

Solar axion detection and the inverse Primakoff process

Jia Liu School of Physics, Peking University

With Christina Gao, Lian-Tao Wang, Xiao-Ping Wang, Wei Xue, Yi-Ming Zhong arXiv: 2006.14598, Phys. Rev. Lett. **125**, 131806

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The axion and the Strong CP problem

$$\mathscr{L} \supset -\frac{\theta g_s^2}{32\pi^2} G\tilde{G} - \left(\bar{u}_L M_u u_R + \bar{d}_L M_d d_R + h.c.\right)$$

- The CKM matrix
 - CP violating phase ~ 1.2 radian
- QCD induced CP violating phase

$$\theta_{\text{QCD}} = \theta + \arg\left[\det\left[M_u M_d\right]\right]$$

• Invariant under quark chiral rotation

According to neutron EDM experiment

$$d_{\text{EDM}}^n \sim \theta \times 10^{-16} \text{ e cm}$$

 $d_{\text{exp}}^n < 10^{-26} \text{ e cm}$

 $\theta_{\rm QCD} \lesssim 1.3 \times 10^{-10}$ radian

Axion: the Strong CP problem and Dark Matter

- Global U(1)_{PQ} symmetry
 - Spontaneous broken leads to massless goldstone (Axion)
- At QCD scale ~ O(1) GeV,
 - Potential from Chiral Lagrangian explicitly breaks the symmetry leads to massive axion
 - Energy stored in coherent oscillation of axion field
 - When mass ~ Hubble, becoming cold dark matter
 - QCD vacuum picks

$$\Theta = \theta_{\rm QCD} + \xi \langle a \rangle / f_a = 0$$





Motivation for axion and axion-like particles

- Strong CP problem
- Dark Matter
- Stellar cooling/ TeV transparency



Experimental approaches

- Dark Matter Axion: haloscopes ...
- Axion independent searches:
 - Rare meson decays
 - Stellar cooling
 - Supernova
 - Helioscopes: solar axion (CAST, IAXO, or DM direct detection searches)
 - Light shining through walls
 - Polarization
 - Fifth force
 - Etc..



1.8

★ RGB stars
▲ AGB stars

RGB track

- RGB fit

11.



Helioscope: solar axion production

• The axion produced in the Sun via photon, electron or nucleon interactions

Atomic recombination and deexcitation, Bremsstrahlung, and Compton (ABC) interactions



Primakoff process



Solar axion production

• The axion produced in the Sun via photon, electron or nucleon interactions



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Solar axion detection

- Axio-electric effect $\mathcal{L} \supset -g_{ae} \frac{\partial_{\mu} a}{2m_e} \bar{e} \gamma^{\mu} \gamma_5 e \frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$
 - keV axions are absorbed by electrons



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The results from axio-electric effect

• Confronting stellar cooling constraints



Inverse Primakoff effect

$$\frac{d\sigma_{a\to\gamma}^{\rm invPrim}}{d\Omega} = \frac{\alpha}{16\pi} g_{a\gamma}^2 \frac{\boldsymbol{q}^2}{\boldsymbol{k}^2} \left(4 - \boldsymbol{q}^2/\boldsymbol{k}^2\right) F_a^2(\boldsymbol{q}^2)$$



Form factor:

$$F_a(q^2) = Zk^2 / (r_0^2 + q^2)$$

Screening effect: $r_0^{-1} = 4 \text{ keV}$



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Photon and electron signal in xenon

- keV photon ionizes Xe.
 - XENON can hardly distinguish keV photon signal from Electron Recoil







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http://nest.physics.ucdavis.edu/benchmark-plots

Inverse Primakoff effect should be included!



2006.14598 Gao, *JL*, Wang, Wang, Xue, Zhong see also 2006.15118 Dent, Dutta, Newstead, Thompson

The effect of Inverse

- An important effect should not be missed
- Alleviated the stellar constraints





The R-parameter

 $R \equiv N_{HB}/N_{RGB} \simeq \tau_{HB}/\tau_{RGB}$



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The R-parameter

$$R \equiv N_{HB}/N_{RGB} \simeq \tau_{HB}/\tau_{RGB}$$

 $R_{\rm th}(g_{a\gamma}, Y) = 6.26 \, Y - 0.41 \, g_{10}^2 - 0.12$

- Used state-of-art stellar evolution models
- Primakoff process emitting axion will decrease R
- sensitive to He abundance Y
- He abundance measurement is challenging
- Choosing observed R from 39 clusters out of 57, biased Robs ?



Stellar cooling anomaly

19 σ?

1406.6053, Ayala, Dominguez, Giannotti, Mirizzi, Straniero

Summary

- The inverse Primakoff effect is important for solar axion
- The stellar cooling tension can be alleviated
- May further alleviate the tension via environment dependent mass (see 2006.15112)
- The uncertainties in the stellar cooling calculation should be checked
- Future experiments like XENONnT, PandaX and LZ can further explore the solar axion properties

Backup slides



2006.15118 Dent, Dutta, Newstead, Thompson