

## "de Sitter Oddities"

work with:

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partially massless  $\rightarrow$  shift symmetric  
 scalars

non scalar bosonic particles in dS:

$$\begin{aligned}
 & \cdot \text{massive} \quad m^2 > s(s-1)H^2 \quad \text{Higuchi bound} \\
 & \left. \begin{array}{l} \cdot \text{massless} \\ \text{gauge} \\ \text{symmetries} \end{array} \right\} \quad m^2 = [s(s-1) - t(t+1)]H^2 \\
 & \quad \text{partially} \\
 & \quad \text{massless} \\
 & \quad \text{depth } t = 0, \dots, s-1 \\
 & m^2 = 0 \iff t = s-1 \\
 & \quad t = 0 \text{ saturates Higuchi} \\
 & \quad \text{bound}
 \end{aligned}$$

e.g.  $s=1$

massive  $m^2 > 0 \quad \pm 1, 0 \text{ helicities}$

massless  $t=0 \quad m^2=0 \quad \pm 1 \text{ helicities}$

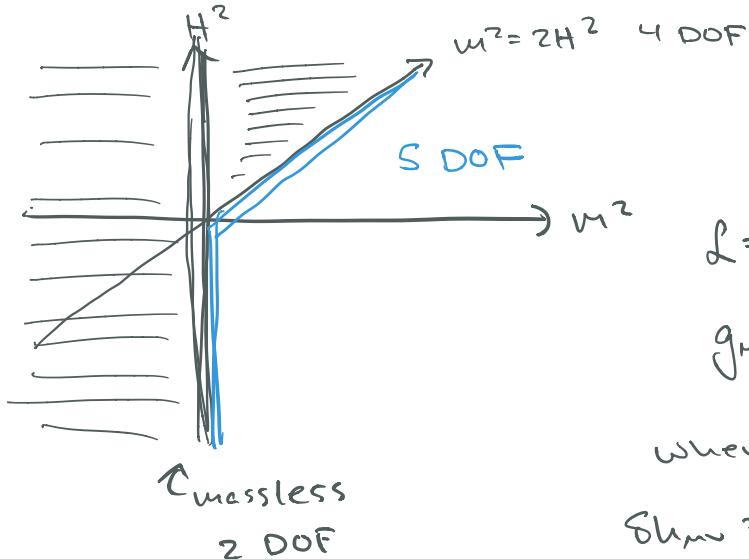
e.g.  $s=2$

massive  $m^2 > 2H^2 \quad \pm 2, \pm 1, 0 \text{ helicities}$

$t=0 \quad m^2=2H^2 \quad \pm 2, \pm 1 \text{ helicities} \leftarrow$  Partially  
 massless

massless  $t=1 \quad m^2=0 \quad \pm 2 \text{ helicities}$

Focus on spin=2



Free theory:

$$\mathcal{L} = \mathcal{L}_{EH}^{(2)} - \frac{m^2}{2} (h_{\mu\nu} h^{\mu\nu} - h^2)$$

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$$

$$\text{when } m^2 = 2H^2$$

$$\delta h_{\mu\nu} = (\nabla_\mu \nabla_\nu + H^2 g_{\mu\nu}) \alpha$$

invariant field strength:

$$F_{\mu\nu\rho} = \nabla_\mu h_{\nu\rho} - \nabla_\nu h_{\mu\rho}$$

$$f = -\frac{1}{4} (F^{\lambda\mu\nu} F_{\lambda\mu\nu} - 2 F^{\lambda\mu}{}_\mu F_{\lambda\mu}{}^\nu)$$

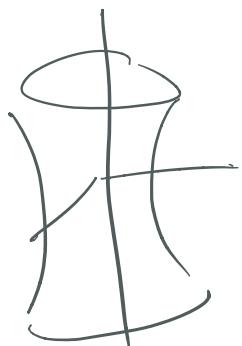
conserved charges  $\longleftrightarrow$  "reducibility parameters"

$$\delta h_{\mu\nu} = 0 = (\nabla_\mu \nabla_\nu + H^2 g_{\mu\nu}) \alpha$$

embedding space  $X^A$

$$\partial_A \partial_B \alpha = 0$$

$$\alpha = C_A X^A$$



pull back:

$$\alpha^A(\kappa) = X^A(\kappa)$$

- 5 conserved currents:  $j_{\mu\nu}^A = F_{\mu\nu}^\lambda \nabla_\lambda \alpha^A$

$$\nabla^\mu j_{\mu\nu}^A = 0$$

- 5 conserved charges  $Q^A = \oint_{S^L} j^A$

invariant charge: take the ds invariant norm  
 2 possibilities depending on whether norm  
 is timelike or spacelike

Flat space limit

$$m^2 = 2H^2 \rightarrow 0$$

$$\mathcal{L}_{PM} \rightarrow \mathcal{L}_{EH}^{(2)} + \mathcal{L}_{EM}$$

$$\pm 2 \quad \pm 1 \quad \nwarrow_{U(1)} \quad \delta A_\mu = \partial_\mu \Lambda$$

$$\mathcal{L} \quad \underline{\underline{g_{\mu\nu} = \partial_\mu \xi_\nu + \partial_\nu \xi_\mu}}$$

reducibility:

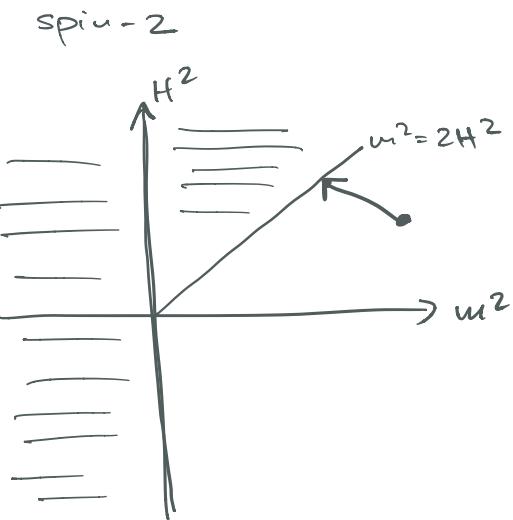
$$(D_\mu D_\nu + H^2 g_{\mu\nu}) \alpha = 0 \rightarrow \partial_\mu \partial_\nu \alpha = 0 \quad \alpha = \text{const}$$

$$\partial_\mu \alpha = \xi_\mu = \text{const}$$

2 invariant conserved charges:

mass, electric charge

$\uparrow$  translations



$$m^2 \rightarrow 2H^2$$

$$L = L_{EH}^{(2)} - \frac{m^2}{2} (h_{\mu\nu} h^{\mu\nu} - h^2)$$

$\leq 0$  DOF

take  $m^2 \rightarrow 2H^2$

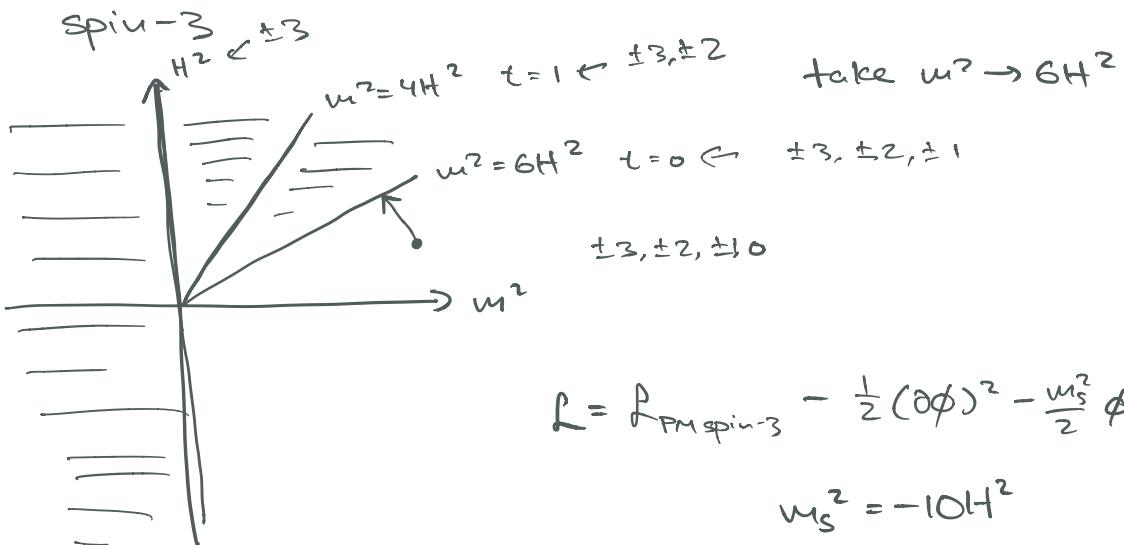
$$L = L_{PM} - \frac{1}{2} (\partial\phi)^2 - \frac{m_s^2}{2} \phi^2$$

$$\begin{matrix} \pm 2 \\ \pm 1 \\ \pm 0 \end{matrix}$$

$$m_s^2 = -4H^2$$

$$\delta g_{\mu\nu} = (\nabla_\mu \nabla_\nu + H^2 g_{\mu\nu}) \alpha$$

embedding space of  $\Phi \rightarrow \delta \Phi = C_A X^A$   
 $\equiv$  embedding space coords



$$L = L_{PM \text{ spin-3}} - \frac{1}{2} (\partial\phi)^2 - \frac{m_s^2}{2} \phi^2$$

$$m_s^2 = -10H^2$$

$$\delta g_{\mu\nu\lambda} = (\nabla_\mu \nabla_\nu \nabla_\lambda + 4H^2 g_{\mu\nu} \nabla_\lambda) \alpha$$

$$\delta \Phi = C_{AB} X^A X^B$$

Generally in de Sitter:

$$\mathcal{L} = -\frac{1}{2}(\partial\phi)^2 - \frac{m_k^2}{2}\phi^2$$

$$m_k^2 = -k(k+3)H^2$$

$$S[\phi] = C_{A_1 \dots A_k} X^{A_1} \dots X^{A_k}$$

↑ const, symm, traceless matrices

$$AdS: m^2 = \frac{1}{L^2} k(k+3)$$

dS:

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Flat spacetime:

$$S = \int d^4x \left( -\frac{1}{2}(\partial\phi)^2 \right) \quad \text{free massless scalar}$$

enjoys infinite number of shift sym:

$$\delta\phi = c + c_\mu x^\mu + c_{\mu\nu} x^\mu x^\nu + \dots$$

$$= \sum_{k=0}^n c_{\mu_1 \dots \mu_k} x^{\mu_1} \dots x^{\mu_k}$$

Add interactions to preserve subset of sym:

$$\bullet k=0 \quad \delta\phi = c$$

$$X = -\frac{1}{2}(\partial\phi)^2 \quad \mathcal{L} = P(X)$$

$$\lim_{P \rightarrow 0} A(P) \propto P^\sigma \quad \sigma = 1$$

- $k = 1 \quad S\phi = C_r x^r$ 
  - galileons  $\rightarrow$  abelian
  - DBI action  $\rightarrow$  nonabelian  $S\phi = C_r x^r + C^r \phi \partial_r \phi$

$$\sigma = 2$$

- $k = 2 \quad S\phi = C_{\mu\nu} x^\mu x^\nu$ 
    - "special galileon" Cheung et al 2014  
Hinterbichler & Joyce 2015
  - $\sigma = 3$
  - $k > 3$
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