

# Effective Theories of Black Hole Dynamics

Roberto Emparan  
ICREA & UBarcelona

*All Things EFT*

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**EFT,**

or

the Search for a Separation of  
Scales

# EFT

- A scheme to systematically parametrize our ignorance about short-distance dynamics
- An approach to simplify the long-distance physics of a known complex theory

# EFT

In this talk

- An approach to simplify the long-distance physics of a known complex theory

*General Relativity*

Not *low-energy EFT*,  
but *long-wavelength EFT*

Quantum EFT: *length*  $\sim \frac{\hbar}{E}$

Gravitational classical EFT: *length*  $\sim GE$

(assume relativistic invce)

## EFT & separation of scales

Integrate out (solve) short-distance dynamics  
to obtain coefficients/functions of long-distance  
state in a systematic expansion

Can be classical or quantum

# Long-wavelength states can differ

Fluctuations around:

- vacuum
- thermal state
- finite-energy localized object (soliton, **black hole**)

# Black holes and Black branes

Two problems amenable to EFT:

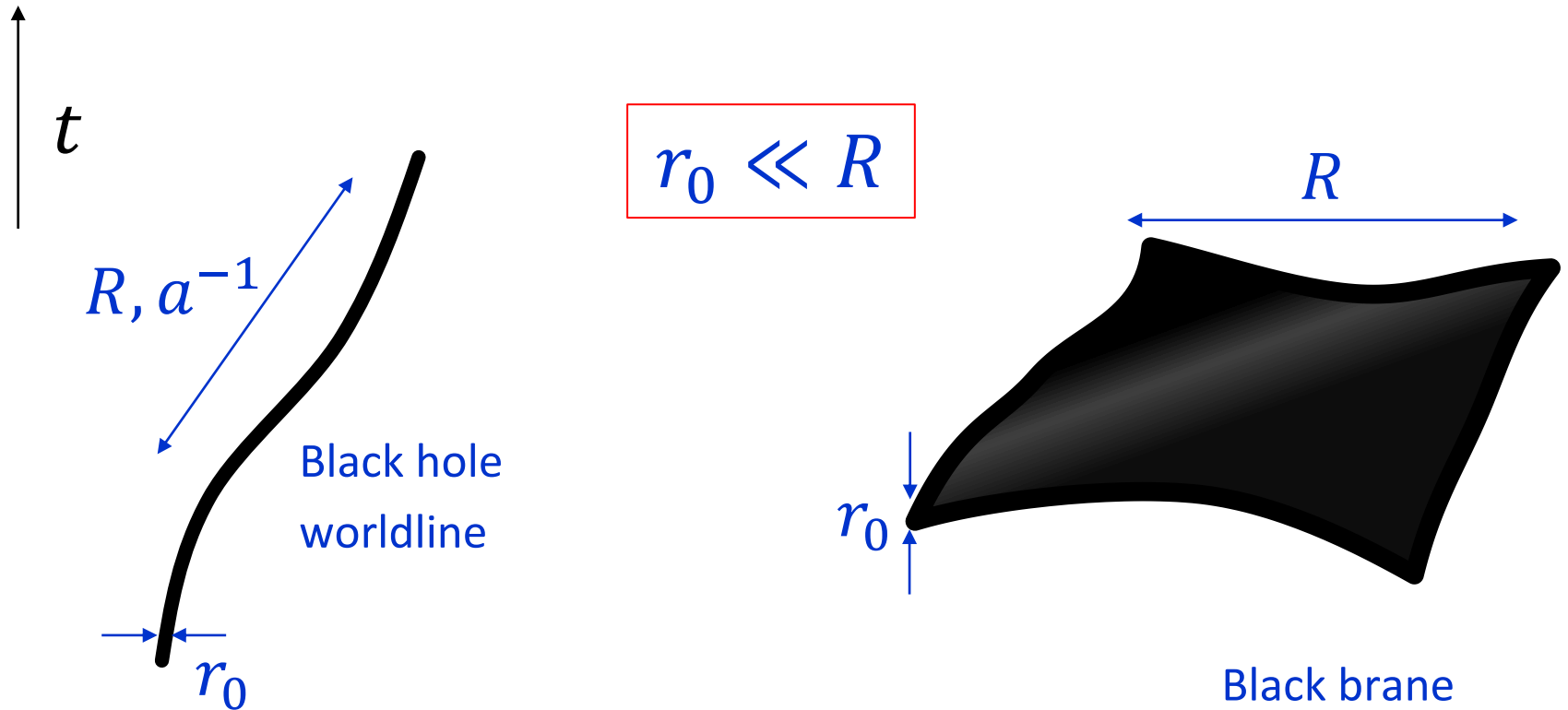
Motion in spacetime

Horizon fluctuations



# Motion in spacetime

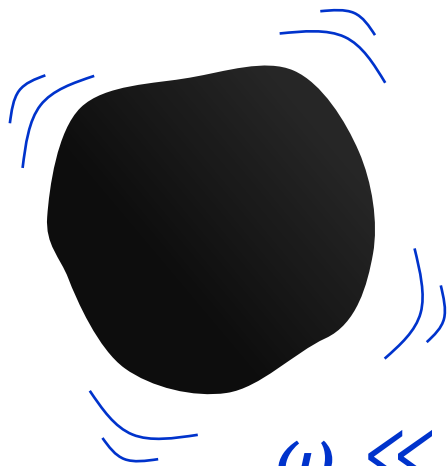
## Worldline/worldvolume dynamics



# Horizon fluctuations

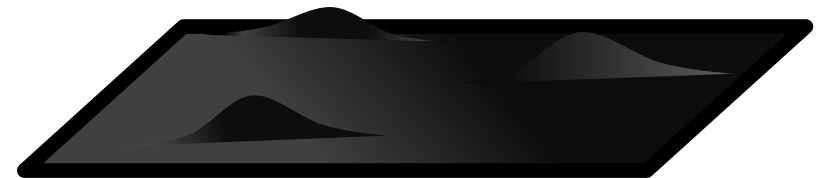
## Quasinormal ringing & Stability

Non-linearly



$$\omega \ll \kappa \sim 1/r_0$$

surface gravity  $\nearrow$  (?)



$$\lambda \gg r_0$$

BHs and Black branes are dissipative

absorption, qnm ringdown, tidal distortion

Effective eqns of mo and stress tensor

rather than effective action

Symmetry and conservation

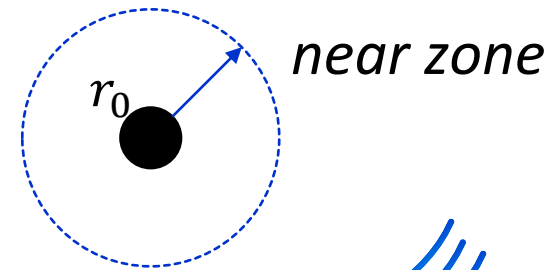
$$\nabla_{\mu} T^{\mu\nu} = 0$$

$T_{\mu\nu}$  : gradient expansion

worldline/worldvolume covariance

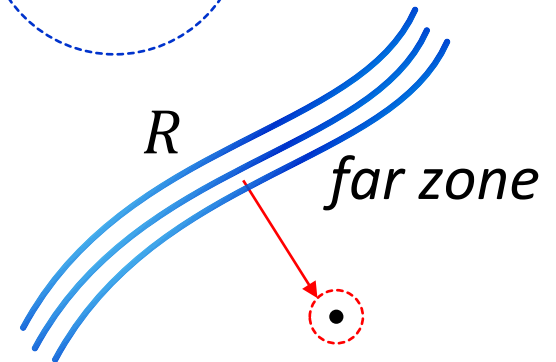
# BH worldline theory

bh →  $r_0 \ll R$  ← background

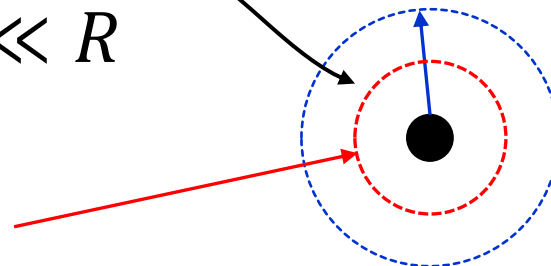


$$ds_{(near)}^2 = ds^2(Schw) + O(r/R)$$

$$ds_{(far)}^2 = ds^2(bckg) + O(r_0/r)$$



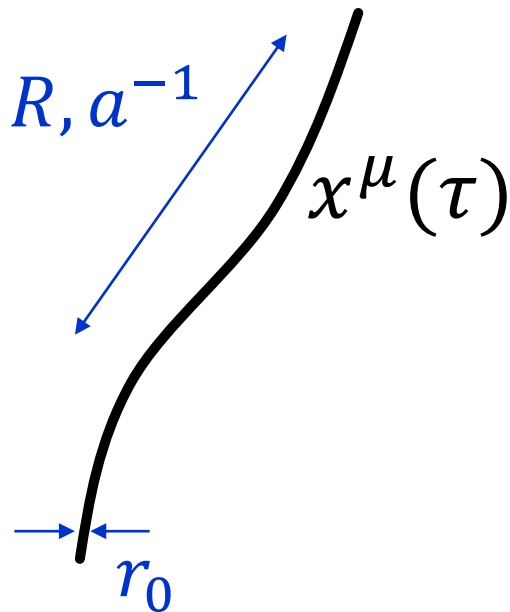
overlap:  $r_0 \ll r \ll R$



In far,  $r \gg r_0$ : BH  $\rightarrow$  worldline  $x^\mu(\tau)$

$$T^{\mu\nu} = c_1(x(\tau)) \dot{x}^\mu(\tau) \dot{x}^\nu(\tau) + \dots$$

(add tetrad for spin)



$$T_{\mu\nu} = c_1(x(\tau)) u_\mu(\tau) u_\nu(\tau) + \dots$$

(minimal coupling)

$$u^\mu = \dot{x}^\mu(\tau)$$

## Symmetries

General coord invce

Worldline reparametrization

Internal Lorentz invariance

$$T^{\mu\nu} = c_1(x(\tau)) u_\mu(\tau) u_\nu(\tau)$$

Effective eom

$$\nabla_\mu T^{\mu\nu} = 0$$

*v* parallel to worldline  $\Rightarrow \partial_\tau c_1 = 0 \Rightarrow c_1 = m$

*v* orthogonal to worldline  $\Rightarrow m a^\mu(\tau) = 0$



*m* is matched to BH mass in overlap zone

Coefficients in effective stress tensor obtained  
by integrating out short distance physics



# BH worldline theory

Derivation from Einstein's equations

- match asymptotics: perturb BH with generic weak asymptotic field (*far* field sources)
- radial (Gauss-Codazzi) constraint  $\rightarrow$  eff eqs

# BH worldline theory

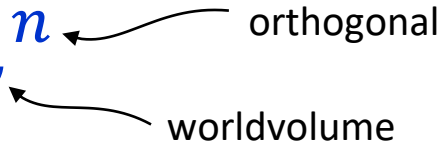
EFT coefficients can also be obtained  
diagrammatically

*Classical* Feynman diagrams with on-shell external legs, no  
virtuality in loops

At lowest orders both methods are similarly  
simple

# Black Brane Worldvolume Theory

$$x^\mu(\tau, \sigma^i)$$

- Derivatives *orthogonal* to worldvolume:  
extrinsic curvature  $K_{\mu\nu}^n$   


- Derivatives *parallel* to worldvolume:  
worldvolume velocities  $u^i$

worldvol Lorentz broken in black brane  
not purely tensional brane, can support wv velocities

# Symmetries

General coord invce

Worldvolume reparametrization

Worldvolume translations + rotations

$$T^{\mu\nu} = c_1(\tau, \sigma^i) u^\mu u^\nu + c_2(\tau, \sigma^i) \gamma^{\mu\nu}$$

↖  
wv induced metric

$$T^{\mu\nu} = c_1(\tau, \sigma^i) u^\mu u^\nu + c_2(\tau, \sigma^i) \gamma^{\mu\nu}$$

Effective eom

$$\nabla_\mu T^{\mu\nu} = 0$$

Orthogonal to  $wv \Rightarrow T^{\mu\nu} K_{\mu\nu}^n = 0$



$$T^{\mu\nu} = c_1(\tau, \sigma^i) u^\mu u^\nu + c_2(\tau, \sigma^i) \gamma^{\mu\nu}$$

Effective eom

$$\nabla_\mu T^{\mu\nu} = 0$$

Parallel to wv  $\Rightarrow$  Relativistic fluid equations

$$\text{with } c_1 = \varepsilon + P \quad c_2 = P$$

$\varepsilon(\tau, \sigma^i)$ ,  $P(\tau, \sigma^i)$ : effective *energy density & pressure*

Coefficients in effective stress tensor  $\varepsilon, P$   
from integrating out short-distance physics

Match *near* to *far* at  $r_0 \ll r \ll R$

⇒ Effective equation of state of black  $p$ -brane

$$P = -\frac{\varepsilon}{D - p - 2}$$

Extrinsic dynamics EFT

Elasticity theory



# Extrinsic dynamics EFT

## Elastic membrane

Example: Dirac-Nambu-Goto pure-tension brane

$$T^{\mu\nu} = -T\gamma^{\mu\nu} \Rightarrow \gamma^{\mu\nu} K_{\mu\nu}{}^n = 0$$

minimal (extremal) surface

*Blackfolds* are elastic, but not minimal surfaces

$$T^{\mu\nu} = (\varepsilon + P)u^\mu u^\nu + P\gamma^{\mu\nu}$$

# Extrinsic dynamics EFT

Elastic solid membrane

Bent brane: Young modulus

*Armas + Camps, Gath, Harmark, Obers*

# Intrinsic dynamics EFT

## Hydrodynamics

Fluid on the brane

$$T_{\mu\nu} = T_{\mu\nu}^{(0)} + T_{\mu\nu}^{(1)} + \dots$$

zero gradient

perfect fluid

one gradient

viscous fluid

# Black Branes as Fluids

## Fluid/gravity correspondence

A classical limit of AdS/CFT

# Black Branes as Fluids

## Fluid/gravity correspondence

Universality of hydrodynamics as EFT

Long-wavelength fluctuations of *thermal state of CFT*:  
hydrodynamics

Long-wavelength fluctuations of *black brane in AdS*:  
hydrodynamics

*Bhattacharyya+Hubeny+Minwalla+Rangamani 2007*

# Fluid/gravity correspondence

can be derived from classical gravity in AdS  
without assuming AdS/CFT

Uniform black brane:  
constant temperature & velocity

Allow them to vary along brane, correcting the  
metric to satisfy Einstein's eqs

Solve Einstein: integrate radial direction  
Obtain eq of state, transport coeffs

# Black branes as fluids: highlights

- AdS branes: Viscosity of “quark-gluon-like” plasma *Policastro+Son+Starinets+Kovtun*
- AF branes: Gregory-Laflamme instability analytically solved as hydro instability
- Correlated stability conjecture *Gubser+Mitra*  
reduces to  
thermodynamic instability  $\Rightarrow$  hydrodynamic instability



# Limitations

Effective description of horizon fluctuations  
requires *long horizon* to support long  
wavelengths

$$\lambda \gg |\partial_r|^{-1}$$

needed for integrating radial dependence

$|\partial_r|^{-1}$  : surface gravity, temperature

$$\lambda \gg \frac{1}{\kappa}, \frac{1}{T}$$

acceleration length, thermal length

For Schw or Kerr (or non-ultraspinning bhs)

Horizon length  $\sim$  Thermal length

$$r_0 \sim |\partial_r|^{-1}, T^{-1}, \kappa^{-1}$$

$$\lambda \simeq r_0$$



$$\omega \ll \kappa \sim 1/r_0$$

Can't have long wavelengths or low frequencies on horizon

No EFT for black hole fluctuations

In more detail

$$T^{-1} = \frac{4\pi}{D-3} r_0$$

$\Rightarrow T^{-1} \sim r_0$  : only scale in system

$\Rightarrow$  ~~A~~ separation of scales, ~~A~~ EFT



In more detail – try again

$$T^{-1} = \frac{4\pi}{D - 3} r_0$$

$$D \gg 1$$

$\Rightarrow T^{-1} \sim r_0/D \ll r_0$  : two scales

$\Rightarrow \exists$  separation of scales,  $\exists$  EFT



*Asnin+Gorbonos+Hadar+Kol+Levi 2007*

*RE+Suzuki+Tanabe 2015*

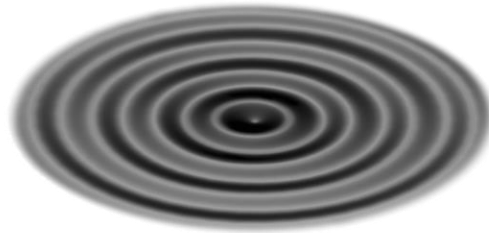
*RE+Herzog RevModPhys 2020*

Study BH fluctuations in an expansion in  $1/D$

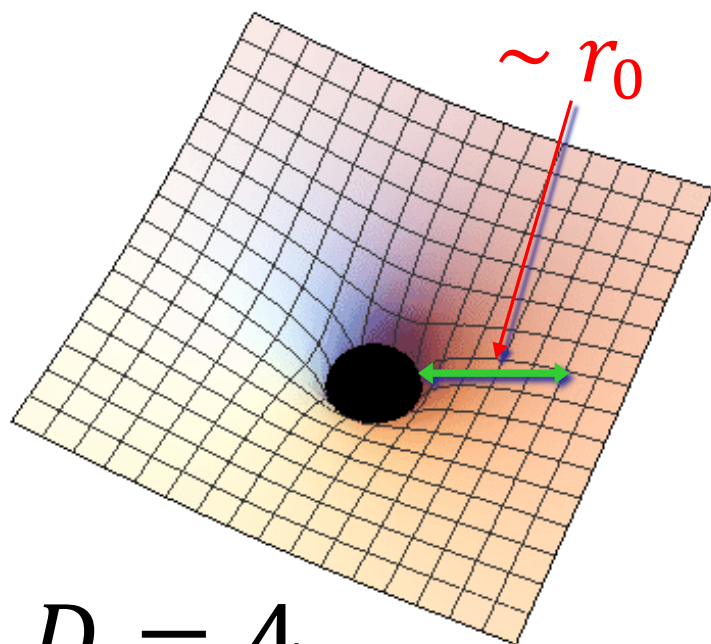
What kind of EFT?

Elasticity theory or Hydrodynamic theory?

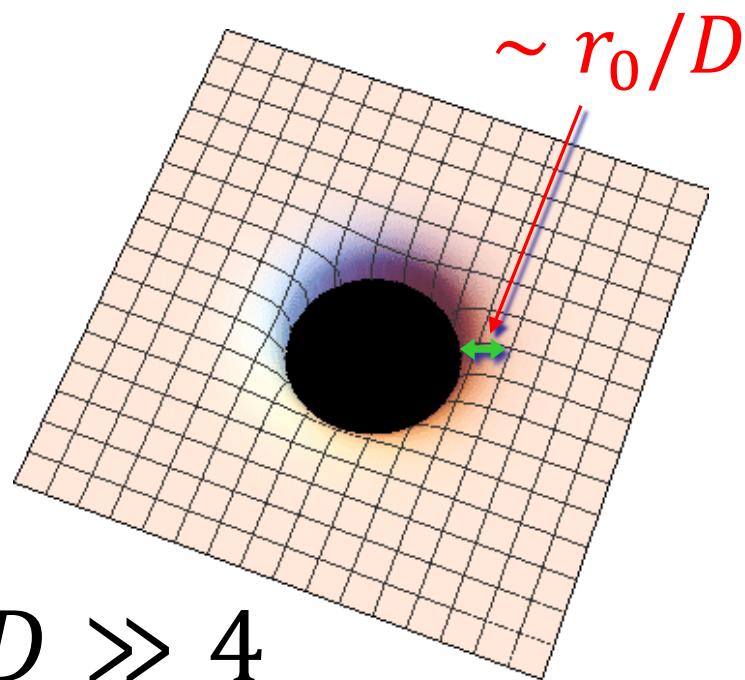
Ripples on a pond, or wrinkles on a membrane?



Not *or*, but *and*



$D = 4$

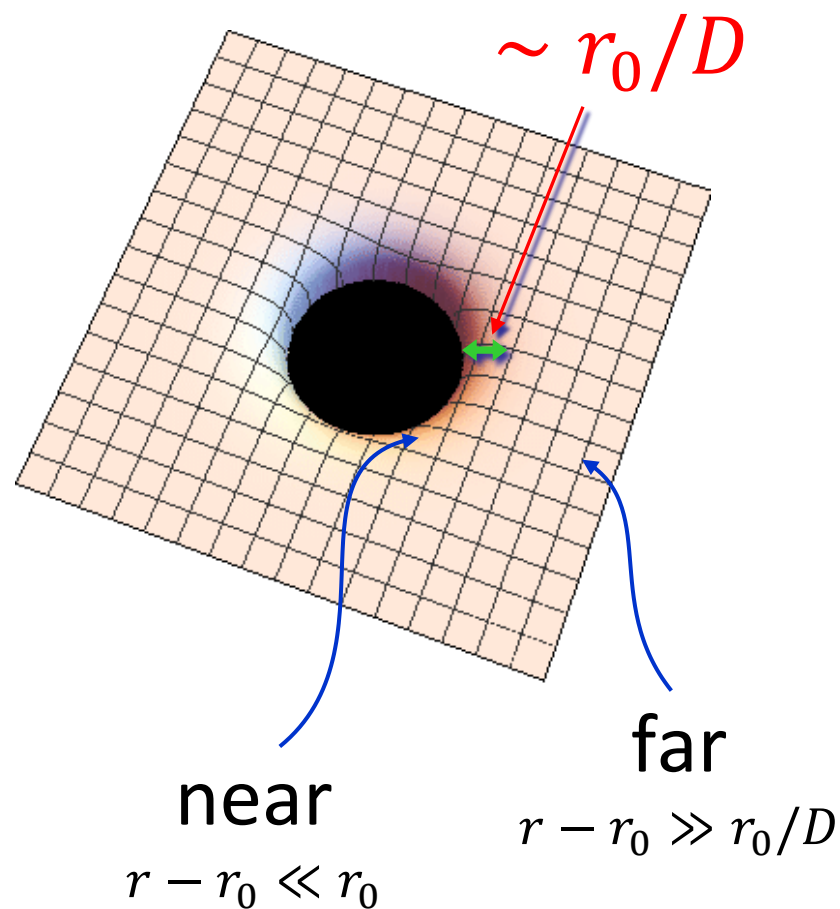


$D \gg 4$

*So near, so far...*

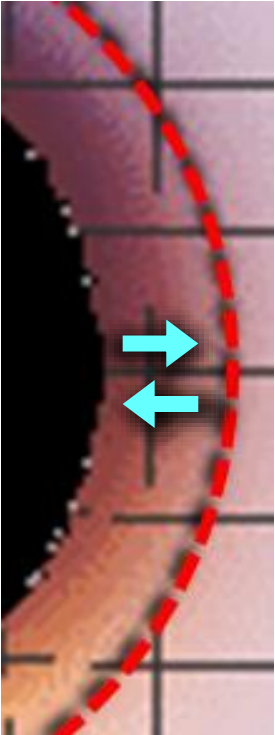
*Near* and *far* defined independently of any background length scale

Just  $r_0 \gg r_0/D$   
intrinsic black hole  
scales





# Slowest fluctuations of BH



$$\omega \sim \frac{1}{r_0} \ll \frac{D}{r_0}$$

Almost static in near-horizon region

Decoupled from far zone

∃ parametrically slow fluctuations

Find **EFT for non-linear slow fluctuations**

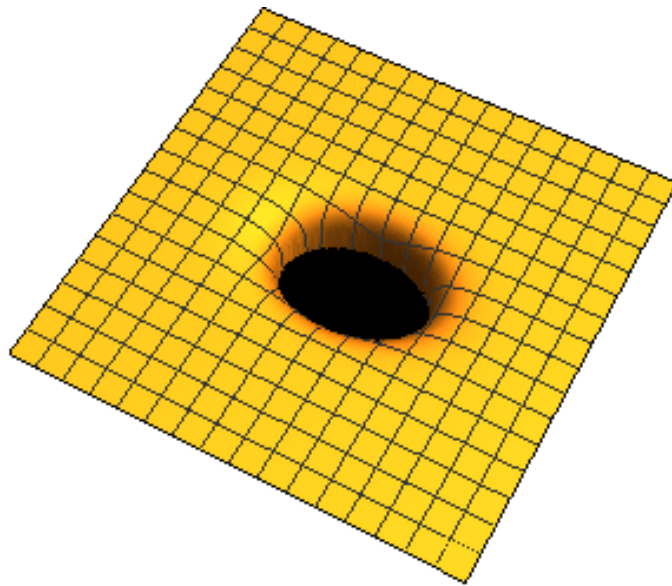
Integrate radial direction away from horizon  
ie, integrate out near-horizon dynamics

# Gradient hierarchy

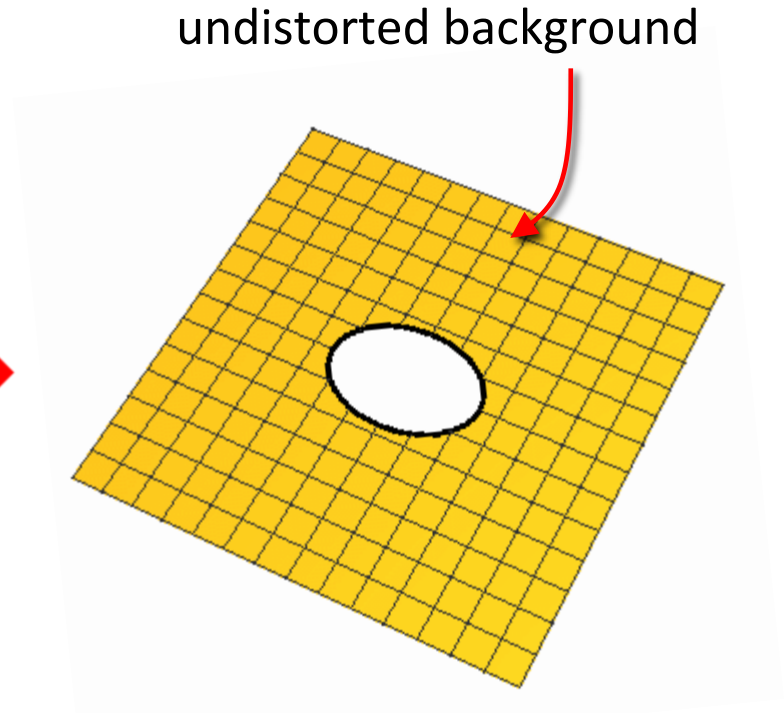
⊥ Horizon:  $\partial_r \sim D$

∥ Horizon:  $\partial_t, \partial_z \sim 1$  (or  $\sim \sqrt{D}$ )

# Replace BH $\rightarrow$ 'effective membrane'

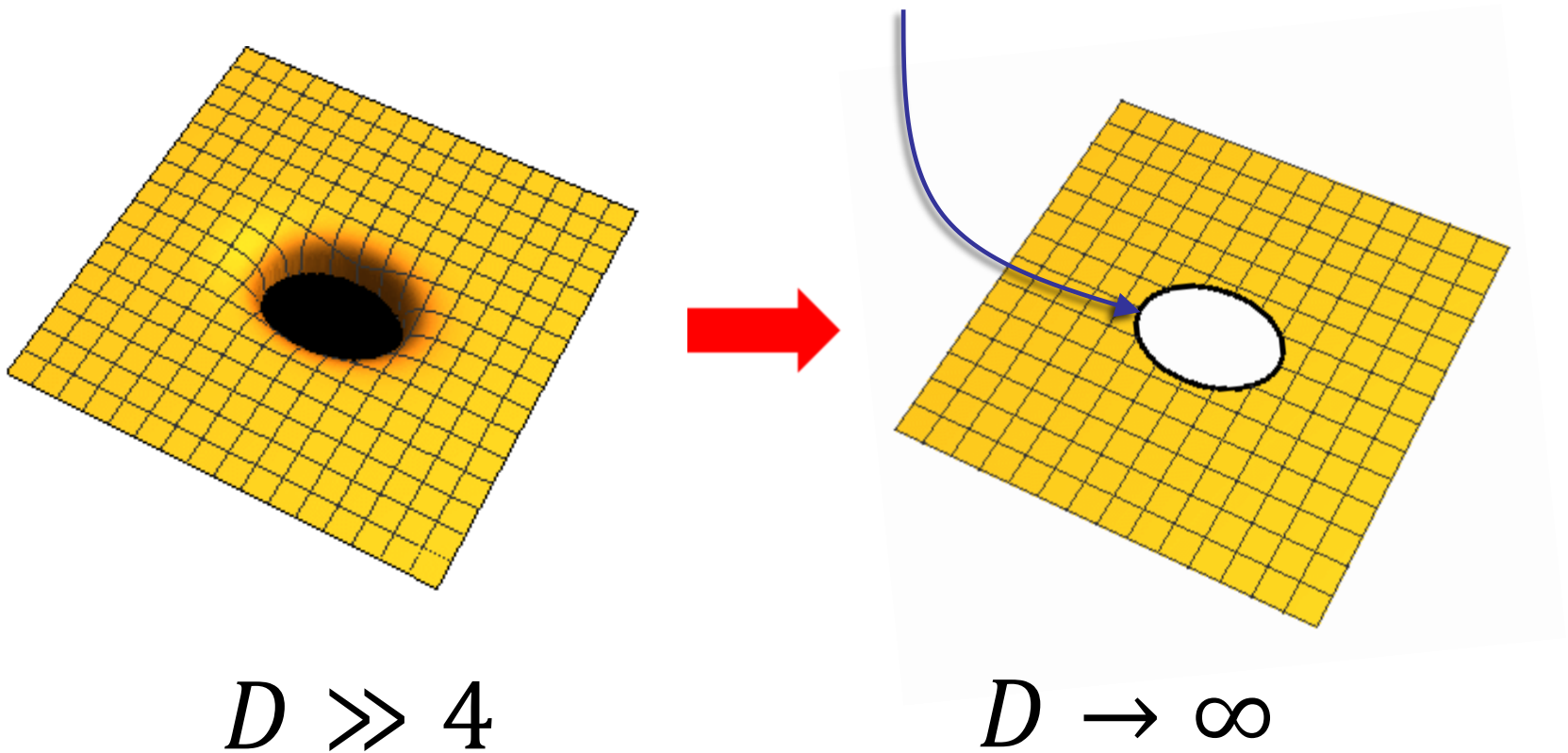


$$D \gg 4$$



$$D \rightarrow \infty$$

What's the dynamics\* of this membrane?



\*what shapes? what evolution?

# Effective fields

- Horizon shape (embedding into background)

$$X(\sigma)$$

w/ finite number of non-trivial horizon directions  $\sigma$

- Horizon velocity

$$u(\sigma)$$

Effective equations  
for  $X(\sigma), u(\sigma)$

Most general & elegant formulation by  
*Bhattacharyya+Minwalla et al 2015*



$$\left( \frac{\nabla^2 u}{\mathcal{K}} - \frac{\nabla \mathcal{K}}{\mathcal{K}} + u \cdot K - (u \cdot \nabla)u \right) \cdot \mathcal{P} = 0$$

$$\mathcal{K} = \eta^{AB} K_{AB}$$

$$\nabla \cdot u = 0, \quad n \cdot u = 0$$

$$\mathcal{P}_{AB} = \eta_{AB} - n_A n_B + u_A u_B$$

$n, K_{AB}$ : normal & **extrinsic curvature** of membrane

$u$ : **velocity** field on membrane

## Simplifies

–conceptually and technically–  
in two important cases:

1. Stationary black holes
2. Black branes, AdS or AF

# Stationary solution

Soap-bubble equation

$$K = 2\gamma\kappa$$

$K$  = trace **extrinsic curvature**  
of membrane

$\gamma$  = **redshift** on membrane

$\kappa$  = **surface gravity**



# Effective equations for fluctuating black brane

*RE+Suzuki+Tanabe 2015*

Effective fields:

$m(t, \sigma^i)$  : mass and area density of black brane

$v_i(t, \sigma^j)$  : velocity along brane

# Effective equations for fluctuating black brane

$$\partial_t m + \partial_i (m v^i) = 0$$

$$\partial_t (m v_i) + \partial^j (\underbrace{\pm m \delta_{ij}}_{\text{pressure}} + \underbrace{m v_i v_j}_{\text{viscosity}} - 2 m \partial_{(i} v_{j)} - \underbrace{m \partial_{ij}^2 \ln m}_{\text{higher transport}}) = 0$$

pressure

viscosity

higher transport

Hydrodynamics truncates exactly

Large D explains unexpected success of hydro?

# Effective equations for fluctuating black brane

$$\partial_t m + \partial_i (m v^i) = 0$$

$$\partial_t (m v_i) + \partial^j (\pm m \delta_{ij} + m v_i v_j - 2 m \partial_{(i} v_{j)} - m \partial_{ij}^2 \ln m) = 0$$

Can be rewritten as Soap Bubble eqn  $K = 2\gamma\kappa$

with  $m(t, \sigma^i)$  local bubble radius

(recall  $\ell \sim GE!$ )

# BH “membrane paradigm”

*Damour*

*Thorne et al*

Not an EFT

Clever & suggestive way of writing boundary conditions for the gravitational field on a null hypersurface

No separation of scales

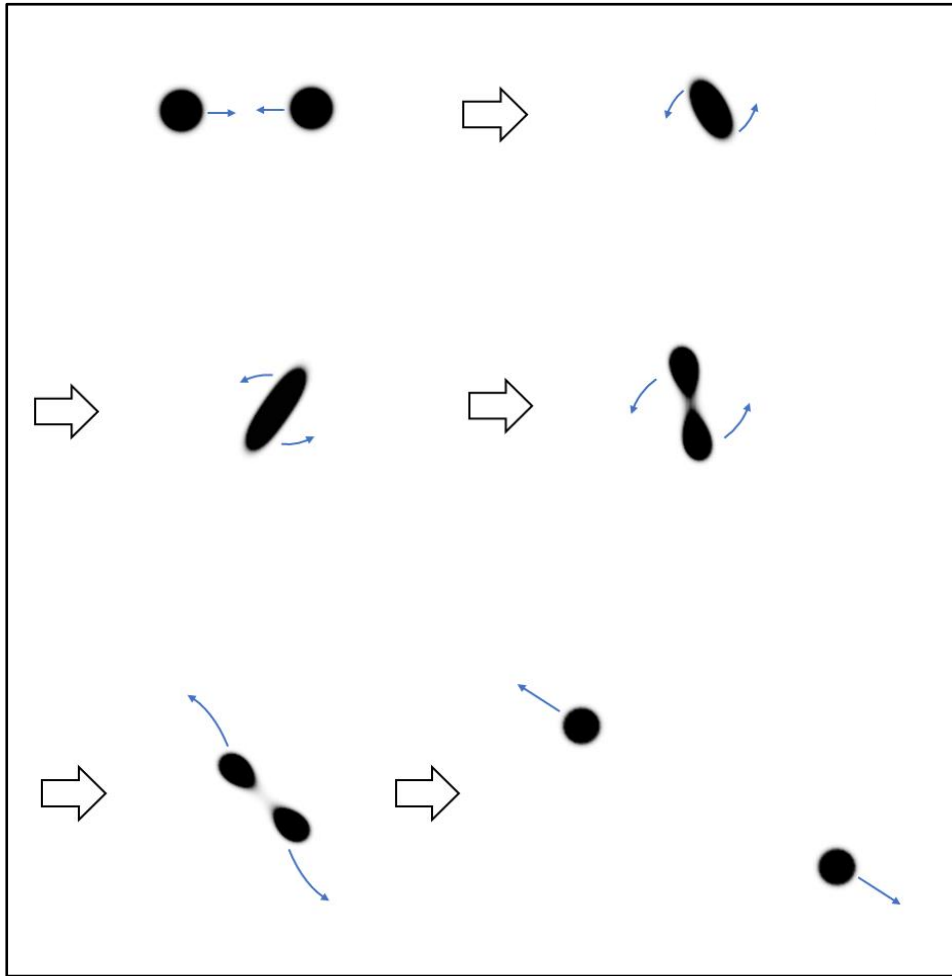
No integration of any short-scale degrees of freedom

Highlight of large  $D$  effective theory:

Black hole collision in higher  $D$  w/  
cosmic censorship violation

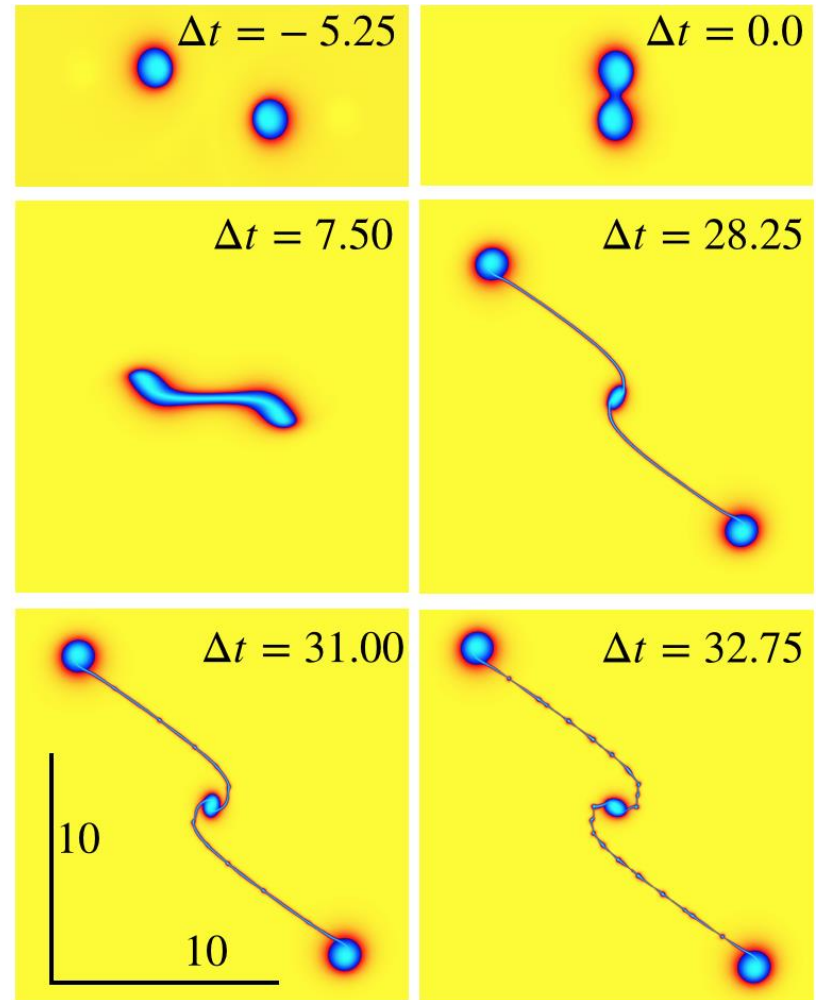






## Large-D EFT

*Andrade+RE+Licht+Luna 2019*



## D=6 NumGR

*Andrade+Figuerras+Sperhake 2020*

# Final remarks

- EFT philosophy and concepts can be very usefully applied to classical theories

Separation of scales, integrating short-distance dynamics, matching conditions, expansion in derivatives...

# Final remarks

- Diagrammatic techniques, renormalization group?
  - So far of little use beyond BH worldline theory
- Matched asymptotics works well enough

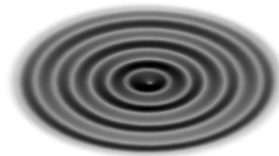
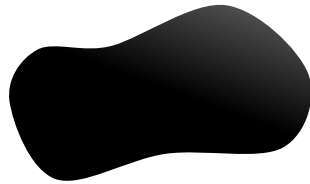
# Final remarks

- Universal EFTs: hydrodynamics, elasticity apply also to black holes

Different EFTs for different dynamical regimes

# Final remarks

EFTs very efficient for isolating and reformulating slow **non-linear dynamics of horizons**



Thank you

## *Brane blobology*

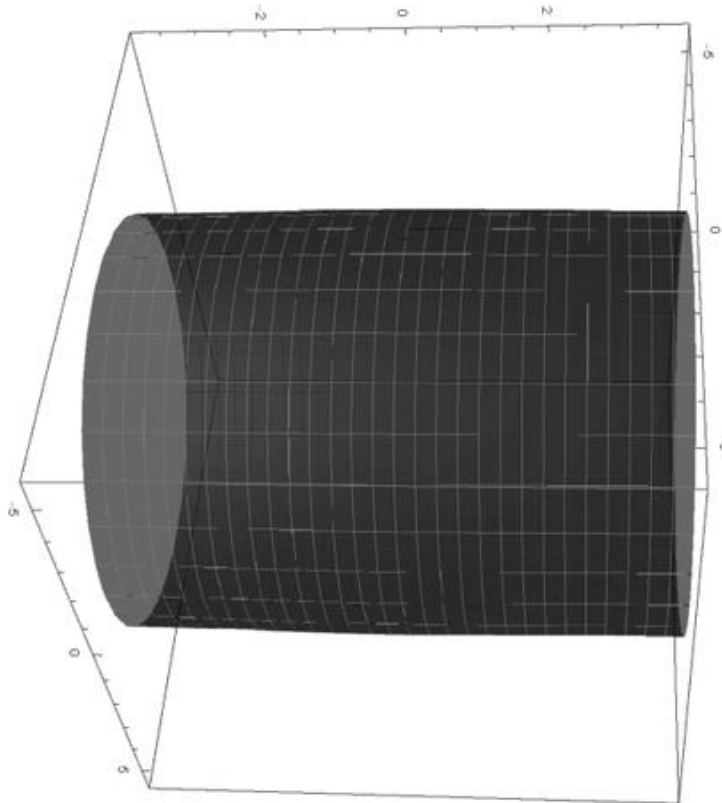
*Very effective* theory of black holes as  
blobs on a brane



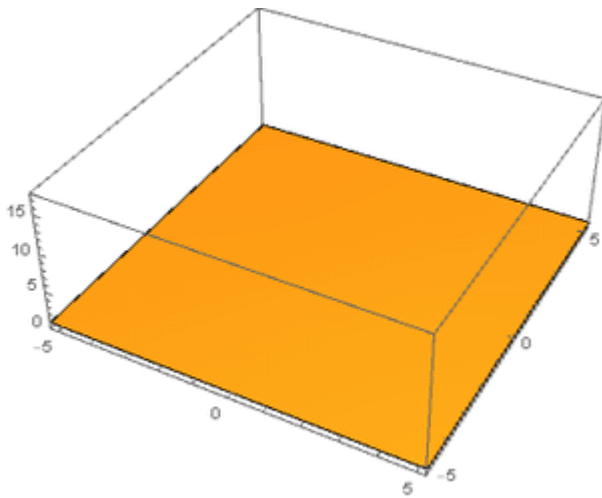
# Black hole as blob on a string or brane

Black string & black brane instability *Gregory-Laflamme 1993*

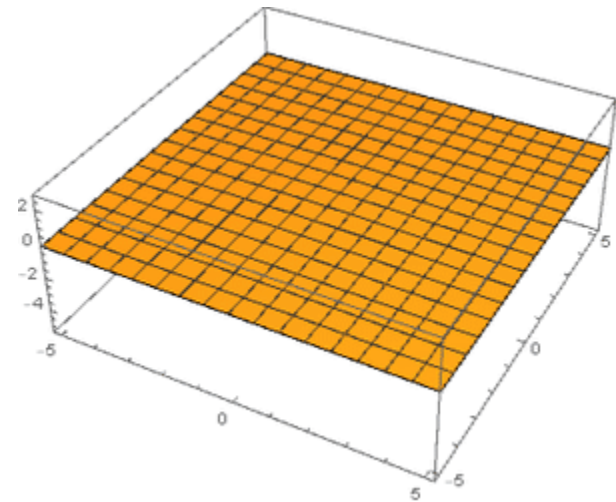
Evolved w/ large- $D$  effective theory



# “Black hole blob” in a black membrane

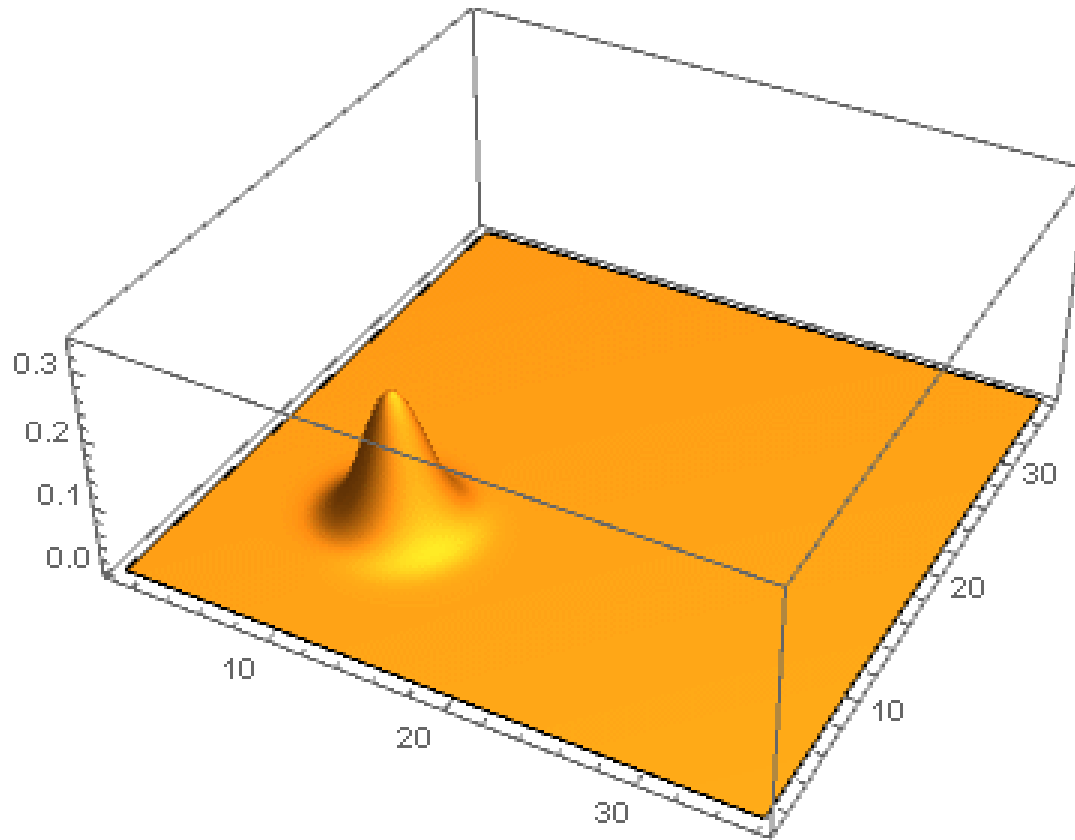


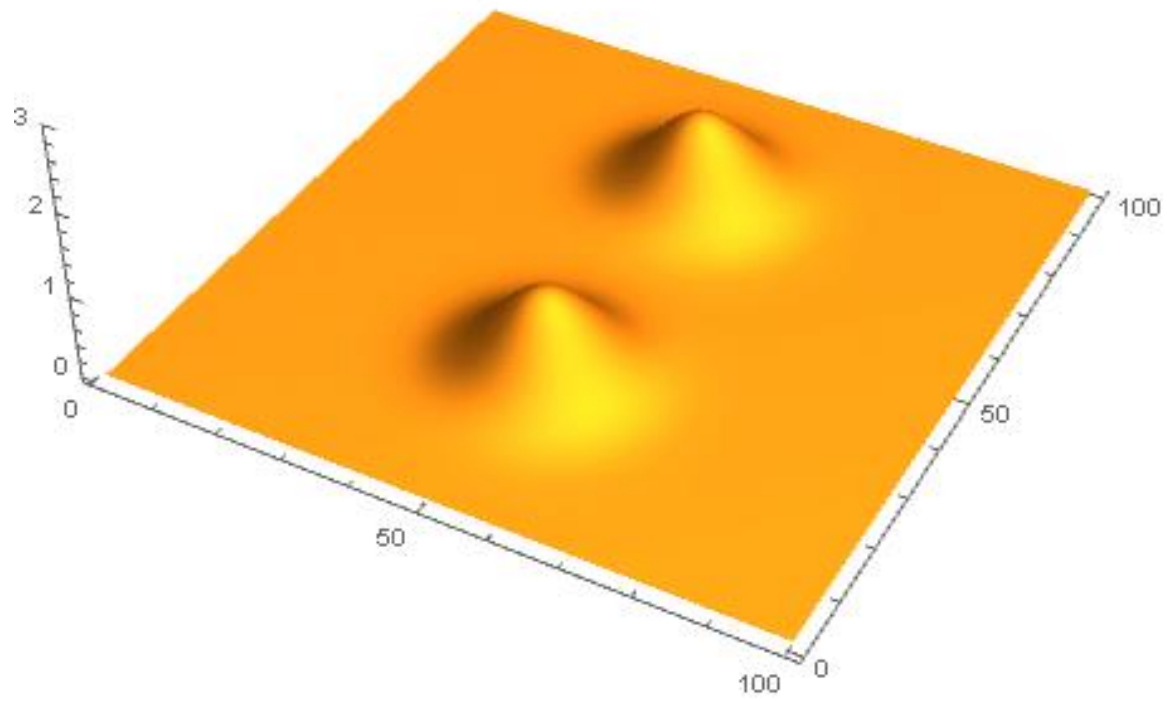
Area density  
 $m$



Area radius  
 $r_H = \ln m$

# Moving black hole





# Collisions of black hole blobs

Brane acts as a “regulator”:  
continuous horizon

BHs never really merge nor split:  
smooth evolution