

陈龙斌 广州大学 2021-5-27

微分方程方法:



A. V. Kotikov, Differential equations method: New technique for massive Feynman diagrams calculation, Phys. Lett. **B254** (1991) 158–164.

A. V. Kotikov, Differential equation method: The Calculation of N point Feynman diagrams, Phys. Lett. **B267** (1991) 123–127. [Erratum: Phys. Lett.B295,409(1992)].

Multiloop integrals in dimensional regularization made simple

Johannes M. Henn (Princeton, Inst. Advanced Study) (Apr 5, 2013)

Published in: *Phys.Rev.Lett.* 110 (2013) 251601 • e-Print: 1304.1806 [hep-th]

Choosing canonical basis (basis with uniform transcendentality)

For random basis g, we may have:

$$\partial_x \vec{g}(x;\epsilon) = B(x,\epsilon) \, \vec{g}(x;\epsilon)$$

We can choose new basis f:

 $\vec{f} = T\vec{g},$



$$d \vec{f}(x,\epsilon) = \epsilon \left(d \tilde{A} \right) \vec{f}(x;\epsilon)$$

$$\tilde{A} = \left[\sum_{k} A_k \log \alpha_k(x)\right] \,.$$

Master Integrals Calculation Example

Two-loop master integrals for heavy-to-light form factors of two different massive fermions

Massive quark decays to massive quark

$$(t \to b + W^+(l + \bar{\nu}), b \to c + l + \bar{\nu})$$

JHEP 02 (2018) 066



Sample of Two-loop Feynman diagrams



Tensor integrals are reduced by IBP (integrateby-parts) method.



$$\begin{split} F_{16} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{16}, \\ F_{17} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{17}, \\ F_{18} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{18}, \\ F_{19} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{19}, \\ F_{20} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{20}, \\ F_{21} &= \epsilon^2 \left(2s \left(\epsilon(M_{19} + 2M_{20}) - m_2^2 M_{21} - 2m_1^2 M_{22}\right)\right) + 2\frac{m_2}{m_1} F_9, \\ F_{22} &= \epsilon^2 \left(2(m_1^2 - m_2^2)(2\epsilon(M_{19} + 2M_{20}) + m_2^2 M_{21} - 2m_1^2 M_{22})\right) + 2\frac{m_1}{m_2} F_7 - 2\frac{m_2}{m_1} F_9, \\ F_{22} &= \epsilon^2 \left(2(m_1^2 - m_2^2)(2\epsilon(M_{19} + 2M_{20}) + (m_1^2 + m_2^2 - s)M_{21} - 4m_1^2 M_{22})\right) \\ &+ 2\frac{m_1}{m_2} F_7 - 2\frac{m_2}{m_1} F_9, \\ F_{23} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{23}, \\ F_{24} &= \epsilon^3 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} M_{24}, \\ F_{25} &= \epsilon^2 \frac{s m_2^2(s - (m_1 + m_2)^2)(s - (m_1 - m_2)^2)}{(s - m_1^2 + m_2^2)} M_{25} \\ &+ \epsilon^3 \frac{(s - (m_1 + m_2)^2)(s - (m_1 - m_2)^2)}{2(s - m_1^2 + m_2^2)} (M_{23} - M_{24}) \\ &- \frac{m_2 s \left(s - m_1^2 - m_2^2\right)}{m_1(s - m_1^2 + m_2^2)} F_9 - \frac{\sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2}}{4(s - m_1^2 + m_2^2)} F_{11} \\ &+ \frac{s m_2^2}{(s - m_1^2 + m_2^2)^2} (F_2 - F_3 - 6F_8 + 2F_{12}), \end{split}$$

Differential Equations In Canonical Form:

$$d\mathbf{F}(x,y;\epsilon) = \epsilon \, \mathrm{d} \, \tilde{A}(x,y) \, \mathbf{F}(x,y;\epsilon)$$

$$\tilde{A}(x,y) = A_1 \ln(x) + A_2 \ln(x+1) + A_3 \ln(x-1) + A_4 \ln(x+y) + A_5 \ln(x-y) + A_6 \ln(xy+1) + A_7 \ln(xy-1) + A_8 \ln(y) + A_9 \ln(y+1) + A_{10} \ln(y-1) + A_{11} \ln(x^2y - 2x + y) + A_{12} \ln(x^2 - 2yx + 1).$$
(4.41)

A_i are rational matrices

$$s = m_1^2 \frac{(x-y)(1-xy)}{x}$$
, and $m_2 = m_1 y$.

Determination of Boundary Conditions

$$F_{10} = \epsilon^2 s M_{10},$$

$$F_{11} = \epsilon^2 \sqrt{s - (m_1 + m_2)^2} \sqrt{s - (m_1 - m_2)^2} (M_{10} + M_{11} + M_{12}),$$

$$F_{12} = \epsilon^2 \left((m_1^2 - m_2^2) (M_{10} + M_{11} + M_{12}) + s (M_{11} - M_{12}) \right),$$

$$\frac{\partial F_{12}}{\partial x} = \epsilon \left(-\frac{F_{11} + F_{12}}{x - y} + y \frac{F_{11} - F_{12}}{x y - 1} + \frac{F_{12}}{x} \right)$$

$$F_{11}|_{x=y} = -F_{12}|_{x=y}.$$
$$F_{11}|_{x=\frac{1}{y}} = F_{12}|_{x=\frac{1}{y}}.$$

Goncharov Polylogarithms

$$G_{a_1,a_2,...,a_n}(x) \equiv \int_0^x \frac{dt}{t-a_1} G_{a_2,...,a_n}(t) ,$$

$$G_{\overrightarrow{0}_n}(x) \equiv \frac{1}{n!} \ln^n x .$$

$$G_{a_1,...,a_m}(x) G_{b_1,...,b_n}(x) = \sum_{c \in a \text{III} b} G_{c_1,c_2,...,c_{m+n}}(x) .$$

A. B. Goncharov, *Multiple polylogarithms, cyclotomy and modular complexes*, Math. Res. Lett. 5, (1998) 497–516, [arXiv:1105.2076].

$$\begin{split} F_{10} &= 2\epsilon^2 \left[G_{0,0}(y) - G_{0,0}(x) \right] + \epsilon^3 \left[\frac{G_{0,0}(y)}{2} (G_{\frac{1}{y}}(x) + G_y(x) + 4G_y(1) - 4G_{\frac{1}{y}}(1) - G_{\frac{1}{y}}(y) \right) \\ &\quad + 4G_0(y) (G_{0,y}(x) - G_{0,\frac{1}{y}}(x) + G_{0,0}(x) + G_{0,\frac{1}{y}}(y) - G_{\frac{1}{y},0}(1) - G_{y,0}(1) + \frac{\pi^2}{3}) \\ &\quad + G_0(x) (\frac{G_{0,0}(y)}{2} - 4(G_{\frac{1}{y}}(1) - G_y(1))G_0(y) - 4G_{\frac{1}{y},0}(1) - 4G_{y,0}(1) + \pi^2) - 6G_{0,0,0}(x) \\ &\quad + 12(G_{0,-1,0}(x) + G_{0,1,0}(x) - G_{0,-1,0}(y) - G_{0,1,0}(y)) - 2(2G_{0,\frac{1}{y},0}(x) + 2G_{0,y,0}(x) \\ &\quad + G_{\frac{1}{y},0,0}(x) + G_{y,0,0}(x)) + 2G_{\frac{1}{y},0,0}(y) + 4G_{0,\frac{1}{y},0}(y) + 6\zeta(3) \right] + \mathcal{O}(\epsilon^4), \\ F_{33} &= \epsilon^3 \left[2G_{0,0,0}(x) - 2G_{0,1,0}(x) - 2G_{0,-1,0}(x) + \frac{1}{6}\pi^2G_0(x) + \zeta(3) \right] + \mathcal{O}(\epsilon^4). \end{split}$$

$$M_{33}^{\text{SecDec}}(-5.4, 1.0, 0.2) = \frac{-0.4466129 \pm 0.0000004}{\epsilon} - 0.507366 \pm 0.000006,$$

$$M_{33}^{\text{FIESTA}}(-5.4, 1.0, 0.2) = \frac{-0.446613 \pm 0.000005}{\epsilon} - 0.507387 \pm 0.000049,$$

$$M_{33}^{\text{Ours}}(-5.4, 1.0, 0.2) = \frac{-0.4466129967 \dots}{\epsilon} - 0.5073683817 \dots$$

Check:

Numerical evaluation of multiple polylogarithms

Ginac

Jens Vollinga (Mainz U., Inst. Phys.), Stefan Weinzierl (Mainz U., Inst. Phys.) (Oct, 2004)

Published in: Comput. Phys. Commun. 167 (2005) 177 • e-Print: hep-ph/0410259 [hep-ph]



Digits = 15: expression = 2*G({0.9393543905251667 + 7.589370620706021e-10*I, 0., 1., -4.696771952625833 + 4.246188272542417e-9*I}, 1) : evalf(expression); quit:

Save as txt file

customer@node01:~/Gnum _ □	×
文件(F) 编辑(E) 查看(V) 搜索(S) 终端(T) 帮助(H) [customer@node01 ^]\$ cd Gnum/ [customer@node01 Gnum]\$ ginsh call.txt ginsh - GiNaC Interactive Shell (GiNaC VI.7.8)	

1. arXiv:1904.07279 [pdf, other] hep-th hep-ph doi 10.1007/JHEP08(2019)135

PolyLogTools - Polylogs for the masses

Authors: Claude Duhr, Falko Dulat

In[1]:= T = 1/(1-x)*G[1/(1-x),0,0,1,1,1/(1+z)];

In[2]:= Ginsh[T, {x->0.3, z->0.45}]

Out[2]:= -0.0294179470484662503367597193416603238279

```
In[4]:= Ginsh[ T, \{x \rightarrow 0.3, z \rightarrow 0.45\}, PrecisionGoal \rightarrow 10]
```

```
In[5]:= Ginsh[ T, {x->0.3, z->0.45}, PrecisionGoal -> 100]
```

Out[4]:= -0.02941794704846625

Out[5]:= -0.02941794704846625033675971934166032382890897 19017828790593790231936752156076091770442309841

11624061172247650989277119

Three-Loop Heavy-to-light Form factors



Phys.Lett. B786, 453

Color-planar diagrams (Leading Color Contribution)



Integrals can be parameterized by



$$\mathcal{D}^d k_i \equiv \frac{m^{2\epsilon}}{\pi^{d/2} \Gamma(1+\epsilon)} d^d k_i, \quad d = 4 - 2\epsilon \; .$$

d=4-2€

All integrals can be reduced to 71 Master Integrals

$$\frac{\partial}{\partial s} = \frac{1}{s - m^2} p_2 \cdot \frac{\partial}{\partial p_2}$$

Canonical Basis:

$$\begin{split} F_{67} &= \epsilon^5 \left(s - m^2 \right) I_{1,1,0,1,1,-1,1,1,0,2,0,0} \,, \\ F_{68} &= \epsilon^5 \left(s - m^2 \right)^2 I_{1,1,0,1,1,0,1,1,0,2,0,0} \,, \\ F_{69} &= \epsilon^6 \left(s - m^2 \right) I_{1,1,0,1,0,0,1,1,1,0,0,1} \,, \\ F_{70} &= \epsilon^6 \left(s - m^2 \right)^2 I_{1,1,0,1,1,0,1,1,1,1,-1,1} \,, \\ F_{71} &= \epsilon^6 \left(s - m^2 \right) I_{1,1,0,1,1,-1,1,1,1,1,-1,1} \,, \\ &+ \frac{1}{12(1 - 2\epsilon)} (12F_2 + 6F_3 + 3F_4 - 2F_7 + 6F_9 \,, \\ &- 18F_{14} + 2F_{24} + 12F_{25}) \,. \end{split}$$

$$\frac{\partial \mathbf{F}(x,\epsilon)}{\partial x} = \epsilon \left(\frac{\mathbf{P}}{x} + \frac{\mathbf{Q}}{x-1}\right) \mathbf{F}(x,\epsilon).$$

P and Q are 71*71 rational matrices

Boundary Conditions

Known $F_{1} = \frac{1}{8},$ $F_{2} = \frac{1}{8} + \epsilon^{2} \frac{\pi^{2}}{12} + \epsilon^{3} \zeta(3) + \epsilon^{4} \frac{4\pi^{4}}{45} + 2\epsilon^{5} \frac{27\zeta(5) + \pi^{2} \zeta(3)}{3} + \epsilon^{6} \left(\frac{229\pi^{6}}{1890} + 4\zeta^{2}(3)\right) + \mathcal{O}(\epsilon^{7}),$

$$x \equiv \frac{s}{m^2}$$
.

$$\frac{\partial F_{38}}{\partial x} = \epsilon \left(\frac{-4F_2 - F_3 + 6F_{11} + 2F_{19} - 6(3F_{38} - 2F_{39})}{6x} + \frac{2F_{38}}{x - 1} \right),$$
Regular at x=0

 $6(3F_{38} - 2F_{39})|_{x=0} = (-4F_2 - F_3 + 6F_{11} + 2F_{19})|_{x=0}$

$$\begin{aligned} F_{38} &= \epsilon^5 \left(s - m^2 \right) I_{0,1,1,1,0,0,1,2,1,0,0,0}, \\ F_{39} &= \epsilon^4 m^2 \left(s - m^2 \right) I_{0,1,2,1,0,0,1,2,1,0,0,0}, \\ \frac{\partial F_{38}}{\partial x} &= \epsilon \left(\frac{-4F_2 - F_3 + 6F_{11} + 2F_{19} - 6(3F_{38} - 2F_{39})}{6x} + \frac{2F_{38}}{x - 1} \right), \\ \frac{\partial F_{39}}{\partial x} &= \epsilon \left(\frac{-20F_2 + 4F_3 + 3F_{11} - 12F_{16} + 30F_{18} + 16F_{19} - 30(3F_{38} - 2F_{39})}{12x} - 2\frac{3F_{39} - 4F_{38}}{x - 1} \right). \end{aligned}$$

$$6(3F_{38} - 2F_{39})|_{x=0} = (-4F_2 - F_3 + 6F_{11} + 2F_{19})|_{x=0},$$

$$30(3F_{38} - 2F_{39})|_{x=0} = (-20F_2 + 4F_3 + 3F_{11} - 12F_{16} + 30F_{18} + 16F_{19})|_{x=0}.$$

$$\begin{split} F_{37}|_{x=0} &= -\frac{11}{240} - \epsilon^2 \frac{\pi^2}{180} + \epsilon^3 \frac{3\zeta(3)}{4} + \epsilon^4 \frac{379\pi^4}{5400} \\ &- \epsilon^5 \left(\frac{4\pi^2 \zeta(3)}{3} - \frac{1107\zeta(5)}{20} \right) \\ &+ \epsilon^6 \left(\frac{901\pi^6}{4725} + \frac{21\zeta^2(3)}{2} \right) + \mathcal{O}(\epsilon^7) \,, \end{split}$$

$$F_{39}|_{x=0} &= -\epsilon^2 \frac{\pi^2}{36} + \epsilon^4 \frac{79\pi^4}{1080} + \epsilon^5 \left(\frac{143\pi^2 \zeta(3)}{18} + \frac{5\zeta(5)}{2} \right) \\ &+ \epsilon^6 \left(\frac{18737\pi^6}{22680} + 48\zeta^2(3) \right) + \mathcal{O}(\epsilon^7) \,, \end{aligned}$$

$$F_{41}|_{x=0} &= \frac{7}{180} - \epsilon^2 \frac{7\pi^2}{270} - \epsilon^3 \frac{89\zeta(3)}{45} - \epsilon^4 \frac{139\pi^4}{900} \\ &- \epsilon^5 \frac{353\pi^2 \zeta(3) + 8469\zeta(5)}{135} \\ &- \epsilon^6 \left(\frac{92077\pi^6}{170100} + \frac{2503\zeta^2(3)}{45} \right) + \mathcal{O}(\epsilon^7) \,, \end{split}$$

Results of Boundary conditions are expressed in terms of zeta functions.

$$\begin{aligned} F_{71} &= \epsilon^4 \left(H_{0,1,0,1}(x) - H_{0,0,1,1}(x) + \frac{\pi^2}{6} H_{0,1}(x) - \frac{\pi^4}{30} \right) \\ &+ \epsilon^5 \left(-2H_{0,0,0,0,1}(x) - 2H_{0,0,0,1,1}(x) - 2H_{0,0,1,0,1}(x) - 10H_{0,0,1,1,1}(x) \right) \\ &+ 2H_{0,1,0,1,1}(x) + 6H_{0,1,1,0,1}(x) - \frac{\pi^2}{6} H_{0,0,1}(x) + \pi^2 H_{0,1,1}(x) + 2\zeta(3)H_{0,1}(x) \\ &- \frac{7\pi^2 \zeta(3)}{6} - \zeta(5) \right) \\ &+ \epsilon^6 \left(- \left(2\zeta(5) + \frac{\pi^2 \zeta(3)}{3} \right) H_1(x) + \frac{9\pi^4}{40} H_{0,1}(x) \right) \\ &+ \chi(3)(-13H_{0,0,1}(x) + 9H_{0,1,1}(x) - 2H_{1,0,1}(x)) - \pi^2 \left(-H_{0,0,0,1}(x) - \frac{5}{6} H_{0,0,1,1}(x) \right) \\ &+ H_{0,1,0,1}(x) + 6H_{0,1,1,1}(x) + \frac{1}{3} H_{1,0,0,1}(x) \right) - 11H_{0,0,0,0,1}(x) - 11H_{0,0,0,0,1,1}(x) \\ &- 20H_{0,0,0,1,0,1}(x) - 20H_{0,0,0,1,1,1}(x) - 16H_{0,0,1,0,0,1}(x) - 26H_{0,0,1,0,1,1}(x) \\ &- 29H_{0,0,1,1,0,1}(x) - 76H_{0,0,1,1,1}(x) - 14H_{0,1,0,0,0,1}(x) - 12H_{0,1,0,0,1,1}(x) \\ &+ 36H_{0,1,1,1,0,1}(x) + 4H_{1,0,0,0,0,1}(x) + 2H_{1,0,0,1,0,1}(x) + 2H_{1,0,1,0,0,1}(x) \\ &- \frac{1219\pi^6}{15120} \right) + \mathcal{O}(\epsilon^7), \end{aligned}$$

H are Harmonic Polylogarithms



Check: (s=-1.3,m=1.0)





https://bitbucket.org/feynmanIntegrals/fiesta/src/master/

fiesta	Tools for Feynman integrals / Feynman integrals fiesta Clone	
Source	EIESTA integrator - Feynman Integral Evaluation by a Sector decomposition Approach	
b Commits	TESTA Integration - regiminar integral evaluation by a sector accomposition Approach	
Properties	By master Files Filter files Q	linata
Branches		insta
Pull requests	■ /	make
Pipelines		make
P Deployments	FIESTA4 2020-04-14 version 4.2	
Issues	README.md 1.62 KB 2020-04-14 information on bitbucket updated	
Jira issues		
🖻 Wiki	README.md	
Downloads		
_	READIVIE	
	FIESTA stands for Feynman Integral Evaluation by a Sector decomposition Approach.	
	FIESTA 4.2 released. Nothing special out there, only bug fixes collected through years and upgrade to new compilers. The only real new feature is the LP=True option to use Lee-Pomeransky representation for region search (https://arxiv.org/abs/1809.04325)	

FIESTA

$$\mathcal{F}(a_1, \dots, a_n) = \int \dots \int \frac{\mathrm{d}^d k_1 \dots \mathrm{d}^d k_l}{E_1^{a_1} \dots E_n^{a_n}}$$

$$d = 4 - 2\epsilon$$

$$E_i \equiv E_i - i0$$

Prefactor

$$\mathbf{i} \pi^{d/2} \mathbf{e}^{-\gamma_{\mathrm{E}} \varepsilon}$$

FIESTA

pySecDec

A GPU compatible quasi-Monte Carlo integrator interfaced to pySecDec

S. Borowka (CERN), G. Heinrich (Munich, Max Planck Inst.), S. Jahn (Munich, Max Planck Inst.), S.P. Jones (CERN and Munich, Max Planck Inst.), M. Kerner (Munich, Max Planck Inst. and Zurich U.) et al. (Nov 28, 2018) Published in: *Comput.Phys.Commun.* 240 (2019) 120-137 • e-Print: 1811.11720 [physics.comp-ph]

Before installing pySECDEC, make sure that recent versions of numpy (http: //www.numpy.org/) and sympy (http://www.sympy.org/) are installed. The pySECDEC program (which includes the QMC integrator library) can be downloaded from https://github.com/mppmu/secdec/releases. To install pySECDEC, perform the following steps

```
tar -xf pySecDec-<version>.tar.gz
cd pySecDec-<version>
make
<copy the highlighted output lines into your .bashrc>
```

DiffExp: 2006.05510

DiffExp, a Mathematica package for computing Feynman integrals in terms of one-dimensional series expansions

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Based on 1907.13234