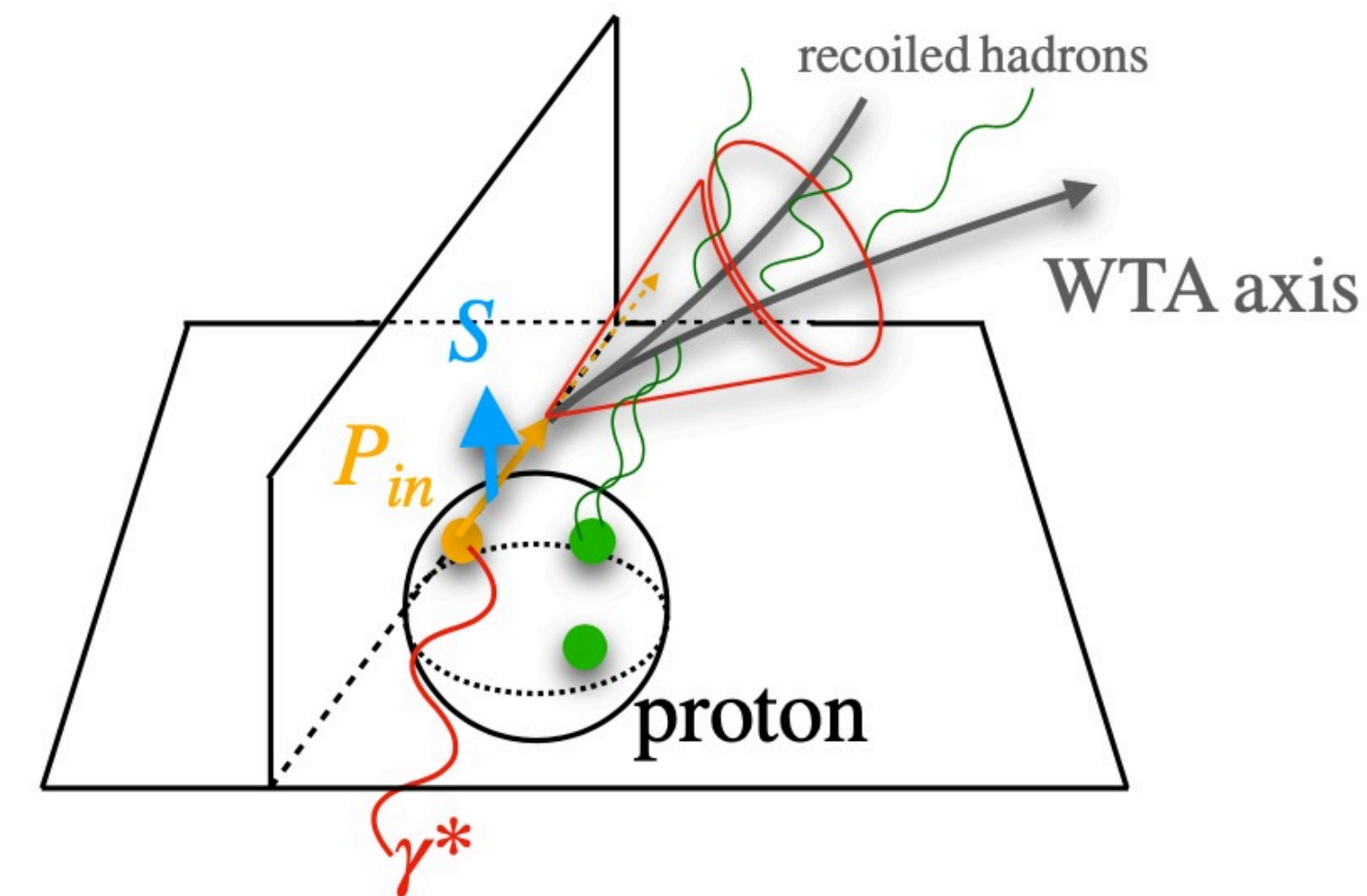




# The time reversal odd side of a jet

**Hongxi Xing**

(South China Normal University)



微扰量子场论研讨会@上海, 5.15-16

# Outline

- Motivation - the role of jets
- The T-odd jet in DIS and  $e^+e^-$  collisions
- Summary

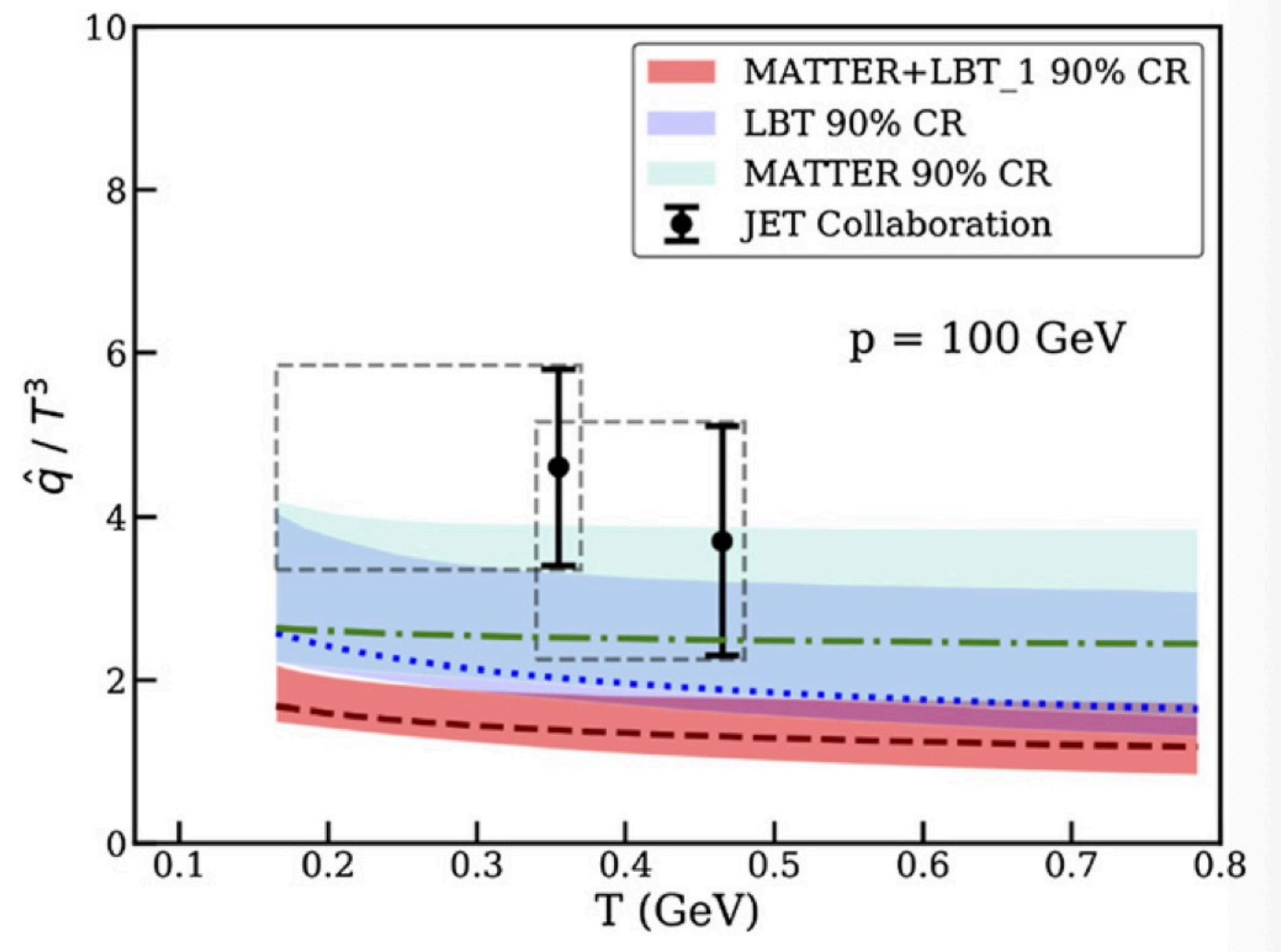
Liu, [HX](#), 2104.03328, 2021

Liu, [HX](#), Zhang, 2021

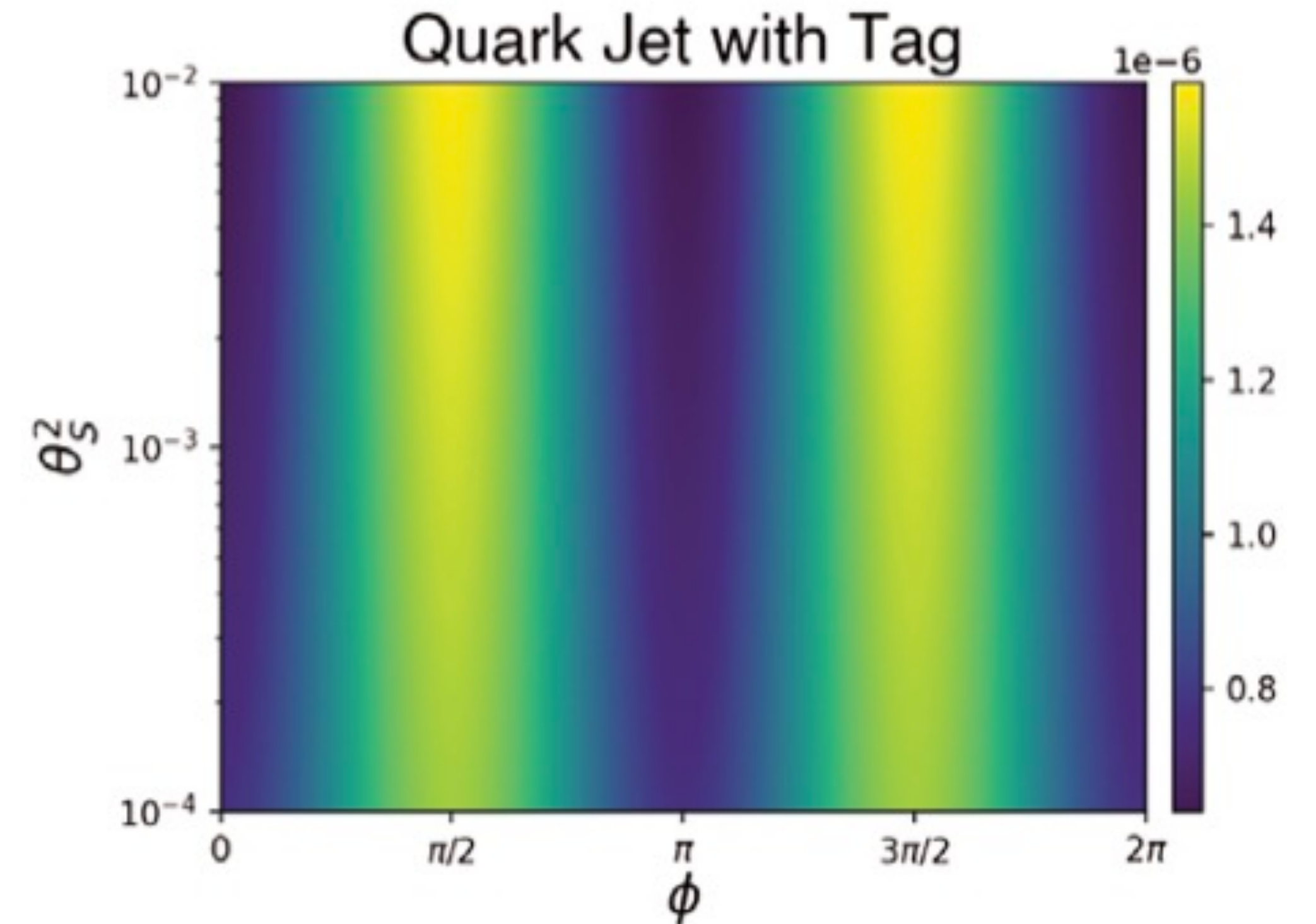
Lai, Liu, Wang, [HX](#), 2021

# Motivation - jet is powerful in many fields

- Jet as a tool for QGP tomography
- Quantum interference in jet substructure



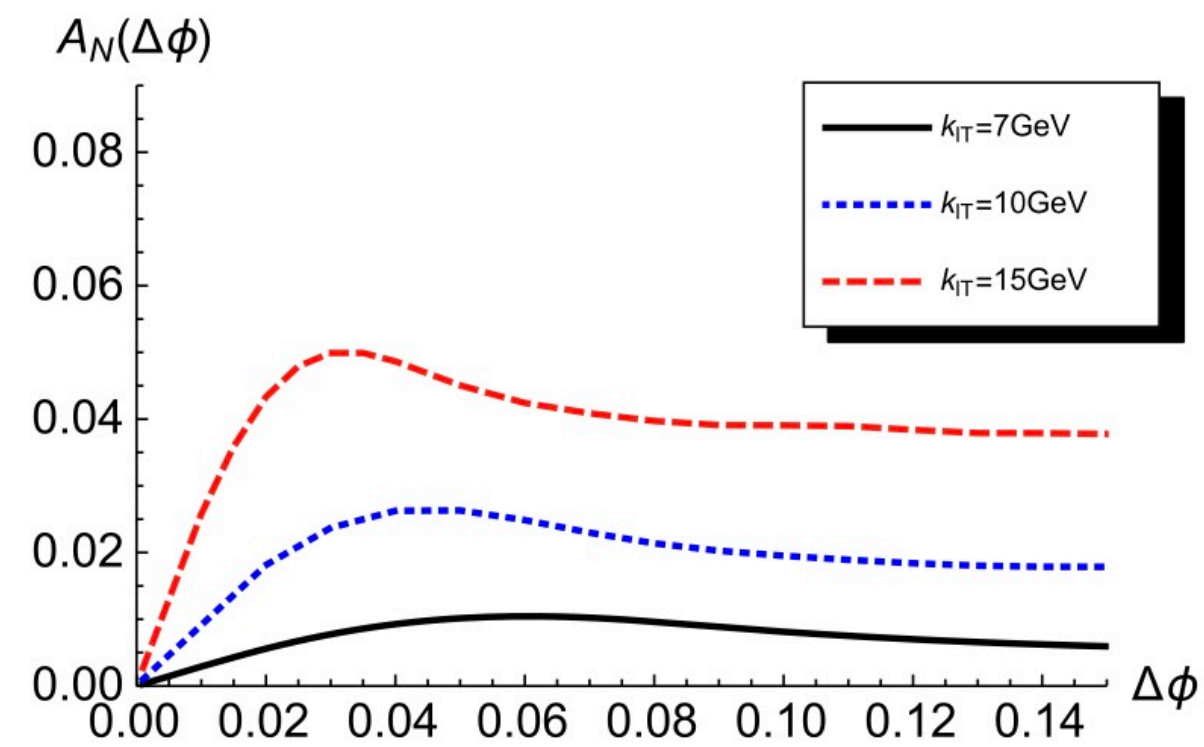
See e.g., review by Cao, Wang, Rept.Prog.Phys 2021



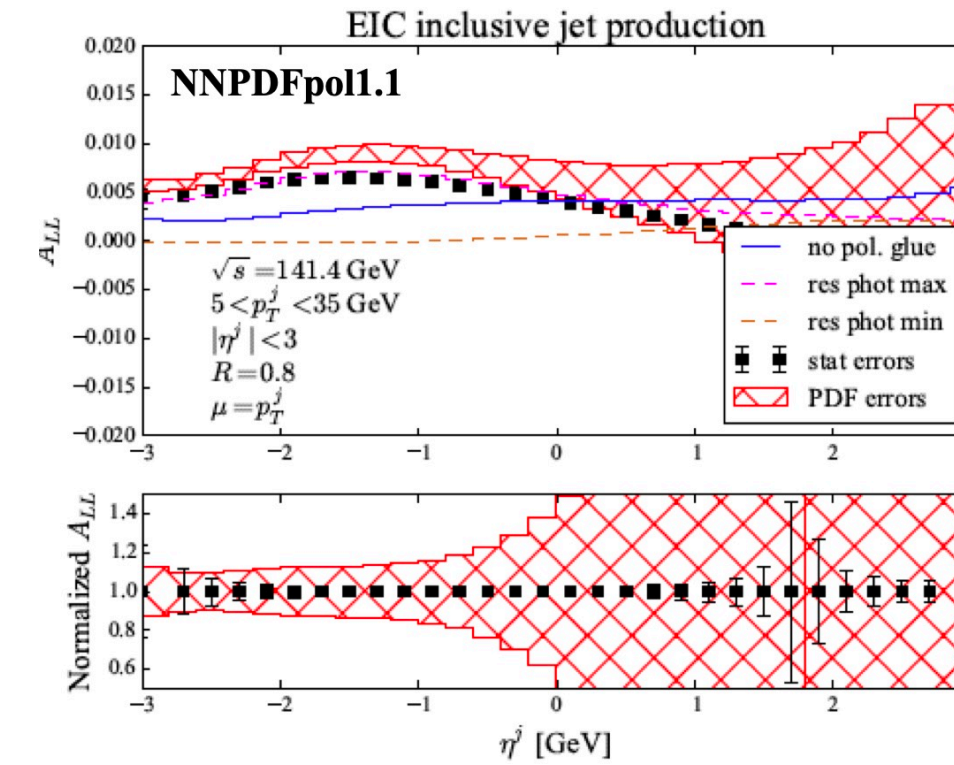
Chen, Moult, Zhu, PRL 2021  
talk by H. Chen

# Jet for spin related physics

- Jet in DIS for Sivers and helicity distributions

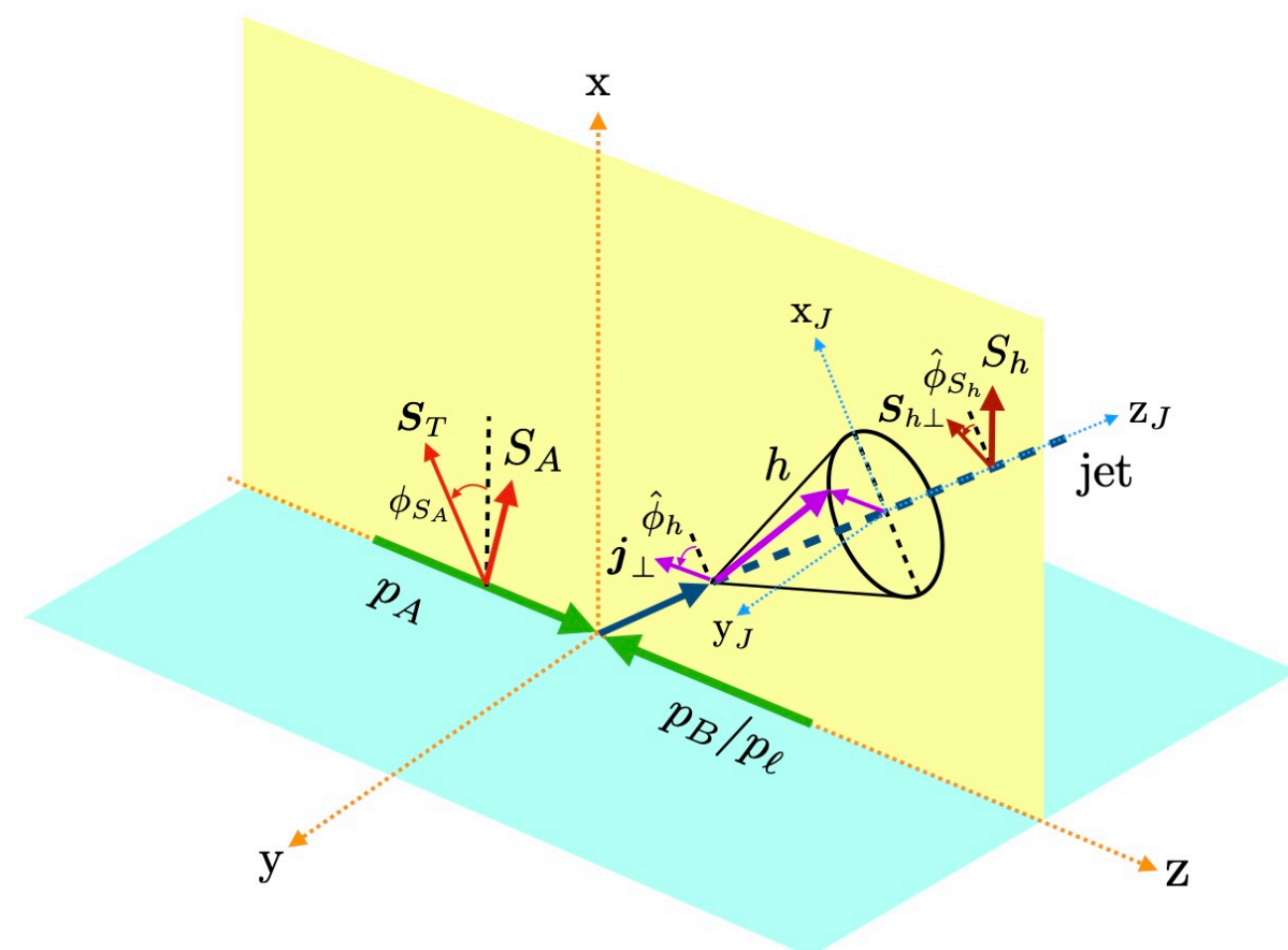


Liu, Ringer, Vogelsang, Yuan, PRL 2019

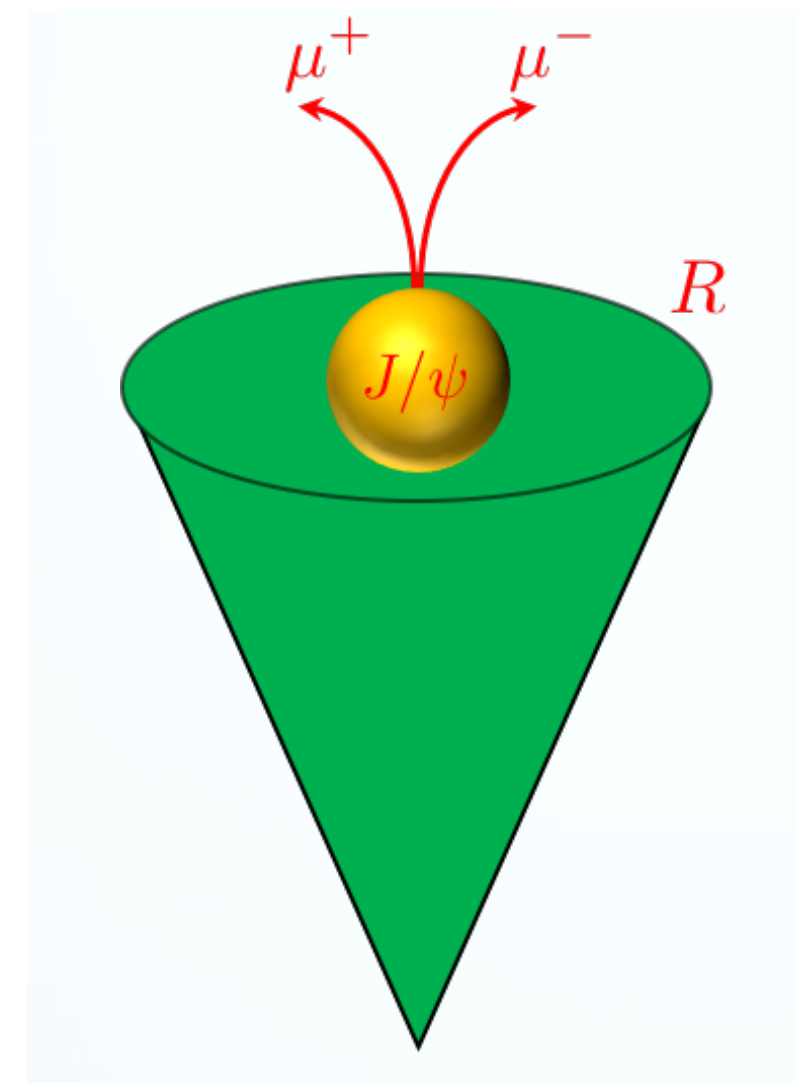


Boughezal, Petriello, **HX**, PRD 2017

- Jet fragmentation -> polarized fragmentation functions ( $J/\psi$ ,  $\Lambda$ )

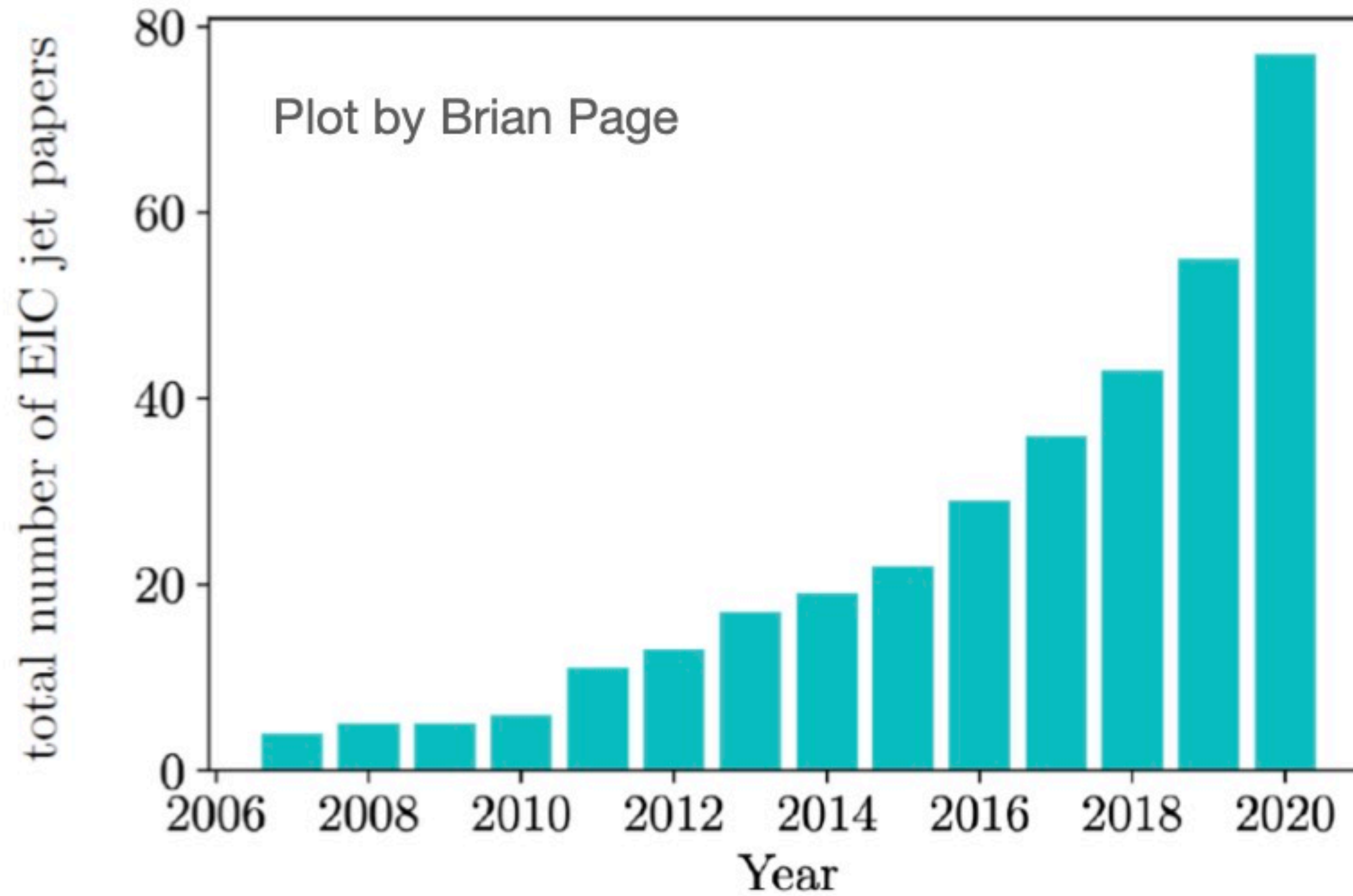


Kang, Lee, Zhao, PLB 2020



Kang, Qiu, Ringer, **HX**, Zhang  
PRL 2017

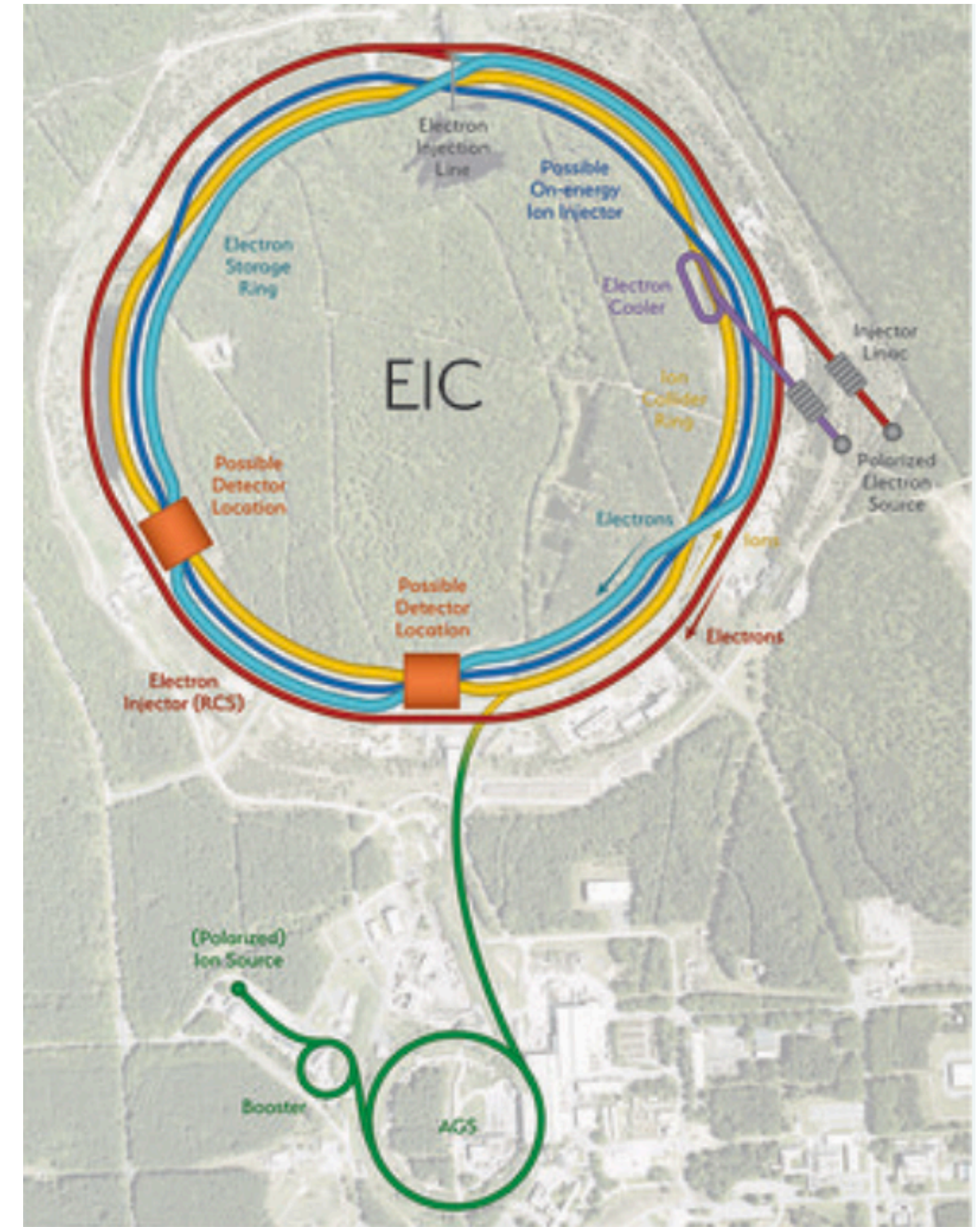
# Rapid growing interests for jets @ EIC



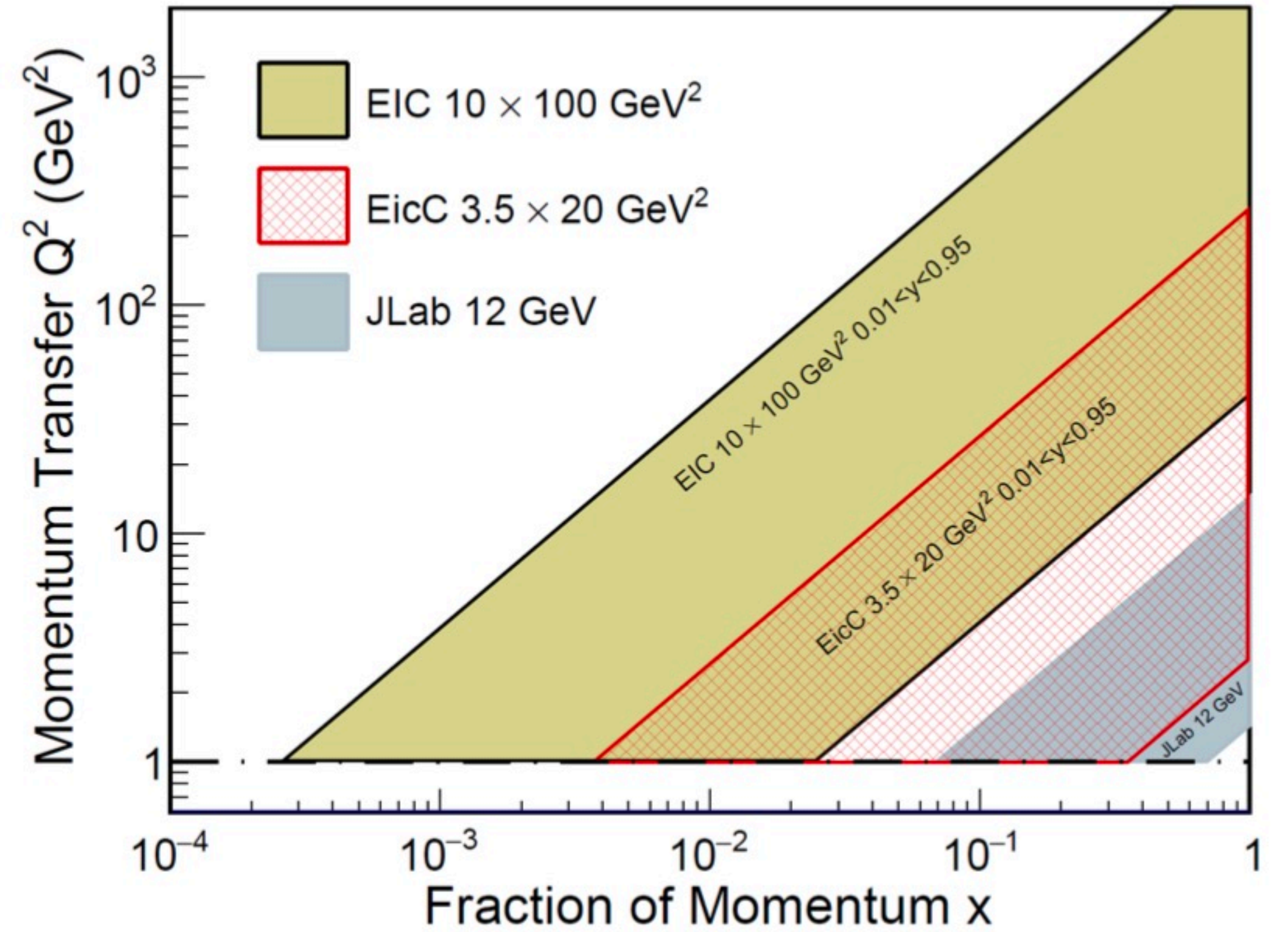
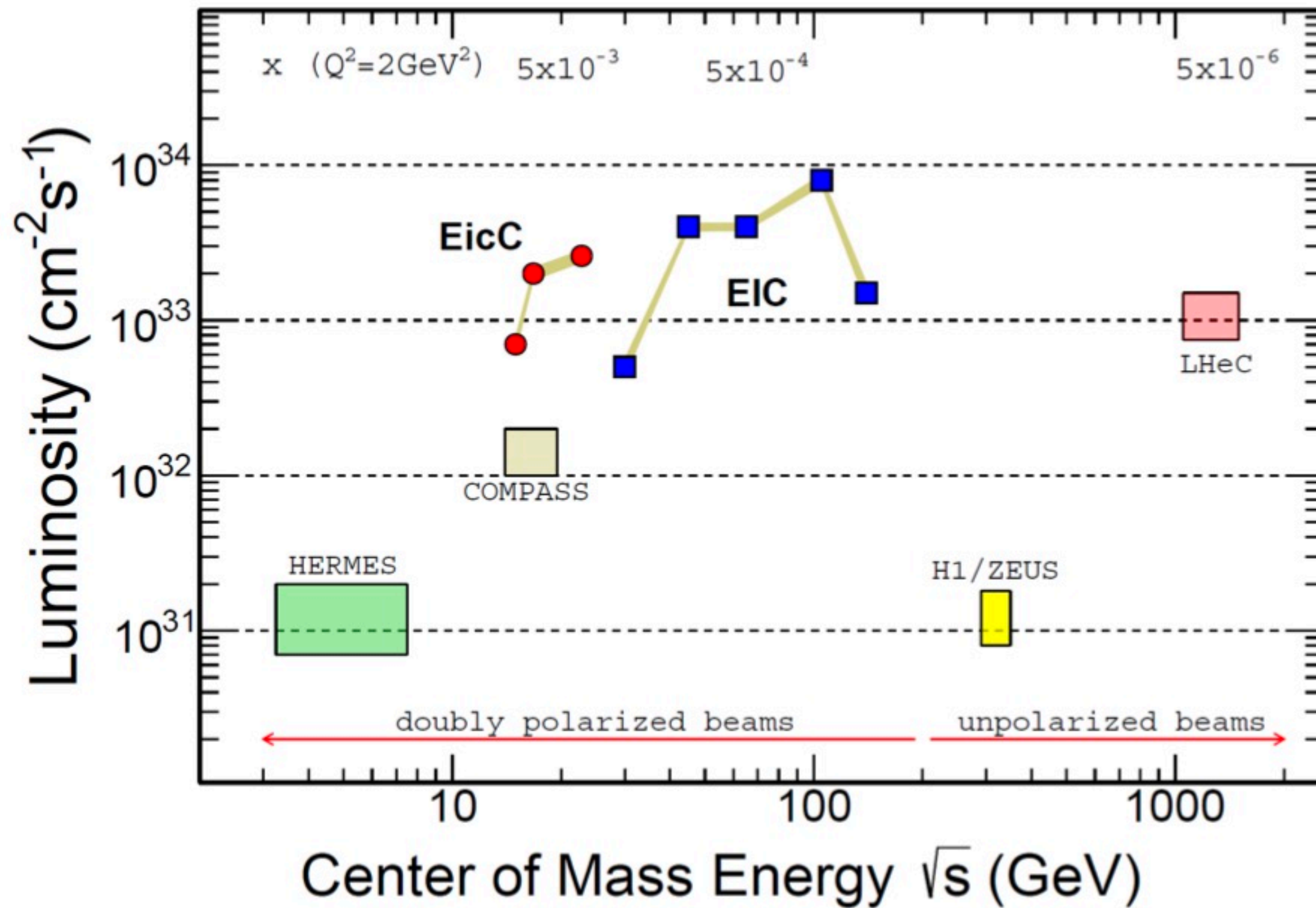


**E**lectron **I**on **C**ollider in **C**hina, EicC

# Electron Ion Collider (US)



# Mapping out the nucleon structure via EICs worldwide



# EicC white paper (arXiv: 2102.09222)

Also in production in the *Frontiers of Physics* Journal

arXiv.org > nucl-ex > arXiv:2102.09222

Search...

Help | Advance

## Nuclear Experiment

[Submitted on 18 Feb 2021]

### Electron-Ion Collider in China

Daniele P. Anderle, Valerio Bertone, Xu Cao, Lei Chang, Ningbo Chang, Gu Chen, Xurong Chen, Zhuojun Chen, Zhufang Cui, Lingyun Dai, Weitian Deng, Minghui Ding, Xu Feng, Chang Gong, Longcheng Gui, Feng-Kun Guo, Chengdong Han, Jun He, Tie-Jiun Hou, Hongxia Huang, Yin Huang, Krešimir Kumerički, L. P. Kaptari, Demin Li, Hengne Li, Minxiang Li, Xueqian Li, Yutie Liang, Zuotang Liang, Chen Liu, Chuan Liu, Guoming Liu, Jie Liu, Liuming Liu, Xiang Liu, Tianbo Liu, Xiaofeng Luo, Zhun Lyu, Boqiang Ma, Fu Ma, Jianping Ma, Yugang Ma, Lijun Mao, Cédric Mezrag, Hervé Moutarde, Jialun Ping, Sixue Qin, Hang Ren, Craig D. Roberts, Juan Rojo, Guodong Shen, Chao Shi, Qintao Song, Hao Sun, Paweł Sznajder, Enke Wang, Fan Wang, Qian Wang, Rong Wang, Ruiru Wang, Taofeng Wang, Wei Wang, Xiaoyu Wang, Xiaoyun Wang, Jiajun Wu, Xinggong Wu, Lei Xia, Bowen Xiao, Guoqing Xiao, Ju-Jun Xie, Yaping Xie, Hongxi Xing, Hushan Xu, Nu Xu, Shusheng Xu, Mengshi Yan, Wenbiao Yan, Wencheng Yan, Xihu Yan, Jiancheng Yang, Yi-Bo Yang, Zhi Yang, Deliang Yao, Peilin Yin, C.-P. Yuan, Wenlong Zhan, Jianhui Zhang, Jinlong Zhang, Pengming Zhang, Chao-Hsi Chang, Zhenyu Zhang, Hongwei Zhao, Kuang-Ta Chao, Qiang Zhao, Yuxiang Zhao, Zhengguo Zhao, Liang Zheng, Jian Zhou, Xiang Zhou, Xiaorong Zhou et al. (2 additional authors not shown)

Lepton scattering is an established ideal tool for studying inner structure of small particles such as nucleons as well as nuclei. As a future high energy nuclear physics project, an Electron-ion collider in China (EicC) has been proposed. It will be constructed based on an upgraded heavy-ion accelerator, High Intensity heavy-ion Accelerator Facility (HIAF) which is currently under construction, together with a new electron ring. The proposed collider will provide highly polarized electrons (with a polarization of  $\sim 80\%$ ) and protons (with a polarization of  $\sim 70\%$ ) with variable center of mass energies from 15 to 20 GeV and the luminosity of  $(2-3) \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$ . Polarized deuterons and Helium-3, as well as unpolarized ion beams from Carbon to Uranium, will be also available at the EicC.

The main foci of the EicC will be precision measurements of the structure of the nucleon in the sea quark region, including 3D tomography of nucleon; the partonic structure of nuclei and the parton interaction with the nuclear environment; the exotic states, especially those with heavy flavor quark contents. In addition, issues fundamental to understanding the origin of mass could be addressed by measurements of heavy quarkonia near-threshold production at the EicC. In order to achieve the above-mentioned physics goals, a hermetical detector system will be constructed with cutting-edge technologies.

This document is the result of collective contributions and valuable inputs from experts across the globe. The EicC physics program complements the ongoing scientific programs at the Jefferson Laboratory and the future EIC project in the United States. The success of this project will also advance both nuclear and particle physics as well as accelerator and detector technology in China.

Comments: EicC white paper, written by the whole EicC working group

Subjects: **Nuclear Experiment (nucl-ex)**; High Energy Physics - Experiment (hep-ex); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th)

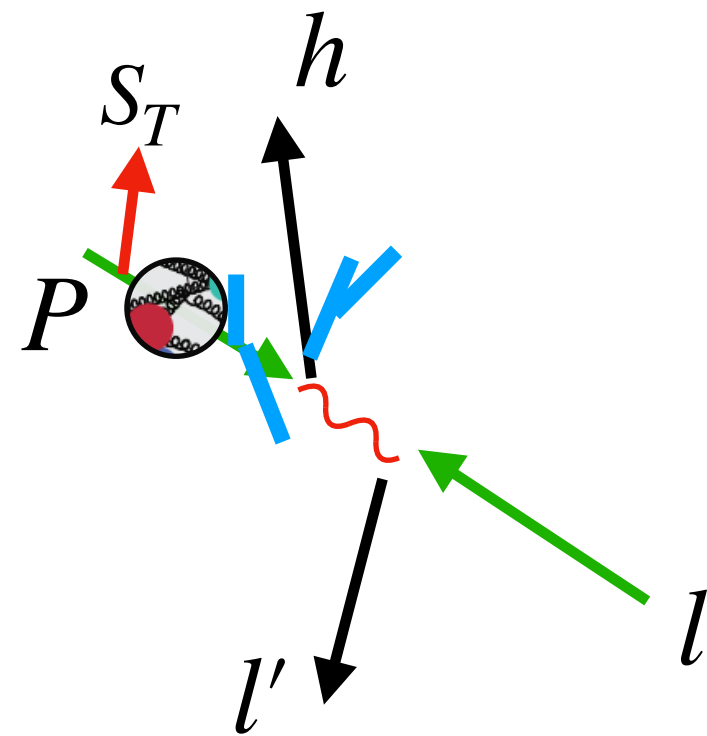
Cite as: [arXiv:2102.09222](https://arxiv.org/abs/2102.09222) [nucl-ex]

(or [arXiv:2102.09222v1](https://arxiv.org/abs/2102.09222v1) [nucl-ex] for this version)

Now we have 46 institutes and >100 physicists



# Jet vs hadron

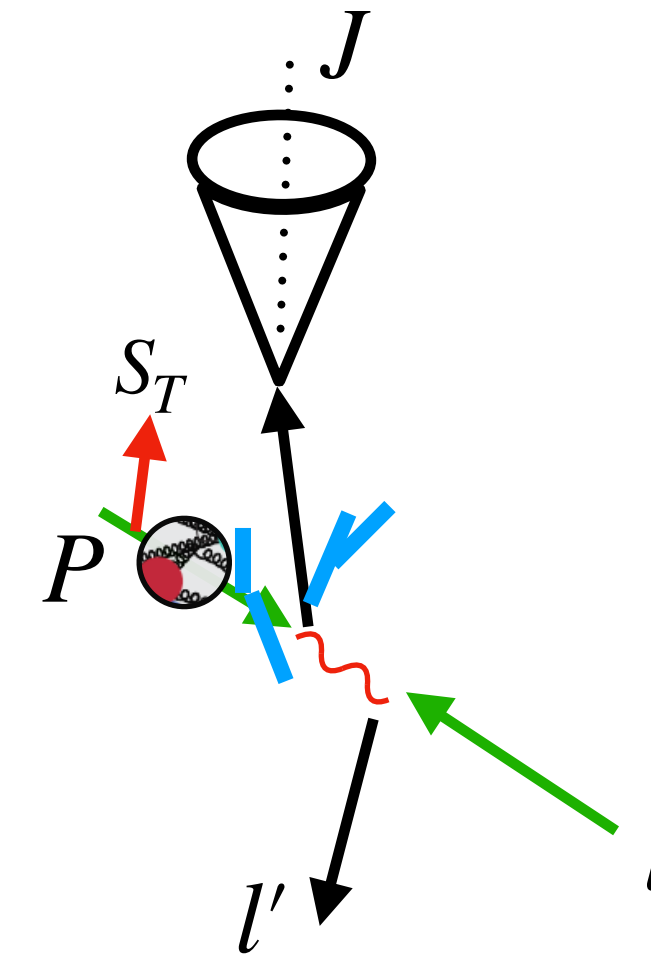


- Semi-inclusive DIS

$$\sigma^h \sim \sum_{i,j} \hat{\sigma}_{ei \rightarrow e'j} f_{ilp}(x, k_T) \otimes D_{j|h}(z, k'_T) \otimes S(k_T)$$

👎 The nonperturbative fragmentation functions -> extra uncertainty

👍 Identified hadrons can be used as flavor separation



- Jet production in DIS

$$\sigma^J = \sum_{i,j} \hat{\sigma}_{ei \rightarrow e'j} f_{ilp}(x, k_T) \otimes J(p_T, R) \otimes S(k_T, p_T R)$$

👍 Jet functions can be calculated perturbatively -> controllable uncertainty

👎 Sum all hadrons within jet, hard for flavor separation

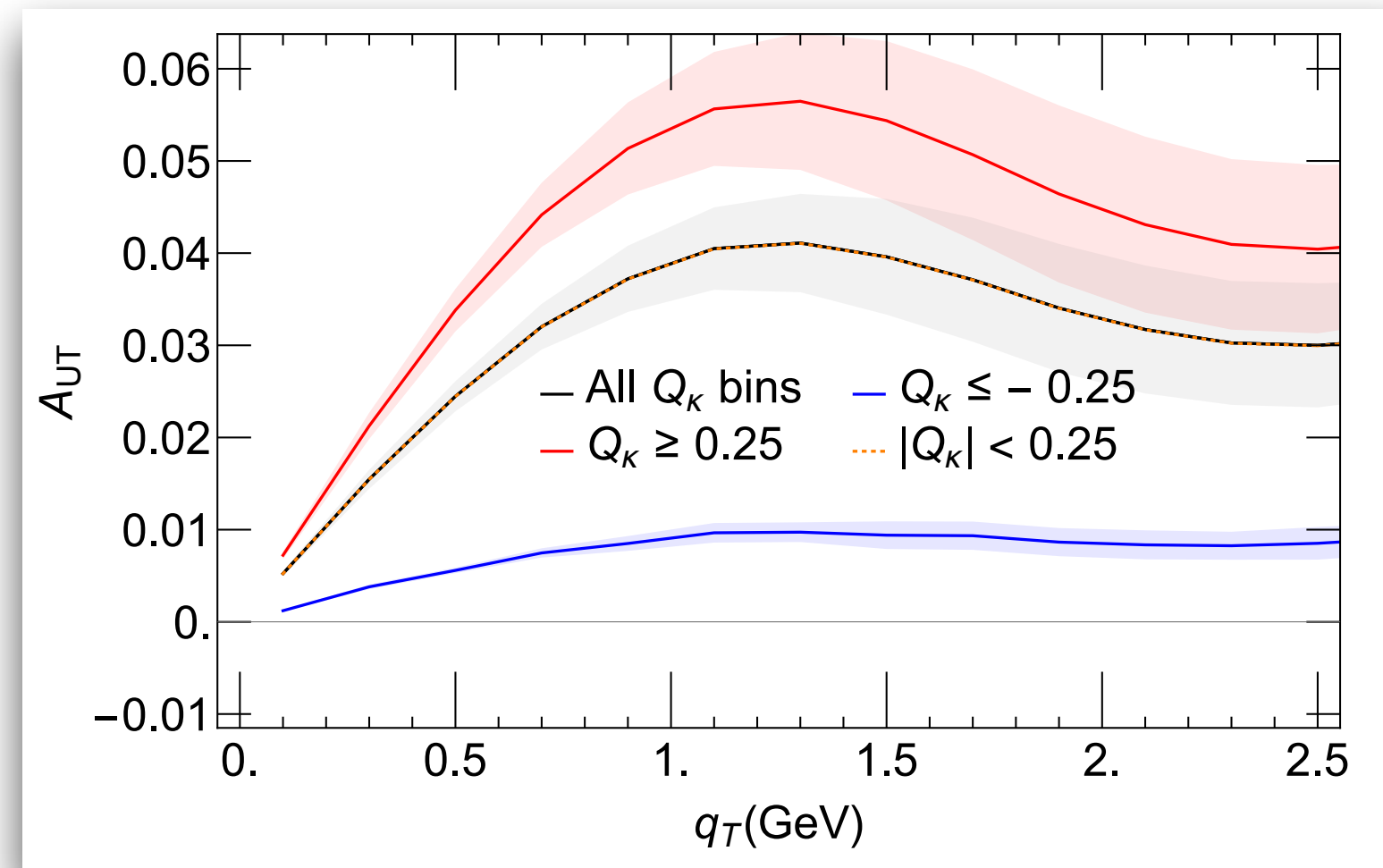
# Jet charge for nucleon/nucleus flavor separation

Definition:

$$Q_\kappa = \sum_h \left( \frac{p_{h,T}}{p_J} \right)^\kappa Q_h$$

Field, Feynman, 1978

Krohn, Schwartz, Lin, Waalewijn, 2013

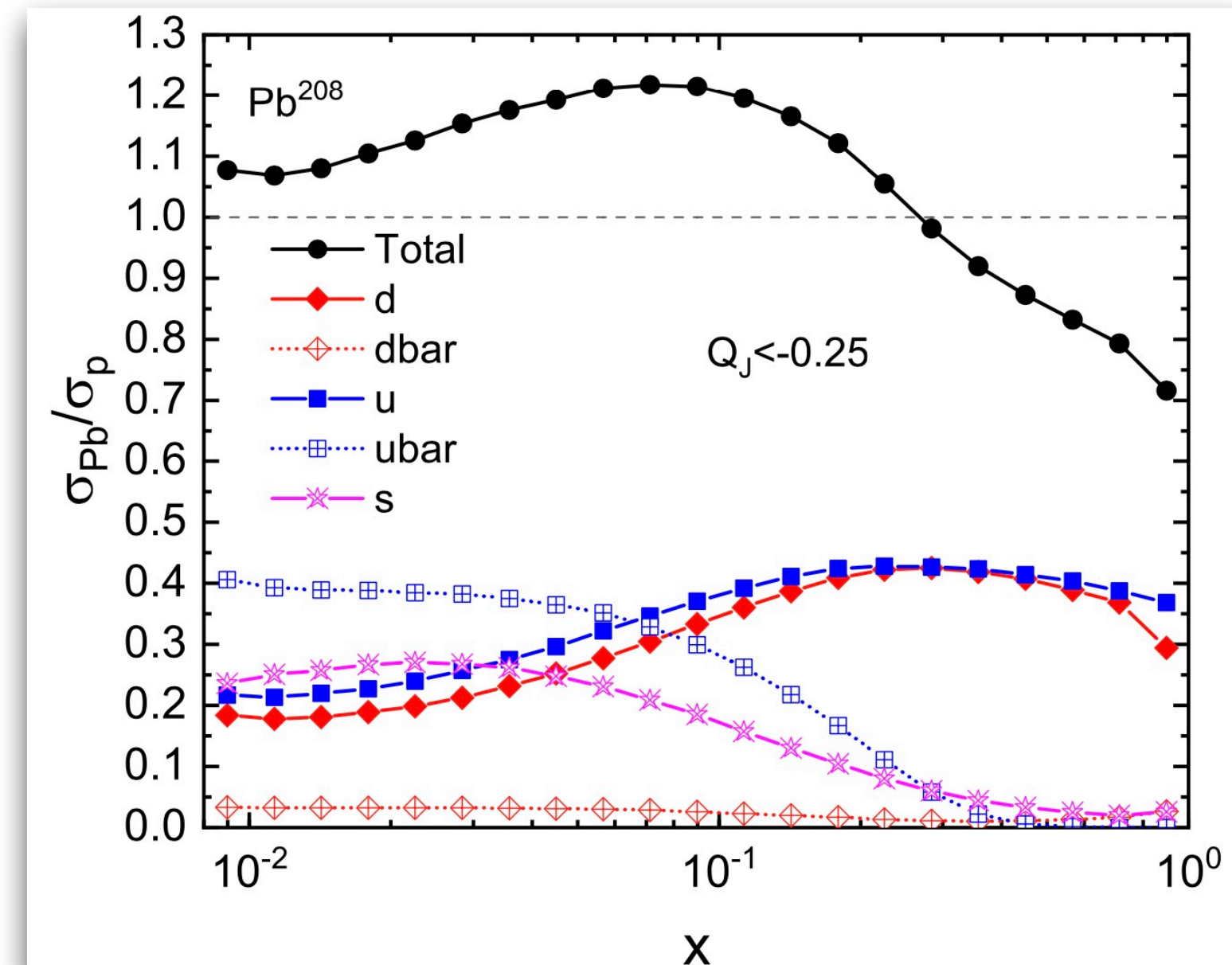
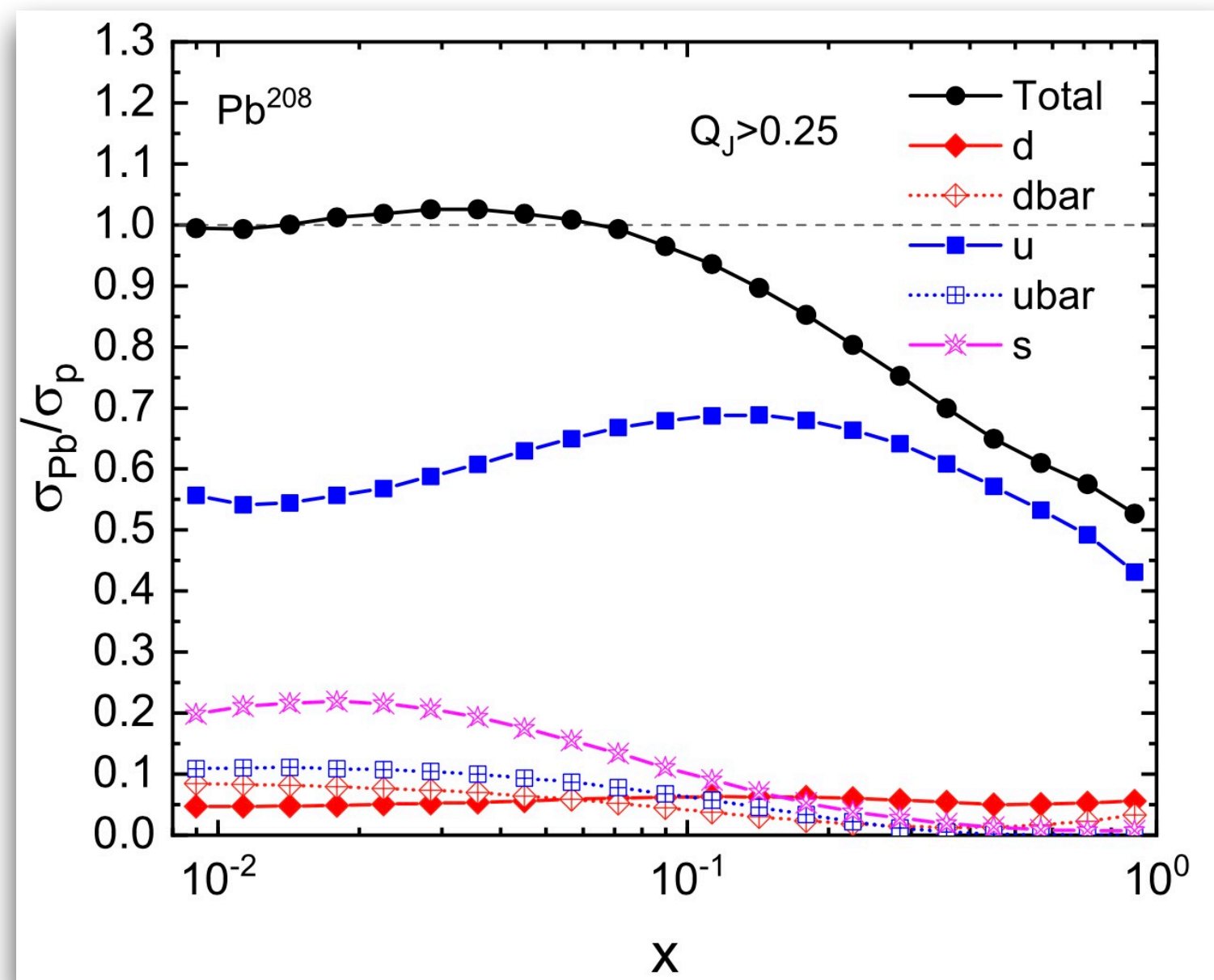


Flavor dependent Sivers effect

Positive charge: u

Negative charge: d

Kang, Liu, Mantry, Shao, PRL 2020



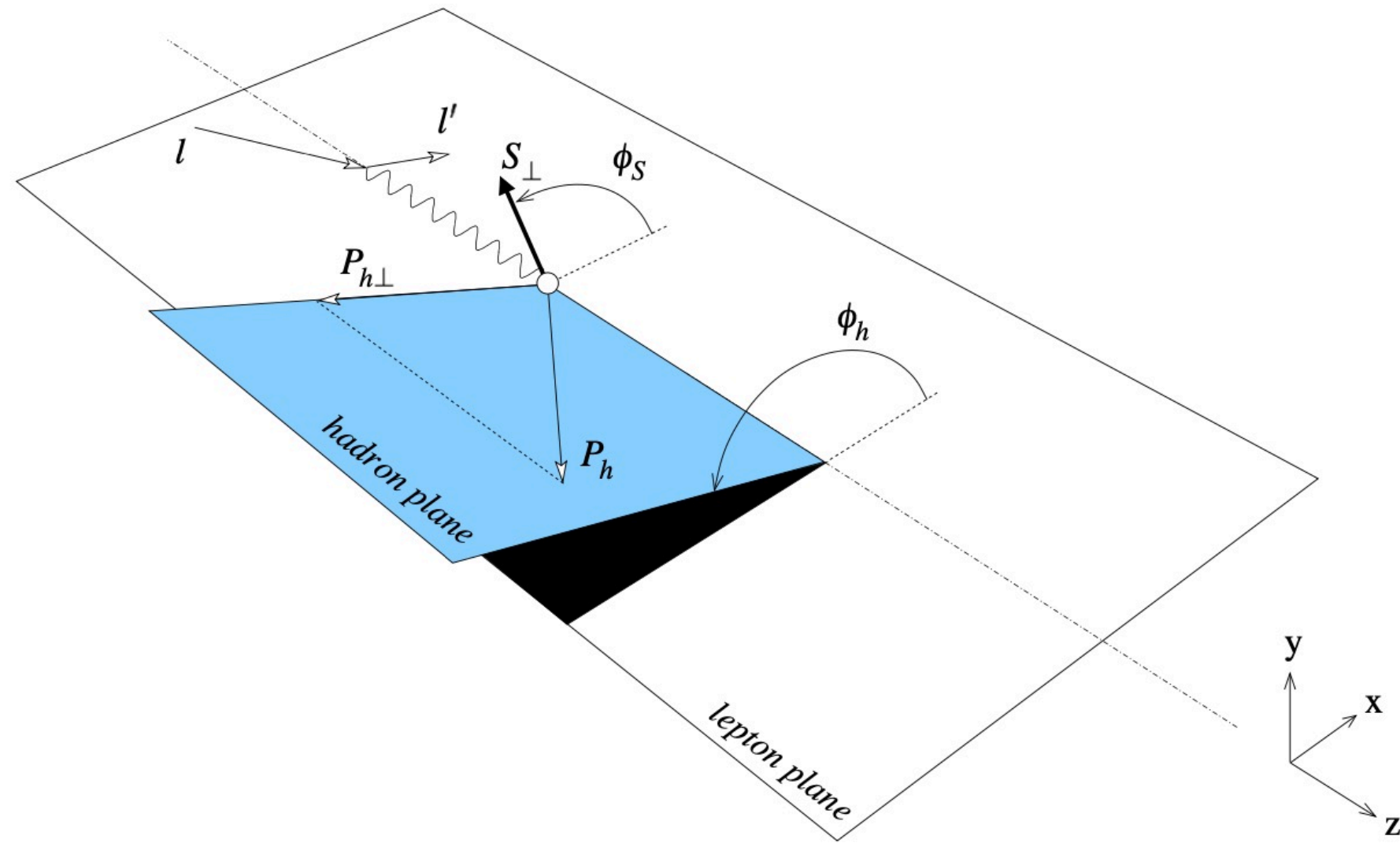
Flavor dependent nuclear PDFs

Positive charge: u

Negative charge: u, d, sea

Liu, HX, Zhang, in preparation, 2021

# The role of T-odd fragmentation functions



$$\begin{aligned}
 \sigma(\phi, \phi_S) &\equiv \frac{d^6 \sigma}{dx dy dz d\phi d\phi_S dP_{KT}^2} = \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \\
 &\left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} \cos\phi F_{UU}^{\cos\phi} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} + \lambda_e \left[ \sqrt{2\epsilon(1-\epsilon)} \sin\phi F_{LU}^{\sin\phi} \right] + \right. \\
 &+ S_L \left[ \sqrt{2\epsilon(1+\epsilon)} \sin\phi F_{UL}^{\sin\phi} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] + S_L \lambda_e \left[ \sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos\phi F_{LL}^{\cos\phi} \right] \\
 &+ |S_T| \left[ \sin(\phi - \phi_S) \left( F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \right. \\
 &\left. + \sqrt{2\epsilon(1+\epsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \right] \\
 &\left. - |S_T| \lambda_e \left[ \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} + \sqrt{2\epsilon(1-\epsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right] \right\} ;
 \end{aligned}$$

Labels for the fragmentation functions in the equation:
 

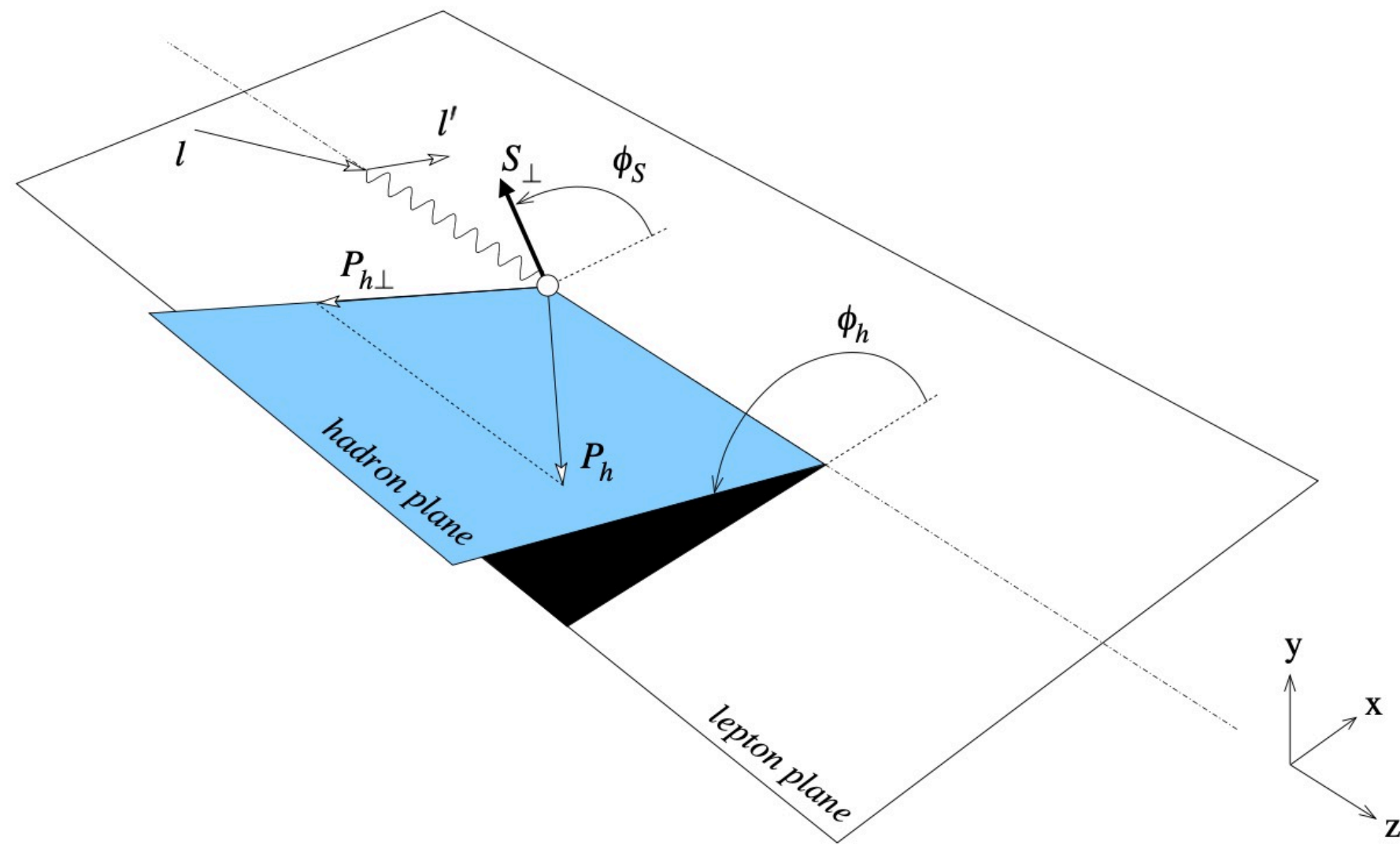
- Cahn-effect + BM ⊗ Collins** (points to the first term)
- Worm-gear (Kotzinian-Mulders) ⊗ Collins** (points to the second term)
- BM ⊗ Collins** (points to the third term)
- Sivers ⊗ D1** (points to the fourth term)
- Worm-gear ⊗ D1** (points to the fifth term)
- Transversity ⊗ Collins** (points to the sixth term)
- Pretzelosity ⊗ Collins** (points to the seventh term)

Bacchetta, Diehl, Goeke, Metz, Mulders, Schlegel, JHEP 2007

QCD requires all physical observables should preserve time reversal invariance!

The T-odd FFs provide us abundant opportunities to probe nucleon structure.

# Limited power of jet probing - conventional wisdom




$$\begin{aligned}
 & \sigma(\phi, \phi_S) \equiv \frac{d^6\sigma}{dx dy dz d\phi d\phi_S dP_{KT}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \\
 & \left\{ F_{UU,T} + \epsilon F_{UU,LL} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi - \phi_S) F_{UU}^{\sin(\phi - \phi_S)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} + \sqrt{2\epsilon(1-\epsilon)} \sin\phi F_{LU}^{\sin\phi} \right\} + \\
 & + S_L \left[ \sqrt{2\epsilon(1+\epsilon)} \sin\phi F_{UL}^{\sin\phi} + \epsilon \sin(2\phi) F_{UL}^{\cos(2\phi)} \right] + S_L \lambda_e \left[ \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LL}^{\cos(\phi - \phi_S)} + \sqrt{2\epsilon(1-\epsilon)} \cos\phi F_{LL}^{\cos\phi} \right] \\
 & + |S_T| \left[ \sin(\phi - \phi_S) \left( F_{UT}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \right. \\
 & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\epsilon(1-\epsilon)} (1+\epsilon) \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \right] \\
 & + |S_T| \lambda_e \left[ \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} + \sqrt{2\epsilon(1-\epsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right] \Big\} ;
 \end{aligned}$$

Labels and annotations in the diagram:
 

- Worm-gear (Kotzinian-Mulders) ⊗ Collins**: Points to the  $F_{UU,T}$  and  $F_{UU,LL}$  terms.
- BM ⊗ Collins**: Points to the  $F_{LU}^{\sin\phi}$  term.
- Sivers ⊗ D1**: Points to the  $F_{UL}^{\sin\phi}$  and  $F_{UL}^{\cos(2\phi)}$  terms.
- Worm-gear**: Points to the  $F_{LT}^{\cos(\phi - \phi_S)}$  term.
- Transversity ⊗ Collins**: Points to the  $F_{LL}^{\cos\phi}$  term.
- Pretzelosity ⊗ Collins**: Points to the  $F_{LL}^{\cos(2\phi)}$  term.

 Red 'X' marks indicate terms that are T-odd and thus vanish in conventional wisdom. Blue checkmarks indicate terms that are T-even and survive.

In conventional wisdom, perturbative jet is T-even, thus many spin structures vanish.

 T-odd distributions involved  
 But Perturbative Jet is (almost) T-even

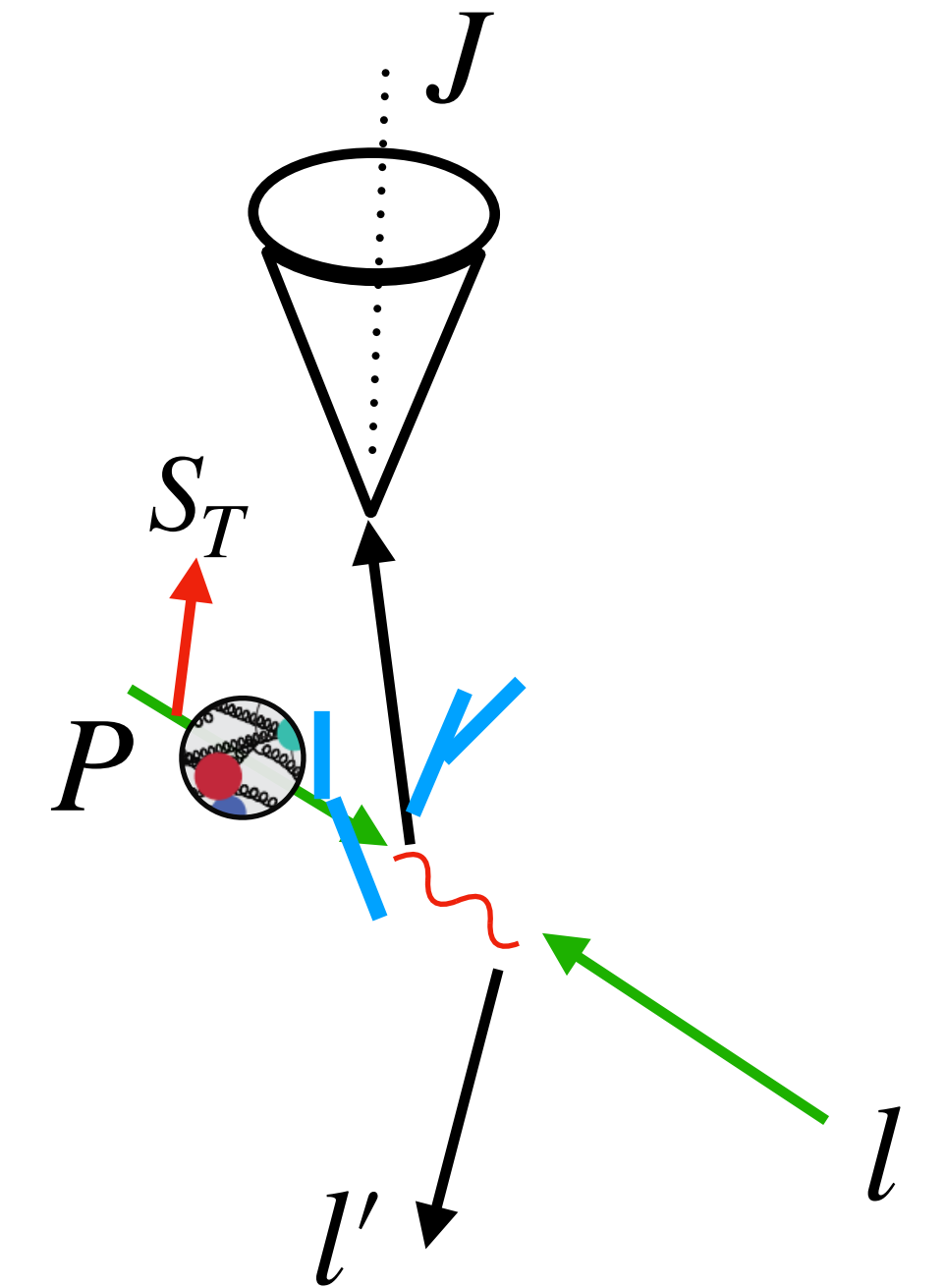
# Revisit jet production in DIS

$$l p(P, s_T) \rightarrow \gamma^*(q) \rightarrow l' J(P_J) + X$$

- Factorization frame

$$P^\mu = \frac{P^-}{2}(1, 0, 0, 1) = \frac{P^-}{2}n^\mu \quad P_J^\mu = \frac{P_J^+}{2}(1, 0, 0, -1) = \frac{P_J^+}{2}\bar{n}^\mu$$

$$q_T = q - \frac{q \cdot n}{2}\bar{n} - \frac{q \cdot \bar{n}}{2}n \quad q_T \ll Q$$



- Winner-take-all jet

$$\text{SJA : } E_{(12)} = E_1 + E_2, \quad \vec{p}_{(12)} = \vec{p}_1 + \vec{p}_2,$$

$$\text{WTA : } E_{(12)} = E_1 + E_2, \quad \vec{p}_{(12)} = E_{(12)} \left[ \frac{\vec{p}_1}{|\vec{p}_1|} \theta(E_1 - E_2) + \frac{\vec{p}_2}{|\vec{p}_2|} \theta(E_2 - E_1) \right]$$

WTA jet TMD has the same RG as TMD FFs

Bertolini, Chan, Thaler, JHEP 2014

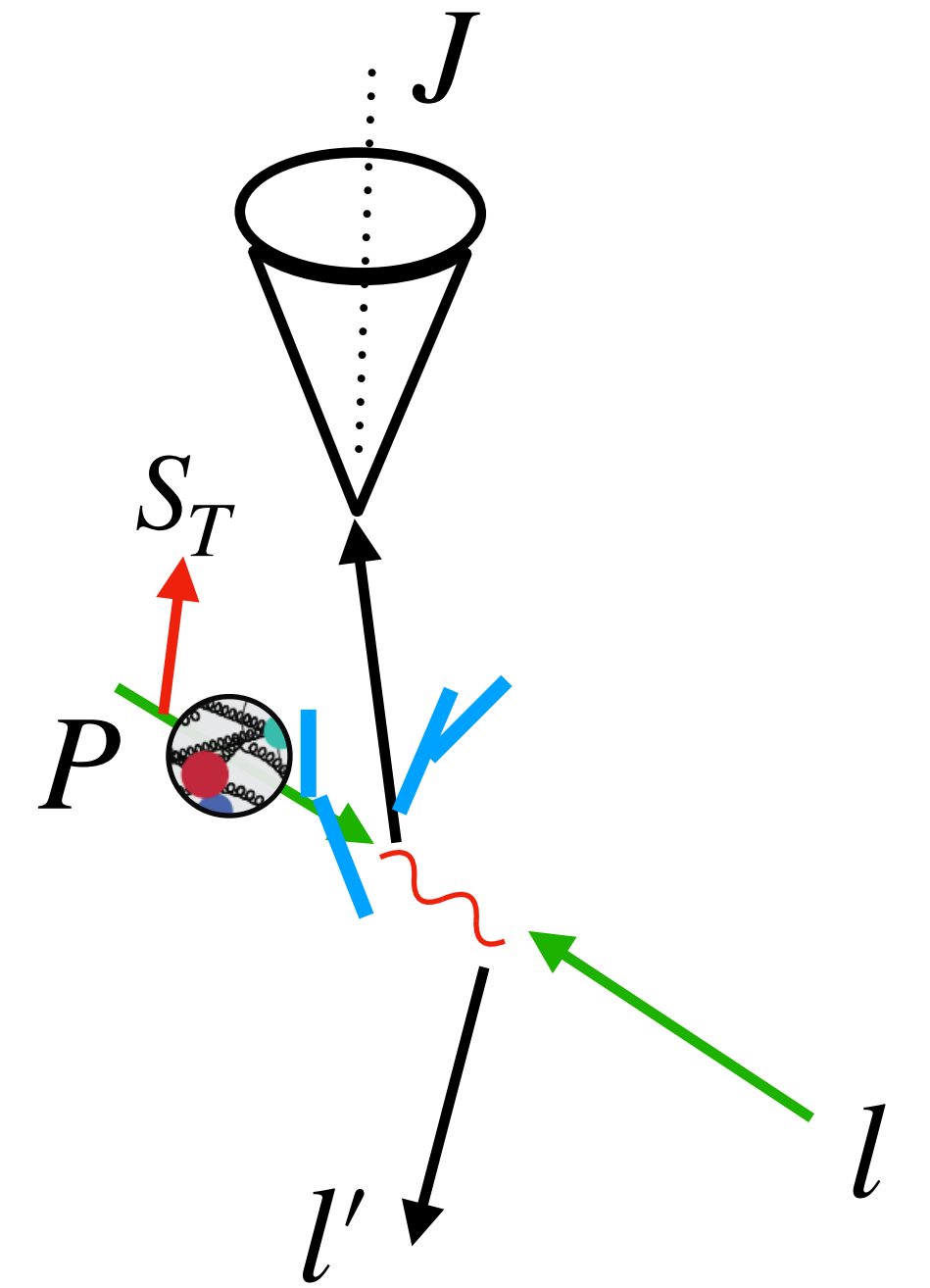
Reyes, Scimemi, Waalewijn, Zoppi, PRL 2018

# Revisit jet production in DIS

$$l p(P, s_T) \rightarrow \gamma^*(q) \rightarrow l' J(P_J) + X$$

In small  $q_T$  limit

$$\begin{aligned} \sigma &= \frac{1}{2s} \int [dl'] \frac{e^4 e_q^2}{Q^4} \frac{1}{2} L_{\mu\nu} \frac{1}{2N_c} \int d^2x_T e^{iq_T \cdot x_T} \gamma_{\alpha\beta}^\mu \gamma_{\beta'\alpha'}^\nu \\ &\times \int dx^- e^{iq^+ x^-} \langle P, s_T | \xi_n(0) \xi_n^\dagger(x^-, x_T) | P, s_T \rangle_{\alpha'\alpha} \\ &\times \int [dP_J] \int dx^+ e^{iq^- x^+} \langle 0 | \xi_{\bar{n}}(x^+, x_T) | J X_{\bar{n}} \rangle_\beta \langle J X_{\bar{n}} | \xi_{\bar{n}}^\dagger(0) | 0 \rangle_{\beta'} \end{aligned}$$




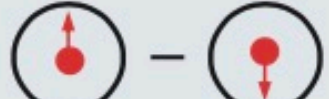




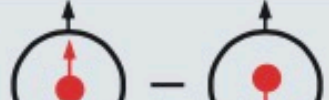
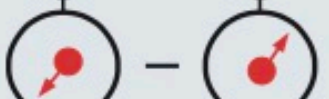
- Only consider photon exchange

SCET quark field:  $\xi_{n(\bar{n})} = W_{n(\bar{n})}^\dagger \chi_{n(\bar{n})}$

- $L_{\mu\nu}$  is leptonic tensor

# Decomposition $\Phi$ into different Dirac structures

$$\Phi(\zeta, p_T) = \frac{1}{2} \left\{ f_1 \not{n} - f_{1T}^\perp \frac{\epsilon_T^{\rho\sigma} p_{T\sigma} s_{T\sigma}}{M} \not{n} + g_{1s} \gamma_5 \not{n} \right. \\ \left. + h_{1T} \frac{[s_T, \not{n}] \gamma_5}{2} + h_{1s}^\perp \frac{[p_T, \not{n}] \gamma_5}{2M} + i h_1^\perp \frac{[p_T, \not{n}]}{2M} \right\} + \dots$$

TMDs		Quark polarization		
		Unpolarized (U)	Longitudinally polarized (L)	Transversely polarized (T)
Nucleon polarization	U	$f_1$  Unpolarized		$h_1^\perp$  Boer-Mulders
	L		$g_{1L}$  Helicity	$h_{1L}^\perp$  Longi-transversity
	T	$f_{1T}^\perp$  Sivers	$g_{1T}$  Trans-helicity	$h_1$  Transversity $h_{1T}^\perp$  Pretzelosity

# Decomposition of $\Delta$ into different Dirac structures

For single hadron production:

Collins function  $\rightarrow$  T-odd

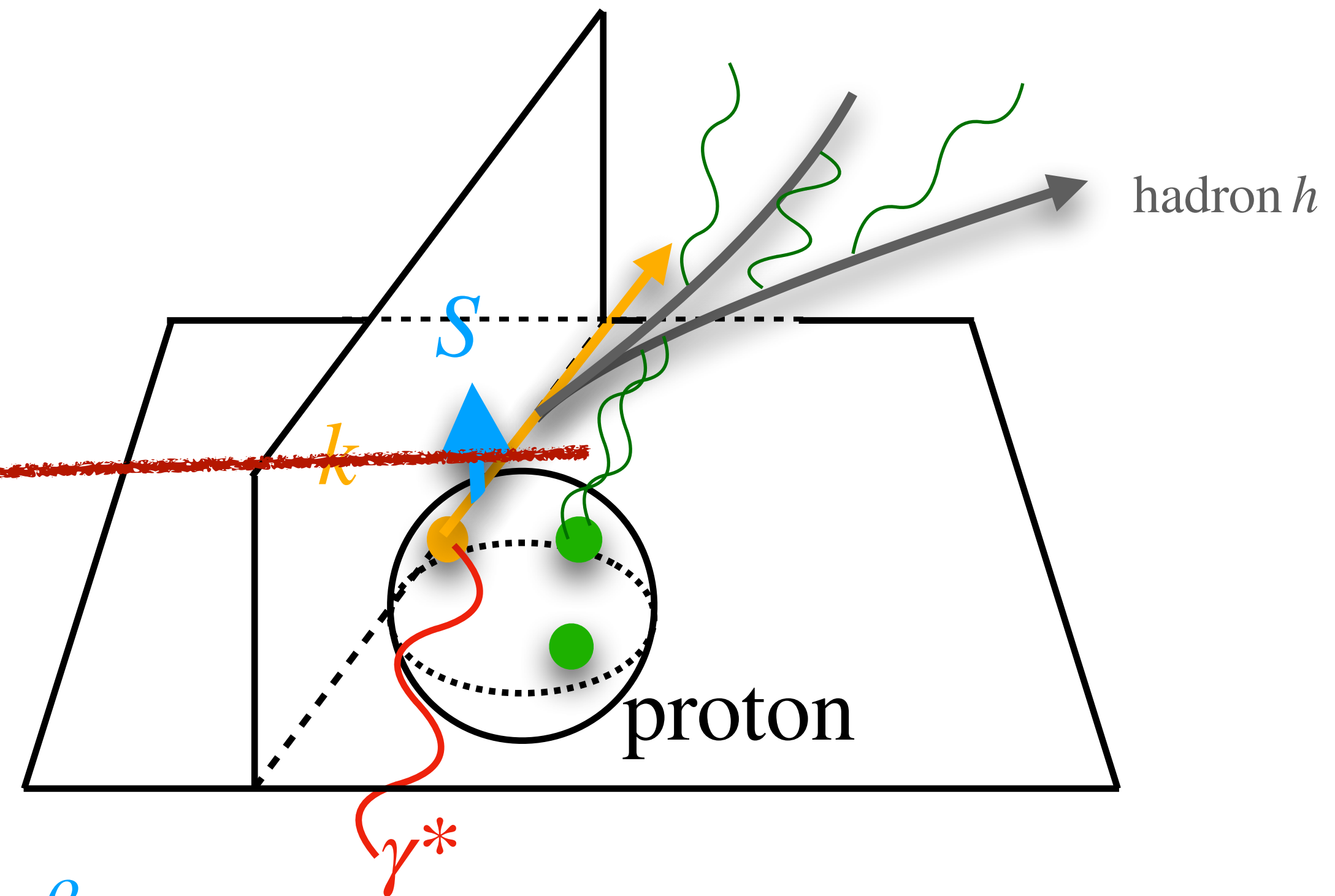
$$\Delta(z, k_T) = \frac{1}{2} \left\{ D_1 \vec{n} + i H_1^\perp \frac{[k_T^\perp, \vec{n}]}{2M_h} \right\} + \dots$$

Final state interactions generate an asymmetry

To observe the asymmetry (Collins 2002):

- h is not in the k-S plane
- Non-perturbative

$$\propto \epsilon_{\alpha\beta\mu\nu} k_\perp^\alpha s_\perp^\beta P_h^\mu n^\nu$$





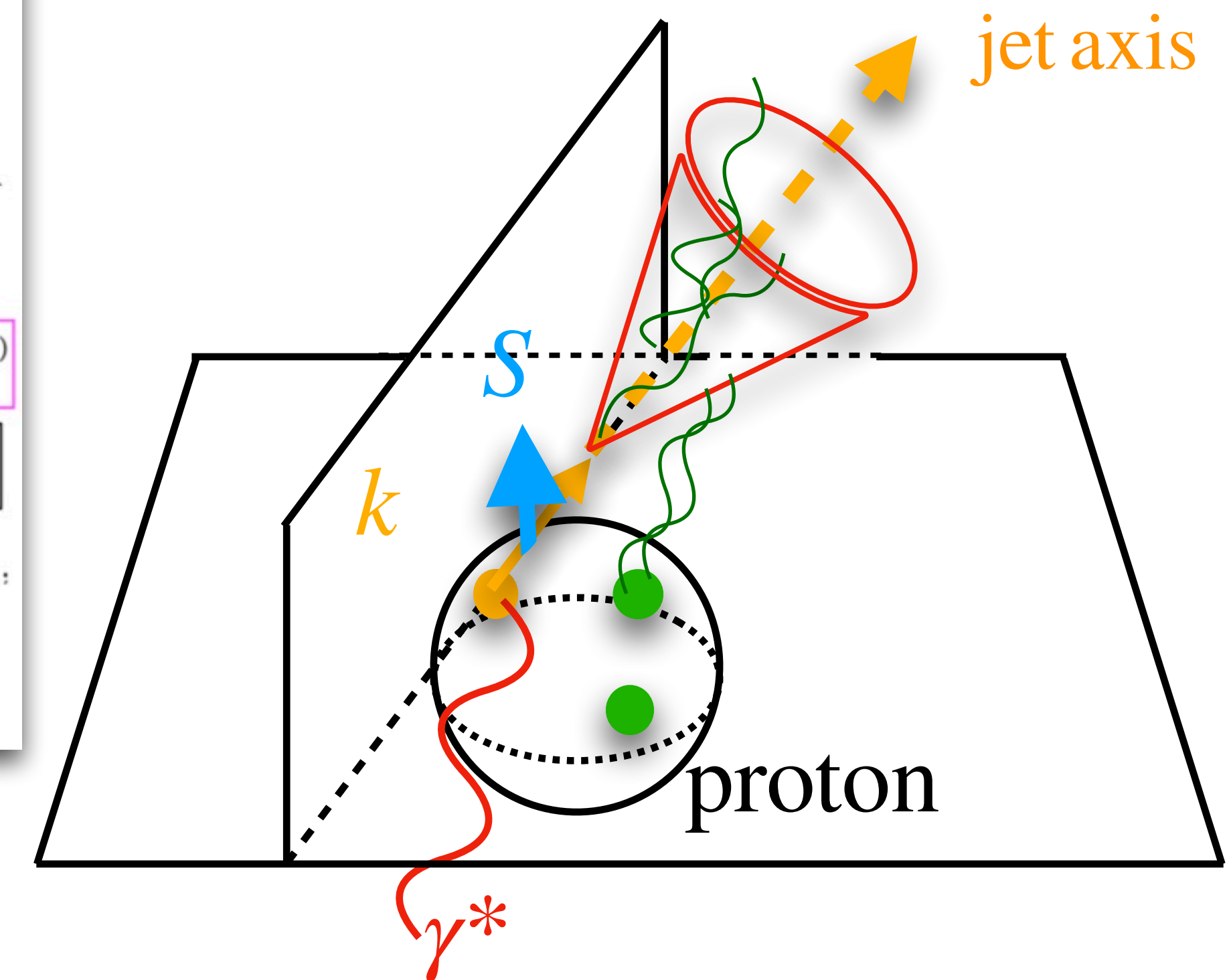
# Decomposition of $\mathcal{J}$ into different Dirac structures

$$\mathcal{J}(z, k_T) = \frac{1}{4} \left\{ J_1 \not{n} + i J_T \frac{[k_T, \not{n}]}{2M_h} \right\} + \dots$$

Diagram illustrating the decomposition of the cross-section  $d^6\sigma$  into various Dirac structures. The cross-section is given by:

$$\sigma(\phi, \phi_S) \equiv \frac{d^6\sigma}{dx dy dz d\phi d\phi_S dP_{KT}^2} = \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\epsilon)} \left( 1 + \frac{\gamma^2}{2x} \right)$$

The decomposition includes terms like  $F_{UU,T} + \epsilon F_{UUL} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi - \phi_S) F_{UT}^{\sin(\phi - \phi_S)}$  and  $F_{UU}^{\cos(2\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)}$ . The diagram uses boxes to group terms into categories: Cahn-effect + BM ⊗ Collins, Worm-gear (Kotzinian-Mulders) ⊗ Collins, BM ⊗ Collins, Sivers ⊗ D1, Worm-gear, Transversity ⊗ Collins, and Pretzelosity ⊗ Collins. Red 'X' marks indicate terms that are not present in the decomposition, while blue checkmarks indicate terms that are present.



- Jet is in the  $k$ - $S$  plane
- perturbative

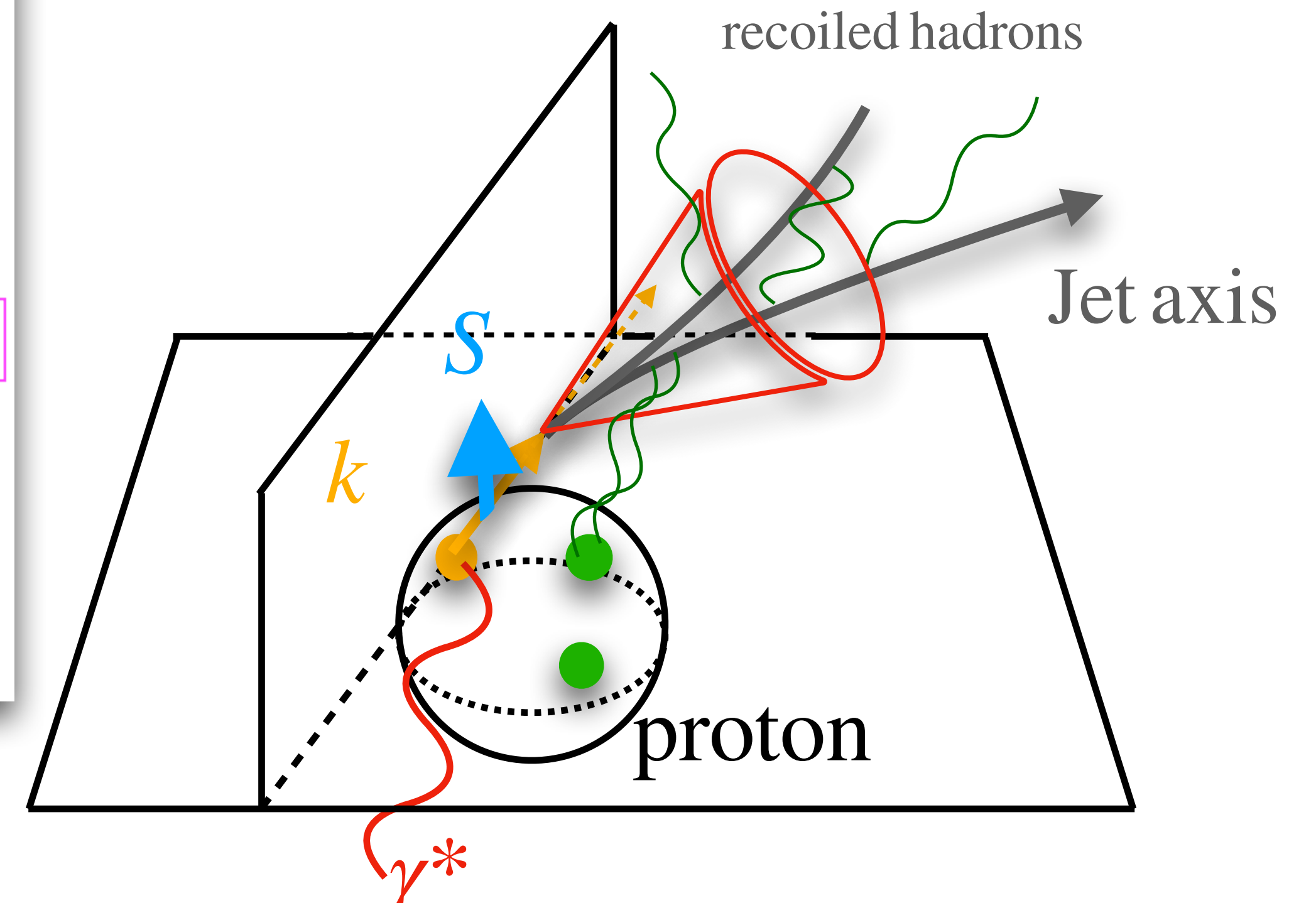
T-even jets limit the power of jet probe!

# Decompose quark field correlation functions into different Dirac structures

$$\mathcal{J}(z, k_T) = \frac{1}{4} \left\{ J_1 \not{n} + i J_T \frac{[k_T', \not{n}]}{2M_h} \right\} + \dots \quad \text{Not true} \rightarrow \text{existence of T-odd jet component}$$

$$d^6\sigma \equiv \frac{d^4\sigma}{dx dy dz d\phi d\phi_S} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left( 1 + \frac{\gamma^2}{2x} \right) \left\{ \begin{aligned} & F_{UU} + F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} \cos\phi F_{UU}^{\cos\phi} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} + \sqrt{2\epsilon(1-\epsilon)} \sin\phi F_{LU}^{\sin\phi} \\ & + S_L \left[ \sqrt{2\epsilon(1-\epsilon)} \sin\phi F_{UL}^{\sin\phi} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] + S_L \lambda_e \left[ \sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos\phi F_{LL}^{\cos\phi} \right] \\ & + |S_T| \left[ \sin(\phi - \phi_S) \left( F_{UT,L}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) + \epsilon \sin(\phi + \phi_S) F_{UT,L}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT,L}^{\sin(3\phi - \phi_S)} \right] \\ & + \sqrt{2\epsilon(1+\epsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \end{aligned} \right\};$$

Labels for Dirac structures in the diagram:  
 - Cahn-effect + BM ⊗ Collins (top left)  
 - Worm-gear (Kotzinian-Mulders) ⊗ Collins (top middle)  
 - BM ⊗ Collins (top right)  
 - Sivers ⊗ D1 (middle left)  
 - Worm-gear ⊗ D1 (bottom left)  
 - Transversity ⊗ Collins (bottom middle)  
 - Pretzelosity ⊗ Collins (bottom right)



Lai, Liu, Wang, HX, 2021

Jet axes can be different with different artificial algorithms

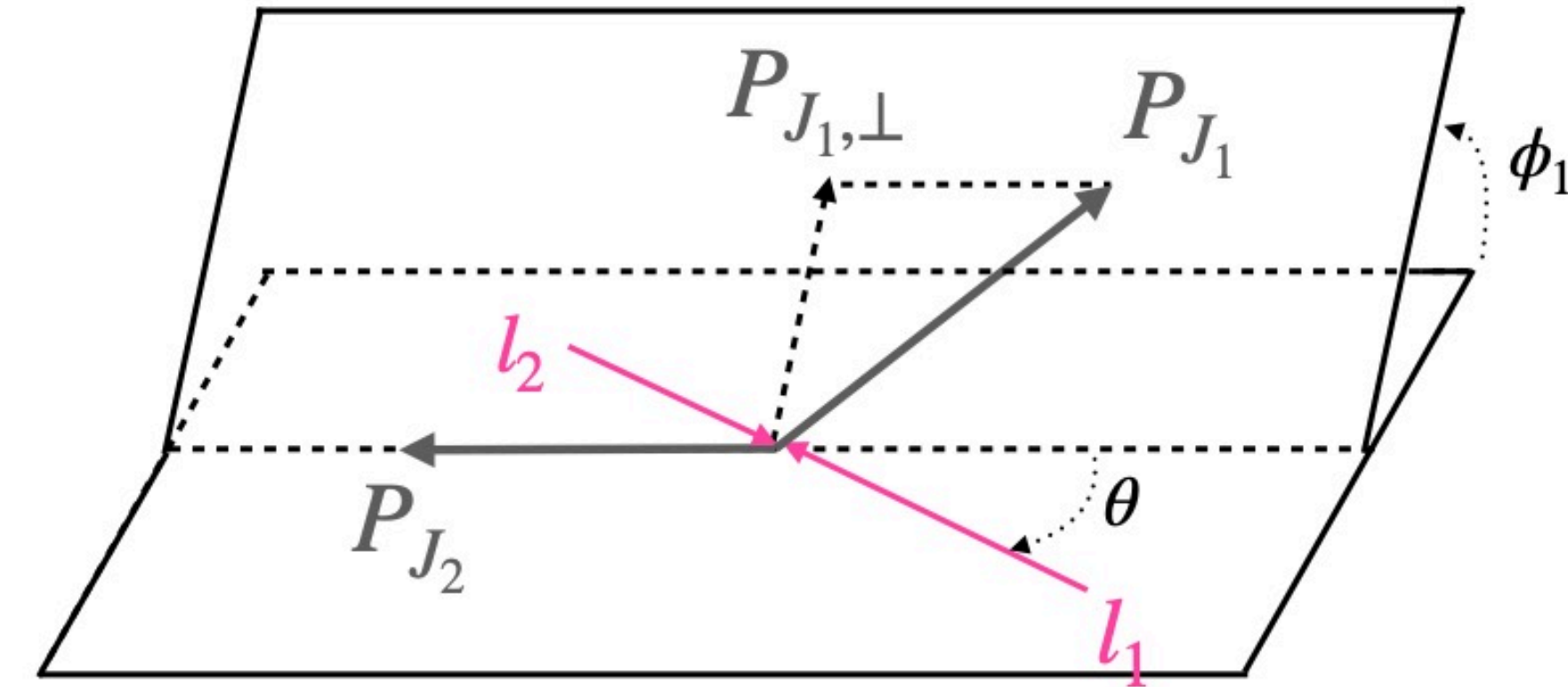
- Jet is not in the k-S plane
- Can be non-perturbative

Cal, Neill, Ringer, Waalewijn, JHEP 2020

# Constrain T-odd jet in $e^+e^-$ collisions

Azimuthal asymmetry:

$$R^{J_1 J_2} = 1 + \cos(2\phi_1) \frac{\sin^2 \theta}{1 + \cos^2 \theta} \frac{F_T(q_T)}{F_U(q_T)}$$



Factorization for WTA jet  $\mathcal{J}_{\beta\beta'}^q(z, k_T, R) = \delta(1 - z) J_{\beta\beta'}^q(k_T) + \mathcal{O}\left(\frac{k_T^2}{p_J^2 R^2}\right)$

$$F_U = q_T \sum_q e_q^2 \int \frac{db b}{2\pi} J_0(q_T b) J^q(b) \bar{J}^q(b)$$

Reyes, Scimemi, Waalewijn, Zoppi, PRL 2018

$$F_T = q_T \sum_q e_q^2 \int \frac{d^2 b}{(2\pi)^2} e^{-iq_T \cdot b} \left( 2 \frac{q_T^\alpha q_T^\beta}{q_T q_T} + g^{\alpha\beta} \right) \partial_{b^\alpha} J_T^q(b) \partial_{b^\beta} \bar{J}_T^q(b)$$

# Azimuthal asymmetry in $e^+e^-$

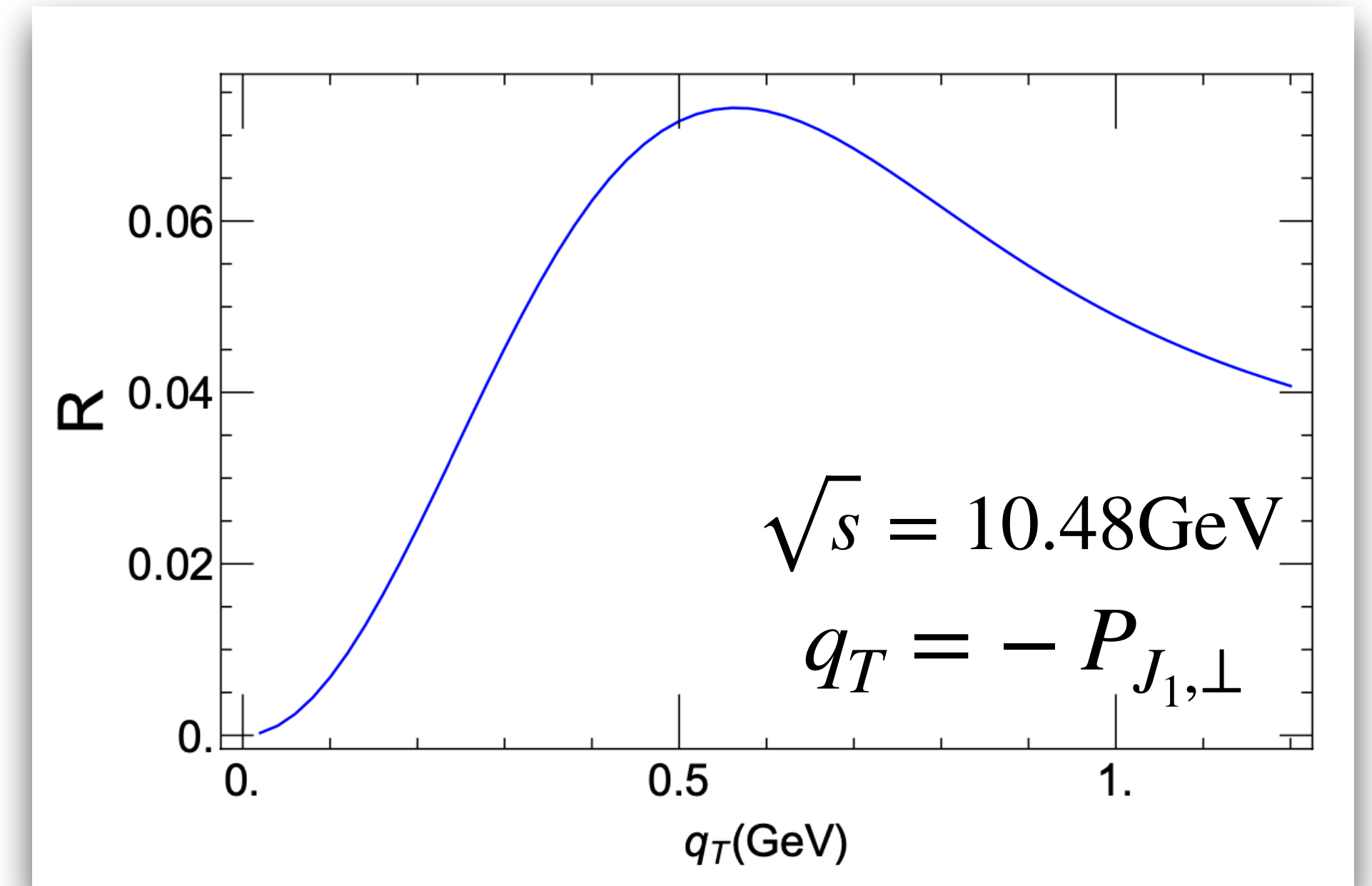
$$R = 2 \int d \cos \theta \frac{d\phi_1}{\pi} \cos(2\phi_1) R^{J_1 J_2}$$

$$J(b) \bar{J}(b) = e^{-S_{pert.} - S_{NP}^J} (1 + \mathcal{O}(\alpha_s)) r_q(Q_{J_{h_1}}) r_{\bar{q}}(Q_{J_{h_2}})$$

$$S_{pert.} = \int_{\mu_b^2}^s \frac{d\mu^2}{\mu^2} \left( A \log \frac{s}{\mu^2} + B \right)$$

$$\partial_{b^\alpha} J_T^q \partial_{b^\beta} \bar{J}_T^{\bar{q}} = e^{-S_{pert.} - S_{NP}^T} \frac{b^\alpha b^\beta}{4} \mathcal{N}_q^h(b) \mathcal{N}_{\bar{q}}^h(b)$$

- Follow the parametrization for Collins TMD (Kang, Prokudin, Sun, Yuan, 2015)
- This is for illustration only, because the non-perturbative T-odd jet functions are yet to be determined.



Can be extracted directly from BELLE or BaBar data

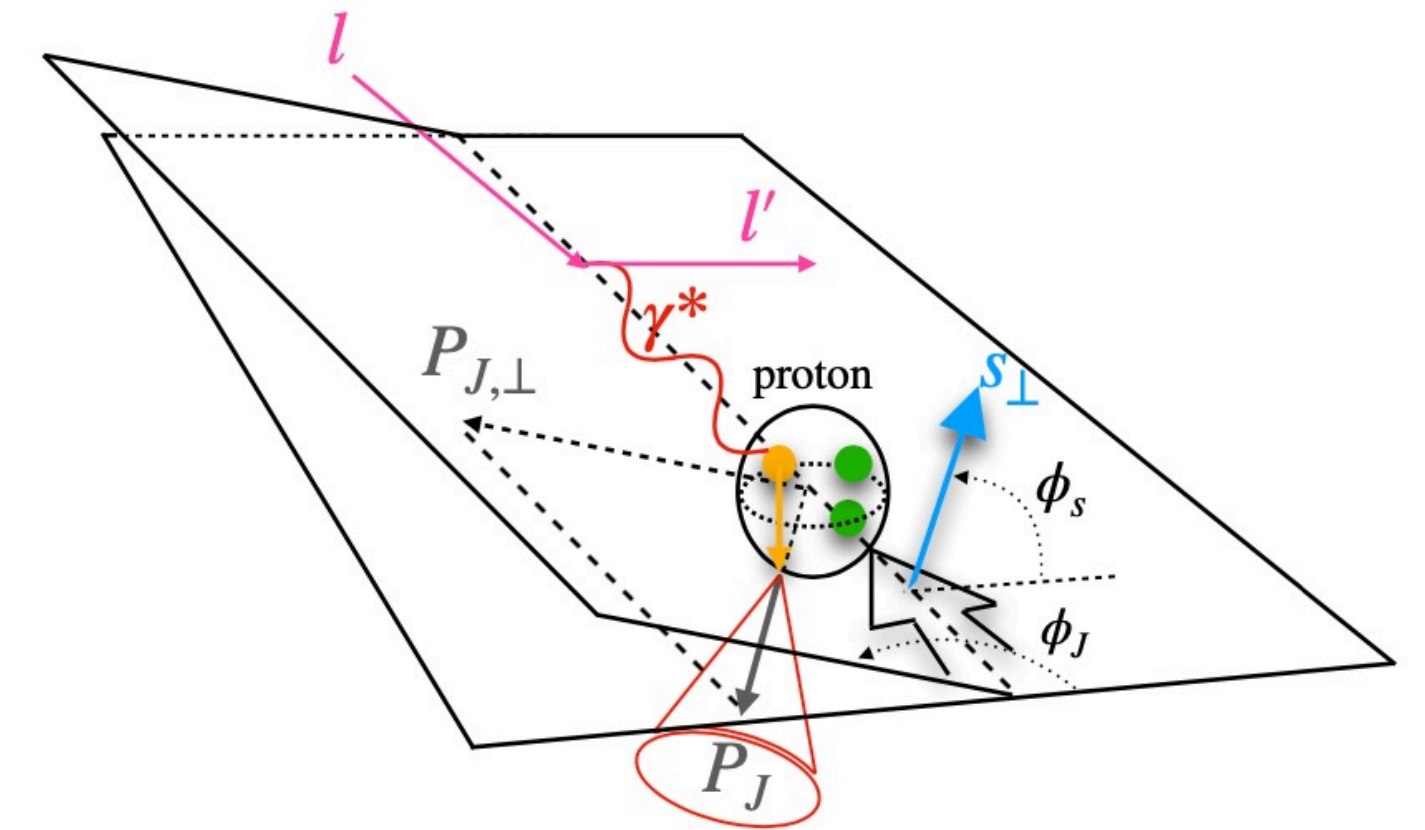
The lower the jet energy the better (if statistics guaranteed)

# Using T-odd jets to constrain the nucleon tensor charge

## Factorization in DIS

$$A(\zeta, y, \phi_s, \phi_J, P_{J\perp}) = 1 + \epsilon |s_\perp| \sin(\phi_J + \phi_s) \frac{F_{UT}}{F_{UU}}$$

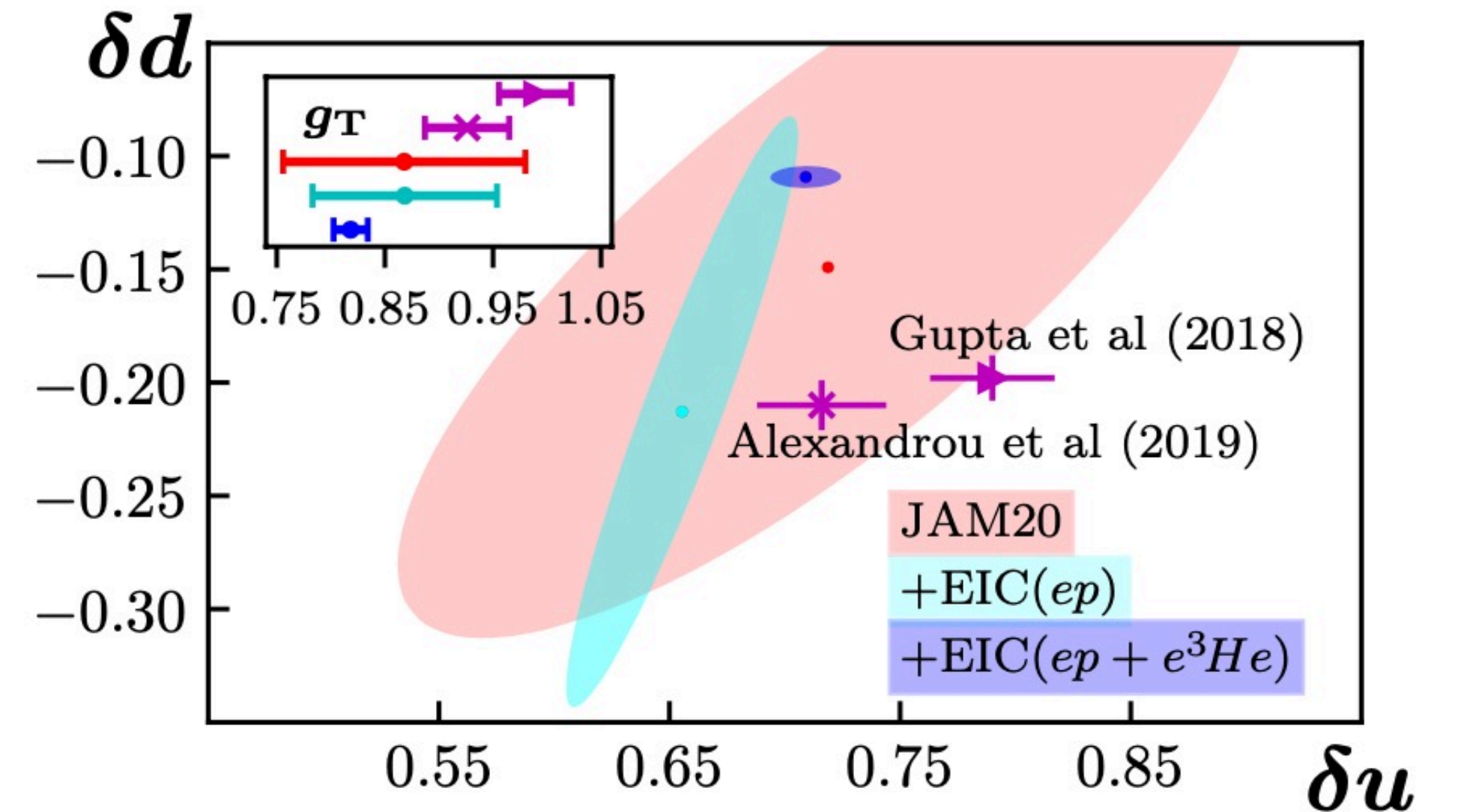
$$F_{UT} = \sum_q e_q^2 \int \frac{d^2b}{4\pi^2} e^{-iP_{J\perp} \cdot b} i \frac{P_{J\perp}^\alpha}{P_{J\perp}} \zeta h_1^q(\zeta, b) \partial_{b\alpha} J_T^q(b)$$



Nucleon tensor charge:  $g_T = \delta u - \delta d$

$$\delta u = \int_0^1 dx (h_1^u(x) - h_1^{\bar{u}}(x)), \quad \delta d = \int_0^1 dx (h_1^d(x) - h_1^{\bar{d}}(x))$$

T-odd jet with jet charge measurement can serve as a fantastic probe for nucleon tensor charge.



Gamberg, Kang, Pitonyak, Prokudin, Sato, PLB 2021

# Conclusion

- Jet is a powerful tool in many fields.
- Two fundamental difficulties in conventional jet

Flavor separation -> Jet charge

Limited power of jet probing on nucleon spin structure -> T-odd jet

- More comprehensive and precise analyses will come soon - stay tuned!

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**Thanks for your attention!**