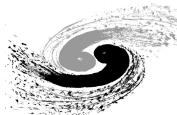


A new method for amplitude with multi fermion line
—The development of FDC program

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- 1 *WHY(Background)*
- 2 *HOW(Method)*
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WHY

Background

- In high energy physics, perturbation theory is widely used to do precise calculation.
- Our FDC system is a complete system of automatic perturbation calculation.

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- Our FDC system is a complete system of automatic perturbation calculation.
- With the development of particle physics experiment technology, especially the increase of collision particle energy, the multi-particle final state process will be more and more important.
 - (1) LEP II: $e^+e^- \rightarrow w^+w^- \rightarrow 4f$, nearly 100 processes
 - (2) Future high energy collider (ILC, CLIC, and etc): $e^+e^- \rightarrow t\bar{t} \rightarrow 6f$In recent years, there are more and more researches for multi-particle final state process

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- It's difficult for existing programs on the market to do this job.
- Helicity amplitude: some trouble in dealing with massive particle

HOW(Improve the fermion line calculation)

Step1: Simplification of spinor

One term for a fermion line in Feynman amplitude can be generally expressed as:

$$M = \bar{u}(n_1, k_1, s_1) B_m u(n_2, k_2, s_2) = u^+(n_1, k_1, s_1) \gamma^0 B_m u(n_2, k_2, s_2), \quad (1)$$

$$u(n, k, s) = \begin{cases} \frac{\hat{k} + m}{\sqrt{k^0 + m}} \mu(+, s) = \hat{k}' \mu(+, s), & n = + \\ \frac{\hat{k} - m}{\sqrt{k^0 + m}} \mu(-, s) = \hat{k}' \mu(-, s), & n = - \end{cases} \quad (2)$$

$$\mu(+, s) = \begin{pmatrix} \chi_s \\ 0 \end{pmatrix}, \quad \mu(-, s) = \begin{pmatrix} 0 \\ \chi_s \end{pmatrix}, \quad \chi_s = \begin{cases} \begin{pmatrix} 1 & 0 \end{pmatrix}^T, & s = 1 \\ \begin{pmatrix} 0 & 1 \end{pmatrix}^T, & s = 2 \end{cases}$$

$$k'^\mu = \left(\sqrt{k^0 + m}, \frac{\vec{k}}{\sqrt{k^0 + m}} \right) \quad (3)$$

HOW(Improve the fermion line calculation)

Step1: Simplification of spinor

Substituting the new form, one can get

$$\begin{aligned} M &= \mu(n_1, s_1)^+ \hat{k}'_1 \gamma_0 B_m \hat{k}'_2 \mu(n_2, s_2) \\ &= \mu(n_1, s_1)^+ \gamma_0 A_n \mu(n_2, s_2) \end{aligned} \quad (4)$$

$$\mu(+, s) = \begin{pmatrix} \chi_s \\ 0 \end{pmatrix}, \quad \mu(-, s) = \begin{pmatrix} 0 \\ \chi_s \end{pmatrix}, \quad \chi_s = \begin{cases} \begin{pmatrix} 1 & 0 \end{pmatrix}^T, & s = 1 \\ \begin{pmatrix} 0 & 1 \end{pmatrix}^T, & s = 2 \end{cases}$$

$$A_n \equiv \hat{k}'_1 B_m \hat{k}'_2, \quad n = m + 2. \quad (5)$$

Meanwhile, if there is a γ_5 in front of B_m ,

$$M = \bar{u}(n_1, k_1, s_1) \gamma_5 B_m u(n_2, k_2, s_2) \quad (6)$$

the final result equals:

$$M = -\mu(n_1, s_1)^+ \gamma_0 \gamma_5 A_n \mu(n_2, s_2). \quad (7)$$

HOW(Improve the fermion line calculation)

Step2: Simplification of Dirac matrices product

In general, one can redefine

$$A_n \equiv \hat{p}_n \cdots \hat{p}_2 \hat{p}_1, \quad n = 2, 3, \dots \quad (8)$$

Form the calculation, the general result with n p_i , whose form has relationship with odevity, and equals:

$$A_n = \hat{p}_n \cdots \hat{p}_2 \hat{p}_1 = \begin{cases} \hat{p}_a \gamma^0 - i \gamma_5 \hat{p}_b \gamma^0, & n \text{ is even} \\ \hat{p}_a - i \gamma_5 \hat{p}_b, & n \text{ is odd} \end{cases} \quad (9)$$

If there is a γ_5 at the front, the general form is

$$\gamma_5 A_n = \gamma_5 \hat{p}_n \cdots \hat{p}_2 \hat{p}_1 = \begin{cases} -i(\hat{p}_b \gamma^0 + i \gamma_5 \hat{p}_a \gamma^0), & n \text{ is even} \\ -i(\hat{p}_b + i \gamma_5 \hat{p}_a), & n \text{ is odd} \end{cases} \quad (10)$$

HOW(Improve the fermion line calculation)

Step3: Simplification of fermion line

The process introduced above should get and use final analytic results.

$$M = \mu(n_1, s_1)^+ \gamma_0 (\hat{p}_a \gamma^0 - i \gamma_5 \hat{p}_b \gamma^0) \mu(n_2, s_2)$$
$$= \mu(n_1, s_1)^+ \begin{pmatrix} p_a^0 - ip_b^3 & -ip_b^1 - p_b^2 & p_a^3 - ip_b^0 & p_a^1 - ip_a^2 \\ -ip_b^1 + p_b^2 & p_a^0 + ip_b^3 & p_a^1 + ip_a^2 & -p_a^3 - ip_b^0 \\ -p_a^3 + ip_b^0 & -p_a^1 + ip_a^2 & -p_a^0 + ip_b^3 & ip_b^1 + p_b^2 \\ -p_a^1 - ip_a^2 & p_a^3 + ip_b^0 & ip_b^1 - p_b^2 & -p_a^0 - ip_b^3 \end{pmatrix} \mu(n_2, s_2).$$

$$\mu(+, s) = \begin{pmatrix} \chi_s \\ 0 \end{pmatrix}, \quad \mu(-, s) = \begin{pmatrix} 0 \\ \chi_s \end{pmatrix}, \quad \chi_s = \begin{cases} \begin{pmatrix} 1 & 0 \end{pmatrix}^T, & s = 1 \\ \begin{pmatrix} 0 & 1 \end{pmatrix}^T, & s = 2 \end{cases}$$

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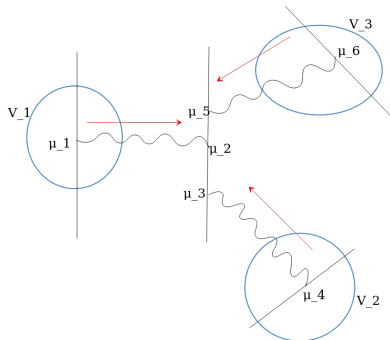
In fact, for a specific fermion line, both sides are particles:

$$M = \begin{cases} p_a^0 - ip_b^3, & s_1 = 1 \quad \text{and} \quad s_2 = 1 \\ -ip_b^1 - p_b^2, & s_1 = 1 \quad \text{and} \quad s_2 = 2 \\ -ip_b^1 + p_b^2, & s_1 = 2 \quad \text{and} \quad s_2 = 1 \\ p_a^0 + ip_b^3, & s_1 = 2 \quad \text{and} \quad s_2 = 2 \end{cases} \quad (11)$$

HOW(Improve the fermion line calculation)

Step4: Simplification of amplitude

- (1) contract to vector boson propagator;
- (2) contract the vector obtained in to other fermion line;
- (3) finally, all fermion lines can be contracted to one fermion line or scalar.

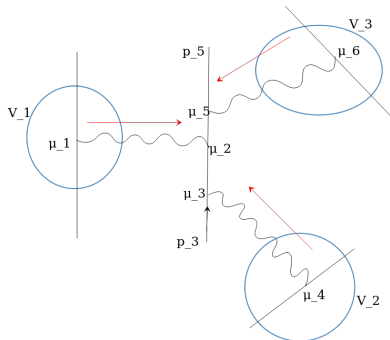


HOW(Improve the fermion line calculation)

Step4: Simplification of amplitude

- (1) contract to vector boson propagator;
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$$M = \bar{u}(p_5, s_5) \hat{V}_3(\hat{p} + m) \hat{V}_1(\hat{p}' + m) \hat{V}_2 u(p_3, s_3)$$



RESULTS

Compare the calculation speed with other programs

process	$e^+ + e^- \rightarrow c + \bar{c} + b + \bar{b}(8)$		$e^+ + e^- \rightarrow c + \bar{c} + c + \bar{c} + c + \bar{c}(576)$	
program	FDC	MadGraph	FDC	MadGraph
\sqrt{s}				
20	0.6	2.26	103.4	4008
50	0.5	2.28	90.4	3936
100	0.5	2.23	111.4	3990
200	0.5	2.24	127.9	4044
500	0.5	2.24	154.2	4002
1000	0.5	2.23	172.8	4002

Table: The expected time for generating 10000 events for a certain process. (\sqrt{s} in unit of GeV and expected time in unit of second.)

RESULTS

Compare the calculation speed with other programs

group	$1(e^+ + e^- \rightarrow c + \bar{c} + b + \bar{b})$		$2(e^+ + e^- \rightarrow c + \bar{c} + c + \bar{c} + c + \bar{c})$	
program	FDC	WHIZARD	FDC	WHIZARD
\sqrt{s}				
20	0.3	4	7.1	319
50	0.3	5	7.0	760
100	0.3	6	6.8	3935
200	0.2	8	6.8	6999
500	0.2	8	7.1	7905
1000	0.2	11	7.0	6277

Table: Group 1 represents the subprocess of $e^+ + e^- \rightarrow c + \bar{c} + b + \bar{b}$, which only contains the 4 diagrams of that. group 2 represents the subprocess of $e^+ + e^- \rightarrow c + \bar{c} + c + \bar{c} + c + \bar{c}$, which only contains 12 diagrams of that.

SUMMARY

- The numerical calculation of the amplitude of multi fermion lines process is improved obviously, and the velocity is increased a lot.
- The improved algorithm has been programmed, and embedded into FDC. Any multi fermion lines process can be calculated.
- The article is in preparation and will be completed soon.

End

The End
Thank you very much

$$A_2 = \begin{pmatrix} p_2 \cdot p_1 - i\varepsilon_{ijk}p_2^i p_1^j \sigma^k & (p_1^0 \vec{p}_2 - p_2^0 \vec{p}_1) \cdot \vec{\sigma} \\ (p_1^0 \vec{p}_2 - p_2^0 \vec{p}_1) \cdot \vec{\sigma} & p_2 \cdot p_1 - i\varepsilon_{ijk}p_2^i p_1^j \sigma^k \end{pmatrix}. \quad (12)$$

one can define two new p donated by p_a, p_b , which are

$$\begin{aligned} p_a &= (p_2 \cdot p_1, -p_2^0 \vec{p}_1 + p_1^0 \vec{p}_2), \\ p_b &= (0, \varepsilon_{ijk}p_2^i p_1^j). \end{aligned} \quad (13)$$

one can rewrite the form of Eq, which equals

$$A_2 = \hat{p}_a \gamma^0 - i\gamma_5 \hat{p}_b \gamma^0. \quad (14)$$

On the basis of the previous,

$$\begin{aligned} A_3 &= \hat{p}_3 \hat{p}_2 \hat{p}_1 = \hat{p}_3 (\hat{p}_a \gamma^0 - i \gamma_5 \hat{p}_b \gamma^0) \\ &= \hat{p}_3 \hat{p}_a \gamma^0 + i \gamma_5 \hat{p}_3 \hat{p}_b \gamma^0 \end{aligned} \quad (15)$$

$$\begin{aligned} \hat{p}_3 \hat{p}_a &= \hat{p}_{a'_1} \gamma^0 - i \gamma_5 \hat{p}_{b'_1} \gamma^0, \\ \hat{p}_3 \hat{p}_b &= \hat{p}_{a'_2} \gamma^0 - i \gamma_5 \hat{p}_{b'_2} \gamma^0 \end{aligned} \quad (16)$$

Hence, one can submit them into formula (15) and continue to calculate:

$$\begin{aligned} A_3 &= (\hat{p}_{a'_1} \gamma^0 - i \gamma_5 \hat{p}_{b'_1} \gamma^0) \gamma^0 + i \gamma_5 (\hat{p}_{a'_2} \gamma^0 - i \gamma_5 \hat{p}_{b'_2} \gamma^0) \gamma^0 \\ &= \hat{p}_{a'_1} - i \gamma_5 \hat{p}_{b'_1} + i \gamma_5 \hat{p}_{a'_2} + \hat{p}_{b'_2} \\ &= (\hat{p}_{a'_1} + \hat{p}_{b'_2}) - i \gamma_5 (\hat{p}_{b'_1} - \hat{p}_{a'_2}). \end{aligned} \quad (17)$$

So one can also define two new p donate by p_{a_2} , p_{b_2} , where

$$\begin{aligned} p_{a_2} &= p_{a'_1} + p_{b'_2}, \\ p_{b_2} &= p_{b'_1} - p_{a'_2}, \end{aligned} \tag{18}$$

and get the final result:

$$A_3 = \hat{p}_{a_2} - i\gamma/5\hat{p}_{b_2}. \tag{19}$$