

# Pion form factor in pQCD at the intermediate and large momentum transfers

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# Overview

- 1 Motivation
- 2 Formulism
- 3 Numerics
- 4 Outlook

"Hitting one constituent of a bound state with a photon initiates a complicated process if the bound state is not to fall apart. The momentum gained must be redistributed between all the constituents so that the whole convey can slew round into the new direction. A price is paid in terms of a reduced interaction strength between the photon and the bound state and this is known as the **electromagnetic form factor** - a function of the square of the (spacelike) 4-momentum,  $q^2$ , transferred from initial to final state."

[HPQCD collaboration, Phys.Rev.D96(2017)054501]

# Motivation

- † Open question: Hadron matrix element  
Factorization approaches: LCDAs  $\otimes$  HK  
Exclusive B,D decays: CPV and NP
- † From factor (FF): The most simple HME  
Pion e.m. FF: The most direct and easy way to study QCD  
Both for pert. and nonpert. information
- † Direct and indirect measurements:  $[0.25, 3] \text{ GeV}^2$ ,  
[CLEO 2005; Jefferson Lab 2001,2006,2008; et.al.]
- † Lattice QCD (LQCD): a few points of small  $Q^2$   
[HPQCD, UKQCD, MILC and Fermilab Lattice Collaborations 2004]  
[QCDSF/UKQCD 2007; ETM 2009; JLQCD and TWQCD 2009; JLQCD 2016, et.al.]  
up to  $1 \text{ GeV}^2$ , [χQCD 2020]
- † QCD-based approaches  
DSE:  $0\text{-}6 \text{ GeV}^2$     Lei Chang's talk;  
QCDSR:  $1\text{-}2 \text{ GeV}^2$ ;    LCSR:  $1\text{-}15 \text{ GeV}^2$ ;    pQCD:  $\gtrsim 10 \text{ GeV}^2$ ;  
[B.L Ioffe 1982]    [P. Ball 1991, V. Braun 1994,2000, SC 2020]    [Li 2001, Li 2012, SC 2014]

# Motivation

† External reason (large  $Q^2$ ) of the small relative distance  $\bar{z}_i \bar{p}_i \sim 1$   
→ OPE in twist(power);

Internal reason ( $M_Z, M_b$ ):  $\bar{z}_i \ll 1/\mu_{\bar{t}}$

† the controversy between LCSRs and pQCD in small-x region,

$\ln(Q^2/(p-q)^2)$

VS  $\ln(Q^2/\Lambda_{QCD}^2)$ ;

$\ln(Q^2/M^2)$  and  $\ln(Q^2/S_0)$

VS  $\ln(Q^2/K_T^2)$ ;

not like to occur

VS Sudakov suppression ;

"soft" + "hard"

VS Hard scattering ← **Leading power;**

† "k<sub>T</sub>" + multiparton

← **High powers;**

† Accuracy:

'q $\bar{q}$ ' + 'qg $\bar{q}$ ', twist four DAs

VS 'q $\bar{q}$ ', 'qg $\bar{q}$ ', twist three

← **LO**

[V. Braun 1994]

[Li 2001]

'q $\bar{q}$ ' + 'qg $\bar{q}$ ', twist six DAs

VS 'q $\bar{q}$ ', 'qg $\bar{q}$ ', twist three

← **NLO-2p**

[V. Braun 2000, J. Bijnens 2002]

[Li 2012, SC 2014]

## Contribution from twist four DAs

† LCSR: a visible enhancement at  $Q^2 > 10 \text{ GeV}^2$ ,  $\sim 40\%$ ,  
the same asymptotic behaviour as twist two  $\mathcal{O}(1/Q^4)$  [J. Bijmans 2002]  
LCSR holds in the intermediate energy region.

† Check the twist expansion in pQCD at large  $Q^2$

† Improve the precision of pQCD prediction

† In turn, with measurements, to determine the nonpert. parameters

† Pion, Kaon,  $SU(3)_f$  breaking

- Factorization formula

$$\begin{aligned}
 & \langle \pi^-(p_2) | J_{\mu}^{e.m.} | \pi^-(p_1) \rangle \\
 = & \int d z_1 d z_2 \langle \pi^-(p_2) | \left\{ \bar{d}_{\gamma}(0) \exp \left( i g_s \int_{z_2}^0 d\sigma_{\nu'} A_{\nu'}(\sigma) \right) u_{\beta}(z_2) \right\}_{kj} | 0 \rangle_{\mu_t} \\
 & \cdot H_{\gamma\beta\alpha\delta,\mu}^{ijkl}(z_2, z_1) \cdot \langle 0 | \left\{ \bar{u}_{\alpha}(z_1) \exp \left( i g_s \int_0^{z_1} d\sigma_{\nu} A_{\nu}(\sigma) \right) d_{\delta}(0) \right\}_{il} | \pi^-(p_1) \rangle_{\mu_t} \quad (1)
 \end{aligned}$$

- Separation: SD (propagator) and LD (Heisenberg operator).
- Hard kernel and nonlocal MEs  $\leftarrow$  spin structures (Fierz trans.).
- nonlocal MEs, defined in terms of LCDAs by twist.
- Truncated (factorizable) scale  $\mu_t$ .

- hard kernel associated with the lowest Fock state

$$H_{\gamma\beta\alpha\delta}^{ijkl}(z_1, z_2) = (-1) [ig_s\gamma_m]_{\alpha\beta} T^{ij} [(ie_q\gamma_\mu)S_0(0 - z_1)(ig_s\gamma_n)]_{\gamma\delta} T^{kl} [-iD_{mn}^0(z_1 - z_2)] , \quad (2)$$

- the free propagators in the coordinate space

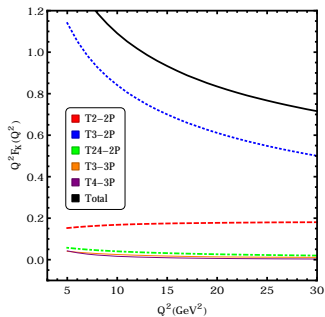
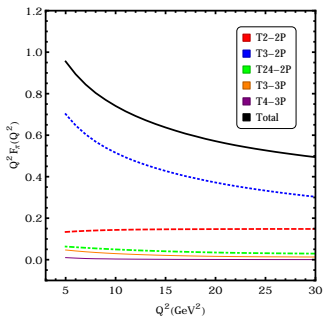
$$S_0(z) = \frac{i}{2\pi} \frac{\not{z}}{z^4} , \quad D_{mn}^0(z) = \frac{1}{4\pi} \frac{g_{mn}}{z^2} . \quad (3)$$

- nonlocal matrix elements in Eq. (1), the mesons breaking-up into a pair of soft quarks, different spin structures

$$\begin{aligned} & \langle 0 | \left\{ \bar{u}_\alpha(z_1) \exp \left( ig_s \int_0^{z_1} d\sigma_\nu A_\nu(\sigma) \right) d_\delta(0) \right\} | \pi^-(p_1) \rangle_{\mu t} \\ &= \frac{\delta_{il}}{3} \left\{ \frac{1}{4} (\gamma_5 \gamma^\rho)_{\delta\alpha} \langle 0 | \bar{u}(z_1) \exp \left( ig_s \int_0^{z_1} d\sigma_\nu A_\nu(\sigma) \right) (\gamma_\rho \gamma_5) d(0) | \pi^-(p_1) \rangle_{\mu t} \right. \\ & \quad + \frac{1}{4} (i\gamma_5)_{\delta\alpha} \langle 0 | \bar{u}(z_1) \exp \left( ig_s \int_0^{z_1} d\sigma_\nu A_\nu(\sigma) \right) (i\gamma_5) d(0) | \pi^-(p_1) \rangle_{\mu t} \\ & \quad + \frac{1}{8} \left( \sigma^{\tau\tau'} \gamma_5 \right)_{\delta\alpha} \langle 0 | \bar{u}(z_1) \exp \left( ig_s \int_0^{z_1} d\sigma_\nu A_\nu(\sigma) \right) (i\sigma_{\tau\tau'} \gamma_5) d(0) | \pi^-(p_1) \rangle_{\mu t} \\ & \quad \left. + \dots \right\} . \quad (4) \end{aligned}$$



- 2p-2p scattering:  $T2 \otimes T2$ ,  $T3 \otimes T3$ ,  $T2 \otimes T4$   
3p-3p scattering:  $T3 \otimes T3$ ,  $T4 \otimes T4$   
~~2p-3p, 3p-2p scattering~~ ← color transparency mechanism
- Gauge invariance:  $\propto k_T$  in 2p-2p cancels with 3p-3p scattering  
four categories:  $N_g = 0, 1, 2$  and the 3g vertex configuration [Li 2001]
- 3p-3p scattering:  
dominated by the gauge invariant configuration with 4g vertex
- Power expansion:  $\mathcal{O}(1) : \mathcal{O}\left(\frac{m_0^2}{Q^2}\right) : \mathcal{O}\left(\frac{\delta_P^2}{Q^2}\right) : \mathcal{O}\left(\frac{f_{3P}^2}{f_P^2 Q^2}\right) : \mathcal{O}\left(\frac{\delta_P^2}{Q^4}\right)$



- 2p-2p scattering:  $T2 \otimes T2$ ,  $T3 \otimes T3$  ← LO+NLO [Li 2012, SC 2014].
- Inputs: LQCD, QCDSR and  $\chi PT$  relation for  $m_0^\pi$ .
- The chiral enhancement at T3, more stronger in Kaon case.
- $SU(3)_f$  breaking  $\sim 30\%$ ,  $R_P(Q^2) \equiv F_P^{T2}/F_P^{T3-2p}$ ,  $A(Q^2) \equiv 1 - R_\pi/R_K$ .

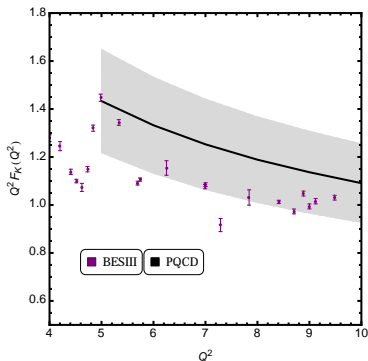
$Q^2(\text{GeV}^2)$	$Q^2 F_\pi^{\text{PQCD}}(Q^2)$	$Q^2 F_\pi^{\text{LCSR}s}(Q^2)$	$Q^2 F_K^{\text{PQCD}}(Q^2)$	$Q^2 F_K^{\text{LCSR}s}(Q^2)$
10	0.75(10)	0.51(15)	1.08(15)	0.76(22)

- Fruitful discussions with BES-III collaboration,

in particular with W.B. Yan and G.X Huang [BES-III 2019]

- $e^+e^- \rightarrow K^+K^-$  at  $\sqrt{s} = 2 - 3.08$  GeV

- Analytical continuum,  $\ln(-Q^2 - i\epsilon) = \ln(Q^2) - i\pi$   
 $\ln(k_T^2 - xQ^2 + i\epsilon) = \ln(k_T^2 - xQ^2) + i\pi\Theta(k_T^2 - xQ^2)$



- Jefferson Lab 12 GeV upgrade program ( $9 \text{ GeV}^2$ )  
to extract the nonpert. paras (besides the LQCD contribution).

*"provide new insights into the structure of the nucleon, the transition between the hadronic and quark/gluon description of matter, and the nature of quark confinement."* [JLAB CDR]

- Dispersion relation: spacelike  $\leftrightarrow$  timelike regions  
resonances (measurements) + high energy tail (pQCD)
- BEPCII, 4.7-4.95 GeV in 2021, STCF (2-7 GeV)

*"Physics highlights: QCD and hadronic physics, Flavor physics and CP violation, New Physics search."*

- nonpert. paras, fact. scale, Nucleons ...

The End, Thanks.