The Discussion of the Singularities in Triangle and Box Single Loop

Jia-Jun Wu

Collaborators: Qi Huang, Chao-Wei Shen

Phys.Rev.D 103 (2021) 1

第二届强子与重味物理理论与实验联合研讨会

2021. 3. 26

兰州大学

Content

- Background of Triangle Singularity
- What is Triangle Singularity ?
- How to search Triangle Singularity?
- How about Box Singularity
- Summary



Background of Triangle Singularity

L. D. Landau, Nucl. Phys. 13, no.1, 181-192 (1960)
S. Coleman, R. E. Norton, Nuovo Cim. 1965, 38, 438 - 442,
R. Karplus, C. M. Sommerfield, E. H. Wichmann, PR 1958, 111, 1187 - 1190.
J. D. Bjorken, Ph. D. Thesis, Stanford University, Stanford, CA, USA, 1959.
C. Schmid, Phys. Rev. 1967, 154, 1363,





Background of Triangle Singularity

L. D. Landau, Nucl. Phys. 13, no.1, 181-192 (1960)
S. Coleman, R.E. Norton, Nuovo Cim. 1965, 38, 438 - 442,
R. Karplus, C.M. Sommerfield, E.H. Wichmann, PR 1958, 111, 1187 - 1190.
J.D. Bjorken, Ph.D. Thesis, Stanford University, Stanford, CA, USA, 1959.
C. Schmid, Phys. Rev. 1967, 154, 1363,





Background of Triangle Singularity

中国科学

University of Chinese Academy of Scienc

L. D. Landau, Nucl. Phys. 13, no.1, 181-192 (1960)
S. Coleman, R.E. Norton, Nuovo Cim. 1965, 38, 438 - 442,
R. Karplus, C.M. Sommerfield, E.H. Wichmann, PR 1958, 111, 1187 - 1190.
J.D. Bjorken, Ph.D. Thesis, Stanford University, Stanford, CA, USA, 1959.
C. Schmid, Phys. Rev. 1967, 154, 1363,



F.K. Guo, X. H. Liu, S. Sakai PPNP 2020, 112, 103757

Structures	Processes	Loops	I/J	F Refs.
2.1 GeV [141]	$\gamma p^+ \rightarrow N^*(2030) \rightarrow K^+\Lambda(1405)$	$K^*\Sigma\pi$	Ι	[142]
$2.1 \ \text{GeV}$	$\pi^- p^+ \rightarrow K^0 \Lambda(1405), \ pp \rightarrow pK^+ \Lambda(140)$	5) $K^* \Sigma \pi$	Ι	[143]
1.88 GeV	$\Lambda_c^+ \to \pi^+ \pi^0 \pi \Sigma$	$\bar{K}^*N\bar{K}$	Ι	$[144, 145]^{a}$
N(1700) [10]	$N(1700) \rightarrow \pi \Delta$	$\rho N \pi$	Ι	[146]
N(1875) [10]	$N(1875) \rightarrow \pi N(1535)$	$\Sigma^*K\Lambda$	Ι	[147]
$\Delta(1700)$ [148–150]	$\gamma p \to \Delta(1700) \to \pi N(1535) \to p \pi^0 \eta$	$\Delta \eta p$	Ι	[151]
2.2 GeV [152]	$\Lambda_c^+ \to \pi^0 \phi p$	$\Sigma^* K^* \Lambda$	F	[153]
1.66 GeV [154, 155]	$\Lambda_c^+ \rightarrow \pi^+ K^- p$	$a_0\Lambda\eta, \Sigma^*\eta\Lambda$	F	[156]
$P_c(4450)$ [35]	$\Lambda_b^0 \to K^- J/\psi p$	$\Lambda(1890)\chi_{c1}$	p F	$[157-160]^{b}$
	0	$N(1900)\chi_{cl}$	р F	[159]
peaks relevant for P_c	$\Lambda^0_h o K^- J/\psi p$	$\bar{D}_{*I}\Lambda_c^{(*)}\bar{D}^{(*)}$) F	[36, 158]
	0 774	00 0		1 , 1
Structures	Processes	Loops	I/F	Refs.
$\rho(1480)$ [78, 79]	$\pi^- p \rightarrow \phi \pi^0 n$	$K^*\bar{K}K$	Ι	[80, 81]
$\eta(1405/1475)$ [82–86]	$\eta(1405/1475) \to \pi f_0$	K^*KK	I	[87–91] ^{a,b}
$f_1(1420)$ [92]	$f_1(1285) \to \pi a_0/\pi f_0$	K^*KK	Ι	[89, 93–95] ^b
$a_1(1420)$ [96, 97]	$a_1(1260) \rightarrow f_0 \pi \rightarrow 3\pi$	K^*KK	I	[97-99]
1.4 GeV [100]	$J/\psi \to \phi \pi^0 \eta / \phi \pi^0 \pi^0$	K^*KK	I	[101] ^b
1.42 GeV	$B^- \to D^{*0} \pi^- f_0(a_0), \tau \to \nu_\tau \pi^- f_0(a_0)$	K*KK	I	[102, 103]
1 (1010) [10]	$D_s^+ \to \pi^+ \pi^0 f_0(a_0), B_s^0 \to J/\psi \pi^0 f_0(a_0)$	K*KK	I	[104, 105]
$f_2(1810)$ [10]	$f_2(1640) \rightarrow \pi \pi \rho$	K*K*K	1	[106]
1.65 GeV	$\tau \to \nu_\tau \pi^- f_1(1285)$	K*K*K	1	[107]
1515 MeV	$J/\psi \to K^+ K^- f_0(a_0)$	φKK	1	[108]
2.85 GeV, 3.0 GeV	$B \rightarrow K \pi D_{s0}^{-}/K \pi D_{s1}$ $B^+ \rightarrow \pi^0 \pi^+ D^0$	$\overline{K}^{*0}D^{(*)*}K^{+}$	I F	[109, 110]
[4.01_4.02] CoV	$\frac{D_c \rightarrow \pi^* \pi^* D_s}{[\bar{D}^{*0} D^{*0}] \rightarrow \alpha Y}$	Λ · <i>D</i> · Λ Δ*0 Δ*0 Δ0	г Т	[111]
4015 MoV	$[D^+D^-] \rightarrow \gamma \chi$ $a^+a^- \rightarrow \gamma \chi$	מים ים ש מים∗ים יים	T	[112]
4015 MeV 4015 MeV	$e \ e \ \rightarrow /\pi$ $B \rightarrow K X \pi \ m/m \bar{n} \rightarrow X \pi + anything$	$D^{*0}\bar{D}^{*0}D^{0}$	ī	[115, 114]
$\Upsilon(11020)$ [117 118]	$e^+e^- \rightarrow Z_{\mu}\pi$	$B_1(5721)\bar{B}B^*$	î	[119, 120]
3.73 GeV	$X \rightarrow \pi^0 \pi^+ \pi^-$	$D^{*0}\bar{D}^0D^0$	F	[121]
[4.22, 4.24] GeV	$e^+e^- \rightarrow \gamma J/\psi \phi/\pi^0 J/\psi \eta$	$D^*_{-0(-1)}\bar{D}^{(*)}_s D^{(*)}_s$	F	[122]
[4.08, 4.09] GeV	$e^+e^- ightarrow \pi^0 J/\psi\eta$	$D^*_{r^{0}(s1)}\bar{D}^{(*)}_s D^{(*)}_s$	F	[122]
$Z_c(3900)$ [31, 32]	$e^+e^- ightarrow J/\psi \pi^+\pi^-$	$D_1 \overline{D} D^*$	F	[119, 123–127] ^c
		$D_0^*(2400)\bar{D}^*D$	F	[128, 129]
$Z_c(4020, 4030)$ [33, 130]	$e^+e^- \rightarrow \pi^+\pi^-h_c(\psi')$	$D_{1(2)}\bar{D}^{(*)}D^{(*)}$	F	[125]
X(4700) [131, 132]	$B^+ \to K^+ J/\psi \phi$	$K_1(1650)\psi'\phi$	F	[133]
$Z_c(4430)$ [30, 134]	$\bar{B}^0 \rightarrow K^- \pi^+ J/\psi$	$\bar{K}^{*0}\psi(4260)\pi^+$	F	[135]
$Z_c(4200)$ [136, 137]	$\bar{B}^0 \rightarrow K^- \pi^+ \psi(2S)$	$\bar{K}_{2}^{*}\psi(3770)\pi^{+}$	\mathbf{F}	[135]
	$\Lambda_b^0 o p \pi^- J/\psi$	$N^*\psi(3770)\pi^-$	\mathbf{F}	[135]
$X(4050)^{\pm}$ [138]	$\bar{B}^0 \rightarrow K^- \pi^+ \chi_{c1}$	$\bar{K}^{*0}X\pi^+$	\mathbf{F}	[139]
$X(4250)^{\pm}$ [138]	$\bar{B}^0 \rightarrow K^- \pi^+ \chi_{c1}$	$\bar{K}_{2}^{*}\psi(3770)\pi^{+}$	\mathbf{F}	[139]
$Z_b(10610)$ [34]	$e^+e^- \rightarrow \Upsilon(1S)\pi^+\pi^-$	$B_J^* \overline{B}^* B$	F	[128]

(b)=



















(2) Particle 3 catch up particle 2.











Why Triangle Singularity interesting ?

1. It is a pure kinematic effect

-> Model independent

2. The effect of Loop

-> Understand hadronic Loop contribution 3.Provide a peak structure

 \rightarrow May mixing with resonance

4. To extract the nature of hadron

 \rightarrow Study the coupling in the special point

5. •••••





1.Threshold

2.Width of the internal particle of the loop





1.Threshold



The corresponding values are listed in Table 1. The ATS peak will then stay close to the normal threshold, as illustrated in Fig. 4(d). In this sense, it would be difficult to distinguish the ATS peak from the pole structure in the invariant mass of the $J/\psi\pi$. We shall come back to the relevant issue later in this Section. It should be

X. H. Liu, M. Oka and Q. Zhao, PLB 753, 297-302 (2016)

2.Width of the internal particle of the loop

(d)





University of Chinese Academy of Science



Our proposal

1.Threshold

2. Width of the internal particle of the loop





Our proposal

1.Threshold

Far away from p η threshold. Singularity point is 1.563 GeV of p $\eta(\pi)$ invariant mass.

2. Width of the internal particle of the loop

All narrow internal particles, $J/\,\psi\,,$ p, η

3.Unknow vertex

All vertices are constrained from experimental data







Our proposal

1.Threshold

Far away from $p\eta$ threshold. Singularity point is 1.563 GeV of $p \eta(\pi)$ invariant mass.

2. Width of the internal particle of the loop

All narrow internal particles, $J/\,\psi\,,$ p, η

3.Unknow vertex

All vertices are constrained from experimental data













































How about Box Singularity









































Summary

We discuss how to detect the Triangle Singularity, we should consider the ${\bullet}$ following three point, 4.0x10 3.5x10-3.2x10 3.0x10⁻⁶ 1.Threshold 3.0x10-6 2.8x10 2.5x10 dl^T/dm 2.0x10-8 2. Width of the internal particle of the loop 1.5x10 1.0x10⁻⁸ 3.Unknow vertex 5.0x10⁻⁹

0.0

1.5

1.6

m_m (GeV)

17

1.8

- We propose a process $\psi(2s) \rightarrow p \,\overline{p} \,\eta$ through J/ψ , p, η loop.
- We need to find new processes to detect Triangle Singularity.
- For the box singularity, it needs more hard work

Backup











1. In the complex plane, the integral routine will be fixed between two singularity.

$$\int \frac{dx}{(x+i\epsilon)^2} \to \text{Convergence} \quad \int \frac{dx}{(x+i\epsilon)(x-i\epsilon)} \to \text{Divergence}$$
$$q_b = q_{on} - i\epsilon'$$

2. For $cos\theta$, it is a one order singularity, thus, it will be convergence except at the edge.

$$\cos\theta = -1 \text{ or } 1$$

PEANSA

University of Chinese Academy of Science



TS in $\psi(2S) \rightarrow p\overline{p}\eta$ process vs Schmid theorem

Effect of the tree diagram $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\bar{p})$



Schmid theorem can't be applied directly here:

1. $p\eta \rightarrow p\eta$ is not purely elastic process, $p\pi$ channel must couples with it, which is effectively included in the imaginary part of N^* propagator.

2. $\psi(2S) \rightarrow \bar{p}(N^* \rightarrow p\eta)$ will modify the amplitude to $\left|t_{J/\psi}^{\text{Tree}} + t_{N^*}^{\text{Loop}} + t_{N^*}^{\text{Tree}}\right|^2$, where $t_{\text{elastic}}^{\text{Loop}}$. 3. Contribution of $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\bar{p})$ can be removed by applying a cut $m_{p\bar{p}} < m_{J/\psi}$.





Improvement: Selection of internal particles

A lesson: Widths of internal particles should not be too narrow !

Or:

1. Peak of TS is too sharp to be detected by experiment. \implies Requirement of the resolution is too high.

