

# Top, Higgs, Diboson and Electroweak Fit to the Standard Model Effective Field Theory

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Based on 2012.02779 J. Ellis, MM, K. Mimasu, V. Sanz, T. You



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# The Standard Model Effective Field Theory

Powerful tool for capturing deviations from the SM and performing indirect searches for new physics.

**Model independent:** assume the BSM physics is heavy

$$E \ll \Lambda$$

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i}{\Lambda^2} \mathcal{O}_i + \dots$$

We restrict to dimension-6 operators.

# The Standard Model Effective Field Theory

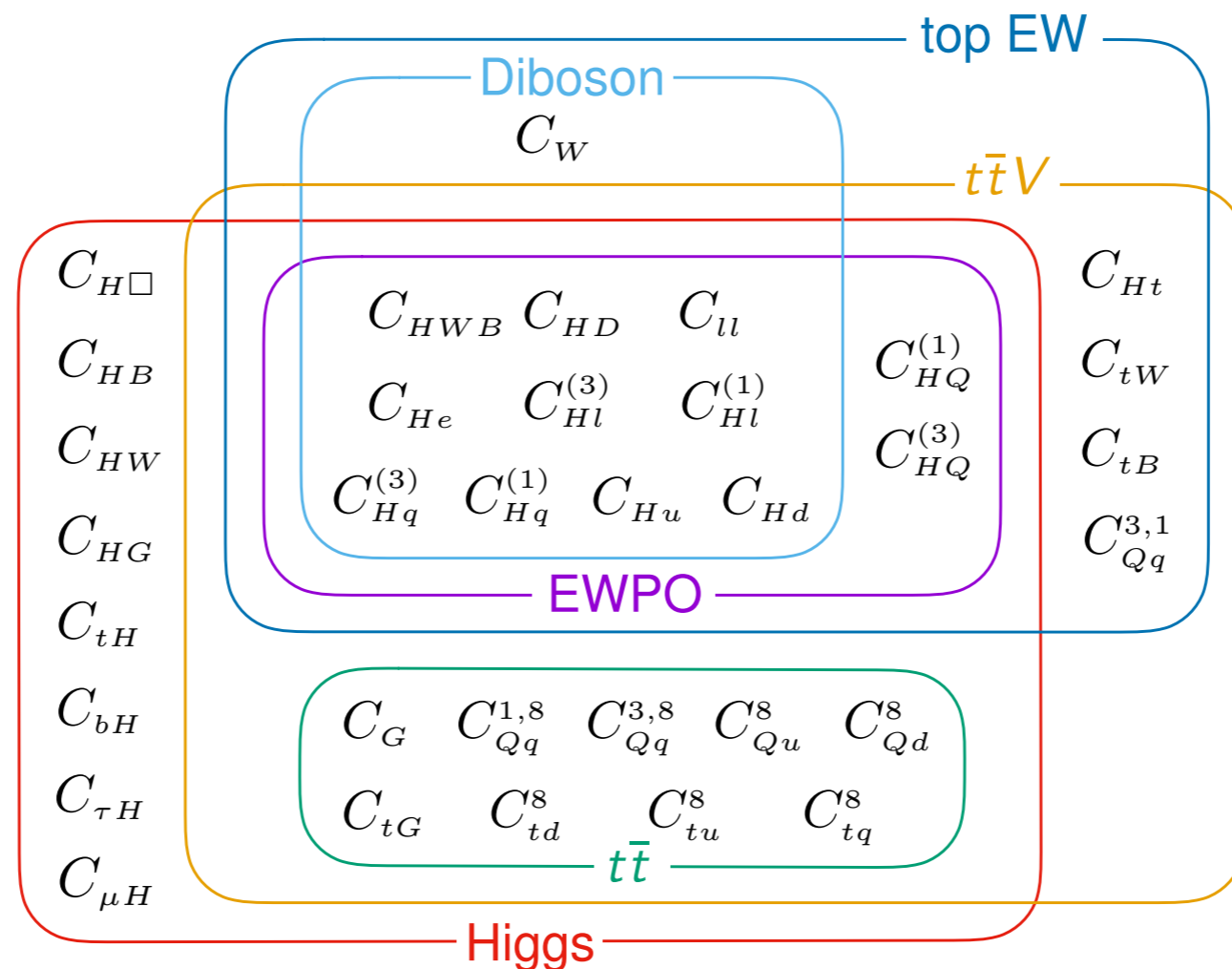
Warsaw basis  
[1008.4884  
Grzadkowski et. al]

$X^3$		$H^6$ and $H^4 D^2$		$\psi^2 H^3$	
$\mathcal{O}_G$	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$\mathcal{O}_H$	$(H^\dagger H)^3$	$\mathcal{O}_{eH}$	$(H^\dagger H)(\bar{l}_p e_r H)$
$\mathcal{O}_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$\mathcal{O}_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	$\mathcal{O}_{uH}$	$(H^\dagger H)(\bar{q}_p u_r \tilde{H})$
$\mathcal{O}_W$	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$\mathcal{O}_{HD}$	$(H^\dagger D^\mu H)^* (H^\dagger D_\mu H)$	$\mathcal{O}_{dH}$	$(H^\dagger H)(\bar{q}_p d_r H)$
$\mathcal{O}_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 H^2$		$\psi^2 XH$		$\psi^2 H^2 D$	
$\mathcal{O}_{HG}$	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	$\mathcal{O}_{eW}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I H W_{\mu\nu}^I$	$\mathcal{O}_{Hl}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$
$\mathcal{O}_{H\tilde{G}}$	$H^\dagger H \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	$\mathcal{O}_{eB}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$\mathcal{O}_{Hl}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$\mathcal{O}_{HW}$	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	$\mathcal{O}_{uG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G_{\mu\nu}^A$	$\mathcal{O}_{He}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{e}_p \gamma^\mu e_r)$
$\mathcal{O}_{H\tilde{W}}$	$H^\dagger H \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	$\mathcal{O}_{uW}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	$\mathcal{O}_{Hq}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$
$\mathcal{O}_{HB}$	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	$\mathcal{O}_{uB}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	$\mathcal{O}_{Hq}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$\mathcal{O}_{H\tilde{B}}$	$H^\dagger H \tilde{B}_{\mu\nu} B^{\mu\nu}$	$\mathcal{O}_{dG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) H G_{\mu\nu}^A$	$\mathcal{O}_{Hu}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$
$\mathcal{O}_{HWB}$	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	$\mathcal{O}_{dW}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H W_{\mu\nu}^I$	$\mathcal{O}_{Hd}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$
$\mathcal{O}_{H\tilde{W}B}$	$H^\dagger \tau^I H \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$\mathcal{O}_{dB}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) H B_{\mu\nu}$	$\mathcal{O}_{Hud}$	$i(\tilde{H}^\dagger D_\mu H)(\bar{u}_p \gamma^\mu d_r)$
$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
$\mathcal{O}_{ll}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	$\mathcal{O}_{ee}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$\mathcal{O}_{le}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$\mathcal{O}_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	$\mathcal{O}_{uu}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	$\mathcal{O}_{lu}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$\mathcal{O}_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$\mathcal{O}_{dd}$	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{ld}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$\mathcal{O}_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	$\mathcal{O}_{eu}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$\mathcal{O}_{qe}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$\mathcal{O}_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$\mathcal{O}_{ed}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$\mathcal{O}_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$\mathcal{O}_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$\mathcal{O}_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$\mathcal{O}_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		$B$ -violating			
$\mathcal{O}_{ledq}$	$(\bar{l}_p^j e_r)(\bar{d}_s^k q_t^j)$	$\mathcal{O}_{duq}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{ijk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^{\gamma j})^T C l_t^k]$		
$\mathcal{O}_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	$\mathcal{O}_{quu}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{ijk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$\mathcal{O}_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	$\mathcal{O}_{quq}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jkn} \varepsilon_{km} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^{\gamma m})^T C l_t^n]$		
$\mathcal{O}_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	$\mathcal{O}_{duu}$	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$\mathcal{O}_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				

# A global fit to the SMEFT

Each operator contributes to multiple datasets

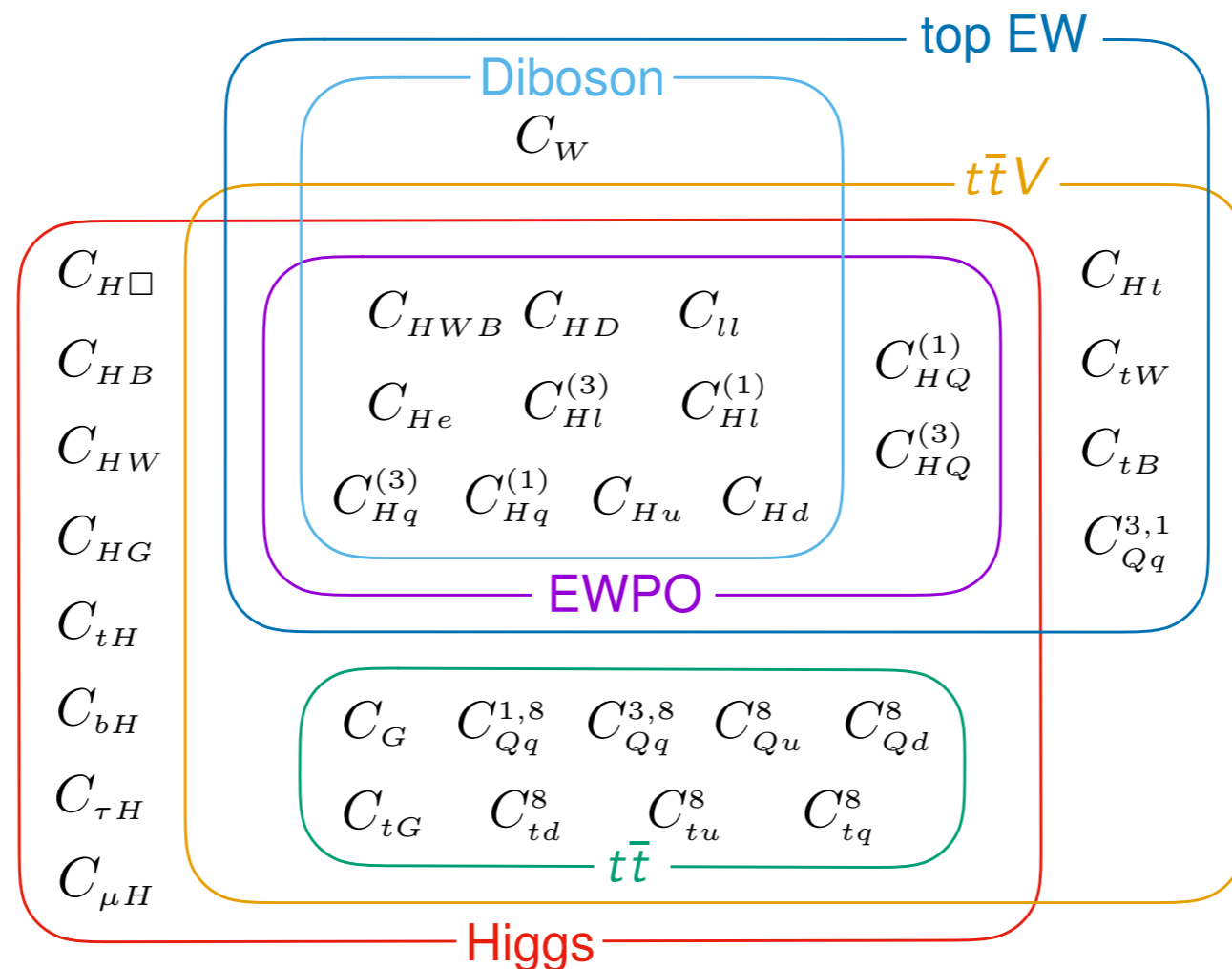
- ▶ expect an interplay between sectors



# A global fit to the SMEFT

This highlights the need for a **global fit** to understand and parameterise the deviations and correlations between operators and sectors.

We include data from top, diboson, Higgs and EWPO in a fit to 34 dim-6 operators.



# SMEFT conventions

- Warsaw basis
- Neglect CP-violating operators
- Two flavour scenarios:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i}{\Lambda^2} \mathcal{O}_i$$

- Flavour universal  $SU(3)^5$
- Top-specific flavour scenario singles out top couplings [1802.07237]

$$SU(3)^5 \rightarrow SU(2)_q \times SU(2)_u \times SU(3)_d \times SU(3)_l \times SU(3)_e$$

In both flavour scenarios we also include 4 Yukawa operators which explicitly break these flavour symmetries:

$$\mathcal{O}_{\tau H}, \mathcal{O}_{\mu H}, \mathcal{O}_{bH}, \mathcal{O}_{tH}$$

# Measurements

- 341 statistically independent measurements
- Correlation information included from published covariance matrices

## Higgs: 72

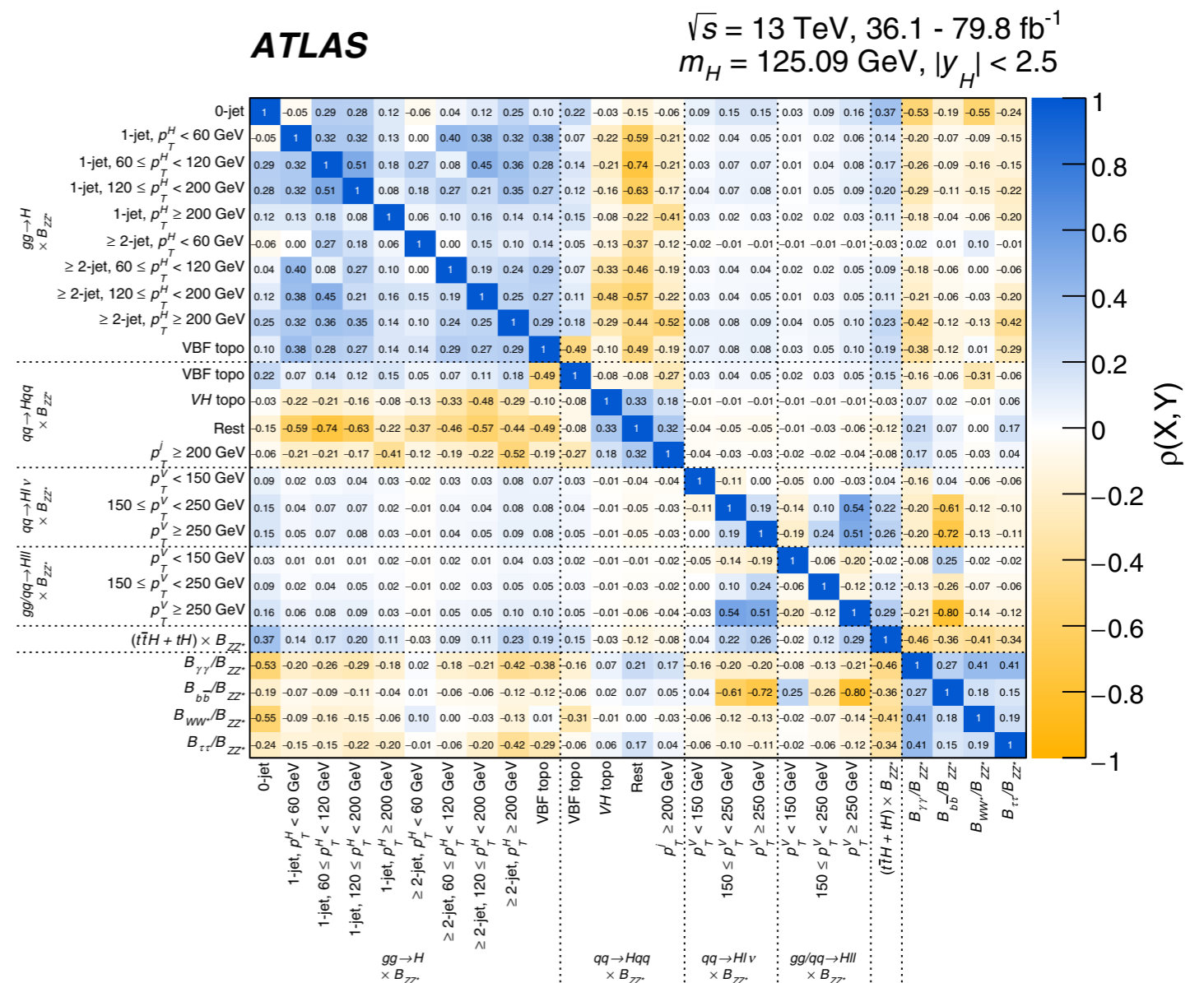
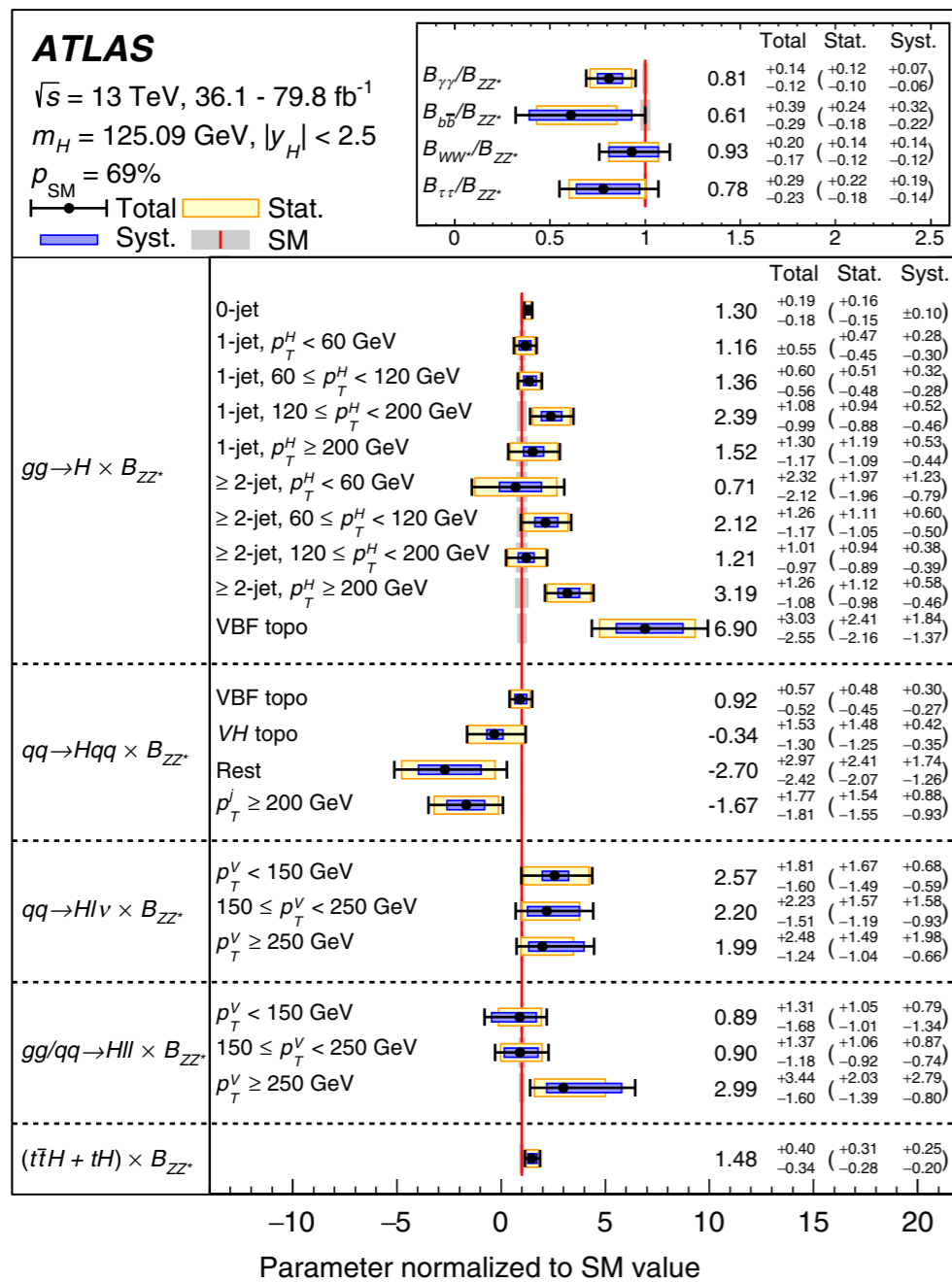
- Signal strength combinations (LHC Run 1 and Run 2)
- STXS combination (LHC Run 2)
- Measurements of

$$H \rightarrow Z\gamma \quad H \rightarrow \mu\mu$$

# Higgs STXS data from LHC Run 2 ATLAS

ATLAS Run 2 STXS combination [HIGG-2018-57, Phys. Rev. D 101 (2020) 012002]

A total of 21 STXS bins with published correlation matrix





# Measurements

- 341 statistically independent measurements
- Correlation information included from published covariance matrices

## Higgs: 72

- Signal strength combinations (LHC Run 1 and Run 2)
- STXS combination (LHC Run 2)
- Measurements of

$$H \rightarrow Z\gamma \quad H \rightarrow \mu\mu$$

## Diboson: 118

- LHC and LEP measurements of

$$WW, WZ, Zjj$$

## EWPO: 14

LEP, Tevatron, LHC measurements

$$\{\Gamma_Z, \sigma_{\text{had.}}^0, R_l^0, A_{FB}^l, A_l, R_b^0, R_c^0, A_{FB}^b, A_{FB}^c, A_b, A_c, M_W\}.$$

## Top: 137

LHC measurements of

$$t\bar{t}, \quad t\bar{t} + V, \quad \text{single top}$$

# Fitting methodology

$$\chi^2(C_i) = (\vec{y} - \vec{\mu}(C_i))^T V^{-1} (\vec{y} - \vec{\mu}(C_i))$$

$\vec{y}$ : vector of observables with covariance matrix  $V$

Predictions:  $\mu_\alpha(C_i) = \mu_\alpha^{SM} + H_{\alpha i} C_i$

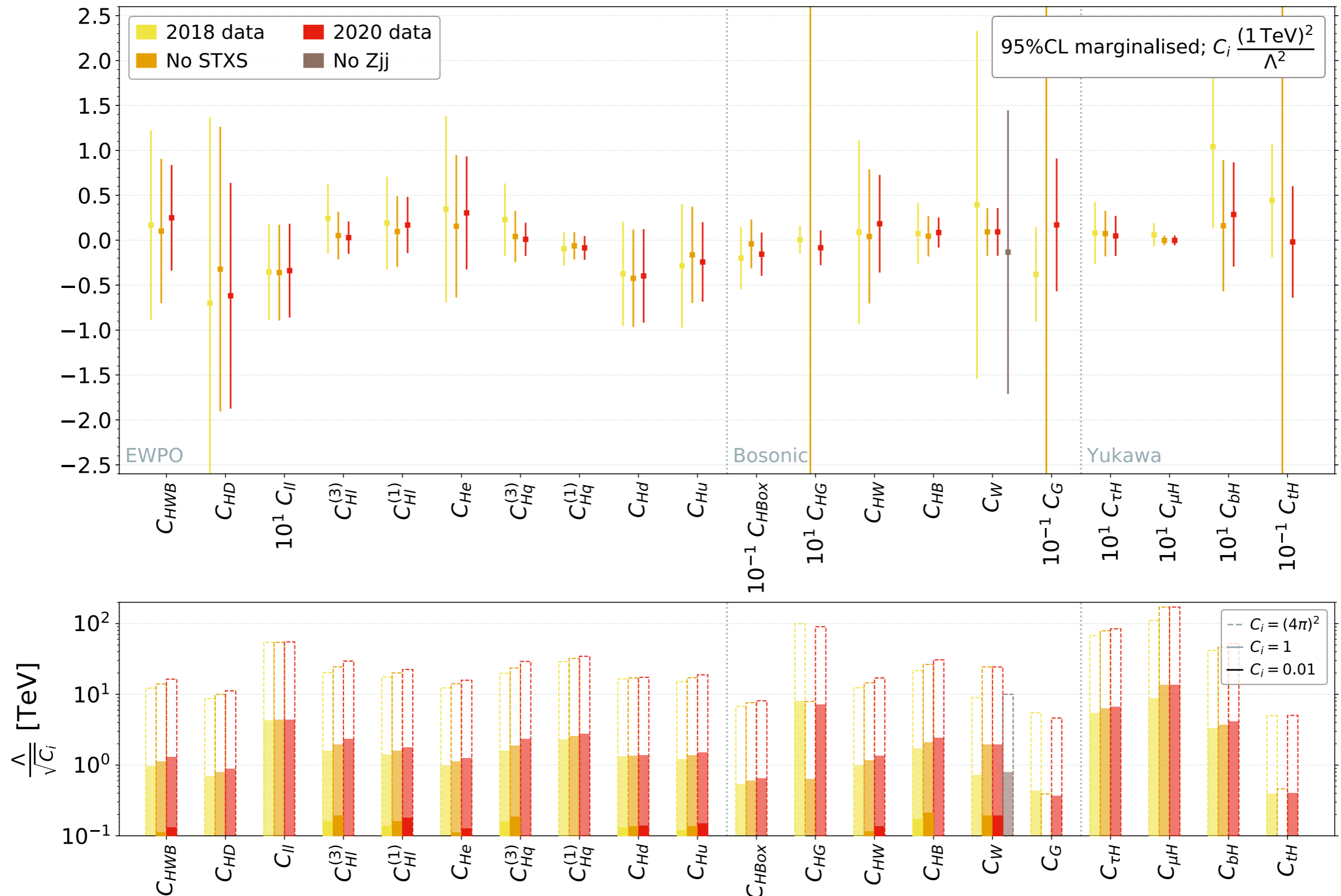
i.e. restricting to  $\mathcal{O}(\Lambda^{-2})$  in the EFT expansion

Best-fit WC:  $\hat{\vec{C}} = (H^T V^{-1} H)^{-1} H^T V^{-1} (\vec{y} - \vec{\mu}^{SM})$

Covariance:  $U = (H^T V^{-1} H)^{-1} = F^{-1}$

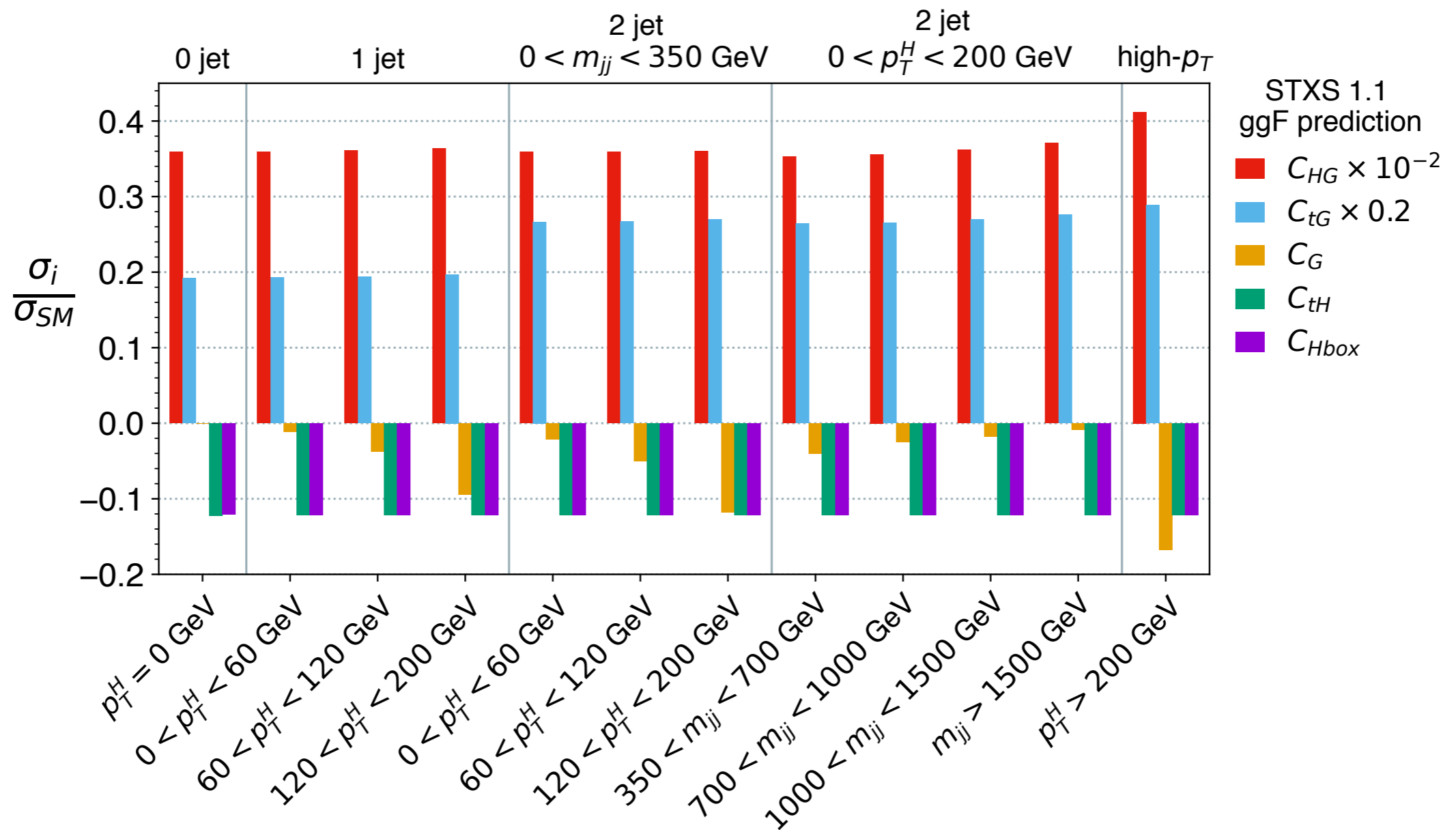
(Fisher information)

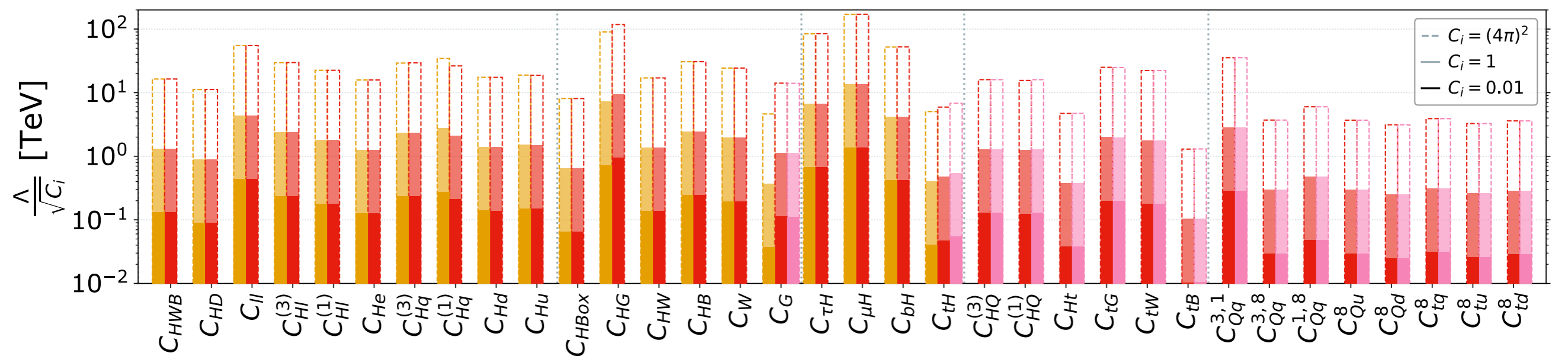
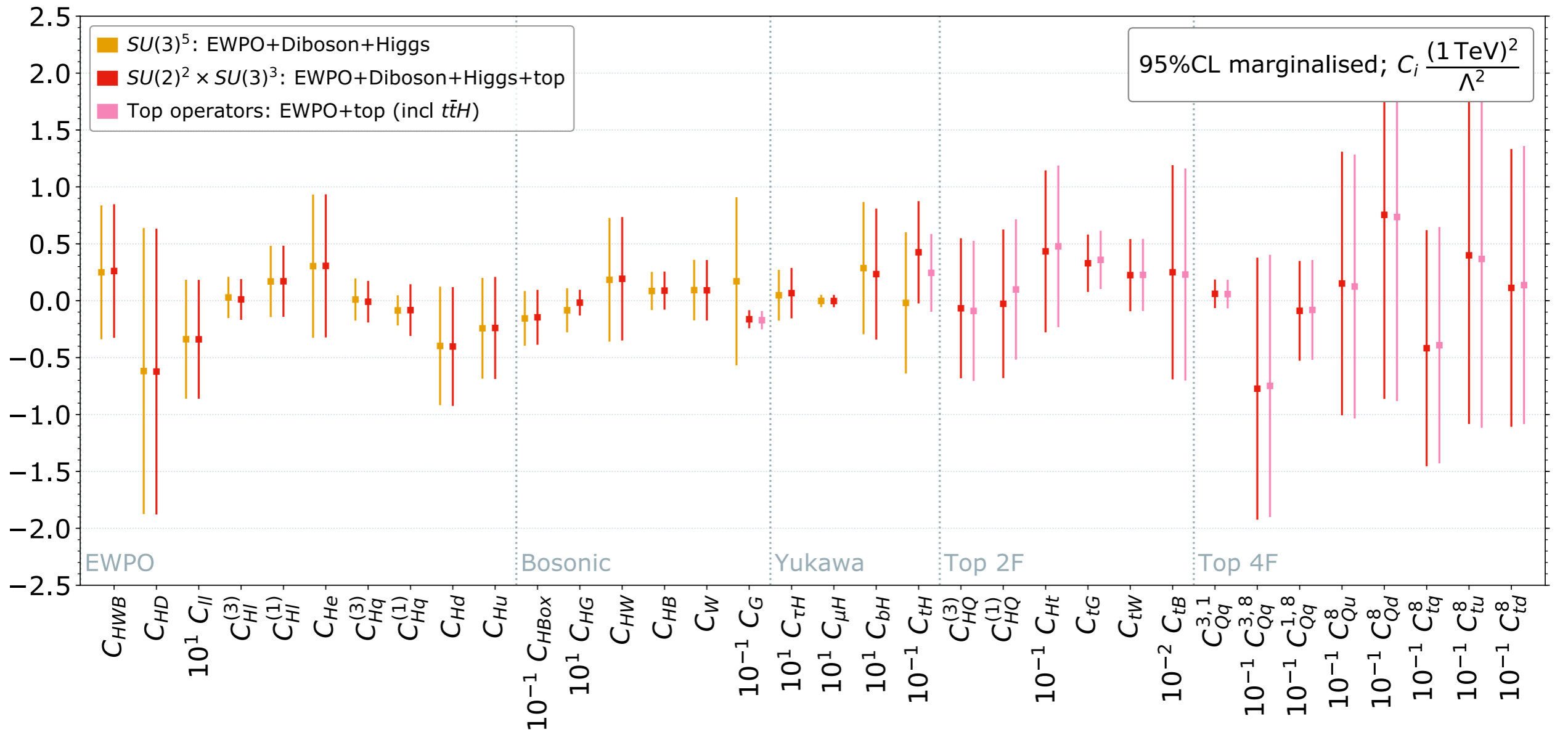
# Fit to Higgs, diboson, electroweak data in the flavour universal scenario:



# STXS measurements for ggF

STXS measurements of gluon gluon fusion **improve sensitivity** and **break degeneracy** between SMEFT operators:



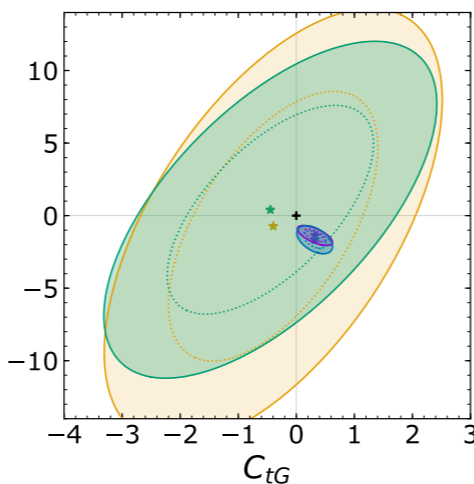
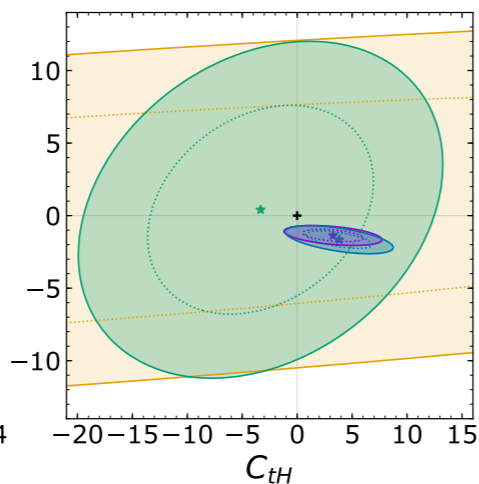
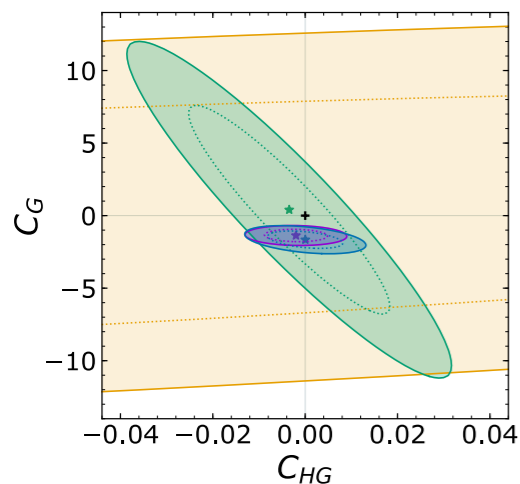
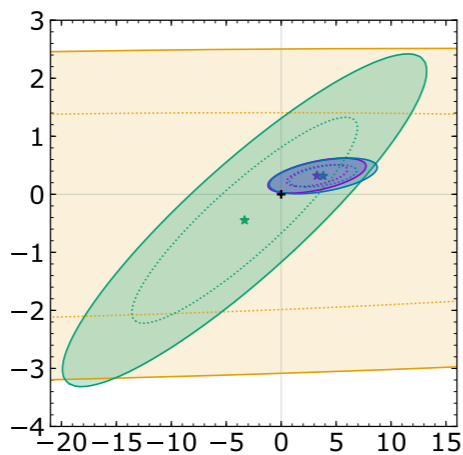
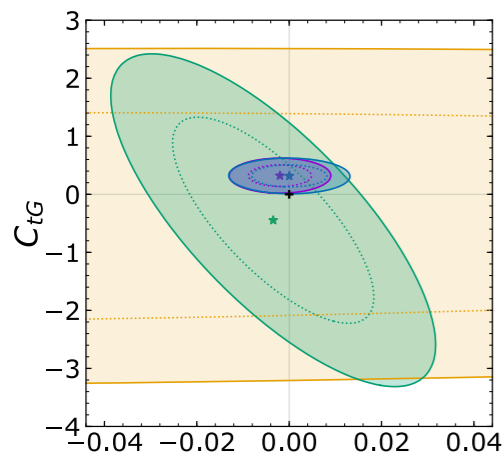
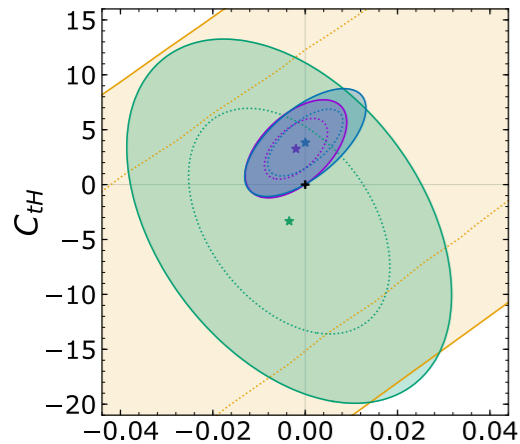


# Top-Higgs interplay

Studying the interplay of Higgs and top data  
in constraining the operators  $\mathcal{O}_{tH}$ ,  $\mathcal{O}_{tG}$ ,  $\mathcal{O}_G$ ,  $\mathcal{O}_{HG}$

while marginalising over

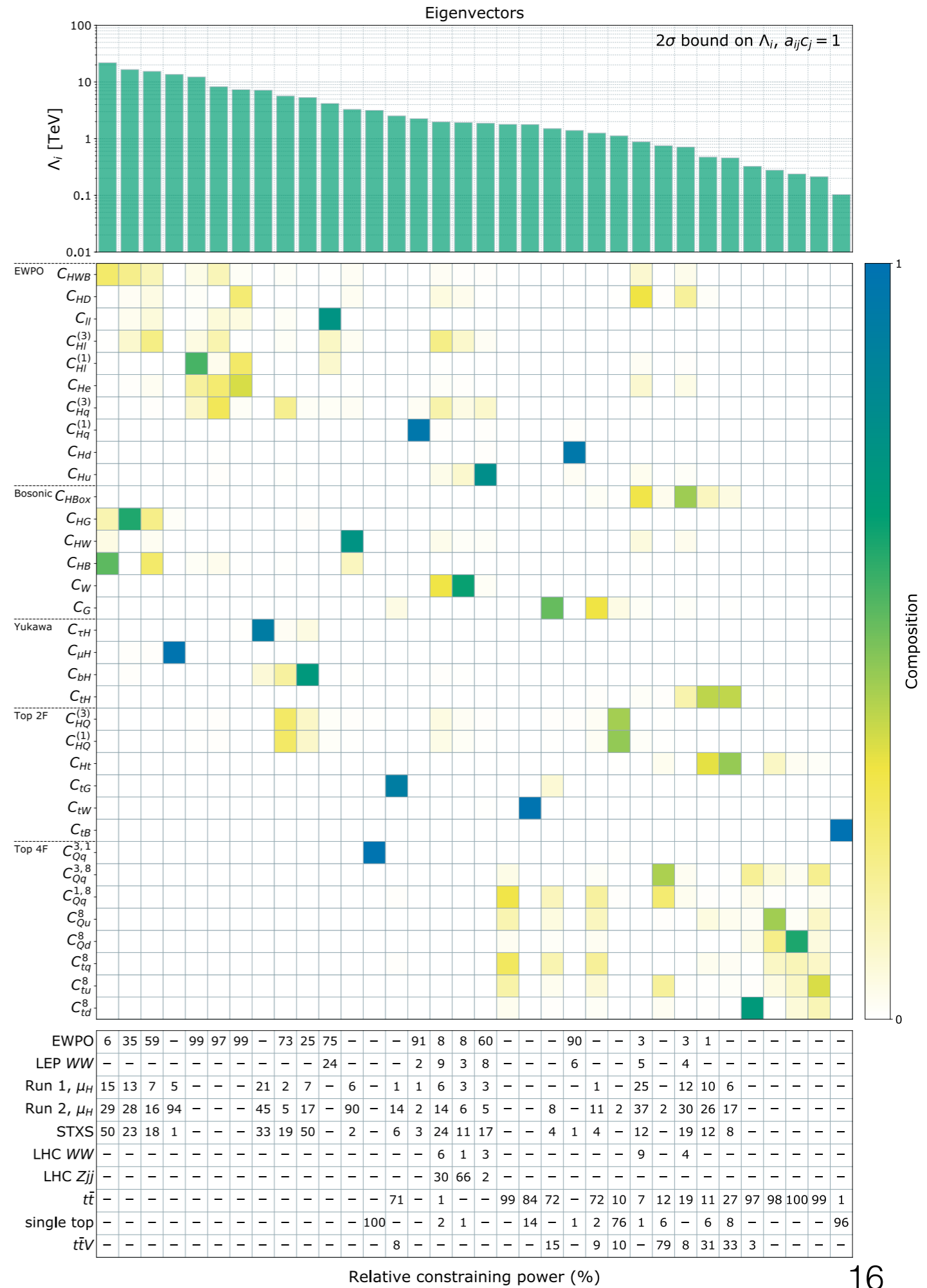
$\mathcal{O}_{H\Box}$ ,  $\mathcal{O}_{HW}$ ,  $\mathcal{O}_{HB}$ ,  $\mathcal{O}_{bH}$ ,  $\mathcal{O}_{\tau H}$ ,  $\mathcal{O}_{\mu H}$  (+4F operators)



*Marginalised 95% C. L.*

- Higgs data (no  $t\bar{t}H$ )
- Higgs data
- Higgs & Top data
- Higgs & Top data (+4F)
- + SM

# Principal Component Analysis



# UV models

We analyse our fit in terms of a set of BSM benchmark models  
from 2009.01249 Marzocca et. al, 1711.10391 de Blas et. al

Name	Spin	SU(3)	SU(2)	U(1)	Name	Spin	SU(3)	SU(2)	U(1)
$S$	0	1	1	0	$\Delta_1$	$\frac{1}{2}$	1	2	$-\frac{1}{2}$
$S_1$	0	1	1	1	$\Delta_3$	$\frac{1}{2}$	1	2	$-\frac{1}{2}$
$\varphi$	0	1	2	$\frac{1}{2}$	$\Sigma$	$\frac{1}{2}$	1	3	0
$\Xi$	0	1	3	0	$\Sigma_1$	$\frac{1}{2}$	1	3	-1
$\Xi_1$	0	1	3	1	$U$	$\frac{1}{2}$	3	1	$\frac{2}{3}$
$B$	1	1	1	0	$D$	$\frac{1}{2}$	3	1	$-\frac{1}{3}$
$B_1$	1	1	1	1	$Q_1$	$\frac{1}{2}$	3	2	$\frac{1}{6}$
$W$	1	1	3	0	$Q_5$	$\frac{1}{2}$	3	2	$-\frac{5}{6}$
$W_1$	1	1	3	1	$Q_7$	$\frac{1}{2}$	3	2	$\frac{7}{6}$
$N$	$\frac{1}{2}$	1	1	0	$T_1$	$\frac{1}{2}$	3	3	$-\frac{1}{3}$
$E$	$\frac{1}{2}$	1	1	-1	$T_2$	$\frac{1}{2}$	3	3	$\frac{2}{3}$
$T$	$\frac{1}{2}$	3	1	$\frac{2}{3}$	$TB$	$\frac{1}{2}$	3	2	$\frac{1}{6}$



# UV models: patterns

We analyse our fit in terms of a set of BSM benchmark models from 2009.01249 Marzocca et. al, 1711.10391 de Blas et. al

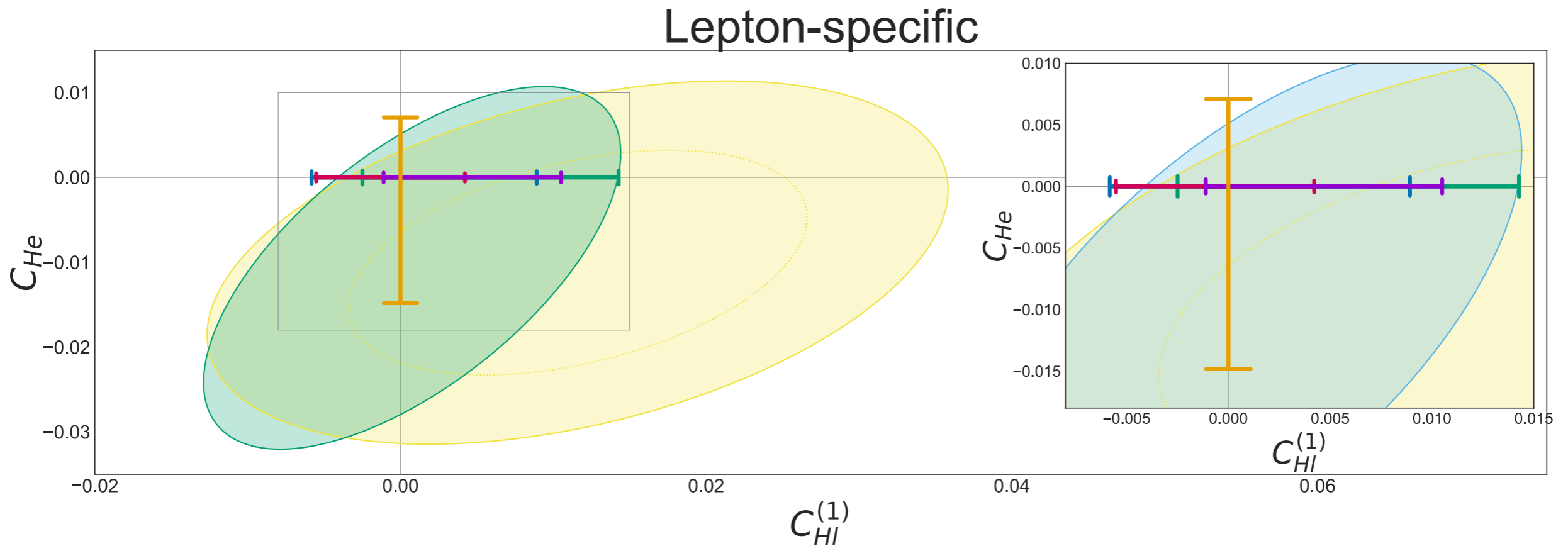
Some models exhibit similar patterns among operators

- Consider models with couplings to leptons:  $N, E, \Delta_1, \Delta_3, \Sigma, \Sigma_1$
- These will generate

$$\mathcal{O}_{Hl}^{(1)}, \mathcal{O}_{Hl}^{(3)}, \mathcal{O}_{He}, \mathcal{O}_{ll}$$

with patterns such as  $C_{Hl}^{(1)} \propto C_{Hl}^{(3)}$  and  $C_{ll} \propto C_{He}$

# UV models: patterns



- |  |  |
|--|--|
| <span style="display: inline-block; width: 15px; height: 15px; background-color: yellow; border: 1px solid black; margin-right: 5px;"></span> $C_{II}, C_{HI}^{(3)}$ marginalised                                      | <span style="display: inline-block; width: 15px; height: 15px; background-color: orange; border: 1px solid black; margin-right: 5px;"></span> $\Delta_1, \Delta_3: C_{HI}^{(1)} = C_{HI}^{(3)} = C_{II} = 0$ |
| <span style="display: inline-block; width: 15px; height: 15px; background-color: lightgreen; border: 1px solid black; margin-right: 5px;"></span> $C_{II} = 0, C_{HI}^{(3)}$ marginalised                              | <span style="display: inline-block; width: 15px; height: 15px; background-color: purple; border: 1px solid black; margin-right: 5px;"></span> $N: C_{HI}^{(3)} = -C_{HI}^{(1)}; C_{II} = C_{He} = 0$         |
| <span style="display: inline-block; width: 15px; height: 15px; background-color: lightblue; border: 1px solid black; margin-right: 5px;"></span> $\Sigma: C_{HI}^{(3)} = \frac{1}{3}C_{HI}^{(1)}; C_{II} = C_{He} = 0$ | <span style="display: inline-block; width: 15px; height: 15px; background-color: red; border: 1px solid black; margin-right: 5px;"></span> $E: C_{HI}^{(3)} = C_{HI}^{(1)}; C_{II} = C_{He} = 0$             |
| <span style="display: inline-block; width: 15px; height: 15px; background-color: teal; border: 1px solid black; margin-right: 5px;"></span> $\Sigma_1: C_{HI}^{(3)} = -\frac{1}{3}C_{HI}^{(1)}; C_{II} = C_{He} = 0$   |  |

See 2012.02779 for quark-specific, top-specific and boson-specific cases

# Conclusions

- ▶ Global fit produced using **Fitmaker**:  
a publicly available python code  
<https://gitlab.com/kenmimasu/fitrepo>  
(Version for public use still to come!)
- ▶ an adaptable, flexible and extensible framework for performing global SMEFT fits.

**Thank you for listening!**

# Backup

# Datasets: Higgs

<b>LHC Run 1 Higgs</b>	$n_{\text{obs}}$	<b>Ref.</b>
ATLAS and CMS LHC Run 1 combination of Higgs signal strengths. Production: $ggF, VBF, ZH, WH$ & $ttH$ Decay: $\gamma\gamma, ZZ, W^+W^-, \tau^+\tau^-$ & $b\bar{b}$	21	[8]
ATLAS inclusive $Z\gamma$ signal strength measurement	1	[9]
<b>LHC Run 2 Higgs (new)</b>	$n_{\text{obs}}$	<b>Ref.</b>
ATLAS combination of signal strengths and stage 1.0 STXS in $H \rightarrow 4\ell$ including ratios of branching fractions to $\gamma\gamma, WW^*, \tau^+\tau^-$ & $b\bar{b}$ Signal strengths coarse STXS bins  fine STXS bins	16 19 25	[10]
CMS LHC combination of Higgs signal strengths. Production: $ggF, VBF, ZH, WH$ & $ttH$ Decay: $\gamma\gamma, ZZ, W^+W^-, \tau^+\tau^-, b\bar{b}$ & $\mu^+\mu^-$	23	[11]
CMS stage 1.0 STXS measurements for $H \rightarrow \gamma\gamma$ . 13 parameter fit   7 parameter fit	13 7	[12]
CMS stage 1.0 STXS measurements for $H \rightarrow \tau^+\tau^-$	9	[13]
CMS stage 1.1 STXS measurements for $H \rightarrow 4\ell$	19	[14]
CMS differential cross section measurements of inclusive Higgs production in the $WW^* \rightarrow l\nu l\nu$ final state. $\frac{d\sigma}{dn_{\text{jet}}}$   $\frac{d\sigma}{dp_H^T}$	5 6	[15]
ATLAS $H \rightarrow Z\gamma$ signal strength.	1	[16]
ATLAS $H \rightarrow \mu^+\mu^-$ signal strength.	1	[17]

<b>EW precision observables</b>	$n_{\text{obs}}$	<b>Ref.</b>
Precision electroweak measurements on the $Z$ resonance. $\Gamma_Z, \sigma_{\text{had.}}^0, R_\ell^0, A_{FB}^\ell, A_\ell(\text{SLD}), A_\ell(\text{Pt}), R_b^0, R_c^0, A_{FB}^b, A_{FB}^c, A_b \& A_c$	12	[1]
Combination of CDF and D0 $W$ -Boson Mass Measurements	1	[6]
LHC run 1 $W$ boson mass measurement by ATLAS	1	[57]
<b>Diboson LEP &amp; LHC</b>	$n_{\text{obs}}$	<b>Ref.</b>
$W^+ W^-$ angular distribution measurements at LEP II.	8	[5]
$W^+ W^-$ total cross section measurements at L3 in the $l\nu l\nu, l\nu qq \& qqqq$ final states for 8 energies	24	[3]
$W^+ W^-$ total cross section measurements at OPAL in the $l\nu l\nu, l\nu qq \& qqqq$ final states for 7 energies	21	[4]
$W^+ W^-$ total cross section measurements at ALEPH in the $l\nu l\nu, l\nu qq \& qqqq$ final states for 8 energies	21	[2]
ATLAS $W^+ W^-$ differential cross section in the $e\nu\mu\nu$ channel, $\frac{d\sigma}{dp_{\ell_1}^T}$ , $p_T > 120$ GeV overflow bin	1	[66]
ATLAS $W^+ W^-$ fiducial differential cross section in the $e\nu\mu\nu$ channel, $\frac{d\sigma}{dp_{\ell_1}^T}$	14	[70]
ATLAS $W^\pm Z$ fiducial differential cross section in the $l^+ l^- l^\pm \nu$ channel, $\frac{d\sigma}{dp_Z^T}$	7	[69]
CMS $W^\pm Z$ normalised fiducial differential cross section in the $l^+ l^- l^\pm \nu$ channel, $\frac{1}{\sigma} \frac{d\sigma}{dp_Z^T}$	11	[67]
ATLAS $Zjj$ fiducial differential cross section in the $l^+ l^-$ channel, $\frac{d\sigma}{d\Delta\varphi_{jj}}$	12	[71]

# Datasets: top

Tevatron & Run 1 top	$n_{\text{obs}}$	Ref.
Tevatron combination of differential $t\bar{t}$ forward-backward asymmetry, $A_{FB}(m_{t\bar{t}})$ .	4	[7]
ATLAS $t\bar{t}$ differential distributions in the dilepton channel. $\frac{d\sigma}{dm_{t\bar{t}}}$	6	[18]
ATLAS $t\bar{t}$ differential distributions in the $\ell$ +jets channel. $\frac{d\sigma}{dm_{t\bar{t}}} \mid \frac{d\sigma}{d y_{t\bar{t}} } \mid \frac{d\sigma}{dp_t^T} \mid \frac{d\sigma}{d y_t }$ .	7 5 8 5	[19]
CMS $t\bar{t}$ differential distributions in the $\ell$ +jets channel. $\frac{d\sigma}{dm_{t\bar{t}}} \mid \frac{d\sigma}{dy_{t\bar{t}}} \mid \frac{d\sigma}{dp_t^T} \mid \frac{d\sigma}{dy_t}$ .	7 10 8  10	[20, 215]
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