Using associated top quark production to probe for new physics within the framework of effective field theory

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Motivation for new physics

The Standard Model (SM) is wonderfully precise, but only accounts for 5% of the universe

Shortcomings include:

- **Dark matter/energy** – Invisible particles? Non-zero vacuum energy?
- **Hierarchy problem** – Why is the Higgs mass $\mathcal{O}(10^2)$ GeV but Plank mass $\mathcal{O}(10^{19})$ GeV?
- **Baryon asymmetry** – Why do we live in a universe devoid of anti-matter?

After discovering the Higgs boson in 2012, the LHC provided no definitive evidence of anything unexpected

Assume $\Lambda_{\mathrm{NP}} > \Lambda_{\mathrm{LHC}}$

How might it appear at the LHC?

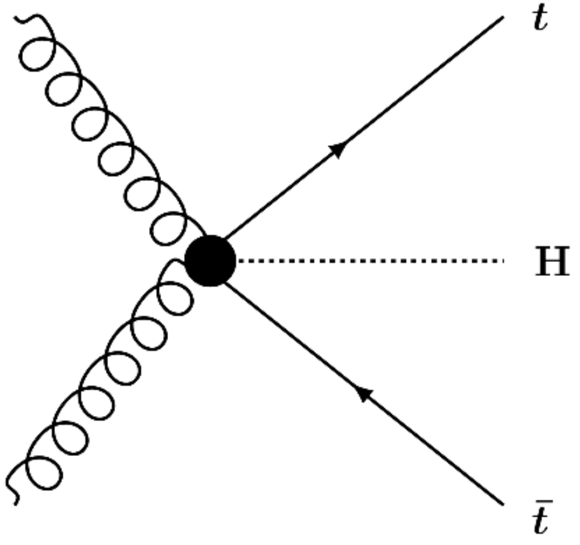
THE OHIO STATE UNIVERSITY

Introduction to effective field theory

New physics at scales beyond what the LHC can directly probe can be approximated by expanding terms of higher dimensional (d) operators $\mathcal{O}$ consisting of SM fields

Operators are suppressed by powers of the energy scale Λ, and the strength is controlled by the Wilson coefficients (WCs) c_i

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_{d,i} \frac{c_i^{(d)}}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}$$



Dimension five violates lepton number

Dimension six is the focus of this analysis

Higher dimensions are suppressed by additional powers of Λ

Analysis overview

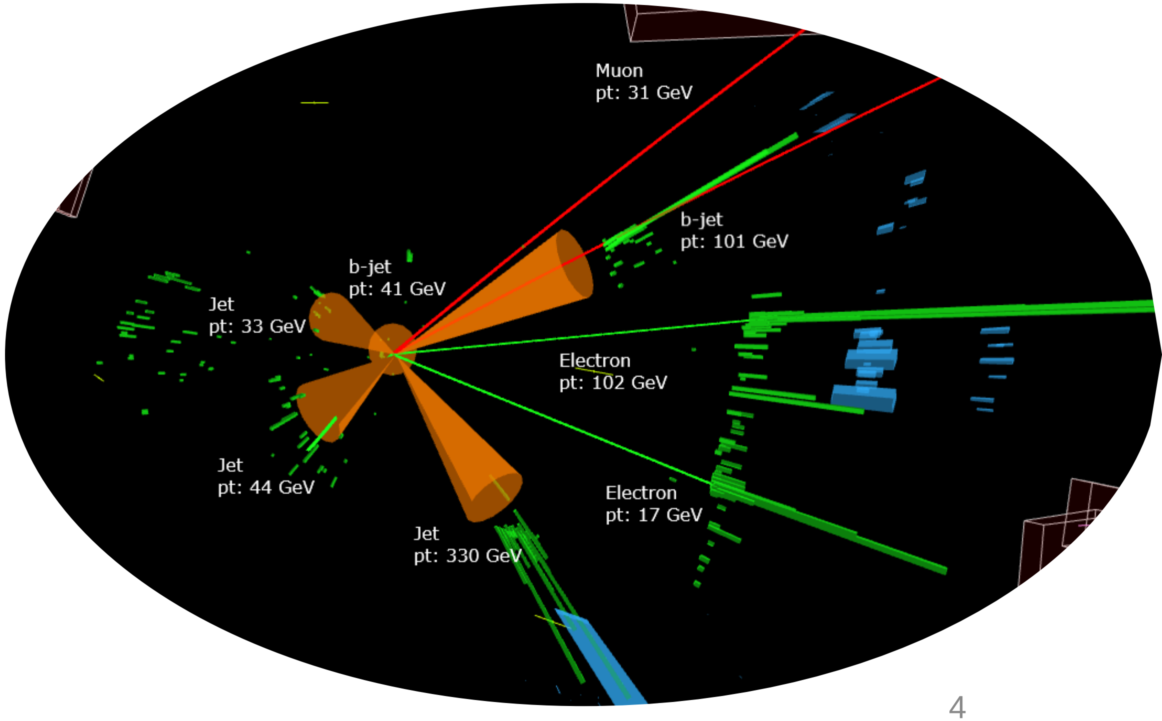
A novel technique to examine data collected in 2017
Performs a global fit across all processes (signal and background)

Probe EFT effects using multiple lepton final states

Procedure helps constrain systematic uncertainties

Correlations rely on data (no assumptions made)

Using channels with $t\bar{t}l\nu$, $t\bar{t}l\bar{l}$, $tl\bar{l}q$, $t\bar{t}H$, and tHq production ($H \rightarrow W^+W^-, ZZ, \tau^+\tau^-$, exclude $H \rightarrow b\bar{b}$)



EFT parameterization

Matrix elements $\mathcal{M}$ split into SM and EFT terms

$$\mathcal{M} = \mathcal{M}_{\mathrm{SM}} + \sum_j \frac{c_j}{\Lambda^2}\mathcal{M}_j$$

Parameterize cross section by WCs:

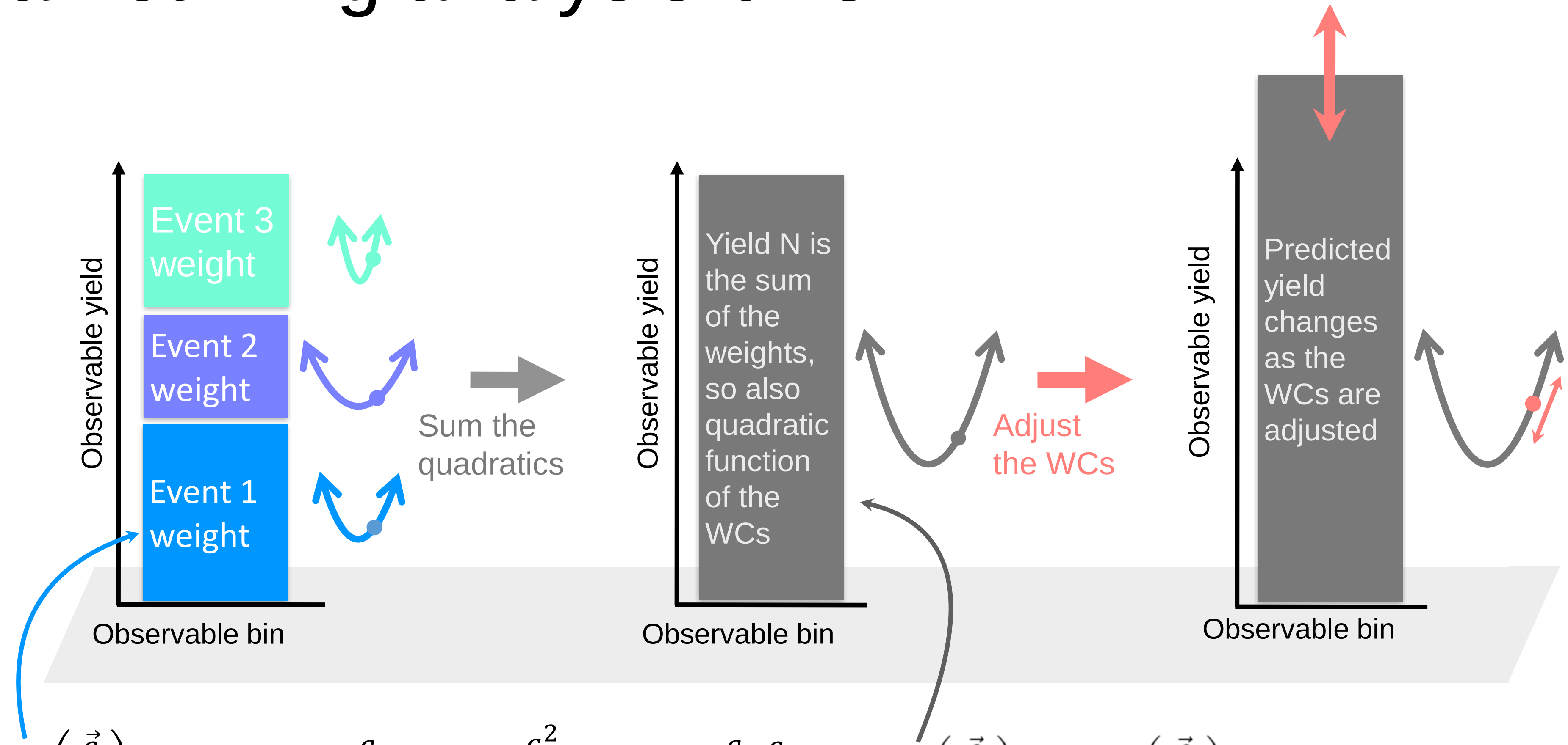
- SM terms (s_{0i})
- Interference terms between the SM and EFT (s_{1ij})
- Pure EFT terms (s_{2ij})
- Interference terms between EFT (s_{3ijk})

$$\sigma \propto \mathcal{M}^2$$

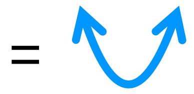
Individual events ($d\sigma$) have weight w_i → can be summed to produce the predicted event yields

$$w_i\left(\frac{\vec{c}}{\Lambda^2}\right) = s_{0i} + \sum_j s_{1ij}\frac{c_j}{\Lambda^2} + \sum_j s_{2ij}\frac{c_j^2}{\Lambda^4} + \sum_{j,k} s_{3ijk}\frac{c_j}{\Lambda^2}\frac{c_k}{\Lambda^2}$$

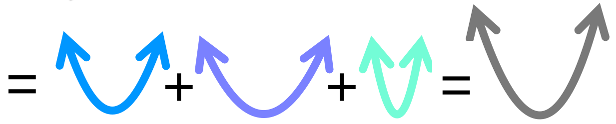
Parametrizing analysis bins



$$w_i\left(\frac{\vec{c}}{\Lambda^2}\right) = s_{0i} + \sum_j s_{1ij}\frac{c_j}{\Lambda^2} + \sum_j s_{2ij}\frac{c_j^2}{\Lambda^4} + \sum_{j,k} s_{3ijk}\frac{c_j}{\Lambda^2}\frac{c_k}{\Lambda^2}$$



$$N\left(\frac{\vec{c}}{\Lambda^2}\right) = \sum_i w_i\left(\frac{\vec{c}}{\Lambda^2}\right)$$



EFT parameterization

MC simulations are generated with non-zero WCs

Extra partons are added when possible

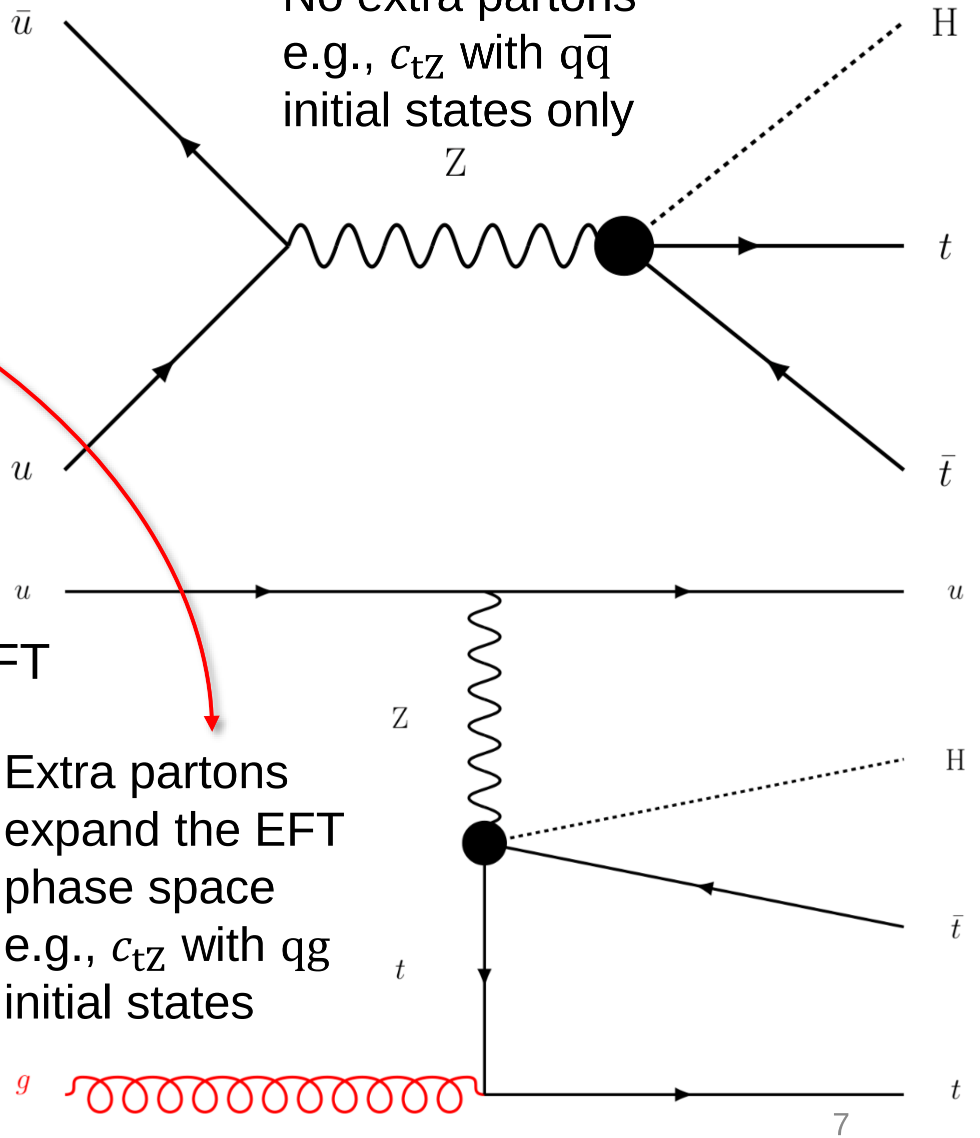
Initial values chosen to include all relevant phase space

and optimize the MC statistical power: $\sigma_{\mathrm{stat}}^2 = \sum w_i^2(\vec{c})$

Weight of each event accounts for variation in yield from EFT

Used to solve for the constant terms in the quadratic fit

This parameterization will be used in the fit



No extra partons
e.g., c_{tZ} with $\mathrm{q\bar{q}}$
initial states only

Extra partons
expand the EFT
phase space
e.g., c_{tZ} with qg
initial states

Dim6TopEFT Model

EFT simulations are generated by MadGraph_aMC@NLO using the dim6TopEFT[1] model

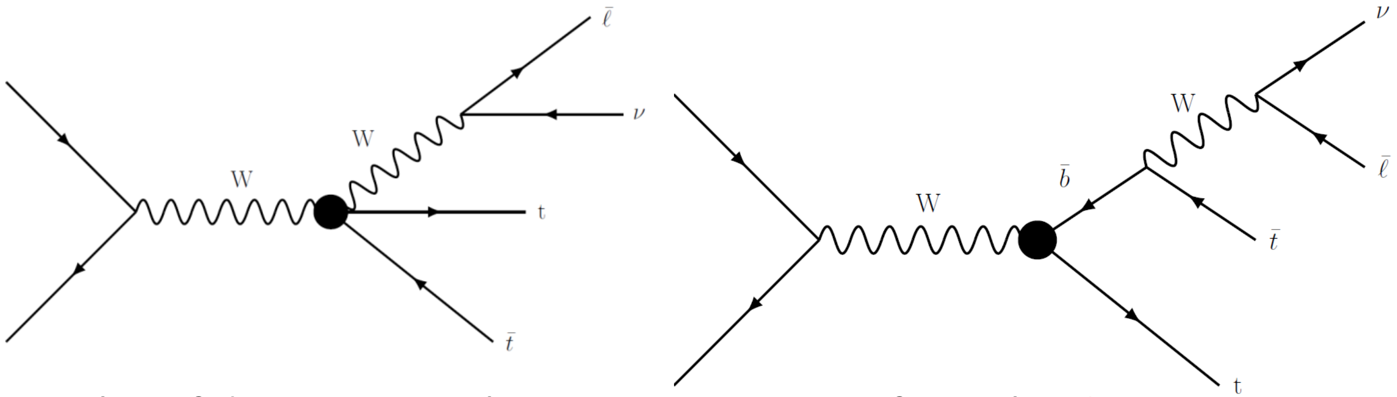
- Warsaw basis of dimension six operators
- $\Lambda = 1\ \mathrm{TeV}$
- CKM matrix is assumed to be a unit matrix
- $\mathrm{u, d, s, c, e, \mu}$ masses all set to zero
- The unitary gauge is used and Goldstone bosons are removed
- Baryon and lepton number violating operators are not included
- Only tree-level simulation is possible
- Lepton universality is assumed (all flavors set to same WCs)

The 16 operators which have the largest impact on the signal processes, and relatively small impact on $\mathrm{t\bar{t}}$ background, are considered

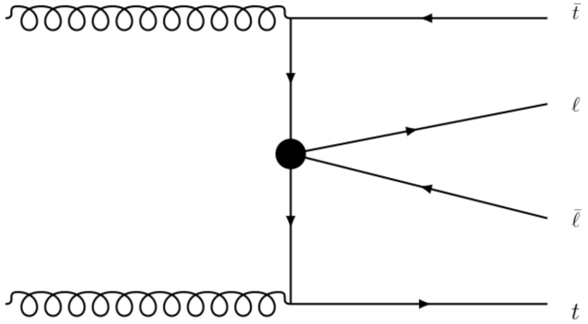
[1] https://arxiv.org/abs/1802.07237

Model operators

Only the real components are considered since the imaginary coefficients lead to CP violation, and are well constrained by EDM experiments and $B \to X_s \gamma$ decays



Examples of diagrams involving vertices arising from the O_{tW} operator



Examples of diagrams involving vertices arising from the $O_{\mathrm{t}\ell}$ operator

Operators involving two quarks and one or more bosons		
Operator	Definition	Wilson coefficient
$\ddagger O_{\mathrm{u}\varphi}^{(ij)}$	$\bar{\mathrm{q}}_i \mathrm{u}_j \tilde{\varphi}\,(\varphi^\dagger \varphi)$	$c_{\mathrm{t}\varphi} + i c_{\mathrm{t}\varphi}^I$
$O_{\varphi\mathrm{q}}^{1(ij)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{\mathrm{q}}_i \gamma^\mu \mathrm{q}_j)$	$c_{\varphi Q}^- + c_{\varphi Q}^3$
$O_{\varphi\mathrm{q}}^{3(ij)}$	$(\varphi^\dagger i \overleftrightarrow{D}{}^I_\mu \varphi)(\bar{\mathrm{q}}_i \gamma^\mu \tau^I \mathrm{q}_j)$	$c_{\varphi Q}^3$
$O_{\varphi\mathrm{u}}^{(ij)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{\mathrm{u}}_i \gamma^\mu \mathrm{u}_j)$	$c_{\varphi\mathrm{t}}$
$\ddagger O_{\varphi\mathrm{ud}}^{(ij)}$	$(\tilde{\varphi}^\dagger i D_\mu \varphi)(\bar{\mathrm{u}}_i \gamma^\mu \mathrm{d}_j)$	$c_{\varphi\mathrm{tb}} + i c_{\varphi\mathrm{tb}}^I$
$\ddagger O_{\mathrm{uW}}^{(ij)}$	$(\bar{\mathrm{q}}_i \sigma^{\mu\nu} \tau^I \mathrm{u}_j)\; \tilde{\varphi} \mathrm{W}_{\mu\nu}^I$	$c_{\mathrm{tW}} + i c_{\mathrm{tW}}^I$
$\ddagger O_{\mathrm{dW}}^{(ij)}$	$(\bar{\mathrm{q}}_i \sigma^{\mu\nu} \tau^I \mathrm{d}_j)\; \varphi \mathrm{W}_{\mu\nu}^I$	$c_{\mathrm{bW}} + i c_{\mathrm{bW}}^I$
$\ddagger O_{\mathrm{uB}}^{(ij)}$	$(\bar{\mathrm{q}}_i \sigma^{\mu\nu} \mathrm{u}_j) \quad \tilde{\varphi} \mathrm{B}_{\mu\nu}$	$(c_{\mathrm{W}} c_{\mathrm{tW}} - c_{\mathrm{tZ}})/s_{\mathrm{W}} + i(c_{\mathrm{W}} c_{\mathrm{tW}}^I - c_{\mathrm{tZ}}^I)/s_{\mathrm{W}}$
$\ddagger O_{\mathrm{uG}}^{(ij)}$	$(\bar{\mathrm{q}}_i \sigma^{\mu\nu} T^A \mathrm{u}_j)\; \tilde{\varphi} G_{\mu\nu}^A$	$c_{\mathrm{tG}} + i c_{\mathrm{tG}}^I$

Operators involving two quarks and two leptons		
Operator	Definition	Wilson coefficient
$O_{\ell\mathrm{q}}^{1(ijkl)}$	$(\bar{\ell}_i \gamma^\mu \ell_j)(\bar{\mathrm{q}}_k \gamma^\mu \mathrm{q}_\ell)$	$c_{Q\ell}^{-(\ell)} + c_{Q\ell}^{3(\ell)}$
$O_{\ell\mathrm{q}}^{3(ijkl)}$	$(\bar{\ell}_i \gamma^\mu \tau^I \ell_j)(\bar{\mathrm{q}}_k \gamma^\mu \tau^I \mathrm{q}_\ell)$	$c_{Q\ell}^{3(\ell)}$
$O_{\ell\mathrm{u}}^{(ijkl)}$	$(\bar{\ell}_i \gamma^\mu \ell_j)(\bar{\mathrm{u}}_k \gamma^\mu \mathrm{u}_\ell)$	$c_{\mathrm{t}\ell}^{(\ell)}$
$O_{\mathrm{e}\bar{\mathrm{q}}}^{(ijkl)}$	$(\bar{\mathrm{e}}_i \gamma^\mu \mathrm{e}_j)(\bar{\mathrm{q}}_k \gamma^\mu \mathrm{q}_\ell)$	$c_{Q\mathrm{e}}^{(\ell)}$
$O_{\mathrm{eu}}^{(ijkl)}$	$(\bar{\mathrm{e}}_i \gamma^\mu \mathrm{e}_j)(\bar{\mathrm{u}}_k \gamma^\mu \mathrm{u}_\ell)$	$c_{\mathrm{te}}^{(\ell)}$
$\ddagger O_{\ell\mathrm{equ}}^{1(ijkl)}$	$(\bar{\ell}_i \mathrm{e}_j)\; \varepsilon\; (\bar{\mathrm{q}}_k \mathrm{u}_\ell)$	$c_{\mathrm{t}}^{S(\ell)} + i c_{\mathrm{t}}^{SI(\ell)}$
$\ddagger O_{\ell\mathrm{equ}}^{3(ijkl)}$	$(\bar{\ell}_i \sigma^{\mu\nu} \mathrm{e}_j)\; \varepsilon\; (\bar{\mathrm{q}}_k \sigma_{\mu\nu} \mathrm{u}_\ell)$	$c_{\mathrm{t}}^{T(\ell)} + i c_{\mathrm{t}}^{TI(\ell)}$

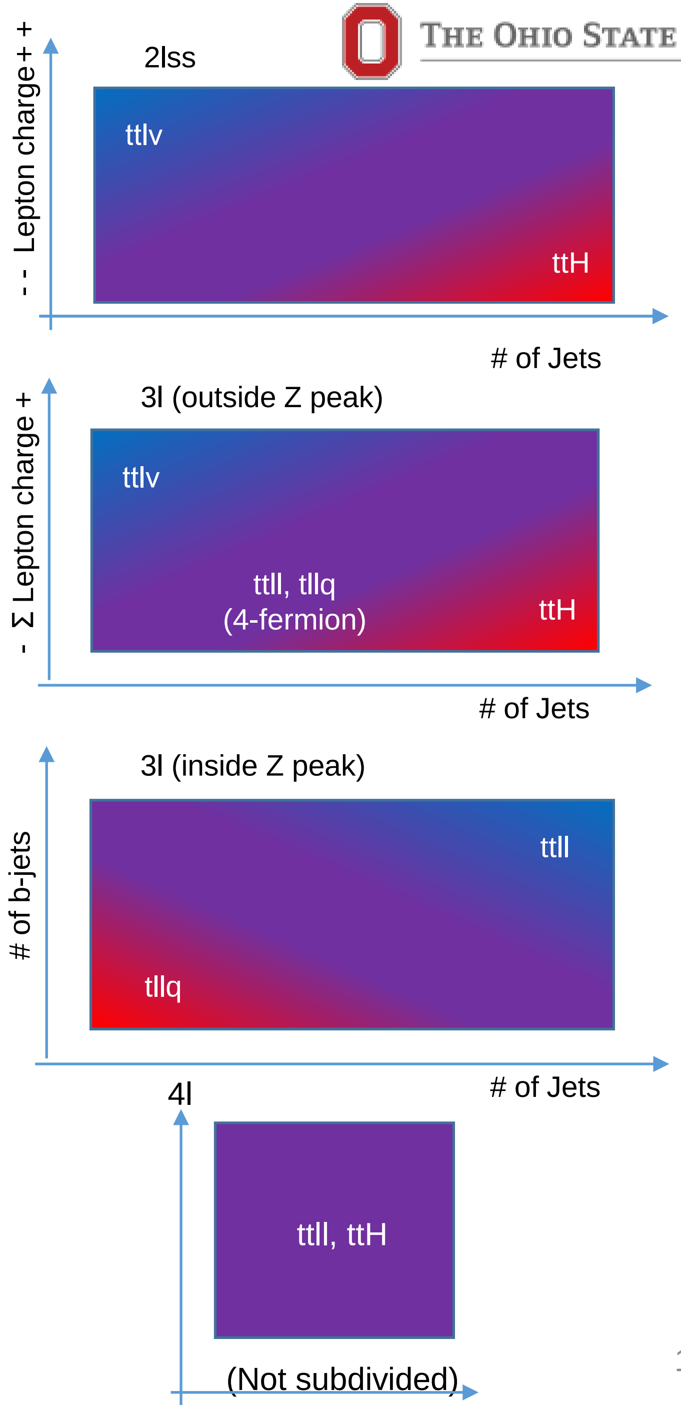
Event selection

The analysis is split into lepton (ℓ) categories as well as jet multiplicity (both light and b-tagged jets)

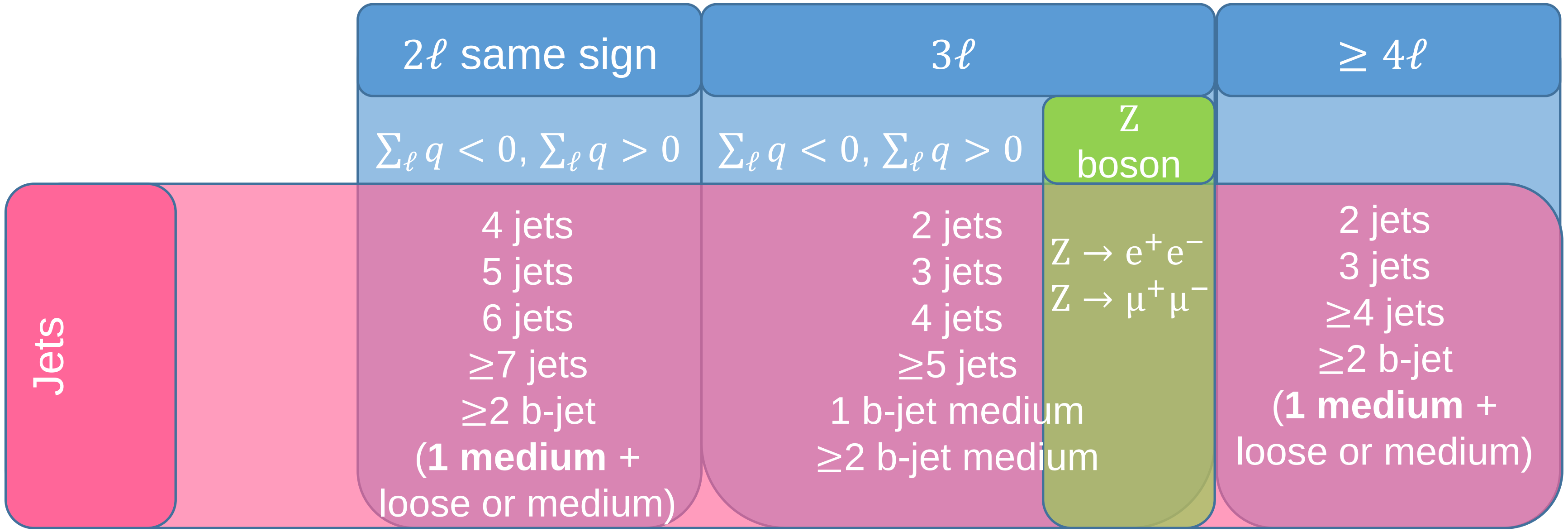
A BDT is applied to separate prompt from non-prompt leptons

Final-state observables are an admixture of processes (the method does not require we separate processes)

- Each analysis bin stores the sum of the quadratic coefficients → event yields are fully parametrized by the WCs

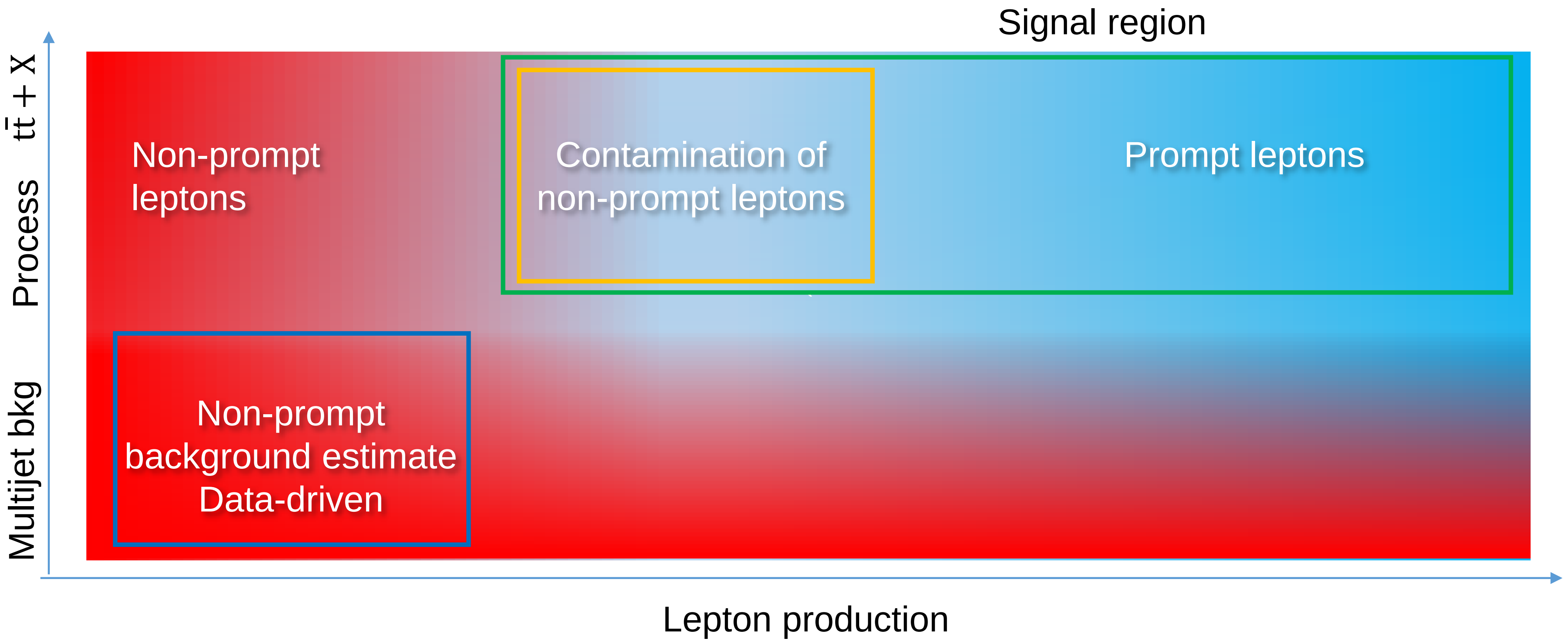


Event categorization



loose = 85% efficiency, 10% light/quark jets
medium = 70% efficiency, 1% light/quark jets

Misidentified lepton background



Probability of a non-prompt lepton passing prompt cuts is measured in a multijet enriched region

Data-driven

Fitting procedure

Each bin is treated as a Poisson experiment with a probability of obtaining the observed data

Profiled likelihood simultaneously fits all bin and extract the 2σ confidence intervals of the various WCs

Two fitting procedures are used:

Scan single WC, other 15 are unconstrained nuisance parameters

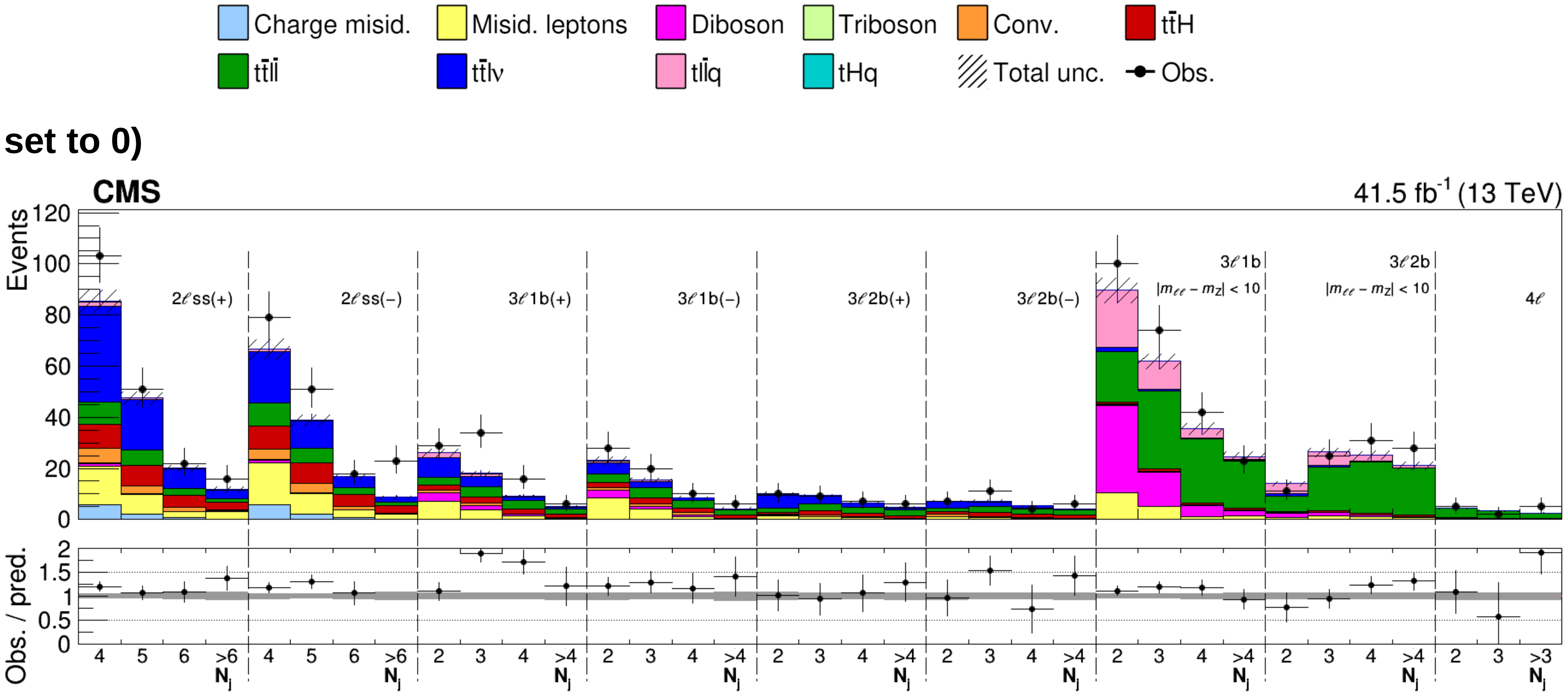
- More physical of the two, no reason for new physics to only favor one WC

Scan a single WC, other 15 are fixed to their SM value of zero

- Extreme scenario where nature has a single WC; the ability to fit one is limited by the lack of knowledge of 15 others

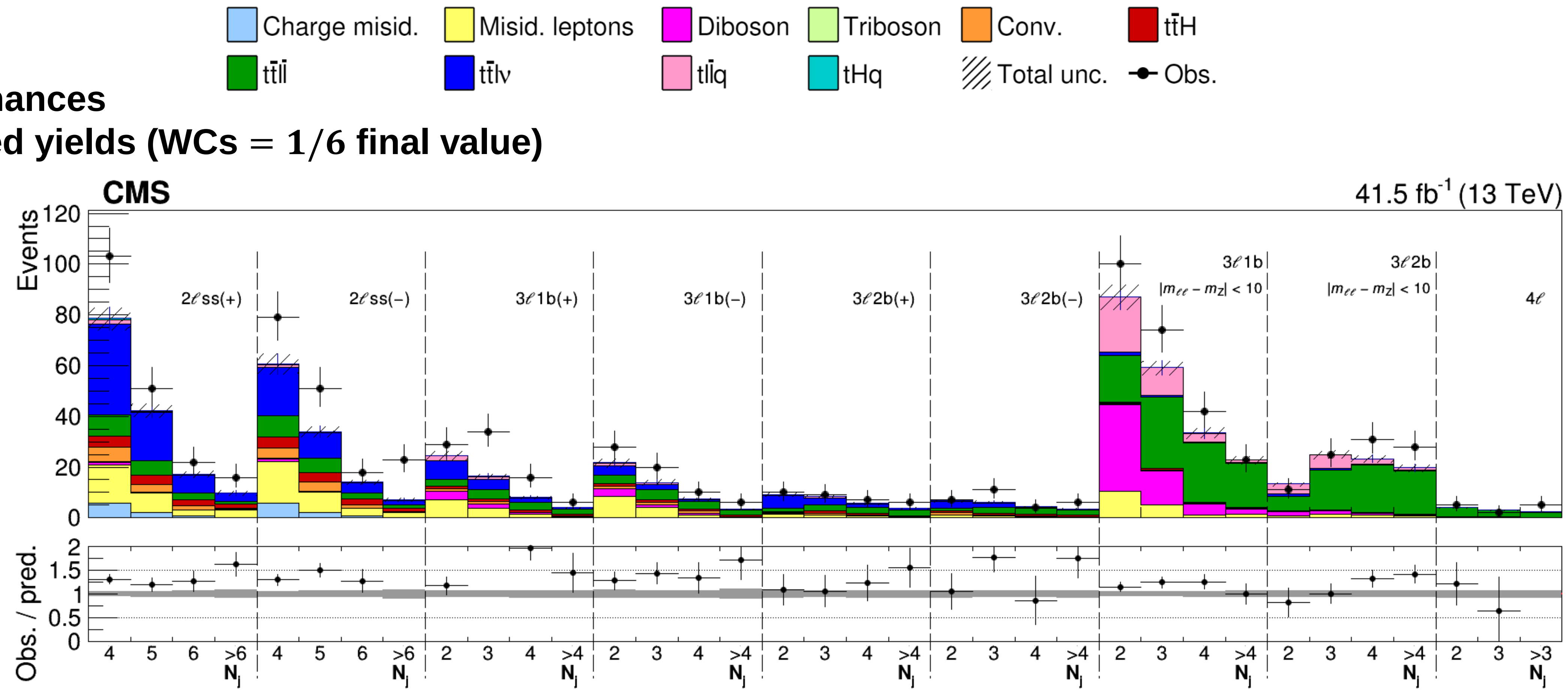
Event yields in analysis bins

SM only
(all WCs set to 0)



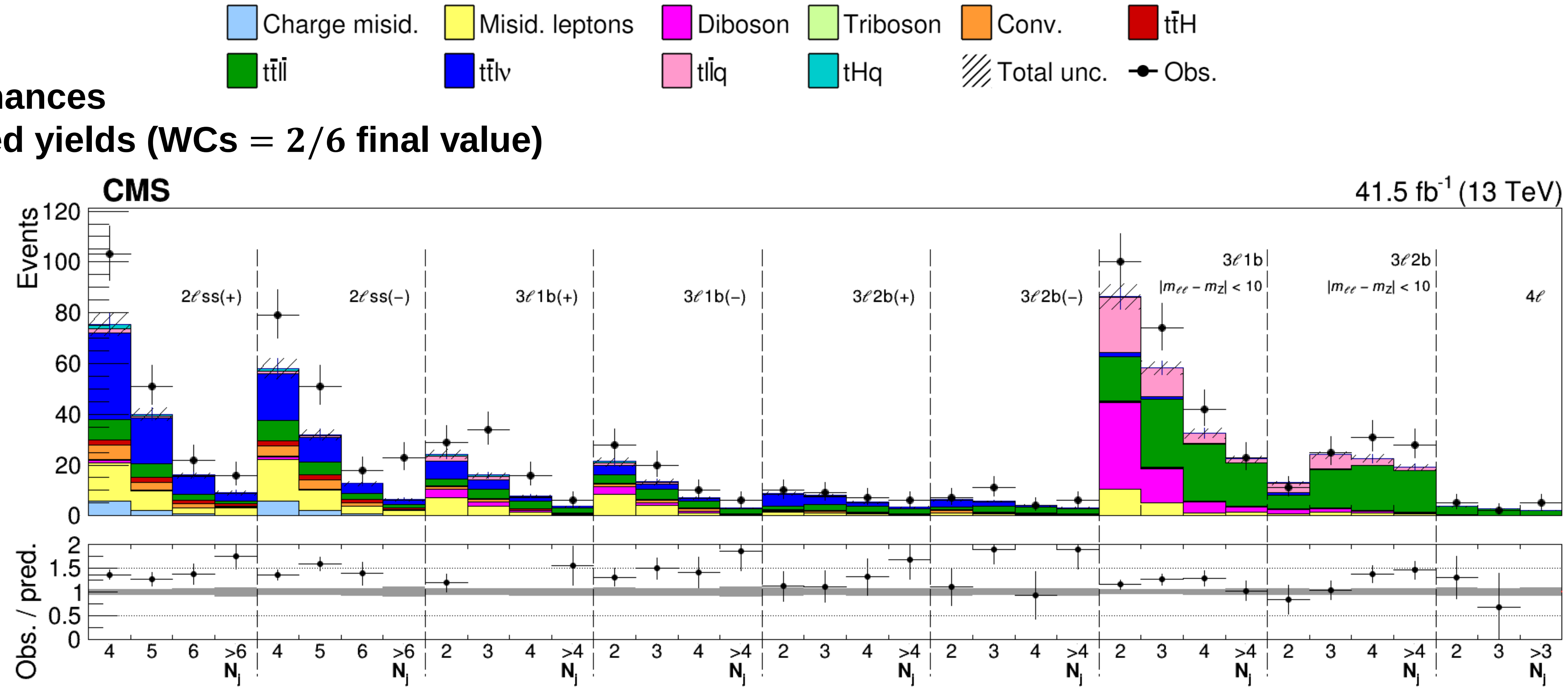
Event yields in analysis bins

EFT enhances predicted yields (WCs $= 1/6$ final value)



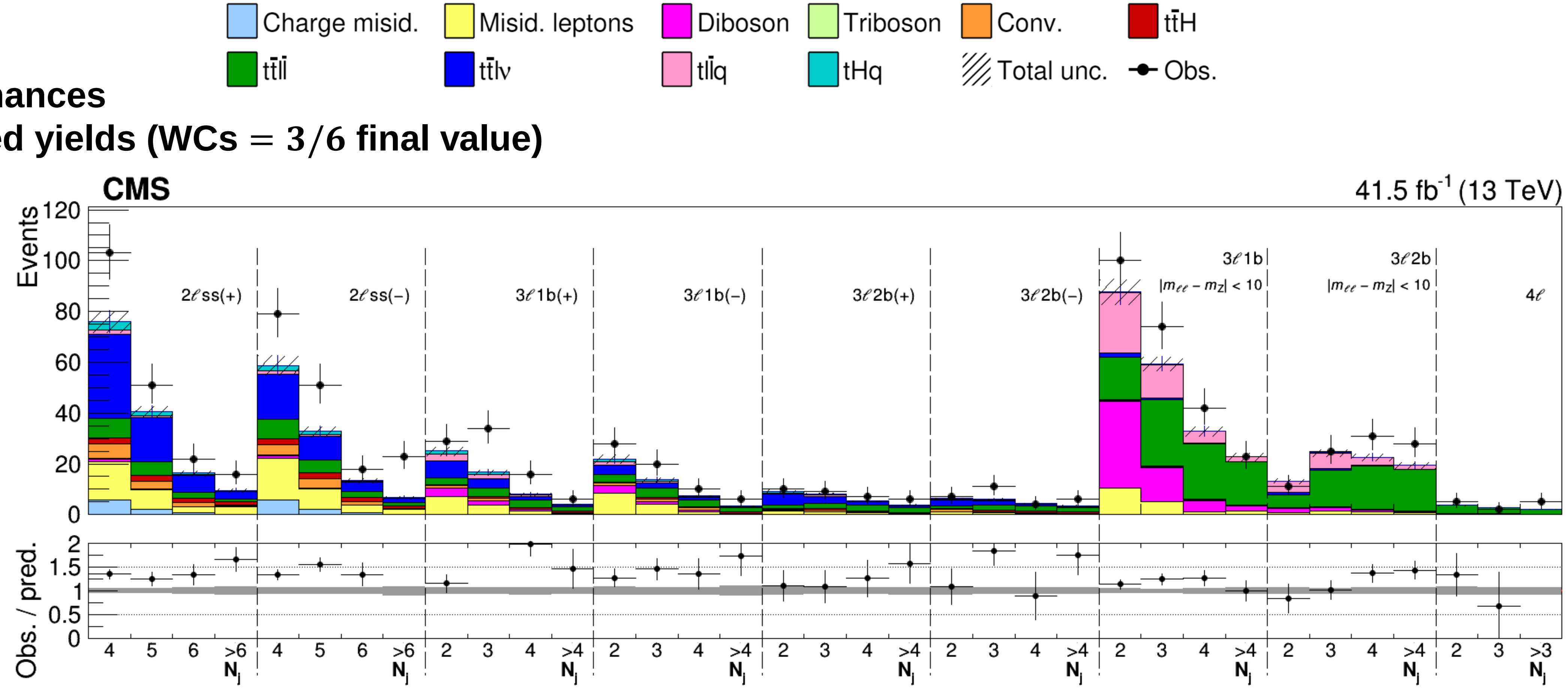
Event yields in analysis bins

EFT enhances predicted yields (WCs $= 2/6$ final value)



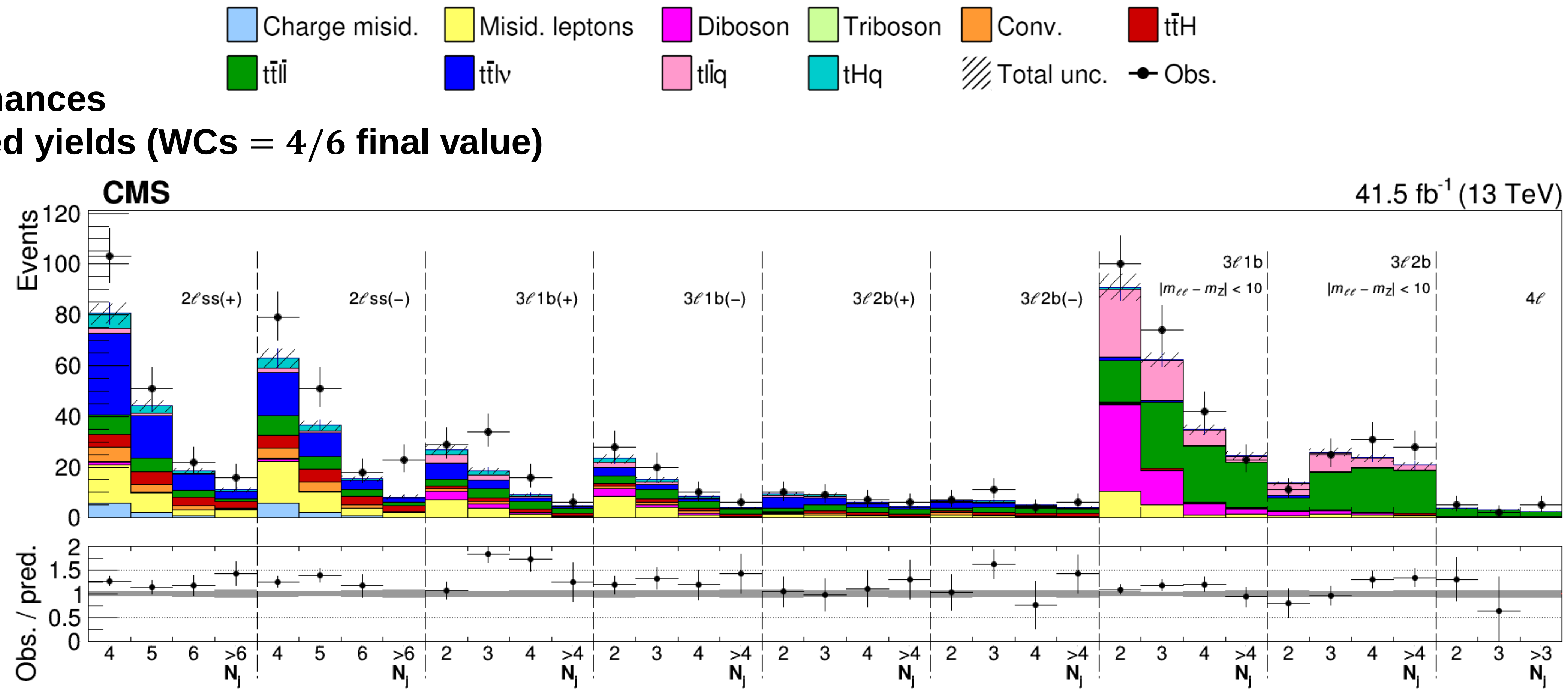
Event yields in analysis bins

EFT enhances
predicted yields (WCs $= 3/6$ final value)



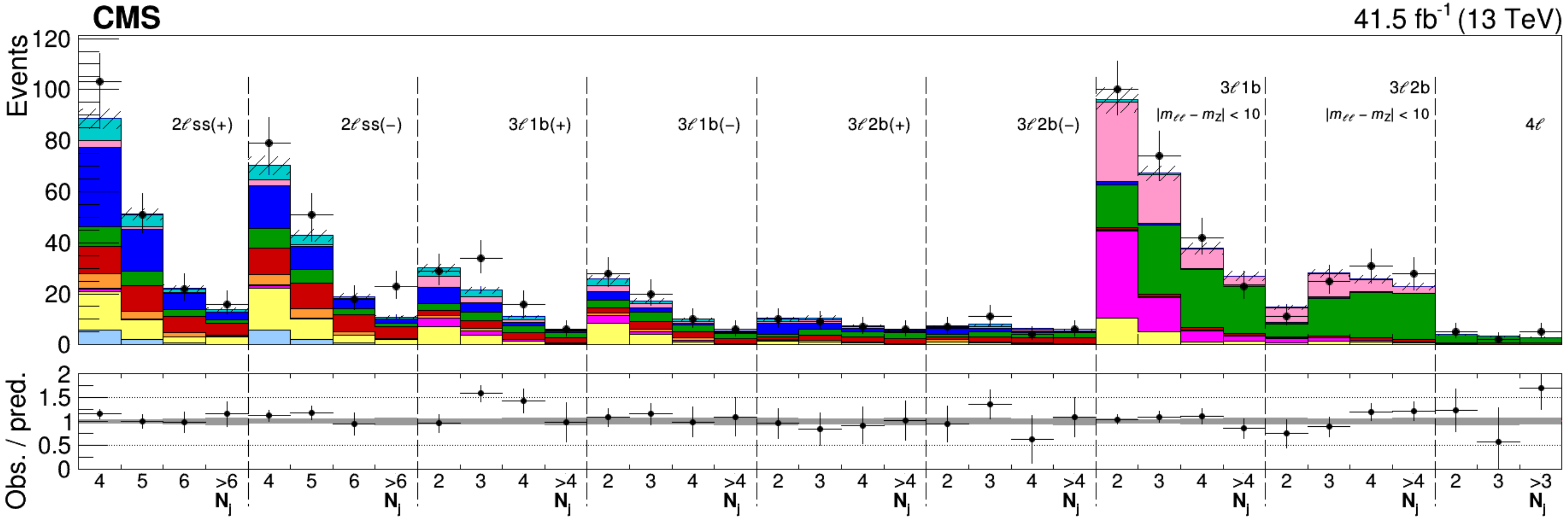
Event yields in analysis bins

EFT enhances predicted yields (WCs $= 4/6$ final value)



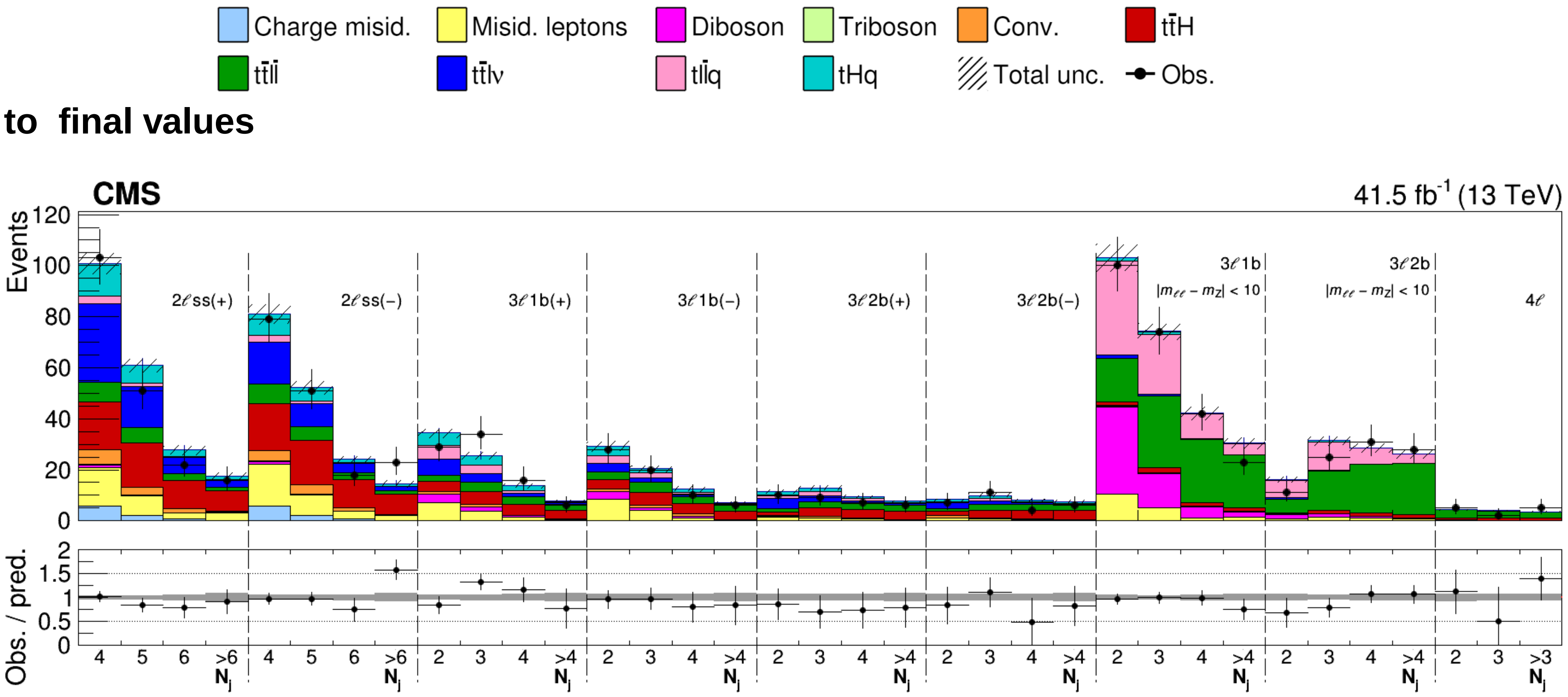
Event yields in analysis bins

EFT enhances
predicted yields (WCs $= 5/6$ final value)



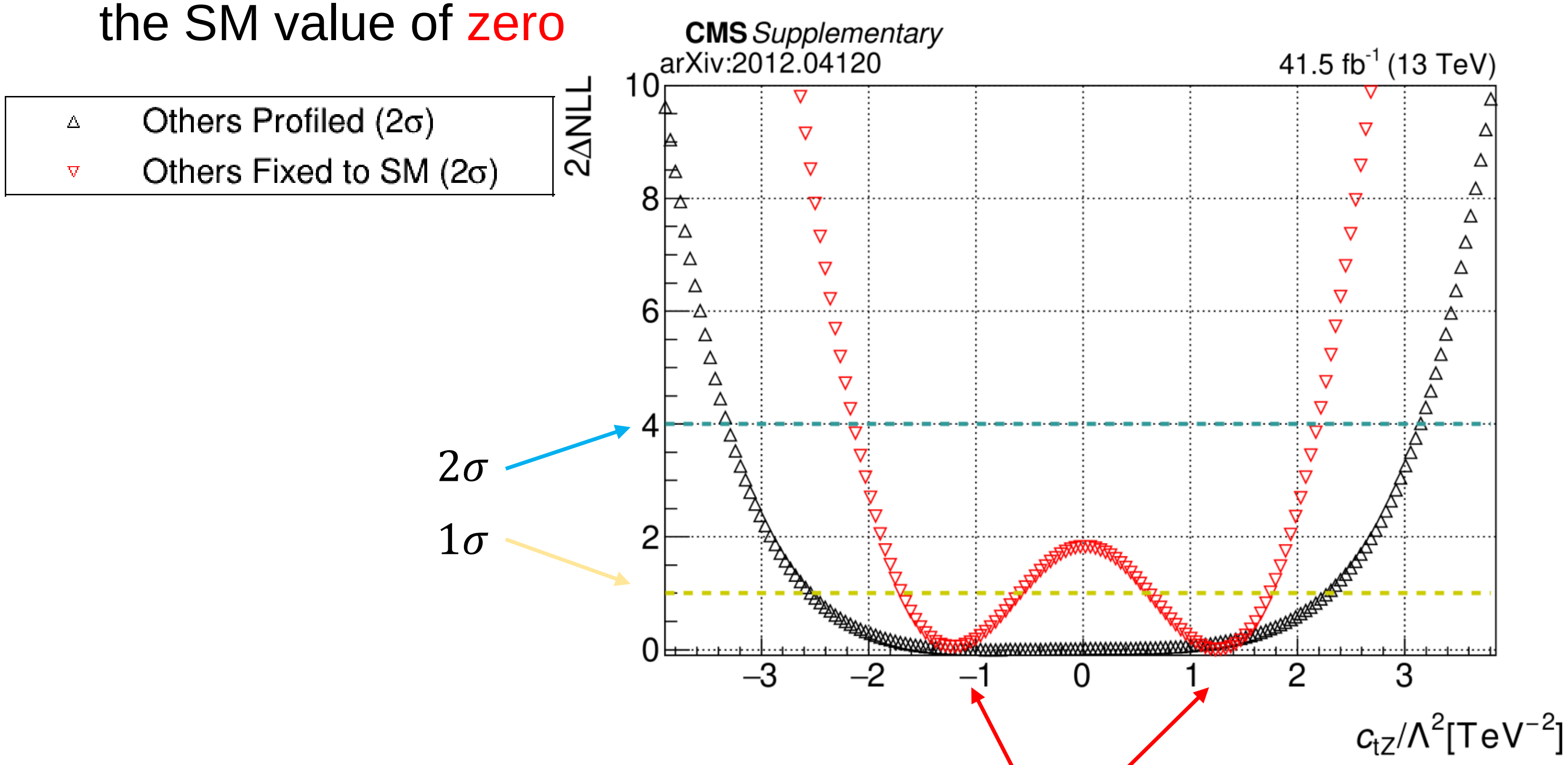
Event yields in analysis bins

WCs set to final values



Single WC scan

Scanning c_{tZ} while treating the other 15 as unconstrained nuisance parameters or fixed to the SM value of zero



Degenerate minima widen the NLL curve

Two-dimensional WC scans

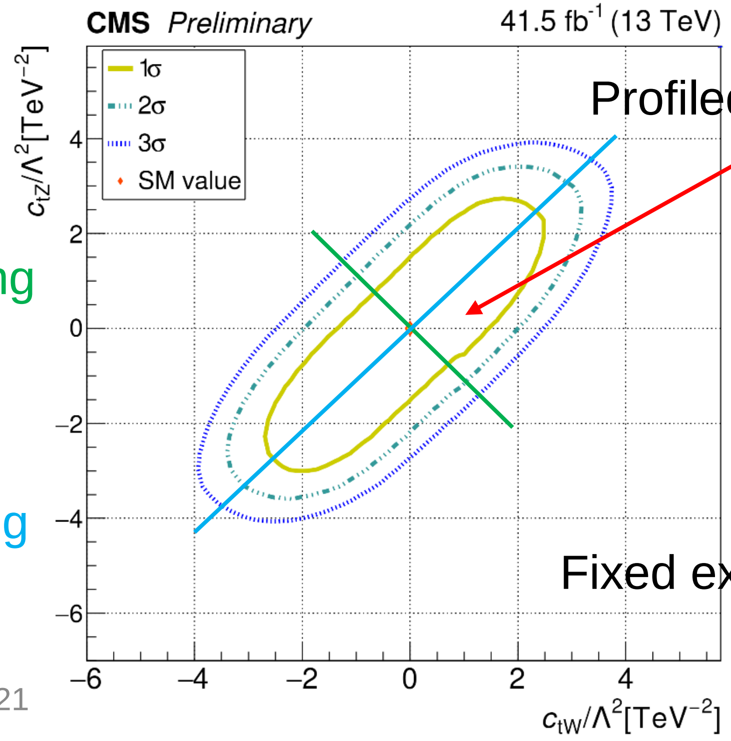
Pairs of WCs are also scanned to help investigate the correlations between WCs, as visualizing the full 16-dimensional hypersurface in not feasible

Other 14 WCs treated as unconstrained nuisance parameters

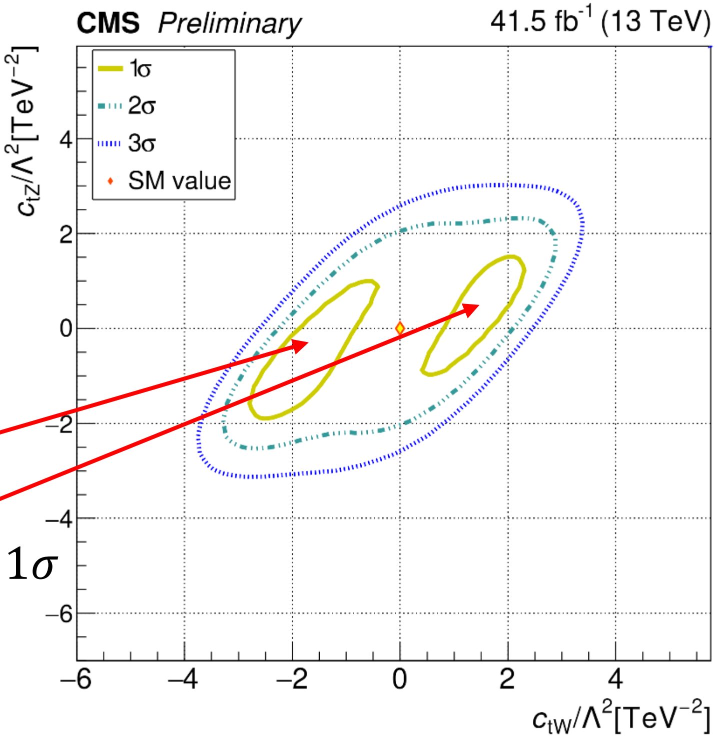
Other 14 WCs are fixed to zero (SM)

Better constraints along the negative vertical axis

Looser constraints along the positive vertical axis



Profiled includes the SM



Fixed excludes the SM at 1σ

2 WCs must compensate for other 14 → less correlation

Important systematic uncertainties

Analysis specific

- **Misidentified lepton rate estimate** – Contamination from non-prompt leptons
  Overcome by examining the analysis side-bands
  Data-driven → statistically limited

Monte Carlo simulation modeling

- **Matrix element parton shower matching** – Matching extra partons to final-state jets
- **Missing parton uncertainty** – Extra partons cannot be added for tHq and tllq
  Compare LO EFT without extra partons to NLO SM simulations, assign uncertainty to cover any discrepancies
  These issues will not be present in SMEFT@NLO, and we are very interested in the development
- **Scale uncertainties** – FSR and ISR

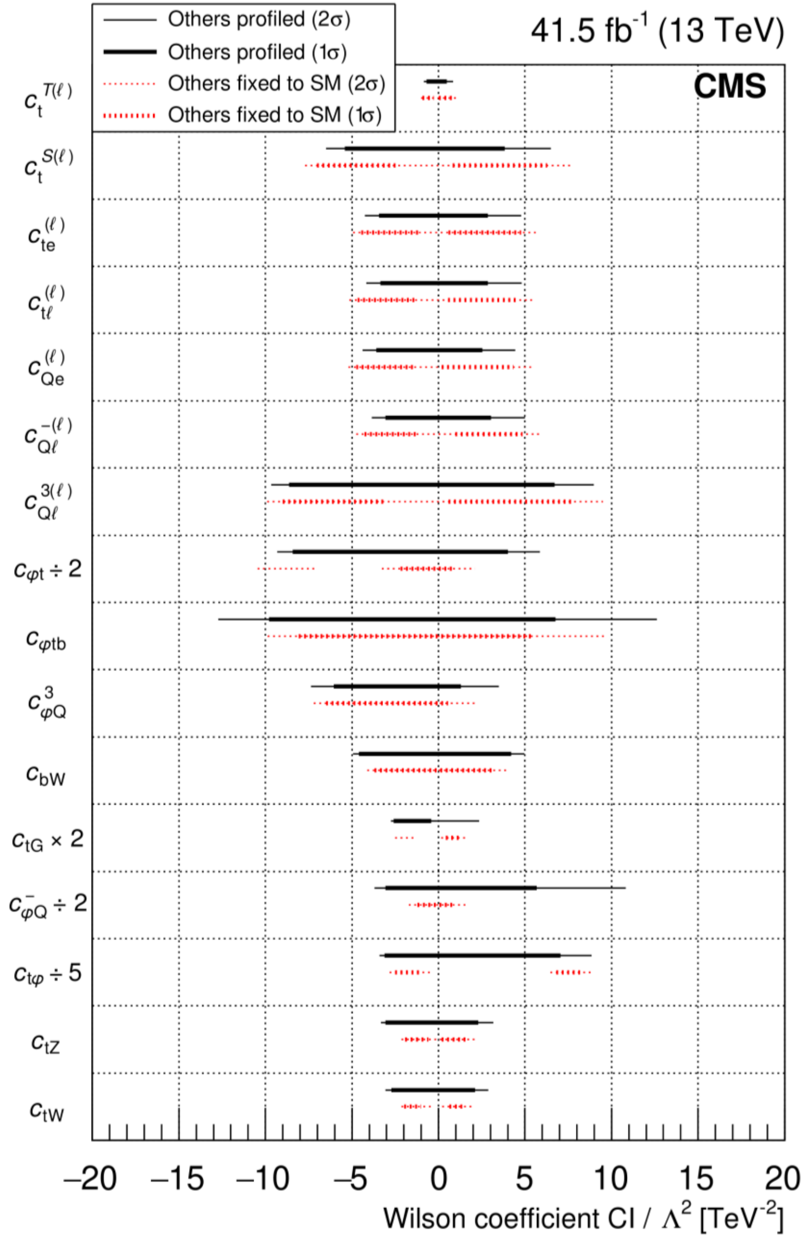
Wilson coefficient CIs

The 1σ and 2σ CIs are given

When the other WCs are fixed to zero, the fit can produce degenerate minima in c_{tW}, $c_{\mathrm{t}\varphi}$, c_{tG}, and $c_{\varphi\mathrm{t}}$

Degenerate minima are due to the quadratic nature of the parameterization

None of the WCs exclude the SM point of zero by a statistically significant amount



Conclusion

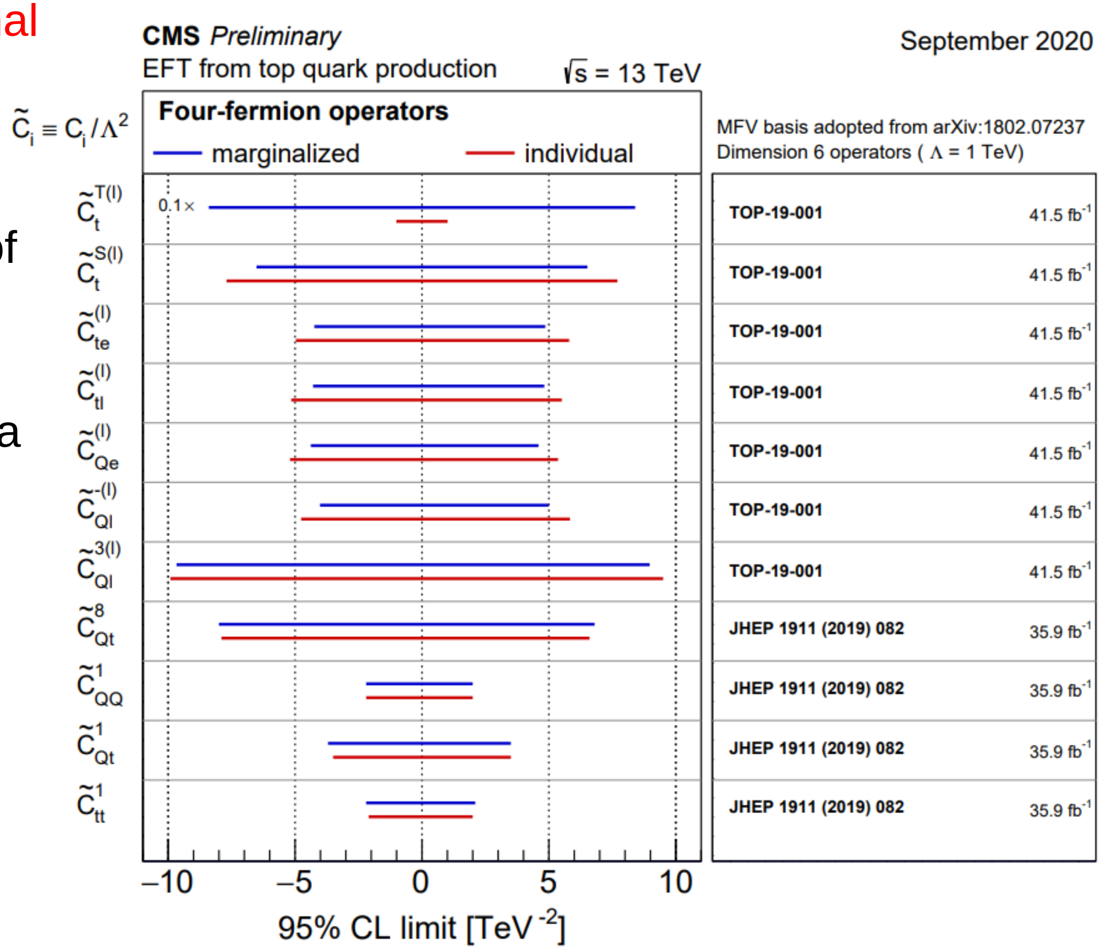
The production of or more t quarks in association with additional leptons were used to measure the confidence intervals of 16 dimension-six EFT operators using data collected in 2017

The EFT yields are parameterized using a quadratic function of event weights

This novel technique allows us to extract EFT from difficult data

The 2σ CIs were extracted for these operators
Intervals are compatible with the SM and other analyses [1]

With the full Run II data set (almost triple the integrated luminosity) more sophisticated analyses may be performed, including using differential distributions



[1] https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsTOPSummaryFigures

Backup

Selecting operators of interest

There are 59 total dimension-6 operators that conserve baryon and lepton number

We only consider 16 operators:

- Operators must appear in signal processes (top + boson) at tree level
- Ignore operators that strongly affect background processes
- Imaginary parts of non-Hermitian operators are set to zero

'Two heavy + boson': $c_{\mathrm{t}\varphi}$, $c_{\varphi\mathrm{Q}}^{-}$, $c_{\varphi\mathrm{Q}}^{3}$, $c_{\varphi\mathrm{t}}$, $c_{\varphi\mathrm{tb}}$, c_{tW}, c_{tZ}, c_{bW}, and c_{tG}

'Two heavy + two lepton': $c_{\mathrm{Ql}}^{3(l)}$, $c_{\mathrm{Ql}}^{-(l)}$, c_{Qe}, c_{tl}, c_{te}, $c_{\mathrm{t}}^{\mathrm{S}(l)}$, and $c_{\mathrm{t}}^{\mathrm{T}(l)}$

- These operators have three copies that couple to the different lepton flavors
- We assume equal coupling to all flavors to reduce the number of operators to these seven

Lepton identification

$$I_\ell = \sum_{\text{charged}} p_\mathrm{T} + \max\left(0, \sum_{\text{neutral}} p_\mathrm{T} - \rho\mathcal{A}\left(\frac{R}{0.3}\right)^2\right)$$

$$R = \begin{cases} 0.05 & \text{if } p_\mathrm{T} > 200\,\mathrm{GeV} \\ 10\,\mathrm{GeV}/p_\mathrm{T} & \text{if } 50 < p_\mathrm{T} < 200\,\mathrm{GeV} \\ 0.20 & \text{if } p_\mathrm{T} < 50\,\mathrm{GeV} \end{cases}$$

Electrons	
Pseudorapidity range	$\mathcal{A}$
$0.0 < \|\eta\| < 1.0$	0.1566
$1.0 < \|\eta\| < 1.479$	0.1626
$1.479 < \|\eta\| < 2.0$	0.1073
$2.0 < \|\eta\| < 2.2$	0.0854
$2.2 < \|\eta\| < 2.3$	0.1051
$2.3 < \|\eta\| < 2.4$	0.1204
$2.4 < \|\eta\| < 2.5$	0.1524

Muons	
Pseudorapidity range	$\mathcal{A}$
$0.0 < \|\eta\| < 0.8$	0.0566
$0.8 < \|\eta\| < 1.3$	0.0562
$1.3 < \|\eta\| < 2.0$	0.0363
$2.0 < \|\eta\| < 2.2$	0.0119
$2.2 < \|\eta\| < 2.5$	0.0064

Inputs to BDT

- Lepton p_{T} and η
- $I_{\ell}^{\mathrm{charged}}$
- $I_{\ell}^{\mathrm{neutral}}$
- $p_{\mathrm{T}}^{\ell}/p_{\mathrm{T}}^{\mathrm{jet}}$
- CSVv2 b-tagging algorithm
- N_{charged} of charged particles within the jet
- $p_{\mathrm{T}}^{\mathrm{rel}} = p_{\ell}\sin(\theta)$
- d_{xy} and d_z w.r.t. the PV
- d/σ_d signed 3D impact parameter significance w.r.t the PV
- Lepton MVA from EGamma POG
- Compatibility of track segments with the muon system

e cuts

Observable	Loose	Fakeable	Tight
p_{T}	$> 7\,\mathrm{GeV}$	$> 10\,\mathrm{GeV}$	$> 10\,\mathrm{GeV}$
$\lvert\eta\rvert$	< 2.5	< 2.5	< 2.5
$\lvert d_{xy}\rvert$	< 0.05 cm	< 0.05 cm	< 0.05 cm
$\lvert d_z\rvert$	< 0.1 cm	< 0.1 cm	< 0.1 cm
d/σ_d	< 8	< 8	< 8
I_e	$< 0.4 \times p_{\mathrm{T}}$	$< 0.4 \times p_{\mathrm{T}}$	$< 0.4 \times p_{\mathrm{T}}$
Missing hits	≤ 1	$= 0$	$= 0$
EGamma POG MVA	See Table 10	> 0.50	See Table 10
Conversion rejection	–	✓	✓
$\sigma_{i\eta i\eta}$	–	$< \{0.011\,/\,0.011\,/\,0.030\}$	$< \{0.011\,/\,0.011\,/\,0.030\}$
H/E	–	< 0.10	< 0.10
1/E - 1/p	–	$> -0.04\frac{1}{\mathrm{GeV}}$	$> -0.04\frac{1}{\mathrm{GeV}}$
$p_{\mathrm{T}}^{\mathrm{e}}/p_{\mathrm{T}}^{\mathrm{j}}$	–	> 0.6	–
DeepCSV(b) of nearby jet	–	< 0.07	< 0.4941
Prompt-e MVA	–	< 0.90	> 0.90

[1] $> \{-0.13\,/\,-0.32\,/\,-0.08\}$ if $p_{\mathrm{T}} < 10\,\mathrm{GeV}$
[2] > 0.50 if prompt-e MVA < 0.90

μ cuts

Observable	Loose	Fakeable	Tight
p_{T}	$> 5\,GeV$	$> 10\,GeV$	$> 10\,GeV$
$\lvert\eta\rvert$	< 2.4	< 2.4	< 2.4
$\lvert d_{xy}\rvert$	< 0.05 cm	< 0.05 cm	< 0.05 cm
$\lvert d_z\rvert$	< 0.1 cm	< 0.1 cm	< 0.1 cm
d/σ_d	< 8	< 8	< 8
I_μ	$< 0.4 \times p_{\mathrm{T}}$	$< 0.4 \times p_{\mathrm{T}}$	$< 0.4 \times p_{\mathrm{T}}$
Loose PF muon	✓	✓	✓
Medium PF muon	–	–	✓
Segment compatibility	–	> 0.3	–
$p_{\mathrm{T}}^{\mu}/p_{\mathrm{T}}^{\mathrm{j}}$	–	> 0.6	–
DeepCSV(b) of nearby jet	–	< 0.07	< 0.4941
Prompt-μ MVA	–	< 0.90	> 0.90

EGamma POG MVA cuts

	$\lvert\eta\rvert < 0.8$	$0.8 < \lvert\eta\rvert < 1.44$	$\lvert\eta\rvert > 1.44$
$p_{\mathrm{T}} < 10\,GeV$	-0.13	-0.32	-0.08
$p_{\mathrm{T}} > 10\,GeV$	-0.86	-0.81	-0.72

Event selection

	2ℓss	3ℓ	$\geq 4\ell$
Tight Leptons	2	3	$\geq$4
Jets	$\geq$4	$\geq$2	$\geq$2
DeepCSV(b) Medium Jets	$\geq$1	1, $\geq$2	$\geq$1
DeepCSV(b) Loose Jets	$\geq$2	-	$\geq$2
nJet Subcategories	4, 5, 6, $\geq$7	2, 3, 4, $\geq$5	2, 3, $\geq$4
Other Subcategories	Lepton Charge	Sign of Net Lepton Charge	-
		"SFOS Z"	-

	2ℓss ("p")	2ℓss ("m")	3ℓ (1b "p")	3ℓ (1b "m")	3ℓ ($\geq$ 2b "p")	3ℓ ($\geq$ 2b "m")	3ℓ (SFOS Z,1b)	3ℓ (SFOS Z,$\geq$ 2b)	$\geq 4\ell$
Diboson	1.6	1.2	5.9	4.7	0.4	0.3	52.1	4.1	0.6
Triboson	0.5	0.5	0.2	0.2	0.0	0.1	3.5	0.6	0.1
Charge Flips	8.5	8.5	0.0	0.0	0.0	0.0	0.0	0.0	0.0
Fakes	25.6	26.8	11.3	13.0	3.3	2.5	16.9	3.8	0.0
Conversions	10.9	9.2	2.3	2.6	1.7	1.9	0.8	0.4	0.0
Sum Background	47.0	46.2	19.7	20.5	5.4	4.8	73.3	8.9	0.7
$t\bar{t}l\nu$	68.7	37.1	14.4	8.0	10.8	5.9	2.9	2.3	0.0
$t\bar{t}l\bar{l}$	19.3	19.0	12.7	13.3	9.1	8.5	95.5	63.2	9.4
$t\bar{t}H$	24.7	24.1	7.9	7.6	5.1	5.2	3.2	2.2	1.0
$tl\bar{l}q$	2.7	1.5	3.5	1.8	1.2	0.6	39.8	13.3	0.0
tHq	0.8	0.4	0.3	0.2	0.2	0.1	0.1	0.1	0.0
Sum Signal	116.2	82.2	38.9	30.8	26.3	20.3	141.6	81.1	10.4
Total Expected	163 ± 20	128 ± 15	59 ± 7	51 ± 6	32 ± 4	25 ± 3	215 ± 25	90 ± 13	11 ± 2
Data	192	171	85	64	32	28	239	95	12

Background Estimation

Misidentified leptons

- Data-driven by dividing data into measurement and application regions (analysis side-bands)
- Measurement region contains QCD multijet background dominated by non-prompt leptons
- A fake rate is derived by comparing looser lepton cuts to tight lepton cuts used in the main analysis
- Fake rate is applied to the application region

Lepton charge mismeasurement

- Also data-driven
- Charge mismeasurement rate is extracted from the $2\ell ss$ region using $Z/\gamma^* \to ee$
- Only applied to the ee region of the analysis

Lepton charge mismeasurement

$2\ell ss$ region using $Z/\gamma^* \to ee$ used to estimate rate at which the CMS detector incorrectly measures lepton charge

Data-driven

Only applied to the ee region of the analysis

Systematic uncertainties

The systematic uncertainties are: standard and analysis specific

Luminosity – vary simulation by integrated luminosity estimate uncertainty

Jet energy scale (JES) – account for pileup, nonuniform detector response, and any residual differences between the data and simulation

b jet tagging scale factors – account for tagging inefficiencies and charm jet contamination

Cross section theoretical uncertainty – vary cross sections in simulation by uncertainties

PDF shape variations – reweighting the spectra according to the 100 replica sets given by the NNPDF31_NLO_as_0118 PDF parameterization

Renormalization and factorization scale – vary scales by $1/2 \; - 2$

Parton shower – vary ISR by 2 and FSR by $\sqrt{2}$

Matching uncertainty – vary the matching scale between MADGRAPH_aMC@NLO and PYTHIA

Muon and electron ID isolation – vary corrections by their uncertainties

Trigger efficiency – vary corrections by their uncertainties

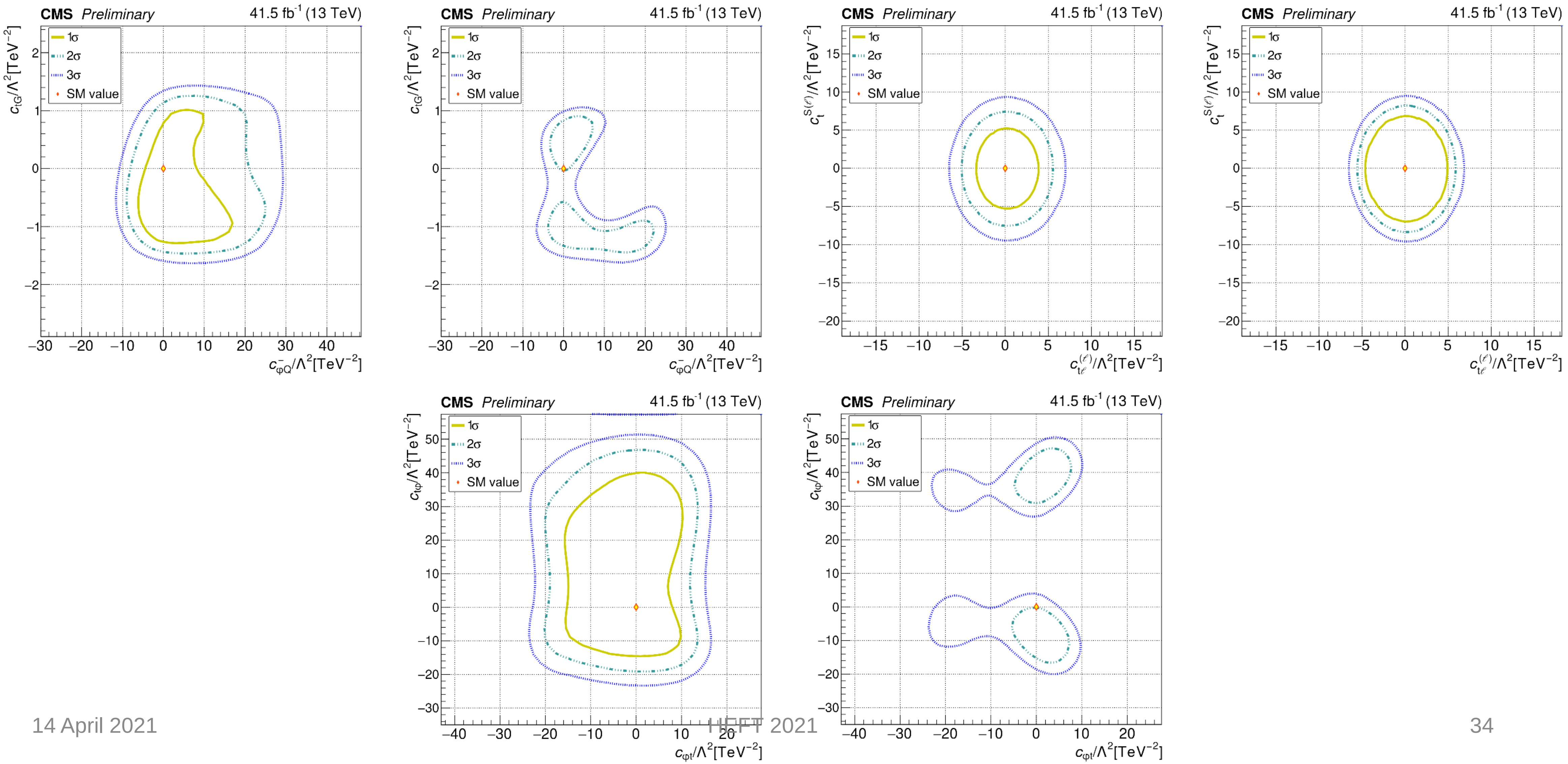
Pileup – vary the pp inelastic cross section by 5%

Missing parton uncertainty – cover differences between LO EFT w/o extra partons and NLO SM

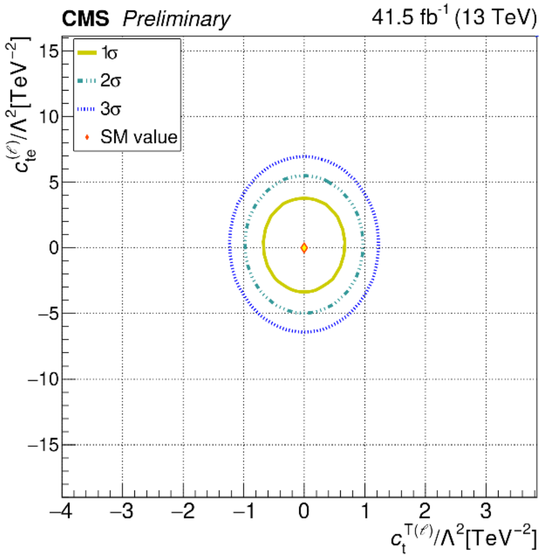
Uncertainty on the misidentified lepton rate estimate (data-driven) – account for non-prompt contamination

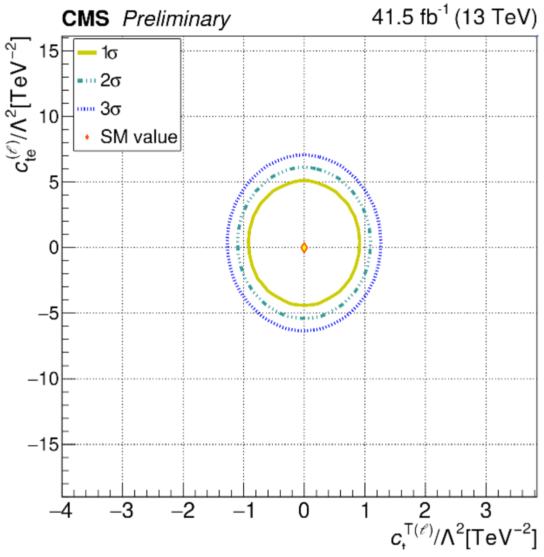
Uncertainty on the charge mismeasurement estimate (data-driven) – account for $\mathrm{Z}/\gamma^* \to \mathrm{e}^{\pm}\mathrm{e}^{\mp}$ becoming $\mathrm{Z}/\gamma^* \to \mathrm{e}^{\pm}\mathrm{e}^{\pm}$

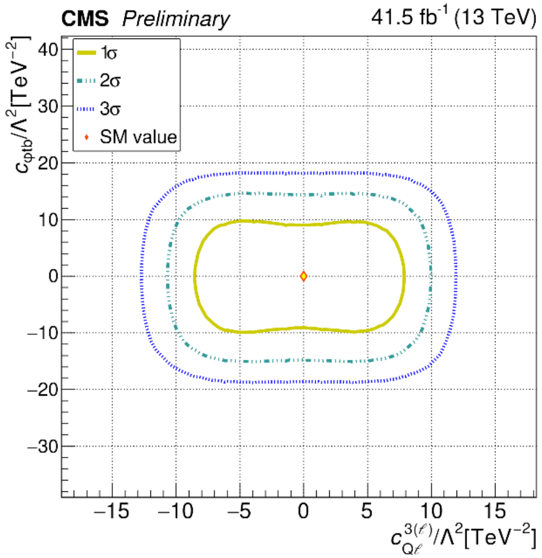
Complete set of 2D contours

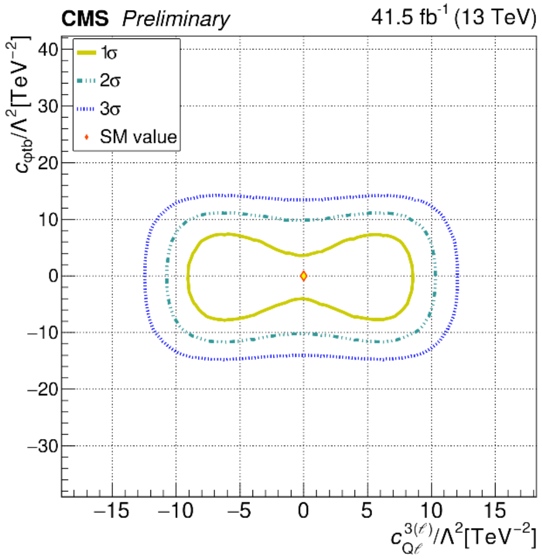


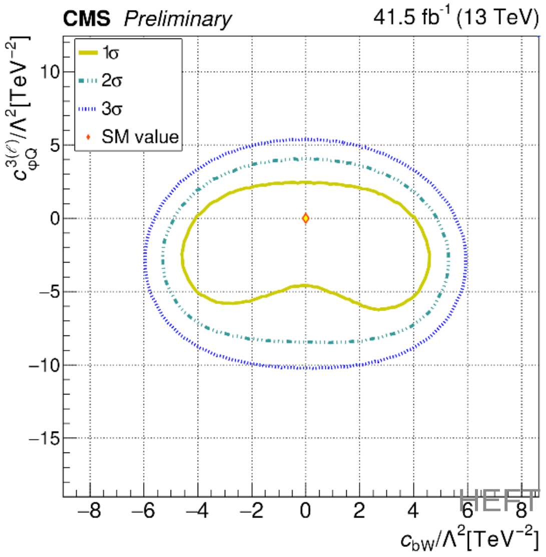
Complete set of 2D contours

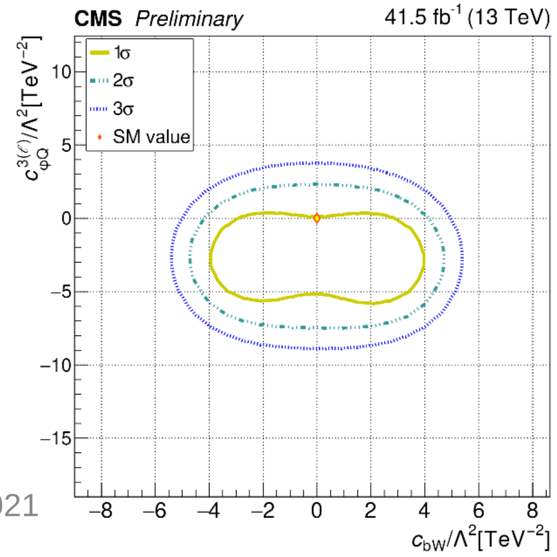












Changes in event yields over the 2σ CI

ttlv ttll ttH tllq tHq |15

Examining the minimum and maximum yield changes within the 2σ CI of various WCs