

Flavor Specific $U(1)_{Bq-L2}$ Gauge Model for Muon $g-2$ and $b \rightarrow s$ Anomalies

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Flavor specific $U(1)_{B2-L2}$ model

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Effects of kinetic mixing

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4 Effects of kinetic mixing

Anomalies



Fermilab

muon anomalous magnetic
dipole moment

$$a_\mu = (g - 2)/2$$



$$\begin{aligned} \Delta a_\mu &= a_\mu^{exp} - a_\mu^{SM} \quad 4.2\sigma \\ &= (251 \pm 59) \times 10^{-11} \end{aligned}$$

Phys. Rev. Lett. 126 (2021) 141801

LHCb

$b \rightarrow s \bar{\mu} \mu$ induced anomalies

$$R_k = \frac{BR(B \rightarrow K \mu^+ \mu^-)}{BR(B \rightarrow K e^+ e^-)}$$



$$R_K = 0.846_{-0.039-0.012}^{+0.042+0.013} \quad 3.1\sigma$$

$$q^2 \in (1.1, 6) \text{ GeV}^2$$

arXiv:2103.11769

New physics beyond the SM

New gauge boson



$$\begin{cases} (g - 2)_\mu \\ b \rightarrow s \bar{\mu} \mu \end{cases}$$



involve muon

2nd-3rd quark mixing



Model with new gauge interactions



Gauge anomaly free $U(1)_{L_i - L_j}, U(1)_{B-L}$

Z'_μ



Flavor changing Z' interaction

Flavor changing Z' interaction



1. New type of quarks

Phys. Rev. D 89 (2014) 095033

2. Non-trivial $U(1)$ quantum number

Phys. Lett. B 774 (2017) 643, arXiv:2103.13991,

Eur. Phys. J. C 81 (2021) 56, Phys. Rev. Lett. 114(2015) 151801

3. $U(1)$ - $U(1)_Y$ kinetic mixing



Flavor specific $U(1)_{B_q-L_2}$ gauge model

Only act on the 2nd generation of fermions in weak basis

$$B_q = B_2 = 1/3, L_\mu = 1$$

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Components



Gauge group $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_{B_2-L_\mu}$

g_s g g' \tilde{g}

Components $H = (1, 2, 1/2)(0)$

$$Q_L = (Q_{L_1}, Q_{L_2}, Q_{L_3}) = (3, 2, 1/6)(0, 1/3, 0)$$

$$U_R = (u_{R_1}, u_{R_2}, u_{R_3}) = (3, 1, 2/3)(0, 1/3, 0)$$

$$D_R = (d_{R_1}, d_{R_2}, d_{R_3}) = (3, 1, -1/3)(0, 1/3, 0)$$

$$L_L = (L_{L_1}, L_{L_2}, L_{L_3}) = (1, 2, -1/2)(0, -1, 0)$$

$$l_R = (l_{R_1}, l_{R_2}, l_{R_3}) = (1, 1, -1)(0, -1, 0)$$

$$\nu_R = (\nu_{R_1}, \nu_{R_2}, \nu_{R_3}) = (1, 1, 0)(0, -1, 0)$$

Interaction with fermions




Z' Interaction Lagrangian

$$L_{int} = \frac{1}{3} \tilde{g} \left(\bar{Q}_{L_2} \gamma^\mu Q_{L_2} + \bar{u}_{R_2} \gamma^\mu u_{R_2} + \bar{d}_{R_2} \gamma^\mu d_{R_2} \right) Z'_\mu \\ - \tilde{g} \left(\bar{L}_{L_2} \gamma^\mu L_{L_2} + \bar{l}_{R_2} \gamma^\mu l_{R_2} + \bar{\nu}_{R_2} \gamma^\mu \nu_{R_2} \right) Z'_\mu$$

Yukawa Interaction and RH neutrino mass term

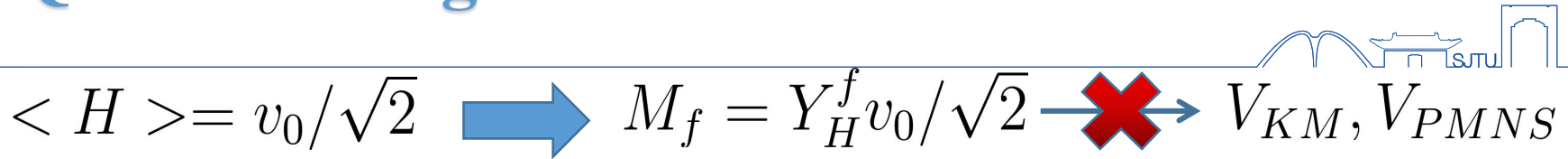
$$-L_{Y-mass} = \bar{Q}_L Y_H^u U_R \tilde{H} + \bar{Q}_L Y_H^d D_R H + \bar{L}_L Y_H^\nu \nu_R \tilde{H} \\ + \bar{L}_L \underbrace{Y_H^l}_{\text{circled}} l_R H + \frac{1}{2} \bar{\nu}_R^c \underbrace{\tilde{M}_R}_{\text{circled}} \nu_R + \text{H.c.}$$

$$Y_H^f = \begin{pmatrix} Y_{11}^f & 0 & Y_{13}^f \\ 0 & Y_{22}^f & 0 \\ Y_{31}^f & 0 & Y_{33}^f \end{pmatrix}$$



$$\tilde{M}_R = \begin{pmatrix} M_{11} & 0 & M_{13} \\ 0 & 0 & 0 \\ M_{13} & 0 & M_{33} \end{pmatrix}$$

Quark mixing

$$\langle H \rangle = v_0/\sqrt{2} \quad \longrightarrow \quad M_f = Y_H^f v_0/\sqrt{2} \quad \not\longrightarrow \quad V_{KM}, V_{PMNS}$$


Introduce another new scalars

$$\left\{ \begin{array}{l} \text{Doublets} \quad H_1^q : (1, 2, 1/2)(1/3), H_2^q : (1, 2, 1/2)(-1/3), \\ \text{Singlets} \quad S_1 : (1, 1, 0)(1), S_2 : (1, 1, 0)(2) \end{array} \right.$$

New added quark and neutrino interactions

$$\begin{aligned} & -\bar{Q}_L (Y_{H_1^q}^u \tilde{H}_1^q + Y_{H_2^q}^u \tilde{H}_2^q) U_R - \bar{Q}_L (Y_{H_1^q}^d H_1^q + Y_{H_2^q}^d H_2^q) D_R \\ & - \frac{1}{2} \bar{\nu}_R^c (Y_{S_1} S_1 + Y_{S_2} S_2) \nu_R + \text{H.c.} \end{aligned}$$

Backup

Mass matrix

$\langle H_i \rangle = \frac{v_i}{\sqrt{2}}$

Quark

$$\begin{cases}
 M_u = \frac{v_0}{\sqrt{2}} Y_H^u + \frac{v_1^q}{\sqrt{2}} Y_{H_1^q}^u + \frac{v_2^q}{\sqrt{2}} Y_{H_2^q}^u, \\
 M_d = \frac{v_0}{\sqrt{2}} Y_H^d + \frac{v_1^q}{\sqrt{2}} Y_{H_1^q}^d + \frac{v_2^q}{\sqrt{2}} Y_{H_2^q}^d,
 \end{cases}$$

Lepton

$$\begin{cases}
 M_l = \frac{v_0}{\sqrt{2}} Y_H^l \\
 M_D = (v_0/\sqrt{2}) Y_H^\nu \\
 M_R = \tilde{M}_R + (v_{S_1}/\sqrt{2}) Y_{S_1} + (v_{S_2}/\sqrt{2}) Y_{S_2}
 \end{cases}$$

$$M_\nu = \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix}$$

No charged lepton mixing



New doublets

Other constraints

Generate mu-tau mixing

Lepton mixing

Doublets $H_1^l : (1, 2, 1/2)(1), H_2^l : (1, 2, 1/2)(-1),$

New added charged lepton interactions

$$- \left(\bar{L}_L (\underline{Y_{H_1^l}^\nu} \tilde{H}_1^l + \underline{Y_{H_2^l}^\nu} \tilde{H}_2^l) \nu_R + \bar{L}_L (\underline{Y_{H_1^l}^l} H_1^l + \underline{Y_{H_2^l}^l} H_2^l) E_R \right) + \text{H.c.}$$

$$Y_{H_1^l}^\nu = \begin{pmatrix} 0 & 0 & 0 \\ Y_{H_1^l}^{\nu 21} & 0 & Y_{H_1^l}^{\nu 23} \\ 0 & 0 & 0 \end{pmatrix}, \quad Y_{H_1^l}^l = \begin{pmatrix} 0 & Y_{H_1^l}^{l12} & 0 \\ 0 & 0 & 0 \\ 0 & Y_{H_1^l}^{l32} & 0 \end{pmatrix}$$

$$Y_{H_2^l}^\nu = \begin{pmatrix} 0 & Y_{H_2^l}^{\nu 12} & 0 \\ 0 & 0 & 0 \\ 0 & Y_{H_2^l}^{\nu 32} & 0 \end{pmatrix}, \quad Y_{H_2^l}^l = \begin{pmatrix} 0 & 0 & 0 \\ Y_{H_2^l}^{l21} & 0 & Y_{H_2^l}^{l23} \\ 0 & 0 & 0 \end{pmatrix}$$

Lepton mass matrix

$$\begin{aligned}
 M_l &= \frac{v_0}{\sqrt{2}} Y_H^l & \text{Modify} & \longrightarrow & M_l &= \frac{v_0}{\sqrt{2}} Y_H^l + \frac{v_1^l}{\sqrt{2}} Y_{H_1^l}^l + \frac{v_2^l}{\sqrt{2}} Y_{H_2^l}^l \\
 M_D &= \frac{v_0}{\sqrt{2}} Y_H^\nu & & & M_D &= \frac{v_0}{\sqrt{2}} Y_H^\nu + \frac{v_1^l}{\sqrt{2}} Y_{H_1^l}^\nu + \frac{v_2^l}{\sqrt{2}} Y_{H_2^l}^\nu
 \end{aligned}$$

Diagonalize mass matrix

$$M_f = V_L^{f\dagger} \hat{M}_f V_R^f, \quad M_\nu = V^\nu T \hat{M}_\nu V^\nu$$

$$V_{KM} = V_L^u V_L^{d\dagger}, \quad V_{PMNS} \approx V_L^l V_{3 \times 3}^{\nu T}$$

Mass-squared matrix (No Z-Z' mixing) $v_1^{q,l} = v_2^{q,l}$

$$m_{\frac{Z}{Z'}}^2 = \frac{g^2 + g'^2}{4} (v_0^2 + (v_1^q)^2 + (v_2^q)^2 + (v_1^l)^2 + (v_2^l)^2)$$

$$m_{\frac{Z}{Z'}}^{\prime 2} = \tilde{g}^2 \left(\frac{1}{(3)^2} ((v_1^q)^2 + (v_2^q)^2) + (v_1^l)^2 + (v_2^l)^2 + v_{S_1}^2 + 4v_{S_2}^2 \right)$$

Z'-q/l interaction in the mass eigen-state

The simplest case

$$N_2 = \text{diag}(0, 1, 0)$$

$$L_{int-f} = \frac{1}{3} \tilde{g} (\bar{U}_L V_L^u N_2 V_L^{u\dagger} \gamma^\mu U_L + \bar{U}_R V_R^u N_2 V_R^{u\dagger} \gamma^\mu U_R \\ + \bar{D}_L V_L^d N_2 V_L^{d\dagger} \gamma^\mu D_L + \bar{D}_R V_R^d N_2 V_R^{d\dagger} \gamma^\mu D_R) Z'_\mu \\ - \tilde{g} (\bar{l}_L V_L^l N_2 V_L^{l\dagger} \gamma^\mu l_L + \bar{l}_R V_R^l N_2 V_R^{l\dagger} \gamma^\mu l_R) Z'_\mu$$

↑ $V_{L22}^l = 1$

The general case

positive g-2

$$L_{int-sb,\mu,\nu} = -\tilde{g} \left((V_{L22}^l V_{L22}^{l*}) \bar{\mu} \gamma^\mu \mu + (V_{L32}^l V_{L32}^{l*}) (\bar{\tau} \gamma^\mu \tau) \right) Z'_\mu \\ - \tilde{g} \left((V_{L22}^l V_{L32}^{l*}) \bar{\mu} \gamma^\mu \tau + (V_{L32}^l V_{L22}^{l*}) \bar{\tau} \gamma^\mu \mu \right) Z'_\mu \\ + \frac{1}{3} \tilde{g} \left(V_{L22}^d V_{L32}^{d*} (\bar{s}_L \gamma^\mu b_L) \right) Z'_\mu \quad \text{LH for } C_9$$

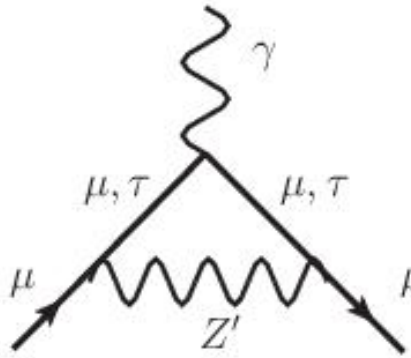
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Muon g-2



$$\Delta a_\mu = (251 \pm 59) \times 10^{-11}$$

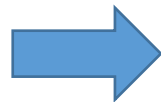
$$\Delta a_\mu^{new} = \frac{\tilde{g}^2 m_\mu^2}{4\pi^2} (V_{L 22}^l V_{L 22}^{l*}) \int_0^1 dx$$

$$\times \left[\frac{(V_{L 32}^l V_{L 32}^{l*})(x-x^2)\left(x + \frac{2m_\tau}{m_\mu} - 2\right)}{m_\mu^2 x^2 + m_{Z'}^2(1-x) + (m_\tau^2 - m_\mu^2)x} + \frac{(V_{L 22}^l V_{L 22}^{l*})x^2(1-x)}{m_\mu^2 x^2 + m_{Z'}^2(1-x)} \right]$$

$$m_{Z'} \gg m_{\tau, \mu} \quad \Downarrow \quad |V_{L 32}^l|^2 = 1 - |V_{L 22}^l|^2$$

$$\frac{\tilde{g}^2}{m_{Z'}^2} = \frac{4\pi^2}{m_\mu^2} \frac{\Delta a_\mu}{|V_{L 22}^l|^2 \left[\left(\frac{m_\tau}{m_\mu} - \frac{2}{3} \right) + |V_{L 22}^l|^2 \left(1 - \frac{m_\tau}{m_\mu} \right) \right]}$$

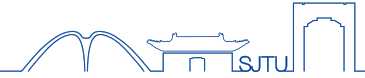
$$\begin{cases} |V_{L 22}^l| = 1 \\ |V_{L 22}^l| = 0.15 \end{cases}$$



$$\tilde{g}^2 / m_{Z'}^2 = (2.66 \pm 0.63) \times 10^{-5} \text{ GeV}^{-2}$$

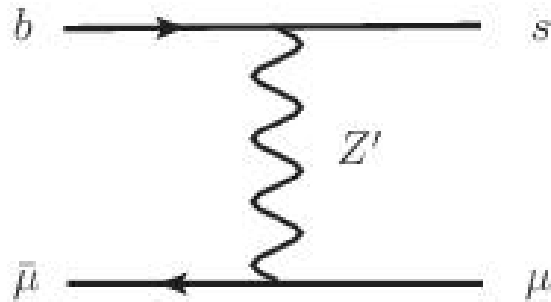
$$\tilde{g}^2 / m_{Z'}^2 = (2.50 \pm 0.59) \times 10^{-5} \text{ GeV}^{-2}$$

b anomaly



Effective Hamiltonian

$$H_{eff} = -\frac{\tilde{g}^2}{3m_{Z'}^2} (V_{L 22}^d V_{L 32}^{d*}) (V_{L 22}^l V_{L 32}^{l*}) (\bar{s}_L \gamma^\mu b_L \bar{\mu} \gamma_\mu \mu)$$



Standard form

$$H_{eff} = -(G_F \alpha_{em} / \sqrt{2} \pi) V_{tb} V_{ts}^* \sum_i C_i O_i$$

$$O_9 = \bar{s}_L \gamma^\mu b_L \bar{\mu} \gamma_\mu \mu,$$

$$C_9 = -0.80 \pm 0.14, 5.7\sigma$$

$$C_9^{new} = \frac{\sqrt{2} \pi \tilde{g}^2}{3m_{Z'}^2 G_F \alpha_{em}} (V_{L 22}^l V_{L 22}^{l*}) \frac{V_{L 22}^d V_{L 32}^{d*}}{V_{tb} V_{ts}^*}$$

arXiv:2103.13370

$$\frac{\tilde{g}^2}{m_{Z'}^2} \frac{V_{L 22}^d V_{L 32}^{d*}}{V_{tb} V_{ts}^*} = \frac{3 G_F \alpha_{em}}{\sqrt{2} \pi} \frac{C_9}{|V_{L 22}^l|^2}$$

$$\begin{cases} |V_{L 22}^l| = 1 \\ |V_{L 22}^l| = 0.15 \end{cases}$$

$$= (-0.46 \pm 0.08) \times 10^{-7} \text{ GeV}^{-2}$$

$$= (-2.04 \pm 0.36) \times 10^{-6} \text{ GeV}^{-2}$$

The simplest case



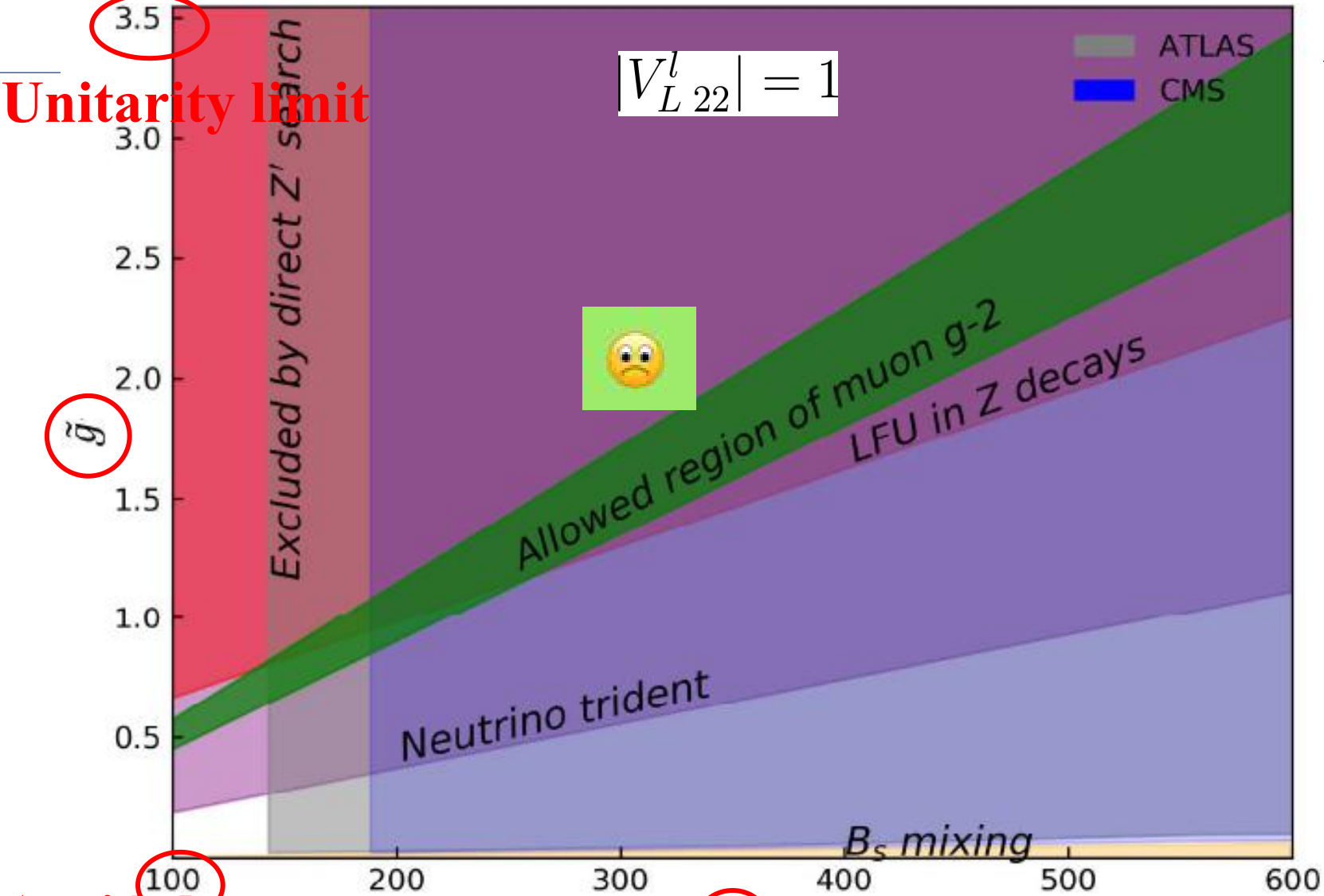
Unitarity limit

$$|V_{L 22}^l| = 1$$

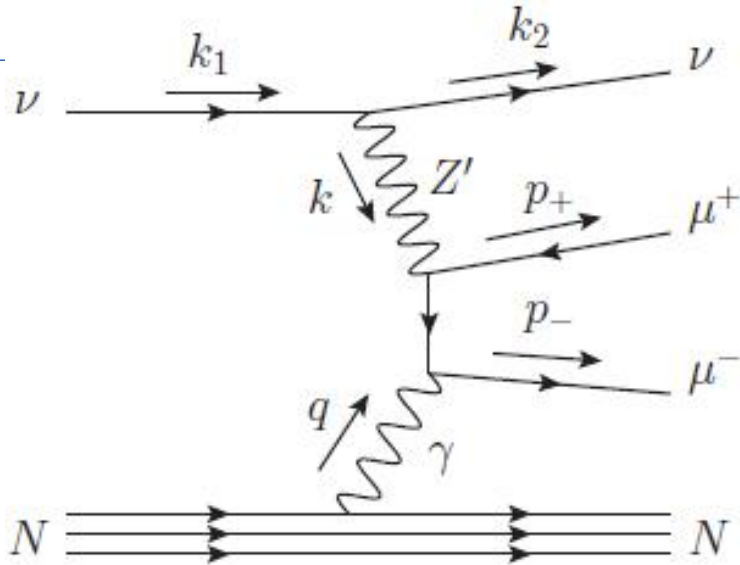
g

Avoid degeneracy

$m_{Z'}$



Neutrino trident production



$$\nu_{\mu} N \rightarrow \nu_{\mu} N \mu^{+} \mu^{-}$$

$$\sigma^{\text{CCFR}} / \sigma^{\text{SM}} = 0.82 \pm 0.28$$

PhysRevLett.66.3117

Suppression factor

$$\frac{\sigma^{NP}}{\sigma^{SM}} = 1 + \frac{(1+4s_W^2) \frac{4v^2 \tilde{g}^2}{m_{Z'}^2} (V_{L22}^l V_{L22}^{l*}) + \frac{4v^4 \tilde{g}^4}{m_{Z'}^4} (V_{L22}^l V_{L22}^{l*})^2}{1 + (1+4s_W^2)^2}$$

decrease

$$|V_{L22}^l| = 1 \quad \longrightarrow \quad |V_{L22}^l| = 0.15 \quad \longrightarrow \quad \text{Weaker constraint}$$

LFU of Z coupling

Lepton flavor universality of Z decay

$Z' - \mu\mu, Z' - \nu\nu \longrightarrow$ **Loop effects**

$$\frac{g_{V\mu}}{g_{Ve}} \simeq \left(\frac{g_{A\mu}}{g_{Ae}} \right) \simeq \left| 1 + \frac{\tilde{g}^2}{(4\pi)^2} (V_{L22}^l V_{L22}^{l*}) k_F(m_Z^2/m_{Z'}^2) \right|$$

Suppression factor

$$\frac{g_{V\nu}}{g_{Ve}} \simeq \frac{g_{A\nu}}{g_{Ae}} \simeq \left| 1 + \frac{\tilde{g}^2}{(4\pi)^2} \frac{1}{3} k_F(m_Z^2/m_{Z'}^2) \right| \quad \text{JHEP08(2011),088}$$

$$g_{Ve} = -0.03817 \pm 0.00047, \quad g_{Ae} = -0.50111 \pm 0.00035$$

$$g_{V\mu} = -0.0367 \pm 0.0023, \quad g_{A\mu} = -0.50120 \pm 0.00054 \quad \text{Strongest}$$

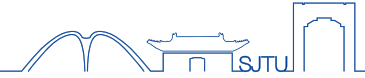
$$g_{V\nu} = g_{A\nu} = 0.5008 \pm 0.0008$$

$$g_{A\mu}/g_{Ae} = 1.00018 \pm 0.00128$$

decrease

$$|V_{L22}^l| = 1 \longrightarrow |V_{L22}^l| = 0.15 \longrightarrow \text{Weaker constraint}$$

LHC direct search



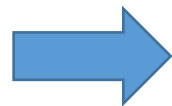
Di-muon pair production

Resonant Z' vector bosons decay into different final states

Excluded space $\left\{ \begin{array}{l} \text{ATLAS : } 150 \text{ GeV} \leq m_{Z'} \leq 5 \text{ TeV} \\ \text{CMS : } 200 \text{ GeV} \leq m_{Z'} \leq 5.5 \text{ TeV} \end{array} \right.$

arXiv:1707.02424 [hep-ex], arXiv:1803.06292 [hep-ex]

$$|V_{L22}^l| = 1$$



Exclude all $m_{Z'} \geq 150 \text{ GeV}$

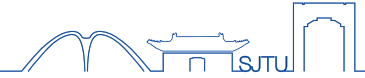
$$|V_{L22}^l| = 0.15$$

Allow to exist $m_{Z'} \geq 150 \text{ GeV}$

$$|V_{L22}^l|$$

The Suppression factor opens more parameter space

B_s mixing



$$B_s - \bar{B}_s$$

$$\frac{M_{12}}{M_{12}^{SM}} = 1 + \frac{\tilde{g}^2}{m_{Z'}^2} \left(\frac{V_{L 22}^d V_{L 32}^{d*}}{V_{tb} V_{ts}^*} \right)^2 \left(\frac{9g^2 S_0}{16\pi^2 v^2} \right)^{-1} \quad S_0 \simeq 2.3$$

arXiv:1008.1593 [hep-ph], arXiv:1211.1896 [hep-ph]

$$\left| \frac{\tilde{g}}{m_{Z'}} \frac{V_{L 22}^d V_{L 32}^{d*}}{V_{tb} V_{ts}^*} \right| = \frac{3g}{4\pi v} \sqrt{\left(\frac{M_{12}}{M_{12}^{SM}} - 1 \right) S_0}$$

$$M_{12}/M_{12}^{SM} = 0.959_{-0.078}^{+0.062}$$

Combine C_9

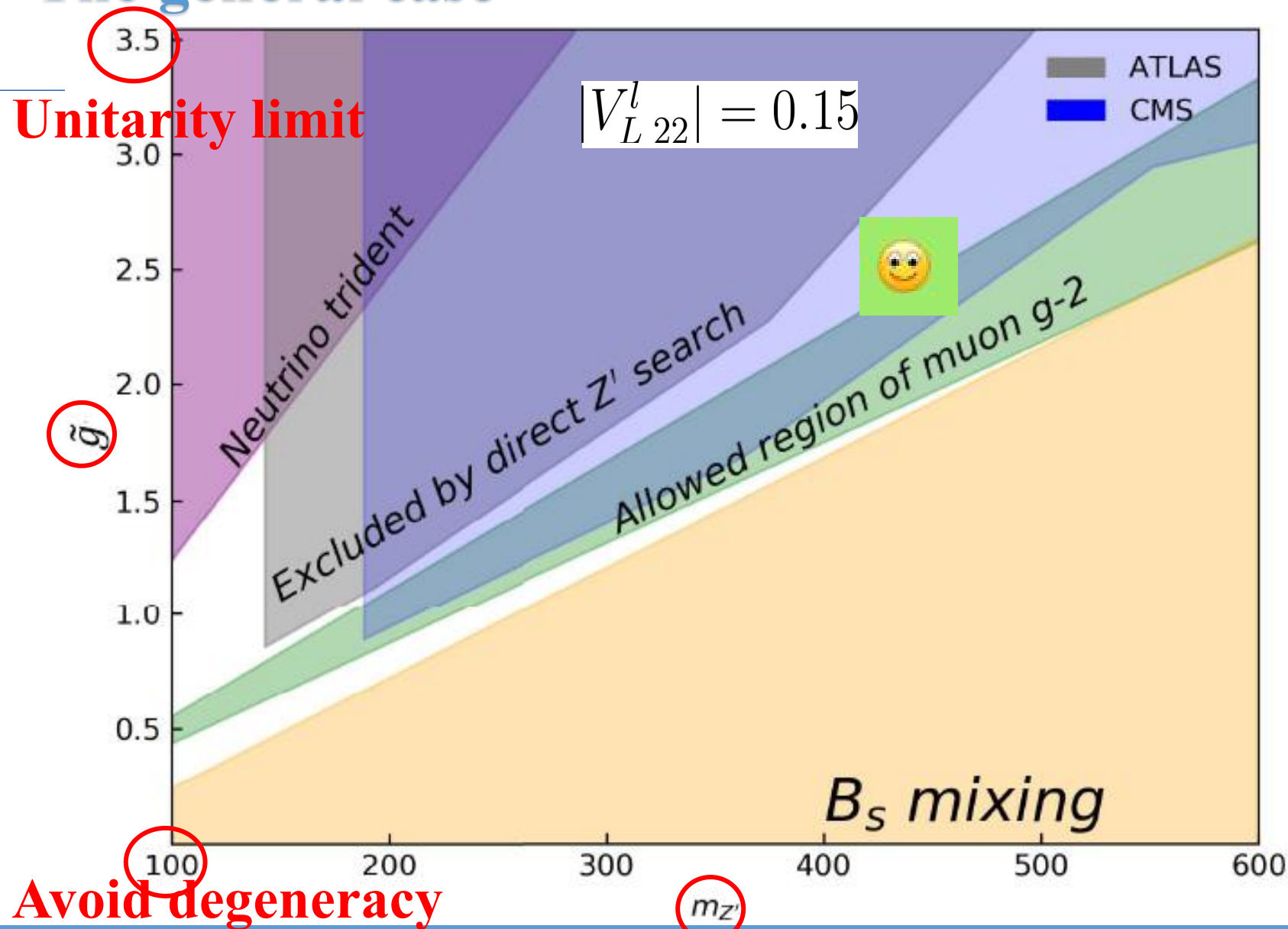
$$\frac{\tilde{g}}{m_{Z'}} = \frac{4v G_F \alpha_{em}}{g} \frac{-C_9}{|V_{L 22}^l|^2} \frac{1}{\sqrt{2S_0 \left(\frac{M_{12}}{M_{12}^{SM}} - 1 \right)}}$$

$$\begin{cases} |V_{L 22}^l| = 1 \\ |V_{L 22}^l| = 0.15 \end{cases}$$



$$\begin{aligned} \tilde{g}/m_{Z'} &\geq 1.08 \times 10^{-4} \\ \tilde{g}/m_{Z'} &\geq 4.81 \times 10^{-3} \end{aligned}$$

The general case



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U(1) kinetic mixing

U(1)_Y-U(1) kinetic mixing $(1/2)\delta B^{\mu\nu}Z'_{\mu\nu}, |\delta| < 1$

$$B = \cos \theta_W A - \sin \theta_W Z$$

Mixing of mass eigen-state

$$\begin{pmatrix} Z^m \\ Z'^m \end{pmatrix} = \begin{pmatrix} \cos \xi & \sin \xi \\ -\sin \xi & \cos \xi \end{pmatrix} \begin{pmatrix} Z \\ Z' \end{pmatrix}$$

$$\tan 2\xi = \frac{2m_Z^2 \delta \sin \theta_W / \sqrt{1-\delta^2}}{m_Z^2 - (m_Z^2 \delta^2 \sin^2 \theta_W + m_{Z'}^2) / (1-\delta^2)}$$



$$\xi \approx \frac{m_Z^2 \delta \sin \theta_W}{m_Z^2 - m_{Z'}^2}$$

Interaction terms



In the mass eigen-state

$$\begin{aligned}
 L_{int} = & J_{em}^\mu A_\mu^m \\
 & + \left(-\frac{s_\xi \delta \cos \theta_W}{\sqrt{1-\delta^2}} J_{em}^\mu + \left(c_\xi + \frac{s_\xi \delta \sin \theta_W}{\sqrt{1-\delta^2}} \right) J_Z^\mu + \frac{s_\xi}{\sqrt{1-\delta^2}} J_{Z'}^\mu \right) Z_\mu^m \\
 & + \left(-\frac{c_\xi \delta \cos \theta_W}{\sqrt{1-\delta^2}} J_{em}^\mu - \left(s_\xi - \frac{c_\xi \delta \sin \theta_W}{\sqrt{1-\delta^2}} \right) J_Z^\mu + \frac{c_\xi}{\sqrt{1-\delta^2}} J_{Z'}^\mu \right) Z_\mu^{Im}
 \end{aligned}$$

Current $\left\{ \begin{array}{l} J_{em}^\mu = -eQ^f \bar{f} \gamma^\mu f \\ J_Z^\mu = -(g/2 \cos \theta_W) \bar{f} \gamma^\mu (g_V^f - g_A^f \gamma_5) f \\ g_V^f = I_3^f - 2Q^f \sin^2 \theta_W, \quad g_A^f = I_3^f \\ J_{Z'}^\mu = -\tilde{g} (V_{L22}^l V_{L22}^{l*}) \bar{\mu} \gamma^\mu \mu \end{array} \right.$

Kinetic mixing effects



g-2

First order correction

$$\Delta a_{\mu}^{mixing}(\delta) = 0$$

Second order correction

$$\frac{\tilde{g}^2}{4\pi^2} \frac{m_{\mu}^2}{m_Z^2} \left(\frac{m_Z^2}{m_Z^2 - m_{Z'}^2} \right)^2 \delta^2 \sin^2 \theta |V_{L 22}^l|^2$$

$$\times \left[|V_{L 32}^l|^2 \left(\frac{m_{\tau}}{m_{\mu}} - \frac{2}{3} \right) + \frac{1}{3} |V_{L 22}^l|^2 \right]$$



Neglect

b anomaly

$$\bar{s}bZ \rightarrow Z - Z' \rightarrow Z' \mu \bar{\mu}$$

$$C_9^{loop} = 2 \cos \theta_W \frac{\tilde{g}(V_{L 22}^l V_{L 22}^{l*})}{e}$$

$$\times \frac{m_Z^2 \delta}{m_Z^2 - m_{Z'}^2} C_0(m_t^2/m_W^2)$$

$$\begin{cases} m_{Z'} \in [100, 600] \text{ GeV} \\ |V_{L 22}^l| = 0.15 \end{cases}$$



$$C_9^{loop} = [-1.74\delta, -0.05\delta]$$

Conclusions



- 1. We construct a flavor specific gauge model to explain muon $g-2$ and b anomalies simultaneously.**
- 2. The original Z' - q_2/l_2 interaction in weak basis can develop mixing in mass eigen-state basis.**
- 3. The introduced substantial mu-tau mixing induces the suppressed factor V_{L22}^l to make other processes satisfied so that some viable parameter space emerges.**
- 4. The model can be tested at a high energy muon collider.**



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Thanks!



Yukawa coefficients

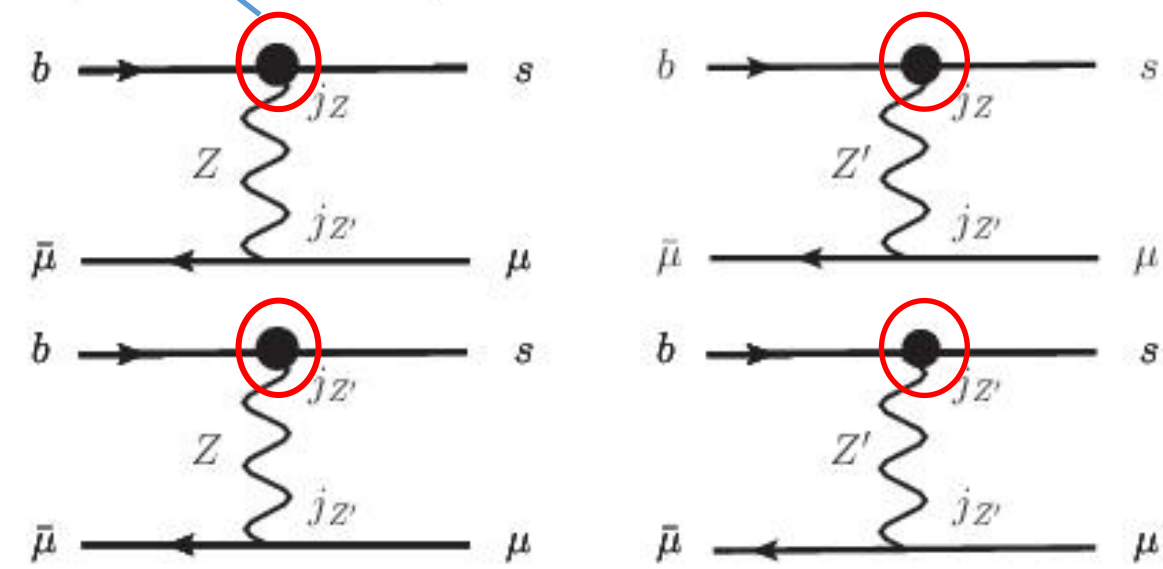
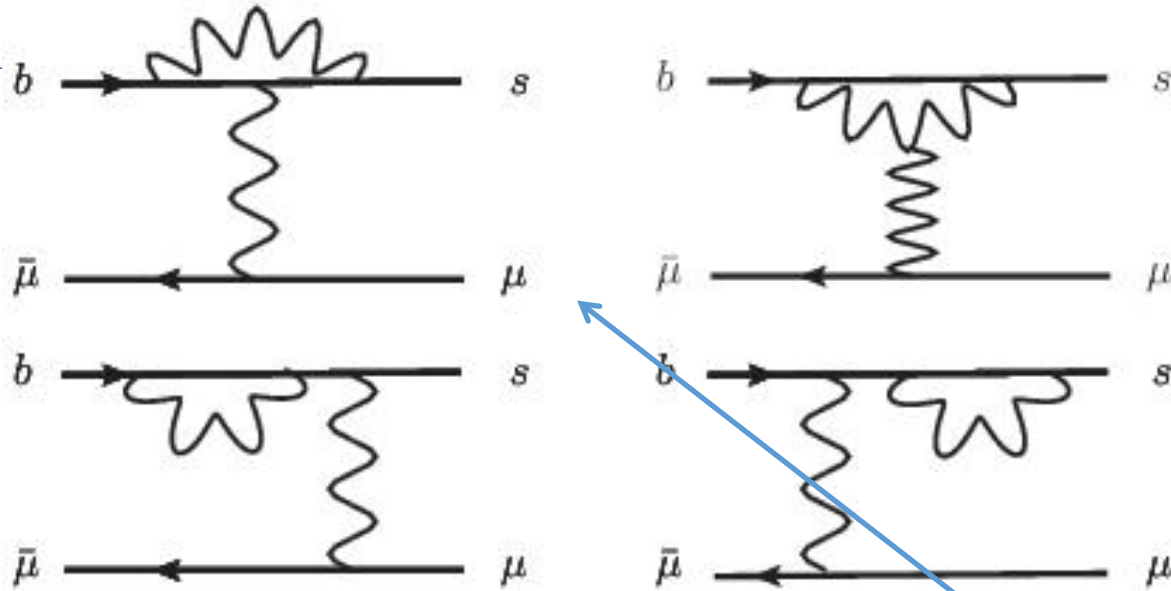


$$Y_{H_1}^u = \begin{pmatrix} 0 & Y_{H_1}^{u12} & 0 \\ 0 & 0 & 0 \\ 0 & Y_{H_1}^{u32} & 0 \end{pmatrix}, \quad Y_{H_1}^d = \begin{pmatrix} 0 & 0 & 0 \\ Y_{H_1}^{d21} & 0 & Y_{H_1}^{d23} \\ 0 & 0 & 0 \end{pmatrix}$$

$$Y_{H_2}^u = \begin{pmatrix} 0 & 0 & 0 \\ Y_{H_2}^{u21} & 0 & Y_{H_2}^{u23} \\ 0 & 0 & 0 \end{pmatrix}, \quad Y_{H_2}^d = \begin{pmatrix} 0 & Y_{H_2}^{d12} & 0 \\ 0 & 0 & 0 \\ 0 & Y_{H_2}^{d32} & 0 \end{pmatrix}$$

$$Y_{S_1} = \begin{pmatrix} 0 & Y_{S_1}^{12} & 0 \\ Y_{S_1}^{12} & 0 & Y_{S_1}^{23} \\ 0 & Y_{S_1}^{23} & 0 \end{pmatrix}, \quad Y_{S_2} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & Y_{S_2}^{22} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

b anomaly correction by kinetic mixing



$$C_0(x) = \frac{x}{8} \times \left(\frac{6-x}{1-x} + \frac{3x+2}{(1-x)^2} \ln x \right)$$