



IR-Safe Observables Extended to Track-Based Measurement

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Based on:

YL, Ian Moult, Solange Schrijnder van Velzen, Wouter Waalewijn, HuaXing Zhu
arXiv:2107.xxxxx

Motivation from experiments

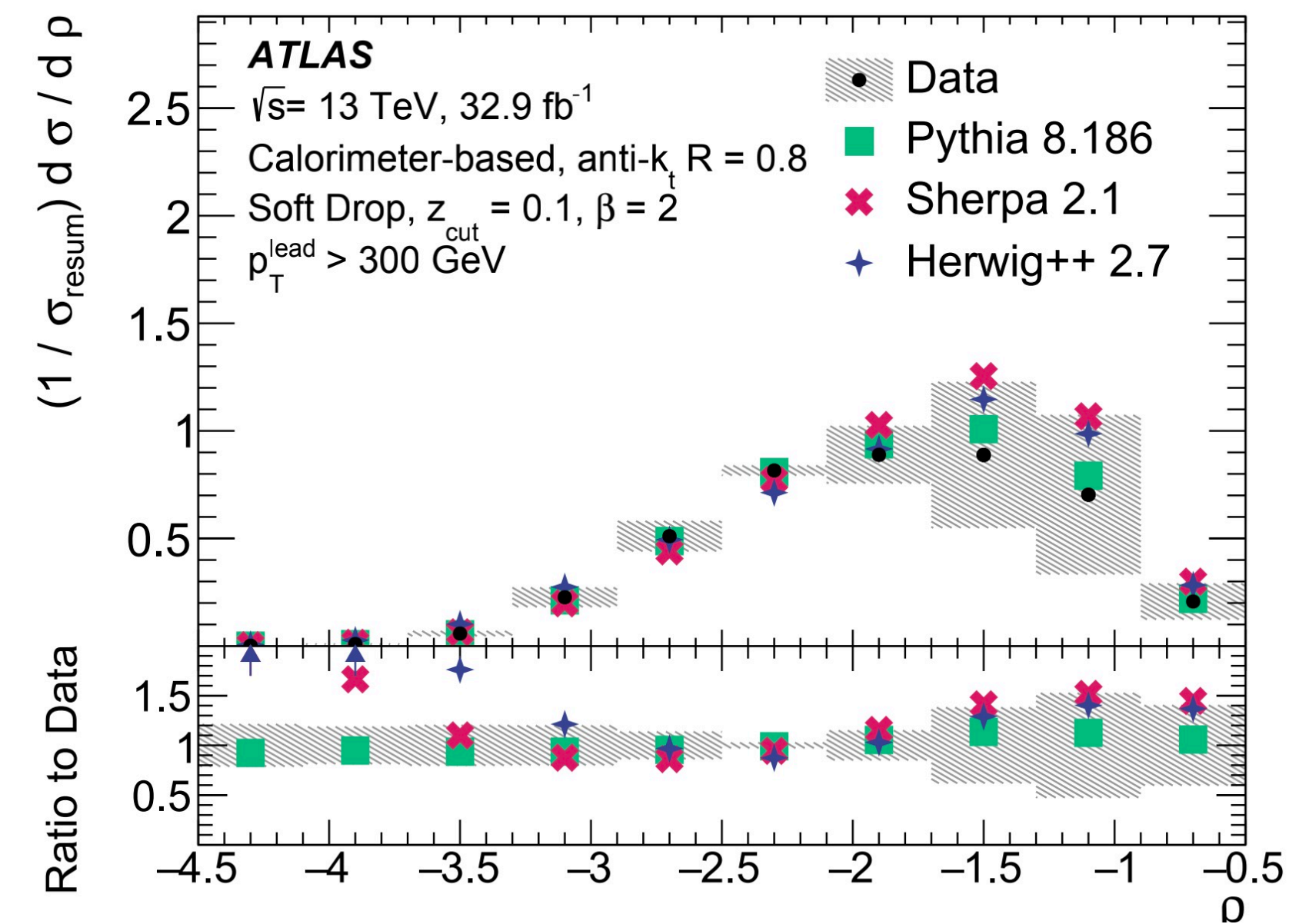
Track-based measurements:

- Superior angular resolution
- To mitigate pileup
- One problem: Track-based calculations are **not** IR safe in perturbative theory.

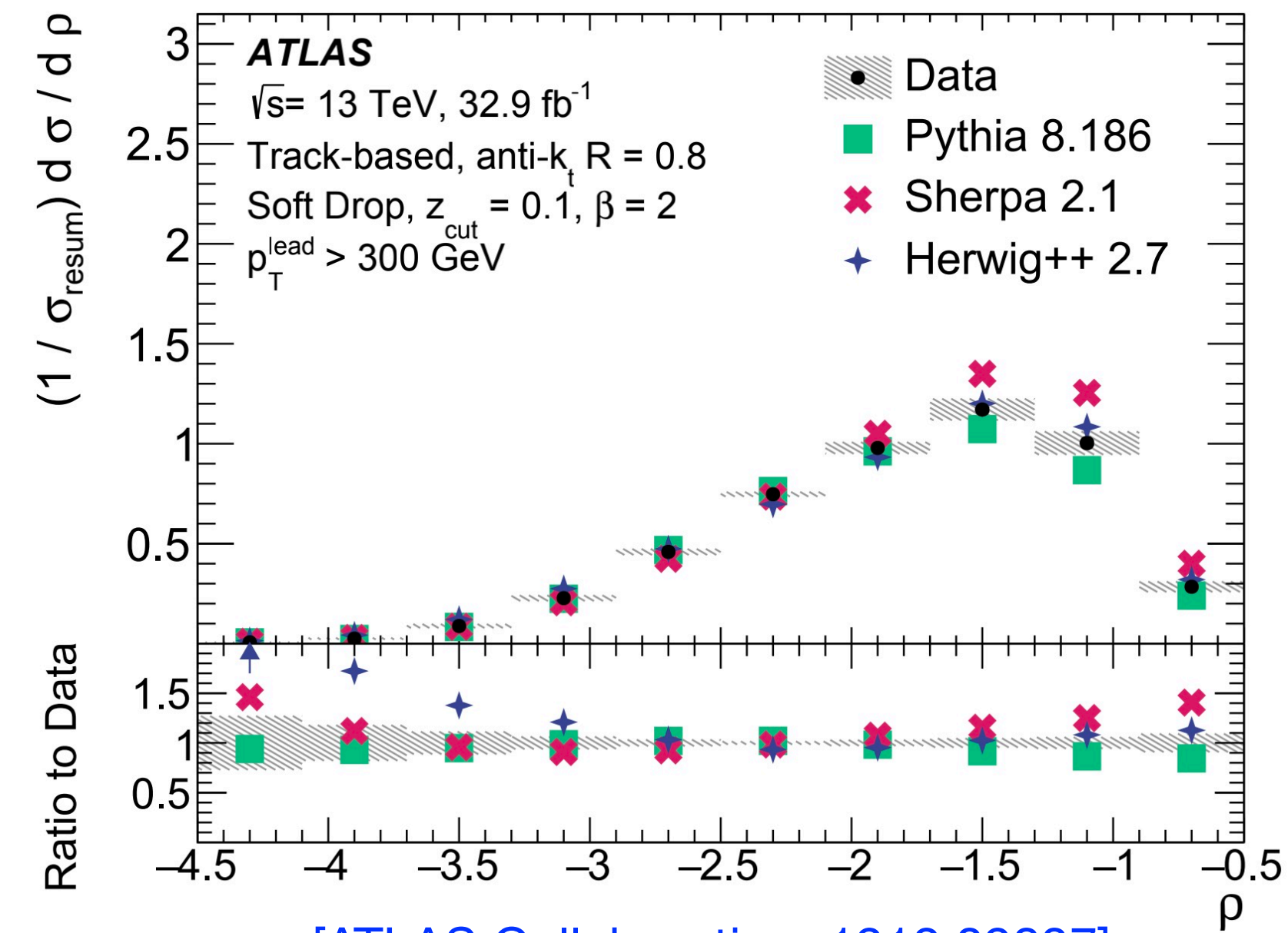
Track Functions [H. Chang, M. Procura, J. Thaler, W. Waalewijn, arXiv: 1303.6637, 1306.6630]

- ▶ IR divergences are absorbed into track functions.

calorimeter-based
(all-particle)



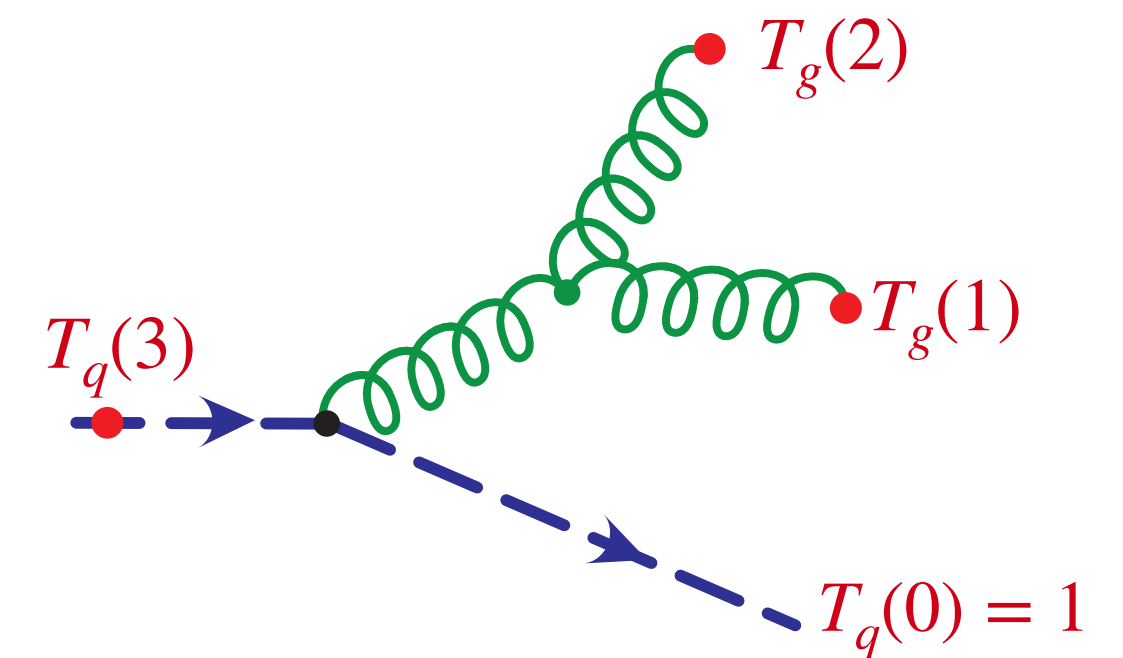
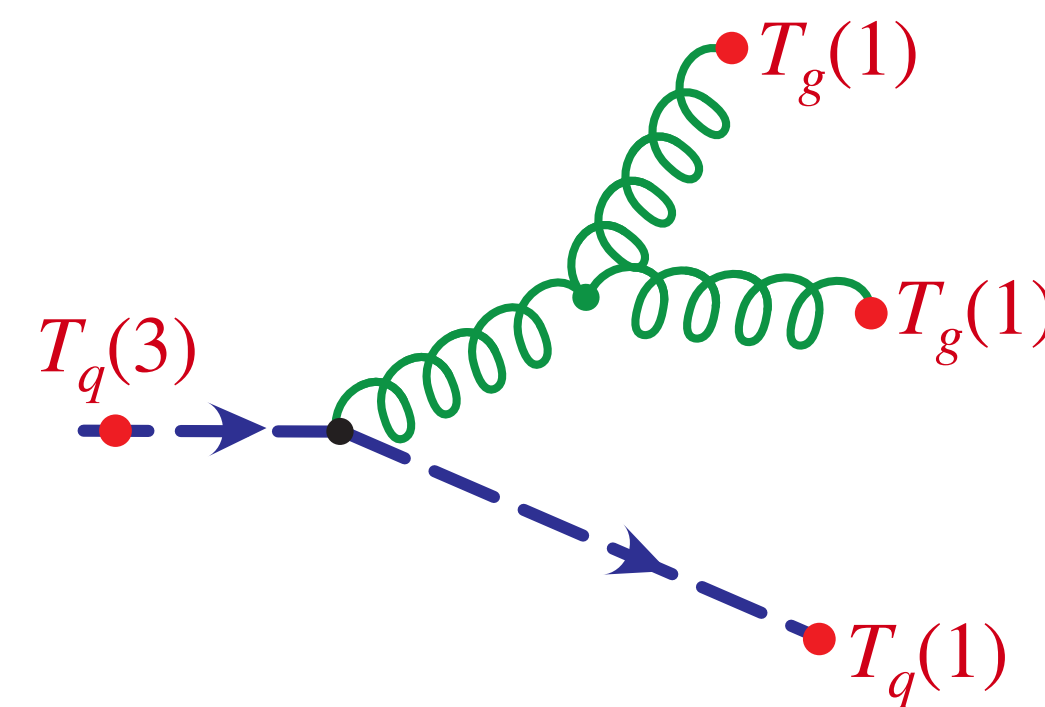
track-based
(charged-particle)



[ATLAS Collaboration, 1912.09837]

Outline

- Introduction
 - Definition of track functions
 - Incorporating tracking information
- Track Function Evolution at $\mathcal{O}(\alpha_s^2)$
 - In Mellin Space
- An Example: Prediction for Track EECs
 - Jet Substructure



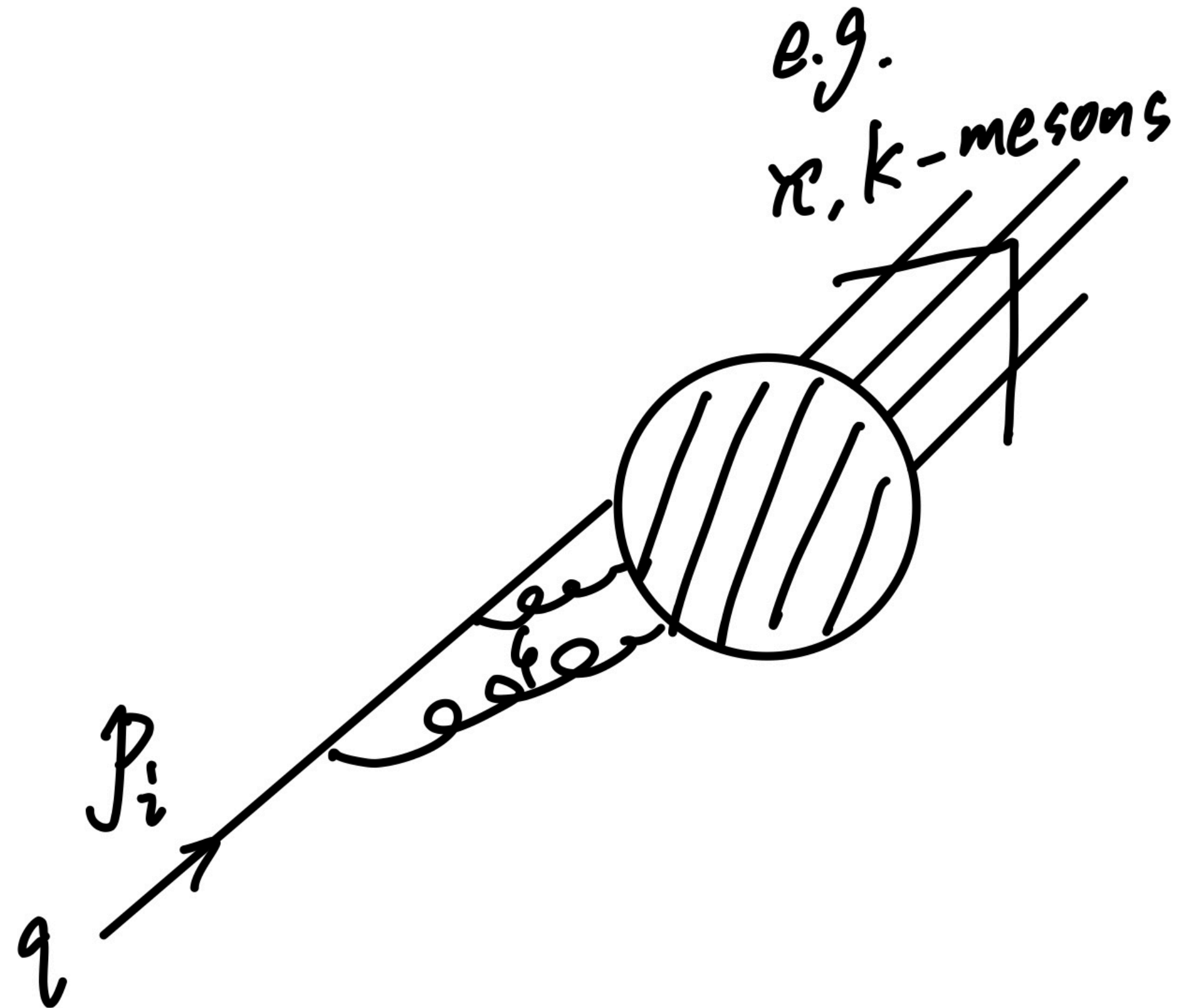
Track Functions $T_i(x, \mu)$

Definition

- The track function $T_i(x, \mu)$ describes the total momentum fraction x of **all** charged particles (tracks) in a jet initiated by a hard parton i (quark/antiquark or gluon).

$$\bar{p}_i^\mu = xp_i^\mu + O(\Lambda_{\text{QCD}}), \quad (0 \leq x \leq 1).$$

A quark with 4-momentum p_i^μ hadronizes into tracks (charged particles) with total 4-momentum \bar{p}_i^μ :



Track Functions

Features

- A generalization of the fragmentation function (FF).
 - Independent of hard process.
 - Fundamentally non-perturbative, with a calculable scale (μ) dependence.

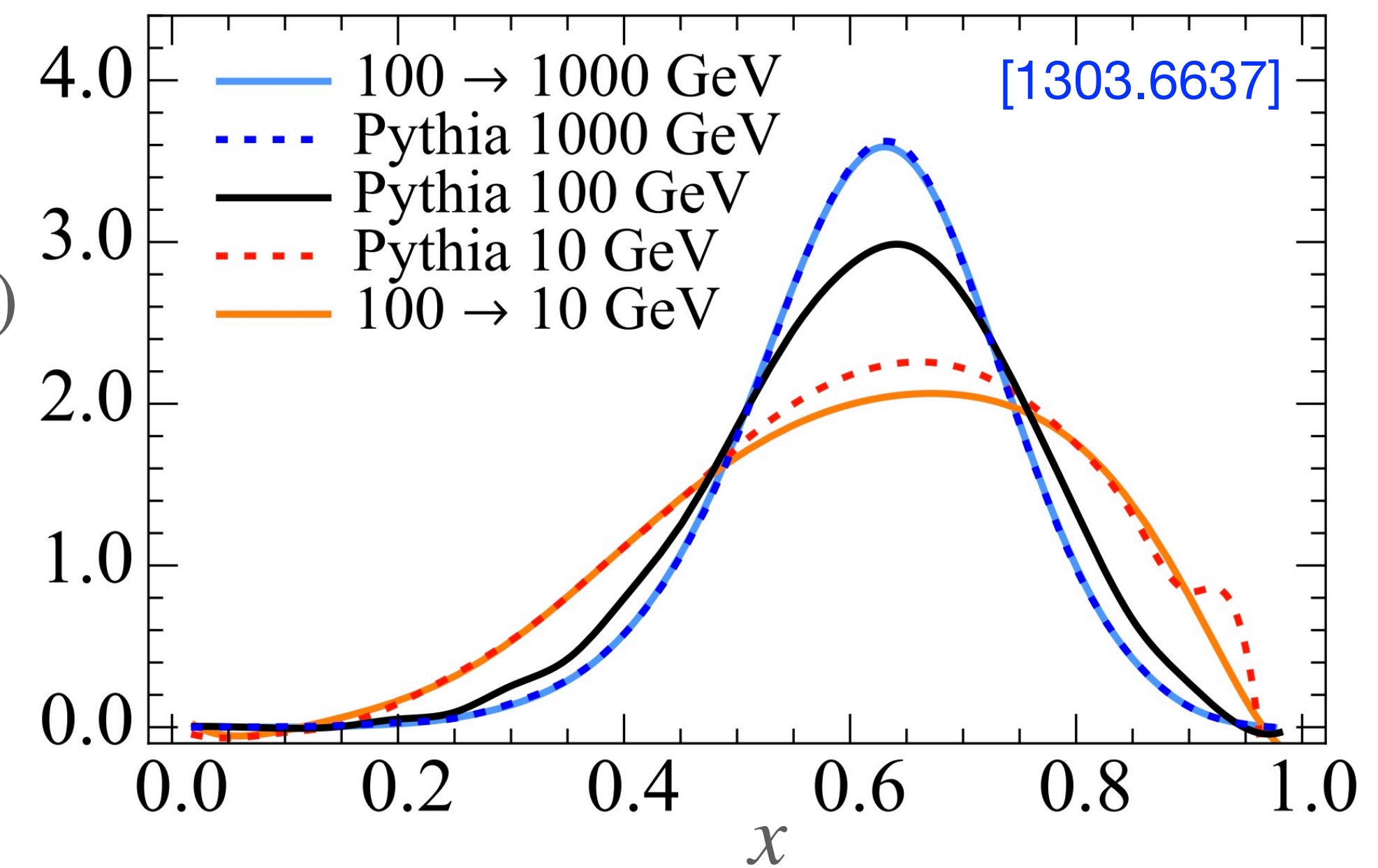
- Sum rule: $\int_0^1 dx T_i(x, \mu) = 1$.

✓ Track functions has already been studied at $\mathcal{O}(\alpha_s)$.

[H. Chang, M. Procura, J. Thaler, W. Waalewijn, arXiv: 1303.6637, 1306.6630]

✓ New in this talk: The track function formalism applies to IR-safe measurements at $\mathcal{O}(\alpha_s^2)$.

$T_g(x, \mu)$



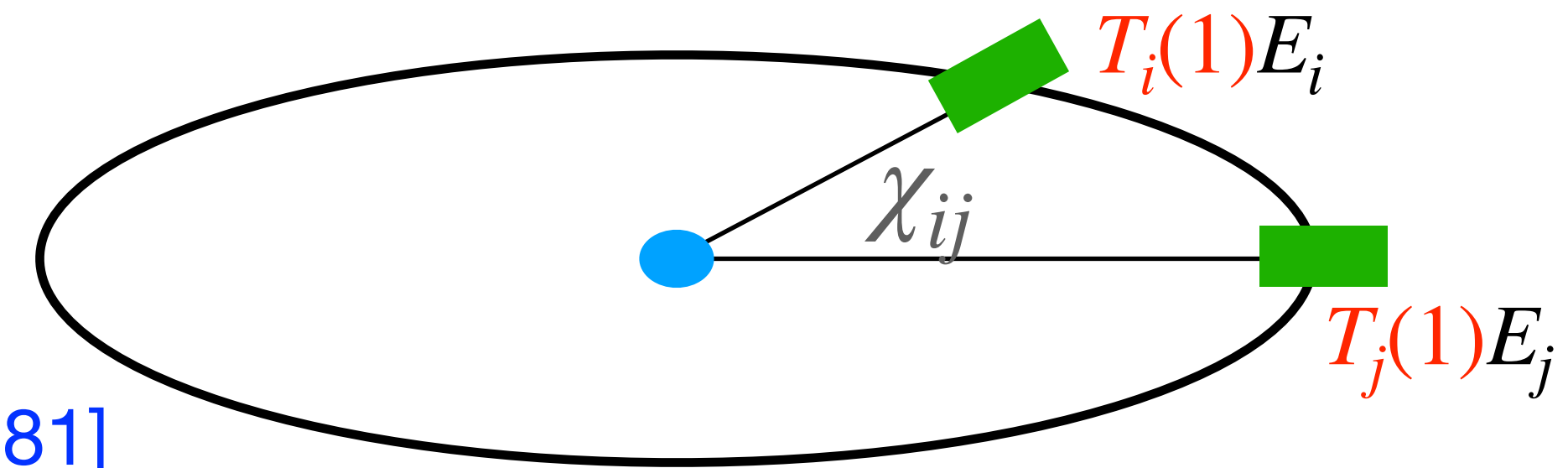
Incorporating Tracks

[1303.6637]

[Chen, Mout, Zhang, Zhu, 2004.11381]

For a δ -function type observable e measured using partons:

$$\frac{d\sigma}{de} = \sum_N \int d\Pi_N \frac{d\sigma_N}{d\Pi_N} \delta \left[e - \hat{e}(p_i^\mu) \right]$$



For an energy correlator at partonic level: e.g. 2-point correlator (EEC)

$$\frac{d\Sigma}{dz} = \sum_{i,j} \int \frac{E_i E_j}{Q^2} \delta \left(z - \frac{1 - \cos \chi_{ij}}{2} \right) d\sigma$$

..... when measured on tracks

Simply

$$E_i^n \rightarrow \int dx_i T_i(x_i) x_i^n E_i^n = T_i(n) E_i^n$$

$$\frac{d\sigma}{d\bar{e}} = \sum_N \int d\Pi_N \frac{d\bar{\sigma}_N}{d\Pi_N} \int \prod_{i=1}^N dx_i T_i(x_i) \delta \left[\bar{e} - \hat{e}(x_i p_i^\mu) \right]$$

full functional form involved

only requiring moments entered as weights

$$\left(\frac{d\Sigma}{dz} \right)_{\text{tr}} = \sum_{i,j} T_i(1) T_j(1) \int \frac{E_i E_j}{Q^2} \delta \left(z - \frac{1 - \cos \chi_{ij}}{2} \right) d\bar{\sigma}$$

Track EEC

- ▶ The (projected) energy correlator is a natural observable to incorporate tracking information conveniently.

Track Function Evolution

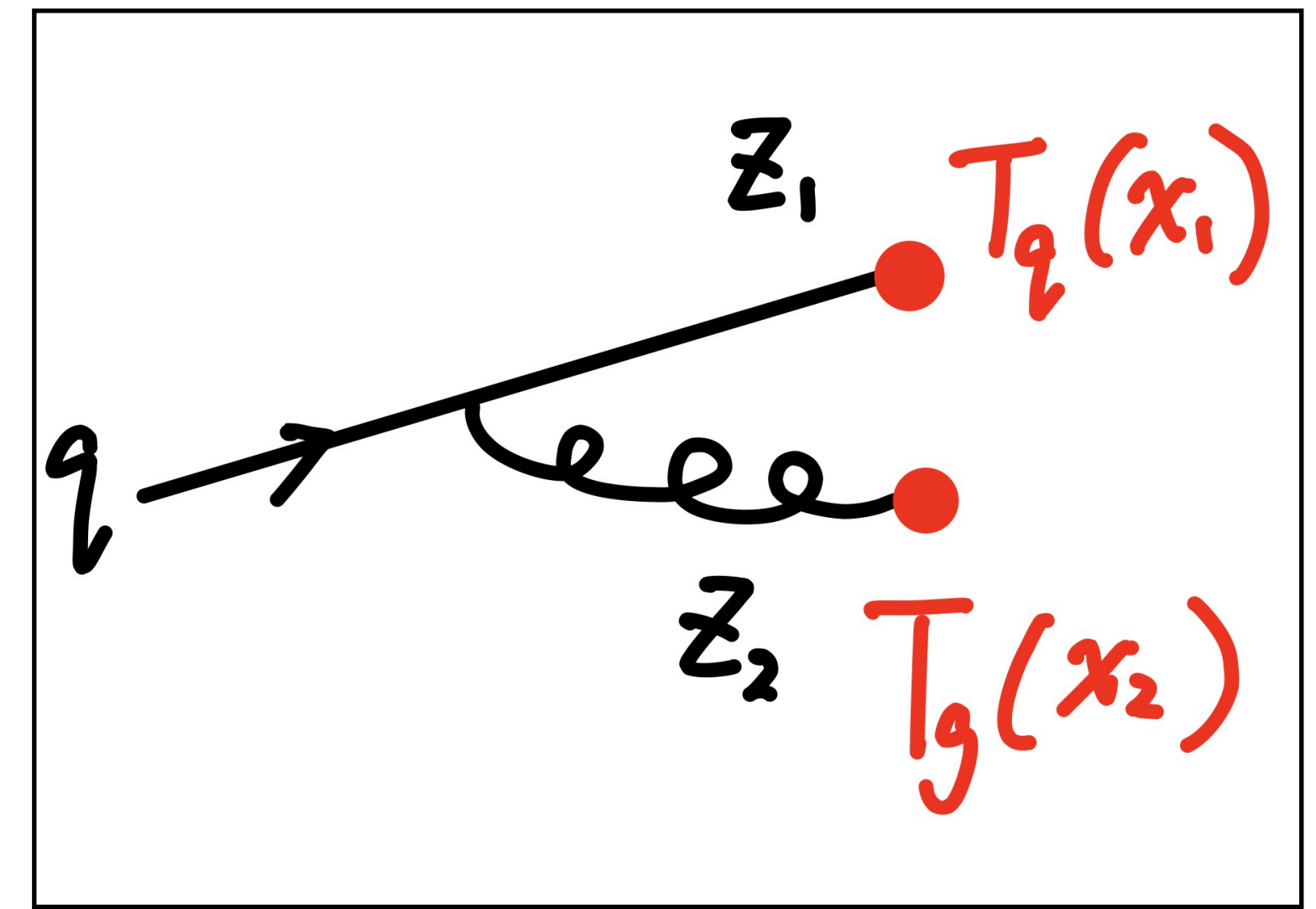
- LO evolution [1303.6637]

$$\frac{d}{d \ln \mu^2} T_i(x, \mu) = a_s(\mu) \sum_{j,k} \int dz dx_1 dx_2 P_{i \rightarrow jk}^{(0)}(z_1, z_2) \delta(1 - z_1 - z_2) \times T_j(x_1, \mu) T_k(x_2, \mu) \delta[x - z_1 x_1 - z_2 x_2] .$$

Nonlinear, involving contributions from both branches of the splitting.

- Beyond leading order: Involving contributions from multiple branchings.
- Of shift symmetry: $x \rightarrow x + a$

$$\frac{d}{d \ln \mu^2} T_i(x + a) = \sum_X \int \left(\prod_m dx_m dz_m T_{i_m}(x_m + a) \right) P_{i \rightarrow i_1 \dots i_m \dots}(\{z_m\}) \delta \left(1 - \sum_m z_m \right) \delta \left(x - \sum_m x_m z_m \right)$$



While for fragmentation functions:

$$\frac{d}{d \ln \mu^2} d_{h/i}(z, \mu) = \sum_j d_{h/lj} \otimes P_{ji}^T(z, \mu)$$

- One of the branches; not depending on correlations between final-state hadrons. \rightarrow Linearity
- Scale invariant $d(y) \rightarrow d(ay)$.

Track Function Evolution

- Two independent approaches to extracting the evolution at $\mathcal{O}(\alpha_s^2)$:

- Jet Functions

←.....→
Agree

- Directly calculating track jet functions. [Ritzmann, Waalewijn, 1407.3272]
- Remaining IR poles cancel when matching onto track functions.

$$\mathcal{G}_i^{(2)} = T_i^{(2)} + \sum J_{i \rightarrow jk}^{(1)} \otimes [T_j^{(1)} T_k^{(0)}] + \sum_{j,k} J_{i \rightarrow jkl}^{(2)} \otimes [T_j^{(0)} T_k^{(0)} T_l^{(0)}]$$

- Taking moments to extract the evolution for $T(n, \mu)$ (avoiding plus distribution with multiple variables).

- (Projected) Energy Correlators $E_i^n \rightarrow T_i(n) E_i^n$

- n -point correlators on tracks involves no higher moments than $T(n)$.

- n -point correlator $\xrightarrow[\text{Pole}]{\text{Cancellation}} z \rightarrow 0$

Evolution for $T(n, \mu)$

- The evolution for the lower moments can be checked at wide-angle $0 < z \leq 1$ region.

checks on the track function formalism

➔ Mellin Space

Why we study moments of track functions?

non-linear equations for the
x-space track functions

- For a δ -function type IR-safe observable, [\[2004.11381\]](#)

$$\frac{d\sigma}{de} = \int d^4x e^{iq \cdot x} \langle 0 | O(x) \delta(e - \hat{e}) O^\dagger(0) | 0 \rangle$$

- **Infinite** number of correlators
→ an **infinite** number of $T(n)$'s
- Complicated dependence on $T(x)$
- Its moments → weighted cross sections.
- For (projected) n -point energy correlators, tracking information can be easily incorporated with a **finite** number of the moments.

- **Linear** RG equations for $\mathbf{T} =$

$$\{T_i(n), T_{i_1}(k)T_{i_2}(n-k), \dots, T_{i_1}(1)\dots T_{i_n}(1)\}^t$$

▶ Matrix form: $\frac{d}{d \ln \mu^2} \mathbf{T} = \mathbb{R} \mathbf{T}$

▶ \mathbb{R} : related to moments of timelike splitting functions.

- Simpler to do resummation
- Give information on full-x evolution equations.

➔ Mellin Space

Evolution for moments of track functions

- For the first moment: the same as that of FFs up to all orders in perturbation theory:

$$\frac{d}{d \ln \mu^2} T_g(1) = -\gamma_{gg}(2) T_g(1) + \sum_q \left(-2\gamma_{qg}(2) \right) T_q(1)$$

$$\frac{d}{d \ln \mu^2} T_q(1) = -\gamma_{gq}(2) T_g(1) + \left(-\gamma_{qq}(2) - \gamma_{\bar{q}q}(2) \right) T_q(1) + \sum_{q' \neq q} \left(-\gamma_{q'q}(2) - \gamma_{\bar{q}'q}(2) \right) T_{q'}(1)$$

For Moments of Fragmentation Functions:

$$\frac{d}{d \ln \mu^2} d_{h/i}(n) = - \sum_j d_{h/j}(n) \gamma_{ji}^T(n+1)$$

- Higher moments involve products of ≥ 2 track functions

$$\frac{d}{d \ln \mu^2} T_i(n) \supset a_s^N T_{i_1}(1) T_{i_2}(1) \cdots T_{i_n}(1) \text{ for } n < N + 1$$

- The expressions for the evolution can be simplified in terms of **shift-invariant** objects (named central moments).

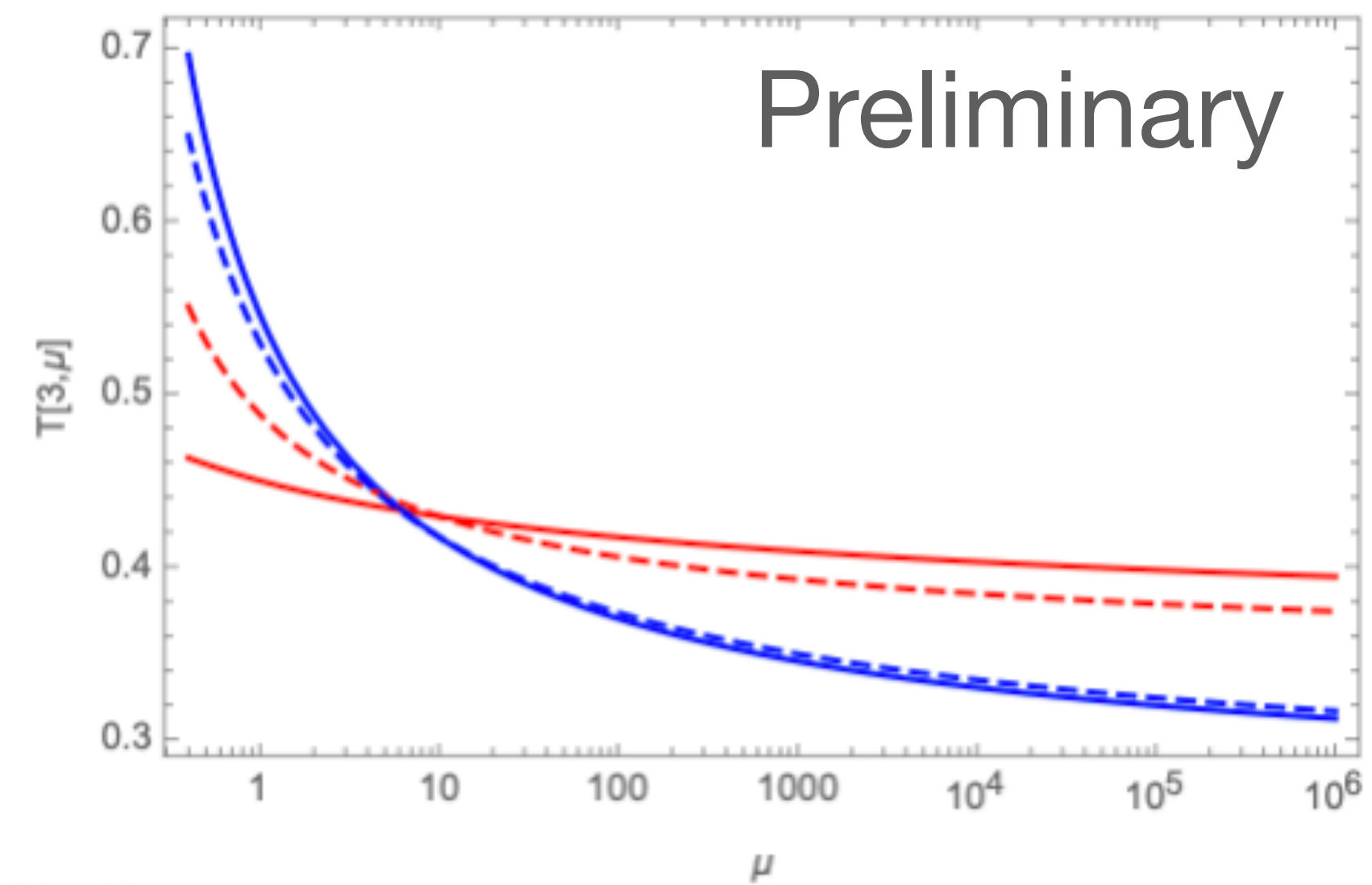
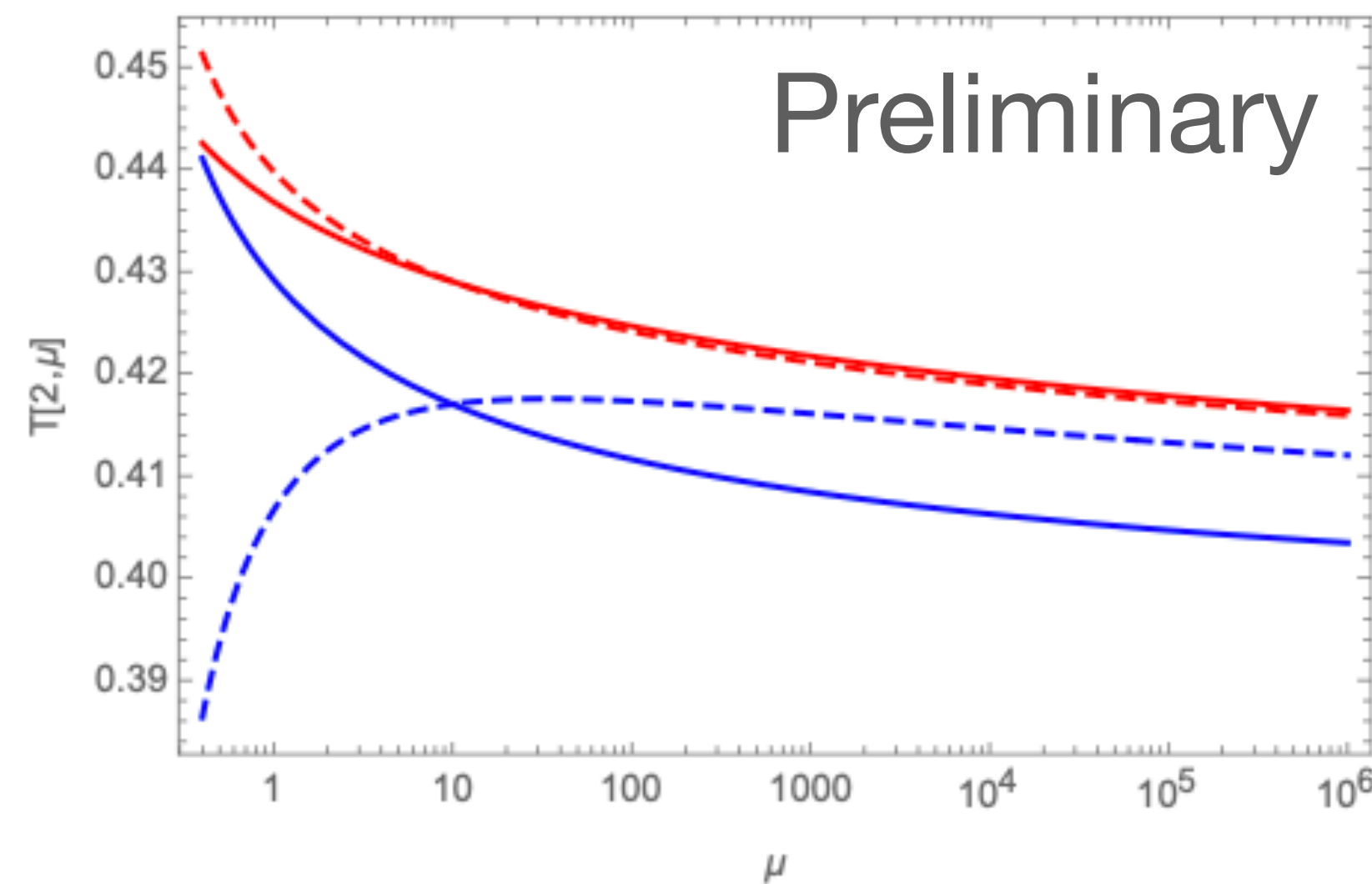
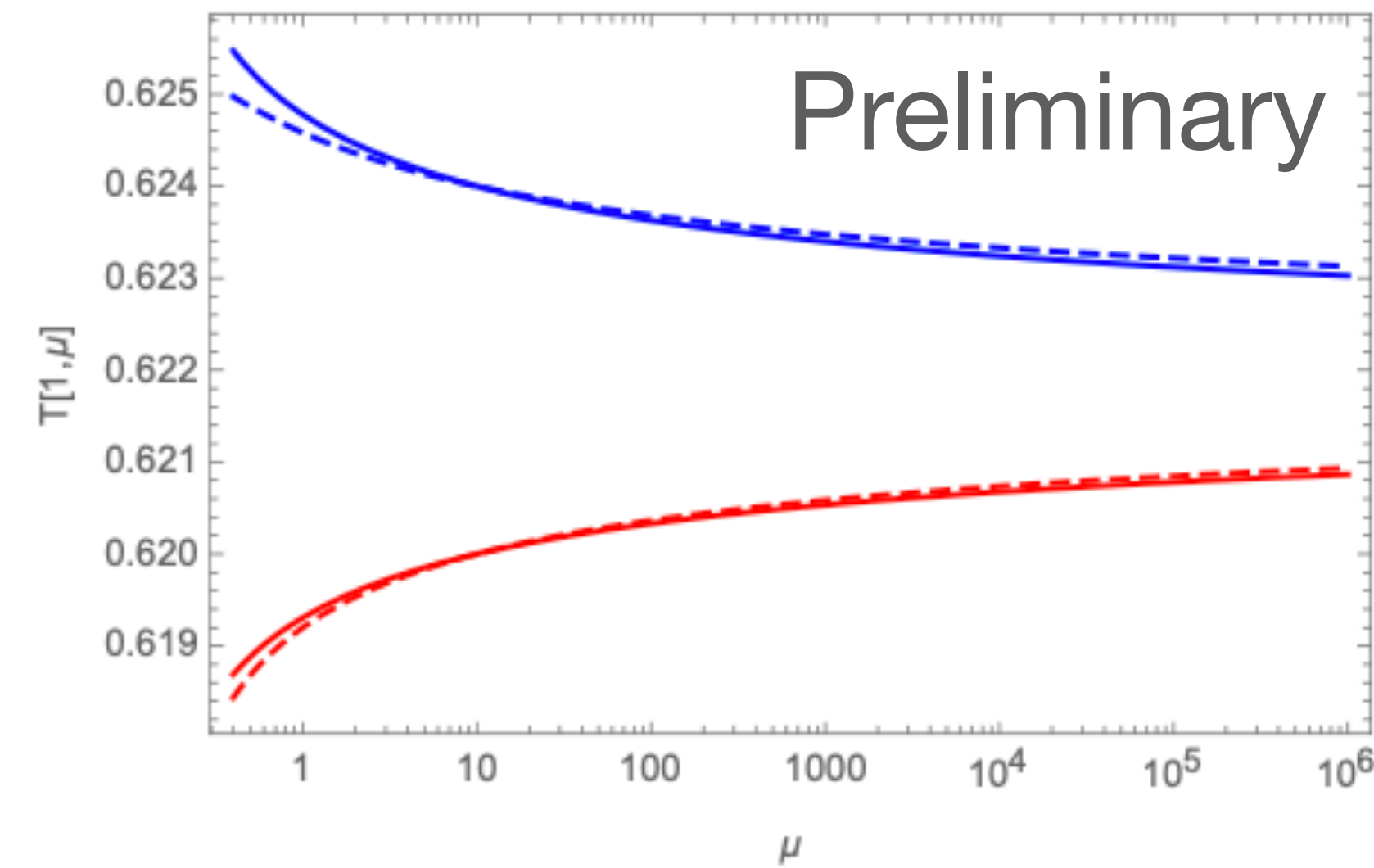
➔ Mellin Space

$\mathcal{O}(\alpha_s^2)$ evolution for moments of track functions

The evolution of $T[1]$

The evolution of $T[2]$

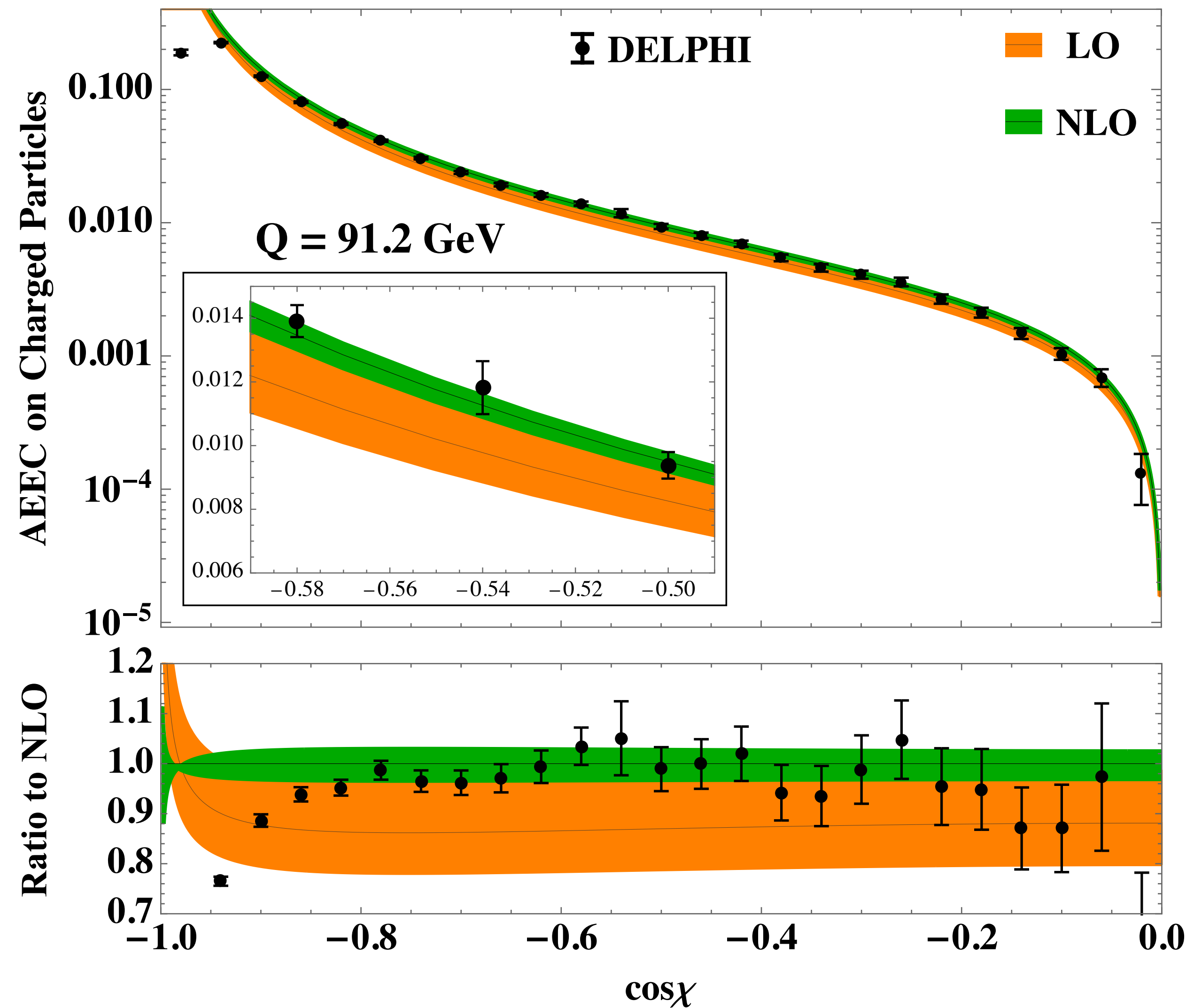
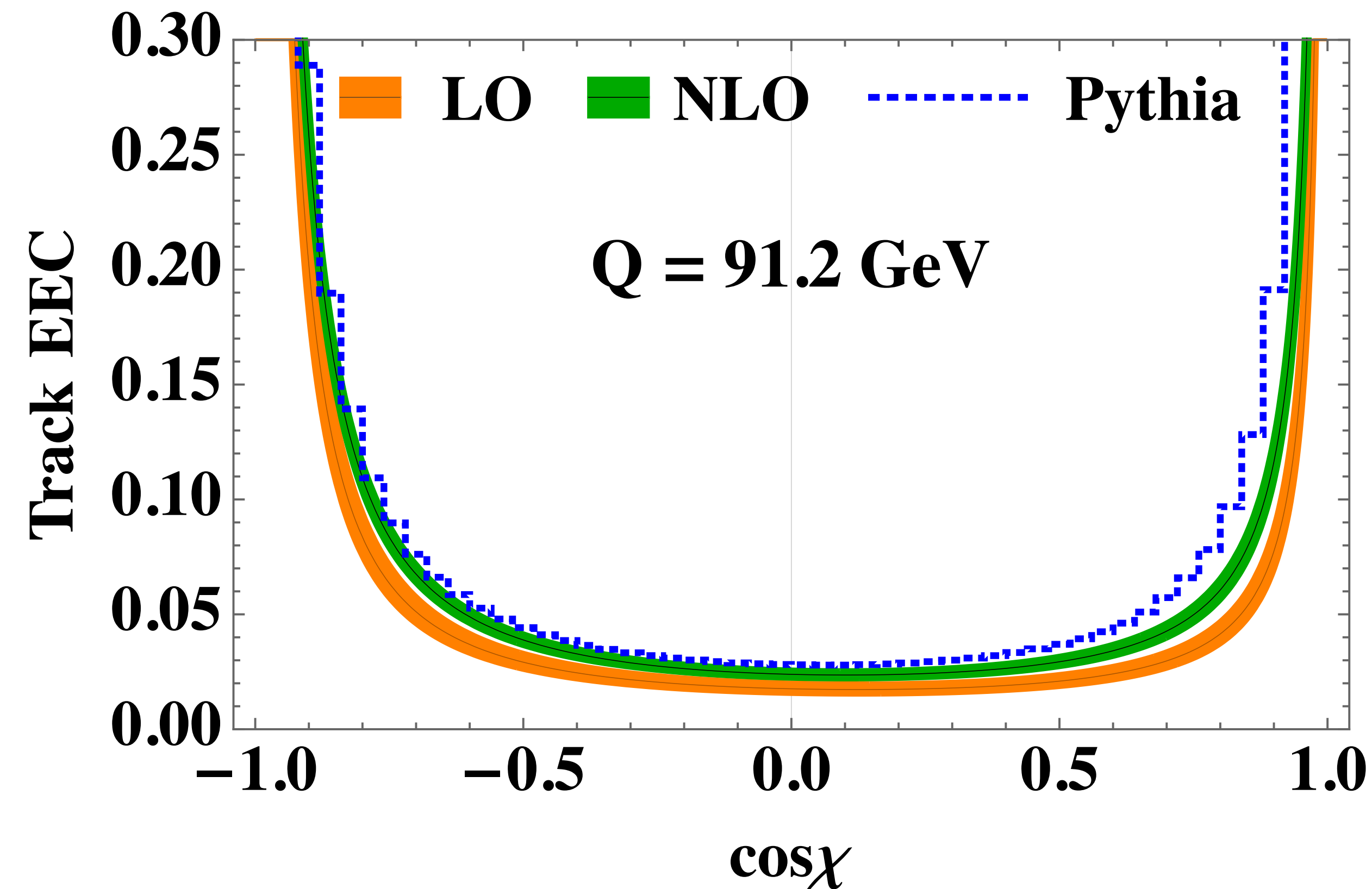
The evolution of $T[3]$



— $Tq[n]$ up to α_s - - - $Tq[n]$ up to α_s^2
— $Tg[n]$ up to α_s - - - $Tg[n]$ up to α_s^2

Prediction for Track EECs

- Here, $\text{EEC} = \frac{1}{\sigma} \frac{d\Sigma}{d\cos\chi}$
- $\text{AEEC}(\cos\chi) = \text{EEC}(\cos\chi) - \text{EEC}(-\cos\chi), \cos\chi \leq 0$



Jet Substructure

In the collinear limit:

[2004.11381]

- The energy correlator is a jet substructure observable.

- Jet function constants (jet functions with the logarithmic dependence excluded):

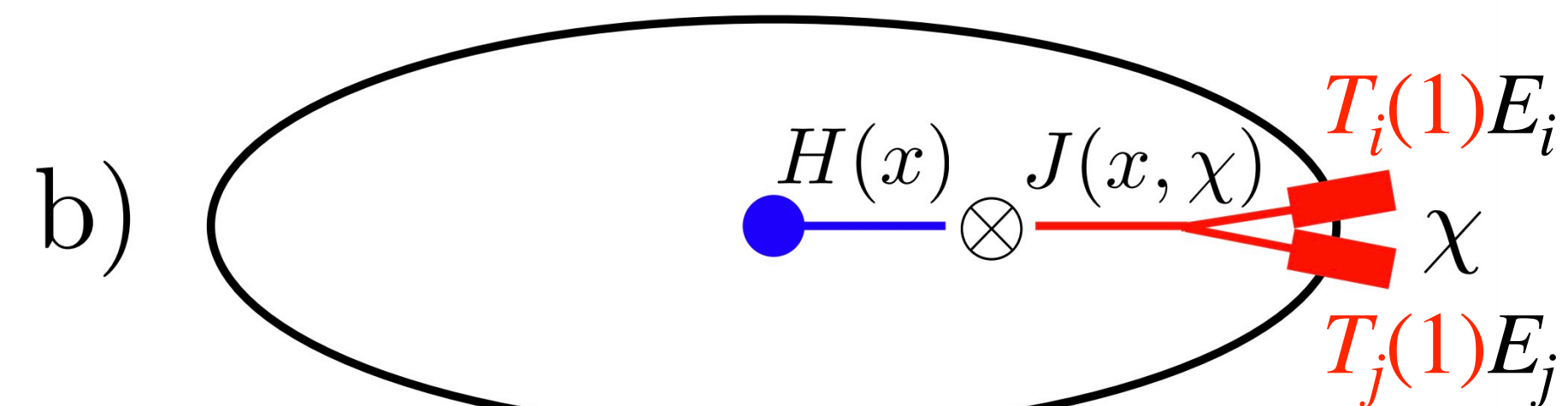
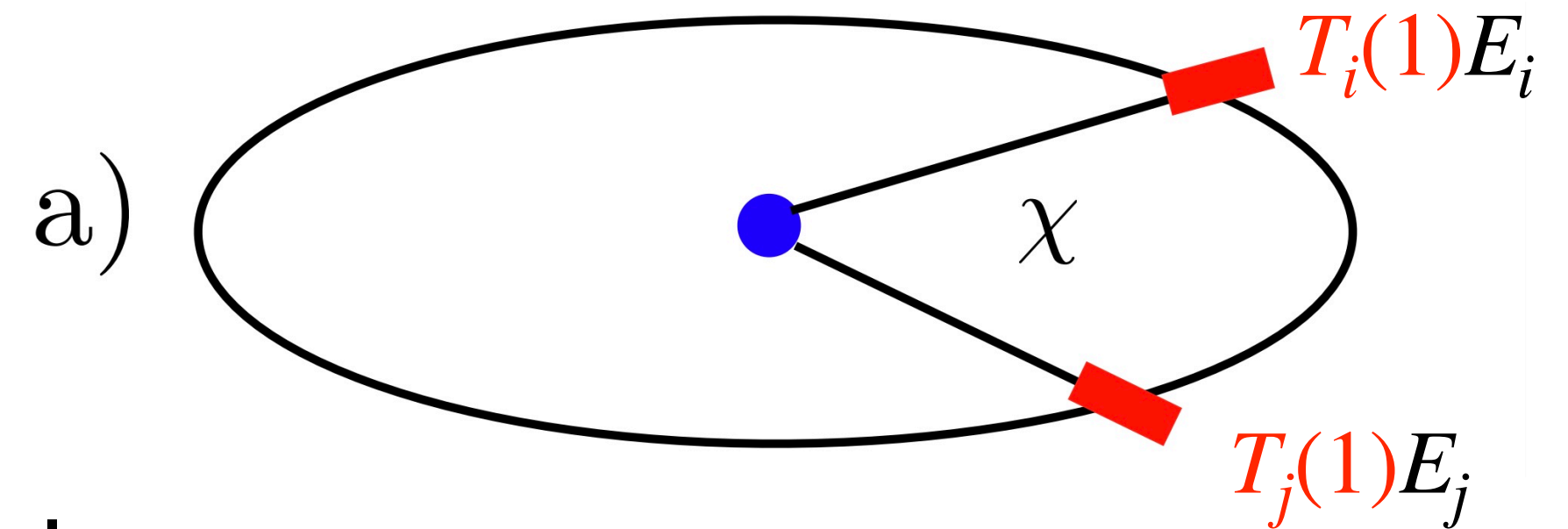
- The moments $T_i(n)$ appear as the coefficients.

- e.g. for EECs, up to $\mathcal{O}(\alpha_s^2)$

$$j_2^g = \frac{1}{4} T_g(2) + a_s \left\{ T_g(1) T_g(1) C_A \left(-\frac{449}{150} \right) + \sum_q T_q(1) T_{\bar{q}}(1) T_F \left(-\frac{7}{25} \right) \right\}$$

$$+ a_s^2 \left\{ T_g(1) T_g(1) \left\{ C_A^2 \left(-\frac{527\zeta(3)}{10} + \frac{133639871}{3240000} - \frac{2159\pi^2}{1800} + \frac{19\pi^4}{90} \right) + C_A n_f T_F \frac{139}{270} \right\} + \sum_q T_q(1) T_{\bar{q}}(1) \dots \right\}$$

- Resummation



Summary & Outlook

- Track functions offer a QFT approach to calculating track-based observables.
 - Superior angular resolution.
 - Reduced effect of pile-up.
 - Track function formalism studied beyond leading order:
 - Evolution for moments of track functions at $\mathcal{O}(\alpha_s^2)$.
 - Two-loop track EEC.
 - The (project) energy correlators interface in a simple manner with tracking information through the moments, allowing for high order calculations.
 - This formalism allows IR-safe observables to be computed on any subset of final-state hadrons specified by some particular quantum numbers.
- ➡ More: Generalized fragmentation functions (GFFs)

- High precision
- Strong check on formalism

[B. Elder, M. Procura, J. Thaler, W. Waalewijn, K. Zhou, 1704.05456]

Thanks!