

Massive Vector Boson Sterile Neutrino Decay

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Sterile Neutrino Decay

- $m_N \gg 200\text{GeV}$, Before EWSB. $N \leftrightarrow h\nu/h^\pm l^\mp$.
- $m_N \ll 50\text{GeV}$, After EWSB, however low temperature. Thermal effects can be neglected!!
- Gap: $m_N \sim 100\text{GeV}$.

Sterile Neutrino Decay

- J.Ghiglieri, M.Laine, “Full” thermal calculations of sterile neutrino production in the early universe. A series of paper.
- $\mathcal{K}^2 = 0$, only for small m_N .
- T.Hambye, D.Teresi, an approximation. When $T < T_c$, replace W_L/Z_L with $G^{0,\pm}$ of the same mass as the dressed H . (arXiv:1606.00017)

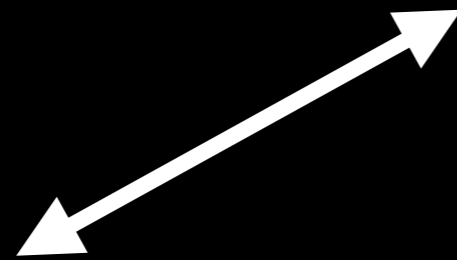
Sterile Neutrino Decay

- What about $m_N \sim 100\text{GeV}$?
- Resolve the massive vector boson's thermal emission at first!

Massive Vector Boson

- $\frac{i}{k^2 - m^2} \left[-g^{\mu\nu} + \frac{k^\mu k^\nu}{k^2 - \xi m^2} (1 - \xi) \right]$

- $\frac{i}{k^2 - \xi m^2}$



**Precise cancellation
of these two poles at
zero temperature!**

Goldstone Equivalence Gauge

- Thanks Junmou Chen.
- Physical Gauge: extend the “polarization vector” with one extra Goldstone degree of freedom:
- $\epsilon_{\pm,L}^{\mu} \rightarrow \epsilon_{\pm,L}^M, M = 0,1,2,3,4.$
- ϵ_{\pm}^M does not change.

Goldstone Equivalence Gauge

$$\epsilon_{Lout}^W(p_2) = \begin{pmatrix} -\frac{\sqrt{p_2^2}}{n_2 \cdot p_2} n_2^\mu \\ -i \frac{m_W(T)}{\sqrt{p_2^2}} \end{pmatrix},$$

$$p_2^\mu \mathcal{M}_{V\mu} = im_V \mathcal{M}^{GS}$$

$$\epsilon_{L,R\xi}^W = \epsilon_{Lin}^W + \begin{pmatrix} \frac{p_2^\mu}{\sqrt{p_2^2}} \\ -i \frac{m_W(T)}{\sqrt{p_2^2}} \end{pmatrix},$$

Advantage: All the Goldstone contributions are separated out!

Finite Temperature

$$D_0^{\text{full},MN}(k) = \frac{i}{k^2 - m_A^2 - \Pi_T(k) + i\epsilon} P_T + \frac{i}{k^2 - m_A^2 - \Pi_L(k) + i\epsilon} P_L \\ + \frac{1}{1 - \frac{\Pi_U(k)}{m_A^2}} \frac{i}{k^2 + i\epsilon} \begin{bmatrix} 0_{4 \times 4} & 0_{4 \times 1} \\ 0_{1 \times 4} & 1 \end{bmatrix}.$$

Transverse, Longitudinal polarizations remain unchanged. However, the Goldstone poles no longer cancel!

$$\Delta_{\text{GS}}^F(k) = \frac{k^2 - \Pi_L(k) + i\epsilon}{k^2 - m_A^2 - \Pi_L(k) + i\epsilon} \frac{i}{k^2 + i\epsilon}$$

Massless Goldstone with the “Renormalization Factor”

$$Z_{\text{GS}}^2 = \frac{\Pi_L(k) + i\epsilon}{m_A^2 + \Pi_L(k) + i\epsilon}!$$

Finite Temperature

- Wait!!!!!!!!!!!!!!!!!!!!

- $\Pi_L(k) = -\frac{2m_E^2 k^2}{\vec{k}^2} \left(1 - \frac{k^0}{|\vec{k}|} Q_0\left(\frac{k^0}{|\vec{k}|}\right) \right)$ in HTL approximation.

- The k^2 cancels the pole in

$$\Delta_{\text{GS}}^F(k) = \frac{k^2 - \Pi_L(k) + i\epsilon}{k^2 - m_A^2 - \Pi_L(k) + i\epsilon} \frac{i}{k^2 + i\epsilon}!$$

Finite Temperature

- $\Pi_L(k) = -\frac{2m_E^2 k^2}{\vec{k}^2} \left(1 - \frac{k^0}{|\vec{k}|} Q_0\left(\frac{k^0}{|\vec{k}|}\right) \right)$ has a branch-cut along $k^0 = (-|\vec{k}|, |\vec{k}|)$.

- $\Delta_{\text{GS}}^F(k) = \frac{k^2 - \Pi_L(k) + i\epsilon}{k^2 - m_A^2 - \Pi_L(k) + i\epsilon} \frac{i}{k^2 + i\epsilon}$ inherit this branch cut in place of the two poles.

- The branch cut peaks significantly at both $k^0 = \pm |\vec{k}|$. I call it a pair of **Quasi-Pole**.

Finite Temperature

- When $T > T_c$, $m_W = 0$, Goldstone and the Longitudinal polarization decouples. $\Delta_{\text{GS}}^F(k) = \frac{i}{k^2 + i\epsilon}$.
- When $T < T_c$, two poles fragment into two branch cut, which is still similar to two poles. I call this a “cadaver” of a Goldstone boson. Longitudinal polarization eats the Goldstone, but could not devour once in a time.
- $T = 0$, the cadaver completely disappear.

Sterile Neutrino 1 \leftrightarrow 2 processes

- T.Hambye, D.Teresi, an approximation. When $T < T_c$, replace W_L/Z_L with $G^{0,\pm}$ of the same mass as the dressed H . (arXiv:1606.00017)
- Fermions dispersion relations approximation according to G.Guidice et.al., 0310123
- Too rough an approximation!

Sterile Neutrino $1 \leftrightarrow 2$ processes

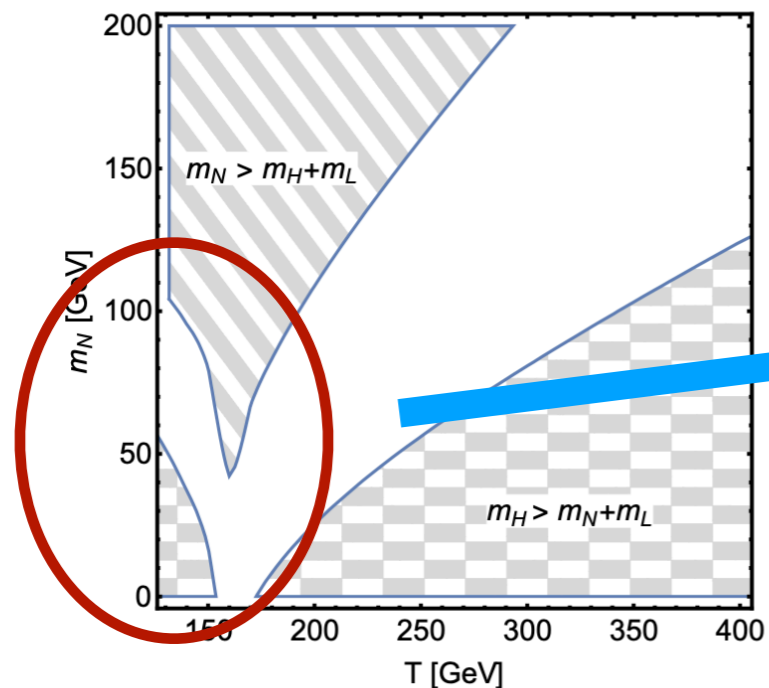


FIG. 1: Values of m_N and T for which the $N \rightarrow LH$ and $H \rightarrow LN$ decays are kinematically allowed.

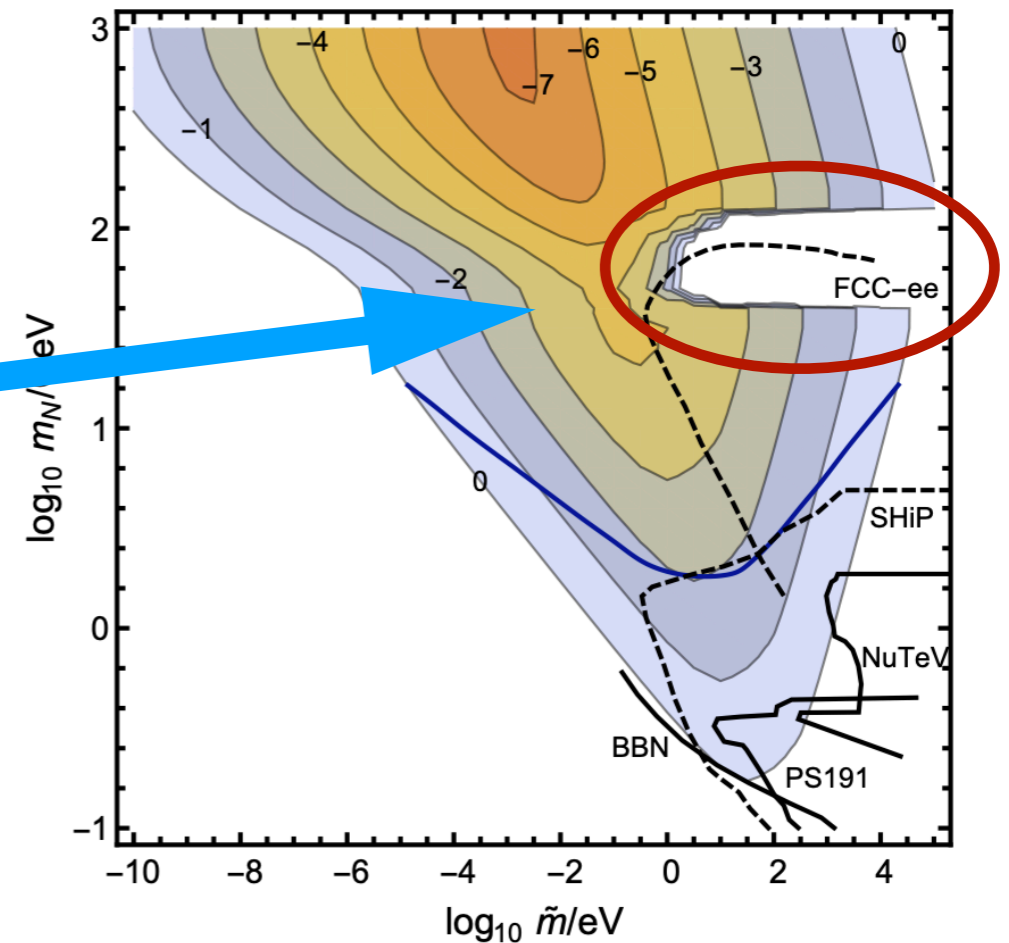
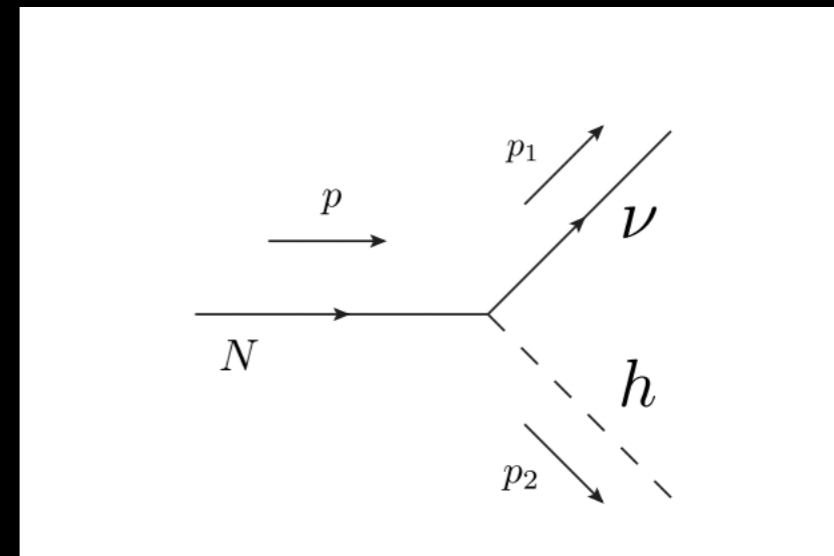
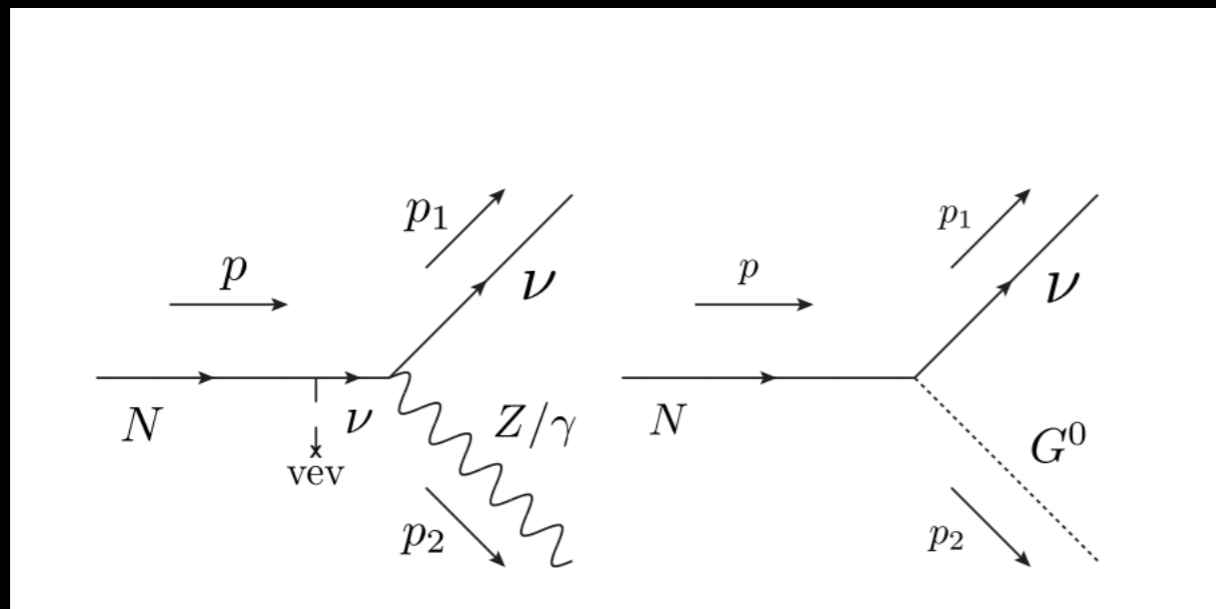
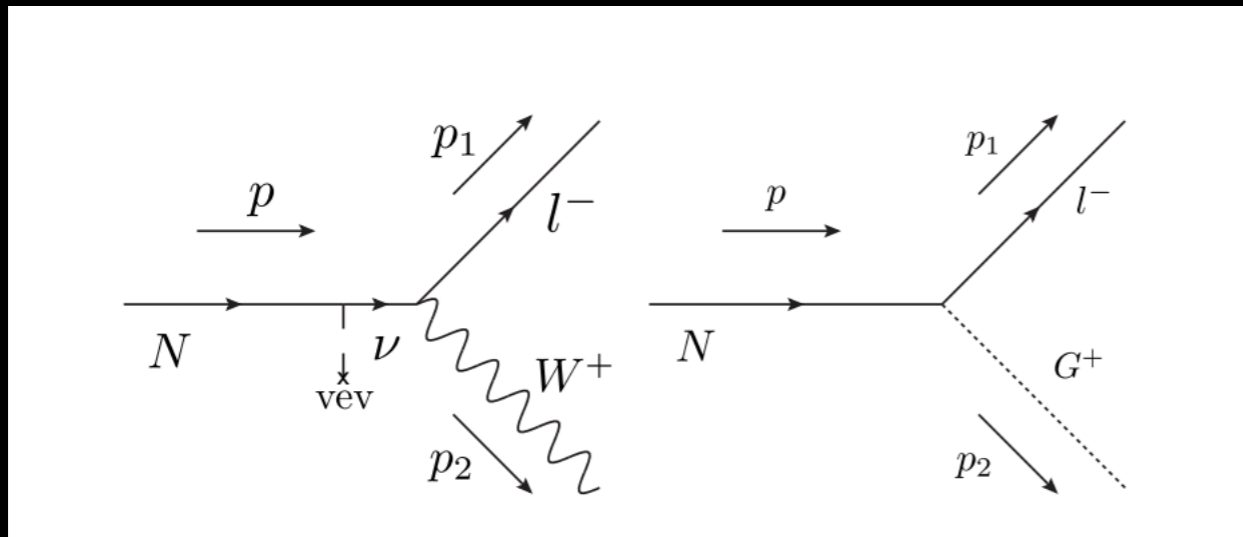


FIG. 3: Logarithm base 10 of the asymmetry ϵ_{CP} needed to obtain successful leptogenesis, with the RH neutrinos initially at thermal equilibrium. We also plot the relevant existing bounds (solid lines) and projected sensitivities of the SHiP [16] and FCC-ee [17] experiments (dashed lines). The area below the thick blue line requires values of ϵ_{CP} which are not reachable for such low m_N .

Sterile Neutrino $1 \leftrightarrow 2$ processes



Sterile Neutrino 1 \leftrightarrow 2 processes

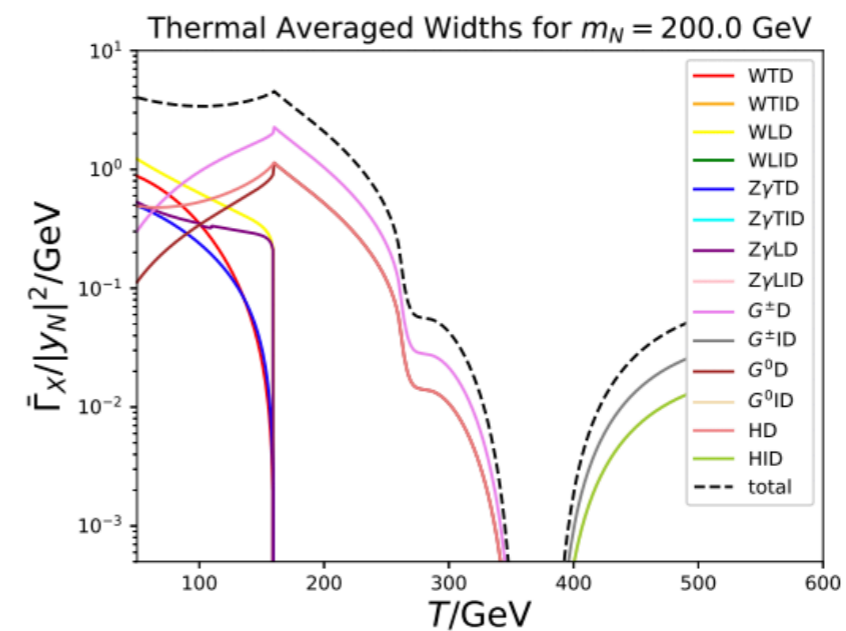
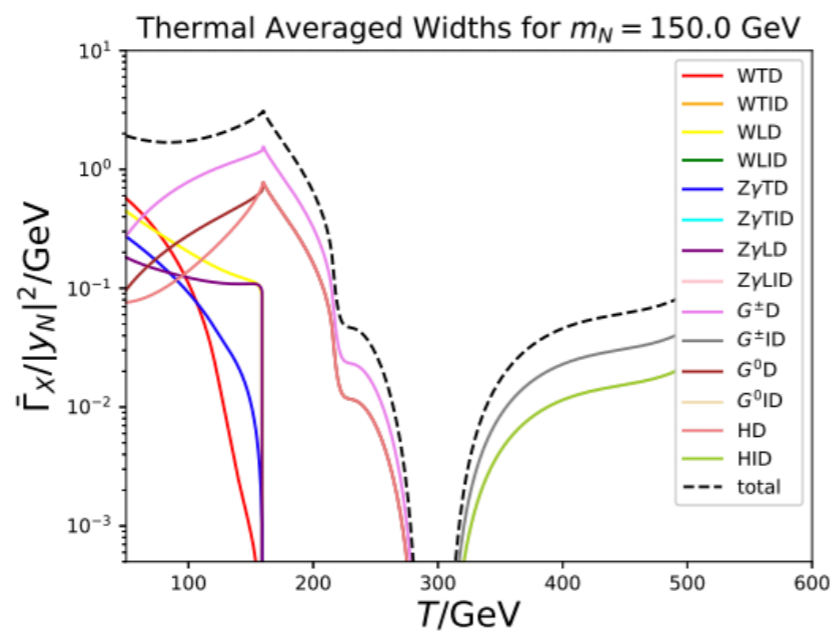
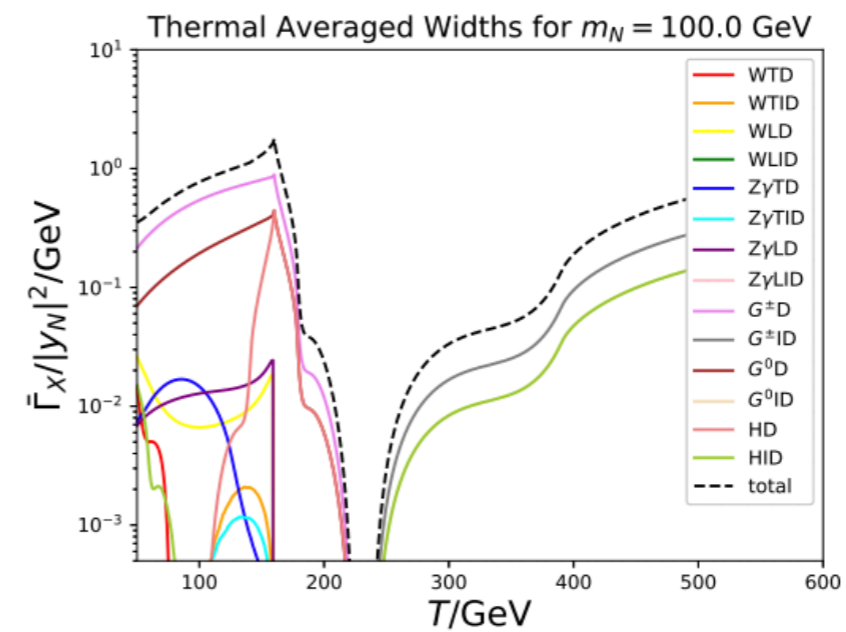
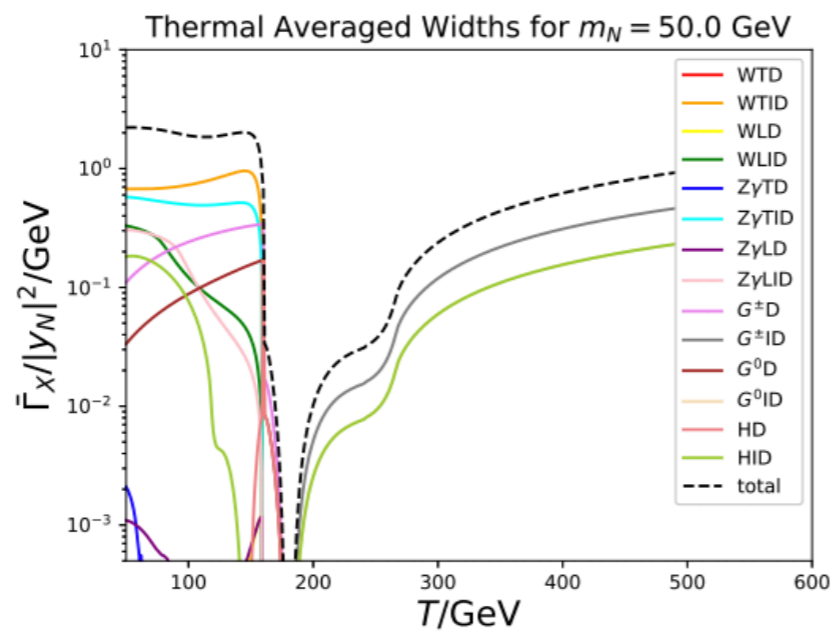
- Lepton: Particle Hole.
- Sterile Neutrino Decay, Inverse Decay of $W/Z/\gamma/h$.
- Transverse, Longitudinal, Goldstone Cadaver.
- 54 channels in total.

Sterile Neutrino $1 \leftrightarrow 2$ processes

Combine all the channels into 14 channels

Alias	Meaning	Alias	Meaning
WTD	$N \leftrightarrow W_T^+ l^-$	Z γ LID	$N\bar{\nu} \leftrightarrow Z_L/\gamma_L$
WTID	$Nl^+ \leftrightarrow W_T^+$	G^\pm D	$N \leftrightarrow G^+ l^-$
WLD	$N \leftrightarrow W_L^+ l^-$	G^\pm ID	$NG^- \leftrightarrow l^-, Nl^+ \leftrightarrow G^+$
WLID	$Nl^+ \leftrightarrow W_L^+$	G^0 D	$NG^0 \leftrightarrow \nu, N\bar{\nu} \leftrightarrow G^0$
Z γ TD	$N \leftrightarrow Z_T/\gamma_T \nu$	G^0 ID	$NG^0 \leftrightarrow \nu, N\bar{\nu} \leftrightarrow G^0$
Z γ TID	$N\bar{\nu} \leftrightarrow Z_T/\gamma_T$	HD	$N \leftrightarrow h\nu$
Z γ LD	$N \leftrightarrow Z_L/\gamma_L \nu$	HID	$Nh \leftrightarrow \nu, N\bar{\nu} \leftrightarrow h$

Sterile Neutrino 1 \leftrightarrow 2 processes



Sterile Neutrino $1 \leftrightarrow 2$ processes

Threshold Effect

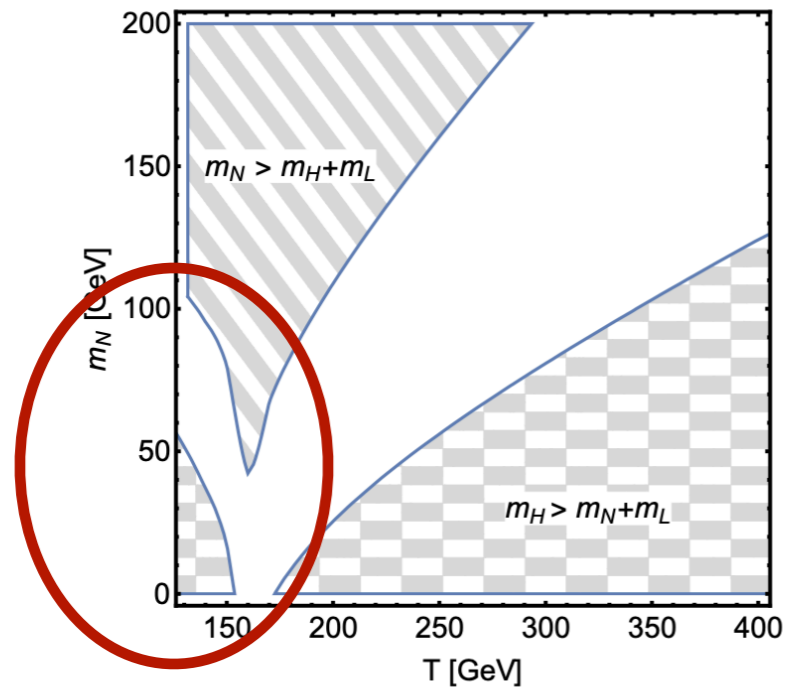
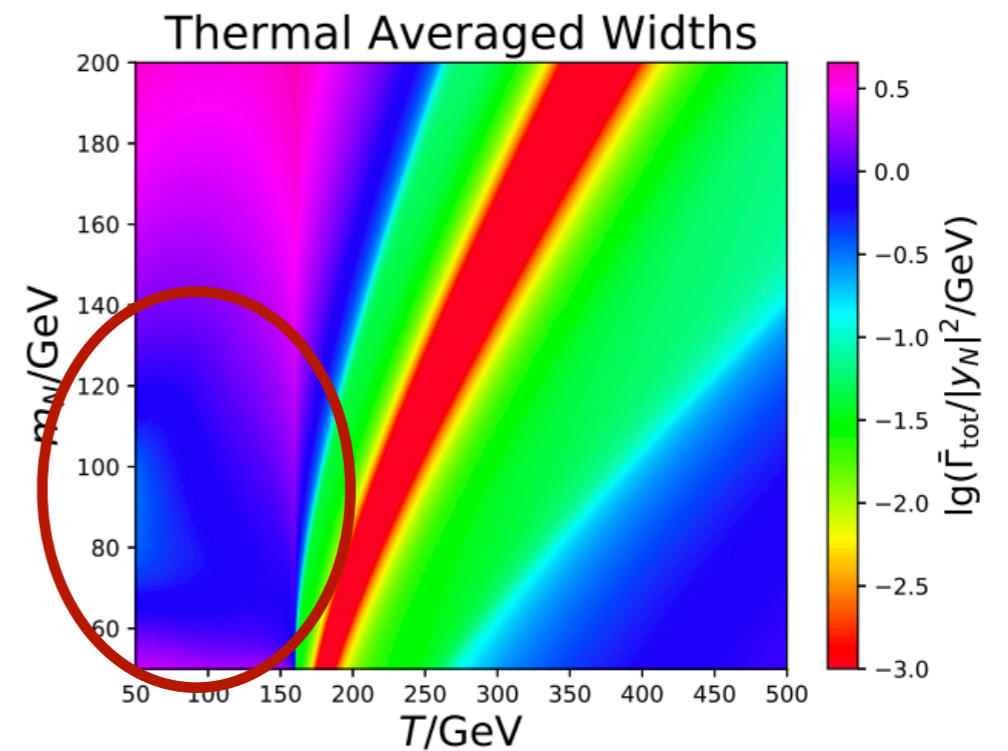


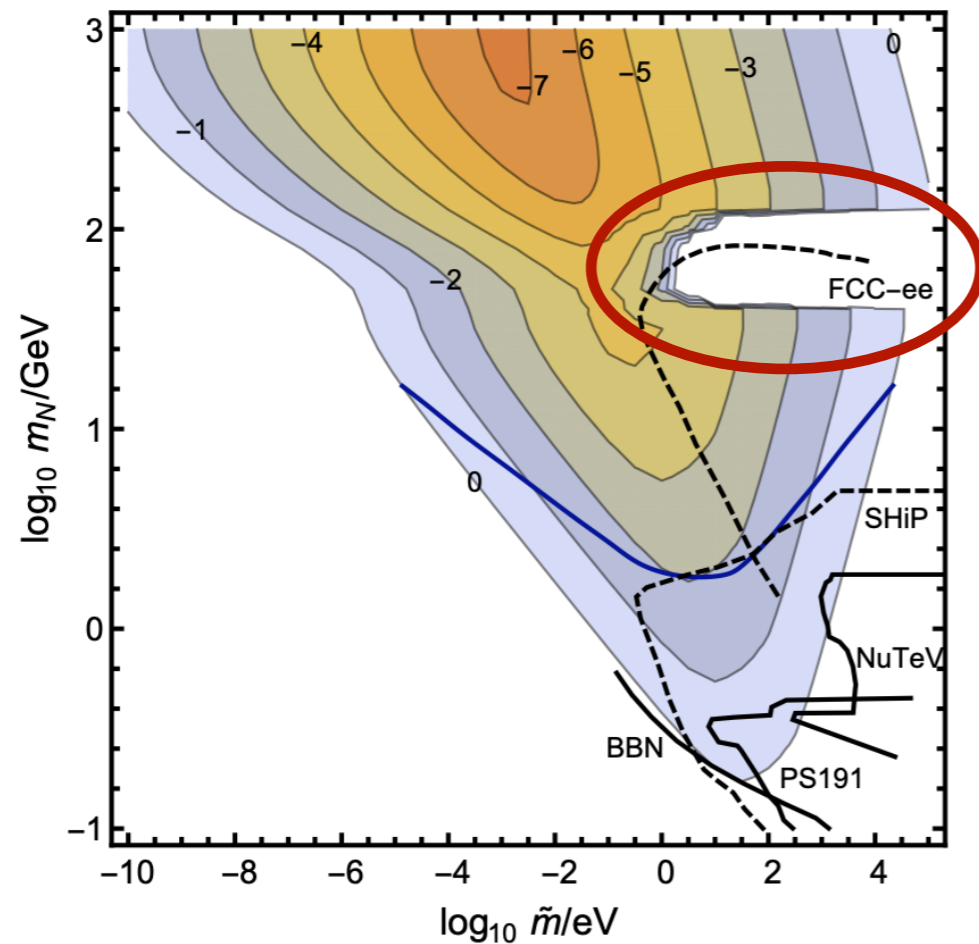
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No Threshold Effect



Leptogenesis

Void existing



Void disappearing

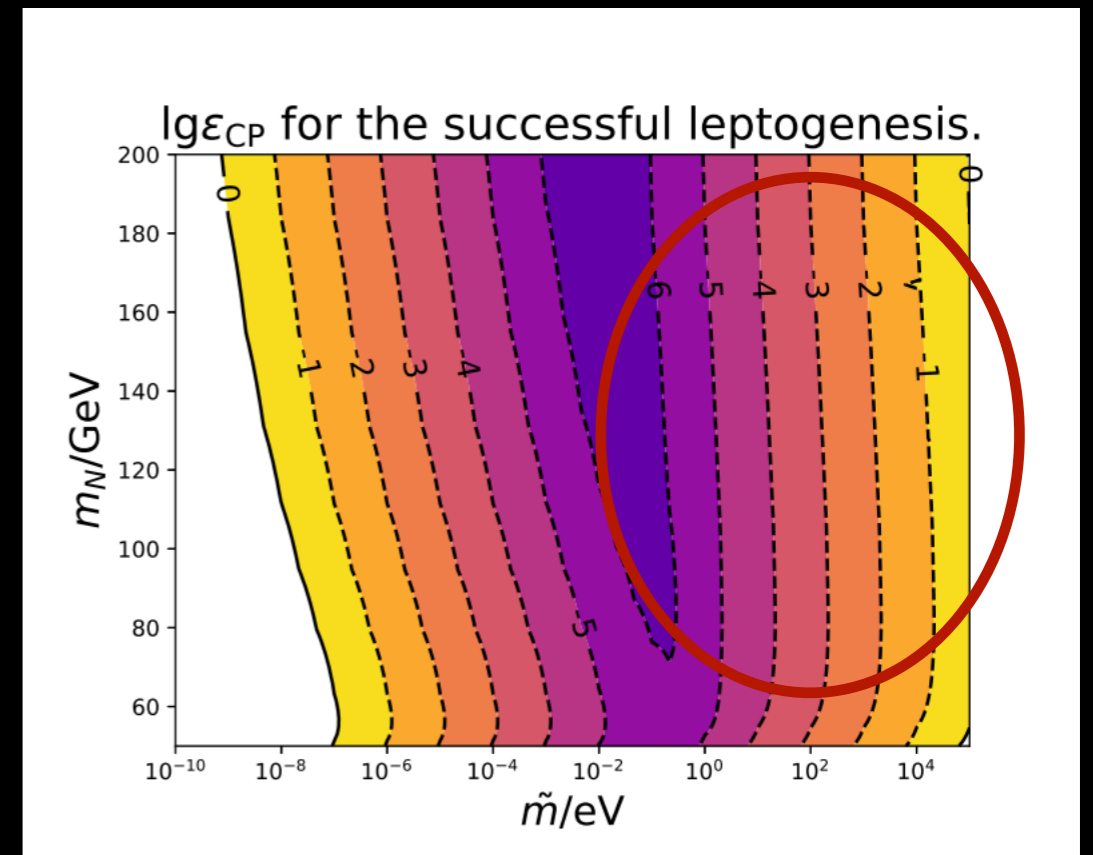


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Future Prospect

- Full calculation of $2 \leftrightarrow 2$ processes. (Thus only initial thermal equilibrium situation)
- LPM resumption.
- Calculation of ϵ_{CP} .
- Dark matter freeze-in involving the massive vector boson.

Thanks!