

J/ψ 光生和电生过程研究

吴佳俊 (中国科学院大学)

合作者: 李正雄, 邹冰松, 杨智, 梁羽铁, 曹须

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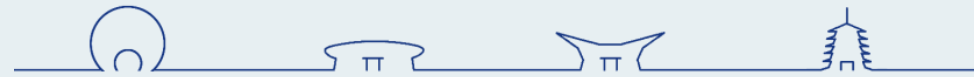


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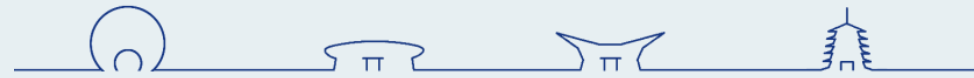
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$\gamma p \rightarrow J/\psi p$ 和 $e p \rightarrow e J/\psi p$ 的物理意义

- 研究胶子场如何传递核子和夸克对之间的相互作用是理解QCD理论的一个重要方向。
- 那么一个重要的过程就是研究 $J/\psi N$ 的相互作用，因为该相互作用领头阶的贡献是 J/ψ 中的正反粲夸克对湮灭以后变成胶子然后和核子 N 相互作用。显然 J/ψ 的光生和电生过程提供了非常好的研究 $J/\psi N$ 相互作用的场所。
- 2015年 P_c 态第一次在LHCb中的，而后在2019年更多的实验数据中确实了 P_c 态的峰结构，并且得到了更加细节的信号。
- 但是至今只有在这个实验中观测到峰结构，这是远远不够的。我们至今不能知道 $P_c \rightarrow J/\psi p$ 的分宽度是多少。所以我们需要更多的实验信息。后面可以看到在 J/ψ 的光生和电生过程中的 $J/\psi p$ 不变质量谱上找到 P_c 信号对理解 P_c 态是非常重要的。



$$\gamma p \rightarrow J/\psi p$$

- 我们在2010年的文章中也提到了 $e p \rightarrow e J/\psi p$ 实验是寻找研究Pc态非常好的地方。随后我在毕业论文里也做了相关计算，但是没有正式发表。
- 去年李老师受到Argonne在Jlab实验组人员委托，希望给他们提供该反应道的建议，所以我们开始从理论估算这个反应道。

[1] Y. Huang, J. He, H.-F. Zhang, and X.-R. Chen, *JPG* **41**, 115004 (2014), 1305.4434.

[2] Q. Wang, X.-H. Liu, and Q. Zhao, *PRD***92**, 034022 (2015), 1508.00339.

[3] V. Kubarovsky and M. B. Voloshin, *PRD***92**, 031502 (2015), 1508.00888.

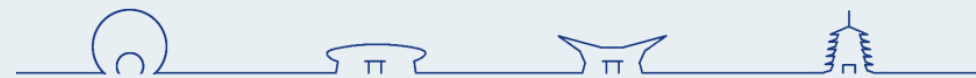
[4] M. Karliner and J. L. Rosner, *PLB***752**, 329 (2016), 1508.01496.

[5] A. N. Hiller Blin, C. Fernandez-Ramirez, A. Jackura, V. Mathieu, etc, *PRD***94**, 034002 (2016), 1606.08912.

[6] E. Ya. Paryev and Yu. T. Kiselev, *NPA***978**, 201 (2018), 1810.01715.

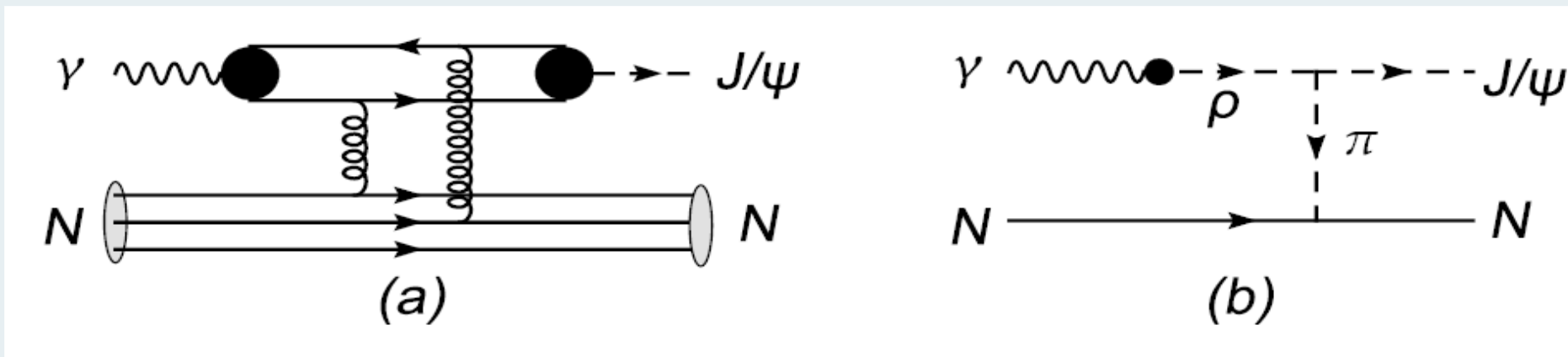
[7] Xu Cao *Eur.Phys.J.C* **81** (2021) 2, 184

[8] Xu Cao, Jian-ping Dai *Phys.Rev.D* **100** (2019) 5, 054033



$\gamma p \rightarrow J/\psi p$ 背景项

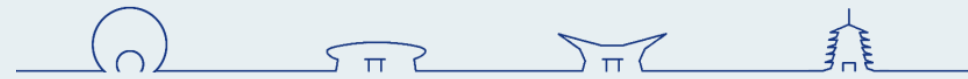
费曼图



公式
$$\frac{d\sigma}{d\Omega} = \frac{1}{(2\pi)^2} \frac{m_N m_B}{4W^2} \frac{1}{4} \sum_{\lambda_\gamma, \lambda_M} \sum_{m_s, m'_s} \left| \bar{u}_p(p', m'_s) \epsilon_\mu^*(q', \lambda'_{J/\psi}) \mathcal{M}^{\mu\nu}(q, p, q', p') u_p(p, m_s) \epsilon_\nu(q, \lambda_\gamma) \right|^2$$

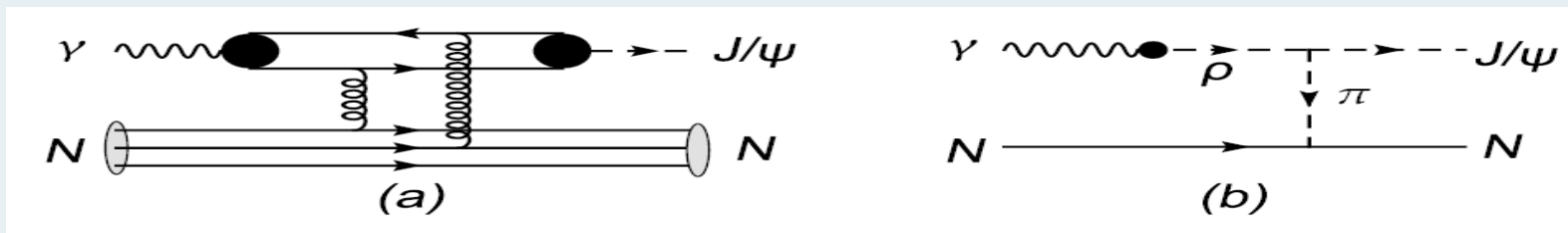
$$\mathcal{M}_P^{\mu\nu}(q, p, q', p') = \left(\frac{s}{s_0}\right)^{\alpha_P(t)-1} \exp\left\{-\frac{i\pi}{2} [\alpha_P(t) - 1]\right\} i 12 e \frac{M_V^2 \beta_q \beta_{q'}}{f_V} \frac{1}{M_V^2 - t} \left(\frac{2\mu_0^2}{2\mu_0^2 + M_V^2 - t}\right) \frac{4M_N^2 - 2.8t}{(4M_N^2 - t)(1 - t/0.71)^2} \{\gamma \cdot q g^{\mu\nu} - q^\mu \gamma^\nu\}$$

$$\mathcal{M}_\pi^{\mu\nu}(q, q', p, p') = \frac{e}{f_\rho} \frac{g_{J/\psi, \rho^0 \pi^0}}{m_{J/\psi}} \frac{f_\pi}{m_\pi} \frac{-m_\rho^2}{q^2 - m_\rho^2 + i\Gamma_\rho m_\rho} \frac{\Lambda_\rho^4}{\Lambda_\rho^4 + (q^2 - m_\rho^2)^2} \frac{1}{t - m_\pi^2} \left(\frac{\Lambda_\pi^2 - m_\pi^2}{\Lambda_\pi^2 - t}\right)^4 \epsilon^{\mu\nu\alpha\beta} q'_\alpha q_\beta (\gamma \cdot (p' - p)) \gamma^5$$

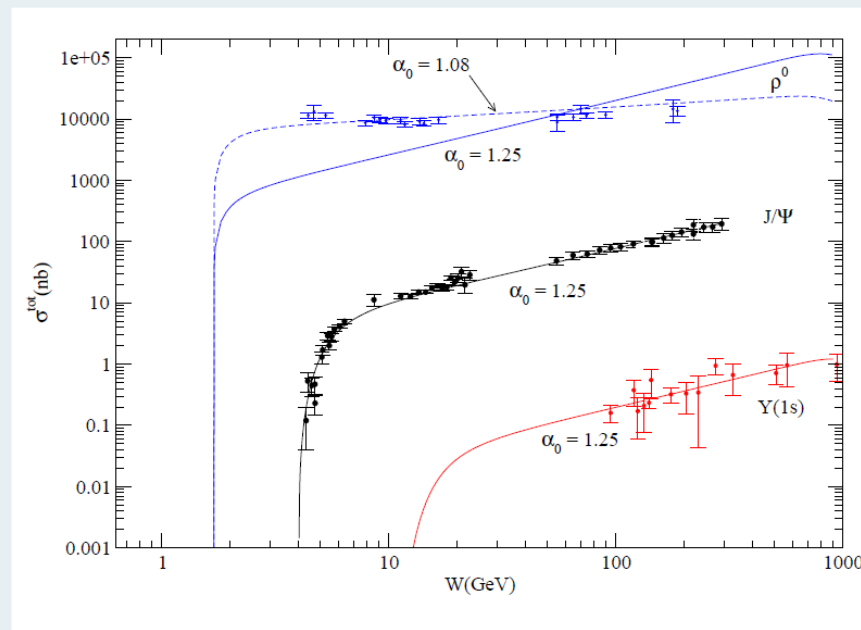
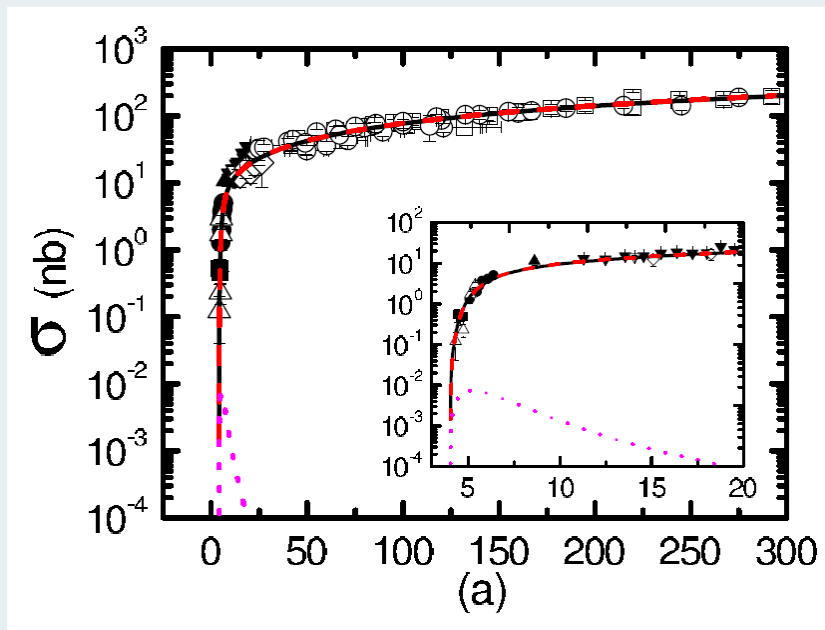


$\gamma p \rightarrow J/\psi p$ 背景项

费曼图



公式



$$\alpha_P(t) = \alpha_0 + \alpha'_p t$$

$$\alpha_0 = 1.08$$

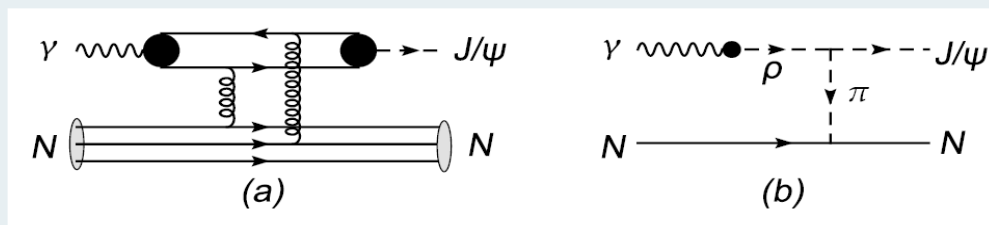
$$\alpha_0 = 1.25$$

$$\mathcal{M}_P^{\mu\nu}(q, p, q', p') = \left(\frac{s}{s_0}\right)^{\alpha_P(t)-1} \exp\left\{-\frac{i\pi}{2}[\alpha_P(t)-1]\right\} i12e \frac{M_V^2 \beta_q \beta_{q'}}{f_V} \frac{1}{M_V^2 - t} \left(\frac{2\mu_0^2}{2\mu_0^2 + M_V^2 - t}\right) \frac{4M_N^2 - 2.8t}{(4M_N^2 - t)(1 - t/0.71)^2} \{\gamma \cdot q g^{\mu\nu} - q^\mu \gamma^\nu\}$$

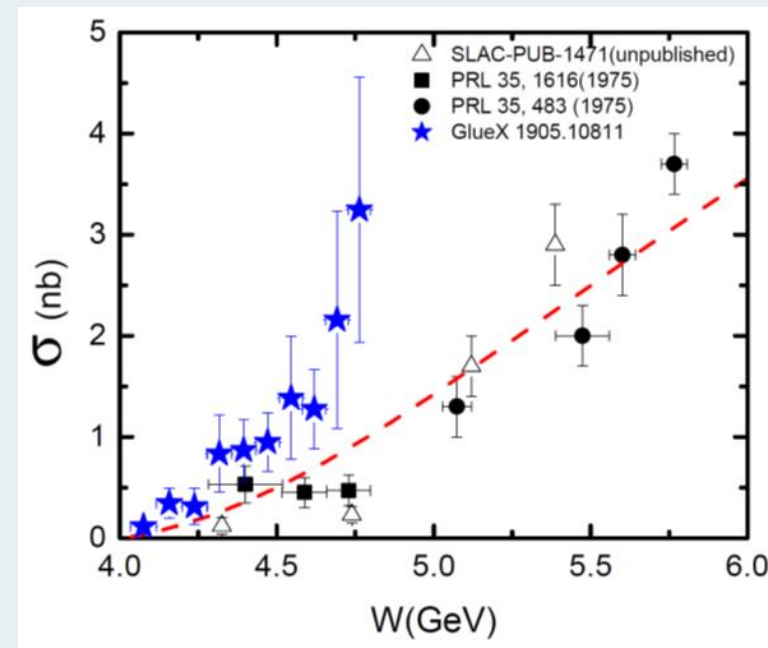
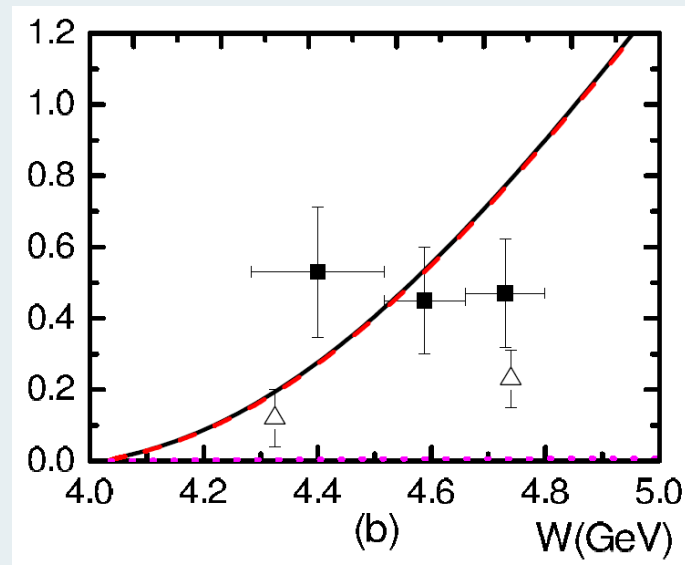
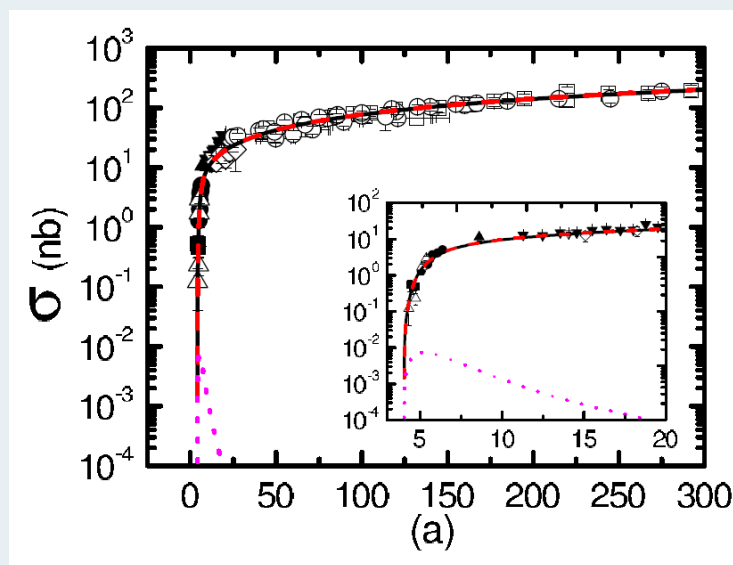


$\gamma p \rightarrow J/\psi p$ 背景项

费曼图



公式



$$\mathcal{M}_P^{\mu\nu}(q, p, q', p') = \left(\frac{s}{s_0}\right)^{\alpha_P(t)-1} \exp\left\{-\frac{i\pi}{2}[\alpha_P(t)-1]\right\} i12e \frac{M_V^2 \beta_q \beta_{q'}}{f_V} \frac{1}{M_V^2 - t} \left(\frac{2\mu_0^2}{2\mu_0^2 + M_V^2 - t}\right) \frac{4M_N^2 - 2.8t}{(4M_N^2 - t)(1 - t/0.71)^2} \{\gamma \cdot q g^{\mu\nu} - q^\mu \gamma^\nu\}$$

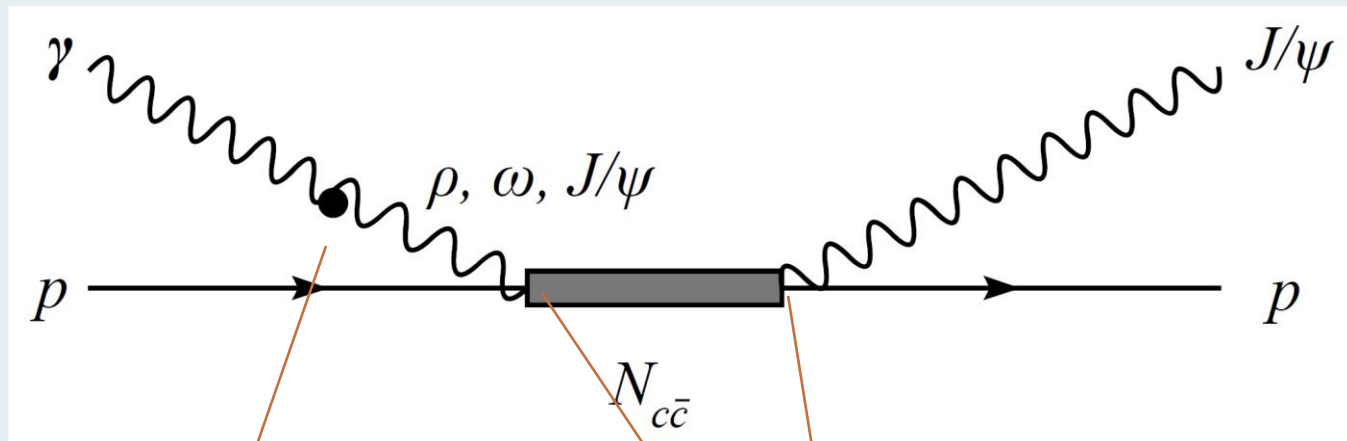
$$\alpha_P(t) = \alpha_0 + \alpha'_p t$$

$$\alpha_0 = 1.08$$

$$\alpha_0 = 1.25$$



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献

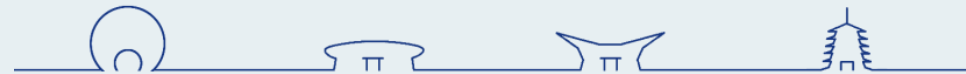
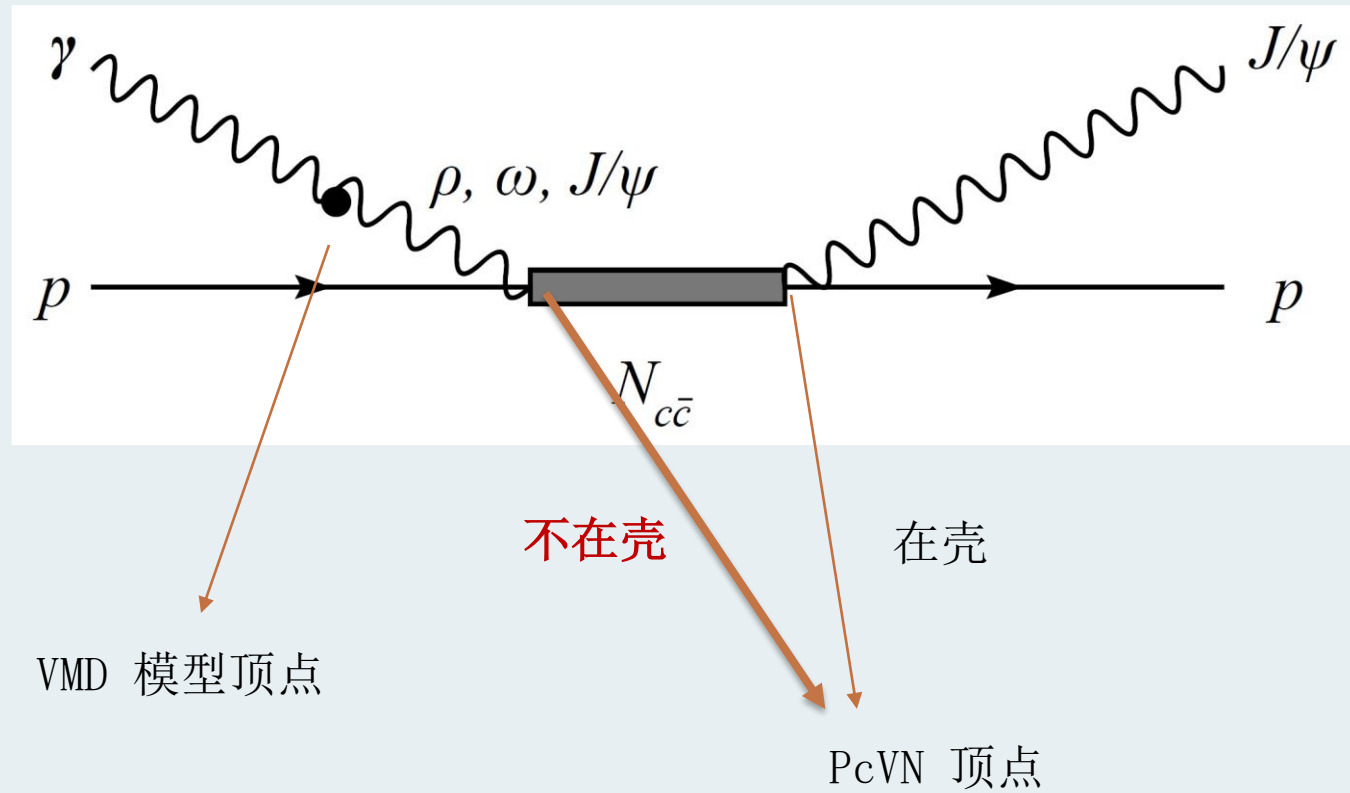


VMD 模型顶点

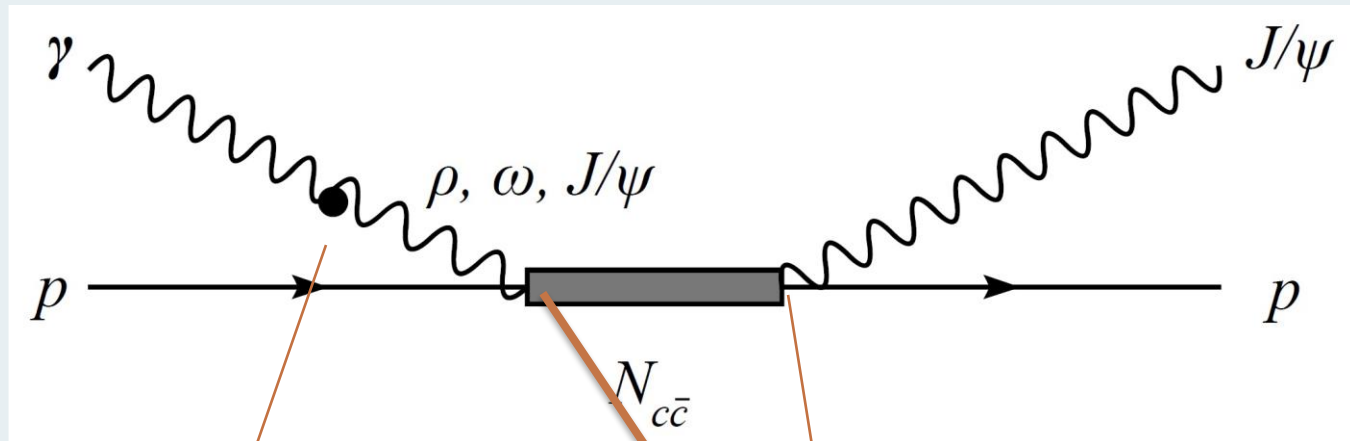
PcVN 顶点



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献



VMD 模型顶点

不在壳

在壳

PcVP 顶点

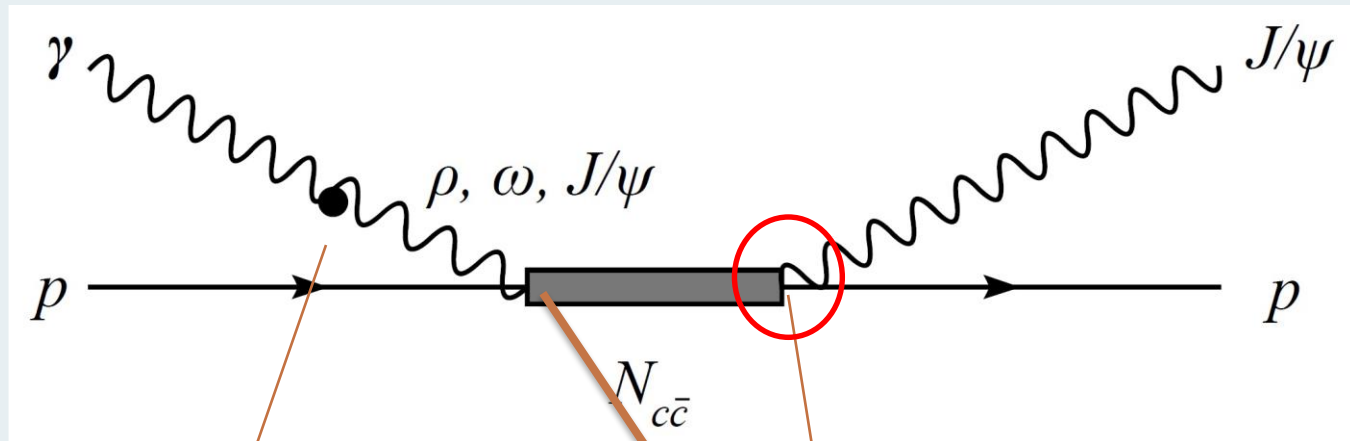
Pc γ P 顶点



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$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献



不在壳

在壳

VMD 模型顶点

PcVP 顶点

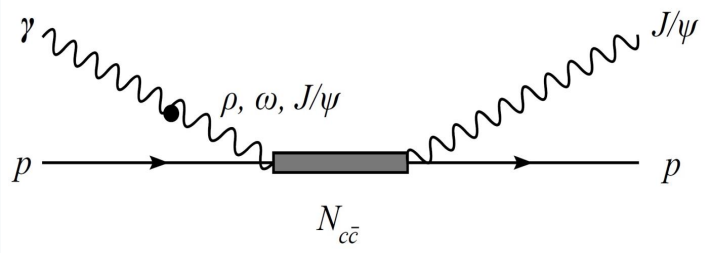
Pc γ P 顶点



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献

• $P_c \rightarrow VB$ 不同的理论模型

No.	J^P	m	Γ	$\Gamma_{J/\psi N}$	$\Gamma_{\rho N}$	$\Gamma_{\omega N}$	$\Gamma_{\bar{D}\Lambda_c}$	$\Gamma_{\bar{D}^*\Lambda_c}$	Main Channel	Ref.
1	$\frac{1}{2}^-$	4262	35.6	10.3	—	—	0.01	—	$\bar{D}\Sigma_c$	[6]
2		4308	—	1.2	—	—	0.02	1.4	$\bar{D}\Sigma_c$	[7]
3		4412	47.3	19.2	3.2	10.4	—	—	$\bar{D}^*\Sigma_c$	[8, 9]
4		4410	58.9	52.5	—	—	0.8	0.7	$\bar{D}^*\Sigma_c$	[6]
5		4460	—	3.9	—	—	1.0	0.3	$\bar{D}^*\Sigma_c$	[7]
6		4481	57.8	14.3	—	—	1.02	0.3	$\bar{D}^*\Sigma_c^*$	[6]
7		4334	38.8	38.0	—	—	—	0.8	$\bar{D}\Sigma_c^*$	[6]
8		4375	—	1.5	—	—	—	0.9	$\bar{D}\Sigma_c^*$	[7]
9		4380	144.3	3.8	1.4	5.3	1.2	131.3	$\bar{D}\Sigma_c^*$	[5]
10		4380	69.9	16.6	0.15	0.6	17.0	35.3	$\bar{D}^*\Sigma_c$	[5]
11	$\frac{3}{2}^-$	4412	47.3	19.2	3.2	10.4	—	—	$\bar{D}^*\Sigma_c$	[8, 9]
12		4417	8.2	4.6	—	—	—	3.1	$\bar{D}^*\Sigma_c$	[6]
13		4450	139.8	16.3	0.14	0.5	41.4	72.3	$\bar{D}^*\Sigma_c$	[5]
14		4450	21.7	0.03	—	—	1.4	6.8	$\bar{D}^*\Sigma_c$	[10]
15		4450	16.2	11	—	—	0.6	4.2	$\Psi'N$	[10]
16		4453	—	1.5	—	—	—	0.3	$\bar{D}\Sigma_c^*$	[7]
17		4481	34.7	32.8	—	—	—	1.2	$\bar{D}^*\Sigma_c^*$	[6]
18	$\frac{5}{2}^+$	4450	46.4	4.0	0.3	0.3	18.8	20.5	$\bar{D}^*\Sigma_c$	[5]
19	$\frac{3}{2}^- / \frac{5}{2}^+$	$4380_{\pm 29}^{+8}$	$205_{\pm 86}^{+18}$	—	—	—	—	—	Exp	[1, 2]
20		$4450_{\pm 3}^{+2}$	$39_{\pm 19}^{+5}$	—	—	—	—	—	Exp	[1, 2]



[5] Lin, Shen, Guo, Zou, PRD95 114017
[6] Xiao, Nieves, Oset, PRD88 056012
[7] Huang, Ping, 1811. 04260
[8,9] Wu, Molina, Oset, Zou, PRL 105 232001, PRC 84 015202
[10] Eides, Petrov, 1811.01691



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献

- $P_c \rightarrow VB$ 相互作用的LS耦合的洛伦兹结构

$$\mathcal{M}_{N^*(\frac{1}{2}^-)NV} = \bar{u}_N \gamma_5 \tilde{\gamma}_\mu u_{N^*} \epsilon_V^*{}_\nu \left(g_{1V} g^{\mu\nu} + f_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right)$$

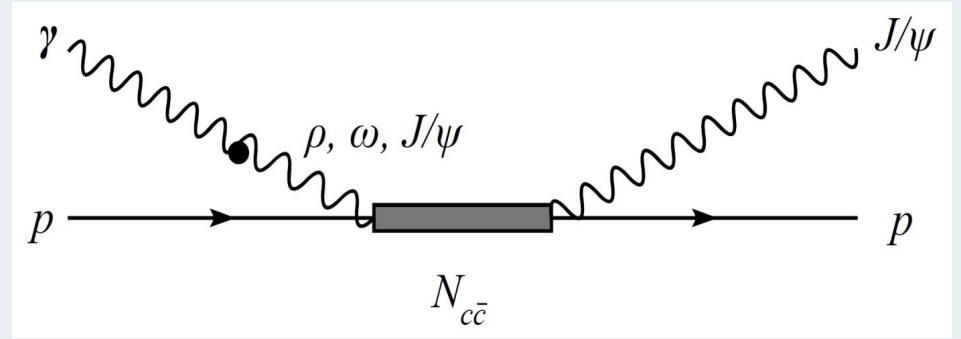
$$\mathcal{M}_{N^*(\frac{3}{2}^-)NV} = \bar{u}_N u_{N^*}{}_\mu \epsilon_V^*{}_\nu \left(g_{3V} g^{\mu\nu} + f_{3V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right)$$

$$+ h_{3V} \epsilon_{\mu\nu\lambda\delta} \bar{u}_N \gamma_5 (\tilde{\gamma}^\mu g_\alpha^\beta + \tilde{\gamma}_\alpha g^{\mu\beta}) u_{N^*}{}_\beta \epsilon_V^*{}_\nu \left(\frac{\tilde{r}^\alpha \tilde{r}^\lambda}{\tilde{r}^2} - \frac{1}{3} \tilde{g}_{N^*}^{\alpha\lambda} \right) \hat{P}^\delta$$

$$\mathcal{M}_{N^*(\frac{5}{2}^+)NV} = \bar{u}_N u_{N^*}{}_{\mu\nu} \epsilon_V^*{}_\alpha \left(\frac{g_{5V}}{m_N} g^{\alpha\mu} \tilde{r}^\nu + \frac{f_{5V}}{m_N} \left(\frac{3}{5} \frac{\tilde{r}^\mu \tilde{r}^\nu \tilde{r}^\alpha}{\tilde{r}^2} - \frac{1}{5} (\tilde{g}_{N^*}^{\mu\nu} \tilde{r}^\alpha + \tilde{g}_{N^*}^{\nu\alpha} \tilde{r}^\mu + \tilde{g}_{N^*}^{\alpha\mu} \tilde{r}^\nu) \right) \right)$$

$$+ \frac{h_{5V}}{m_N} \epsilon_{\mu\nu\lambda\delta} \bar{u}_N \gamma_5 (\tilde{\gamma}^\mu g_{\xi\alpha} g_{\sigma\beta} + \tilde{\gamma}_\xi g_{\sigma\beta} g_{\mu\beta} + \tilde{\gamma}_\sigma g_{\mu\beta} g_{\xi\beta}) u_{N^*}^{\alpha\beta} \epsilon_V^*{}^\mu$$

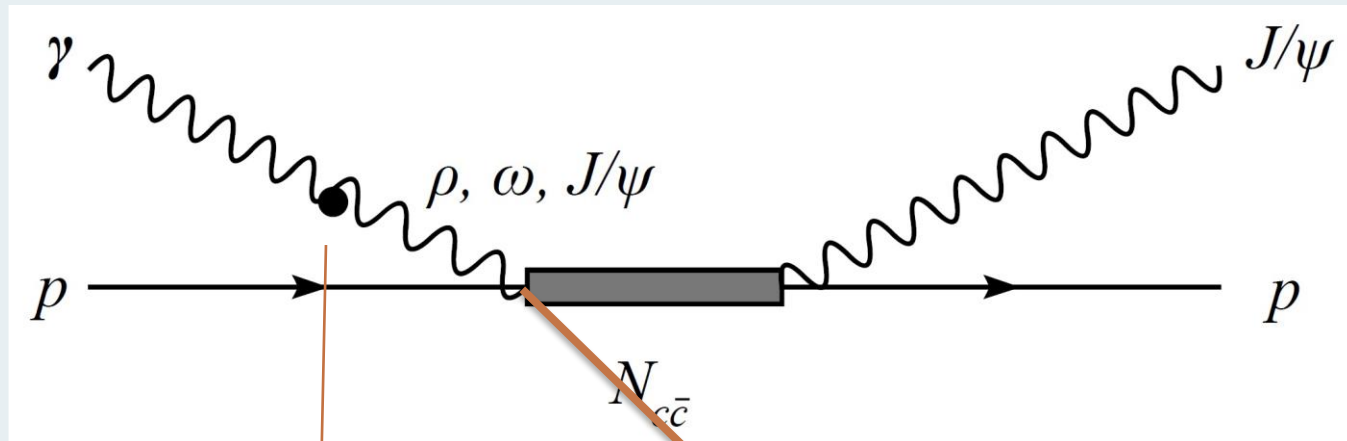
$$\times \left(\frac{\tilde{r}^\xi \tilde{r}^\lambda \tilde{r}^\sigma}{\tilde{r}^2} - \frac{1}{3} (\tilde{g}_{N^*}^{\xi\sigma} \tilde{r}^\lambda + \tilde{g}_{N^*}^{\sigma\lambda} \tilde{r}^\xi + \tilde{g}_{N^*}^{\lambda\xi} \tilde{r}^\sigma) \right) \hat{P}^\delta$$



No.	J^P	m	Γ_{tot}	$\frac{g_V}{J/\Psi p}$
1	$\frac{1}{2}^-$	4262	35.6	0.39
2		4308	—	0.13
3		4412	47.3	0.46
4		4410	58.9	0.75
5		4460	—	0.20
6		4481	57.8	0.37
7	$\frac{3}{2}^-$	4334	38.8	1.19
8		4375	—	0.23
9		4380	144.3	0.36
10		4380	69.9	0.75
11		4412	47.3	0.79
12		4417	8.2	0.39
13		4450	139.8	0.71
14		4450	21.7	0.030
15		4450	16.2	0.58
16		4453	—	0.21
17		4481	34.7	0.98
18	$\frac{5}{2}^+$	4450	46.4	0.35



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点



VMD 模型顶点

$$\mathcal{L}_{VDM} = \frac{iem_V^2}{f_V} A_\mu V^\mu$$

不在壳

$$F_V(q^2) = \frac{\Lambda^4}{\Lambda^4 + (q^2 - m_V^2)^2}$$

问题2. off-shell的影响

PcVP 顶点

$$\mathcal{L}_{N^*(\frac{1}{2}^-)NV} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N V_\nu \left(g_{1V} g^{\mu\nu} + f_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$

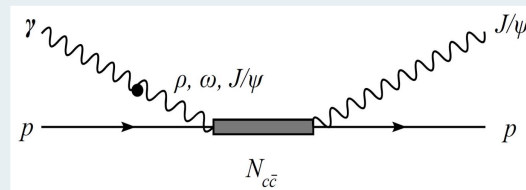
问题1. 光子的
规范不变性

PcγP 顶点



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点

- $\gamma p \rightarrow P_c$ 顶点的 **光子的规范不变性**



$$\mathcal{L}_{VDM} = \frac{iem_V^2}{f_V} A_\mu V^\mu$$

$$\mathcal{L}_{N^*(\frac{1}{2}^-)NV} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N V_\nu \left(g_{1V} g^{\mu\nu} - f_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$

$$\mathcal{T}_{N^*(\frac{1}{2}^-) \rightarrow NV \rightarrow N\gamma} = \bar{u}_N \mathcal{M}^\nu u_{N^*} \epsilon_\nu^*,$$

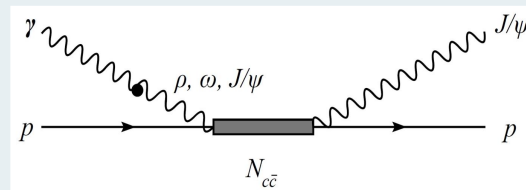
$$\mathcal{M}^\nu = \frac{ie}{f_V} \frac{-m_V^2}{q^2 - m_V^2 + i\Gamma_V m_V} \gamma_5 \tilde{\gamma}_\mu \left(g_{1V} g_{\mu\nu'} - f_{1V} \left(\frac{3}{2} \frac{\tilde{r}_\mu \tilde{r}_{\nu'}}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu'} \right) \right) \tilde{g}_V^{\nu'\nu}(q) \times F_V(q^2)$$

$$\mathcal{M}^\nu q_\nu \sim (g_{1V} - f_{1V}) \neq 0 \quad \text{破坏了光子的规范不变性}$$



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点

- $\gamma p \rightarrow P_c$ 顶点的光子的规范不变性



$$\mathcal{L}_{VDM} = \frac{iem_V^2}{f_V} A_\mu V^\mu$$

$$\mathcal{L}_{N^*(\frac{1}{2}^-)NV} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N V_\nu \left(g_{1V} g^{\mu\nu} - f_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$

$$\mathcal{T}_{N^*(\frac{1}{2}^-) \rightarrow NV \rightarrow N\gamma} = \bar{u}_N \mathcal{M}^\nu u_{N^*} \epsilon_\nu^*,$$

$$\mathcal{M}^\nu = \frac{ie}{f_V} \frac{-m_V^2}{q^2 - m_V^2 + i\Gamma_V m_V} \gamma_5 \tilde{\gamma}_\mu \left(g_{1V} g_{\mu\nu'} - f_{1V} \left(\frac{3}{2} \frac{\tilde{r}_\mu \tilde{r}_{\nu'}}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu'} \right) \right) \tilde{g}_V^{\nu'\nu}(q) \times F_V(q^2)$$

$$\mathcal{M}^\nu q_\nu \sim (g_{1V} - f_{1V}) \neq 0 \quad \text{破坏了光子的规范不变性}$$

$$\mathcal{L}_{N^*(\frac{1}{2}^-)N\gamma} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N A_\nu \left(\underline{g}_{1\gamma} g^{\mu\nu} - \underline{g}_{1\gamma} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$

$$\text{Dulat, Wu, Zou, PRD83, 094032} \quad J/\psi \rightarrow \bar{B} B \gamma$$

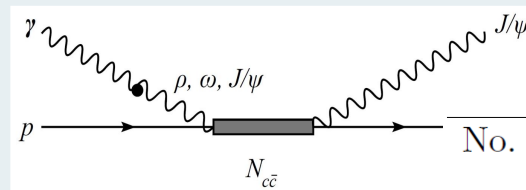
一个简单粗暴的方案是：（不是说我们解决了这个问题）

$$\mathcal{L}_{N^*(\frac{1}{2}^-)NV} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N V_\nu \left(\underline{\tilde{g}}_{1V} g^{\mu\nu} - \underline{\tilde{g}}_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点

- $\gamma p \rightarrow P_c$ 顶点的光子的规范不变性



No.	J^P	m	Γ_{tot}	g_V $J/\Psi p$	$\tilde{g}_V$ $J/\Psi p$	$\tilde{g}_V$ ρp	$\tilde{g}_V$ ωp
1	$\frac{1}{2}^-$	4262	35.6	0.39	0.32	—	—
2		4308	—	0.13	0.11	—	—
3		4412	47.3	0.46	0.38	0.078	0.14
4		4410	58.9	0.75	0.62	—	—
5	$\frac{3}{2}^-$	4460	—	0.20	0.16	—	—
6		4481	57.8	0.37	0.31	—	—
7		4334	38.8	1.19	0.98	—	—
8		4375	—	0.23	0.19	—	—
9		4380	144.3	0.36	0.30	0.090	0.17
10		4380	69.9	0.75	0.62	0.039	0.059
11		4412	47.3	0.79	0.65	0.14	0.24
12		4417	8.2	0.39	0.32	—	—
13		4450	139.8	0.71	0.58	0.028	0.053
14		4450	21.7	0.030	0.025	—	—
15		4450	16.2	0.58	0.48	—	—
16		4453	—	0.21	0.18	—	—
17		4481	34.7	0.98	0.81	—	—
18	$\frac{5}{2}^+$	4450	46.4	0.35	0.27	0.016	0.016

$$\mathcal{L}_{VDM} = \frac{iem_V^2}{f_V} A_\mu V^\mu$$

$$\mathcal{L}_{N^*(\frac{1}{2}^-)NV} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N V_\nu \left(g_{1V} g^{\mu\nu} - f_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$

$$\mathcal{T}_{N^*(\frac{1}{2}^-) \rightarrow NV \rightarrow N\gamma} = \bar{u}_N \mathcal{M}^\nu u_{N^*} \epsilon_\nu^*,$$

$$\mathcal{M}^\nu = \frac{ie}{f_V} \frac{-m_V^2}{q^2 - m_V^2 + i\Gamma_V m_V} \gamma_5 \tilde{\gamma}_\mu \left(g_{1V} g_{\mu\nu'} - f_{1V} \left(\frac{3}{2} \frac{\tilde{r}_\mu \tilde{r}_{\nu'}}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}{}_{\mu\nu'} \right) \right) \tilde{g}_V^{\nu'\nu}(q),$$

$$\mathcal{M}^\nu q_\nu \sim (g_{1V} - f_{1V}) \neq 0 \quad \text{破坏了光子的规范不变性}$$

$$\mathcal{L}_{N^*(\frac{1}{2}^-)N\gamma} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N A_\nu \left(\textcolor{red}{g}_{1\gamma} g^{\mu\nu} - \textcolor{red}{g}_{1\gamma} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$

$$\text{Dulat, Wu, Zou, PRD83, 094032} \quad J/\psi \rightarrow \bar{B} B \gamma$$

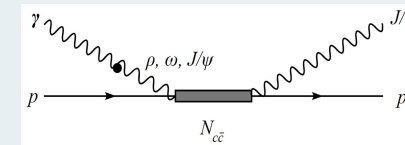
一个简单粗暴的方案是：（不是说我们解决了这个问题）

$$\mathcal{L}_{N^*(\frac{1}{2}^-)NV} = \bar{N}^* \gamma_5 \tilde{\gamma}_\mu N V_\nu \left(\textcolor{red}{\tilde{g}}_{1V} g^{\mu\nu} - \textcolor{red}{\tilde{g}}_{1V} \left(\frac{3}{2} \frac{\tilde{r}^\mu \tilde{r}^\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^*}^{\mu\nu} \right) \right) + h.c.$$



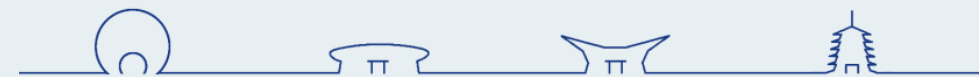
$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点

- $\gamma p \rightarrow P_c$ 顶点的 **矢量介子 off-shell** 的影响



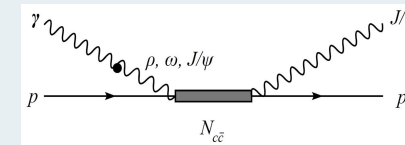
$$\mathcal{T}_{N^*(\frac{1}{2}^-) \rightarrow NV \rightarrow N\gamma} = \bar{u}_N \mathcal{M}^\nu u_{N^*} \epsilon_\nu^*,$$

$$\mathcal{M}^\nu = \frac{ie}{f_V} \frac{-m_V^2 \tilde{g}_{1V}}{q^2 - m_V^2 + i\Gamma_V m_V} \gamma_5 \tilde{\gamma}_\mu \left(g_{\mu\nu'} - \left(\frac{3}{2} \frac{\tilde{r}_\mu \tilde{r}_\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^* \mu\nu'} \right) \right) \tilde{g}_V^{\nu'\nu}(q) \times F_V(q^2)$$



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点

- $\gamma p \rightarrow P_c$ 顶点的 **矢量介子off-shell** 的影响



$$\mathcal{T}_{N^*(\frac{1}{2}^-) \rightarrow NV \rightarrow N\gamma} = \bar{u}_N \mathcal{M}^\nu u_{N^*} \epsilon_\nu^*,$$

$$\mathcal{M}^\nu = \frac{ie}{f_V} \frac{-m_V^2 \tilde{g}_{1V}}{q^2 - m_V^2 + i\Gamma_V m_V} \gamma_5 \tilde{\gamma}_\mu \left(g_{\mu\nu'} - \left(\frac{3}{2} \frac{\tilde{r}_\mu \tilde{r}_\nu}{\tilde{r}^2} - \frac{1}{2} \tilde{g}_{N^* \mu\nu'} \right) \right) \tilde{g}_V^{\nu'\nu}(q) \times F_V(q^2)$$

$$F_V(q^2) = \frac{\Lambda^4}{\Lambda^4 + (q^2 - m_V^2)^2} = \frac{\Lambda^4}{\Lambda^4 + m_V^4}$$

$D^* \rightarrow D\gamma$ 用VMD 模型来处理, $D^* \rightarrow D(\rho \ \& \ \omega \ \& \ J/\psi) \rightarrow D\gamma$ 来估算 Λ 的大小, 这个值其实挺小的, 650 MeV左右。这个结果给我们一个非常强烈的暗示, $N^*N\gamma$ 的强度可能被这个矢量介子的off-shell极大的压低了, 所以我们在随后的计算中, 尽量选取符合现有实验数据的结果, 避免出现预言非常大的共振态贡献。最后我们取 $\Lambda = 550$ MeV。



$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献中 $\gamma p \rightarrow P_c$ 顶点

• γ

$\mathcal{M}^\nu$

$D^* \rightarrow D\gamma$
这个结果
后的计算

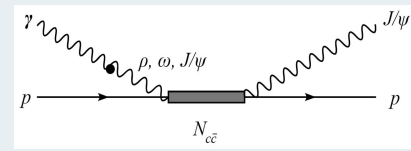
No.	J^P	m	Γ_{tot}	g_V $J/\Psi p$	$\tilde{g}_V$ $J/\Psi p$	$\tilde{g}_V$ ρp	$\tilde{g}_V$ ωp	$\Gamma_{p\gamma}$ (keV)
1	$\frac{1}{2}^-$	4262	35.6	0.39	0.32	—	—	3.9×10^{-5}
2		4308	—	0.13	0.11	—	—	4.5×10^{-6}
3		4412	47.3	0.46	0.38	0.078	0.14	1.14
4		4410	58.9	0.75	0.62	—	—	1.5×10^{-4}
5		4460	—	0.20	0.16	—	—	1.1×10^{-5}
6		4481	57.8	0.37	0.31	—	—	3.8×10^{-5}
7	$\frac{3}{2}^-$	4334	38.8	1.19	0.98	—	—	1.3×10^{-4}
8		4375	—	0.23	0.19	—	—	4.6×10^{-6}
9		4380	144.3	0.36	0.30	0.090	0.17	0.53
10		4380	69.9	0.75	0.62	0.039	0.059	0.060
11		4412	47.3	0.79	0.65	0.14	0.24	1.1
12		4417	8.2	0.39	0.32	—	—	1.4×10^{-5}
13		4450	139.8	0.71	0.58	0.028	0.053	0.054
14		4450	21.7	0.030	0.025	—	—	8.4×10^{-8}
15		4450	16.2	0.58	0.48	—	—	3.1×10^{-5}
16		4453	—	0.21	0.18	—	—	4.2×10^{-6}
17		4481	34.7	0.98	0.81	—	—	8.8×10^{-5}
18	$\frac{5}{2}^+$	4450	46.4	0.35	0.27	0.016	0.016	8.3×10^{-2}

影响

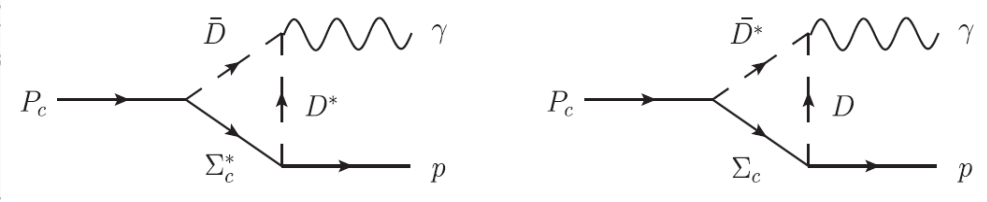
$\bar{u}$

$\bar{c}$

$\Lambda^4 + m_V^4$



No.	J^P	m	Γ_{tot}	$\Gamma_{J/\Psi p}$	$\Gamma_{p\gamma}$ (keV)	$\sigma^{(tot)}$ (nb)	Ref.
9	$\frac{3}{2}^-$	4380	144.3	3.8	0.53	0.11	This work
		4380	144.3	3.8	0.70	0.15	[19]
18	$\frac{5}{2}^+$	4450	46.4	4.0	0.083	0.25	This work
		4450	46.4	4.0	1.13	3.4	[19]

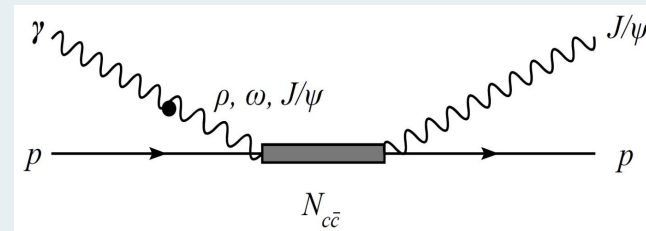


[19] Lin, Shen, Guo, Zou, PRD95 114017

$D\gamma$ 来估算 Λ 的大小, 这个值其实挺小的, 650 MeV 左右。
个矢量介子的 off-shell 极大的压低了, 所以我们在随
预言非常大的共振态贡献。最后我们取 $\Lambda = 550$ MeV。

$\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的贡献

- $\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的振幅



$$\frac{d\sigma}{d\Omega} = \frac{1}{(2\pi)^2} \frac{m_N m_B}{4W^2} \frac{1}{4} \sum_{\lambda_\gamma, \lambda_M} \sum_{m_s, m'_s} \left| \bar{u}_p(p', m'_s) \epsilon_\mu^*(q', \lambda'_{J/\psi}) \mathcal{M}^{\mu\nu}(q, p, q', p') u_p(p, m_s) \epsilon_\nu(q, \lambda_\gamma) \right|^2$$

$$\mathcal{M}_{N^*(\frac{1}{2}^-)}^{\mu\nu}(q, p, q', p') = \sum_{V=J/\Psi, \rho, \omega} g_{1J/\Psi} \gamma_5 \tilde{\gamma}_\alpha \tilde{g}^{\alpha\mu}(q) \frac{\gamma \cdot (q+p) + m_{N_{c\bar{c}}^*}}{W^2 - m_{N_{c\bar{c}}^*}^2 + i\Gamma_{N_{c\bar{c}}^*} m_{N_{c\bar{c}}^*}} F_V(0) \times \frac{ie}{f_V - m_V^2 + i\Gamma_V m_V} \frac{-m_V^2 \tilde{g}_{1V}}{m_V} \gamma_5 \tilde{\gamma}_\beta \left(g^{\beta\nu} - \frac{3}{2} \frac{\tilde{r}^\beta \tilde{r}^\nu}{\tilde{r}^2} + \frac{1}{2} \tilde{g}_{N^*}^{\beta\nu} \right)$$

$$\mathcal{M}_{N^*(\frac{3}{2}^-)}^{\mu\nu}(q, p, q', p') = \sum_{V=J/\Psi, \rho, \omega} g_{3J/\Psi} g^{\mu\alpha} \frac{(\gamma \cdot (q+p) + m_{N_{c\bar{c}}^*}) P_{\alpha\beta}^{\frac{3}{2}}(p+q)}{W^2 - m_{N_{c\bar{c}}^*}^2 + i\Gamma_{N_{c\bar{c}}^*} m_{N_{c\bar{c}}^*}} F_V(0) \frac{ie}{f_V - m_V^2 + i\Gamma_V m_V} \frac{-m_V^2 \tilde{g}_{3V}}{m_V} \left(g^{\beta\nu} - \frac{3}{2} \frac{\tilde{r}^\beta \tilde{r}^\nu}{\tilde{r}^2} + \frac{1}{2} \tilde{g}_{N^*}^{\beta\nu} \right),$$

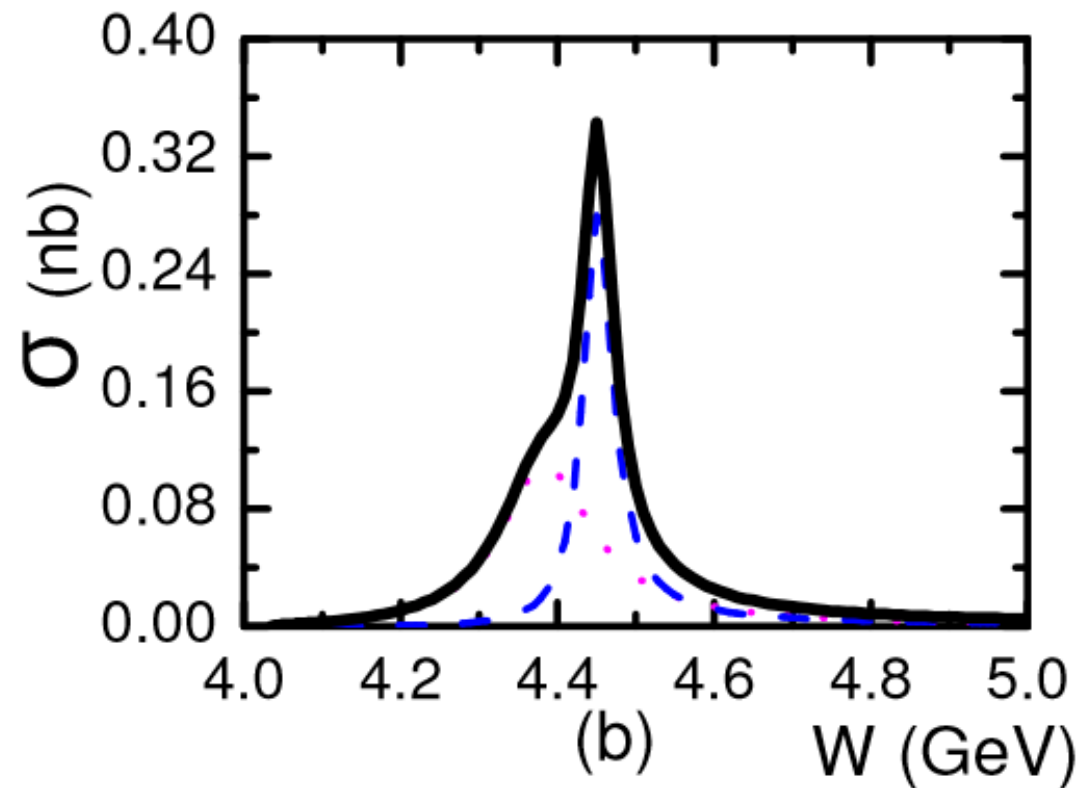
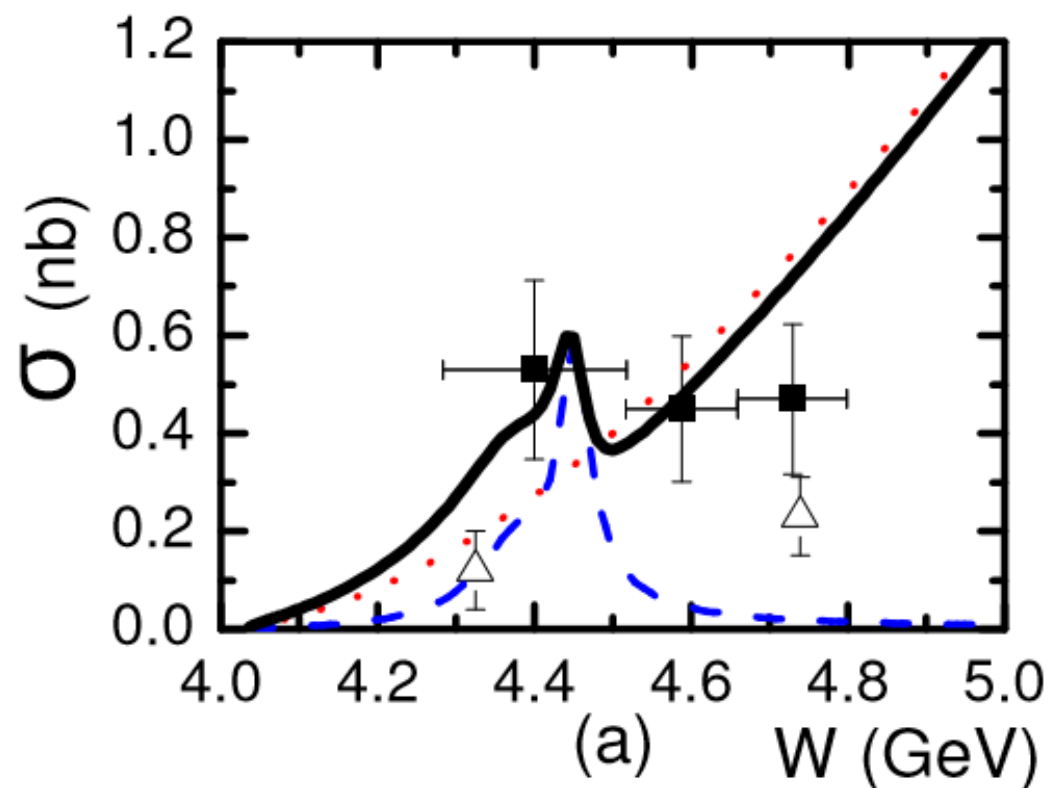
$$\mathcal{M}_{N^*(\frac{5}{2}^+)}^{\mu\nu}(q, p, q', p') = \sum_{V=J/\Psi, \rho, \omega} \frac{g_{5J/\Psi}}{m_N} g^{\mu\alpha} \tilde{r}^\beta \frac{(\gamma \cdot (q+p) + m_{N_{c\bar{c}}^*}) P_{\alpha\beta\alpha'\beta'}^{\frac{5}{2}}(p+q)}{W^2 - m_{N_{c\bar{c}}^*}^2 + i\Gamma_{N_{c\bar{c}}^*} m_{N_{c\bar{c}}^*}} F_V(0) \times \frac{ie}{f_V - m_V^2 + i\Gamma_V m_V} \frac{-m_V^2 \tilde{g}_{5V}/m_N}{m_V} \left(g^{\nu\alpha'} \tilde{r}^{\beta'} - \frac{5}{3} \frac{\tilde{r}^\nu \tilde{r}^{\alpha'} \tilde{r}^{\beta'}}{\tilde{r}^2} + \frac{1}{3} \left(\tilde{g}_{N^*}^{\nu\alpha'} \tilde{r}^{\beta'} + \tilde{g}_{N^*}^{\nu\beta'} \tilde{r}^{\alpha'} + \tilde{g}_{N^*}^{\alpha'\beta'} \tilde{r}^\nu \right) \right)$$

$$P_{\alpha\beta}^{\frac{3}{2}}(p) = -g_{\alpha\beta} + \frac{1}{3} \gamma_\mu \gamma_\nu + \frac{2}{3} \frac{p_\mu p_\nu}{m_{N^*}^2} + \frac{1}{3m_{N^*}} (\gamma_\mu p_\nu - \gamma_\nu p_\mu)$$

$$P_{\frac{5}{2}}^{\alpha\alpha'\beta\beta'}(p) = \frac{1}{2} (\tilde{g}_{N^*}^{\alpha\alpha'} \tilde{g}_{N^*}^{\beta\beta'} + \tilde{g}_{N^*}^{\alpha\beta'} \tilde{g}_{N^*}^{\beta\alpha'}) - \frac{1}{5} \tilde{g}_{N^*}^{\alpha\beta} \tilde{g}_{N^*}^{\alpha'\beta'} - \frac{1}{10} \left(\tilde{\gamma}^\alpha \tilde{\gamma}^{\alpha'} \tilde{g}_{N^*}^{\beta\beta'} + \tilde{\gamma}^\alpha \tilde{\gamma}^{\beta'} \tilde{g}_{N^*}^{\beta\alpha'} + \tilde{\gamma}^\beta \tilde{\gamma}^{\alpha'} \tilde{g}_{N^*}^{\alpha\beta'} + \tilde{\gamma}^\beta \tilde{\gamma}^{\beta'} \tilde{g}_{N^*}^{\alpha\alpha'} \right)$$



$$\gamma p \rightarrow P_c \rightarrow J/\psi p$$



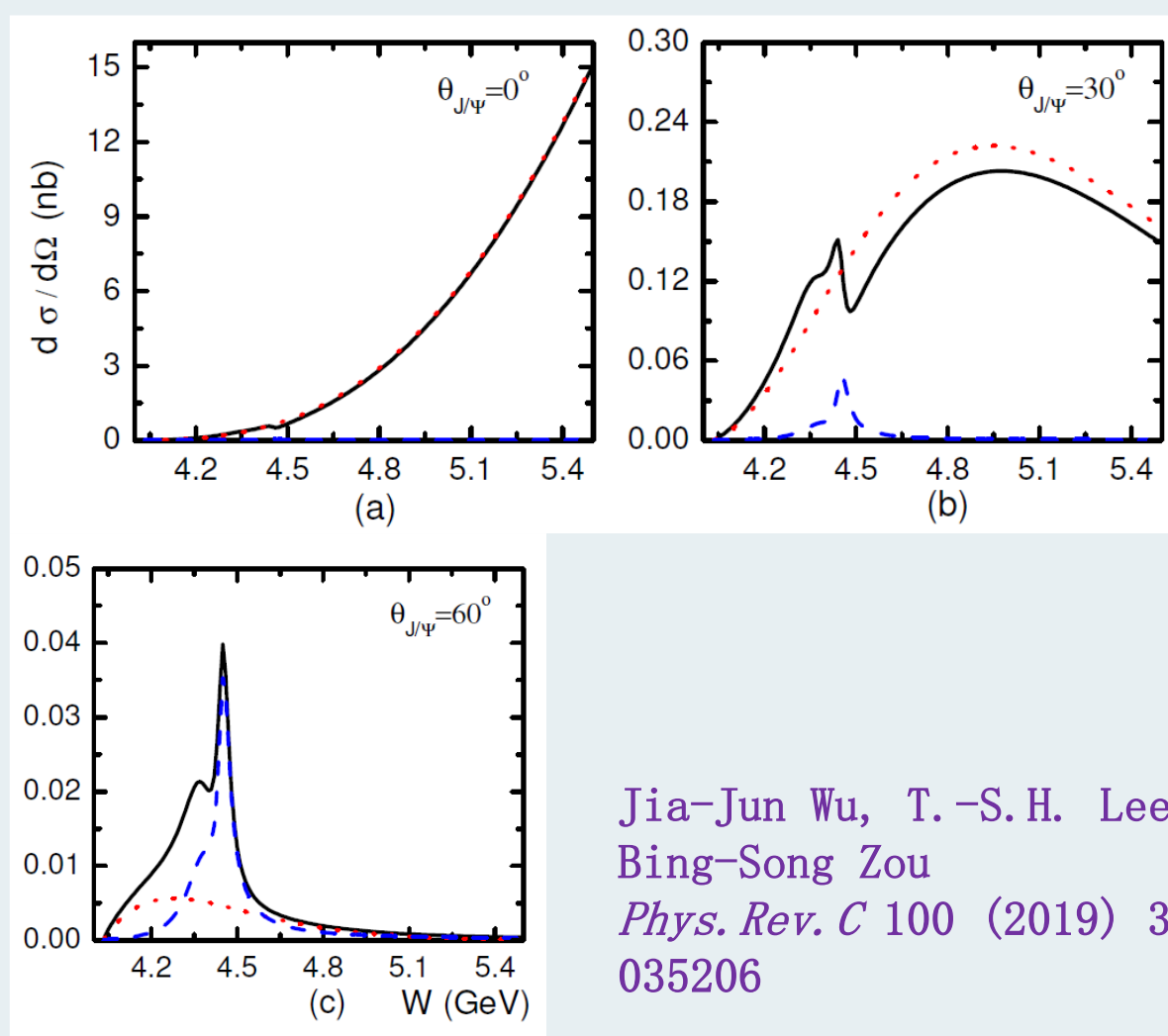
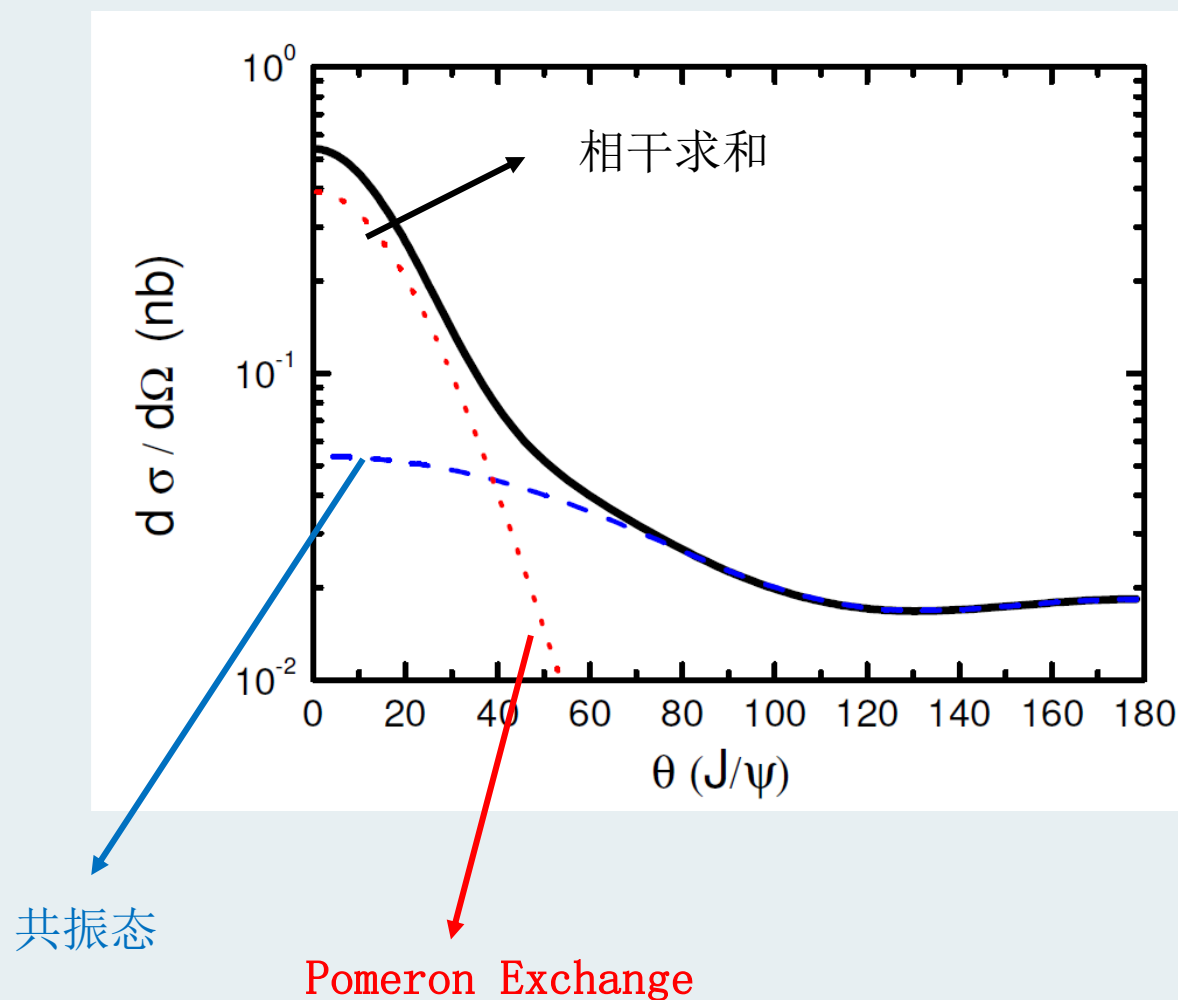
Black: total
 Red: Pomeron-exchange
 Blue: P_c



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如何从J/ψ p末态中尽可能抽取 P_C 的峰结构?

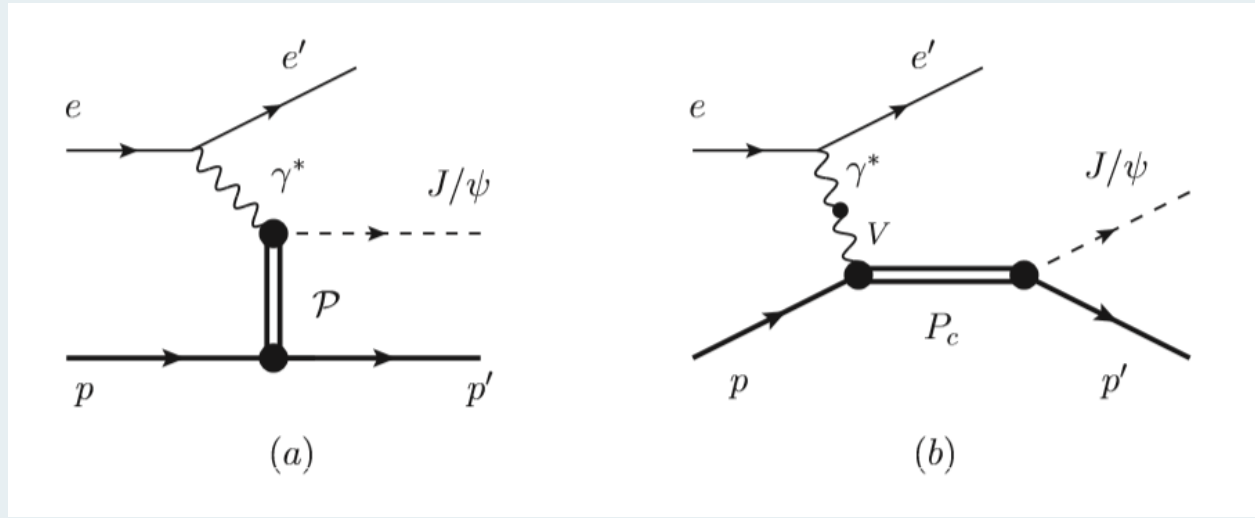


Jia-Jun Wu, T.-S.H. Lee,
Bing-Song Zou
Phys. Rev. C 100 (2019) 3,
035206



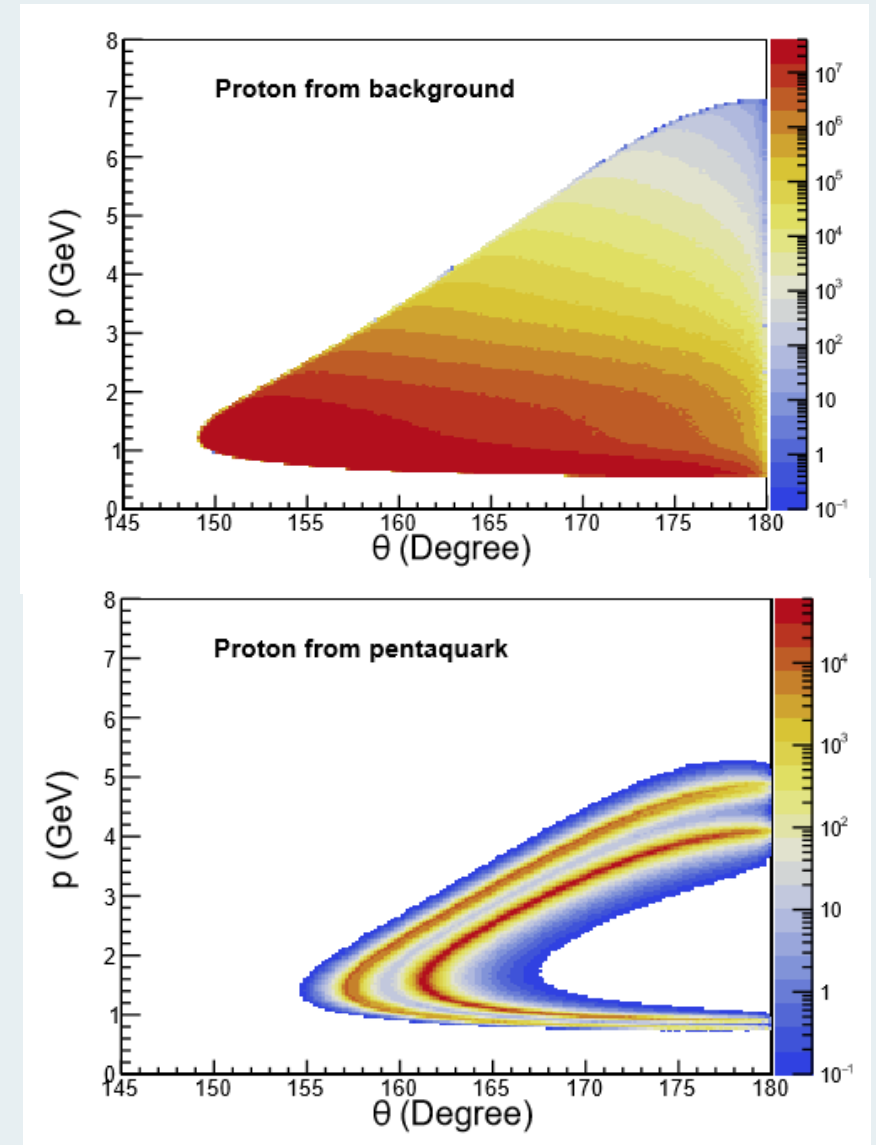
如何从J/ψ p末态中尽可能抽取 P_c 的峰结构?

$$e p \rightarrow e J/\psi p$$



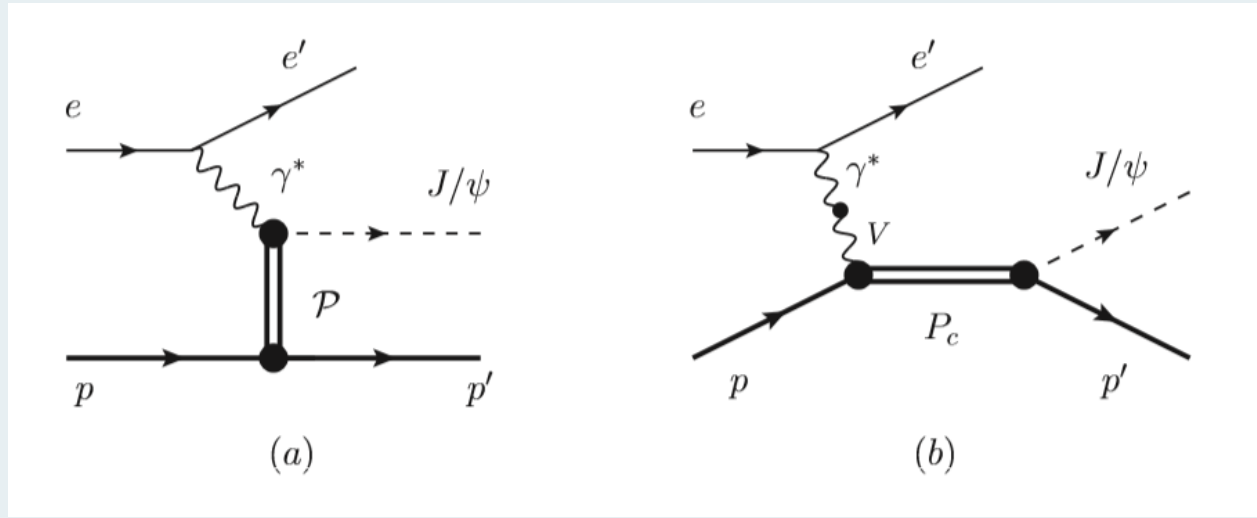
The distribution of momentum versus scattering angles of final proton in the laboratory frame at the **JLab energy configuration, with 12 GeV** electron beam projectile on rest proton.

Zhi Yang, Xu Cao, Yu-Tie Liang, Jia-Jun Wu
Chin. Phys. C 44 (2020) 8, 084102



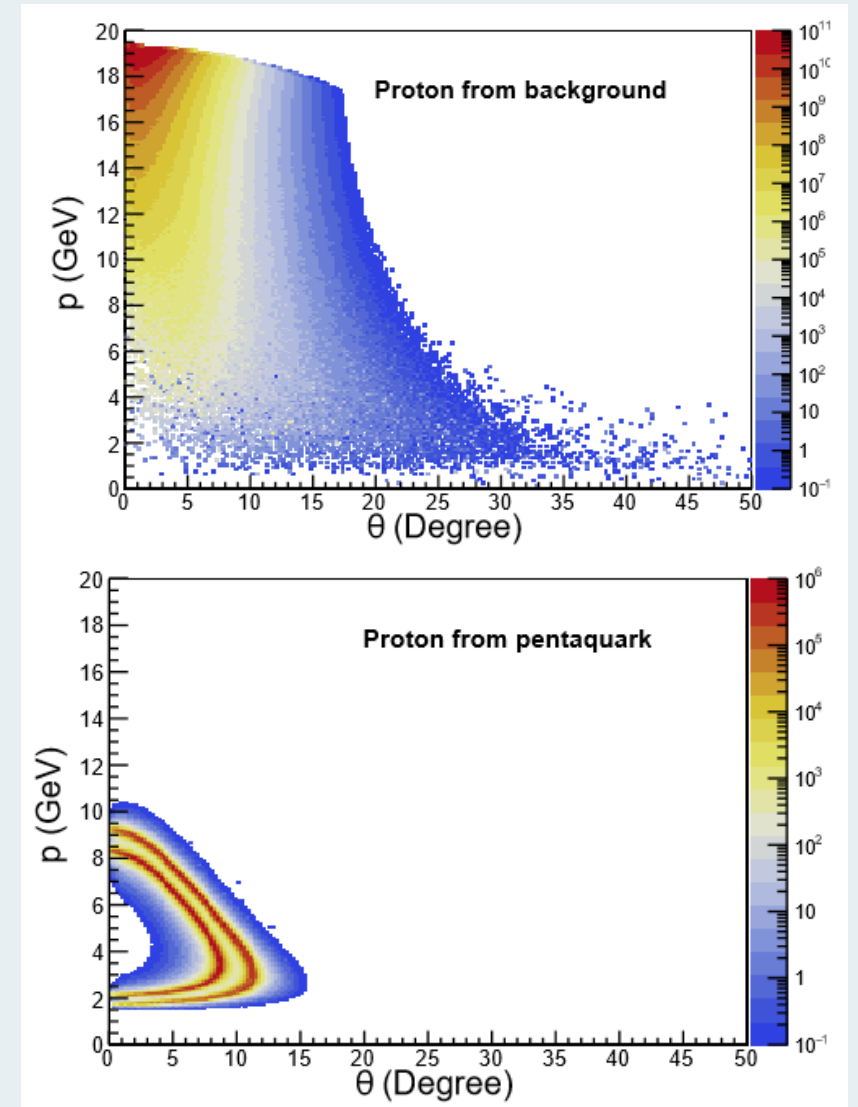
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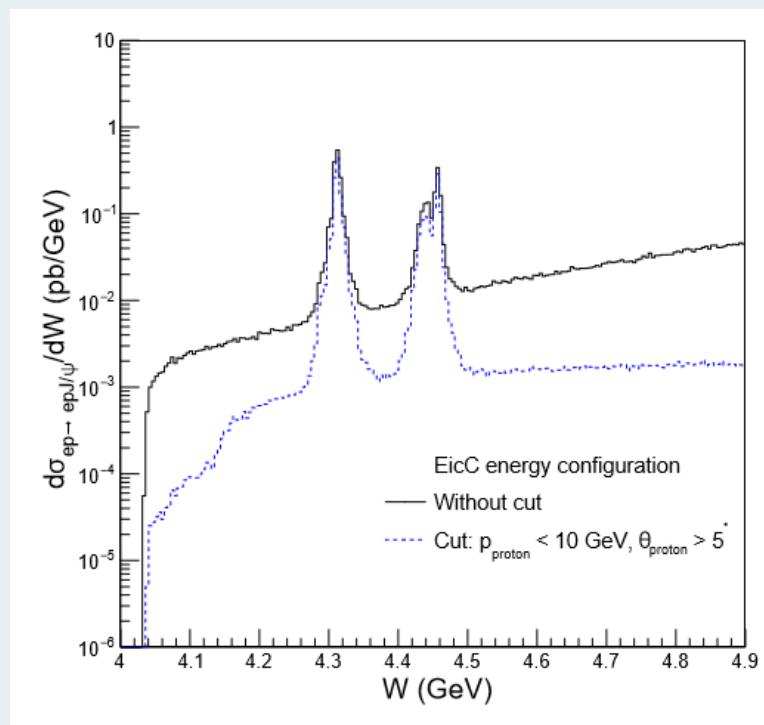
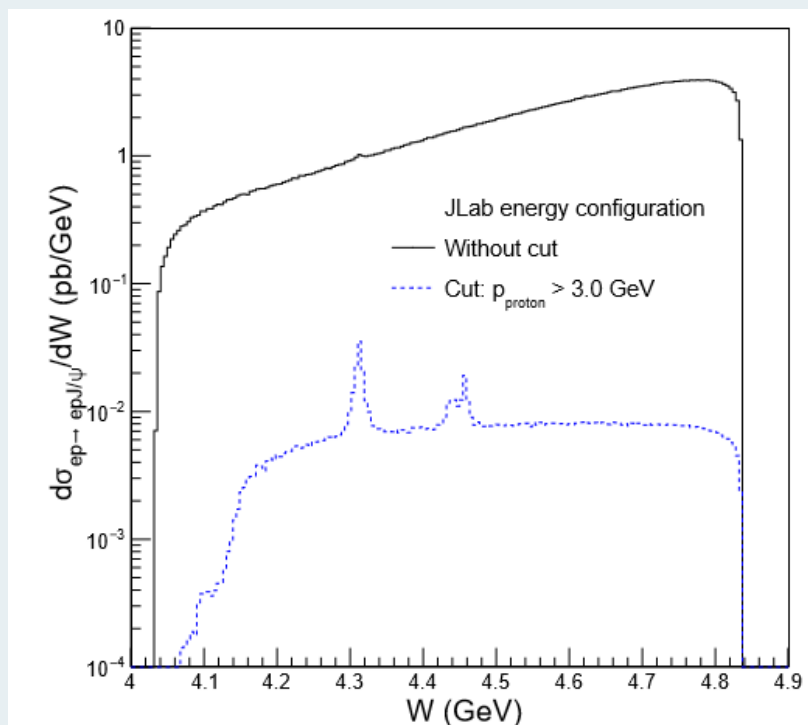
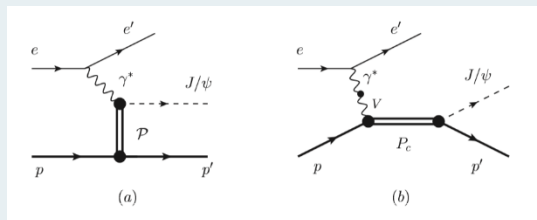
The distribution of momentum versus scattering angles of final proton in the laboratory frame at the but at the **EicC energy configuration, which is 20 GeV proton colliding with 3.5 GeV electron.**

Zhi Yang, Xu Cao, Yu-Tie Liang, Jia-Jun Wu
Chin. Phys. C 44 (2020) 8, 084102



如何从J/ψ p末态中尽可能抽取 P_c 的峰结构?

$$e p \rightarrow e J/\psi p$$



Jlab 能量下可能就是信号道非常小。

而在EicC可能会大一些。

希望将来实验能够进一步研究。

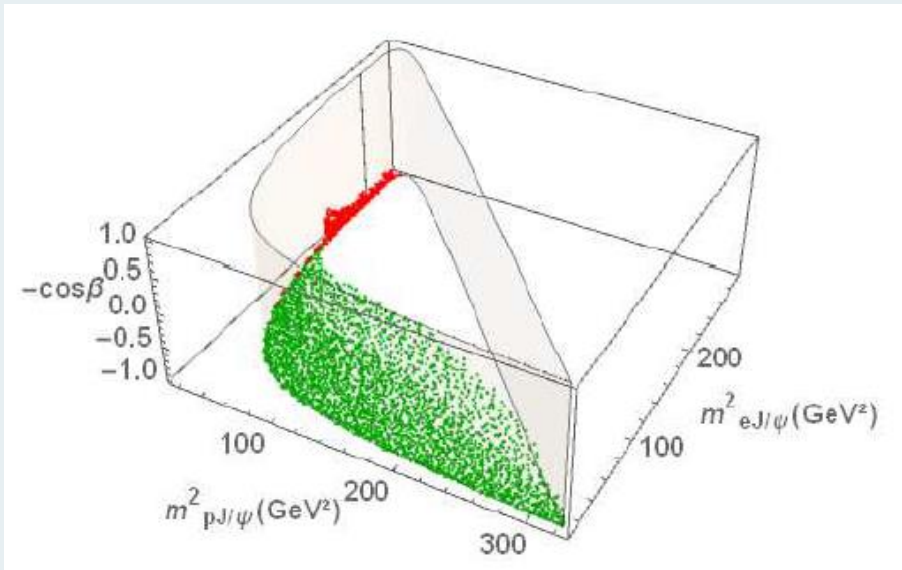
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如何从J/ψ p末态中尽可能抽取 P_c 的峰结构?

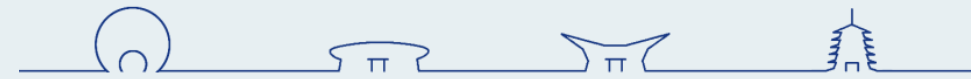


$$e p \rightarrow e J/\psi p$$

三体，四体相空间的解析表达

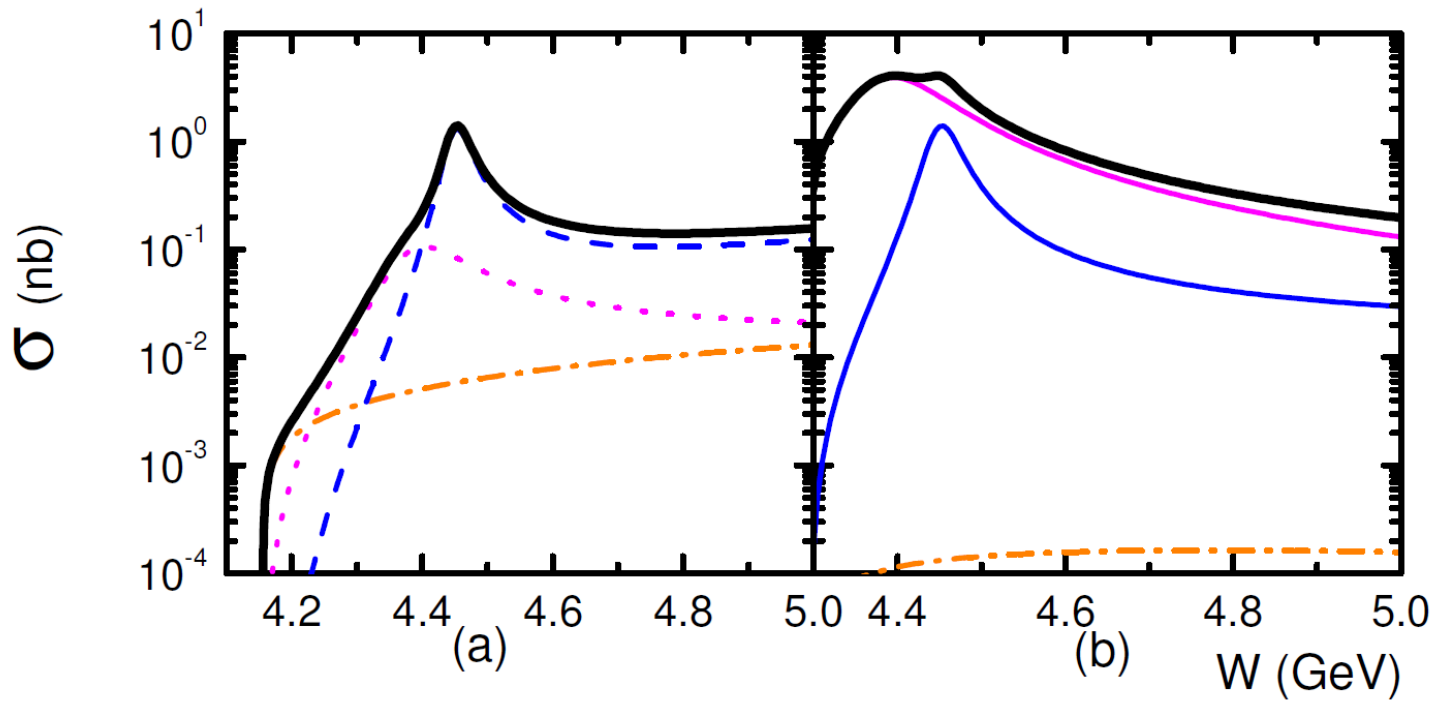
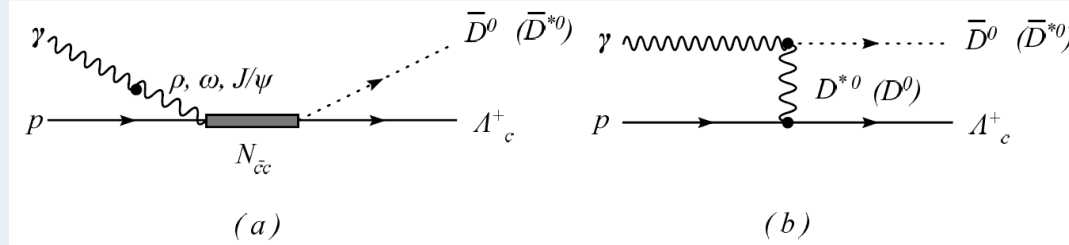
contribution number	(3; 7; 3220)
kinematic variables chosen	$m_{12}^2, m_{13}^2, m_{23}^2, \cos \theta_2^*, \cos \theta_3^*$
phase space factor	$\frac{A_g \mathbf{p}_2^* \mathbf{p}_3^* \lambda(m, m_{123}, m_4)}{(2\pi)^{12} 2^9 m^3 m_{123}^2 \mathbf{p}_1 \mathbf{p}_2 \mathbf{p}_3 }$
intrinsic quantities	$p_1^0 = \frac{1}{2m} (m^2 + m_{123}^2 - m_4^2) - 2m p_2^0 - 2m p_3^0$ $p_2^0 = \gamma_L \sqrt{ \mathbf{p}_2^* ^2 + m_2^2} - \gamma_L \beta_L \mathbf{p}_2^* \cos \theta_2^*$ $p_3^0 = \gamma_L \sqrt{ \mathbf{p}_3^* ^2 + m_3^2} - \gamma_L \beta_L \mathbf{p}_3^* \cos \theta_3^*$ $m_{23}^2 = m_{123}^2 + m_1^2 + m_2^2 + m_3^2 - m_{12}^2 - m_{13}^2$
parameters	$ \mathbf{p}_2^* = \frac{1}{2m_{123}} \lambda^{\frac{1}{2}}(m_{123}, m_{13}, m_2) \quad \mathbf{p}_3^* = \frac{1}{2m_{123}} \lambda^{\frac{1}{2}}(m_{123}, m_{12}, m_3)$ $\gamma_L = \frac{1}{2mm_{123}} (m^2 + m_{123}^2 - m_4^2) \quad \gamma_L \beta_L = \frac{1}{2mm_{123}} \lambda^{\frac{1}{2}}(m, m_{123}, m_4)$

$$d\phi_4 = \delta \left(\sqrt{s} - \sum_{i=1}^4 p_i^0 \right) \delta \left(\sum_{i=1}^4 \vec{p}_i \right) \frac{d^3 \vec{p}_1}{(2\pi)^3 2E_1} \frac{d^3 \vec{p}_2}{(2\pi)^3 2E_2} \frac{d^3 \vec{p}_3}{(2\pi)^3 2E_3} \frac{d^3 \vec{p}_4}{(2\pi)^3 2E_4}$$



$\gamma p \rightarrow \text{other final states}$

- $\gamma p \rightarrow \Lambda_c^+ \bar{D}^0 (\bar{D}^{*0})$



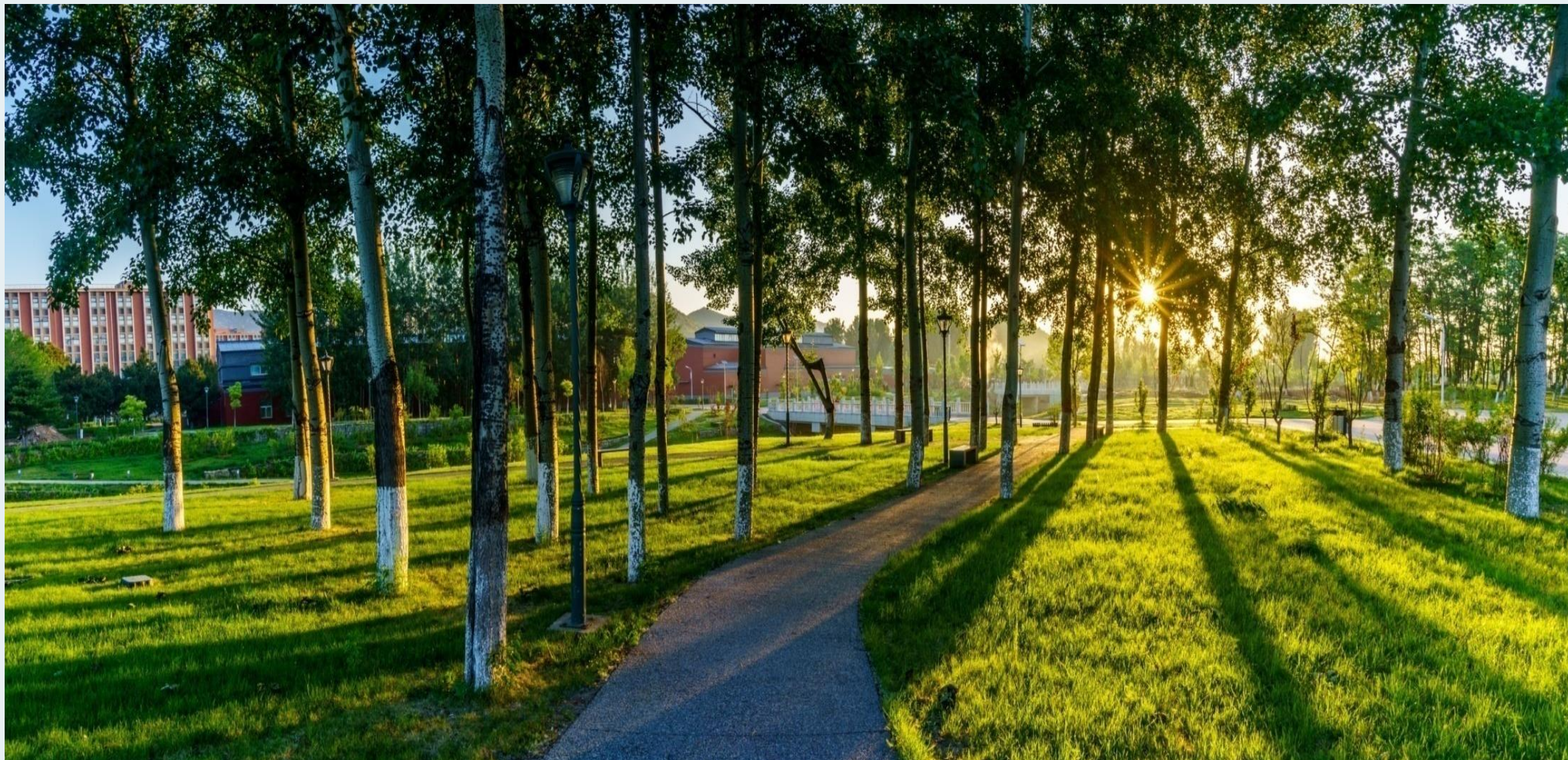
Jia-Jun Wu, T.-S.H. Lee,
Bing-Song Zou
Phys. Rev. C 100 (2019) 3,
035206



小结

- 详细讨论了 $\gamma p \rightarrow P_c \rightarrow J/\psi p$ 的信号以及背景贡献。
- 讨论了如何从 $J/\psi p$ 末态中尽可能抽取 P_c 的峰结构。





谢谢!



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