

# Single-Pion Electroproduction

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[Jülich-Bonn-Washington (JBW) collaboration]

[arXiv: 2104.07312](https://arxiv.org/abs/2104.07312) [nucl-th], Phys. Rev. C, in print



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HPC support by JSC grant *jikp07*

[several slides by  
D. Rönchen and M. Mai]

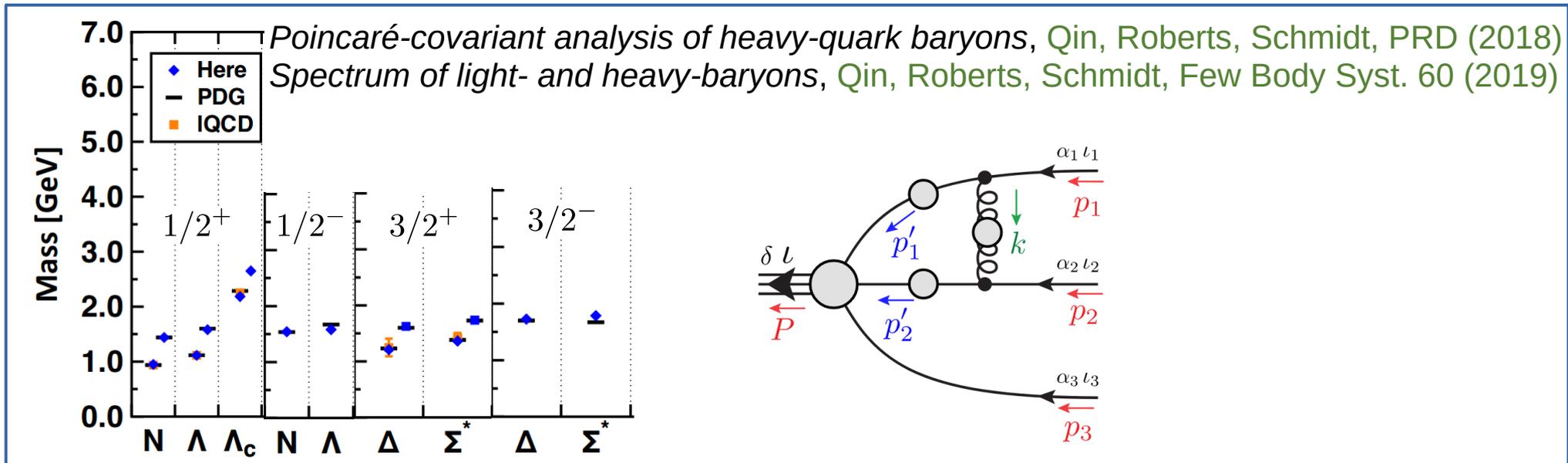
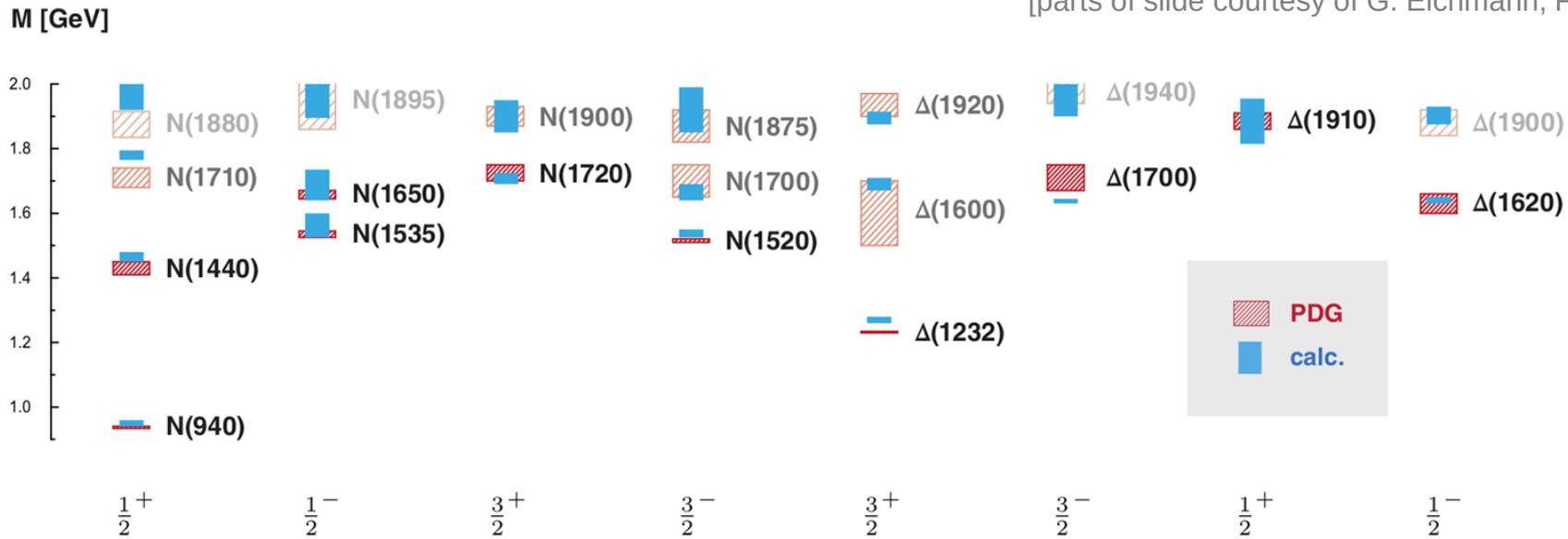
Degrees of freedom: Quarks or hadrons?



# Results in dynamical quark picture

Quark-diquark with reduced pseudoscalar + vector diquarks: [GE, Fischer, Sanchis-Alepuz, PRD 94 \(2016\)](#)

[parts of slide courtesy of G. Eichmann, Few Body 2018]



# Single-meson photoproduction with JuBo

A boundary condition for electroproduction analysis

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e.g.: D. Ronchen et al., EPJA (2018), arXiv: [1801.10458](https://arxiv.org/abs/1801.10458)

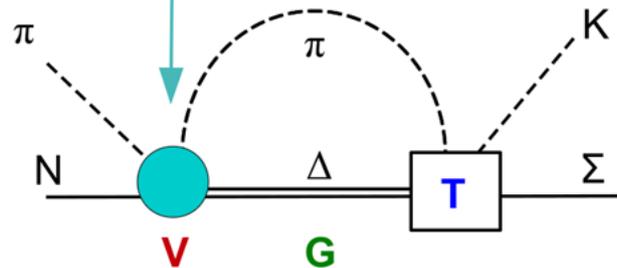
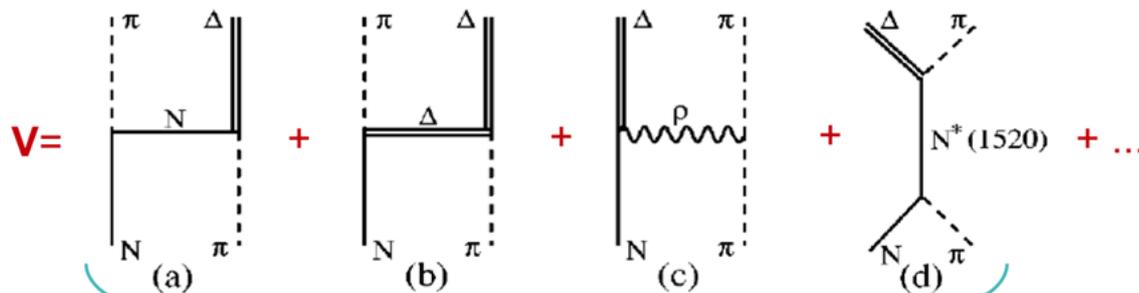
# The Julich-Bonn Dynamical Coupled-Channel Approach

e.g. EPJ A 49, 44 (2013)

**Dynamical coupled-channels (DCC): simultaneous** analysis of different reactions

The scattering equation in partial-wave basis

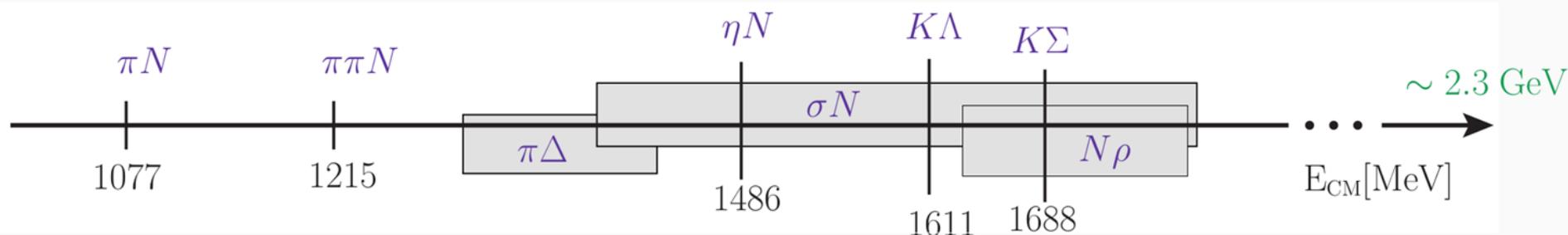
$$\langle L' S' p' | T_{\mu\nu}^{JJ} | L S p \rangle = \langle L' S' p' | V_{\mu\nu}^{JJ} | L S p \rangle + \sum_{\gamma, L'' S''} \int_0^\infty dq q^2 \langle L' S' p' | V_{\mu\gamma}^{JJ} | L'' S'' q \rangle \frac{1}{E - E_\gamma(q) + i\epsilon} \langle L'' S'' q | T_{\gamma\nu}^{JJ} | L S p \rangle$$



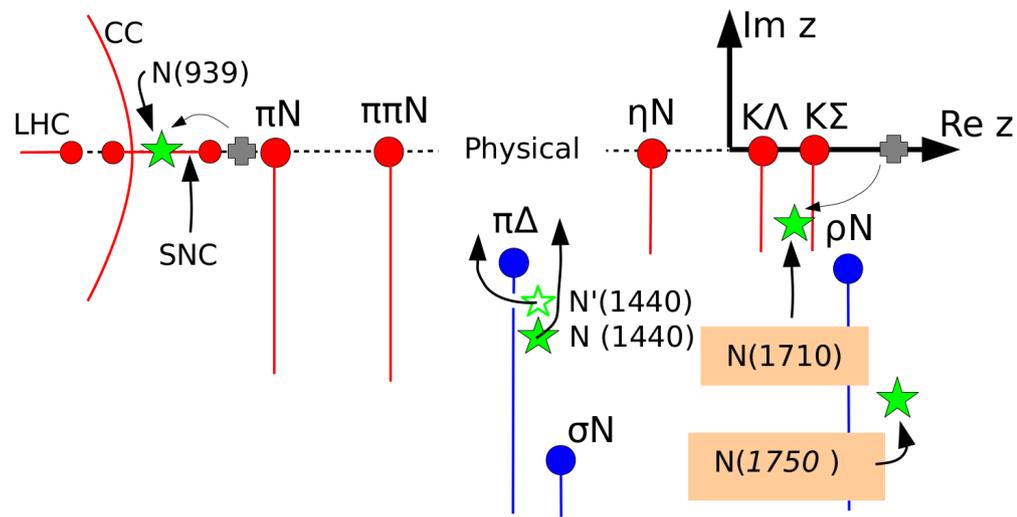
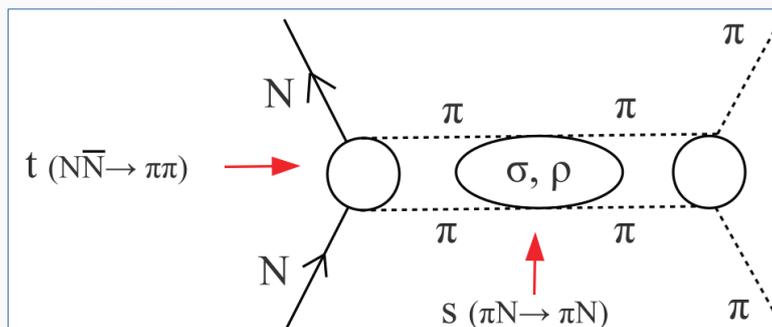
- potentials  $V$  constructed from effective  $\mathcal{L}$
- s-channel diagrams:  $T^P$   
genuine resonance states
- t- and u-channel:  $T^{NP}$   
dynamical generation of poles  
partial waves strongly correlated

# JuBo: Channels and Analytic Structure

Channels included:



- (2-body) unitarity and analyticity respected
  - 3-body  $\pi\pi N$  channel:
    - parameterized effectively as  $\pi\Delta$ ,  $\sigma N$ ,  $\rho N$
    - $\pi N/\pi\pi$  subsystems fit the respective phase shifts
- ↳ branch points move into complex plane



# JuBo: Photoproduction Data base

- $\pi N \rightarrow X$ : > 7,000 data points ( $\pi N \rightarrow \pi N$ : GW-SAID WI08 (ED solution))
- $\gamma N \rightarrow X$ :

Reaction	Observables (# data points)	p./channel
$\gamma p \rightarrow \pi^0 p$	$d\sigma/d\Omega$ (18721), $\Sigma$ (2927), $P$ (768), $T$ (1404), $\Delta\sigma_{31}$ (140), $G$ (393), $H$ (225), $E$ (467), $F$ (397), $C_{x'}$ (74), $C_{z'}$ (26)	25,542
$\gamma p \rightarrow \pi^+ n$	$d\sigma/d\Omega$ (5961), $\Sigma$ (1456), $P$ (265), $T$ (718), $\Delta\sigma_{31}$ (231), $G$ (86), $H$ (128), $E$ (903)	9,748
$\gamma p \rightarrow \eta p$	$d\sigma/d\Omega$ (9112), $\Sigma$ (403), $P$ (7), $T$ (144), $F$ (144), $E$ (129)	9,939
$\gamma p \rightarrow K^+ \Lambda$	$d\sigma/d\Omega$ (2478), $P$ (1612), $\Sigma$ (459), $T$ (383), $C_{x'}$ (121), $C_{z'}$ (123), $O_{x'}$ (66), $O_{z'}$ (66), $O_x$ (314), $O_z$ (314),	5,936
$\gamma p \rightarrow K^+ \Sigma^0$	$d\sigma/d\Omega$ (4271), $P$ (422), $\Sigma$ (280), $T$ (127), $C_{x',z'}$ (188), $O_{x,z}$ (254)	5,542
$\gamma p \rightarrow K^0 \Sigma^+$	$d\sigma/d\Omega$ (242), $P$ (78)	320
	in total	<b>57,027</b>

A new web interface [<https://jbw.phys.gwu.edu/>]

# Pion Electroproduction

A first step towards a coupled-channel photo- and electroproduction analysis

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[M. Mai et al., 2104.07312 \[nucl-th\]](#), Phys. Rev. C, in print

# Single-meson electroproduction to reveal resonance structure

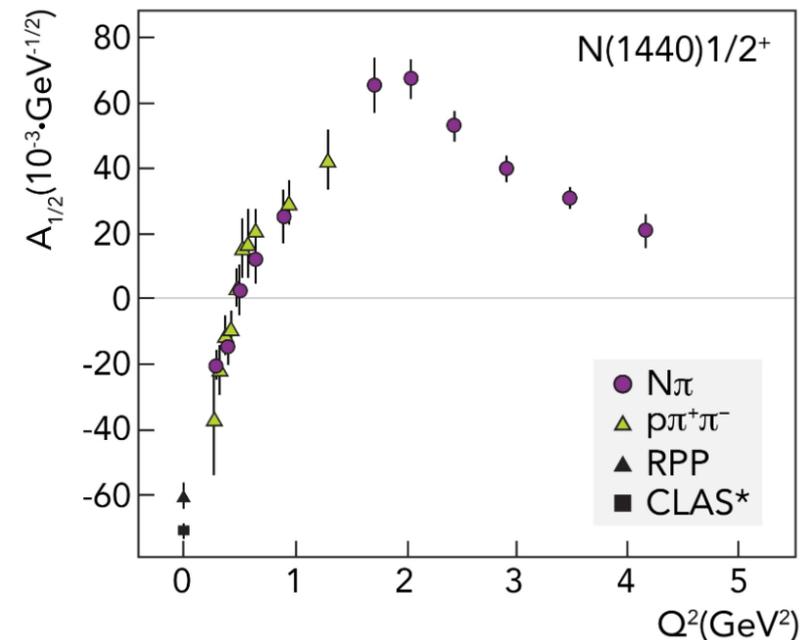
See talks by **D. Carman** and **V. Mokeev**

- **ANL-Osaka PRC 80**, 025207 (2009), Few-Body Syst. 59, 24 (2018),...
- **Aznauryan, Burkert, Mokeev et al.**, PRC 80, 055203 (2009), Int. J. Mod. Phys. E22, 1330015 (2013),...
- **EtaMAID2018**, EPJA 54 (2018), 210
- **MAID2007**, EPJA 34 (2007) 69
- **SAID**, PiN Newsletter 16, 150 (2002)
- **Gent group** Phys. Rev. C 89, 065202 (2014),...

## Highlights:

- Simultaneous description of pion photo- and electroproduction (MAID)
- Consistent extraction of the Roper form factor from single and double pion electroproduction
- New resonance in electroproduction claimed  
Mokeev et al., PLB (2020) [2004.13531 \[nucl-ex\]](#)

Burkert, Roberts, [1710.02549](#)



# Needed: Coupled-channel electroproduction analysis

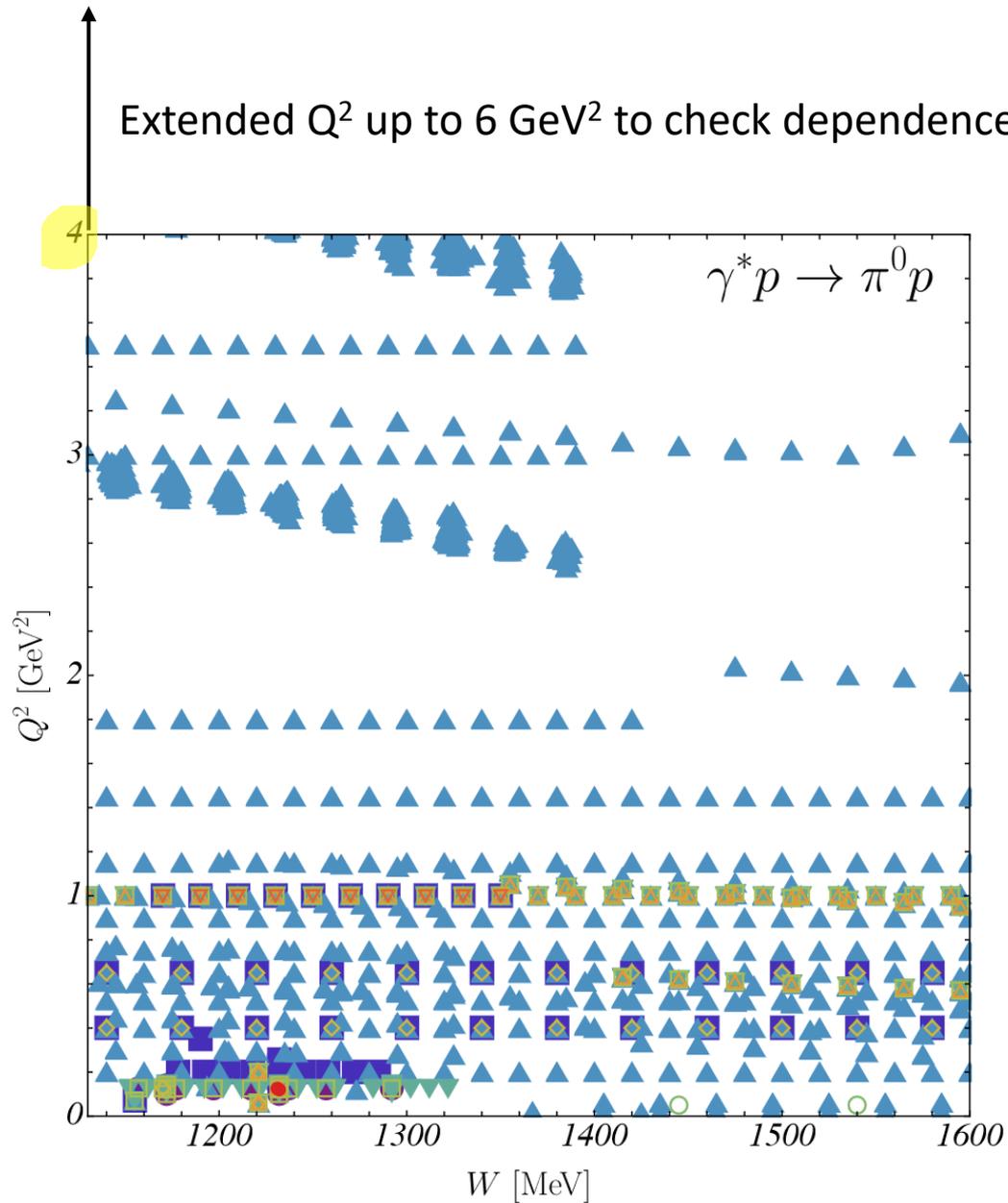
Take advantage of multi-channel approach  
 → analyze simultaneously final states  $\pi N$ ,  $\eta N$ ,  $K\Lambda$   
 ~ $10^6$  pion electroproduction data;  $\eta N$ ,  $K\Lambda$  :

Reaction	Observable	$Q^2$ [GeV]	W [GeV]	Ref.
$ep \rightarrow e'p'\eta$	$\sigma_U, \sigma_{LT}, \sigma_{TT}$	1.6 – 4.6	2.0 – 3.0	[132]
	$\sigma_U, \sigma_{LT}, \sigma_{TT}$	0.13 – 3.3	1.5 – 2.3	[137]
	$d\sigma/d\Omega$	0.25 – 1.5	1.5 – 1.86	[138]
$ep \rightarrow e'K^+\Lambda$	$P_N^0$	0.8 – 3.2	1.6 – 2.7	[139]
	$\sigma_U, \sigma_{LT}, \sigma_{TT}, \sigma_{LT}'$	1.4 – 3.9	1.6 – 2.6	[140]
	$P'_x, P'_z$	0.7 – 5.4	1.6 – 2.6	[141]
	$\sigma_T, \sigma_L, \sigma_{LT}, \sigma_{TT}$	0.5 – 2.8	1.6 – 2.4	[142]
	$P'_x, P'_z$	0.3 – 1.5	1.6 – 2.15	[143]

**Table 1:** Overview of  $\eta p$  and  $K^+\Lambda$  electroproduction data measured at CLAS for different photon virtualities  $Q^2$  and total energy  $W$ . Based on material provided by courtesy of D. Carman (JLab) and I. Strakovsky (GW).

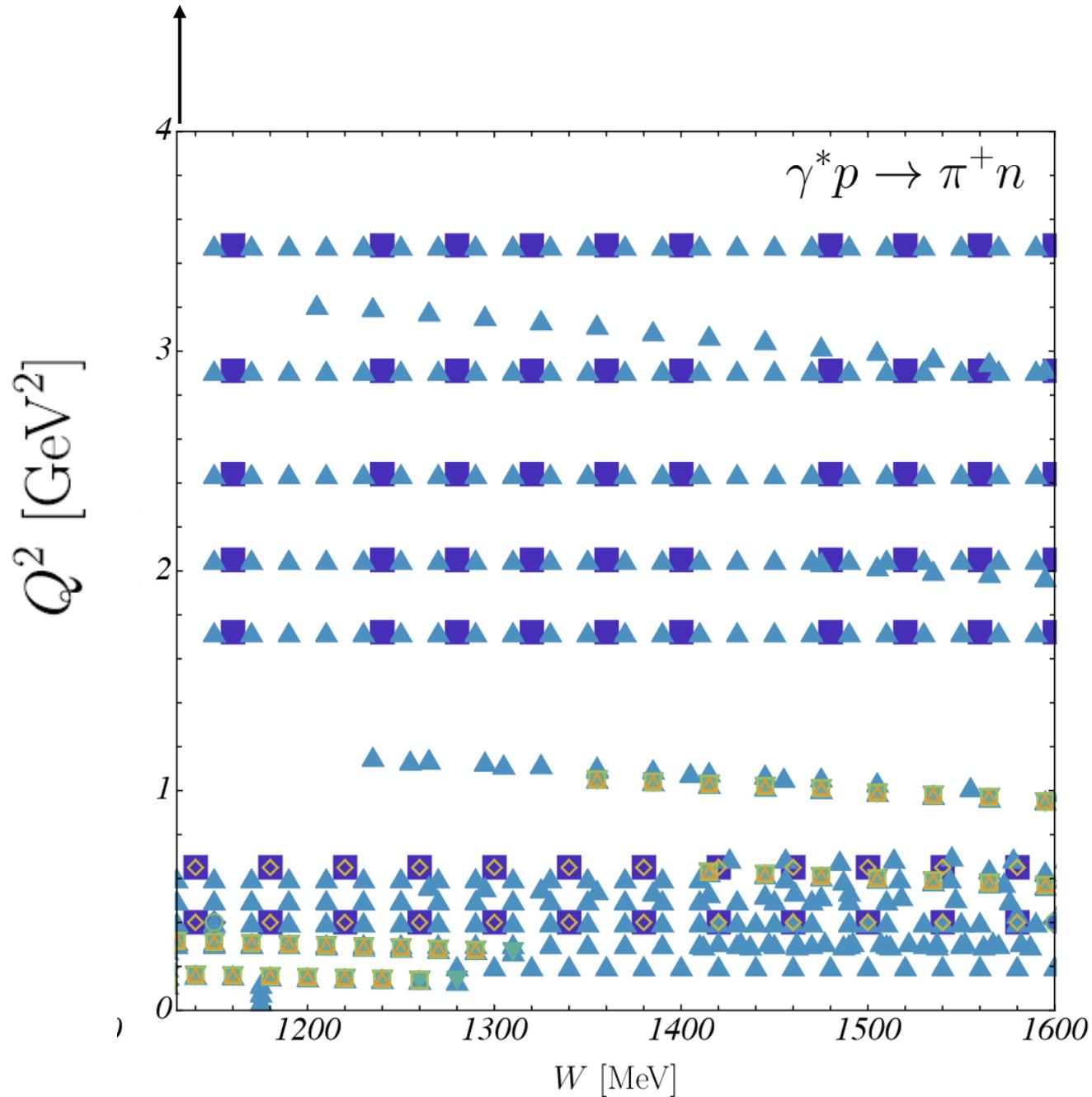
- Many of these (and similar) data await analysis.
- Many more data to emerge at Jlab ( $Q^2 = 5 - 12 \text{ Gev}^2$ )  
 e.g.: Carman, Joo, Moiseev, *Few Body Syst.* 61, 29 (2020)
- Approved Jlab experiments to study
  - Higher-lying nucleon resonances
  - Hybrid baryons
  - Transition regime between nonperturbative and perturbative regions

# Pion Electroproduction – data base



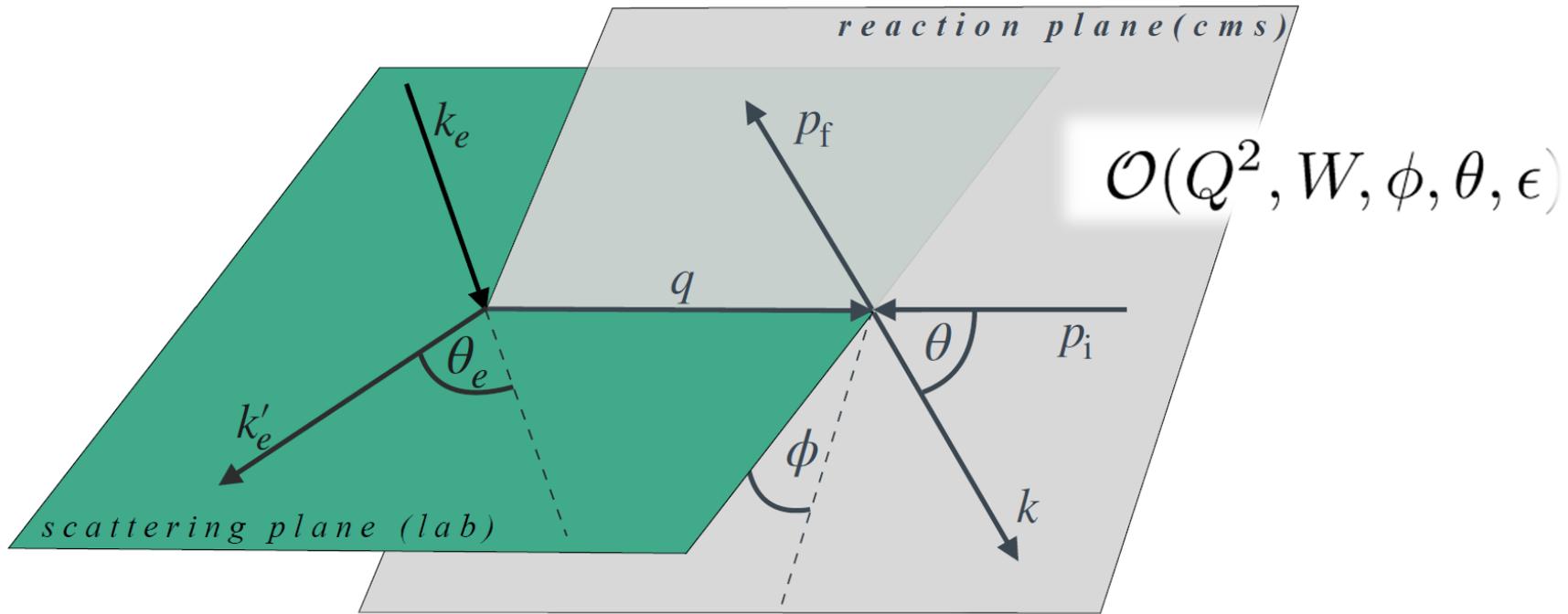
Type	$N_{\text{data}}$
$\rho_{LT}$	45
$\rho_{LT'}$	2644
$\sigma_L$	–
$d\sigma/d\Omega$	39942
$\sigma_T + \epsilon\sigma_L$	318
$\sigma_T$	10
$\sigma_{LT}$	312
$\sigma_{LT'}$	198
$\sigma_{TT}$	266
$K_{D1}$	1527
$P_Y$	2

# Pion Electroproduction – data base



Type	$N_{\text{data}}$
<span style="color: purple;">●</span> $\rho_{LT}$	–
<span style="color: blue;">■</span> $\rho_{LT'}$	4354
<span style="color: blue;">◆</span> $\sigma_L$	2
<span style="color: blue;">▲</span> $d\sigma/d\Omega$	32813
<span style="color: teal;">▼</span> $\sigma_T + \epsilon\sigma_L$	144
<span style="color: cyan;">○</span> $\sigma_T$	2
<span style="color: lightgreen;">□</span> $\sigma_{LT}$	106
<span style="color: yellow;">◇</span> $\sigma_{LT'}$	192
<span style="color: orange;">△</span> $\sigma_{TT}$	91
<span style="color: orange;">▽</span> $K_{D1}$	–
<span style="color: red;">●</span> $P_Y$	–

# Kinematics



$$\frac{d\sigma}{d\Omega_f dE_f d\Omega} = \left( \frac{\alpha}{2\pi^2} \frac{E_f}{E_i} \frac{q_L}{Q^2} \frac{1}{1-\epsilon} \right) \frac{d\sigma^v}{d\Omega},$$

(Un)polarized differential cross section:

$$\begin{aligned} \frac{d\sigma^v}{d\Omega} = & \sigma_T + \epsilon\sigma_L + \sqrt{2\epsilon(1+\epsilon)}\sigma_{LT} \cos\phi \\ & + \epsilon\sigma_{TT} \cos 2\phi + h\sqrt{2\epsilon(1-\epsilon)}\sigma_{LT'} \sin\phi \end{aligned}$$

$$\epsilon = 1 + 2 \frac{q_L^2}{Q^2} \tan^2 \frac{\theta_e}{2}$$

# Polarized Observables

- CLAS: Structure functions  $\sigma_{LT'}$   
 K. Joo et al. [CLAS], [Phys. Rev. C 68 \(2003\)](#),  
 K. Joo et al. [CLAS], [Phys. Rev. C 70 \(2004\)](#).
- Jlab-Hall A for  $K_{1D} = \{K_{1D}^X | X = A, B, \dots, T\}$   
 J. J. Kelly, [Phys. Rev. Lett. 95 \(2005\)](#).
- Response functions (R)  $\Leftrightarrow$  Kelly notation (RL, RT, ...)  $\Leftarrow$  Helicity amplitudes  $H \Leftrightarrow$  CGNL amplitude. For example:

$$\sigma_T = \frac{k}{q_\gamma} R_T^{00}, \quad \sigma_L = \frac{k}{q_\gamma} \frac{Q^2}{\omega^2} R_L^{00}, \quad \sigma_{TT} = \frac{k}{q_\gamma} R_{TT}^{00}$$

$$\sigma_{LT} = \frac{k}{q_\gamma} \frac{\sqrt{Q^2}}{\omega} R_{LT}^{00}, \quad \sigma_{LT'} = \frac{k}{q_\gamma} \frac{\sqrt{Q^2}}{\omega} R_{LT'}^{00}.$$

$$P_Y = -\sqrt{2\epsilon(1+\epsilon)} \frac{\omega}{\sqrt{Q^2}} \frac{R_{LT}^{0Y}}{R_T^{00} + \epsilon\omega^2/Q^2 R_L^{00}}$$

$$\rho_{LT} = \sqrt{2\epsilon(1+\epsilon)} \frac{R_{LT}^{00}}{R_T^{00} + \epsilon(R_L^{00} + R_{TT}^{00})},$$

$$\rho_{LT'} = \sqrt{2\epsilon(1-\epsilon)} \sin \phi \frac{\sigma_{LT'}}{d\sigma^v/d\Omega},$$

# Parameterization

- Photoproduction solution as constraint
- Constraints from (Pseudo)-threshold:

$$\begin{aligned}
 (E_{l+}^I, L_{l+}^I) &\rightarrow k^l q^l \quad (l \geq 0) \\
 (M_{l+}^I, M_{l-}^I) &\rightarrow k^l q^l \quad (l \geq 1) \\
 (L_l^I) &\rightarrow kq \quad (l = 1) \\
 (E_{l-}^I, L_{l-}^I) &\rightarrow k^{l-2} q^l \quad (l \geq 2)
 \end{aligned}
 \quad
 \begin{aligned}
 k = |\mathbf{k}| &= \frac{\sqrt{\left((W - M_N)^2 + Q^2\right) \left((W + M_N)^2 + Q^2\right)}}{2W} \\
 q = |\mathbf{q}| &= \frac{\sqrt{\left((W - M_N)^2 - M_m^2\right) \left((W + M_N)^2 - M_m^2\right)}}{2W}
 \end{aligned}$$

- Siegert's theorem at pseudo-threshold:

$$\frac{E_{l+}}{L_{l+}} \rightarrow 1, \quad \frac{E_{l-}}{L_{l-}} \rightarrow \frac{-l}{l-1}$$

Amaldi, Fubini, Furlan,  
 Springer Tracts Mod. Phys. 83, 1 (1979)  
 Tiator, Few-body Systems 57, 1087 (2016)

- Watson's theorem, multi-channel unitarity

$$\begin{aligned}
 M_{\mu\gamma^*}(q, W, Q^2) &= V_{\mu\gamma^*}(q, W, Q^2) + \sum_{\kappa} \int dp p^2 T_{\mu\kappa}(q, p, W) G_{\kappa}(p, W) V_{\nu\gamma^*}(p, W, Q^2) \\
 V_{\mu\gamma^*}(p, W, Q^2) &= \alpha_{\mu\gamma^*}^{NP}(p, W, Q^2) + \sum_i \frac{\gamma_{\mu;i}^a(p) \gamma_{\gamma^*;i}^c(W, Q^2)}{W - m_i^b}
 \end{aligned}$$

# Parameterization (2)

- Up to  $D$ -waves included (photoproduction part includes up to  $J=9/2$ )
- Energy range up to  $W \approx 1.6$  allows to include  $\eta N$  electro-production without much extra effort, but  $KY$  electroproduction requires additional work
- Final state interaction given by JuBo/JBW model such that pole positions and hadronic branching ratios (pole residues) are universal as required by reaction dynamics

- $Q^2$ -dependence: Several analytic forms tested; settled for:

$$\tilde{F}(Q^2) = \tilde{F}_D(Q^2) e^{-\beta_0 Q^2/m^2} P^N(Q^2/m^2)$$

where P<sup>N</sup>: Polynomial

$$\tilde{F}_D(Q^2) = \frac{1}{(1 + Q^2/b^2)^2} \frac{1 + e^{-Q_r^2/Q_w^2}}{1 + e^{(Q^2 - Q_r^2)/Q_w^2}}$$

- Some multipoles difficult to determine (longitudinal more difficult than E and M; sometime not even Siegert's condition helps because corresponding electric multipole does not exist)
- But: No model-dependent input from (photonic) Feynman diagrams to model longitudinal multipoles

# Parameterization Dependence

- Can parametrization dependence be avoided? Not if the data is far from being complete enough to represent even a truncated complete electroproduction experiment

L. Tiator et al. Phys. Rev. C (2017), [arXiv: 1702.08375](#)

- Future: Bias-variance tradeoff: Different statistical criteria (Akaike, Bayesian) to find sweet spot between no. of parameters or no. of partial waves and predictivity (model selection)

J. Landay et al., Phys.Rev.C (2017), [arXiv: 1610.07547](#)

- Future: Single- $Q^2$  analysis can decrease parametrization-independence but not remove it (discrete & continuous ambiguities).
- Towards complete data: CLAS/Kelly data provides unique opportunity to confront parametrization with different polarization data at given  $W$  and  $Q^2$ .

J. J. Kelly, [Phys. Rev. Lett. 95 \(2005\)](#).

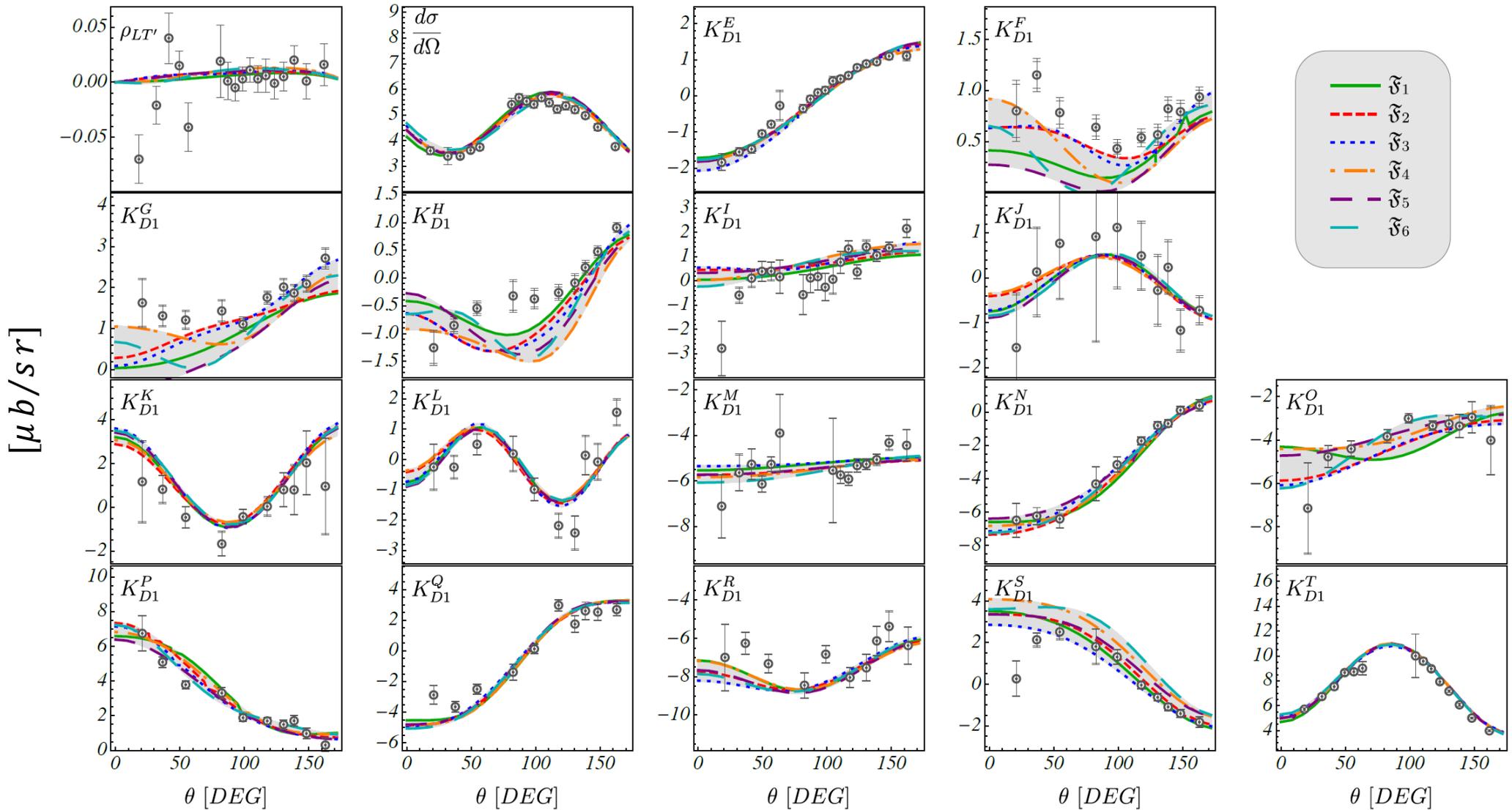
# Results (1): Fit Strategies

- Six different fit strategies:
  - Avoid fitting structure function if corresponding cross sections can be fitted (respect data correlations)
  - Sequential  $S \rightarrow S+P \rightarrow S+P+D$  waves;
  - Subsets of data until full data set reached
  - Simultaneous fit all parameters (209) set to zero without any (!) guidance
  - Extend data range from  $0 < Q^2 < 4 \text{ GeV}^2$  to  $0 < Q^2 < 6 \text{ GeV}^2$  to check for stability

Fit	$\sigma_L$		$d\sigma/d\Omega$		$\sigma_T + \epsilon\sigma_L$		$\sigma_T$		$\sigma_{LT}$		$\sigma_{LT'}$		$\sigma_{TT}$		$K_{D1}$		$P_Y$		$\rho_{LT}$		$\rho_{LT'}$		$\chi^2_{\text{dof}}$
	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	$\pi^0 p$	$\pi^+ n$	
$\mathfrak{F}_1$	-	9	65355	53229	870	418	87	88	1212	133	862	762	4400	251	4493	-	234	-	525	-	3300	10294	1.77
$\mathfrak{F}_2$	-	4	69472	55889	1081	619	65	78	1780	150	1225	822	4274	237	4518	-	325	-	590	-	3545	10629	1.69
$\mathfrak{F}_3$	-	8	66981	54979	568	388	84	95	1863	181	1201	437	3934	339	4296	-	686	-	687	-	3556	9377	1.81
$\mathfrak{F}_4$	-	22	63113	52616	562	378	153	107	1270	146	1198	1015	4385	218	5929	-	699	-	604	-	3548	11028	1.78
$\mathfrak{F}_5$	-	20	65724	53340	536	528	125	81	1507	219	1075	756	4134	230	5236	-	692	-	554	-	3580	11254	1.81
$\mathfrak{F}_6$	-	18	71982	58434	1075	501	29	68	1353	135	1600	1810	3935	291	5364	-	421	-	587	-	3932	11475	1.78

$\chi^2$

# Results (2): Kelly data



$\pi^0 p$ ,  $Q^2=1 \text{ GeV}^2$ ,  $W=1.23 \text{ GeV}$ ,  $\phi=15^\circ$

J. J. Kelly, [Phys. Rev. Lett. 95 \(2005\)](#).

# Results (3): Structure Functions (Selection)

$[\mu b/sr]$

$W = 1.14$  GeV

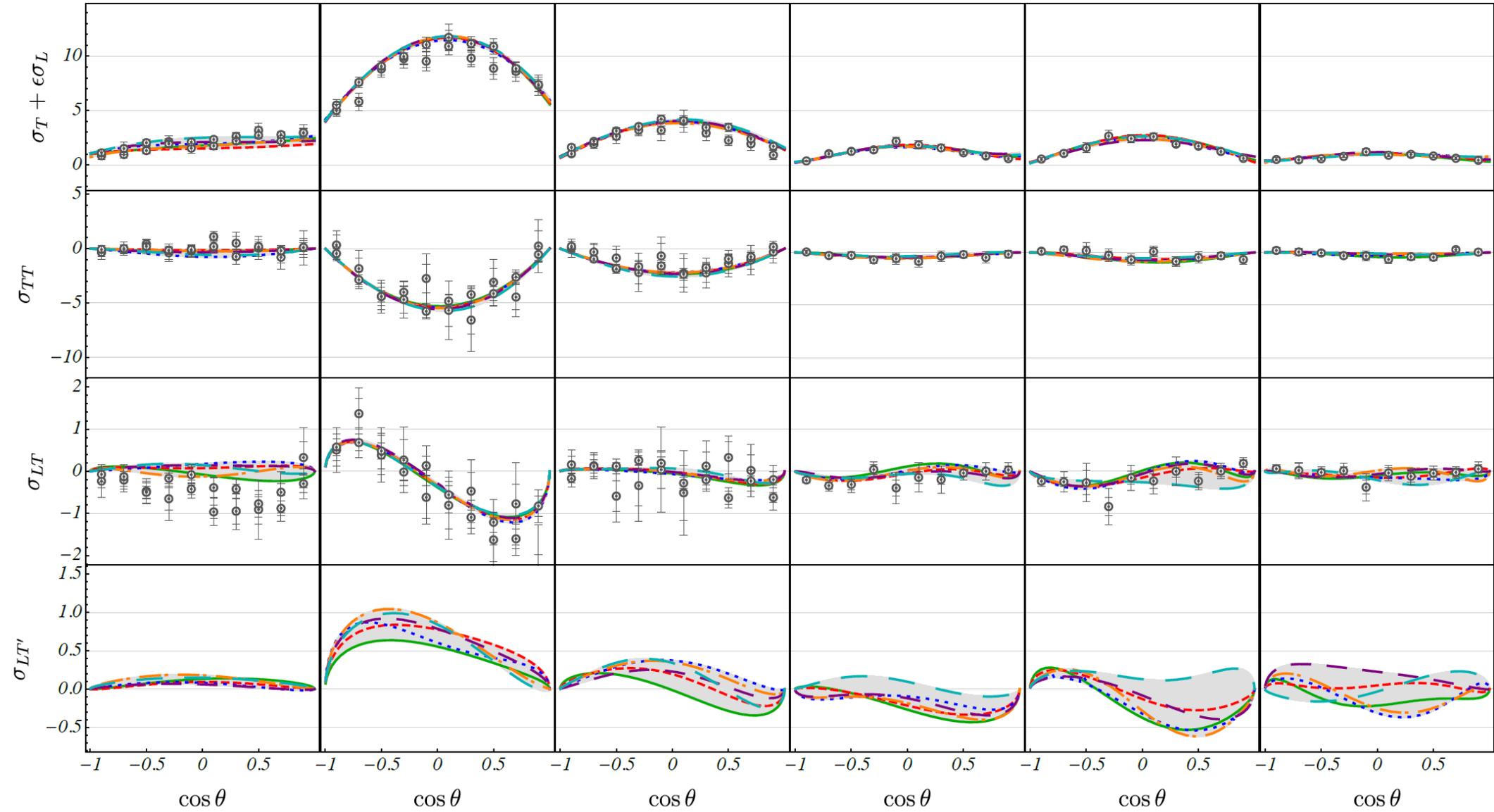
$W = 1.22$  GeV

$W = 1.30$  GeV

$W = 1.42$  GeV

$W = 1.50$  GeV

$W = 1.58$  GeV

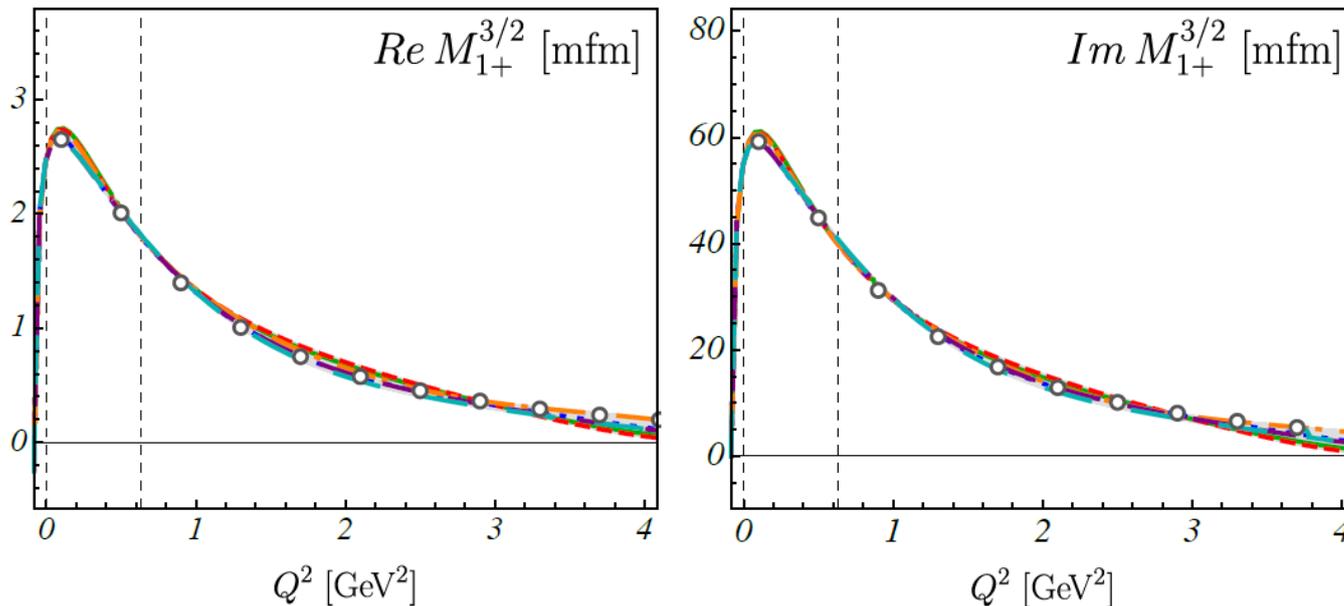


$Q^2 = 0.9 \text{ GeV}^2, \pi^0 p$

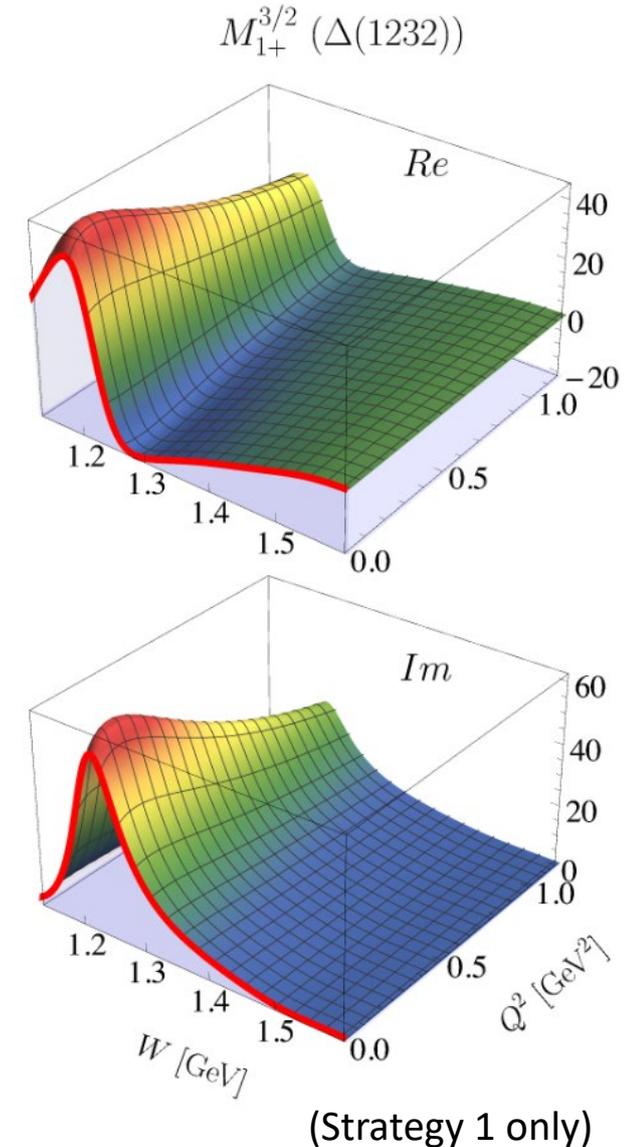
data: CLAS, Phys. Rev. C (2003) [0301012 \[nucl-ex\]](#), Phys. Rev. Lett. (2002) [0110007 \[hep-ex\]](#)

# Results (4): Large Multipoles

**Prominent multipoles are well determined**, even with significantly different fit strategies (e.g., all parameters initially set to zero, no guidance for fit!)

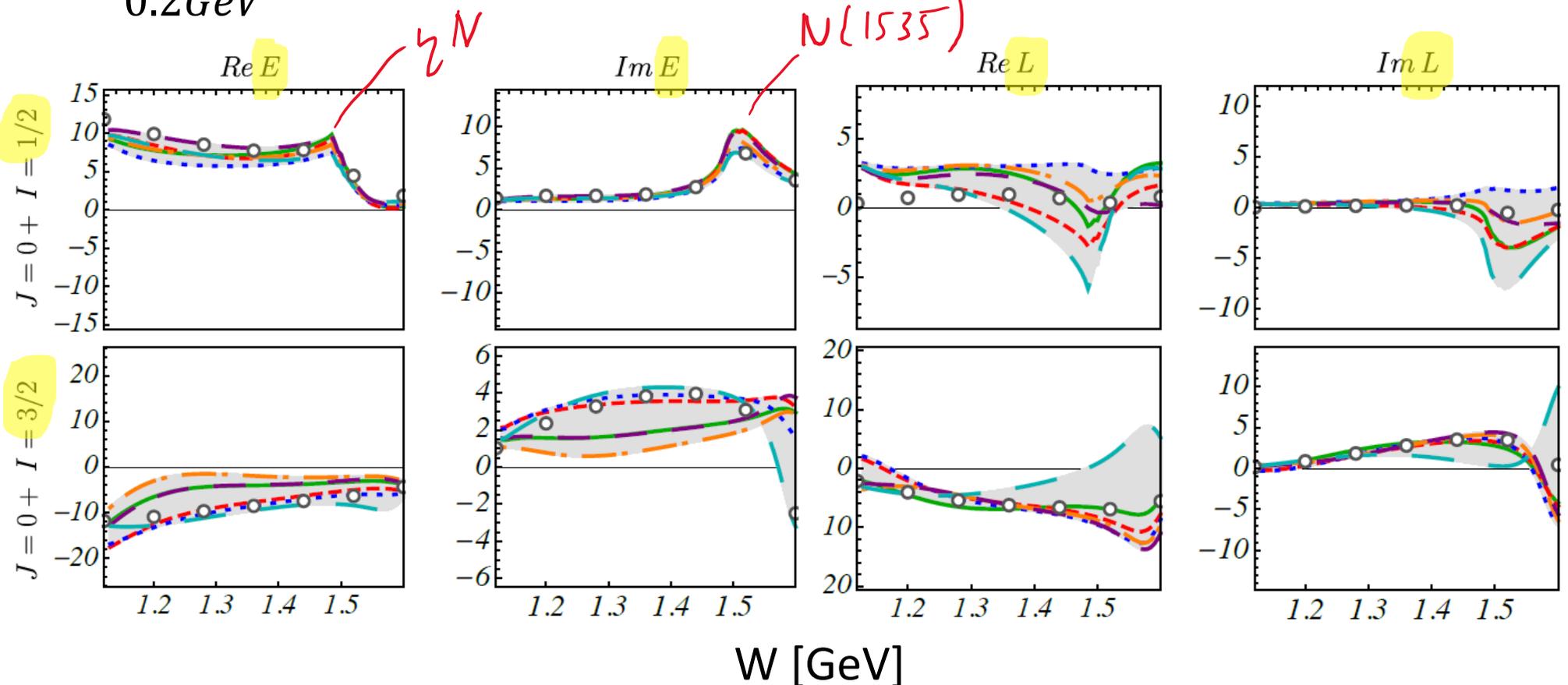


Fit strategies 1-6 together with MAID (open dots) for the magnetic multipole of the  $\Delta(1232)$  [Drechsel et al., EPJA \(2007\) 0710.0306 \[nucl-th\]](#)



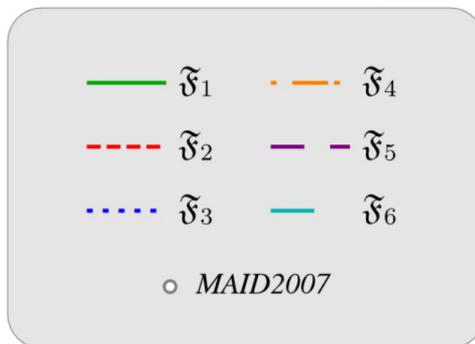
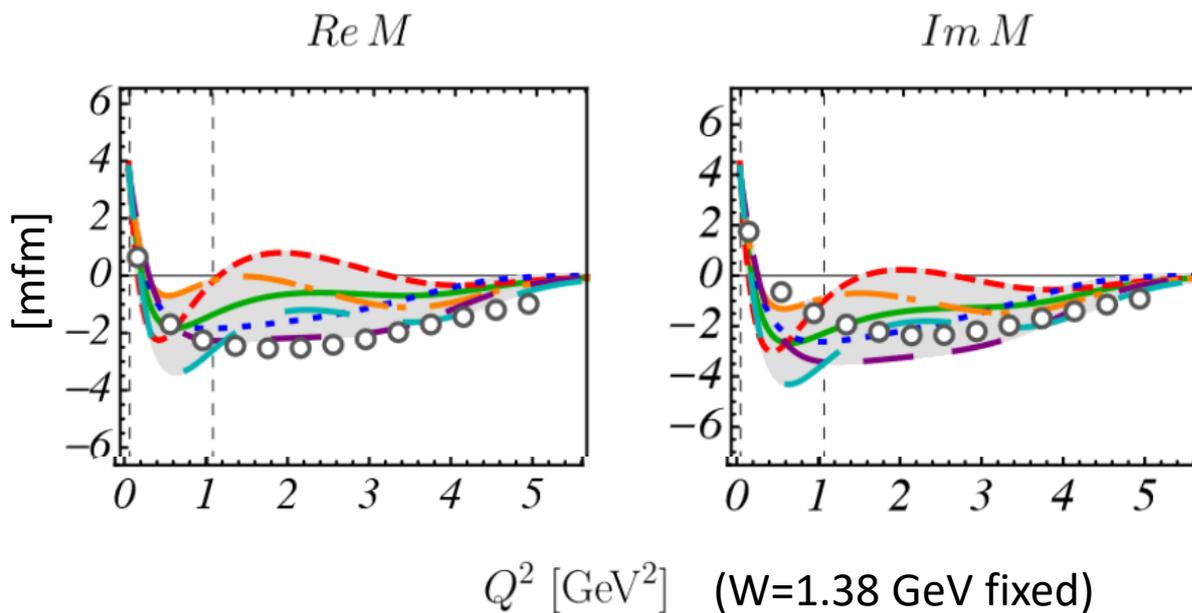
# Results (5): Other multipoles

- Less prominent multipoles are sometimes less well determined
- Overall: solutions are still surprisingly close together given vastly different strategies
- Differences from various strategies (different local  $\chi^2$  minima) much larger than statistical uncertainties; larger than typical MAID uncertainties.
- **Example:** S-wave multipoles [*mfm*] as function of energy  $W$  at fixed  $Q^2 = 0.2\text{GeV}^2$

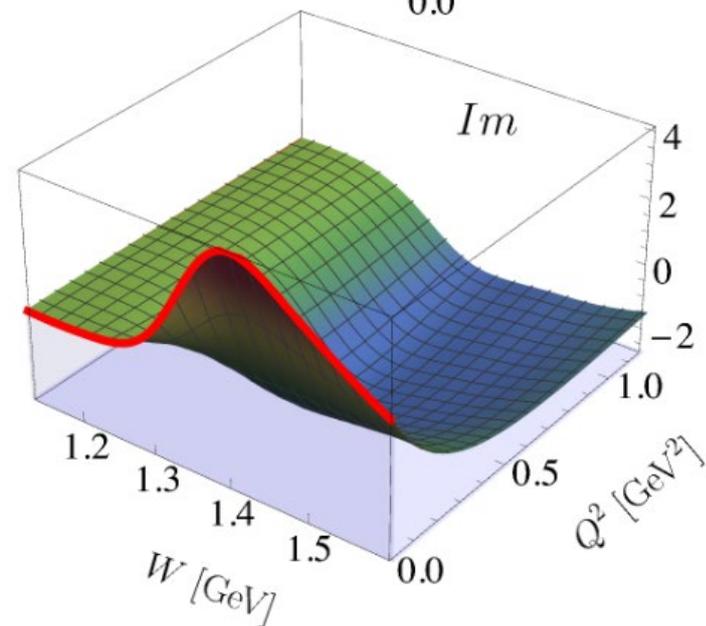
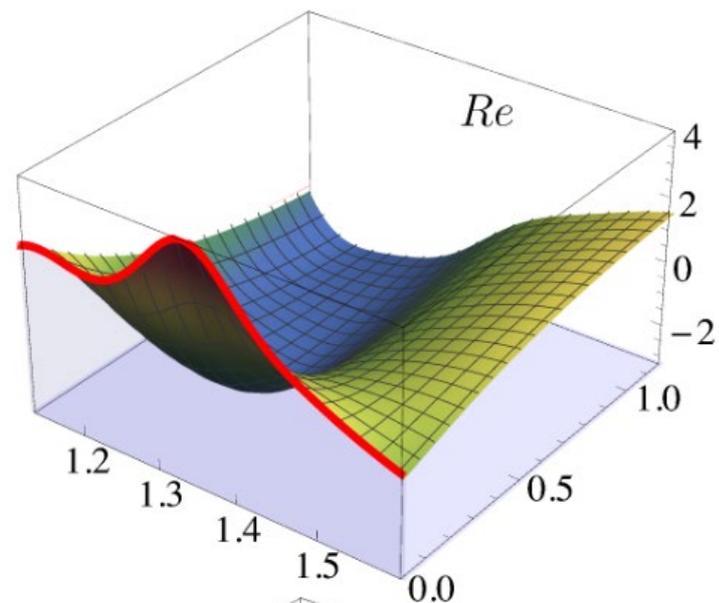


# Results (6): Roper Multipole

- Non-trivial structure
- Zero transition
- Helicity coupling still to be extracted



$M_{1-}^{1/2} (N(1440))$



(Strategy 1 only)

# Summary

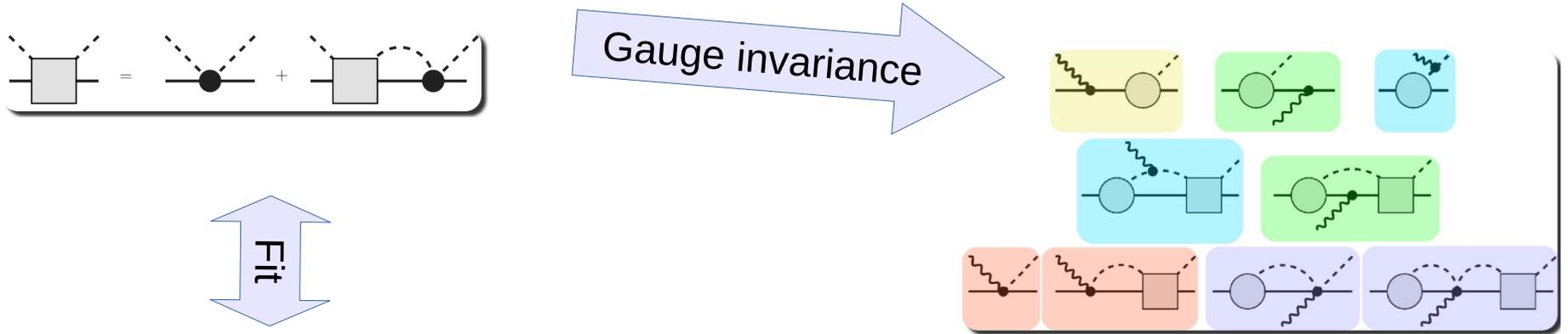
- JBW model: Phenomenology of excited baryons through coupled-channels, two- and three-body effects
  - Analysis finds/confirms new states in analysis of photo-production data, renewed effort to explore additional reaction channels
  - Pion electroproduction analysis performed
    - Exploration of parameter space through different fit strategies reveals different local minima leading to significantly different multipole content.
    - Yet, prominent multipole well determined, albeit with uncertainties larger than in other analyses.
- 
- Extraction of helicity couplings and fixed- $Q^2$  analysis planned
  - Upgrade to  $\eta$  and KY electroproduction straightforward (existing and future JLab data; photoproduction solution exists)
  - Statistical upgrade: How to find a minimal resonance spectrum through model selection J. Landay et al., Phys.Rev.D (2019), [1810.00075 \[nucl-th\]](#)

(spare slides)

Using ONLY meson-baryon degrees of freedom (no explicit quark dynamics):

# Manifestly gauge invariant approach based on full BSE solution

[Ruic, M. Mai, U.-G. Meissner PLB 704 (2011)]



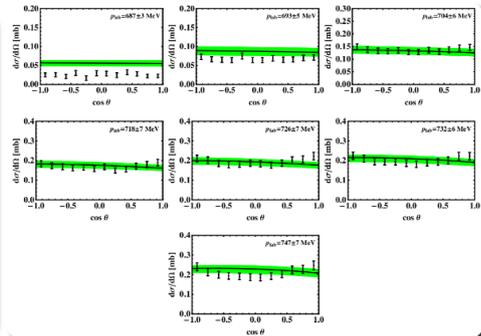
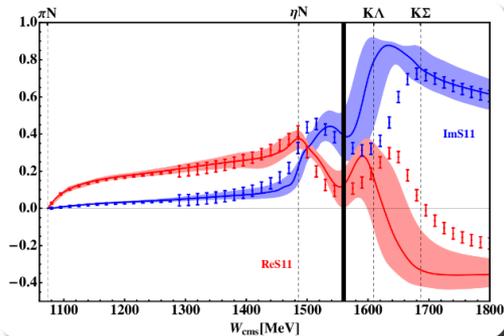
► Exact unitary meson-baryon scattering amplitude  $T$  with parameters, fixed to reproduce:

►  $\pi N$ -partial wave  $S_{11}$  and  $S_{31}$  for  $\sqrt{s} < 1560$  MeV

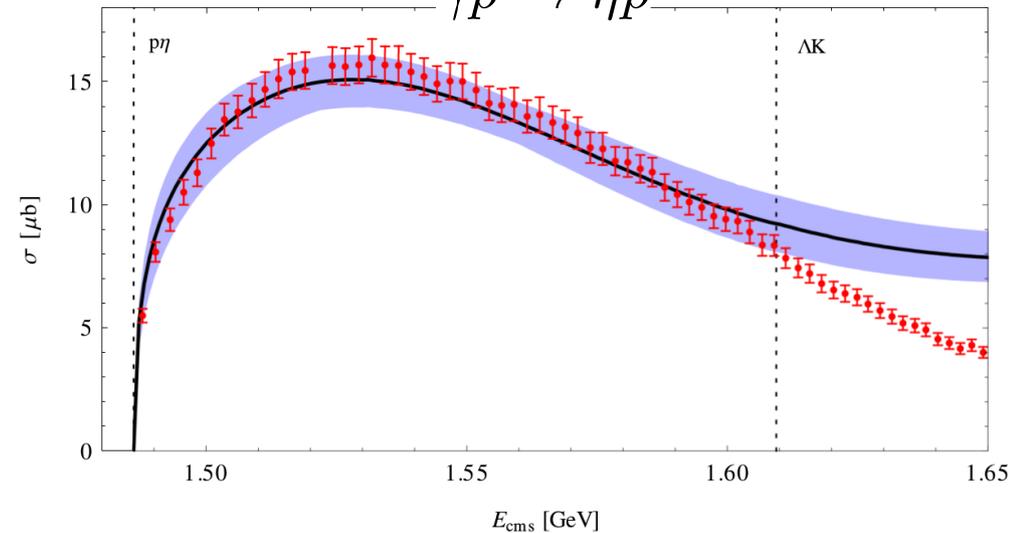
Arndt et al. (2012)

►  $\pi^- p \rightarrow \eta n$  differential cross sections

Prakhov et al. (2005)



$\gamma p \rightarrow \eta p$

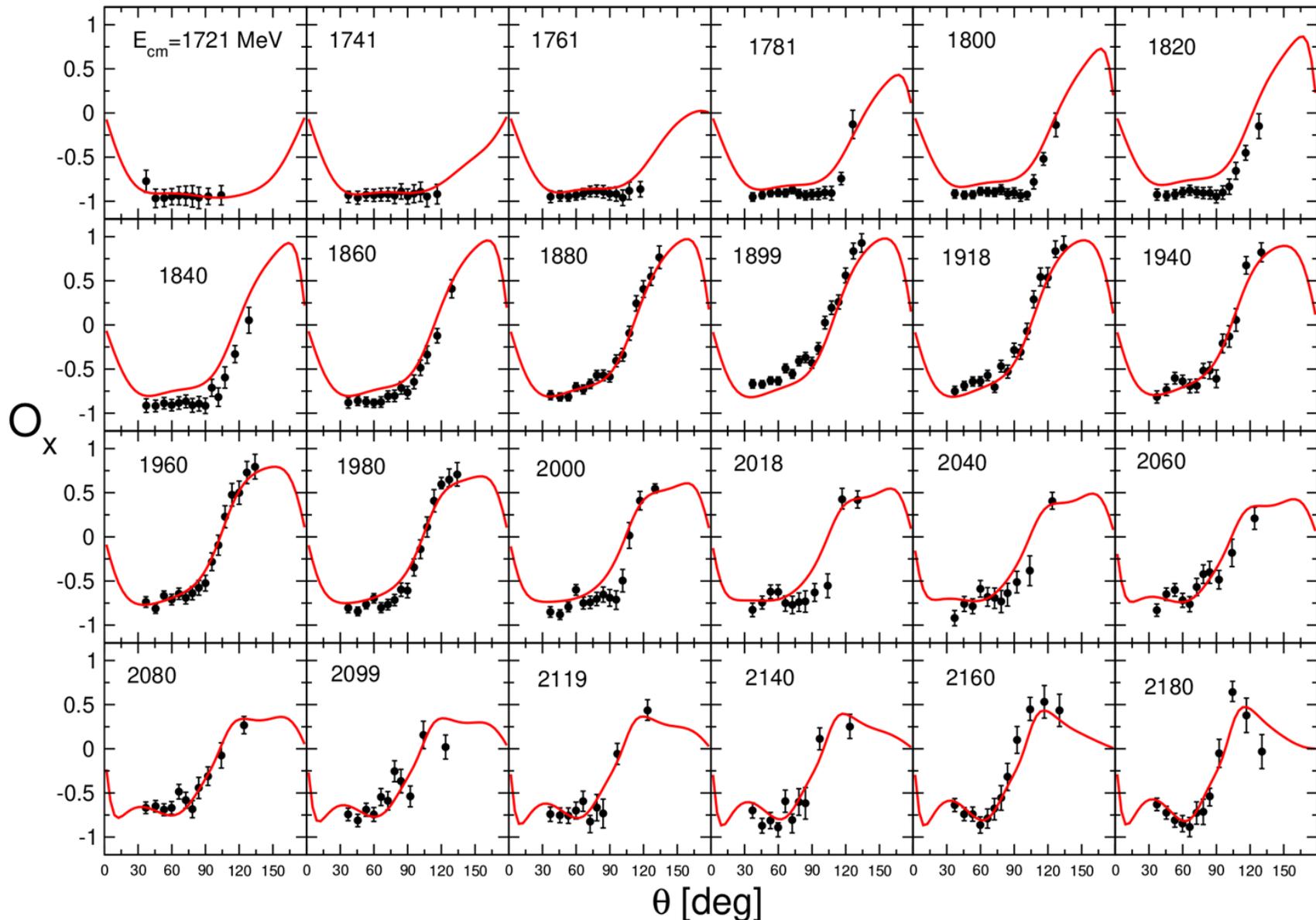


→ Making the “Missing resonance problem” worse ?!

# Selected Fit Results (I)

●  $\gamma p \rightarrow K^+ \Lambda$ :

<http://collaborations.fz-juelich.de/ikp/meson-baryon/main>



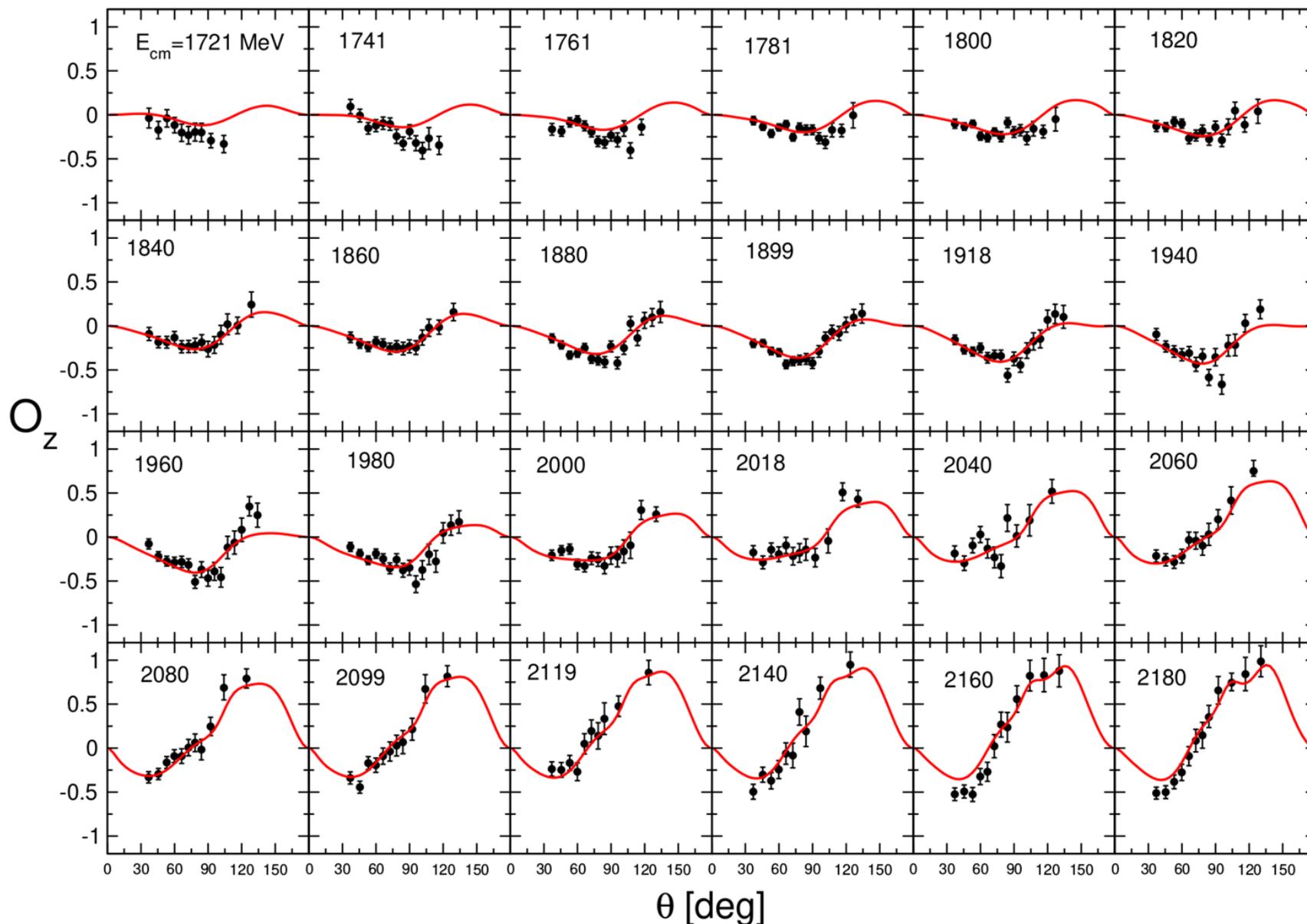
[D. Roenchen, M. D., U.-G. Meißner, EPJ A 54, 110 (2018)]

data: Paterson (CLAS) PRC 93, 065201 (2016), red line: fit JüBo2019

# Selected Fit Results (II)

●  $\gamma p \rightarrow K^+ \Lambda$ :

<http://collaborations.fz-juelich.de/ikp/meson-baryon/main>



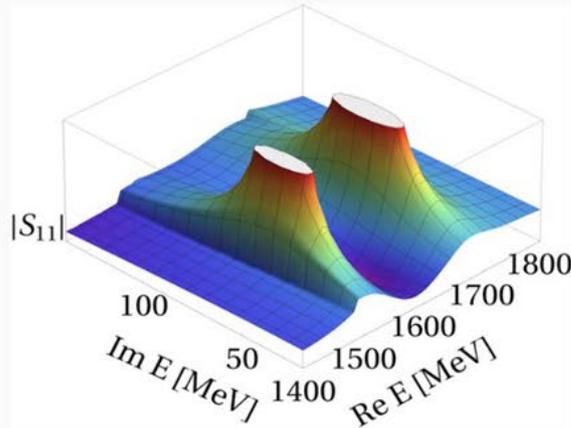
[D. Roenchen, M. D., U.-G. Meißner, EPJ A 54, 110 (2018)]

data: Paterson (CLAS) PRC 93, 065201 (2016), red line: fit JüBo2019

# Resonance Couplings

Resonance states: Poles in the  $T$ -matrix on the 2<sup>nd</sup> Riemann sheet

[D. Roenchen, M. D., U.-G. Meißner, EPJ A 54, 110 (2018)]



- $\text{Re}(E_0)$  = “mass”,  $-2\text{Im}(E_0)$  = “width”
- elastic  $\pi N$  residue ( $|r_{\pi N}|, \theta_{\pi N \rightarrow \pi N}$ ), normalized residues for inelastic channels ( $\sqrt{\Gamma_{\pi N} \Gamma_{\mu}} / \Gamma_{\text{tot}}, \theta_{\pi N \rightarrow \mu}$ )
- photocouplings at the pole:  $\tilde{A}_{\text{pole}}^h = A_{\text{pole}}^h e^{i\vartheta^h}$ ,  $h = 1/2, 3/2$

Inclusion of  $\gamma p \rightarrow K^+ \Lambda$  in JüBo (“JuBo2017-1”): 3 additional states

	$z_0$ [MeV]	$\frac{\Gamma_{\pi N}}{\Gamma_{\text{tot}}}$	$\frac{\Gamma_{\eta N}}{\Gamma_{\text{tot}}}$	$\frac{\Gamma_{K\Lambda}}{\Gamma_{\text{tot}}}$
N(1900)3/2 <sup>+</sup>	1923 – $i$ 108.4	1.5 %	0.78 %	2.99 %
N(2060)5/2 <sup>-</sup>	1924 – $i$ 100.4	0.35 %	0.15 %	13.47 %
$\Delta(2190)$ 1/2 <sup>+</sup>	2191 – $i$ 103.0	33.12 %		

- N(1900)3/2<sup>+</sup>: s-channel resonances, seen in many other analyses of kaon photoproduction (BnGa), 3 stars in PDG
- N(2060)5/2<sup>-</sup>: dynamically generated, 2 stars in PDG, seen e.g. by BnGa
- $\Delta(2190)$  1/2<sup>+</sup>: dyn. gen., no equivalent PDG state