Study of the X(6900): from partial wave decomposition

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Outlines

1 Introduction

2 Formalism

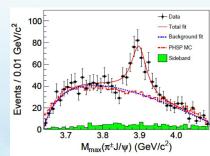
Nature of the X(6900)

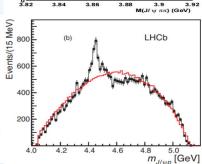
4 Summary

Introduction

- X(3872) by Belle, in 2003
- Z+c(3900) by BESIII and Belle, in 2013
- Pc states by LHCb, in2015
-
- Their nature?
 - Quantum number?
 - The inner structure?



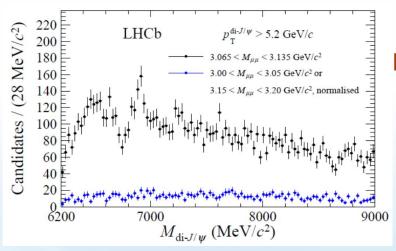


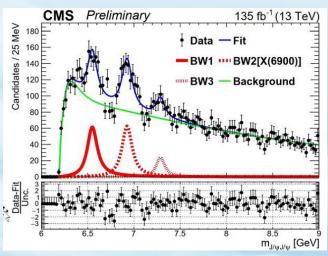


- Only one physical amplitude!
- It should satisfy the fundamental QFT principles
- It should be compatible with the data

X(6900)

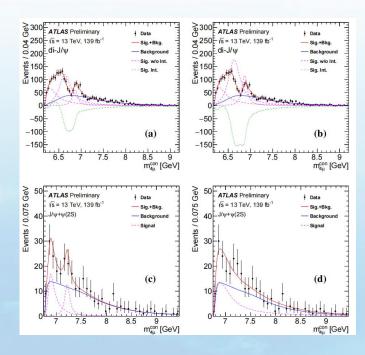
Fully heavy tetra-quark state(s)?





CMS, Kai Yi's talk at ICHEP 2022.

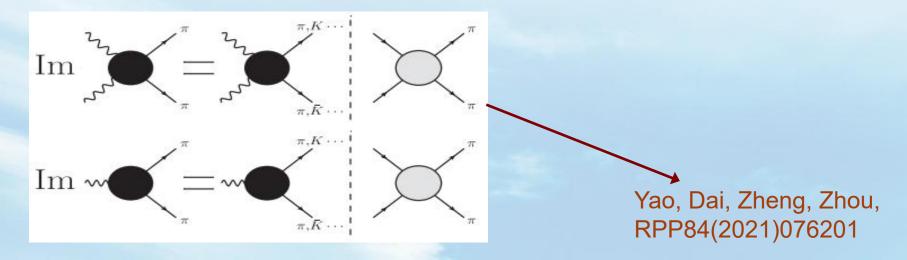
LHCb, Sci.Bull. 65 (2020) 23, 1983



ATLAS, ATLAS-CONF-2022-040

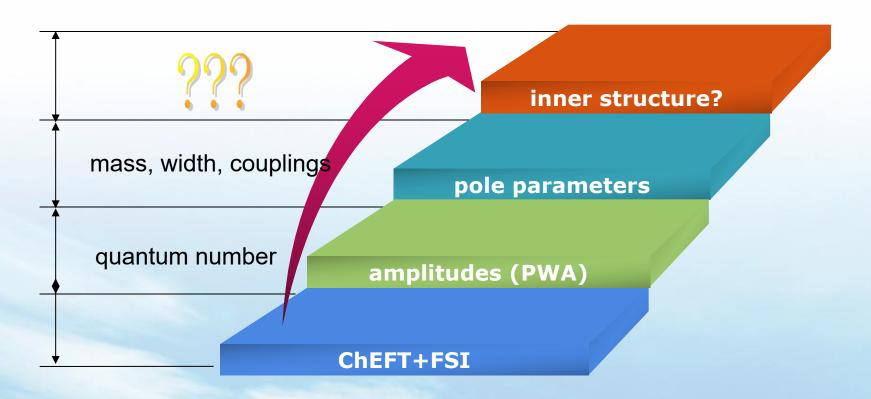
why FSI?

- FSI needs to be taken into account to perform an amplitude analysis
- Methods: KM, N/D, AMP, Roy equation, PKU, Pade, LSE, BSE, ChEFT, et.al.
- Fixed scattering, decaying amplitudes: extracting resonance information



Inner structure?

Interpretation these amplitudes, poles?



Pole counting rule

- Pole counting rule: distinguish molecule and BW resonance.
- At the very beginning, applied to light mesons

Morgan NPA543 (1992) 632.

Dai, Wang and Zheng
CTP57 (2012) 841, CTP58 (2012) 410

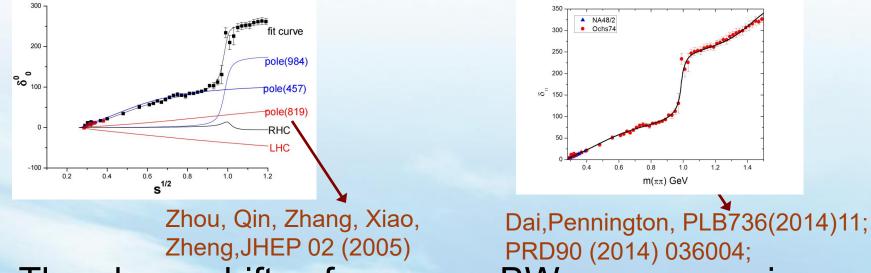
Applied to heavier mesons

Any tools else?

Dai, Sun, Kang, Szczepaniak, Yu, PRD 105 (2022) 5, L051507 Kuang, Dai, Kang, Yao, EPJC 80 (2020) 5, 433 Dai, Shi, Tang, Zheng, PRD 92 (2015) 1, 014020

phase shifts?

- Phase shifts help to study hadronic scatterings and resonances therein
- Successful in study of light scalars: PKU,Roy equation, Omnes representations, etc.



 The phase shifts of a narrow BW resonance rise from 0 to 180 degrees, crossing 90 degrees at E=M

2. Formalism

 coupled channels: J/ψJ/ψ、 ψ(2s)J/ψ、 ψ(3770)J/ψ

J^{PC}	J/ψ $1^{}$	$\psi(2s)$ $1^{}$	$\psi(3770)$ 1	$\frac{\eta_c}{0^{-+}}$	h_c 1^{+-}	χ_{c0} 0^{++}	χ_{c1} 1^{++}	$\frac{\chi_{c2}}{2^{++}}$
质量								
(MeV)	3096.9	3686.1	3773.7	2983.9	3525.4	3414.7	3510.7	3556.2

η_cη_c h_ch_c channels are supressed by HQS

Energy region: [6.2,7.2] GeV

Dong, Baru, Guo, Hanhart, Nefediev, PRL126, 132001(2021) Gong, Du, Zhao, PRD, 106(5):054011, 2022

Effective Lagrangians

Four vector interactions, similar to HGLS formalism

Geng, Molina, Oset. CPC41(12):124101, 2017.

$$\mathcal{L} = c_1 V_{\mu} V_{\alpha} V^{\mu} V^{\alpha} + c_2 V_{\mu} V_{\alpha} V^{\mu} V'^{\alpha} + c_3 V_{\mu} V_{\alpha}' V^{\mu} V'^{\alpha}$$

$$+ c_4 V_{\mu} V'^{\mu} V_{\alpha} V'^{\alpha} + c_5 V_{\mu} V_{\alpha} V^{\mu} V''^{\alpha} + c_6 V_{\mu} V_{\alpha}'' V^{\mu} V''^{\alpha}$$

$$+ c_7 V_{\mu} V''^{\mu} V_{\alpha} V''^{\alpha} + c_8 V_{\mu} V_{\alpha}' V^{\mu} V''^{\alpha} + c_9 V_{\mu} V'^{\mu} V_{\alpha} V''^{\alpha}$$

Zhou, Guo, Kuang, Yang, Dai, arxiv:2207.07537 [hep-ph]

Feynman diagrams up to NLO

Building amplitudes

PWA

$$T^{ij}_{\mu_1\mu_2;\mu_3\mu_4}(s,z_s) = 16\pi N_{ij} \sum_{I} (2J+1) T^{J,ij}_{\mu_1\mu_2;\mu_3\mu_4}(s) d^J_{\mu\mu'}(z_s) .$$

- Time reversal and parity invariance reduce the amplitudes
- Tranfer amplitudes from helicity to the JMLS representation

$$\begin{split} T^{ij}_{\mu_1\mu_2;\mu_3\mu_4}(s,z_s) &= 16\pi N_{ij} \sum_{J=M}^{\infty} (2J+1) d^J_{\mu\mu'}(z_s) \sum_{LS,L'S'} \frac{\sqrt{(2L+1)(2L'+1)}}{2J+1} \left\langle LS0\mu|J\mu\right\rangle \left\langle J\mu'|L'S'0\mu'\right\rangle \\ & \left\langle s_1s_2\mu_1, -\mu_2|S\mu\right\rangle \left\langle S'\mu'|s_3s_4\mu_3, -\mu_4\right\rangle T^{J,ij}_{LS,L'S'} \,. \end{split}$$

$$\begin{split} T^{ij}_{^{3}P_{2}}(s) &= \frac{2}{5}T^{2}_{++++}(s) + \frac{\sqrt{3}}{5}T^{2}_{+++0}(s) + \frac{\sqrt{3}}{5}T^{2}_{+0++}(s) + \frac{\sqrt{3}}{5}T^{2}_{++0+}(s) + \frac{\sqrt{3}}{5}T^{2}_{0+++}(s) - \frac{2}{5}T^{2}_{++--}(s) \\ &- \frac{\sqrt{3}}{5}T^{2}_{++-0}(s) - \frac{\sqrt{3}}{5}T^{2}_{-0++}(s) - \frac{\sqrt{3}}{5}T^{2}_{++0-}(s) - \frac{\sqrt{3}}{5}T^{2}_{0-++}(s) + \frac{3}{10}T^{2}_{+0+0}(s) + \frac{3}{10}T^{2}_{+00+}(s) \\ &+ \frac{3}{10}T^{2}_{0++0}(s) + \frac{3}{10}T^{2}_{0+0+}(s) - \frac{3}{10}T^{2}_{0+0-}(s) - \frac{3}{10}T^{2}_{0-0+0}(s) - \frac{3}{10}T^{2}_{0-+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) + \frac{3}{10}T^{2}_{0+0+}(s) - \frac{3}{10}T^{2}_{0+0-}(s) - \frac{3}{10}T^{2}_{0-0+0}(s) - \frac{3}{10}T^{2}_{0-0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) + \frac{3}{10}T^{2}_{0+0+}(s) - \frac{3}{10}T^{2}_{0+0-}(s) - \frac{3}{10}T^{2}_{0-0+0}(s) - \frac{3}{10}T^{2}_{0-0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) + \frac{3}{10}T^{2}_{0+0+}(s) - \frac{3}{10}T^{2}_{0+0-}(s) - \frac{3}{10}T^{2}_{0-0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) + \frac{3}{10}T^{2}_{0+0+}(s) - \frac{3}{10}T^{2}_{0+0-}(s) - \frac{3}{10}T^{2}_{0+0-}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) + \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0-}(s) - \frac{3}{10}T^{2}_{0+0-}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) + \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) - \frac{3}{10}T^{2}_{0+0}(s) \\ &+ \frac{3}{10}T^{2}_{0+0}(s) - \frac{3$$

PW amplitudes

Five PWs at last

$\overline{}$	S = 0	S = 1	S=2				
0	$0^{++} \ (^1S_0)$	1+-	2^{++} ($^{5}\mathbf{S_{2}}$)				
1	1	$0^{-+} \; (^3P_0) 1^{-+} \; (^3P_1) 2^{-+} \; (^3P_2)$	1 2 3				
2	2++	1^{+-} 2^{+-} 3^{+-}	0^{++} 1^{++} 2^{++} 3^{++} 4^{++}				
:	i i	<u>:</u>	<u>;</u>				

- Ignoring PWs with L>=2
- No S-D coupled channel scatterings

Unitarization

unitarity: ignore the S-D coupled waves

$$\langle L'S'|T^J - T^{J\dagger}|LS \rangle = i \frac{4|\vec{p}''|}{E''_{cm}} \sum_{L''S''} \langle L'S'|T^{J\dagger}|L''S'' \rangle \langle L''S''|T^J|LS \rangle$$

$$\mathbf{Im} T^{ij}_{JLS} = \sum_{k=1}^a T^{ik}_{JLS} \ \rho_k \ T^{kj \ *}_{JLS}$$

 Pade approximation: successful in confirming existence of light scalars, restore unitarity and perturbative amplitudes up to NLO

$$T = T^{LO} \cdot [T^{LO} - T^{NLO}]^{-1} \cdot T^{LO}$$

Dai, Wang and Zheng CTP57 (2012) 841, CTP58 (2012) 410

Invariant mass spectra

• Each channel contributes a ratio α_i

$$\frac{d \text{ Events}^{1,2}}{d\sqrt{s}} = \tilde{N}_{1,2} \ p_{cm}(s) \sum_{\mu_1 \mu_2 \mu_3 \mu_4} \int_{-1}^{1} dz_s |\sum_{i=1}^{a} \alpha_i T_{\mu_1 \mu_2 \mu_3 \mu_4}^{i1,i2}(s,z_s)|^2$$

Into the forms of PWs

$$\sum_{\mu_1\mu_2\mu_3\mu_4} \int_{-1}^{1} |\sum_{i=1}^{a} \alpha_i T^{i1,i2}_{\mu_1\mu_2\mu_3\mu_4}(s,z_s)|^2 dz_s = 512\pi^2 \left[|F^{1,2}_{^{1}S_0}(s)|^2 + 5|F^{1,2}_{^{5}S_2}(s)|^2 + |F^{1,2}_{^{3}P_0}(s)|^2 + 3|F^{1,2}_{^{3}P_1}(s)|^2 + 5|F^{1,2}_{^{3}P_2}(s)|^2 \right]$$

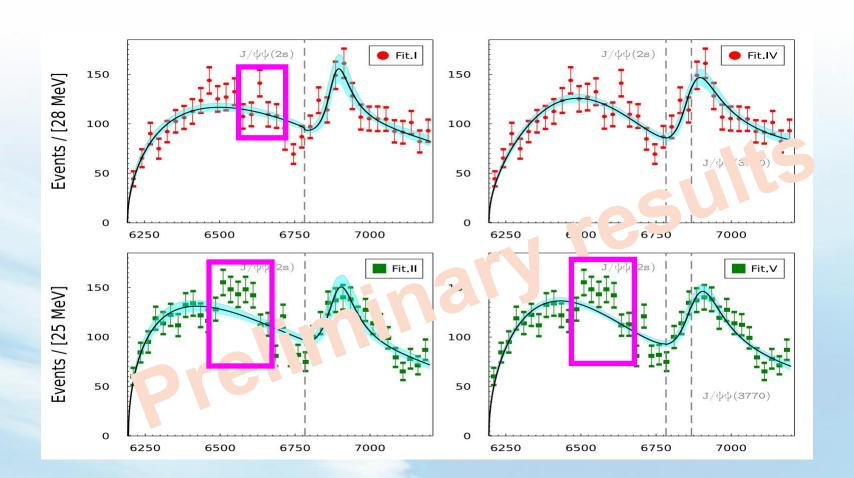
- Compatible with AMP method
 - absorbing l.h.c. and distant r.h.c.
 - Fulfill FSI theorem

$$F_{JLS}^{1,2}(s) = \sum_{i=1}^{a} \alpha_i N_i T_{JLS}^{i1,i2}(s)$$

Dai, Pennington, PLB736(2014)11; PRD90 (2014) 036004; **Dai**, Shi, Tang, Zheng, PRD 92 (2015) 1, 014020

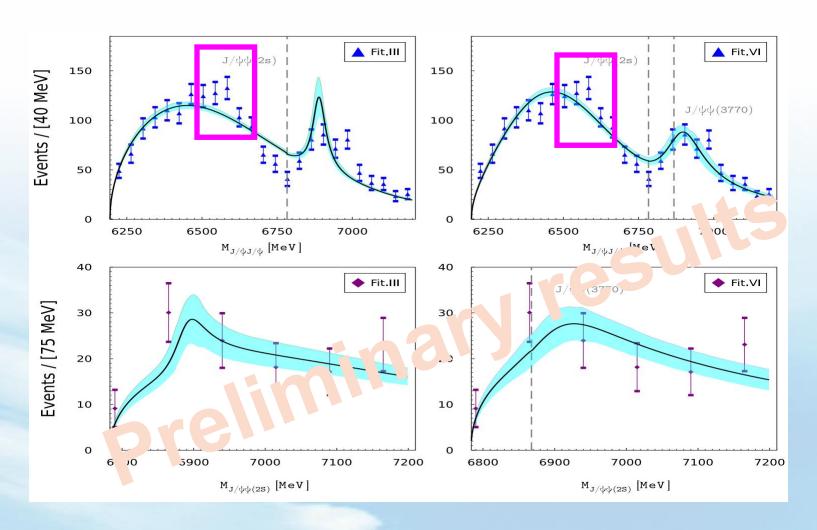
3. Nature of the X(6900)

High quality Fits for all the data-sets



Fits

Coupled channe fits on the data-sets of ATLAS



Couplings

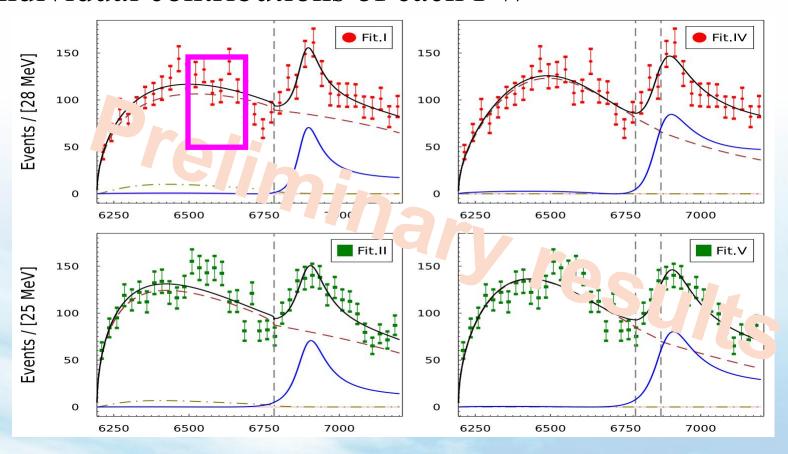
- Couplings are reasonable
- So far, so good....

parameter	$\mathbf{Fit.I}(\mathbf{LHCb})$	Fit.II(CMS)	Fit.III(ATLAS)
c_1	$-0.1232^{+0.0001}_{-0.0001}$	$-0.1504^{+0.0001}_{-0.0002}$	$-0.0617^{+0.0001}_{-0.0001}$
()	$-0.5359^{+0.0021}_{-0.0001}$	$-0.6203^{+0.0001}_{-0.0001}$	$-0.3370^{+0.0001}_{-0.0001}$
c_3	$-0.3250^{+0.0171}_{-0.0001}$	$-0.3492^{+0.0004}_{-0.0001}$	$-0.3230^{+0.0005}_{-0.0003}$
c_4	+0.0234	$-0.6835^{+0.0004}_{-0.0001}$	$-0.5426^{+0.0009}_{-0.0001}$
$c_4 \ ilde{ ilde{N_1}}$	$1.2585 \frac{1}{1.6} \cdot \frac{0.002}{1.85}$	$0.5336^{+0.1404}_{-0.0193}$	$0.0490^{+0.0020}_{-0.0092}$
$ ilde{N_2}$		qa,	$0.0746^{+0.0021}_{-0.0098}$
α_1	$0.3691^{+0.0104}_{-0.0052}$	$0.3510^{-3.001}$	$0^{124}_{-0.0001}^{+0.0001}$
$lpha_2$	$-0.9294^{+0.0089}_{-0.0022}$	$-0.9364^{+0.0001}_{-0.0009}$	U.S. 1. 10.70 1
$\chi^2_{d.o.f.}$	1.29	1.77	2.73

Lack of angular distributions!

Fits

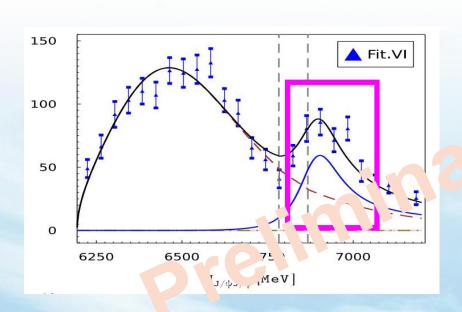
Individual contributions of each PW

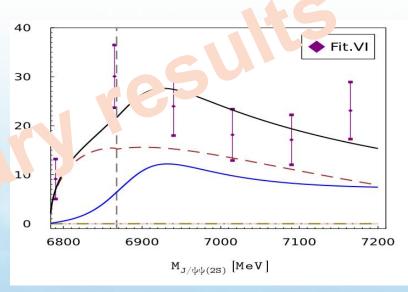


■ ¹S₀ dominates around 6900 MeV!

Fits

 ¹S₀ dominates around 6900 MeV for both J/ψJ/ψ \ ψ(2S)J/ψ channels





⁵S₂ Contributes a large background

Pole analysis

- One resonace found in ${}^{1}S_{0}$ wave: X(6900)---- 0^{++}
- Couple channels case: a pair of accompanying poles

Triple channels case: Four poles in unphysical sheets, implying

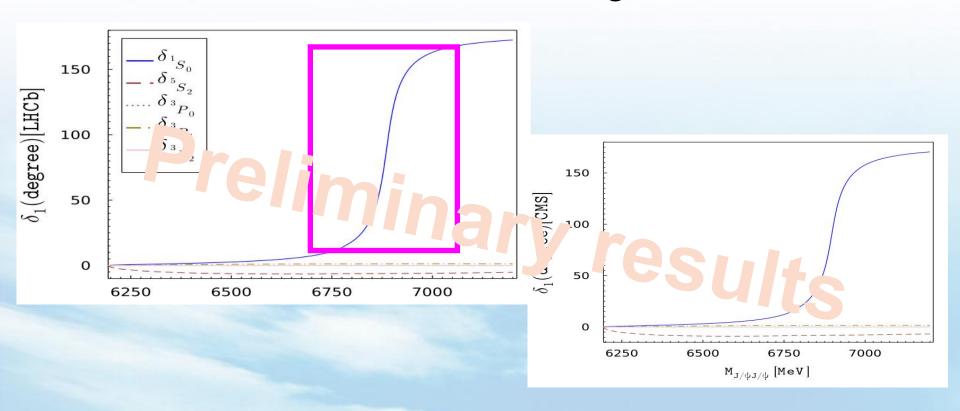
the BW origin.

pole counting rule:
Morgan,
NPA543 (1992) 632.
Dai, Wang and Zheng
CTP57 (2012) 841

Dete	RS	pole location (MeV)	$g_{J/\psi J/\psi} = g e^{i\varphi}$		$g_{J/\psi\psi(2S)} = g e^{i\varphi}$		$g_{J/\psi\psi(3770)} = g e^{i\varphi}$	
Data			$ g_1 (\text{MeV})$	$\varphi_1(^{\circ})$	$ g_2 (\text{MeV})$	$\varphi_2(^\circ)$	$ g_3 (\text{MeV})$	φ ₃ (°)
	II(- + +)	$6872.7^{+6.0}_{-8.6}$ - $i46.5^{+2.1}_{-1.0}$	$1352.7^{+28.1}_{-11.7}$	$86.2^{+0.1}_{-0.1}$	$946.5^{+18.8}_{-8.0}$	$84.9^{+0.1}_{-0.1}$	$14.4^{+0.9}_{-0.3}$	$3.3^{+1.7}_{-0.6}$
LHCb(Fit.IV)	III(+)	$6861.0_{-8.8}^{+6.3}$ $-i64.5_{-1.7}^{+2.8}$	$1326.2^{+24.7}_{-11.3}$	$81.0^{+0.2}_{-0.1}$	$917.9_{-7.9}^{+15.1}$	$78.0^{+0.3}_{-0.}$	0 .8	$0^{+1.2}_{-0.7}$
EHOS(Fit.FV)	IV()	$6861.0^{+6.3}_{-8.8}$ - $i64.5^{+2.8}_{-1.7}$	$1322.9^{+26.2}_{-14.6}$	$81.0^{+0.3}_{-0.1}$	915.6+16.0	.8.6 0.4	16.1+0.8	$-70.7^{+0.9}_{-0.8}$
	VII(- + -)	$6872.7^{+6.0}_{-8.6}$ $-i46.5^{+2.1}_{-1.0}$	$1349.7^{+29.1}_{-10.8}$	86.2+0.1	9 4 19.9	$5.0^{+0.1}_{-0.1}$	$14.5^{+0.9}_{-0.3}$	$-72.6^{+1.5}_{-0.4}$
	II(- + +)	$6888.6^{+11.3}_{-7.2}$ - $i58.5^{+1.7}_{-0.5}$	$1450.0^{+23.1}_{00000000000000000000000000000000000$	85.6 ⁺⁰ ₋₀ .	$796.3_{-4.3}^{+12.2}$	$83.4^{+0.1}_{-0.1}$	$40.9^{+2.1}_{-0.1}$	$81.9^{+0.3}_{-0.1}$
CMS(Fit.V)	III(+)	$6879.2^{+11.3}_{-7.4}$ $-i72.3^{+2}_{-1}$	$42 5^{+29}_{-5.7}$	$32.1_{-0.1}^{+0.1}$	$775.1_{-4.2}^{+15.5}$	$77.9^{+0.2}_{-0.1}$	$38.5^{+2.2}_{-0.1}$	$65.1^{+1.6}_{-0.4}$
	IV()	$6879.2_{-4}^{+11.3}$ $72{-1}^{+2}$	$424.0^{+18.8}_{-5.0}$	$82.1^{+0.1}_{-0.1}$	$772.4_{-3.1}^{+8.7}$	$77.9^{+0.2}_{-0.1}$	$38.7^{+2.1}_{-0.1}$	$65.1^{+1.6}_{-0.4}$
	VII.	$ \begin{array}{ccccccccccccccccccccccccccccccccc$	$1450.1_{-5.6}^{+24.4}$	$85.6^{+0.1}_{-0.1}$	$796.1^{+13.6}_{-3.4}$	$83.3^{+0.1}_{-0.1}$	$41.6^{+2.2}_{-0.1}$	$83.1^{+0.4}_{-0.2}$
40	1. ++)	$6896.9_{-4.3}^{+19.1}$ - $i50.8_{-0.2}^{+0.9}$	$1410.1_{-1.9}^{+12.0}$	86.3 ^{+0.1} _{-0.1}	$997.4^{+8.8}_{-1.8}$	$85.1^{+0.1}_{-0.1}$	$9.2^{+0.1}_{-0.1}$	$70.0^{+0.8}_{-0.3}$
A LAS Fit.VI)	íII(+)	$6883.1_{-4.0}^{+18.3} -i73.2_{-0.7}^{+2.8}$	$1379.8^{+7.3}_{-2.7}$	$80.9^{+0.1}_{-0.1}$	$964.5^{+5.6}_{-1.3}$	$77.6^{+0.1}_{-0.2}$	$9.6^{+0.1}_{-0.1}$	$41.3^{+1.1}_{-1.0}$
		$3883.1^{+18.3}_{-4.0}$ $-i73.2^{+2.8}_{-0.7}$	$377.3^{+10.0}_{-2.0}$	$80.8^{+0.1}_{-0.1}$	$962.8^{+7.1}_{-1.2}$	$77.6^{+0.1}_{-0.1}$	$9.7^{+0.1}_{-0.1}$	$41.7^{+1.1}_{-1.0}$
	VII(- + -)	$6896.9_{-4.3}^{+19.1} -i50.8_{-0.2}^{+0.9}$	$1412.1^{+10.4}_{-2.0}$	$86.2^{+0.1}_{-0.1}$	$998.9^{+7.4}_{-2.3}$	$85.1^{+0.1}_{-0.1}$	$9.3^{+0.2}_{-0.1}$	$70.9^{+0.9}_{-0.2}$

Phase shifts

- Phase shifts of ¹S₀ is very likely to the one generated by a narrow BW resonacne: tetra-quark
- Other waves should contribute backgrounds



4. Summary

PWA

One needs PWA to extract pole parameters for the resonance: quantum number, mass, width, residues...

FSI

Amplitude analysis connects QFT principles and Exp. FSI needs to be considered when performing amplitude analysis.

X6900

It is very likely to be a compact tetraquark, and has quantum number 0++

Next?

We lack enough data on the angular distribution, which would be essential to separate the partial waves. Our framework can be applied to other scatterings.

Thank You For your patience!