

2022年第七届手征有效场论研讨会，2022年10月15-17日

# **Hyperon dynamics and hypernuclear formation in heavy-ion collisions**

**重离子碰撞中超子产生和超核的形成**

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# 报告内容

- 奇异性核物理研究现状
- 重离子碰撞中奇异粒子产生
- 原子核碎裂反应和超核形成
- 总结

# 一、奇异性核物理现状介绍

◆ 强子的内部结构

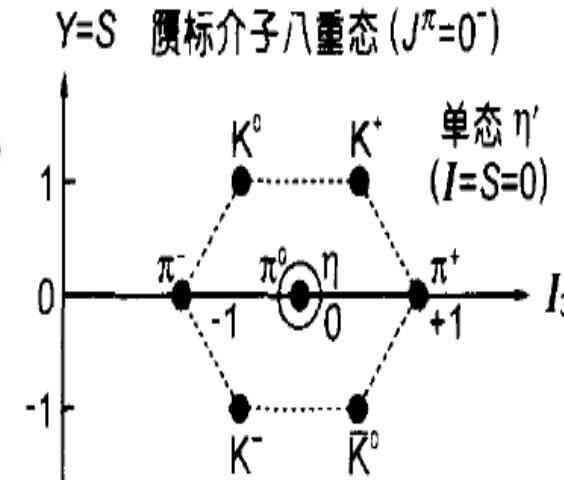


Normal baryon

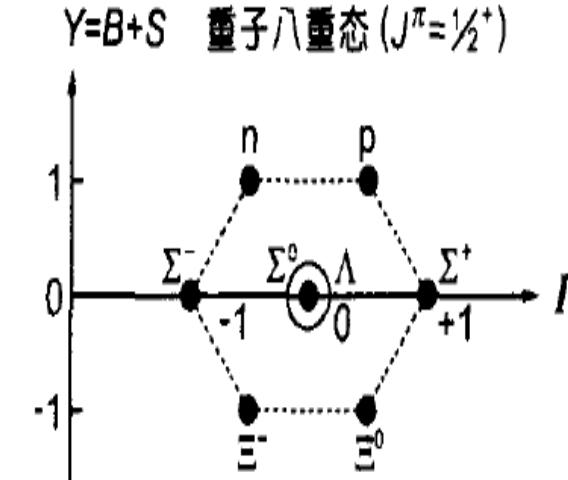


Normal meson

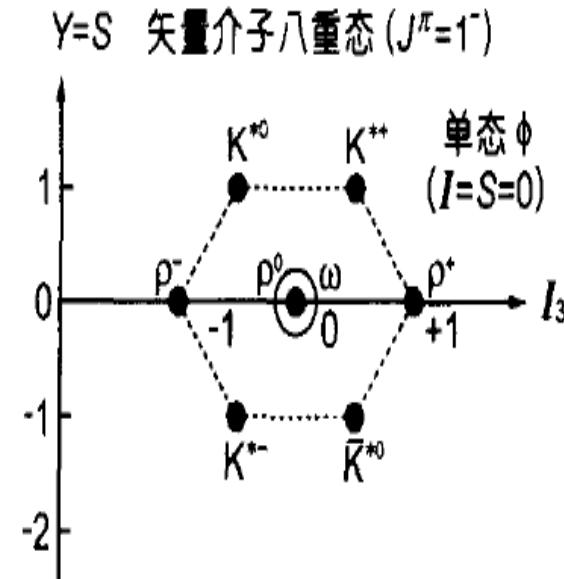
$\Upsilon=S$  质标介子八重态 ( $J^\pi=0^-$ )



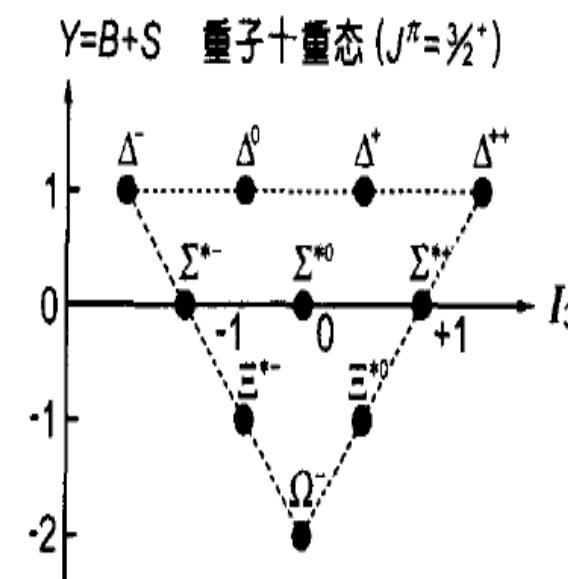
$\Upsilon=B+S$  重子八重态 ( $J^\pi=\frac{1}{2}^+$ )



$\Upsilon=S$  矢量介子八重态 ( $J^\pi=1^-$ )



$\Upsilon=B+S$  重子十重态 ( $J^\pi=\frac{3}{2}^+$ )

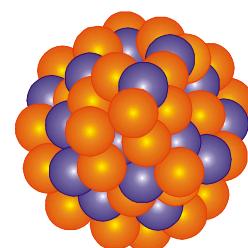


◆ 核子和核子共振态：八(十)重态重子

◆ 介子八重态：矢量介子和赝标量介子

◆ 奇异粒子(含s夸克)主要指介子K( $K^0, K^+$ )和  
 $\bar{K}$ ( $\bar{K}^0, K^-$ )，超子 $\Lambda$ 、 $\Sigma$ 、 $\Xi$ 和 $\Omega$

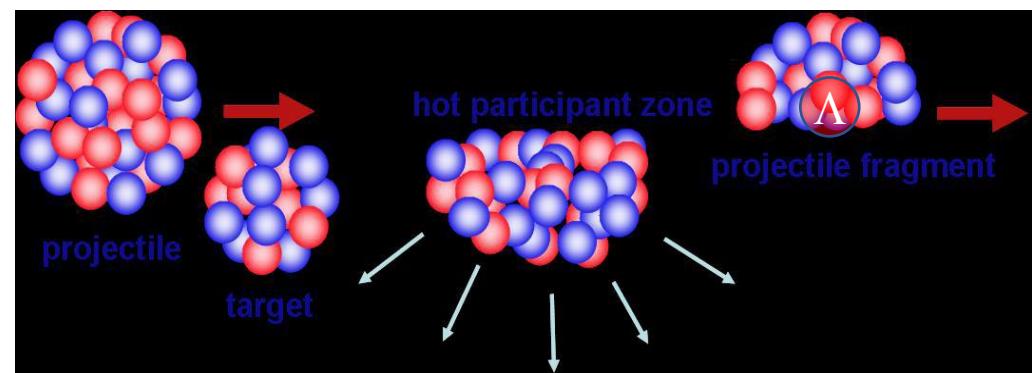
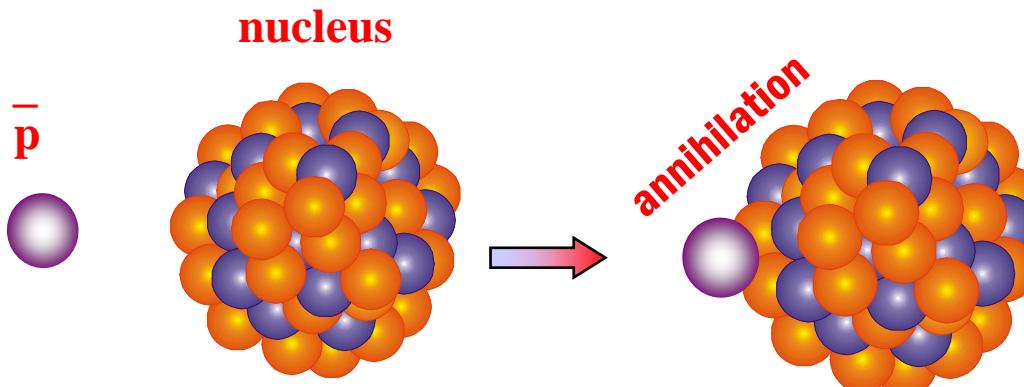
◆ 含有奇异粒子的原子核-超核



# 奇异粒子产生：

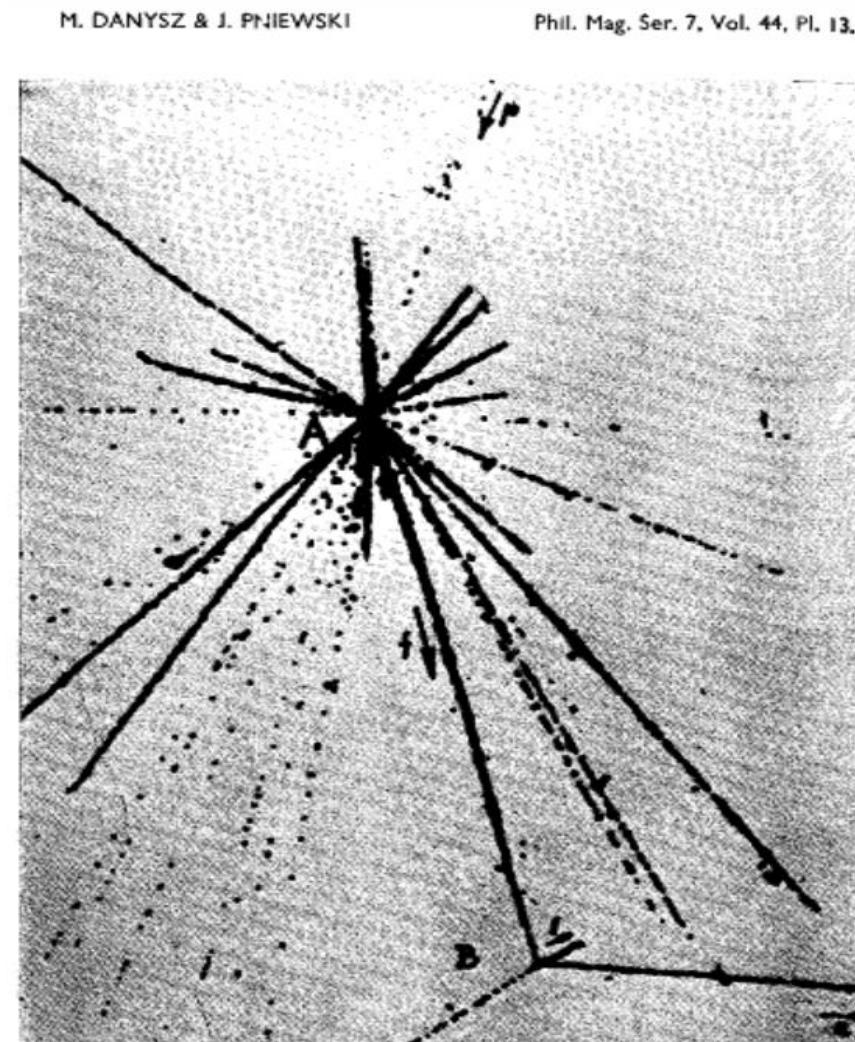
- 中高能重离子碰撞
- 强子(质子, 反质子, 介子)引起的核反应
- 高能电子轰击原子核
- 光核反应

...

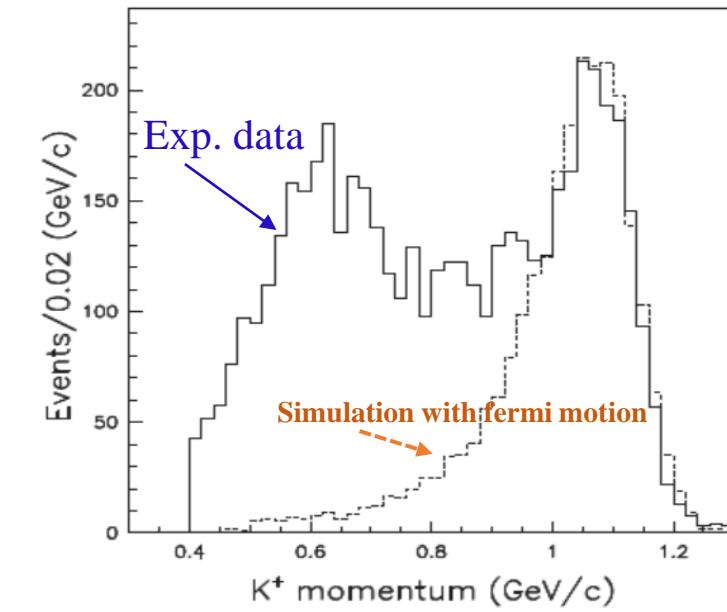
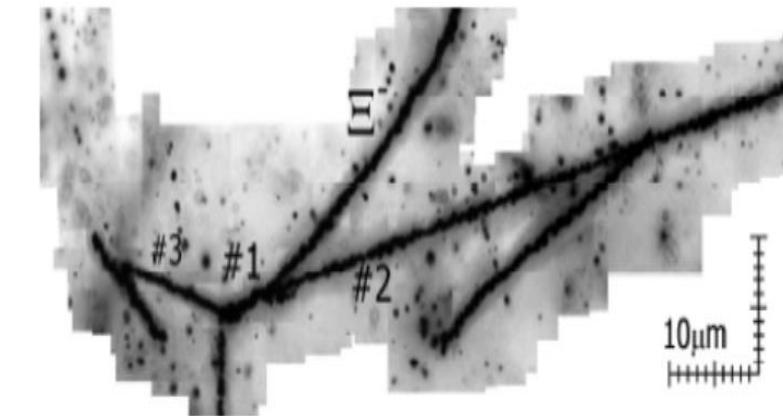


# 1. 超核实验观测和进展

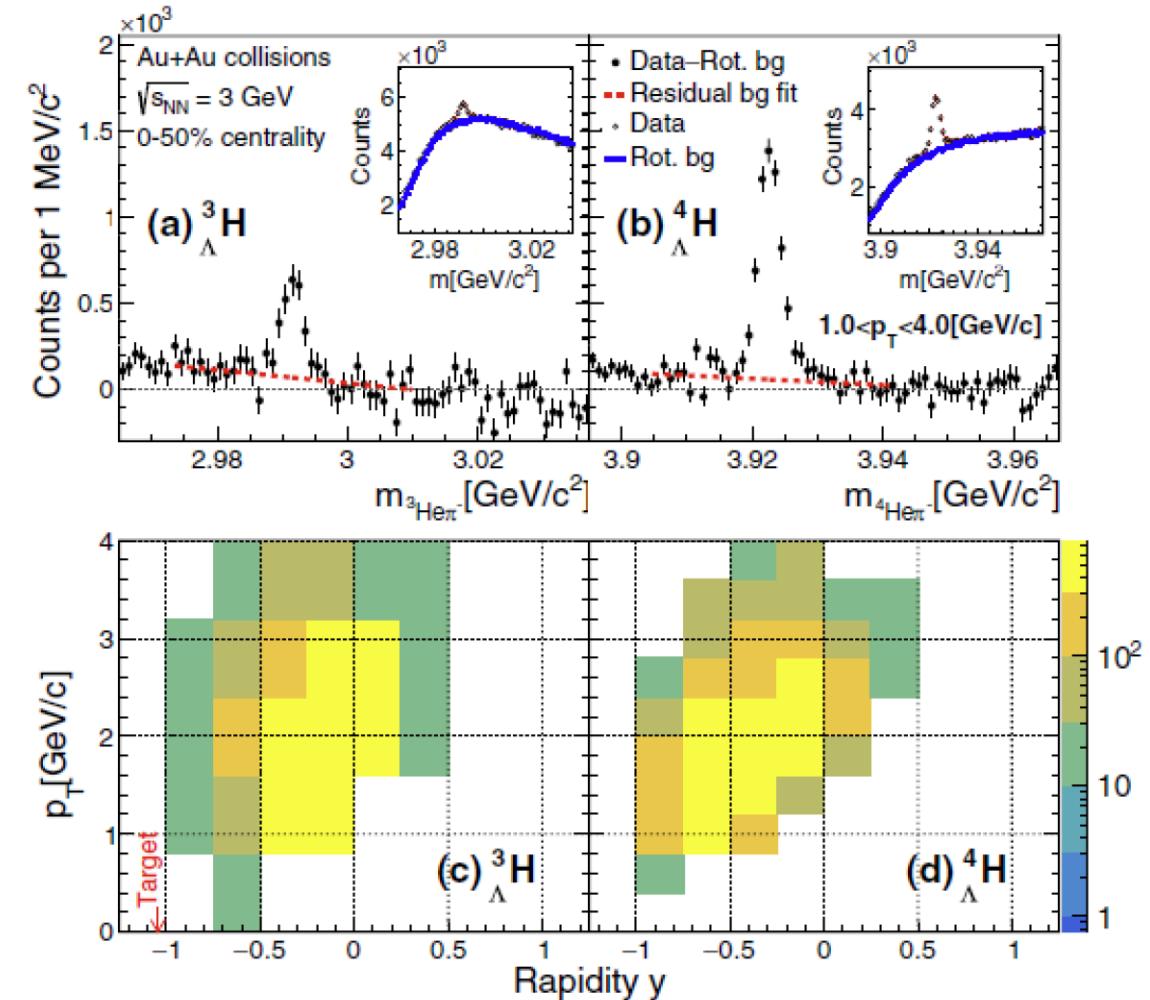
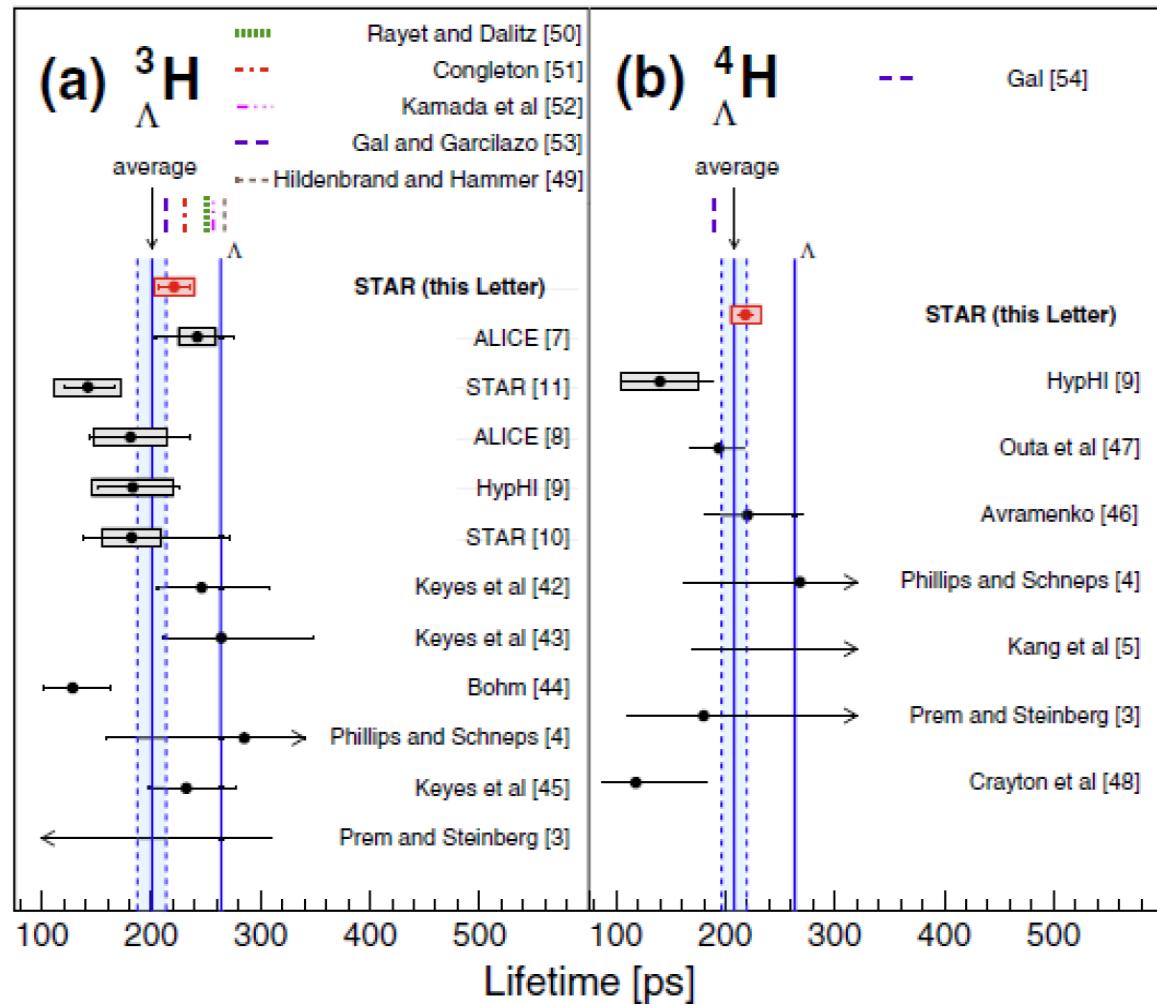
1953年波兰物理学家M. Danysz和J. Pniewski  
在宇宙线乳胶实验中首次发现 $\Lambda$ 超核



利用  $(K^-, K^+)$  产生 $\Xi$ 超核实验观测  
Kazuma Nakazawa et al, J. Phys.: Conf. Ser. 569 (2014) 012082

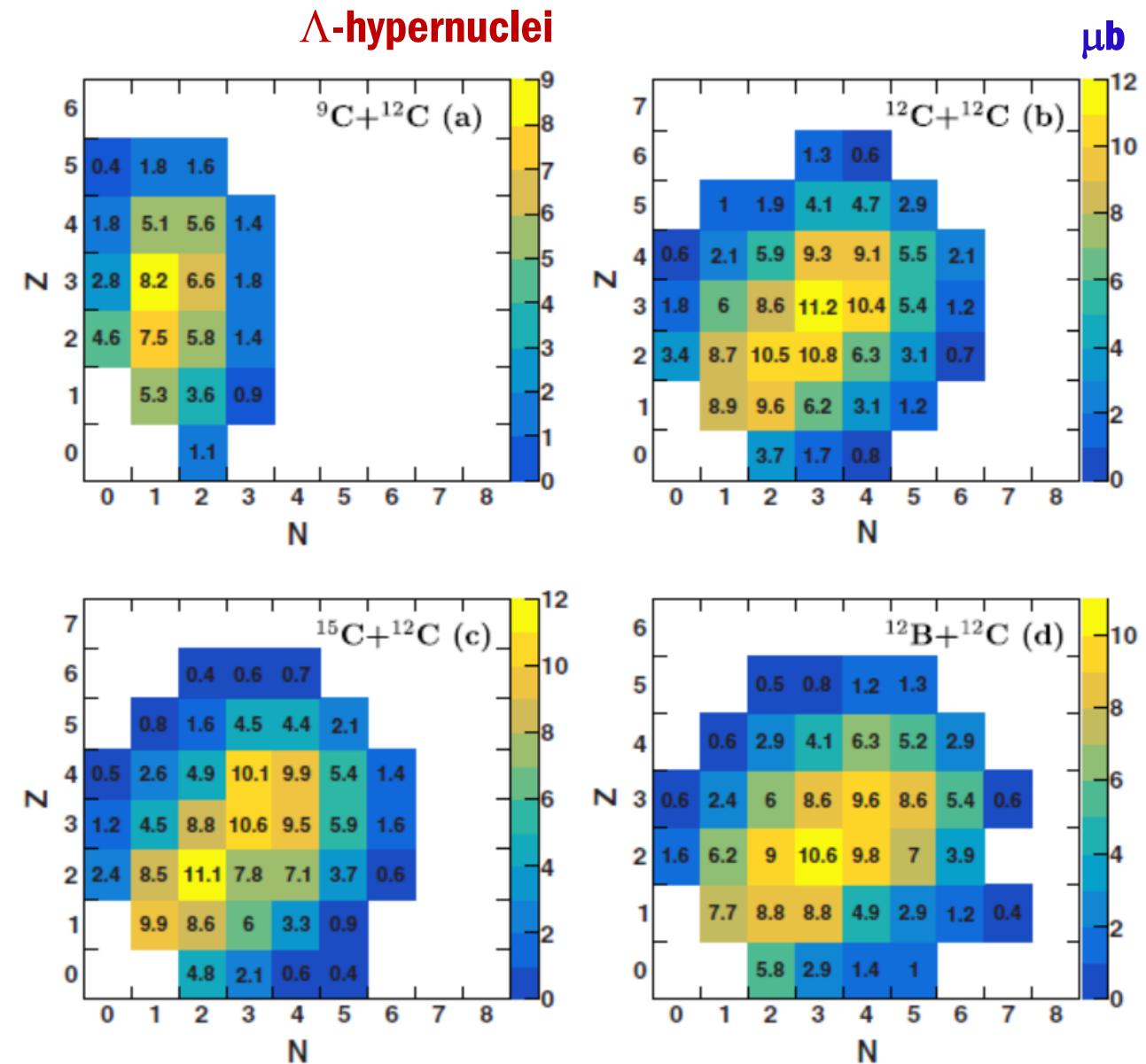
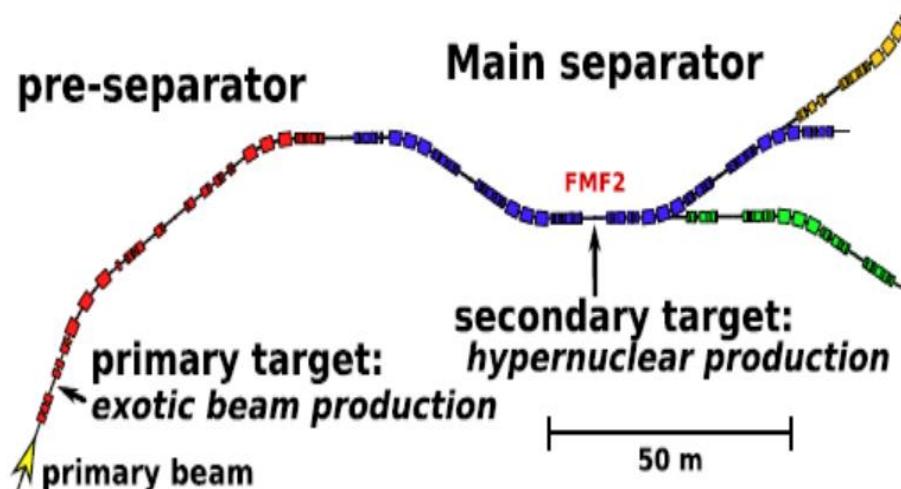


# Hypernuclide ${}^3_{\Lambda}\text{H}$ and ${}^4_{\Lambda}\text{H}$ measured by STAR Collaboration (Phys. Rev. Lett. 128, 202301 (2022))



# 德国GSI产生丰质子/丰质子超核 实验条件分析

C. Rappold, J. López-Fidalgo, PRC 94,  
044616 (2016)

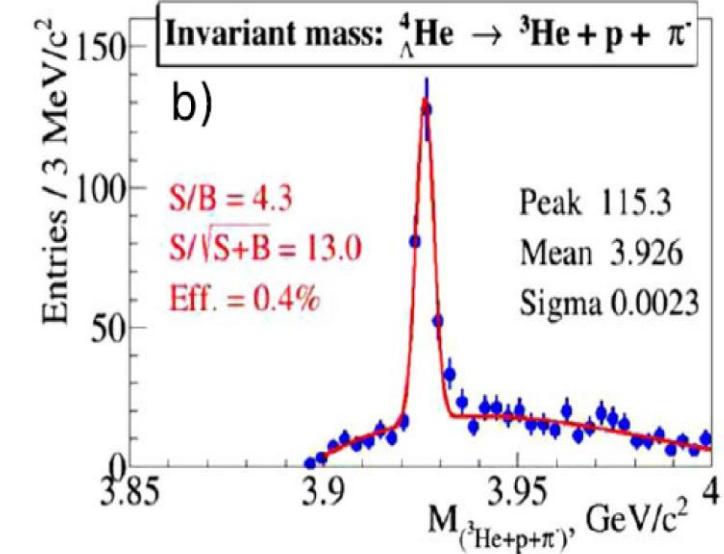
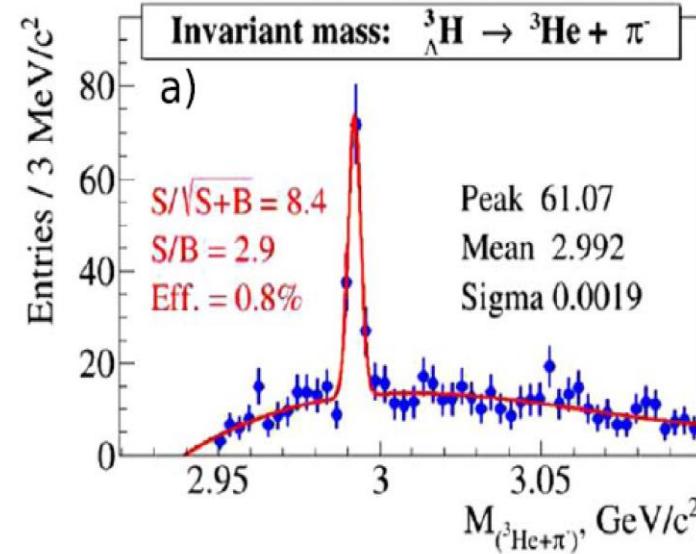
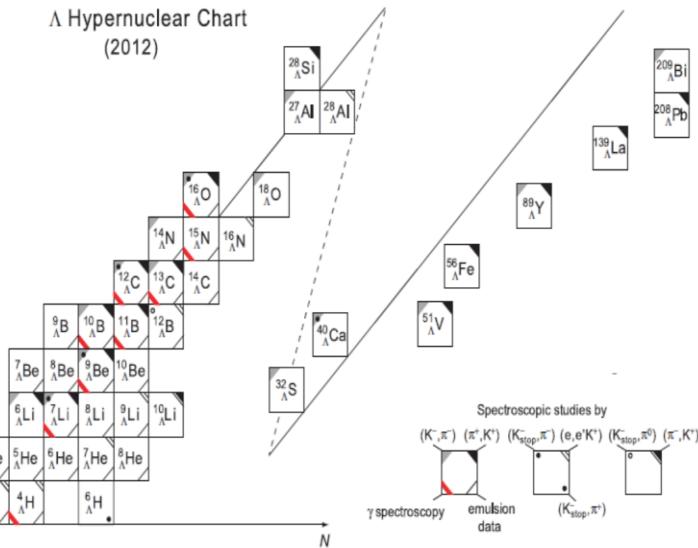
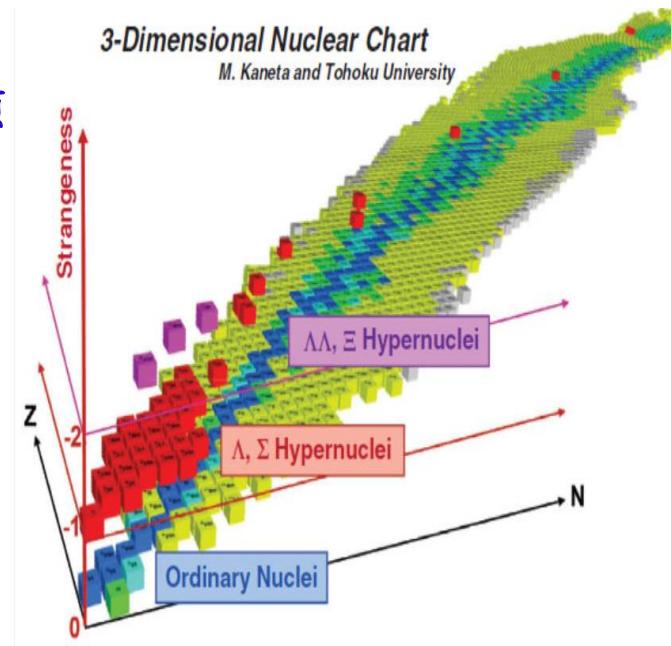
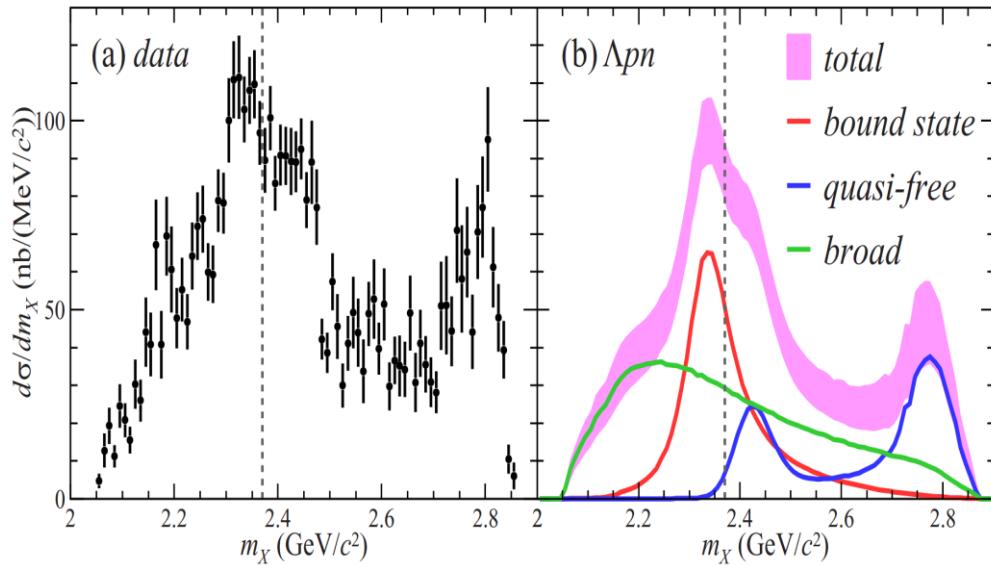


# 重离子碰撞产生超核的优势

1. 极端丰中子或丰质子超核产生和谱学性质
2. 奇特超核产生 ( $s=-2$ )  $_{\Lambda\Lambda}X$  和  $_{\Xi}X$
3. 核物质中  $\Lambda-\Lambda$  和  $\Xi-N$  相互作用

PHYSICAL REVIEW C 102, 044002 (2020)

Observation of a  $\bar{K}NN$  bound state in the  ${}^3\text{He}(K^-, \Lambda p)n$  reaction

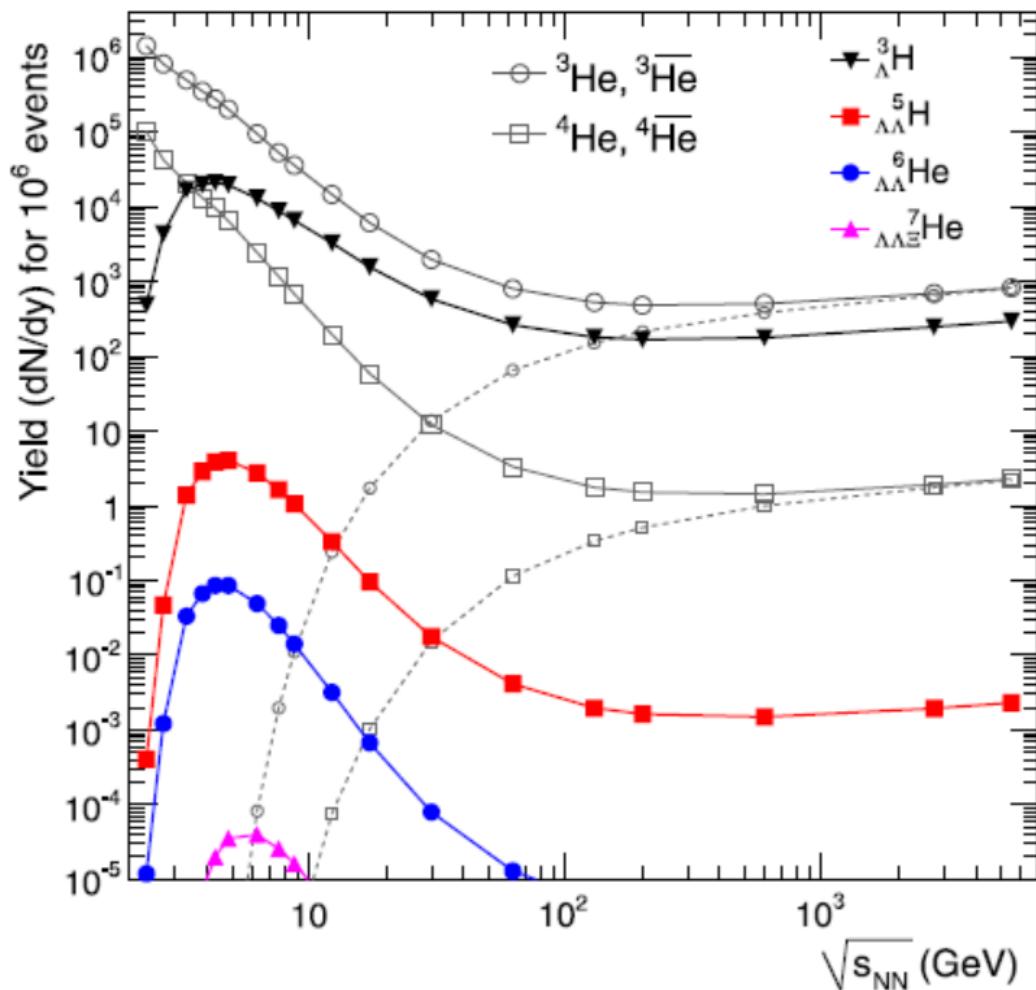


## 2. 中高能重离子碰撞产生超核理论研究—统计理论

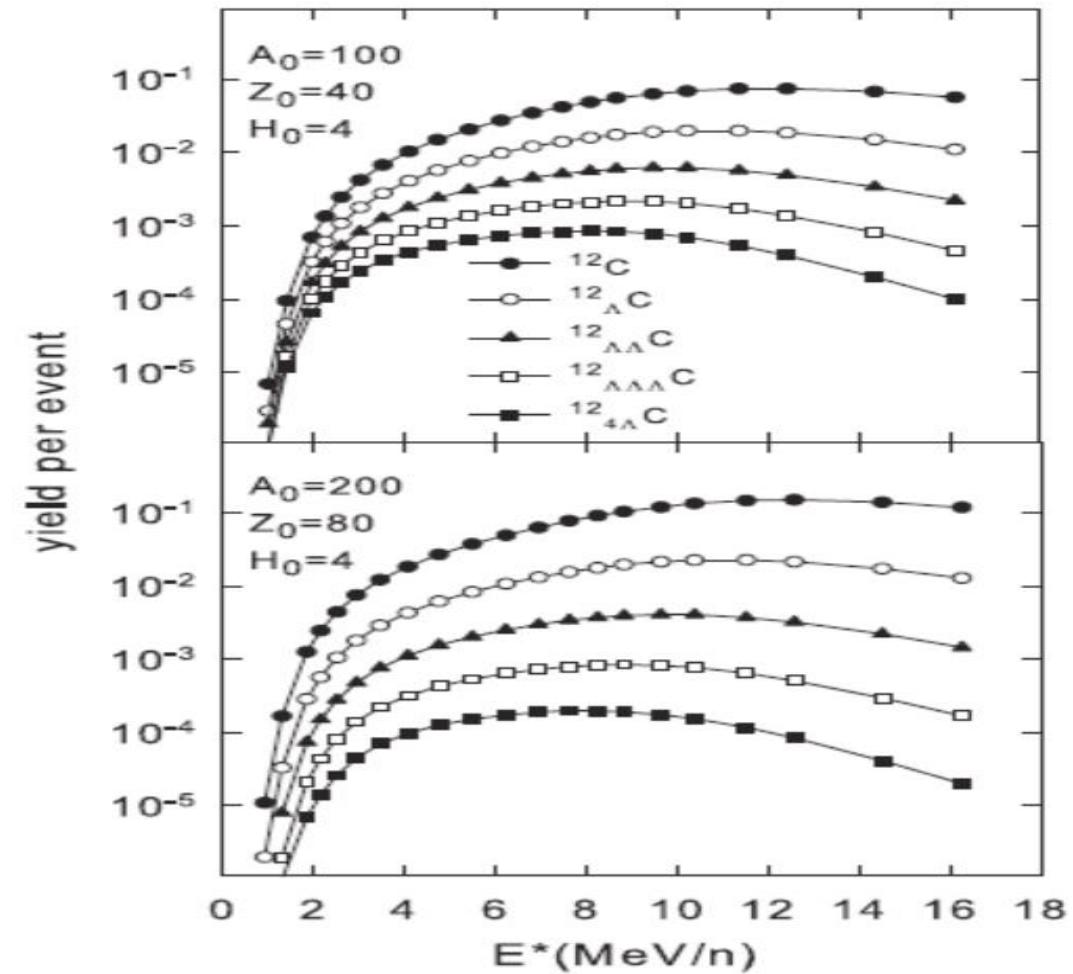
A. Andronic, P. Braun-Munzinger, J. Stachel, H. Stöcker,  
Physics Letters B 697 (2011) 203–207

N. Buyukcizmeci, R. Ogul, A. S. Botvina, M.  
Bleicher, Phys. Scr. 95 075311 (2020)

Pb+Pb

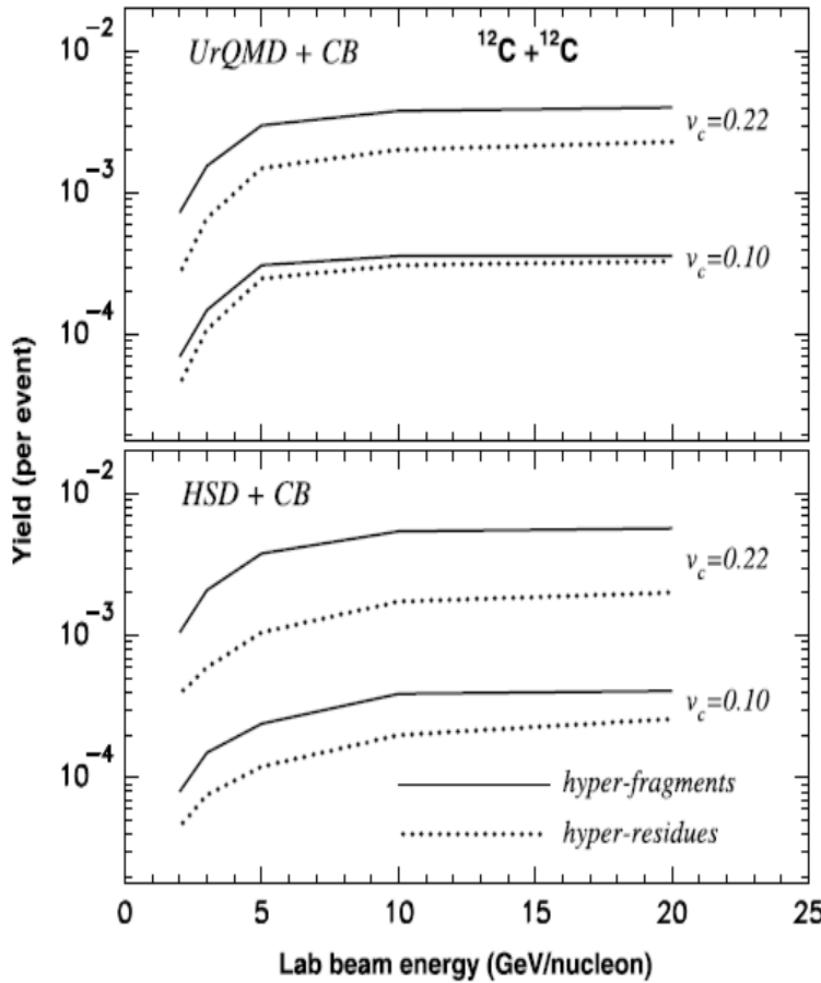


Statistical multifragmentation model (SMM)



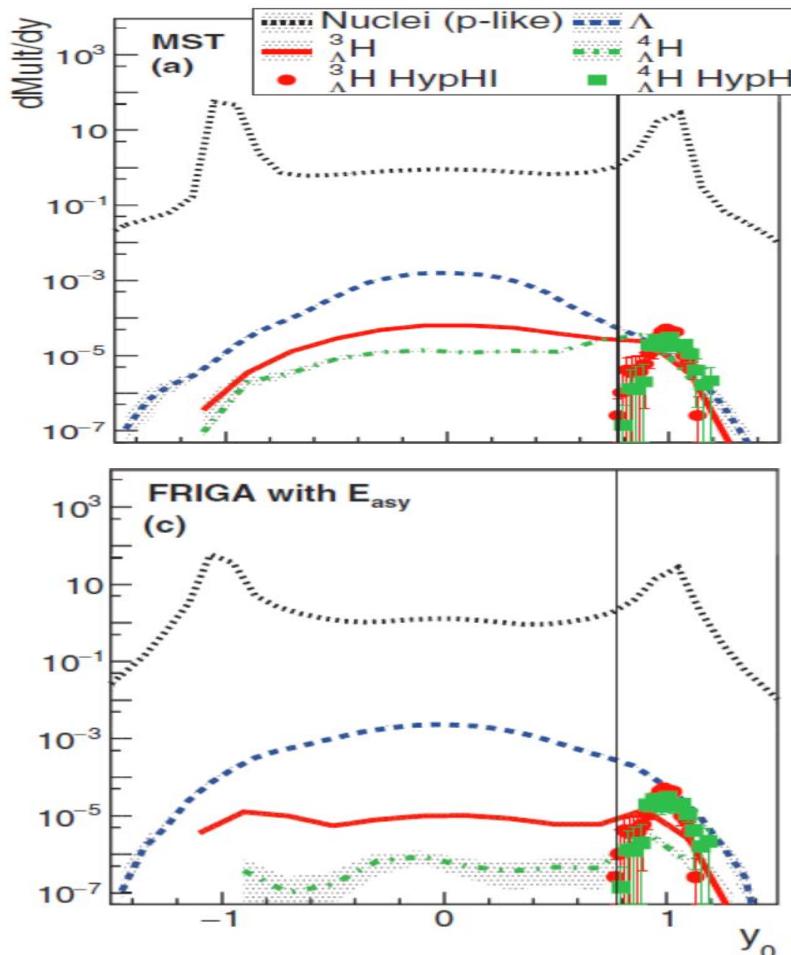
### 3. 中高能重离子碰撞产生超核理论研究方法 — 输运理论+并合模型

A.S. Botvina, J. Steinheimer, E.Bratkovskaya et al., Physics Letters B 742 (2015) 7–14



J. Aichelin, E. Bratkovskaya, A. Le Fèvre et al.,  
Physical Review C 101, 044905 (2020)  
A. Le Fèvre, J. Aichelin, C. Hartnack and Y. Leifels 100, Physical Review C 034904 (2019)

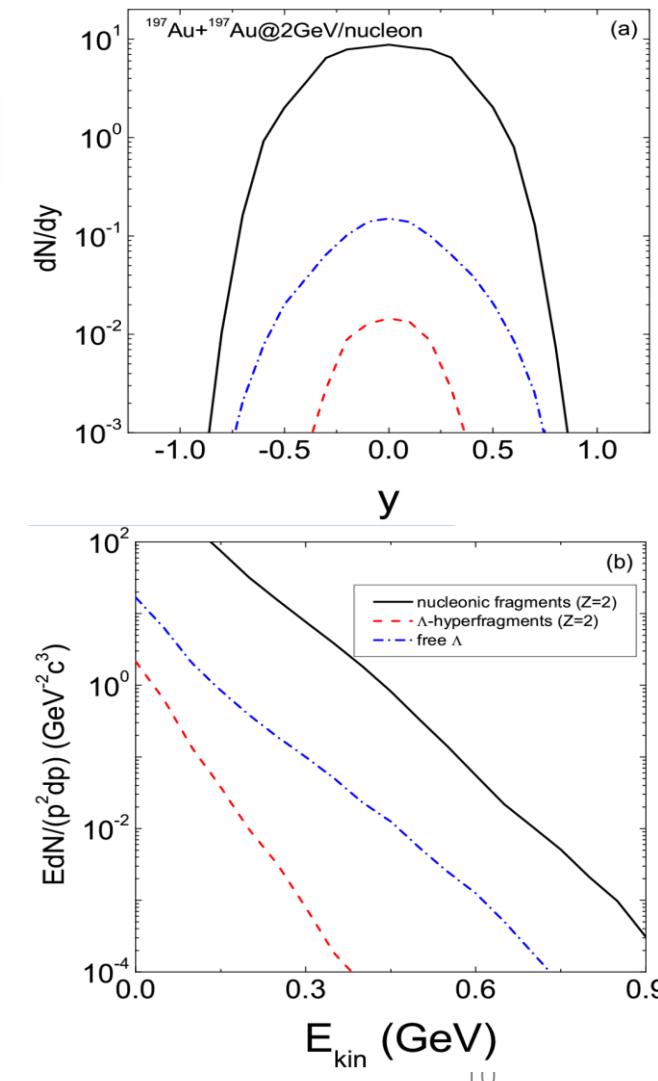
$^{6}\text{Li} + ^{12}\text{C} @ 2\text{A GeV}$



中高能重离子碰撞中奇异粒子产生和超核形成机制

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E-mail: fengzq@impcas.ac.cn



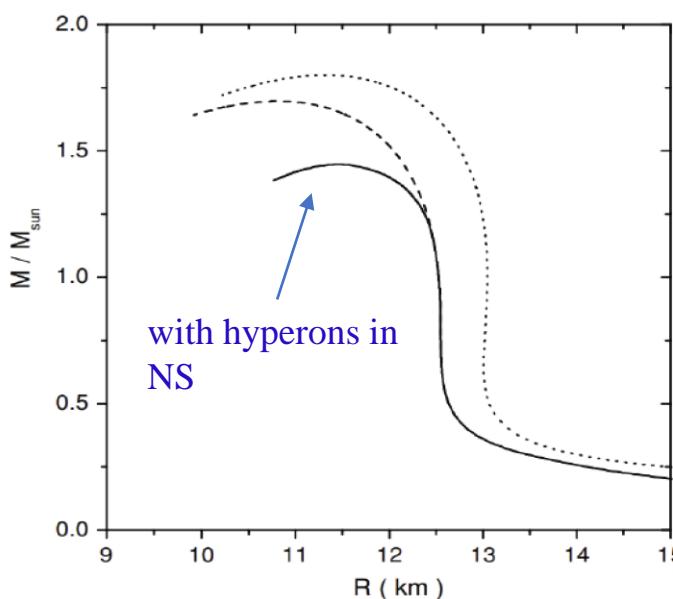
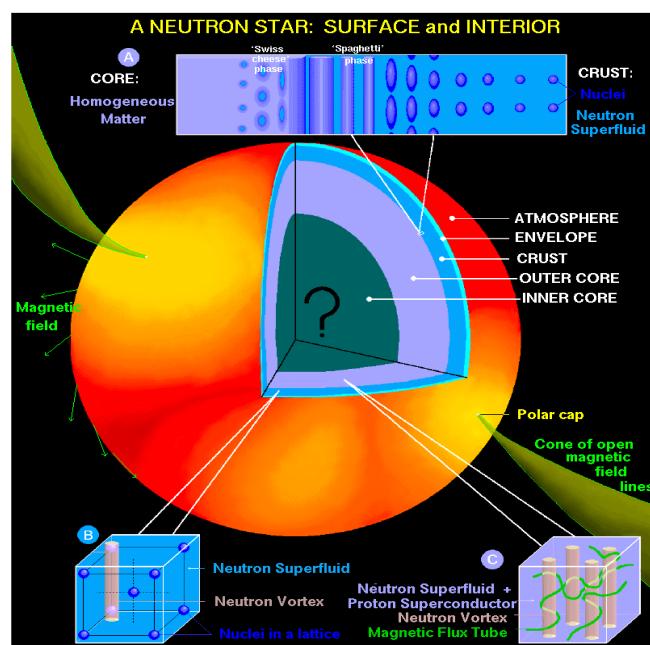
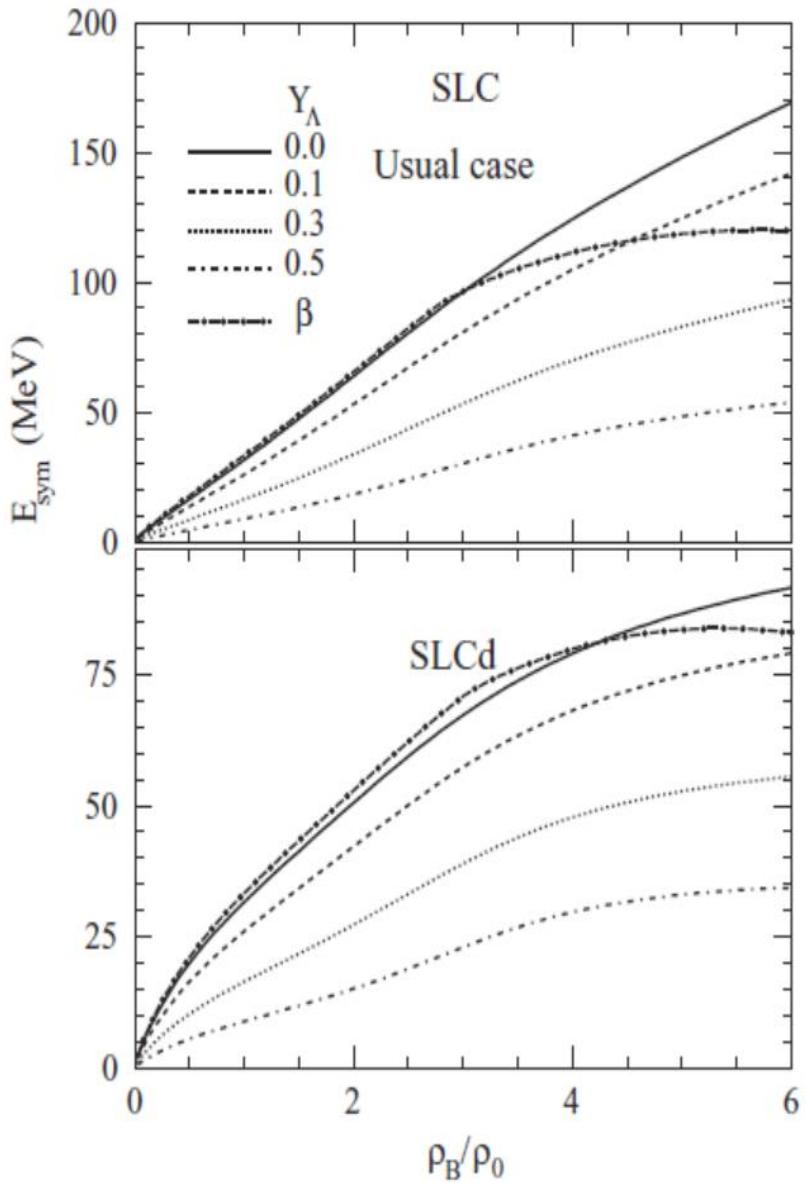
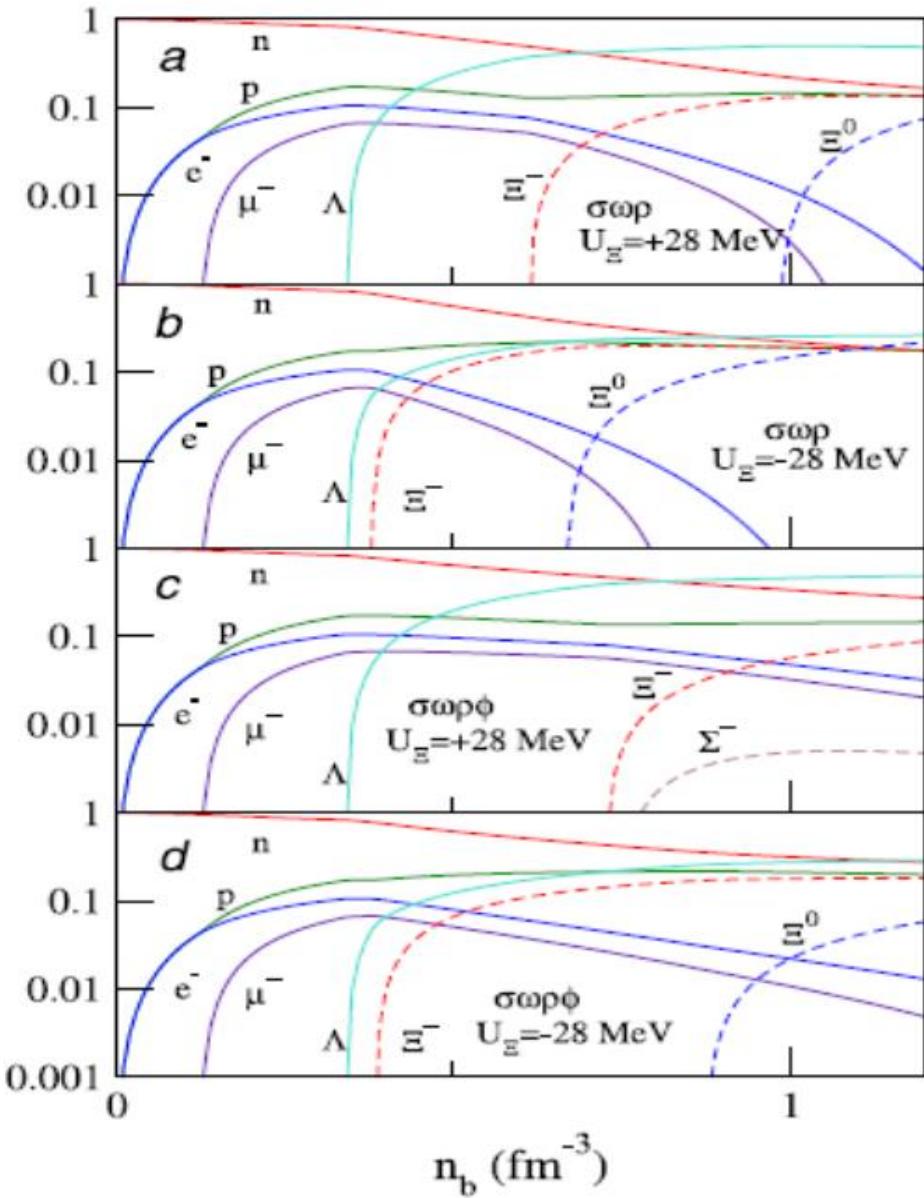
## 4. 致密物质中超子成分的影响

S. Weissenborn, D. Chatterjee, J. Schaffner-Bielich, Nucl. Phys. A 881, 62 (2012)

W. Z. Jiang, R. Y. Yang, and D. R. Zhang, Phys. Rev. C 87, 064314 (2013)

J. M. Lattimer and M. Prakash, Science 304, 536 (2004)

Particle fractions



## 核物质对称能

$$E_{\text{sym}}(\rho) = E_{\text{sym}}(\rho_0) + \frac{L}{3} \left( \frac{\rho - \rho_0}{\rho_0} \right) + \frac{K_{\text{sym}}}{18} \left( \frac{\rho - \rho_0}{\rho_0} \right)^2, \quad (\rho \sim \rho_0)$$

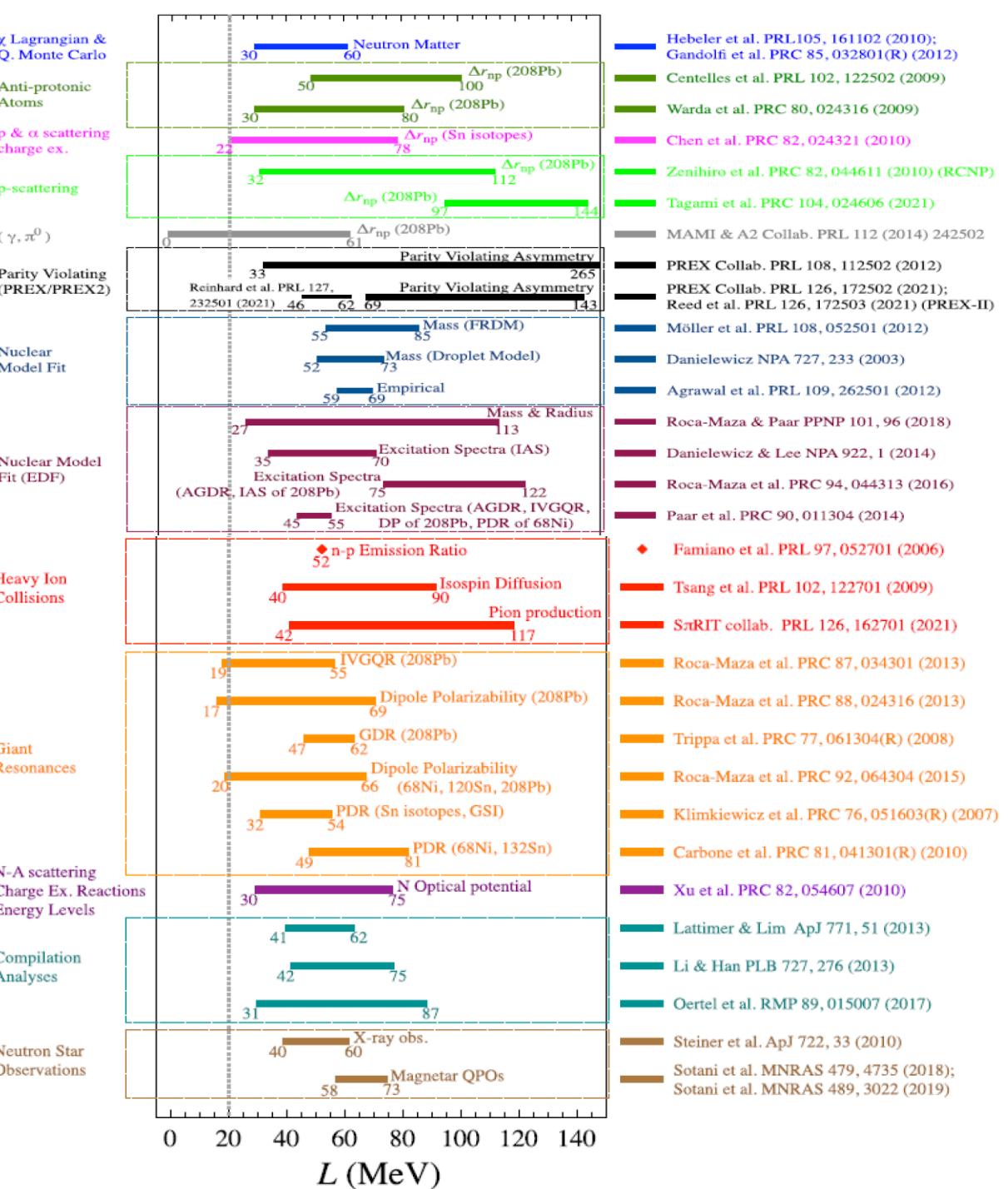
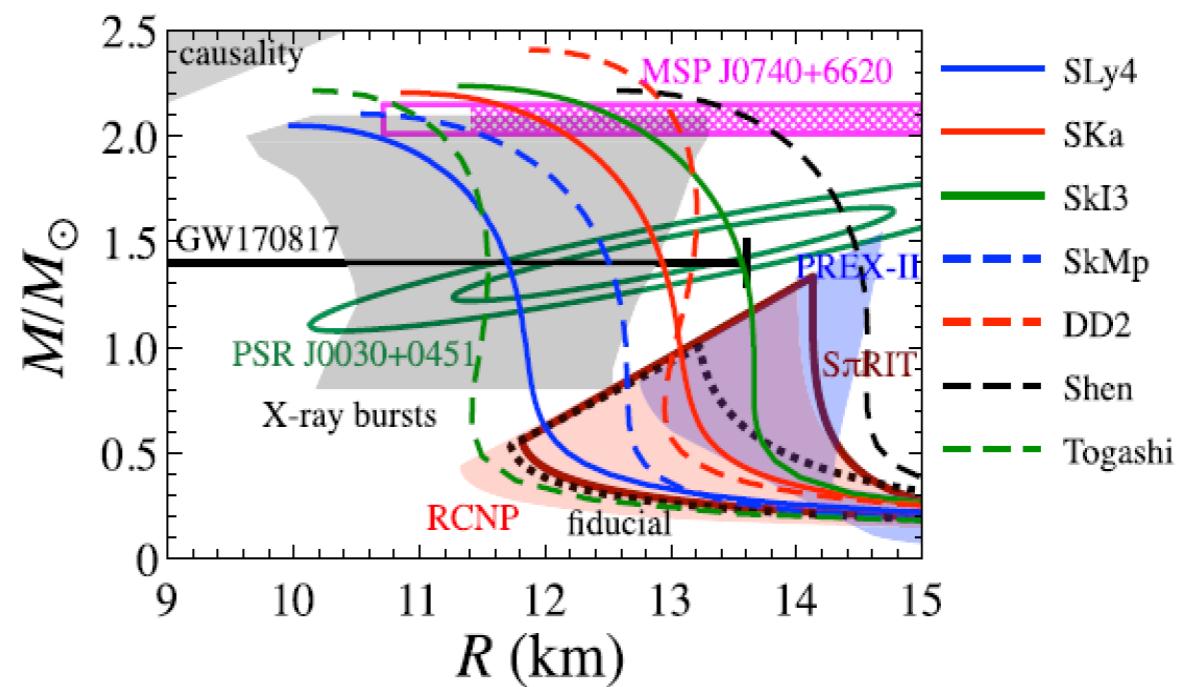
# 液滴模型

$$E_{\text{sym}}(\rho_0) \approx 32 \text{ MeV}$$

## 斜率和曲率参数

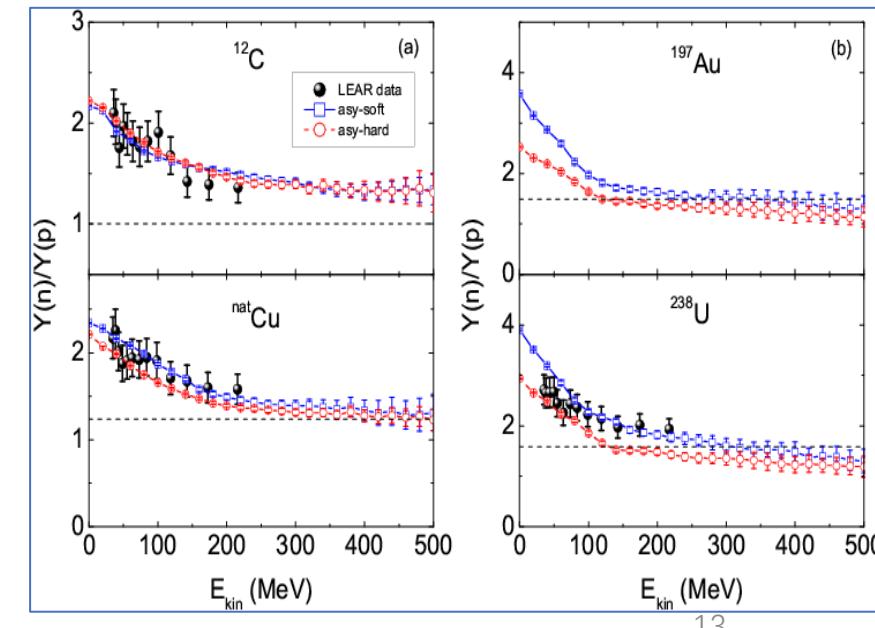
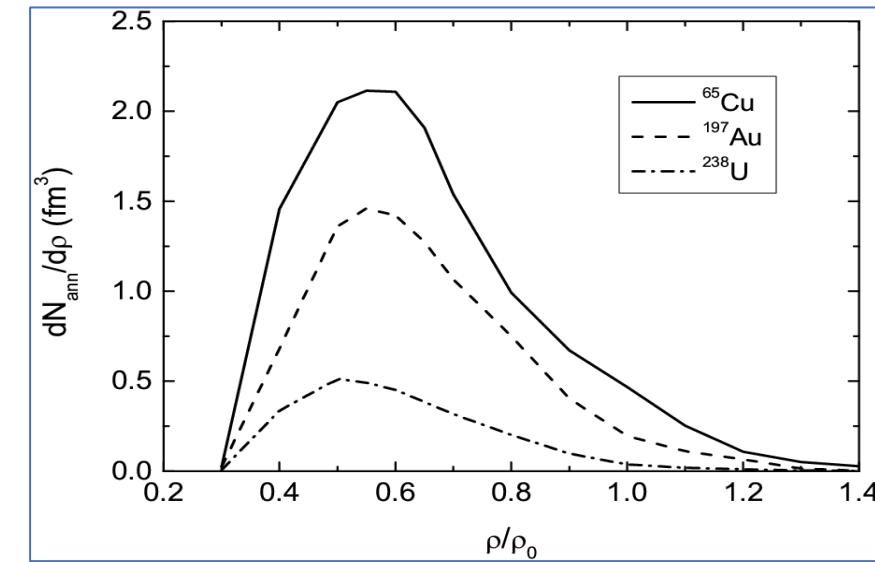
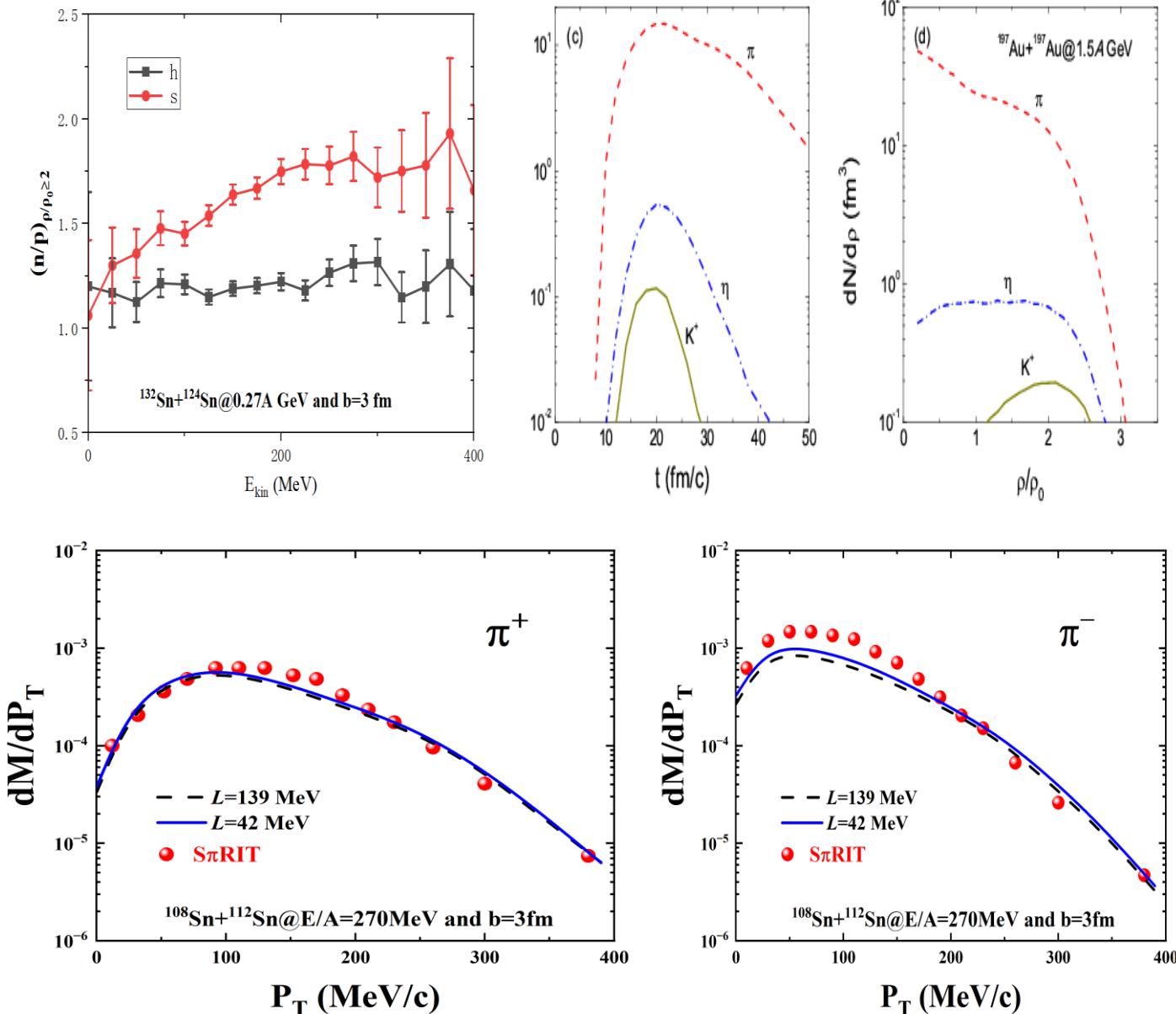
$$L \equiv 3\rho_0 \frac{\partial E_{\text{sym}}(\rho)}{\partial \rho} \Bigg|_{\rho=\rho_0} \quad K_{\text{sym}} \equiv 9\rho_0^2 \frac{\partial^2 E_{\text{sym}}(\rho)}{\partial \rho^2} \Bigg|_{\rho=\rho_0}$$

Hajime Sotani, Nobuya Nishimura, and Tomoya Naito, Prog. Theor. Exp. Phys. 041D01 (2022)



# 反质子轰击原子核中子和质子产额比值 $\gamma_s=0.5$ , L=53 MeV

## 重离子碰撞中同位旋耗散和介子产生



## Article

## Constraining neutron-star matter with microscopic and macroscopic collisions

<https://doi.org/10.1038/s41586-022-04750-w>

Received: 13 July 2021

Accepted: 11 April 2022

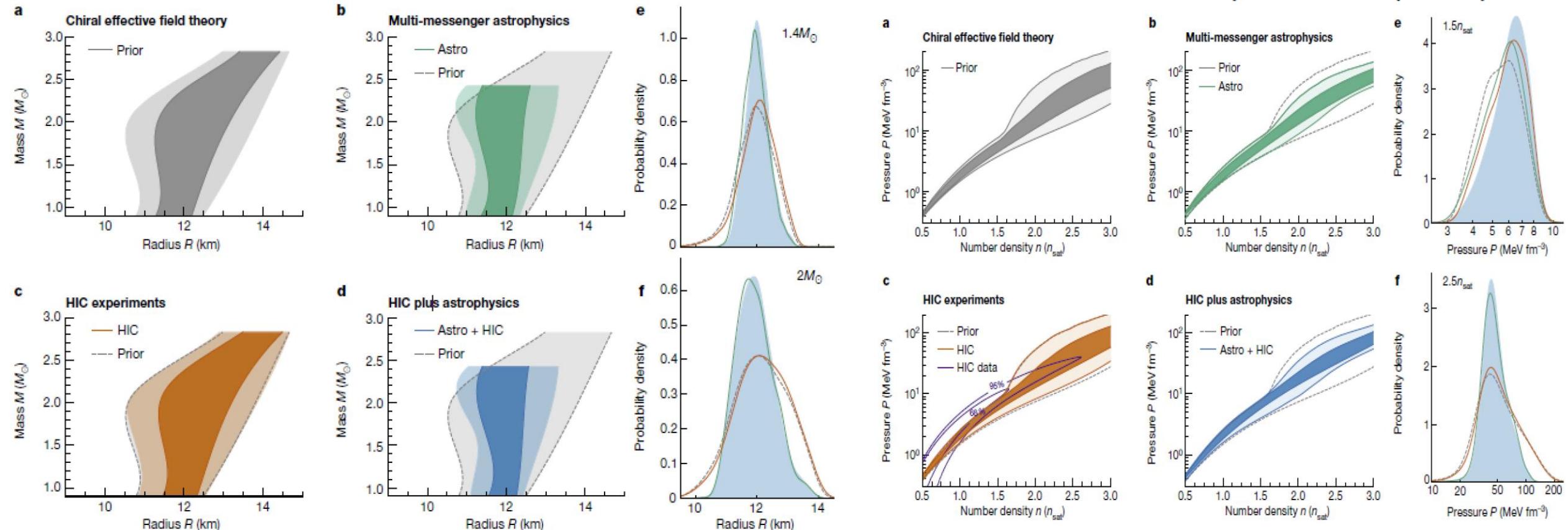
Published online: 8 June 2022

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Sabrina Huth<sup>1,2,13</sup>, Peter T. H. Pang<sup>3,4,13</sup>, Ingo Tews<sup>5</sup>, Tim Dietrich<sup>6,7</sup>, Arnaud Le Fèvre<sup>8</sup>, Achim Schwenk<sup>1,2,9</sup>, Wolfgang Trautmann<sup>8</sup>, Kshitij Agarwal<sup>10</sup>, Mattia Bulla<sup>11</sup>, Michael W. Coughlin<sup>12</sup> & Chris Van Den Broeck<sup>3,4</sup>

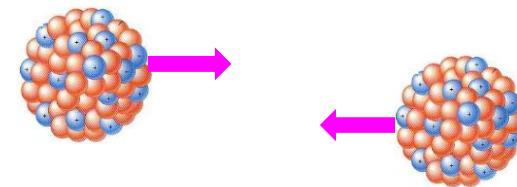
Interpreting high-energy, astrophysical phenomena, such as supernova explosions or neutron-star collisions, requires a robust understanding of matter at supranuclear densities. However, our knowledge about dense matter explored in the cores of neutron stars remains limited. Fortunately, dense matter is not probed only in



## 二、量子分子动力学模型 (LQMD)

Nuclear dynamics from 5 MeV/nucleon – 10 GeV/nucleon for HICs, antiproton (proton,  $\pi$ , K, etc)

- **Lanzhou quantum molecular dynamics** (Skyrme interaction, Walecka model with  $\sigma, \omega, \rho, \delta$ )
- **Isospin physics at intermediate energies** (constraining nuclear symmetry energy at sub- and supersaturation densities in HICs and probing isospin splitting of nucleon effective mass from HICs)
- **In-medium properties of hadrons in dense nuclear matter from heavy-ion collisions** (extracting optical potentials, i.e.,  $\Delta(1232)$ ,  $N^*(1440)$ ,  $N^*(1535)$ ), hyperons ( $\Lambda, \Sigma, \Xi, \Omega$ ) and mesons ( $\pi, K, \eta, \rho, \omega, \phi, \dots$ ), hypernucleus dynamics)
- **Hadron (antiproton, proton,  $\pi^\pm$ ,  $K^\pm$ ) induced reactions** (hypernucleus production, e.g.,  $\Lambda(\Sigma)X$ ,  $\Lambda\Lambda X$ ,  $\Xi X$ ,  $\bar{\Lambda}X(S=1)$ , in-medium modifications of hadrons, cold QGP)



# 1. 基于Skyrme相互作用的量子分子动力学模型 (LQMD-Skyrme)

PHYSICAL REVIEW C 84, 024610 (2011)

$$H_B = \sum_i \sqrt{\mathbf{p}_i^2 + \mathbf{m}_i^2} + U_{\text{int}} + U_{\text{mom}}$$

Momentum dependence of the symmetry potential and its influence on nuclear reactions

Zhao-Qing Feng\*

Institute of Modern Physics, Chinese Academy of Sciences, Lanzhou 730000, People's Republic of China

(Received 11 July 2011; published 19 August 2011)

$$U_{loc} = \int V_{loc}(\rho(\mathbf{r})) d\mathbf{r}$$

$$V_{loc}(\rho) = \frac{\alpha}{2} \frac{\rho^2}{\rho_0} + \frac{\beta}{1+\gamma} \frac{\rho^{1+\gamma}}{\rho_0^\gamma} + E_{sym}^{loc}(\rho) \rho \delta^2 + \frac{g_{sur}}{2\rho_0} (\nabla \rho)^2 + \frac{g_{sur}^{iso}}{2\rho_0} [\nabla(\rho_n - \rho_p)]^2,$$

$$U_{mom} = \frac{1}{2\rho_0} \sum_{i,j,j \neq i} \sum_{\tau,\tau'} C_{\tau,\tau'} \delta_{\tau,\tau_i} \delta_{\tau',\tau_j} \iiint d\mathbf{p} d\mathbf{p}' d\mathbf{r} f_i(\mathbf{r}, \mathbf{p}, t) \\ \times [\ln(\epsilon(\mathbf{p} - \mathbf{p}')^2 + 1)]^2 f_j(\mathbf{r}, \mathbf{p}', t).$$

$$E_{sym}(\rho) = \frac{1}{3} \frac{\hbar^2}{2m} \left( \frac{3}{2} \pi^2 \rho \right)^{2/3} + E_{sym}^{loc}(\rho) + E_{sym}^{mom}(\rho).$$

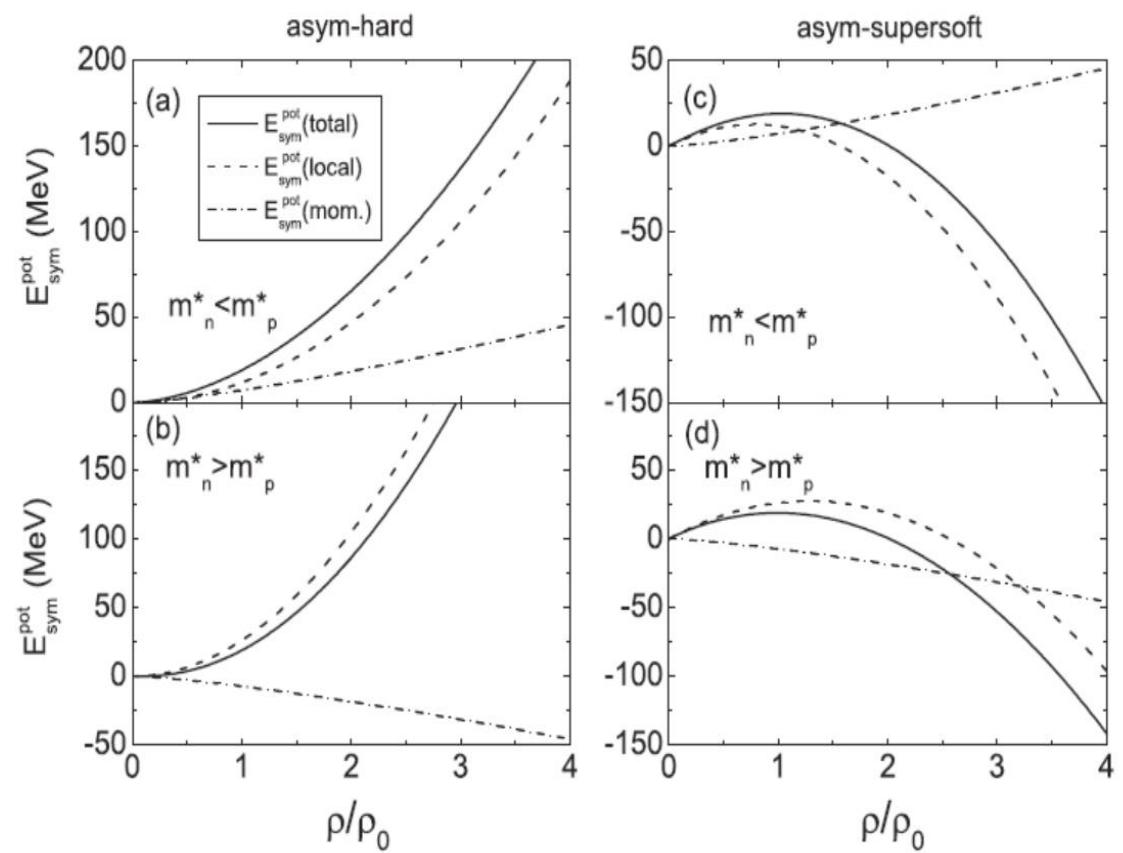
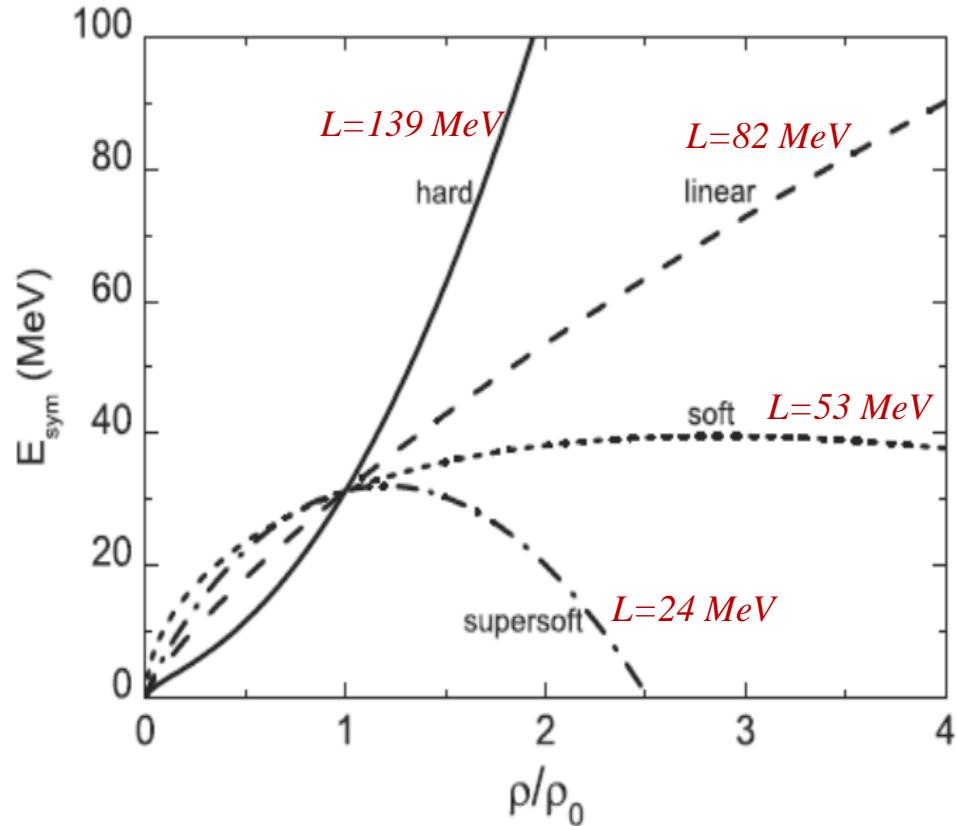
$$E_{sym}^{loc}(\rho) = \frac{1}{2} C_{sym} (\rho / \rho_0)^{\gamma_s}$$

$$E_{sym}^{loc}(\rho) = a_{sym} (\rho / \rho_0) + b_{sym} (\rho / \rho_0)^2.$$

**C<sub>sym</sub>**= 38 MeV  
**a<sub>sym</sub>**= 37.7 MeV  
**b<sub>sym</sub>**= -18.7 MeV

Table 1: The parameters and properties of isospin symmetric EoS used in the LQMD model at the density of  $0.16 \text{ fm}^{-3}$ .

Parameters	$\alpha$ (MeV)	$\beta$ (MeV)	$\gamma$	$C_{mom}$ (MeV)	$\epsilon$ ( $c^2/\text{MeV}^2$ )	$m_\infty^*/m$	$K_\infty$ (MeV)
PAR1	-215.7	142.4	1.322	1.76	$5 \times 10^{-4}$	0.75	230
PAR2	-226.5	173.7	1.309	0.	0.	1.	230



## 2. 相对论量子分子动力学(LQMD-RMF)

$$\begin{aligned}
L = & \bar{\psi} [i\gamma_\mu \partial^\mu - (M_N - g_\sigma \varphi - g_\delta \vec{\tau} \cdot \vec{\delta}) - g_\omega \gamma_\mu \omega^\mu - g_\rho \gamma_\mu \vec{\tau} \cdot \vec{b}^\mu] \psi \\
& + \frac{1}{2} (\partial_\mu \varphi \partial^\mu \varphi - m_\sigma^2 \varphi^2) - U(\varphi) + \frac{1}{2} (\partial_\mu \vec{\delta} \partial^\mu \vec{\delta} - m_\delta^2 \vec{\delta}^2) \\
& + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{b}_\mu \vec{b}^\mu - \frac{1}{4} \vec{G}_{\mu\nu} \vec{G}^{\mu\nu}
\end{aligned}$$

$$F_{\mu\nu} = \partial_\mu \omega_\nu - \partial_\nu \omega_\mu,$$

$$G_{\mu\nu} = \partial_\mu \vec{b}_\nu - \partial_\nu \vec{b}_\mu,$$

$$U(\varphi) = \frac{g_2}{3} \varphi^3 + \frac{g_3}{4} \varphi^4$$

能量密度：

$$\varepsilon = \sum_{i=n,p} 2 \int \frac{d^3 k}{(2\pi)^3} \sqrt{k^2 + M_i^{*2}} + \frac{1}{2} m_\sigma^2 \varphi^2 + U(\varphi) + \frac{1}{2} m_\omega^2 \omega_0^2 + \frac{1}{2} m_\rho^2 b_0^2 + \frac{1}{2} m_\delta^2 \delta_0^2$$

核子的相空间演化：

$$\begin{aligned}
\dot{\mathbf{x}} = & \frac{\mathbf{p}_i^*}{p_0^*} + \sum_{i \neq j}^N \left\{ \frac{g_v^2}{2m_v^2} z_j^{*\mu} u_{i,\mu} B_i B_j \frac{\partial \rho_{ij}}{\partial \mathbf{p}_i} + \frac{g_v^2}{2m_v^2} z_i^{*\mu} u_{j,\mu} B_i B_j \frac{\partial \rho_{ji}}{\partial \mathbf{p}_i} + \frac{g_v^2}{2m_v^2} z_j^{*\mu} \rho_{ji} B_i B_j \frac{\partial u_{i,\mu}}{\partial \mathbf{p}_i} \right. \\
& + z_j^{*\mu} \frac{B_i B_j \bar{g}_v^2}{2m_v^2} \left[ \frac{\rho_{ij}}{1 - p_{T,ij}^2/\Lambda_v^2} \frac{\partial u_{i,\mu}}{\partial \mathbf{p}_i} + \frac{u_{i,\mu}}{1 - p_{T,ij}^2/\Lambda_v^2} \frac{\partial \rho_{ij}}{\partial \mathbf{p}_i} + u_{i,\mu} \rho_{ij} \frac{\partial [1/(1 - p_{T,ij}^2/\Lambda_v^2)]}{\partial \mathbf{p}_i} \right] \\
& + z_i^{*\mu} \frac{B_i B_j \bar{g}_v^2}{2m_v^2} \left[ \frac{u_{j,\mu}}{1 - p_{T,ji}^2/\Lambda_v^2} \frac{\partial \rho_{ji}}{\partial \mathbf{p}_i} + u_{j,\mu} \rho_{ji} \frac{\partial [1/(1 - p_{T,ji}^2/\Lambda_v^2)]}{\partial \mathbf{p}_i} \right] \\
& \left. - \frac{m_j^*}{p_j^{*0}} \frac{\partial S_j}{\partial \mathbf{p}_i} - \frac{m_i^*}{p_i^{*0}} \frac{\partial S_i}{\partial \mathbf{p}_i} \right\},
\end{aligned}$$

$$\begin{aligned}
\dot{\mathbf{p}} = & - \sum_{i \neq j}^N \left\{ \frac{g_v^2}{2m_v^2} z_j^{*\mu} u_{i,\mu} B_i B_j \frac{\partial \rho_{ij}}{\partial \mathbf{r}_i} + \frac{g_v^2}{2m_v^2} z_i^{*\mu} u_{j,\mu} B_i B_j \frac{\partial \rho_{ji}}{\partial \mathbf{r}_i} \right. \\
& + z_j^{*\mu} \frac{B_i B_j \bar{g}_v^2}{2m_v^2} \frac{u_{i,\mu}}{1 - p_{T,ij}^2/\Lambda_v^2} \frac{\partial \rho_{ij}}{\partial \mathbf{r}_i} \\
& + z_i^{*\mu} \frac{B_i B_j \bar{g}_v^2}{2m_v^2} \frac{u_{j,\mu}}{1 - p_{T,ji}^2/\Lambda_v^2} \frac{\partial \rho_{ji}}{\partial \mathbf{r}_i} \\
& \left. - \frac{m_j^*}{p_j^{*0}} \frac{\partial S_j}{\partial \mathbf{r}_i} - \frac{m_i^*}{p_i^{*0}} \frac{\partial S_i}{\partial \mathbf{r}_i} \right\},
\end{aligned}$$

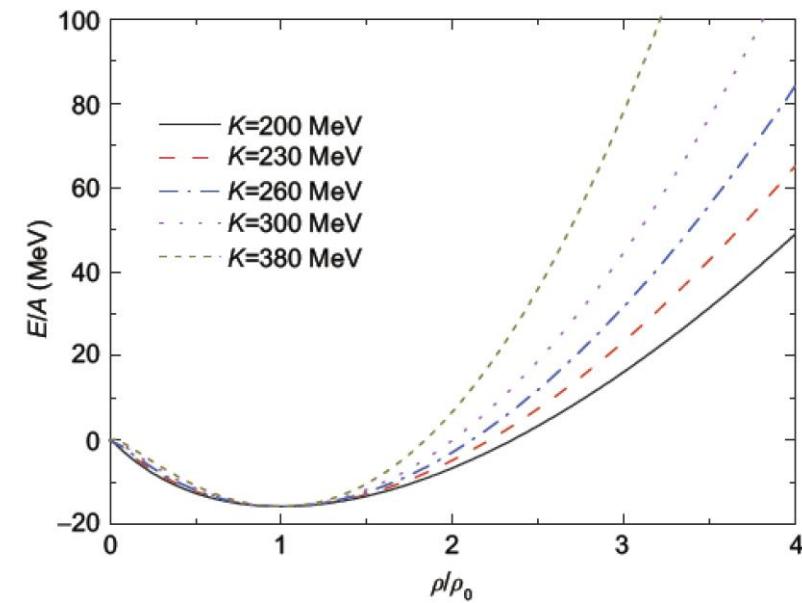
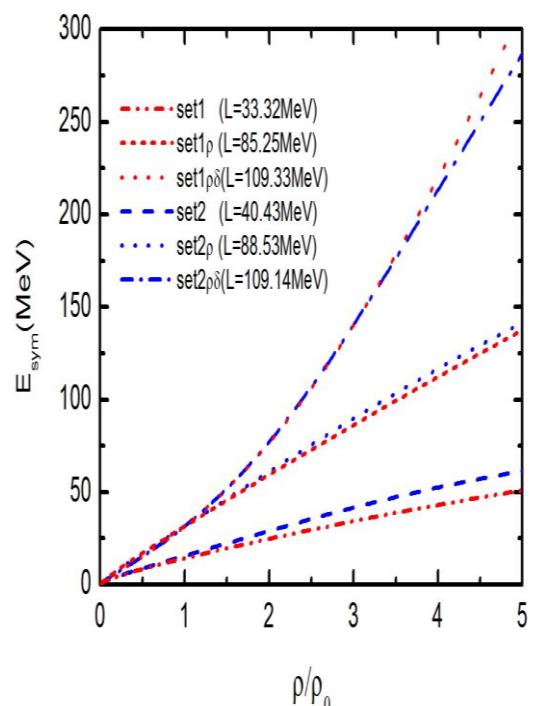
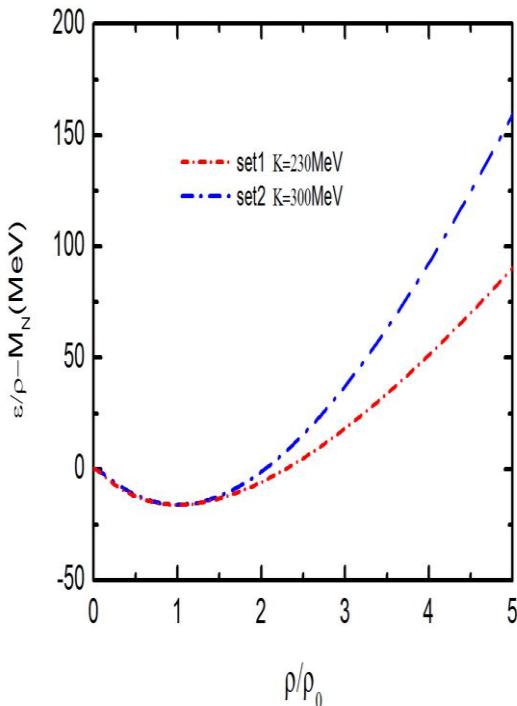
TABLE I: Parameters for the relativistic mean-field theory with non-linear scalar interaction for a binding energy of  $E_B = \epsilon/\rho - M_N = -16$  MeV and for normal nuclear matter density of  $\rho_0 = 0.16 fm^{-3}$ . The  $\sigma$  meson mass  $m_\sigma$ , the  $\omega$  meson mass  $m_\omega$ , the  $\rho$  meson mass  $m_\rho$  and the  $\delta$  meson mass  $m_\delta$  are set to be 550, 783, 763 and 500 MeV, respectively.

Model	$g_\sigma$	$g_\omega$	$g_2(fm^{-1})$	$g_3$	$g_\rho$	$g_\delta$	$M^*/M_N$	$E_B(\text{MeV})$	$K(\text{MeV})$	$E_{sym}(\text{MeV})$	$L(\text{MeV})$
set1	8.145	7.570	31.900	21.800	-	-	0.813	-16.0	230	14.3	33.3
set1 $\rho$	8.145	7.570	31.900	21.800	4.049	-	0.813	-16.0	230	31.6	85.3
set1 $\rho\delta$	8.145	7.570	31.900	21.800	8.673	5.347	0.813	-16.0	230	31.6	109.3
set2	8.830	9.500	11.310	13.750	-	-	0.738	-16.0	300	15.6	40.4
set2 $\rho$	8.830	9.500	11.310	13.750	3.897	-	0.738	-16.0	300	31.6	88.5
set2 $\rho\delta$	8.830	9.500	11.310	13.750	7.219	4.280	0.738	-16.0	300	31.6	109.4

# 核物质对称能

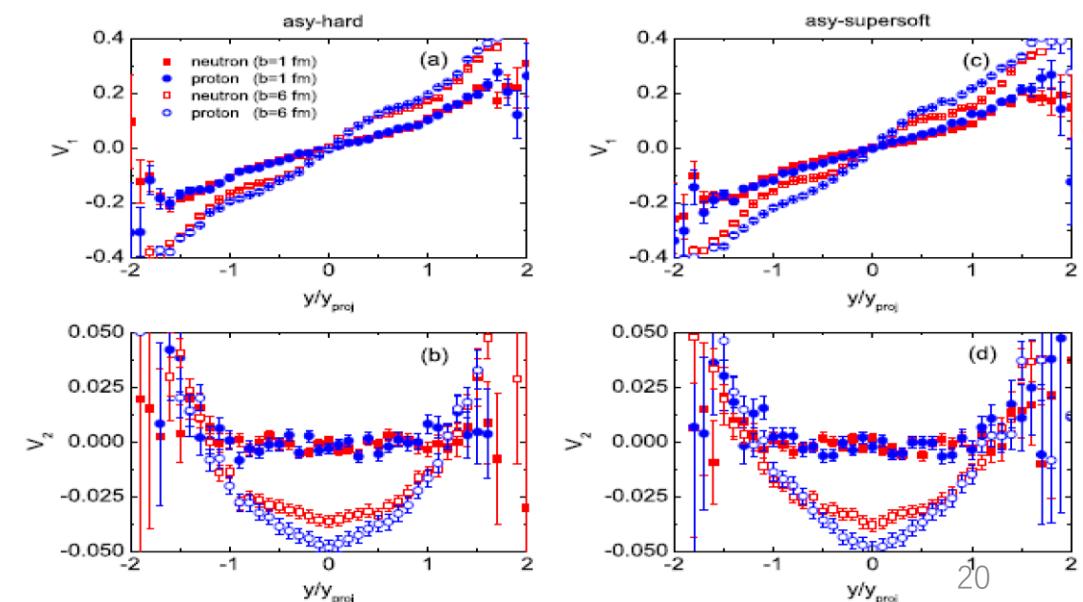
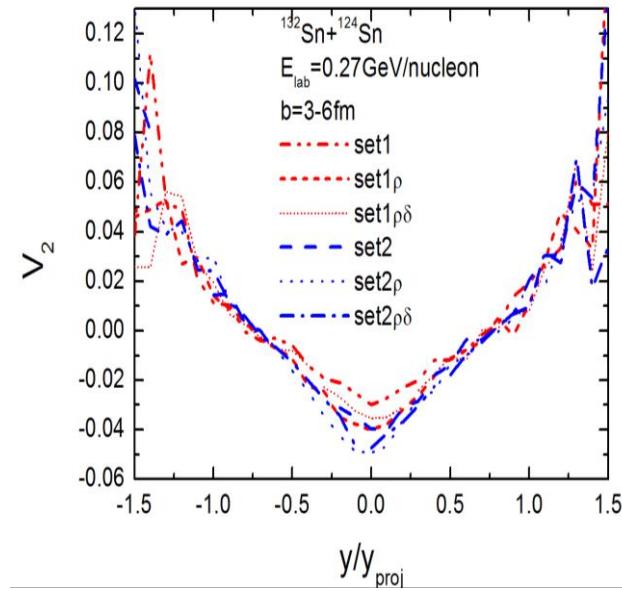
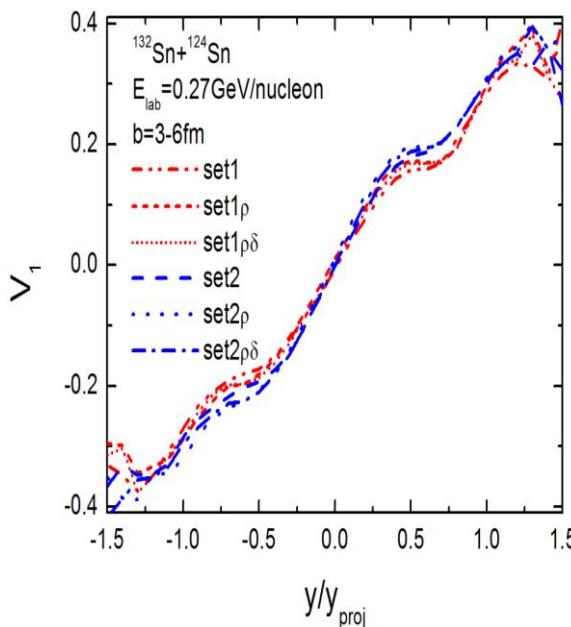
$$E_{sym} = \frac{1}{6} \frac{k_F^2}{E_F^*} + \frac{1}{2} \left[ f_\rho - f_\delta \left( \frac{M^*}{E_F^*} \right) \right] \rho$$

$$f_{\rho,\delta} = g_{\rho,\delta}/m_{\rho,\delta}$$



Z.-Q. Feng / Nuclear Physics A 878 (2012) 3–13

9



### 3. 粒子产生反应道

$\pi$  and resonances ( $\Delta(1232)$ ,  $N^*(1440)$ ,  $N^*(1535)$ , ...) production:

$$\begin{aligned} NN &\leftrightarrow N\Delta, \quad NN \leftrightarrow NN^*, \quad NN \leftrightarrow \Delta\Delta, \quad \Delta \leftrightarrow N\pi, \\ N^* &\leftrightarrow N\pi, \quad NN \leftrightarrow NN\pi(s-state), \quad N^*(1535) \leftrightarrow N\eta \end{aligned}$$

Collisions between resonances,  $NN^* \leftrightarrow N\Delta$ ,  $NN^* \leftrightarrow NN^*$

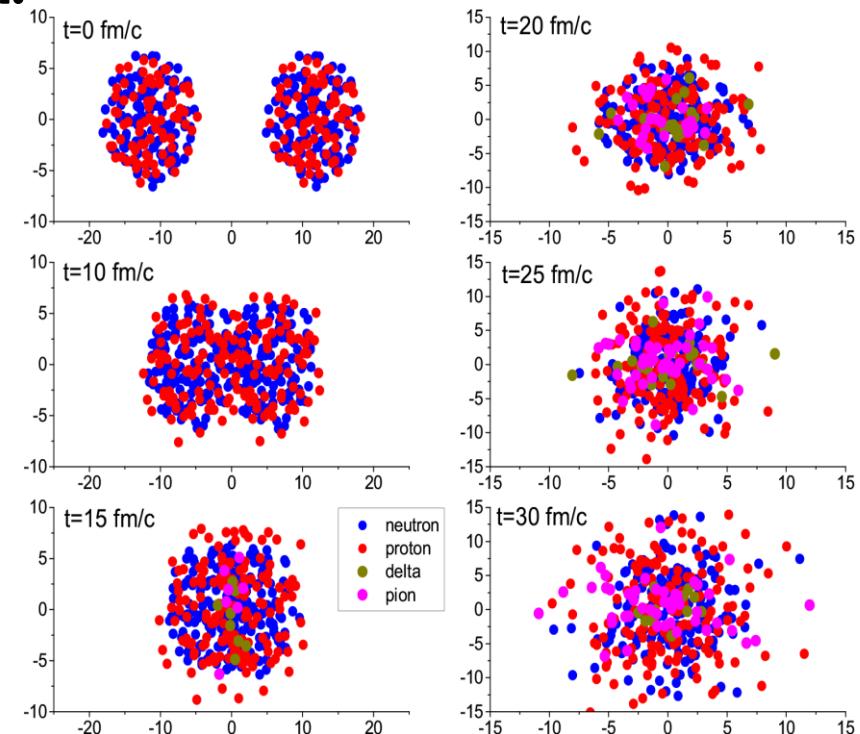
Strangeness channels:

$$\begin{aligned} BB &\rightarrow BYK, BB \rightarrow BB\bar{K}, B\pi(\eta) \rightarrow YK, YK \rightarrow B\pi, \\ B\pi &\rightarrow NK\bar{K}, Y\pi \rightarrow B\bar{K}, \quad B\bar{K} \rightarrow Y\pi, \quad YN \rightarrow \bar{K}NN, \\ BB &\rightarrow B\Xi KK, \bar{K}B \leftrightarrow K\Xi, YY \leftrightarrow N\Xi, \bar{K}Y \leftrightarrow \pi\Xi. \end{aligned}$$

Reaction channels with antiproton:

$$\begin{aligned} \bar{p}N &\rightarrow \bar{N}N, \quad \bar{N}N \rightarrow \bar{N}N, \quad \bar{N}N \rightarrow \bar{B}B, \quad \bar{N}N \rightarrow \bar{Y}Y \\ \bar{N}N &\rightarrow \text{annihilation}(\pi, \eta, \rho, \omega, K, \bar{K}, K^*, \bar{K}^*, \phi) \end{aligned}$$

The PYTHIA and FRITIOF code are used for baryon(meson)-baryon and antibaryon-baryon collisions at high invariant energies

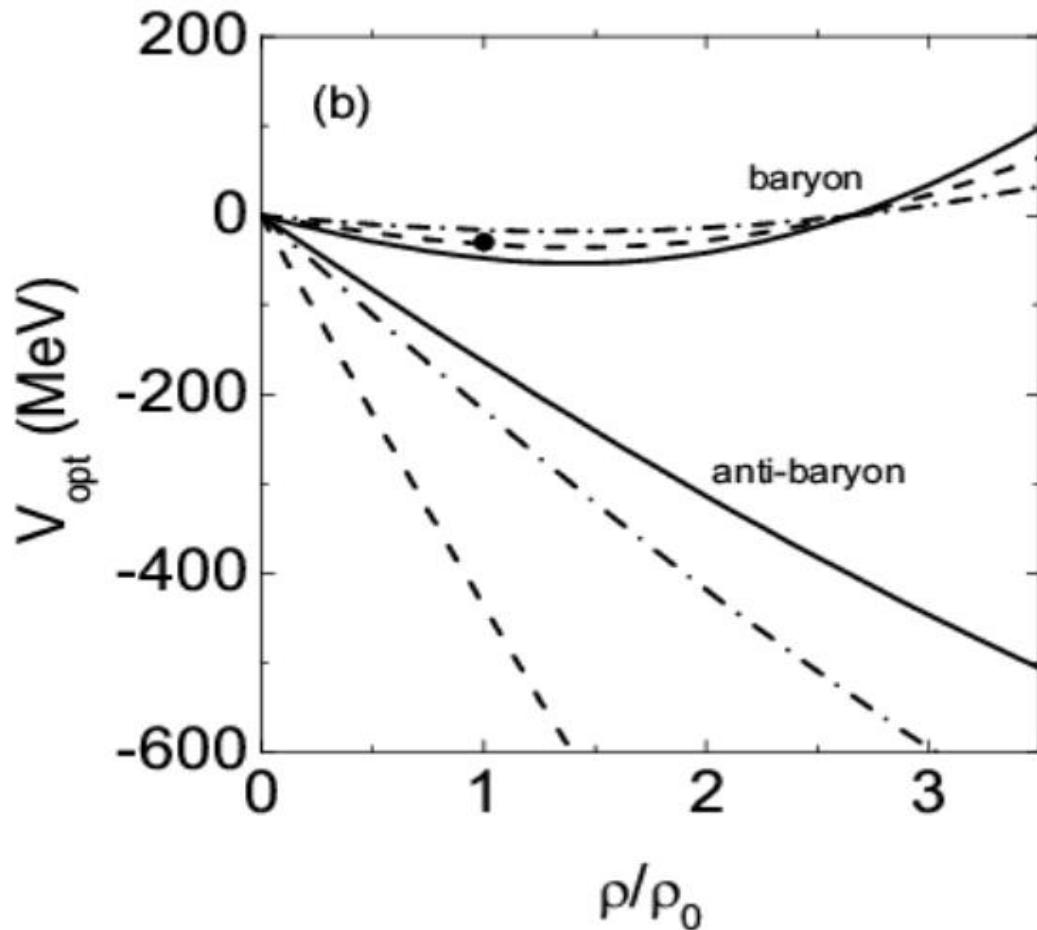
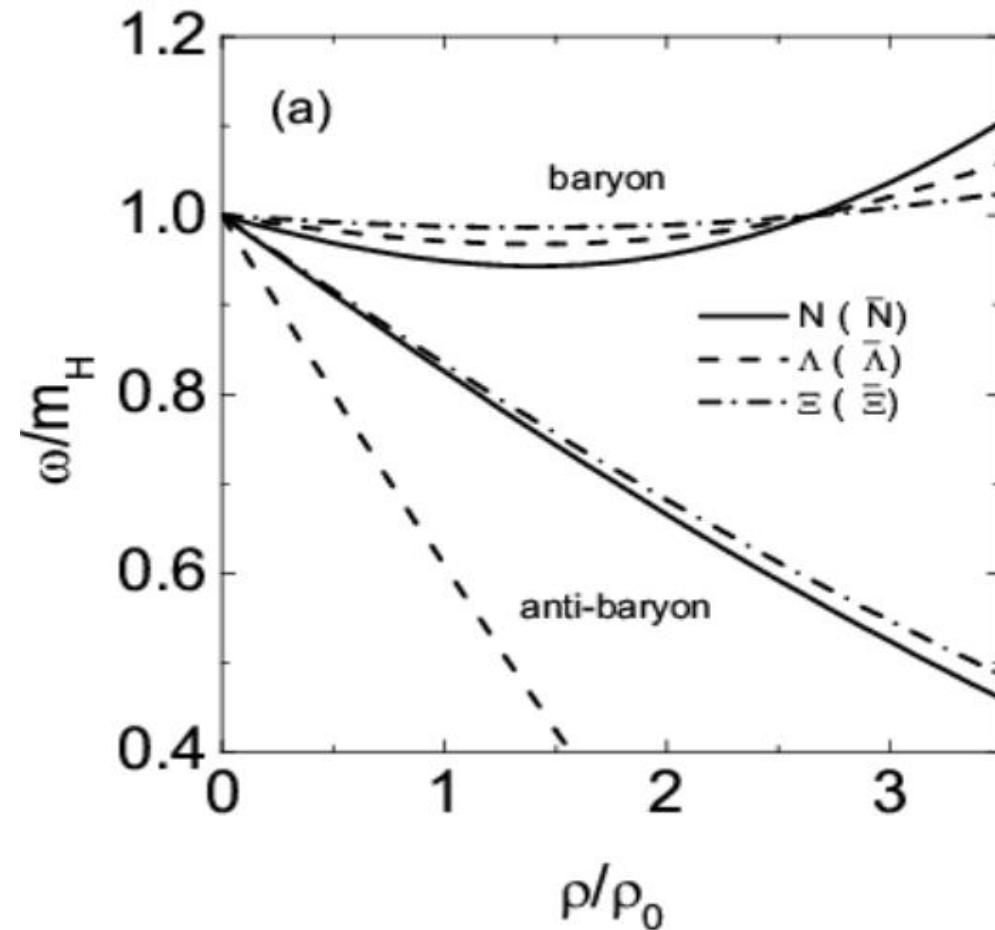


Statistical model with SU(3)  
symmetry for annihilation  
(E.S. Golubeva et al., Nucl. Phys. A 537, 393  
(1992))

## 4. 超子(反超子)平均场

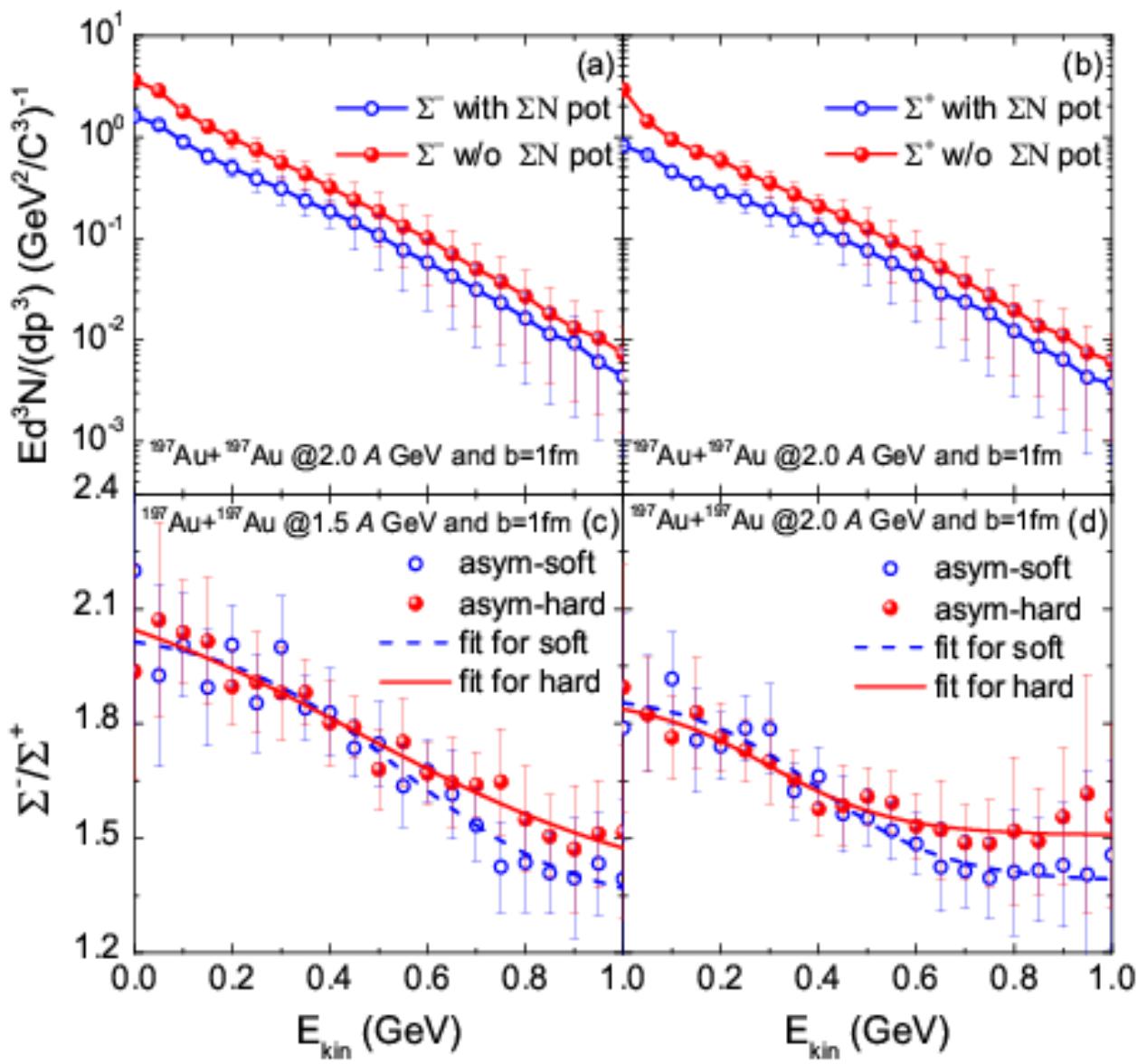
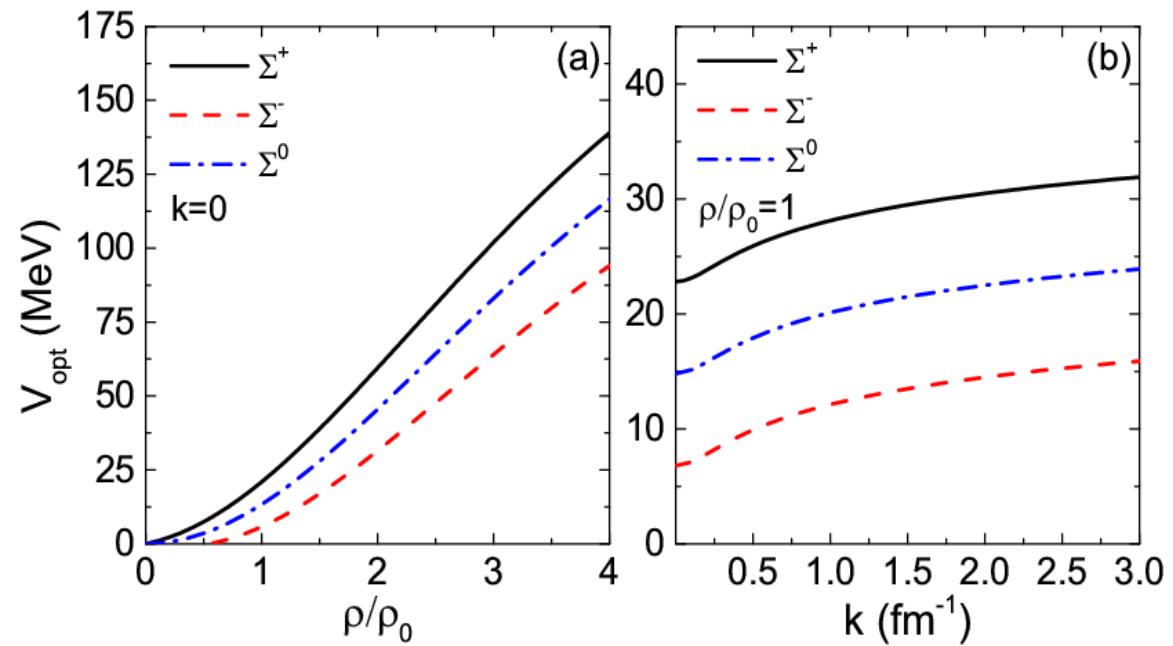
A factor  $\xi$  is introduced in evaluating self-energies of the antinucleon, e.g.,  $\xi=0.25$  for  $V_{\bar{N}N} = -160 \text{ MeV}$  at  $\rho = \rho_0$

$$H_M = \sum_{i=1}^{N_M} (V_i^{\text{Coul}} + \omega(\mathbf{p}_i, \rho_i)) \quad \omega(\mathbf{p}_i, \rho_i) = \sqrt{(m_H + \Sigma_S^H)^2 + \mathbf{p}_i^2} + \Sigma_V^H \quad V_{opt}(\mathbf{p}, \rho) = \omega(\mathbf{p}, \rho) - \sqrt{\mathbf{p}^2 + m^2}$$



Ding-Chang Zhang, Hui-Gan Cheng and Zhao-Qing Feng.  
*Chinese Physics Letters* 38 (2021) 092501.  
(arXiv: 2107.00277, editor's suggestion)

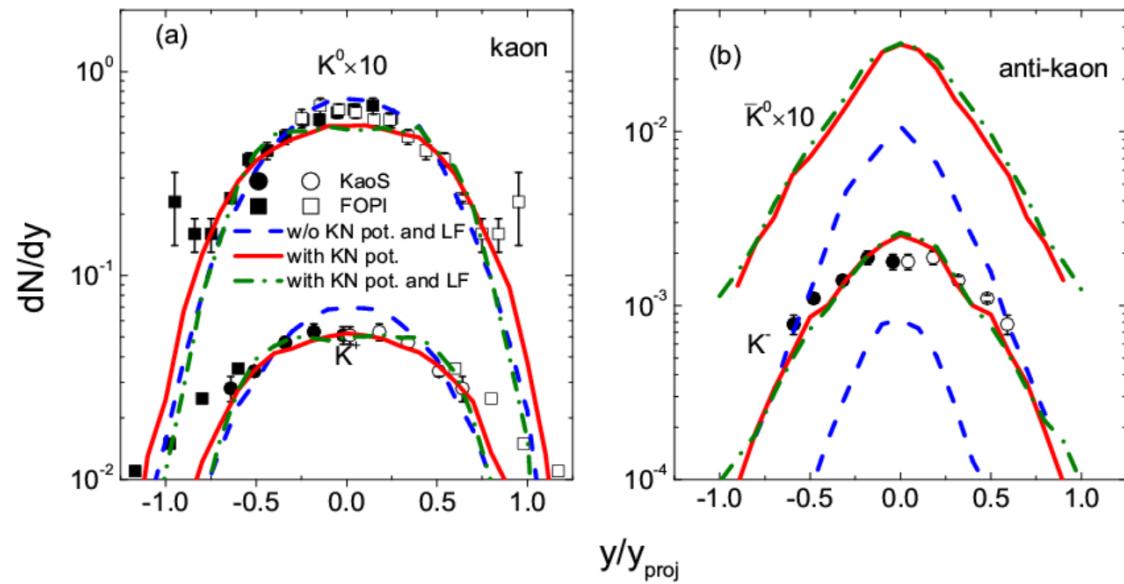
$$V_{\text{opt}}^{\Sigma}(p_i, \rho_i) = V_0(\rho_i/\rho_0)^{\gamma_s} + V_1(\rho_n - \rho_p)t_{\Sigma}\rho_i^{\gamma_s-1}/\rho_0^{\gamma_s} + C_{\text{mom}}\rho_i \ln(\epsilon p_i^2 + 1)$$



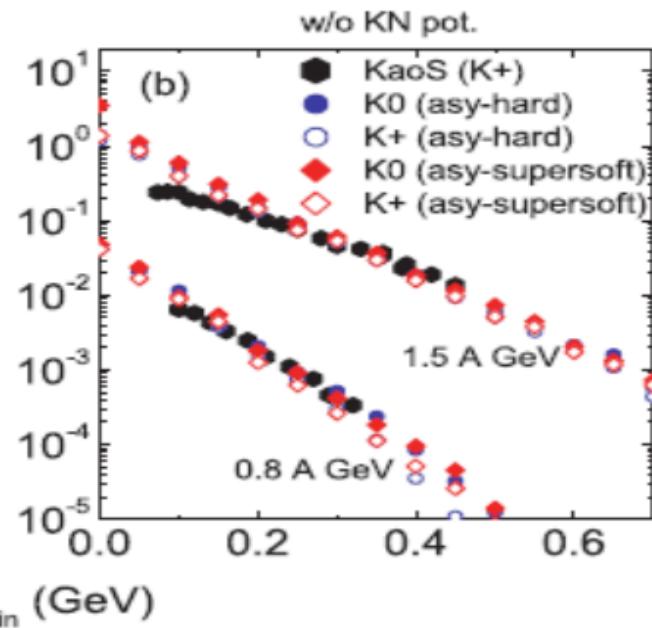
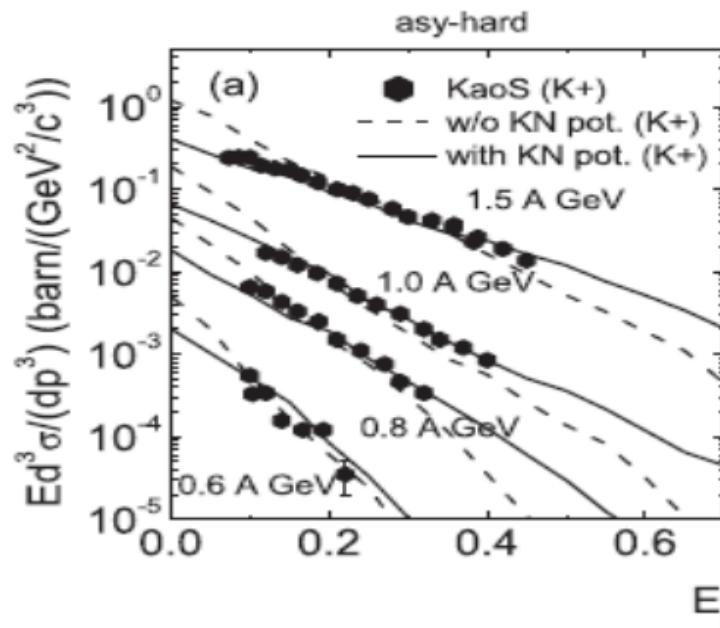
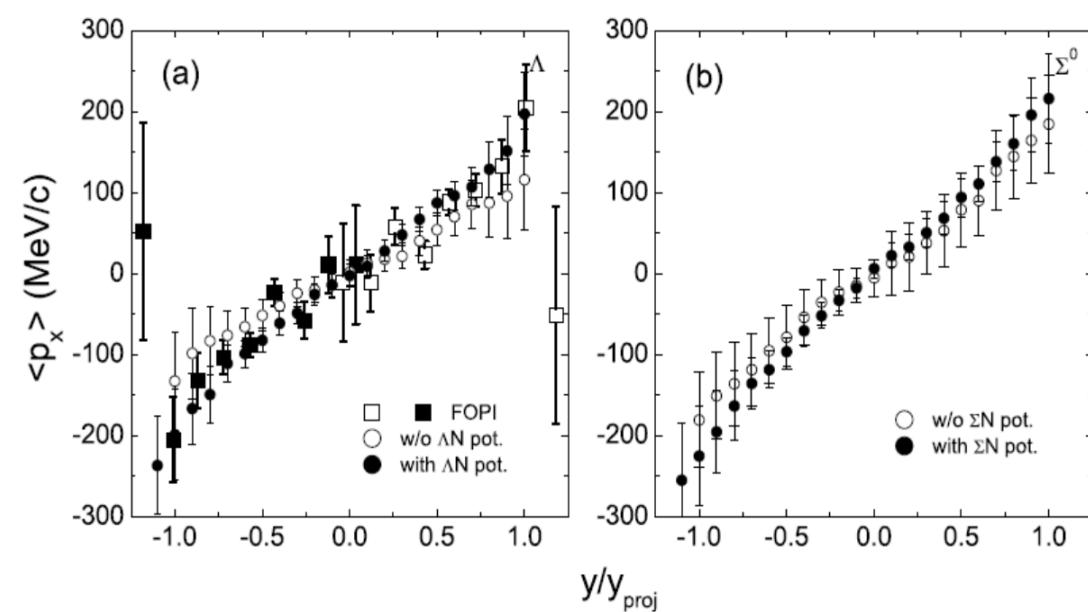
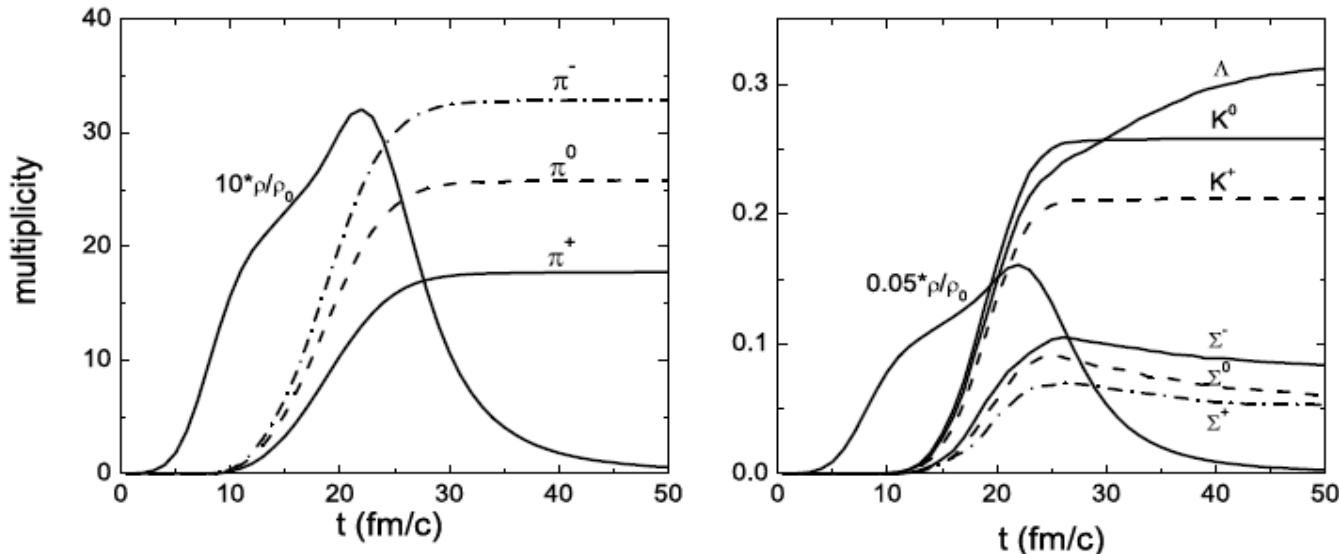
# 5. 重离子碰撞中K介子和超子产生

Phys. Rev. C 82 (2010) 057901; Phys. Rev. C 87, 064605 (2013); Nuclear Physics A 919 (2013) 32-45

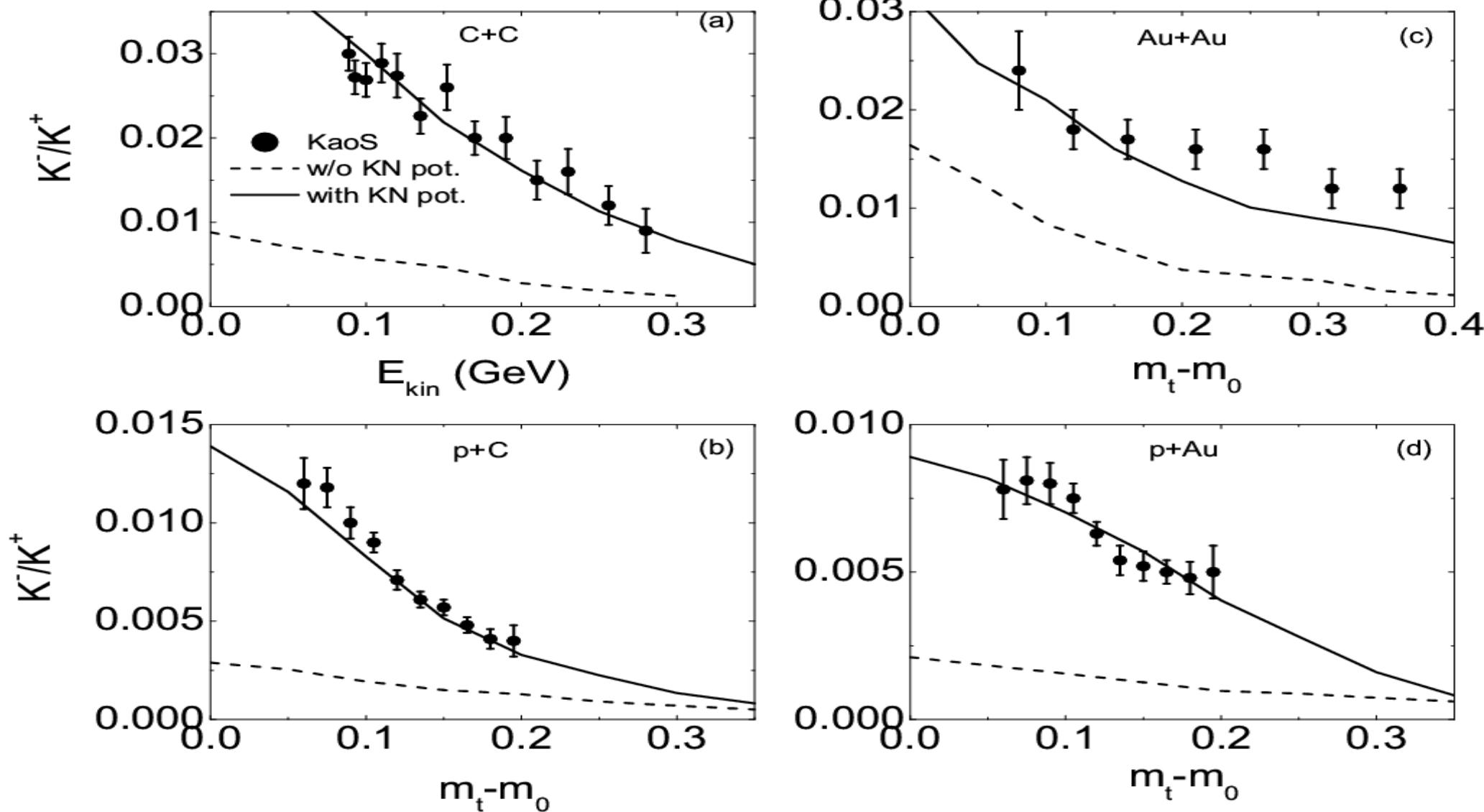
$^{58}\text{Ni} + ^{58}\text{Ni}$  @ 1.93A GeV



$^{197}\text{Au} + ^{197}\text{Au}$  @ 1.5A GeV



## The $K^-/K^+$ ratio in collisions of $^{12}\text{C} + ^{12}\text{C}$ at 1.84 GeV and protons on $^{12}\text{C}$ and $^{197}\text{Au}$ with 2.5 GeV

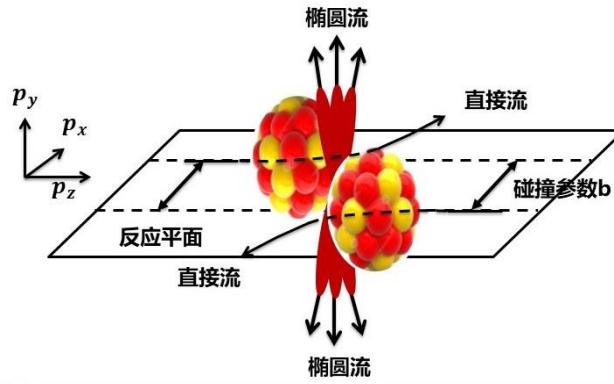
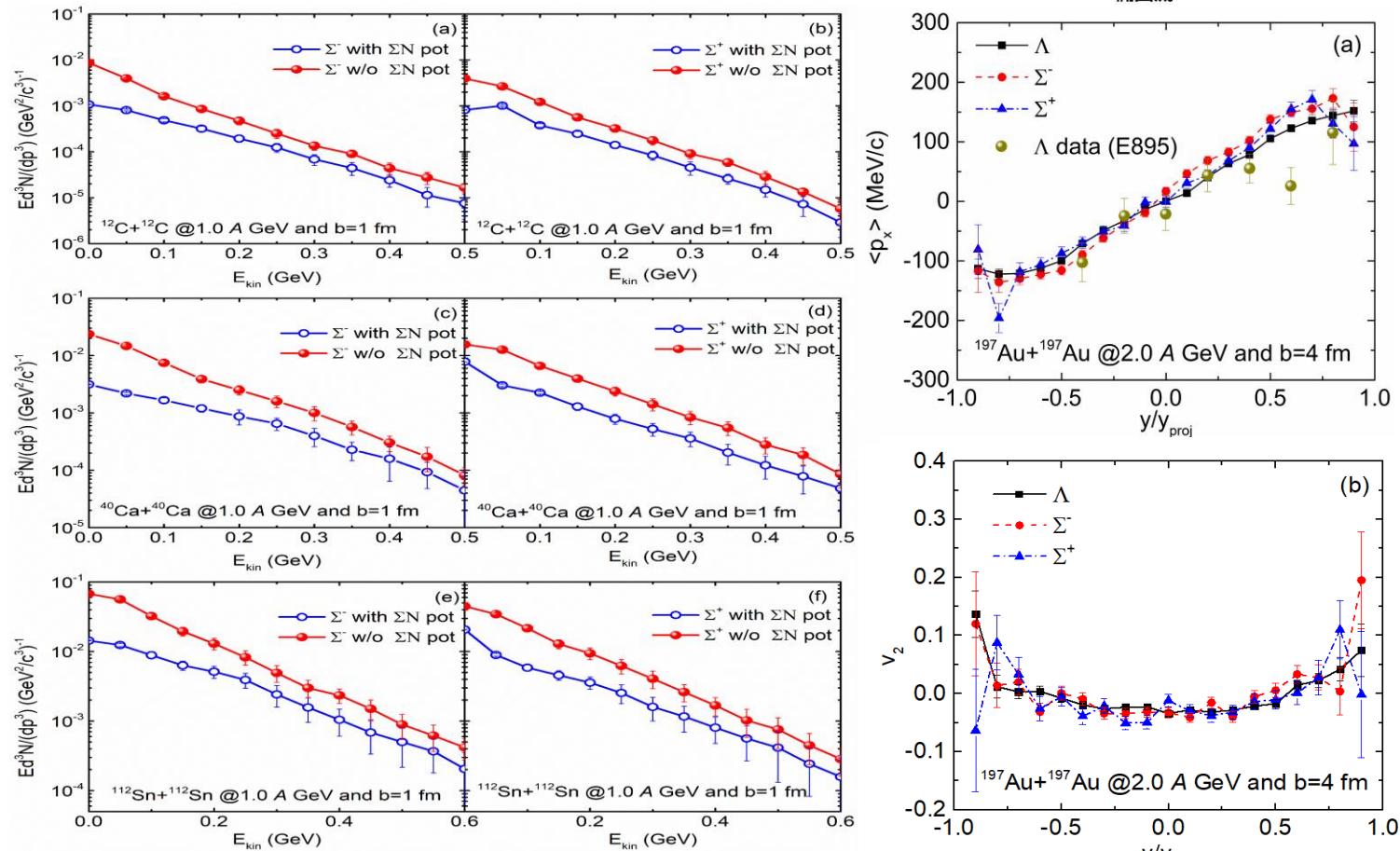
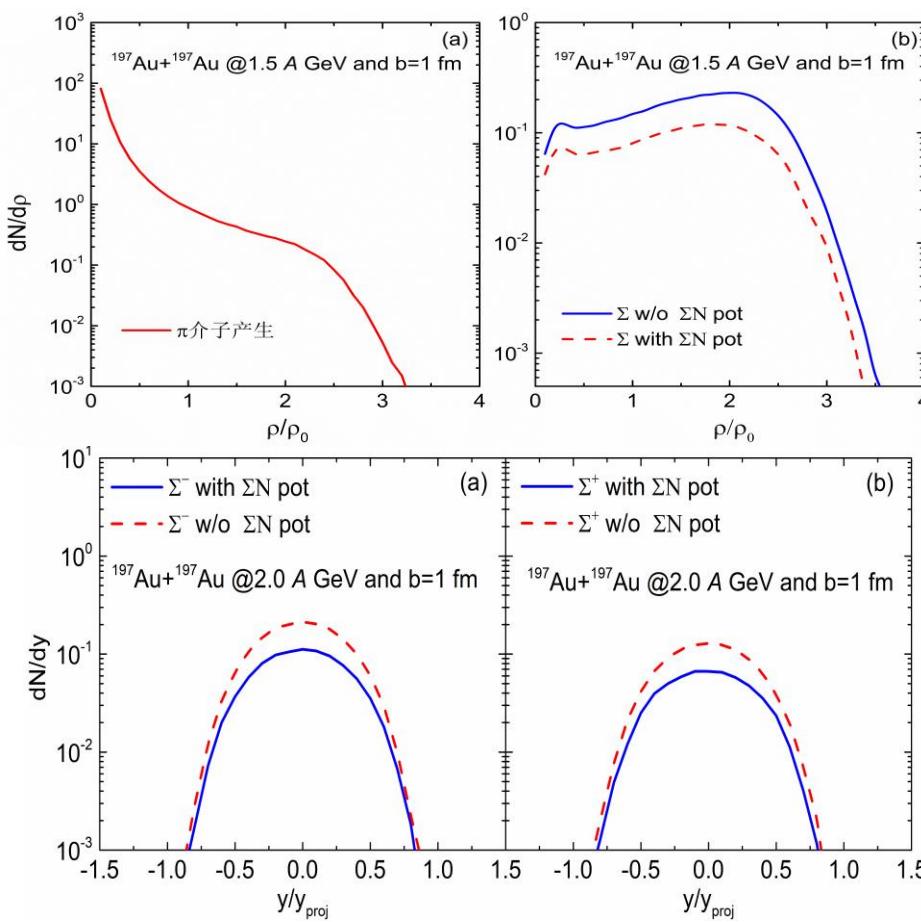


$$V_{K^+}(\rho_0) = 28 \text{ MeV}, \quad V_{K^-}(\rho_0) = -100 \text{ MeV}$$

重离子碰撞中 $\Sigma$ 超子分布 (Ding-Chang Zhang, Hui-Gan Cheng and Zhao-Qing Feng,  
Chinese Physics Letters 38 (2021) 092501)

出射粒子的角分布可按傅里叶展开表示为:  $\frac{dN}{d\phi}(y, p_t) = N_0[1 + 2v_1(y, p_t)\cos\phi + 2v_2(y, p_t)\cos 2\phi]$

其中 $\phi$ 为出射粒子的方位角 ( $\tan\phi = p_y/p_x$ )





# 三、原子核碎裂反应和超核形成

Physics Reports 510 (2012) 119–200



Contents lists available at SciVerse ScienceDirect

Physics Reports

journal homepage: [www.elsevier.com/locate/physrep](http://www.elsevier.com/locate/physrep)



Physics Reports 512 (2012) 1–124



Contents lists available at SciVerse ScienceDirect

Physics Reports

journal homepage: [www.elsevier.com/locate/physrep](http://www.elsevier.com/locate/physrep)



Strangeness production close to the threshold in proton–nucleus and heavy-ion collisions

Christoph Hartnack<sup>a</sup>, Helmut Oeschler<sup>b,\*</sup>, Yvonne Leifels<sup>c</sup>, Elena L. Bratkovskaya<sup>d,e</sup>, Jörg Aichelin<sup>a</sup>



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Progress in Particle and Nuclear Physics 56 (2006) 1–103

Progress in  
Particle and  
Nuclear Physics

[www.elsevier.com/locate/ppnp](http://www.elsevier.com/locate/ppnp)

Review

Kaon production in heavy ion reactions at intermediate energies

Christian Fuchs\*

Institut für Theoretische Physik der Universität Tübingen, Auf der Morgenstelle 14,  
D-72076 Tübingen, Germany

Transport-theoretical description of nuclear reactions

O. Buss, T. Gaitanos, K. Gallmeister, H. van Hees, M. Kaskulov, O. Lalakulich, A.B. Larionov<sup>1</sup>, T. Leitner, J. Weil, U. Mosel\*

Institut für Theoretische Physik, Universität Giessen, Germany



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Progress in Particle and Nuclear Physics 53 (2004) 225–237

Progress in  
Particle and  
Nuclear Physics

[www.elsevier.com/locate/pinp](http://www.elsevier.com/locate/pinp)

Review

Strangeness dynamics in relativistic nucleus–nucleus collisions

E.L. Bratkovskaya<sup>a</sup>, M. Bleicher<sup>a</sup>, W. Cassing<sup>b,\*</sup>, M. van Leeuwen<sup>c,d</sup>, M. Reiter<sup>a</sup>, S. Soff<sup>a</sup>, H. Stöcker<sup>a</sup>, H. Weber<sup>a</sup>

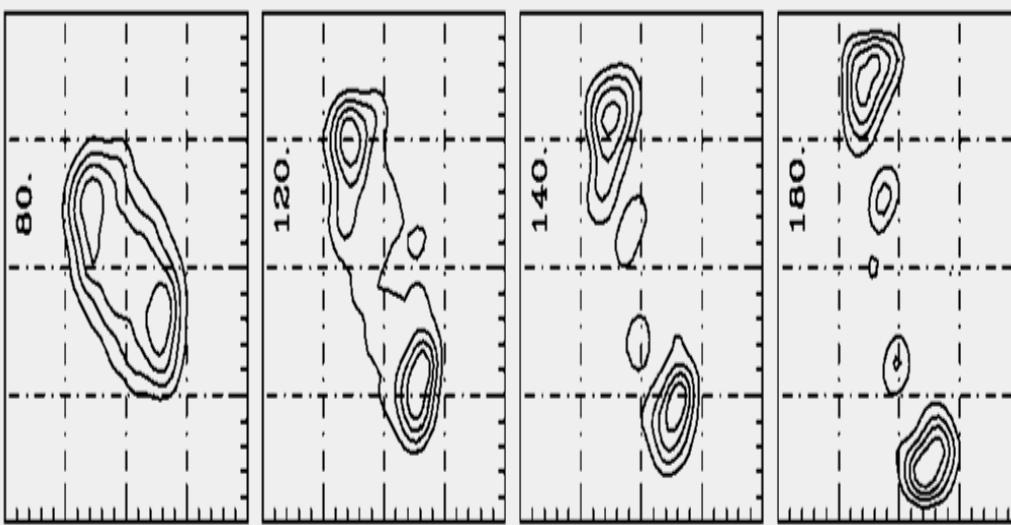
# 1. 费米能区(10A–100A MeV)重离子碰撞中的同位旋效应

Experiments:

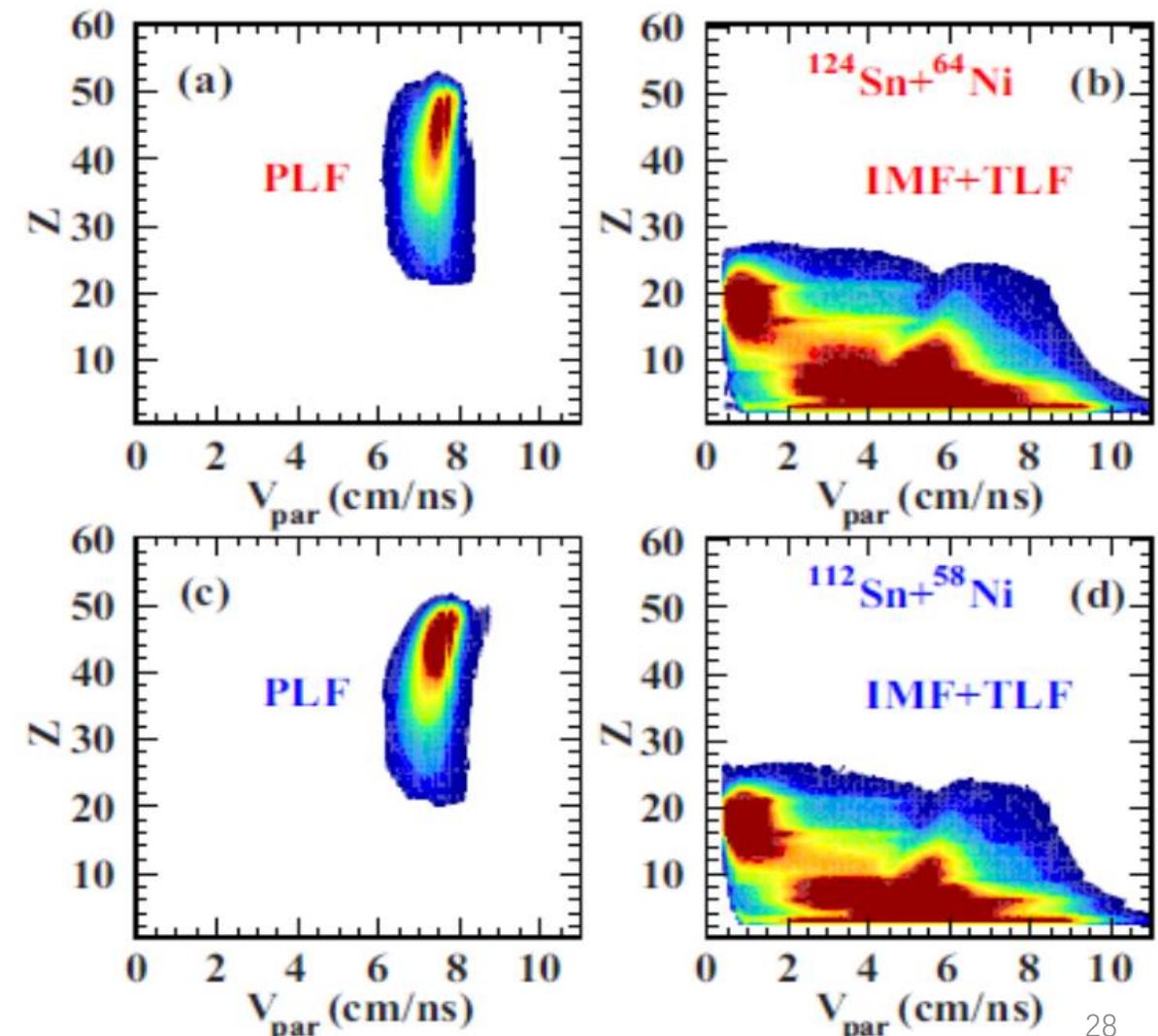
**INDRA (GANIL), CHIMERA (LNS), NSCL (MSU)**

**SSC (HIRFL) ...**

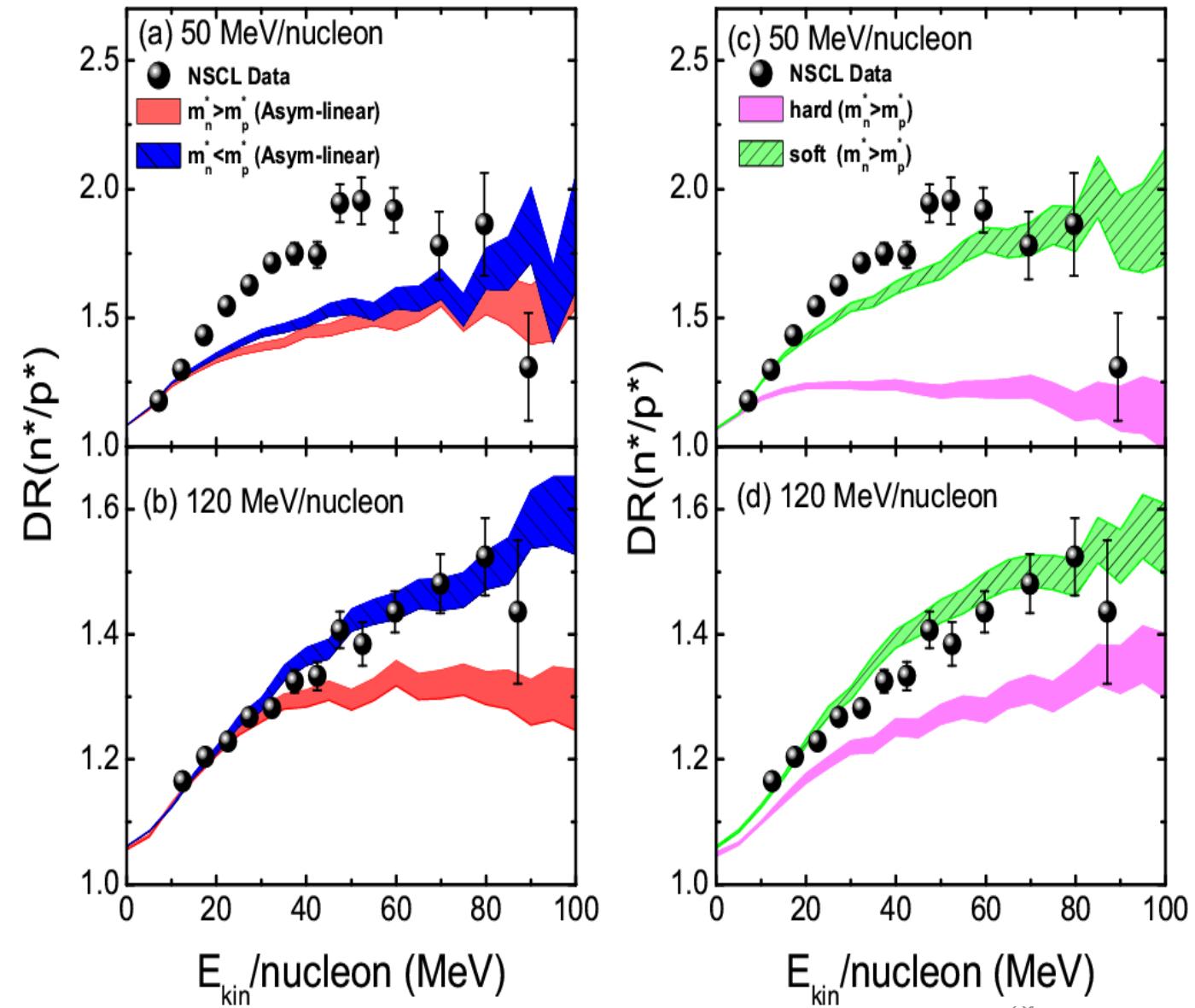
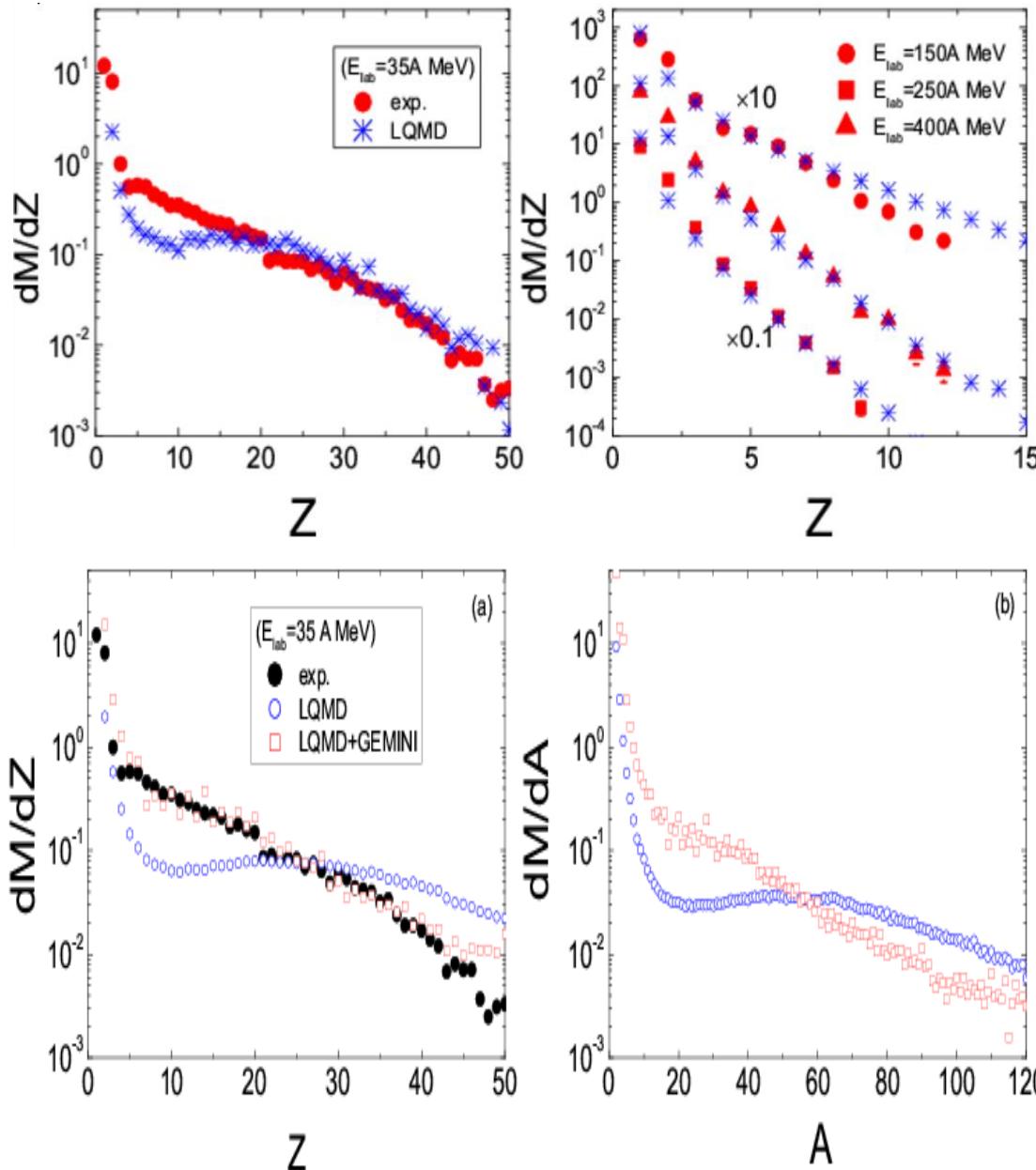
密度演化等高图



P. Russotto *et al.*, PRC 91, 014610 (2015)



# $^{197}\text{Au} + ^{197}\text{Au}$ 碰撞中碎裂分布 (Phys. Rev. C 82, 044615 (2010); 94, 014609 (2016), Chin. Phys. C 41 (2017) 104104 )

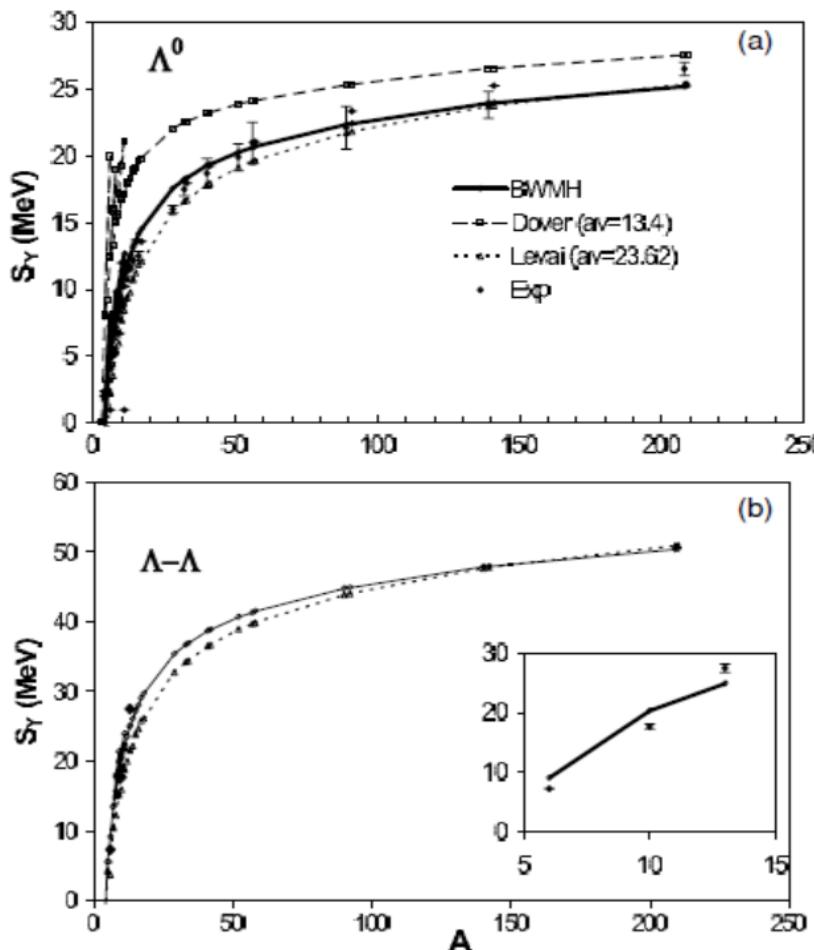


## 2. 超核碎片构造和动力学分析

1) Classical coalescence approach in phase space for nuclides of  $Z>2$  combined with the GEMINI decay code (minimum spanning tree (MST) procedure)

$$|\mathbf{r}_i - \mathbf{r}_j| \leq 3 \text{ fm}, |\mathbf{r}_i - \mathbf{r}_Y| \leq 4.5 \text{ fm}, |\mathbf{p}_i - \mathbf{p}_j| \leq 0.3 \text{ GeV/c}$$

C. Samanta et al, J. Phys. G: Nucl. Part. Phys. 32 (2006) 363



$$\text{Binding energy: } E_B(Z_i, N_i) = \sum_j \sqrt{p_j^2 + m_j^2} - m_j$$

$$+ \frac{1}{2} \sum_{j,k,k \neq j} \int f_j(\mathbf{r}, \mathbf{p}, t) f_k(\mathbf{r}', \mathbf{p}', t)$$

$$v(\mathbf{r}, \mathbf{r}', \mathbf{p}, \mathbf{p}') d\mathbf{r} d\mathbf{r}' d\mathbf{p} d\mathbf{p}'$$

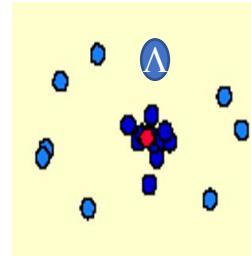
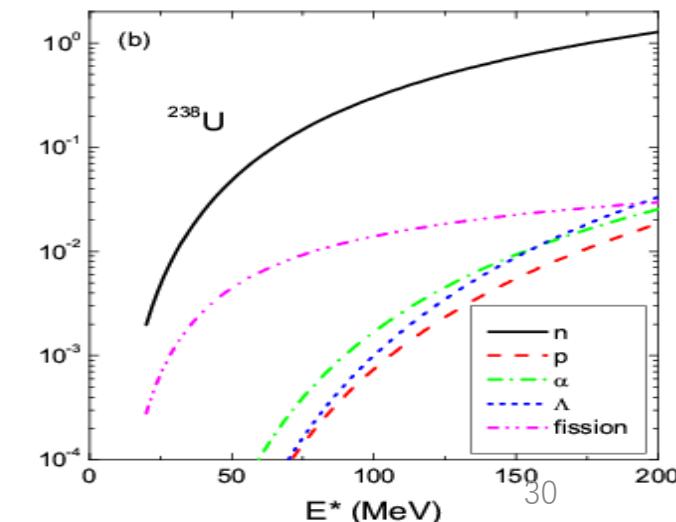
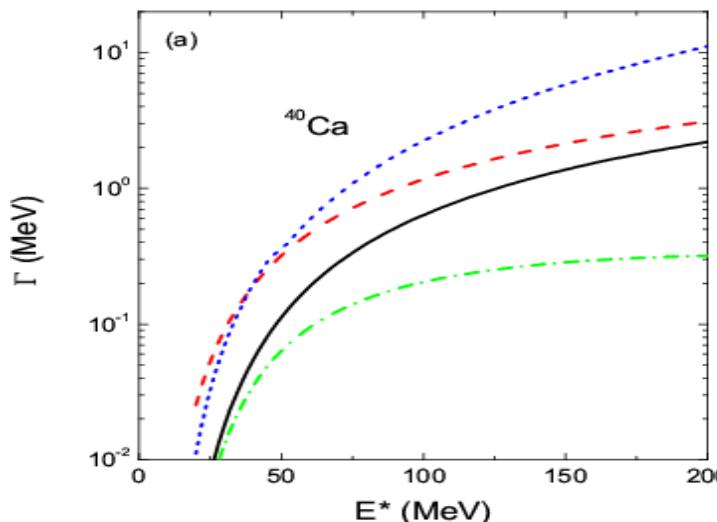
$$+ \frac{1}{6} \sum_{j,k,l} \sum_{k \neq j, k \neq l, j \neq l} \int f_j(\mathbf{r}, \mathbf{p}, t) f_k(\mathbf{r}', \mathbf{p}', t)$$

$$v(\mathbf{r}'', \mathbf{p}'', t) v(\mathbf{r}, \mathbf{r}', \mathbf{r}'', \mathbf{p}, \mathbf{p}', \mathbf{p}'') d\mathbf{r} d\mathbf{r}' d\mathbf{r}'' d\mathbf{p} d\mathbf{p}' d\mathbf{p}'' ,$$

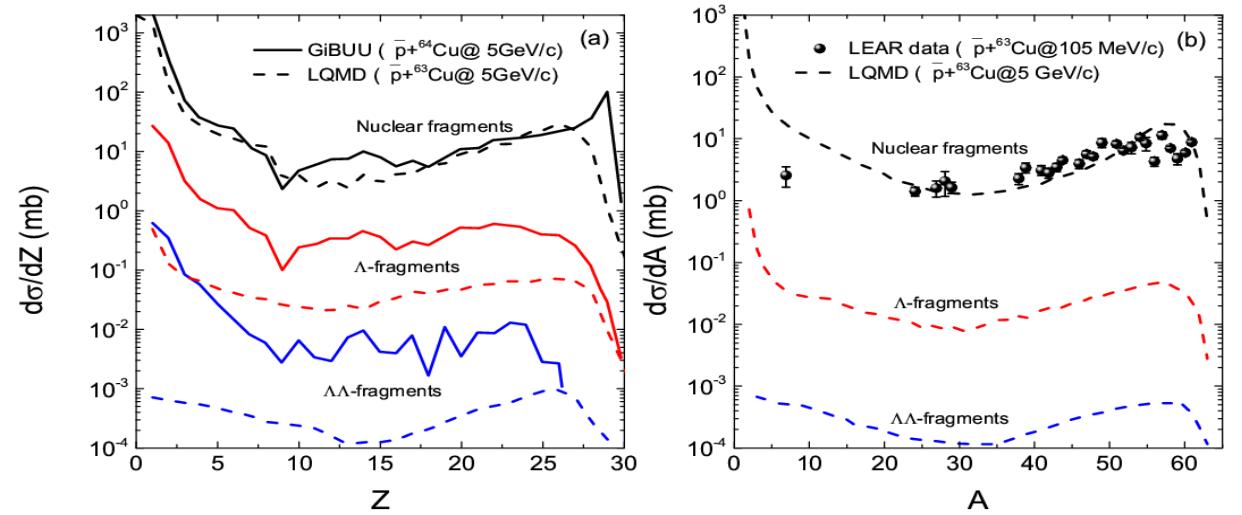
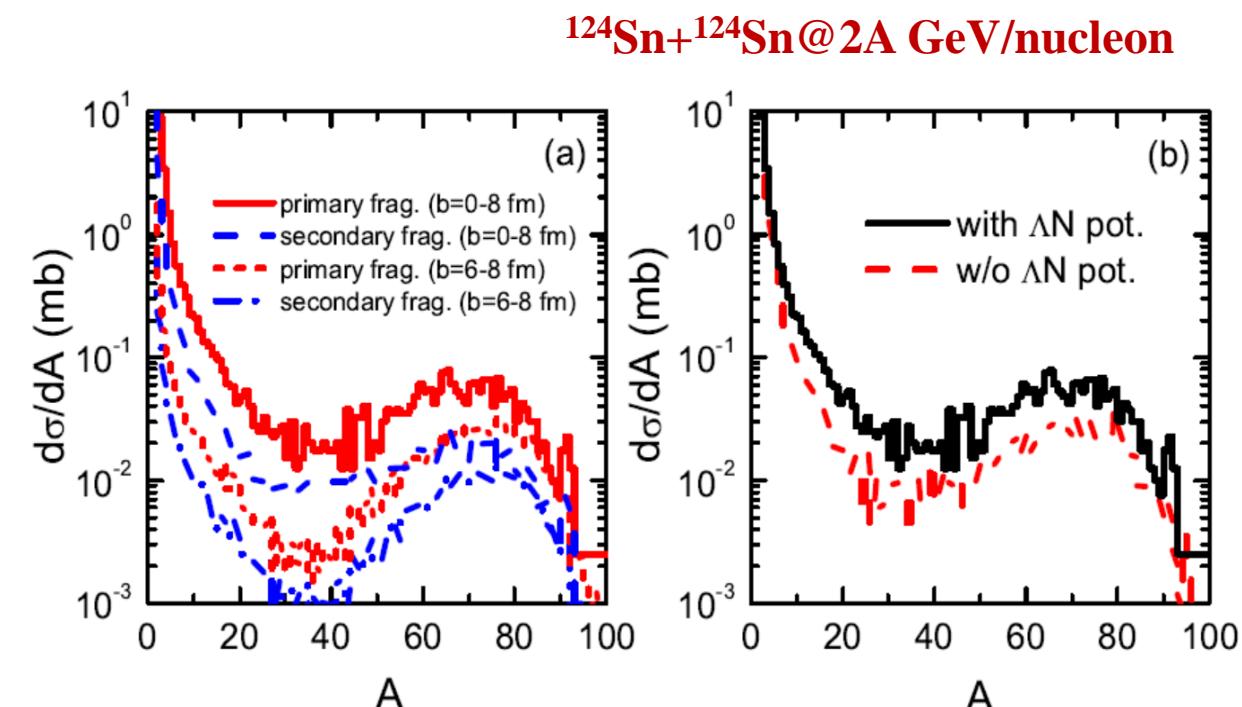
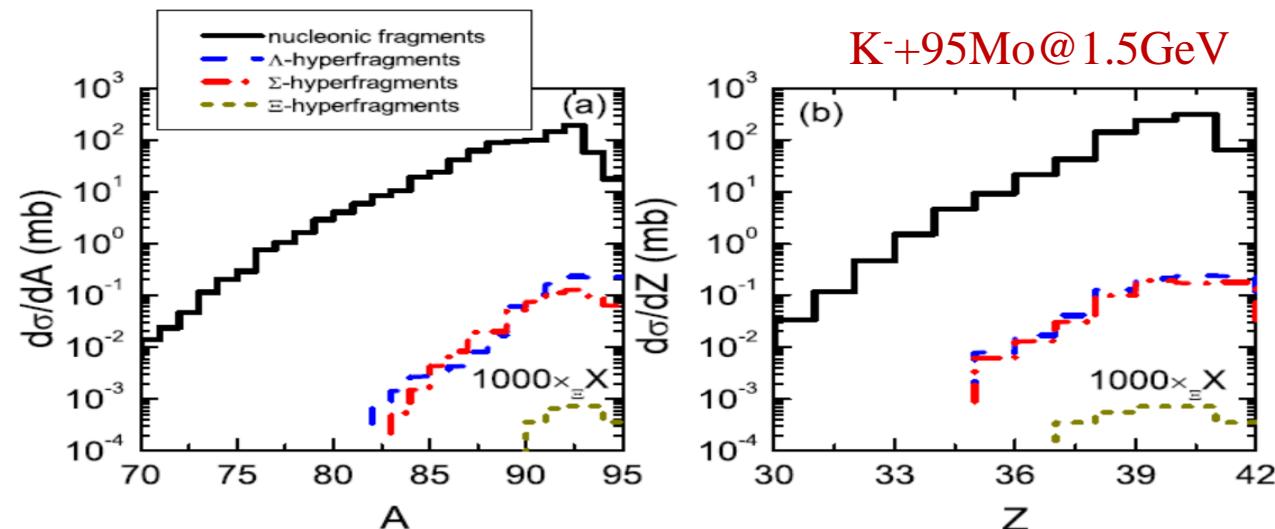
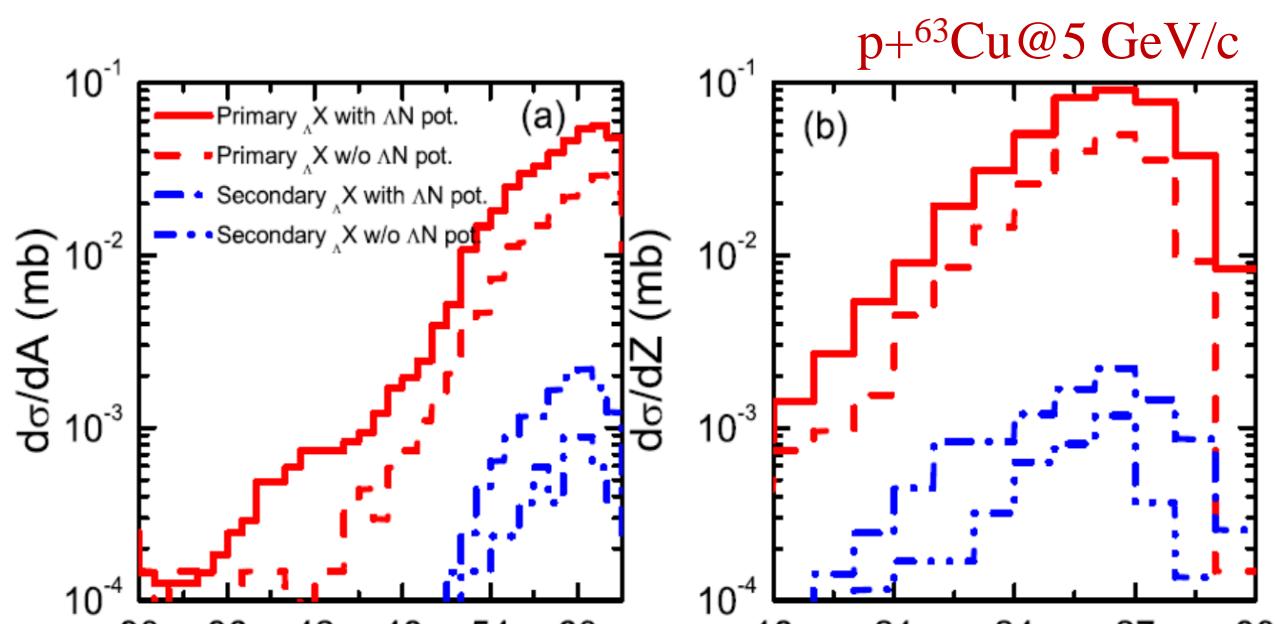
**Excitation energy  $E^*(Z_\nu, N_\nu, nY)$**

$$= E_B(Z_\nu, N_\nu, nY) - E_{LD}(Z_\nu, N_\nu, nY)$$

**The decay of excited hypernucleus is described by the GEMINI code!**



Influence of the statistical decay and hyperon-nucleon potential on the hyperfragment production induced by proton, K- and antiproton ([Physical Review C 101, 064601 \(2020\)](#); [101, 014605 \(2020\)](#); [101, 064601 \(2020\)](#))

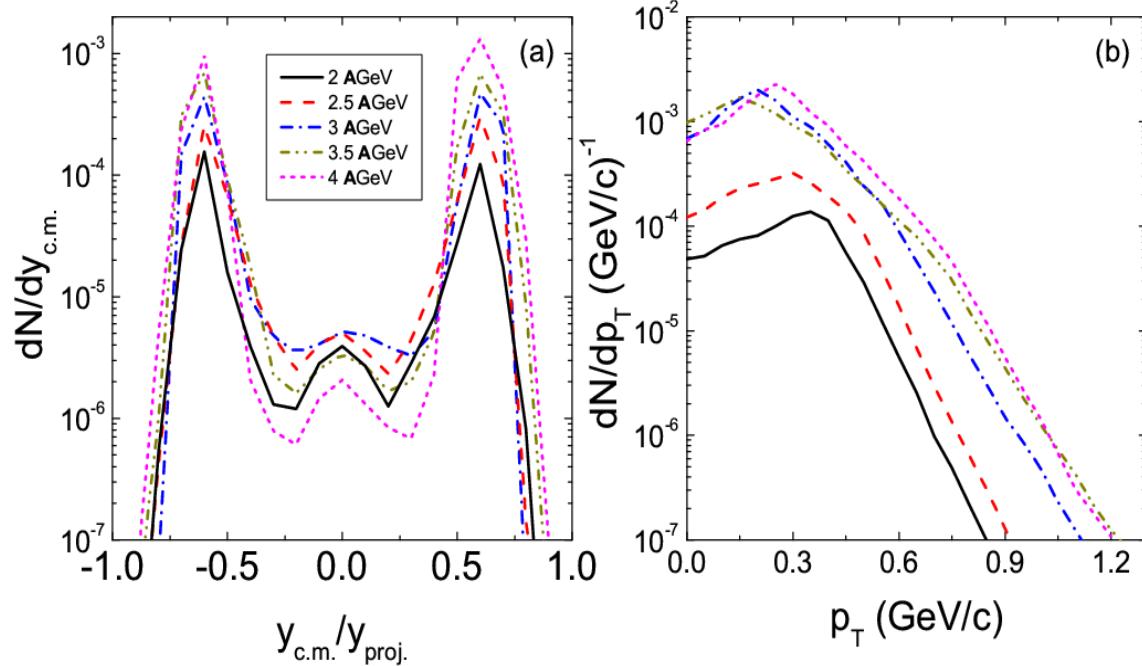


## 2) Wigner density approach for $Z \leq 2$

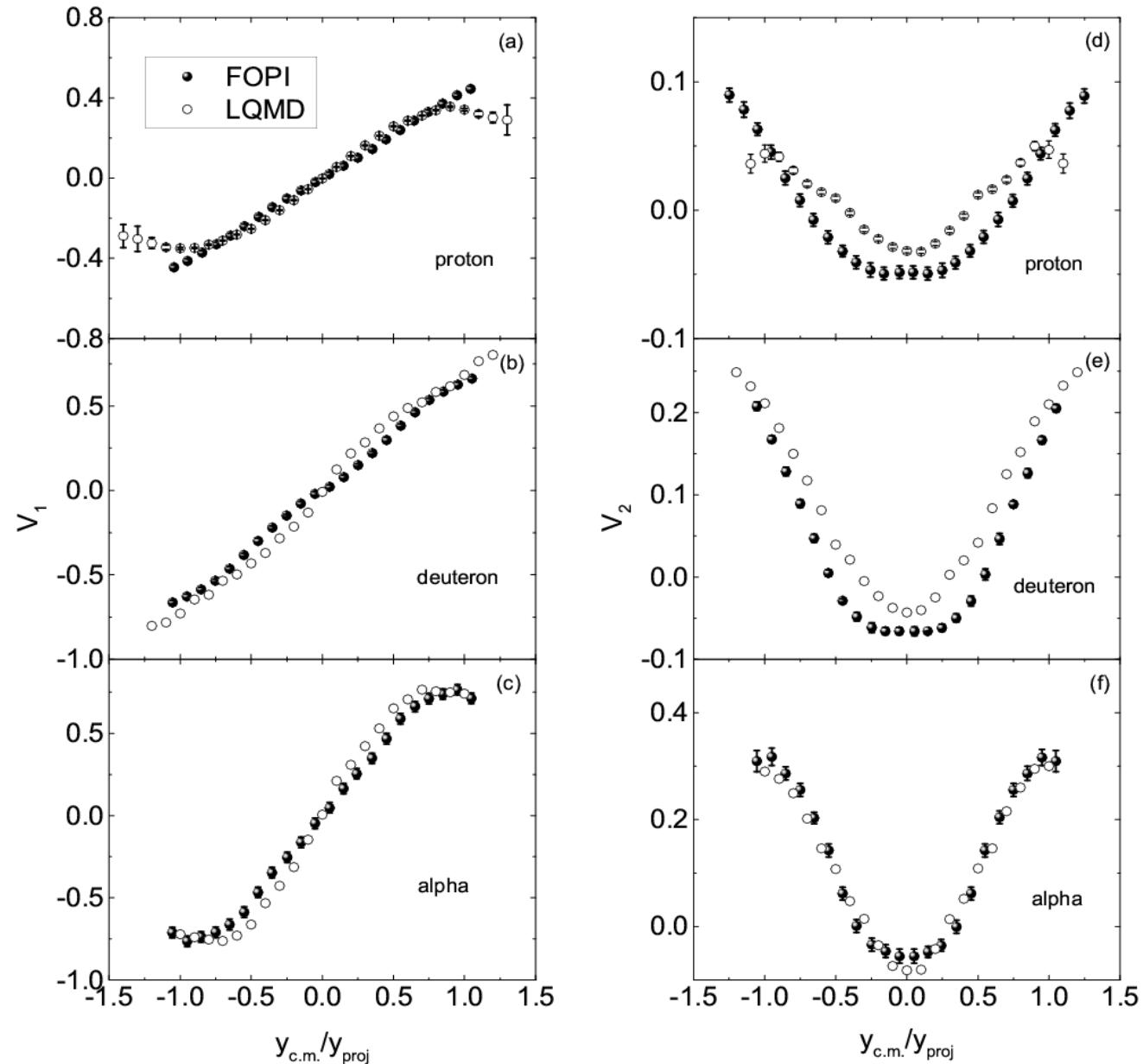
R. Mattiello et al., Phys. Rev. C 55, 1443 (1997)

$$\frac{dN_M}{d^3P} = G_M \binom{A}{M} \binom{M}{Z} \frac{1}{A^M} \int \prod_{i=1}^Z f_p(\mathbf{r}_i, \mathbf{p}_i) \prod_{i=Z+1}^M f_n(\mathbf{r}_i, \mathbf{p}_i) \times \rho^W(\mathbf{r}_{k_1}, \mathbf{p}_{k_1}, \dots, \mathbf{r}_{k_{M-1}}, \mathbf{p}_{k_{M-1}}) \delta(\mathbf{P} - (\mathbf{p}_1 + \dots + \mathbf{p}_M)) d\mathbf{r}_1 d\mathbf{p}_1 \dots d\mathbf{r}_M d\mathbf{p}_M$$

$^3\Lambda$ H via  $^{197}\text{Au} + ^{197}\text{Au}$



Cal: Eur. Phys. J. A, 57 (2021) 18; FOPI data, Nucl. Phys. A 876, 1 (2012)



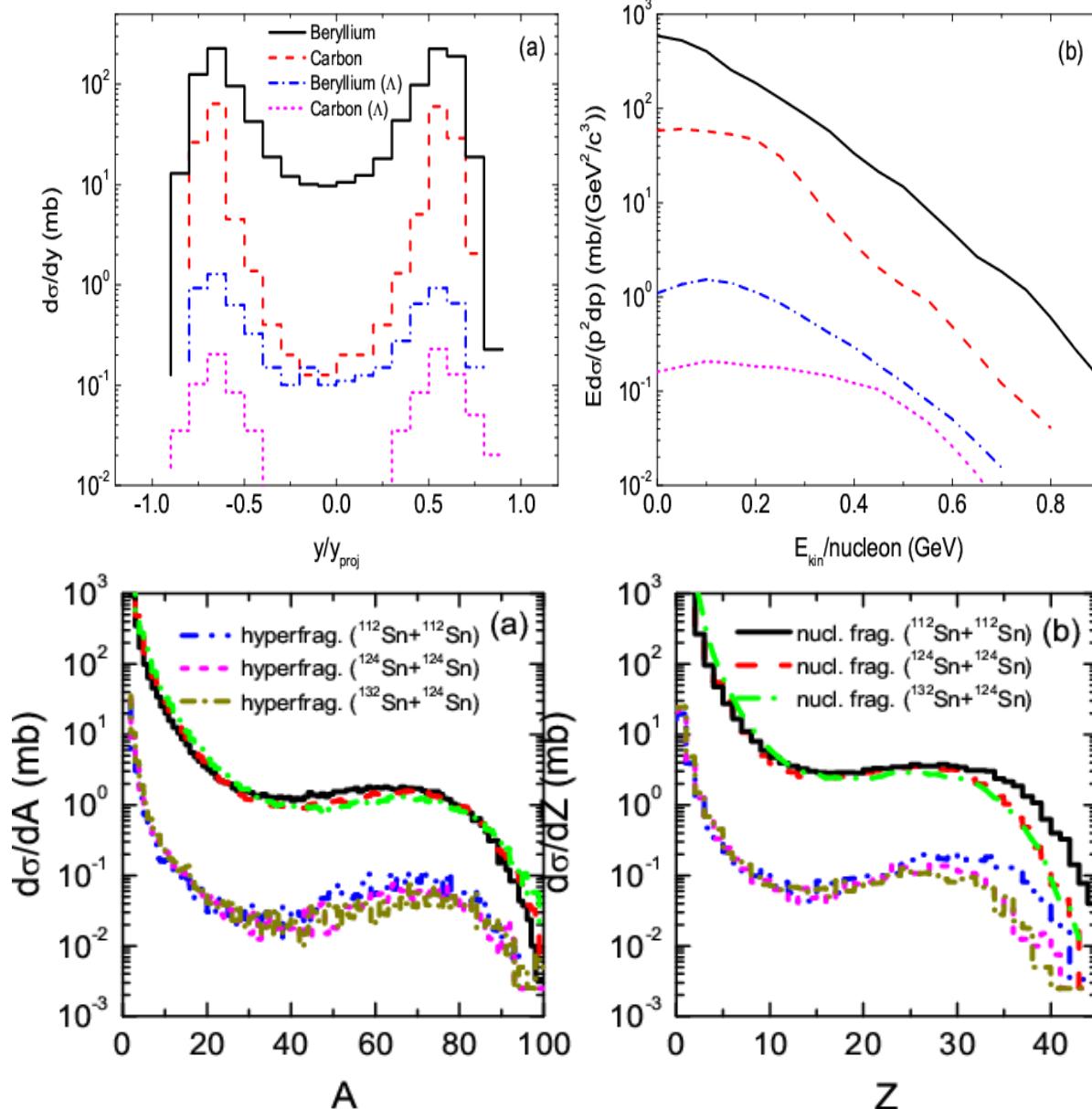
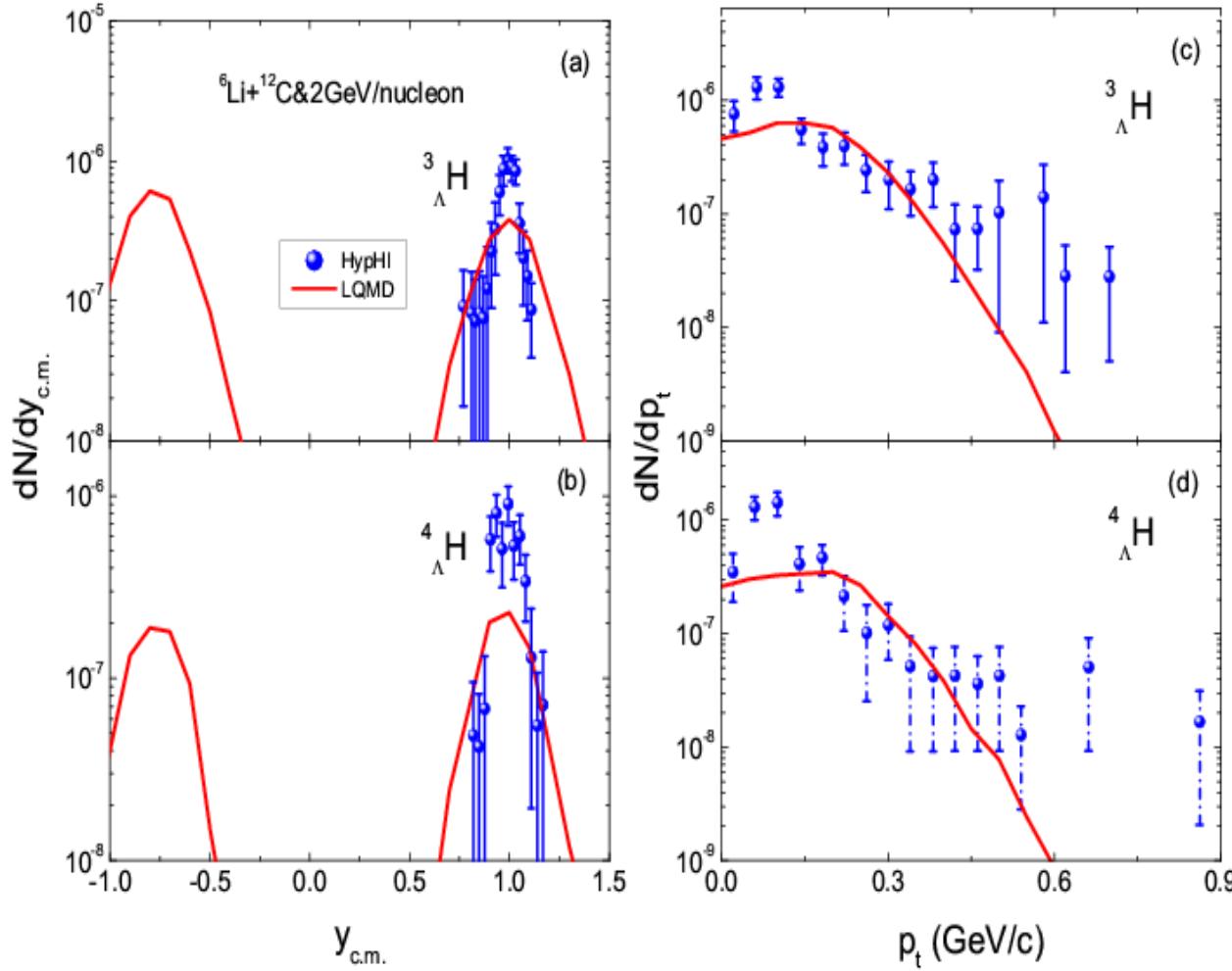
### 3. 重离子碰撞中超核的形成

$^{124}\text{Sn} + ^{124}\text{Sn}$  @ 2A GeV

Z. Q. Feng, Phys. Rev. C 102, 044604 (2020)

Data: C. Rappold et al., (HypHI collaboration)

Phys. Lett. B 747, 129 (2015).



# Multi-strangeness hypernuclide production

H.G. Cheng, Z. Q. Feng, Phys. Lett. B 824 (2022) 136849

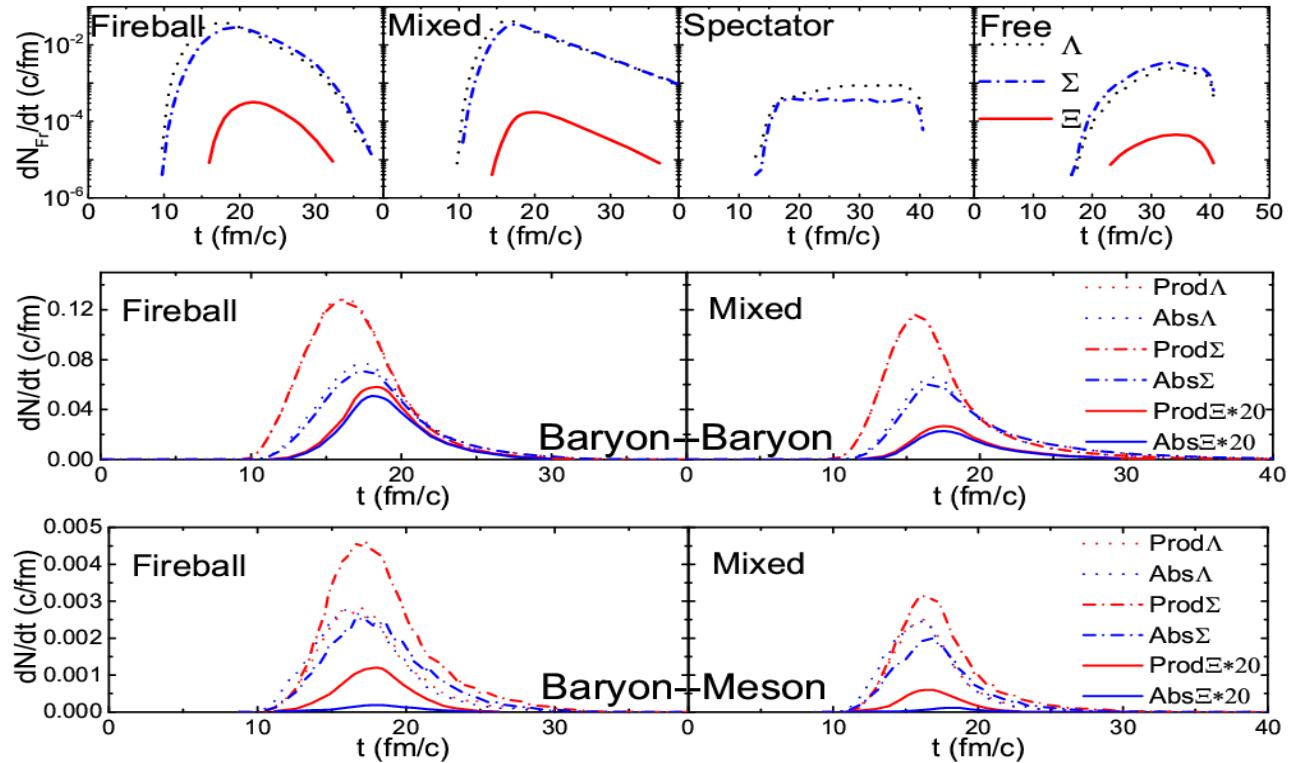
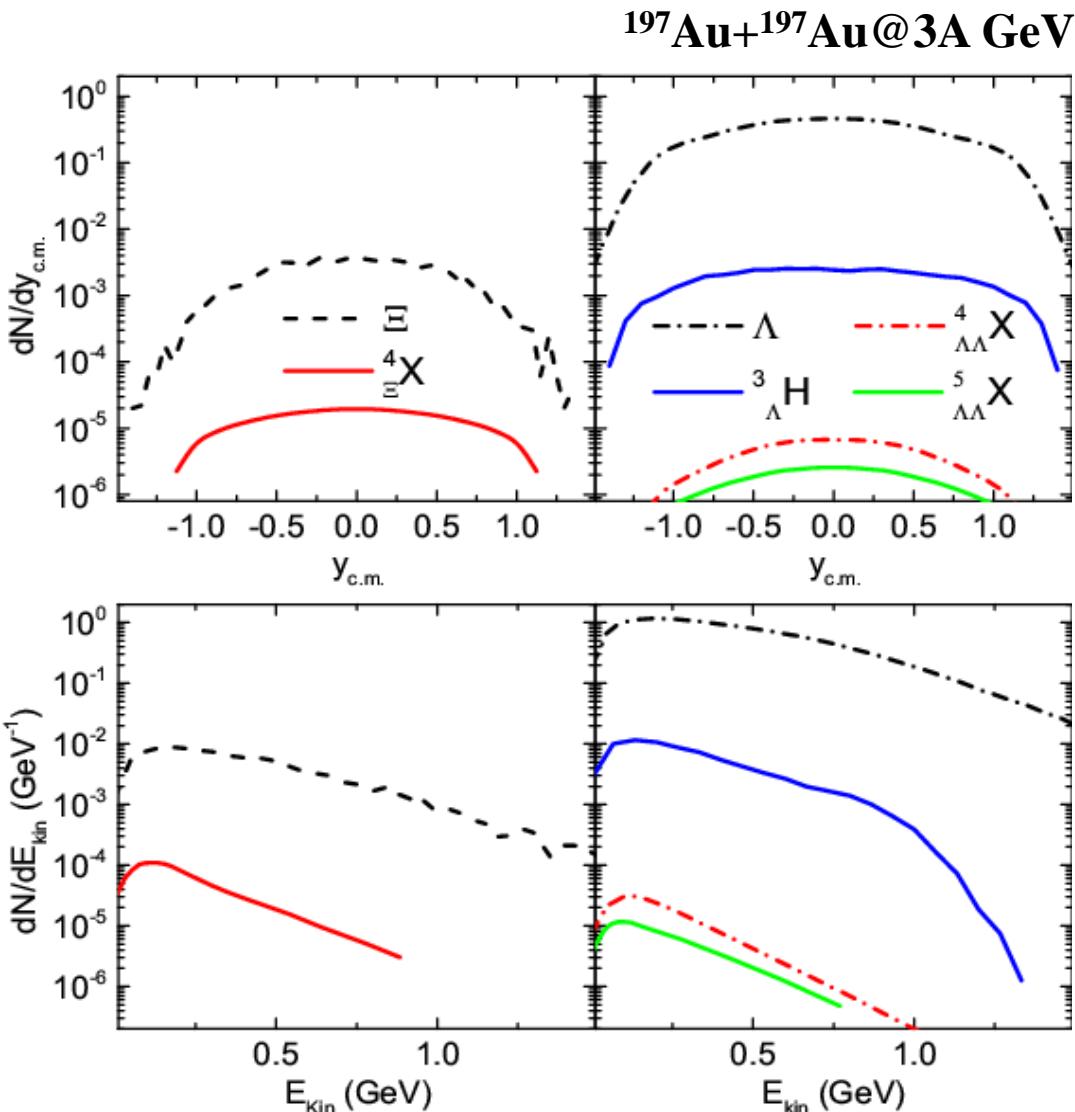
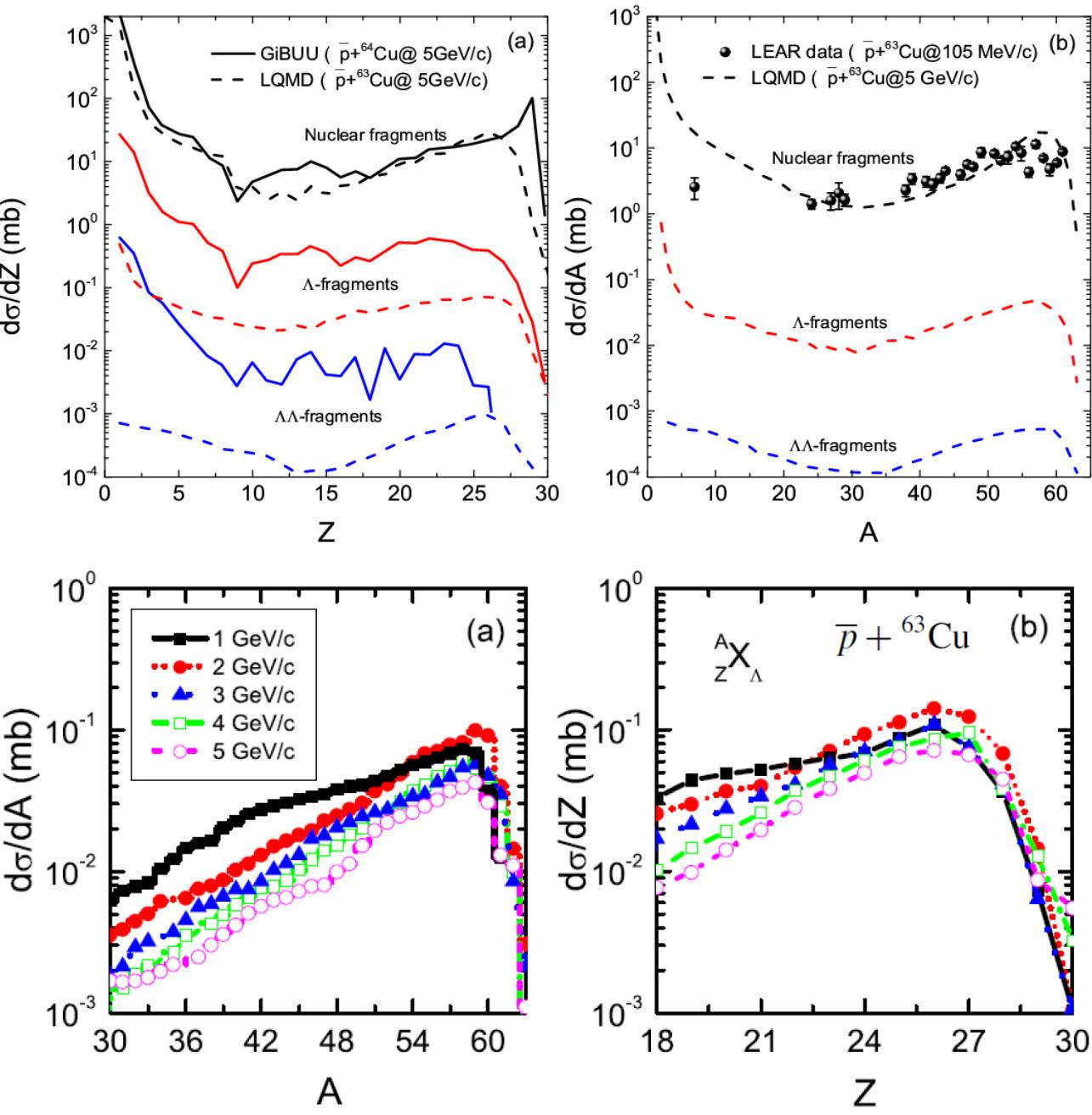
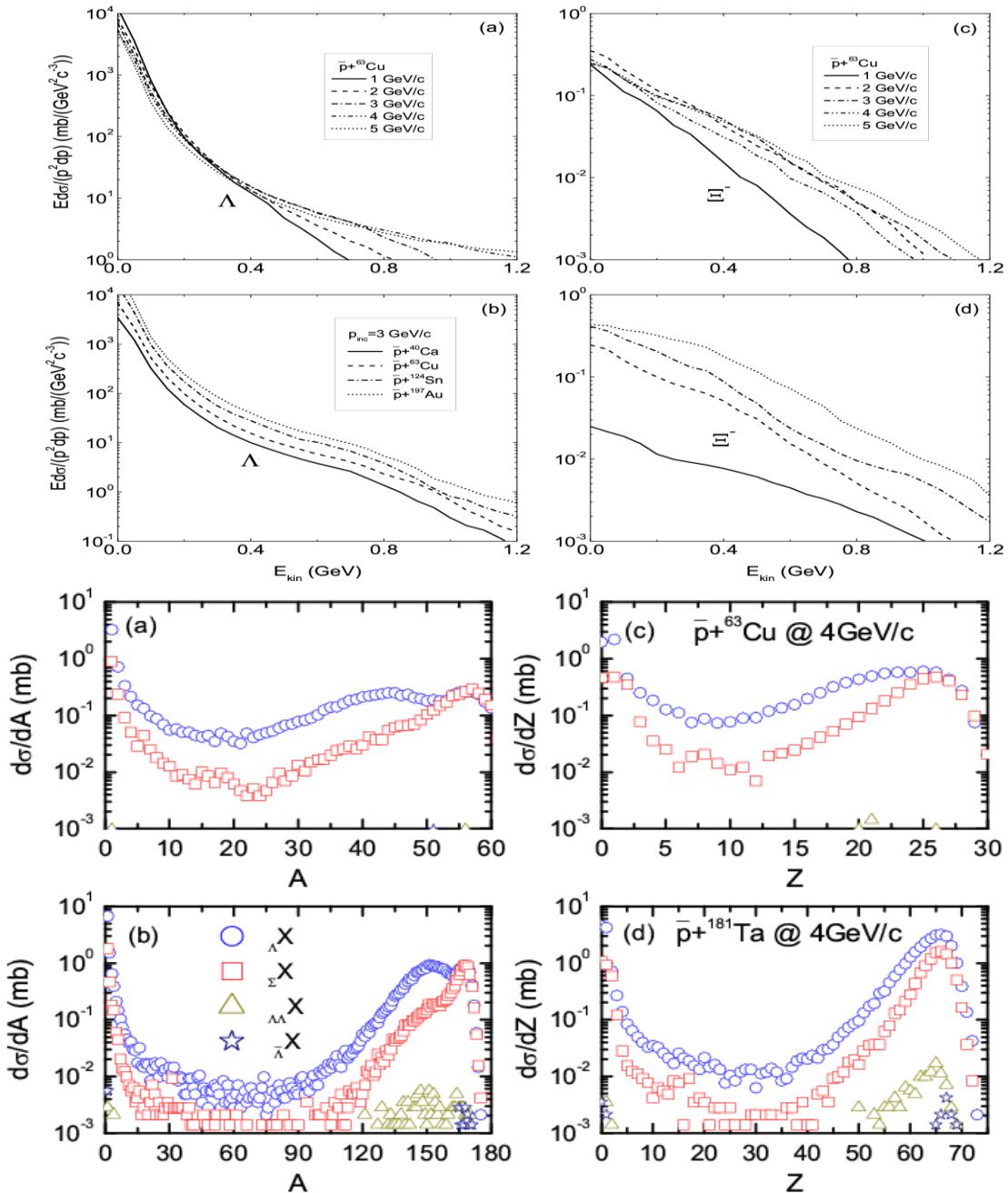


TABLE I. Comparison between cross sections of double lambda hypernuclei calculated with  $r_0 = 3.5$  fm for  $\Lambda$  in  $^{197}\text{Au} + ^{197}\text{Au}$  and  $^{40}\text{Ca} + ^{40}\text{Ca}$  collisions at 3A GeV

Hypernuclei	Cross sections (mb)	
	$^{197}\text{Au} + ^{197}\text{Au}$	$^{40}\text{Ca} + ^{40}\text{Ca}$
$^4\Lambda\Lambda H$	$2.6 \times 10^{-2}$	$1.0 \times 10^{-4}$
$^4\Lambda\Lambda He$	$1.0 \times 10^{-2}$	$\sim 10^{-5}$
$^5\Lambda\Lambda H$	$5.9 \times 10^{-3}$	$\sim 10^{-5}$
$^5\Lambda\Lambda He$	$5.1 \times 10^{-3}$	$\sim 10^{-5}$
$^5\Lambda\Lambda Li$	$1.4 \times 10^{-3}$	$\sim 10^{-6}$
$^6\Lambda\Lambda He$	$2.2 \times 10^{-3}$	$\sim 10^{-6}$
$^7\Lambda\Lambda He$	$6.8 \times 10^{-4}$	$\lesssim 10^{-6}$

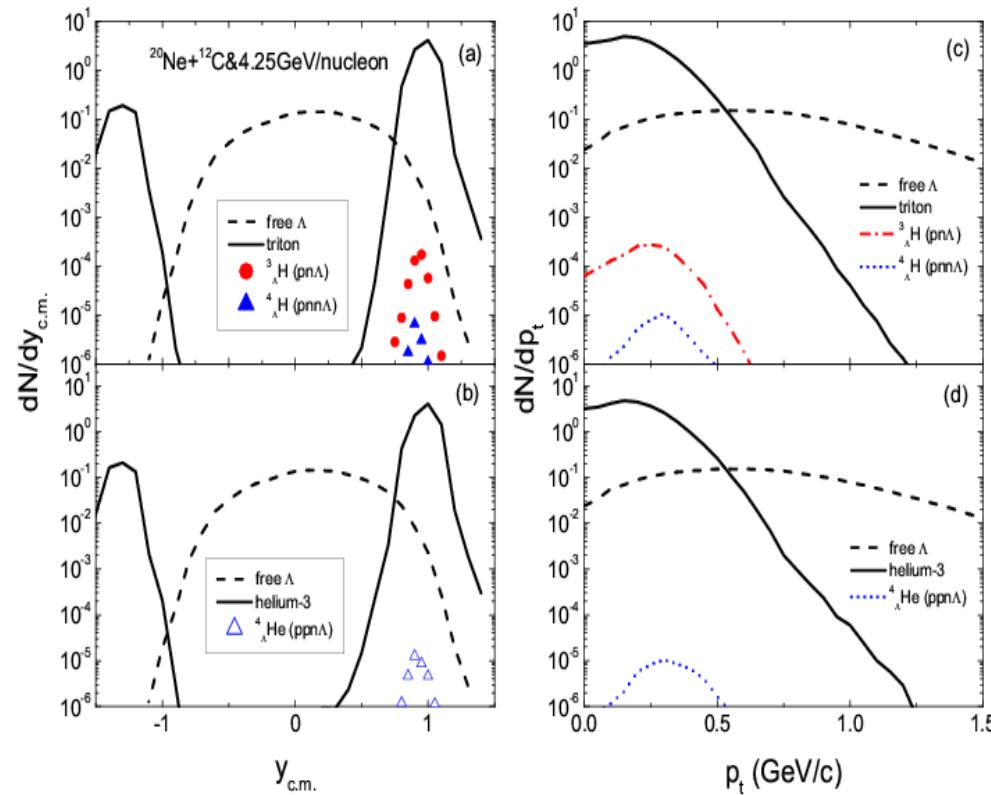
# 4. 反质子引起的核反应中超子的产生和超核形成

Z. Q. Feng, Phys. Rev. C 89, 044617 (2014); Phys. Rev. C 93, 041601(R) (2016); Phys. Rev. C 101, 064601 (2020)



## 四、总结

- 碰撞区域产生的超子被“旁观者”(旁观核子)俘获后形成超核，轻质量超核可以在类弹、类靶和中心快度区域产生，重质量超核只能在类弹(靶)区域产生
- 重离子碰撞可以产生极端丰中子/丰质子超核、多奇异性超核。入射能量4.25GeV/核子 $^{20}\text{Ne}+^{12}\text{C}$ 反应可以在HIAF上做超核研究测试实验。
- 问题：核子-核子碰撞中涉及超子的三体和四体碰撞直接产生的超核(轻质量、高动量、奇异性，如 $\text{nn}\Lambda$ ,  $\text{nn}\Lambda\Lambda$ ,  $_{\Lambda}^5\text{H}$ ,  $_{\Lambda}^6\text{H}$ )还没有考虑！



谢谢大家！

