



Compton Polarimeter for CEPC

Shanhong CHEN

On behalf of the BEMS team

- Motivation of CEPC Z-pole polarized beam program
- Compton Polarimeter for CEPC
- The measurements of beam transverse polarization(Z pole)
- The toy MC simulations
- The fit results of transverse polarization
- Discussion & Summary

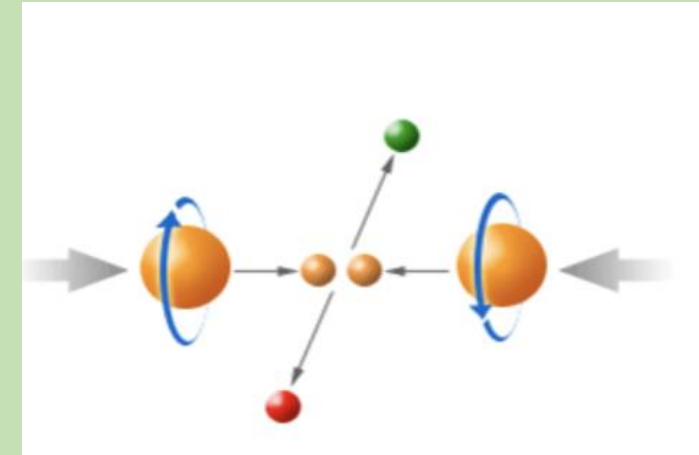
Motivation of CEPC Z-pole polarized beam program

- Transversely polarized beams in the ARC

- Beam energy calibration via the resonant depolarization technique(Accuracy 10^{-6})
- Essential for precision measurements of Z and W properties
- At least 5% ~ 10% transverse polarization, for both e⁺ and e⁻ beams

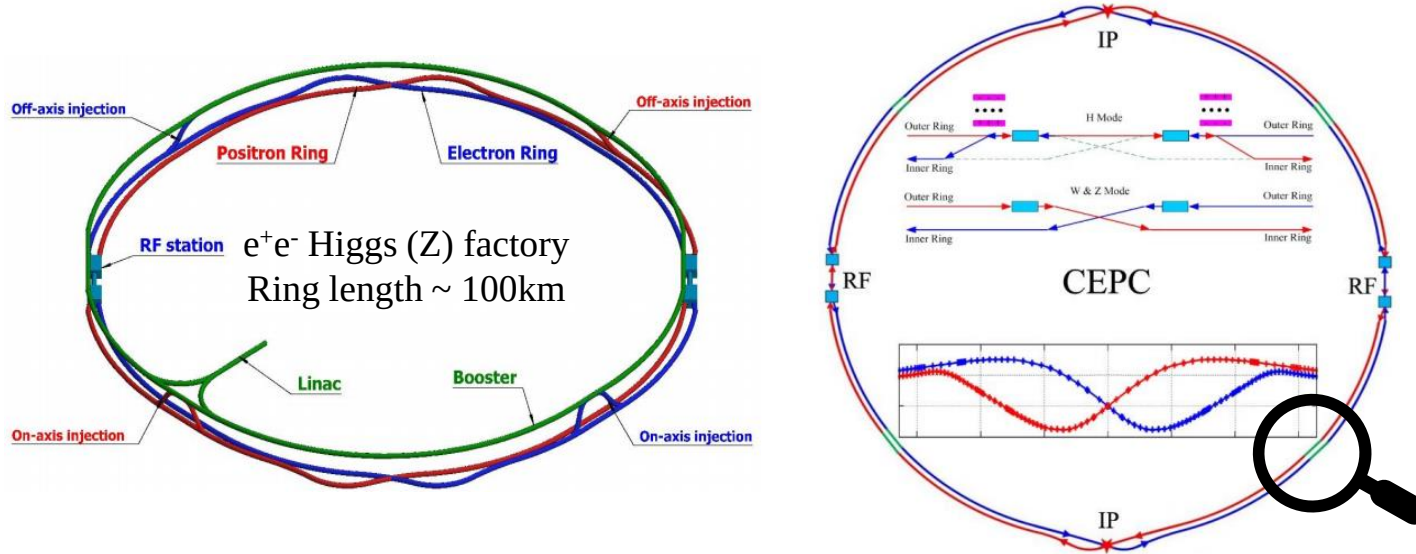
- Longitudinally polarized beams at IPs

- Beneficial to colliding beam physics programs at Z, W and Higgs(weak interactions & spin structure of particle)
- ~50% or more longitudinal polarization is desired, for one beam, or both beams



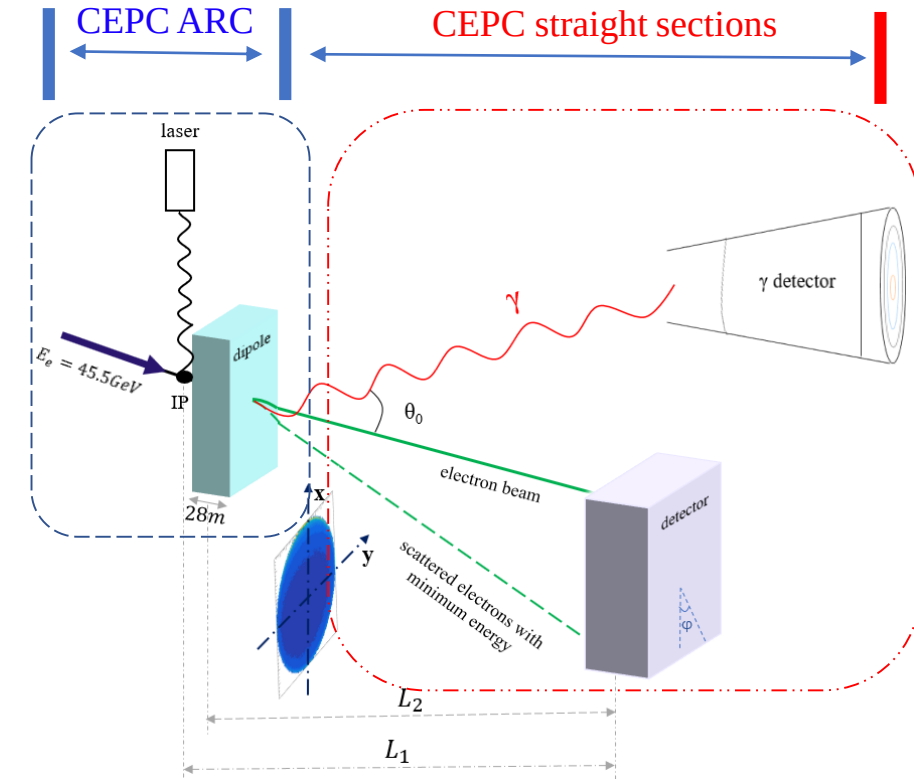
Compton Polarimeter at CEPC(Z pole)

- Discuss the system layout(for transverse polarization)



- **Prospects for the Compton polarimeter system:**
- Arrange our device in straight sections on the ring.(ensuring not affect the normal operation of beam on the ring)
- The last dipole in the ARC can be used as our bending magnet for our Compton polarimeter.
- **Note that:** the layout above is only aimed at measuring transverse polarization, the longitudinal polarimeter must located between two spin rotators(will discuss in the future)

△ Compton Polarimeter



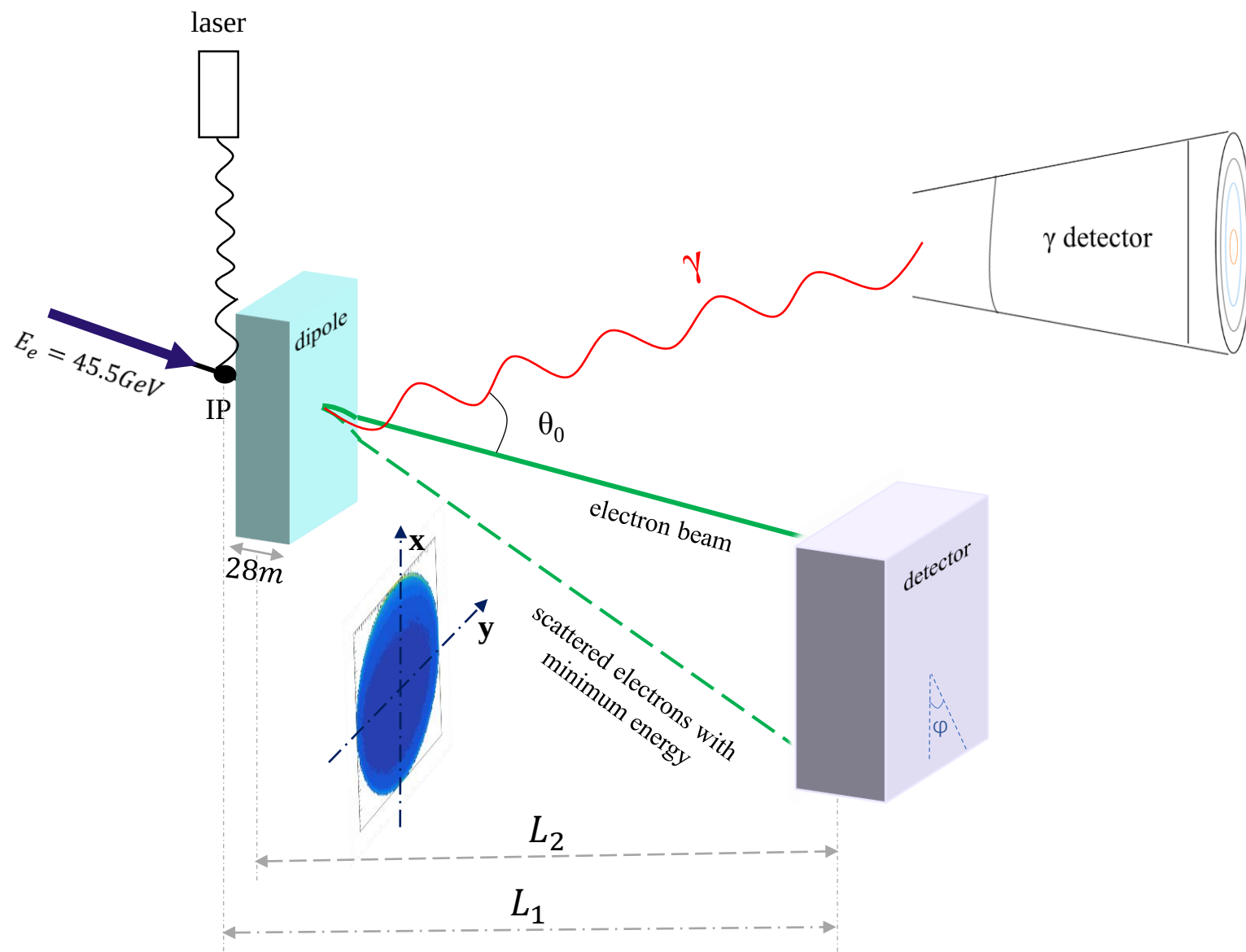
Dipole parameter(CDR): $B = 0.01867T$ $l = 28m$

Draft distance(Our Preliminary design):

- $L_1 = 60m$; $L_2 = 40m$
- Meaning that after dipole the scattered particle will drift $40m - 28/2 = 16m$

Compton Polarimeter at CEPC(Z pole)

Parameters	
E	45.5 [GeV]
Laser	1.24 [eV]
Dipole	$B = 0.01867[T]$
	$l = 28 [m]$
	$\theta_0 = 3.4467[mrad]$
Drift distance	$L_1 = 60 [m]$
	$L_2 = 40 [m]$
Beam angle (at IP of laser and electron beam)	$\beta_x = 121m$
	$\beta_y = 15m$
	$\epsilon_x = 0.18nm$
	$\epsilon_y = 0.004nm$
	$\sigma'_x = 1.2197 [\mu rad]$
	$\sigma'_y = 0.5164 [\mu rad]$



The distribution of scattered electrons

● A toy MC

$$X_e = \sigma'_x L_1 + \frac{L_1}{\gamma} \sqrt{u(\kappa - u)} \cos \varphi + u \theta_0 L_2$$

$$Y_e = \sigma'_y L_1 + \frac{L_1}{\gamma} \sqrt{u(\kappa - u)} \sin \varphi$$

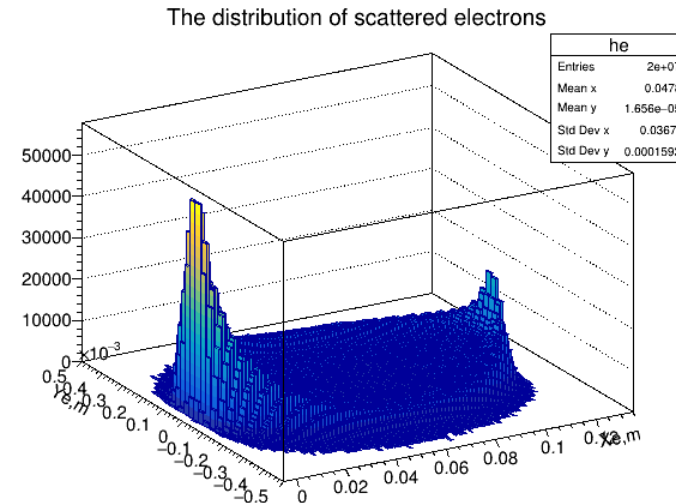
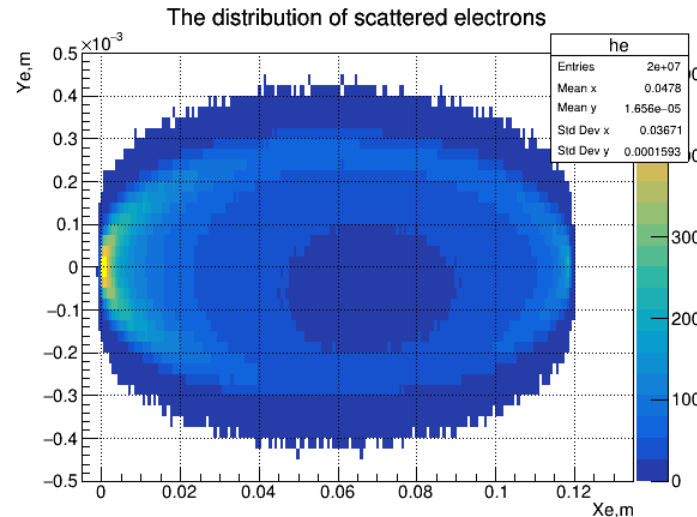
Where σ'_x / σ'_y are horizontal and vertical electron angle ;

L_1 : Distance between IP and detector;

L_2 : Distance between Dipole and detector;

$u = \frac{\omega}{\varepsilon}$ (the ratio of energy of scattered photons and scattered electrons);

$u \in [0, \kappa]$



The detector resolution has been considered in MC simulation

$$X \sim \text{smear}(X_e, 202.07 \mu\text{m})$$

$$Y \sim \text{smear}(Y_e, 7.2168 \mu\text{m})$$

The active dimension of detector : $140\text{mm} \times 1\text{mm}$

The pixel size : $700\mu\text{m} \times 25\mu\text{m}$

The pixel number : 200×40

The resolution of the detector ($\sqrt{12}$) : $202.07\mu\text{m} \times 7.2168\mu\text{m}$

Compton Scattering cross section

● The fit function-----Analyzing power $\Pi(Xe)$

- The distribution of scattered electrons depends on electron polarization(ζ) and laser polarization(ξ)

$$d\sigma_0 = \frac{r_e^2}{(1+u)^3 \sqrt{1-x^2-y^2}} \left(1 + (1+u)^2 - 4 \frac{u}{\kappa} (1+u) \right) dx dy$$

$$d\sigma_{\parallel} = \frac{\xi_{\parallel} \zeta_{\parallel} r_e^2}{\kappa (1+u)^3 \sqrt{1-x^2-y^2}} u(u+2) \left(1 - 2 \frac{u}{\kappa} \right) dx dy$$

$$d\sigma_{\perp} = - \frac{\xi_{\perp} \zeta_{\perp} r_e^2}{(1+u)^3 \sqrt{1-x^2-y^2}} u y dx dy$$

Ref: Nickolai Muchnoi, "FCC-ee polarimeter," arXiv e-prints, arXiv:1803.09595 (2018).

$$x = \frac{(X_e - \sigma'_x L_1) - \frac{\kappa}{2} \theta_0 L_2}{\frac{\kappa}{2} \sqrt{\left(\frac{L_1}{\gamma}\right)^2 + (\theta_0 L_2)^2}} \quad y = \frac{Y_e - \sigma'_y L_1}{\frac{L_1}{\gamma} \frac{\kappa}{2}} \quad < u > = \frac{\theta_0 L_2}{\sqrt{\left(\frac{L_1}{\gamma}\right)^2 + (\theta_0 L_2)^2}} \frac{\kappa}{2} x + \frac{\kappa}{2}$$

Note that: In our simulation, we consider laser beam has no linear polarization and electron beam has no longitudinal polarization

- The average y is given by:

$$< y > = \frac{\int y \frac{d\sigma}{dx dy} dy}{\int \frac{d\sigma}{dx dy} dy}$$

- Use it to obtain the transverse polarization:

$$\Delta < y > |_{Xe} = \frac{< y > |_{left} - < y > |_{right}}{2} = \zeta_{\perp} \Pi(Xe)$$

- Analyzing power Π :

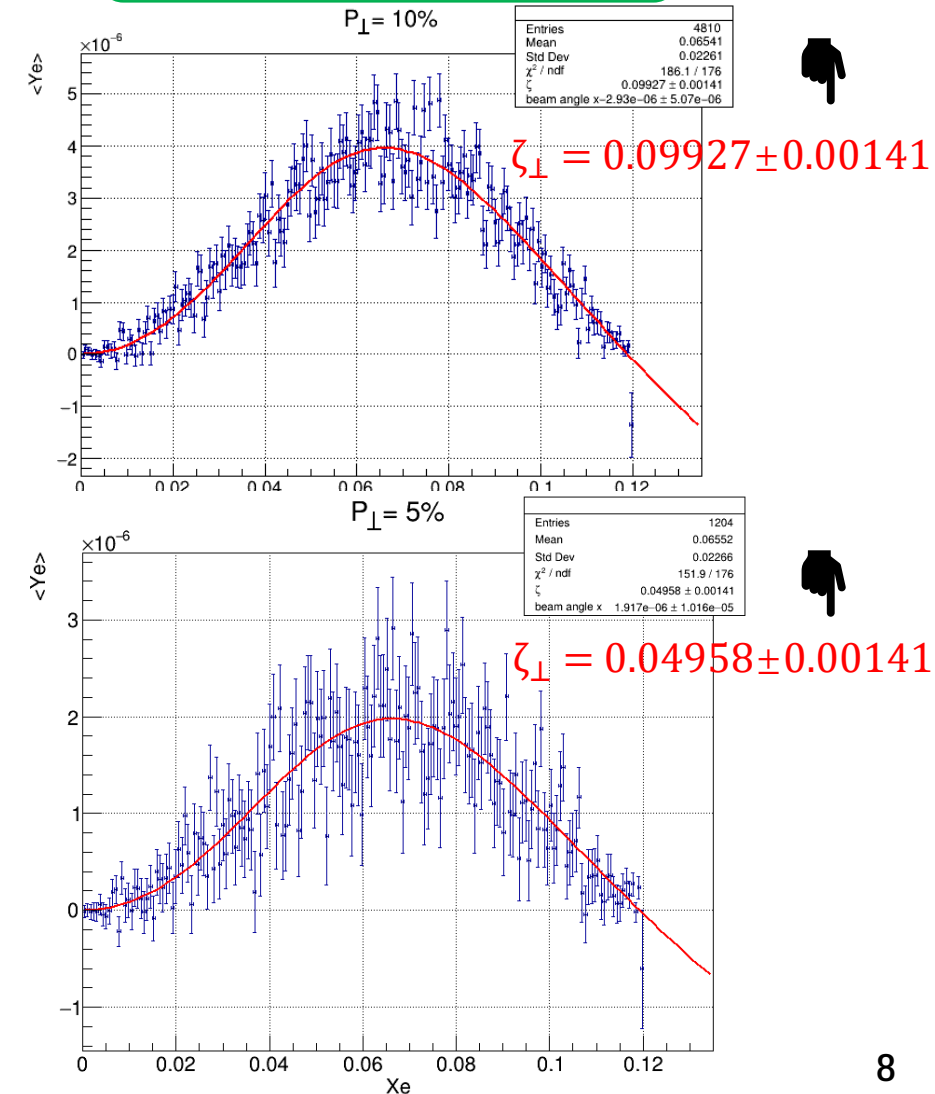
the Analyzing Power of the polarimeter is its value for a 100% polarization.

Polarization Measurement via Compton Polarimeter

● The method

- ▶ 1. Obtain the 2D distribution scattered electrons
- ▶ 2. Projection
 - Projection onto the X-axis and add up the Y coordinates of each X bin
- ▶ 3. we fit the asymmetry:
 - $\Delta \langle y \rangle |_{Xe} = \zeta_{\perp} \Pi(Xe)$
 - Where the transverse polarization $\zeta_{\perp}$ is fit parameter

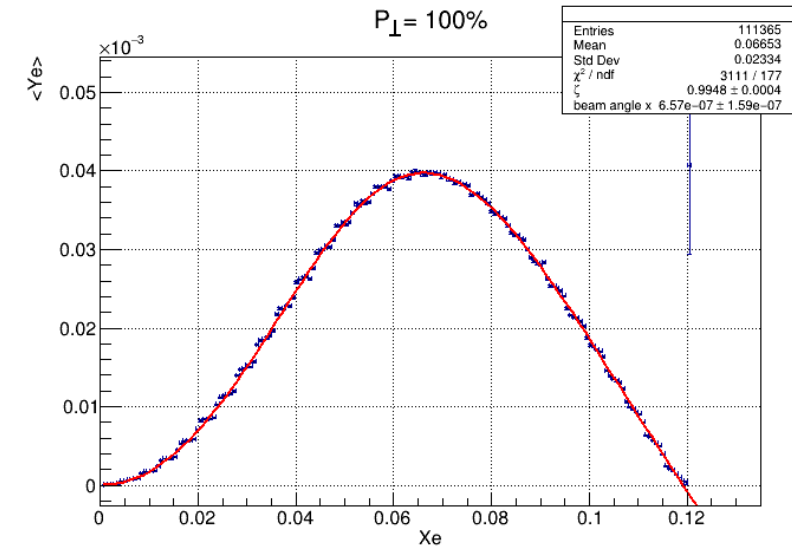
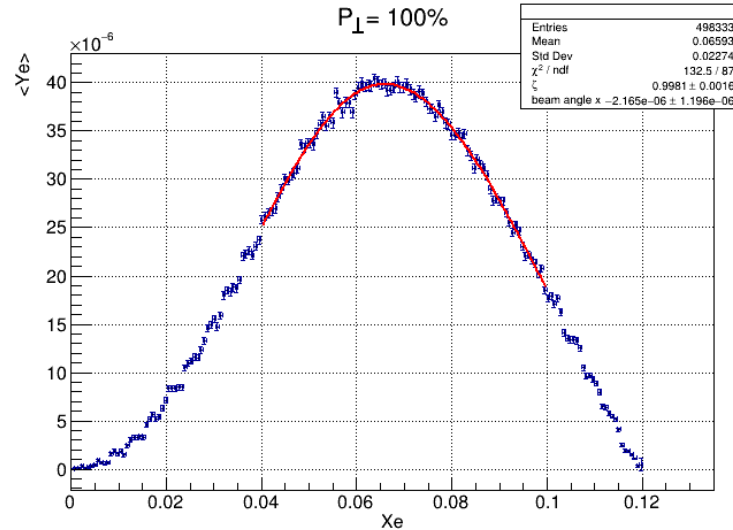
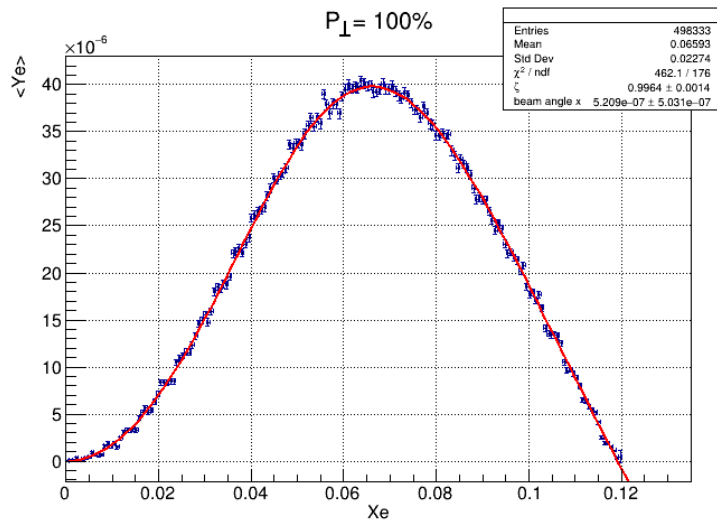
The fit result of
Polarization = 10% and 5%



Polarization Measurement via Compton Polarimeter

● About the statistics error and fit range

	Entries	Fit range	Fit result
P = 100%	Entries = 2E7	[0.00 , 0.12]	0.9964 ± 0.0014
	Entries = 2E7	[0.04 , 0.10]	0.9981 ± 0.0016
	Entries = 2E8	[0.00 , 0.12]	0.9948 ± 0.0004



- For statistic error, 0.15% for 2E7, 0.04% for 2E8.



The deviation between the fitting result and the true value: ①The value of the $\langle Y \rangle$ is obtain by each bin; ②smear due to detector resolution

Polarization Measurement via Compton Polarimeter

● The luminosity for continuous lasers

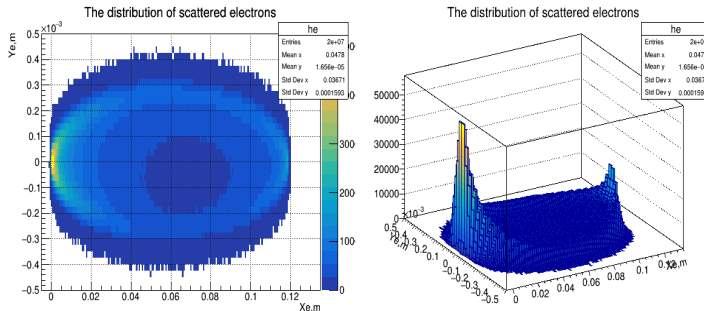
$$\mathfrak{L} = \frac{1 + \cos\theta_0}{\sqrt{2\pi}} \frac{I_e}{e} \frac{P_L \lambda}{hc^2} \frac{1}{\sqrt{\sigma_e^2 + \sigma_\gamma^2}} \frac{1}{\sin\theta_0} \approx \frac{2}{\sqrt{2\pi}} \frac{I_e}{e} \frac{P_L \lambda}{hc^2} \frac{1}{\sqrt{\sigma_e^2 + \sigma_\gamma^2}} \frac{1}{\theta_0}$$

Par.	Unit	Physics meaning	CEPC(Z pole)
I_e	A=C/s	mean electron current	461mA
P_L	W=J/s	power of the laser	1.0W
λ	m	Wavelength of laser	1.002 μ m(1.24eV)
σ_e/σ_γ	m	Rms beam size	$\sigma_\gamma = 160\mu m$
h	J/s	Planck constant	$6.626 \times 10^{-34} J \cdot s$
c	m/s	Light speed	$3 \times 10^8 m \cdot s$
$\mathfrak{L}$	$cm^{-2}s^{-1}$	luminosity	$2.56 \times 10^{30} cm^{-2} \cdot s^{-1}$
σ	barn	Cross section	393.5mb
Max.rate	s^{-1}	Compton scattering event rate	$1.0088 \times 10^6 s^{-1}$

- A 5%/~10% polarization level is expected to be measured in a 20s with an relative accuracy of about 1.4%.

Polarization Measurement via Compton Polarimeter

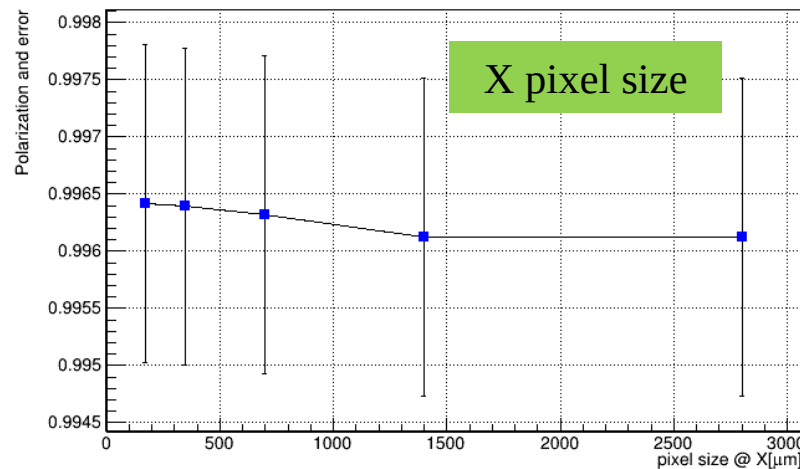
● About the active detector dimension



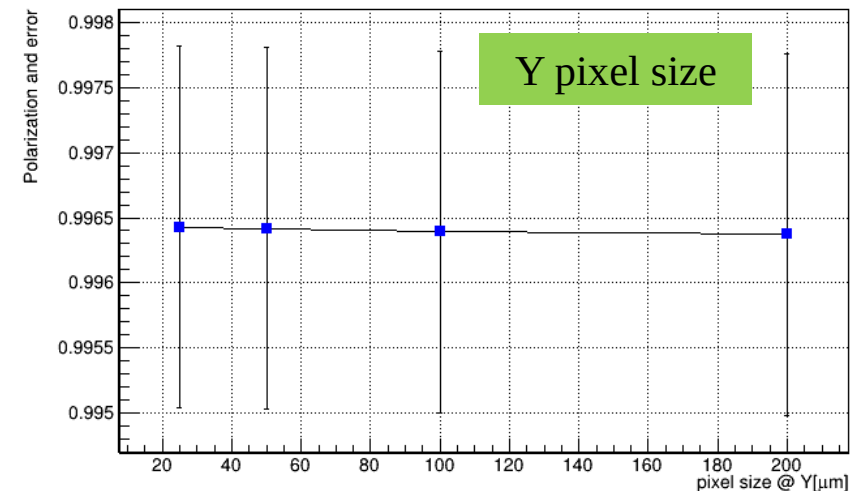
The active dimension of detector :
140mm × 1mm

The relationship between pixel size and uncertainty					
X			Y		
Pixel size	number	Fit result	Pixel size	number	Fit result
175μm*25μm	800*40	0.99642±0.00139	700μm*25μm	200*40	0.99643±0.00139
350μm*25μm	400*40	0.99639±0.00139	700μm*50μm	200*20	0.99642±0.00139
700μm*25μm	200*40	0.99632±0.00139	700μm*100μm	200*10	0.99639±0.00139
1400μm*25μm	100*40	0.99612±0.00139	700μm*200μm	200*5	0.99637±0.00139
2800μm*25μm	50*40	0.99530±0.00139	*****	*****	*****

The fit result vs pixel size



The fit result vs pixel size



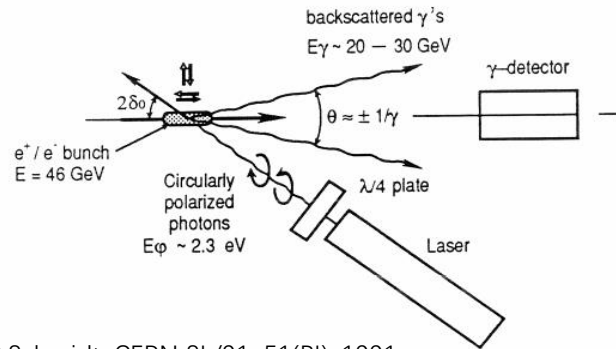
Discussion & Summary

- Compton Polarimeter is the clear technique of choice for electron polarization at CEPC
 - High precision ($\sim 1.4\%$) has been achieved in 20s time
 - The resolution of detector and the beam angle have been considered in our Toy MC
 - The measurement of polarization has low requirements for detector (different from calibrate the beam energy)
- The systematic error may be from : ① the uncertainty of beam energy; ② layout of drift distance and magnet; ③ The Angle of collision between laser beam and electron beam; ④ error from detector (A full simulation by Geant4 would be desirable to study the systematics uncertainty).
- The measurement of longitudinal polarization by Compton polarimeter are going on~

Backup

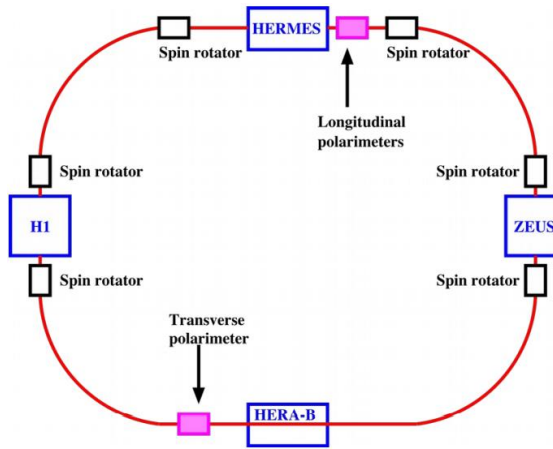
Compton polarimeter

CERN LEP 46GeV



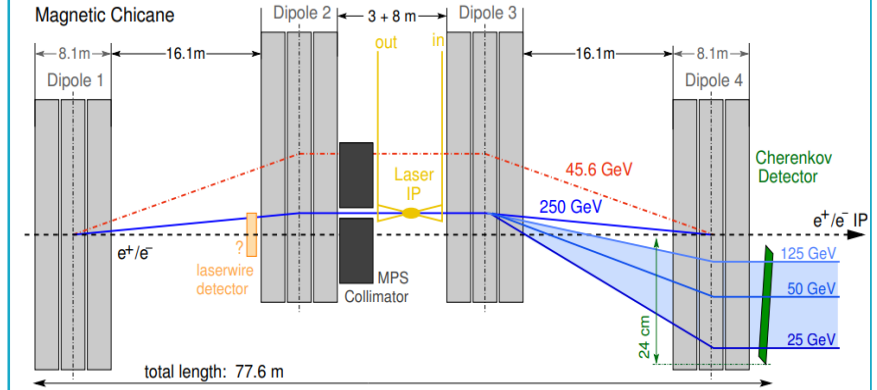
R.Schmidt. CERN SL/91-51(BI), 1991.

HERA 27GeV



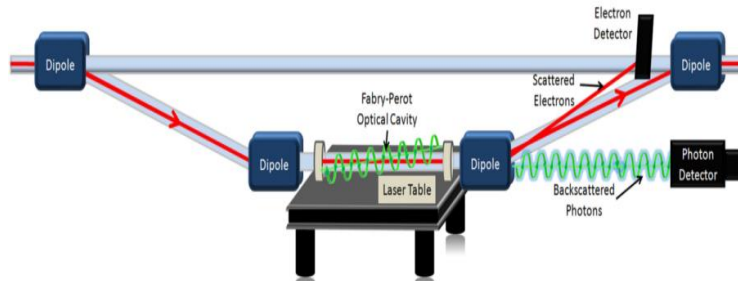
S.Schmitt, HERA polarimeter review

ILC 45.6GeV



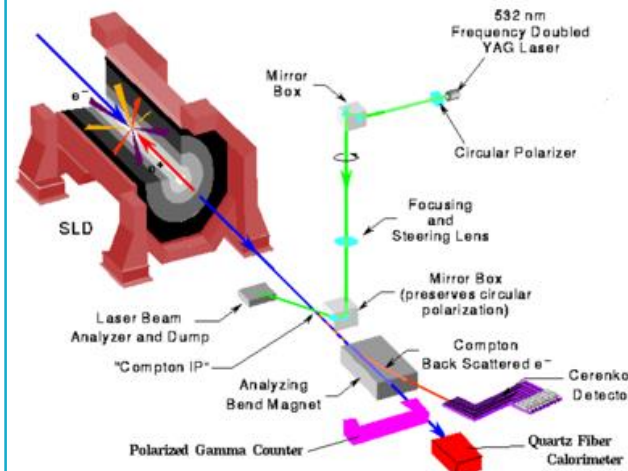
Jenny List, ILC Polarimetry, 2020

JLab Hall C 1.16GeV



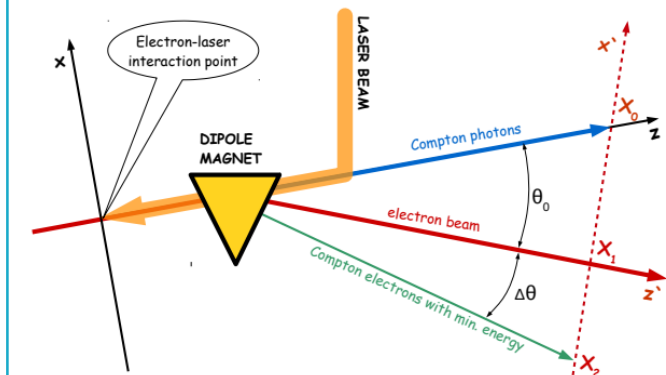
Donald Jones, Hall C Compton Polarimetry, PSTP 2013.

SLD at SLAC 45.6GeV



M.Wood. SLAC, 2003.

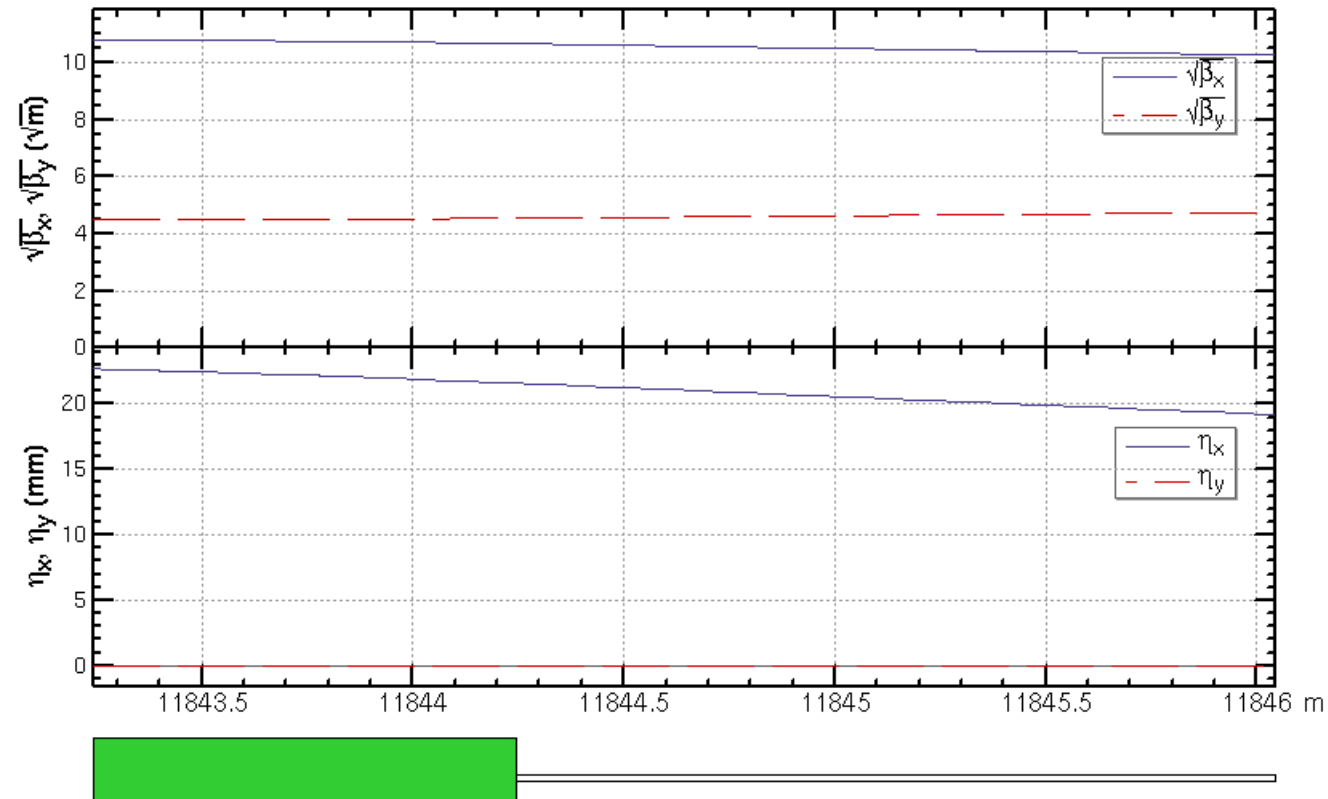
FCC-ee 45.6GeV



Nickolai Muchnoi, 2018

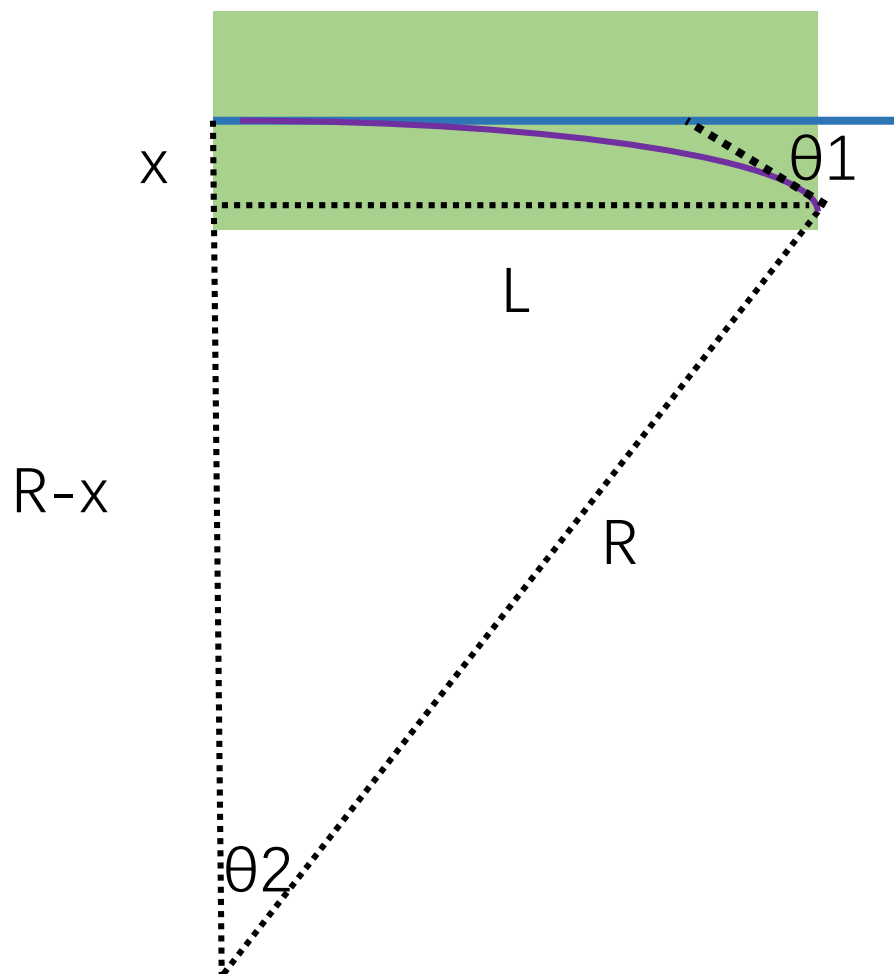
β function

$$\beta_x = 121m$$
$$\beta_y = 25m$$



About bending angle

- 推导偏转角和偏移量



- 偏转角等于圆心角

$$\text{bending angle: } \theta_1 = \theta_2 = \frac{l}{R}$$

- 偏移量 x : $R^2 = L^2 + (R - x)^2$

- $Bqv = mv^2/R$

$$E = mc^2 \rightarrow mc = E/c$$

$$R = \frac{mv^2}{Bqv} = \frac{mv}{Be} = \frac{E}{Bec}$$
$$= \frac{45.5 * 10^9 eV}{0.01867 T \times e \times 3 * 10^8 m/s}$$

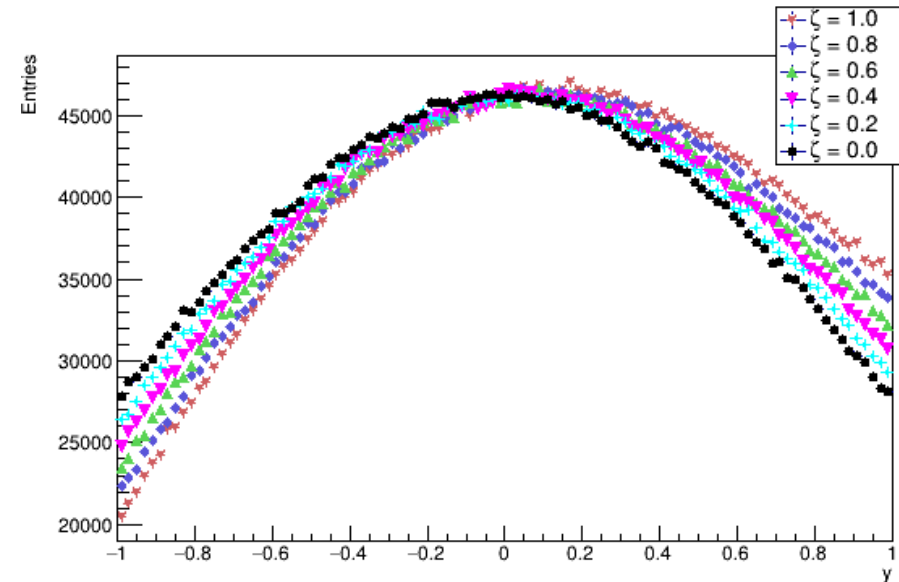
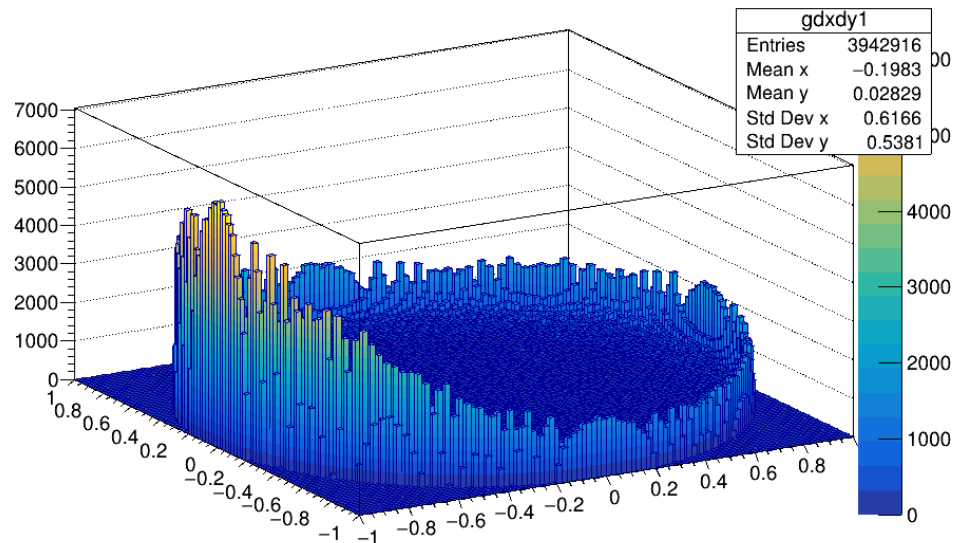
- $\theta = \frac{l}{R} = 3.4467 mrad$

Inverse Compton scattering cross section

$$d\sigma_0 = \frac{r_e^2}{\kappa^2(1+u)^3} \left(\kappa(1 + (1+u)^2) - 4\frac{u}{\kappa}(1+u)(\kappa-u) \left[1 - \xi_{\perp} \cos(2(\varphi - \varphi_{\perp})) \right] \right) du d\varphi,$$

$$d\sigma_{\parallel} = \frac{\xi_{\parallel} \zeta_{\parallel} r_e^2}{\kappa^2(1+u)^3} u(u+2)(\kappa-2u) du d\varphi,$$

$$d\sigma_{\perp} = -\frac{\xi_{\parallel} \zeta_{\perp} r_e^2}{\kappa^2(1+u)^3} 2u\sqrt{u(\kappa-u)} \cos(\varphi - \phi_{\perp}) du d\varphi.$$



Decide the dimension of detector

$$Xe_{max} \approx \kappa \theta_2 L_2 = \frac{4\omega_0 E_e}{m_e^2} \theta_2 L_2 = \frac{4\omega_0}{m_e^2} (\hbar B e c) L_2$$

Laser energy
Drift distance

Dipole length

$$Ye_{max} \approx \frac{L_1}{\gamma} \times \frac{\kappa}{2} = \frac{L_1 * \frac{4\omega_0 E_e}{m_e^2}}{2 \frac{E_e}{m_e}} = \frac{2\omega_0 L_1}{m_e}$$

$$\Pi(Xe)$$

$$\Pi(x) = \frac{\int y \frac{d\sigma_{\perp}}{dx dy} dy}{\int \frac{d\sigma_0}{dx dy} dy}$$

we use the $y = \sqrt{1-x^2} \sin \theta, \theta \in [-\frac{\pi}{2}, \frac{\pi}{2}]$

$$\int y \frac{d\sigma_{\perp}}{dx dy} dy = \int_{-\sqrt{1-x^2}}^{\sqrt{1-x^2}} \frac{r_e^2}{(1+u)^3 \sqrt{1-x^2-y^2}} u y^2 dy = \frac{u r_e^2}{(1+u)^3} \int_{-\sqrt{1-x^2}}^{\sqrt{1-x^2}} \frac{y^2}{\sqrt{1-x^2-y^2}} dy = \frac{(1-x^2) u r_e^2}{(1+u)^3} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \sin^3 \theta dy = \frac{\pi(1-x^2) u r_e^2}{2(1+u)^3}$$

$$\int \frac{d\sigma_0}{dx dy} dy = \int_{-\sqrt{1-x^2}}^{\sqrt{1-x^2}} \frac{r_e^2}{(1+u)^3 \sqrt{1-x^2-y^2}} \left(1 + (1+u)^2 - 4 \frac{u}{\kappa} (1+u) \left(1 - \frac{u}{\kappa} \right) \right) dy = \frac{\pi r_e^2 \left(1 + (1+u)^2 - 4 \frac{u}{\kappa} (1+u) \left(1 - \frac{u}{\kappa} \right) \right)}{(1+u)^3}$$

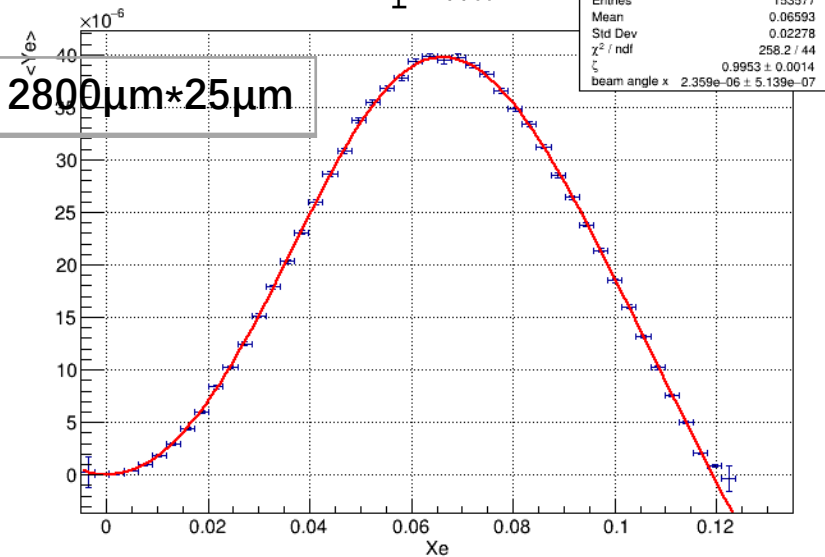
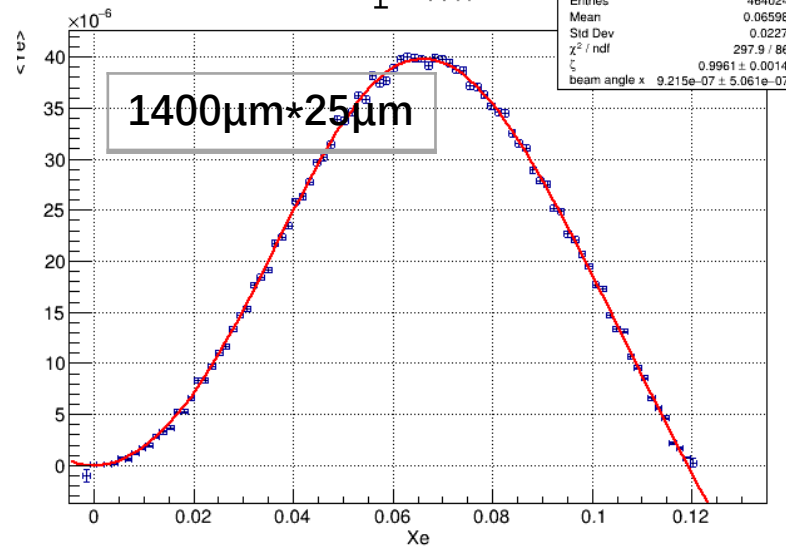
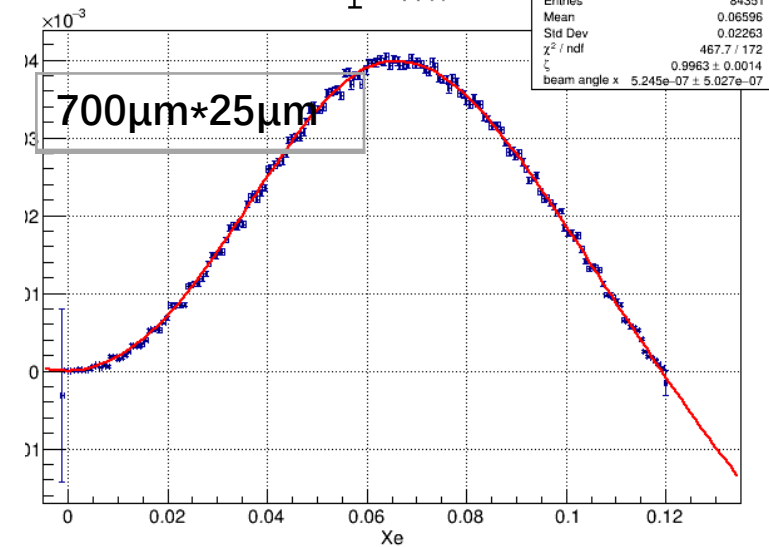
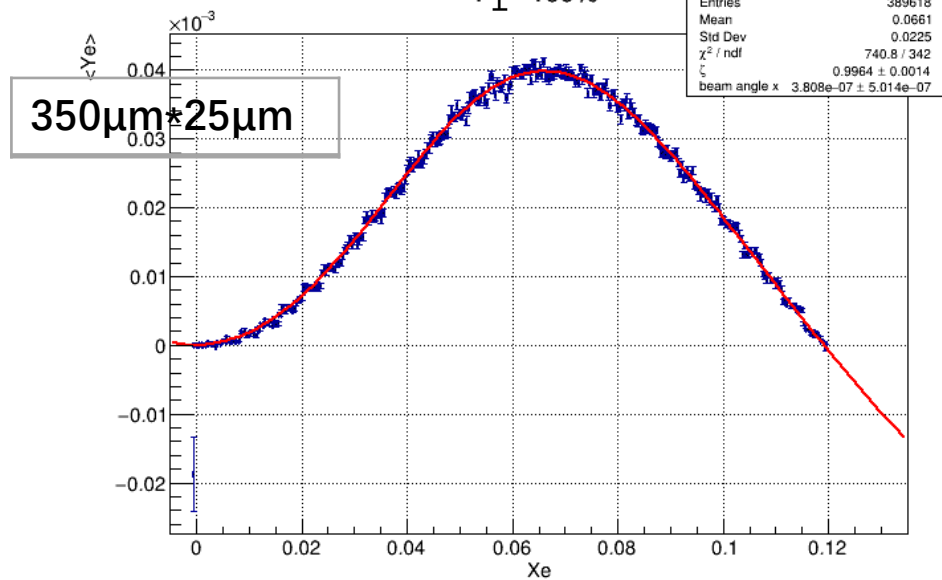
- $\Pi(x) = \frac{(1-x^2)u}{2 \left(1 + (1+u)^2 - 4 \frac{u}{\kappa} (1+u) \left(1 - \frac{u}{\kappa} \right) \right)}$
- $x = \frac{(X_e - \sigma'_x L_1) - \frac{\kappa}{2} \theta_0 L_2}{\frac{\kappa}{2} \sqrt{\left(\frac{L_1}{\gamma} \right)^2 + (\theta_0 L_2)^2}}$
- $y = \frac{Y_e - \sigma'_y L_1}{\frac{L_1 \kappa}{\gamma^2}}$
- $u(Xe) = \frac{2(X_e - \sigma'_x L_1) \theta_0 L_2 + \kappa \left(\frac{L_1}{\gamma} \right)^2}{2 \left(\left(\frac{L_1}{\gamma} \right)^2 + (\theta_0 L_2)^2 \right)}$

$$\Pi(Xe) = \frac{L_1 \kappa}{\gamma^2} \frac{\left(1 - \left(\frac{(X_e - \sigma'_x L_1) - \frac{\kappa}{2} \theta_0 L_2}{\frac{\kappa}{2} \sqrt{\left(\frac{L_1}{\gamma} \right)^2 + (\theta_0 L_2)^2}} \right)^2 \right) u(Xe)}{2 \left(1 + (1+u)^2 - 4 \frac{u}{\kappa} (1+u) \left(1 - \frac{u}{\kappa} \right) \right)}$$

Table 4.3.3.5: Parameters of the dual aperture dipole.

Beam center separation [mm]		350
Magnetic length [m]		28.686
Magnetic strength [Gs]		373.4
Gap [mm]		70
Coil	Number	
	Shape	
	Material	Aluminum
	Conductor specs. [mm]	30×54
Current [A]		1058
Current density [A/mm ²]		0.67
Resistance [mΩ]		2.44
Voltage [V]		2.58
Power consumption [kW]		2.73
Cooling water	Loop number	1
	Pressure drop [kg/cm ²]	6
	Velocity [m/s]	1.75
	Flux [l/s]	0.138
	Temperature rise [°C]	4.7

The magnetic strength of the dipole located in last is $\frac{1}{2}B = \frac{373.4Gs}{2} = 0.01867T$

$P_{\perp} = 100\%$  $P_{\perp} = 100\%$  $P_{\perp} = 100\%$  $P_{\perp} = 100\%$  $P_{\perp} = 100\%$ 