

## Outline

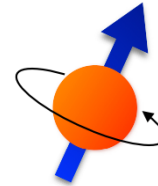
- Brief historical introduction to electron and muon magnetic anomaly
- Dispersion-relation based prediction to muon magnetic anomaly
- Alternative way used to evaluate the prediction
- Summary and perspectives

Work in collaboration with M. Davier, A. Hoecker, B. Malaescu (DHMZ) and others

# (Anomalous) Magnetic Moment of a Charged Lepton

For a charged lepton with charge  $q_\ell$  and mass  $m_\ell$ , its magnetic moment is connected to its spin by a  $g_\ell$  factor:

$$\vec{\mu}_\ell = g_\ell \left( \frac{q_\ell}{2m_\ell} \right) \vec{S}_\ell$$



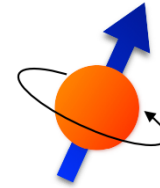
Dirac predicted  $g_e = 2$  for electron in 1928 [Proc. Roy. Soc. Lond. A 118, 351 (1928) ]

It was confirmed to 0.1% by Kinsler & Houston in 1934 through studying the Zeeman effect in neon [ Phys. Rev. 46, 533 (1934) ]

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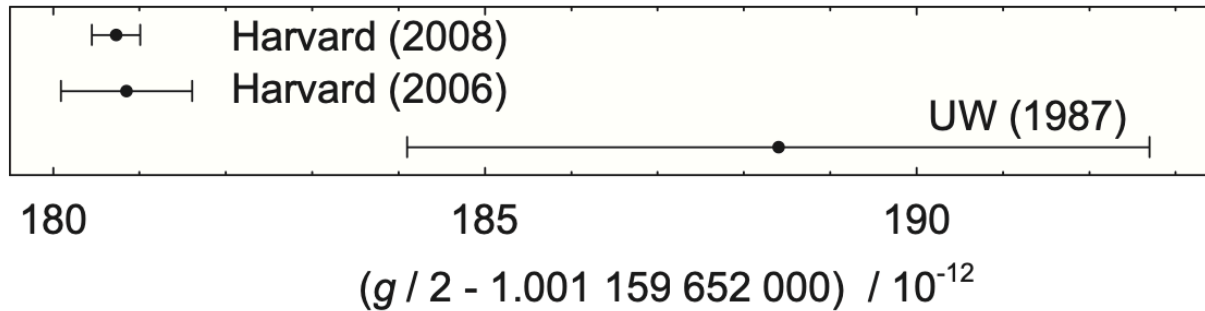
A deviation from  $g_e = 2$  was established by Nafe, Nels & Rabi only in 1947 by comparing the hyperfine structure of hydrogen and deuterium spectra [ Phys. Rev. 71, 914 (1947) ]

A first precision measurement of  $g_e = 2.00344 \pm 0.00012$  (wrong: 2.00232...!) was made by Kusch & Foley in 1947 using Rabi's atomic beam magnetic resonance technique [ Phys. Rev. 72, 1256 (1947) ]

The anomalous magnetic moment  $a_\ell$  was introduced to quantify the deviation from 2:

$$a_\ell = \frac{g_\ell - 2}{2}$$

# The $a_e$ Measurements



$$a_e = 1\ 159\ 652\ 180.73(28) \times 10^{-12} \quad [ \text{Phys. Rev. Lett. 100, 120801 (2008), Phys. Rev. A 83, 052122 (2011)} ]$$

24 ppb precision

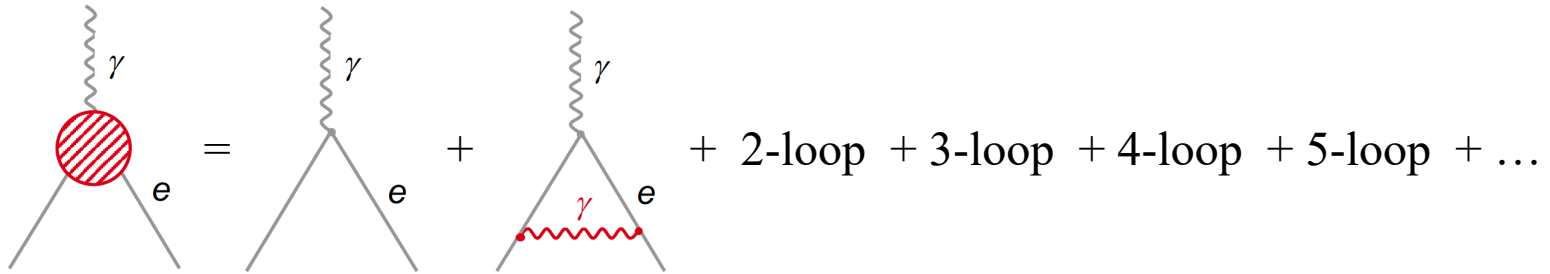
using a one-electron quantum cyclotron in a cylindrical Penning trap cavity, a technique invented by Dehmelt in 70's who was awarded with a Nobel Prize in Physics in 1989

This is the most precisely measured quantity in particle physics

Do we have a prediction with a comparable precision to compare with?



# The $a_e$ Prediction



$$a_e = 0 + \frac{\alpha}{2\pi} + A_2 \left(\frac{\alpha}{\pi}\right)^2 + A_3 \left(\frac{\alpha}{\pi}\right)^3 + A_4 \left(\frac{\alpha}{\pi}\right)^4 + A_5 \left(\frac{\alpha}{\pi}\right)^5 + \dots$$

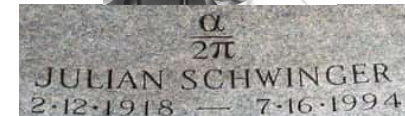
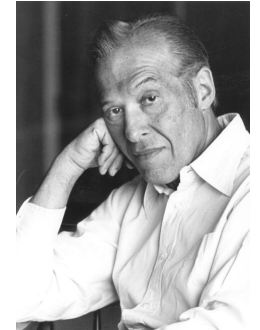
(Dirac:  $g_e = 2$ ) (Schwinger 1947)

$A_2, A_3$  known analytically,  
 $A_2$ – $A_4$ : cross-checked by different groups using different methods  
 $A_5$ : calculated by one group with numerical means, only small portions double-checked

Aoyama, Hayakawa, Kinoshita, Nio(2012-2019)



Paul Dirac  
Nobel prize 1933



Nobel prize 1965

# The $a_e$ Prediction

$$a_e = 0 + \frac{\alpha}{2\pi} + A_2 \left(\frac{\alpha}{\pi}\right)^2 + A_3 \left(\frac{\alpha}{\pi}\right)^3 + A_4 \left(\frac{\alpha}{\pi}\right)^4 + A_5 \left(\frac{\alpha}{\pi}\right)^5 + \dots$$

Using the latest  $\alpha$  measurement: 137.035 999 206(11) [ Morel et al., Nature 588, 61 (2020) ]  
one gets

$$a_e \text{ (SM prediction)} = 1\,159\,652\,180.252(95) \times 10^{-12}$$

which includes a tiny contribution of

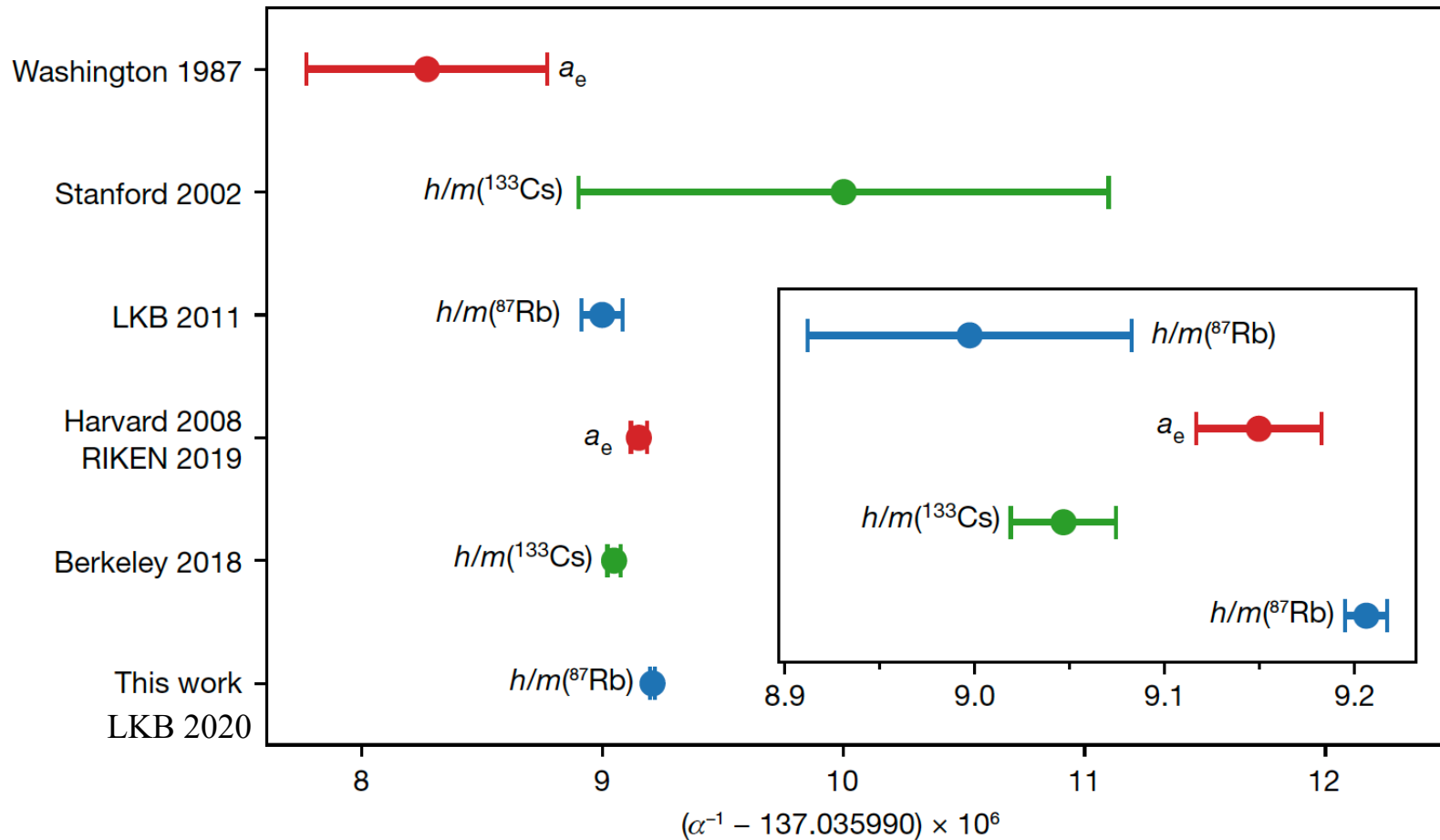
$$\begin{aligned} a_e \text{ (hadron)} &= 1.6927 (120) \times 10^{-12} \\ a_e \text{ (weak)} &= 0.03052 (23) \times 10^{-12} \end{aligned} \quad [ \text{Jegerlehner, 1711.06089} ]$$

Thus

$$a_e \text{ (exp)} - a_e \text{ (SM prediction)} = (4.8 \pm 3.0) \times 10^{-13} [+1.6\sigma] \rightarrow \text{Great success of QED and the SM!?!}$$

# Summary on $a_e$ and $\alpha$

[ Morel et al., Nature 588, 61 (2020) ]



The latest the  $\alpha$  measurement from Rb differs from that from Cs by more than 5 standard deviations!

# Why Are We Interested in $a_\mu$ ?

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Contrary to the electron, the muon is unstable ( $2.2\mu\text{s}$ )

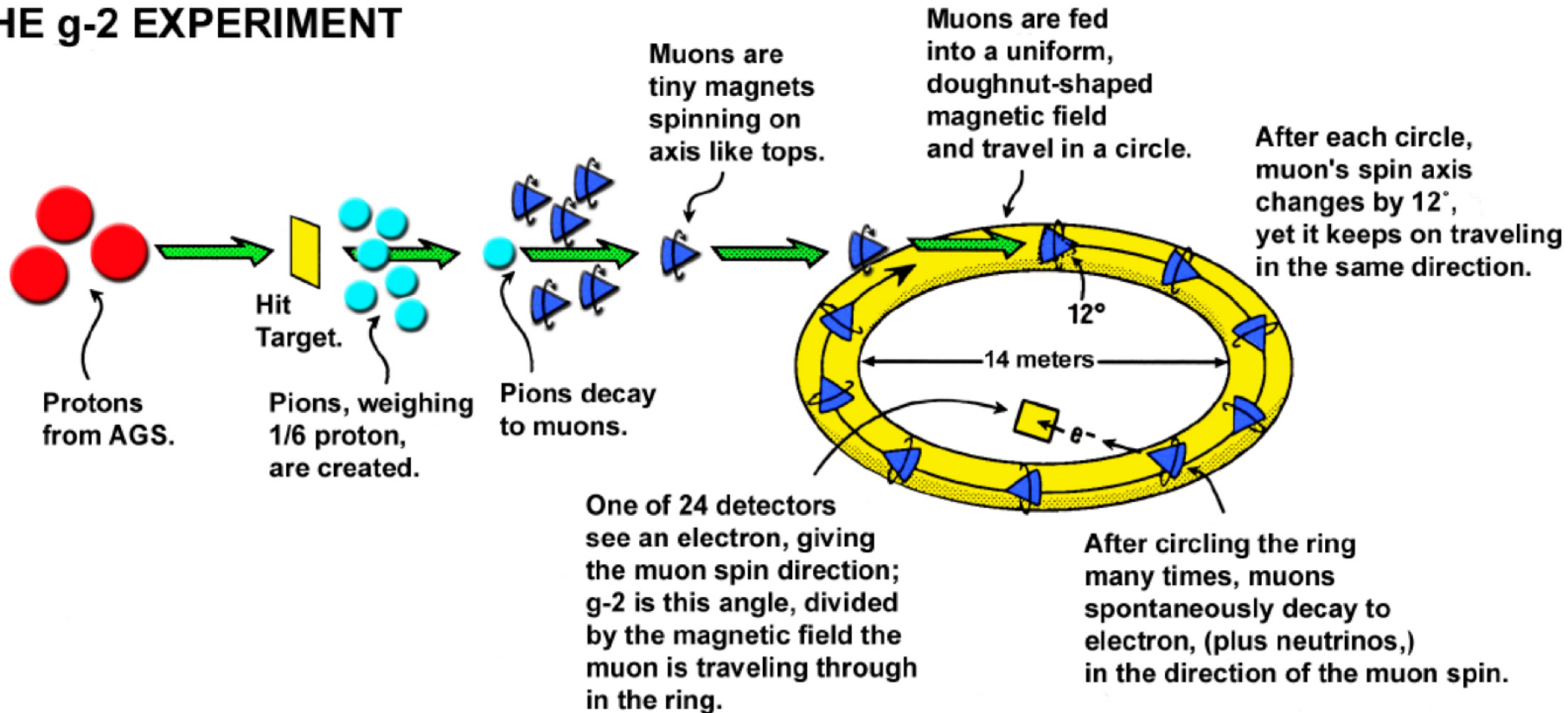
so it is more difficult to measure (and predict), nevertheless

$a_\mu$  receives sizeable contributions from all three sectors of the SM

and due to its heavier mass, its sensitivity to new physics is  $\sim (m_\mu/m_e)^2 \sim 43\,000$  larger

# Overview of a Muon $g-2$ Experiment

## LIFE OF A MUON: THE $g-2$ EXPERIMENT



The key elements:

1. Parity violation in pion decay: polarised muons
2. Parity violation in muon decay: positron emitted in the direction of muon spin

# The Measurement Concept

Measure “anomalous” frequency difference between spin precession and cyclotron frequencies:

$$\vec{\omega}_a \equiv \vec{\omega}_s - \vec{\omega}_c = \frac{e}{m_\mu c} \left[ a_\mu \vec{B} - \left( a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} \right]$$

One defines a magic  $\gamma$  value of 29.3 (i.e.  $p_\mu=3.09$  GeV) to remove the 2nd term

The magnetic field B is determined from  $\omega_p = 2\mu_p B$

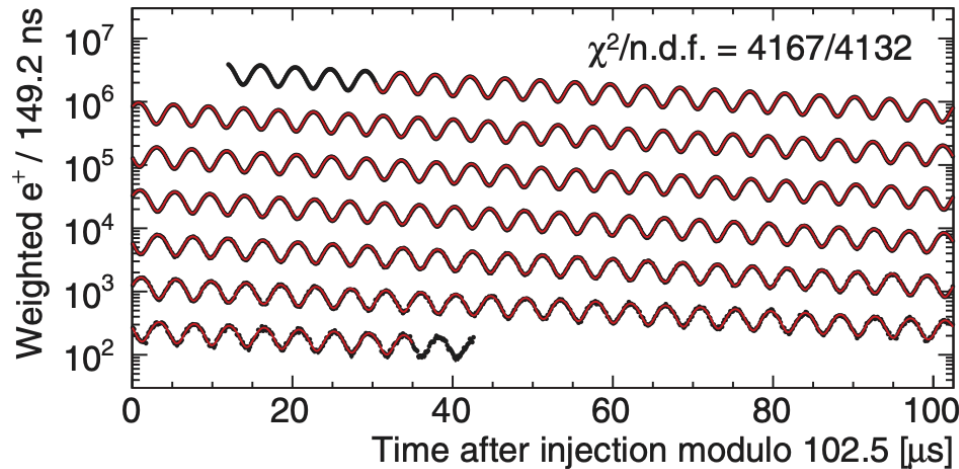
Thus

$$a_\mu = \frac{\omega_a}{\omega_p} \left[ \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2} \right]$$

The 1st ratio term is measured in a double blinded way

The other ratio term within the brackets is known with high precision  $\pm 25$  ppb

# FNAL Muon g-2 Result



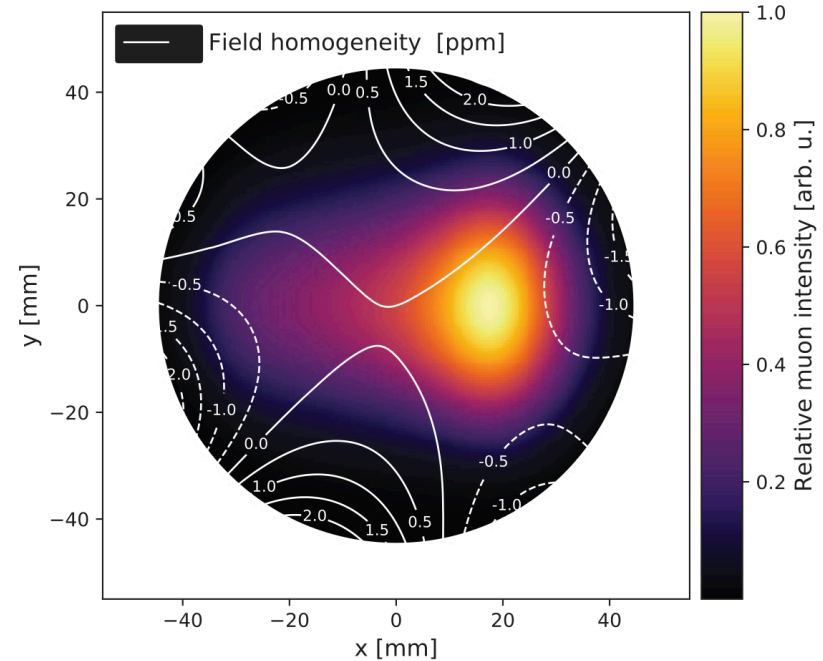
Positron rate measured in the calorimeters is fitted with

$$N(t) = N_0 e^{-t/\gamma\tau} [1 - A \cdot \cos(\omega_a t + \phi)]$$

to determine  $\omega_a$  with a stat and syst error of 434, 56 ppb


$$a_\mu(\text{FNAL}) = 116\,592\,040(54) \times 10^{-11} \quad (0.46 \text{ ppm})$$

The total uncertainty of 0.46 ppm is dominated by statistical one of 0.43 ppm at Run 1



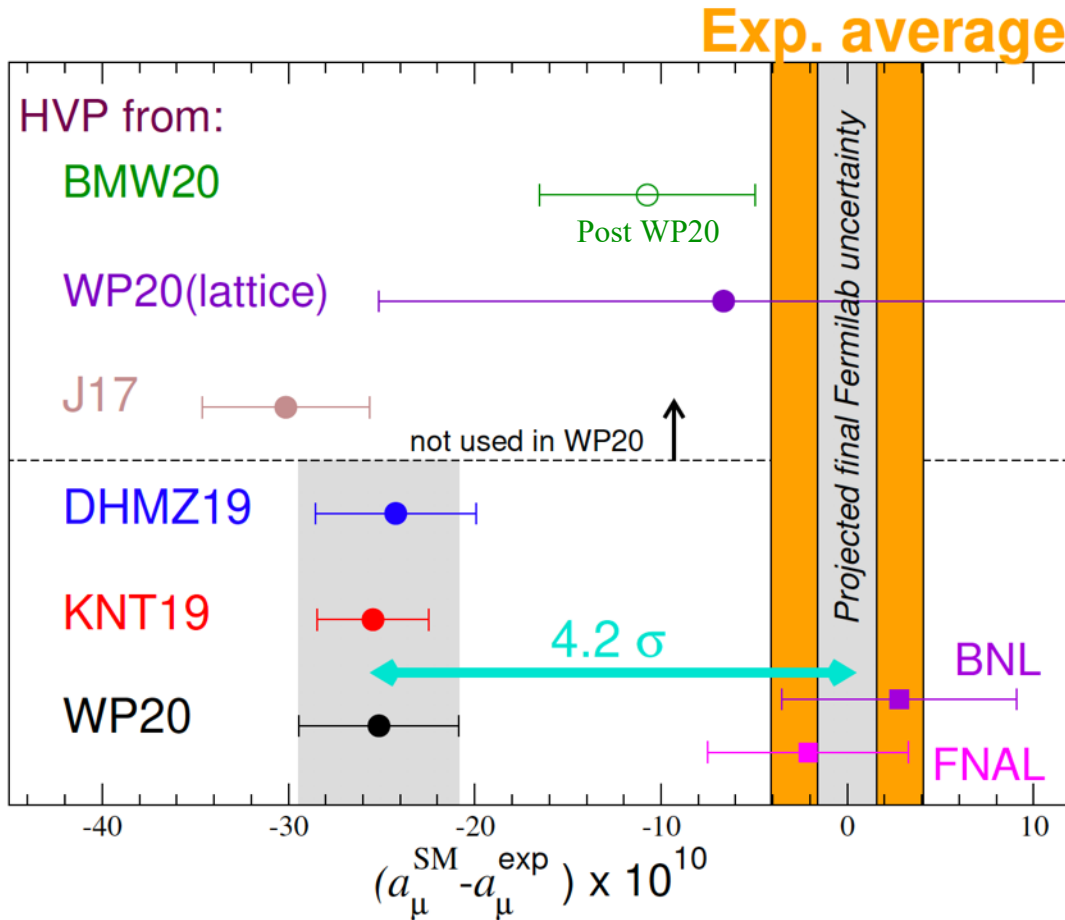
The magnetic field is measured with a syst error of 56 ppb

# More than 60 years of measurements

Experiment	Beam	Measurement	$\delta a_\mu/a_\mu$	Required th. terms
Columbia-Nevis (57)	$\mu^+$	$g=2.00\pm 0.10$		$g=2$
Columbia-Nevis (59)	$\mu^+$	0.001 13(+16)(-12)	12.4%	$\alpha/\pi$
CERN 1 (61)	$\mu^+$	0.001 145(22)	1.9%	$\alpha/\pi$
CERN 1 (62)	$\mu^+$	0.001 162(5)	0.43%	$(\alpha/\pi)^2$
CERN 2 (68)	$\mu^+$	0.001 166 16(31)	265 ppm	$(\alpha/\pi)^3$
CERN 3 (75)	$\mu^\pm$	0.001 165 895(27)	23 ppm	$(\alpha/\pi)^3 + \text{had}$
CERN 3 (79)	$\mu^\pm$	0.001 165 911(11)	7.3 ppm	$(\alpha/\pi)^3 + \text{had}$
BNL E821 (00)	$\mu^+$	0.001 165 919 1(59)	5 ppm	$(\alpha/\pi)^3 + \text{had}$
BNL E821 (01)	$\mu^+$	0.001 165 920 2(16)	1.3 ppm	$(\alpha/\pi)^4 + \text{had} + \text{weak}$
BNL E821 (02)	$\mu^+$	0.001 165 920 3(8)	0.7 ppm	$(\alpha/\pi)^4 + \text{had} + \text{weak} + ?$
BNL E821 (04)	$\mu^-$	0.001 165 921 4(8)(3)	0.7 ppm	$(\alpha/\pi)^4 + \text{had} + \text{weak} + ?$
 FNAL Run1 (21)	$\mu^+$	0.001 165 920 40(54)	0.46 ppm	$(\alpha/\pi)^4 + \text{had} + \text{weak} + ?$



# Current Situation



WP20: White Paper published in 2020  
*Phys. Rept.* 887 (2020) 1 ([link](#))

An outcome after several dedicated workshops since 2017

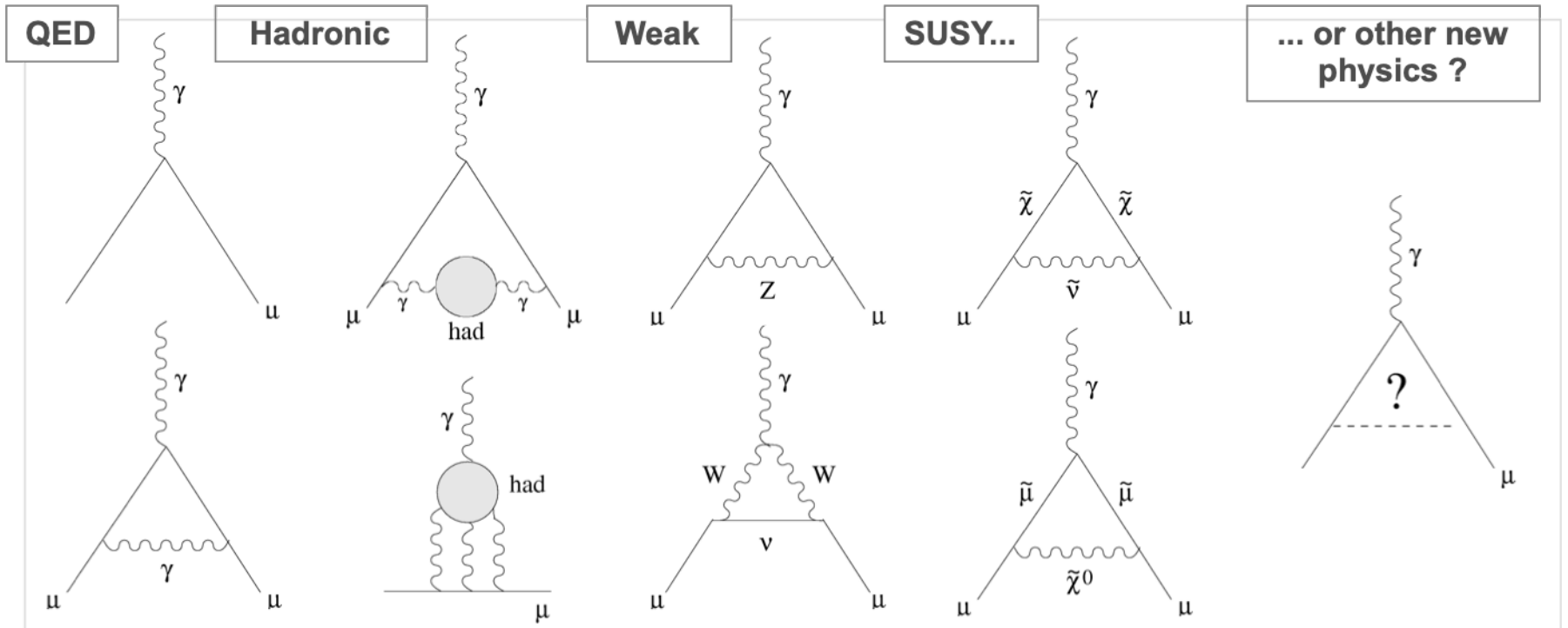
In the following, I shall review the prediction using our DHMZ19 as an example

**A discrepancy of  $4.2\sigma$**   
**→ Strong evidence for new physics?**

# Theoretical Contributions

$$a_{\mu}^{\text{th}} = a_{\mu}^{\text{SM}} + a_{\mu}^{\text{BSM}}$$

$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{had}} + a_{\mu}^{\text{Weak}} \quad a_{\mu}^{\text{BSM}}$$



# Contributions of the SM components

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QED:  $116\,584\,718.9 (0.1) \times 10^{-11}$  [0.001 ppm]

Weak:  $153.6 (1.0) \times 10^{-11}$  [ 0.01 ppm]

Hadronic ...

... Vacuum polarisation (HVP):  $6845.0 (40.0) \times 10^{-11}$  [ 0.37 ppm]

... Light-by-Light (HLbL):  $92.0 (18.0) \times 10^{-11}$  [ 0.15 ppm]

HVP has the largest uncertainty to the prediction

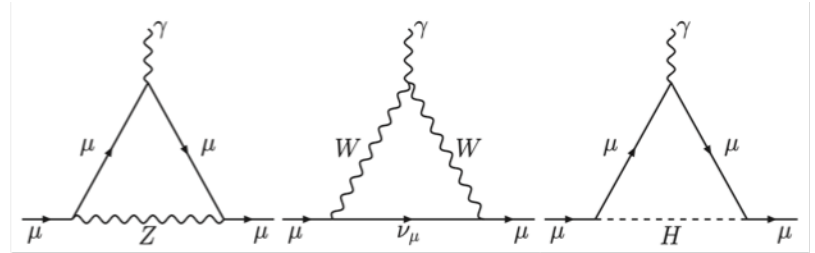
In the following, unless stated otherwise, the numbers are from WP20

# Prediction QED

$$\begin{aligned} a_{\mu}^{\text{QED}} &= \frac{\alpha}{2\pi} + A_2 \left(\frac{\alpha}{\pi}\right)^2 + A_3 \left(\frac{\alpha}{\pi}\right)^3 + A_4 \left(\frac{\alpha}{\pi}\right)^4 + A_5 \left(\frac{\alpha}{\pi}\right)^5 && \text{No of Feynman} \\ & && \text{diagrams:} \\ &= && 116\,140\,793.321 \quad (23) \times 10^{-11} && \mathcal{O}(\alpha) && 1 \\ &+ && 413\,217.626 \quad (7) \times 10^{-11} && \mathcal{O}(\alpha^2) && 7 \\ &+ && 30\,141.902 \quad (33) \times 10^{-11} && \mathcal{O}(\alpha^3) && 72 \\ &+ && 381.004 \quad (17) \times 10^{-11} && \mathcal{O}(\alpha^4) && 891 \\ &+ && 5.078 \quad (6) \times 10^{-11} && \mathcal{O}(\alpha^5) && 12\,672 \\ &= && 116\,584\,718.931 \quad (104) \times 10^{-11} && && \end{aligned}$$

# Prediction Weak

The weak contributions are defined as all SM contributions that are not contained in the pure QED, the HVP, or the HLbL contributions



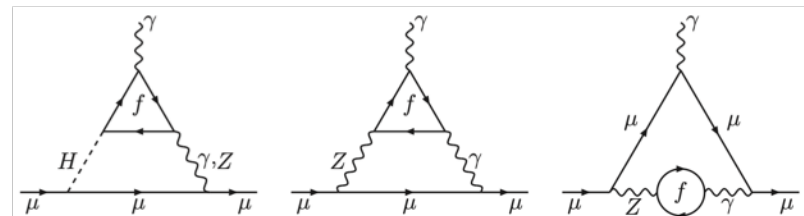
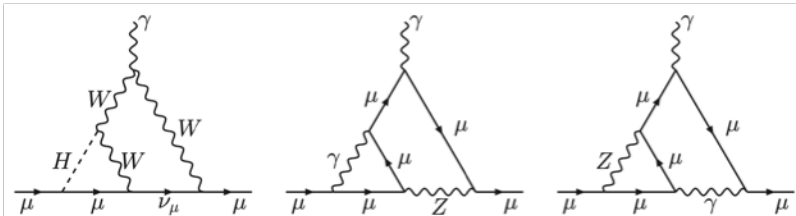
$$a_{\mu}^{\text{Weak}(1)} = \frac{G_F m_{\mu}^2}{\sqrt{2} 8\pi^2} \left[ \frac{5}{3} + \frac{1}{3} (1 - 4 \sin^2_W)^2 + \mathcal{O}\left(\frac{m_{\mu}^2}{m_W^2}\right) + \mathcal{O}\left(\frac{m_{\mu}^2}{m_H^2}\right) \right]$$

$$= 194.79(1) \times 10^{-11}$$

$$a_{\mu}^{\text{Weak}(2)} = a_{\mu}^{\text{bos}} + a_{\mu}^{e,\mu,u,c,d,s} + a_{\mu}^{\tau,t,b} + a_{\mu}^H + a_{\mu}^{\text{no } H}$$

$$= [-19.96(1) - 6.91(20)(30) - 8.21(10) - 1.51(1) - 4.64(10)] \times 10^{-11}$$

$$a_{\mu}^{\text{Weak}(\geq 3)} = 0(0.20) \times 10^{-11}$$

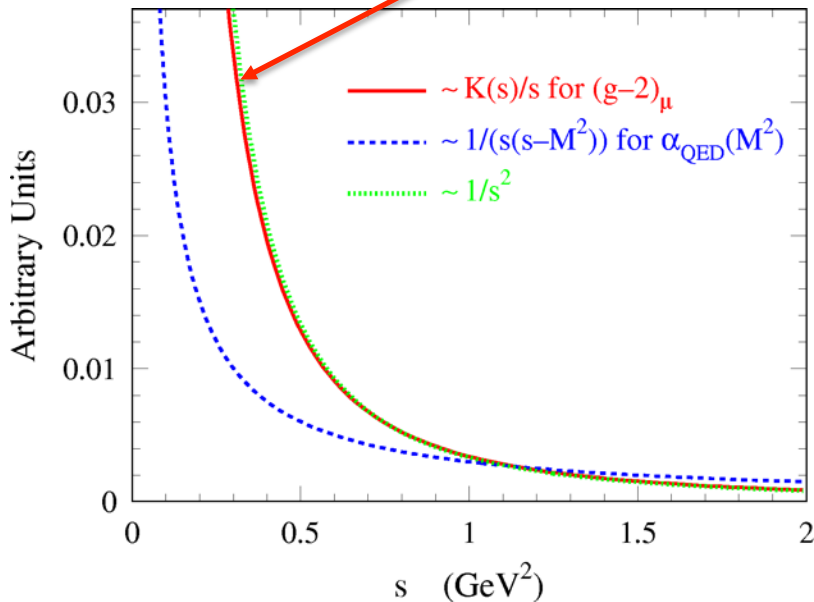


# Prediction Hadronic

$$a_{\mu}^{\text{had}} = a_{\mu}^{\text{had,LO}} + a_{\mu}^{\text{had,HO}} + a_{\mu}^{\text{had,LBL}}$$

Based on analyticity and unitarity, the LO HVP contribution can be calculated using the dispersion relation [1] over  $e^+e^- \rightarrow \text{hadrons}$  cross sections

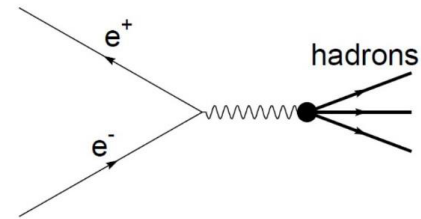
$$a_{\mu}^{\text{had,LO}} = \frac{\alpha^2}{3\pi^2} \int_{4m_{\pi}^2}^{\infty} ds \frac{K(s)}{s} R(s)$$



Born:  $\sigma^{(0)}(s) = \sigma(s)(\alpha/\alpha(s))^2$

$$12\pi \text{Im}\Pi_{\gamma}(s) = \frac{\sigma^0 [e^+e^- \rightarrow \text{hadrons} (\gamma_{\text{FSR}})]}{\sigma_{pt}} \equiv R(s)$$

$$\text{Im}[\text{Diagram}] \propto |\text{Diagram} \rightarrow \text{hadrons}|^2$$



The QED kernel  $K(s)$  [2] has such an  $s$  dependence that low energy data contribute most:

$e^+e^- \rightarrow \pi^+\pi^-$  contributes  $\sim 73\%$  (58% in uncertainty)

**→ The precision is data-driven!**

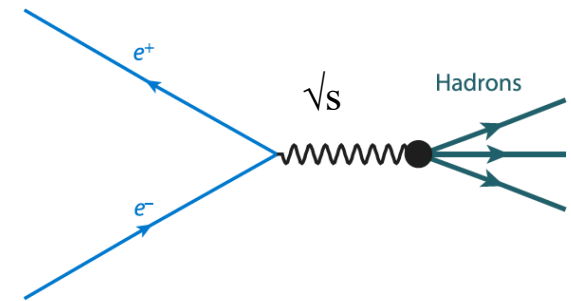
[1] Bouchiat and Michel, 1961

[2] Brodsky, de Rafael, 1968

# Two Types of $\sigma(e^+e^- \rightarrow \text{hadrons})$ Measurements

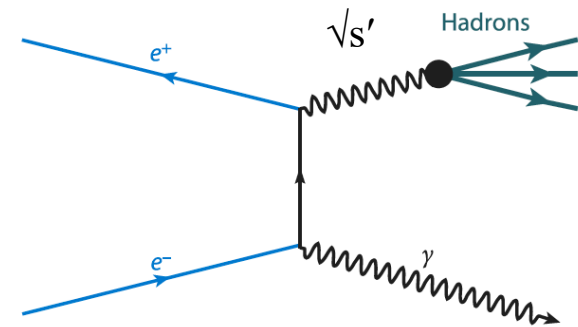
1. The scan method: e.g. CMD-2/3, SND at Novosibirsk

- Advantages:
  - Well defined  $\sqrt{s}$
  - Good energy resolution  $\sim 10^{-3}\sqrt{s}$
- Disadvantages:
  - Energy gap between two scans
  - Low luminosity at low energies
  - Limited  $\sqrt{s}$  range of a given experiment



2. The ISR approach: e.g. BABAR, BES, CLEO-c, KLOE

- Advantages:
  - Continuous cross section measurement over a broad energy range down to threshold
  - Large acceptance for hadrons if ISR detected at large angle
  - $\sigma(e^+e^- \rightarrow \text{hadrons})$  may be measured over  $\sigma(e^+e^- \rightarrow \mu^+\mu^-)$  thus reducing some syst uncertainties
- Disadvantages:
  - Require high luminosity to compensate higher order in  $\alpha$  (that's why the  $\sqrt{s}$  is chosen at a resonance)



$$s' = (1-x)/s$$
$$x = 2E_\gamma/\sqrt{s}$$

# A Large Number of Exclusive Processes below 1.8GeV

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Channel

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$\pi^0\gamma$   
 $\eta\gamma$   
 $\pi^+\pi^-$   
 $\pi^+\pi^-\pi^0$   
 $2\pi^+2\pi^-$   
 $\pi^+\pi^-2\pi^0$   
 $2\pi^+2\pi^-\pi^0$  ( $\eta$  excl.)  
 $\pi^+\pi^-3\pi^0$  ( $\eta$  excl.)  
 $3\pi^+3\pi^-$   
 $2\pi^+2\pi^-2\pi^0$  ( $\eta$  excl.)  
 $\pi^+\pi^-4\pi^0$  ( $\eta$  excl., isospin)  
 $\eta\pi^+\pi^-$   
 $\eta\omega$   
 $\eta\pi^+\pi^-\pi^0$  (non- $\omega$ ,  $\phi$ )  
 $\eta 2\pi^+2\pi^-$   
 $\omega\eta\pi^0$   
 $\omega\pi^0$  ( $\omega \rightarrow \pi^0\gamma$ )  
 $\omega 2\pi$  ( $\omega \rightarrow \pi^0\gamma$ )  
 $\omega$  (non- $3\pi$ ,  $\pi\gamma$ ,  $\eta\gamma$ )  
 $K^+K^-$   
 $K_S K_L$   
 $\phi$  (non- $K\bar{K}$ ,  $3\pi$ ,  $\pi\gamma$ ,  $\eta\gamma$ )  
 $K\bar{K}\pi$   
 $K\bar{K}2\pi$   
 $K\bar{K}\omega$   
 $\eta\phi$   
 $\eta K\bar{K}$  (non- $\phi$ )  
 $\omega 3\pi$  ( $\omega \rightarrow \pi^0\gamma$ )  
 $7\pi$  ( $3\pi^+3\pi^-\pi^0$  + estimate)

List of 30 channels evaluated in DHMZ19 [ *Eur. Phys. J. C* 80 (2020) 3, 241, [link](#) ]

DHMZ group involved in HVP evaluation since 1997 with more than 10 publications and over 3500 citations ([link](#))

Result used as reference for the Brookhaven experiment: comparison revealed a deficit in the prediction at  $\sim 2$ - $3\sigma$  level, hence our motivation to continue this effort for a more precise prediction

In the following, a few examples of measurements from different experiments in the dominant channels will be shown

Then we discuss

- the combination of different measurements of a given channel
- comparison and tension between different measurements

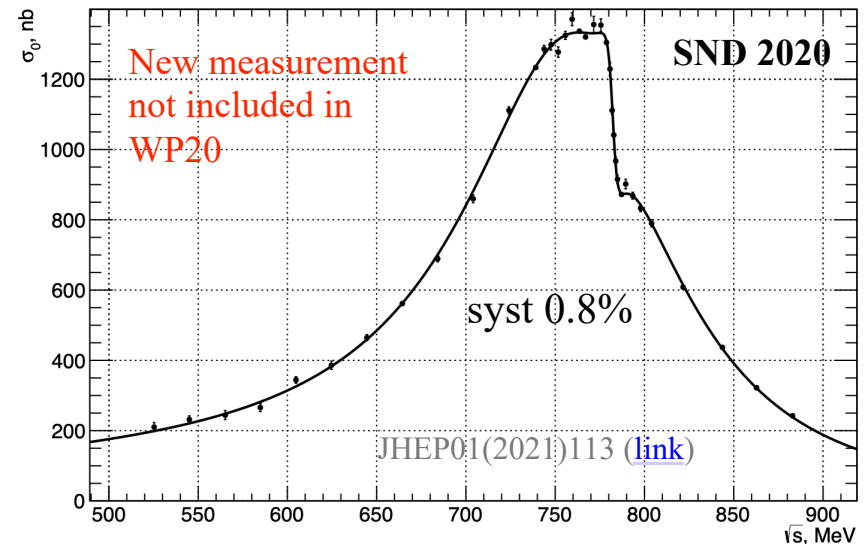
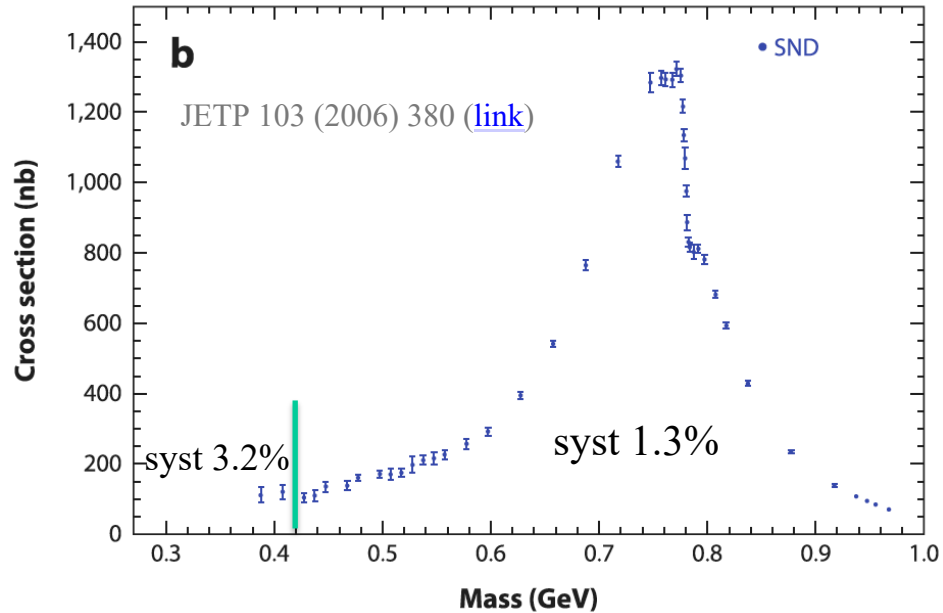
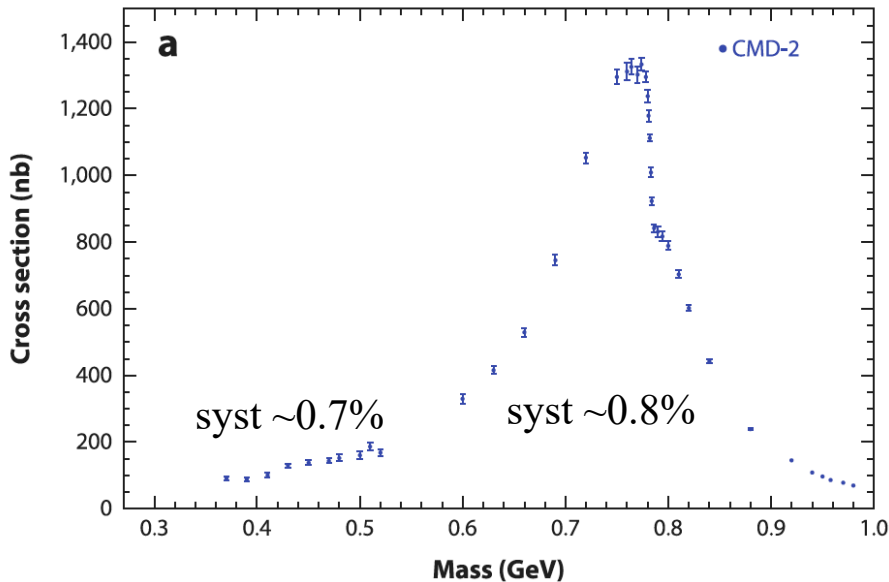


# Measurements of $2\pi$ Channel: CMD-2, SND

CMD-2 (2006)

- Energy 0.35-0.52 GeV [ JETP Lett. 84:413-417, 2006 ([link](#)) ]

- Energy 0.6-1.0 GeV [ Phys. Lett. B 648: 28-38, 2007 ([link](#)) ]

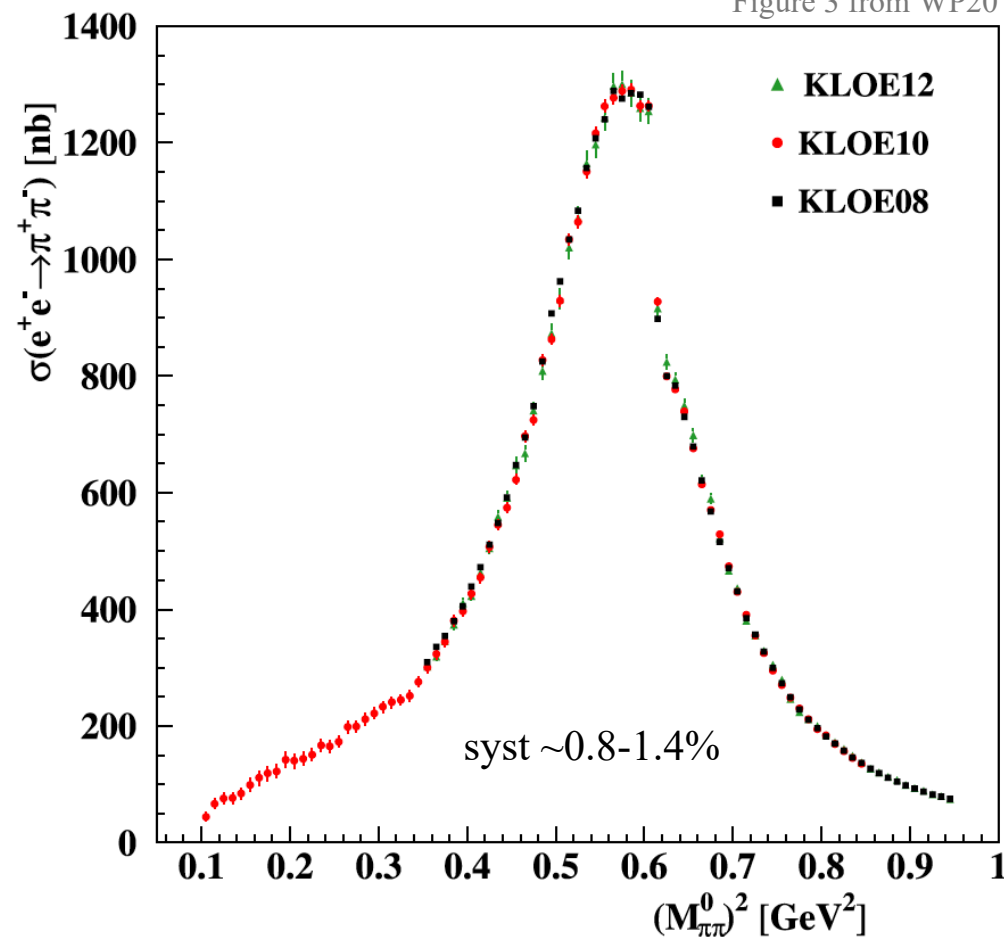


Figures a, b from M. Davier, Ann. Rev. Nucl. Part. Sci. 63 (2013) 407 ([link](#))

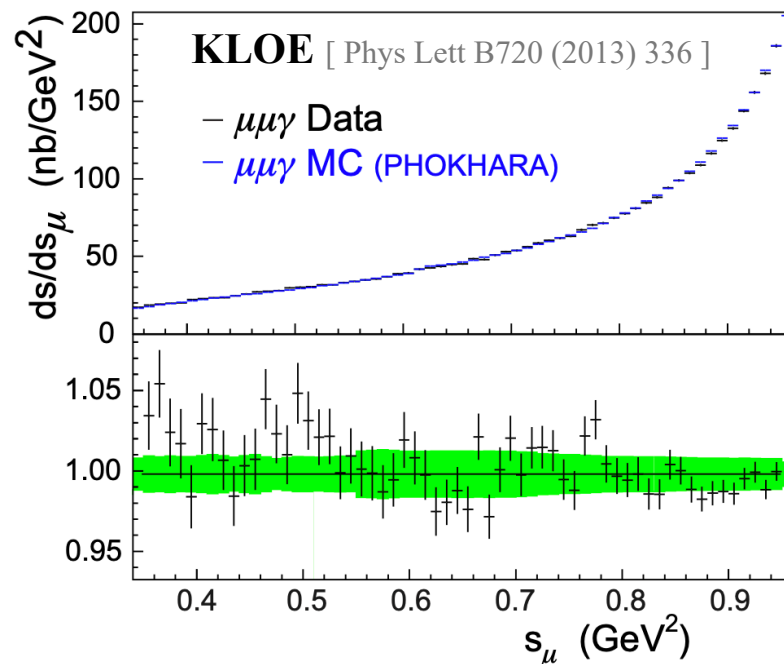
# Measurements of $2\pi$ Channel: KLOE 08,10,12

$\sqrt{s}=1.02$  GeV  $\Rightarrow$  Soft ISR photons

Figure 3 from WP20



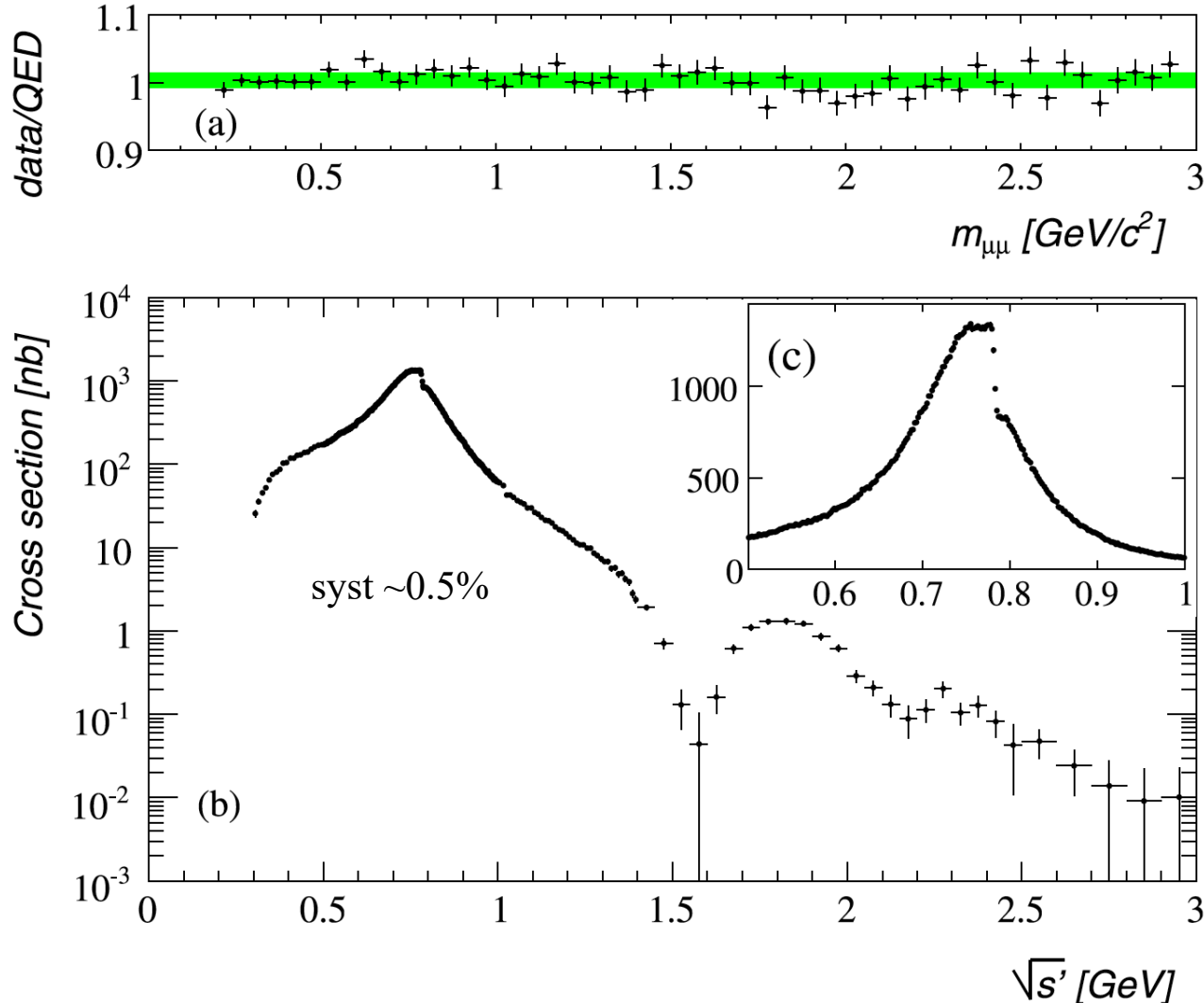
- KLOE12: photon at small angle and undetected, radiator function from measured  $\mu^+\mu^-(\gamma)$  events
- KLOE10: photon at large angle and detected, radiator function from NLO QED
- KLOE08: photon at small angle and undetected, radiator function from NLO QED



# Measurements of $2\pi$ Channel: BaBar 09

$\sqrt{s}=10.58$  GeV  $\Rightarrow$  Hard ISR photons

Figure 3 from WP20



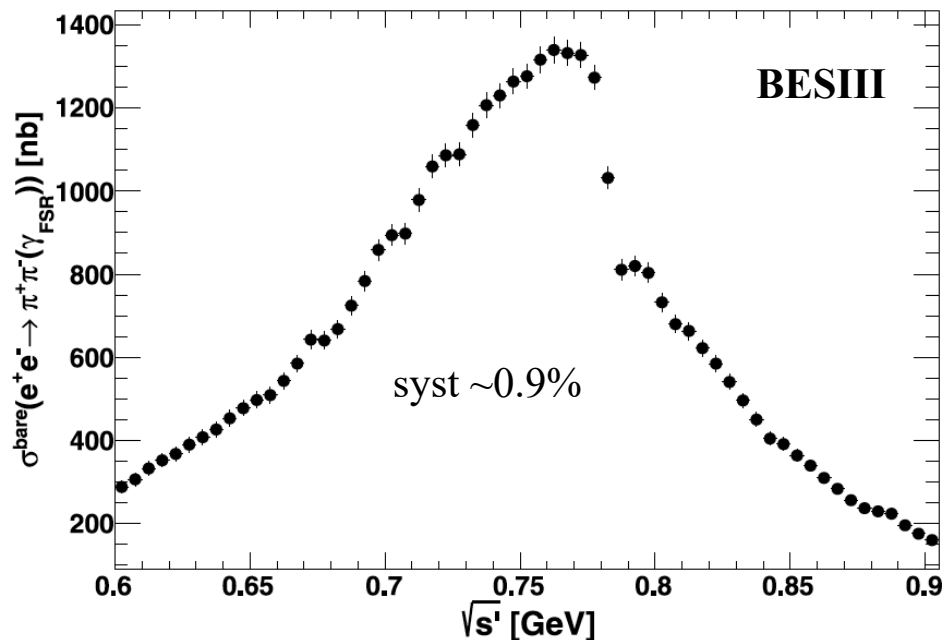
BABAR measurement covers a huge mass range from threshold to 3 GeV!

In BABAR, the ISR photon is detected at large angle

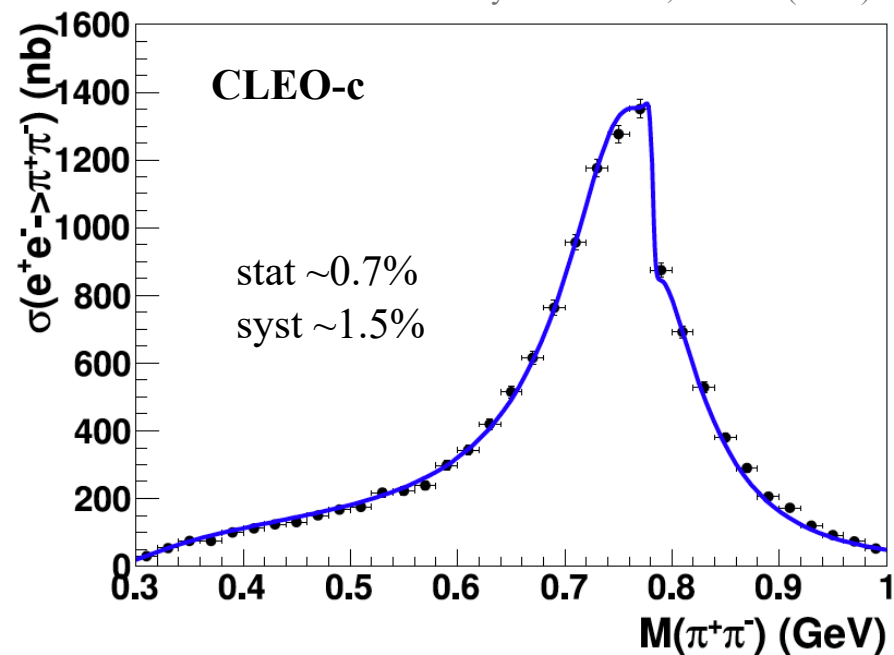
Both pion and muon pairs are measured and the ratio  $\pi\pi(\gamma)/\mu\mu(\gamma)$  directly provides the  $\pi\pi(\gamma)$  cross section

# Measurements of $2\pi$ Channel: BESIII, CLEO-c

Original publication: Phys. Lett. B 753, 629 (2016)  
Erratum: Phys. Lett. B 812, 135982 (2021) for updating  
the stat covariance matrix

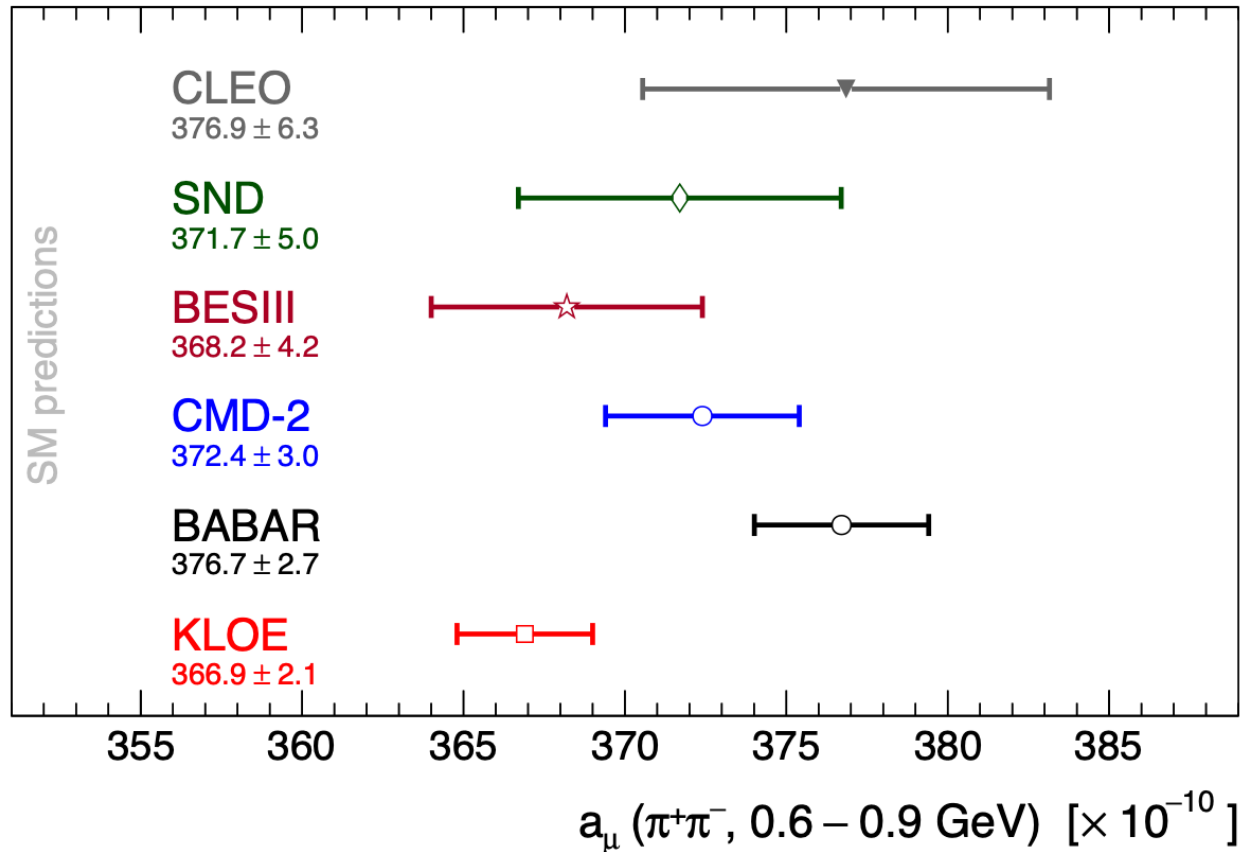


Phys. Rev. D 97, 032012 (2018)



# Comparison

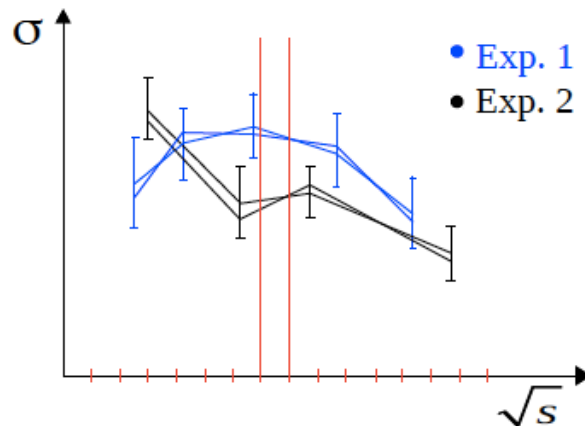
Figure from DHMZ, EPJC80 (2020) 241



BABAR and KLOE most precise but in clear discrepancy  
Combination needs special treatment (see later)

# Data Combination Using HVPTools\*

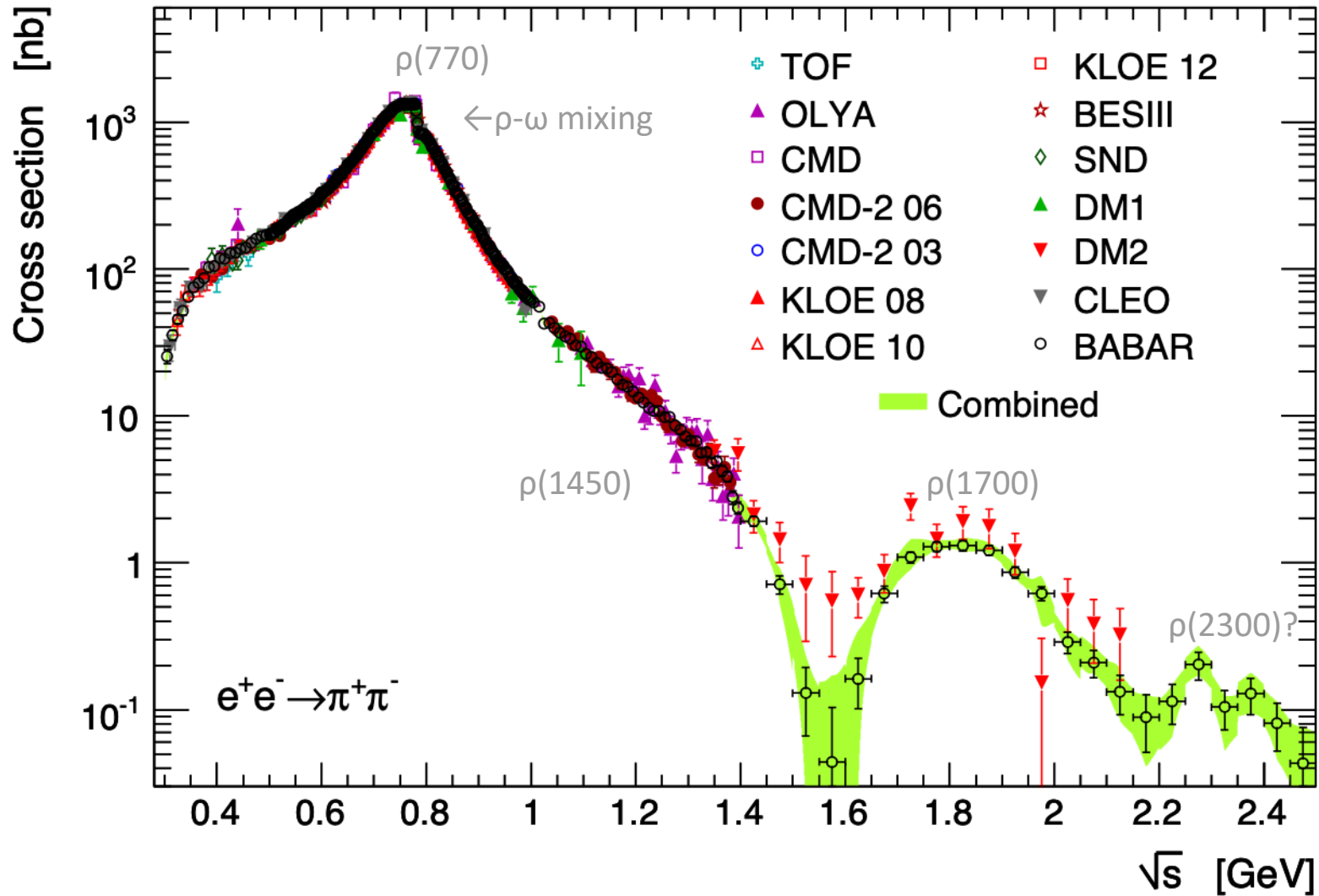
- Combine experimental spectra with arbitrary spacing/binning
- Properly propagate uncertainties and correlations
  - ▶ Between measurements (data points/bins) of a given experiment
  - ▶ Between experiments
  - ▶ Between different channels
- Linear/quadratic splines to interpolate between points/bins of each experiment
  - ▶ For binned measurements: preserve integral inside each bin
- Fluctuate data points taking into account correlations and redo the splines for each (pseudo)experiment
  - ▶ Each uncertainty fluctuated coherently for all points/bins that it impacts
  - ▶ Eigenvector decomposition for (stat & syst) covariance matrices
- Resulting combination shown in fine binning
  - ▶ Local error inflation following PDG prescription ( $\sqrt{\chi^2/\text{dof}}$ ) to take better into account data tension



\* HVPTools: Davier et al.,  
EPJC66 (2010) 127

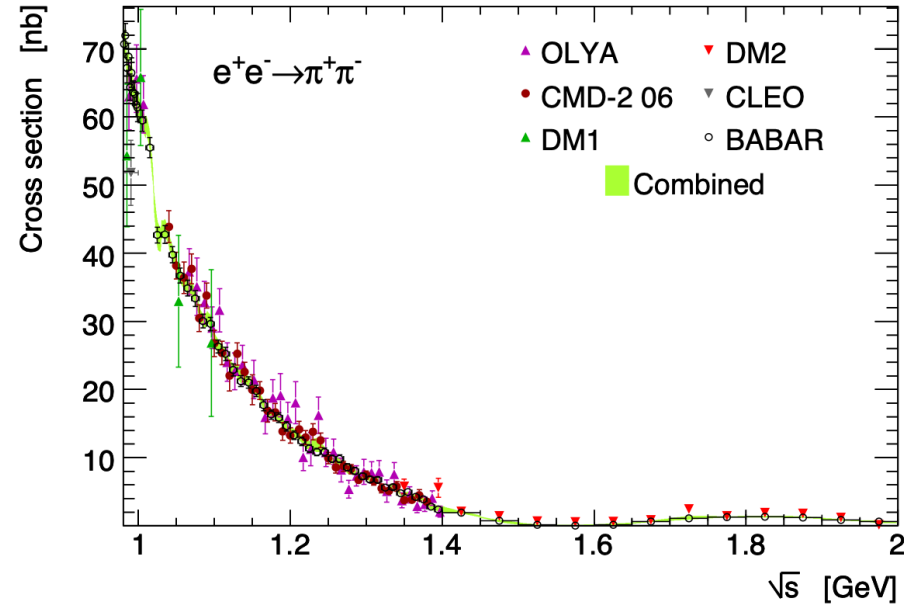
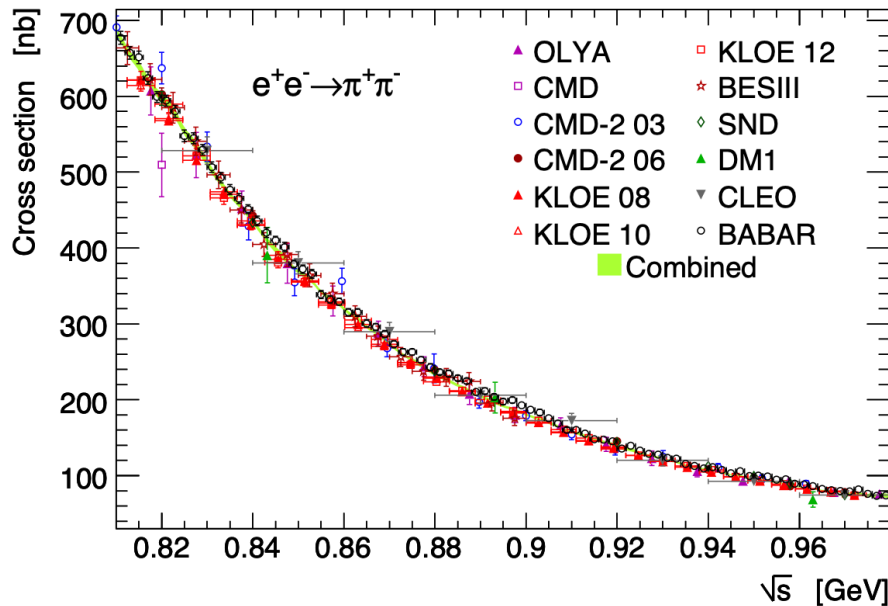
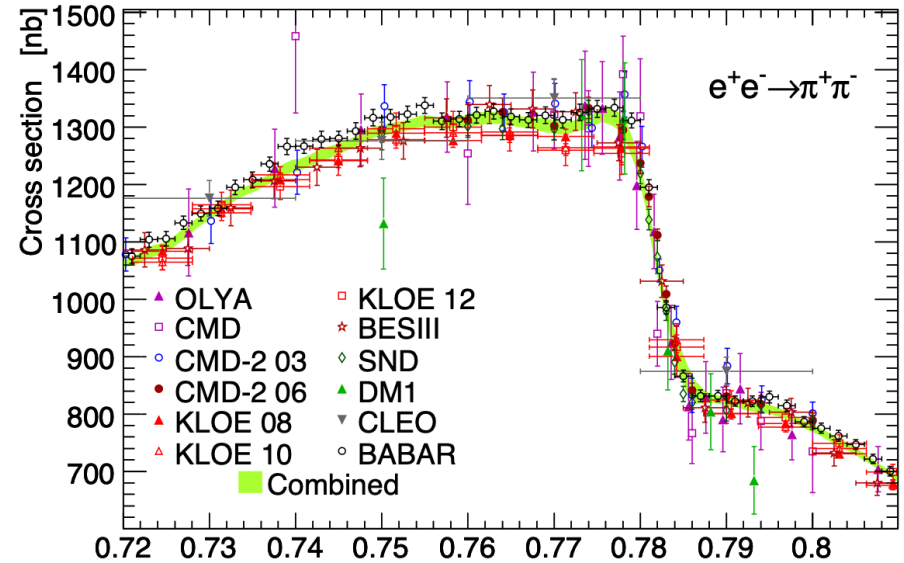
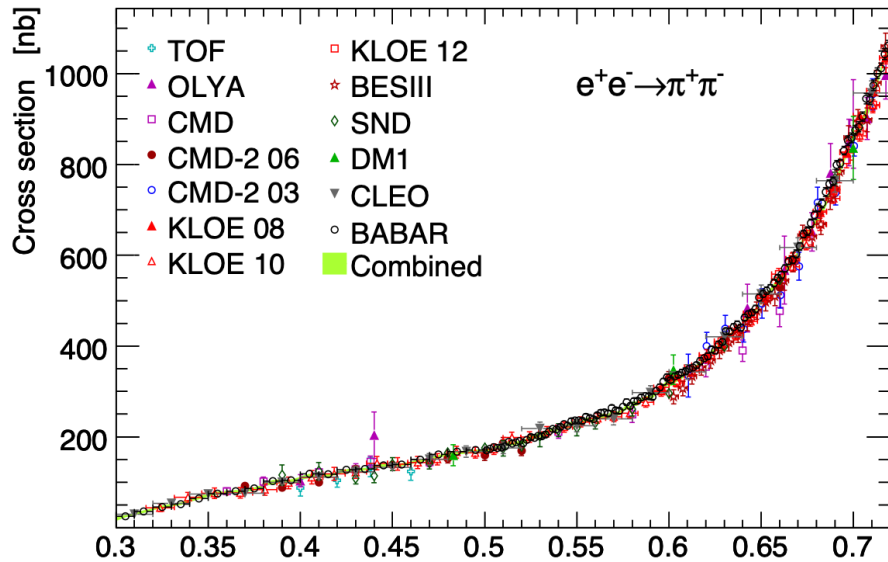
# Combined $2\pi$ vs Individual Measurements

Figures from DHMZ, EPJC80 (2020) 241



# Combined $2\pi$ vs Individual Measurements

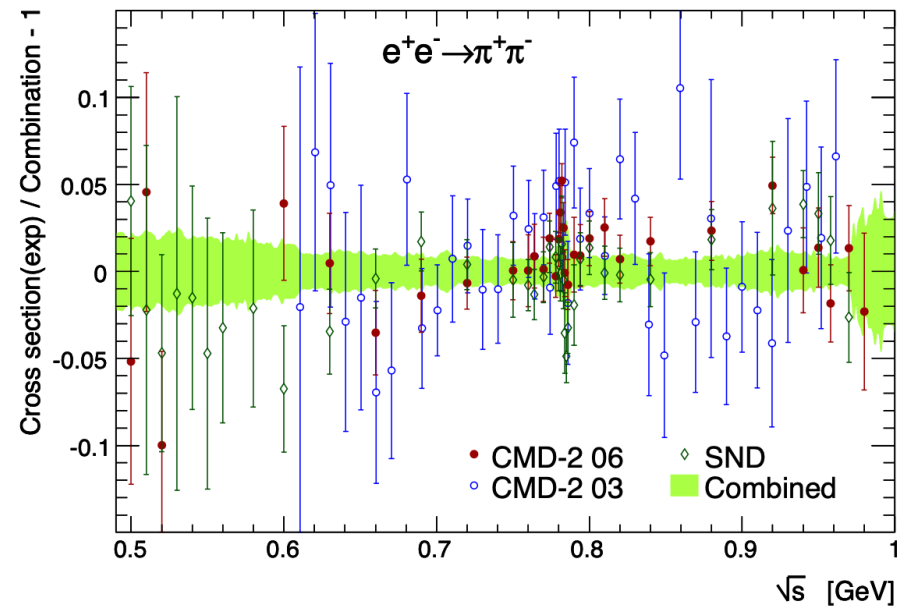
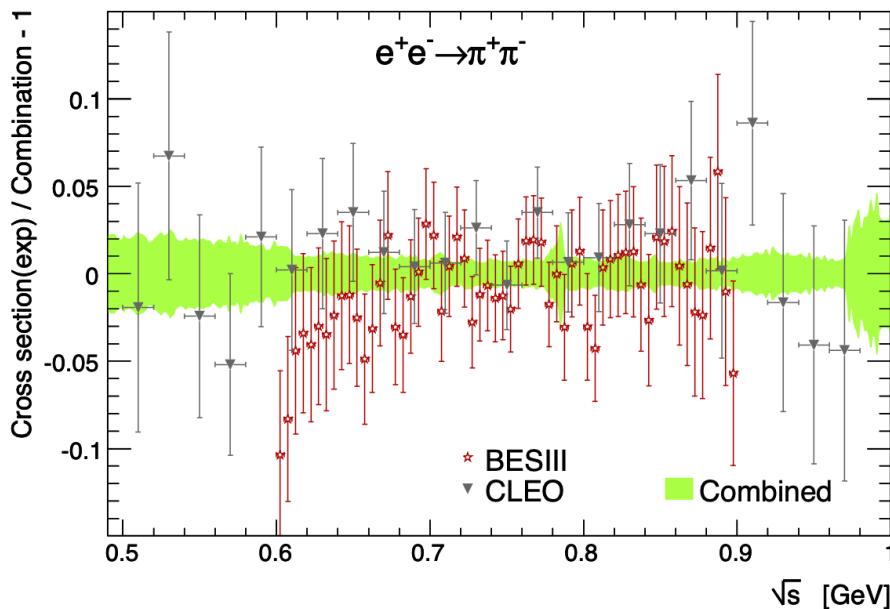
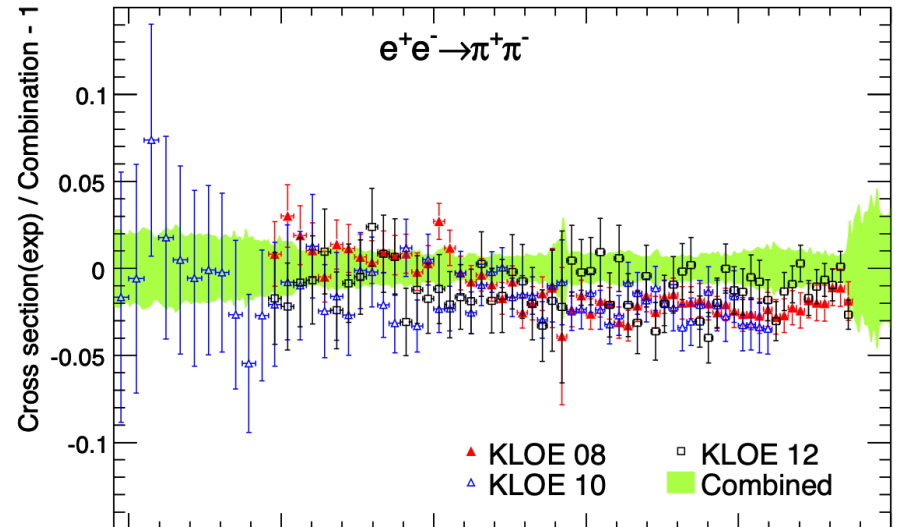
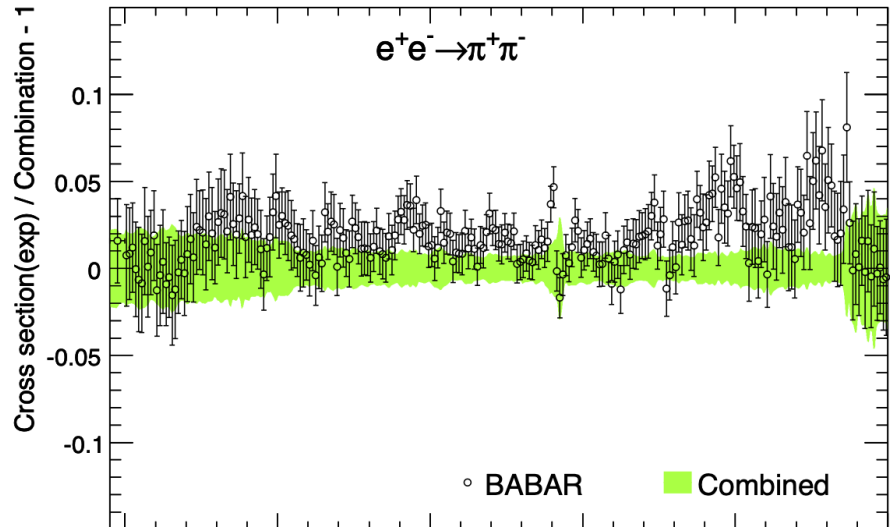
Figures from DHMZ, EPJC80 (2020) 241





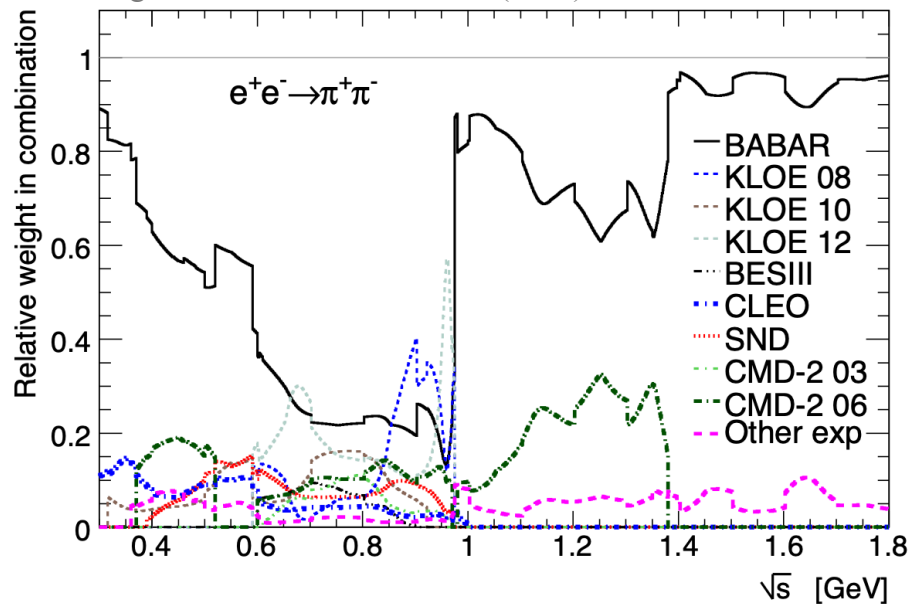
# Combined $2\pi$ vs Individual Measurements

Figures from DHMZ, EPJC80 (2020) 241



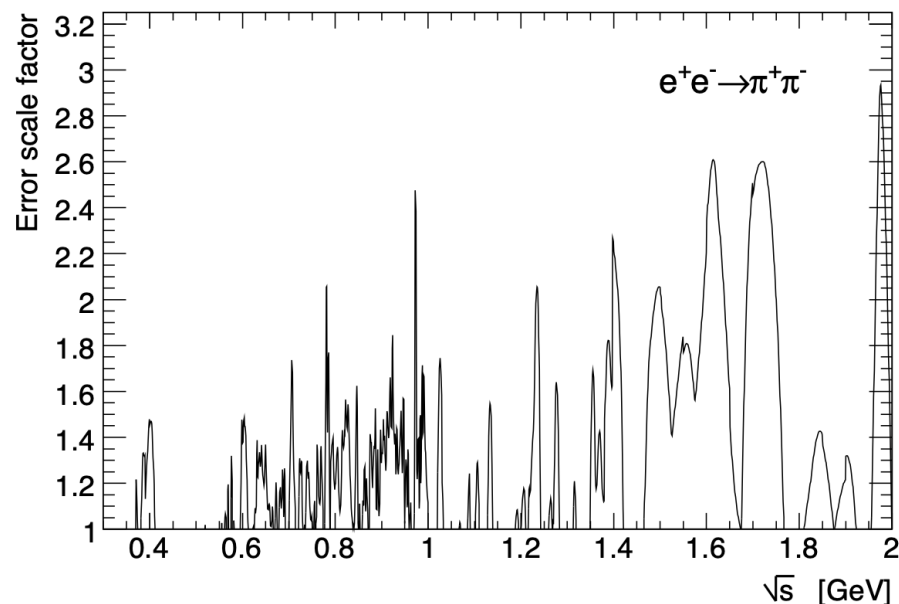
# Relative Weights and Tension

Figures from DHMZ, EPJC80 (2020) 241



In the energy range of [0.6-0.9] GeV, BABAR and KLOE dominate, elsewhere BABAR has better precision

Inconsistency between measurements reflected by large local scale factor (SF) defined by  $\sqrt{\chi^2/\text{dof}}$  a la PDG ( $\sim 15\%$  increase on  $\delta a_\mu^{\text{had}}$ )



# Fit at Low Energy Based on Analyticity and Unitarity

Pion form factor  $F_\pi^0$  extracted from  $\pi^+\pi^-$  bare cross sections as in [1810.00007]

$$|F_\pi^0|^2 = |G(s) \times J(s)|^2$$

$$G(s) = 1 + \alpha_V s + \frac{\kappa s}{m_\omega^2 - s - im_\omega \Gamma_\omega}$$

$$J(s) = e^{1 - \frac{\delta_1(s_0)}{\pi}} \left(1 - \frac{s}{s_0}\right)^{\left[1 - \frac{\delta_1(s_0)}{\pi}\right] \frac{s_0}{s}} \left(1 - \frac{s}{s_0}\right)^{-1} e^{\frac{s}{\pi} \int_{4m_\pi^2}^{s_0} dt \frac{\delta_1(t)}{t(t-s)}}$$

$$\cot \delta_1(s) = \frac{\sqrt{s}}{2k^3} (m_\rho^2 - s) \left[ \frac{2m_\pi^3}{m_\rho^2 \sqrt{s}} + B_0 + B_1 \omega(s) \right]$$

$$k = \frac{\sqrt{s - 4m_\pi^2}}{2}$$

$$\omega(s) = \frac{\sqrt{s} - \sqrt{s_0 - s}}{\sqrt{s} + \sqrt{s_0 - s}}$$

$$\sqrt{s_0} = 1.05 \text{ GeV}$$

$G(s)$  from [1611.09359]

$J(s)$  from [hep-ph/0106025, 0402285]

$\cot \delta_1(s)$  from [1102.2183]

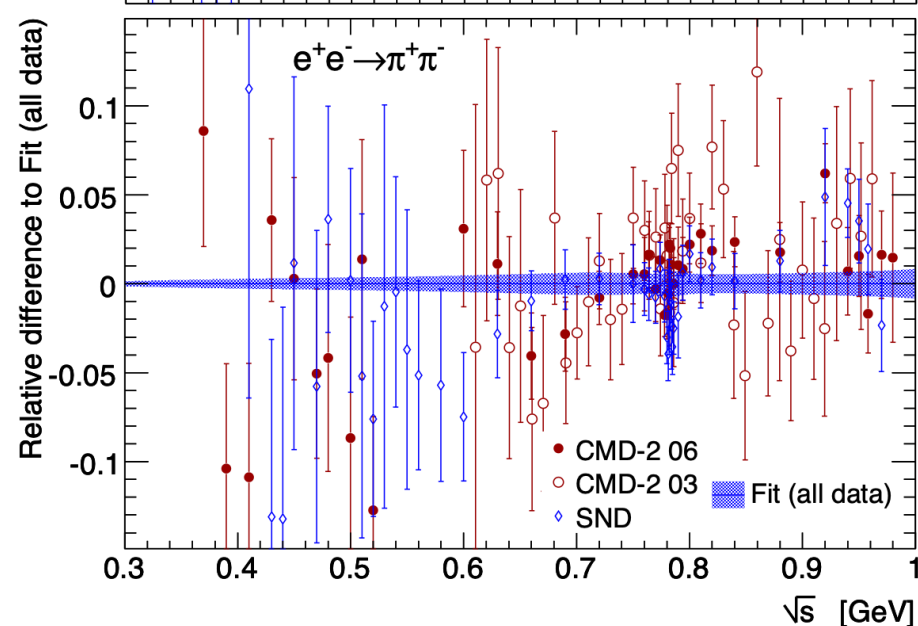
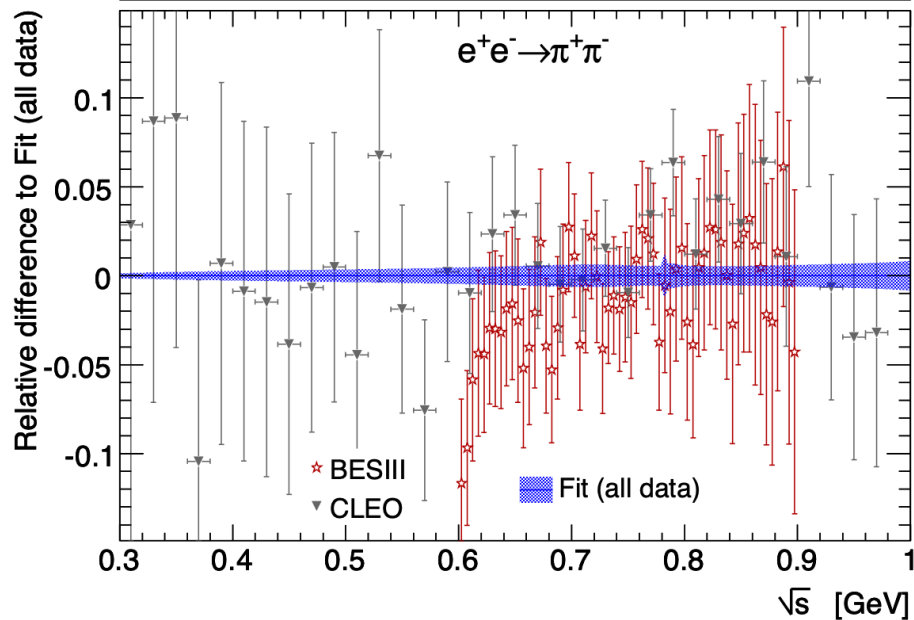
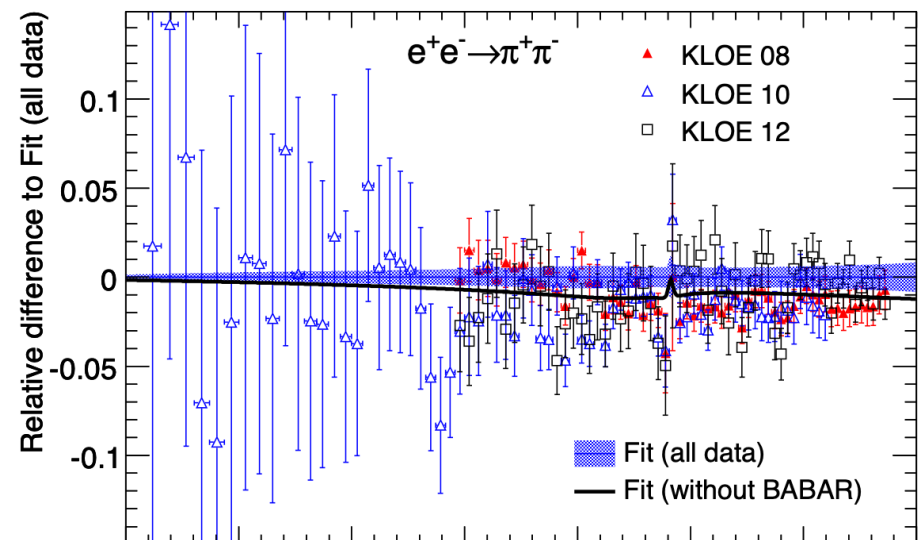
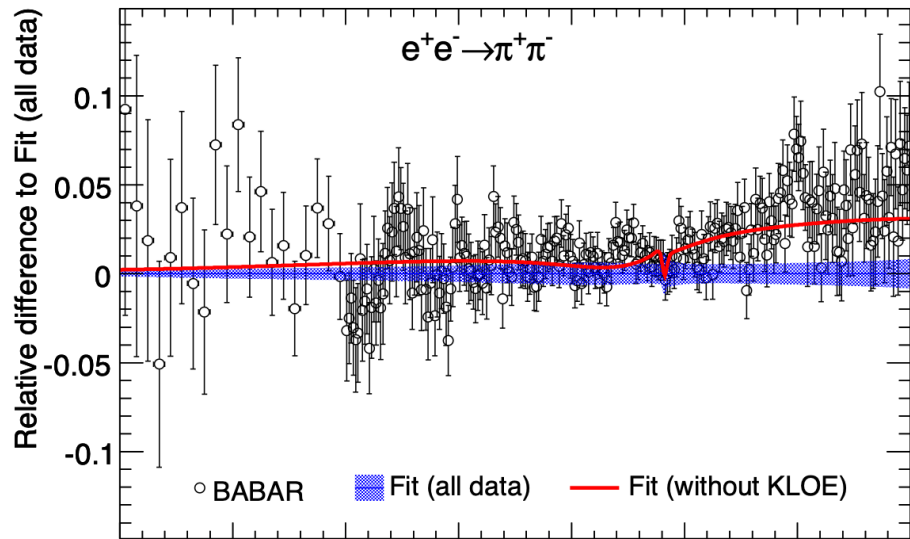
Six free parameters to fit:

$\alpha_V, \kappa, m_\omega, m_\rho, B_0, B_1$

( $\Gamma_\omega$  fixed to PDG value)

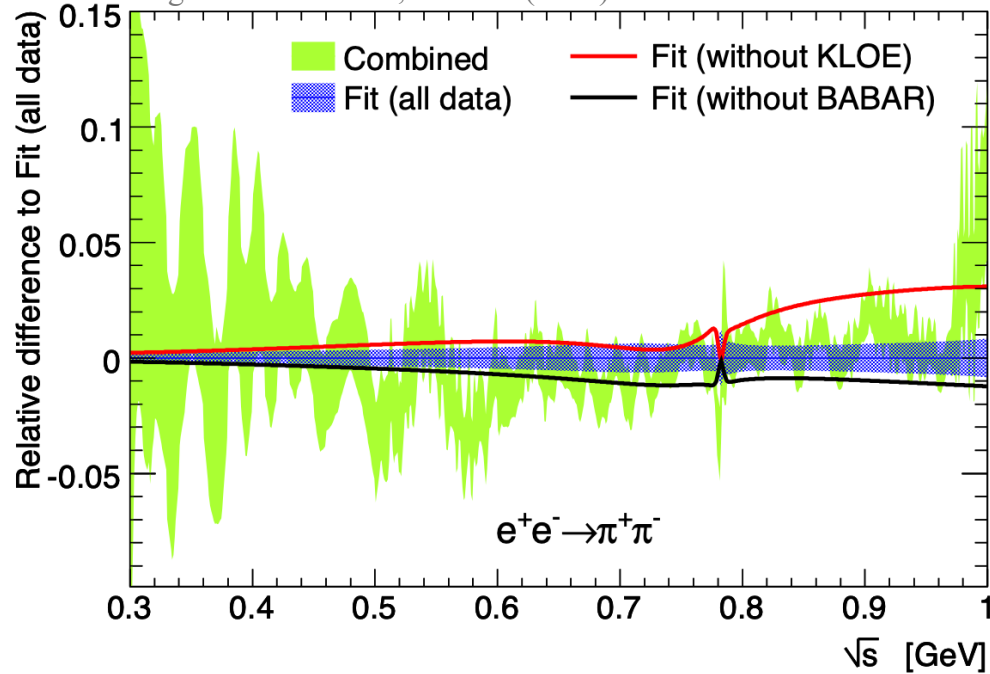
# Fit Performed to 1 GeV, Results Used to 0.6 GeV

Figures from DHMZ, EPJC80 (2020) 241



# Comparison Fit and Data Integration

Figure from DHMZ, EPJC80 (2020) 241



□ Use fit only below 0.6 GeV for  $a_{\mu}^{\text{had}}$  integral

- ▶ Where the data are less precise and scarce
- ▶ Less affected by potential uncertainties of inelastic contribution at high energy

$\sqrt{s}$ range [GeV]	$a_{\mu}^{\text{had}} [10^{-10}]$ Fit	$a_{\mu}^{\text{had}} [10^{-10}]$ Data Integration
0.3 - 0.6	$109.8 \pm 0.4_{\text{exp}} \pm 0.4_{\text{para}}^*$	$109.6 \pm 1.0_{\text{exp}}$

⇒ The difference  $0.2 \pm 0.8$  (correlation accounted for)

⇒ The fit improves the precision by  $\sim 2$

\* Parameter uncertainty corresponds to variations by removing the  $B_l$  term in the phase shift formula and by varying  $\sqrt{s_0}$  from 1.05 GeV to 1.3 GeV

# Combined Results Fit [ $<0.6$ GeV] + Data [0.6-1.8 GeV]

Take into account the correlation of 62% (based on pseudo-data samples) of the two regions

$\sqrt{s}$ range [GeV]	$a_{\mu}^{\text{had}} [10^{-10}]$ All data	$a_{\mu}^{\text{had}} [10^{-10}]$ All but BABAR	$a_{\mu}^{\text{had}} [10^{-10}]$ All but KLOE
threshold - 1.8	$506.9 \pm 1.9_{\text{total}}$	$505.0 \pm 2.1_{\text{total}}$	$510.6 \pm 2.2_{\text{total}}$

⇒ The difference “All but BABAR” and “All but KLOE” = 5.6 to be compared with 1.9 uncertainty with “All data”

- The local error inflation is not sufficient to amplify the uncertainty
- Global tension (normalisation/shape) not previously accounted for
- Potential underestimated uncertainty in at least one of the measurements?
- Other measurements not precise enough and are in agreement with BABAR or KLOE

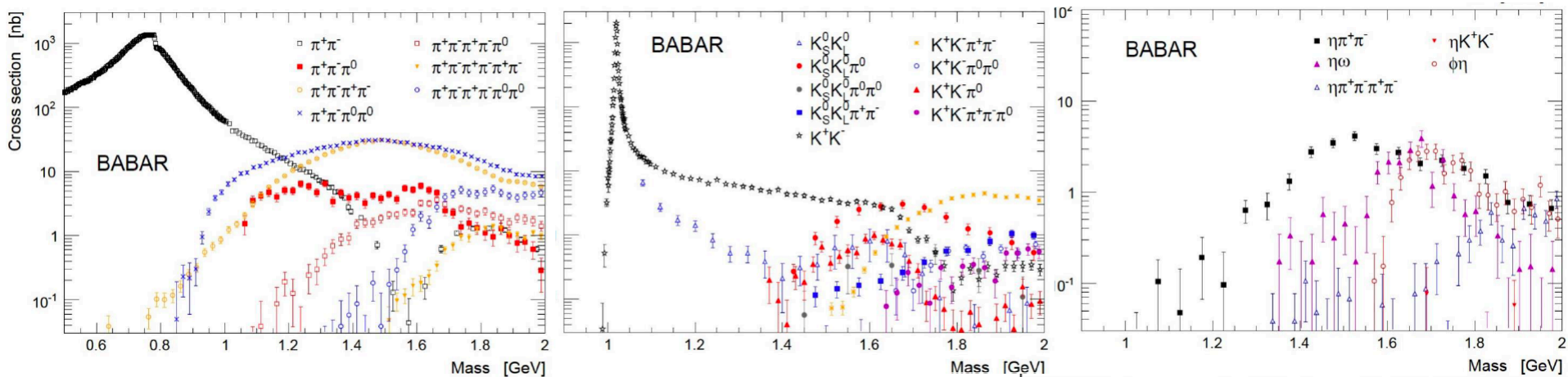
⇒ Given the fact we do not know which dataset is problematic, we decide to

- Add half of the discrepancy (2.8) as an additional uncertainty (correcting the local PDG inflation to avoid double counting)
- Take the mean value “All but BABAR” and “All but KLOE” as our central value

# Other Channels e.g. Those Measured by BABAR

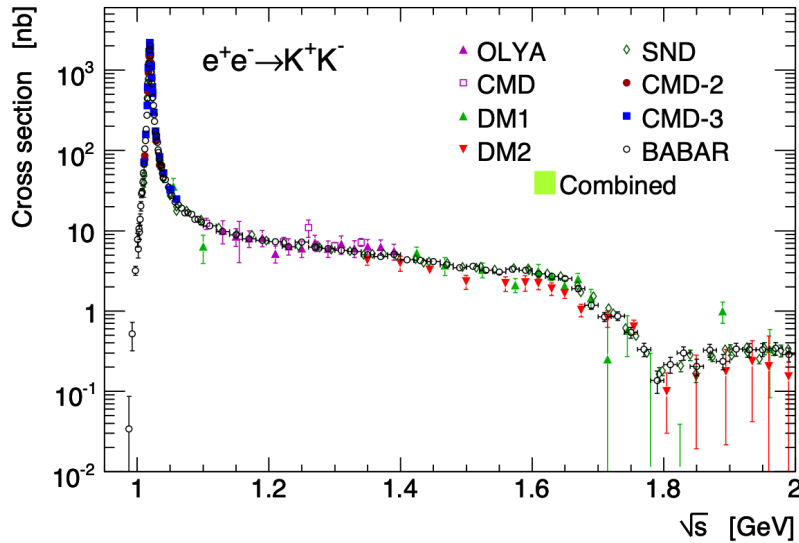
There are many exclusive channels ( $\sim 40$  processes) contributing to HVP

Here are some example measurements from BABAR

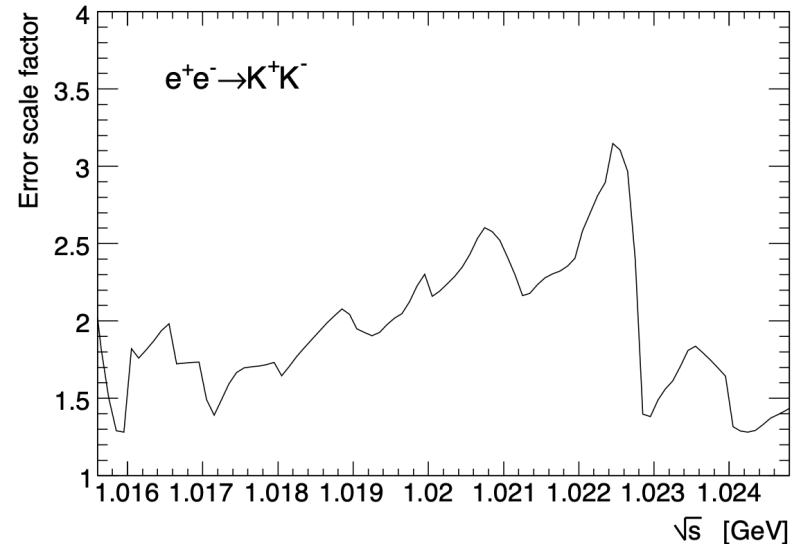
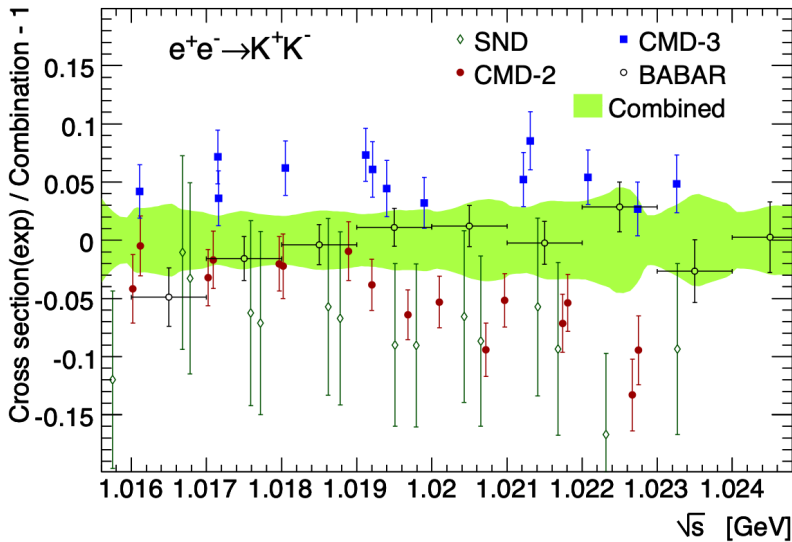


# Tension in Other Channel (e.g. KK)

Figures from DHMZ, EPJC80 (2020) 241

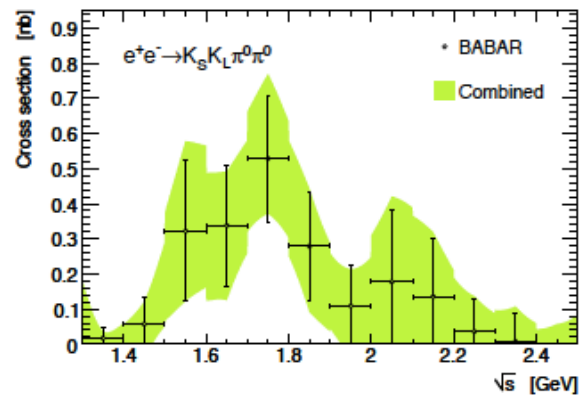
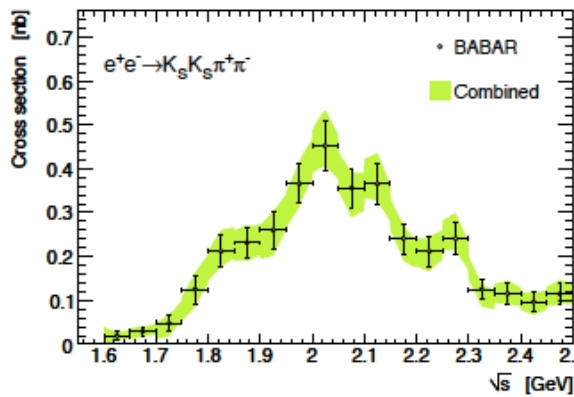
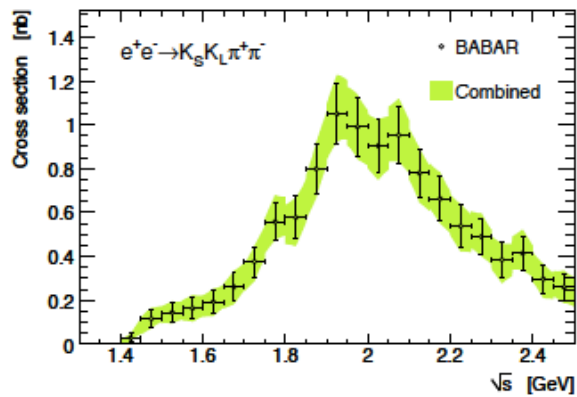
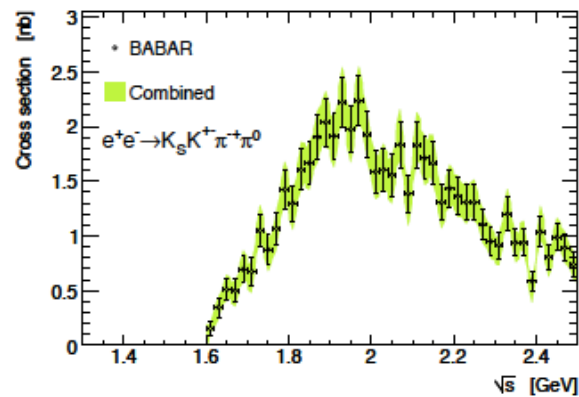
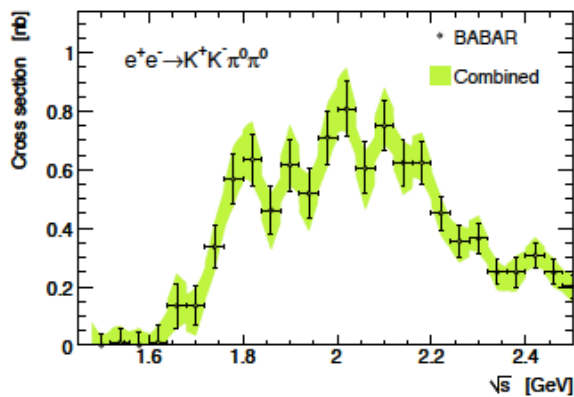
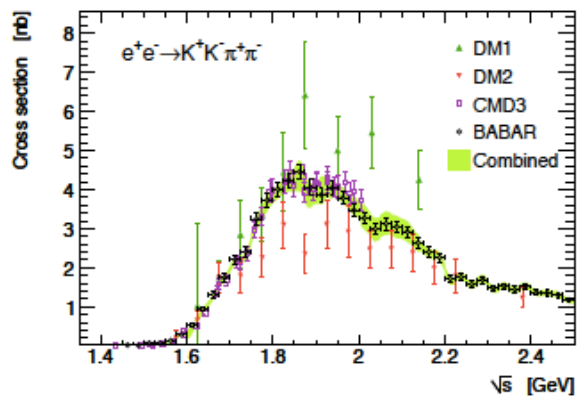
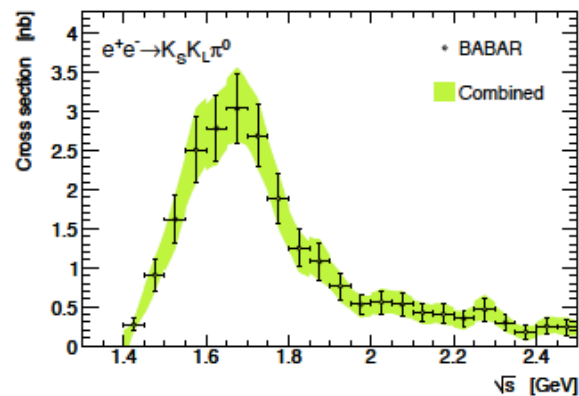
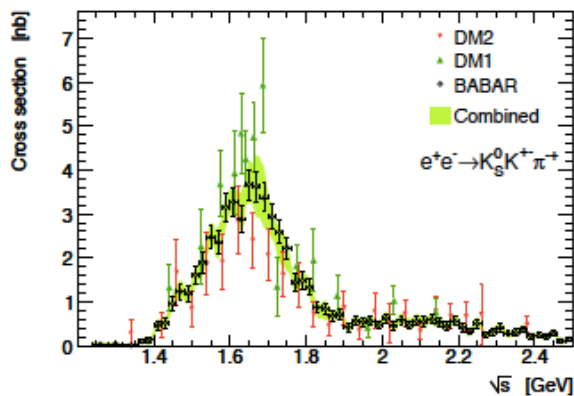
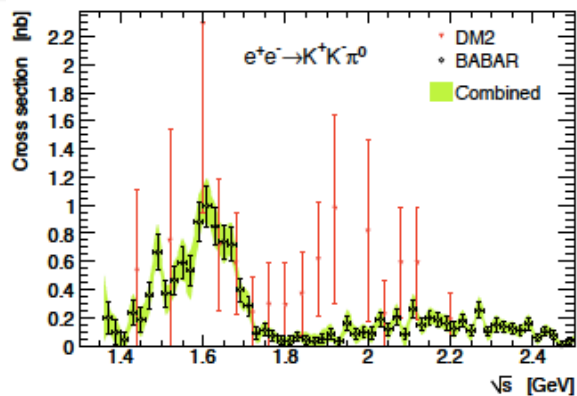


Several measurements with different precisions, CMD-2 and CMD-3 do not agree within the quoted uncertainties!  
→ Large error scaling factors  
though this channel only contributes 3.3% to LO HVP and 1.2% to uncertainty-squared.

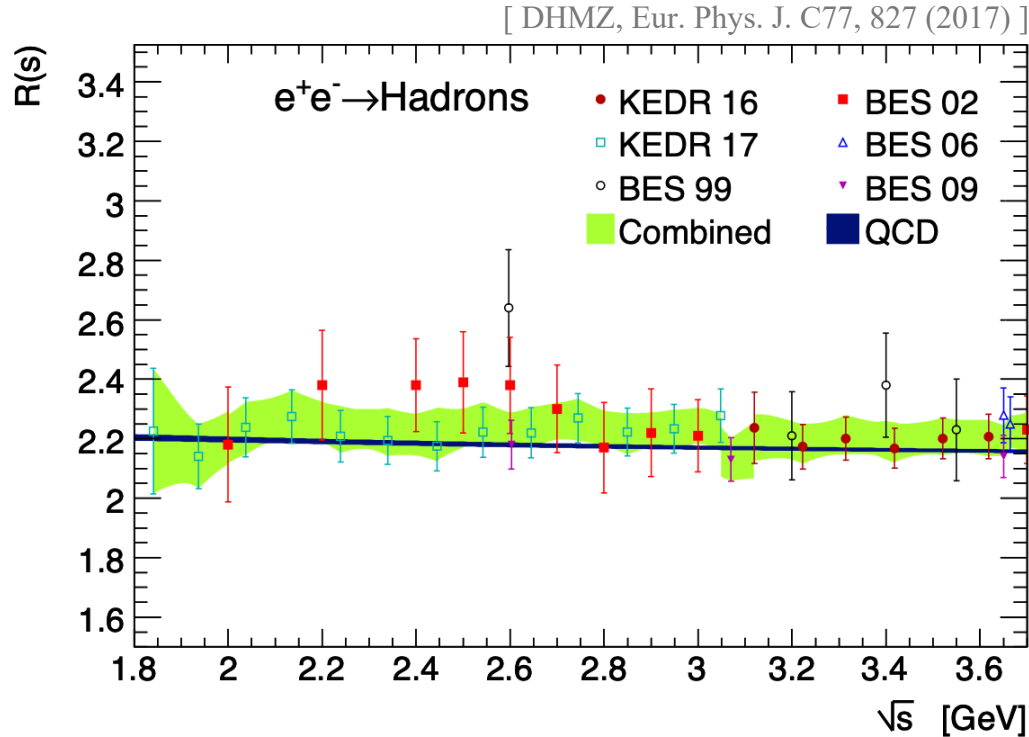




# KKbar+ $\pi$ 's Channels [ DHMZ, Eur. Phys. J. C77, 827 (2017) ]



# Contributions in the Region 1.8-3.7 GeV

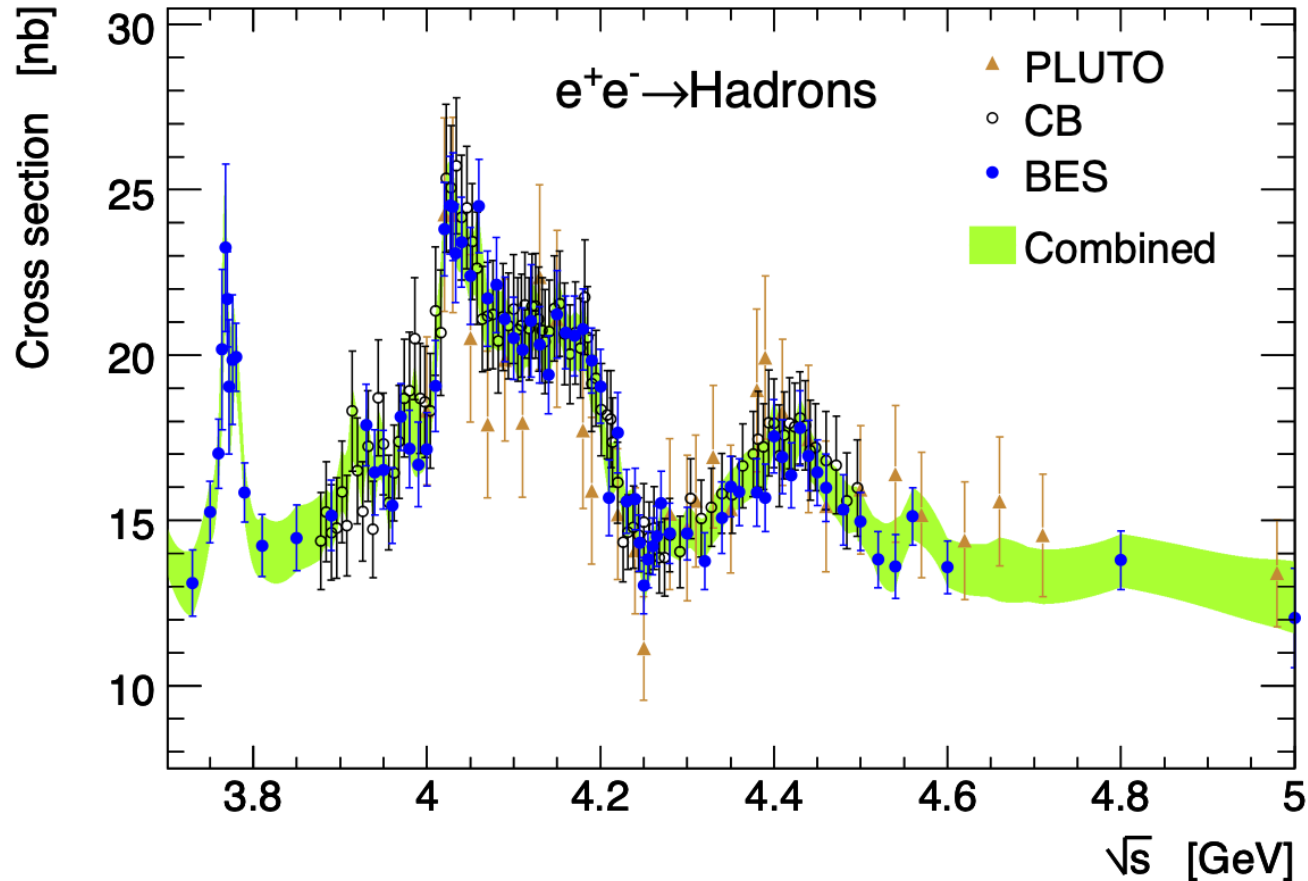


Energy range [GeV]	1.8 - 2.0 [2020]	2.0 - 3.7 [2017]
Data	$7.65 \pm 0.31$	$25.82 \pm 0.61$
pQCD	$8.30 \pm 0.09$	$25.15 \pm 0.19$
Difference	$0.65 \rightarrow \text{dual}$	agree $< 1\sigma$

pQCD evaluated from 4 loops +  $O(\alpha_s^2)$  quark mass corrections  
 Uncertainties:  $\alpha_s$ , truncation, FOPT/CIPT,  $m_q$

# Contributions from Charm Resonance Region

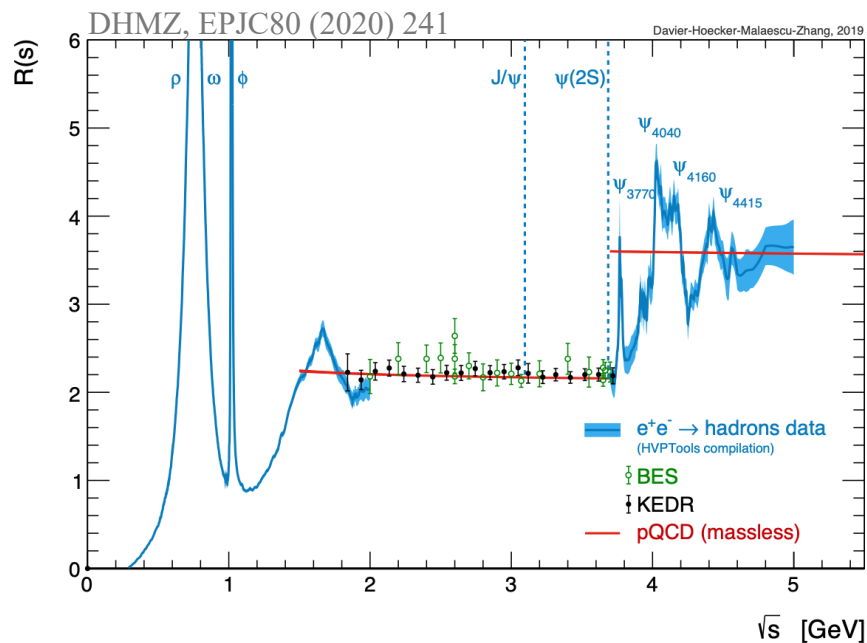
[ DHMZ, Eur. Phys. J. C77, 827 (2017) ]



$$7.29 \pm 0.05 \pm 0.30 \pm 0.00 \Rightarrow 1.05\% \text{ of } a_\mu^{\text{had, LO}}$$

stat    sys    cor

# A Big Picture in Terms of R(s)

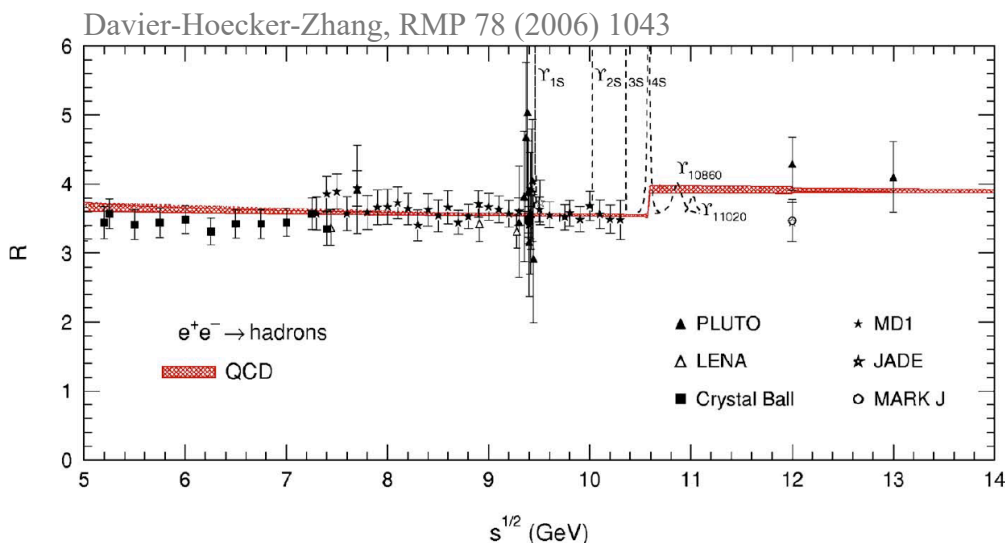


- $[\pi^0\gamma-1.8\text{GeV}]$ 
  - sum about 22  $\rightarrow$  37 exclusive channels
  - estimate unmeasured channels using isospin relations (now  $< 0.1\%$ )

- $[1.8-3.7]$  GeV
  - good agreement between data and pQCD calculation  $\rightarrow$  use 4-loop pQCD
  - $J/\psi, \psi(2s)$ : Breit-Wigner integral

- $[3.7-5]$  GeV
  - use data

- $>5\text{GeV}$ 
  - use 4-loop pQCD calculation



# Overall Results

Channel	$a_\mu^{\text{had, LO}} [10^{-10}]$	$\Delta\alpha(m_Z^2) [10^{-4}]$
$\pi^0\gamma$	$4.29 \pm 0.06 \pm 0.04 \pm 0.07$	$0.35 \pm 0.00 \pm 0.00 \pm 0.01$
$\eta\gamma$	$0.65 \pm 0.02 \pm 0.01 \pm 0.01$	$0.08 \pm 0.00 \pm 0.00 \pm 0.00$
$\pi^+\pi^-$	$507.80 \pm 0.83 \pm 3.19 \pm 0.60$	$34.49 \pm 0.06 \pm 0.20 \pm 0.04$
$\pi^+\pi^-\pi^0$	$46.20 \pm 0.40 \pm 1.10 \pm 0.86$	$4.60 \pm 0.04 \pm 0.11 \pm 0.08$
$2\pi^+2\pi^-$	$13.68 \pm 0.03 \pm 0.27 \pm 0.14$	$3.58 \pm 0.01 \pm 0.07 \pm 0.03$
$\pi^+\pi^-2\pi^0$	$18.03 \pm 0.06 \pm 0.48 \pm 0.26$	$4.45 \pm 0.02 \pm 0.12 \pm 0.07$
$2\pi^+2\pi^-\pi^0$ ( $\eta$ excl.)	$0.69 \pm 0.04 \pm 0.06 \pm 0.03$	$0.21 \pm 0.01 \pm 0.02 \pm 0.01$
$\pi^+\pi^-3\pi^0$ ( $\eta$ excl.)	$0.49 \pm 0.03 \pm 0.09 \pm 0.00$	$0.15 \pm 0.01 \pm 0.03 \pm 0.00$
$3\pi^+3\pi^-$	$0.11 \pm 0.00 \pm 0.01 \pm 0.00$	$0.04 \pm 0.00 \pm 0.00 \pm 0.00$
$2\pi^+2\pi^-2\pi^0$ ( $\eta$ excl.)	$0.71 \pm 0.06 \pm 0.07 \pm 0.14$	$0.25 \pm 0.02 \pm 0.02 \pm 0.05$
$\pi^+\pi^-4\pi^0$ ( $\eta$ excl., isospin)	$0.08 \pm 0.01 \pm 0.08 \pm 0.00$	$0.03 \pm 0.00 \pm 0.03 \pm 0.00$
$\eta\pi^+\pi^-$	$1.19 \pm 0.02 \pm 0.04 \pm 0.02$	$0.35 \pm 0.01 \pm 0.01 \pm 0.01$
$\eta\omega$	$0.35 \pm 0.01 \pm 0.02 \pm 0.01$	$0.11 \pm 0.00 \pm 0.01 \pm 0.00$
$\eta\pi^+\pi^-\pi^0$ (non- $\omega, \phi$ )	$0.34 \pm 0.03 \pm 0.03 \pm 0.04$	$0.12 \pm 0.01 \pm 0.01 \pm 0.01$
$\eta 2\pi^+2\pi^-$	$0.02 \pm 0.01 \pm 0.00 \pm 0.00$	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega\eta\pi^0$	$0.06 \pm 0.01 \pm 0.01 \pm 0.00$	$0.02 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega\pi^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.94 \pm 0.01 \pm 0.03 \pm 0.00$	$0.20 \pm 0.00 \pm 0.01 \pm 0.00$
$\omega(\pi\pi)^0$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.07 \pm 0.00 \pm 0.00 \pm 0.00$	$0.02 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega$ (non- $3\pi, \pi\gamma, \eta\gamma$ )	$0.04 \pm 0.00 \pm 0.00 \pm 0.00$	$0.00 \pm 0.00 \pm 0.00 \pm 0.00$
$K^+K^-$	$23.08 \pm 0.20 \pm 0.33 \pm 0.21$	$3.35 \pm 0.03 \pm 0.05 \pm 0.03$
$K_S K_L$	$12.82 \pm 0.06 \pm 0.18 \pm 0.15$	$1.74 \pm 0.01 \pm 0.03 \pm 0.02$
$\phi$ (non- $K\bar{K}, 3\pi, \pi\gamma, \eta\gamma$ )	$0.05 \pm 0.00 \pm 0.00 \pm 0.00$	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$K\bar{K}\pi$	$2.45 \pm 0.05 \pm 0.10 \pm 0.06$	$0.78 \pm 0.02 \pm 0.03 \pm 0.02$
$K\bar{K}2\pi$	$0.85 \pm 0.02 \pm 0.05 \pm 0.01$	$0.30 \pm 0.01 \pm 0.02 \pm 0.00$
$K\bar{K}3\pi$ (estimate)	$-0.02 \pm 0.01 \pm 0.01 \pm 0.00$	$-0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$\eta\phi$	$0.33 \pm 0.01 \pm 0.01 \pm 0.00$	$0.11 \pm 0.00 \pm 0.00 \pm 0.00$
$\eta K\bar{K}$ (non- $\phi$ )	$0.01 \pm 0.01 \pm 0.01 \pm 0.00$	$0.00 \pm 0.00 \pm 0.01 \pm 0.00$
$\omega K\bar{K}$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$	$0.00 \pm 0.00 \pm 0.00 \pm 0.00$
$\omega 3\pi$ ( $\omega \rightarrow \pi^0\gamma$ )	$0.06 \pm 0.01 \pm 0.01 \pm 0.01$	$0.02 \pm 0.00 \pm 0.00 \pm 0.00$
$7\pi$ ( $3\pi^+3\pi^-\pi^0$ + estimate)	$0.02 \pm 0.00 \pm 0.01 \pm 0.00$	$0.01 \pm 0.00 \pm 0.00 \pm 0.00$
$J/\psi$ (BW integral)	$6.28 \pm 0.07$	$7.09 \pm 0.08$
$\psi(2S)$ (BW integral)	$1.57 \pm 0.03$	$2.50 \pm 0.04$
$R$ data [3.7 – 5.0] GeV	$7.29 \pm 0.05 \pm 0.30 \pm 0.00$	$15.79 \pm 0.12 \pm 0.66 \pm 0.00$
$R_{\text{QCD}}$ [1.8 – 3.7 GeV] <sub>uds</sub>	$33.45 \pm 0.28 \pm 0.65_{\text{dual}}$	$24.27 \pm 0.18 \pm 0.28_{\text{dual}}$
$R_{\text{QCD}}$ [5.0 – 9.3 GeV] <sub>udsc</sub>	$6.86 \pm 0.04$	$34.89 \pm 0.17$
$R_{\text{QCD}}$ [9.3 – 12.0 GeV] <sub>udscb</sub>	$1.21 \pm 0.01$	$15.56 \pm 0.04$
$R_{\text{QCD}}$ [12.0 – 40.0 GeV] <sub>udscb</sub>	$1.64 \pm 0.00$	$77.94 \pm 0.12$
$R_{\text{QCD}}$ [ $> 40.0$ GeV] <sub>udscb</sub>	$0.16 \pm 0.00$	$42.70 \pm 0.06$
$R_{\text{QCD}}$ [ $> 40.0$ GeV] <sub>t</sub>	$0.00 \pm 0.00$	$-0.72 \pm 0.01$
<b>Sum</b>	<b><math>693.9 \pm 1.0 \pm 3.4 \pm 1.6 \pm 0.1_\psi \pm 0.7_{\text{QCD}}</math></b>	<b><math>275.42 \pm 0.15 \pm 0.72 \pm 0.23 \pm 0.09_\psi \pm 0.55_{\text{QCD}}</math></b>

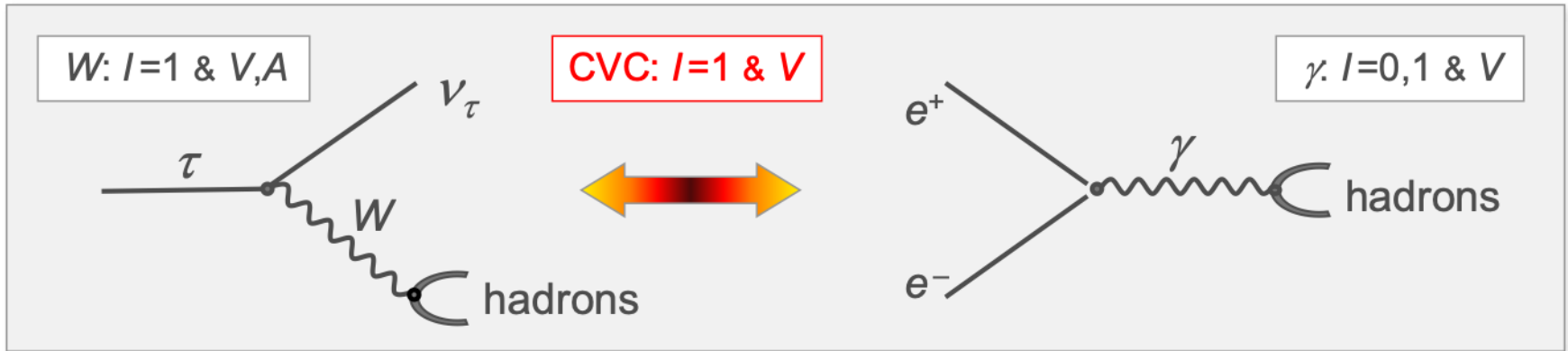
More than 30  
exclusive  
channels  
( $< 1.8$  GeV)  
evaluated

Estimation for  
missing modes  
based on isospin  
constraints  
becomes  
negligible  
(0.016%)

Table taken from  
DHMZ, EPJC80  
(2020) 241

# An Alternative Way Used to Evaluate HVP

Proposed by Alemany-Davier-Hoecker, EPJC 2 (1998) 123



Hadronic physics factorises in **Spectral Functions**:

Isospin symmetry connects  $I=1$   $e^+e^-$  cross section to vector  $\tau$  spectral functions

**Fundamental ingredient relating long distance (resonances) to short distance description (QCD)**

$$\sigma^{(I=1)}[e^+e^- \rightarrow \pi^+\pi^-] = \frac{4\pi\alpha^2}{s} \nu[\tau^- \rightarrow \pi^-\pi^0\nu_\tau]$$

$$\nu[\tau^- \rightarrow \pi^-\pi^0\nu_\tau] \propto \frac{\text{BR}[\tau^- \rightarrow \pi^-\pi^0\nu_\tau]}{\text{BR}[\tau^- \rightarrow e^-\bar{\nu}_e\nu_\tau]} \cdot \frac{1}{N_{\pi\pi^0}} \frac{dN_{\pi\pi^0}}{ds} \cdot \frac{m_\tau^2}{(1-s/m_\tau^2)^2 (1+s/m_\tau^2)}$$

Branching fractions

Mass spectrum

Kinematic factors (PS)

# Known Isospin Breaking Corrections

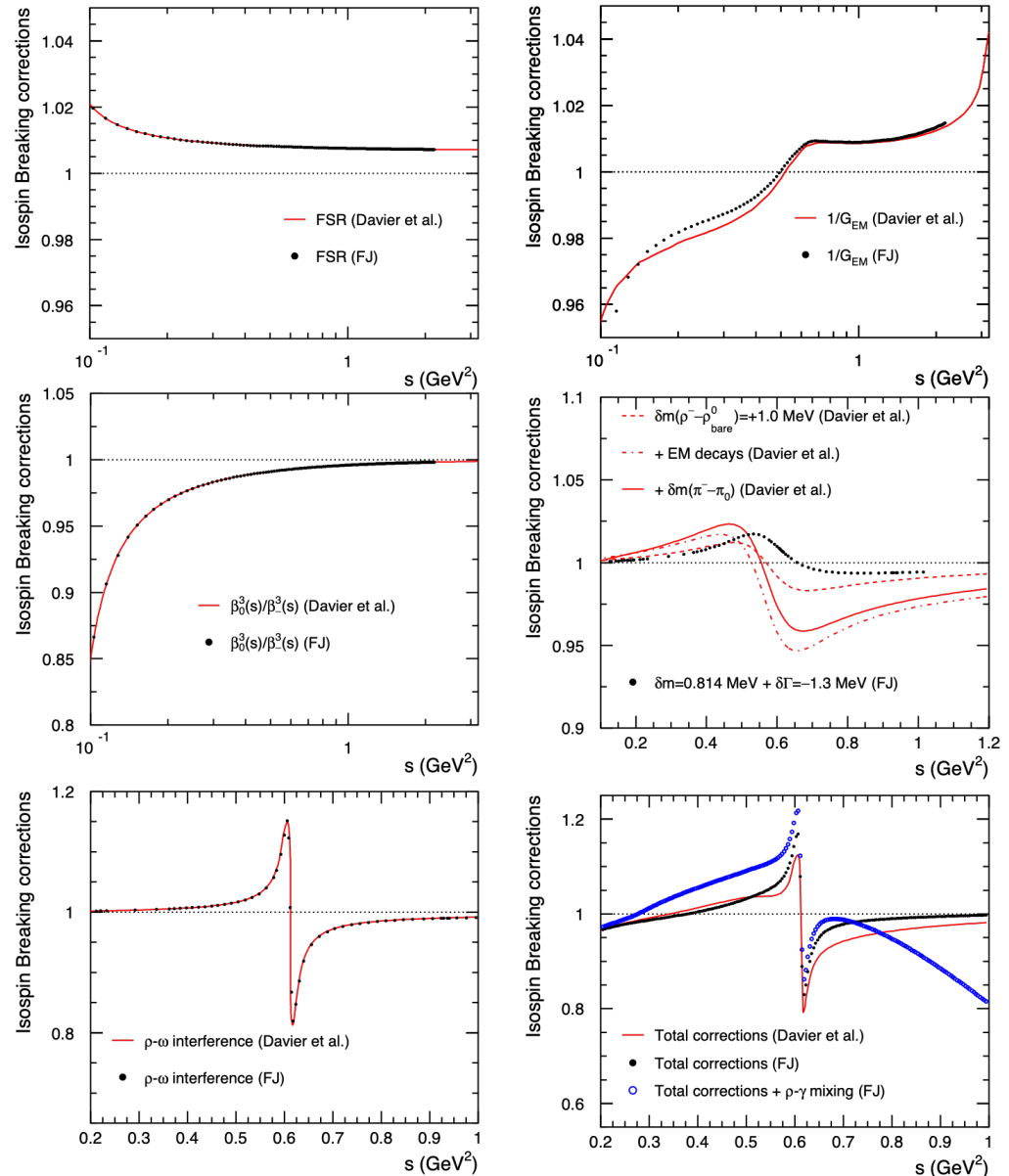
Davier et al., EPJC66 (2010) 127

$$v_{1,X^-}(s) = \frac{m_\tau^2}{6|V_{ud}|^2} \frac{B_{X^-}}{B_e} \frac{1}{N_X} \frac{dN_X}{ds} \times \left(1 - \frac{s}{m_\tau^2}\right)^{-2} \left(1 + \frac{2s}{m_\tau^2}\right)^{-1} \frac{R_{IB}(s)}{S_{EW}},$$

$$R_{IB}(s) = \frac{FSR(s)}{G_{EM}(s)} \frac{\beta_0^3(s)}{\beta_-^3(s)} \left| \frac{F_0(s)}{F_-(s)} \right|^2.$$

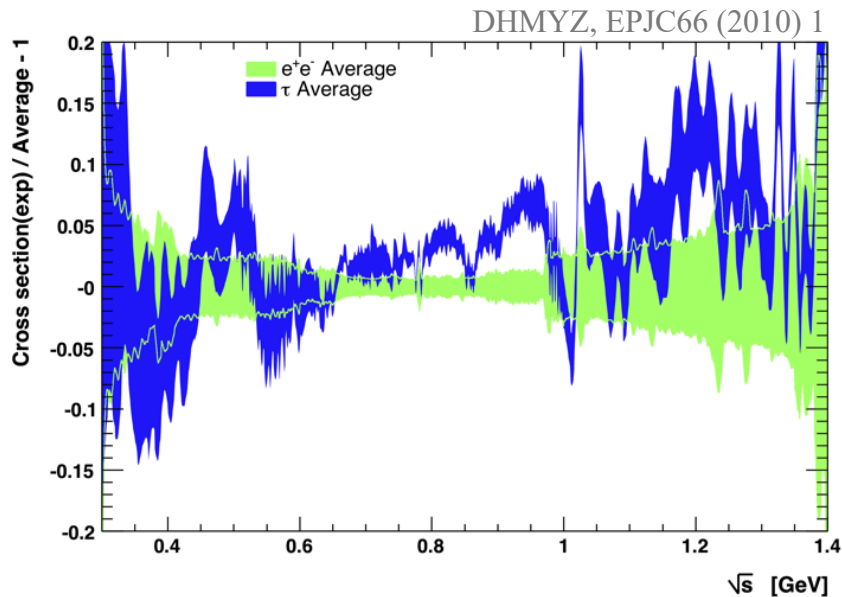
Good agreement between Davier et al. and FJ for most of the isospin breaking components

Figure 19 from WP20  
Studies initiated in Davier et al., EPJC66 (2010) 127

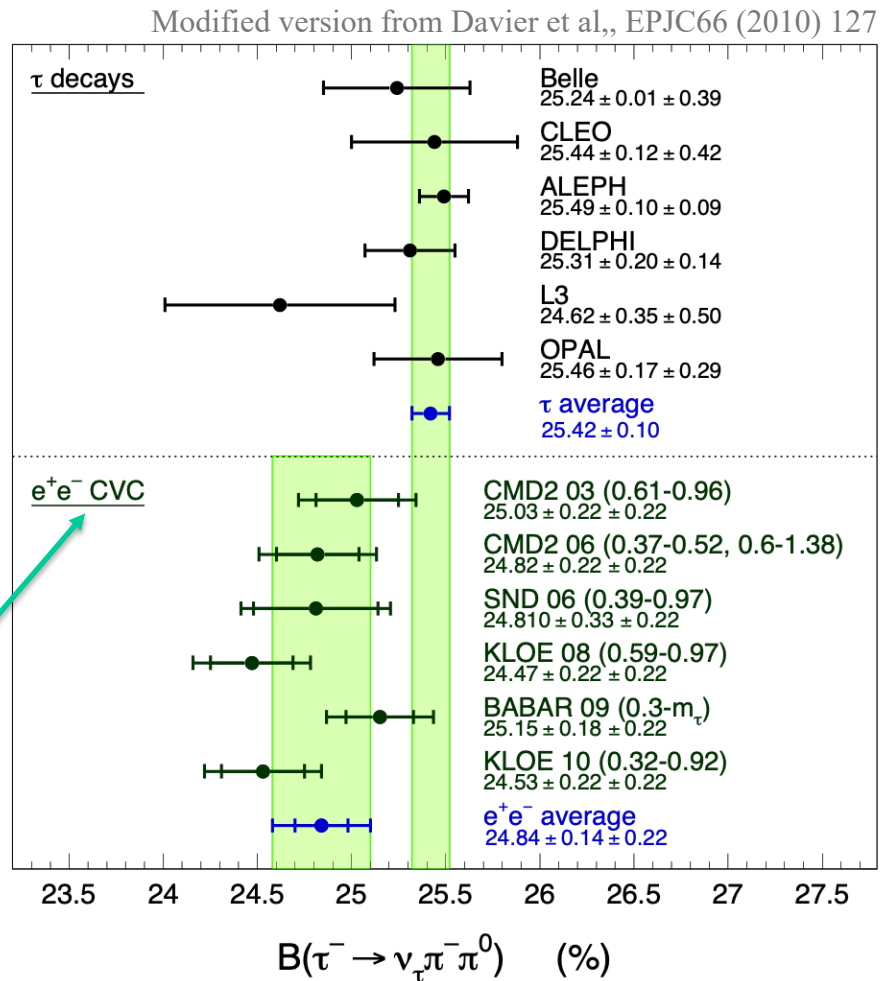


# Open Issue in $2\pi$ Channel

Take into account all known isospin breaking corrections except for the  $\rho$ - $\gamma$  mixing correction



$$B_X^{\text{CVC}} = \frac{3}{2} \frac{B_e |V_{ud}|^2}{\pi \alpha^2 m_\tau^2} \int_{s_{\min}}^{m_\tau^2} ds s \sigma_{X^0} \left(1 - \frac{s}{m_\tau^2}\right)^2 \left(1 + \frac{2s}{m_\tau^2}\right)$$

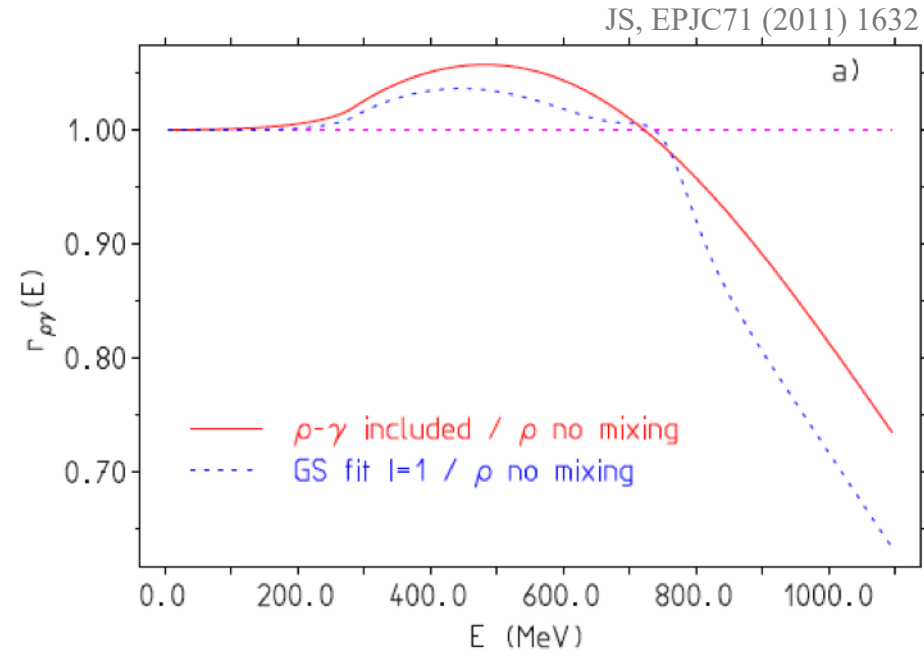
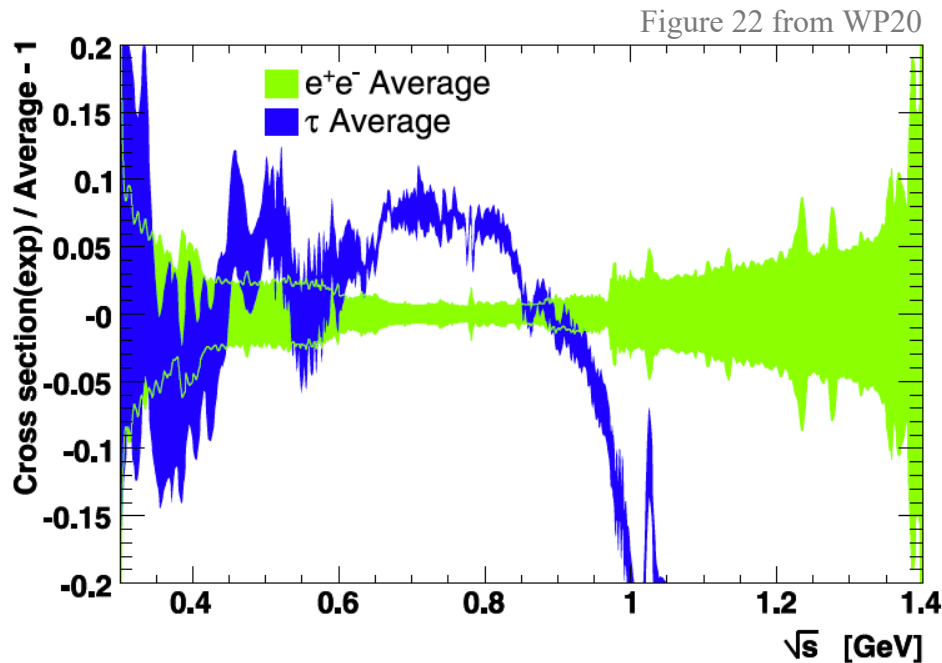


Clear difference in shape and in BR between  $e^+e^-$  and  $\tau$  average



# Additional EFT Based $\rho$ - $\gamma$ Mixing Correction

Jegerlehner and Szafron proposed to use the missing  $\rho$ - $\gamma$  mixing in  $\tau$  data to explain the remaining  $e^+e^-$  and  $\tau$  difference

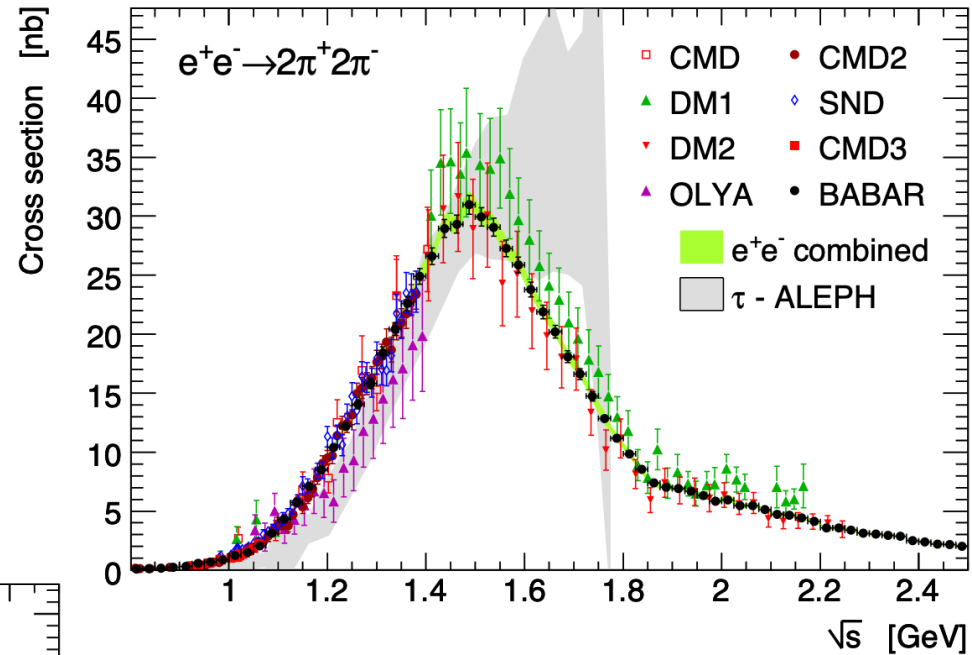
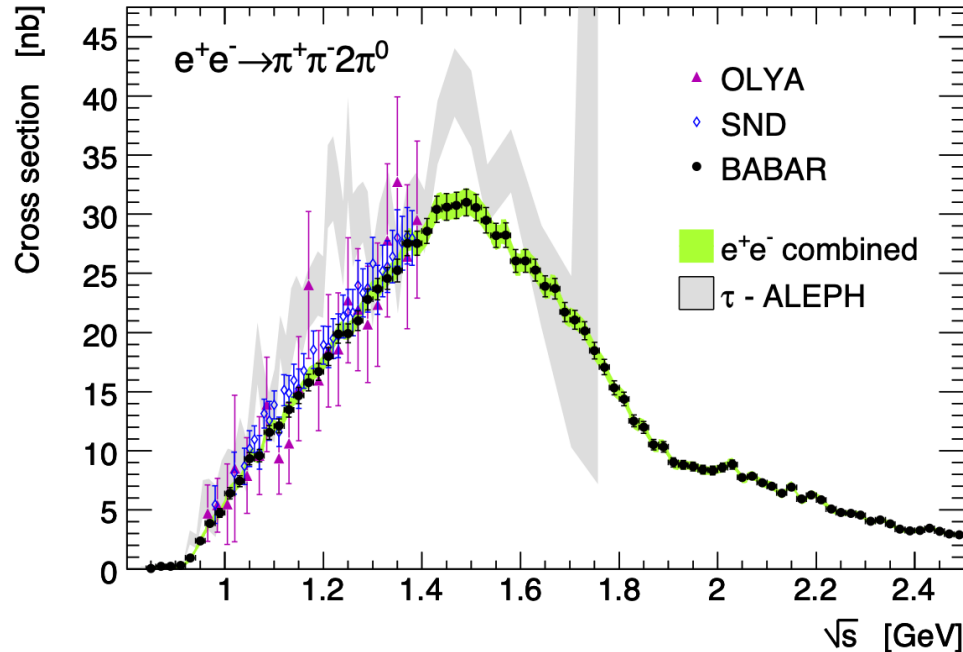


Applying the  $\rho$ - $\gamma$  mixing correction makes the  $e^+e^-$  and  $\tau$  difference worse in some of the mass range

# Comparison of $4\pi$ Channels

The precision of  $e^+e^-$  data increased over time, a factor of 1.7-2.3 between 2011 and 2017

Figures from DHMZ, *EPJC* 77 (2017) 827



In comparison,  $\tau$  data are now less precise

We no longer pursue the  $\tau$  data-based evaluation due to the open issue in the  $2\pi$  + less precise  $\tau$  data

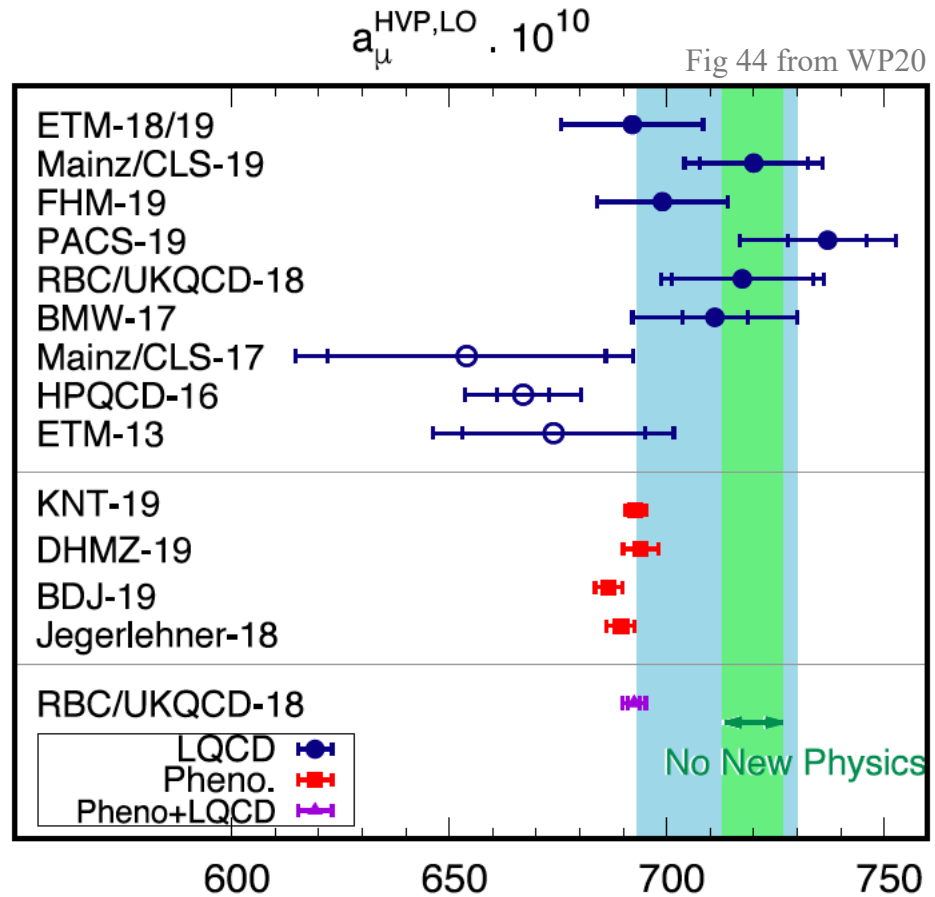
# Alternative LO HVP Prediction from Lattice QCD

Lattice QCD allows to directly compute the real part of the two-point correlation function without invoking the resonances occurring on the imaginary axis

Several groups provided predictions, however their uncertainties are still large so that they were not used in providing the LO HVP prediction in the WP20

Recently, BMW has provided a prediction varying  
from  $v1: 712.4 (4.5) \times 10^{-10}$   
to  $v3: 707.5 (5.5) \times 10^{-10}$   
reaching 0.8% close to 0.6% (dispersive).

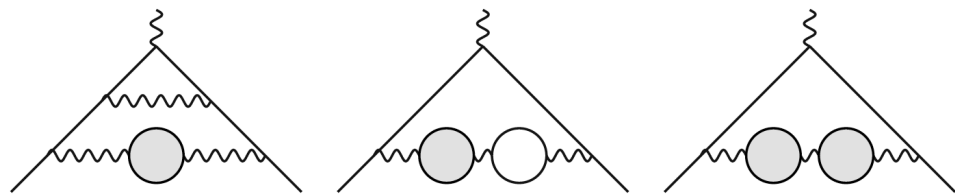
This prediction needs to be confirmed by other lattice groups with comparable precision



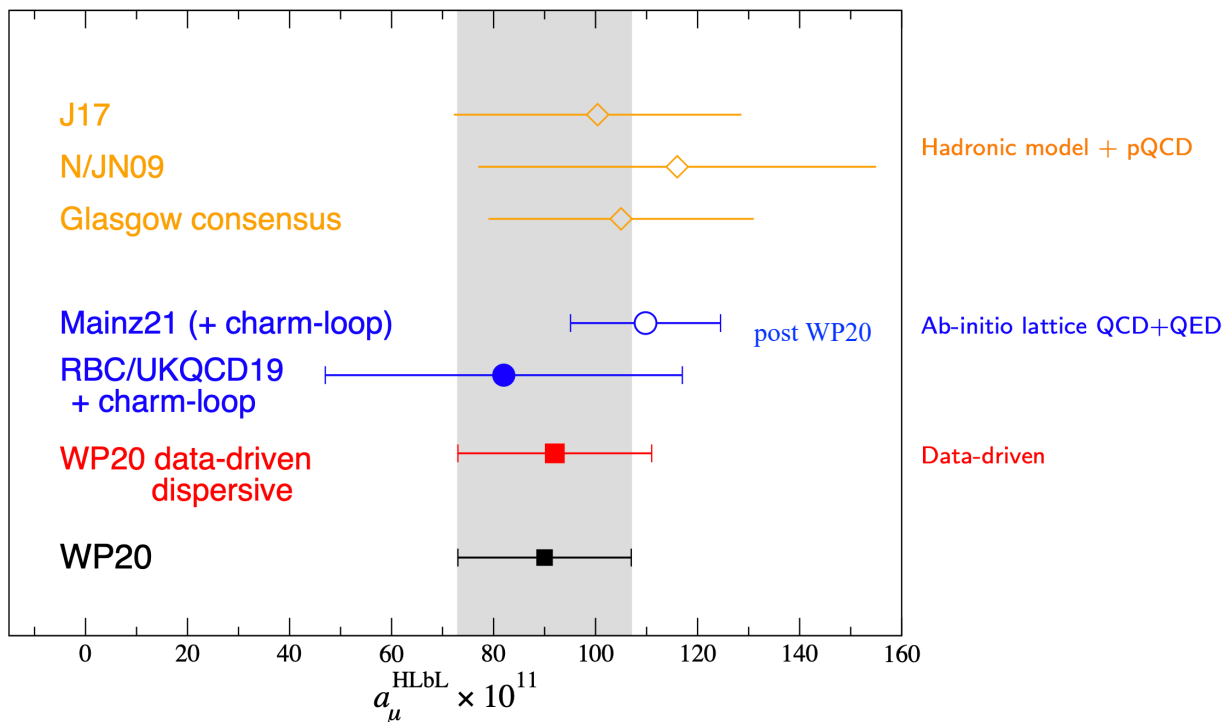
# High Order HVP + HLbL Predictions

$$a_{\mu}^{\text{HVP, NLO}} = -9.83(7) \times 10^{-10}$$

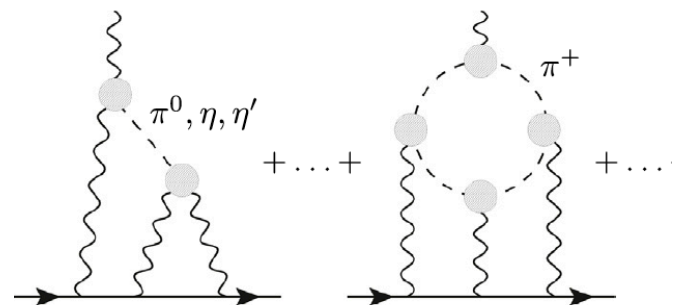
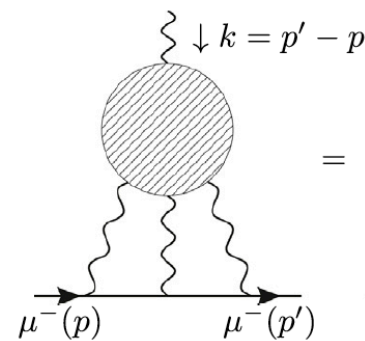
$$a_{\mu}^{\text{HVP, NNLO}} = 1.24(1) \times 10^{-10}$$



## Status of hadronic light-by-light contribution



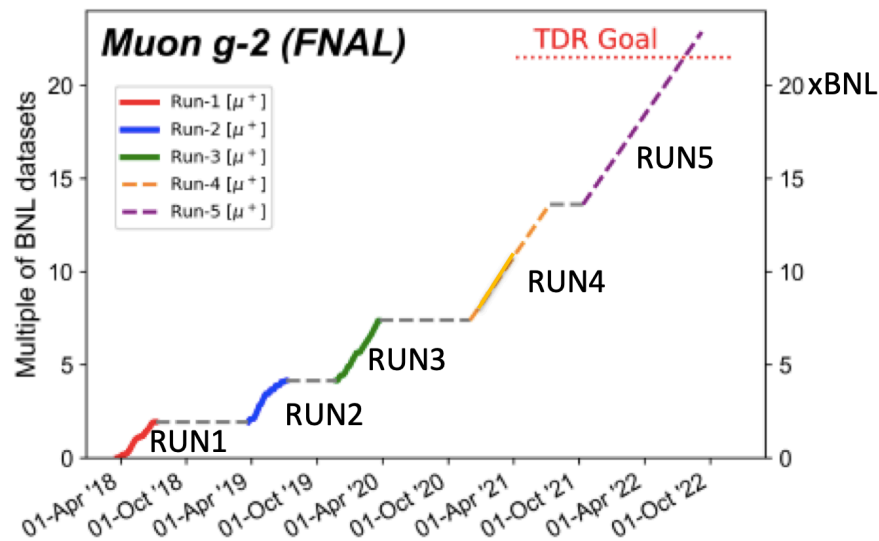
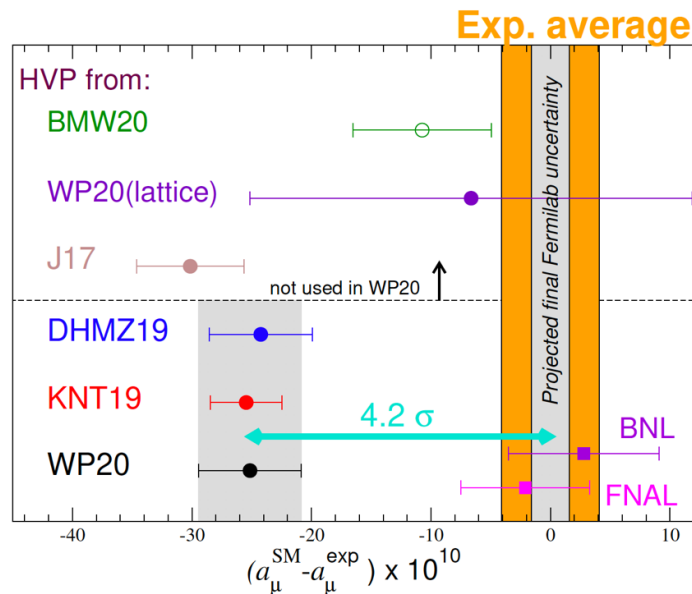
Values, graphs from WP20  
Left figure from C. Lehner's  
CERN seminar talk



$$a_{\mu}^{\text{HLbL}}(\text{phenomenology} + \text{lattice QCD}) + a_{\mu}^{\text{HLbL, NLO}} = 92(18) \times 10^{-11}$$

# Summary and Perspectives

- The current precision of the dispersive prediction  $\sim$  direct measurement
- The latter will be improved in the next years by a factor  $\sim 4$
- On the theory side, the situation is less clear
  - We do expect more measurements (e.g. in  $2\pi$ ) from BABAR, Belle 2, BESIII, CMD-3, ...
  - However the measurements have to be very precise in order to resolve the BABAR-KLOE discrepancy which prevented us from improving further the precision in the data combination
  - If the BMW prediction is confirmed, we also need to understand the difference between the dispersive and lattice predictions



# List of publications

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1. ADH 1998, [Eur.Phys.J.C 2 \(1998\) 123](#) [330 citations\*]
  2. DH 1998, [Phys.Lett.B 419 \(1998\) 419](#) [219 citations]
  3. DH 1998, [Phys.Lett.B 435 \(1998\) 427](#) [292 citations]
  4. DEHZ 2003, [Eur.Phys.J.C 27 \(2003\) 497](#) [394 citations]
  5. DEHZ 2003, [Eur.Phys.J.C 31 \(2003\) 503](#) [430 citations]
  6. DHMZ+ 2010, [Eur.Phys.J.C 66 \(2010\) 127](#) [157 citations]
  7. DHMYZ 2010, [Eur.Phys.J.C 66 \(2010\) 1](#) [209 citations]
  8. DHMZ 2011, [Eur.Phys.J.C 71 \(2011\) 1515](#) [866 citations]
  9. DHMZ 2017, [Eur.Phys.J.C 77 \(2017\) 827](#) [259 citations]
  10. DHMZ 2019, [Eur.Phys.J.C 80 \(2020\) 241](#) [169 citations]
  11. Theory initiative WP 2020, [Phys.Rept. 887 \(2020\) 1](#) [171 citations]
- Total number of citations: ~3500

\* Status of April 9, 2021

# Efforts on the prediction side

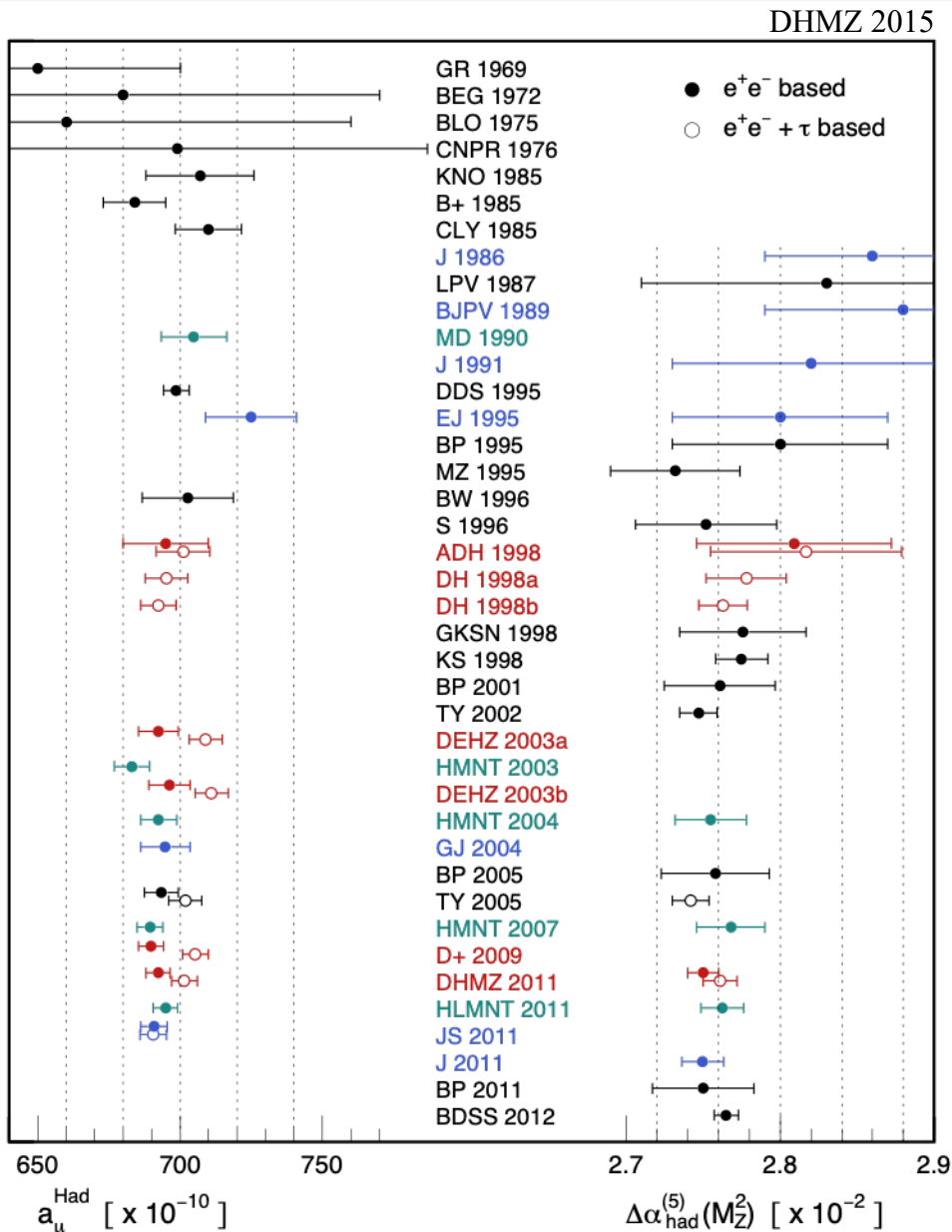


Fig. 3 prepared by Davier, Hoecker, Malaescu, Zhang for “Standard theory essays in the 60th anniversary of CERN”

Our efforts started by Michel with a first publication with Andreas and Richard in 1998 (ADH 1998)

Since then ZZ(02), Bogdan(09) and others joined the efforts

In total we have published 10 highly cited articles ([link](#))

Our prediction has been one main reference used for comparing with the direct measurement

The precision of the HVP prediction is data-driven

It depends on

- the precision of  $e^+e^-$  annihilation (& tau) data
- state of the art techniques (HVPTools) for data interpolation, combination and error correlation treatment