



# Combined fitting of the weak mixing angle and the proton structure parameters using the forward-backward charge asymmetry

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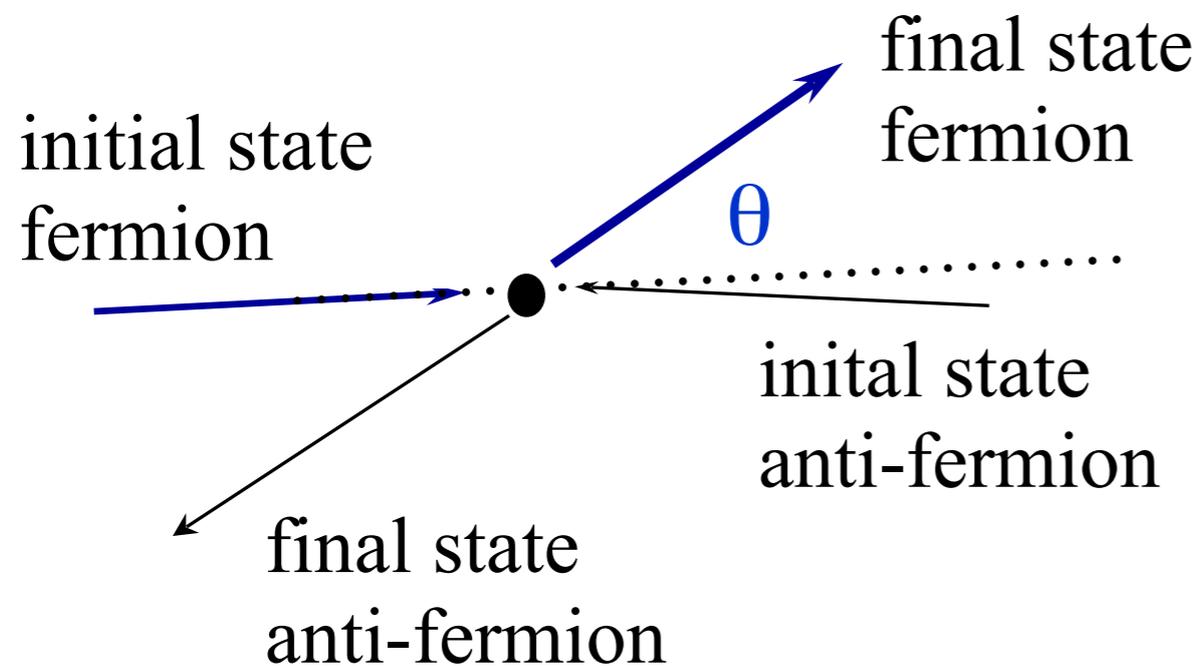
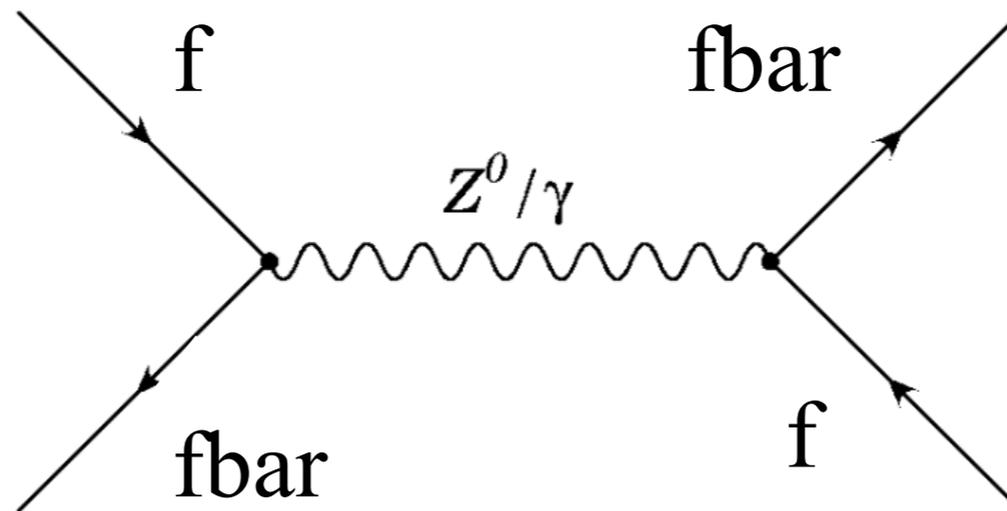
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# AFB from Drell-Yan process

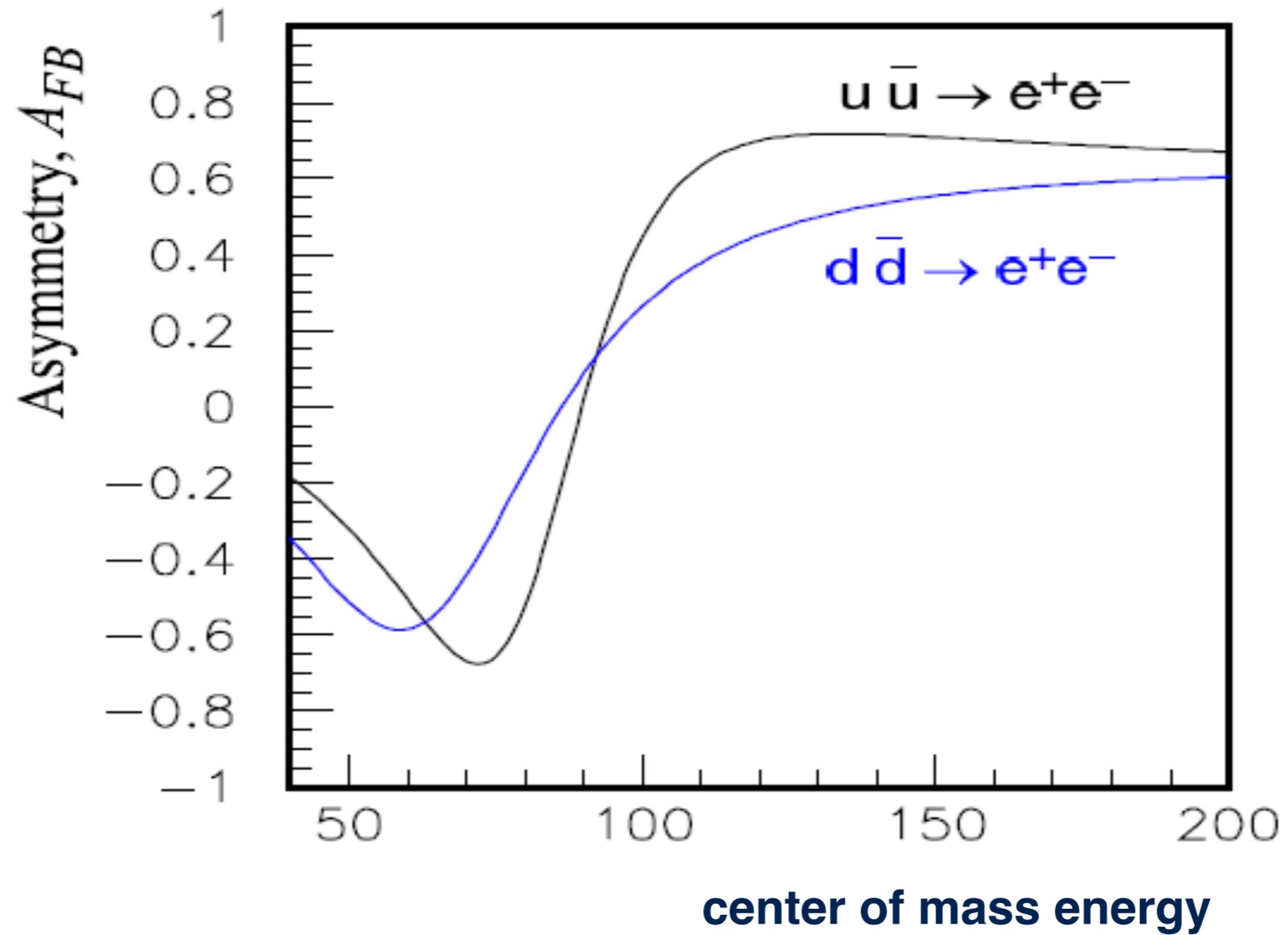


**$\cos\theta > 0$ , forward**

**$\cos\theta < 0$ , backward**

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = A_{FB}(\sin^2 \theta_{\text{eff}}^f)$$

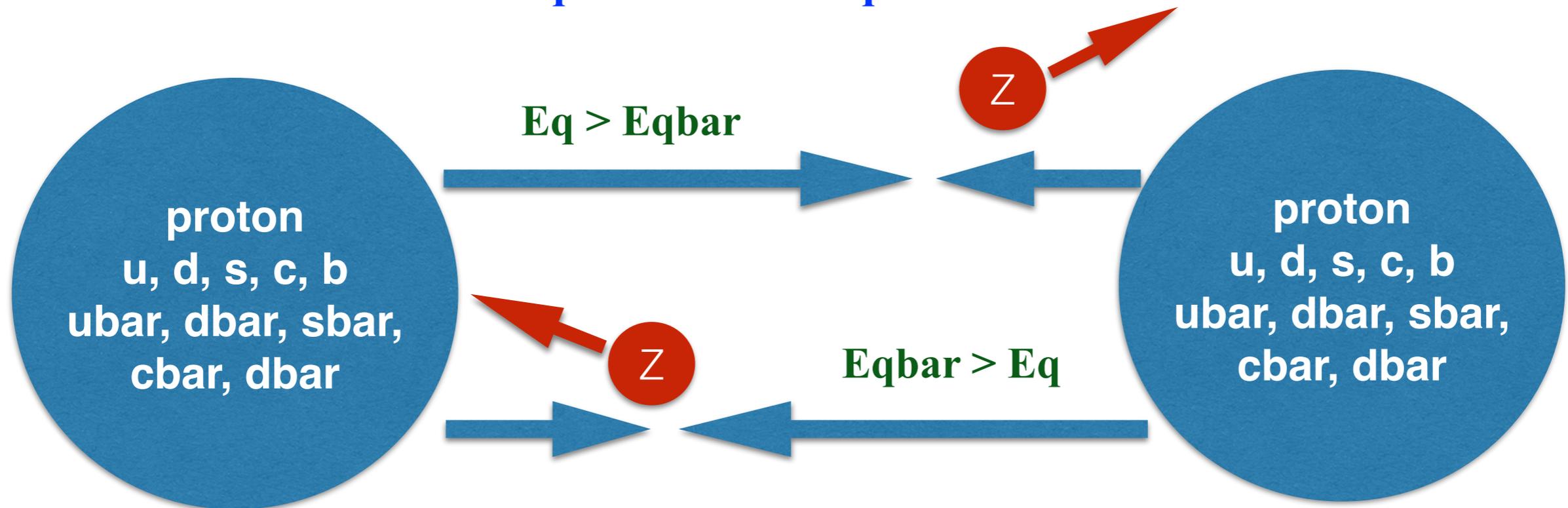
# AFB from Drell-Yan process



# AFB at hadron colliders

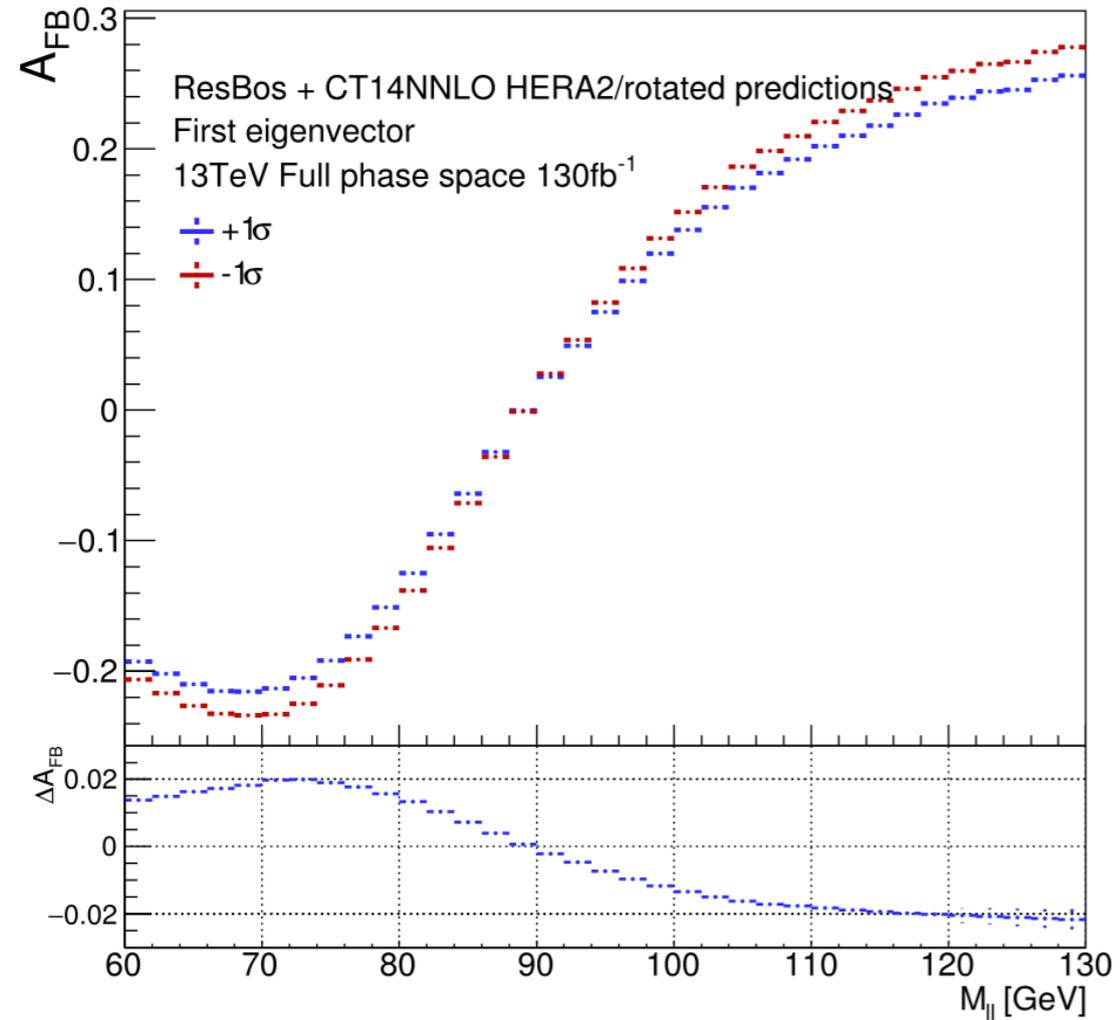
## Dilution effect

- Symmetrical initial state of proton-proton
- Asymmetry redefined according to the Z boson boost
- Relative difference between quarks and antiquarks

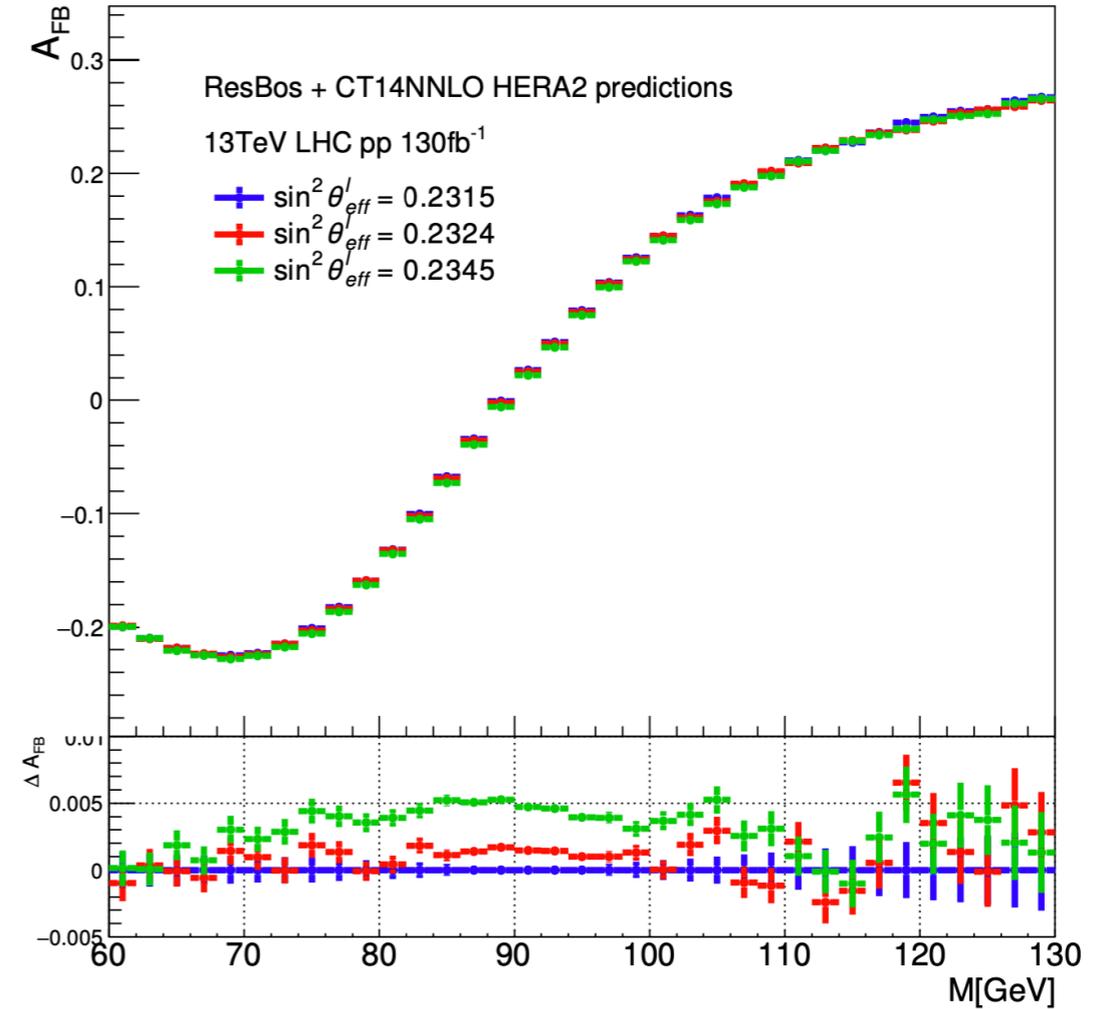


$$A_{FB} \rightarrow A_{FB}^h$$

# AFB governed by EW and dilution effect

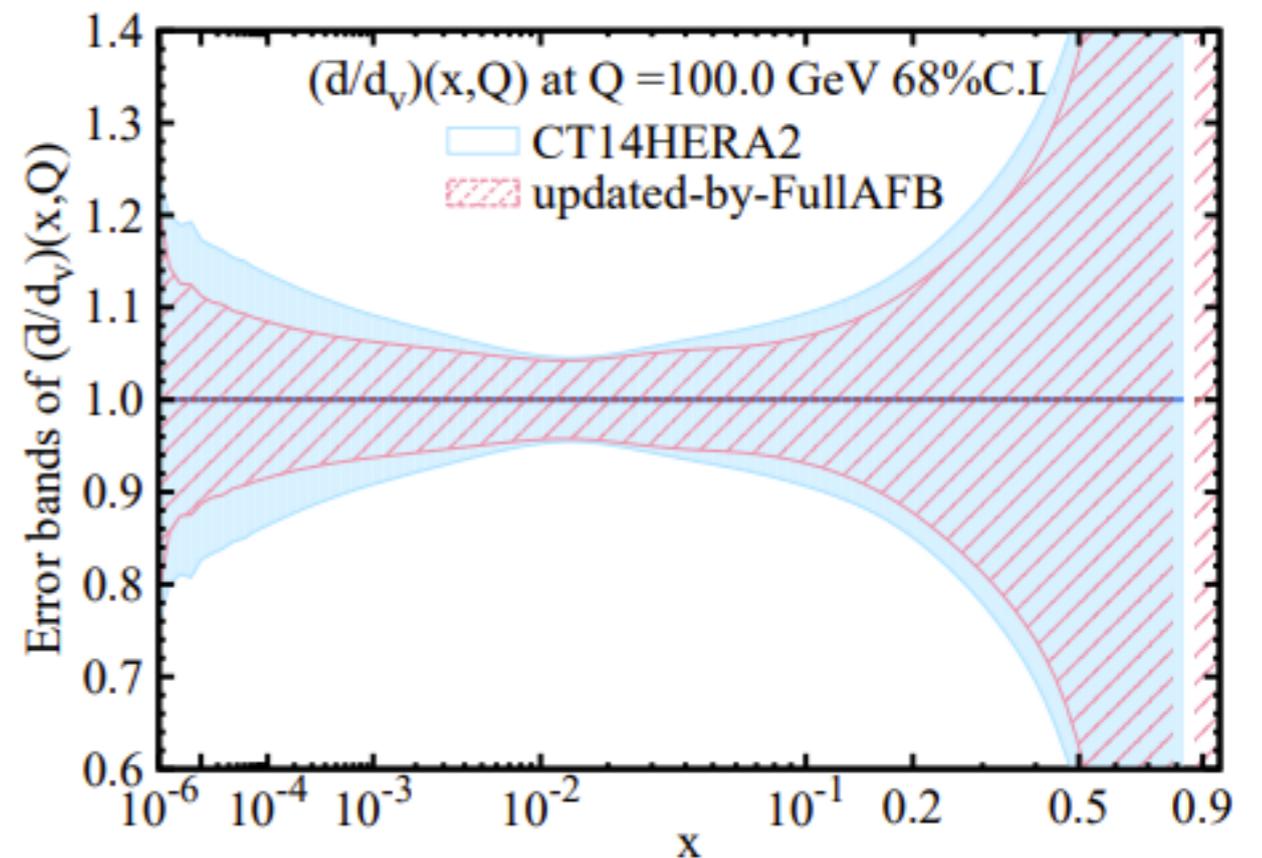
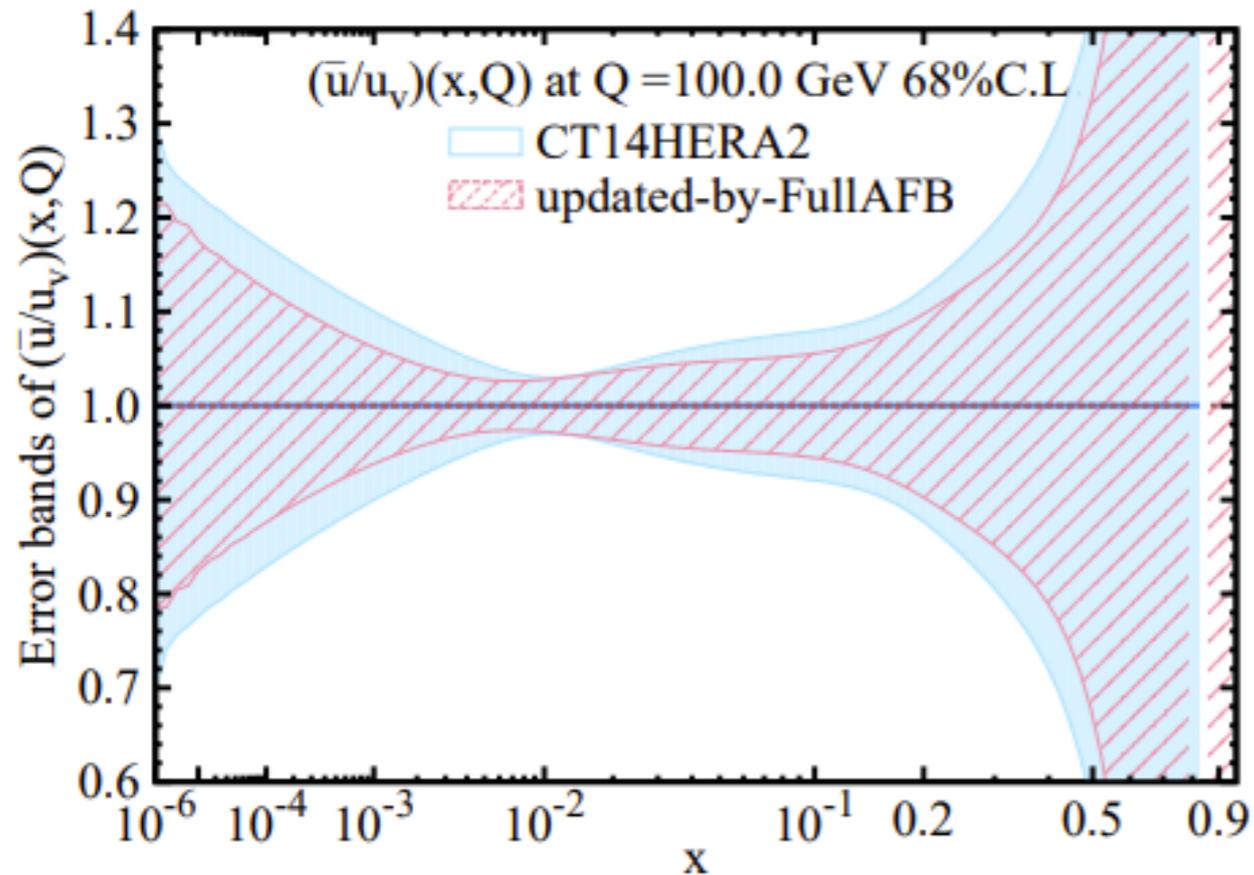


AFB governed by PDF,  
estimated by CT14



AFB governed by weak  
mixing angle

# AFB used in PDF global fitting



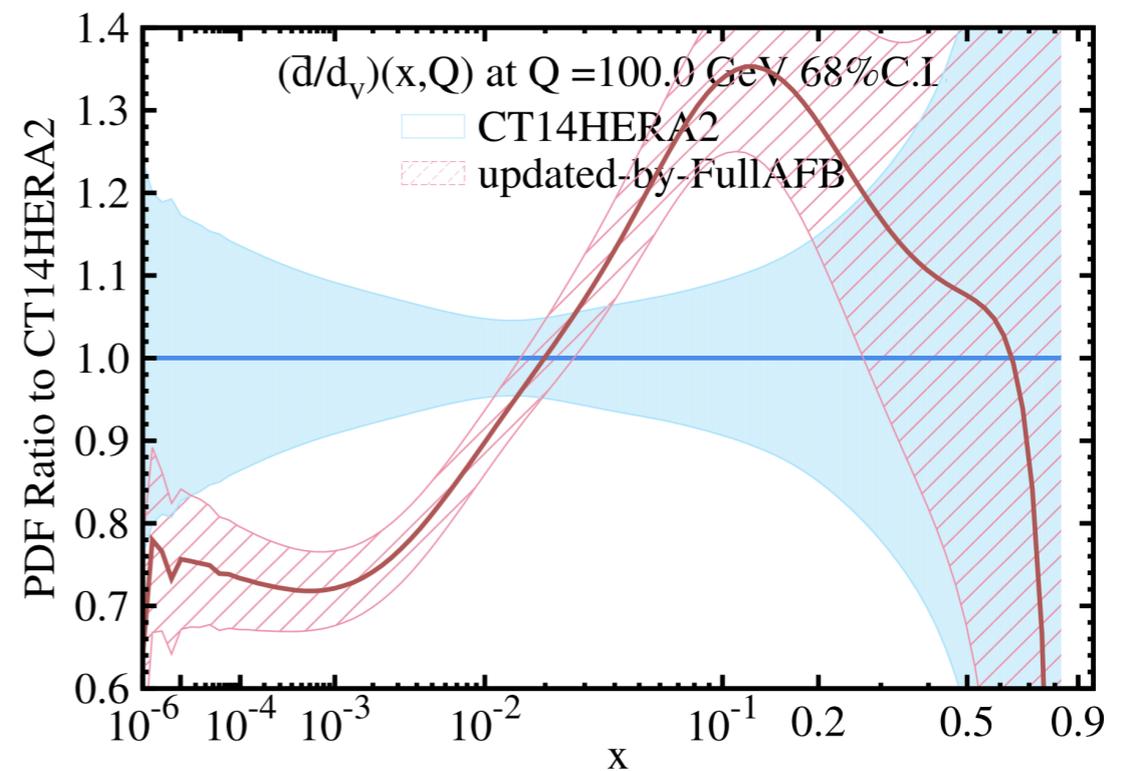
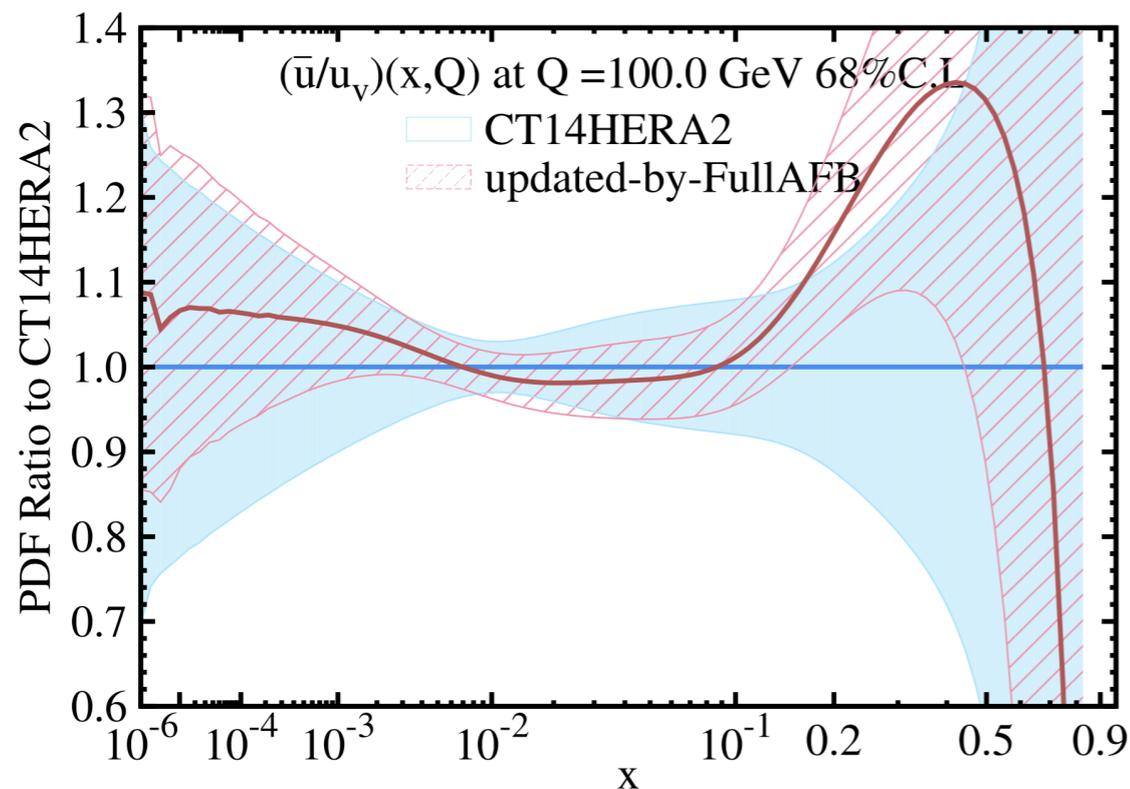
**Using AFB in the PDF global fitting, corresponding to  $130 \text{ fb}^{-1}$  data collected by ATLAS or CMS (LHC Run 2)**

**AFB generated using ResBos+CT14HERA2**

# Correlations between EW and proton structure

## Correlations

- Observation on proton structure highly correlated with determination on  $\sin^2\theta_W$
- Proton structure: need PDF-independent measurement of  $\sin^2\theta_W$  as input (LEP, SLC, or Tevatron)
- Measurement of  $\sin^2\theta_W$ : limited by the current modeling of proton structure



Using AFB in the PDF global fitting. The input AFB is generated with its  $\sin^2\theta_W$  value different from the PDF global fitting theory (0.2315 vs 0.2324)

# Determination on $\sin^2\theta_w$ and proton structure

## PDF global fitting strategy: not a good strategy

- Not supporting a combined fitting on  $\sin^2\theta_w$  and proton structure
- Massive work to provide a unified EW calculation in extensive experimental results to treat  $st_w$  and a floating parameter
- PDF global fitting is a combination of all data under various SM assumptions. Determinations on proton structure and  $st_w$  **are not single-data independent measurement**

# Factorization on AFB

## Collins-soper frame

- Center of mass frame
- z-axis defined as bisector of the angle formed by the direction of one incoming hadron beam (HA) and the negative direction of the other incoming hadron beam (HB)

**Collins-soper frame 1:**

**HA: according to “quark direction”**

**HB: according to “antiquark direction”**

$$\cos \theta_q$$

**Collins-soper frame 2:**

**HA: according to Z boson boost**

**HB: negative direction to HA**

$$\cos \theta_h$$

**if  $E_q > E_{qbar}$ :**

$$\cos \theta_h = \cos \theta_q$$

**if  $E_{qbar} > E_q$ :**

$$\cos \theta_h = -\cos \theta_q$$

# Factorization on AFB

**Differential cross section  
of DY process in  $\cos\theta_q$ :**

$$\begin{aligned} \frac{d\sigma}{d\cos\theta_q dY dM dQ_T} &= \sum_f \alpha_f(Y, M, Q_T) \times \left\{ (1 + \cos^2\theta_q) \right. \\ &+ A_0^f(Y, M, Q_T)(1 - 3\cos^2\theta_q) \\ &\left. + A_4^f(Y, M, Q_T)\cos\theta_q \right\}, \end{aligned}$$

**Differential cross section  
of DY process in  $\cos\theta_h$ :**

$$\begin{aligned} \frac{d\sigma}{d\cos\theta_h dY dM dQ_T} &= \sum_f \alpha_f(Y, M, Q_T) \\ &\times \left\{ (1 + \cos^2\theta_h) + A_0^f(Y, M, Q_T)(1 - 3\cos^2\theta_h) \right. \\ &\left. + [1 - 2D_f(Y, M, Q_T)]A_4^f(Y, M, Q_T)\cos\theta_h \right\}, \end{aligned}$$

$$D_f(Y, M, Q_T)$$

**probability of  $\cos\theta_h = -\cos\theta_q$  in a  
specific  $f\bar{f}$  subprocess, i.e.  
the probability of anti-quarks carries  
higher energy than quarks**

# Factorization on AFB

$$A_{FB}^h(Y, M, Q_T) = \frac{\sum_f [1 - 2D_f(Y, M, Q_T)] \alpha_f(Y, M, Q_T) A_{FB}^f(Y, M, Q_T)}{\sum_f \alpha_f(Y, M, Q_T)}$$

$$A_{FB}^h(Y, M, Q_T) = [\Delta_u(Y, M, Q_T) + P_0^u(Y, Q_T)] \cdot A_{FB}^u(Y, M, Q_T; \sin^2 \theta_{\text{eff}}^\ell) + [\Delta_d(Y, M, Q_T) + P_0^d(Y, Q_T)] \cdot A_{FB}^d(Y, M, Q_T; \sin^2 \theta_{\text{eff}}^\ell)$$

**AFB<sup>f</sup> is independent with proton structure. Fully calculated in EW theory !!**

$$\begin{aligned} P_0^u(Y, Q_T) &= \int \frac{1 - 2D_u(Y, M, Q_T) \alpha_u(Y, M, Q_T)}{\sum_f \alpha_f(Y, M, Q_T)} dM / \int dM \\ P_0^d(Y, Q_T) &= \int \frac{1 - 2D_d(Y, M, Q_T) \alpha_d(Y, M, Q_T)}{\sum_f \alpha_f(Y, M, Q_T)} dM / \int dM \\ \Delta_u(Y, M, Q_T) &= \frac{1 - 2D_u(Y, M, Q_T) \alpha_u(Y, M, Q_T)}{\sum_f \alpha_f(Y, M, Q_T)} - P_0^u(Y, Q_T) \\ \Delta_d(Y, M, Q_T) &= \frac{1 - 2D_d(Y, M, Q_T) \alpha_d(Y, M, Q_T)}{\sum_f \alpha_f(Y, M, Q_T)} - P_0^d(Y, Q_T) \end{aligned} \quad (5)$$

# Parton parameters

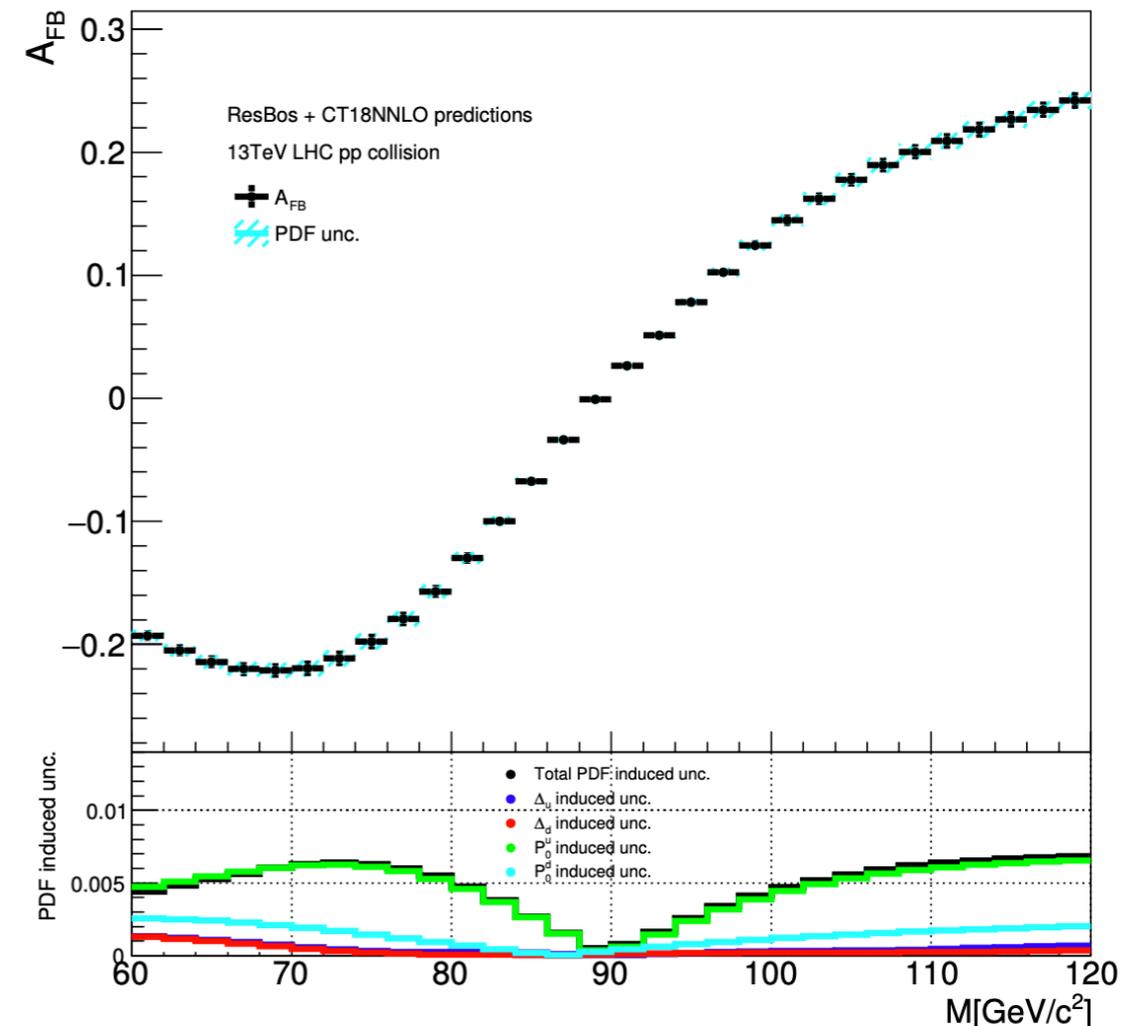
$$A_{FB}^h(Y, M, Q_T) = [\Delta_u(Y, M, Q_T) + P_0^u(Y, Q_T)] \cdot A_{FB}^u(Y, M, Q_T; \sin^2 \theta_{\text{eff}}^\ell) + [\Delta_d(Y, M, Q_T) + P_0^d(Y, Q_T)] \cdot A_{FB}^d(Y, M, Q_T; \sin^2 \theta_{\text{eff}}^\ell)$$

## P0 parameters

- Combined information on **dilution effect and relative cross section**, averaged in a given mass range (or x). Dominating part. P0 parameters and weak mixing angle will be the parameters in a combined fitting using AFB

## $\Delta$ parameters

- Mass dependence (or x dependence). It is difficult to be determined using AFB due to lack of sensitivity.
- However it introduces small uncertainties, because AFB is observed in a very small mass region with respect to the O(10) TeV level collision energy



# Parton parameters

## $P_{0u}$ vs $P_{0d}$

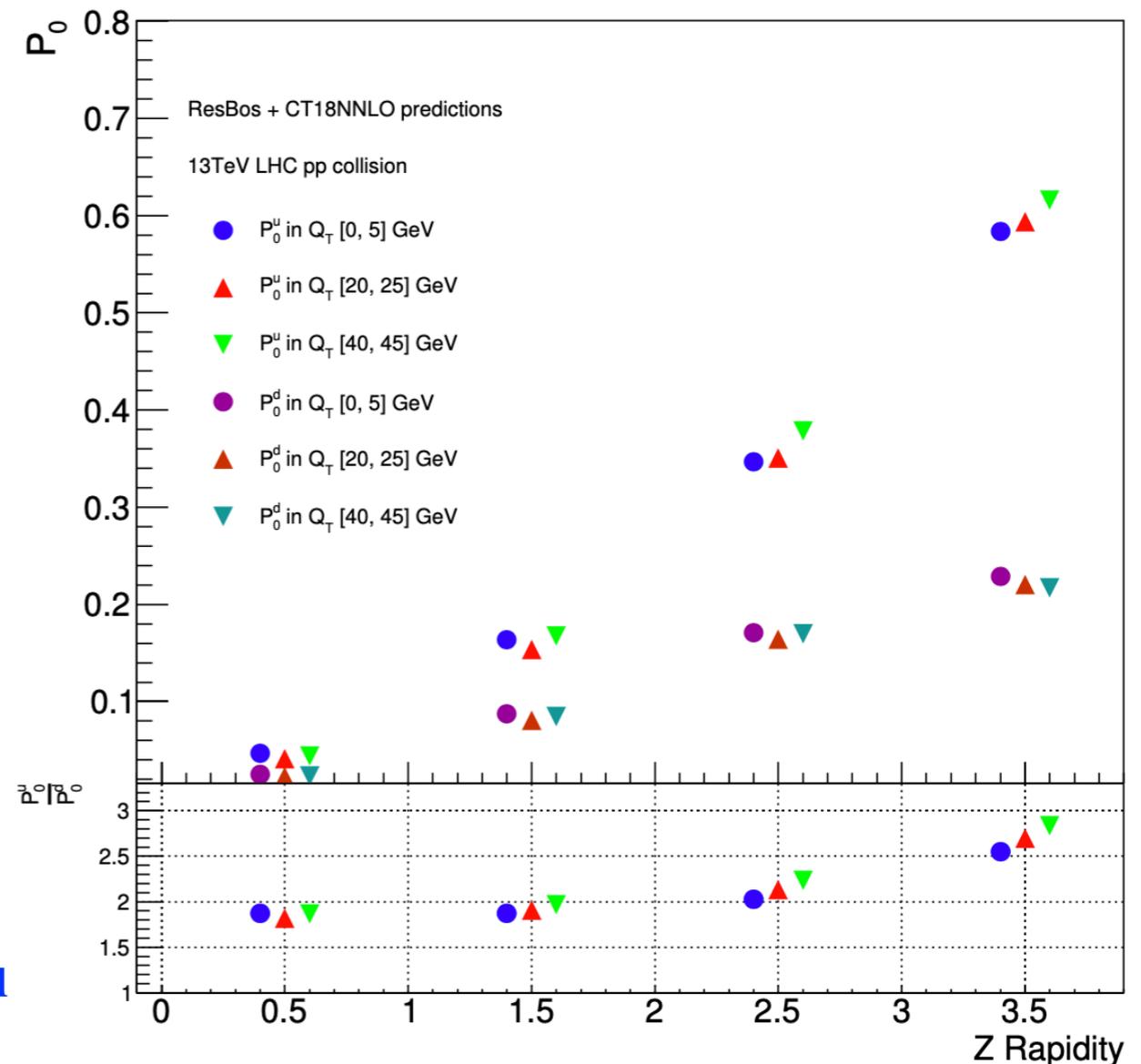
- $P_{0u}$  generally larger than  $P_{0d}$ , due to a smaller dilution effect in  $u\bar{u}$  subprocess

## As a function of $ZY$

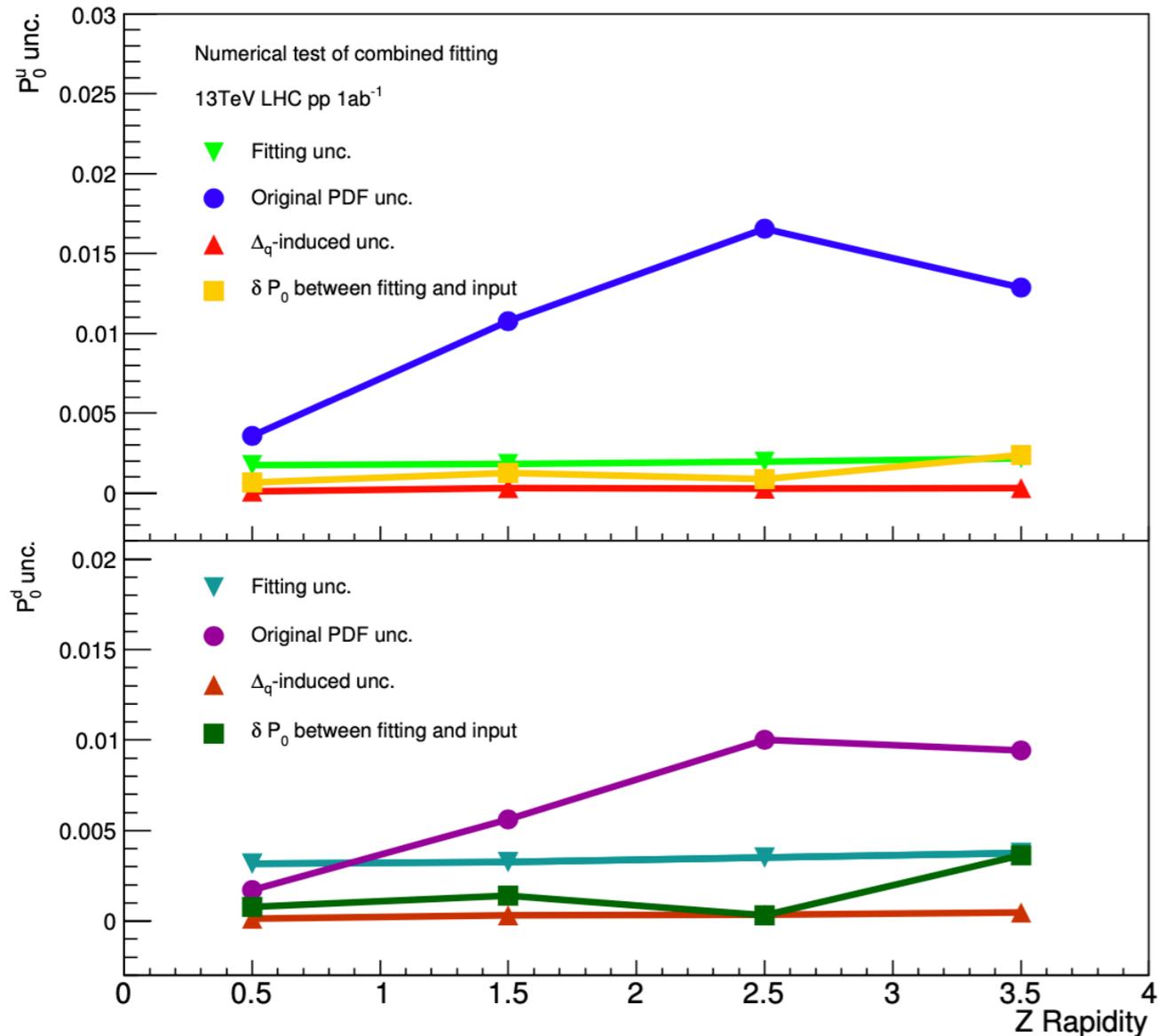
- Smaller  $ZY$  == larger dilution effect,  $P_0 \rightarrow 0$
- Larger  $ZY$  == smaller dilution effect,  $P_0 \rightarrow$  relative cross section

## Ratio of $P_{0u} / P_{0d}$

- $s$ ,  $c$  and  $b$  contributions cancel out
- Pure-observation on the difference between  $u$  and  $d$  quarks



# A combined fitting procedure and closure



**A proton structure observation  
using 1000 fb-1 data at the LHC**

|  |         |
|--|---------|
| CT18 induced uncertainty on $\sin^2\theta_W$       | 0.00038 |
| Combined fitting uncertainty on $\sin^2\theta_W$   | 0.00005 |
| $\Delta$ -induced uncertainty on $\sin^2\theta_W$  | 0.00008 |
| Input value of $\sin^2\theta_W$ in the pseudo-data | 0.23150 |
| Fitted value of $\sin^2\theta_W$                   | 0.23156 |

# Summary

## Background review

- **AFB@LHC are sensitive to both EW and proton structure**
- **However strongly correlated, both of them are limited**

## Factorization on AFB

- **We report a method which factorize the observed AFB into EW calculations and well-defined proton structure parameters**
- **The protons structure parameters can be experimentally observed, universally derived in theory calculation. Its definition is model independent**
- **Relative difference between  $q$  and  $qbar$ ; Relative difference between  $u$  and  $d$**

## For future measurement

- **Independent measurement using single AFB observation**
- **Correlation between EW and proton structure well considered**
- **Precision significantly improved for determination on both weak mixin angle and proton structure**

**Backup**

# An example of pre-work

Phys. Rev. Lett. 126, 041801

Emanuele Bagnaschi and Alessandro Vicini

$$\chi_k^2 = \sum_{i \in \text{bins}} \frac{[(\mathcal{T}_{0,k})_i - (\mathcal{D}^{\text{exp}})_i - \sum_{r \in \mathcal{R}} \alpha_r (\mathcal{S}_{r,k})_i]^2}{\sigma_i^2} + \sum_{r \in \mathcal{R}} \alpha_r^2,$$

- $(\mathcal{T}_0, k)_i$ : theory prediction of an observable
- $(\mathcal{D}^{\text{exp}})_i$ : data measurement
- $\alpha_r$ : nuisance parameters
- $(\mathcal{S}_r, k)_i$ :  $(\mathcal{T}_r, k)_i - (\mathcal{T}_0, k)_i$  representing PDF uncertainties

## Nuisance parameter fitting:

- Only for EW parameter determination. Nuisance parameters can't be used as observation on proton structure
- Model-dependent
- Assumption: the data used for measurement is well-consistent with current theory predictions both on EW and PDFs