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Updated Result on Nucleon Transversity PDF from Lattice QCD

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CLHCP, 2021/11/25

OUTLINE

- **Introduction**
- **Lattice calculation**
- **Numerical result**
- **Summary and outlook**

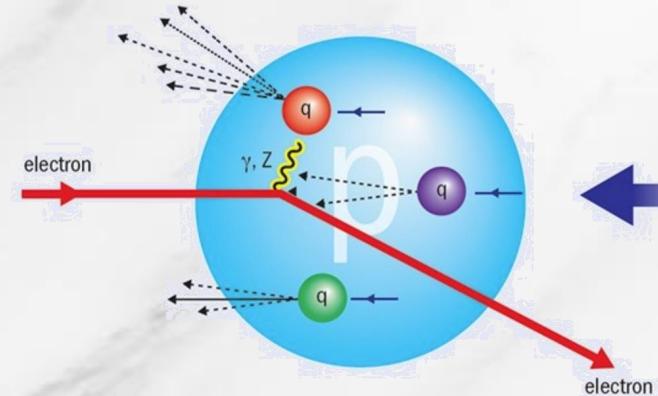


PART.1

Introduction

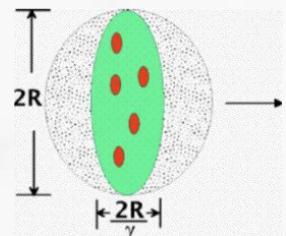
Introduction

□ Deep inelastic scattering

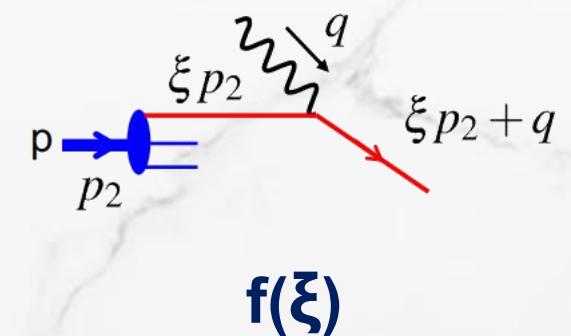
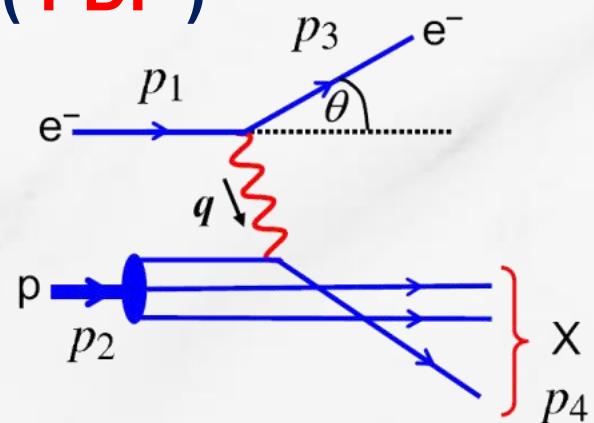
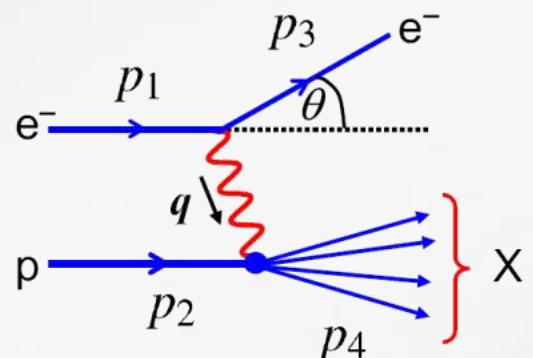


- proton contains **point-like objects**
- Feynman's parton model

□ Nucleons are Lorentz contracted along the direction of motion

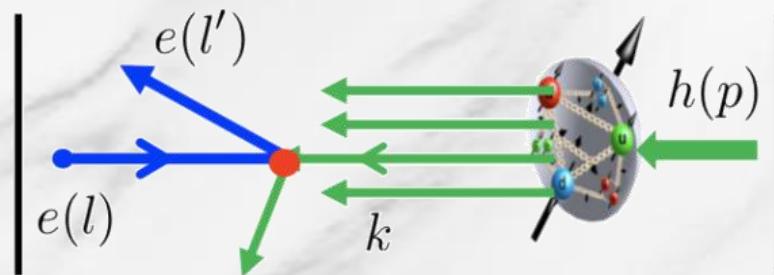


□ Parton distribution function (PDF)

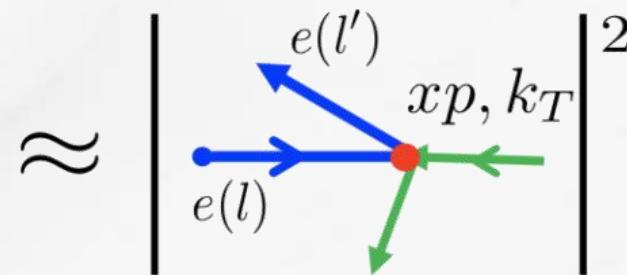


Introduction

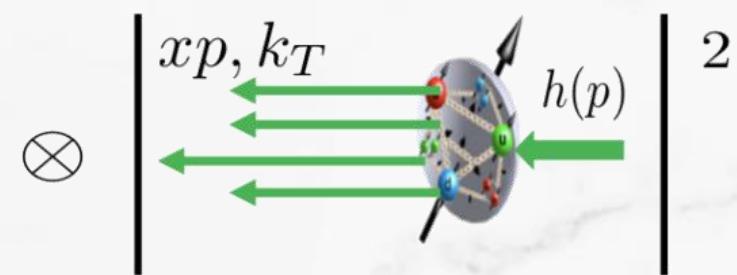
□ QCD factorization $\sigma = \sum_a \hat{H}_{l+a \rightarrow X} \otimes f_{a/p}$



Cross section

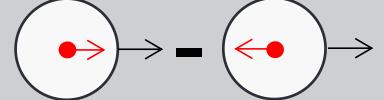


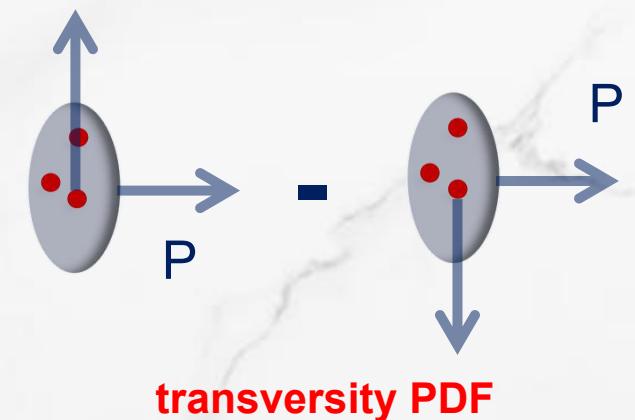
Hard part
(Perturbative)



Parton distribution
(Nonperturbative)

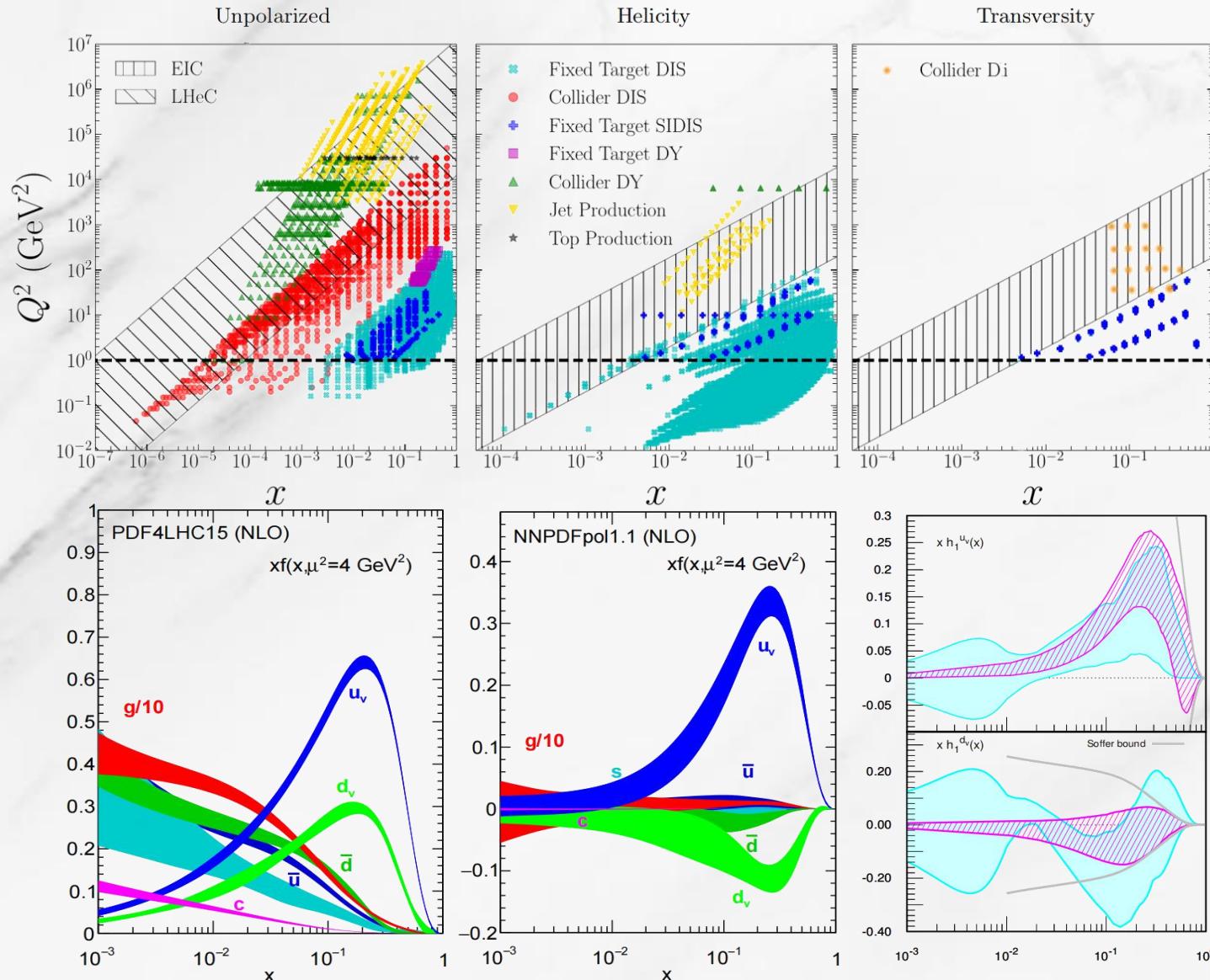
□ Polarized PDF

Unpolarized	Longitudinally polarized	Transversely polarized
	 helicity	 transversity



Motivation

□ Studying PDF can understand inner structure of nucleons in detail

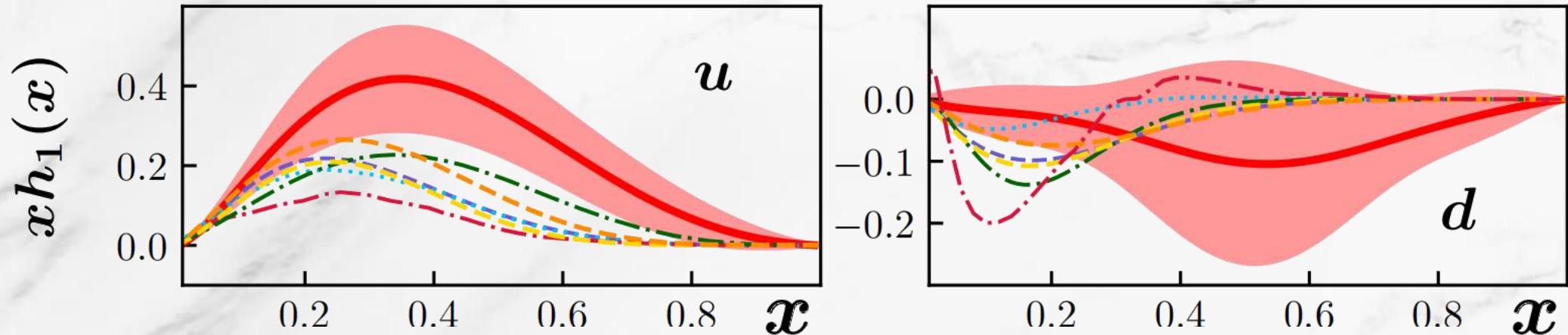


- There are more experimental data of unpolarized and helicity PDF than transversity PDF
- Wider error band

Lin et al—Prog.Part.Nucl.Phys. 121
(2021) 103908

Motivation

- Global fit



JAM Collaboration — PRD.102 (2020) 5, 054002

- Large error compared with unpolarized and helicity PDF
- The fitted antiquark transversity is zero within **large uncertainties**
- Lattice calculation can provide information complementary to experimental data

Theoretical framework

□ The light-cone transversity PDF Definitions:

$$\delta q(x) = \int \frac{d\lambda}{2\pi} e^{ix\lambda} \langle P, s_\perp | \bar{\psi}(\lambda n) i\gamma^x \gamma^t \gamma^5 L(\lambda n, 0) \psi(0) | P, s_\perp \rangle \quad L(\lambda n, 0) = e^{ig \int_0^\lambda d\lambda' n \cdot A(\lambda' n)}$$

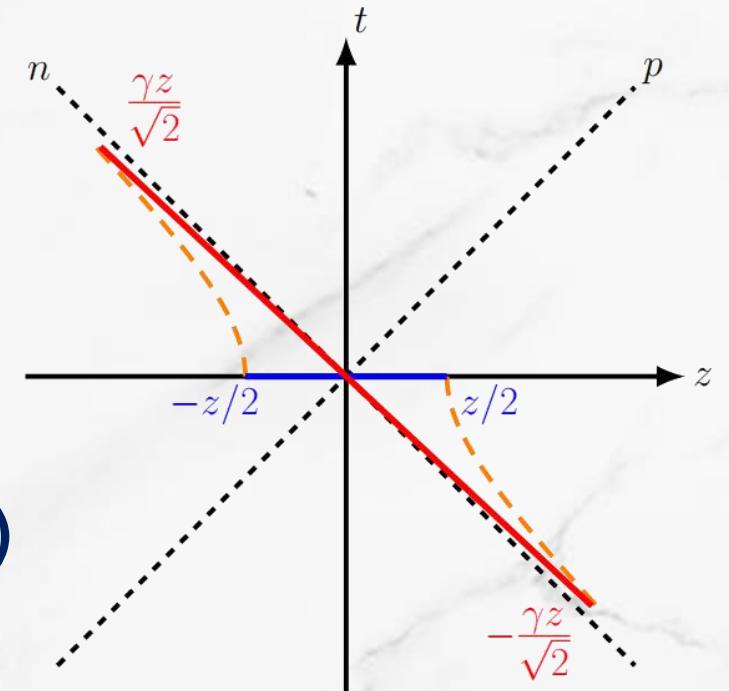
□ Lattice QCD calculate physical observable from path integral in Euclidean space

$$t \rightarrow -i\tau, e^{iS_M} \rightarrow e^{-S_E}$$

Recover physical limit

$$m_\pi \rightarrow m_\pi^{phys}, a \rightarrow 0, L \rightarrow \infty$$

□ Can be used to compute Euclidean (quasi-LF) correlations at large momentum in LaMET



Ji, PRL 110 (2013) 262002



Theoretical framework

- **quasi-LF correlation in Euclidean space**

$$\tilde{h}(P, z) = \langle P, s | \bar{\psi}(z) i\gamma^x \gamma^t \gamma^5 L(z, 0) \psi(0) | P, s \rangle$$

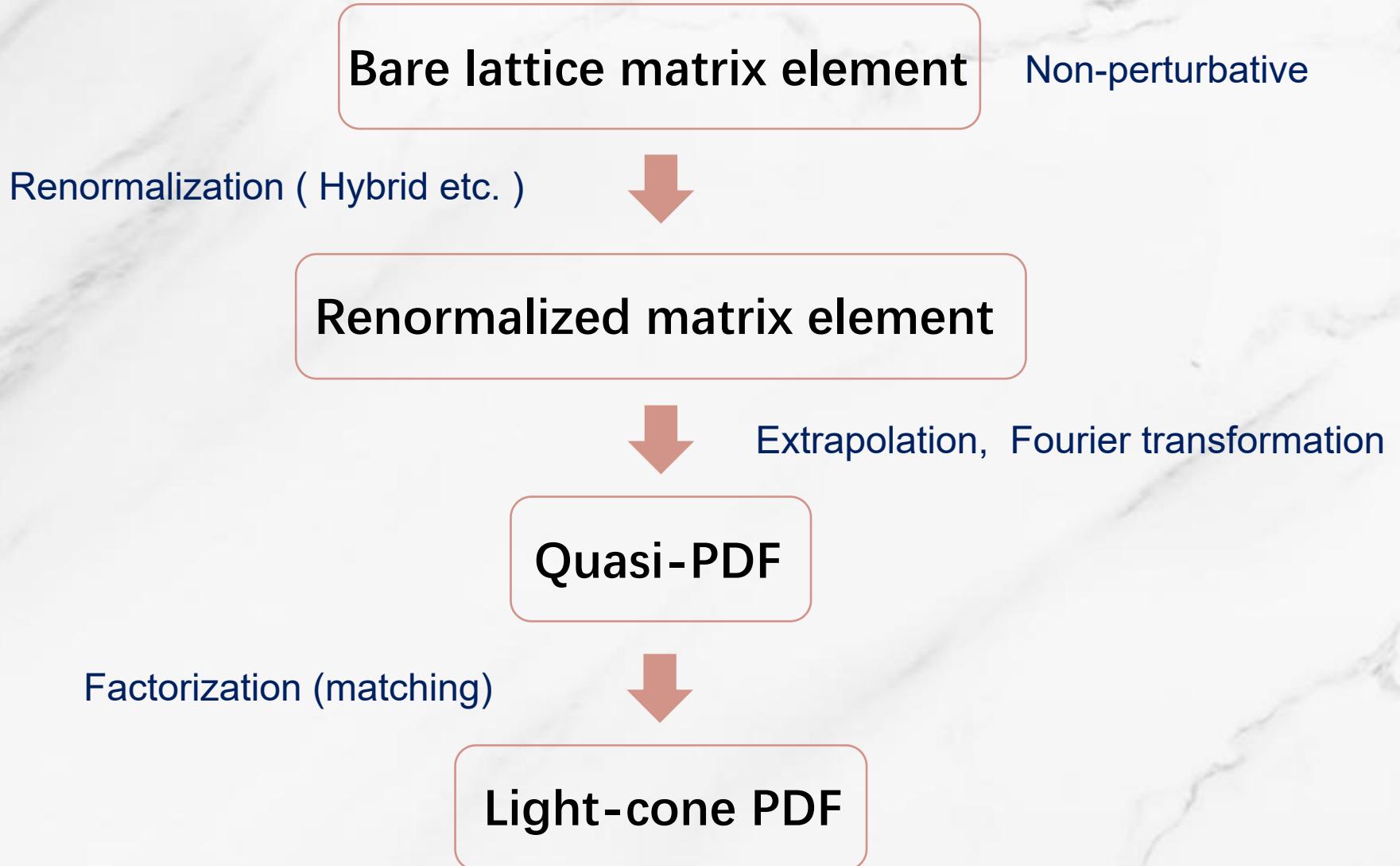
- **The light-cone transversity PDF can be related to the quasi-PDF by the following factorization formula**

$$\boxed{\tilde{\delta}q(x, P^z)} = \int_0^1 \frac{dy}{|y|} \delta C\left(\frac{x}{y}, \frac{\mu}{y P_z}\right) \delta q(y, \mu) + \mathcal{O}(\Lambda_{QCD}^2/(y P^z)^2, \Lambda_{QCD}^2/((1-y) P^z)^2)$$

$$\boxed{\tilde{\delta}q\left(x, \frac{\mu}{P^z}\right) = \int_{-\infty}^{\infty} \frac{d\zeta}{2\pi} e^{ix\zeta} \tilde{h}\left(\zeta, \frac{\mu^2 \zeta^2}{P_z^2}\right) \quad \zeta = z p_z}$$



Theoretical framework





PART.2

Lattice calculation



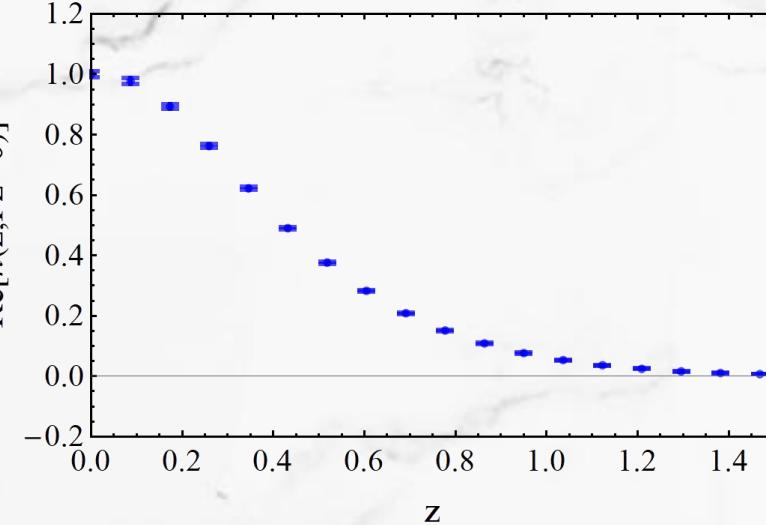
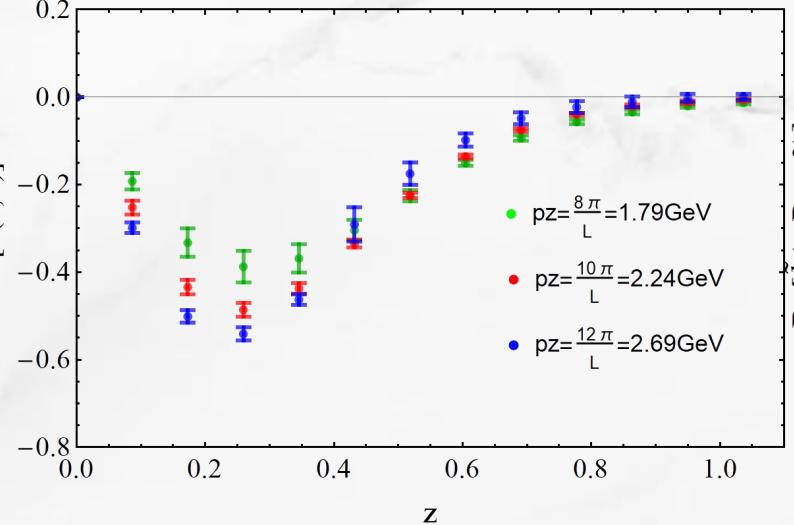
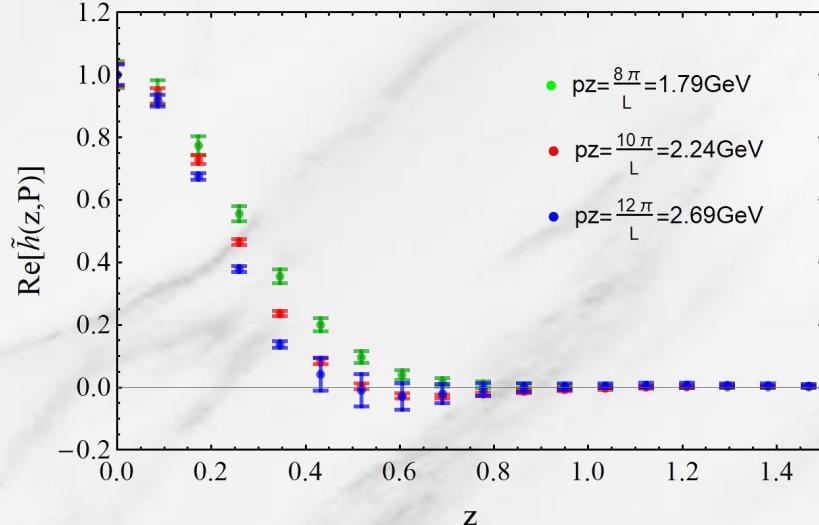
CLS ensembles

ensembles	a(fm)	L ³ ×T	m _π (MeV)
H102	0.08636	32 ³ ×96	354
N203	0.06426	48 ³ ×128	345
N302	0.04981	48 ³ ×128	346

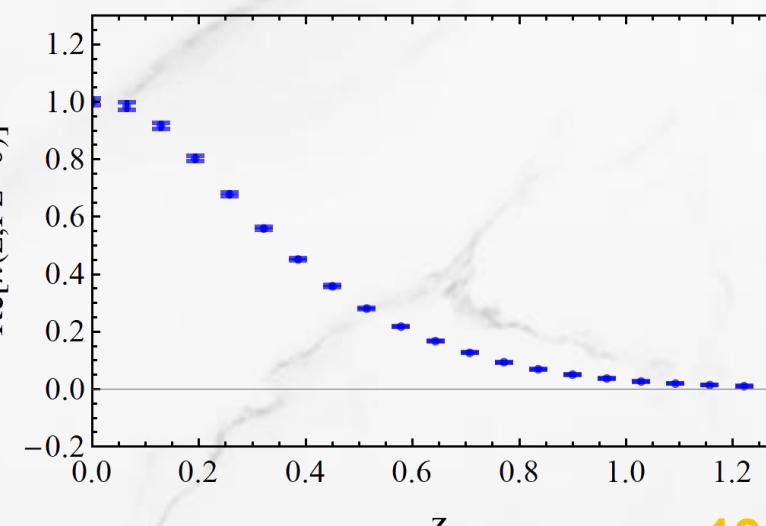
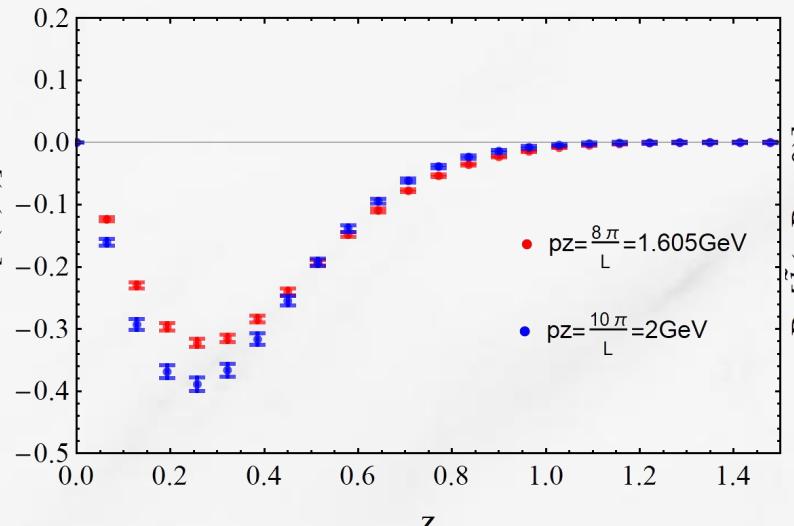
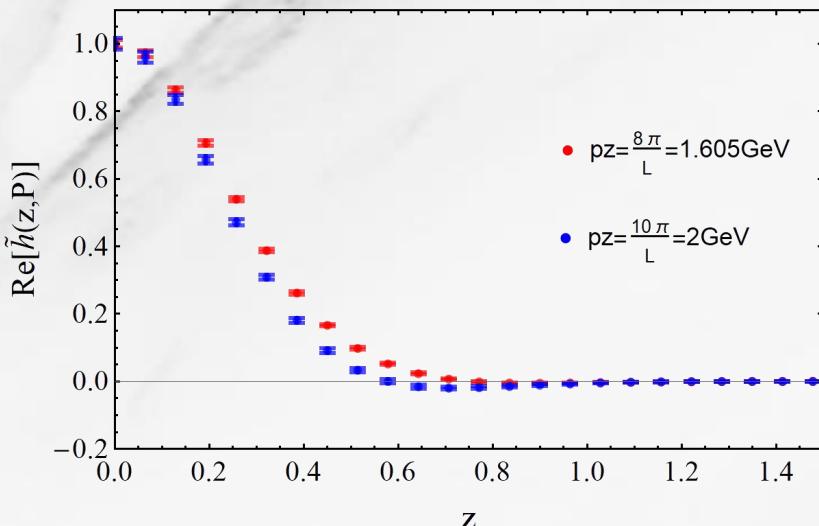
- H102: $P_z = \{4, 5, 6\}2\pi/L \approx \{1.79, 2.24, 2.69\}\text{GeV}$
- N203: $P_z = \{4, 5\}2\pi/L \approx \{1.6, 2\}\text{GeV}$
- N302: $P_z = \{4, 5\}2\pi/L \approx \{2.07, 2.59\}\text{GeV}$

Bare matrix elements

a=0.08636fm



a=0.06426fm

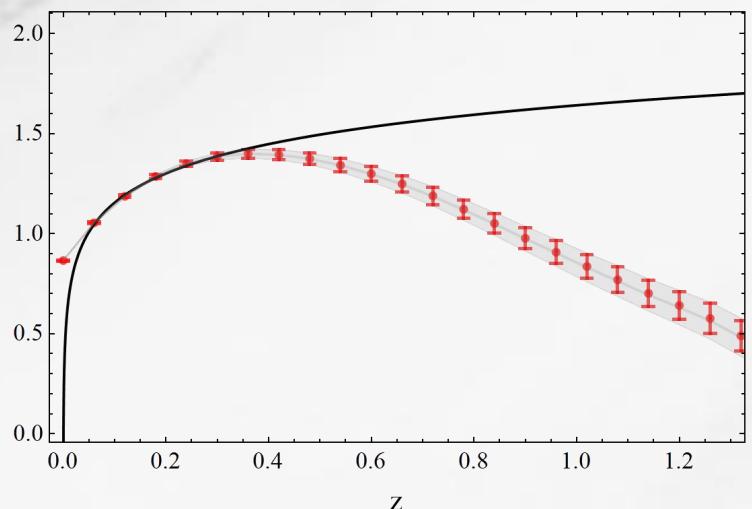


Renormalization in Hybrid scheme

- The quasi-LF correlation operator is **multiplicatively renormalized**

$$[\bar{\psi}(z)\Gamma W(z,0)\psi(0)]_B = e^{\delta m|z|} Z [\bar{\psi}(z)\Gamma W(z,0)\psi(0)]_R \quad \text{Ji, Zhang, Zhao, PRL 120, 112001 (2018)}$$

- Ratio scheme : Divide by the same operator matrix element at zero momentum Radyushkin, PRD 98 (2018) 1, 014019
- Problem: introduce undesired IR effects at large distances



- Hybrid renormalization scheme**

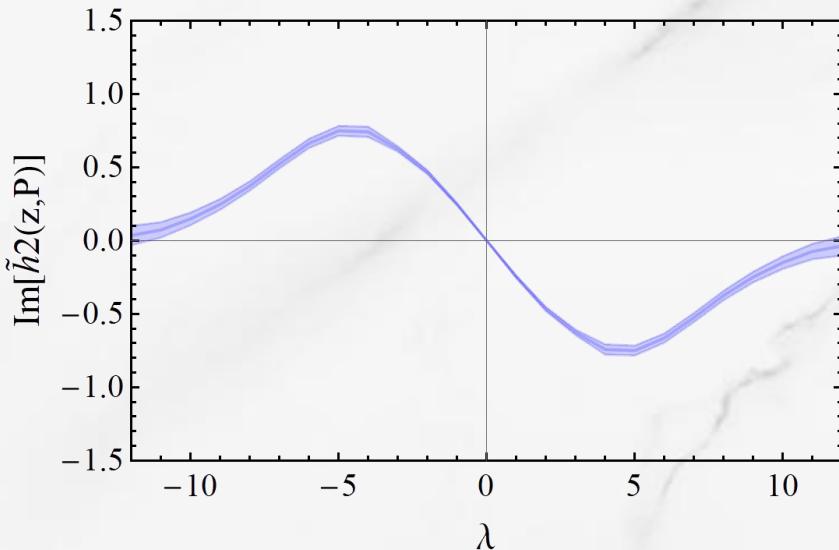
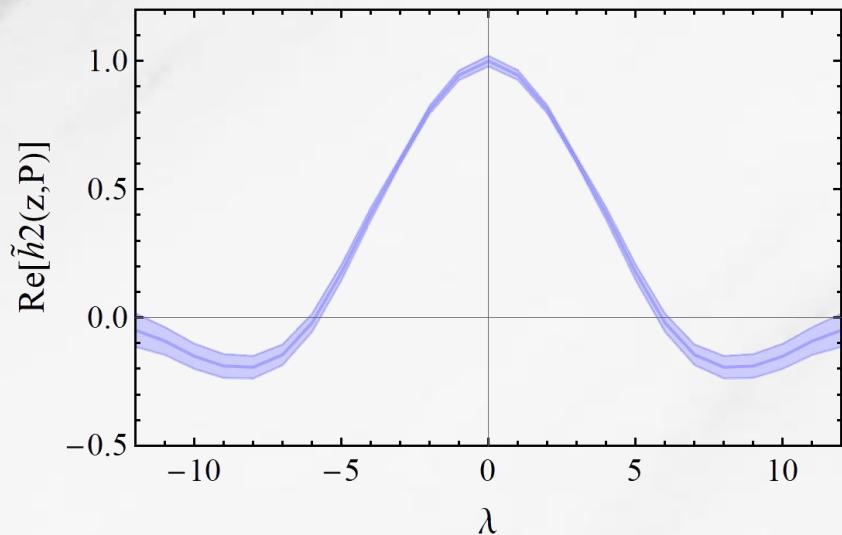
Ji, JHZ et al, Nucl.Phys.B 964 (2021) 115311

$$\tilde{h}^R(z, a, P_z) = \frac{\tilde{h}(z, a, P_z)}{Z_X(z, a)} \theta(z_S - |z|) + \tilde{h}(z, a, P_z) e^{-\delta m|z|} Z_{\text{hybrid}}(z_S, a) \theta(|z| - z_S)$$

Renormalization in Hybrid scheme

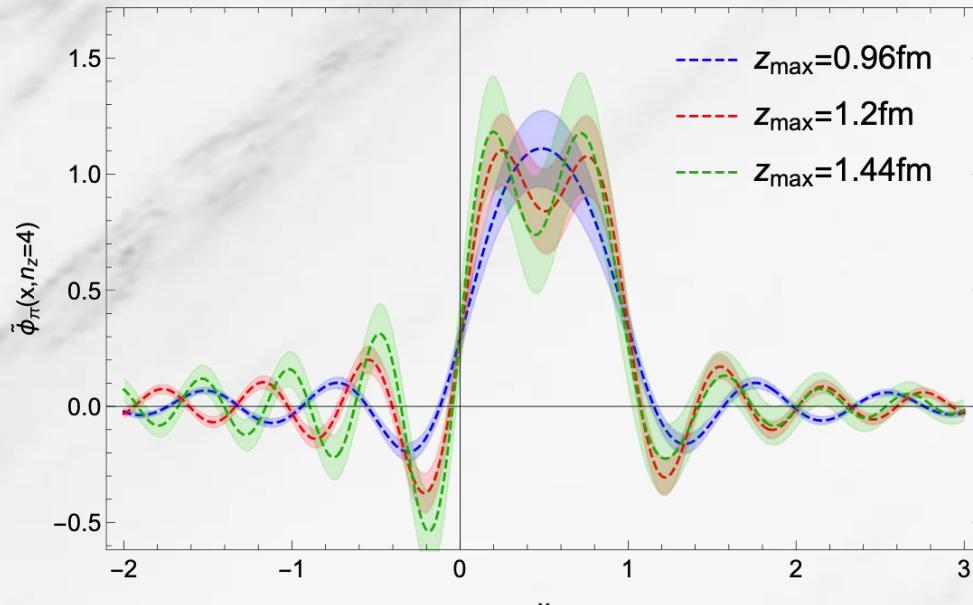
$$\tilde{h}^R(z, a, P_z) = \frac{\tilde{h}(z, a, P_z)}{Z_X(z, a)} \theta(z_S - |z|) + \tilde{h}(z, a, P_z) e^{-\delta m |z|} Z_{\text{hybrid}}(z_S, a) \theta(|z| - z_S) \quad z_S = 0.3 \text{ fm}$$

- For $|z| \leq z_S$ (perturbative region), do the renormalization in the ratio scheme
- For $|z| > z_S$, use Wilson line mass renormalization that avoids introducing extra nonperturbative effects at large distances



Extrapolation and Fourier transformation

- Missing large λ information
- Unphysical oscillation



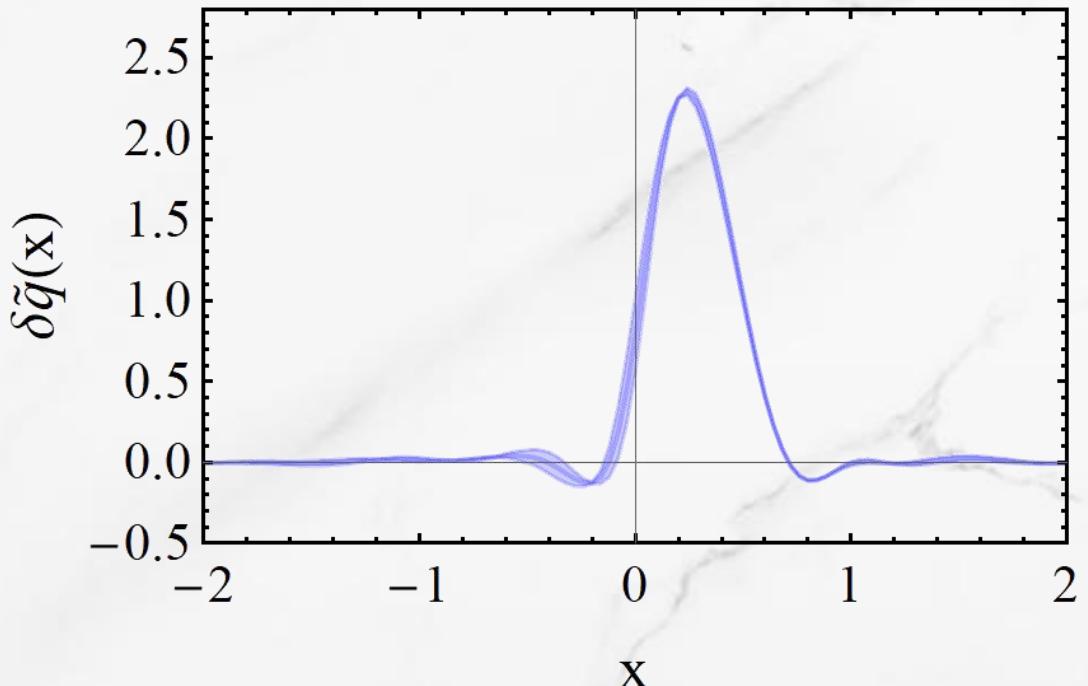
R. Zhang et al —PRD.102.094519(2020)

- **Extrapolation**

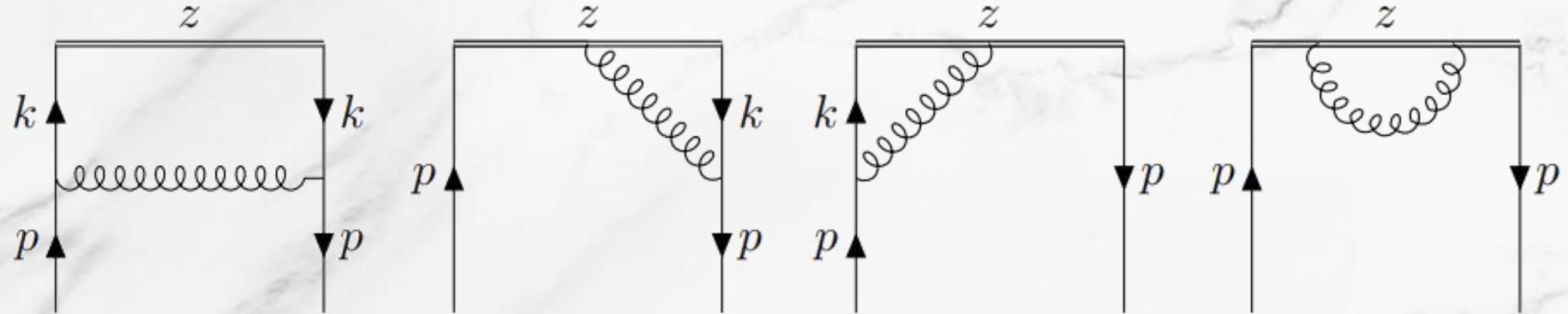
$$\tilde{h}(z, P_z) = \left[\frac{c_1}{(-i\lambda)^a} + e^{i\lambda} \frac{c_2}{(i\lambda)^b} \right] e^{-\frac{\lambda}{\lambda_0}}$$

- **Fourier transformation**

$$\delta \tilde{q}(x) = \int \frac{d\lambda}{2\pi} e^{ix\lambda} \tilde{h}(\lambda)$$



Matching coefficient (In momentum space)



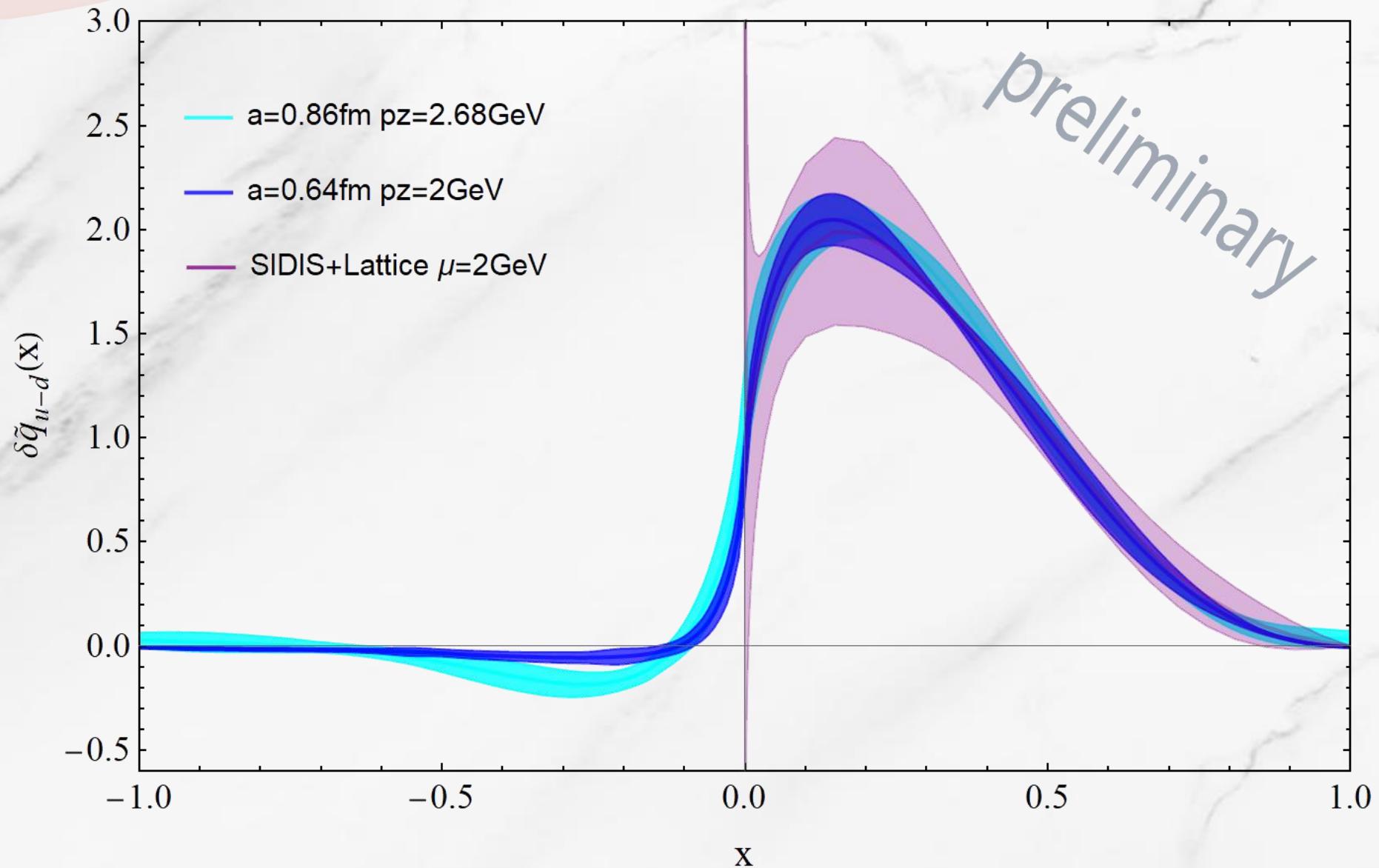
Ratio scheme

$$\delta C_r \left(x, \frac{\mu}{p_z} \right) = \delta(1-x) + \frac{\alpha_s C_F}{2\pi} \begin{cases} \left[\frac{2x}{1-x} \ln \frac{x}{x-1} - \frac{2}{1-x} \right]_+ & x > 1 \\ \left[\frac{2x}{1-x} \left(\ln \frac{4p_z^2}{\mu^2} + \ln x(1-x) \right) - \frac{2x}{1-x} + \frac{2}{1-x} \right]_+ & 0 < x < 1 \\ \left[-\frac{2x}{1-x} \ln \frac{x}{x-1} + \frac{2}{1-x} \right]_+ & x < 0 \end{cases}$$

Hybrid scheme

$$\delta C_h \left(x, \frac{\mu}{p_z}, z_s^2 \mu^2 \right) = \delta C_r \left(x, \frac{\mu}{p_z} \right) + \frac{\alpha_s C_F}{2\pi} 2 \left[-\frac{1}{|1-x|} + \frac{2Si((1-x)\lambda_s)}{\pi(1-x)} \right]_+$$

PART.3 Numerical result





Summary

- PDF plays an important role in the detailed understanding of the internal structure of nucleon
- We calculate the Nucleon Transversity PDF on lattice using LaMET
- The numerical results agree with the global fit results within the error band
- Follow-up
 - Improving error analysis
 - physical extrapolation



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Thank you for listening !