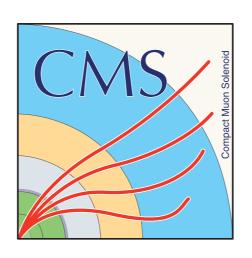
Observation of the electroweak production of Zy and two jets at 13 TeV and constraints on EFTs

Ying An, Peking University CLHCP meeting

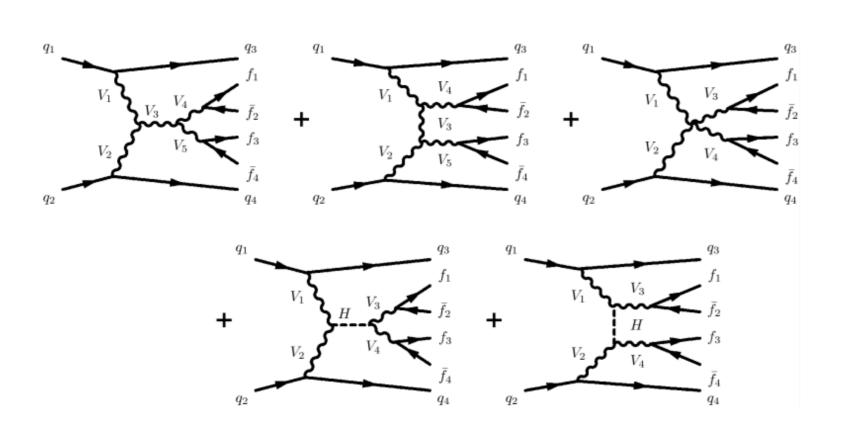


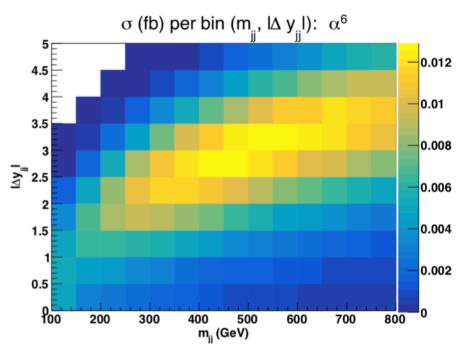


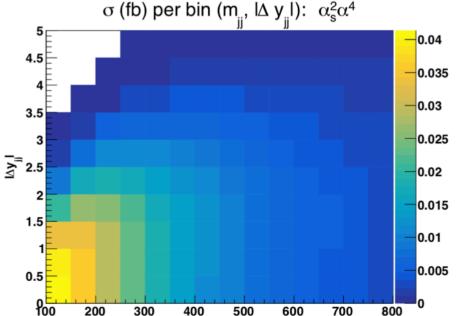


Introduction & Motivation









m_{ii} (GeV)

- **Signal:** six fermions final state at leading order $\mathcal{O}(\alpha^6)$
- Irreducible background: QCD-induced $\mathcal{O}(\alpha^4\alpha_s^2)$
- Interference: between EW and QCD $\mathcal{O}(\alpha^5\alpha_s)$
- Reducible background due to mis-ID of final state particles
- Significant systematic uncertainties from jet energy reconstruction and background modeling



Introduction & Motivation



Important process to investigate electroweak symmetry breaking(EWSB)

- Probe the nature of EW symmetry breaking
- Unitarity preservation visible only in VV scattering

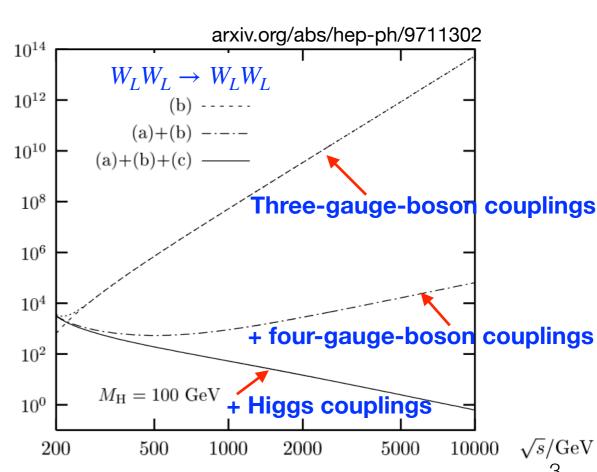
Complimentary to direct Higgs Boson measurement

The perturbative cross section of longitudinal VBS ($V_L V_L \to V_L V_L$) diverges, if there was no Higgs boson or a similar mechanics

Sensitive to anomalous coupling

Triple and quartic gauge coupling

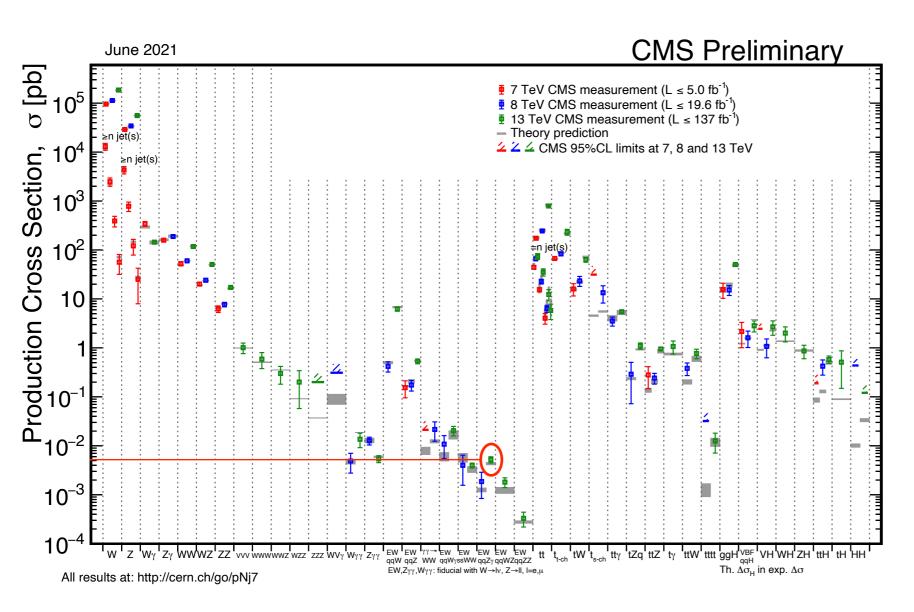
$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_{i} \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}^{(6)} + \frac{c_i^{(8)}}{\Lambda^2} \mathcal{O}^{(8)} + \dots$$



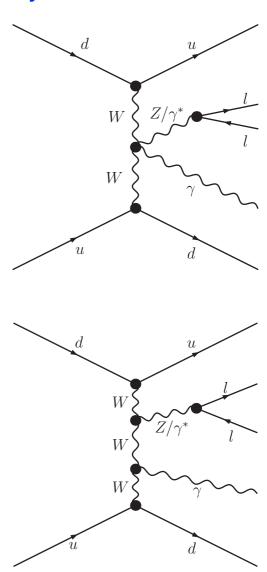


Introduction & Motivation





Phys. Rev. D 104, 072001



Final states: Z to ee/µµ plus a photon with two additional jets.

Vector boson scattering (VBS) signature: large dijet mass and large η separation between the jets.

Main results:

- √ Signal significance
- √ Fiducial cross section
- ✓ Unfolded differential cross section
- √ Limits on anomalous couplings



Sample & Selection



Data: collected from 2016 to 2018 with integrated luminosity: 137 fb⁻¹

Signal: EW Zγjj

- MadGraph_aMC@NLO (MG5) at LO
- Pythia8 with CP5 (CUETP8M1 for 2016)
- · NNPDF 3.1(3.0 for 2016)
- $m_{II} > 50 \text{ GeV}$

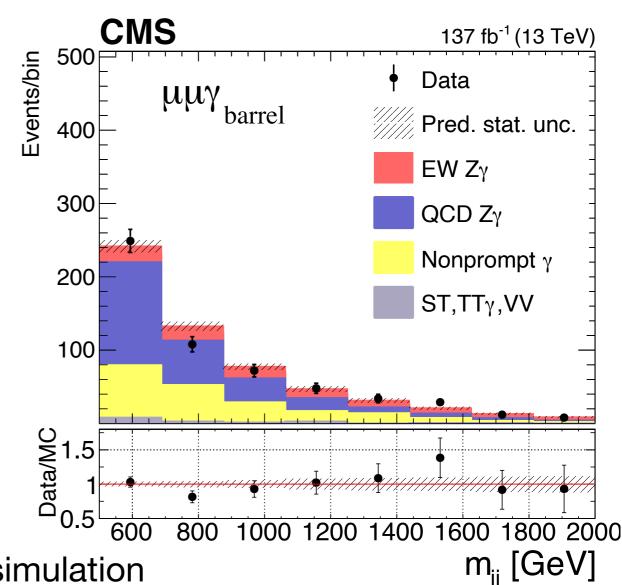
Backgrounds:

- * **Z**γ **plus QCD jets** from simulation
 - MG5 with FxFx jet merging scheme at NLO
 - Pythia8 with CP5 (CUETP8M1 for 2016)
 - NNPDF 3.1(3.0 for 2016)
- Nonprompt photon from data
- EW/QCD Interference from simulation





- $\cdot t\bar{t}\gamma$: MG5 at NLO with FxFx jet merging scheme
- single top: POWHEG at NLO



Phys. Rev. D 104, 072001 (FIG. 2)



Tight

Sample & Selection





Medium Loose

a series of variables reflecting the properties of the particle are optimized to identify the particle.

High quality High efficiency

Good Muon

- Tight muon WP
- Relative PF-isolation (0.4 cone) < 0.15
- $p_T > 20 \text{ GeV}, |\eta| < 2.4$

Veto Muon

- Loose muon WP
- Relative PF-isolation (0.4 cone) < 0.25
- $p_T > 20 \text{ GeV}$, |n| < 2.4

Veto Electron

- Loose electron WP
- p_T > 20 GeV, |η| < 2.5, |η| < 1.4442 or
 1.566 < |η| < 2.5 For third lepton veto

Good Electron

- Medium electron WP
- $p_T > 25$ GeV, $|\eta| < 2.5$

Good Photon

- Medium photon WP
- Electron veto
- p_T > 20 GeV and |η| < 1.4442 or
 1.566 < |η| < 2.5

Jets

- Particle-flow jets and AK4CHS (0.4 cone; charged particles from pileup are removed)
- Tight jet WP and pileup jet WP (p_T < 50 GeV)
- p_T>30 GeV
- $|\eta| < 4.7$



Sample & Selection



- Two same-flavor opposite-sign tight leptons *
- Double muon/electron HLT paths
- Third lepton veto
- 70 GeV $< m_{ll} <$ 110 GeV \Rightarrow
- One good photon in barrel/endcap
- Two jets with $p_T > 30$ GeV, $|\eta| < 4.7$
- $m_{ll\gamma}$ > 100 GeV
- 150 GeV $< m_{jj} < 500$ GeV
- $m_{jj} > 500 \text{ GeV } +$
- $\Delta \eta_{ij} > 2.5$
- $p_T^{\gamma} > 120 \text{ GeV}$
- $zepp = |\eta_{Z\gamma} (\eta_{j1} + \eta_{j2})/2| < 2.4$
- dphi = $|\phi_{Z\gamma} (\phi_{j1} + \phi_{j2})| > 1.9$

Basic event selection

Suppress FSR

Low m_{jj} control region

VBS Signal region

Special cut added for aQGC

EW signal extraction for signal significance

Selection with the generator-level defines the fiducial volume

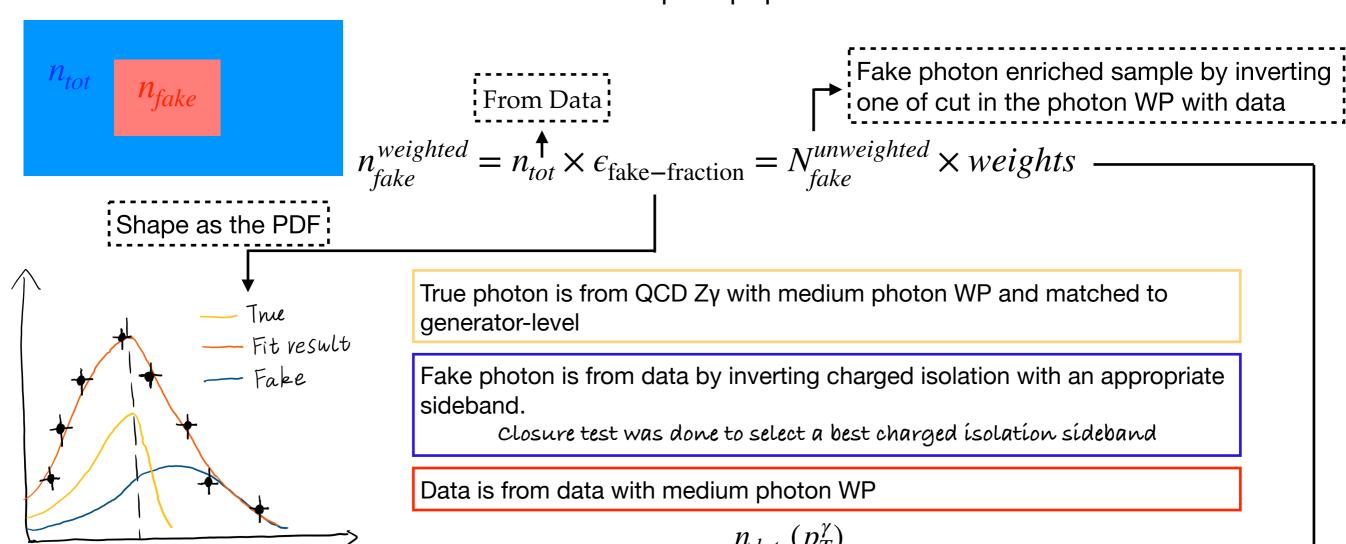


A shower shape variable

Background estimation



- Background processes estimated from simulation are normalized to the best theoretical cross section prediction.
- Irreducible background QCD Zγ normalization is constrained by data in the low m_{ij} control region.
- A data-driven method is used to estimate nonprompt photon contribution.



 $\mathbf{weight}(p_T^{\gamma}) = \frac{n_{data}(p_T^{\gamma})}{N_{fake}^{unweighted}(p_T^{\gamma})} \times \epsilon_{\text{fake-fraction}}(p_T^{\gamma})$



Systematic uncertainties



QCD Factorization and renormalization scale uncertainty

- Exclude the two variations where $(2\mu_0, 0.5\mu_0)$ and $(0.5\mu_0, 2\mu_0)$. μ_0 is the nominal scale.
- Nuisance parameter 1: μ_F only, $(2\mu_0,\mu_0)$ and $(0.5\mu_0,1\mu_0)$

Theoretical

- Nuisance parameter 2: μ_R only, $(\mu_0, 2\mu_0)$ and $(1\mu_0, 0.5\mu_0)$
- Nuisance parameter 3: $\mu_R + \mu_F$ fully correlated, $(2\mu_0, 2\mu_0)$ and $(0.5\mu_0, 0.5\mu_0)$
- Calculated bin-by-bin, correlated between bins, categories, and years

PDF uncertainty

- Standard deviation of the around 100 NNPDF PDF set variations
- Calculated bin-by-bin, correlated between bins, categories, and years

Jet energy resolution&scale uncertainty

Experimental

• Calculated bin-by-bin, correlated between bins and categories, uncorrelated between years

Nonprompt photon uncertainty

- Closure test + Sideband choice + True template choice
- Calculated bin-by-bin, correlated between bins, uncorrelated between categories and years

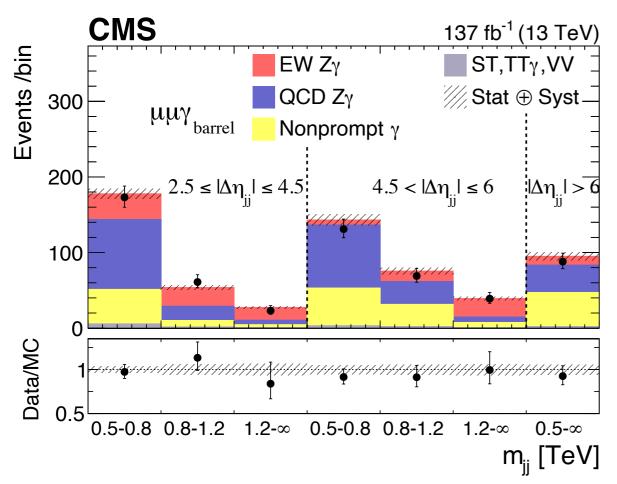
MC Statistical uncertainty

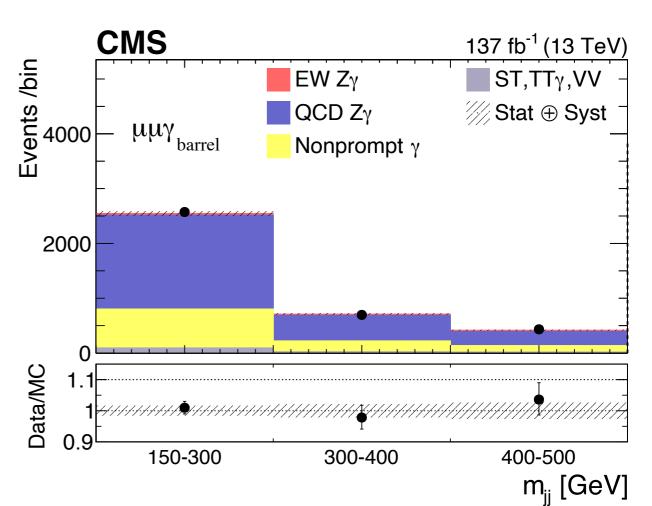
Efficiencies of lepton/photon ID/ISO/Reco, HLT, pileup, L1prefiring and luminosity.



Signal significance







Phys. Rev. D **104**, 072001 (FIG. 4 and 6)

The significance is calculated using a **simultaneous fit** in the signal region with **2D mjj-Δηjj** binning and the control region with **1D mjj binning** in 4 categories for muon/electron choice and barrel photon/endcap photon choice.

• The observed (expected) significance is 9.4 σ (8.5 σ).



Fiducial cross section



$$\sigma_{fiducial-region} = \sigma_{generator} \cdot \mu_{signal-strength}$$

• $\mu_{\text{signal-strength}}$ is the best-fit signal strength, representing the ratio of observed to expected signal yields, which is

$$\square \mu = 1.20^{+0.12}_{-0.12} \text{ (stat)} ^{+0.14}_{-0.12} \text{ (syst)} = 1.20^{+0.18}_{-0.17} \text{ for EW}$$

$$\mathbf{M} = 1.11^{+0.06}_{-0.06} \text{(stat)} ^{+0.10}_{-0.09} \text{(syst)} = 1.11^{+0.12}_{-0.11} \text{ for EW+QCD}.$$

• $\sigma_{generator}$ is the cross section computed by the generator (MadGraph5_aMC@NLO) in the fiducial region which is

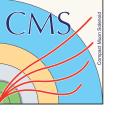
$$\sigma_{\text{generator}} = 4.34 \pm 0.26 \text{ (scale)} \pm 0.06 \text{ (PDF)}$$
 fb for EW

$$\sigma_{\text{generator}} = 13.3 \pm 1.72 \text{ (scale)} \pm 0.10 \text{ (PDF)}$$
 fb for EW+QCD

• $\sigma_{\text{fiducial-region}}$ and its uncertainty is the calculated

$$\sigma_{fid} = 5.21 \pm 0.52 \text{ (stat)} \pm 0.56 \text{ (syst)} = 5.21 \pm 0.76 \text{ fb for EW}$$

$$\sigma_{fid} = 14.7 \pm 0.80 \text{ (stat)} \pm 1.26 \text{ (syst)} = 14.7 \pm 1.53 \text{ fb for EW+QCD}$$



Unfolded differential cross section



Similar with the fiducial cross section measurement, 'unfolding' was performed to revert the 'detector smearing' on the data to get the 'True' distribution.

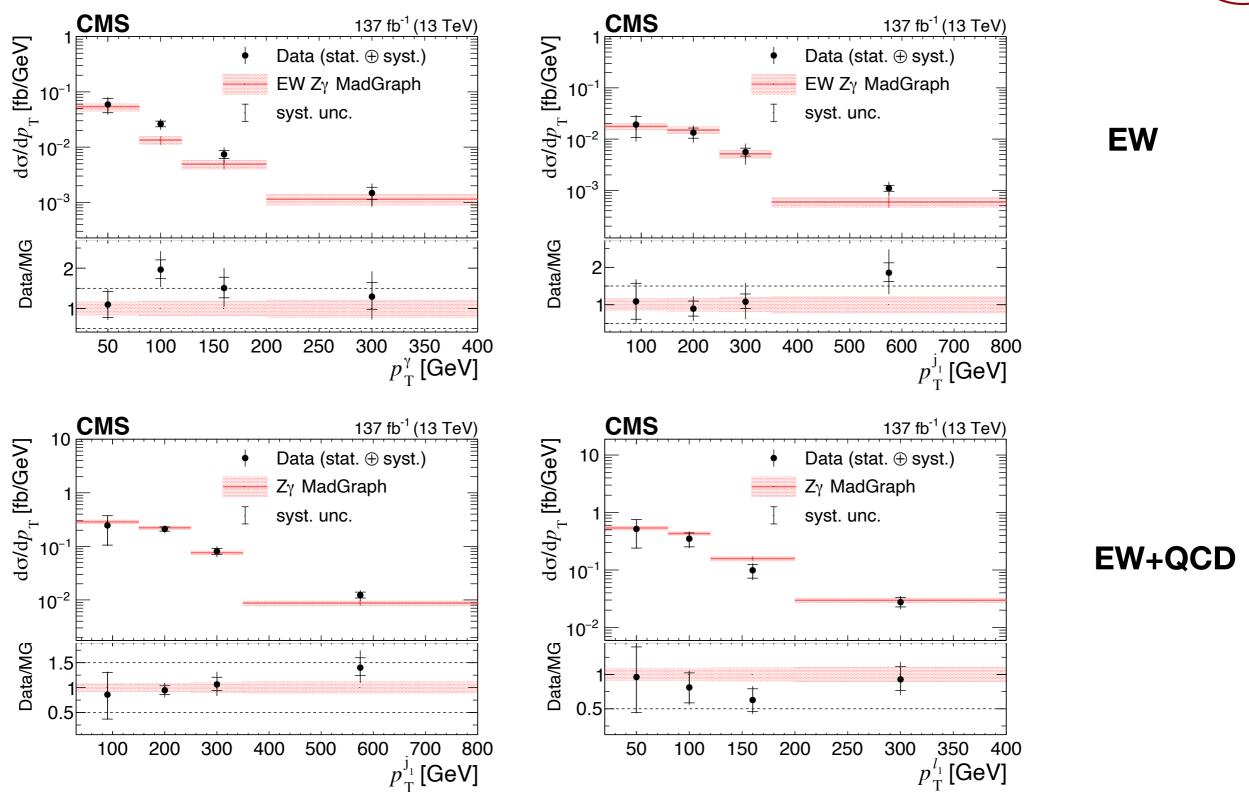
$$\mathcal{L}(\overrightarrow{\mu}; \overrightarrow{\theta}) = \prod_{j} \mathbf{Poisson}(n_{j}; \sum_{i} R_{ji}(\overrightarrow{\theta}) \mu_{i} s_{i}(\overrightarrow{\theta}) + b_{j}(\overrightarrow{\theta})) \cdot \mathcal{N}(\overrightarrow{\theta})$$

- Each reconstructed bin (j) describes the contribution from each truth bin (i) this is the R_{ji} (response matrix).
 - $^{\text{d}}$ Condition number of the R is smaller than about 10, so the regularization is not needed
- Same uncertainties with significance measurement are applied
- 1D variables of leading lepton, photon and jet, and 2D variable $m_{jj} \Delta \eta_{jj}$ are measured



Unfolded differential cross section





Phys. Rev. D 104, 072001 (FIG. 8)

Within the uncertainties, the measurements agree with the SM predictions.



aQGC limits



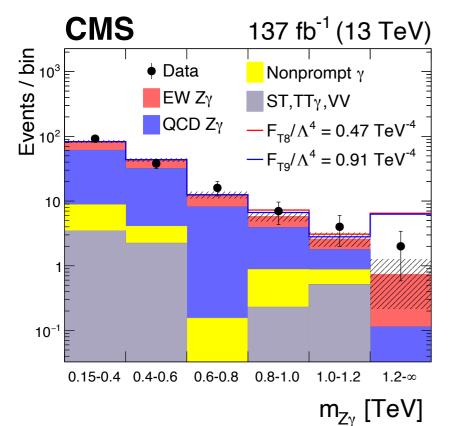
SM Lagrangian can be extended with higher dimensional operators maintaining SU(2)×U(1) gauge symmetry:

$$L_{EFT} = L_{SM} + \sum_{i} \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}^{(6)} + \frac{c_i^{(8)}}{\Lambda^2} \mathcal{O}^{(8)} + \dots$$

mmetry:
$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_{i} \frac{c_{i}^{(6)}}{\Lambda^{2}} \mathcal{O}^{(6)} + \frac{c_{i}^{(8)}}{\Lambda^{2}} \mathcal{O}^{(8)} + \dots.$$
 Test statistic $t_{\alpha_{test}} = -2ln \frac{\mathcal{L}(\alpha, \hat{\theta})}{\mathcal{L}(\hat{\alpha}, \hat{\theta})}$: follows χ^{2} distribution;

Extract the limits directly using the profiling log likelihood ratio $\Delta NLL = t_{\alpha_{test}}/2$;

The 95% CL limit corresponds $2\Delta NLL=3.84$.



The most stringent limit for operator T_9



aQGC limits



As the sensitivity on the T_i operators of VBS Z γ , we show the comparison of the limits of T_i from recent public VBS results with the full Run2 data

Operator	SMP-20-016 VBS Zγ	SMP-20-001 VBS ZZ	SMP-19-012 VBS W±W±
f _{T0}	-0.64 , 0.57	-0.24 , 0.22	-0.28 , 0.31
f _{T1}	-0.81 , 0.90	-0.31 , 0.31	-0.12 , 0.15
f _{T2}	-1.68 , 1.54	-0.63 , 0.59	-0.38 , 0.50
f _{T5}	-0.58 , 0.64		
f _{T6}	-1.30 , 1.33		
f _{T7}	-2.15 , 2.43		
f _{T8}	-0.47 , 0.47	-0.43 , 0.43	
f _{T9}	-0.91 , 0.91	-0.92 , 0.92	

Similar sensitivity on T_8 and T_9 between VBS Z γ and VBS ZZ, which is expected, as the T_8 and T_9 give rise to QGCs only containing the neutral gauge bosons.



Summary



- ✓ Overall significance is far more 5σ .
- √ Fiducial cross section measurement reported
- ✓ Unfolded differential cross section as functions of leading lepton/jet/ photon pT and m_{jj}- $\Delta\eta_{jj}$
- \checkmark AQGC limits for operator $M_{0-7},\,T_{0-2},\,{\rm and}\,\,T_{5-9}$.
 - \checkmark Limit for T_9 is the most stringent limit to date



Backup



variables	2016	2017	2018
p_T^{γ}	1.08	1.12	1.21
$p_T^{j_1}$	1.35	1.41	1.44
$p_T^{l_1}$	1.09	1.09	1.11
m_{jj} - $\Delta \eta jj$	1.87	1.97	1.95

variables	2016	2017	2018
p_T^{γ}	1.16	1.41	1.37
$p_T^{j_1}$	1.33	1.41	1.39
$p_T^{l_1}$	1.10	1.35	1.16
m_{jj} - $\Delta \eta jj$	1.93	2.32	2.09

Condition Number of R for EW

Condition Number of R for EW+QCD

If the condition number is small (~10), then the problem is well-conditioned and can most likely be solved using the unregularized maximum likelihood estimate (MLE). This happens when the resolution effects are small and R is almost diagonal. If on the other hand, the condition number is large (~10⁵) then the problem is ill-conditioned and the unfolded estimator needs to be regularized.



Backup



Building blocks:

- $D_{\mu}\Phi$: Higgs doublet field, affects the coupling of longitudinal modes of the gauge bosons.
- + $\hat{W}_{\mu
 u}$, $\hat{B}_{\mu
 u}$: Field strength tensors

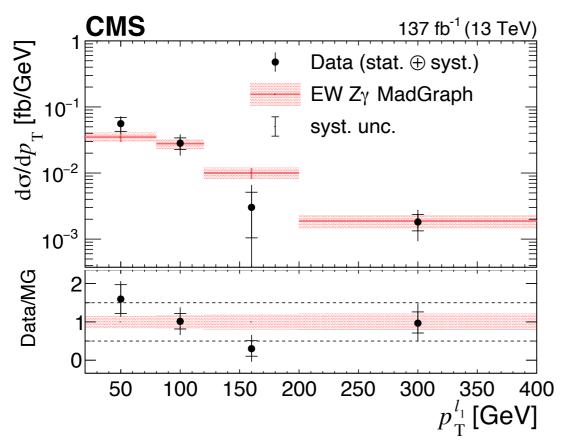
Dimension-8 operators (only field strength/mixed)

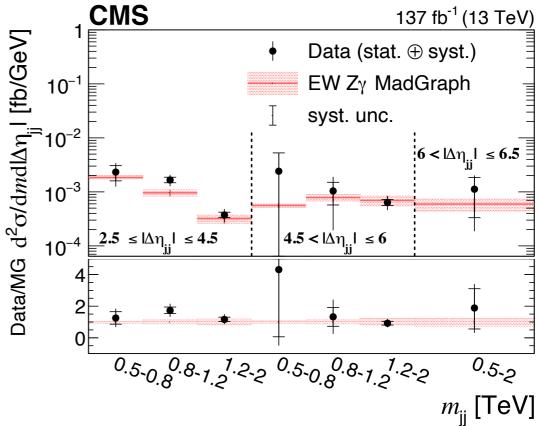
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\mathcal{O}_{M,0} = \operatorname{Tr}\left[W_{\mu\nu}W^{\mu\nu}\right] \cdot \left[(D_{\beta}\Phi)^{\dagger}D^{\beta}\Phi\right]
\mathcal{O}_{T,0} = \operatorname{Tr}\left[W_{\mu\nu}W^{\mu\nu}\right] \cdot \operatorname{Tr}\left[W_{\alpha\beta}W^{\alpha\beta}\right],
                                                                                                                                    \mathcal{O}_{M,1} = \operatorname{Tr}\left[W_{\mu\nu}W^{\nu\beta}\right] \cdot \left[(D_{\beta}\Phi)^{\dagger}D^{\mu}\Phi\right]
\mathcal{O}_{T,1} = \operatorname{Tr} \left[ W_{\alpha \nu} W^{\mu \beta} \right] \cdot \operatorname{Tr} \left[ W_{\mu \beta} W^{\alpha \nu} \right] ,
                                                                                                                                   \mathcal{O}_{M,2} = \left[ B_{\mu\nu} B^{\mu\nu} \right] \cdot \left[ (D_{\beta} \Phi)^{\dagger} D^{\beta} \Phi \right] ,
\mathcal{O}_{T,2} = \operatorname{Tr}\left[W_{\alpha\mu}W^{\mu\beta}\right] \cdot \operatorname{Tr}\left[W_{\beta\nu}W^{\nu\alpha}\right],
                                                                                                                                   \mathcal{O}_{M,3} = \left[ B_{\mu\nu} B^{\nu\beta} \right] \cdot \left[ (D_{\beta} \Phi)^{\dagger} D^{\mu} \Phi \right] ,
\mathcal{O}_{T.5} = \operatorname{Tr} \left[ W_{\mu\nu} W^{\mu\nu} \right] \cdot B_{\alpha\beta} B^{\alpha\beta} ,
                                                                                                             \mathcal{O}_{M,4} = \left[ (D_{\mu}\Phi)^{\dagger} W_{\beta\nu} D^{\mu}\Phi \right] \cdot B^{\beta\nu} ,
\mathcal{O}_{T,6} = \operatorname{Tr}\left[W_{\alpha\nu}W^{\mu\beta}\right] \cdot B_{\mu\beta}B^{\alpha\nu},
                                                                                                              \mathcal{O}_{M,5} = \left[ (D_{\mu} \Phi)^{\dagger} W_{\beta \nu} D^{\nu} \Phi \right] \cdot B^{\beta \mu} ,
\mathcal{O}_{T,7} = \operatorname{Tr}\left[W_{\alpha\mu}W^{\mu\beta}\right] \cdot B_{\beta\nu}B^{\nu\alpha},
\mathcal{O}_{T,8} = B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta}
                                                                                                                                   \mathcal{O}_{M,6} = \left[ (D_{\mu}\Phi)^{\dagger} W_{\beta\nu} W^{\beta\nu} D^{\mu} \Phi \right] ,
 \mathcal{O}_{T,9} = B_{\alpha\mu} B^{\mu\beta} B_{\beta\nu} B^{\nu\alpha} .
                                                                                                                                     \mathcal{O}_{M,7} = \left[ (D_{\mu} \Phi)^{\dagger} W_{\beta \nu} W^{\beta \mu} D^{\nu} \Phi \right] ,
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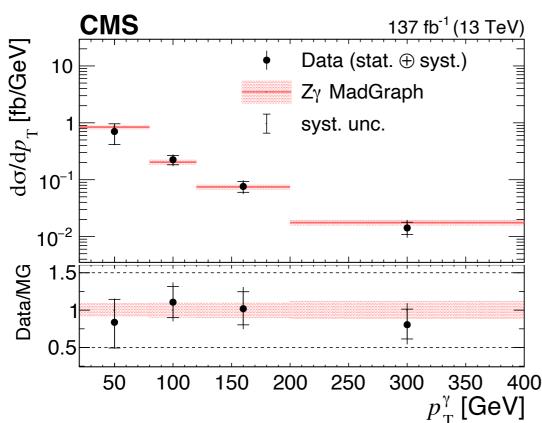


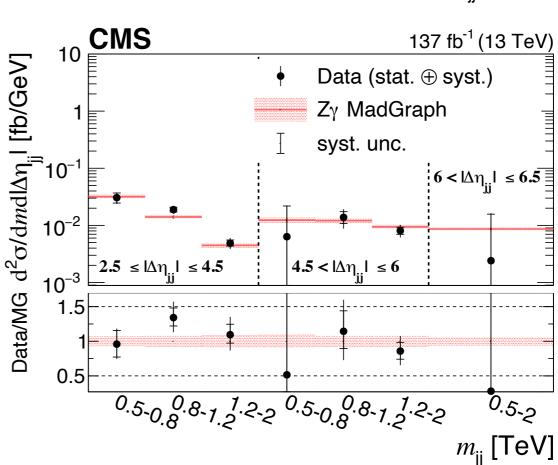
Backup— Unfolding for EW+QCD













Backup



sigmalEtalEta is the log energy weighted RMS of the shower in units of crystals

$$- \sigma_{i\eta i\eta} = \sqrt{\left(\frac{\Sigma_i^{5\times5} w_i (\eta_i - \overline{\eta}_{5\times5})^2}{\Sigma_i^{5\times5} w_i}\right)}$$

$$- w_i = 4.7 + \ln \frac{E_i}{E_{5 \times 5}}$$

- this is effectively a noise cut, each crystal needs to have > 0.9% of 5x5 energy
- means that very low energy electrons are sensitive to noise as 0.9% of a small number brings it below noise threshold
- E_i = energy of crystal, $E_{5\times5}$ energy of 5x5
 - likewise for η
- η is in units of crystals, not absolute η
 - endcap uses $(ix^2 + iy^2)^{1/2}$ to get η in terms of crystals
- normalised to 0.01745 in barrel and 0.0447 in endcap
- cut effectively means that all the energy is within two crystals