



Collective expansion in pp collisions using the Tsallis statistics

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This work will appear in arXiv soon.

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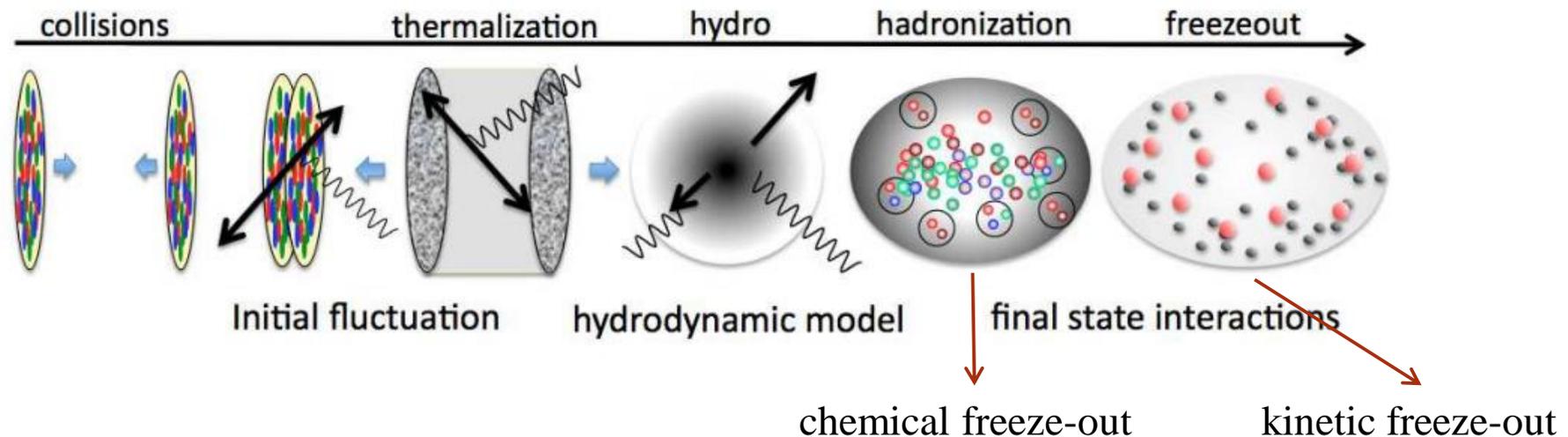
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Background and motivation

- Quantum-Chromodynamics (QCD) predicts that at high temperature and energy density there exists a hot and dense strongly interacting matter denoted as quark-gluon plasma(QGP).

[E.V. Shuryak, *Phys. Rept.* 61, 71 (1980)]

- QGP is expected to be produced in ultrarelativistic heavy-ion collisions.





Background and motivation

- The transverse momentum (p_T) spectra of identified particles are significant observables in high energy collisions.
- They are utilized to investigate the dynamics of particle productions.
- In the low p_T region, particle productions are governed by soft physics and described by some nonperturbative theory or model, such as the Boltzmann-Gibbs blast-wave model (BGBW).

$$\frac{d^2N}{2\pi p_T dp_T dy} \propto m_T \int_0^{R_0} r dr K_1 \left(\frac{m_T \cosh \rho}{T} \right) I_0 \left(\frac{p_T \sinh \rho}{T} \right) \quad \beta_r = \beta_s \left(\frac{r}{R_0} \right)^n \quad \rho = \tanh^{-1}(\beta_r)$$



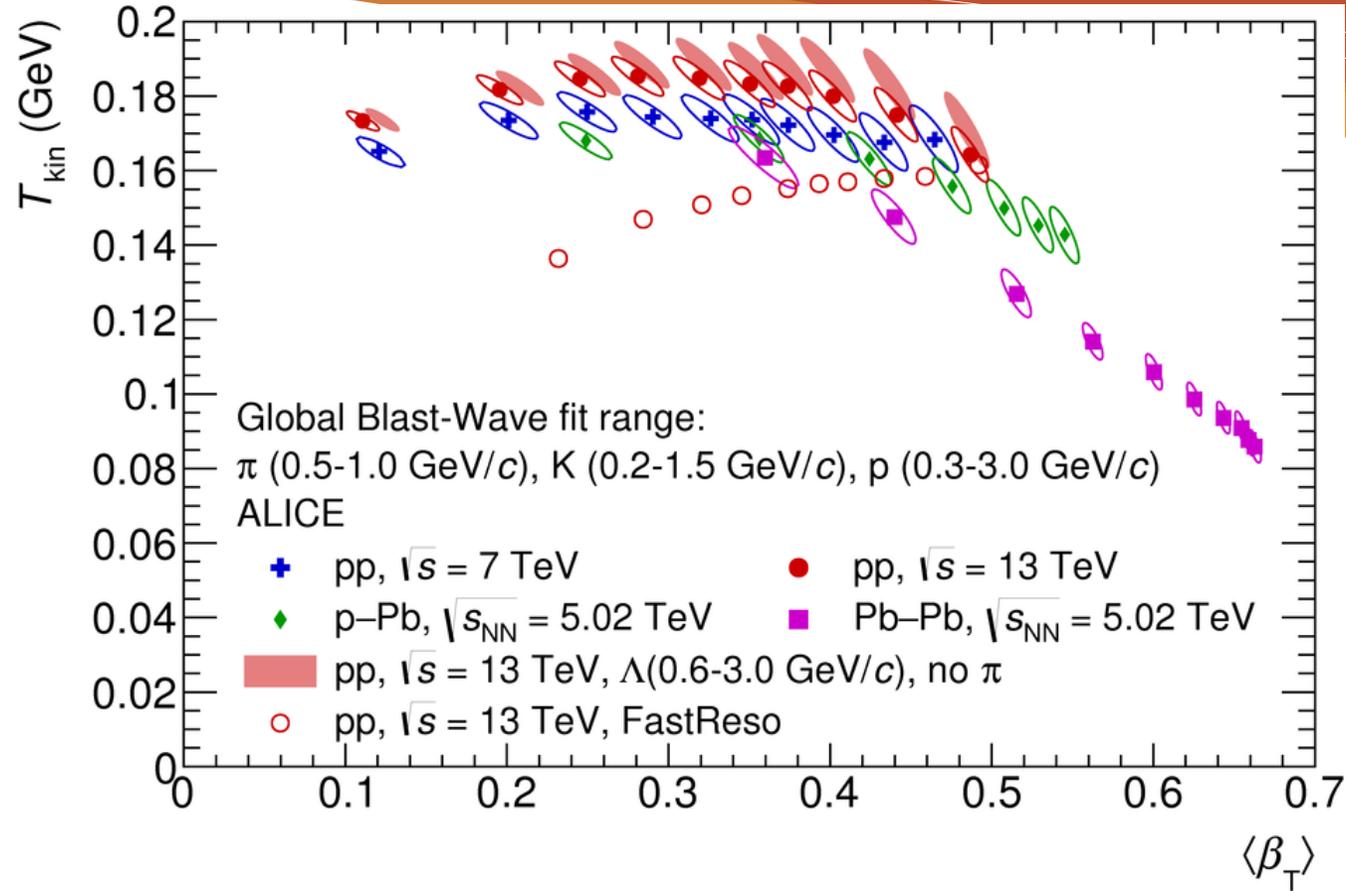
Background and motivation

- In pp collisions at 7 and 13 TeV, the p_T spectra of identified particles get harder with the increase of the charged-particle multiplicity, with the effect being more obvious for particles with larger mass.
- This trend is highly similar to that observed in the evolution of the spectra in p-Pb and Pb-Pb collisions.
- Moreover, double-ridge structures have also been observed in high-multiplicity pp collisions. [G. Aad *et al.* *Phys. Rev. Lett.* 116, 172301 (2016)]
- These collective phenomena are reminiscent to observations attributed to the creation of QGP in Pb-Pb collisions.



Background and motivation

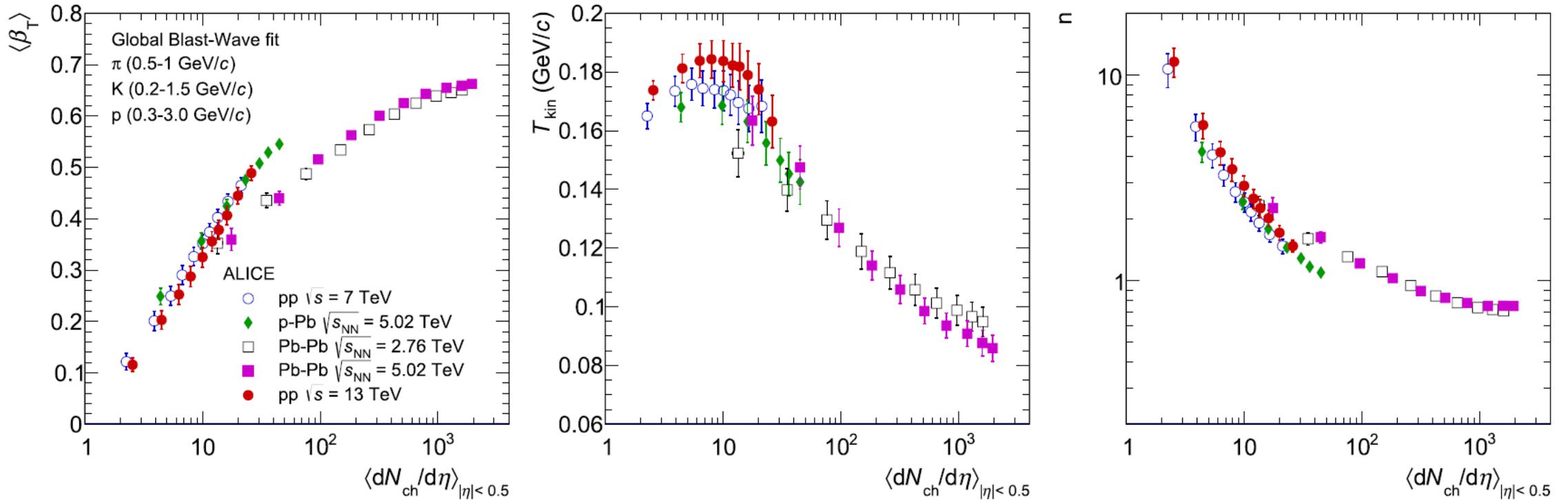
pp, $\sqrt{s} = 7$ and 13 TeV
Global Blast-wave fit
Different multiplicity



S. Acharya *et al.* (ALICE Collaboration), Eur. Phys. J. C (2020) 80:693



Background and motivation



S. Acharya *et al.* (ALICE Collaboration), *Eur. Phys. J. C* (2020) 80:693

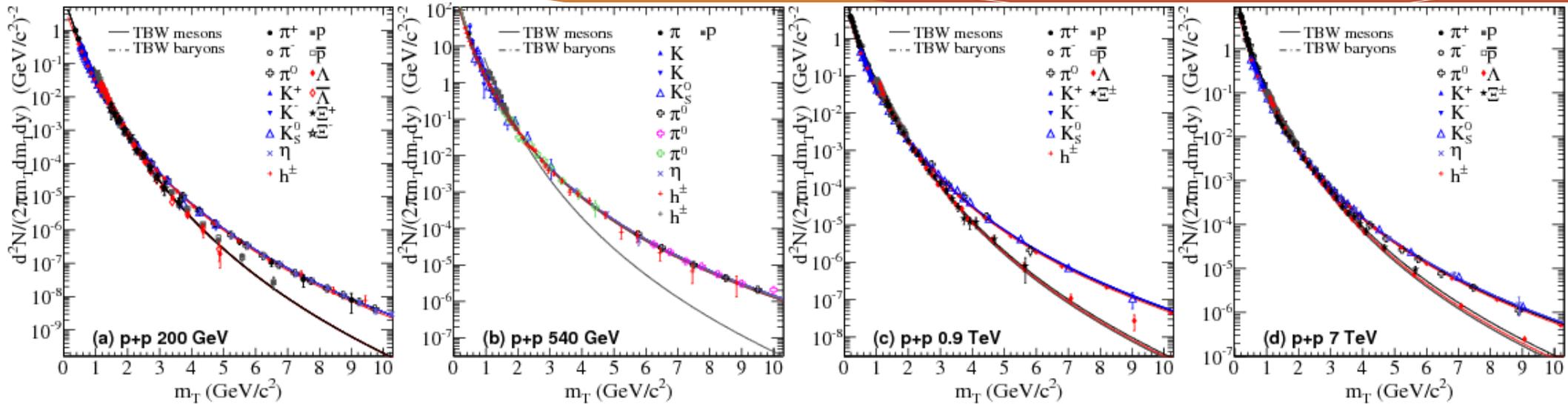


Background and motivation

- In the BGBW model, there is a strong assumption that the system will reach a local thermal equilibrium at some instant of time and then undergoes the hydrodynamic evolution.
- However, in fact the initial condition for the hydrodynamic evolution fluctuates from event to event.
- In order to take this fluctuation into account, the authors have changed the sources of particle emission in the BGBW model from the Boltzmann distribution to the Tsallis distribution (TBW). *Zebo Tang et al. , Phys. Rev. C 79, 051901 (2009)*



Background and motivation



\sqrt{s}	$\langle\beta\rangle$	T (MeV)	$q_M - 1$	$q_B - 1$	χ^2/ndf
7 TeV	0.320 ± 0.005	70.3 ± 0.8	0.1314 ± 0.0003	0.1035 ± 0.0008	490/431
900 GeV	0.264 ± 0.005	74.6 ± 0.5	0.1127 ± 0.0003	0.0827 ± 0.0008	545/501
540 GeV	$0.000^{+0.105}_{-0.000}$	81.8 ± 0.6	0.1158 ± 0.0007	0.0841 ± 0.0036	205/168
200 GeV	$0.000^{+0.124}_{-0.000}$	92.3 ± 2.7	0.0946 ± 0.0006	0.0743 ± 0.0015	268/268

K. Jiang *et al.*, Phys. Rev. C **91**, 024910 (2015)



Background and motivation

- ALICE and CMS Collaboration have published the identified particle transverse momentum spectra in pp collisions at $\sqrt{s_{NN}} = 0.9, 2.76, 5.02, 7, 13$ TeV.
- The ALICE collaboration has also presented the identified particle spectra in pp collisions at 7 and 13 TeV with different charged-particle multiplicities.

Motivations:

- Investigate the dependence of the radial flow $\langle\beta\rangle$, Tsallis temperature T and the degree of off-equilibrium q on the collision energy in pp collisions.
- Predict the identified particle spectra in pp collisions at 8 and 14 TeV.
- Investigate the dependence of the parameters $\langle\beta\rangle$, T and q on the multiplicity.
- Shed some light on the possible underlying mechanism for particle productions.



Method

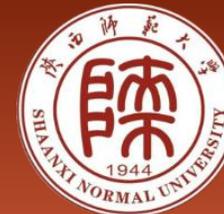
Tsallis blast-wave model

$$\frac{d^2 N}{2\pi p_T dp_T dy} \propto m_T \int_{-y_b}^{+y_b} \exp\left(\sqrt{y_b^2 - y_s^2}\right) \cosh(y_s) dy_s \int_{-\pi}^{+\pi} d\phi$$

$$\times \int_0^{R_0} r dr \left[1 + \frac{q-1}{T} (m_T \cosh(y_s) \cosh(\rho) - p_T \sinh(\rho) \cos(\phi))\right]^{\frac{-1}{q-1}}$$

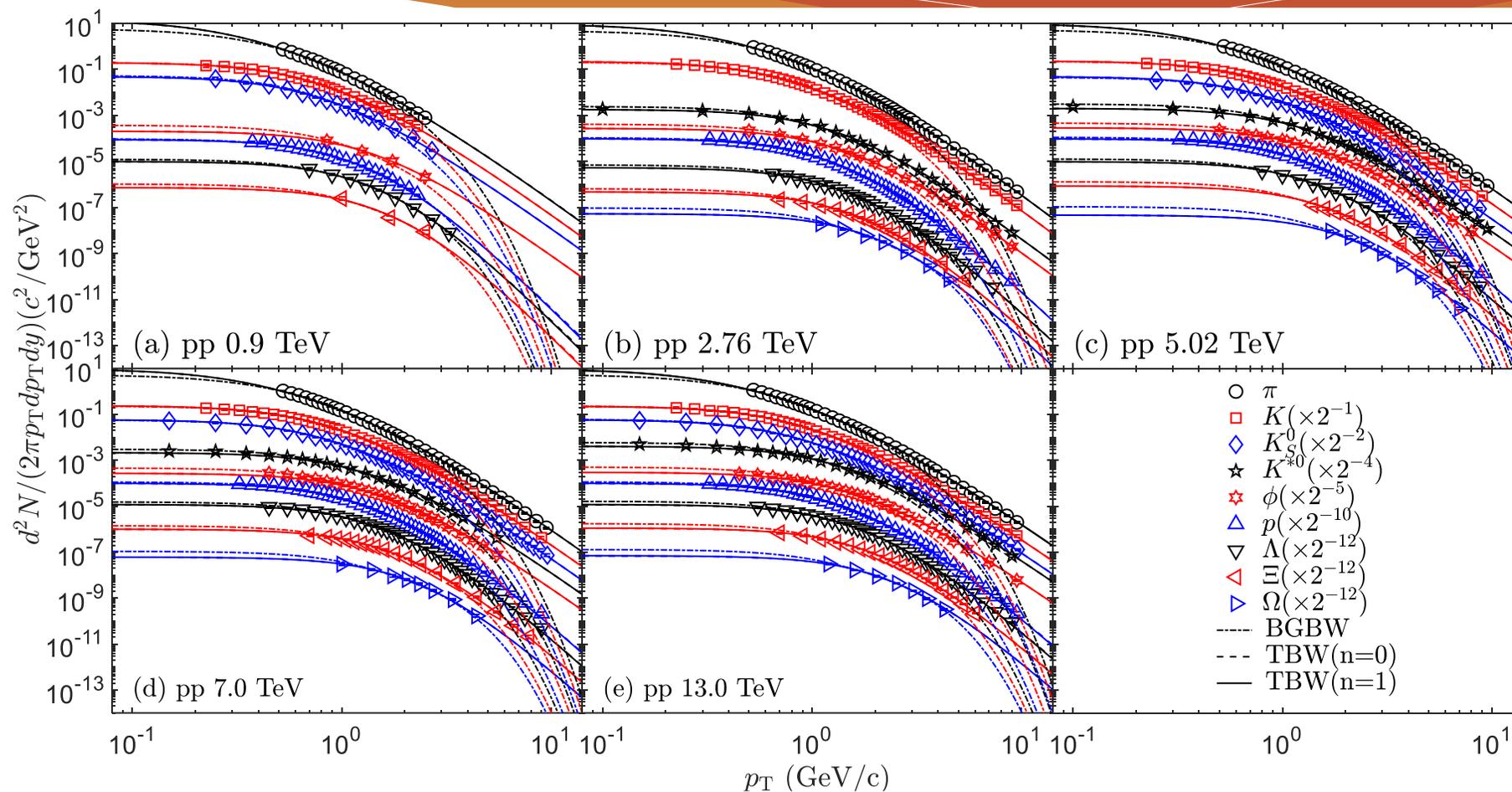
$$\rho_0 = \tanh^{-1}\langle\beta\rangle \quad \langle\beta\rangle = \frac{2}{n+2} \beta_s \quad \beta_r = \beta_s \left(\frac{r}{R_0}\right)^n$$

- In our research, we consider two kinds of velocity profiles for TBW model:
Linear profile, $n=1$; Constant profile, $n=0$.
- Free parameters: non-extensive parameter q , temperature T and average velocity $\langle\beta\rangle$.
- $\langle\beta\rangle$ and T are common for all particles. While q_M (q_B) are the same for all of the mesons (baryons).



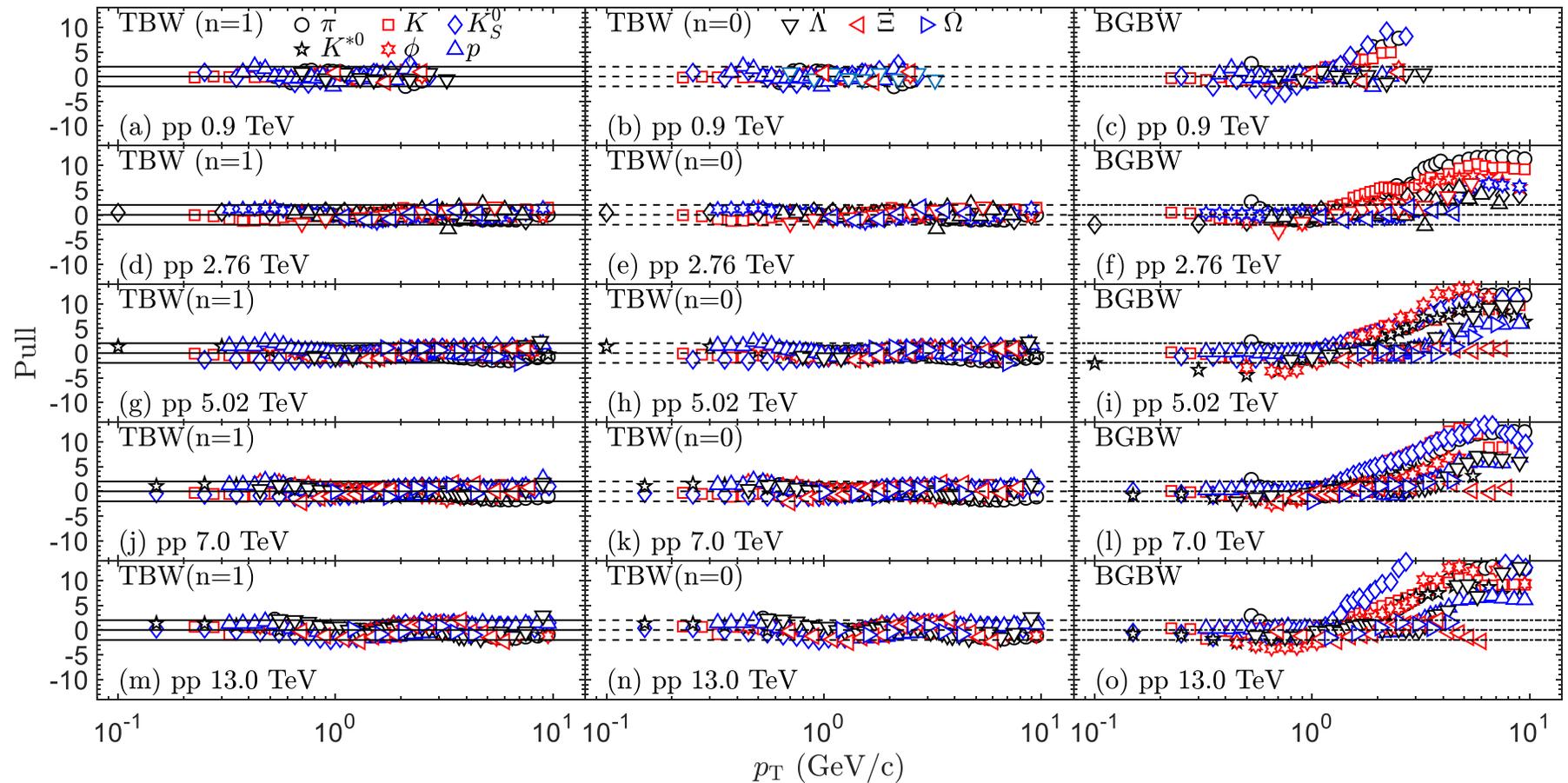
Results and discussions

pp 0.9, 2.76,
5.02, 7.0, 13.0 TeV
Transverse momentum
spectra





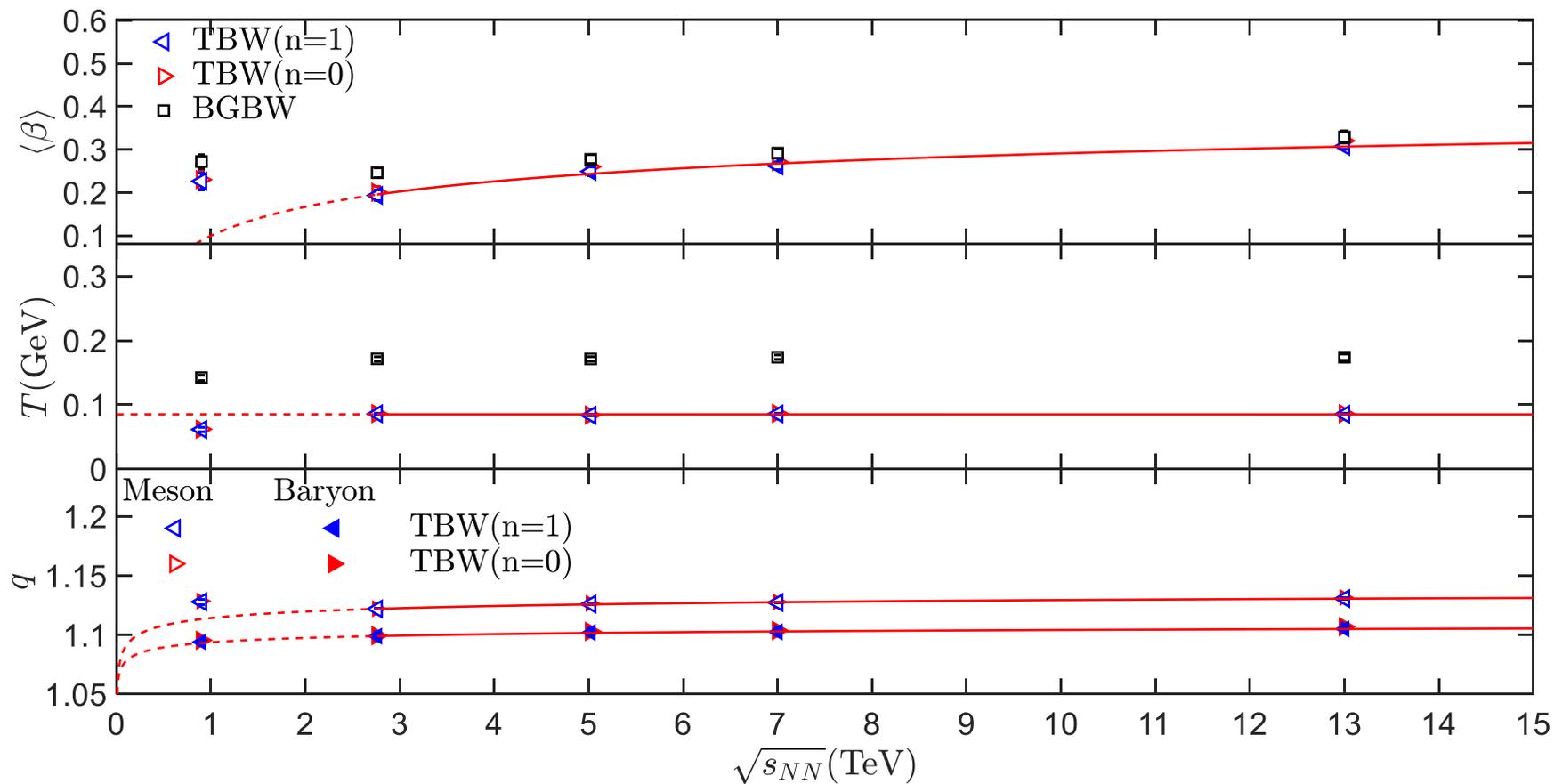
Results and discussions



$$\text{Pull} = \frac{\text{data} - \text{fit}}{|\Delta \text{data}|}$$



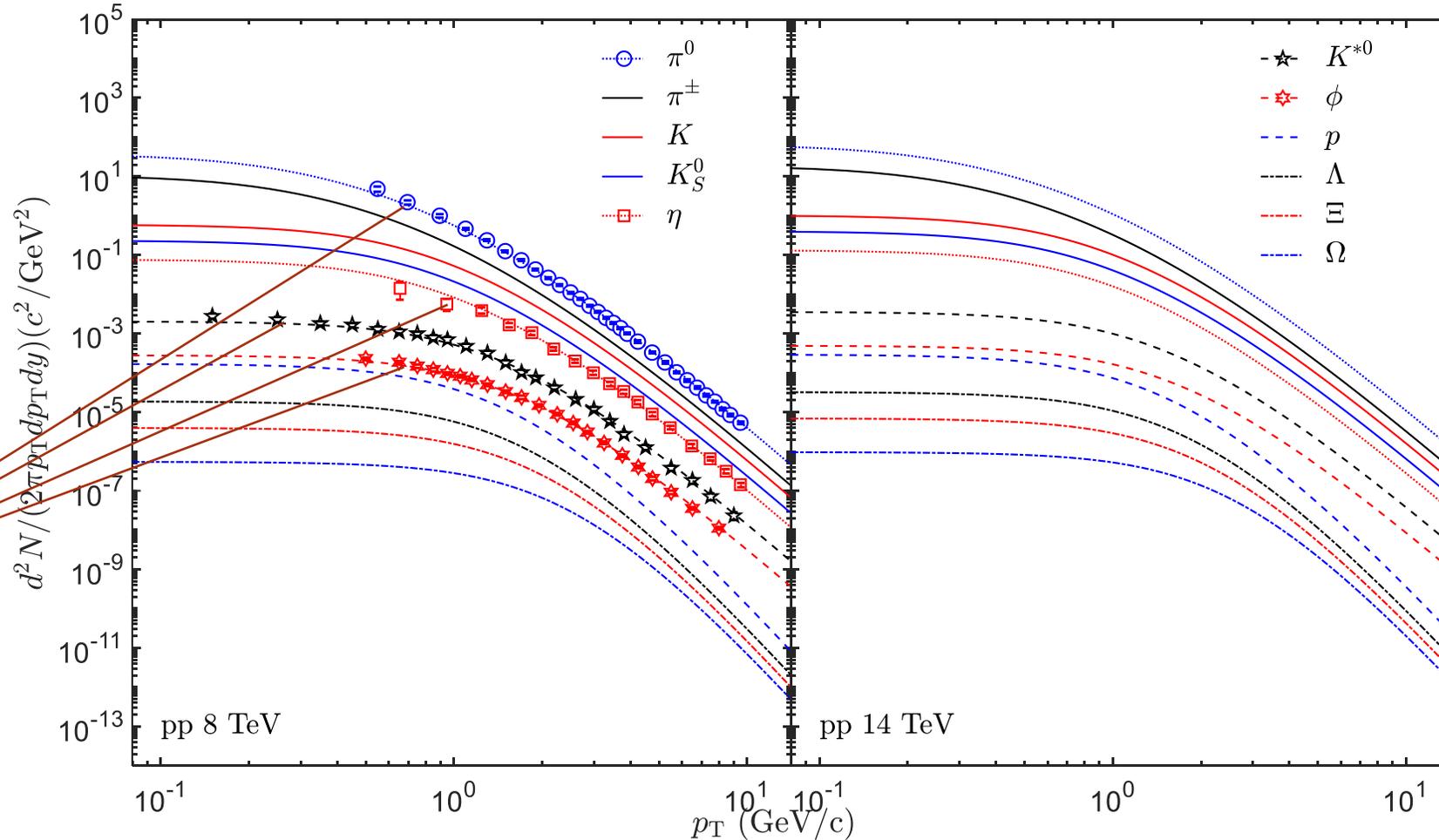
Results and discussions





Results and discussions

Predicted spectra
in pp collisions
at 8 and 14 TeV.



Confirm the prediction

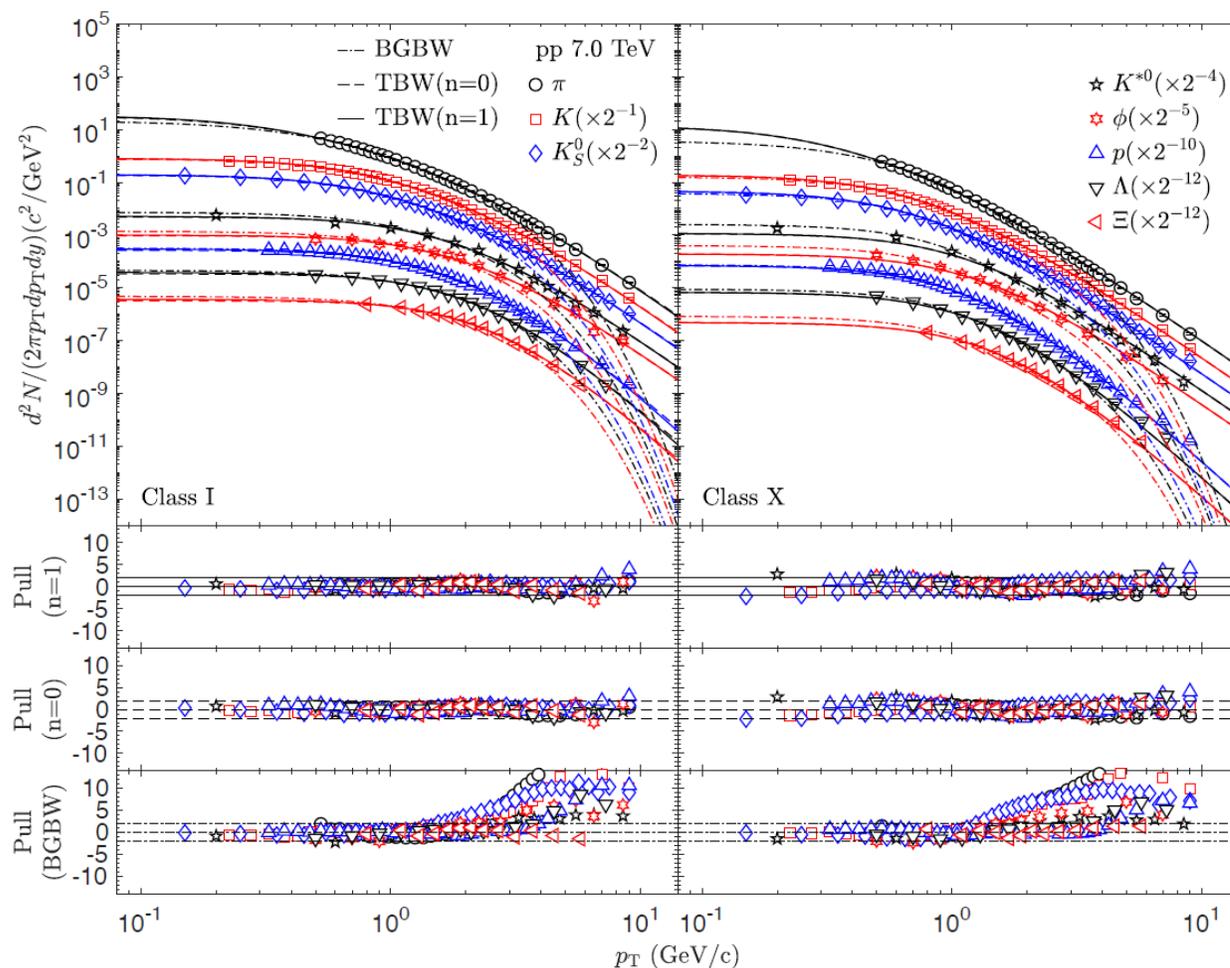
Predicted value: At 8 TeV, $\langle\beta\rangle = 0.2762$, $T = 85.4$ MeV, $q_M = 1.1282$, $q_B = 1.1033$
 At 14 TeV, $\langle\beta\rangle = 0.3108$, $T = 85.4$ MeV, $q_M = 1.1308$, $q_B = 1.1051$

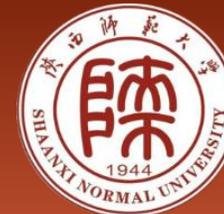


Results and discussions

pp 7.0 TeV Class I, X

	$\langle dN_{ch}/d\eta \rangle$
Class I	21.3 ± 0.6
Class X	2.26 ± 0.12

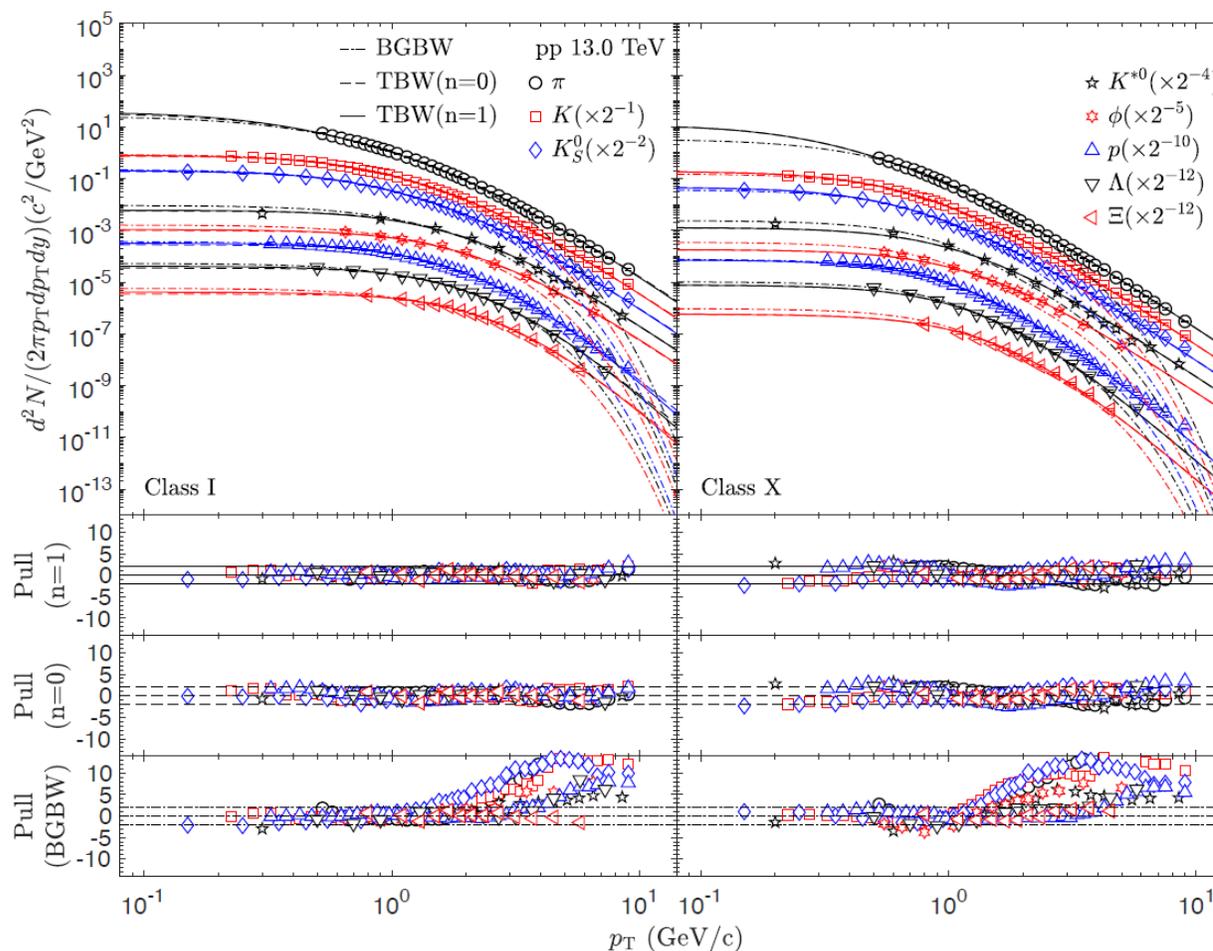




Results and discussions

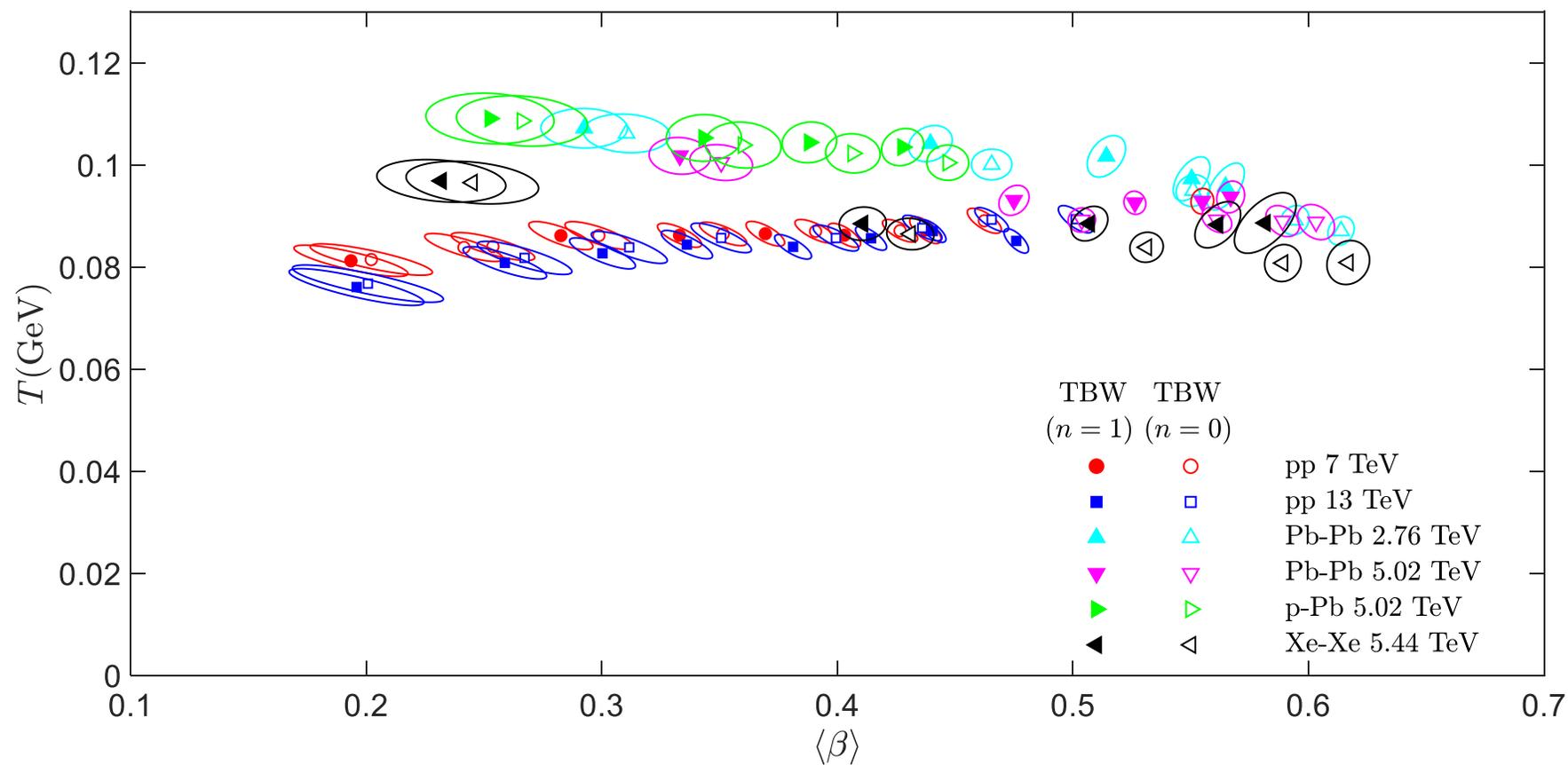
pp 13.0 TeV Class I, X

	$\langle dN_{ch}/d\eta \rangle$
Class I	26.02 ± 0.35
Class X	2.55 ± 0.04



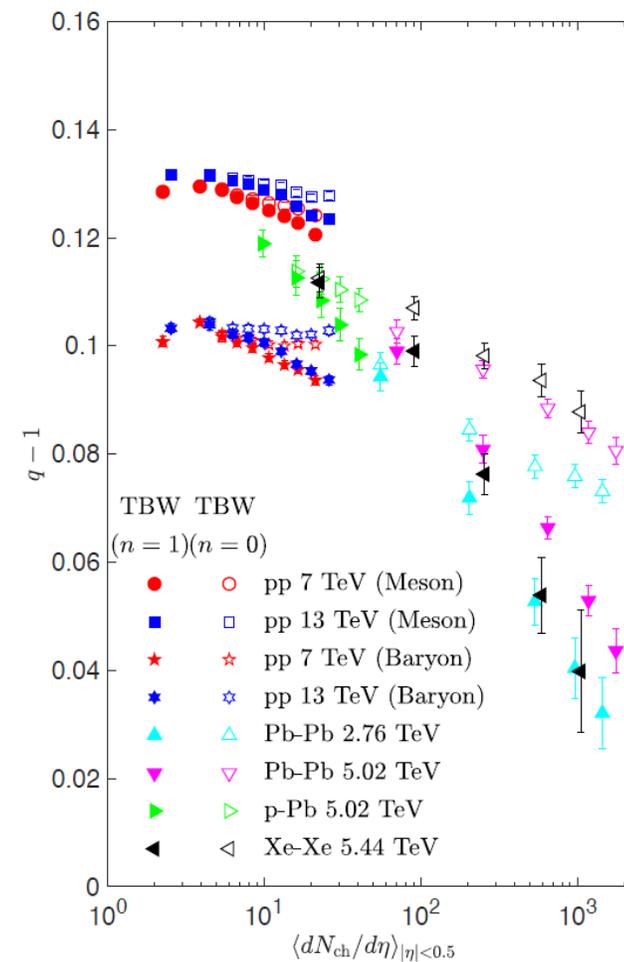
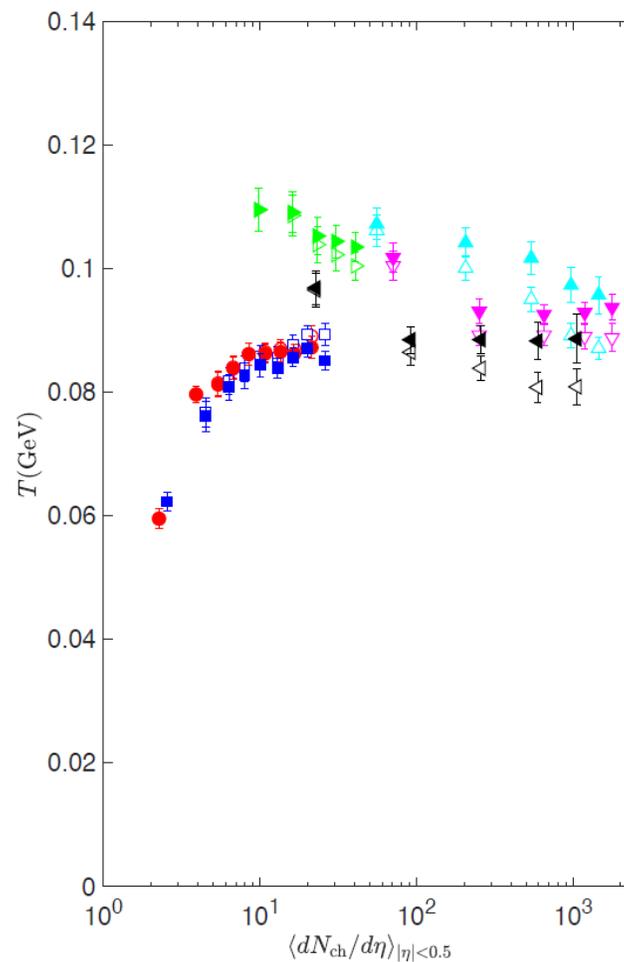
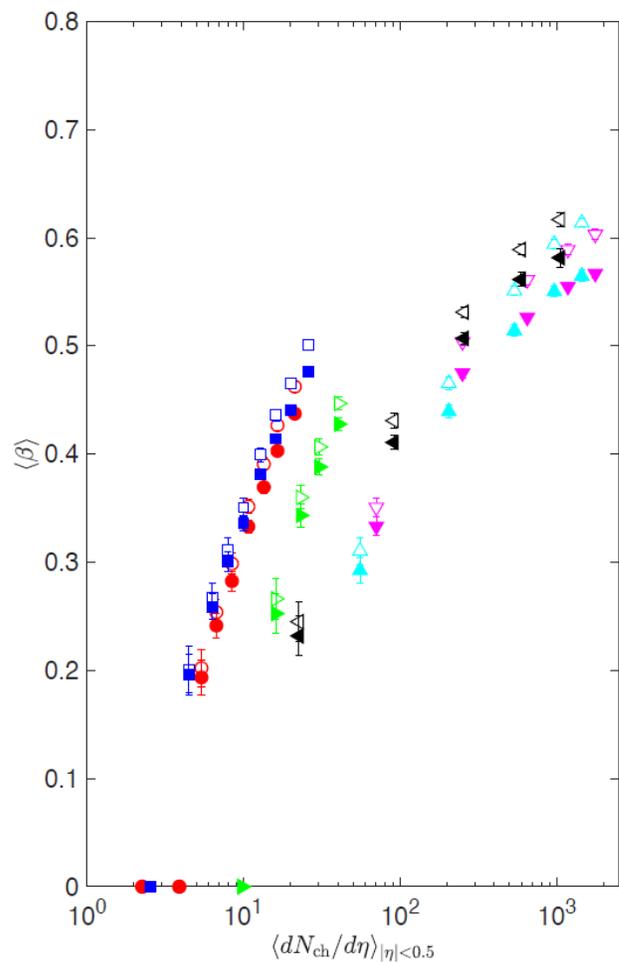
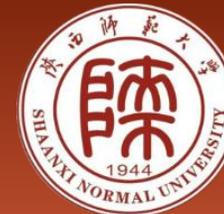


Results and discussions



Guorong Che *et al.* J. Phys. G: Nucl. Part. Phys. 48 095103(2021)

Results and discussions





Summary

- The $\langle\beta\rangle$ and the q increases with \sqrt{s} , while the T first increases with energy up to 2.76 TeV and then saturates around 85 MeV.
- In pp collisions both at 7 and 13 TeV, with the increase of multiplicities, T grows with $\langle\beta\rangle$ until it saturates at around 89 MeV.
- At both energies, $\langle\beta\rangle$ increases with the multiplicity while the q shows the opposite behavior, which is similar to that observed in Pb-Pb, Xe-Xe and p-Pb collisions at the LHC energy regime. However, T increases with the multiplicity, which is opposite to the trend observed in Pb-Pb, Xe-Xe and p-Pb collisions.
- At similar charged-particle multiplicities, $\langle\beta\rangle$ in pp collisions is larger than those in Pb-Pb, Xe-Xe, which indicates that the size of the colliding system might have significant effects on the final state particle dynamics.

Thanks!



Back up

$$\lambda_f \propto \frac{1}{\pi d^2 n_f},$$

For $R_f \gg R_0$ and $\tau_f \gg \tau_0$, we get $R_f = \langle \beta \rangle \tau_f$,

$$n_f = \frac{N}{V_f} = \frac{N}{4/3\pi R_f^3};$$

$$\lambda_f \propto \langle \beta \rangle^3 \tau_f^3 N^{-1}.$$

$$R_f = R_0 + \int_{\tau_0}^{\tau_f} \beta_s dt.$$

$$T_f = T_0 \left(\frac{\tau_0}{\tau_f} \right)^{1/3};$$

$$\int_{\tau_0}^{\tau_f} \beta_s dt = \beta' \cdot (\tau_f - \tau_0) \propto \langle \beta \rangle \cdot (\tau_f - \tau_0).$$

$$T_f \propto T_0 \left(\frac{\langle \beta \rangle^3}{N \lambda_f} \right)^{1/9}$$



Back up

