

Dispersive Analysis of the Primakoff Reaction

$$\gamma K \rightarrow K\pi$$

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Chiral Dynamics 2021



HISKP (Theorie)
Bonn University



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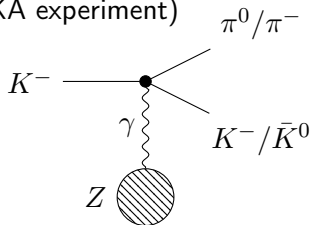


[Dax, DS and Kubis; Eur. Phys. J. C **81** (2021) 221]



Motivation

- pion production in the **Coulomb field** of a heavy nucleus
- $\gamma^{(*)}\pi \rightarrow \pi\pi$ investigated [Hoferichter et al., 2012, 2017], [Niehus et al., 2021] [Niehus, Talk today at 10.50pm]
- **COMPASS++ experiment** planned (+ OKA experiment)
- upgrade from **pion** beam to **kaon** beam

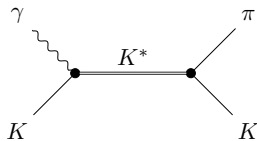
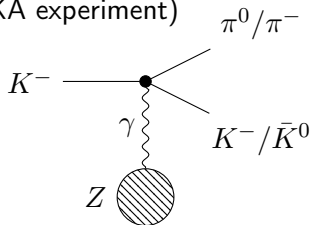


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- upgrade from **pion** beam to **kaon** beam
- combine knowledge about
 - **chiral anomaly** at $s = t = u = 0$

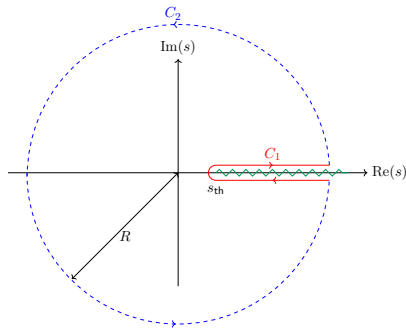
$$F_{KK\pi\gamma} = \frac{e}{4\pi^2 F_\pi^3}$$

- **resonances** ($K^*(892)$) at higher energies (**radiative couplings**)



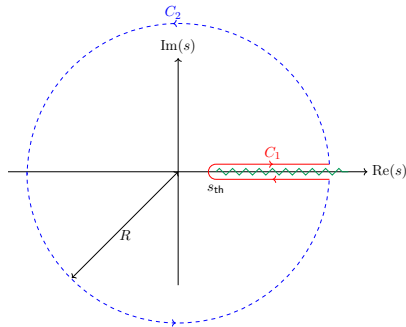
- **analyticity** (\simeq **causality**) & **Cauchy's integral formula**

$$f(s) = \frac{1}{2\pi i} \oint_{\partial U} \frac{f(s')}{s' - s} ds'$$



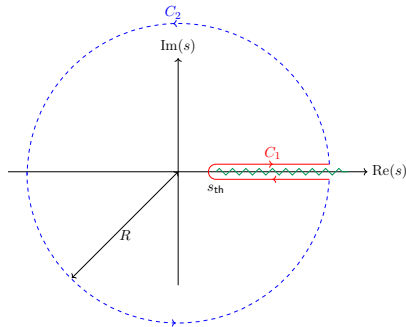
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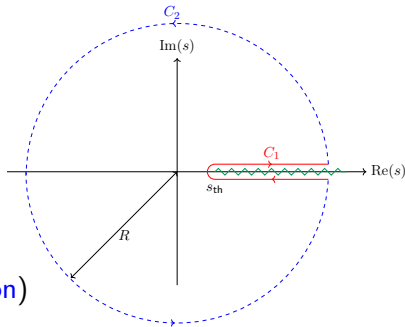
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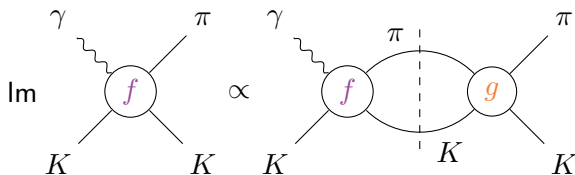
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- unitarity relation (\simeq prob. conservation)

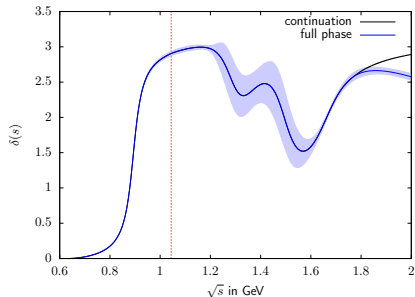
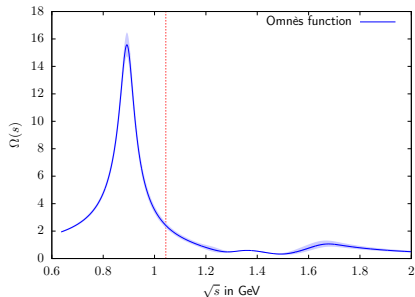
$$\text{Im} f(s) \propto f(s) \cdot g^*(s)$$



- obeys Watson's final state theorem
[Watson, 1954]

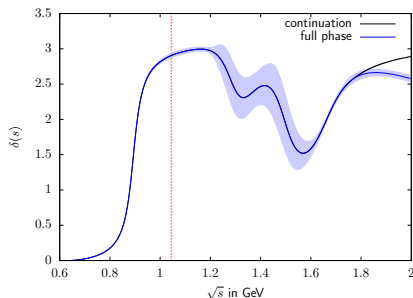
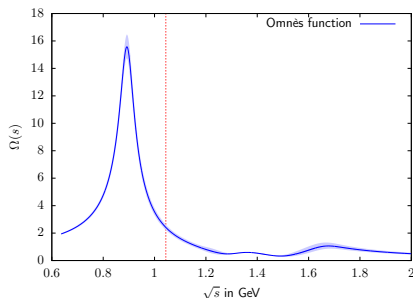
Homogeneous Omnès problem

- $K\pi$ P -wave phase shift from [Peláez and Rodas, 2016]
- $I = 1/2$ phase shift contains $K^*(892)$, $K^*(1410)$ and $K^*(1680)$
- very well constrained up to the $K\eta$ threshold
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- $I = 3/2$ phase shift is $|\delta(s)| < 3^\circ$ for $s < (1.74 \text{ GeV})^2$
- approximate it with $\delta(s) = 0 \Rightarrow \Omega(s) = 1$

Reconstruction Theorem

- separate **kinematic prefactor** $\mathcal{M} = i\varepsilon_{\mu\nu\alpha\beta}\epsilon^\mu p_1^\nu p_2^\alpha p_0^\beta \mathcal{F}(s, t, u)$
- decompose scalar amplitude $\mathcal{F}(s, t, u)$ using isospin and $s \leftrightarrow u$ symmetry into **single variable amplitudes**

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- D -wave a_2 resonance via tensor meson dominance

Inhomogeneous Omnès problem

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- solve this with a separation ansatz

$$\mathcal{F}(s) = \Omega(s) \left(P_{n-1}(s) + \frac{s^n}{\pi} \int_{s_{\text{th}}}^{\infty} \frac{ds'}{s'^n} \frac{\widehat{\mathcal{F}}(s') \sin(\delta(s'))}{|\Omega(s')|(s' - s)} \right)$$

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- isovector and isoscalar part of the **photon** decouple

Basis Functions

- solution depends on **subtraction polynomials** linearly
- construct **basis functions** that correspond to one **subtraction constant**

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- reduces computational effort dramatically
- fit/matching can be done using the **basis functions**

- using $\delta^{(3/2)} = 0$ and $\delta^{(1/2)} = \delta$ we find

$$\mathcal{F}^{(0,1/2)}(s) = \Omega(s) \left(P_{n-1}^{(0,1/2)}(s) + \frac{s^n}{\pi} \int_{s_{\text{th}}}^{\infty} \frac{ds'}{s'^n} \frac{\widehat{\mathcal{F}}^{(0,1/2)}(s') \sin(\delta(s'))}{|\Omega(s')|(s' - s)} \right)$$

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- $n = 2$: twice subtracted scheme $a_2^{(1/2)}$, $b_2^{(1/2)}$; $a_2^{(0)}$, $b_2^{(0)}$; $a_2^{(3/2)}$

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- remove the $I = 3/2$ component with the **ambiguity**
- two/four **free parameters**

- solve KT-equations [Khuri and Treiman, 1960] iteratively

$$\mathcal{F} \left[\widehat{\mathcal{F}} \right] (s) = \Omega(s) \left(P_{n-1}(s) + \frac{s^n}{\pi} \int_{s_{\text{th}}}^{\infty} \frac{ds'}{s'^n} \frac{\widehat{\mathcal{F}}(s') \sin(\delta(s'))}{|\Omega(s')|(s' - s)} \right)$$

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- separate t - and u -channel contributions in the inhomogeneity

$$\widehat{\mathcal{F}} [\mathcal{F}] (s) = \widehat{\mathcal{F}}_{\text{fix}}(s) + \widehat{\mathcal{F}}_{\text{it}} [\mathcal{F}] (s)$$

- calculate chiral **Wess–Zumino–Witten** [1971,1983] anomaly

$$\mathcal{F}^{-0/00}(s=0, t=0, u=0) = F_{KK\pi\gamma} = \frac{e}{4\pi^2 F_\pi^3}$$

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$$a_n^{(1/2)} = 1.0(1.3) \text{ GeV}^{-3}$$

$$a_n^{(0)} = 0.9(1.2) \text{ GeV}^{-3}$$

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Matching: Coupling

Analytic continuation to the K^* pole

- starting from the unitarity condition we can connect the amplitudes on the **first (I)** and **second (II)** Riemann sheet

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Analytic continuation to the K^* pole

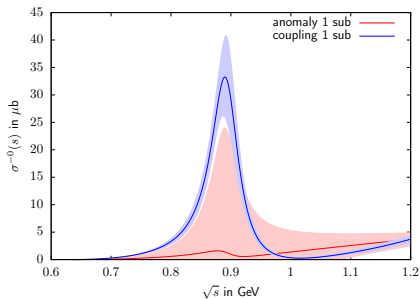
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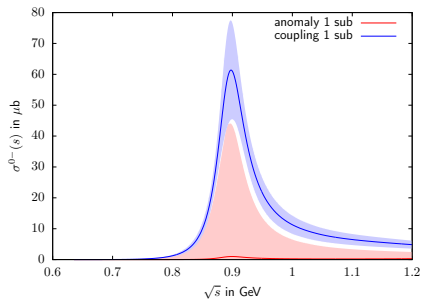
- approximate amplitudes on **second Riemann sheet** with **Breit–Wigner** parametrisations
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- $K\pi$ coupling** and **K^* pole position** from Roy–Steiner analysis [Ruiz de Elvira, 2019]
- $f(s_{K^*})$ on the **first Riemann sheet** is calculated depending on the subtraction constants via the **kernel method**
- do this for both **isospin components** separately

Matching: Anomaly or Coupling

- minimal subtraction scheme
- fully determined by the **anomaly** or **coupling**



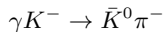
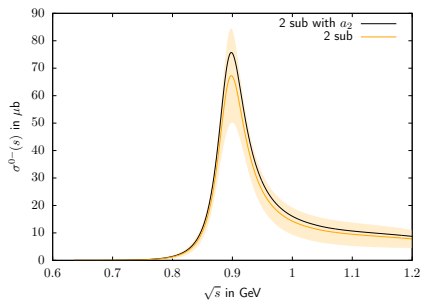
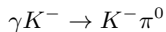
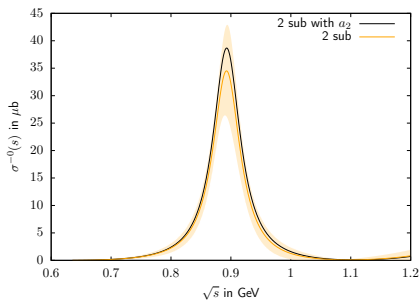
$$\gamma K^- \rightarrow K^- \pi^0$$



$$\gamma K^- \rightarrow \bar{K}^0 \pi^-$$

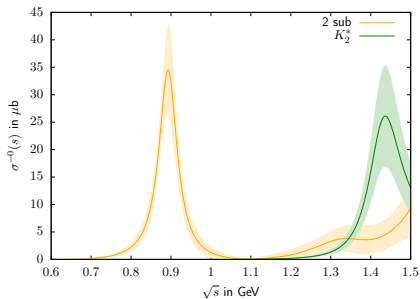
Matching: Anomaly and Coupling

- twice subtracted scheme with and without a_2 resonance

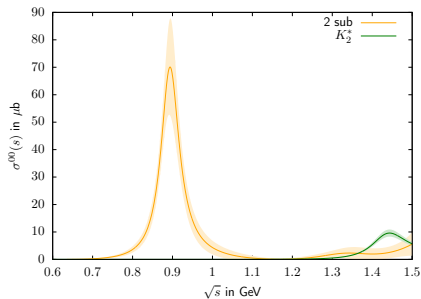


D-wave $K_2^*(1430)$

- use Lagrangians [Ecker and Zauner, 2007], [Plenter and Kubis, 2015]
- for neutral channel the radiative coupling is only an upper limit
 $\Gamma_{K_2^* \rightarrow K^0 \gamma} < 5.4 \text{ keV}$ [Alavi-Harati et al. (KTeV), 2001]



$$\gamma K^- \rightarrow K^- \pi^0$$



$$\gamma \bar{K}^0 \rightarrow \bar{K}^0 \pi^0$$

- D-wave relevant in charged channels above 1.35 GeV [Bacho, 2021]

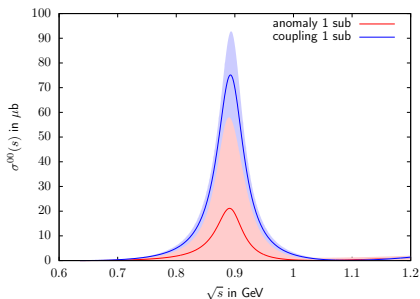
Summary and Outlook

- constructed a **dispersive solution** for the Primakoff reaction $\gamma K \rightarrow K\pi$ for all charge configurations
- **input**: fixed t -channels and $K\pi$ phase shift
- using the basis functions a fit to COMPASS++ (or OKA) data is possible to determine the free parameters
- **matching**: use **radiative couplings** and **chiral anomaly** to predict the free parameters (using the fit, extract these quantities)
- reduce error on $a_i^{(1/2)}$:
 - next-to-leading-order correction to the anomaly
 - $\omega \rightarrow K\bar{K}$ coupling: space-like kaon form factor

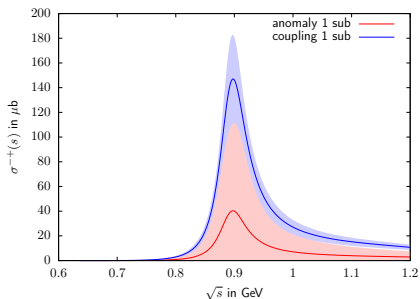
Spares

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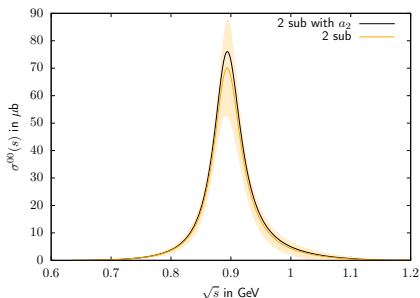
$$\gamma \bar{K}^0 \rightarrow \bar{K}^0 \pi^0$$



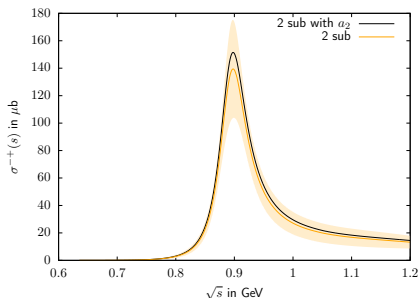
$$\gamma \bar{K}^0 \rightarrow K^- \pi^+$$

Matching: Anomaly and Coupling

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$$\gamma \bar{K}^0 \rightarrow \bar{K}^0 \pi^0$$



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