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NNLO Positivity Bounds on χ PT for a General Number of Flavours

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FACULTY OF
SCIENCE



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New Bounds on χ PT

2-flavour NLO

2-flavour NNLO

3-flavour NLO

3-flavour NLO

Higher flavour

Conclusions

The curse of χ PT (and other EFTs)

Many LECs — Limited observables	LO LECs (2):	high precision
	NLO LECs (10):	%-level precision
	NNLO LECs (90):	educated guesses
	N ³ LO LECs (1233):	unknown

(J. Bijnens and G. Ecker, *Ann. Rev. Nucl. Part. Sci.* 64 (2014) 149)



Whence bounds?

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Guaranteed
for “ordinary”
QFT

Analyticity

Unitarity

Crossing symmetry

Perturbativity

Not necessarily
consistent for EFT
 \Rightarrow bounds!





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■ Principles go back to the 60's; renewed interest in the 90's

A. Martin, Springer-Verlag, 1 ed., 1969

B. Ananthanarayan, D. Toublan and G. Wanders, *Phys. Rev. D* 51 (1995) 1093 [hep-ph/9410302]
etc., etc.

■ Our method is based on work by Manohar & Mateu

A. V. Manohar and V. Mateu, *Phys. Rev. D* 77 (2008) 094019 [0801.3222]

V. Mateu, *Phys. Rev. D* 77 (2008) 094020 [0801.3627]

■ This talk is based on ongoing work by Alvarez, Bijmens & MS (extension of Benjamin Alvarez' 2018 master thesis)

B. Alvarez, J. Bijmens and M. Sjö, [2111.XXXXX]

■ Some recent (2020) extensions in similar directions

Y.-J. Wang, F.-K. Guo, C. Zhang and S.-Y. Zhou, *JHEP* 07 (2020) 214 [2004.03992]

A. J. Tolley, Z.-Y. Wang and S.-Y. Zhou, *JHEP* 05 (2021) 255 [2111.02400]



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Decomposition and Crossing

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(Generalised) Isospin Decomposition

$$T(s, t) = \sum_{\mathcal{J}} a_{\mathcal{J}} T^{\mathcal{J}}(s, t)$$

$$\mathcal{J} : \begin{cases} \text{isospin } 0, 1, 2 & \text{in 2-flavour,} \\ \text{representation } I, A, S, AS, SS, (AA) & \text{in 3(4+)-flavour} \end{cases}$$

Example: $\pi^{\pm}\pi^{\pm} \rightarrow \pi^{\pm}\pi^{\pm}$ is purely isospin 2 (or SS)

$s \leftrightarrow u$ Crossing Symmetry

$$T^I(u, t) = C_u^{I\mathcal{J}} T^{\mathcal{J}}(s, t)$$

Matrix $C_u^{I\mathcal{J}}$ determined entirely by group structure



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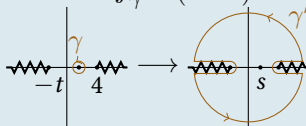
3-flavour NNLO

Higher flavour

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Dispersion relations (following Manohar & Mateu)

$$a_J \frac{d^k}{ds^k} T^J(s, t) = \frac{k!}{2\pi i} \oint_{\gamma} dz \frac{a_J T^J(z, t)}{(z-s)^{k+1}} \xrightarrow{\text{Contour manipulation}}$$



$$= \frac{k!}{\pi} \int_4^{\infty} dz \left(\frac{a_J}{(z-s)^{k+1}} + \frac{(-1)^k a_I C_u^{IJ}}{(z-u)^{k+1}} \right) \text{Im } T^J(z + \varepsilon i, t)$$

Partial-wave expansion, optical theorem

$$\text{Im } T^J(s, t) = \sum_{\ell=0}^{\infty} (2\ell+1) s \beta(s) \sigma_{\ell}^I P_{\ell} \left(1 + \frac{2t}{s-4} \right) \geq 0$$

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$$a_J \frac{d^k}{ds^k} T^J(s, t) \geq 0$$

Linear in all LECs up to NNLO

Conditions

$k \geq 2$ (convergence), k even in 2-flavour

$$a_I \left\{ \delta^{IJ} \left[\frac{z-u}{z-s} \right]^{k+1} + (-1)^k C_u^{IJ} \right\} \geq 0 \quad \text{for all } z \geq 4$$

$$t \in [0, 4], \quad s \in [-t, 4]$$

Conditions on s, t, u

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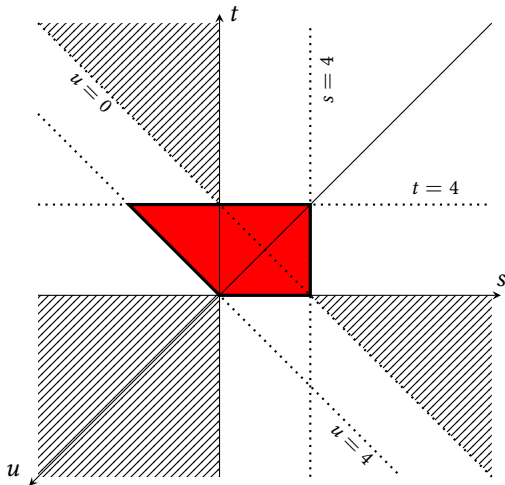
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Conditions on a_j

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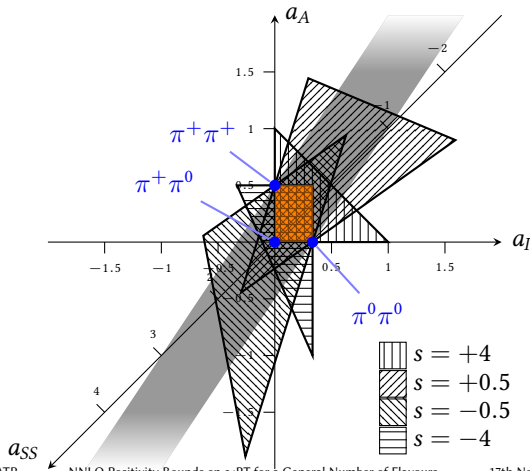
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No need to restrict to **mass eigenstates**

No need to consider bounds valid **for all s, t**

— specific values is enough

(also sufficient to check $z = 4$ and $z = \infty$)



Improving Bounds

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Conclusions

- Move up to NNLO amplitude at equal meson masses

J. Bijnens and J. Lu, *JHEP* 03 (2011) 028 [1102.0172]

- Most general choice of a_J (previous slide)
- Integration above threshold

$$a_J \frac{d^k}{ds^k} T^J = \int_4^\infty dz[\dots] \text{Im } T^J$$

$$\rightarrow a_J \left(\frac{d^k}{ds^k} T^J - \int_4^\lambda dz[\dots] \text{Im } T^J \right) = \int_\lambda^\infty dz[\dots] \text{Im } T^J$$

- Strengthens bounds
- Reliant on low-energy approximation
- Breakdown scale is $\lambda \sim 70/n$ at n flavours
R. S. Chivukula, M. J. Dugan and M. Golden, *Phys. Rev. D* 47 (1993) 2930 [hep-ph/9206222]
- We have performed analytic integral up to NNLO
- Mathematical framework for reducing sets of bounds





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Definition of a Linear Constraint

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Linear Constraints

Parameters b_i , constants α_i, c

$$\alpha_1 b_1 + \alpha_2 b_2 + \dots + \alpha_N b_N - c \geq 0 \quad \Leftrightarrow \quad \boldsymbol{\alpha} \cdot \mathbf{b} \geq c$$

Expressed as

“The (linear) constraint $\langle \boldsymbol{\alpha}, c \rangle$ is satisfied by \mathbf{b} ”



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Combination and Comparison

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Combined constraints

$\langle \alpha, c \rangle + \langle \beta, d \rangle$ satisfied if *both* $\langle \alpha, c \rangle$ and $\langle \beta, d \rangle$ satisfied.

Generally:

$$\Omega = \sum_i \langle \alpha_i, c_i \rangle$$

Stronger and weaker constraints

If Ω' satisfied by all points satisfying Ω :

$$\Omega' \leq \Omega$$

Basic examples

$$\langle \alpha, -1 \rangle \leq \langle \alpha, 0 \rangle \leq \langle \alpha, 1 \rangle, \quad \Omega \leq (\Omega + \Omega') \geq \Omega'$$



Deep Results

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Result 1

$\langle \alpha, c \rangle$ weaker than $\Omega = \sum_i \langle \alpha_i, c_i \rangle$ if and only if

$$\alpha = \sum_i \lambda_i \alpha_i, \quad \sum_i \lambda_i c_i \geq c, \quad \lambda_i \geq 0$$

Result 2

Equivalently, if and only if α satisfies $\sum_j \langle \mathbf{n}_j, r_j \rangle$

Straightforward algorithm to generate $\langle \mathbf{n}_j, r_j \rangle$ based on convex hulls.

Result 3

Of all possible sets \mathcal{S} such that $\Omega = \sum_{\langle \alpha, c \rangle \in \mathcal{S}} \langle \alpha, c \rangle$, there is a (nearly) unique *smallest* such set.

Obtained as side-effect of above algorithm.





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2-flavour NLO

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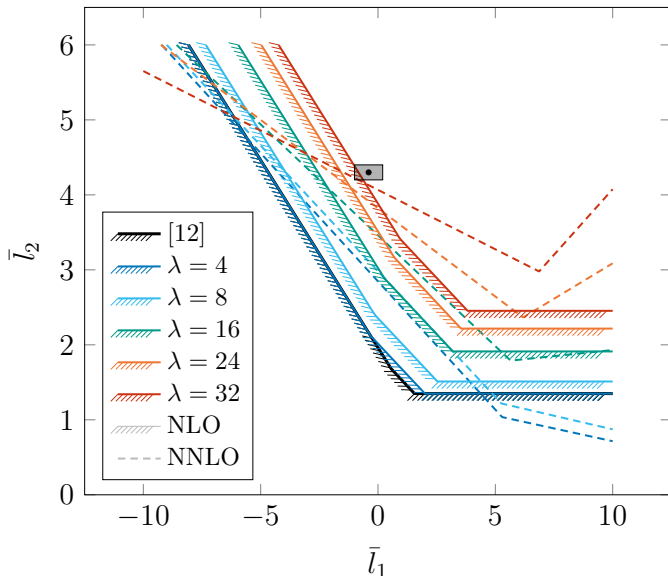
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2-flavour NNLO

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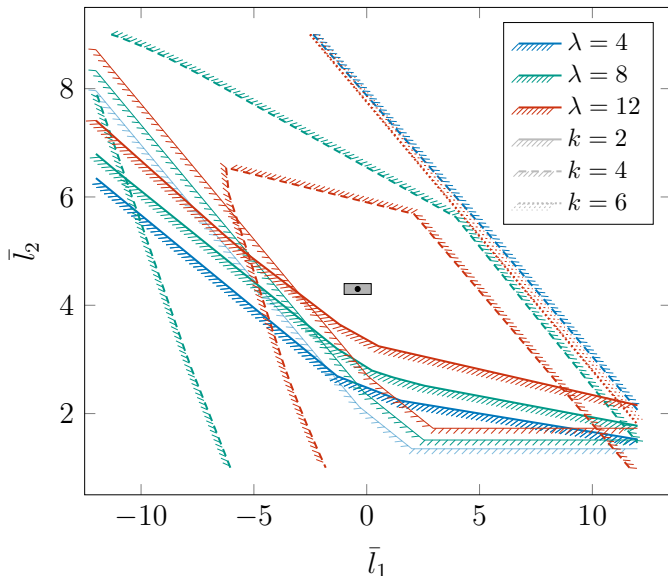
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2-flavour NNLO

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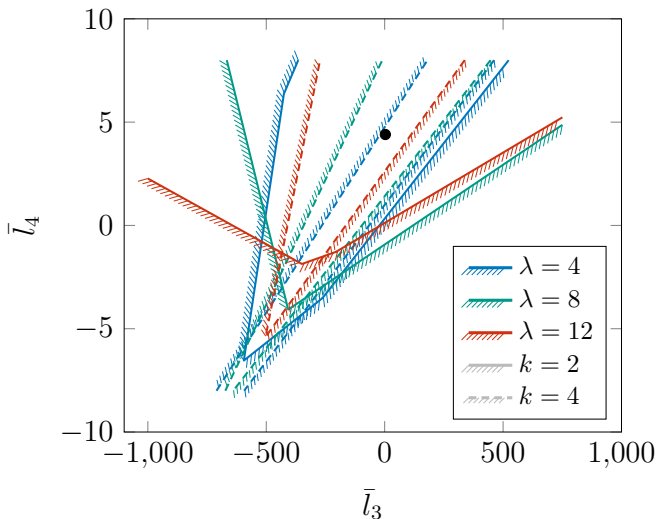
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3-flavour NLO

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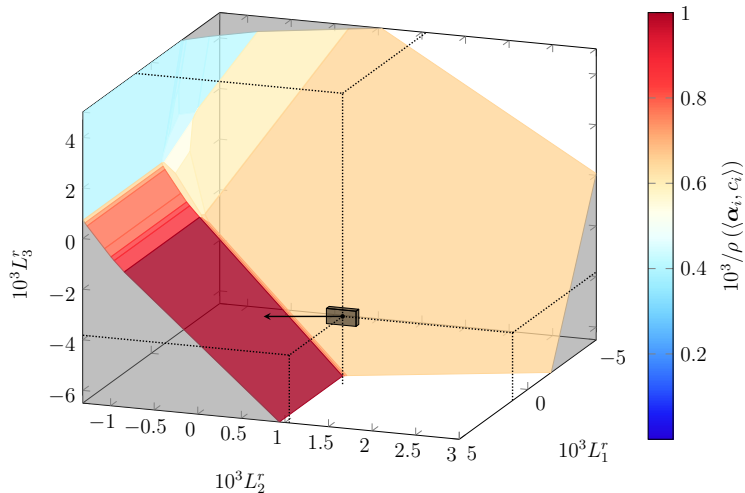
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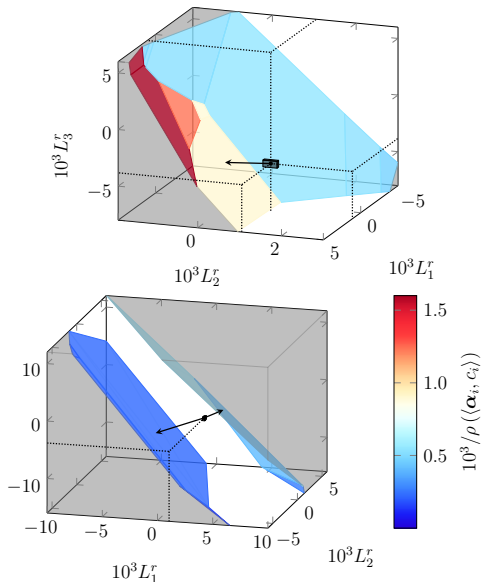
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3-flavour NNLO

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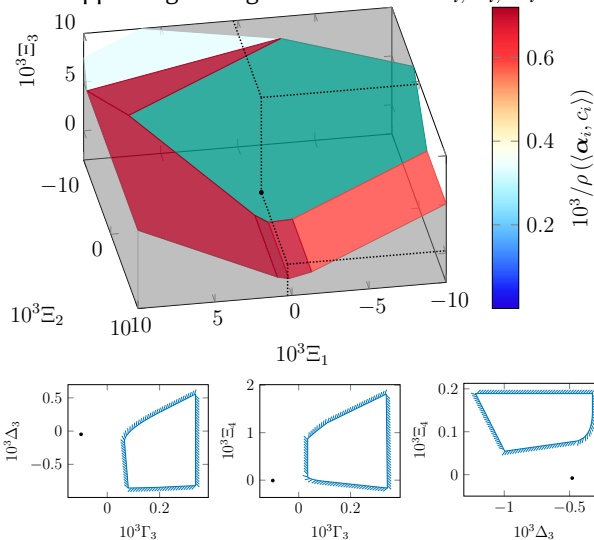
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NNLO LECs appearing through combinations $\Xi_i, \Gamma_i, \Delta_i$



Higher Flavour

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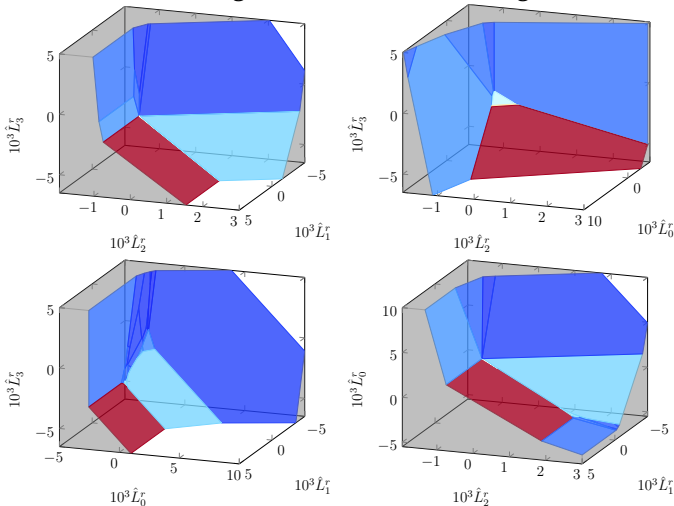
3-flavour NLO

3-flavour NLO

Higher flavour

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4-flavour NLO, no great difference at even higher flavour



Summary & Outlook

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Summary

- First full NNLO bounds
- Stronger and more general bounds produced
- Powerful mathematical framework

Outlook

- Multiple improvements possible (general masses, etc.)
- Mostly plug-and-play for new amplitudes
 - N^3 LO, higher-point, etc.
 - Other EFTs

Thank you!

