

Patterns of C - and CP -violation in hadronic η and η' three-body decays

Hakan Akdag

16th November 2021

HISKP (Theory), University of Bonn

presentation based on: Akdag, Isken, Kubis [arxiv:2111.02417 [hep-ph]] (2021)

Asymmetry between matter and antimatter

- Origin of matter from baryogenesis presumes C - and CP -violation [*Sakharov 1967*]
- In SM: violation from weak interaction is not sufficient to create observed asymmetry

Asymmetry between matter and antimatter

- Origin of matter from baryogenesis presumes C - and CP -violation [*Sakharov 1967*]
- In SM: violation from weak interaction is not sufficient to create observed asymmetry

Search for new sources of CP -violation:

- Mostly neglected since 1960's: T -odd P -even (TOPE) operators in strong interactions
- Consider an eigenstate of C , we focus on the η meson
 \hookrightarrow Can investigate CP -violation in absence of weak interaction

Asymmetry between matter and antimatter

- Origin of matter from baryogenesis presumes C - and CP -violation [*Sakharov 1967*]
- In SM: violation from weak interaction is not sufficient to create observed asymmetry

Search for new sources of CP -violation:

- Mostly neglected since 1960's: T -odd P -even (TOPE) operators in strong interactions
- Consider an eigenstate of C , we focus on the η meson
 \hookrightarrow Can investigate CP -violation in absence of weak interaction

Ideal stage to investigate TOPE forces:

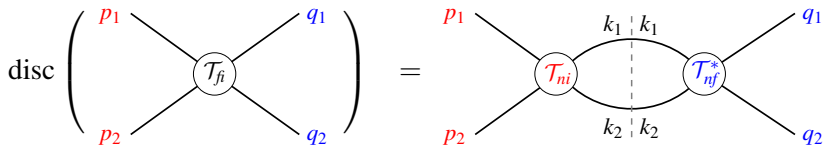
- $\eta \rightarrow \pi^0 \pi^+ \pi^-$
 - $\eta' \rightarrow \eta \pi^+ \pi^-$
- } Dalitz plots sensitive to mirror symmetry breaking

Evaluate three-particle decay in dispersive (*Khuri-Treiman*) framework:

Model independent and non-perturbative re-summation of final state interactions, based on

1 Unitarity (\sim probability conservation) gives rise to *optical theorem*:

$$\mathcal{T}_{fi} - \mathcal{T}_{if}^* = i \sum_n \int d\Pi_n (2\pi)^4 \delta^4\left(\sum_{i,n} p_i - k_n\right) \mathcal{T}_{ni} \mathcal{T}_{nf}^*$$



2 Analyticity (\sim causality)

Dispersion relations: reconstruct whole amplitude with knowledge about discontinuity

Idea: derive $2 \rightarrow 2$ scattering amplitude and analytically continue to realm of $1 \rightarrow 3$ decay

$$\eta \rightarrow \pi^0 \pi^+ \pi^-$$

- $\eta \rightarrow \pi^+ \pi^- \pi^0$ breaks G-parity; in the **Standard Model** consider transition with $\Delta I = 1$

$$\mathcal{M}(s, t, u) = \xi \mathcal{M}_1^C(s, t, u)$$

[Colangelo et al. 2018; Albaladejo et al. 2017; Guo et al. 2017; ...]

- $\eta \rightarrow \pi^+ \pi^- \pi^0$ breaks G-parity; in the **Standard Model** consider transition with $\Delta I = 1$
- For **C-violating** parts consider $C = -(-1)^{\Delta I}$, i.e. need even *total* isospin [*Gardner, Shi 2020*]

$$\mathcal{M}(s, t, u) = \xi \mathcal{M}_1^C(s, t, u) + \mathcal{M}_0^{\mathcal{C}}(s, t, u) + \mathcal{M}_2^{\mathcal{C}}(s, t, u)$$

- $\eta \rightarrow \pi^+ \pi^- \pi^0$ breaks G-parity; in the **Standard Model** consider transition with $\Delta I = 1$
- For **C-violating** parts consider $C = -(-1)^{\Delta I}$, i.e. need even *total* isospin [*Gardner, Shi 2020*]

$$\mathcal{M}(s, t, u) = \xi \mathcal{M}_1^C(s, t, u) + \mathcal{M}_0^{\mathcal{C}}(s, t, u) + \mathcal{M}_2^{\mathcal{C}}(s, t, u)$$

- Bose symmetry: odd (even) $\pi\pi$ -isospin must have odd (even) partial wave
- Reconstruction theorem: expand for fixed *two-body* isospin and partial wave

$$\mathcal{M}_1^C(s, t, u) = \mathcal{F}_0(s) + (s - u)\mathcal{F}_1(t) + (s - t)\mathcal{F}_1(u) + \mathcal{F}_2(t) + \mathcal{F}_2(u) - \frac{2}{3}\mathcal{F}_2(s)$$

$$\mathcal{M}_0^{\mathcal{C}}(s, t, u) = (t - u)\mathcal{G}_1(s) + (u - s)\mathcal{G}_1(t) + (s - t)\mathcal{G}_1(u)$$

$$\mathcal{M}_2^{\mathcal{C}}(s, t, u) = 2(u - t)\mathcal{H}_1(s) + (u - s)\mathcal{H}_1(t) + (s - t)\mathcal{H}_1(u) - \mathcal{H}_2(t) + \mathcal{H}_2(u)$$

- **C-even** terms are **symmetric** and **C-odd** ones **antisymmetric** under $t \leftrightarrow u$
- Note: \mathcal{F}_I , \mathcal{G}_I and \mathcal{H}_I are completely independent

- Single variable amplitudes $\mathcal{A} \in \{\mathcal{F}, \mathcal{G}, \mathcal{H}\}$ obey discontinuity relation

$$\text{disc} \mathcal{A}_I(s) = 2i \theta(s - 4M_\pi^2) [\mathcal{A}_I(s) + \hat{\mathcal{A}}_I(s)] \sin \delta_I(s) e^{-i\delta_I(s)}$$

- Single variable amplitudes $\mathcal{A} \in \{\mathcal{F}, \mathcal{G}, \mathcal{H}\}$ obey discontinuity relation

$$\text{disc} \mathcal{A}_I(s) = 2i \theta(s - 4M_\pi^2) [\mathcal{A}_I(s) + \hat{\mathcal{A}}_I(s)] \sin \delta_I(s) e^{-i\delta_I(s)}$$

- Homogeneous solution

$$\mathcal{A}_I(s) = P(s) \Omega_I(s), \quad \Omega_I(s) = \exp \left(\frac{s}{\pi} \int_{4M_\pi^2}^{\infty} \frac{\delta_I(x)}{x(x-s)} dx \right)$$

- Single variable amplitudes $\mathcal{A} \in \{\mathcal{F}, \mathcal{G}, \mathcal{H}\}$ obey discontinuity relation

$$\text{disc} \mathcal{A}_I(s) = 2i \theta(s - 4M_\pi^2) [\mathcal{A}_I(s) + \hat{\mathcal{A}}_I(s)] \sin \delta_I(s) e^{-i\delta_I(s)}$$

- Homogeneous solution

$$\mathcal{A}_I(s) = P(s) \Omega_I(s), \quad \Omega_I(s) = \exp \left(\frac{s}{\pi} \int_{4M_\pi^2}^{\infty} \frac{\delta_I(x)}{x(x-s)} dx \right)$$

- Inhomogeneities defined via partial wave (include left-hand cut contribution)

$$a_I(s) = \mathcal{A}_I(s) + \hat{\mathcal{A}}_I(s)$$

- Full solution

$$\mathcal{A}_I(s) = \Omega_I(s) \left(P_{n-1}(s) + \frac{s^n}{\pi} \int_{4M_\pi^2}^{\infty} \frac{dx}{x^n} \frac{\sin \delta_I(x) \hat{\mathcal{A}}_I(s)}{|\Omega_I(x)| (x-s)} \right)$$

- Subtraction polynomial P_{n-1} fixed by asymptotics imposed on \mathcal{A}_I and δ_I and by data

Regression to Dalitz plot [\[KLOE-2, 2016\]](#)

The SM amplitude \mathcal{M}_1^C :

- Minimal subtraction scheme 3 dof: $\chi_{\text{red}}^2 \approx 1.054$
 - Observables agree with current literature
 - 1 Taylor invariants [\[Colangelo et al. 2018\]](#)
 - 2 $\text{BR}(\eta \rightarrow 3\pi^0)/\text{BR}(\eta \rightarrow \pi^0\pi^+\pi^-)$ [\[PDG 2020\]](#)
 - 3 Dalitz plot parameters [\[Colangelo et al. 2018, PDG 2020\]](#)
- \implies subtraction scheme justified, apply also to $\mathcal{M}_{0,2}^C$

Regression to Dalitz plot [KLOE-2, 2016]

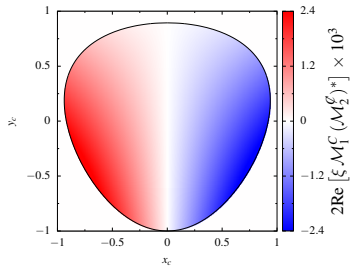
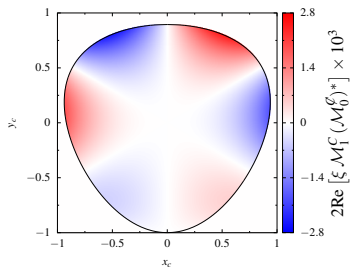
The SM amplitude \mathcal{M}_1^C :

- Minimal subtraction scheme 3 dof: $\chi_{\text{red}}^2 \approx 1.054$
- Observables agree with current literature
 - 1 Taylor invariants [Colangelo et al. 2018]
 - 2 $\text{BR}(\eta \rightarrow 3\pi^0)/\text{BR}(\eta \rightarrow \pi^0\pi^+\pi^-)$ [PDG 2020]
 - 3 Dalitz plot parameters [Colangelo et al. 2018, PDG 2020]

\Rightarrow subtraction scheme justified, apply also to $\mathcal{M}_{0,2}^C$

The BSM amplitude $\mathcal{M} = \xi \mathcal{M}_1^C + \mathcal{M}_0^C + \mathcal{M}_2^C$:

- Fix $\mathcal{M}_{0,2}^C$ by just one complex normalization each
- Full amplitude 7 dof: $\chi_{\text{red}}^2 \approx 1.048$
- Upper limit on TOPE effects in per mille level
- All C - and CP -violating signals vanish within 1-2 σ



- Effective BSM operators

$$X_0^\ell \sim g_0(s-t)(t-u)(u-s) + \mathcal{O}(p^8)$$

$$X_2^\ell \sim g_2(t-u) + \mathcal{O}(p^4)$$

- Obtain couplings by a Taylor expansion of $\mathcal{M}_0^\ell, \mathcal{M}_2^\ell$:

$$g_0 = (-3.3(3.6) + 6.7(12.6) i) \text{ GeV}^{-6}$$

$$g_2 = (0.001(15) - 0.006(42) i) \text{ GeV}^{-2}$$

- Relative deviation $|g_0/g_2| \approx 10^3 \text{ GeV}^{-4} \implies$ fit compensates kinematic suppression of X_0^ℓ

■ Effective BSM operators

$$X_0^\ell \sim g_0(s-t)(t-u)(u-s) + \mathcal{O}(p^8)$$

$$X_2^\ell \sim g_2(t-u) + \mathcal{O}(p^4)$$

■ Obtain couplings by a Taylor expansion of $\mathcal{M}_0^\ell, \mathcal{M}_2^\ell$:

$$g_0 = (-3.3(3.6) + 6.7(12.6)i) \text{ GeV}^{-6}$$

$$g_2 = (0.001(15) - 0.006(42)i) \text{ GeV}^{-2}$$

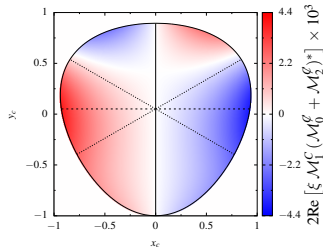
■ Relative deviation $|g_0/g_2| \approx 10^3 \text{ GeV}^{-4} \Rightarrow$ fit compensates kinematic suppression of X_0^ℓ

■ Dalitz-plot asymmetries in 10^{-4} (g_0 in GeV^{-6} , g_2 in 10^3 GeV^{-2}):

$$A_{LR} = -0.943 \text{ Re } g_0 - 0.300 \text{ Im } g_0 - 2.493 \text{ Re } g_2 - 0.936 \text{ Im } g_2 = -7.6(4.7)$$

$$A_Q = 0.479 \text{ Re } g_0 + 0.443 \text{ Im } g_0 + 0.536 \text{ Re } g_2 + 0.336 \text{ Im } g_2 = 4.1(4.3)$$

$$A_S = -2.971 \text{ Re } g_0 - 0.850 \text{ Im } g_0 - 0.057 \text{ Re } g_2 - 0.043 \text{ Im } g_2 = 3.7(4.3)$$



Can not repeating the high-precision analysis for $\eta' \rightarrow \pi^0 \pi^+ \pi^-$

- $\eta' \rightarrow \pi^0 \pi^+ \pi^-$ rare decay mode
- Precise investigation of Dalitz plot [BESIII 2017] not possible yet... [Isken et al. 2021 (in preperation)]

What happens for an increased phase space ($M_\eta \rightarrow M_{\eta'}$)?

- Both decays are driven by same $\Delta I = 0, 2$ operators
- Assume *same* couplings g_0 and g_2 to predict C -odd amplitudes
- \mathcal{M}_0^ℓ dominates \mathcal{M}_2^ℓ by 10^2
 \implies TOPE forces in $\eta' \rightarrow \pi^0 \pi^+ \pi^-$ are more sensitive to **isoscalar** transitions

$$\eta' \rightarrow \eta \pi^+ \pi^-$$

- $\eta' \rightarrow \eta\pi^+\pi^-$ conserves G-parity; in the **Standard Model** consider transition with $\Delta I = 0$

$$\mathcal{M}(s, t, u) = \mathcal{M}_0^G(s, t, u)$$

[Isken et al. 2017, Isken et al. 2021 (in preparation)]

- $\eta' \rightarrow \eta \pi^+ \pi^-$ conserves G-parity; in the **Standard Model** consider transition with $\Delta I = 0$
- **C-violating** driven by isospin $\Delta I = 1$

$$\mathcal{M}(s, t, u) = \mathcal{M}_0^C(s, t, u) + \mathcal{M}_1^{\mathcal{C}}(s, t, u)$$

- $\eta' \rightarrow \eta \pi^+ \pi^-$ conserves G-parity; in the Standard Model consider transition with $\Delta I = 0$
- C-violating driven by isospin $\Delta I = 1$

$$\mathcal{M}(s, t, u) = \mathcal{M}_0^C(s, t, u) + \mathcal{M}_1^{\mathcal{C}}(s, t, u)$$

- Sensitive to a different class of BSM operators!

- $\eta' \rightarrow \eta \pi^+ \pi^-$ conserves G-parity; in the **Standard Model** consider transition with $\Delta I = 0$
- **C-violating** driven by isospin $\Delta I = 1$

$$\mathcal{M}(s, t, u) = \mathcal{M}_0^C(s, t, u) + \mathcal{M}_1^{\mathcal{C}}(s, t, u)$$

- Sensitive to a different class of BSM operators!
- Reconstruction theorem: two different intermediate states

$$\mathcal{M}_0^C(s, t, u) = \mathcal{F}_{\pi\pi}(s) + \mathcal{F}_{\eta\pi}(t) + \mathcal{F}_{\eta\pi}(u)$$

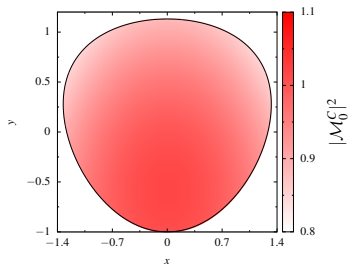
$$\mathcal{M}_1^{\mathcal{C}}(s, t, u) = (t - u) \mathcal{G}_{\pi\pi}(s) + \mathcal{G}_{\eta\pi}(t) - \mathcal{G}_{\eta\pi}(u)$$

- $\mathcal{F}_{\pi\pi}$, $\mathcal{F}_{\eta\pi}$ and $\mathcal{G}_{\eta\pi}$ in S -waves, $\mathcal{G}_{\pi\pi}$ in P -wave
- Solution analogous to $\eta \rightarrow 3\pi$

Regression to Dalitz plot [\[BESIII, 2018\]](#)

The SM amplitude \mathcal{M}_0^C :

- Minimal subtraction scheme fails to describe data accurately
- Need at least 4 dof: $\chi_{\text{red}}^2 \approx 0.994$



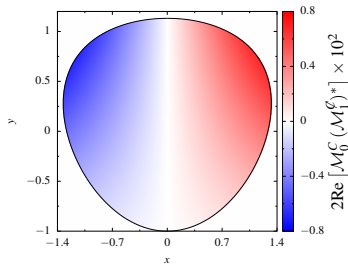
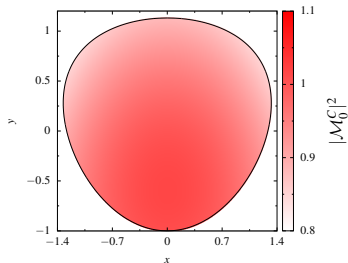
Regression to Dalitz plot [BESIII, 2018]

The SM amplitude \mathcal{M}_0^C :

- Minimal subtraction scheme fails to describe data accurately
- Need at least 4 dof: $\chi_{\text{red}}^2 \approx 0.994$

Apply same subtraction scheme for BSM amplitude:

- Fix \mathcal{M}_1^C by two complex coefficients with the same phase
- Full amplitude 7 dof: $\chi_{\text{red}}^2 \approx 0.994$
- Upper limit on TOPE effects in per cent level
- All C - and CP -violating signals vanish within $< 1.5\sigma$



- Effective BSM operator

$$X_1^{\mathcal{Q}} \sim g_1 (t - u) (1 + s \delta g_1) + \mathcal{O}(p^6)$$

- Obtain coupling and s -dependent correction by a Taylor expansion of $\mathcal{M}_1^{\mathcal{Q}}$:

$$g_1 = (0.17(27) - 0.3(5.7) i) \text{ GeV}^{-2}$$

$$\delta g_1 \approx -4(99) \text{ GeV}^{-2} \quad \text{imaginary part negligibly small}$$

- Dalitz-plot asymmetry in 10^{-3} (g_1 and δg_1 in GeV^{-2}):

$$A_{LR} = -31.0 \text{ Re } g_1 (1 + 0.09 \delta g_1) - 6.6 \text{ Im } g_1 (1 + 0.10 \delta g_1) = -2.3(1.7)$$

The dispersive framework for TOPE forces in $\eta \rightarrow \pi^0 \pi^+ \pi^-$ and $\eta' \rightarrow \eta \pi^+ \pi^-$

- Based on fundamental principles of analyticity, unitarity and crossing
- Derived C - and CP -odd contributions driven by $\Delta I = 0, 2$ and $\Delta I = 1$ transitions
- Extracted effective BSM couplings g_0, g_2 and g_1
- Current experimental precision:
 - Upper limit on C -odd signals in the relative per mille and cent level
 - In both cases BSM signals vanish within $1\text{-}2\sigma$

The dispersive framework for TOPE forces in $\eta \rightarrow \pi^0 \pi^+ \pi^-$ and $\eta' \rightarrow \eta \pi^+ \pi^-$

- Based on fundamental principles of analyticity, unitarity and crossing
- Derived C - and CP -odd contributions driven by $\Delta I = 0, 2$ and $\Delta I = 1$ transitions
- Extracted effective BSM couplings g_0 , g_2 and g_1
- Current experimental precision:
 - Upper limit on C -odd signals in the relative per mille and cent level
 - In both cases BSM signals vanish within $1\text{-}2\sigma$

Future theoretical interest:

- Amplitudes can be used to calculate/predict different TOPE processes
- Couplings g_0 , g_2 and g_1 may be used to match future effective field theories

The dispersive framework for TOPE forces in $\eta \rightarrow \pi^0 \pi^+ \pi^-$ and $\eta' \rightarrow \eta \pi^+ \pi^-$

- Based on fundamental principles of analyticity, unitarity and crossing
- Derived C - and CP -odd contributions driven by $\Delta I = 0, 2$ and $\Delta I = 1$ transitions
- Extracted effective BSM couplings g_0 , g_2 and g_1
- Current experimental precision:
 - Upper limit on C -odd signals in the relative per mille and cent level
 - In both cases BSM signals vanish within $1-2\sigma$

Future theoretical interest:

- Amplitudes can be used to calculate/predict different TOPE processes
- Couplings g_0 , g_2 and g_1 may be used to match future effective field theories

From experimental point of view:

- *JLab Eta Factory* (JEF)
- *Rare Eta Decays with a TPC for Optical Photons* (REDTOP)

Thank you very much for your attention!