



**University of New Hampshire**  
**Nuclear & Particle Physics Group**

# The $g_2p$ Experiment: A Measurement of the Proton's Spin Structure Functions

**David Ruth**

**10<sup>th</sup> International Workshop on Chiral Dynamics**

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# Acknowledgements

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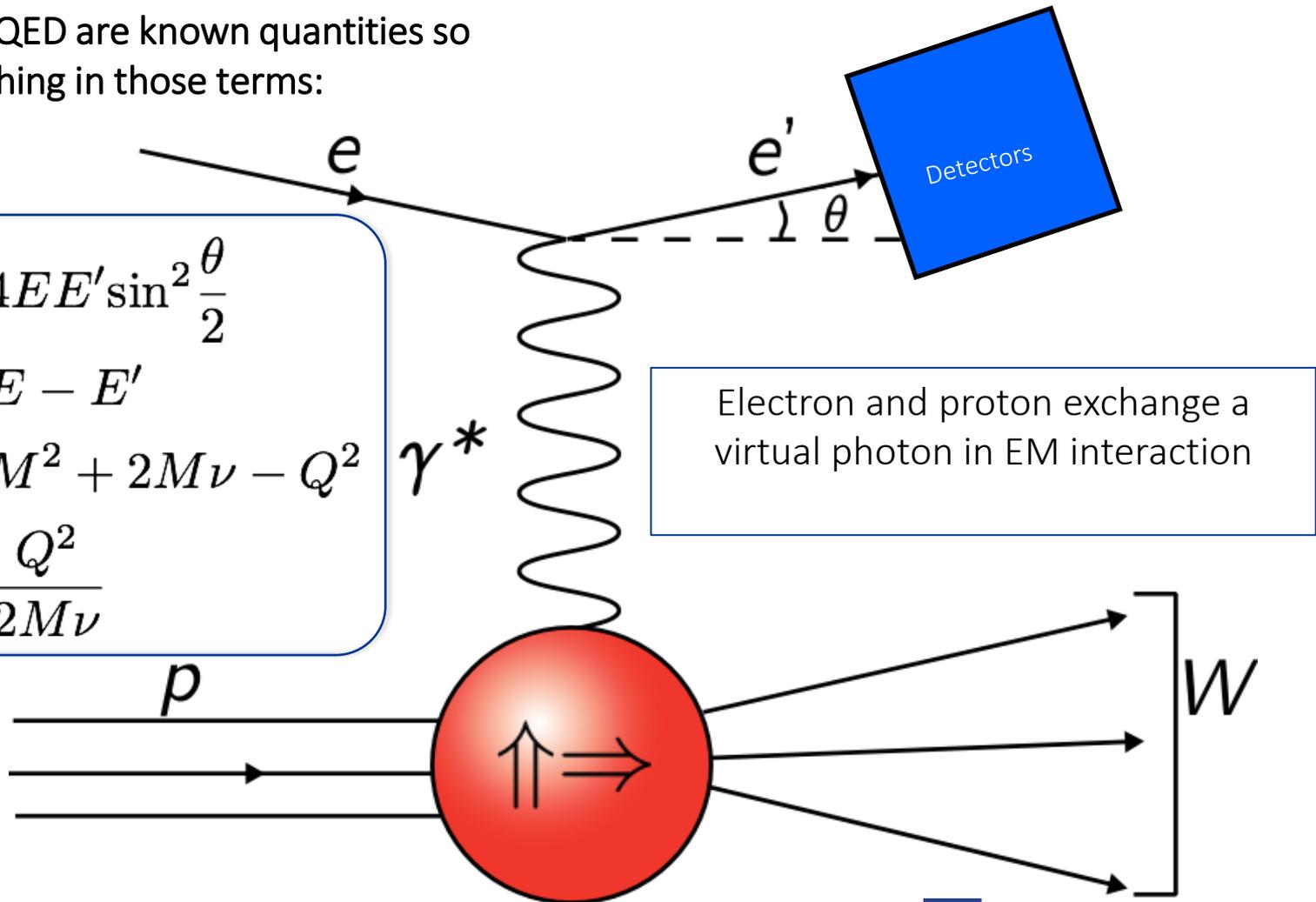
Special thanks again to Ryan Zielinski for allowing me to adapt some of his slides and figures for this presentation.



# Essential Quantities in $ep$ Scattering

Electron and QED are known quantities so define everything in those terms:

$$Q^2 = 4EE' \sin^2 \frac{\theta}{2}$$
$$\nu = E - E'$$
$$W^2 = M^2 + 2M\nu - Q^2$$
$$x = \frac{Q^2}{2M\nu}$$



# Inclusive $ep$ Scattering Cross Sections describe normalized interaction rate

Elastic scattering: target remains in the ground state after interaction

$$E'_{\text{elas}} = \frac{E}{1 + \frac{2E}{M} \sin^2 \frac{\theta}{2}}$$

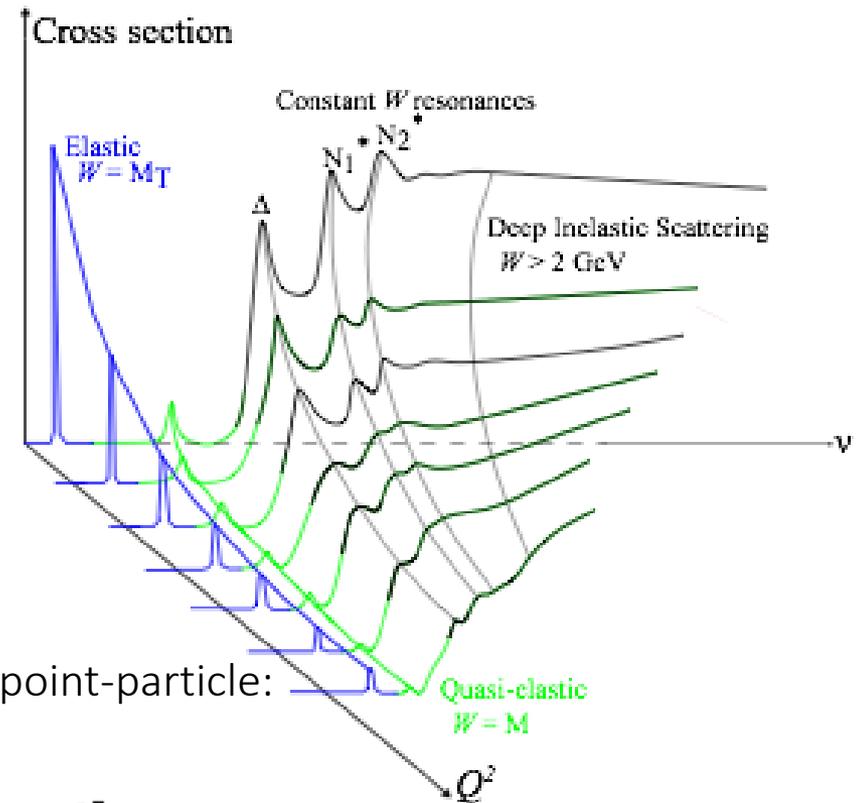
Mott cross section describes scattering from point-particle:

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} = \frac{\alpha^2}{4 E^2 \sin^4 \frac{\theta}{2}} \cos^2 \frac{\theta}{2}$$

Rosenbluth cross section describes deviation from point-particle:

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} \left[ \frac{G_E^2(Q^2) + \tau G_M^2(Q^2)}{1 + \tau} + 2\tau G_M^2(Q^2) \tan^2 \frac{\theta}{2} \right]$$

$G_E$  and  $G_M$  related to charge and current distributions



# Inclusive $ep$ Scattering Cross Sections describe normalized interaction rate

Inelastic scattering: Target is in excited state after interaction

Structure Functions:

Inclusive *unpolarized* cross sections

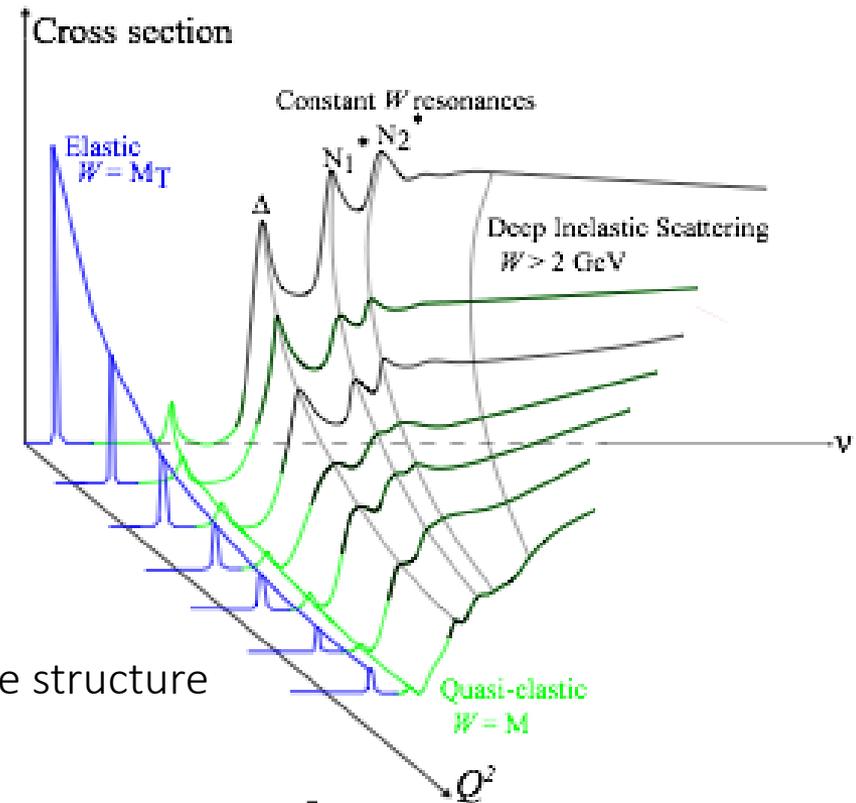
$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{\text{Mott}} \left[ \frac{1}{\nu} F_2(x, Q^2) + \frac{2}{M} F_1(x, Q^2) \tan^2 \frac{\theta}{2} \right]$$

$F_1$  and  $F_2$  related to quark/gluon distribution

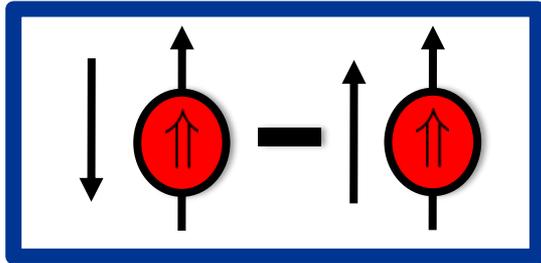
Adding a *polarized* beam and target adds two more structure functions

$$\frac{d^2\sigma^\pm}{d\Omega dE'} = \sigma_{\text{Mott}} \left[ \alpha F_1(x, Q^2) + \beta F_2(x, Q^2) \pm \gamma g_1(x, Q^2) \pm \delta g_2(x, Q^2) \right]$$

$g_1$  and  $g_2$  related to spin distribution



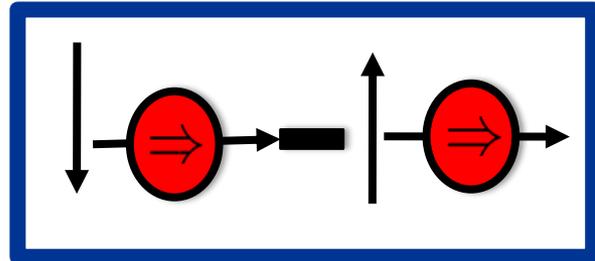
# Extracting Spin Structure by Looking at Cross Section Differences



Parallel

Inclusive *polarized* cross sections

$$\frac{d^2\sigma^{\uparrow\uparrow}}{dE'd\Omega} - \frac{d^2\sigma^{\downarrow\uparrow}}{dE'd\Omega} = \frac{4\alpha^2}{M\nu Q^2} \frac{E'}{E} \left[ g_1(x, Q^2) \{E + E' \cos\theta\} - \frac{Q^2}{\nu} g_2(\nu, Q^2) \right]$$



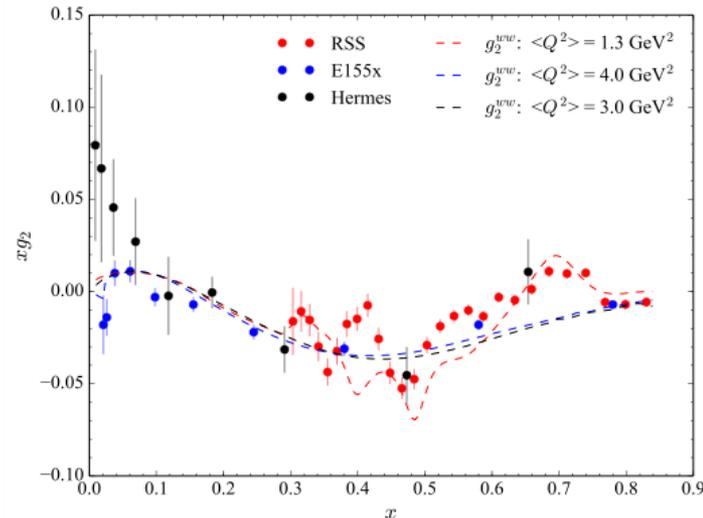
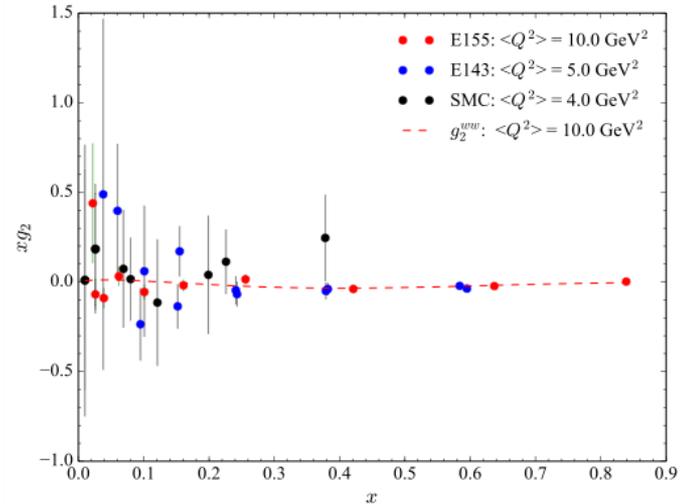
Perpendicular

$$\frac{d^2\sigma^{\uparrow\Rightarrow}}{dE'd\Omega} - \frac{d^2\sigma^{\downarrow\Rightarrow}}{dE'd\Omega} = \frac{4\alpha^2}{M\nu Q^2} \frac{E'^2}{E} \sin\theta \left[ \nu g_1(x, Q^2) + 2E g_2(\nu, Q^2) \right]$$

*Two equations, two unknowns...*

# Motivation: Measure a fundamental spin observable ( $g_2$ ) in the region $0.02 < Q^2 < 0.20 \text{ GeV}^2$ for the first time

- Measurements at Jefferson Lab:
  - RSS – medium  $Q^2$  ( 1-2  $\text{GeV}^2$ ) (published)
  - SANE – high  $Q^2$  (2-6  $\text{GeV}^2$ ) (analysis)
  - $g_2\text{p}$  – low  $Q^2$  (0.02-0.20  $\text{GeV}^2$ ) (analysis)
- Low  $Q^2$  is difficult:
  - Electrons strongly influenced by target field
  - Strong kinematic dependence on observables
- Low  $Q^2$  is useful:
  - Test predictions of Chiral Perturbation Theory ( $\chi\text{PT}$ )
  - Test sum rules and measure moments of  $g_2$
  - Study finite size effects of the proton
- $g_2\text{p}$  experiment ran spring 2012 in Hall A

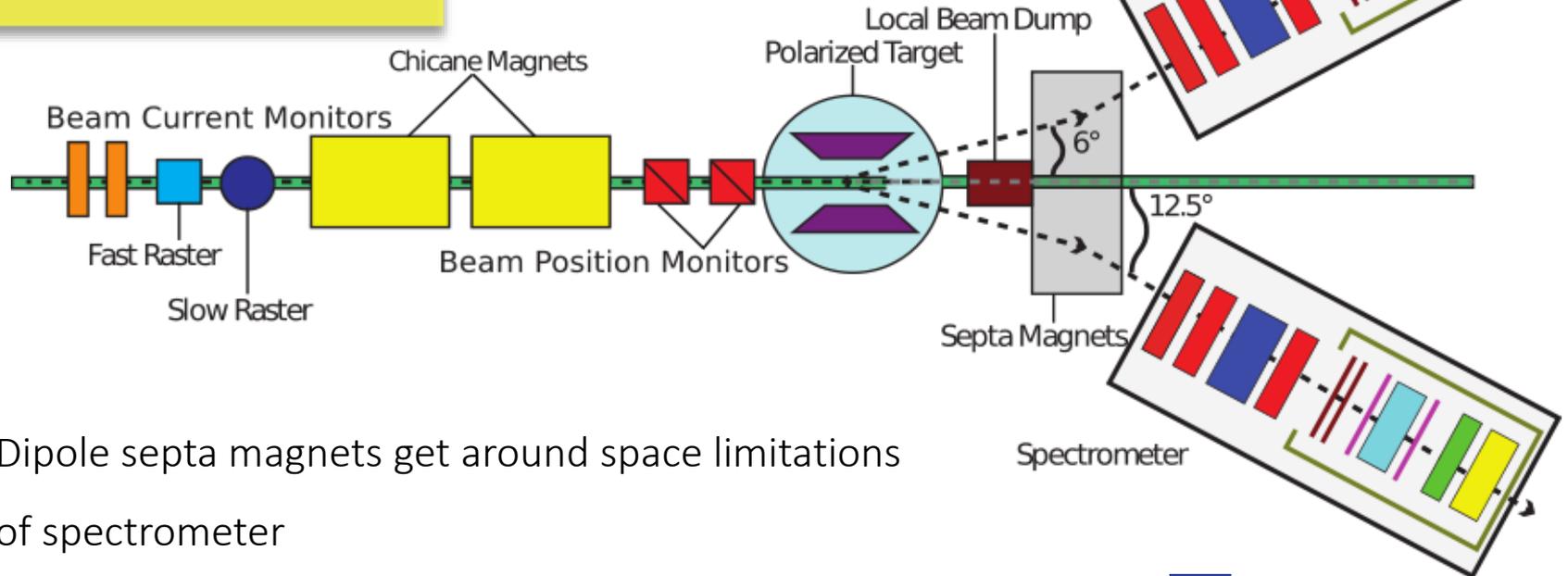


# Hall A Experimental Setup:

Measuring  $g_2^p$

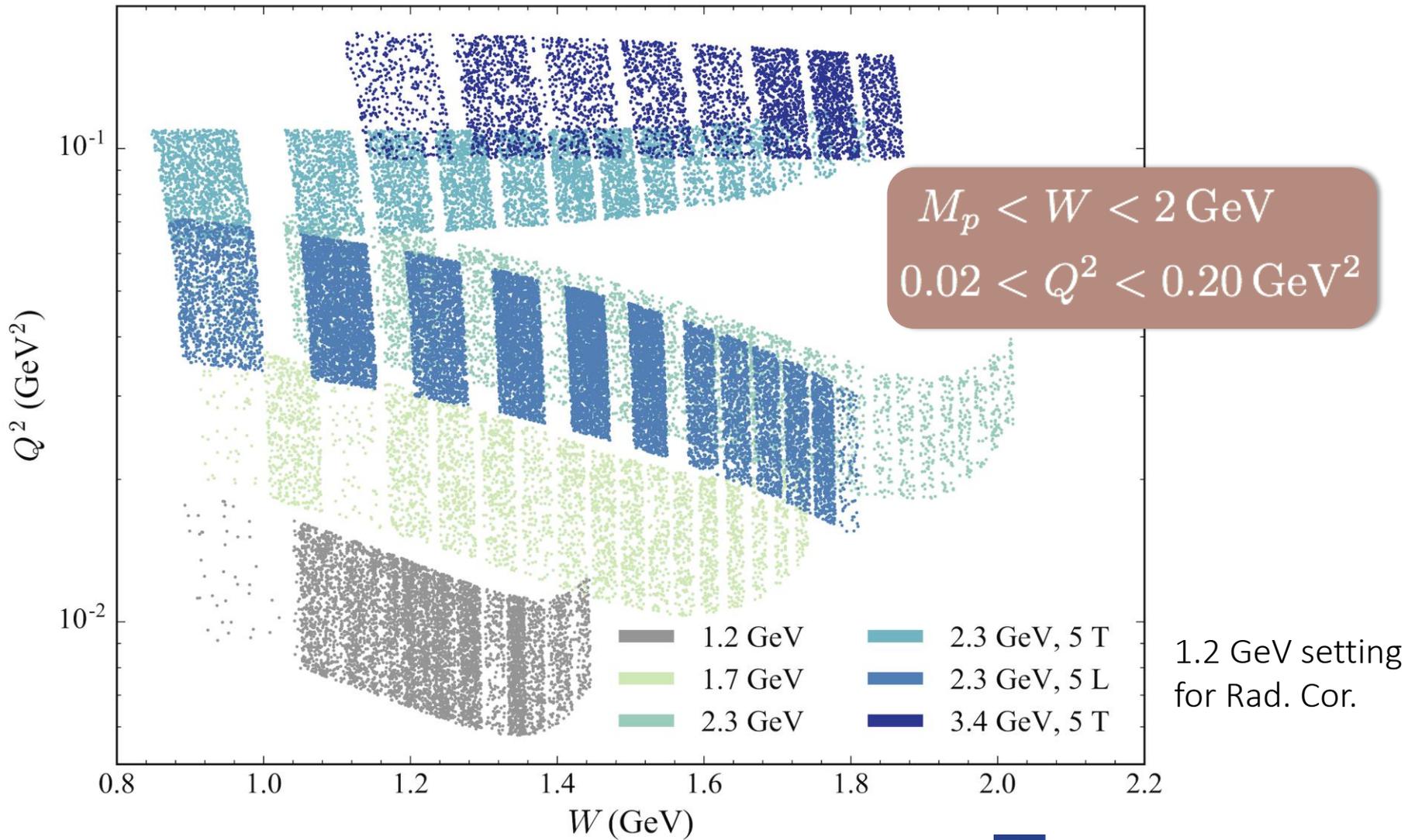
- Electron Beam
- Polarized Proton Target
- Spectrometer/Detectors
- Small Scattering Angle

- Transverse polarized  $NH_3$  target (2.5/5.0T) provided by UVA / JLab target groups
- Dipole chicane magnets help compensate for target field bending of beam



- Dipole septa magnets get around space limitations of spectrometer

# $g_2p$ Kinematic Coverage



# MEASURING $g_{1,2}$ from data

## What can we measure?

1. Helicity dependent asymmetries
2. Unpolarized cross sections
3. Polarized cross sections

1.  $A_{\perp} = \frac{\frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow - \uparrow\Rightarrow)}{\frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow + \uparrow\Rightarrow)}$

2.

$$\sigma_0 = \frac{1}{2} \frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow + \uparrow\Rightarrow)$$

3.

$$\Delta\sigma_{\perp} = \frac{d^2\sigma}{d\Omega dE'}(\downarrow\Rightarrow - \uparrow\Rightarrow) = 2 \cdot A_{\perp} \sigma_0$$

Similar equation for parallel polarized cross section

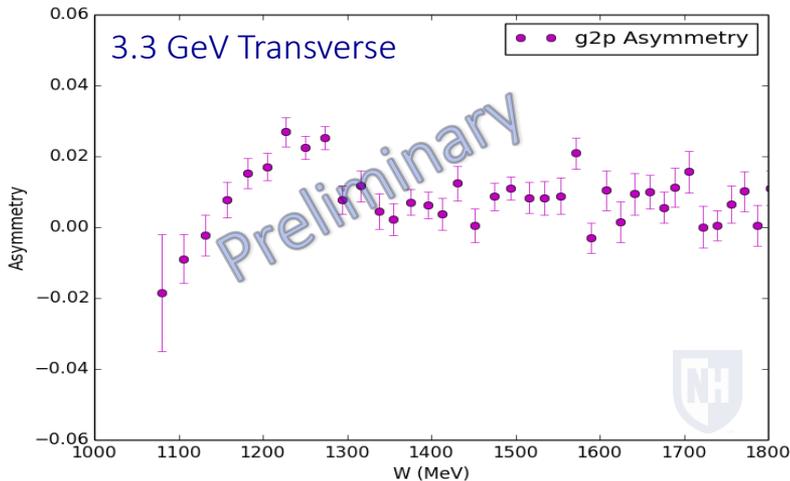
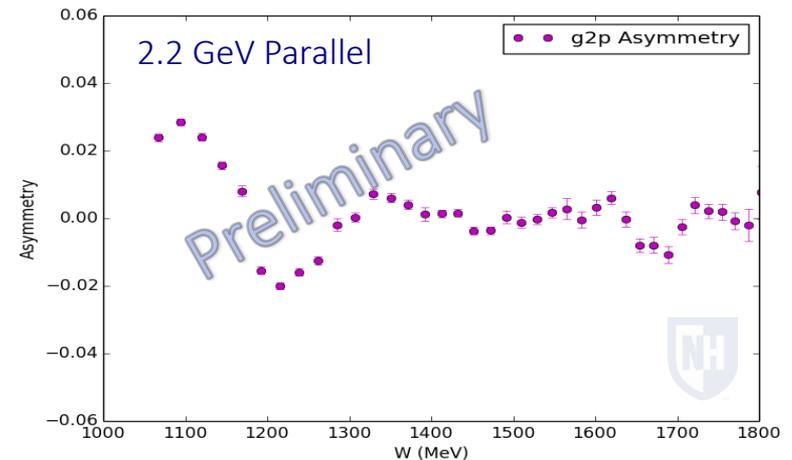
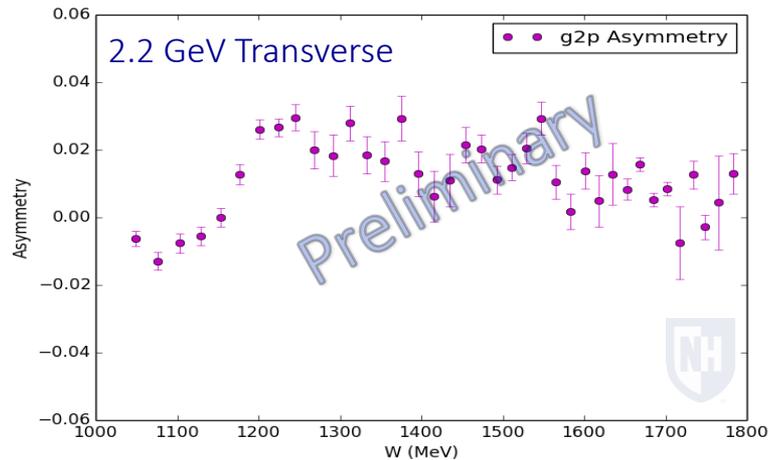
## Why do it this way?

- Asymmetries are easy to measure
- Lots of data on unpolarized cross sections so models are a possibility

Need to be mindful of contributions from scattering from anything other than protons



# 5T Proton Asymmetries



Raw Counts:

$$Y_{\pm} = \frac{N_{\pm}}{LT_{\pm}Q_{\pm}}$$

Measured Asymmetries:

$$A^{\text{raw}} = \frac{Y_{+} - Y_{-}}{Y_{+} + Y_{-}},$$

Combine both  
HRS for best  
statistics!

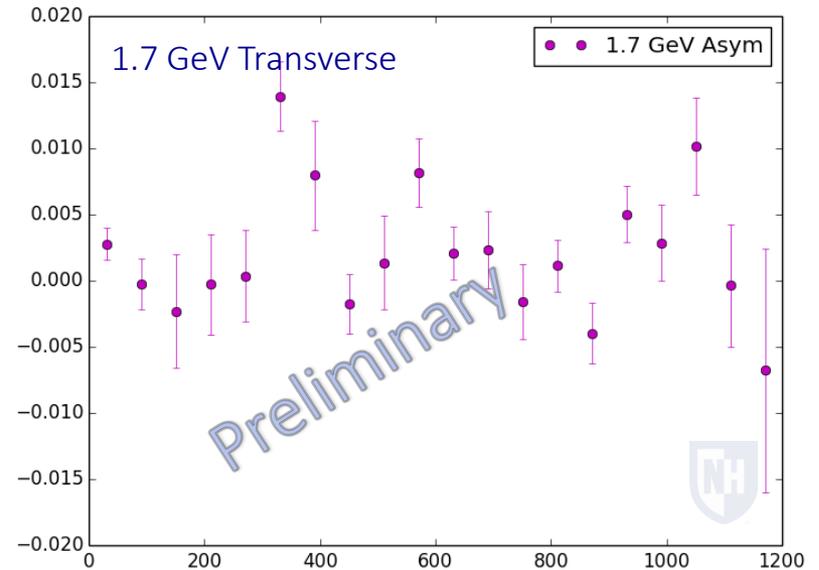
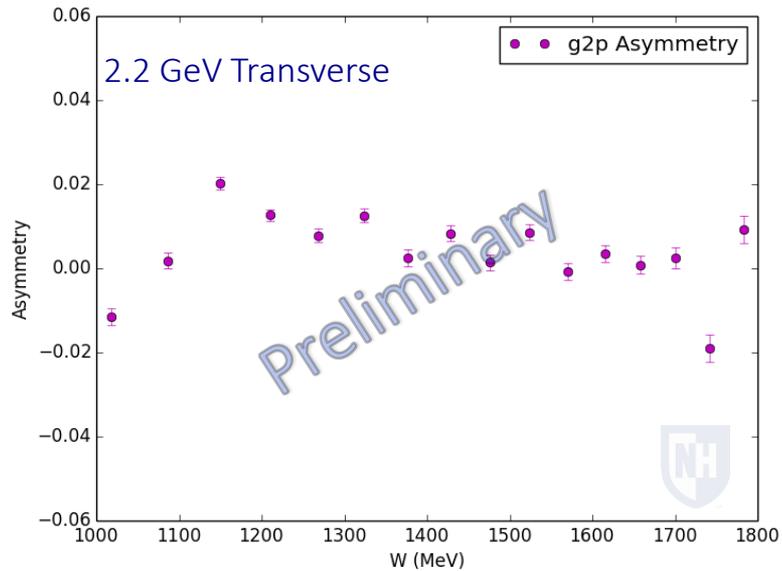
$$A^{\text{exp}} = \frac{1}{f \cdot P_t \cdot P_b} A^{\text{raw}}$$

dilution factor

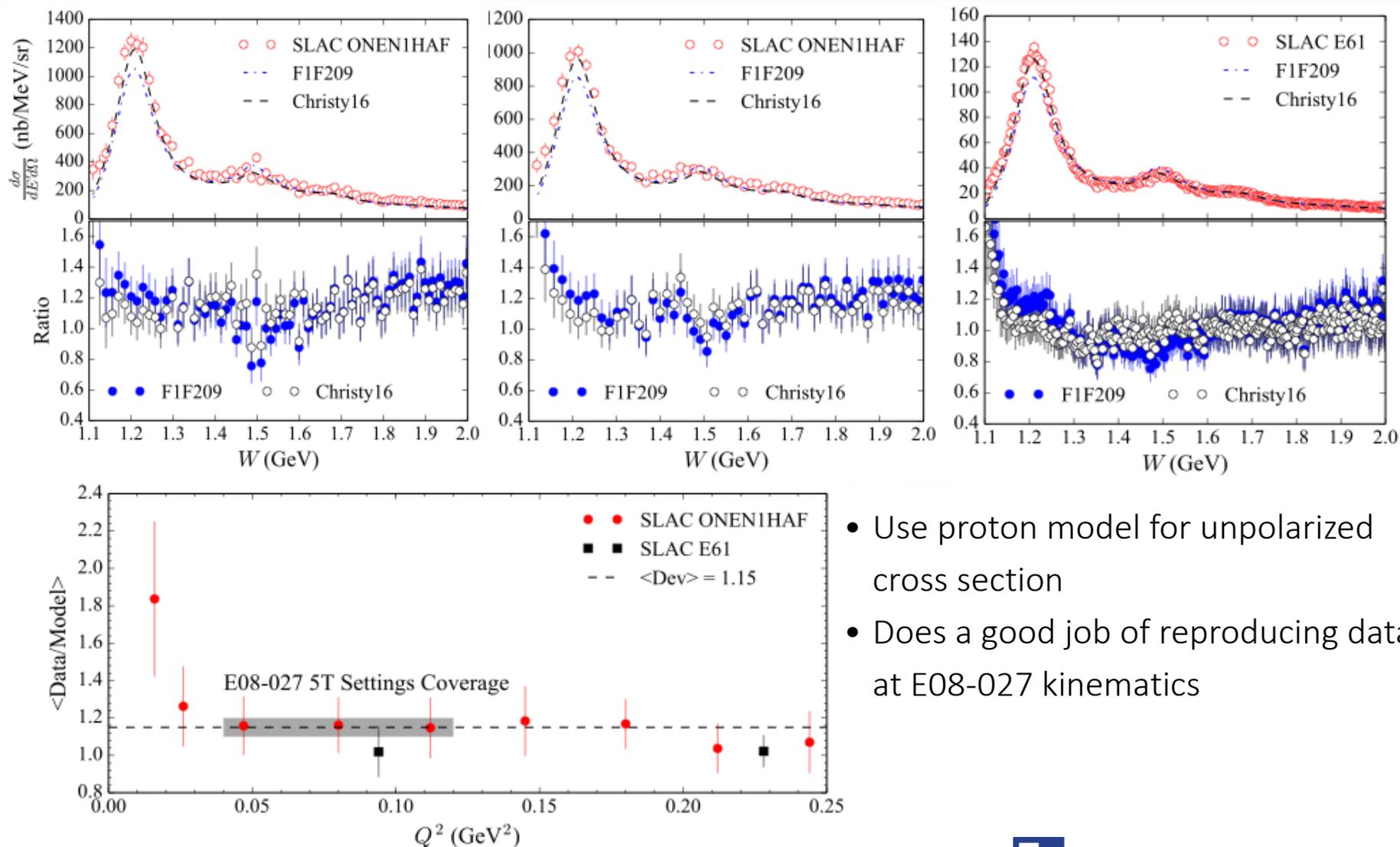
beam/target pol



# 2.5T Proton Asymmetries

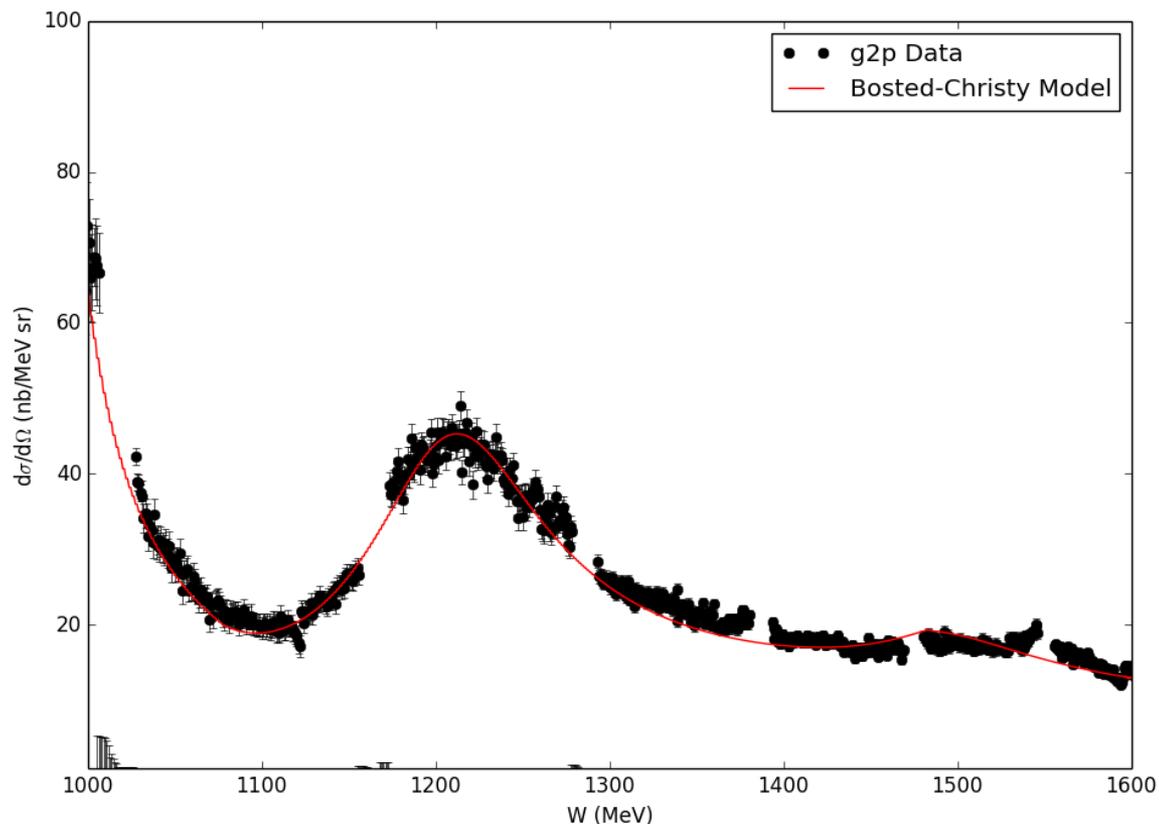


# Model Cross Section



- Use proton model for unpolarized cross section
- Does a good job of reproducing data at E08-027 kinematics

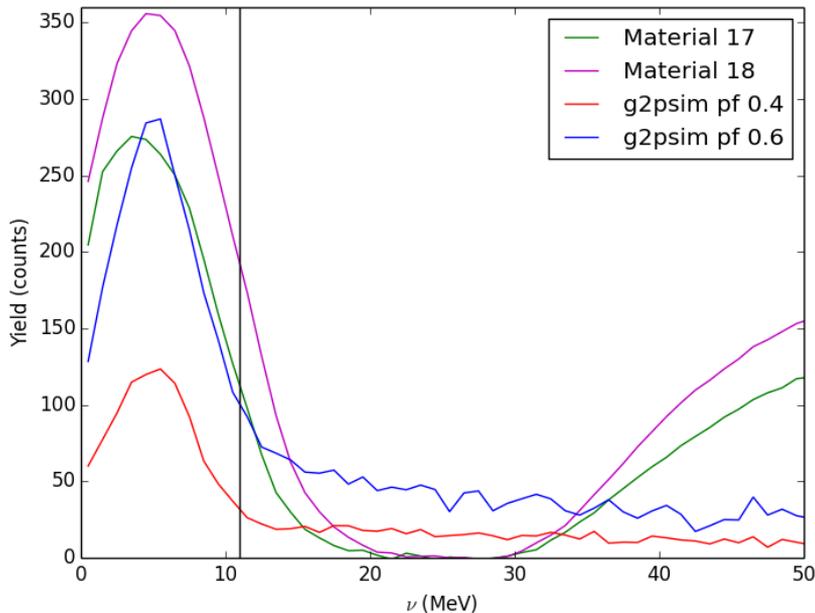
# Model Cross Section



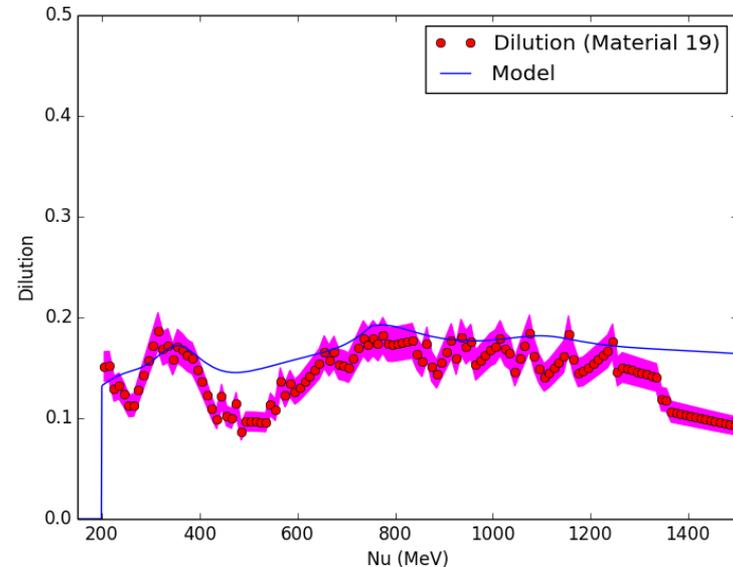
- Direct comparison to g2p Longitudinal cross section yields very similar comparison results
- Necessary to scale model by 1.15 to obtain good agreement for both SLAC and g2p
- Systematic impact on the moments is very small

# Packing Fraction & Dilution Analysis

- Packing fraction describes how much material is in the target cell, important for calculating dilution factor
- Previous packing fraction and dilution analysis yielded unrealistic results, in February I concluded a lengthy re-analysis of both
- Packing Fraction Analysis re-done with Oscar Rondon's method from RSS



- Dilution approximates how much of data comes from other materials
- $$f = \frac{\sigma_{Proton}}{\sigma_{Prod}} = 1 - \frac{Y_N + Y_{He} + Y_{Al}}{Y_{Prod}}$$
- Acceptance effects on edge of momentum settings and BPM calibration issues complicated this analysis



# Extracting the Spin Structure Functions

Model driven procedure for unmeasured part

$$g_2(x, Q^2) = \frac{K_1 y}{2} \left[ \Delta\sigma_{\perp} \left( K_2 + \tan\frac{\theta}{2} \right) \right] - \frac{g_1(x, Q^2) y}{2}$$

$$K_1 = \frac{MQ^2}{4\alpha} \frac{y}{(1-y)(2-y)}$$

$$K_2 = \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta}$$

$$g_1(x, Q^2) = K_1 \left[ \Delta\sigma_{\parallel} \left( 1 + \frac{1}{K_2} \tan\frac{\theta}{2} \right) \right] + \frac{2g_2(x, Q^2)}{K_2 y} \tan\frac{\theta}{2}$$

$$K_1 = \frac{MQ^2}{4\alpha} \frac{y}{(1-y)(2-y)}$$

$$K_2 = \frac{1 + (1-y)\cos\theta}{(1-y)\sin\theta}$$

## Adjusting to a constant $Q^2$

$$\delta_{\text{evolve}} = g_{1,2}^{\text{mod}}(x_{\text{data}}, Q_{\text{data}}^2) - g_{1,2}^{\text{mod}}(x_{\text{const}}, Q_{\text{const}}^2),$$

$$x_{\text{const}} = Q_{\text{const}}^2 / (W^2 - M^2 + Q_{\text{const}}^2),$$

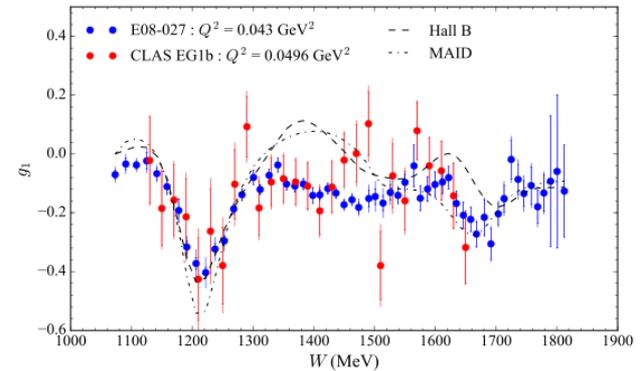
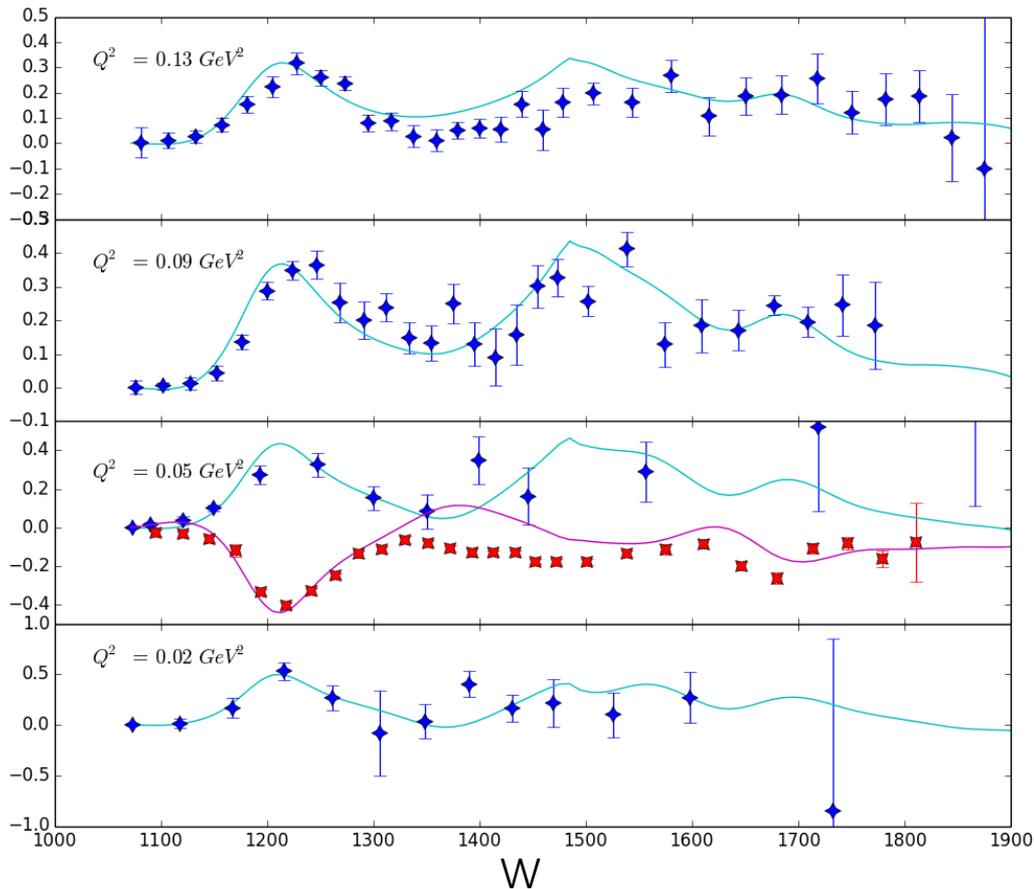
Small effect at the transverse settings



# Structure Function Results

Blue Stars –  $g_2$  (Transverse Setting)

Red Xs –  $g_1$  (Longitudinal Setting)

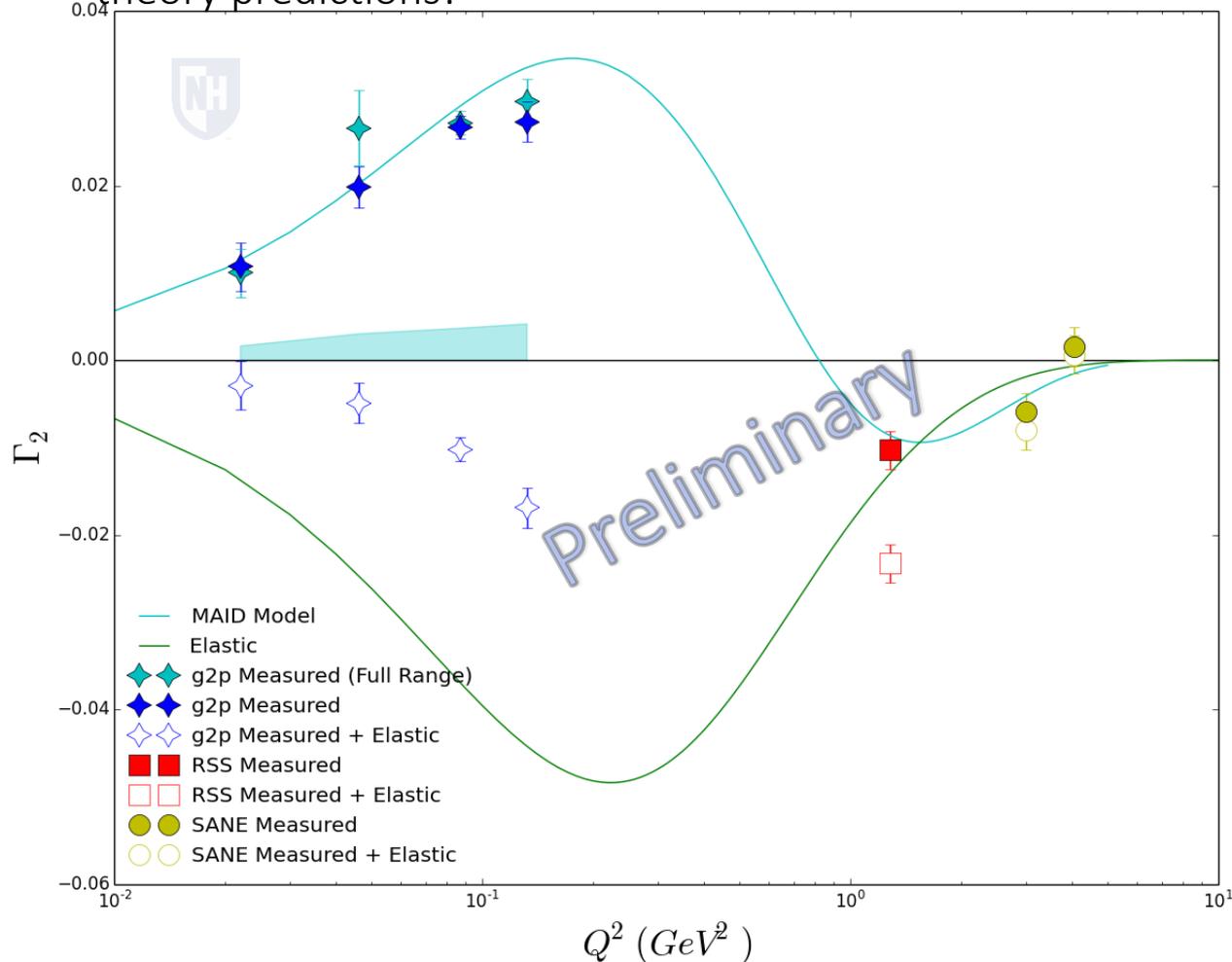


- E08-027 data is consistent with previously published data from CLAS
- But with much better statistics!

# First Moment of $g_2(x, Q^2)$

$$\Gamma_2 = \int_0^{x_{th}} g_2(x, Q^2) dx$$

Moments provide a useful quantity that can be related back to theory predictions!



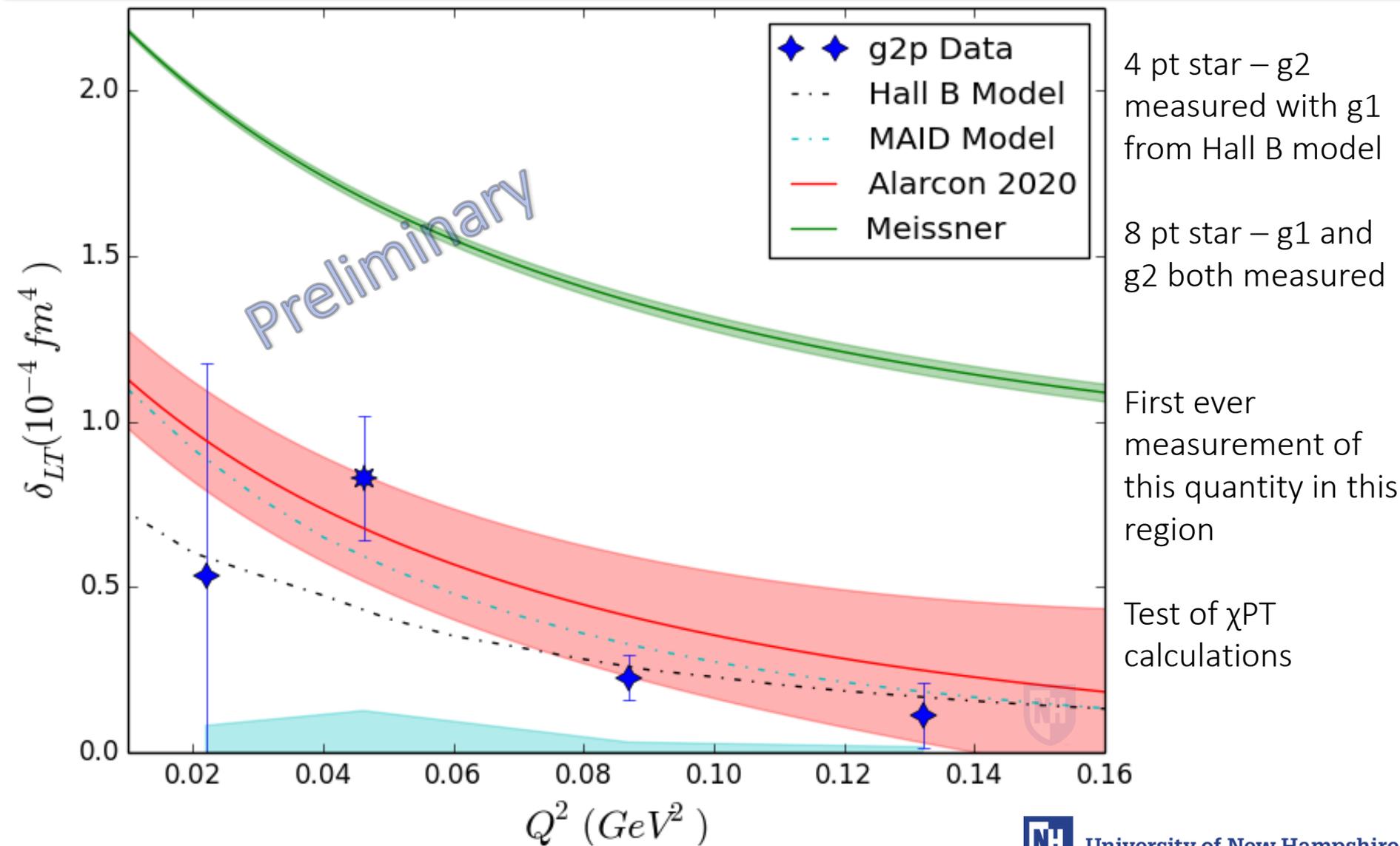
Burkhardt-Cottingham Sum rule says this moment should be zero everywhere...

Unmeasured, low  $x$  part difficult to calculate accurately at low  $Q^2$

Distance between Measured+elastic and zero can be taken as measurement of this hard to measure region if BC sum rule is followed

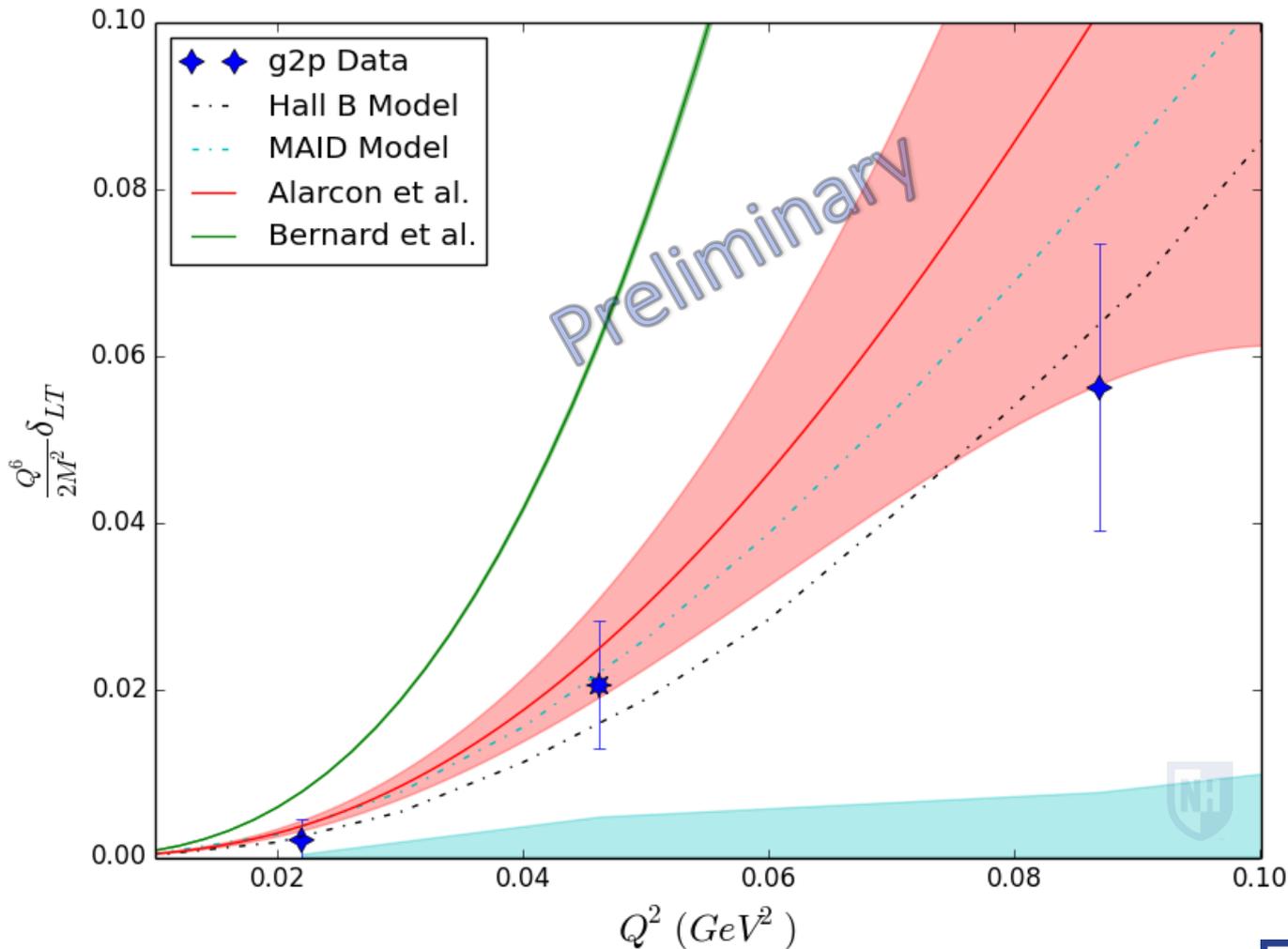
# Transverse-Longitudinal Spin Polarizability

$$\delta_{LT} = \frac{16\alpha M^2}{Q^6} \int_0^{x_{th}} x^2 [g_1(x, Q^2) + g_2(x, Q^2)] dx$$



# Transverse-Longitudinal Spin Polarizability

$$\delta_{LT} = \frac{16\alpha M^2}{Q^6} \int_0^{x_{th}} x^2 [g_1(x, Q^2) + g_2(x, Q^2)] dx$$



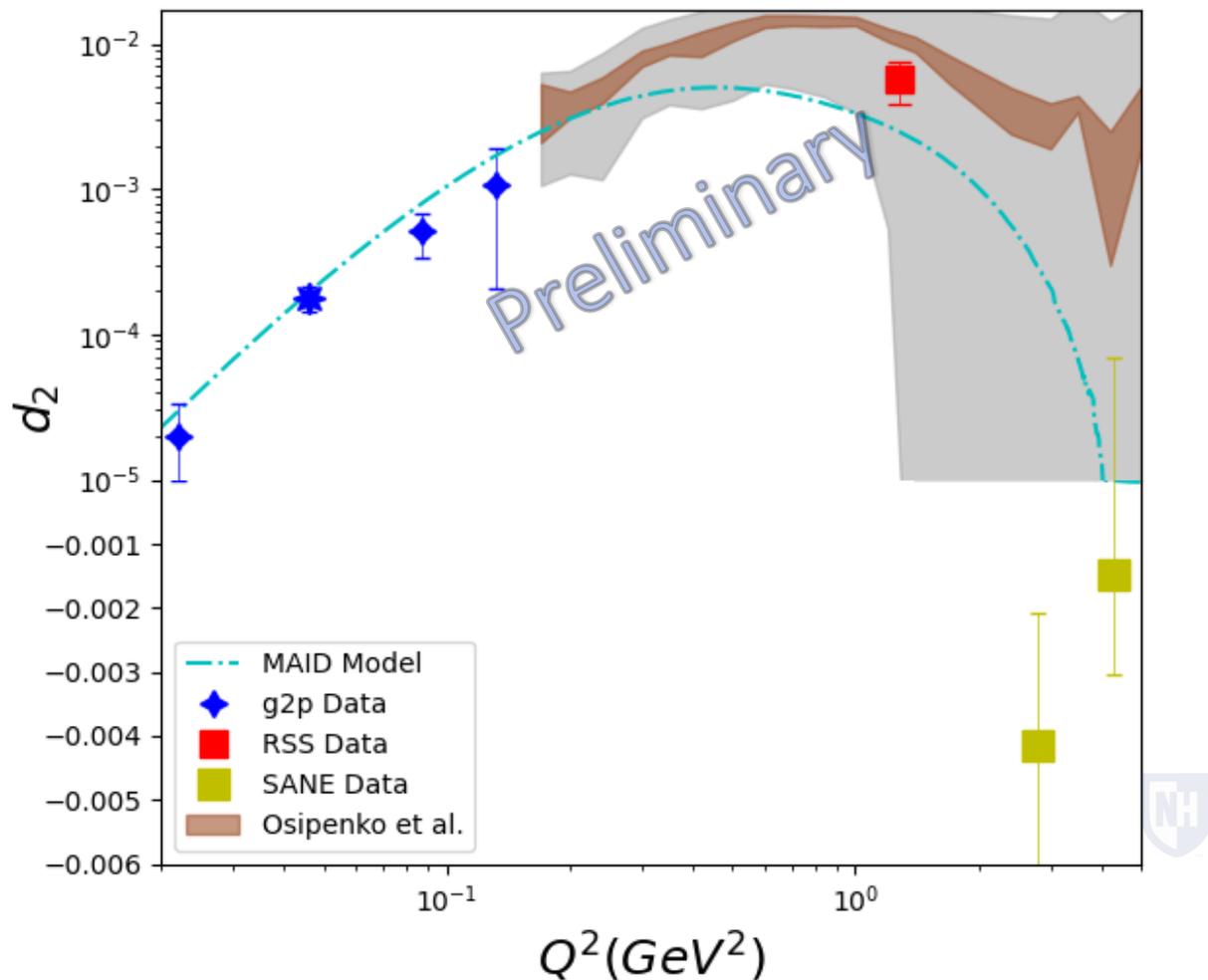
Known difference between calculations –  $\delta$ -Power-Counting vs.  $\epsilon$ -Power-Counting

Low  $Q^2$  data is heavily weighted so many look at a unitless  $Q^6$  scaling of the moment



# Higher Moment $d_2$

$$d_2 = \int_0^{x_{th}} x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)] dx$$

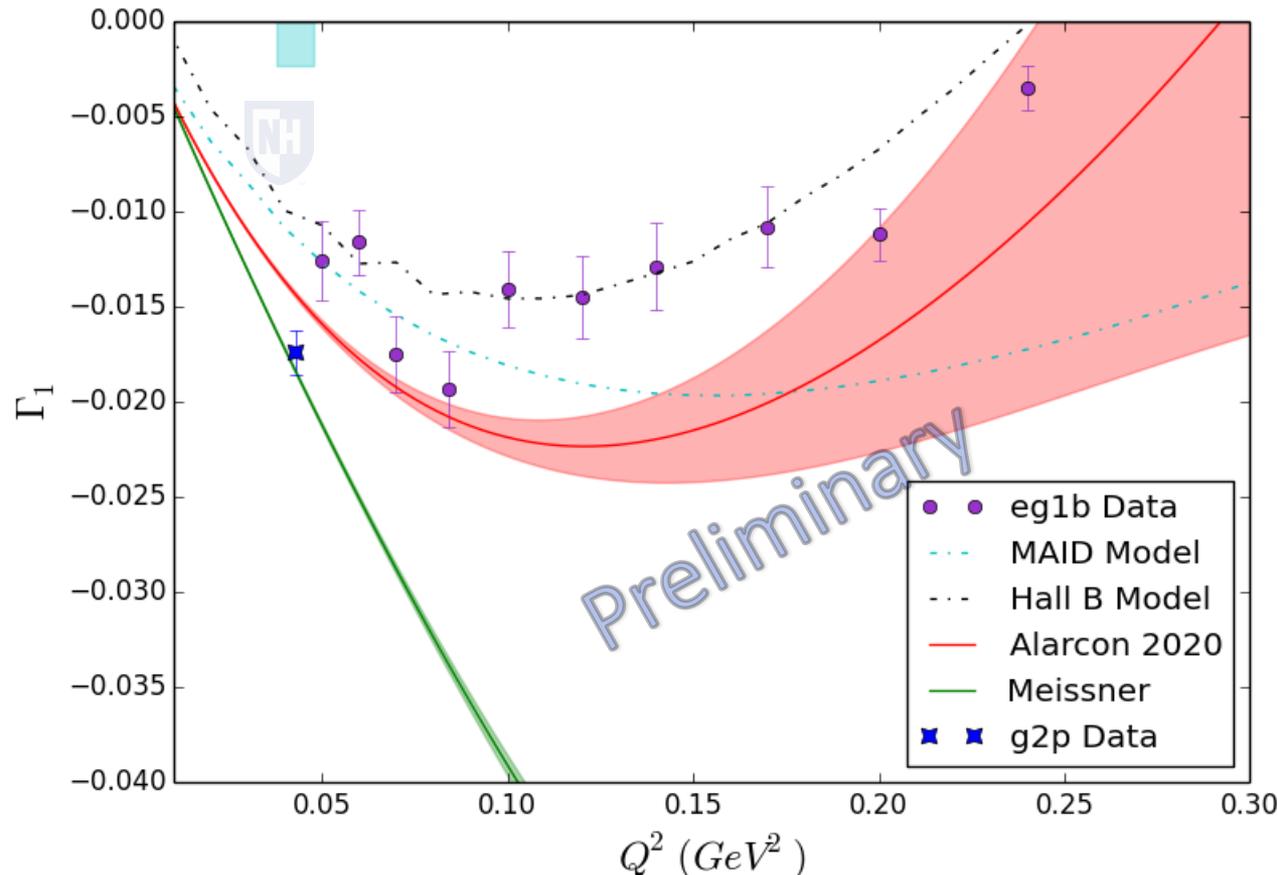


g2p agrees well with phenomenological MAID model

SANE measured anomalously negative moment at high  $Q^2$

# First Moment of $g_1(x, Q^2)$

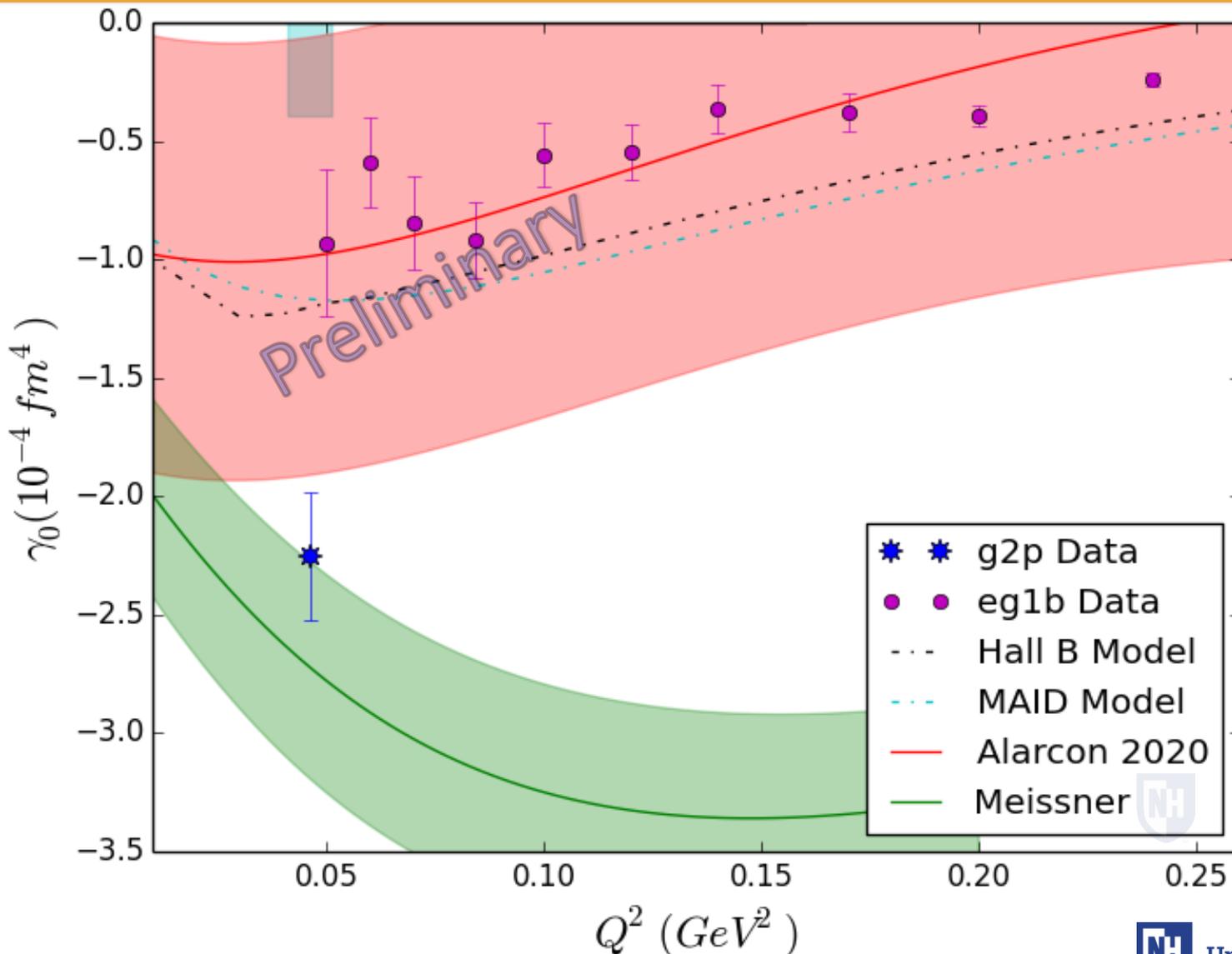
$$\Gamma_1(Q^2) = \int_0^{x_{th}} g_1(x, Q^2) dx$$



Another recently published experiment obtained lower  $Q^2$  data for this moment (EG4, X. Zheng et al., Nature Physics 17-6)

# Generalized Forward Spin Polarizability

$$\gamma_0 = \frac{16\alpha M^2}{Q^6} \int_0^{x_{th}} x^2 g_1(x, Q^2) - \frac{4M^2}{Q^2} x^4 g_2(x, Q^2) dx$$



Strong disagreement with eg1b data

$g_2^p$  data includes measured data for  $g_2$ ...

And goes closer to threshold

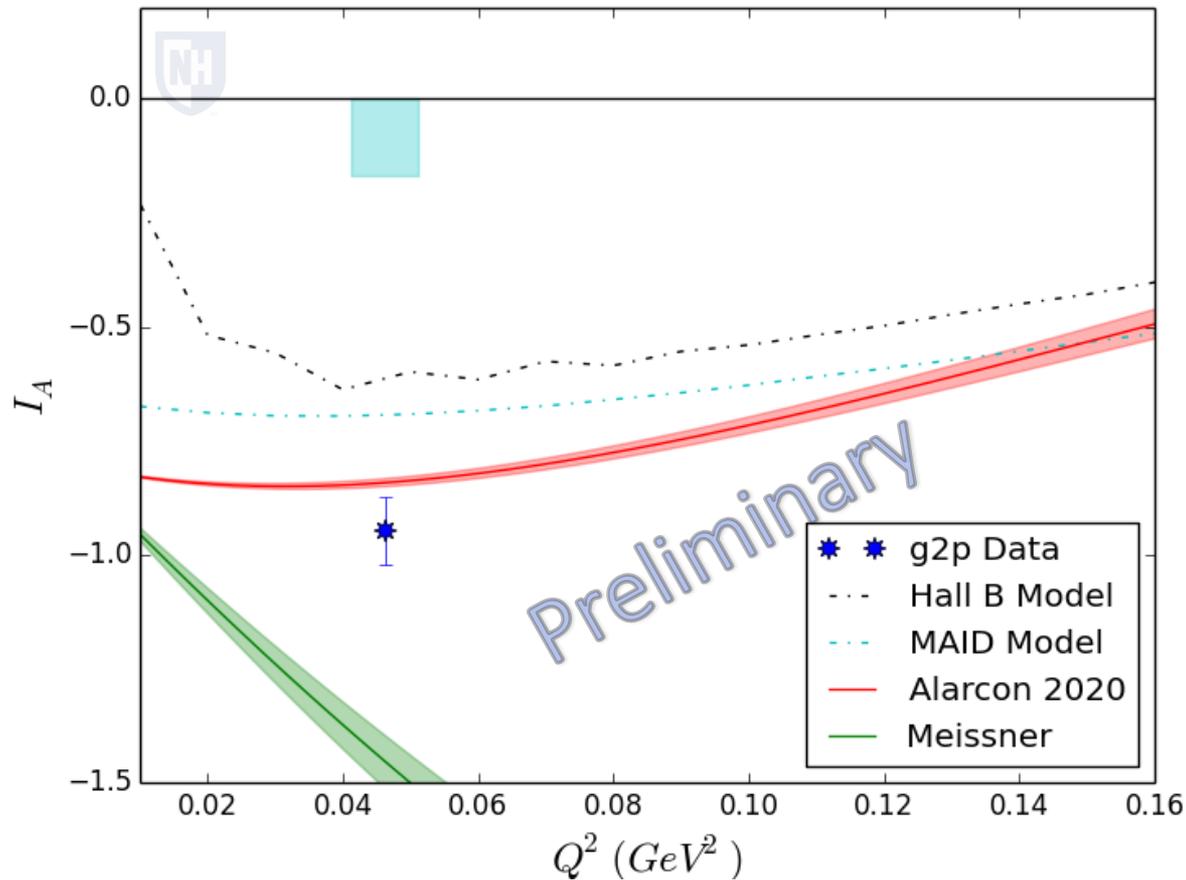
Another recently published experiment obtained lower  $Q^2$  data for this moment (EG4, X. Zheng et al., Nature Physics 17-6)

- \* \*  $g_{2p}$  Data
- ●  $eg_{1b}$  Data
- - - Hall B Model
- · · MAID Model
- Alarcon 2020
- Meissner



# Gerasimov-Drell-Hearn Sum Rule

$$I_A = \frac{2M^2}{Q^2} \int_0^{x_{th}} g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) dx$$



# First publication nearly finished

- Results split into two intended publications: a paper focused on the transverse results which we intend to submit to Nature Physics, following the successful publications in that journal of EG4 and Small-Angle GDH, and a shorter paper focused on the longitudinal results to be submitted as a Physical Review C Rapid Publication.
- The transverse-focused paper is almost finished, we intend to circulate it to all our collaborators for comments after several more rounds of revisions!

## Transverse Proton Spin Structure for $0.021 < Q^2 < 0.13 \text{ GeV}^2$

Author Name<sup>1</sup>  
(The E08-027 Collaboration)

<sup>1</sup>*Author's Institution*  
(Dated: July 6, 2021)

We have extracted the polarized spin structure functions of the proton,  $g_1$  and  $g_2$ , in the nucleon resonance region at four momentum transfer of  $Q^2 = 0.021, 0.045, 0.086$  and  $0.13 \text{ GeV}^2$ . This paper will discuss the  $g_2$  results obtained by measuring transversely polarized yields and using them to form asymmetries and polarized cross section differences  $\Delta\sigma_{\parallel}(\nu, Q^2)$  and  $\Delta\sigma_{\perp}(\nu, Q^2)$ . These asymmetries were formed using measurements from the Jefferson Lab polarized electron beam, a solid polarized target and the Hall A high resolution spectrometers. The structure functions were used to calculate the  $\bar{\Gamma}_2(Q^2)$ ,  $\delta_{LT}(Q^2)$  and  $d_2(Q^2)$  moments of the proton with high precision. Current chiral perturbation theory calculations disagree in the measured region, and our data seems to show a strong preference for one of these calculations over the other. This data represents the first determination of  $\delta_{LT}(Q^2)$  for the proton, which had previously been subject to the “ $\delta_{LT}(Q^2)$  puzzle” in the neutron.

PACS numbers: 11.55.Hx, 25.30.Bf, 29.25.Pj, 29.27.Hj



# Conclusion

- Experimental measurements of proton structure are key to understanding the proton!
- The  $g_2p$  experiment was a precision measurement of proton  $g_2$  in low  $Q^2$  region **for the first time!**
- Analysis is complete!
- Two publications in progress: Transverse-focused paper almost finished!



Extra Slides

# Polarized Protons Created with Dynamic Nuclear Polarization (DNP)

## Creating initial polarization:

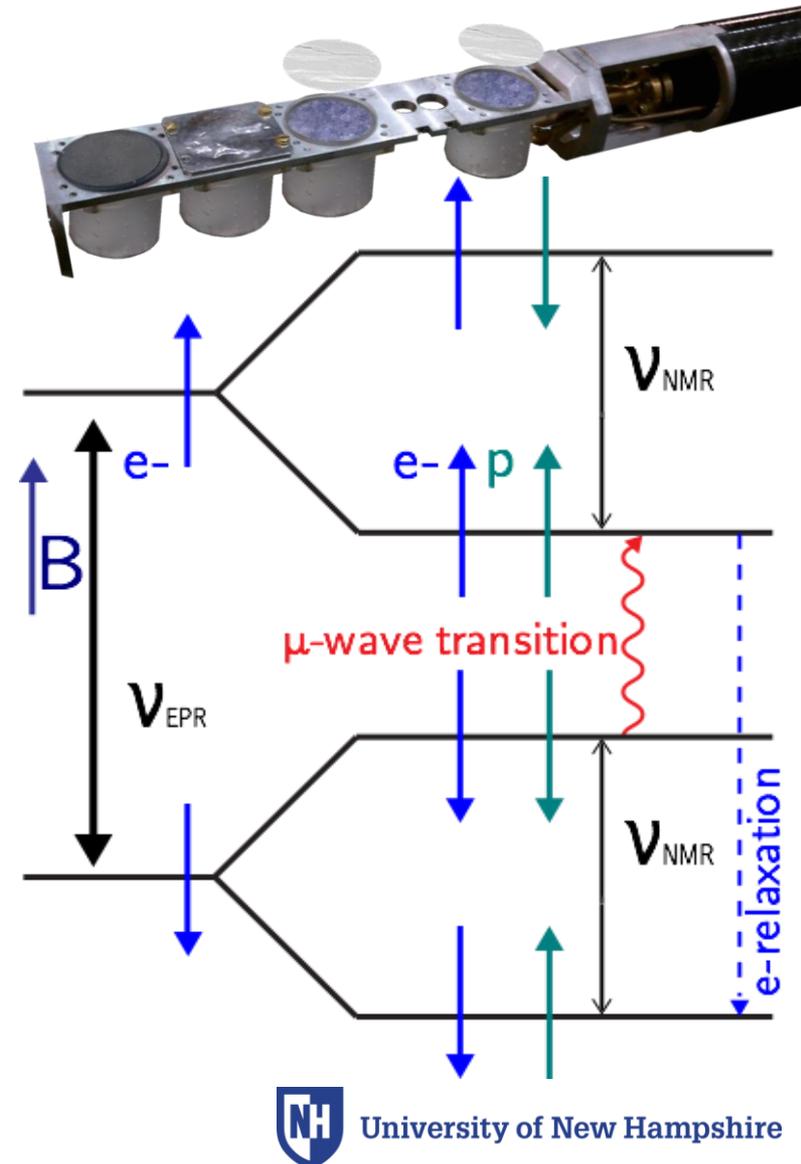
- Align spins in large B and low T
  - 5.0 T/ 2.5 T @ 1 K

$$P_{TE} = \frac{N^\uparrow - N^\downarrow}{N^\uparrow + N^\downarrow} = \frac{e^{\frac{\mu B}{kT}} - e^{-\frac{\mu B}{kT}}}{e^{\frac{\mu B}{kT}} + e^{-\frac{\mu B}{kT}}}$$

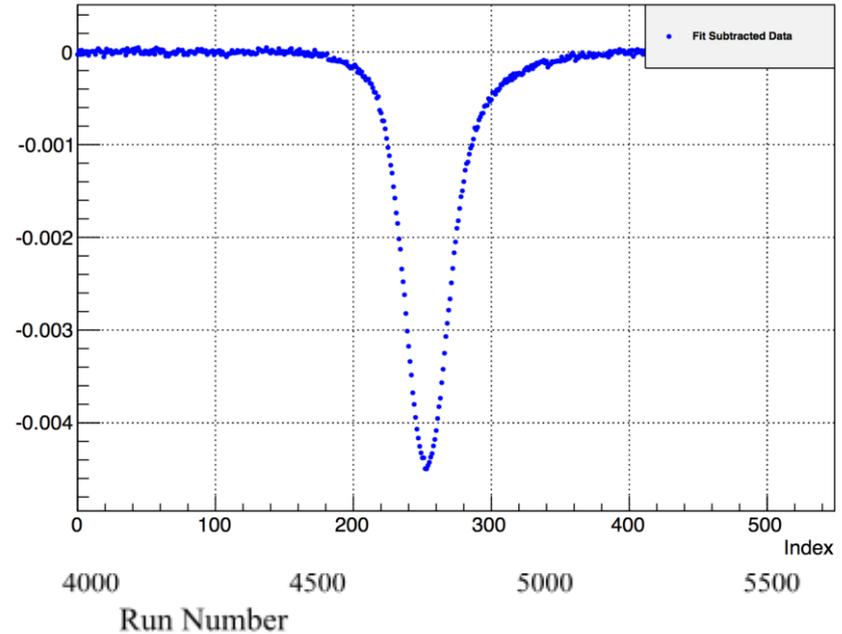
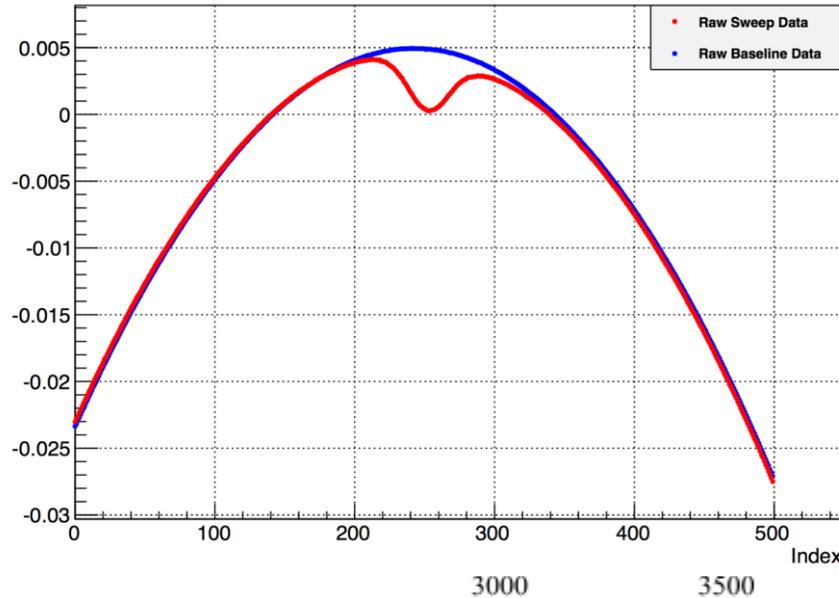
- Large  $\mu_e$  ( $\sim 660\mu_p$ ) creates large electron polarization ( $\sim 99\%$  at 5T/1K)

## Enhancing initial polarization:

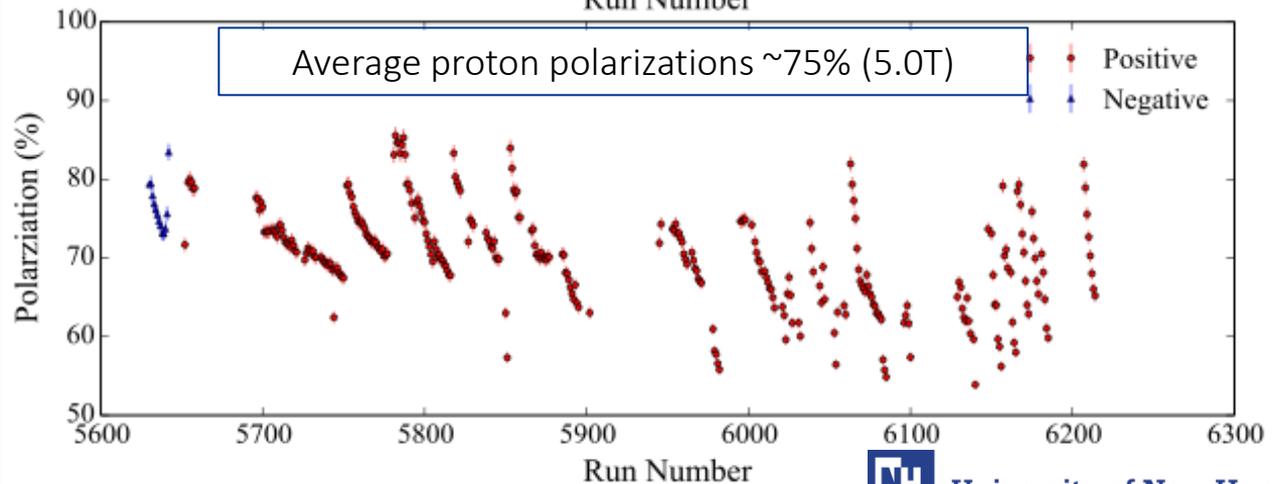
- Proton pol. much smaller ( $\sim 0.5\%$  5T) at TE
- $ep$  spin coupling and microwaves drive pol.
- Electrons relax much quicker than protons so polarization is sustained



# Proton Polarization Measured with Q-Meter



- LRC circuit where proton spin's couple with and change inductance



# EG4 Moment Results

