

The evaluation of the magnetometry using AOC-related NMOR at Cesium D1 Line



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INTRODUCTION

Optical atomic magnetometers based on alkali atom are one of the most sensitive magnetic sensors. The polarization of alkali ensemble plays a crucial role in revealing the property of magneto-optical rotation that further affects the sensitivity of the optical atomic magnetometers. Alignment-to-orientation conversion (AOC) involves the evolution of atomic polarization induced by an applied electric field or by intense laser light ^[1], providing alternative method to modify the polarization state of atom. Therefore, we theoretically analyzed and evaluated the AOC-related nonlinear magneto-optical rotation (NMOR) at ¹³³Cs D1 transition (895 nm) experimentally and theoretically, aiming to explore its value in magnetometry.



METHODS

The experimental apparatus ^[2] is a single beam scheme, as shown in Fig. 1. The temperature of the paraffin coating cell is maintained at 37 °C. The nonperturbative density matrix theory^[3] is used to evaluate the AOC-related NMOR via the AtomicDensityMatrix package ^[4].



Fig. 2. The evaluation of AOC-related NMOR at the 6 ${}^{2}S_{1/2}$ $F = 3 \rightarrow 6 {}^{2}P_{1/2}$ F = 4 transition. (a) Numerical calculations of the polarization rotation vary as a function of the magnetic field. (b) The measured amplitude and linewidth of AOC-related NMOR versus laser power (c) Response of the magnetometer to a sinusoidal magnetic field at different B_z offsets.(d) The amplitude and linewidth of the AOC-related NMOR versus laser power.

An anomaly reversal at the $6^2S_{1/2}$ $F=4 \rightarrow 6^2P_{1/2}$ F=3 transition was experimentally observed at laser power above 3 mW (Fig. 3a). Such reversal at $F \rightarrow F-1$ transition shows some differences from the results of rubidium D2 line, where the flip of the rotation sign causing by AOC only occurred at the $F \rightarrow F+1$ transition ^[1]. We attributed this phenomenon to the hyperfine sublevels effect and decomposed the rotation of laser

Fig. 1. The diagram of the experimental setup. Some abbreviations are used as follows: wavelength meter (WM), noise eater (NE), mirror (M), function generator (FG), attenuator (A), polarizer (P), Wollaston prism (WP).

RESULTS and DISCUSSIONS

Numerical simulations at the $6^2S^{1/2}$ F=3 \rightarrow 6²P_{1/2} F=4 transition were performed to interpret the AOC-related NMOR at high intensity. A "twist" structure (red curve) is clearly visible in NMOR curve (Fig.2a). The angularmomentum probability surface (AMPS)^[5] shows a superposition of the alignment and orientation (the insets in Fig.2a), indicating that the "twist" structure is caused by AOC effect. Subsequently, the experiments were performed. A similar "twist" curve was observed at 0.5 mW (Fig. 2b). The amplitude-to-linewidth ratio (ALR) for the AOCrelated NMOR is 171 µrad/Hz. The response to AC magnetic signals (38 pT, 80 Hz) were measured (Fig. 2c). The maximum response value in the AOC region is 0.428 mrad/nT, which is a 1.7 fold enhancement of the value in the coherence region ($B_7=11$ nT). Further experiments with different laser power (Fig. 2d) show that this AOC-related NMOR maintains an available ALR (>30 µrad/Hz) in the high light intensity region (3~50 mW), benefiting from narrow linewidth (~30 Hz) and weak power broadening (1.37 Hz/mW).

polarization φ_{total} to the sum of the $\varphi_{Fg,Fe}$ (contributed from the phase shift of σ + and σ - between the groundand excited-states hyperfine sublevels) in the Figs. 3b-3d. Only at high saturation parameter κ =10⁷ and low B_z , the contributions from other hyperfine sublevels emerge (Fig. 3c), resulting in a negative rotation.



Fig. 3. The investigations of the reversal of laser polarization rotation at the $6^2S_{1/2}$ $F=4\rightarrow 6^2P_{1/2}$ F=3 transition. (a) Experimental results of the polarization rotation variation as a function of the magnetic field at different laser power. (b-d) Calculated results of the hyperfine sublevels contributed polarization rotation vary with laser detuning at (b) $\kappa=10^3$, $B_z=0.1$ nT, (c) $\kappa=10^7$, $B_z=0.1$ nT, and (d) $\kappa=10^7$, $B_z=200$ nT, respectively. The detuning $\Delta = 0$ GHz corresponds to the center of the cesium D1 line.

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