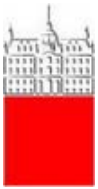


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University
of Ljubljana



“Jožef Stefan”
Institute



Part I

1. Introduction
2. Mixing phenomenology
3. Mixing measurements

Part II

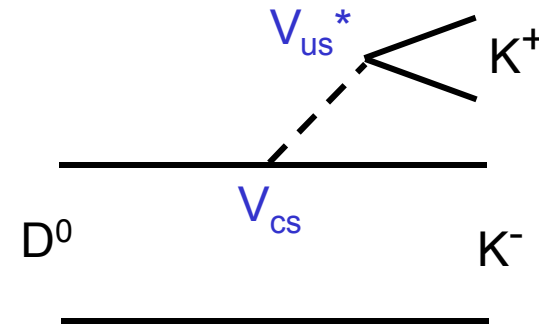
1. CPV phenomenology
2. CPV measurements
3. Constraints on NP
4. Outlook

Topical Seminars on Frontier of Particle Physics,
Hu Yu Village, Aug 27 – Aug 31, 2010

CP violation in charm

Magnitude of CPV
 ...is small. Why?

CPV: complex CKM matrix phase;
 D^0 (and other processes involving charm hadrons):
 first two quark generations;
 CKM elements \approx real;



$$\arg\left(\frac{\langle f | D^0 \rangle}{\langle \bar{f} | \bar{D}^0 \rangle}\right) = \arg\left(\frac{V_{cs} V_{us}^*}{V_{cs}^* V_{us}}\right) = 2 \arg[V_{cs} V_{us}^*]$$

$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + \frac{1}{2}A^2\lambda^5[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4(1 + 4A^2) & A\lambda^2 \\ A\lambda^3[1 - (1 - \frac{1}{2}\lambda^2)(\rho + i\eta)] & -A\lambda^2 + \frac{1}{2}A\lambda^4[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}A^2\lambda^4 \end{pmatrix}$$

CPV of $\mathcal{O}(10^{-3})$ is (just) below current exp. sensitivity;
 larger CPV signals New Physics

$$2 \arg[V_{cs} V_{us}^*] \underset{\text{CKM unitarity}}{\approx} 2 \arg[-V_{cd} V_{ud}^* - V_{cb} V_{ub}^*] \approx -2 \frac{A^2 \lambda^5 \eta}{\lambda} = -2 A^2 \lambda^4 \eta = 1.15 \cdot 10^{-3}$$

CP violation in charm Parametrization

...is sometimes messy

$R_D \neq 1$: Cabibbo suppression

3 types of CPV:

$A_D \neq 0$: CPV in decay

$A_M \neq 0$: CPV in mixing

$\phi \neq 0$: CPV in interference

quantity appearing in
 decay rates (" λ_f "):

$$\left| \frac{\langle \bar{f} | D^0 \rangle}{\langle f | D^0 \rangle} \right| \equiv \sqrt{R_D},$$

$$\left| \frac{\langle f | D^0 \rangle}{\langle \bar{f} | \bar{D}^0 \rangle} \right| \equiv 1 + \frac{A_D}{2},$$

$$\frac{q}{p} \equiv \left(1 + \frac{A_M}{2}\right) e^{i\phi}$$

$$\frac{q}{p} \frac{\langle f | \bar{D}^0 \rangle}{\langle f | D^0 \rangle} \equiv - \frac{(1 + A_M / 2) \sqrt{R_D}}{1 + A_D / 2} e^{-i(\delta_f - \phi)}$$

n.b.:

$$(R_D), A_D, A_M, \phi \ll 1$$

CP violation in charm

Parametrization direct CPV

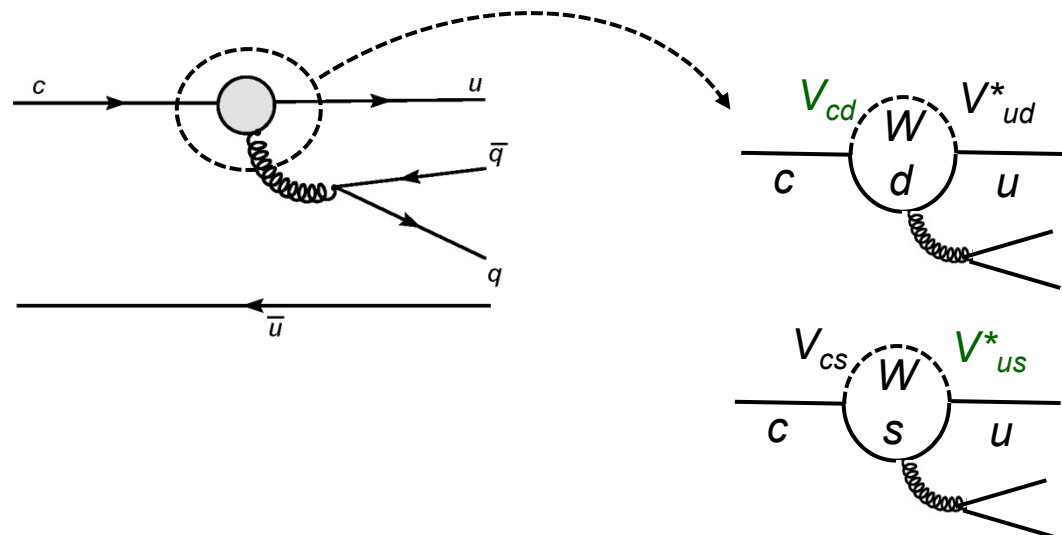
for direct CPV two amplitudes with different strong and weak (CKM) phases are necessary;

in D meson decays this is only possible in CS decays with contribution of penguin decays (beside tree contrib.)

$$A_f = a_1 + a_2 = |a_1| e^{i(\delta_1 + \varphi_1)} + |a_2| e^{i(\delta_2 + \varphi_2)}$$

$$A_{CP} = \frac{\Gamma(M \rightarrow f) - \Gamma(\bar{M} \rightarrow \bar{f})}{\Gamma(M \rightarrow f) + \Gamma(\bar{M} \rightarrow \bar{f})} = \frac{|A_f / \bar{A}_{\bar{f}}|^2 - 1}{|A_f / \bar{A}_{\bar{f}}|^2 + 1} =$$

$$= \dots = \frac{2 |a_1 a_2| \sin(\delta_2 - \delta_1) \sin(\varphi_2 - \varphi_1)}{|a_1|^2 + |a_2|^2 + 2 |a_1 a_2| \cos(\delta_2 - \delta_1) \cos(\varphi_2 - \varphi_1)}$$



CP violation in charm

Observables

master formulas p. 1/18, 19
 still valid, need to keep
 q/p and write amplitudes
 A_f etc. in accordance with
 parametrization on previous
 slide

$$\frac{dN(D^0 \rightarrow f)}{dt} \propto e^{-\bar{\Gamma}t} \left| A_f + \frac{q}{p} \frac{ix+y}{2} \bar{A}_f t \right|^2$$

uncorrelated production

$$\begin{aligned} \Gamma(V \rightarrow D^0 \bar{D}^0 \rightarrow f_1 f_2) &= \\ &= \frac{1}{2} |a_-|^2 \left(\frac{1}{1-y^2} + \frac{1}{1+x^2} \right) + \frac{1}{2} |b_-|^2 \left(\frac{1}{1-y^2} - \frac{1}{1+x^2} \right) \end{aligned}$$

$$a_- = A_{f_1} \bar{A}_{f_2} - \bar{A}_{f_1} A_{f_2}; \quad b_- = \frac{p}{q} A_{f_1} A_{f_2} - \frac{q}{p} \bar{A}_{f_1} \bar{A}_{f_2}$$

coherent production with $C=-1$

Decays to CP eigenstates

$$f = \bar{f}; \quad A_f = A_{\bar{f}}; \quad \bar{A}_f = \bar{A}_{\bar{f}}; \quad \left| \frac{A_f}{\bar{A}_f} \right| = 1 + \frac{A_D}{2}$$

Principle

t-dependent method

$$\frac{|\langle f | P^0(t) \rangle|^2}{|A_f|^2 e^{-t}} = \left[1 + \left(1 + \frac{A_M}{2} - \frac{A_D}{2}\right)(x \sin \phi - y \cos \phi) t \right]$$

measuring lifetime
 in $K^+K^-/\pi^+\pi^-$ state
 separately for D^0
 and \bar{D}^0

$$\frac{|\langle f | \bar{P}^0(t) \rangle|^2}{|A_f|^2 e^{-t}} = \left[1 - A_D - \left(1 - \frac{A_M}{2} - \frac{A_M}{2}\right)(x \sin \phi + y \cos \phi) t \right]$$

following derivation on p. I/26 and keeping CPV parameters

$$\tau_{KK} = \tau / (1 + y_{CP}); \quad y_{CP} = (\tau / \tau_{KK}) - 1 = y \cos \phi - (A_M / 2) x \sin \phi$$

no CPV: $y_{CP} = y$

$$\tau_{KK} \approx \tau \left[1 + \left(1 + \frac{A_M}{2} - \frac{A_D}{2}\right)(x \sin \phi - y \cos \phi) \right],$$

$$\bar{\tau}_{KK} \approx \tau \left[1 - \left(1 - \frac{A_M}{2} - \frac{A_D}{2}\right)(x \sin \phi + y \cos \phi) \right]$$

following same derivation
 separately for D and \bar{D}

$$A_\Gamma \equiv \frac{\bar{\tau}_{KK} - \tau_{KK}}{\bar{\tau}_{KK} + \tau_{KK}} = \frac{A_M}{2} y \cos \phi - \left(1 - \frac{A_D}{2}\right) x \sin \phi \approx \frac{A_M}{2} y \cos \phi - x \sin \phi$$

Decays to CP eigenstates

Results

t-dependent method

Belle, PRL 98, 211803 (2007), 540fb⁻¹

$$A_{\Gamma} = (0.01 \pm 0.30 \pm 0.15)\%$$

BaBar, PRD78, 011105 (2008), 384fb⁻¹

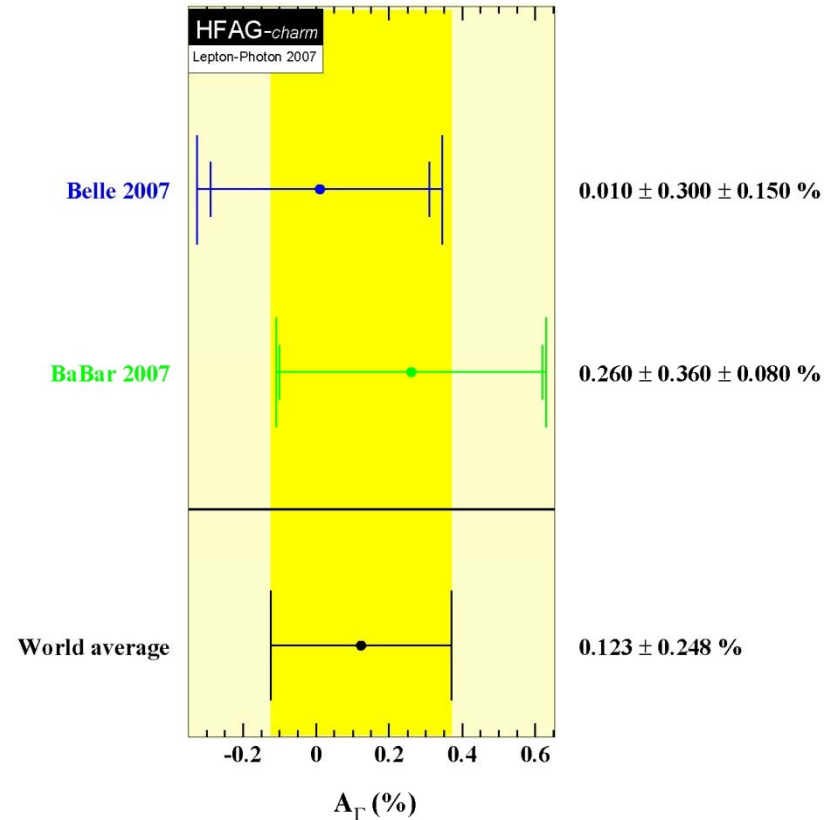
$$A_{\Gamma} = (0.26 \pm 0.36 \pm 0.08)\%$$

dominant syst.: same as for y_{CP}

$$A_{\Gamma} = (A_M/2)y \cos \phi - x \sin \phi$$

CPV in mixing and interference

HFAG, <http://www.slac.stanford.edu/xorg/hfag/>



no CPV at $\sim 3 \cdot 10^{-3}$

Decays to CP eigenstates

Principle

t-integrated method

asymmetry of
 t-integrated rates;
 CPV in decay, mixing and
 interference;

However.....
 measuring absolute rates
 instead of decay-t distrib.
 involves sensitivity to
 acceptance

$$\frac{\Gamma(D^0 \rightarrow K^+ K^-)}{|A_f|^2} = 1 + \left(1 + \frac{A_M}{2} - \frac{A_D}{2}\right)(x \sin \varphi - y \cos \varphi)$$

$$\frac{\Gamma(\bar{D}^0 \rightarrow K^+ K^-)}{|A_f|^2} = 1 - A_D - \left(1 - \frac{A_M}{2} - \frac{A_D}{2}\right)(x \sin \varphi + y \cos \varphi)$$

$$A_{CP}^{KK} = \frac{\Gamma(D^0 \rightarrow KK) - \Gamma(\bar{D}^0 \rightarrow KK)}{\Gamma(D^0 \rightarrow KK) + \Gamma(\bar{D}^0 \rightarrow KK)} \approx$$

$$\approx \frac{A_D}{2} + x \sin \varphi - \frac{A_M}{2} y \cos \varphi$$

n.b.: A_M and ϕ universal among various
 decay modes; A_D is decay mode specific

Decays to CP eigenstates

$$A^{meas} = \frac{N(D^0 \rightarrow KK) - N(\bar{D}^0 \rightarrow KK)}{N(D^0 \rightarrow KK) + N(\bar{D}^0 \rightarrow KK)} =$$

$$= A_{\epsilon}^{\pi} + A_{FB} + A_{CP}^{KK}$$

Principle

t-integrated method

A_{ϵ}^{π} : π^+ / π^- detection eff. asymmetry
 $D^{*+/-} \rightarrow D^0 \pi^{+/-}$; e.g. due to different $\pi^{+/-}$ interactions on detect. material

comparison of tagged/untagged
 $(D^{*+} \rightarrow D^0 \pi^+), D^0 \rightarrow K^- \pi^+$

$$A^{untag} = A_{FB}^{D^0} + A_{CP}^{K\pi} + A_{\epsilon}^{K\pi}$$

$$A^{tag} = A_{FB}^{D^{*+}} + A_{CP}^{K\pi} + A_{\epsilon}^{K\pi} + A_{\epsilon}^{\pi_{slow}}$$

A_{FB} : forward-backward asymmetry

$$\gamma^*/Z^0 \rightarrow c \bar{c};$$

A_{FB} is an odd function of θ_D (in CMS);
 vanishes if integrated over θ_D ;
 since working in bins of θ_{π} (correlated with θ_D) need to correct for it

assuming $A_{FB}^{D^*} = A_{FB}^D \Rightarrow$

$$A^{tag} - A^{untag} = A_{\epsilon}^{\pi};$$

need to perform meas. in bins of p_{π}, θ_{π}

$$A_{FB} = \frac{A^{meas}(\cos \theta_D) - A^{meas}(-\cos \theta_D)}{2}$$

A_{CP}^f : physical CPV asymmetry

$$A_{CP} = \frac{A^{meas}(\cos \theta_D) + A^{meas}(-\cos \theta_D)}{2}$$

Decays to CP eigenstates

Results

t-integrated method

BaBar, PRL 100, 061803 (2007), 386fb⁻¹

$$A_{CP}^{KK} = (0.00 \pm 0.34 \pm 0.13)\%$$

Belle, PLB670, 190 (2008), 540fb⁻¹

$$A_{CP}^{KK} = (-0.43 \pm 0.30 \pm 0.11)\%$$

stat. precision of $\pi\pi$ somewhat worse;
 dominant syst.: stat. uncertainty of A_{ε}^{π}

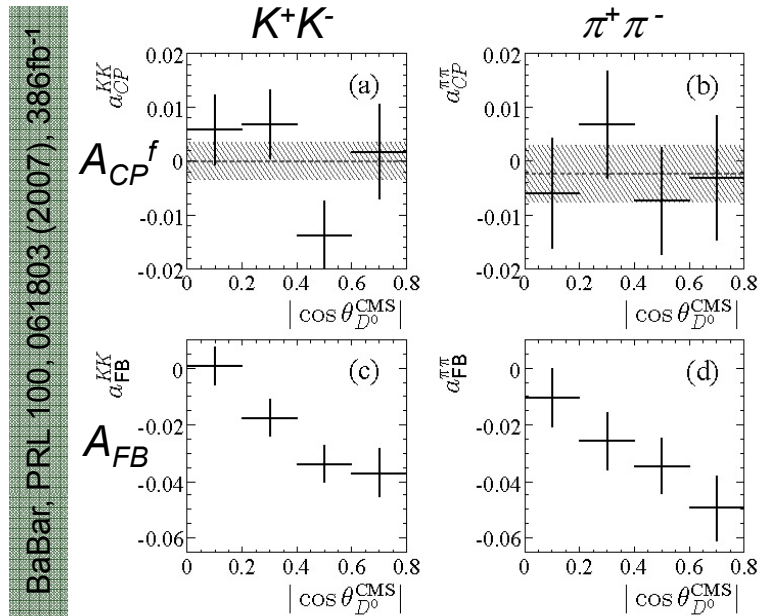
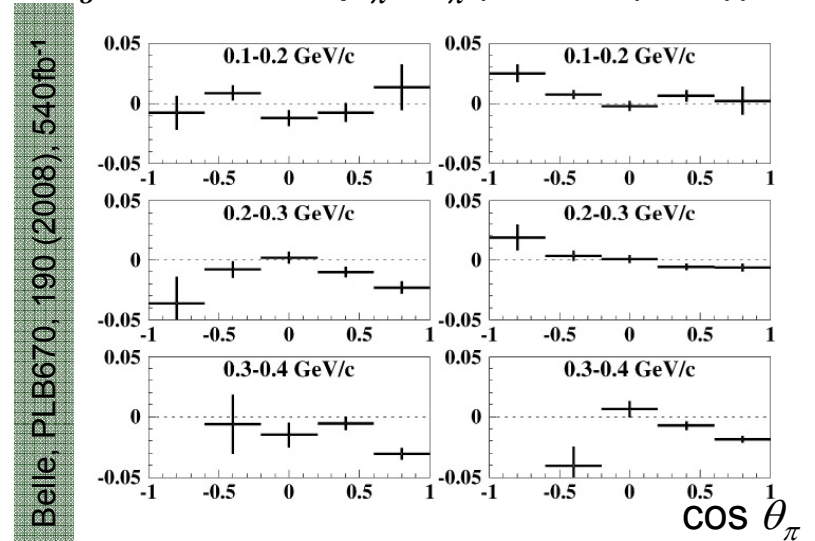
world average:

$$A_{CP}^{KK} = (-0.16 \pm 0.23)\%$$

HFAG,
<http://www.slac.stanford.edu/xorg/hfag/>

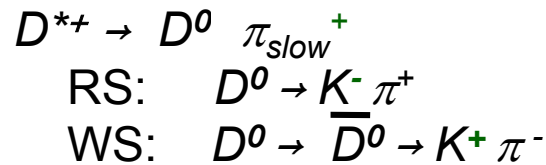
no CPV at $\sim 2 \cdot 10^{-3}$

A_{ε}^{π} in bins of p_{π}, θ_{π} (n.b.: $\mathcal{O}(10^{-2})$)



WS 2-body decays

Principle



equivalent measurement
 separately for D^0, \bar{D}^0
 $(x'^2, y', R_D) \rightarrow (x'^{\pm 2}, y'^{\pm}, R_D^{\pm});$

CPV in decay and mixing

$$\left| \langle K^+ \pi^- | D^0(t) \rangle \right|^2 \propto \left[\underbrace{R_D^+}_{\text{DCS}} + \underbrace{\sqrt{R_D^+} y'^+ t}_{\text{interf.}} + \underbrace{\frac{x'^{+2} + y'^{+2}}{4} t^2}_{\text{mix}} \right] e^{-t}$$

$$\left| \langle K^- \pi^+ | \bar{D}^0(t) \rangle \right|^2 \propto \left[\underbrace{R_D^-}_{\text{DCS}} + \underbrace{\sqrt{R_D^-} y'^- t}_{\text{interf.}} + \underbrace{\frac{x'^{-2} + y'^{-2}}{4} t^2}_{\text{mix}} \right] e^{-t}$$

$$R_M^{\pm} = \frac{x'^{\pm 2} + y'^{\pm 2}}{2} \quad A_M = \frac{R_M^+ - R_M^-}{R_M^+ + R_M^-}$$

$$A_D = \frac{R_D^+ - R_D^-}{R_D^+ + R_D^-}$$

WS 2-body decays

Results

$$D^0 \rightarrow K^+ \pi^-$$

BaBar, PRL 98, 211802 (2007), 384fb⁻¹

R_D	$3.03 \pm 0.16 \pm 0.10$
A_D	$-21 \pm 52 \pm 15$
x'^{2+}	$-0.24 \pm 0.43 \pm 0.30$
y'^{+}	$9.8 \pm 6.4 \pm 4.5$
x'^{2-}	$-0.20 \pm 0.41 \pm 0.29$
y'^{-}	$9.6 \pm 6.1 \pm 4.3$

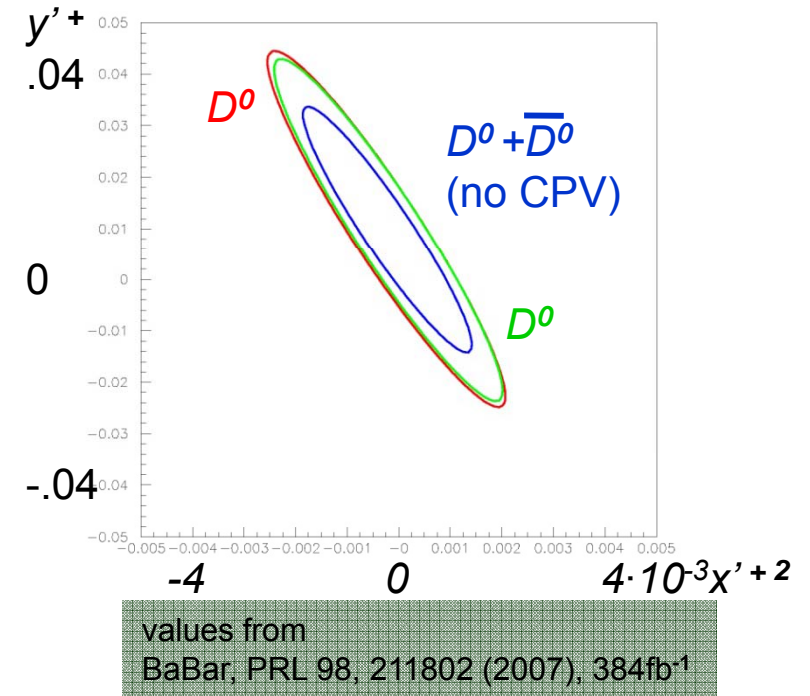
Belle, PRL 96, 151801 (2006), 400fb⁻¹

$$A_D = (23 \pm 47) \cdot 10^{-3}$$

$$A_M = (670 \pm 1200) \cdot 10^{-3}$$

CDF, PRL 100, 121802 (2008), 1.5fb⁻¹

does not fit for CPV param.



Multi-body self conjugated states

same as eqs. on p. 1/42
 but including q/p

Principle

$$D^0 \rightarrow K_S \pi^+ \pi^-$$

t-depedent matrix
 elements $\mathcal{M}, \overline{\mathcal{M}}$
 are in case of CPV
 not trivially related;

$$\mathcal{M}(m_-^2, m_+^2, t) \equiv \langle K_S \pi^+ \pi^- | D^0(t) \rangle = \frac{1}{2} \mathcal{A}(m_-^2, m_+^2) [e^{-i\lambda_1 t} + e^{-i\lambda_2 t}] + \frac{1}{2} \frac{q}{p} \overline{\mathcal{A}}(m_-^2, m_+^2) [e^{-i\lambda_1 t} - e^{-i\lambda_2 t}]$$

$$\overline{\mathcal{M}}(m_-^2, m_+^2, t) \equiv \langle K_S \pi^+ \pi^- | \overline{D}^0(t) \rangle = \frac{1}{2} \overline{\mathcal{A}}(m_-^2, m_+^2) [e^{-i\lambda_1 t} + e^{-i\lambda_2 t}] + \frac{1}{2} \frac{p}{q} \mathcal{A}(m_+^2, m_-^2) [e^{-i\lambda_1 t} - e^{-i\lambda_2 t}]$$

no CPV: $\frac{q}{p} = 1, \overline{\mathcal{A}}(m_-^2, m_+^2) = \mathcal{A}(m_+^2, m_-^2) \Rightarrow \overline{\mathcal{M}}(m_-^2, m_+^2, t) = \mathcal{M}(m_+^2, m_-^2, t)$

CPV:

$\overline{a}_p \overline{\phi}_r \neq a_p \phi_r$: direct CPV
 (in decay)

$|q/p| \neq 1, \phi \neq 0$: indirect CPV
 (in mixing and
 interference)

$$\mathcal{A}(m_-^2, m_+^2) = \sum a_r e^{i\Phi_r} B(m_-^2, m_+^2) + a_{NR} e^{i\Phi_{NR}}$$

$$\overline{\mathcal{A}}(m_-^2, m_+^2) = \sum \overline{a}_r e^{i\overline{\Phi}_r} B(m_+^2, m_-^2) + \overline{a}_{NR} e^{i\overline{\Phi}_{NR}}$$

$$\overline{\mathcal{M}}(m_-^2, m_+^2, t) \neq \mathcal{M}(m_+^2, m_-^2, t)$$

Multi-body self conjugated states

Results

$$D^0 \rightarrow K_S \pi^+ \pi^-$$

Belle, PRL 99, 131803 (2007), 540fb⁻¹

$$|q/p| = 0.86 \pm 0.30 \pm 0.10$$

$$0.29 \pm 0.09$$

$$\varphi = (-0.24 \pm 0.28 \pm 0.09) \text{ rad}$$

$$0.30 \pm 0.09$$

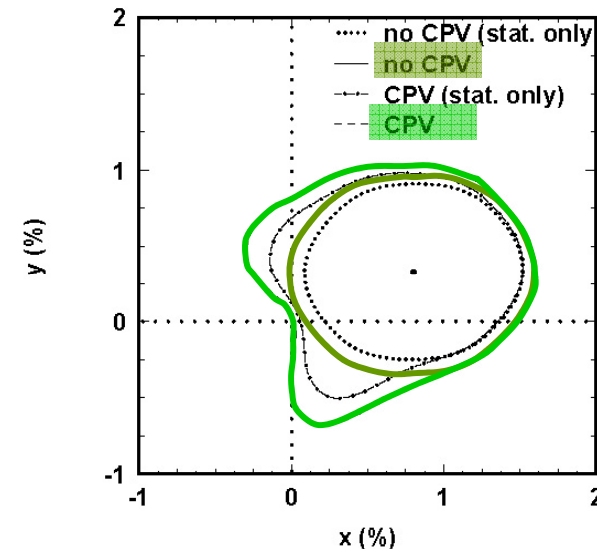
no evidence of direct CPV;
 x, y almost unchanged w.r.t. no CPV fit

BaBar, arXiv:1004.5053, 470 fb⁻¹

$$K_S \pi^+ \pi^- / K_S K^+ K^-$$

does not fit for CPV param.

95% C.L. contour



Other states

Principle

various decay modes
 of $D_{(s)}^{0(+)}$

t-integrated asymmetry;

example: $D_{(s)}^+ \rightarrow K_S h^+$
 $h=K, \pi$

charged mesons: CPV in decay only;

corrections for detector induced
 asymmetries and A_{FB} (n.b.: p. II/9);

example:

$$A_{rec}(D \rightarrow K_S \pi^+) - A_{rec}(D_s \rightarrow \phi \pi^+) \\
 \Rightarrow A_{CP}(D \rightarrow K_S \pi^+)$$

(technically much more involved
 due to $p_\pi, \theta_\pi, \theta_D$ dependence, see p. II/9)

$$A_{CP} = \frac{\Gamma(M \rightarrow f) - \Gamma(\bar{M} \rightarrow \bar{f})}{\Gamma(M \rightarrow f) + \Gamma(\bar{M} \rightarrow \bar{f})}$$

$$D^+ \rightarrow K_S \pi^+ \quad \text{CF}$$

$$D^+ \rightarrow K_S K^+ \quad \text{CS}$$

$$D_s^+ \rightarrow K_S K^+ \quad \text{CF}$$

$$D_s^+ \rightarrow K_S \pi^+ \quad \text{CS}$$

$$A_{rec}^{D \rightarrow K_S^0 \pi^+} = A_{CP}^{D \rightarrow K_S^0 \pi^+} + A_{FB}^D + A_\epsilon^{\pi^+}$$

$$A_{rec}^{D \rightarrow K_S^0 K^+} = A_{CP}^{D \rightarrow K_S^0 K^+} + A_{FB}^D + A_\epsilon^{K^+}$$

$$A_{rec}^{D^{*+} \rightarrow D^0 \pi_s^+} = A_{CP}^{D^0 \rightarrow K_S^0 P^0} + A_{FB}^{D^{*+}} + A_\epsilon^{\pi_s^+}$$

$$A_{rec}^{D_s^+ \rightarrow \phi \pi^+} = A_{FB}^{D_s^+} + A_\epsilon^{\pi^+}$$

$$A_{rec}^{untagged D^0 \rightarrow K^- \pi^+} = A_{FB}^{D^0} + A_\epsilon^{K^-} + A_\epsilon^{\pi^+}$$

$$A_{rec}^{tagged D^0 \rightarrow K^- \pi^+} = A_{FB}^{D^{*+}} + A_\epsilon^{K^-} + A_\epsilon^{\pi^+} + A_\epsilon^{\pi_s^+}$$

Other states

Principle

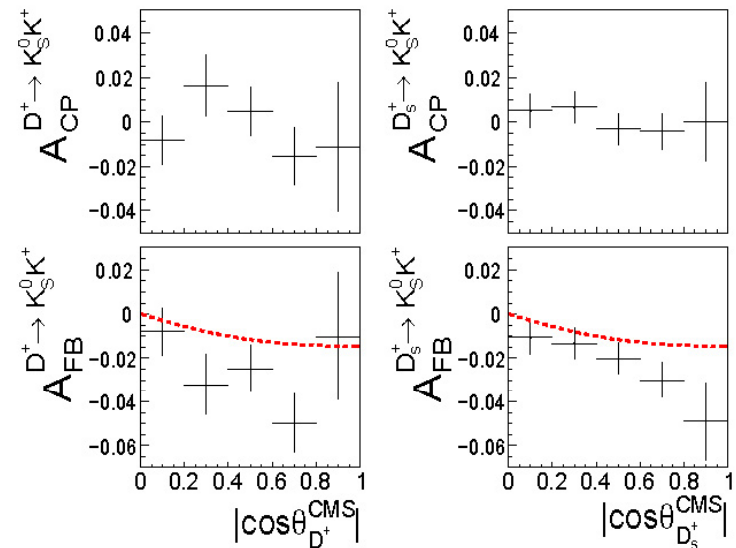
K_S in final state not a CP eigenstate itself;
 in weak D decays K^0, \bar{K}^0 produced;
 CPV in K^0 system \Rightarrow
 even in absence of CPV in D system
 some asymmetry expected

$$A_K = \frac{|\langle \pi\pi | K^0 \rangle|^2 - |\langle \pi\pi | \bar{K}^0 \rangle|^2}{|\langle \pi\pi | K^0 \rangle|^2 + |\langle \pi\pi | \bar{K}^0 \rangle|^2} \approx \frac{1 - |(p/q)_K|^2}{1 + |(p/q)_K|^2} = \frac{2 \operatorname{Re}(\varepsilon)}{1 + |\varepsilon|^2} = 0.332\%$$

Results Belle, PRL 104, 181602 (2010), 673fb⁻¹

	A_K
$A_{CP}^{D^+ \rightarrow K_S^0 \pi^+}$	$-0.71 \pm 0.19 \pm 0.20 \quad -0.332^\dagger$
$A_{CP}^{D_s^+ \rightarrow K_S^0 \pi^+}$	$+5.45 \pm 2.50 \pm 0.33 \quad +0.332$
$A_{CP}^{D^+ \rightarrow K_S^0 K^+}$	$-0.16 \pm 0.58 \pm 0.25 \quad -0.332$
$A_{CP}^{D_s^+ \rightarrow K_S^0 K^+}$	$+0.12 \pm 0.36 \pm 0.22 \quad -0.332^\dagger$

2.6 σ from 0, consistent with A_K



no CPV at $\geq 3 \cdot 10^{-3}$

Averages Results

same fit as for the
 mixing parameters,
 + CPV parameters

HFAG,
<http://www.slac.stanford.edu/xorg/hfag/>

$x \quad (0.59 \pm 0.20)\%$

$y \quad (0.80 \pm 0.13)\%$

$\delta \quad (27.6^{+11.2}_{-12.2})^\circ$

$\delta_{K\pi\pi} \quad (23.2^{+22.3}_{-23.3})^\circ$

$R_D \quad (0.332 \pm 0.008)\%$

$A_D \quad (-2.0 \pm 2.4)\%$

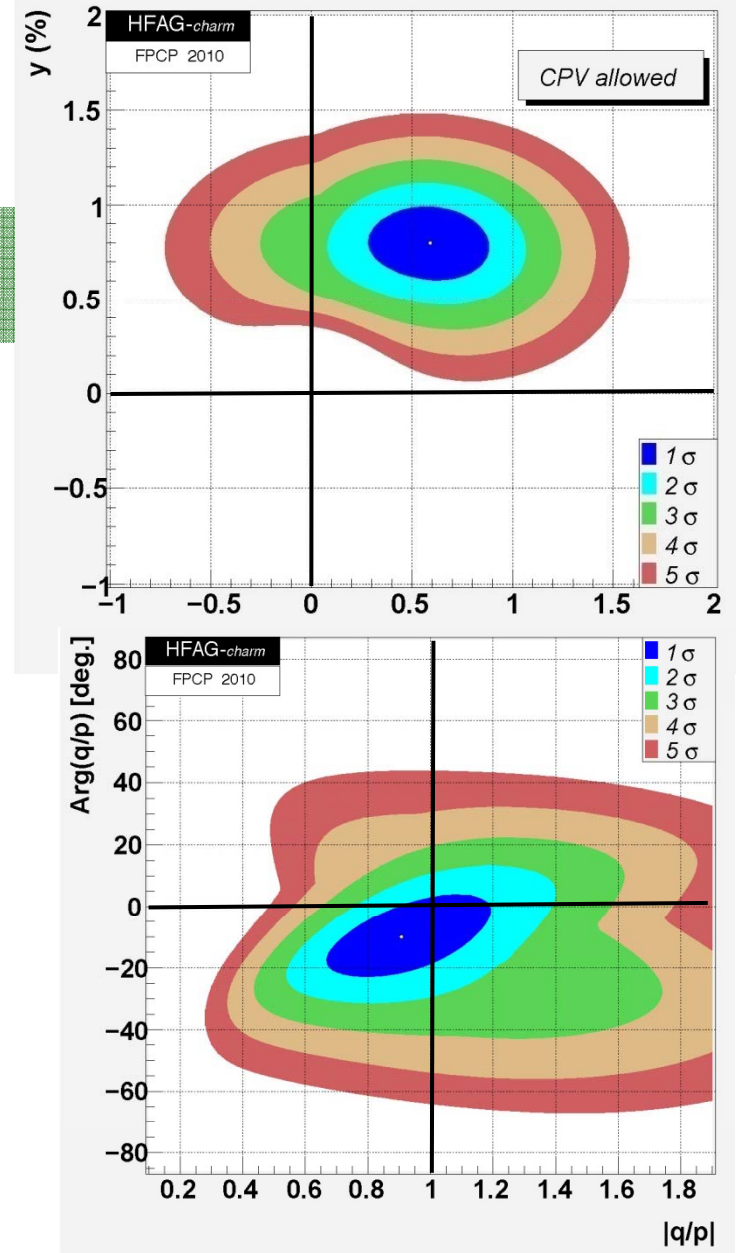
CPV {

$|q/p| \quad 0.91^{+0.19}_{-0.16}$

$\phi \quad (-10.0^{+9.3}_{-8.7})^\circ$

$\chi^2/n.d.f. =$
 $31.9/(30-8) =$
 $31.9/22$

(largest cont.
 from $K^+ \pi^- \pi^0$
 and Cleo-c)



Constraints from mixing (examples from

E. Golowich et al., PRD76, 095009 (2007)

4th generation of fermions

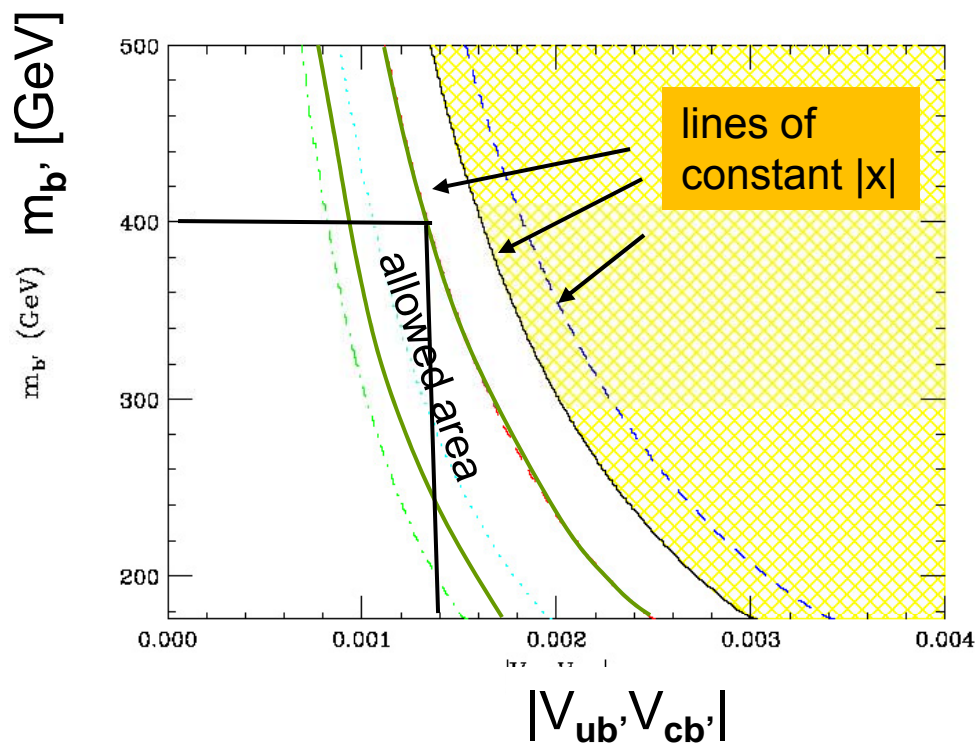
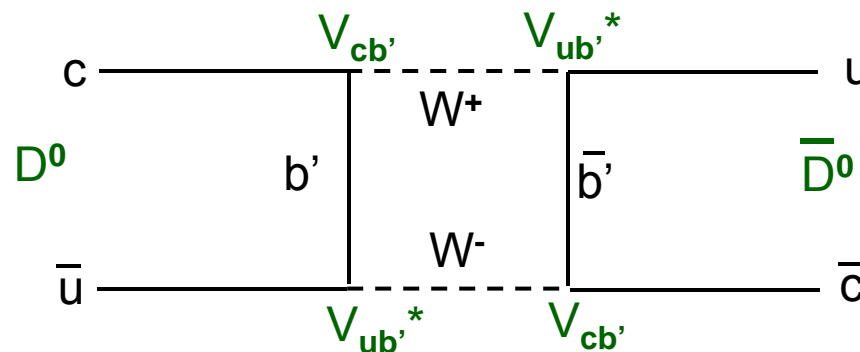
b' beside d,s,b exchanged
 in loop;

$$|V_{ub'}V_{cb'}| < 1.4 \cdot 10^{-3}$$

for $m_{b'} > 400$ GeV

more severe constraints
 than from CKM unitarity

complementarity of
 such constraints w.r.t.
 down-like FCNC obvious

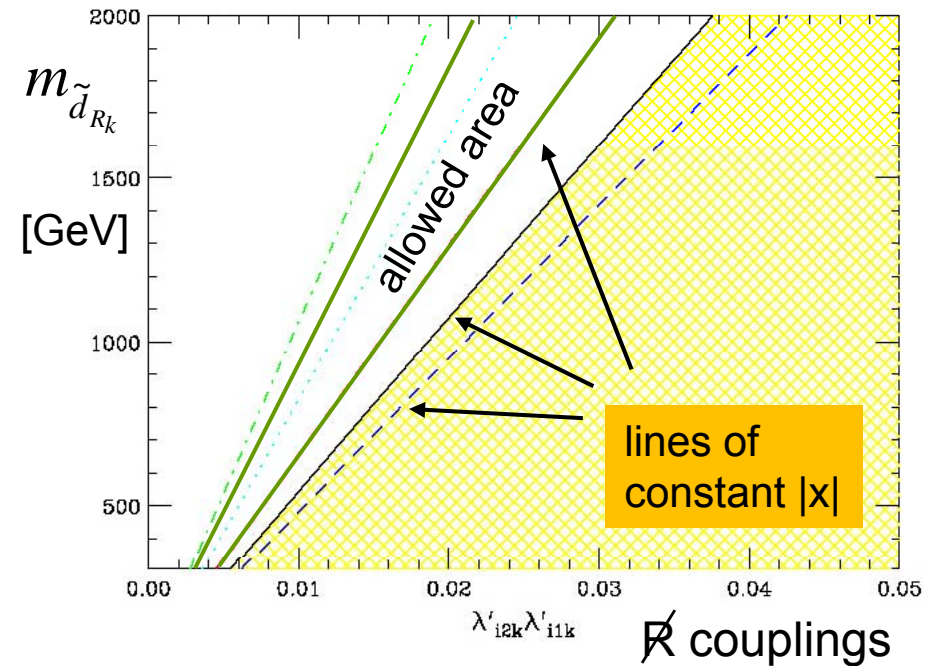
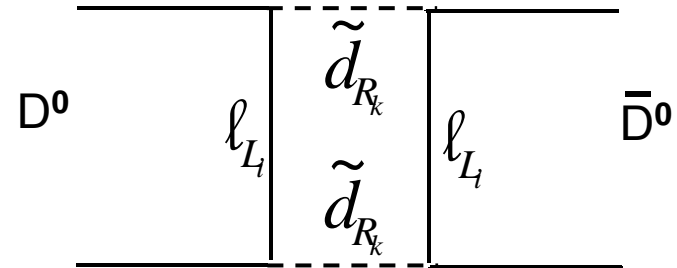


Constraints from mixing

E. Golowich et al., PRD76, 095009 (2007)

R-parity violating SUSY

squark-lepton (or vice versa)
 exchange in loop;



Constraints from CPV

SCS decays (probing penguins);
 small in SM, small DCPV;
 SUSY: squark-gluino loops;
 could lead to CPV

Y. Grossman et al., PRD75, 036008 (2007)

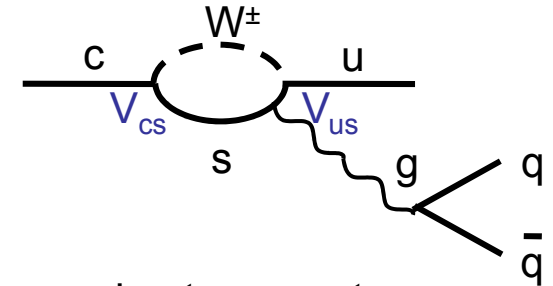
$$\underbrace{A_\Gamma}_{t\text{-dep.}} + \underbrace{A_{CP}}_{t\text{-integr.}} = A_D / 2 \quad \text{see eqs. on p. II/6, 8}$$

$$A_D = 4r_f \sin \varphi_f \sin \delta_f$$

$$r_f = \left| A_f^{NP} / A_f^{SM} \right| \quad A_D^{KK} = (-0.08 \pm 0.68)\%$$

contrib. to r_f of $\mathcal{O}(1\%)$ allowed
 using current constraints from mixing;

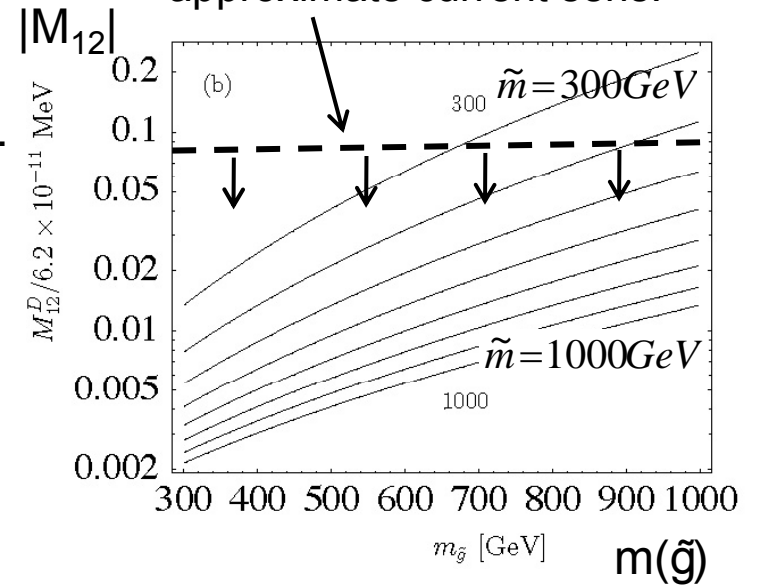
if $\phi_f, \delta_f \sim \mathcal{O}(1)$ (i.e. if no similar CPV suppression in NP like in SM) A_D constraint excludes all shown values



for M_{12} relation to observables
 see p. II/27, 28

approximate current sens.

$|M_{12}|$ vs. $m(\tilde{g})$ for $r_f=1\%$



$$r_f = \frac{A_D}{4 \sin \varphi_f \sin \delta_f} \sim \frac{A_D}{4}$$

Near future facilities

Charm-factories

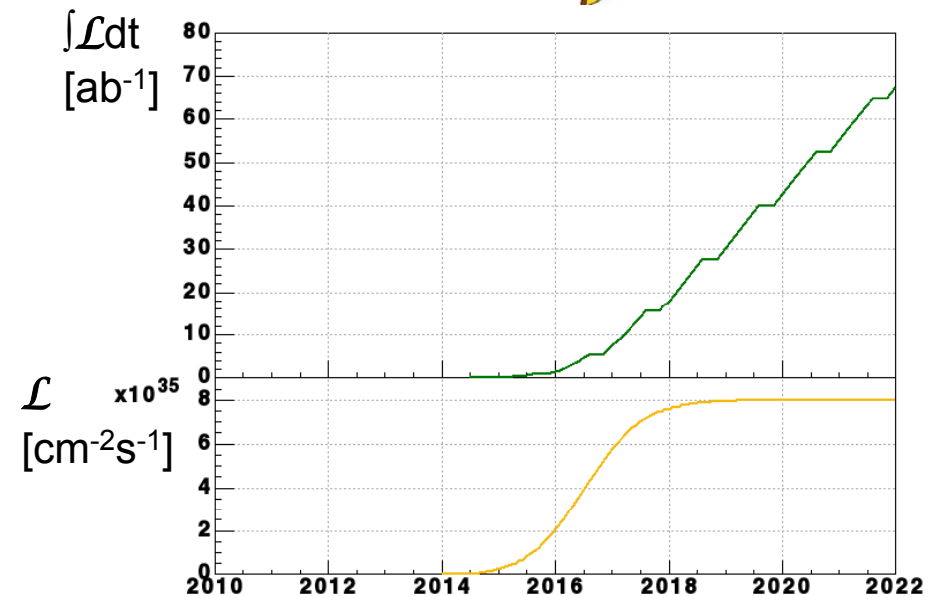
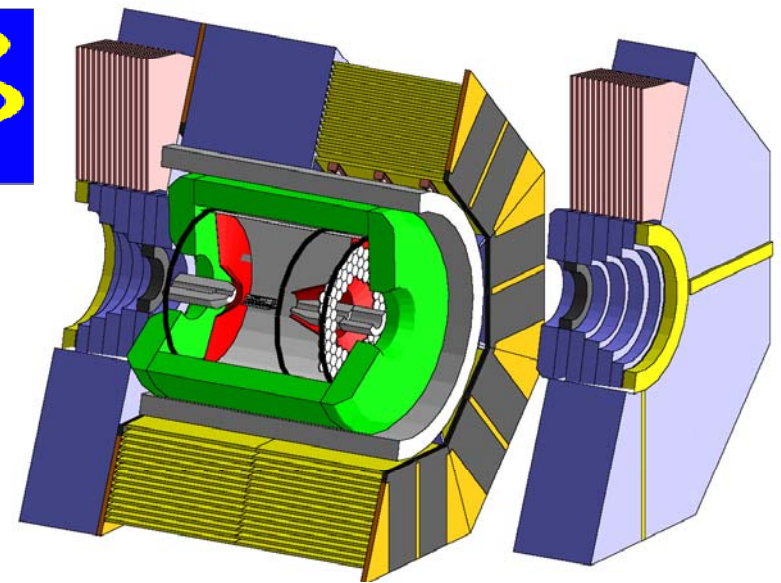
main results from BES-III
expected in the near future;

LHCb

great hopes for nice results,
although 1 fb^{-1} (by end of 2011)
may not be enough

Super B-factories

- Super KEKB, 5 ab^{-1} in 2016;
- SuperB, Frascati



Illustrative expected sensitivities

Mixing parameters

Charm-factories

BES III, arXiv:0809.1869

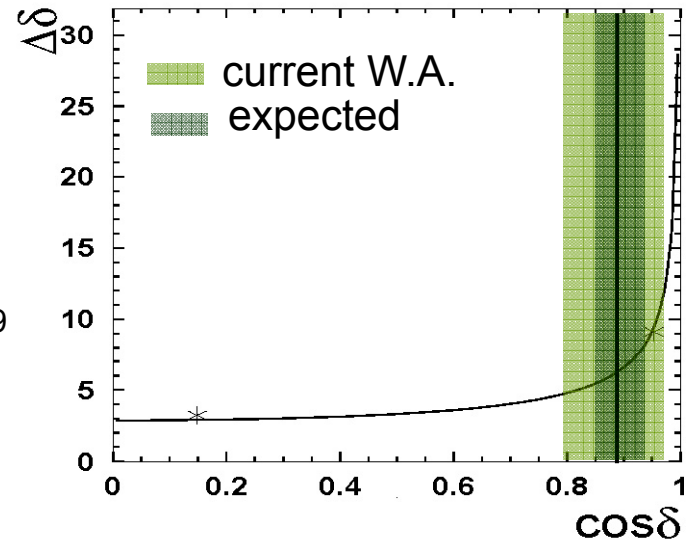
$20 \text{ fb}^{-1}: \sigma(R_M) \sim 1 \cdot 10^{-4}$

(from $\psi(3770) \rightarrow K^-\pi^+, K^-\pi^+$) n.b.: p. I/39

(n.b.: $R_M \sim 1 \cdot 10^{-4}$);

$\sigma(y) \sim 0.3\%$

$\sigma(\cos \delta) \sim 0.04$



LHCb

G. Wilkinson et al., public note LHCb-2007-49

$10 \text{ fb}^{-1}: \sigma(x) \sim 0.25\%, \sigma(y) \sim 0.05\%$

(estimated from precision on x'^2, y' and y_{CP})

Prediction is very difficult, especially of the future.

N. Bohr (1885 - 1962)

Super B-factories

A.G. Akeroyd et al., arXiv:1002.5012

Super-KEKB: $5 \text{ ab}^{-1}: \sigma(x), \sigma(y) \sim 0.1\%$

(combined from $K\pi, KK, K_S\pi\pi$)

Illustrative expected sensitivities

CPV parameters

$$r \equiv \frac{\Gamma(S_+ S_+)}{\Gamma(S_+ X)} = 2R_M Br(S_+) \sin^2 \varphi$$

Charm-factories

several possibilities;

decays to same sign CP states;

not sensitive due to R_M

suppression;

C=+1 initial state ($D^0 D^0 \gamma$);

20 fb⁻¹: $\sigma(A_\Gamma) \sim 0.6\%$ (stat. only)

BES III, arXiv:0809.1869

$$A_{CP}^{C=+1} = \frac{\Gamma(S_+ e^-) - \Gamma(S_+ e^+)}{\Gamma(S_+ e^-) + \Gamma(S_+ e^+)} \approx$$

$$\approx y \frac{A_M}{2} \cos \varphi - x \sin \varphi = A_\Gamma$$

$$n.b.: A_{CP}^{C=-1} \approx R_M A_M$$

(neglecting direct CPV,
 i.e. $A_D=0$; see p. 11/29
 for other asymmetries)

Super B-factories

Super-KEKB:

5 ab⁻¹: $\sigma(\phi) \sim 5^\circ$

$\sigma(A_\Gamma) \sim 0.1\%-0.2\%$

A.G. Akeroyd et al., arXiv:1002.5012

Illustrative expected sensitivities

CPV parameters

Belle II, 50 ab⁻¹

$$x = (0.832 \pm 0.095)\%$$

$$y = (0.813 \pm 0.064)\%$$

$$\delta_{K\pi} = 24.6^\circ \pm 4.9^\circ$$

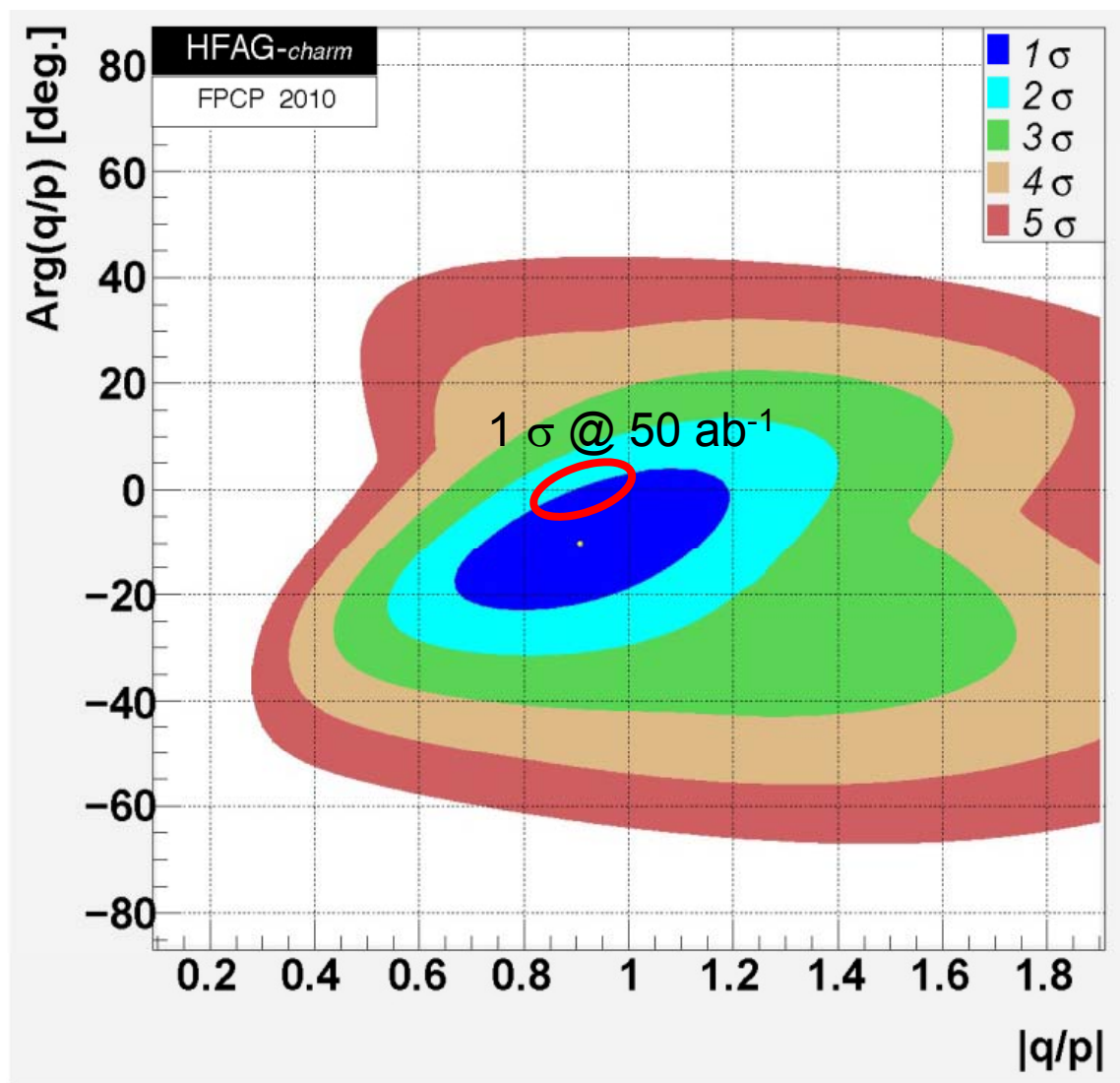
$$R_D = (0.336 \pm 0.003)\%$$

$$\frac{|q|}{|p|} = 0.894 \pm 0.054$$

$$\varphi = -0.004 \pm 0.049 \text{ rad}$$

$$A_D = (-0.1 \pm 0.8)\%$$

only KK/ππ, Kπ and K_sππ
 projected sensitivities included



- entering precision era in D^0 mixing and CPV (mixing only estab. in 2007)
- provide unique constraints/searches of NP in u-like FCNC

Today:

- B-factories (and Tevatron) still to say the final word

Tomorrow:

- Charm-factories, LHCb and Super-B factories
- will be able to search for NP effects (in CPV) in whole range down to SM predictions

D mixing rate in exclusive approach

J.F. Donoghue et al., PRD33, 179 (1986)

$$(M - i\frac{\Gamma}{2})_{ij} = \frac{\langle D_i | H_{eff} | D_j \rangle}{2M_D} = M_D \delta_{ij} + \frac{1}{2M_D} \langle \bar{D}^0 | H_w^{\Delta C=-2} | D^0 \rangle +$$

dominant contrib. to mixing;

take as example *PP* final states:

$$+ \frac{1}{2M_D} \sum_n \frac{\langle \bar{D}^0 | H_w^{\Delta C=-1} | n \rangle \langle n | H_w^{\Delta C=-1} | D^0 \rangle}{M_D - E_n + i\epsilon}$$

state	$ \langle \bar{D}^0 H_w^{\Delta C=-1} n \rangle ^2 \propto$	$\langle \bar{D}^0 H_w^{\Delta C=-1} n \rangle \langle n H_w^{\Delta C=-1} D^0 \rangle \propto$	measured Br	contrib. to mixing
$K^-\pi^+$	1	$-\lambda^2$	r_1	$-\sqrt{(r_1 r_4)} \lambda^2$
K^-K^+	λ^2	λ^2	$r_2 \lambda^2$	$r_2 \lambda^2$
$\pi^-\pi^+$	λ^2	λ^2	$r_3 \lambda^2$	$r_3 \lambda^2$
$K^+\pi^-$	λ^4	$-\lambda^2$	$r_4 \lambda^4$	$-\sqrt{(r_1 r_4)} \lambda^2$

minus sign due to relative sign of V_{us}/V_{cd} (see p. 1/32);
 summing over this states $\Rightarrow 0$ (GIM mechanism);

$$\Sigma = \lambda^2 (r_2 + r_3 - 2\sqrt{r_1 r_4}) \neq 0$$

however, in D^0 decays SU(3) is broken;
 in other words, measured Br's with $r_i \neq 1$

Mixing parameters

B mesons:

calculate M_{12} , Γ_{12} from
 box diagram; from that
 calculate Δm , $\Delta\Gamma$

$$M_{12} = -\frac{G_F^2 m_W^2 \eta_B m_{Bq} B_{Bq} f_{Bq}^2}{12\pi^2} S_0(m_t^2 / m_W^2) (V_{tq}^* V_{tb})^2$$

$$\Gamma_{12} = \frac{G_F^2 m_b^2 \eta_B' m_{Bq} B_{Bq} f_{Bq}^2}{8\pi} (V_{tq}^* V_{tb})^2$$

q: d (B_d) or s (B_s)

B_{Bq} : bag parameter, $\langle B_q^0 | b\gamma^\mu(1-\gamma^5)q | B_q^0 \rangle$

f_{Bq} : decay constant

$\eta_B^{(\prime)}$: QCD corr. $\mathcal{O}(1)$

$S_0(x_t)$: known kinematic function

$$\varphi_{12} = \arg \frac{M_{12}}{\Gamma_{12}} = \pi + \mathcal{O}(m_c^2 / m_b^2)$$

$$\left| \frac{\Gamma_{12}}{M_{12}} \right| \approx \frac{3\pi}{2} \frac{m_b^2}{m_W^2} \frac{1}{S_0(m_t^2 / m_W^2)} \sim \mathcal{O}(m_b^2 / m_t^2)$$

$$\left| \frac{q}{p} \right|^2 = 1 + \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin \varphi_{12} + \mathcal{O}(|\Gamma_{12} / M_{12}|^2)$$

Mixing parameters

B mesons:

$$|\Gamma_{12}| \ll |M_{12}|$$

measured $x_d = 0.776 \pm 0.008$

$$x_s = 25.5 \pm 0.6$$

from $M_{12}/\Gamma_{12} \Rightarrow y_d < 1\%$

$$y_s \sim 10\%$$

D mesons:

$$\Delta m = 2 |M_{12}| (1 + \dots)$$

$$\Delta \Gamma = -2 |\Gamma_{12}| (1 + \dots)$$

$$\dots \rightarrow \mathcal{O}(|\Gamma_{12}| / |M_{12}|)$$

$$|M_{12}^D| = \frac{\bar{\Gamma} x}{2} \sqrt{1 + (A_M x / 2y)^2}$$

G. Raz, PRD66, 057502 (2002)

Asymmetries at charm-factories

untagged asymmetry

$$A_{CP} = \frac{\Gamma(D^0 \rightarrow K^- \pi^+) + \Gamma(\bar{D}^0 \rightarrow K^- \pi^+) - \Gamma(D^0 \rightarrow K^+ \pi^-) - \Gamma(\bar{D}^0 \rightarrow K^+ \pi^-)}{\Gamma(D^0 \rightarrow K^- \pi^+) + \Gamma(\bar{D}^0 \rightarrow K^- \pi^+) + \Gamma(D^0 \rightarrow K^+ \pi^-) + \Gamma(\bar{D}^0 \rightarrow K^+ \pi^-)}$$

$$A_{CP} \approx 2\sqrt{R_D} \sin \delta \left[y \sin \varphi + \frac{A_M}{2} x \cos \varphi \right] \approx 2\sqrt{R_D} \sin \delta y \sin \varphi$$

semileptonic asymmetry

$$A_{CP} = \frac{\Gamma(\bar{D}^0 D^0 \rightarrow e^+ e^+) - \Gamma(\bar{D}^0 D^0 \rightarrow e^- e^-)}{\Gamma(\bar{D}^0 D^0 \rightarrow e^+ e^+) + \Gamma(\bar{D}^0 D^0 \rightarrow e^- e^-)}$$

$$A_{CP} \approx -2A_M$$

direct CPV

$$A_{CP} = \frac{\Gamma(\bar{D}^0 D^0 \rightarrow S_+ e^-) - \Gamma(\bar{D}^0 D^0 \rightarrow S_+ e^+)}{\Gamma(\bar{D}^0 D^0 \rightarrow S_+ e^-) + \Gamma(\bar{D}^0 D^0 \rightarrow S_+ e^+)} \approx \frac{A_D}{2}$$

(n.b.: A_D is decay mode specific, in this case represents direct CPV in S_+ decay mode)