

Genuine CP-Odd Observables and ZZH Coupling

Tao Han *

Univ. of Wisconsin–Madison

(LCWS 2010, March 28, 2010)

*Collaborators: Jing Jiang; Neil Christensen, Y.-C. Li

Genuine CP-Odd Observables and ZZH Coupling

Tao Han *

Univ. of Wisconsin–Madison

(LCWS 2010, March 28, 2010)

General ZZH Vertex

CP-Odd Observables at ILC

CP-Odd Observables at LHC

*Collaborators: Jing Jiang; Neil Christensen, Y.-C. Li

General ZZH Vertex

The most important coupling for EWSB:

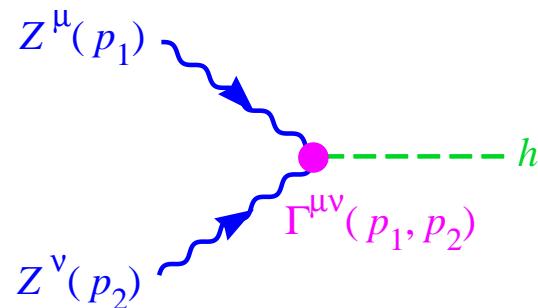
The Higgs boson to gauge bosons, in particular ZZh ,

General ZZH Vertex

The most important coupling for EWSB:

The Higgs boson to gauge bosons, in particular ZZh ,

The most general vertex function for ZZh



$$\Gamma^{\mu\nu}(p_1, p_2) = i \frac{2}{v} h [a M_Z^2 g^{\mu\nu} + b (p_1^\mu p_2^\nu - p_1 \cdot p_2 g^{\mu\nu}) + \tilde{b} \epsilon^{\mu\nu\rho\sigma} p_{1\rho} p_{2\sigma}]$$

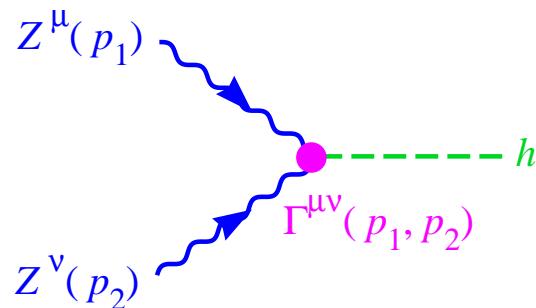
$a = 1$, $b = \tilde{b} = 0$ for SM; a , b terms: CP-even; \tilde{b} term: CP-odd.

General ZZH Vertex

The most important coupling for EWSB:

The Higgs boson to gauge bosons, in particular ZZh ,

The most general vertex function for ZZh



$$\Gamma^{\mu\nu}(p_1, p_2) = i \frac{2}{v} h [a M_Z^2 g^{\mu\nu} + b (p_1^\mu p_2^\nu - p_1 \cdot p_2 g^{\mu\nu}) + \tilde{b} \epsilon^{\mu\nu\rho\sigma} p_{1\rho} p_{2\sigma}]$$

$a = 1$, $b = \tilde{b} = 0$ for SM; a , b terms: CP-even; \tilde{b} term: CP-odd.

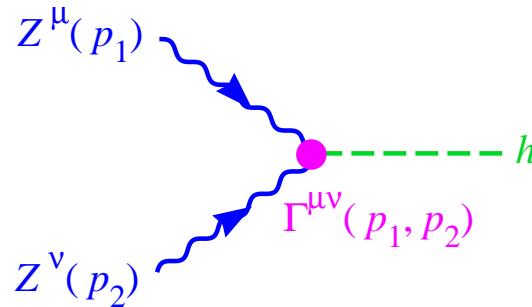
In general, a , b , \tilde{b} complex “form factors”, from loops.

General ZZH Vertex

The most important coupling for EWSB:

The Higgs boson to gauge bosons, in particular ZZh ,

The most general vertex function for ZZh



$$\Gamma^{\mu\nu}(p_1, p_2) = i \frac{2}{v} h [a M_Z^2 g^{\mu\nu} + b (p_1^\mu p_2^\nu - p_1 \cdot p_2 g^{\mu\nu}) + \tilde{b} \epsilon^{\mu\nu\rho\sigma} p_{1\rho} p_{2\sigma}]$$

$a = 1$, $b = \tilde{b} = 0$ for SM; a , b terms: CP-even; \tilde{b} term: CP-odd.

In general, a , b , \tilde{b} complex “form factors”, from loops.

In an “effective field theory”, operators $\frac{g^2}{\Lambda^2} HHW^{\mu\nu}W_{\mu\nu}$, $\frac{g^2}{\Lambda^2} HHW^{\mu\nu}\tilde{W}_{\mu\nu}$, natural size: a , b , $\tilde{b} \sim \mathcal{O}(\frac{1}{16\pi^2} \sim 1)$.

Consider direct production:

$$e^-(p_1) \ e^+(p_2) \rightarrow e^-(q_1) \ e^+(q_2) \ h(q_3)$$

via both Zh (Bjorken/Higgs-strahlung) and ZZ (Fusion).

Consider direct production:

$$e^-(p_1) e^+(p_2) \rightarrow e^-(q_1) e^+(q_2) h(q_3)$$

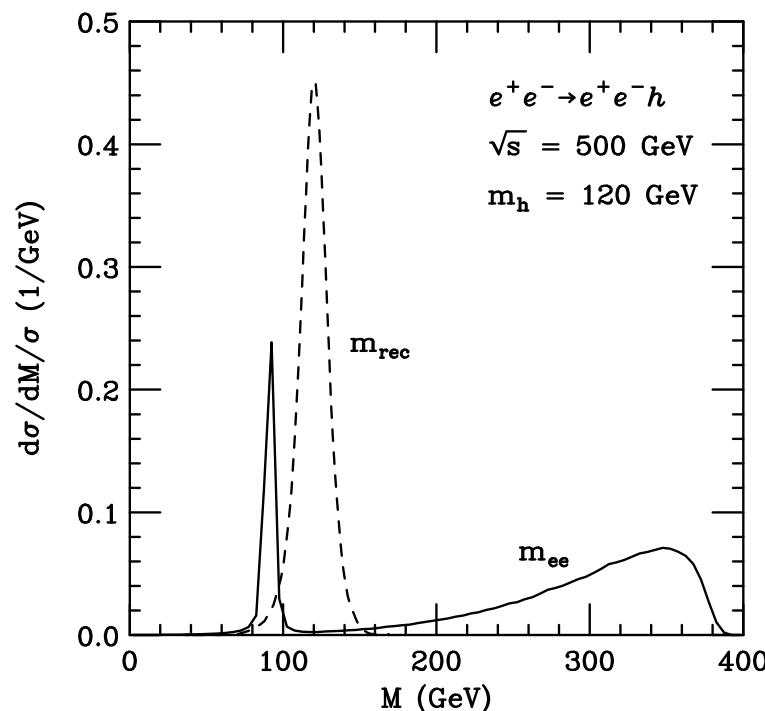
via both Zh (Bjorken/Higgs-strahlung) and ZZ (Fusion).

- Recoil mass, regardless the Higgs decay mode:

$$(p_1 + p_2 - q_1 - q_2)^2 = m_h^2$$

- Distinctive $m_{e^+e^-}$ spectrum:

Bjorken process and Fusion can be separated.



Cross section rates:

Zh Bjorken process:^{*} $\sigma \propto \frac{1}{s}$; ZZ fusion:[†] $\sigma \propto \log^2(\frac{s}{M_Z^2})$.

^{*}K. Hagiwara and M. Stong; V. Barger et al.; A. Skjold and P. Osland; ...

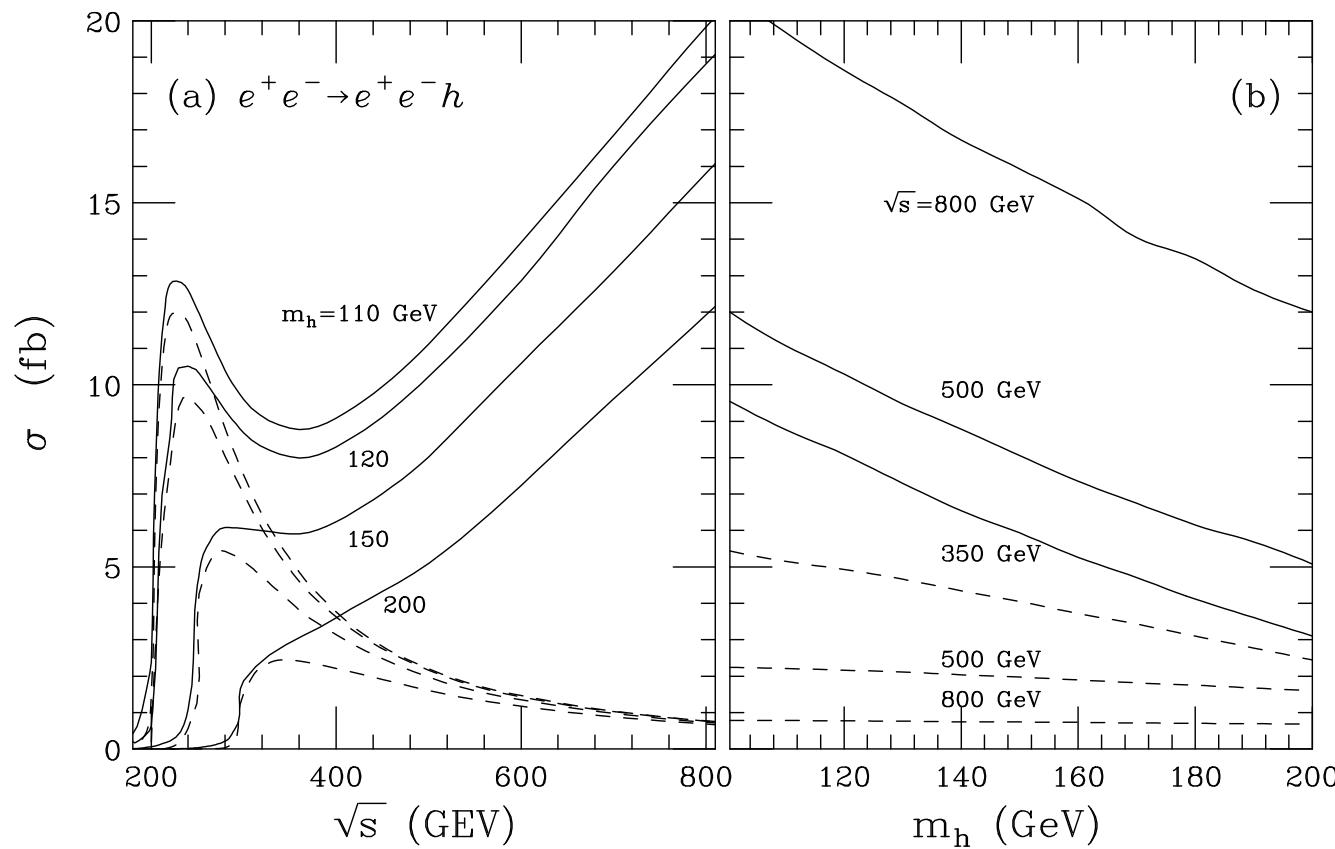
[†]TH and J. Jiang

Cross section rates:

$$Zh \text{ Bjorken process:}^* \sigma \propto \frac{1}{s}; \quad ZZ \text{ fusion:}^\dagger \sigma \propto \log^2\left(\frac{s}{M_Z^2}\right).$$

For instance: for $\sqrt{s} = 500 \text{ GeV}$, $m_h = 120 \text{ GeV}$:

$$\sigma(fusion) > 2\sigma(Bjorken, e^+e^- + \mu^+\mu^-) \approx 10 \text{ fb}^{-1}.$$



^{*}K. Hagiwara and M. Stong; V. Barger et al.; A. Skjold and P. Osland; ...

[†]TH and J. Jiang

CP-odd observables at ILC

$$e^-(p_1) \ e^+(p_2) \rightarrow e^-(q_1) \ e^+(q_2) \ h(q_3)$$

With longitudinally polarized beams, under CP:

$$\mathcal{M}_{--}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{++}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1), \quad (1)$$

$$\mathcal{M}_{-+}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{-+}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1). \quad (2)$$

CP-odd observables at ILC

$$e^-(p_1) \ e^+(p_2) \rightarrow e^-(q_1) \ e^+(q_2) \ h(q_3)$$

With longitudinally polarized beams, under CP:

$$\mathcal{M}_{--}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{++}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1), \quad (1)$$

$$\mathcal{M}_{-+}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{-+}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1). \quad (2)$$

Based on Eq. (1), construct a LR asymmetry:

$$\mathcal{A}_{hel} = \sigma_{--} - \sigma_{++}$$

Based on Eq. (2), construct angular FB asymmetries:

$$\mathcal{A}_{FB} = \sigma_{-+}^F - \sigma_{-+}^B.$$

CP-odd observables at ILC

$$e^-(p_1) \ e^+(p_2) \rightarrow e^-(q_1) \ e^+(q_2) \ h(q_3)$$

With longitudinally polarized beams, under CP:

$$\mathcal{M}_{--}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{++}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1), \quad (1)$$

$$\mathcal{M}_{-+}(\vec{p}_1, \vec{p}_2; \vec{q}_1, \vec{q}_2) \Rightarrow \mathcal{M}_{-+}(\vec{p}_1, \vec{p}_2; -\vec{q}_2, -\vec{q}_1). \quad (2)$$

Based on Eq. (1), construct a LR asymmetry:

$$\mathcal{A}_{hel} = \sigma_{--} - \sigma_{++}$$

Based on Eq. (2), construct angular FB asymmetries:

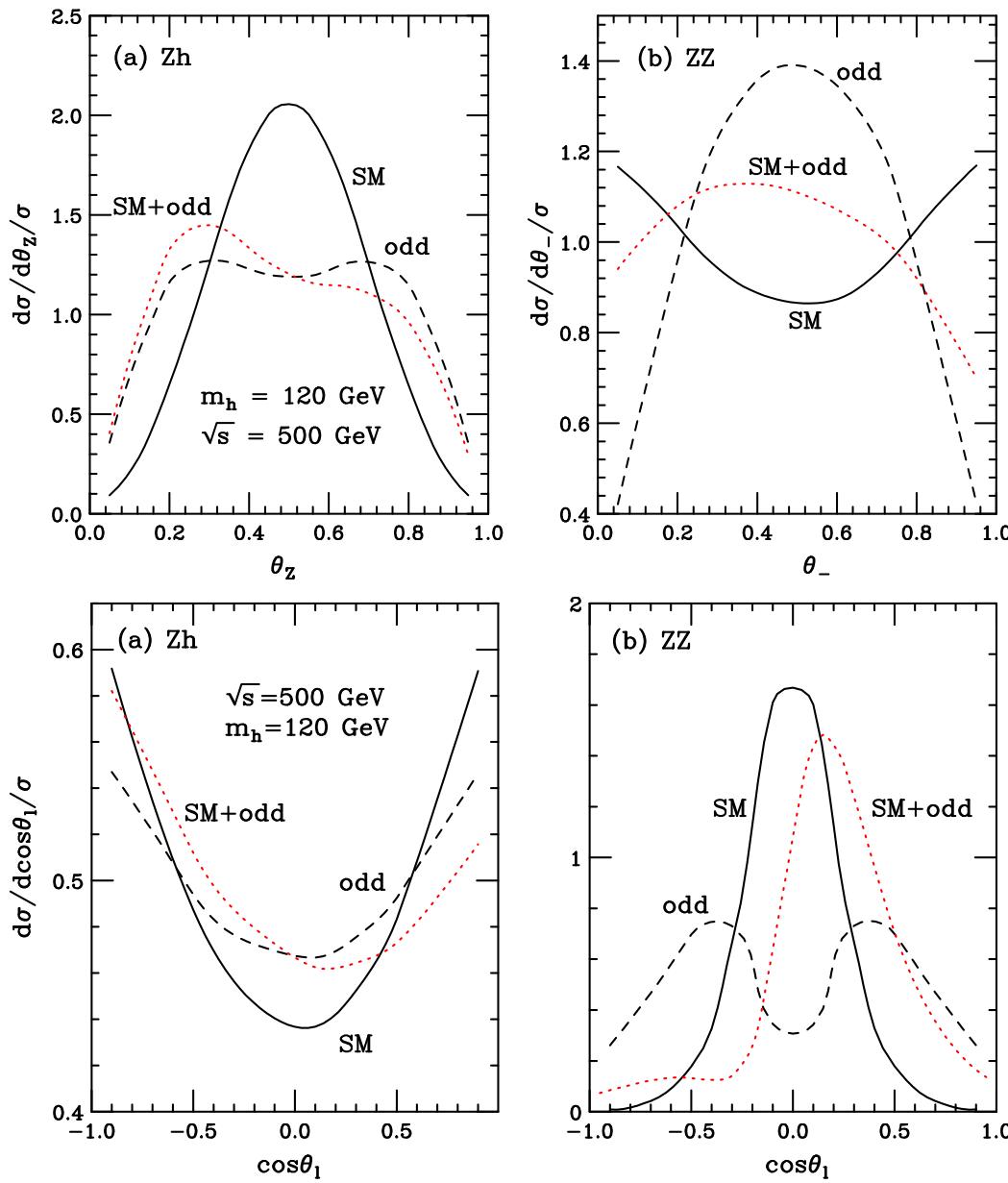
$$\mathcal{A}_{FB} = \sigma_{-+}^F - \sigma_{-+}^B.$$

Such as CP-odd observables:

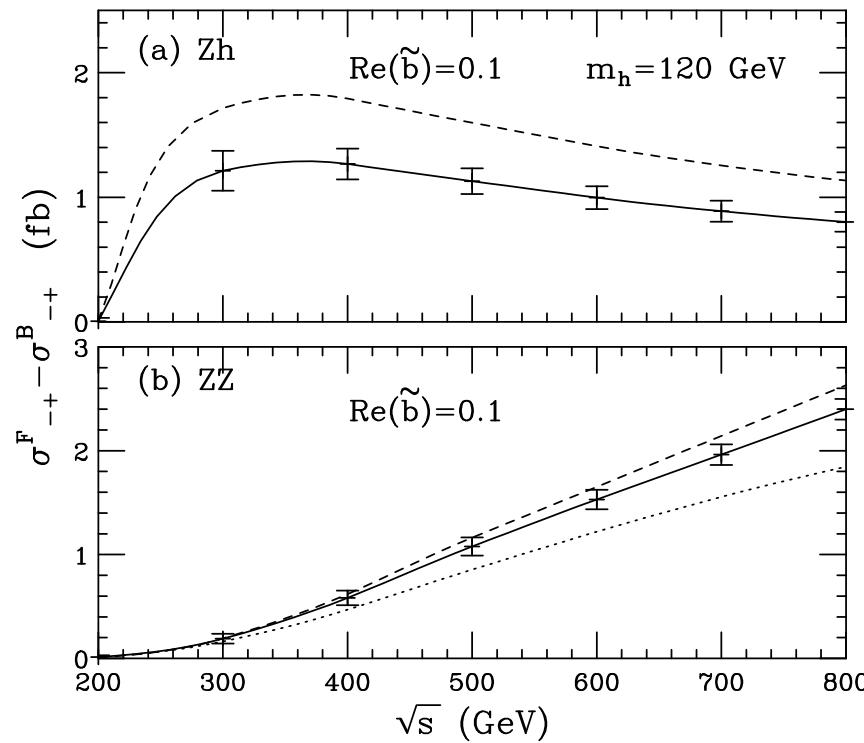
$$\cos \theta_Z \sim \hat{z} \cdot \vec{q}_+, \quad \cos \theta_\ell \sim \hat{z} \cdot (\vec{q}_1 \times \vec{q}_2), \quad \sin \theta_- \sim (\hat{z} \times \vec{q}_-) \cdot (\vec{q}_1 \times \vec{q}_2),$$

where $\vec{q}_\pm = \vec{q}_1 \pm \vec{q}_2$.

Angular distributions for the asymmetries: $\text{Im}(\tilde{b}) \sim 1$



Asymmetries versus c.m. energy:



	\sqrt{s} (GeV)	500	500	800	800
	\mathcal{L} (fb^{-1})	500	1000	500	1000
$Im(\tilde{b})$	$\mathcal{A}_{\theta_z}^{FB}(Zh)$ [-+]	0.0028	0.0022	0.0043	0.0032
	$\mathcal{A}_{\theta_z}^{FB}(Zh)$ [unpol.]	0.019	0.013	0.025	0.019
	$\mathcal{A}_{\theta_{\perp}}^{FB}(ZZ)$ [-+]	0.21	0.16	0.19	0.13
	$\mathcal{A}_{LR}(ZZ)$	0.071	0.045	0.065	0.041
$Re(\tilde{b})$	$\mathcal{A}_{\theta_t}^{FB}(Zh)$ [-+]	0.023	0.018	0.019	0.014
	$\mathcal{A}_{\theta_t}^{FB}(ZZ)$ [-+]	0.021	0.017	0.014	0.009
	$\mathcal{A}_{\theta_t}^{FB}(ZZ)$ [unpol.]	0.024	0.018	0.016	0.010

CP-odd observables at LHC

$pp \rightarrow \ell^- \ell^+ H X$
or $q\bar{q} \rightarrow ZH \rightarrow \ell^- \ell^+ H.$

CP-odd observables at LHC

$$pp \rightarrow \ell^- \ell^+ H X$$

or $q\bar{q} \rightarrow ZH \rightarrow \ell^- \ell^+ H.$

Problems:

- pp is NOT a CP eigen-state (comparing with $\bar{p}\bar{p}$?)
⇒ Most of the subprocesses are NOT from CP eigen-states, $W^\pm\dots$
- Even $q\bar{q}$, gg only CP eigen-state in c.m. frame
- Directions of q and \bar{q} randomly symmetrized.

CP-odd observables at LHC

$$pp \rightarrow \ell^- \ell^+ H X$$

or $q\bar{q} \rightarrow ZH \rightarrow \ell^- \ell^+ H.$

Problems:

- pp is NOT a CP eigen-state (comparing with $\bar{p}\bar{p}$?)
⇒ Most of the subprocesses are NOT from CP eigen-states, $W^\pm\dots$
- Even $q\bar{q}$, gg only CP eigen-state in c.m. frame
- Directions of q and \bar{q} randomly symmetrized.

The naive triplet product variables:

$$(\vec{p}_f \times \vec{p}_{\bar{f}}) \cdot \hat{z}$$

will NOT work, because the above • • •

Must find new variables suitable to LHC!

$$\mathcal{O}_1 \equiv p_T^+ - p_T^- \quad \text{or} \quad E_T^+ - E_T^-,$$
$$p_T = \sqrt{p_x^2 + p_y^2}, \quad E_T = \sqrt{p_T^2 + m_f^2}.$$

It is CP-odd but \hat{T} -even.

Must find new variables suitable to LHC!

$$\mathcal{O}_1 \equiv p_T^+ - p_T^- \quad \text{or} \quad E_T^+ - E_T^-,$$
$$p_T = \sqrt{p_x^2 + p_y^2}, \quad E_T = \sqrt{p_T^2 + m_f^2}.$$

It is CP-odd but \hat{T} -even.

$$\mathcal{O}_2 \equiv (\vec{p}_f \times \vec{p}_{\bar{f}}) \cdot \hat{z} \operatorname{sgn}((\vec{p}_f - \vec{p}_{\bar{f}}) \cdot \hat{z}),$$

It is CP-odd and \hat{T} -odd.

Must find new variables suitable to LHC!

$$\mathcal{O}_1 \equiv p_T^+ - p_T^- \quad \text{or} \quad E_T^+ - E_T^-,$$
$$p_T = \sqrt{p_x^2 + p_y^2}, \quad E_T = \sqrt{p_T^2 + m_f^2}.$$

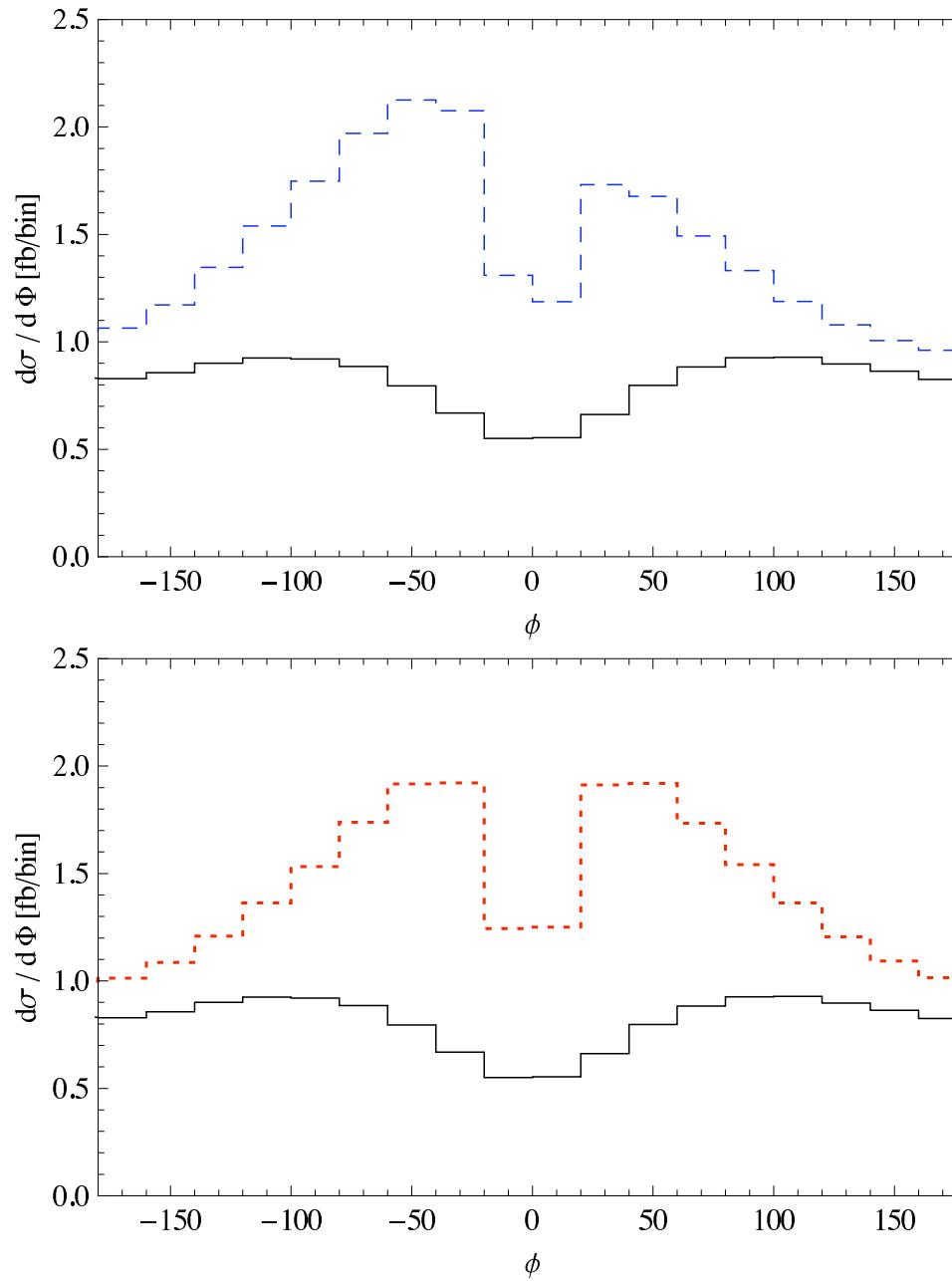
It is CP-odd but \hat{T} -even.

$$\mathcal{O}_2 \equiv (\vec{p}_f \times \vec{p}_{\bar{f}}) \cdot \hat{z} \operatorname{sgn}((\vec{p}_f - \vec{p}_{\bar{f}}) \cdot \hat{z}),$$

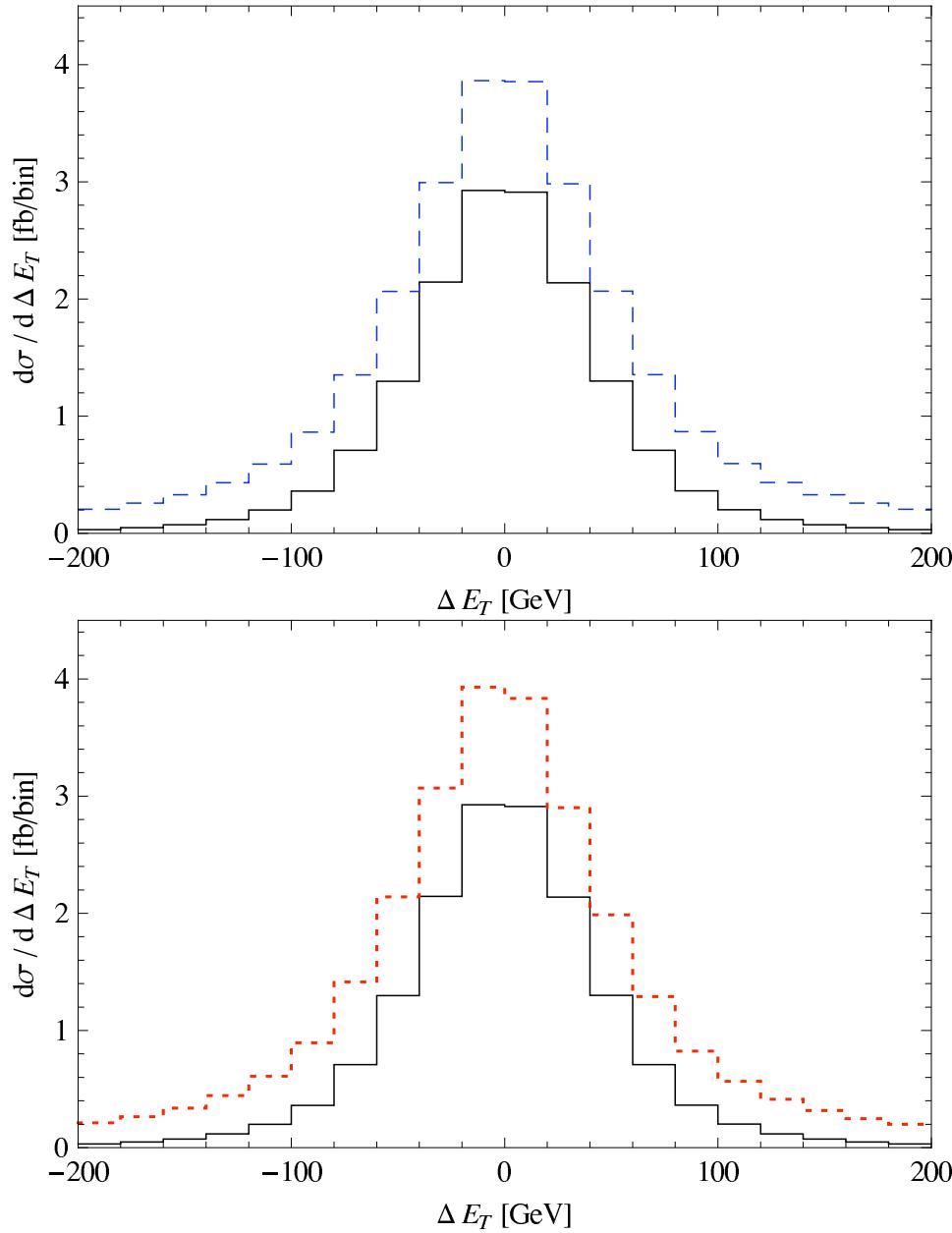
It is CP-odd and \hat{T} -odd.

Note that a CP-odd+ \hat{T} -even $\sim \sin \delta$ (a CP-conserving strong phase);
while a CP-odd+ \hat{T} -odd $\sim \cos \delta$.

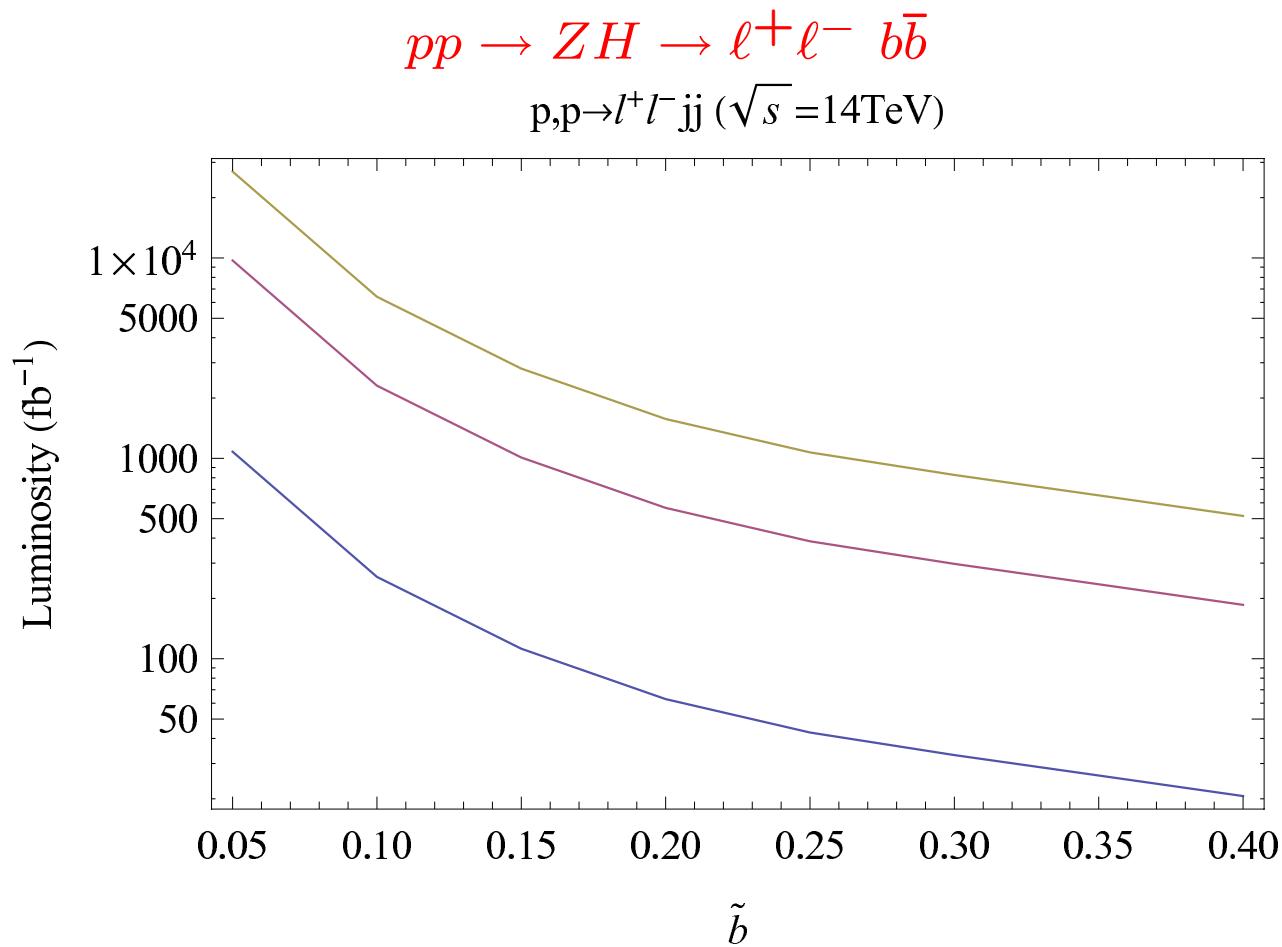
Angular asymmetries at the LHC: $\text{Re}(\tilde{b})$, $\text{Im}(\tilde{b})$



Energy asymmetries at the LHC: $\text{Re}(\tilde{b})$, $\text{Im}(\tilde{b})$



Integrated luminosity needed for 1, 3, 5 σ



Final remarks:

Final remarks:

- New sources of CP-violation is needed.
Higgs sector BSM is likely the place.

Final remarks:

- New sources of CP-violation is needed.
Higgs sector BSM is likely the place.
- Genuine CP-odd observables important.

Final remarks:

- New sources of CP-violation is needed.
Higgs sector BSM is likely the place.
- Genuine CP-odd observables important.
- ILC: Good opportunity to search for CPv,
reaching $\tilde{b} \sim 10^{-3}$.

Final remarks:

- New sources of CP-violation is needed.
Higgs sector BSM is likely the place.
- Genuine CP-odd observables important.
- ILC: Good opportunity to search for CPv,
reaching $\tilde{b} \sim 10^{-3}$.
- LHC: Very challenging, but possible (only recently),
reaching $\tilde{b} \sim 0.25$.