



The WIMP dark matter paradigm and beyond

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高能理论论坛 (High Energy Theory Forum)@IHEP

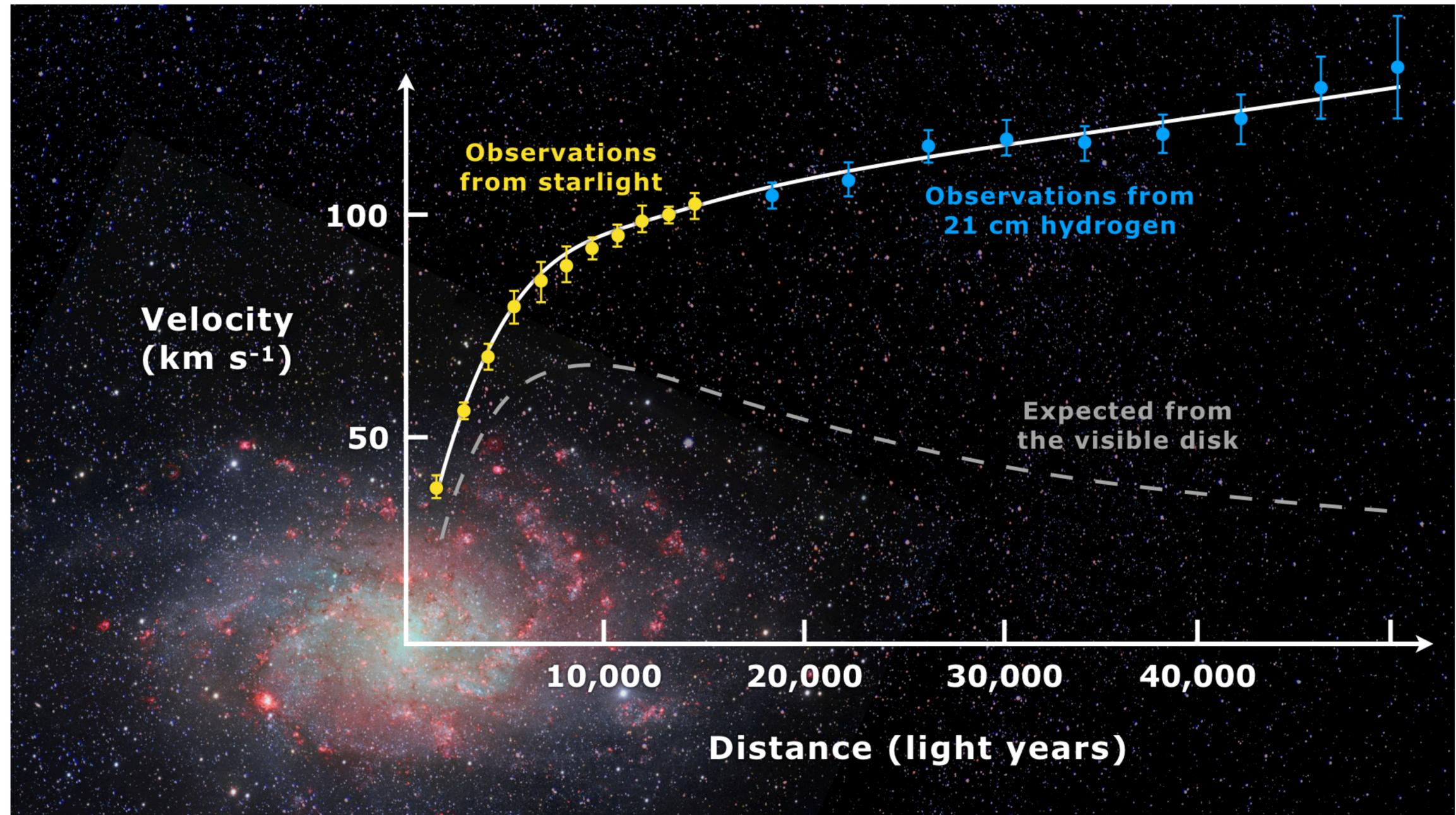
2021-10-19

Outlines

- The dark matter in astrophysics perspective
- The dark matter in particle physics perspective
 - The WIMP crisis from direct detection
 - The DM limits from indirect detection
 - A WIMP variant from cosmological evolution
- Summary

Observational evidence for DM

- Galaxy rotation curves
- Bullet cluster
- Gravitational Lensing
- Structure formation
- Cosmic Microwave Background



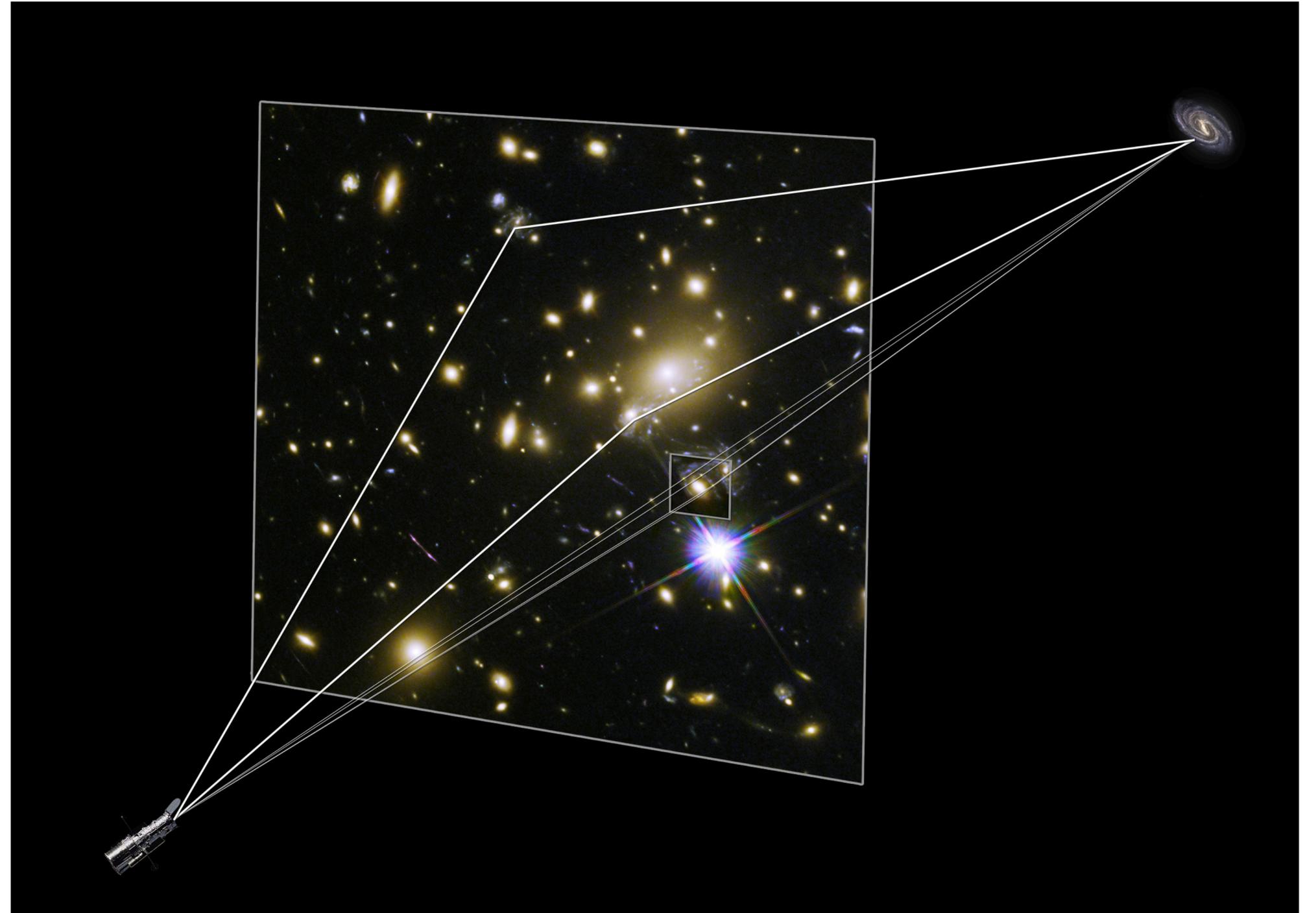
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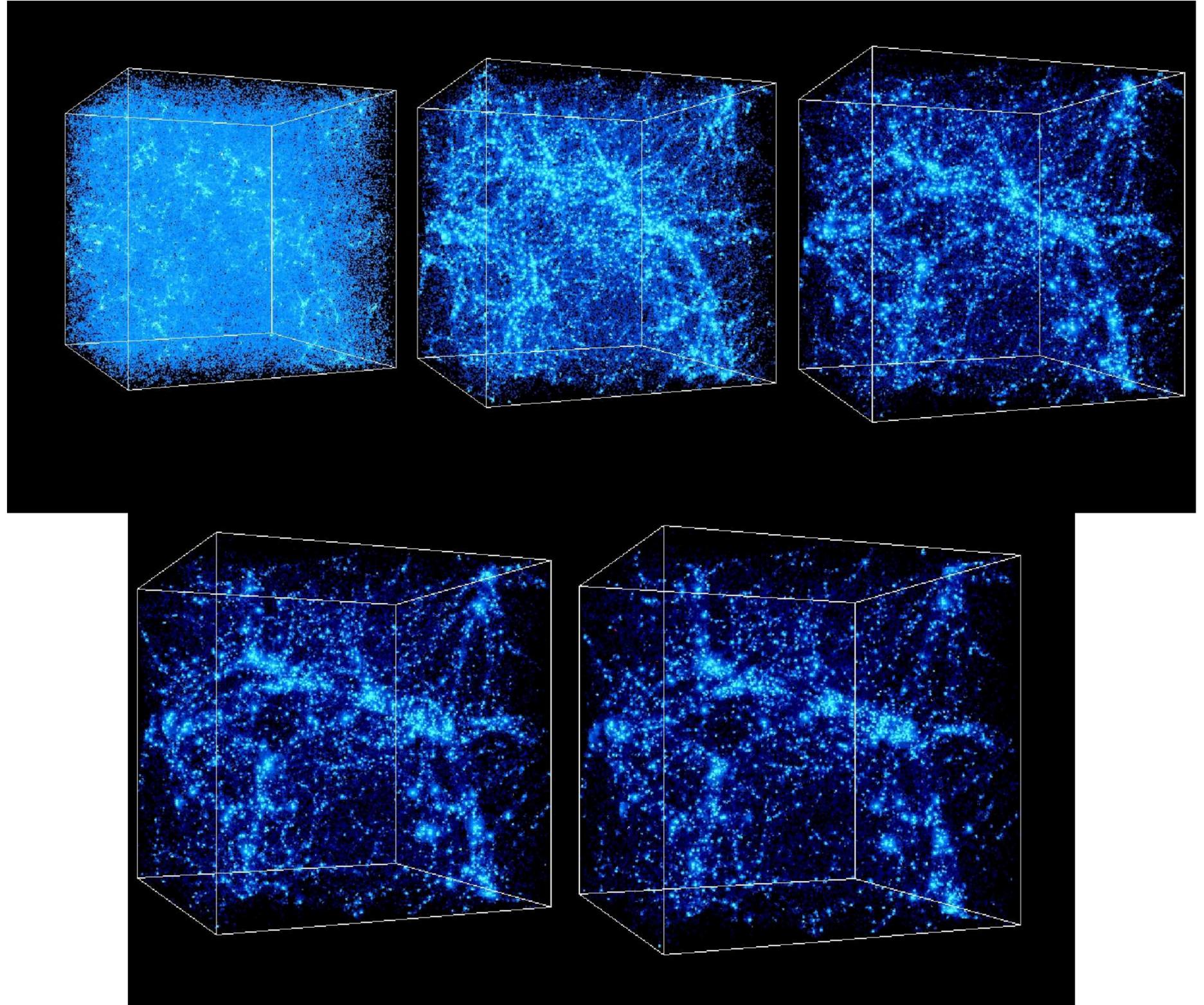
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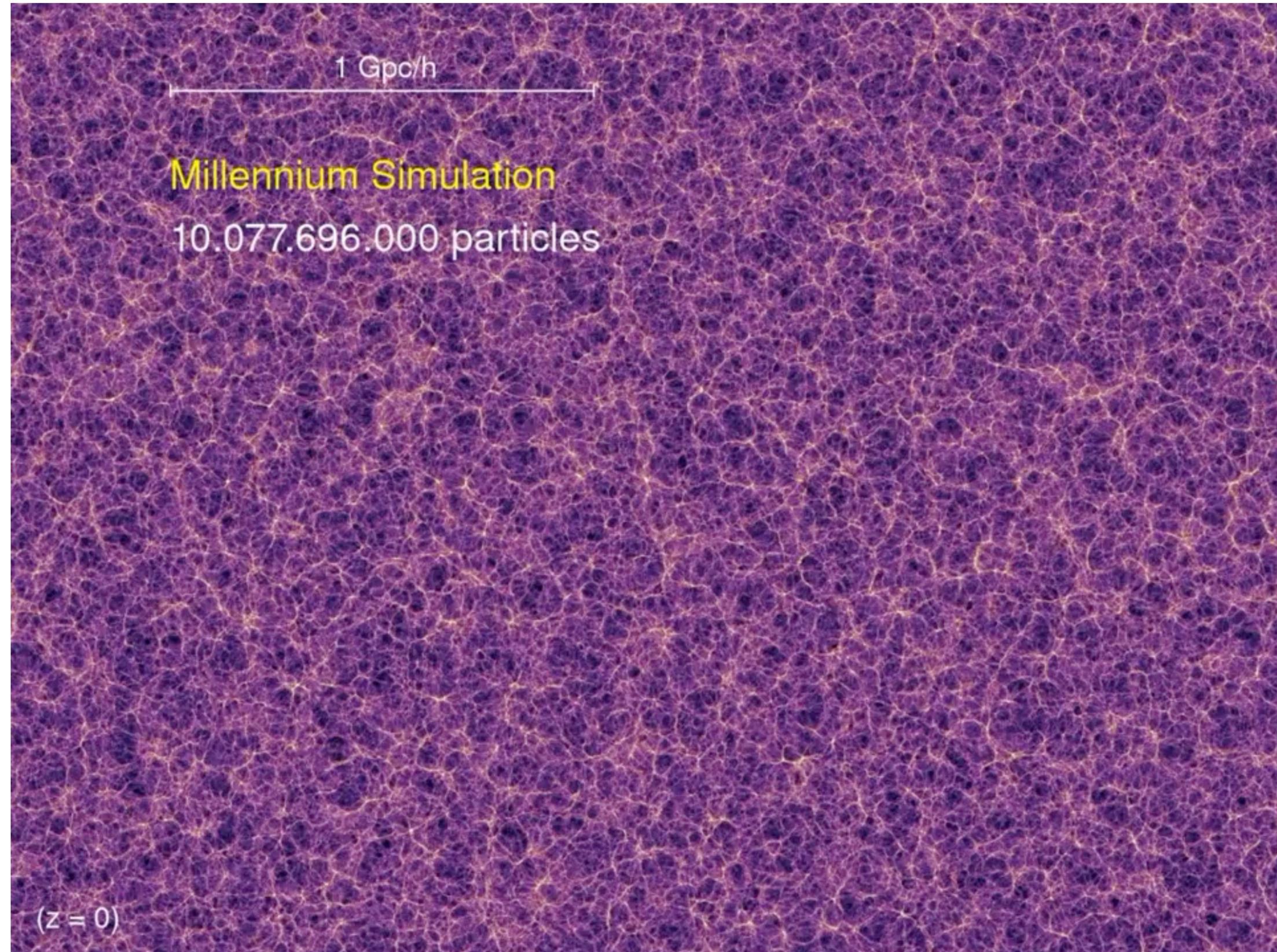
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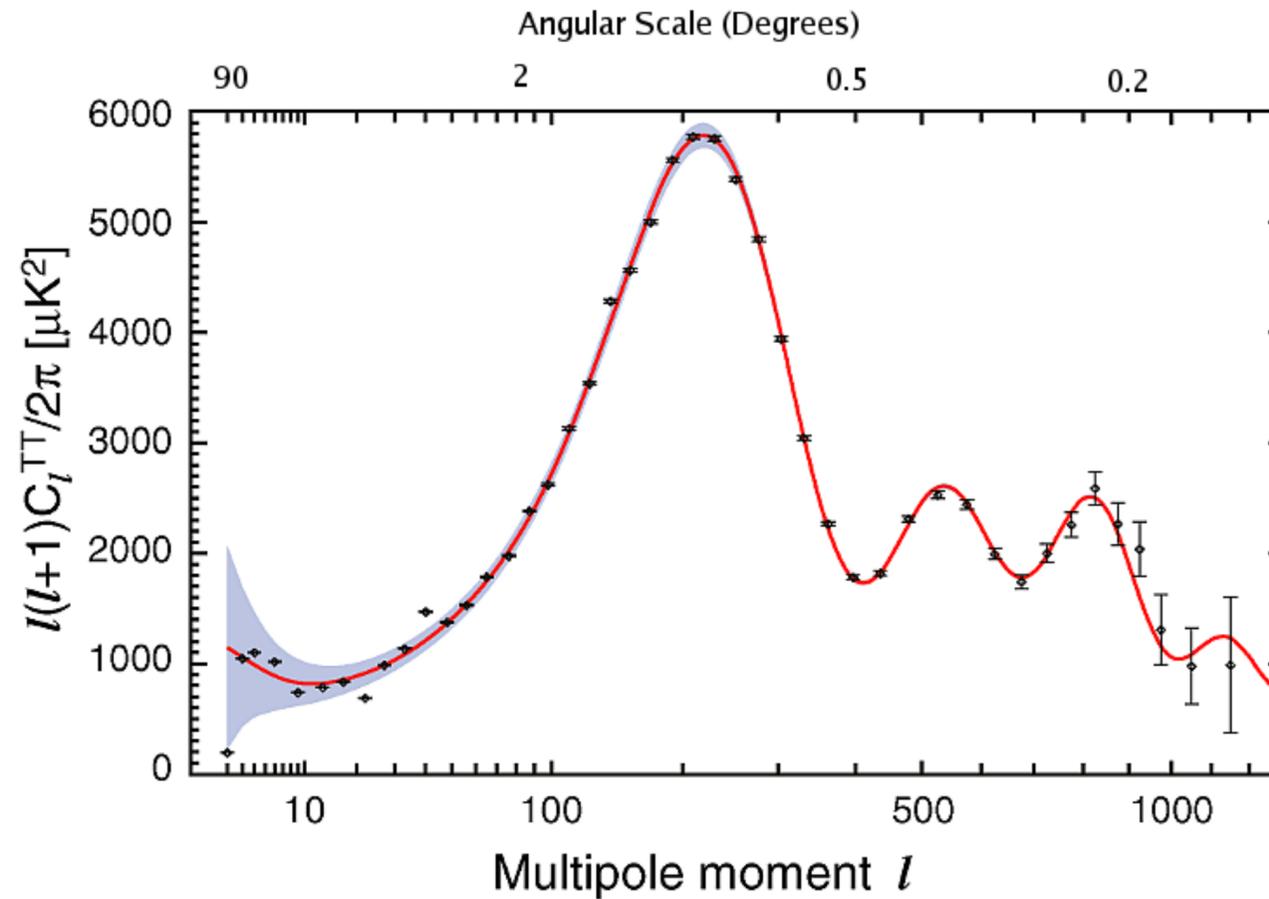
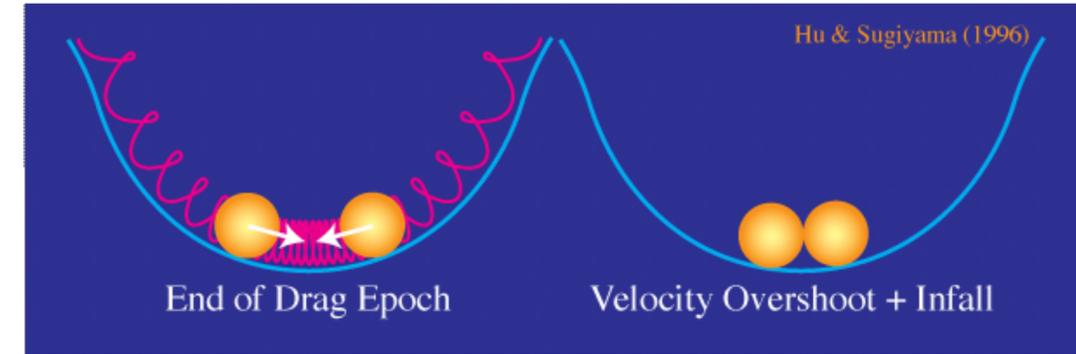
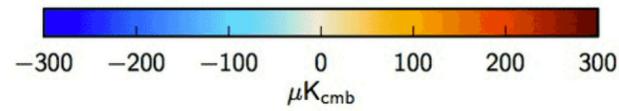
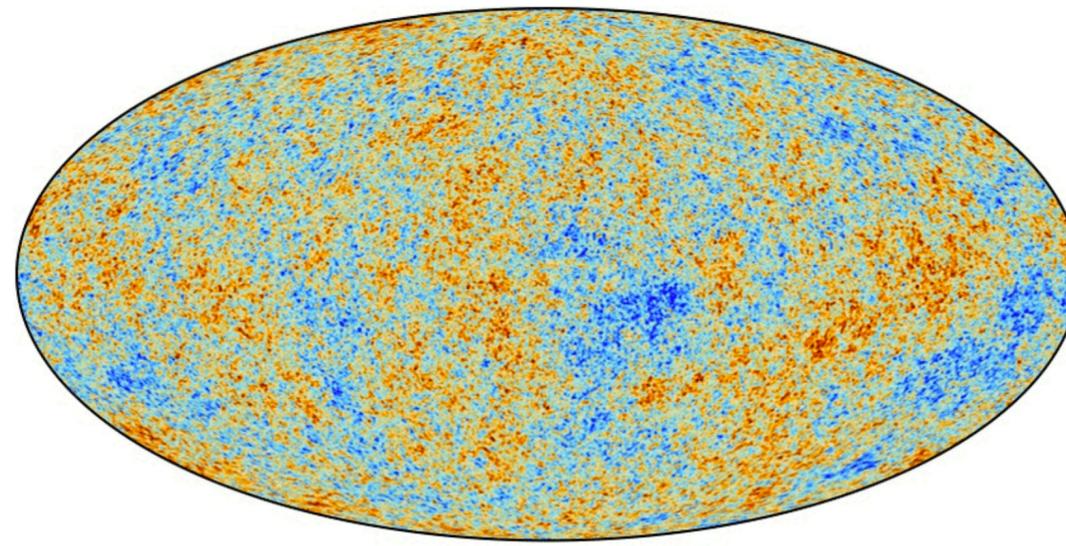
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The *Millennium Run* used more than 10 billion particles to trace the evolution of the matter distribution in a cubic region of the Universe over 2 billion light-years on a side.



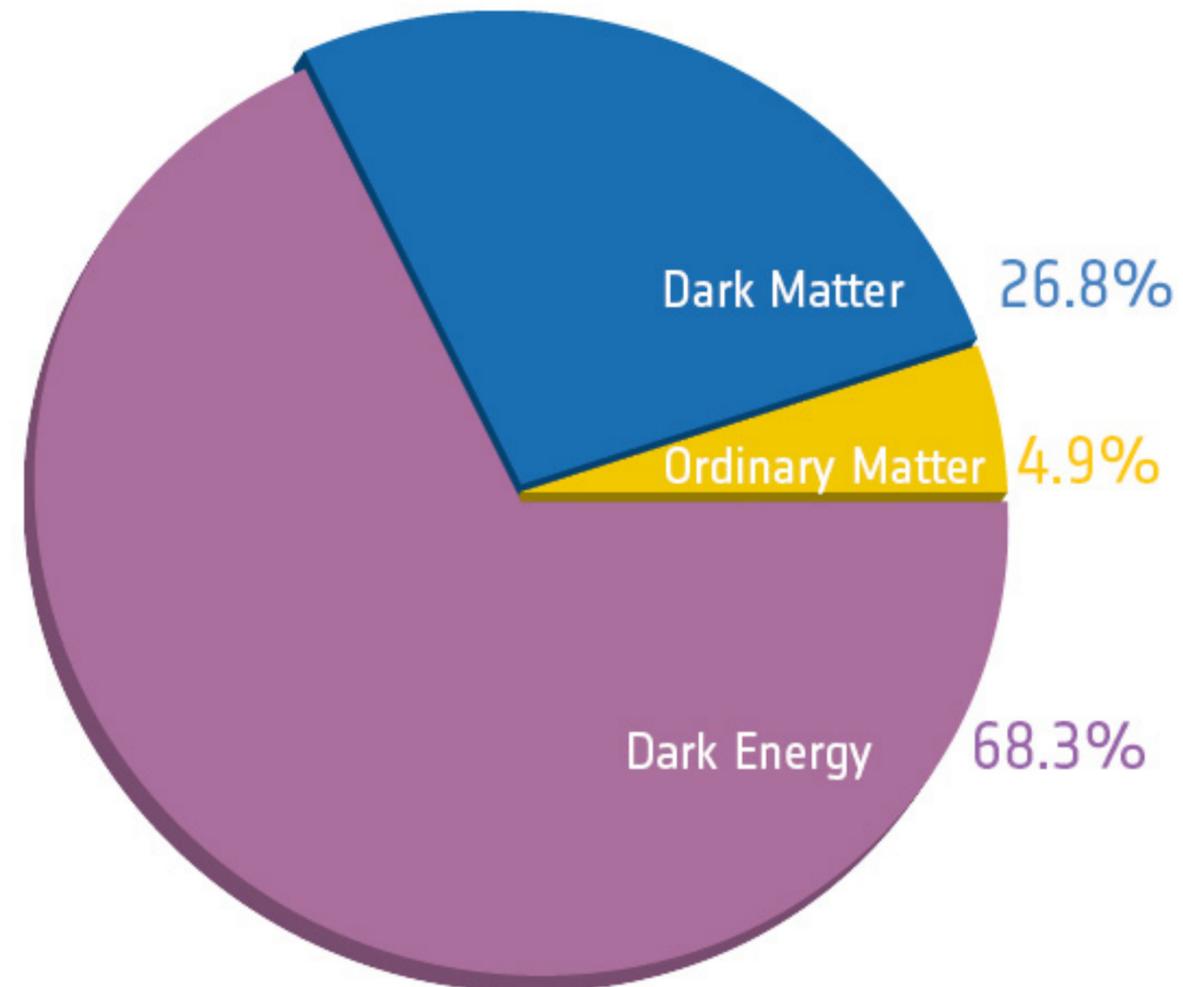
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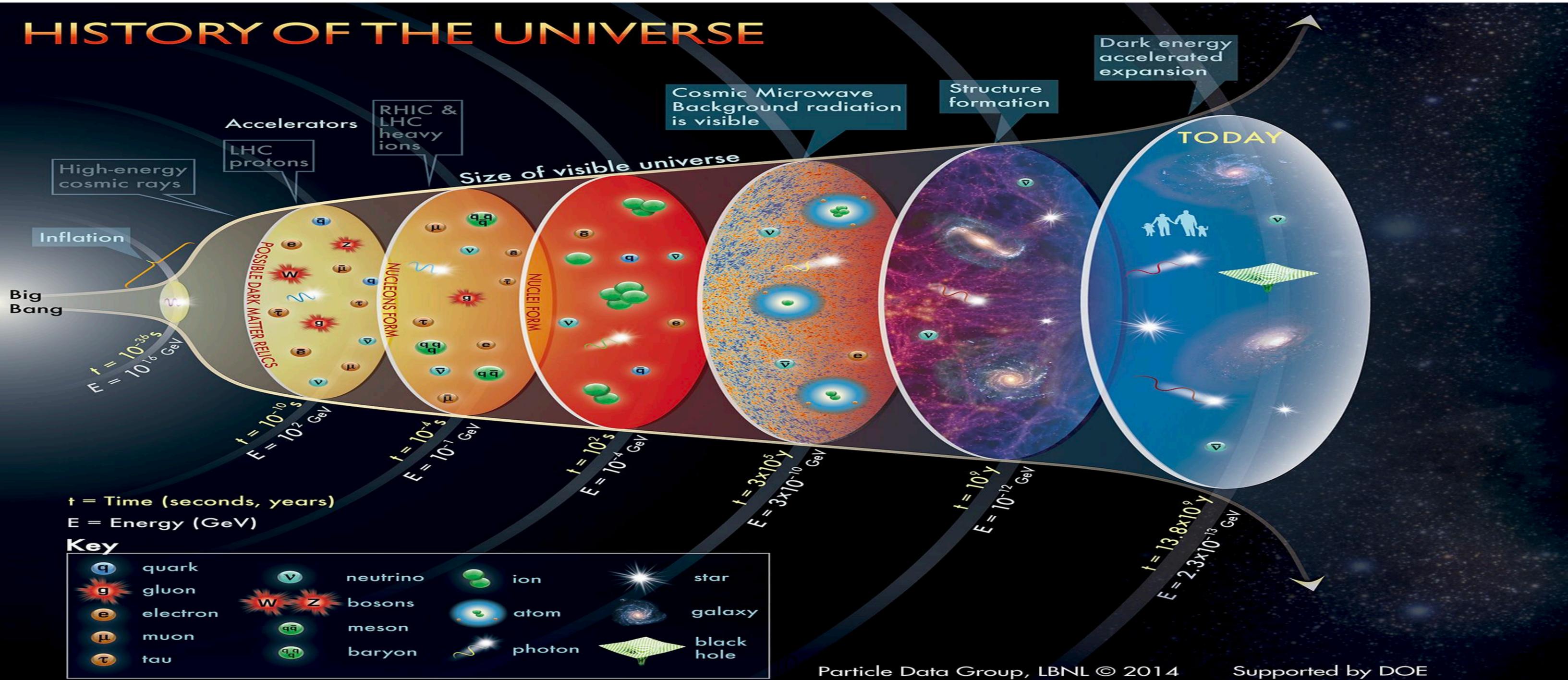
The success of the Lambda cold dark matter Model

- The standard model of Big Bang cosmology
- Λ , dark energy; CDM, cold dark matter; Matter, SM particles



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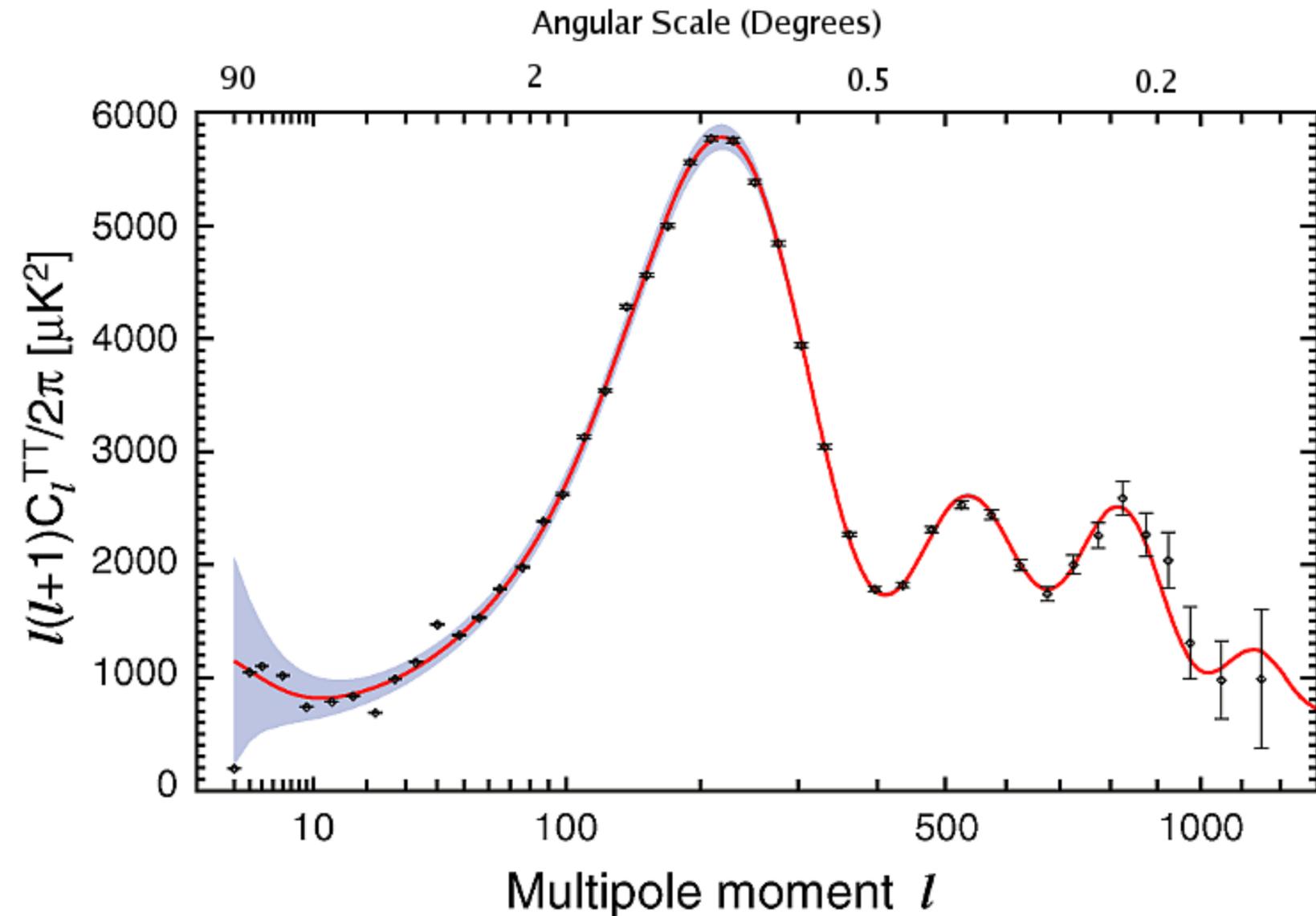


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- The standard model of Big Bang cosmology
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- 6 parameter for the Universe: Baryon matter density, DM density, lifetime of the Universe ...
 - Explain the structure of the CMB
 - Large-scale structure in the distribution of the galaxies
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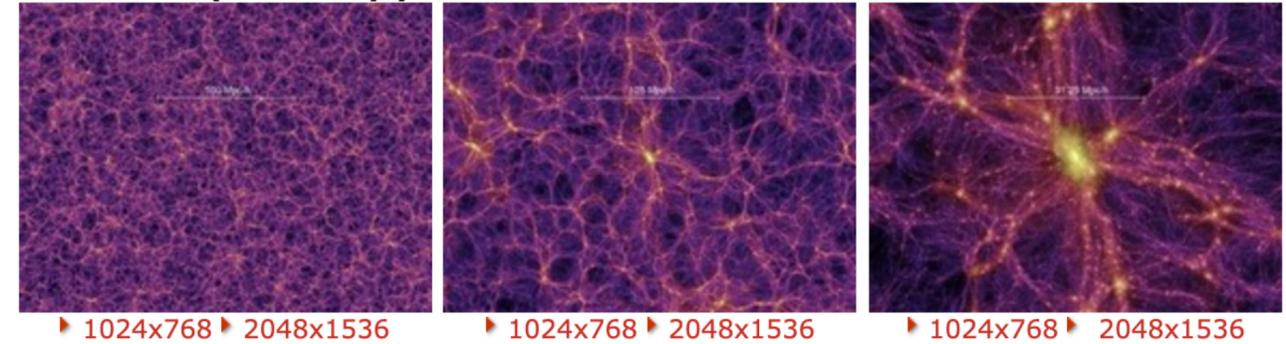
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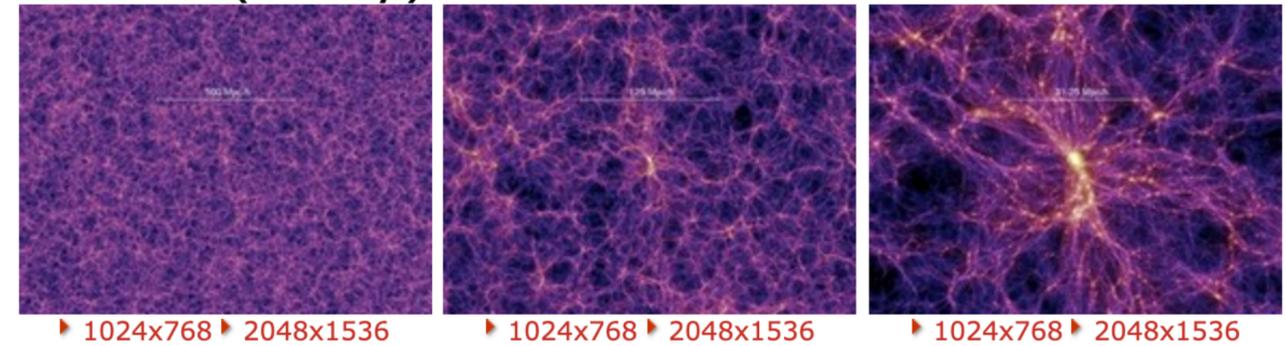
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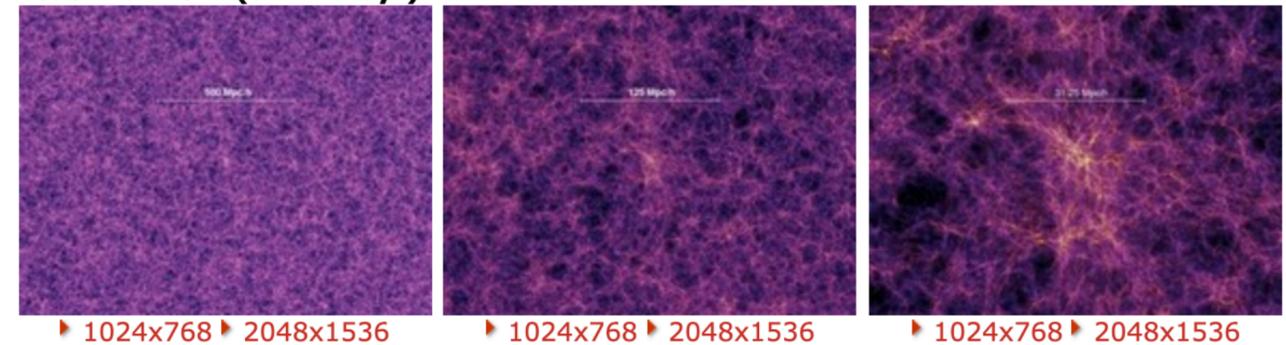
Redshift $z=0$ ($t = 13.6$ Gyr):



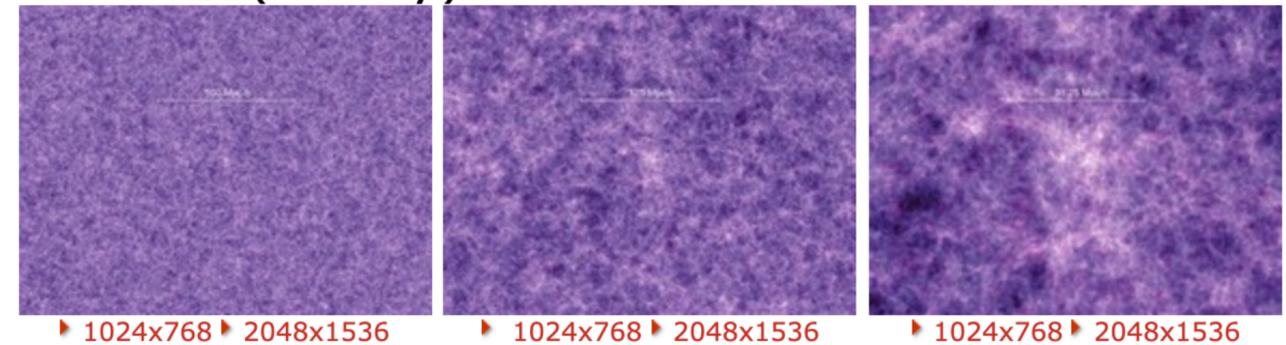
Redshift $z=1.4$ ($t = 4.7$ Gyr):



Redshift $z=5.7$ ($t = 1.0$ Gyr):

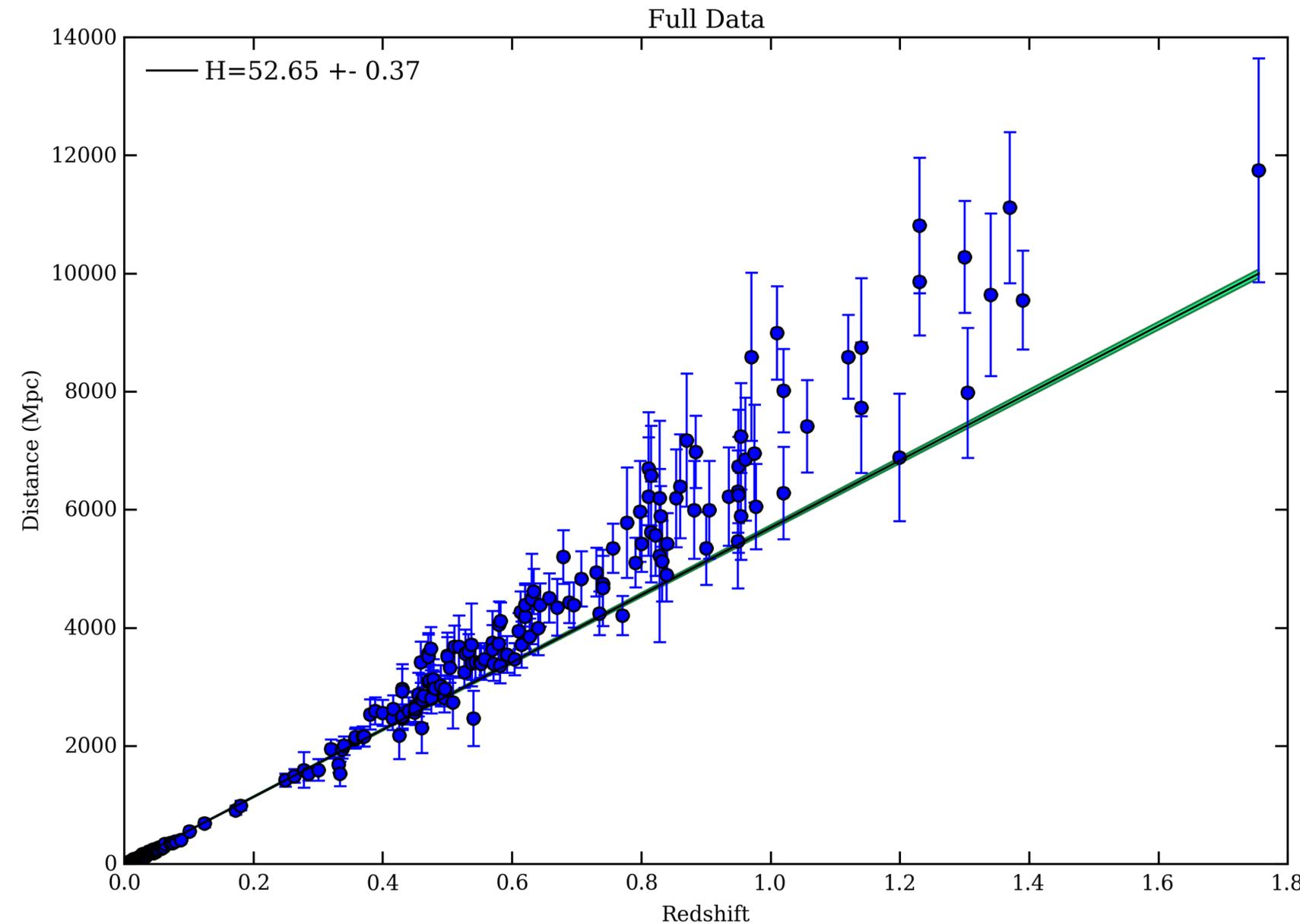


Redshift $z=18.3$ ($t = 0.21$ Gyr):



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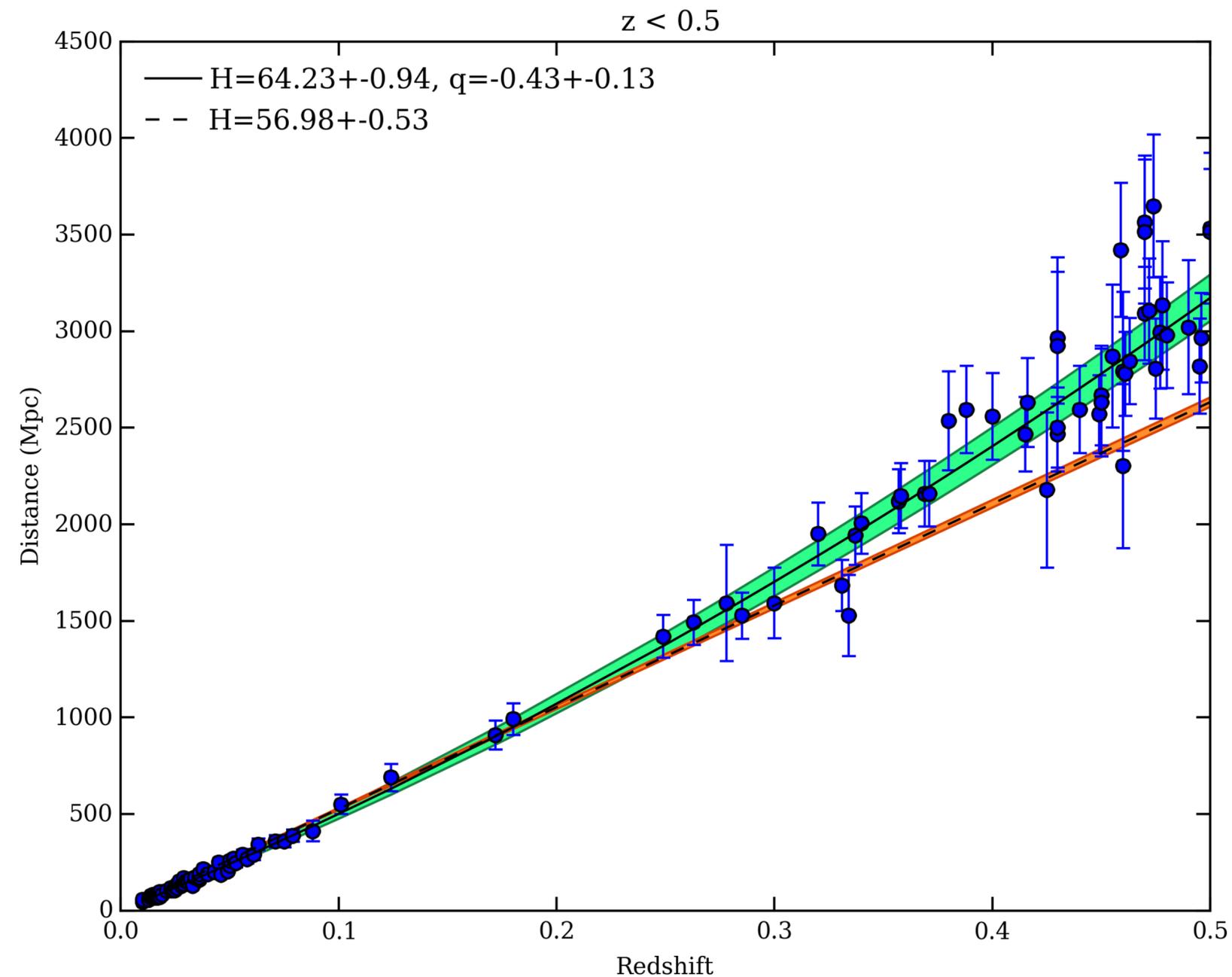


(Re)Discovering Dark Energy and the Expanding Universe

Credit: 2016, Adam Dempsey. CIERA, Northwestern University

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The cosmology frontier of Particle Physics



The dark matter in astrophysics/cosmology

- Energy density scales as $\rho \propto a^{-3}$, others $\rho_r \propto a^{-4}$, $\rho_{cc} \propto a^0$
- Massive, interacting gravitationally
- Neutral, not quite interacting with others, collision-less
- Stable

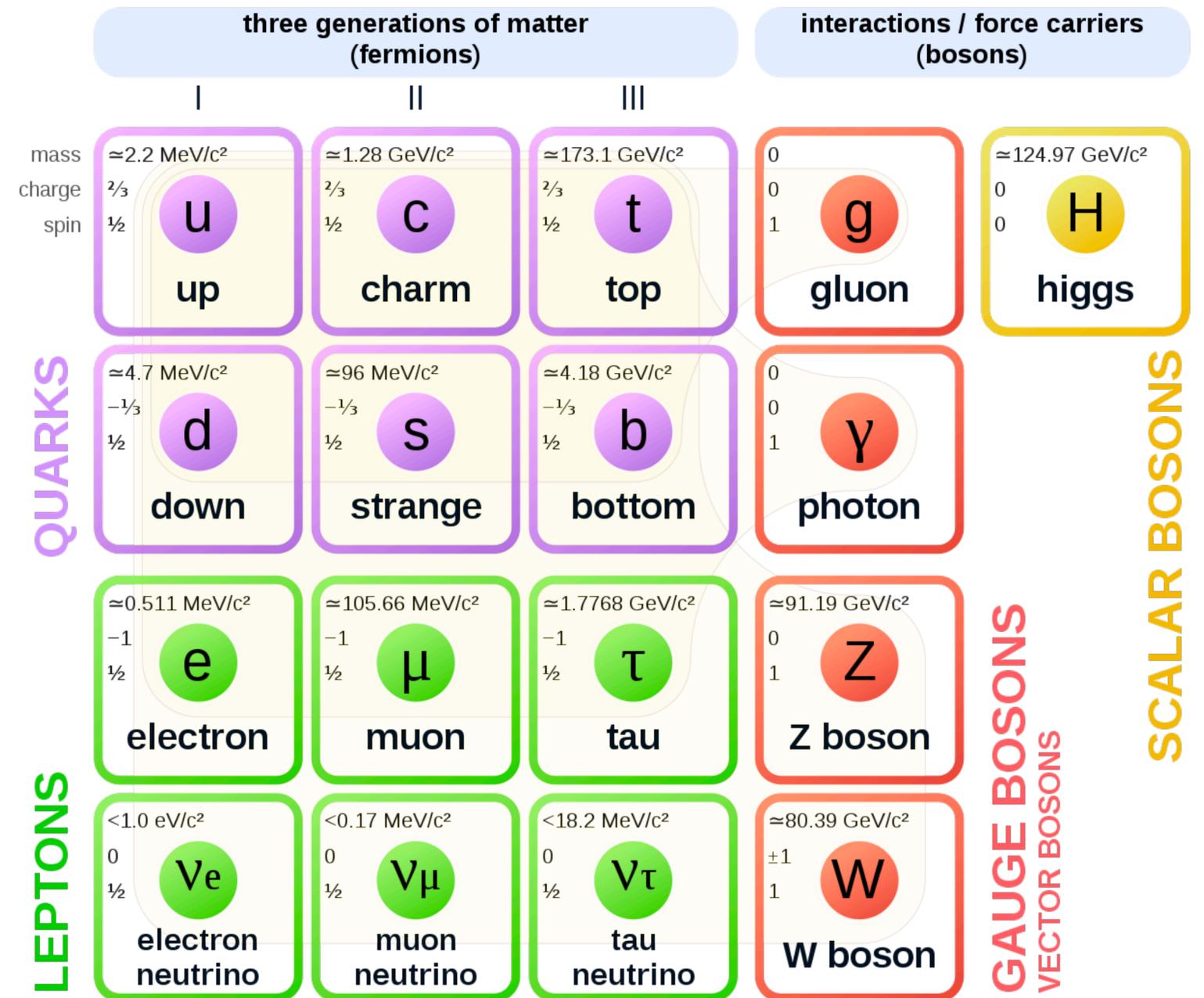
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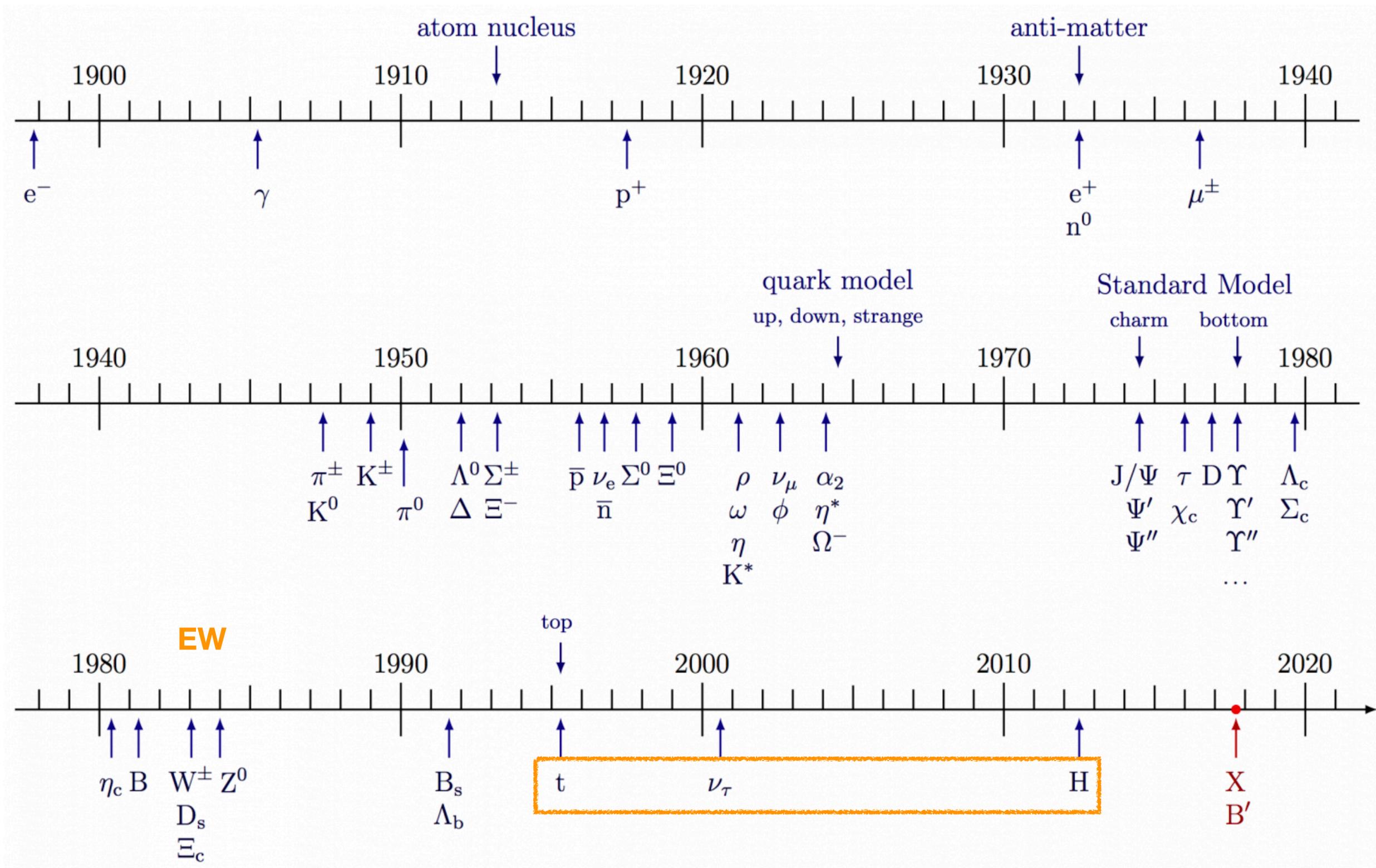
The dark matter in particle physics perspective

- No body knows what DM is
- Not in Standard Model
- There are good guesses

Standard Model of Elementary Particles



Chronicle of particle discoveries



The menu of the Standard Model



Top quark

The heaviest fermion
3rd gen up-type quark
1995



Tau neutrino

The last discovered neutrino
3rd gen neutrino
2000



Higgs

It gives masses to other particles
The heaviest scalar particle
2012

- We always hold a menu
- What about next?

The new menu from Supersymmetry

- SUSY model MSSM guided the phenomenology study for a long time
 - Sizable coupling to SM sector
 - Collider searches
 - Direct searches
 - Indirect searches
- Neutralino DM is well-motivated
- A role-model for Weakly Interacting Massive Particle

至尊私房菜
菜品价目表

炸菜类

- 新的时空对称性 x
- 玻色子和费米子对称性
- 超对称伴子

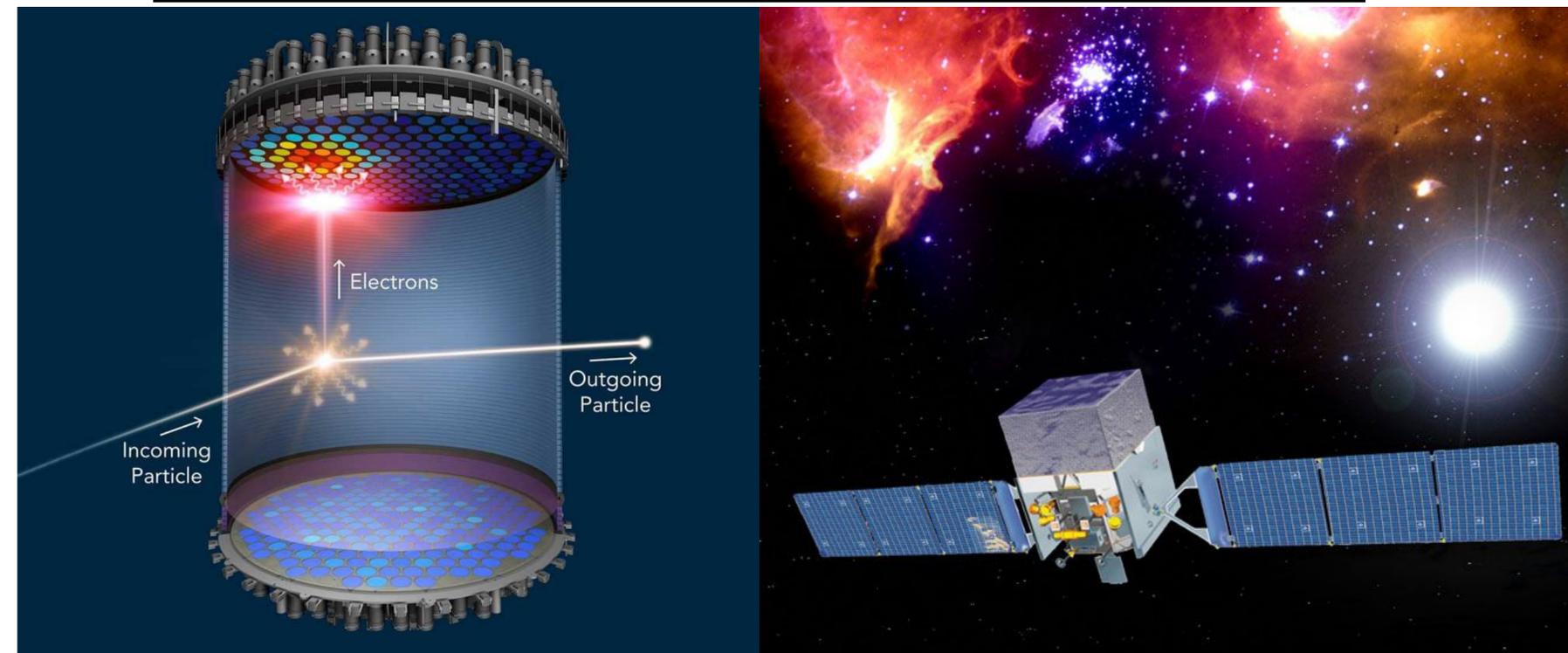
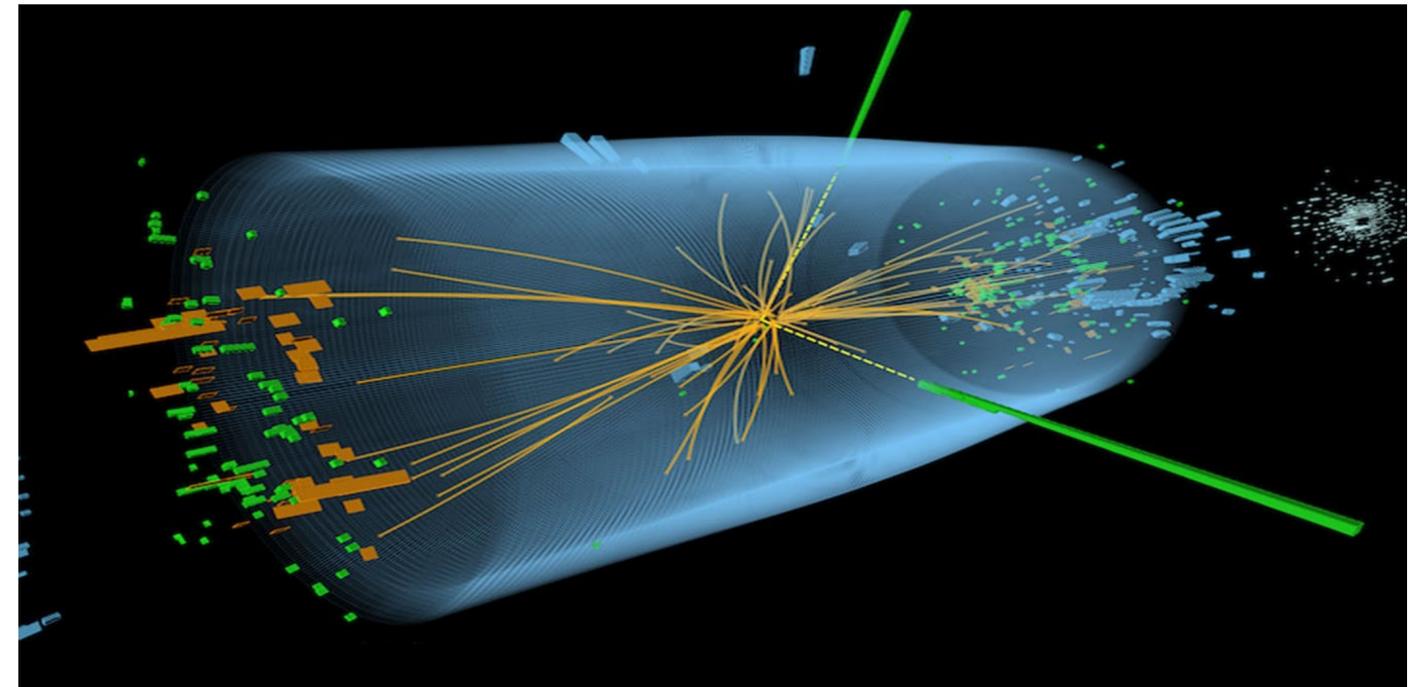
特色菜

- 希格斯质量问题 x
- Neutralino暗物质候选者
- 力的统一和MSSM

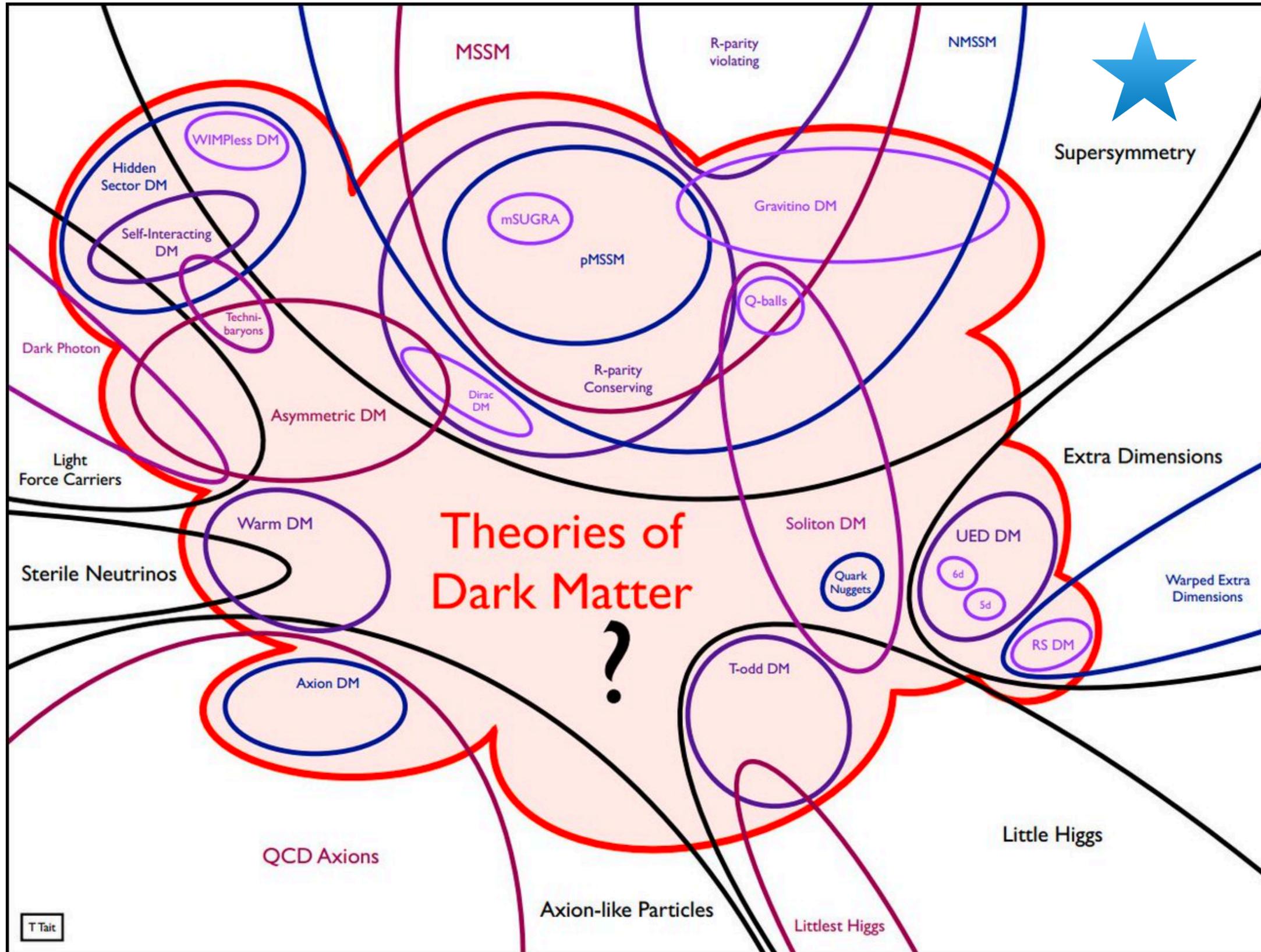
酸菜鱼 RM
野味菌 RM

The new menu from Supersymmetry

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The menu for dark matter models



The wide ranges of dark matter searches

- In this talk:

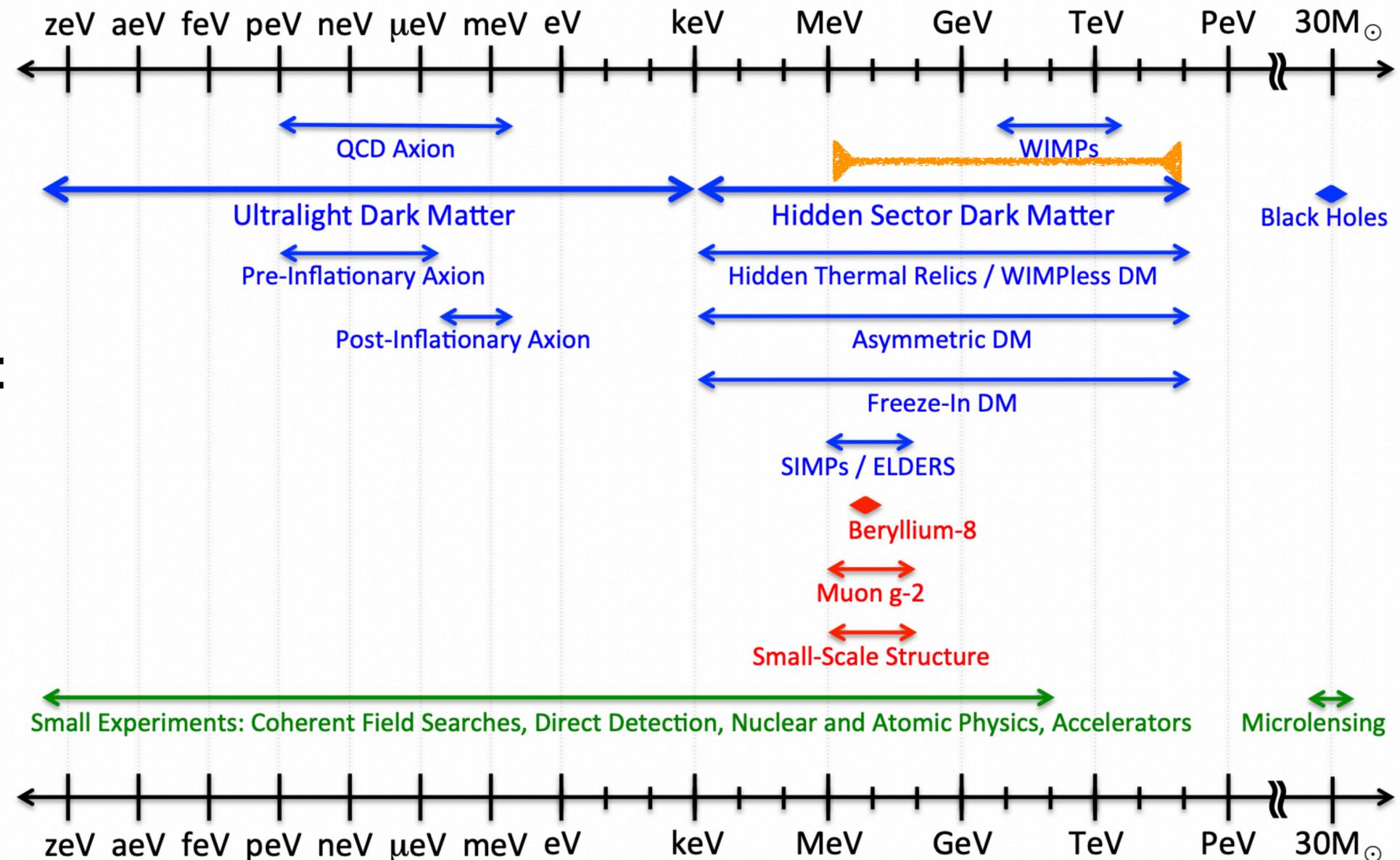
- Skip very small mass: DM is wave-like rather than particle like

- Skip very small interaction rate: DM is accumulated through thermal leakage of SM fields

- We focus on mass range \sim [MeV (Neff bound), 100 TeV (unitary bound)]

- DM starts in thermal equilibrium

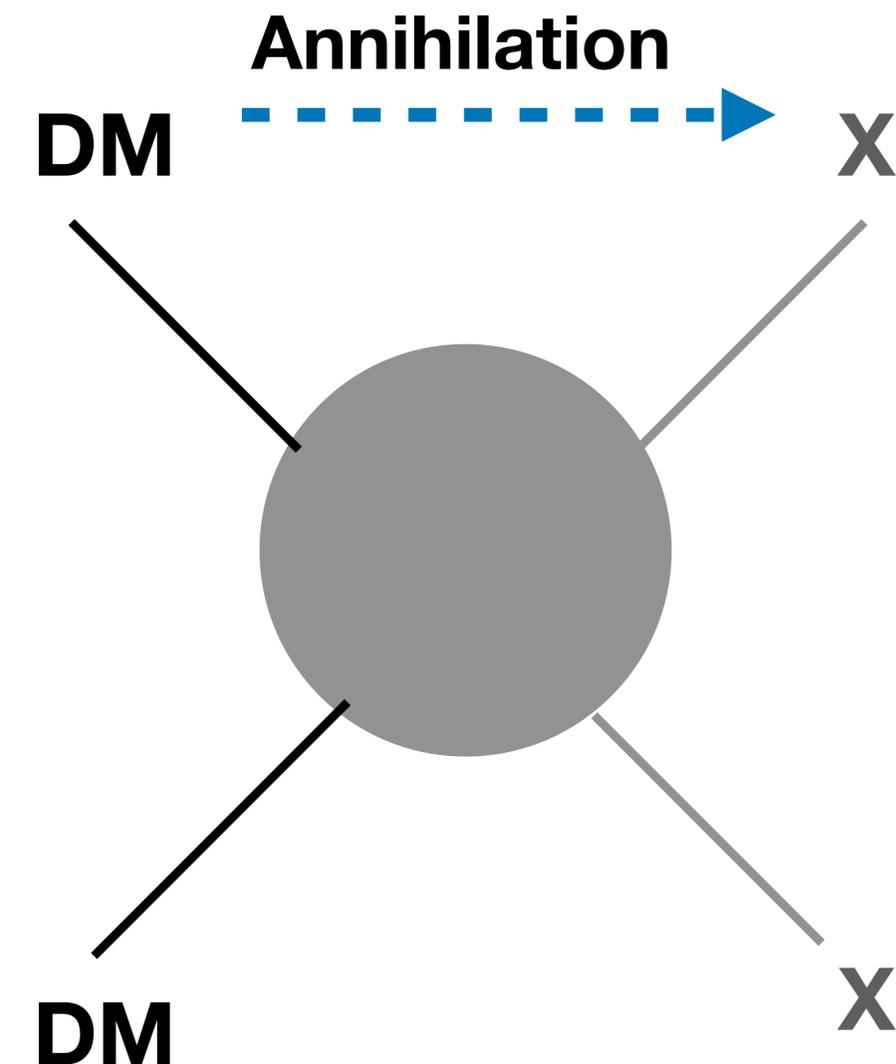
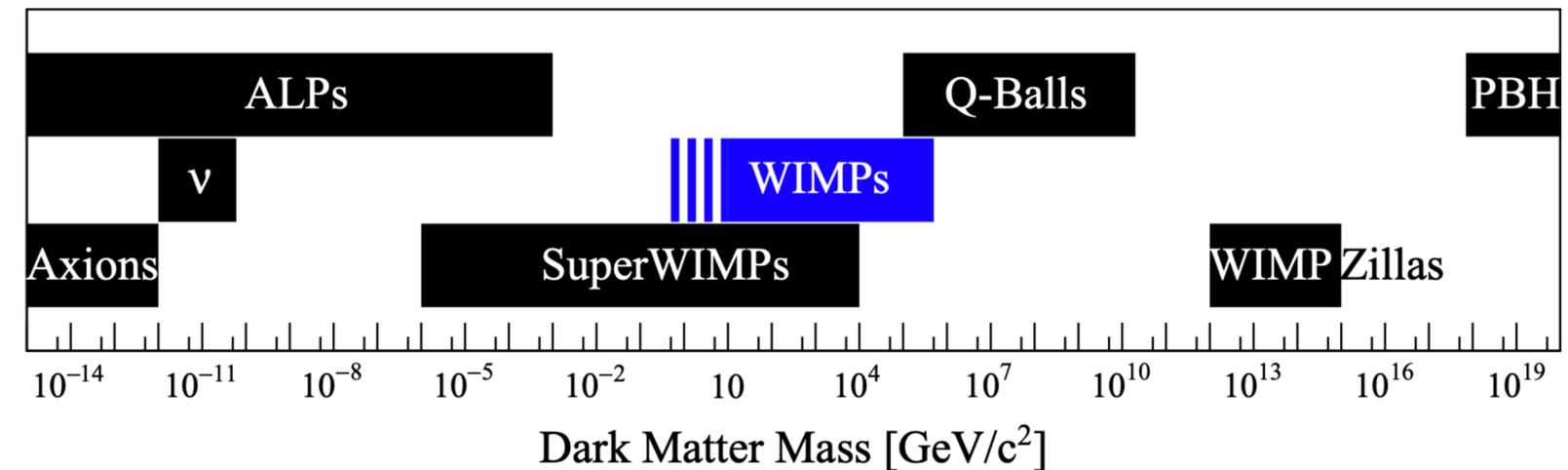
Dark Sector Candidates, Anomalies, and Search Techniques



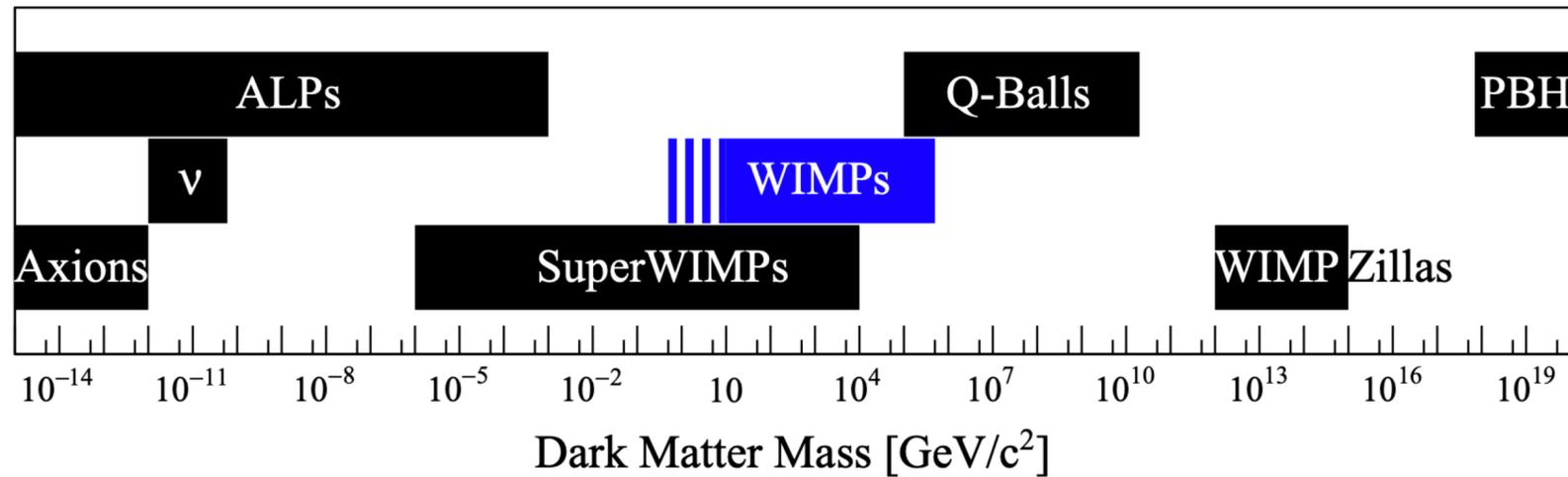
Cosmic Visions [arXiv:1707.04591]

The Weakly Interacting Massive Particle paradigm

- DM is a massive elementary particle
- DM has an electroweak-scale coupling
- DM starts with thermal distribution
- Relic abundance is determined by freeze-out mechanism
- **DM Annihilation into**
 - **X = Standard Model particles (direct coupling)**
 - **X = Dark Sector particles (secluded DM models)**



The freeze-out of WIMP DM

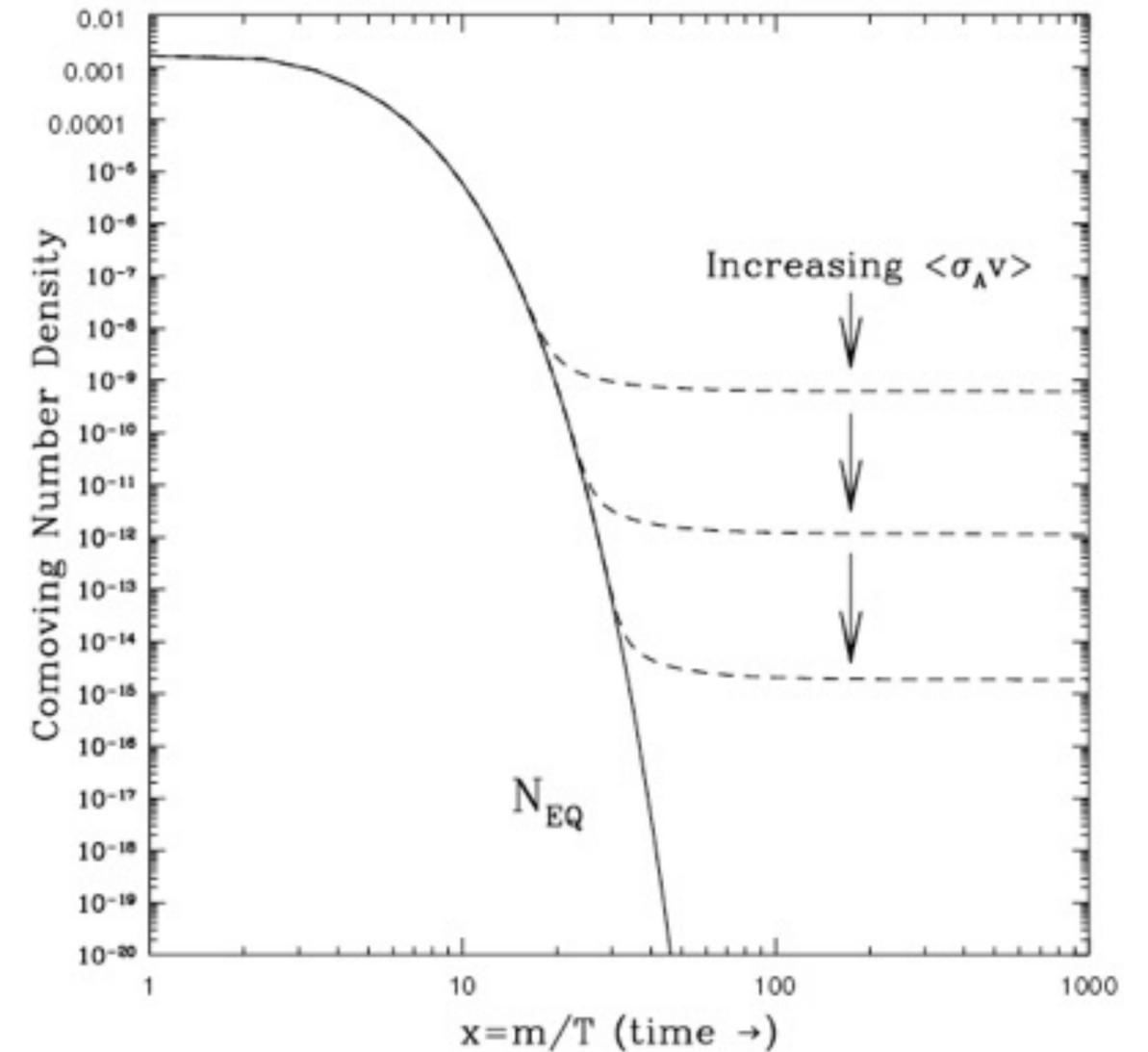


- Thermal cross-section

$$\langle \sigma v \rangle \sim \frac{\alpha^2}{m_W^2} \sim 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$$

- DM Annihilation cross-section

$$\langle \sigma v \rangle \sim \frac{g^4}{m_{\text{DM}}^2} \Rightarrow g \sim \sqrt{\frac{m_{\text{DM}}}{10 \text{TeV}}}$$

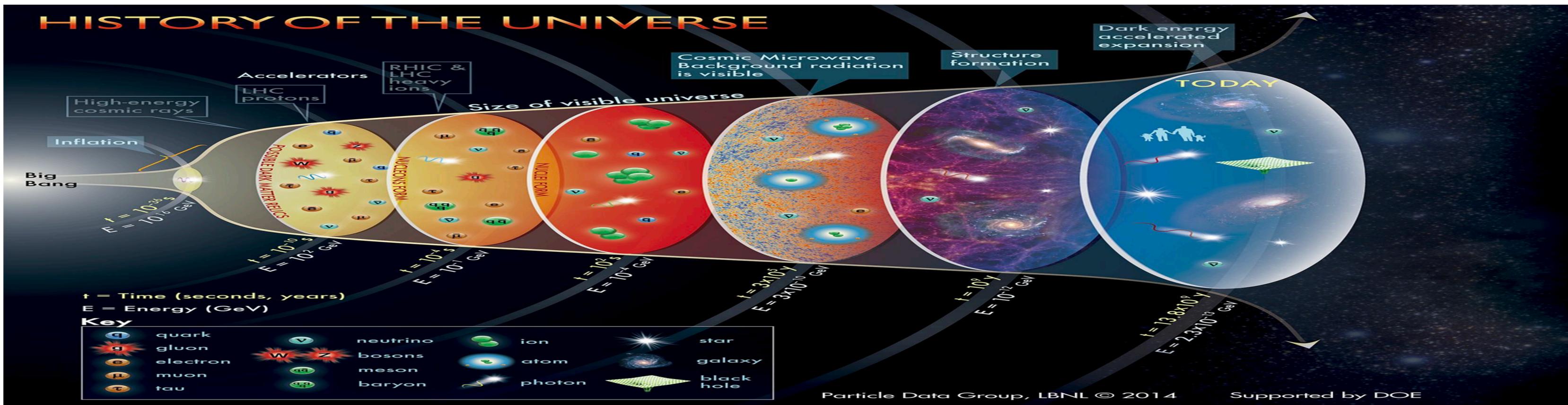


Jungman *et al* hep-ph/9506380

This is called WIMP miracle!

The WIMP DM and freeze-out

- DM relic abundance
 - No further UV info needed (started with a thermal distribution)
 - Electroweak scale annihilation cross-section
 - Similar stories in SM (ν decoupling, n_p/n_n ratio, nuclear elements)
 - Leads to collider/direct/indirect signal as well

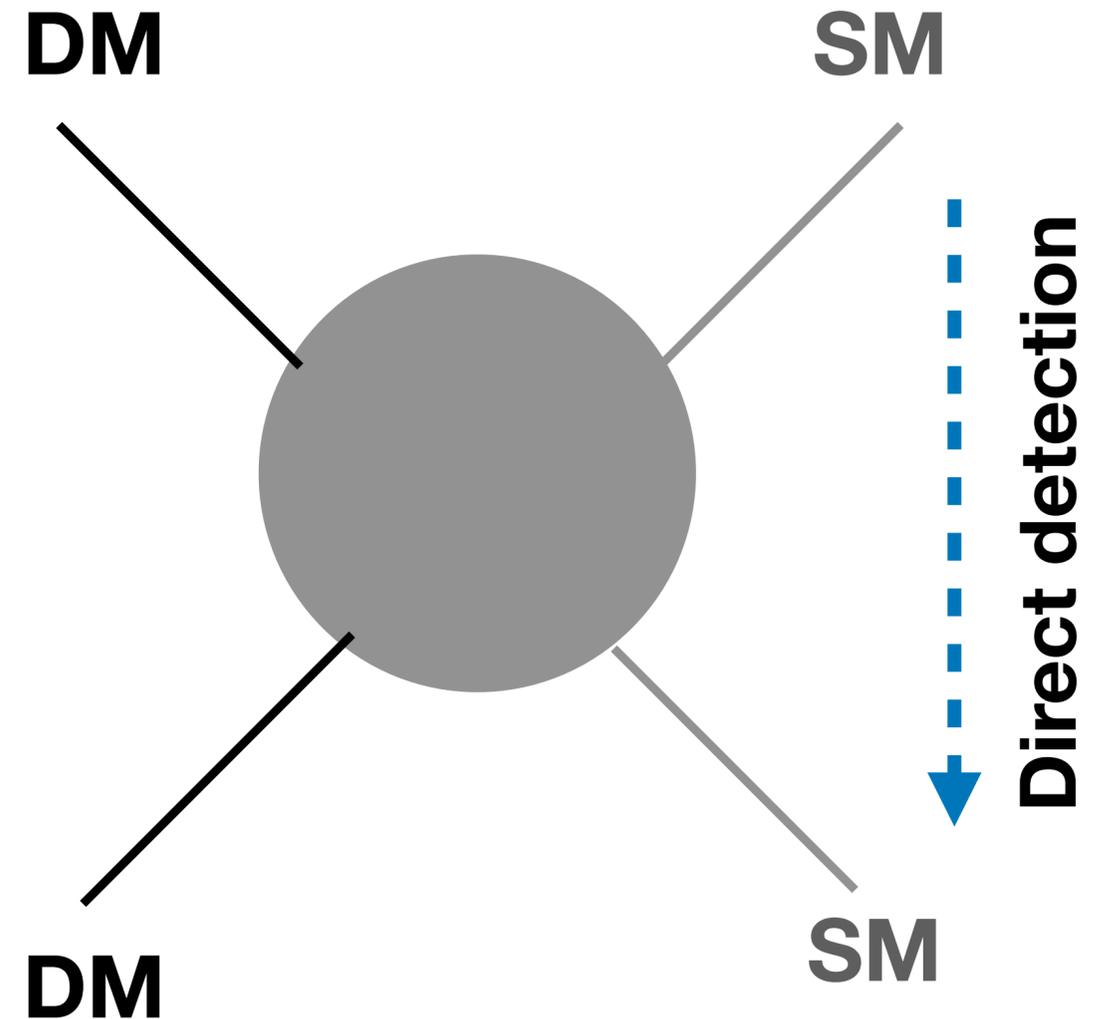
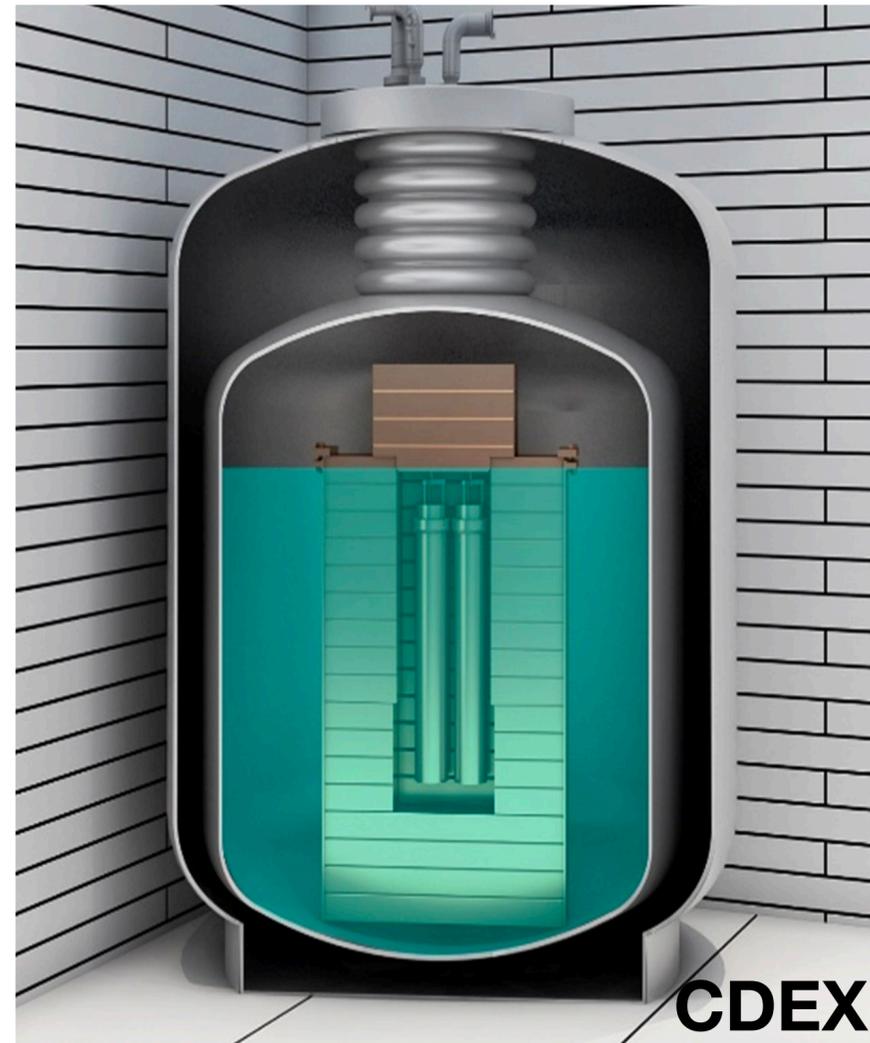
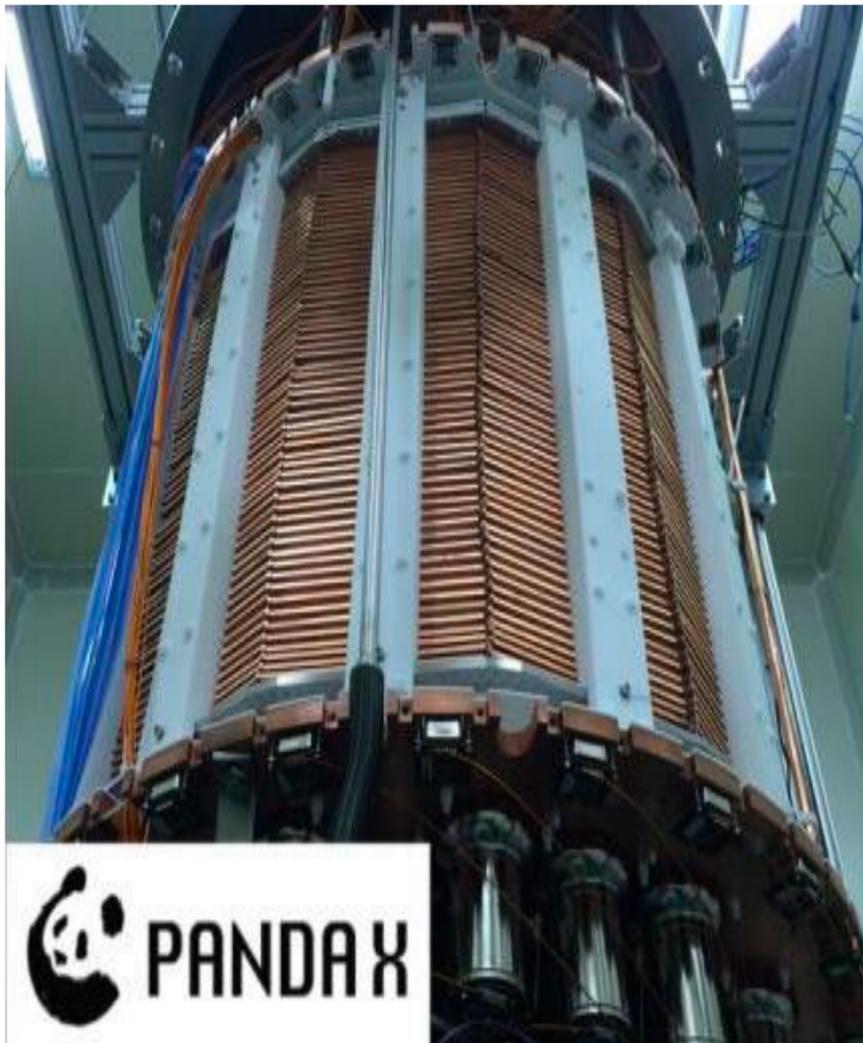


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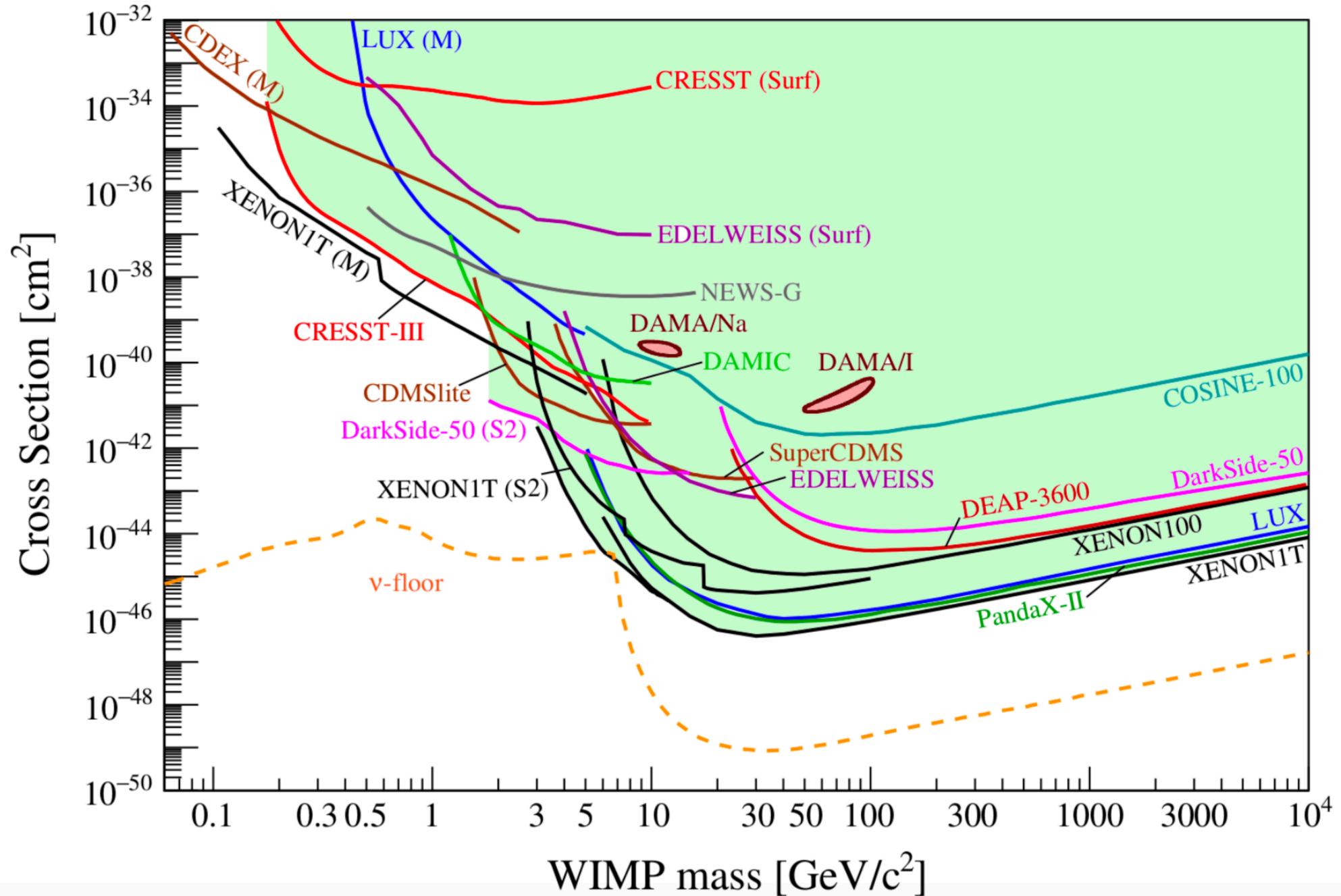
The WIMP crisis from direct detection

- Weakly Interacting Massive Particle
- The sizable coupling of DM to SM particles predicts sizable scattering cross-section



The WIMP crisis from direct detection

- Null result from direct detection
- Maybe discovery in the corner?
- Neutrino floor and beyond: directional ..
- The rise of light dark matter ($\lesssim 10$ GeV)
- We focus on EW scale ($\gtrsim 10$ GeV)



The WIMP crisis from direct detection

- Null result from direct detection

Together with the fact that, we have not seen SUSY either.

- Maybe discovery in the corner? 转角遇到?

- Neutrino floor and beyond: directional ..

- The rise of light dark matter ($\lesssim 10$ GeV)

- We focus on EW scale ($\gtrsim 10$ GeV)

ATLAS SUSY Searches* - 95% CL Lower Limits

June 2021

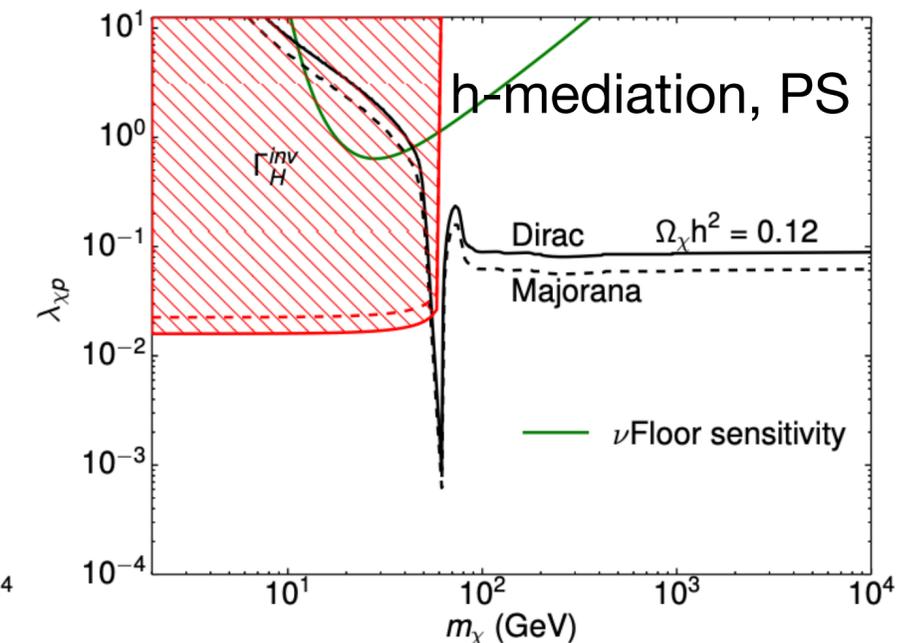
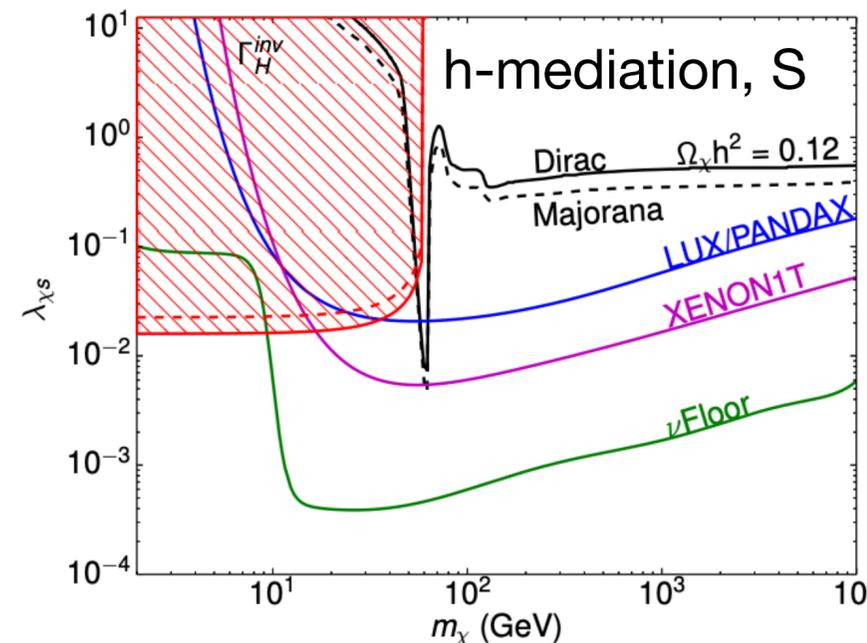
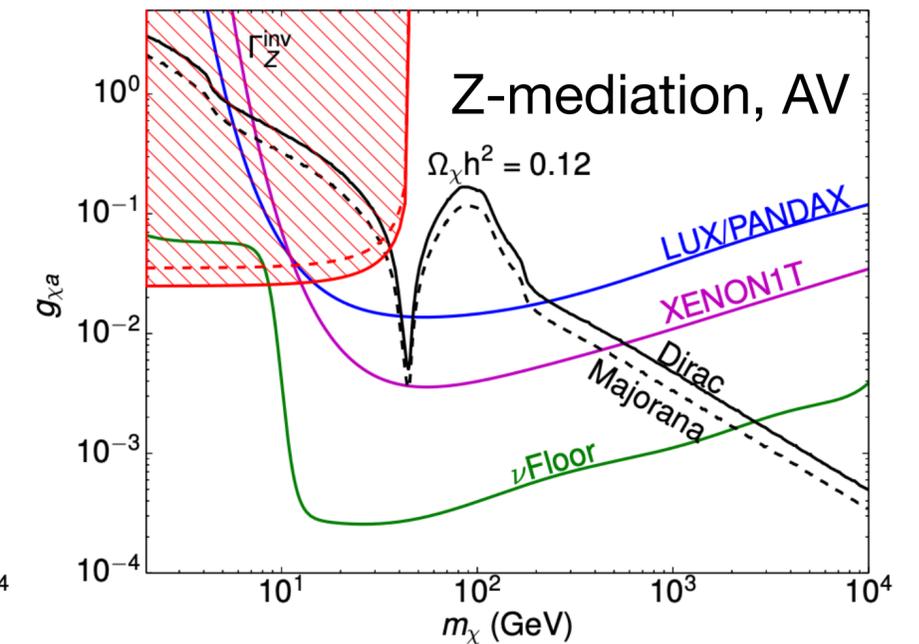
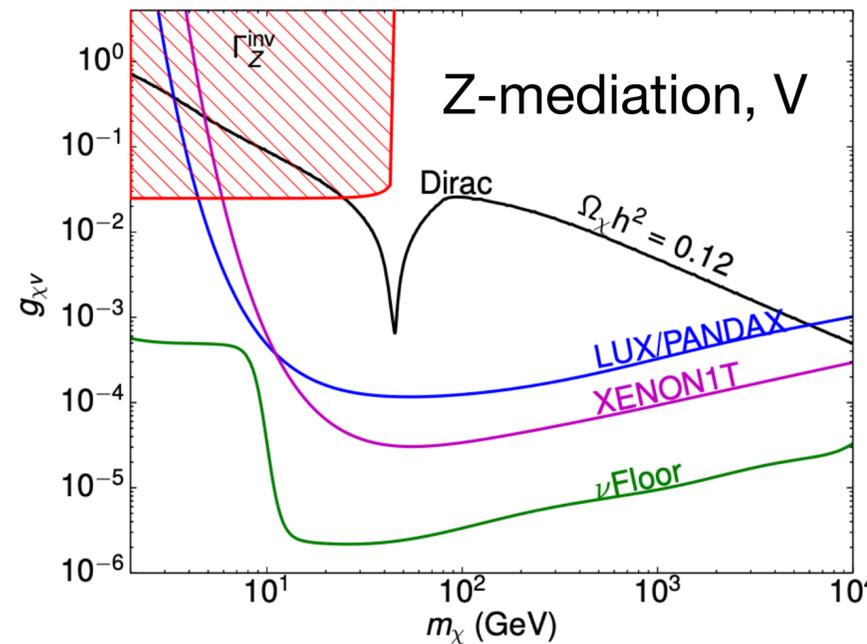
Model	Signature	$\int \mathcal{L} dt$ [fb $^{-1}$]	Mass limit	
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	E_T^{miss} 139 E_T^{miss} 36.1	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ 2-6 jets	E_T^{miss} 139	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}W\tilde{\chi}_1^0$	1 e, μ 2-6 jets	E_T^{miss} 139	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	$ee, \mu\mu$ 2 jets	E_T^{miss} 36.1	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0 e, μ SS e, μ	E_T^{miss} 139 6 jets	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ SS e, μ	E_T^{miss} 79.8 6 jets	
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1$	0 e, μ 2 b	E_T^{miss} 139	
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0 \rightarrow bh\tilde{\chi}_1^0$	0 e, μ 2 τ	E_T^{miss} 139 E_T^{miss} 139	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ ≥ 1 jet	E_T^{miss} 139	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	1 e, μ 3 jets/1 b	E_T^{miss} 139	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$	1-2 τ 2 jets/1 b	E_T^{miss} 139	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0 e, μ 0 e, μ	E_T^{miss} 36.1 mono-jet	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$	1-2 e, μ 1-4 b	E_T^{miss} 139	
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ 1 b	E_T^{miss} 139	
	EW direct	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ	Multiple ℓ /jets $ee, \mu\mu$	E_T^{miss} 139 E_T^{miss} 139
		$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via WW	2 e, μ	E_T^{miss} 139
$\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh		Multiple ℓ /jets	E_T^{miss} 139	
$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$		2 e, μ	E_T^{miss} 139	
$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$		2 τ	E_T^{miss} 139	
$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$		2 e, μ $ee, \mu\mu$	E_T^{miss} 139 E_T^{miss} 139	
$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$		0 e, μ 4 e, μ 0 e, μ	E_T^{miss} 36.1 E_T^{miss} 139 E_T^{miss} 139	
		≥ 3 b 0 jets ≥ 2 large jets	E_T^{miss} 36.1 E_T^{miss} 139 E_T^{miss} 139	

屋漏偏逢连夜雨

The WIMP crisis from direct detection

- SM Higgs and Z mediated scenarios are highly constrained
- Other mediators without DD suppression is also highly constrained, e.g. A'
 - Unless in the resonant region

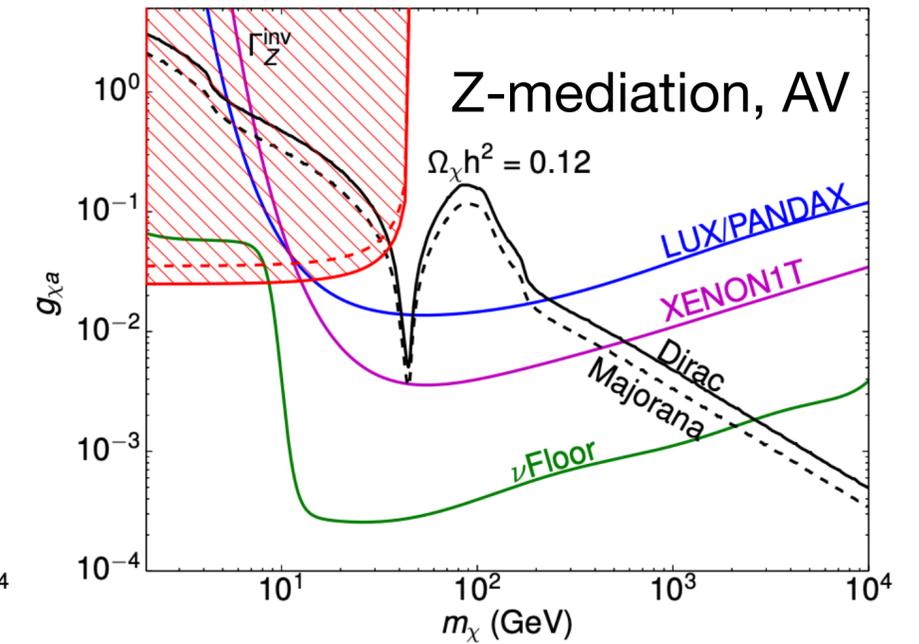
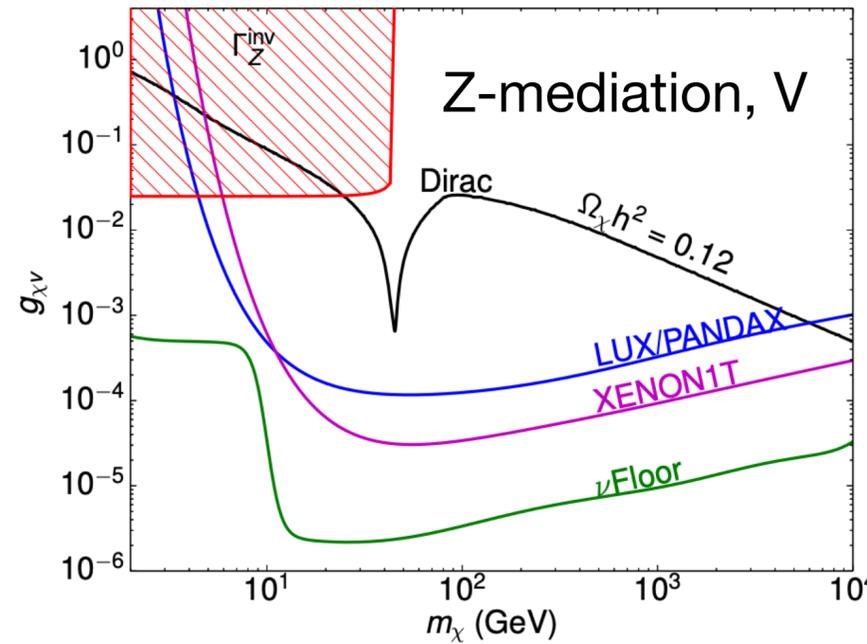
Toward (Finally!) Ruling Out Z and Higgs Mediated Dark Matter Models
 Hooper et al, ArXiv: 1609.09079, JCAP



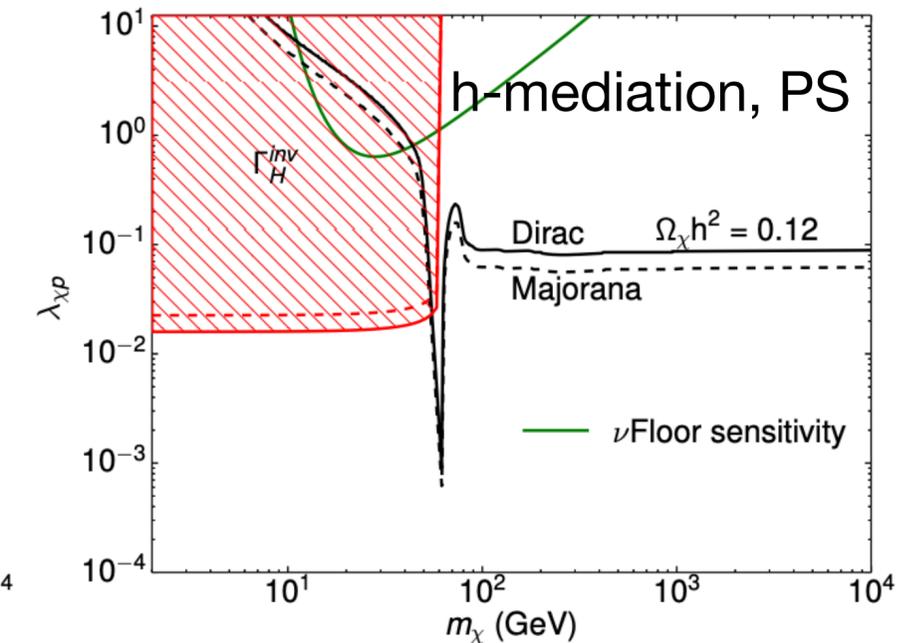
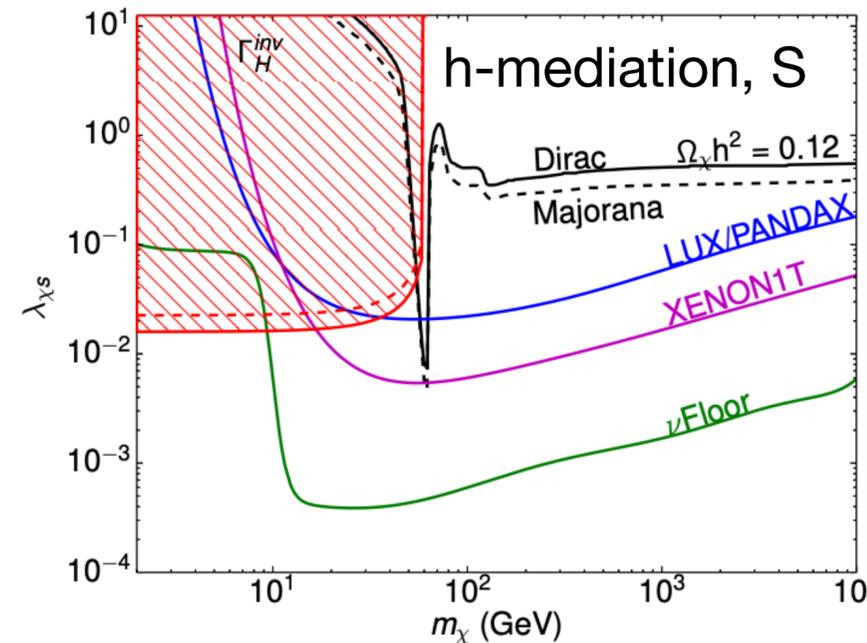
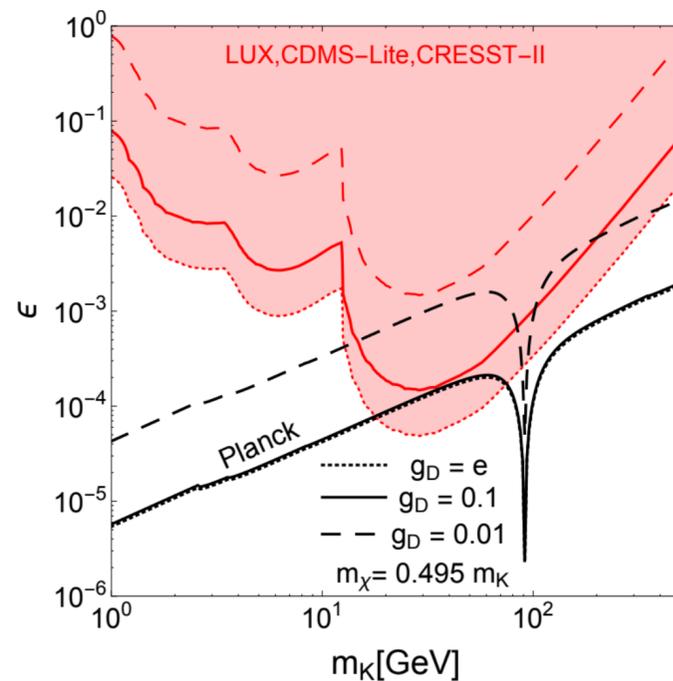
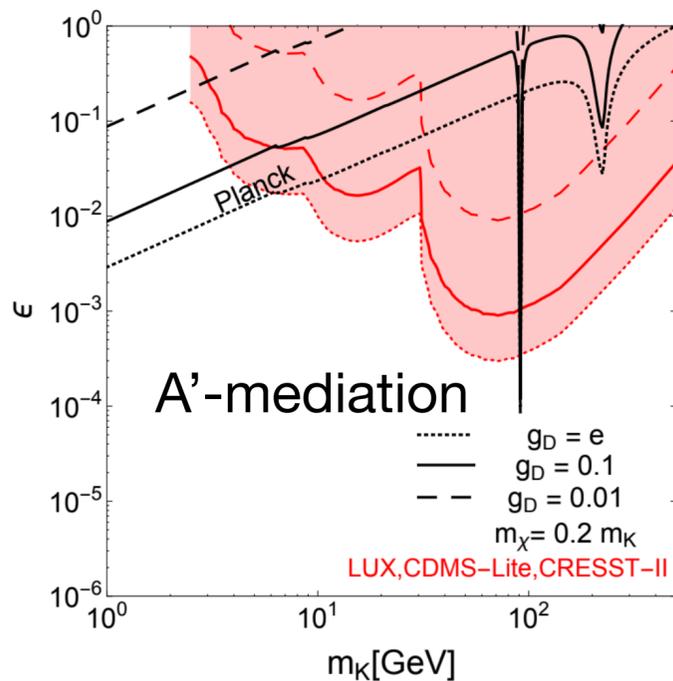
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Hooper et al, ArXiv: 1609.09079, JCAP

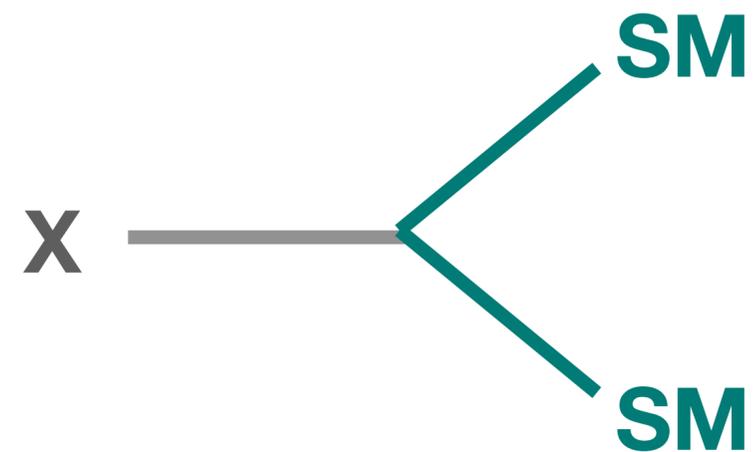
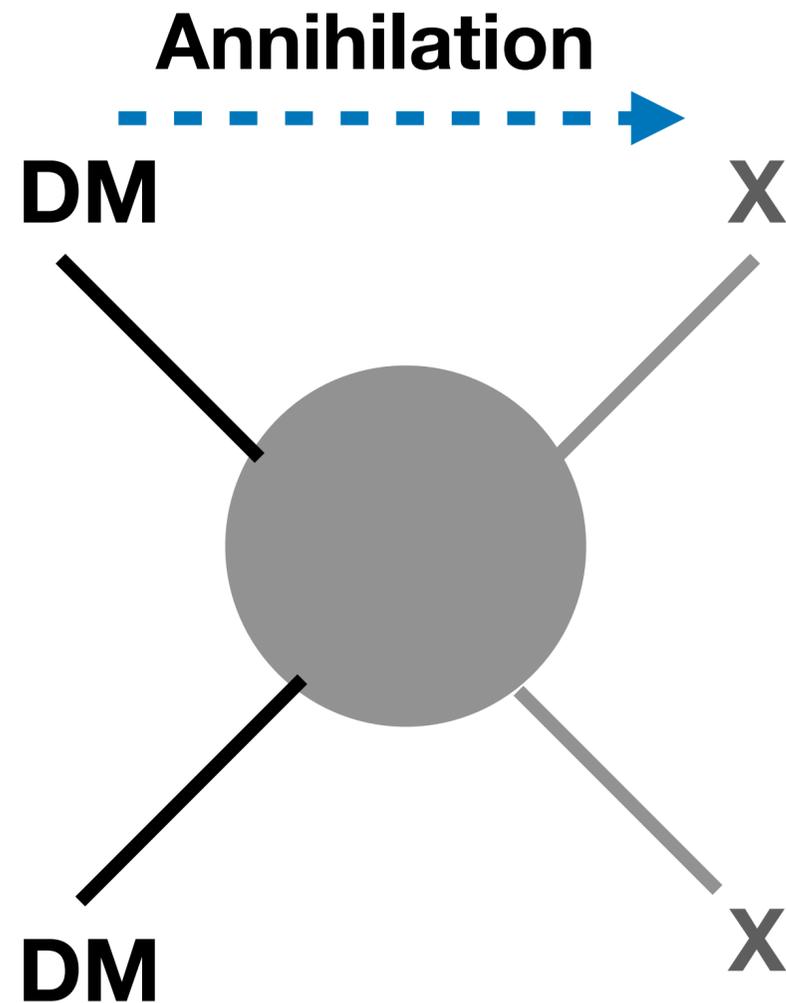


JL, X.P. Wang, F. Yu, 1704.00730, JHEP



The way-out from direct detection limits

- 1. Very small coupling:
 - 1.1 Secluded dark matter (dark sector)



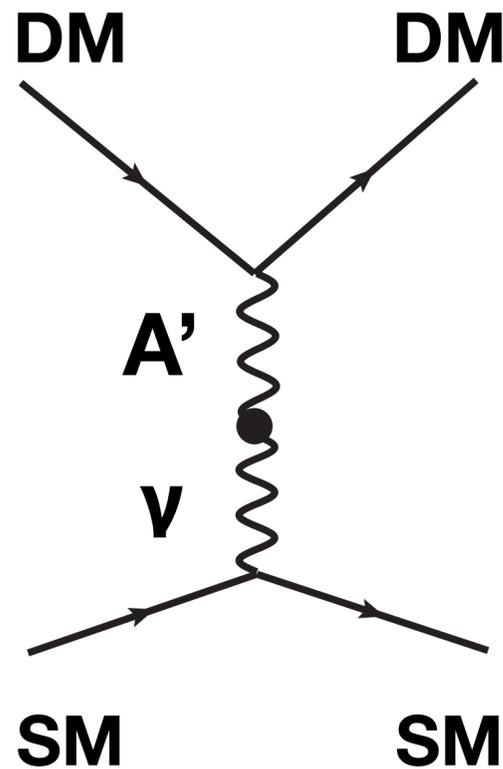
**Dark mediator
with very small coupling to SM**

The way-out from direct detection limits

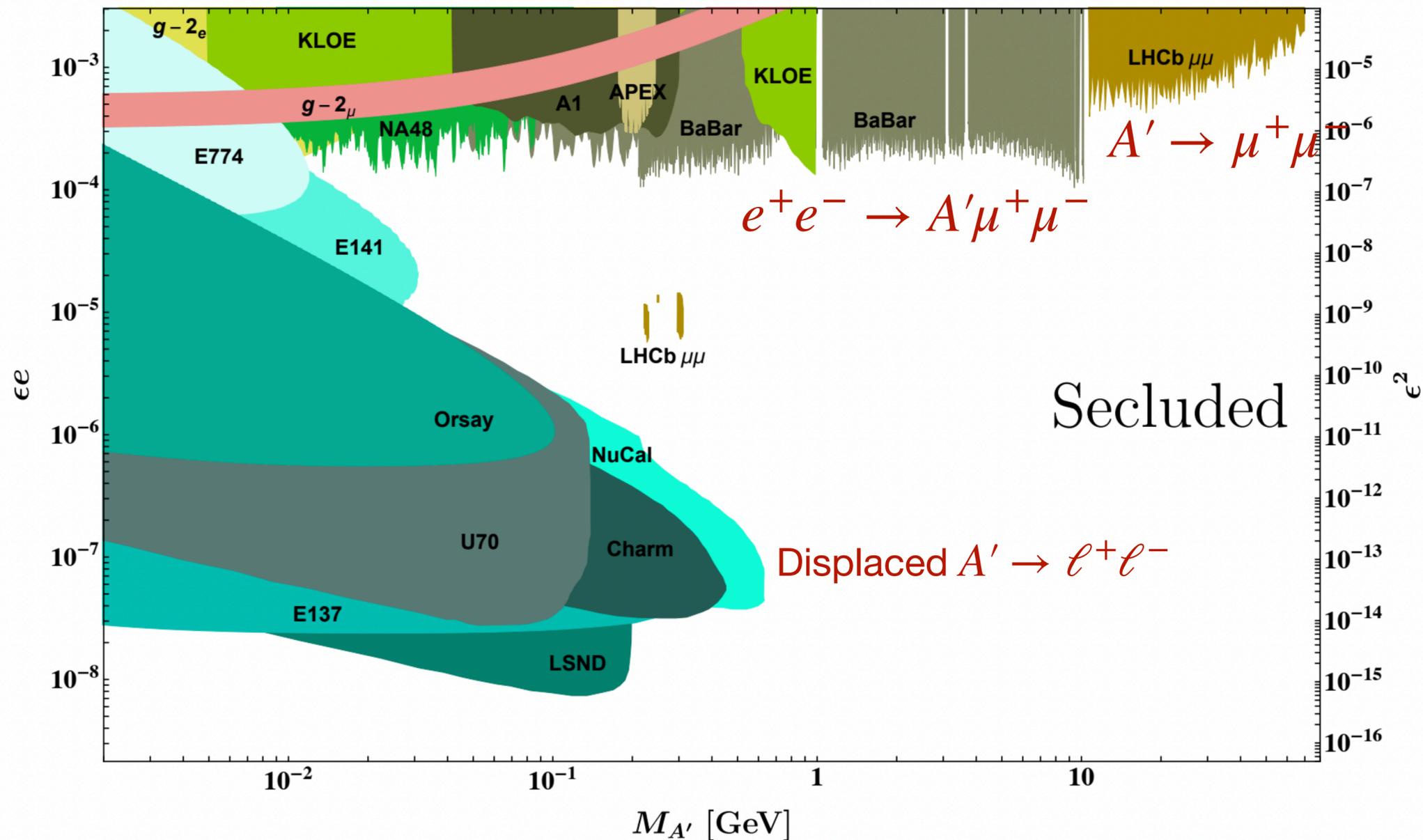
- 1.1 Secluded dark matter (dark sector)
 - Looking for mediator X is easier than DM

Dark photon A' example: visible

$$\epsilon F'_{\mu\nu} B^{\mu\nu} : A' \rightarrow \ell^+ \ell^-$$



Bauer et al: 1803.05466 (JHEP)



The way-out from direct detection limits

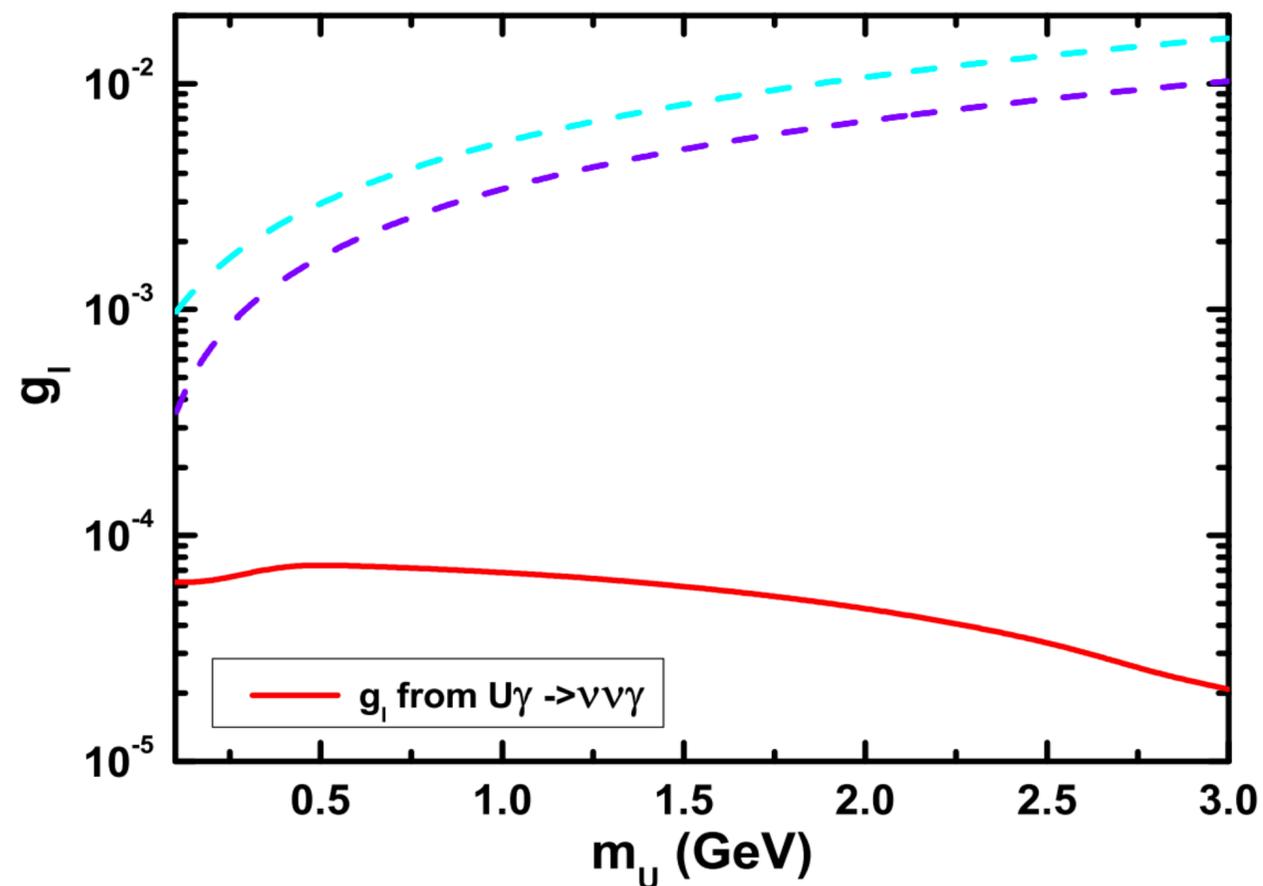
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Dark photon A' example: invisible

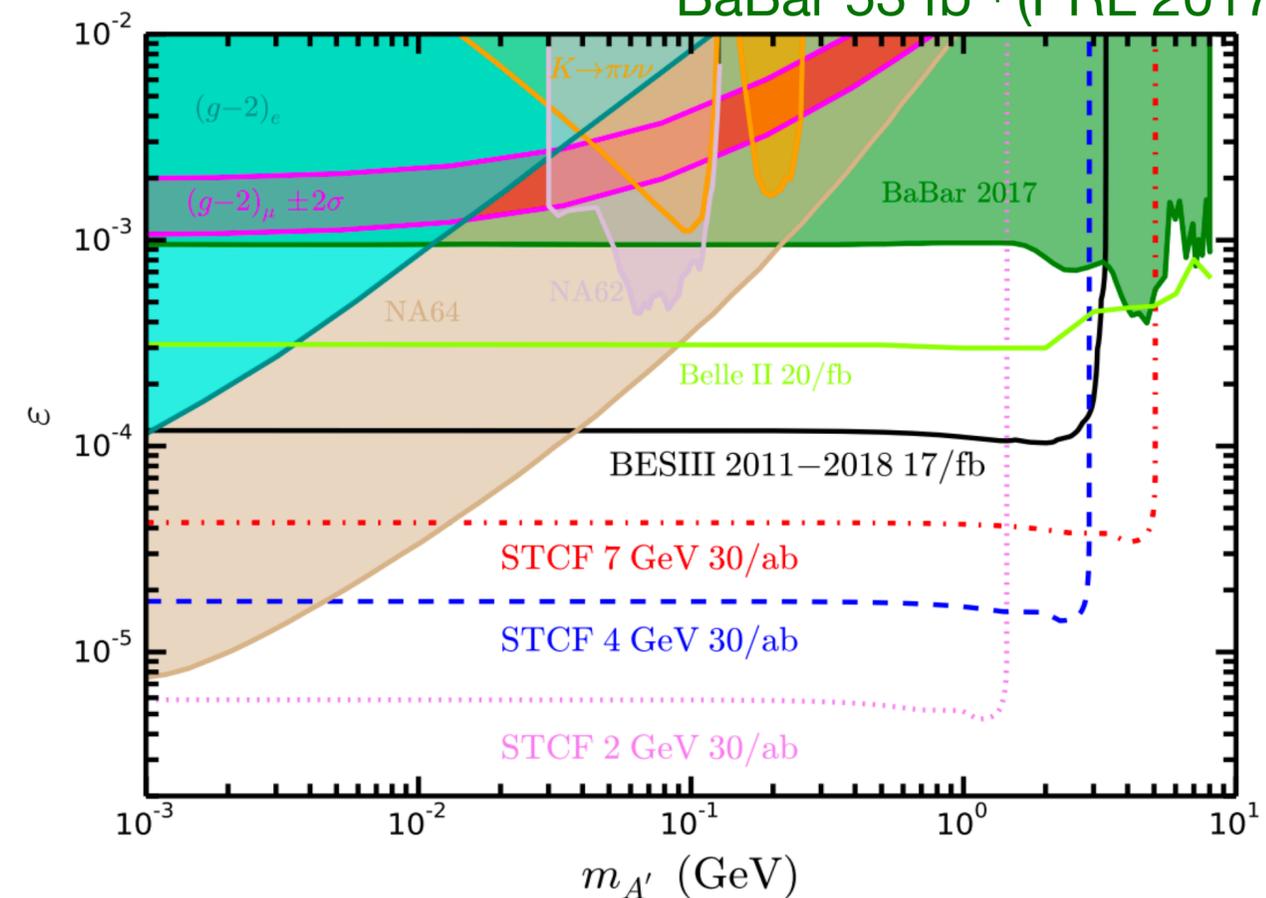
$$A' \rightarrow \bar{\nu}\nu, \bar{\chi}\chi$$

NA64: e beam dump, but for invisible final states

BaBar 53 fb⁻¹ (PRL 2017)



PF Yin, JL, SH Zhu: 0904.4644 (PRD)



BESIII: 1907.07046 (PRD)

The way-out from direct detection limits

- 2. Suppressed scattering cross-section:

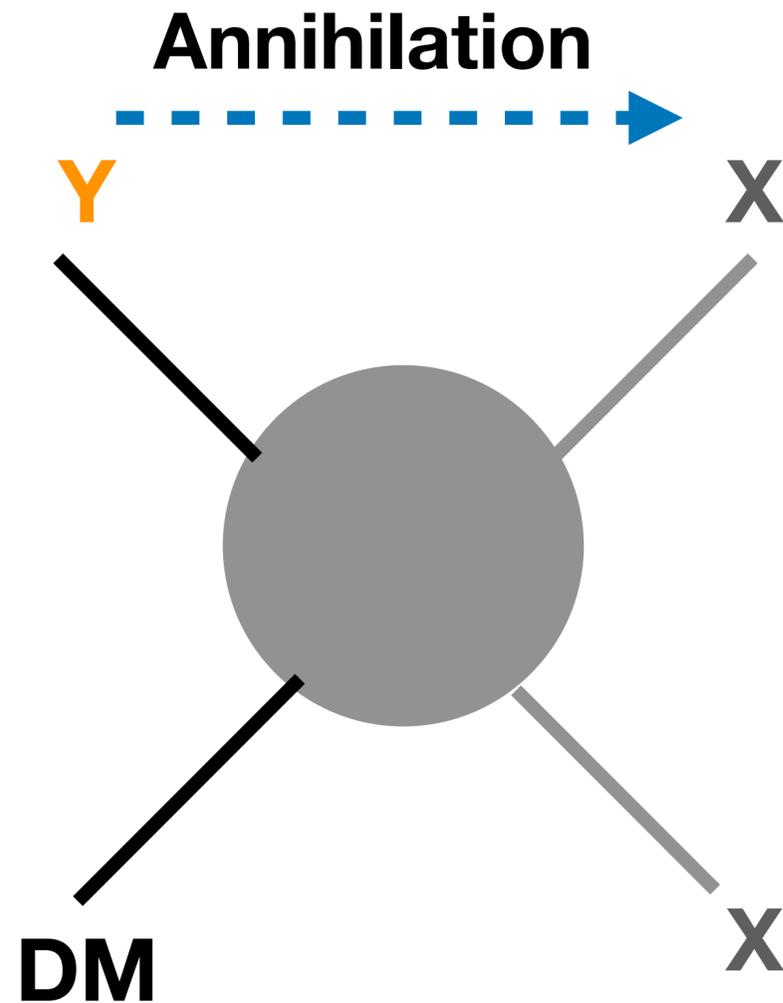
- By velocity or momentum transfer

Case for Fermionic DM
Kumar & Marfatia:1305.1611 (PRD)

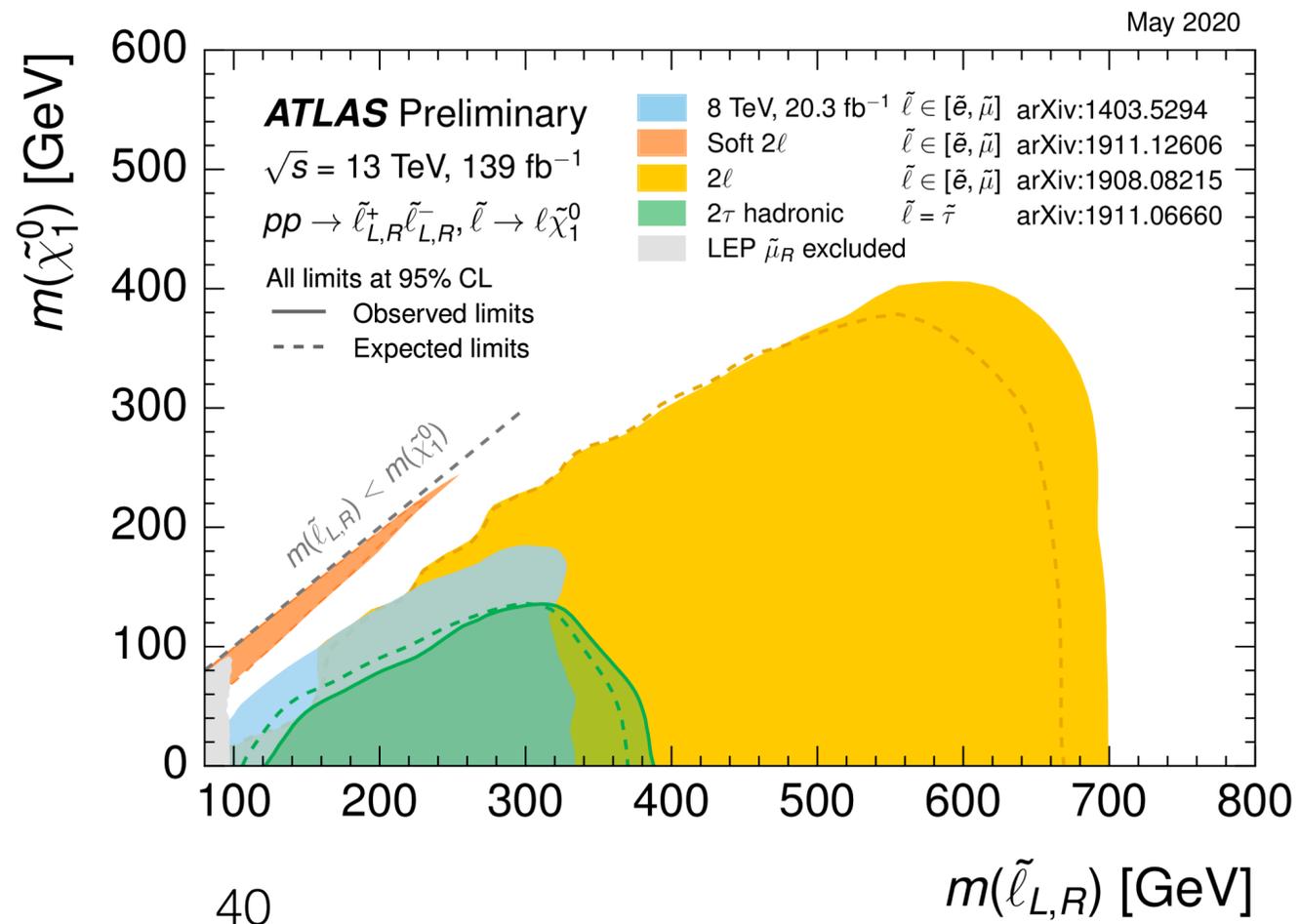
	Name	Interaction Structure	σ_{SI} suppression	σ_{SD} suppression	<i>s</i> -wave?
Scalar	F1	$\bar{X} X \bar{q} q$	1	$q^2 v^{\perp 2}$ (SM)	No
	F2	$\bar{X} \gamma^5 X \bar{q} q$	q^2 (DM)	$q^2 v^{\perp 2}$ (SM); q^2 (DM)	Yes
	F3	$\bar{X} X \bar{q} \gamma^5 q$	0	q^2 (SM)	No
Pseudoscalar	F4	$\bar{X} \gamma^5 X \bar{q} \gamma^5 q$	0	q^2 (SM); q^2 (DM)	Yes
Vector	F5	$\bar{X} \gamma^\mu X \bar{q} \gamma_\mu q$ (vanishes for Majorana X)	1	$q^2 v^{\perp 2}$ (SM) q^2 (SM); q^2 or $v^{\perp 2}$ (DM)	Yes
	Anapole	F6	$\bar{X} \gamma^\mu \gamma^5 X \bar{q} \gamma_\mu q$	$v^{\perp 2}$ (SM or DM)	q^2 (SM)
F7		$\bar{X} \gamma^\mu X \bar{q} \gamma_\mu \gamma^5 q$ (vanishes for Majorana X)	$q^2 v^{\perp 2}$ (SM); q^2 (DM)	$v^{\perp 2}$ (SM) $v^{\perp 2}$ or q^2 (DM)	Yes
F8		$\bar{X} \gamma^\mu \gamma^5 X \bar{q} \gamma_\mu \gamma^5 q$	$q^2 v^{\perp 2}$ (SM)	1	$\propto m_f^2 / m_X^2$
F9		$\bar{X} \sigma^{\mu\nu} X \bar{q} \sigma_{\mu\nu} q$ (vanishes for Majorana X)	q^2 (SM); q^2 or $v^{\perp 2}$ (DM) $q^2 v^{\perp 2}$ (SM)	1	Yes
	F10	$\bar{X} \sigma^{\mu\nu} \gamma^5 X \bar{q} \sigma_{\mu\nu} q$ (vanishes for Majorana X)	q^2 (SM)	$v^{\perp 2}$ (SM) q^2 or $v^{\perp 2}$ (DM)	Yes

The way-out from direct detection limits

- 3. Coannihilation mechanism

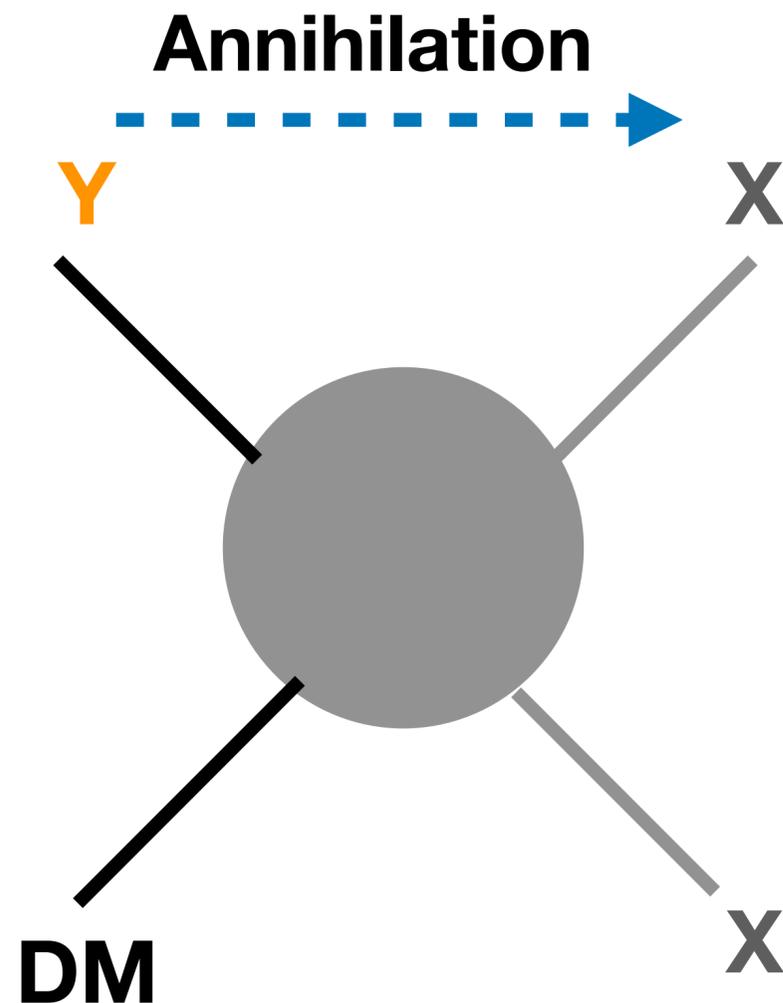


- Y has a close mass with DM
- Y is not populated today due to decay
- Charged Y: near degenerate spectrum of SUSY, AMSB; EW multiplet DM $(2n+1, 0)$ ($\delta m \sim 166$ MeV)



The way-out from direct detection limits

• 3. Coannihilation mechanism



- **Y** has a close mass with DM
- **Y** is not populated today due to decay
- Charged **Y**: near degenerate spectrum of SUSY, AMSB
- Neutral **Y**: Inelastic Dark Matter

Fermionic DM with kinetic mixing A' mediator

$$\mathcal{L} = \bar{\psi} i \gamma_{\mu} D^{\mu} \psi + m \bar{\psi} \psi + \delta \bar{\psi}^c \psi / 2$$

$$\bar{\psi} \gamma_{\mu} \psi \simeq i(\bar{\chi}_1 \bar{\sigma}_{\mu} \chi_2 - \bar{\chi}_2 \bar{\sigma}_{\mu} \chi_1) + \frac{\delta}{2m} (\bar{\chi}_2 \bar{\sigma}_{\mu} \chi_2 - \bar{\chi}_1 \bar{\sigma}_{\mu} \chi_1).$$

$$m_{\chi_1} = m - \delta; \quad m_{\chi_2} = m + \delta$$

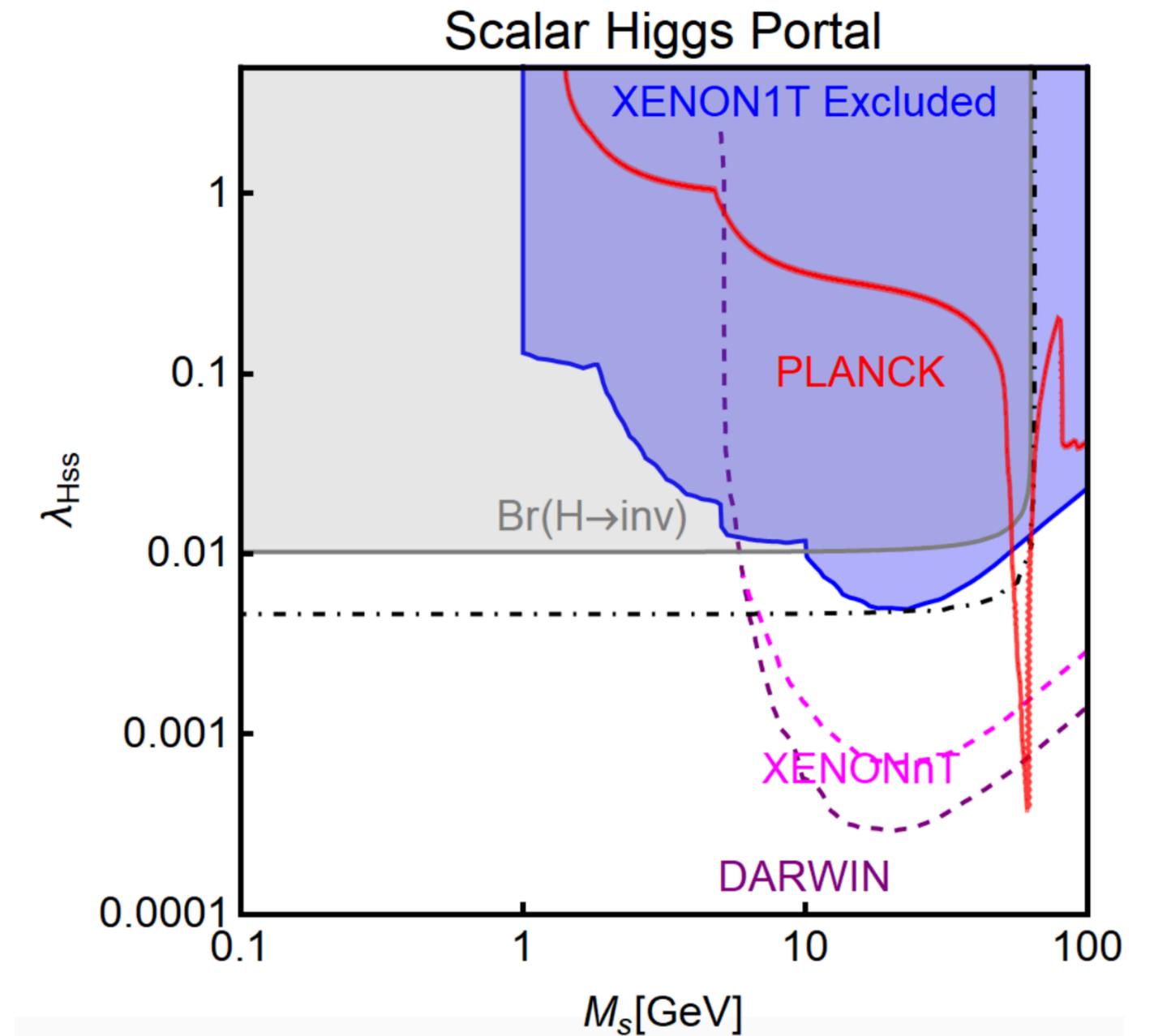
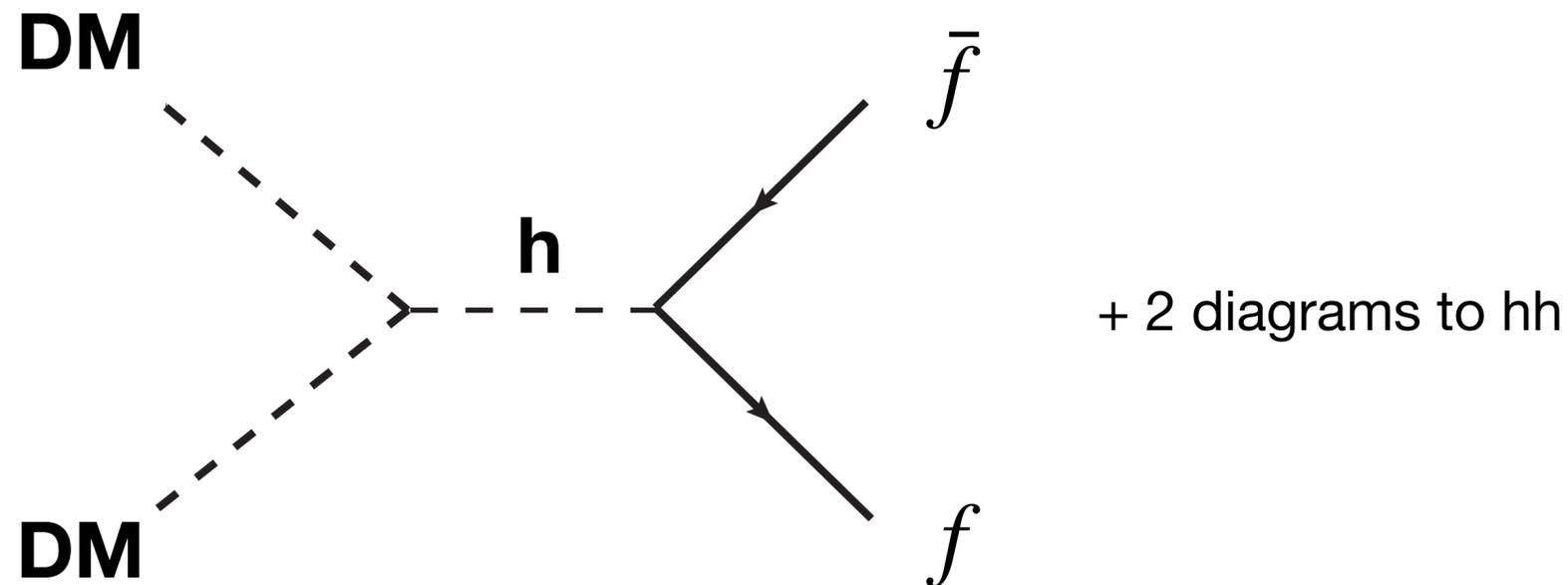
The way-out from direct detection limits

- 4. Resonant annihilation

- $2m_{\text{DM}} \approx m_X$

Scalar DM (s) with a Higgs portal coupling

$$\Delta\mathcal{L}_s = -\frac{1}{2}m_s^2 s^2 - \frac{1}{4}\lambda_s s^4 - \frac{1}{4}\lambda_{Hss}\phi^\dagger\phi s^2$$



Arcadi et al: 2101.02507

The way-out from direct detection limits

- 5. Cancellation effect in scattering cross-section

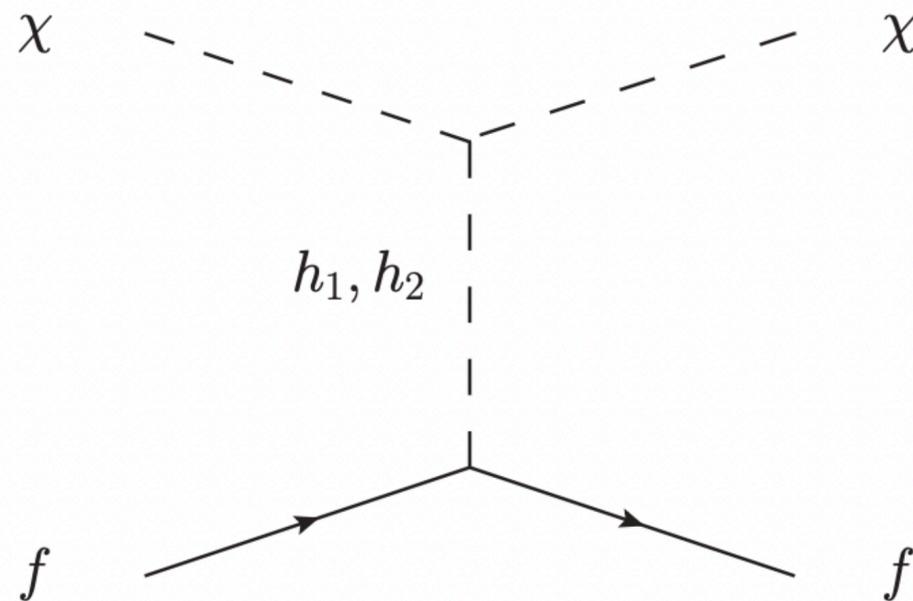
- SM Higgs - Dark scalar mediator cancellation

Gross, Lebedev1, Toma: 1708.02253 (PRL)

$$V_0 = -\frac{\mu_H^2}{2} |H|^2 - \frac{\mu_S^2}{2} |S|^2 + \frac{\lambda_H}{2} |H|^4 + \lambda_{HS} |H|^2 |S|^2 + \frac{\lambda_S}{2} |S|^4$$

$$V_{\text{soft}} = -\frac{\mu_S'^2}{4} S^2 + \text{h.c.} \quad \text{symmetry : } S \leftrightarrow S^*$$

$$S = (v_s + s + i\chi)/\sqrt{2} \quad \text{Pseudoscalar DM}$$



CP-even scalar mixing (s, h) → (h1, h2)

$$\mathcal{L} \supset -(h_1 \cos \theta + h_2 \sin \theta) \sum_f \frac{m_f}{v} \bar{f} f \quad \mathcal{L} \supset \frac{\chi^2}{2v_s} (m_{h_1}^2 \sin \theta h_1 - m_{h_2}^2 \cos \theta h_2)$$

$$\mathcal{A}_{dd}(t) \propto \sin \theta \cos \theta \left(\frac{m_{h_2}^2}{t - m_{h_2}^2} - \frac{m_{h_1}^2}{t - m_{h_1}^2} \right) \simeq \sin \theta \cos \theta \frac{t (m_{h_2}^2 - m_{h_1}^2)}{m_{h_1}^2 m_{h_2}^2} \simeq 0$$

See JL, XP Wang and F Yu 1704.00730 (JHEP), for cancellation between A' - Z boson in kinetic mixing dark photon model

The amplitude is suppressed by q² from pseudo-goldstone nature

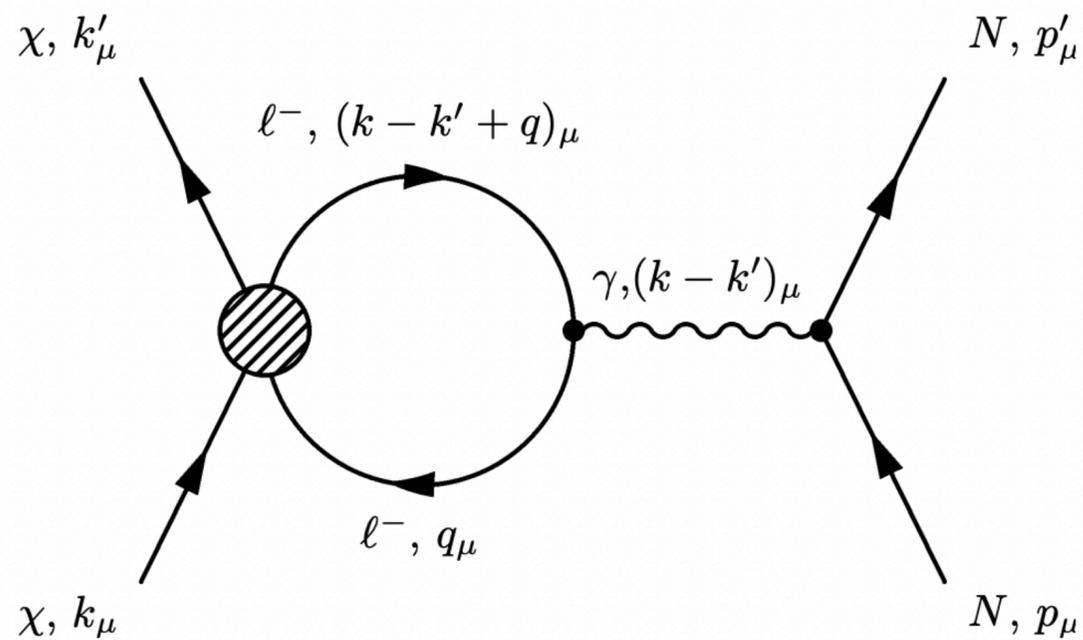
The way-out from direct detection limits

- 6. Leptophilic models

- Only couples to electrons, couples to nucleons at 1-loop

- For light DM, e-DM recoils can have stringent limits (e.g. XENON1T, PANDAX, CDEX)

- For heavy DM, nucleus-DM recoils wins over e-DM recoil



$$R^{\text{WAS}} : R^{\text{WES}} : R^{\text{WNS}} \sim \epsilon_{\text{WAS}} : \epsilon_{\text{WES}} \frac{m_e}{m_N} : \left(\frac{\alpha_{\text{em}} Z}{\pi} \right)^2 \sim 10^{-17} : 10^{-10} : 1$$

WAS = e kicked out

WES = e to higher energy level

WNS = nucleus recoil

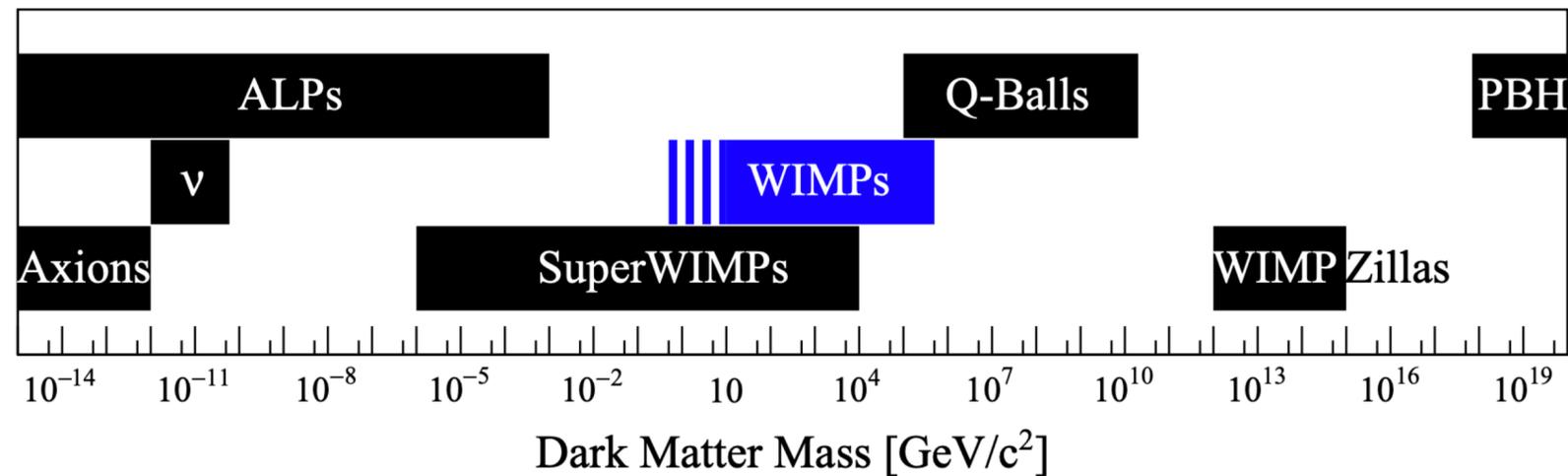
The probability to find a high p electron in the wave function is highly suppressed!

Kopp et al: 0907.3159 (PRD)

Outlines

- The dark matter in astrophysics perspective
- The dark matter in particle physics perspective
 - The WIMP crisis from direct detection
 - The DM limits from indirect detection
 - A WIMP variant from cosmological evolution
- Summary

The indirect detection limits from DM annihilation



- DM starts with thermal distribution
- DM has electroweak-scale coupling
- Relic abundance is determined by freeze-out mechanism
- DM Annihilation into
 - X = Standard Model particles (direct coupling)
 - X = Dark Sector particles (secluded DM models)



The entropy of DM goes into SM sector most of the time!
(Secluded X → SM + SM)

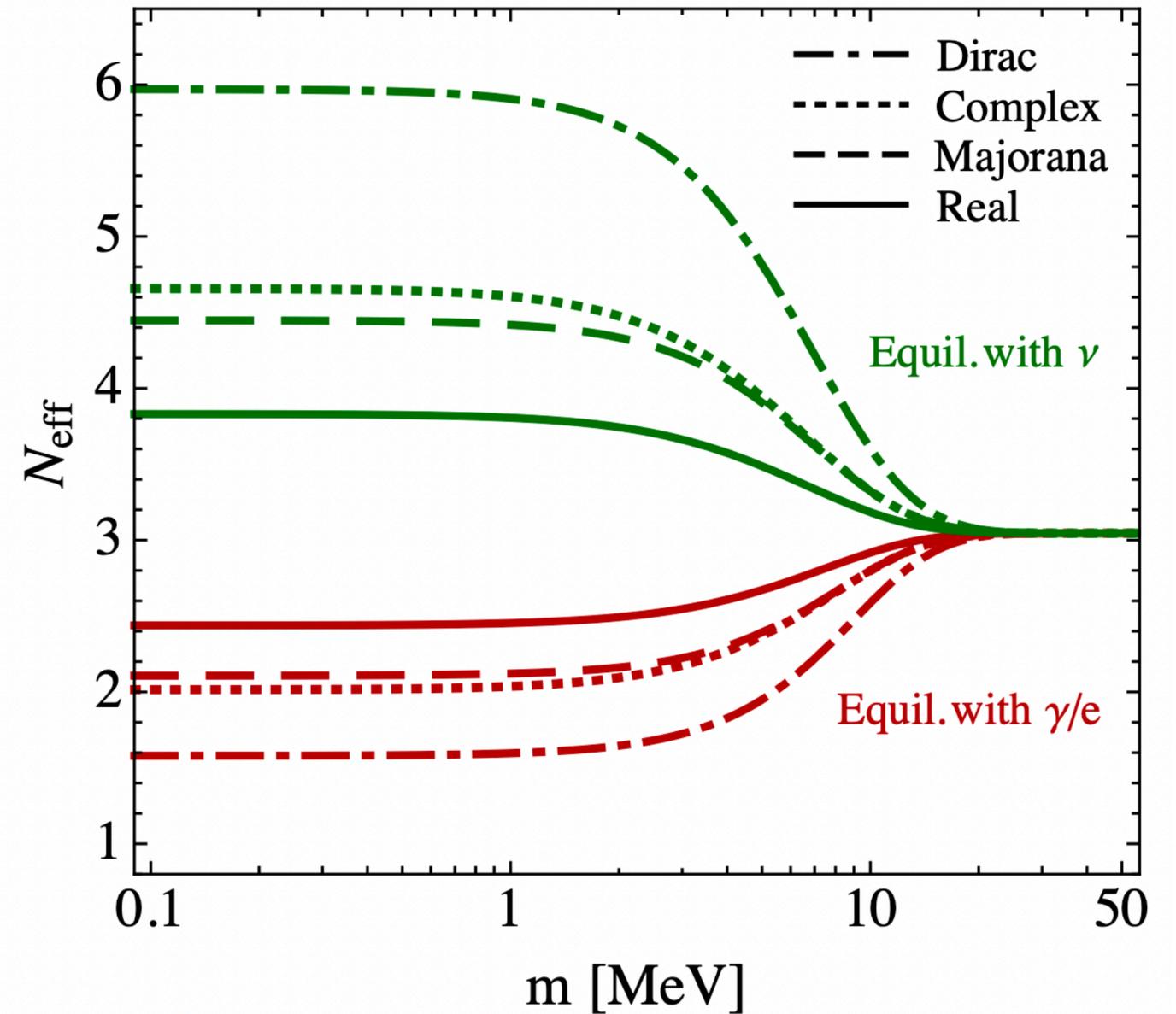
Lower mass bound for thermal DM

- Lower bound from N_{eff} at CMB
- Light DM freeze-out after neutrino decoupling at $T_D \approx 2.3 \text{ MeV}$
- Normally $T_{fo} \sim m_{\text{DM}}/20$
- DM entropy goes into neutrinos or e/γ , will modify T_ν/T_γ



Lower mass bound for thermal DM

- Lower bound from N_{eff} at CMB
- Light DM freeze-out after neutrino decoupling at $T_D \approx 2.3$ MeV
- Normally $T_{fo} \sim m_{\text{DM}}/20$
- DM entropy goes into neutrinos or e/γ , will modify T_ν/T_γ
- DM mass $\gtrsim 5$ MeV, depending on d.o.f.



Boehm et al: 1303.6270 (JCAP)

The annihilation cross-section expansion

- Expansion over velocity
 - S-wave
 - P-wave (L=1)
 - D-wave (L=2), due to extra chiral suppression

$$\sigma v \sim \sigma_s + \sigma_p v^2 + \sigma_d v^4 + \dots$$

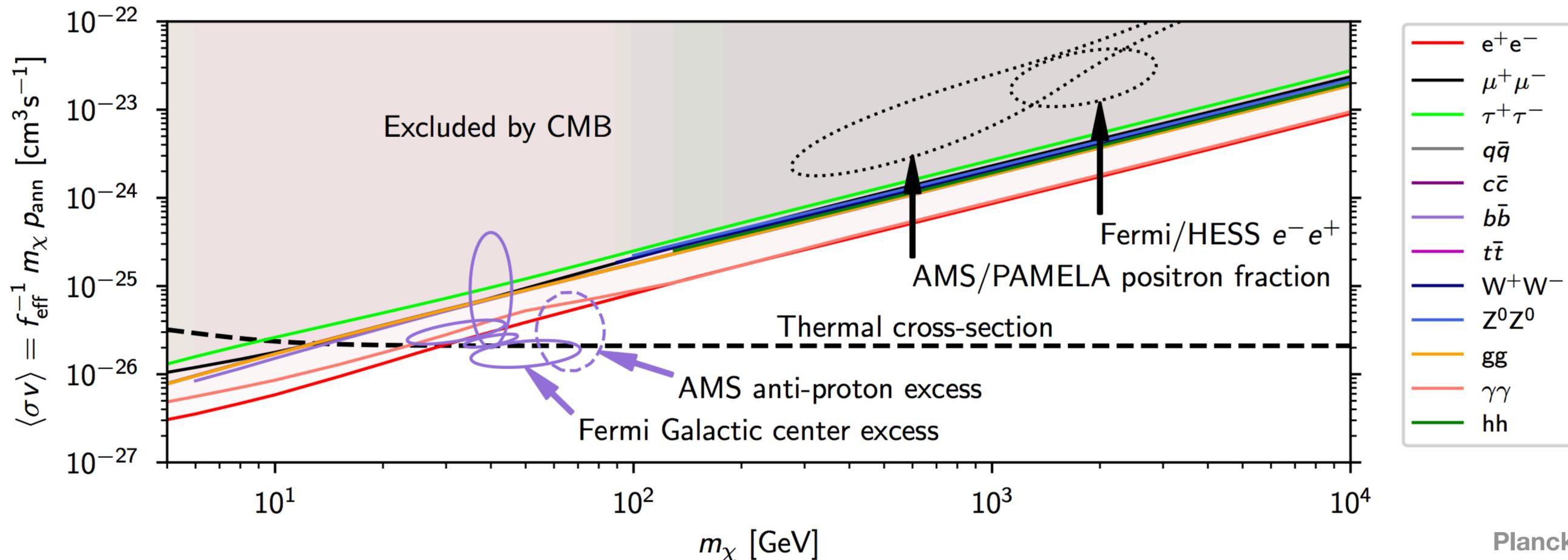
- The value of velocities at different time
 - Freeze-out: $v^2 \sim 0.25$
 - CMB: $v^2 \sim \text{eV}/m_{\text{DM}} \sim 10^{-5}$
 - Today: $v \sim 10^{-3}c$

Annihilation constraints from CMB

- The annihilation: $DM + DM \rightarrow SM + SM$
- The rate DM energy density converted into EM energy

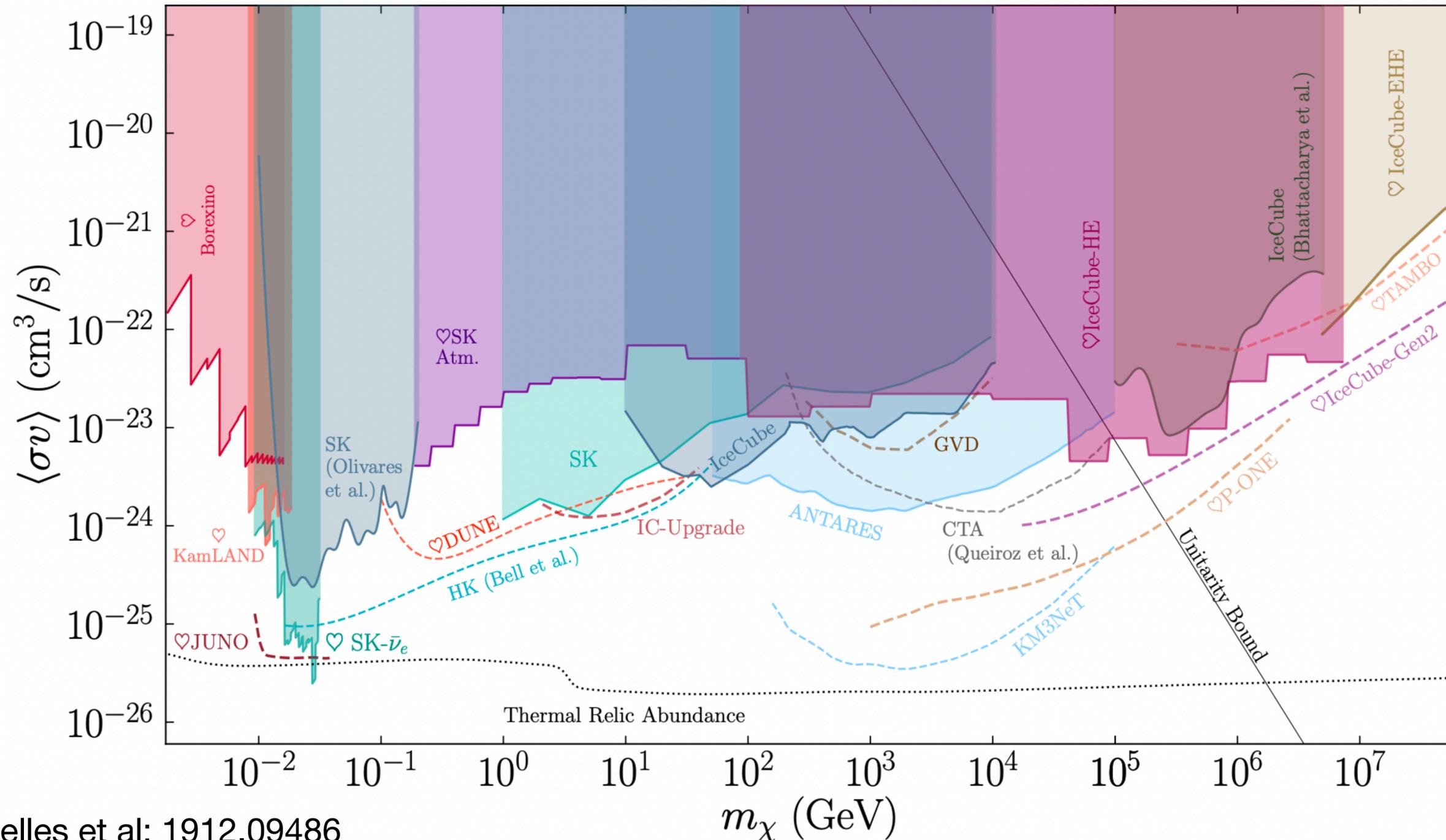
$$\frac{d\rho_{DM}}{dt} = m_{DM} n_{DM}^2 \langle \sigma v \rangle \times f_{\text{eff}}$$

- f_{eff} : the efficiency with which the energy released in DM annihilation is absorbed by the primordial plasma



How to escape CMB constraints?

- 1. Annihilation to neutrinos ($2\text{DM} \rightarrow \bar{\nu}\nu$): $f_{\text{eff}} = 0$



How to escape CMB constraints?

- 2. P-wave annihilation or no annihilation (asymmetric DM) but no indirect detection signal

- Expansion over velocity

$$\sigma v \sim \sigma_s + \sigma_p v^2 + \sigma_d v^4 + \dots$$

- S-wave

- P-wave (L=1)

- D-wave (L=2), due to extra chiral suppression

- Linear v dependence?

- Final state phase space suppression

$(m_{\text{DM}} \approx m_X)$ from symmetry reason

- The value of velocities at different time

- Freeze-out: $v^2 \sim 0.25$

- CMB: $v^2 \sim \text{eV}/m_{\text{DM}} \sim 10^{-5}$

- Today: $v \sim 10^{-3}c$

How to escape CMB constraints?

- 2+. Linear v suppression
 - How about cross-section linear in v ? ($\sigma v \propto v$)
 - For CMB, linear v is enough to be safe
 - For indirect detection
 - Cluster, $v \sim 1000$ km/s $\sim 3 \times 10^{-3}$
 - Galaxy, $v \sim 220$ km/s $\sim 1 \times 10^{-3}$
 - Dwarfs, $v \sim 10$ km/s $\sim 3 \times 10^{-5}$
 - Detectable in Cluster and Galaxy, not in Dwarfs

Linear v to escape CMB limits

- Cross-section linear in v

$$\text{DM} + \text{DM} \rightarrow \text{X} + \text{X} \quad \langle \sigma v \rangle = \frac{1}{4m_{\text{DM}}^2} \int dPS_2 |\mathbf{M}|^2$$

- If $m_{\text{MD}} = m_{\text{X}}$, then the two-body phase space

$$\int dPS_2 = \frac{1}{8\pi} v$$

- For s-wave annihilation, this gives

$$\langle \sigma v \rangle \approx \frac{1}{2} \sigma_0 v$$

- In practice, not exact degenerate

$$\Delta = m_{\text{DM}} - m_{\text{X}}$$

$$\langle \sigma v_{\text{rel}} \rangle \simeq \sigma_0 \sqrt{\frac{v_{\text{rel}}^2}{4} + \frac{2\Delta}{m_{\text{DM}}}}$$

- Model building for $\Delta \ll m_{\text{DM}}$

- Symmetry reason

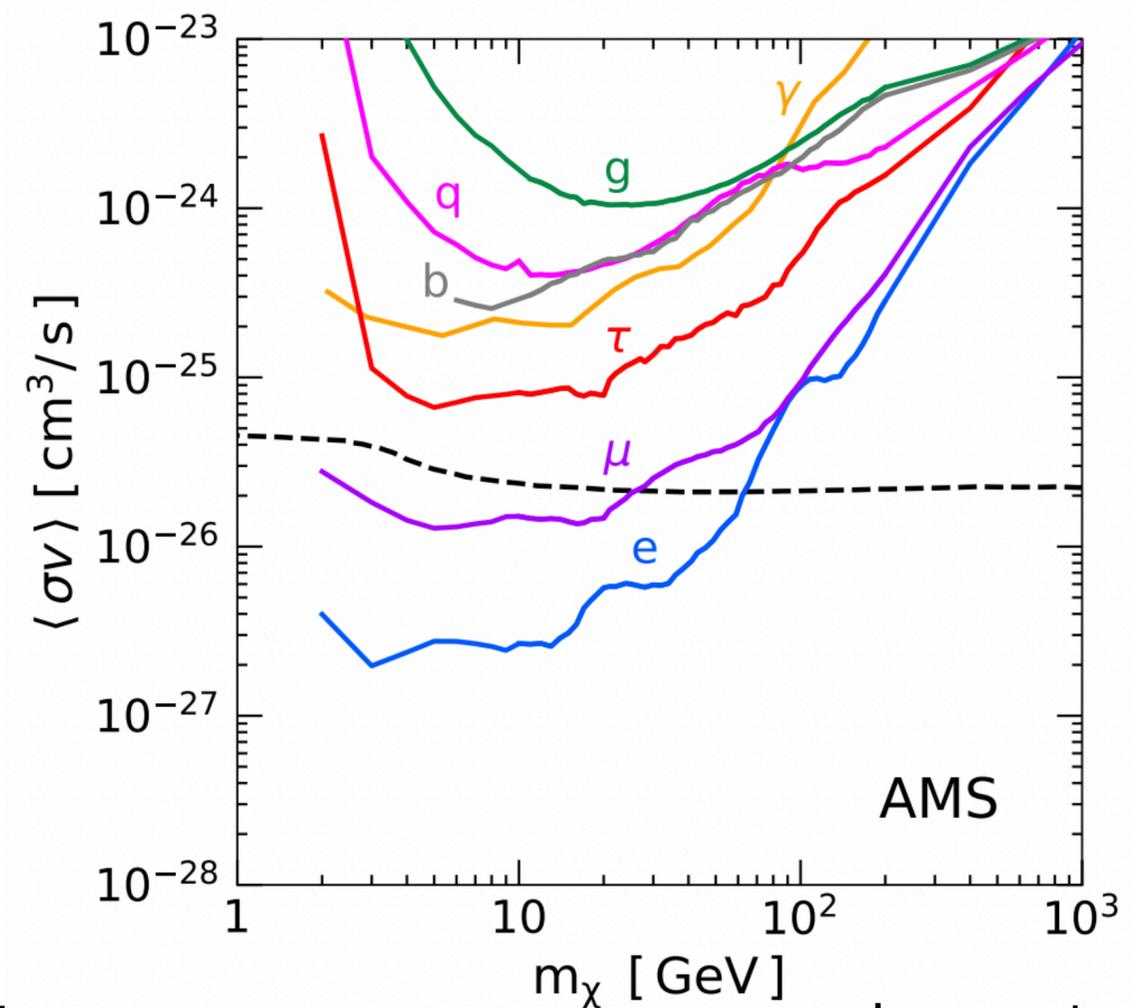
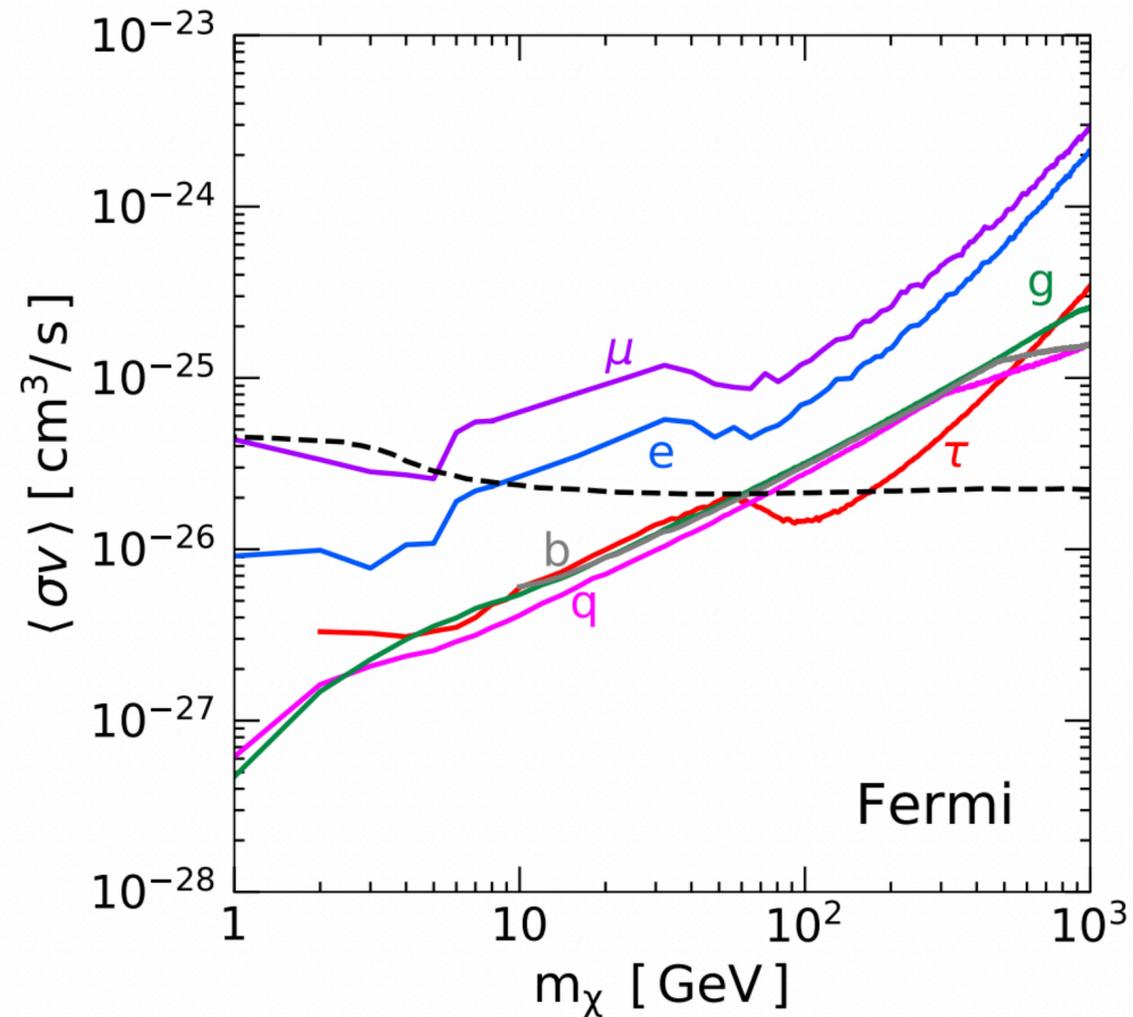
- Custodial symmetry: dark SU(2) vector DM $\Delta < 0$

- Chiral symmetry: dark pion DM $\Delta > 0$

- Supersymmetry: NMSSM setup 1901.02018

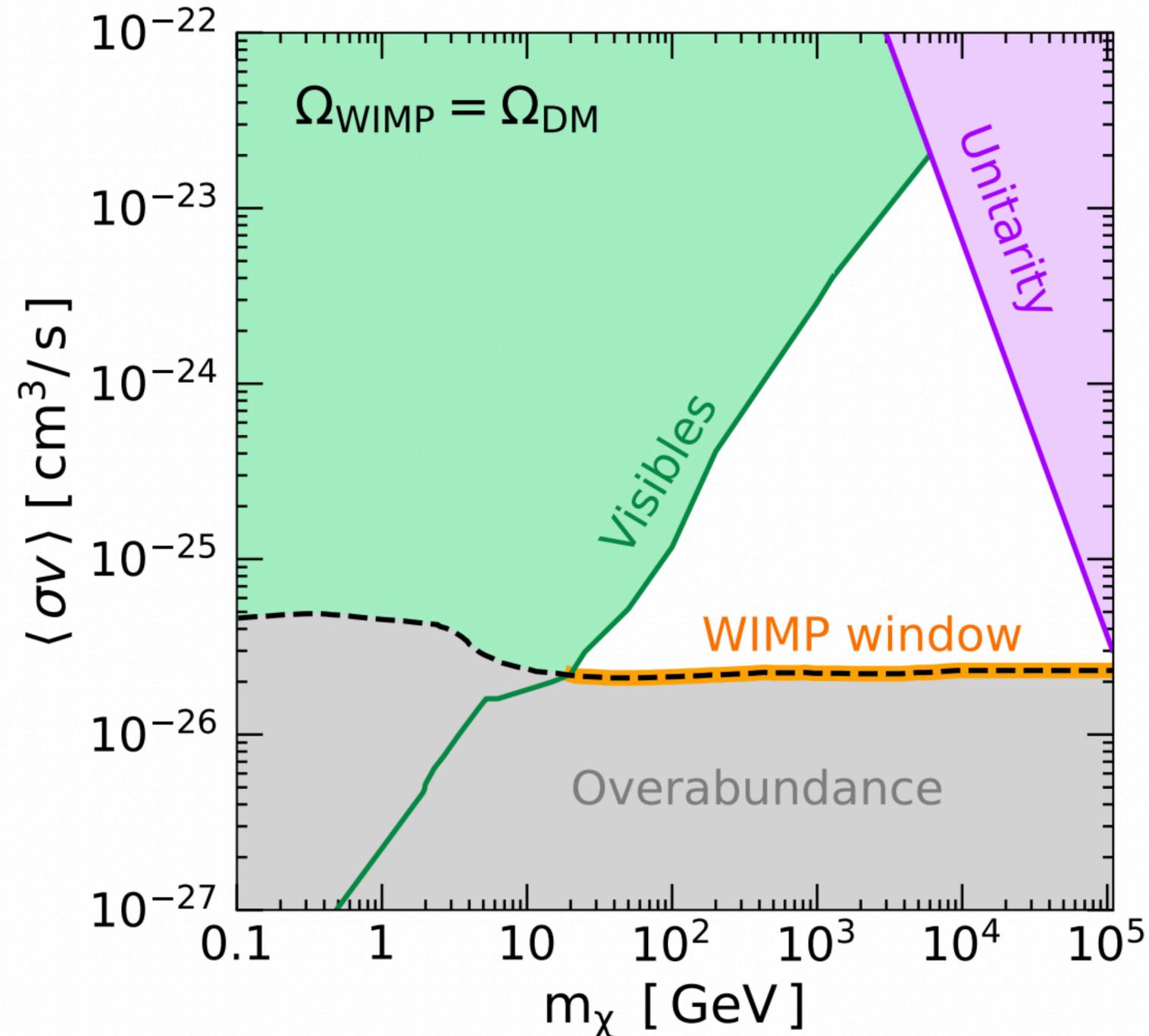
Other indirect limits

- CMB limits only works for DM mass $\lesssim 10$ GeV
- Indirect limits from AMS-02, DAMPE, Fermi-LAT



The WIMP limits from indirect detection

- WIMP mass $\gtrsim 10$ GeV is still viable



GeV-Scale Thermal WIMPs: Not Even Slightly Dead

Leane et al: 1805.10305 (PRD)

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 - **A WIMP variant from cosmological evolution**
- **Summary**

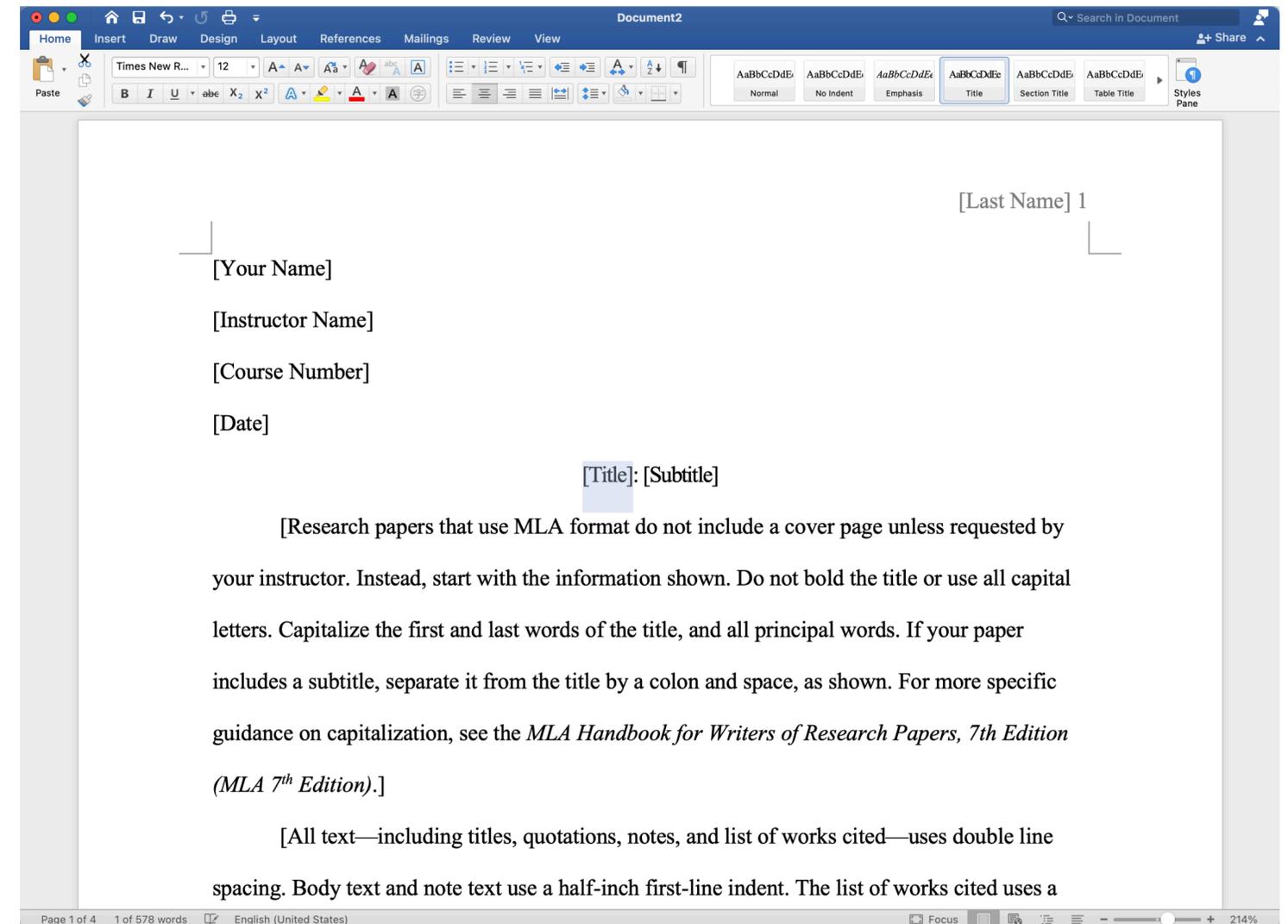
Dark matter transient annihilations in the early Universe
Katsuya Hashino, Jia Liu, Xiao-Ping Wang, and Ke-Pan Xie
ArXiv: 2109.07479

DM properties and cosmological evolution

- DM evolution can be deeply affected by the thermal history of the Universe
- DM properties at freeze-out may be different from today
- DM mass, stability, interaction couplings, decay and annihilation channels, rates

T. Cohen et al, 0808.3994
M. Baker, J. Kopp et al, 1608.07578, 1712.03962, 1811.03101
Kobakhidze and Schmidt et al, 1712.05170, 1910.01433
Hektor et al, 1801.06184
L. Bian and Y.L. Tang, 1810.03172
L. Bian and X. Liu, 1811.03279
L. Heurtier et al, 1912.02828
H. Murayama et al, 2012.15284
B. Batell et al, 2109.04476

...



Word, WPS 等办公软件符合“所见即所得”
WYSIWYG, “What You See Is What You Get”

A WIMP variant: DM with transient annihilations

- 1. Massive gauge boson has a varying mass in the early universe
- 2. If it is the DM-SM mediator, and the mass variation happens near DM freeze-out, what happens?

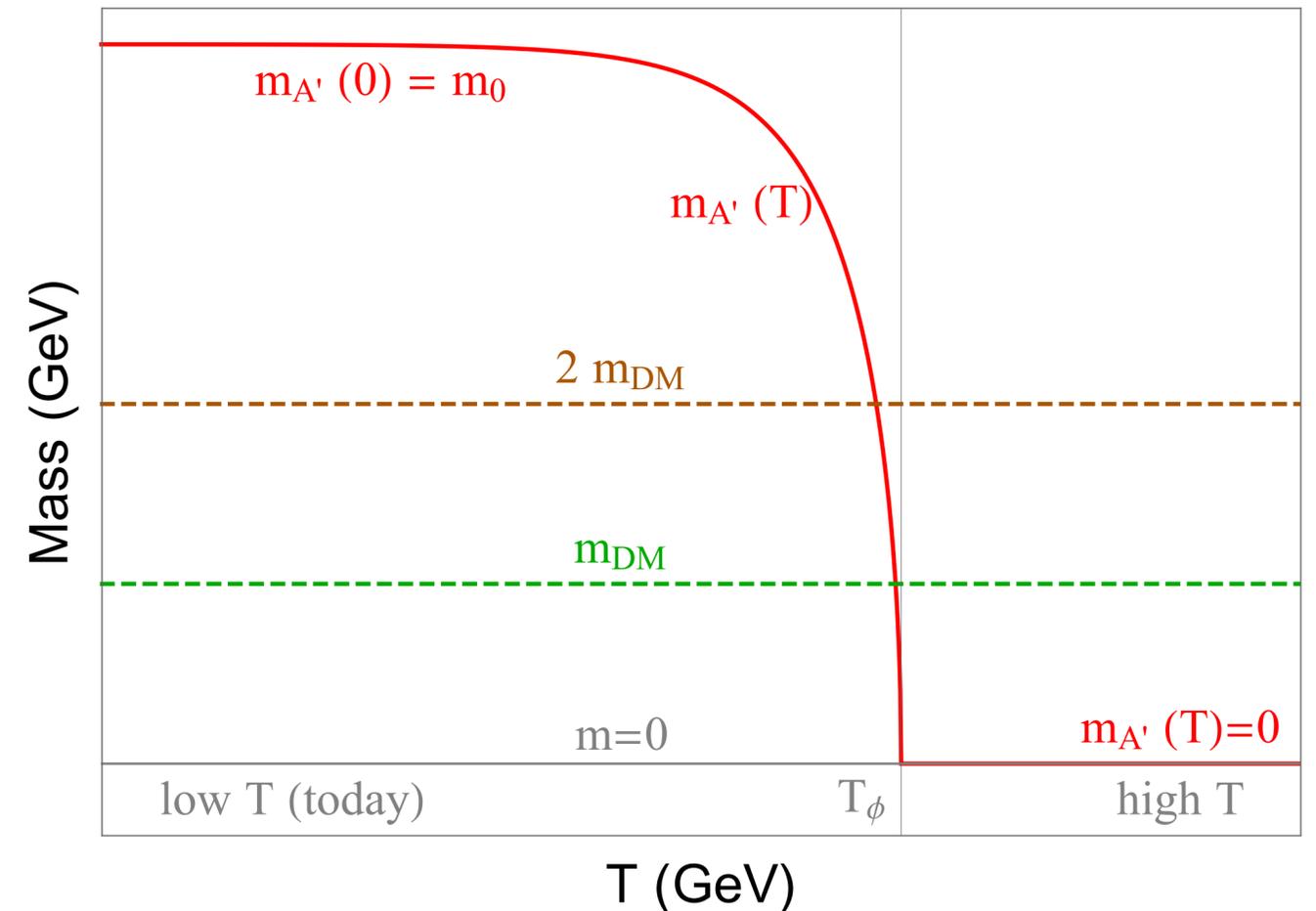
$$m_{A'}^2(T) = \begin{cases} 0 & T > T_\phi, \\ m_{A',0}^2 - \kappa m_\psi^2 \left(\frac{T}{m_\psi}\right)^n & T < T_\phi \end{cases}$$

$$\mathcal{L}_d = \bar{\psi} (i\not{D} - m_\psi) \psi - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + \epsilon e A'_\mu J_{\text{em}}^\mu$$

Transient secluded: $(\bar{\psi}\psi \rightarrow A'A')$ $m_{A'} = m_\psi,$

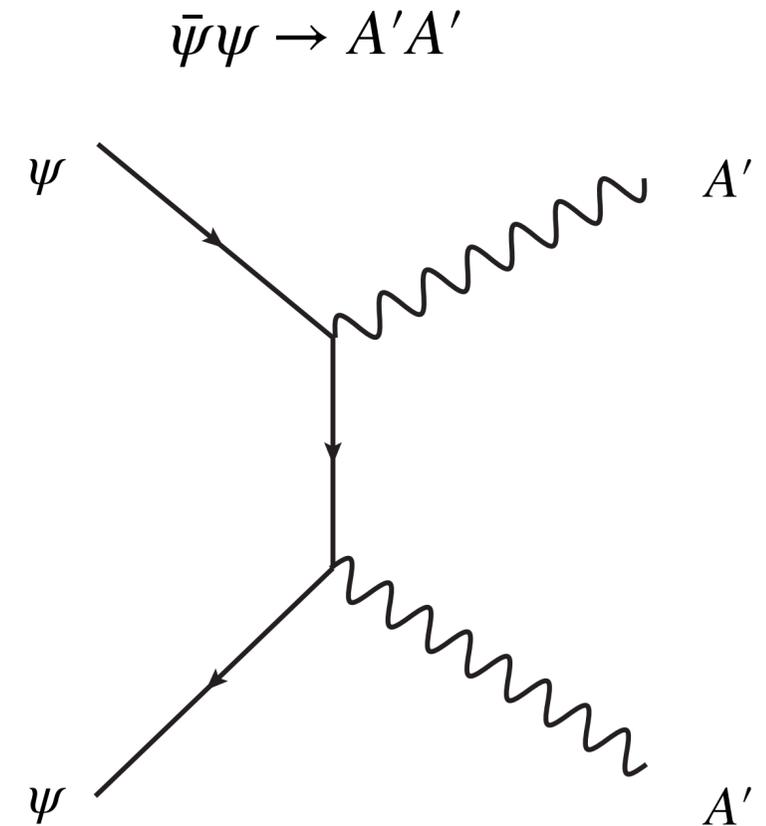
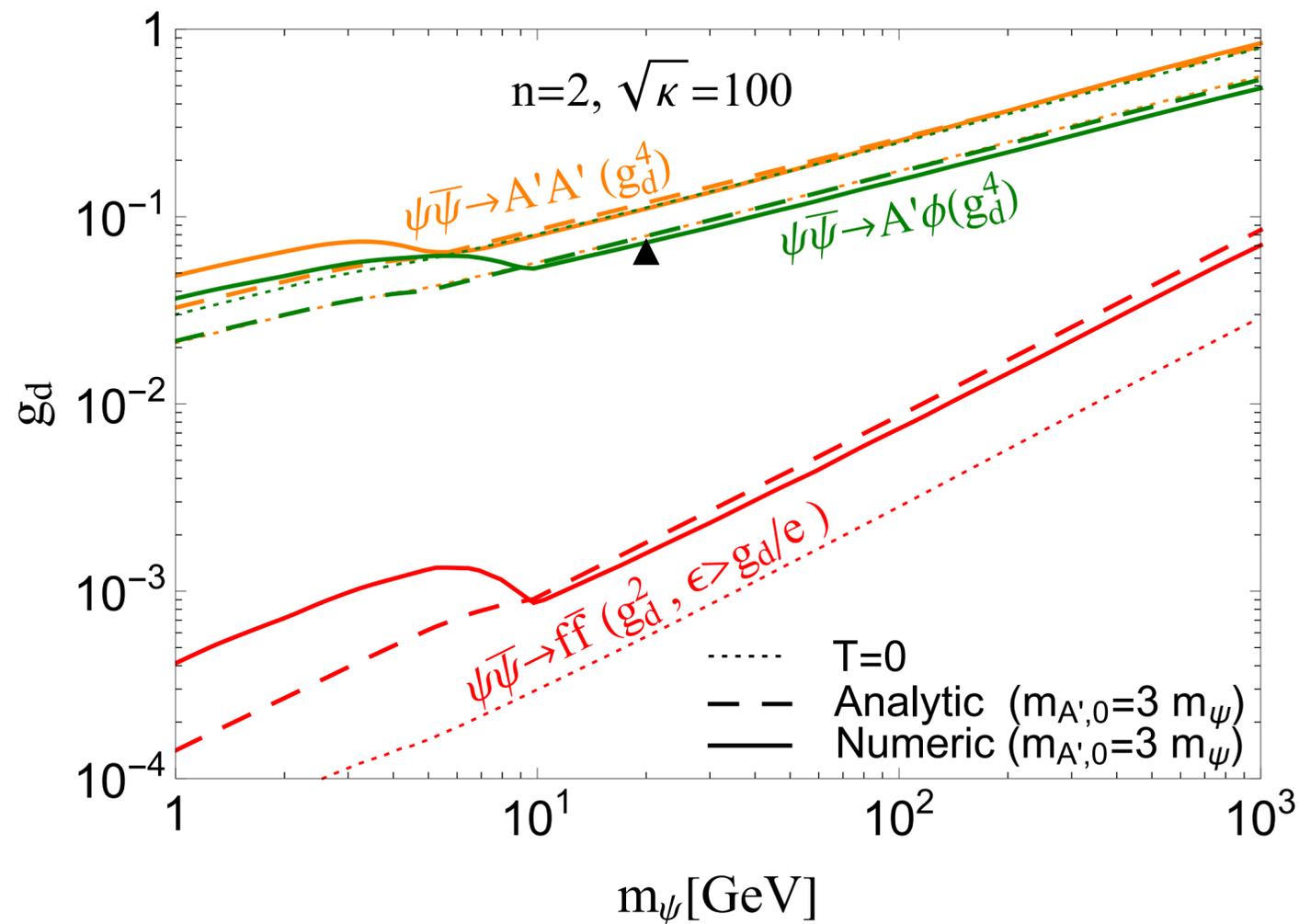
$(\bar{\psi}\psi \rightarrow A'\phi)$ $m_{A'} = 2m_\psi - m_\phi,$

Transient resonant: $(\bar{\psi}\psi \rightarrow \bar{f}f)$ $m_{A'} = 2m_\psi.$



Features for DM with transient annihilations

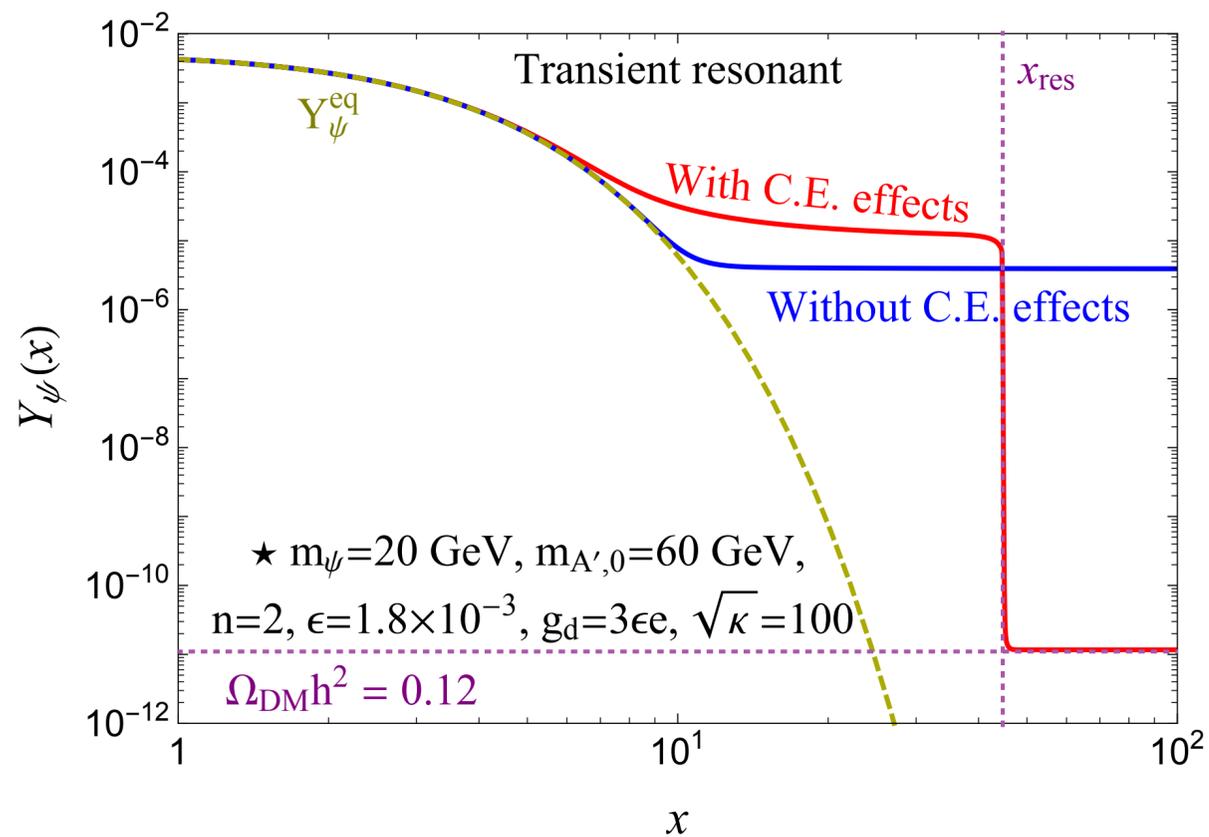
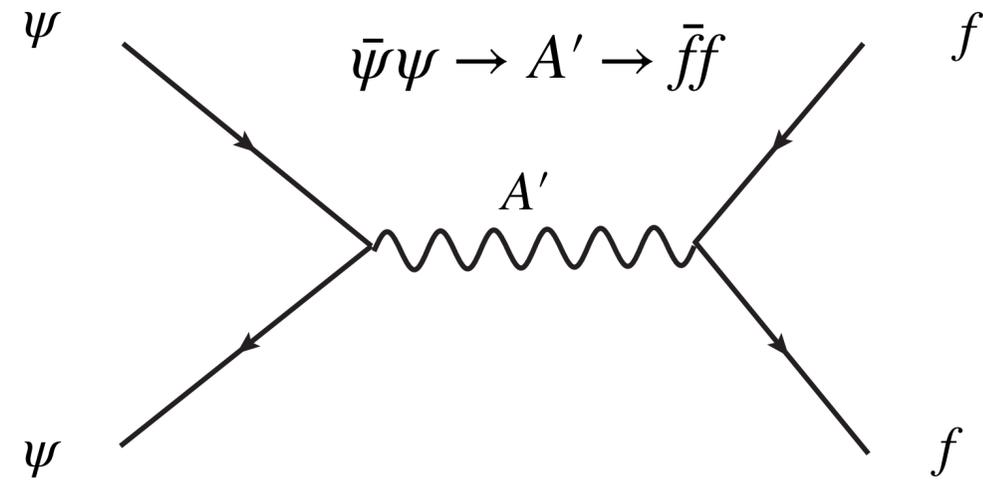
- Transient secluded annihilation only happens in the early universe
- is forbidden today
- No indirect constraints



$$\langle \sigma v_{A'A'} \rangle \approx \frac{g_d^4}{16\pi m_\psi^2} (1-r^2)^{3/2} (1-r^2/2)^{-2}$$

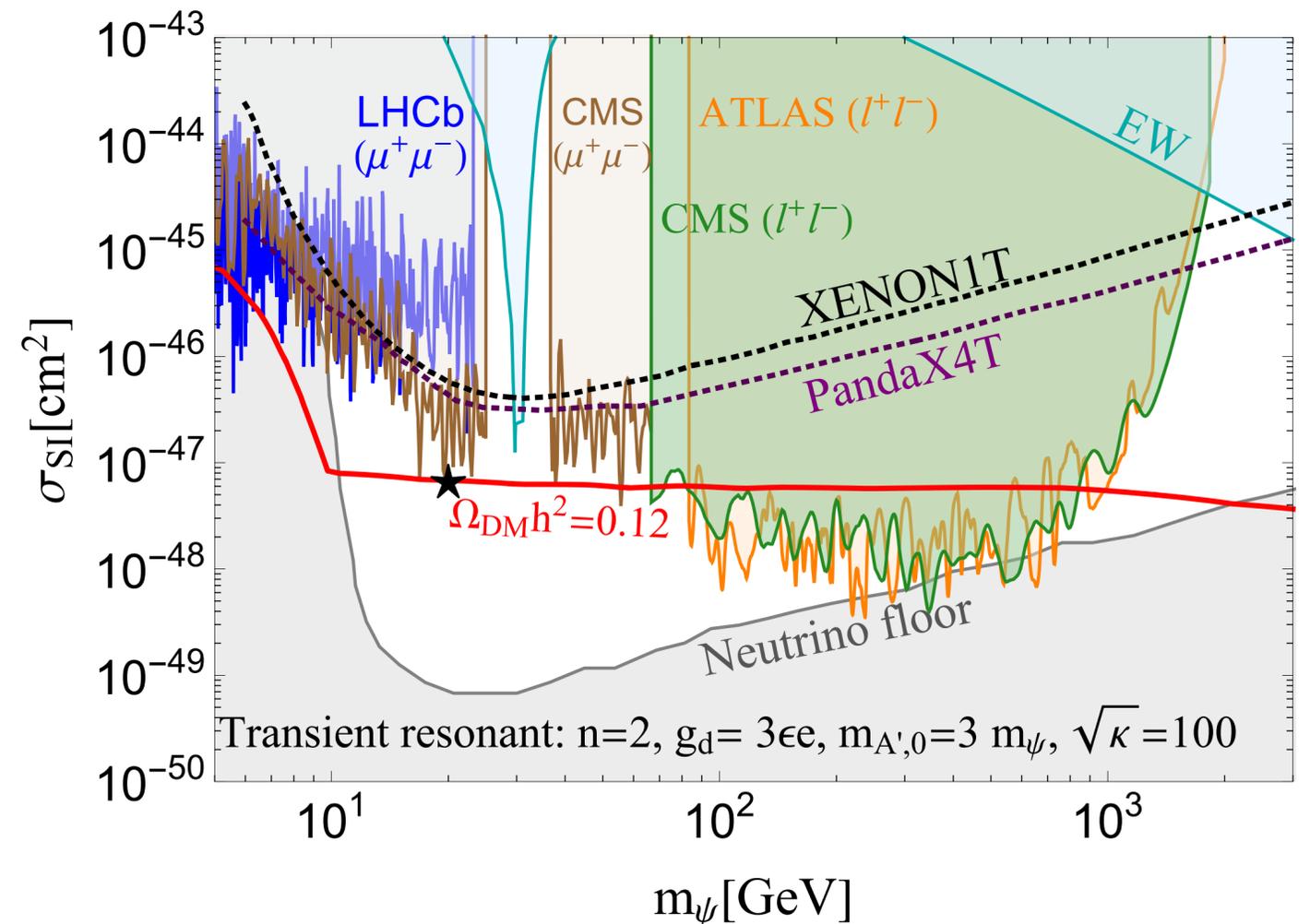
$$r \equiv m_{A'}/m_\psi$$

Features for DM with transient annihilations



$$\langle \sigma v \rangle_{\bar{f}f}^{\text{res}} \approx \frac{g_d^2 \epsilon^2 e^2 (2 + r^2) x}{48 \sqrt{2\pi} m_\psi \Gamma_{A'}} \sqrt{r(r^2 - 4) x} e^{-(r-2)x}$$

- Transient resonant annihilation only happens in the early universe
- No indirect constraints
- Collider and direct detection constraints are evaded
- Can be soon tested in the future



Summary

- WIMP DM has significant coupling to SM model
- Direct detection sets strong limits, but there are at least six ways to escape the limits
- Indirect detection sets strong limits, less way to escape the limits.
But it leaves open for DM mass $\gtrsim 10$ GeV
- GeV-Scale Thermal WIMPs: Not Even Slightly Dead
- A variant of WIMP model from cosmic evolution: transient annihilation DM, evading DD, collider and indirect searches but can be tested soon

Thank you!

Backup slides