Synchrotron sideband spin resonances centered on an integer

Xia Wenhao & Duan Zhe 2021. 12. 06

CEPC 中的一阶自旋共振

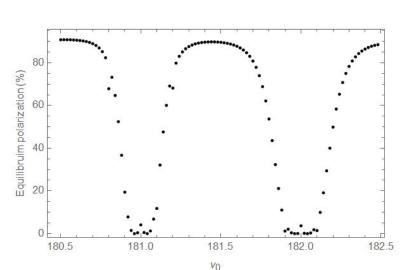
- 含磁铁误差,以**整数自旋共振及其纵向边带共振**为主;
- SLIM模拟结果, Δv₀=0.02;
- 理论解释:

Equilibruim polarization (%) 00 00 00

102.5

$$P \approx \frac{P_{\infty}}{1 + \frac{\tau_{sk}}{\tau_{sk}}}$$

$$P_{\infty} = \frac{-8}{5\sqrt{3}} \frac{\oint ds \left\langle \frac{1}{|\rho|^3} \widehat{b} \cdot (\widehat{n} - \gamma \frac{\partial \widehat{n}}{\partial \gamma}) \right\rangle}{\oint ds \left\langle \frac{1}{|\rho|^3} \right\rangle} \approx 92\%$$



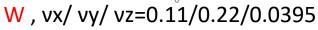
103.5

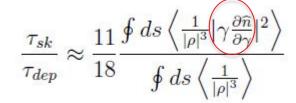
103.0

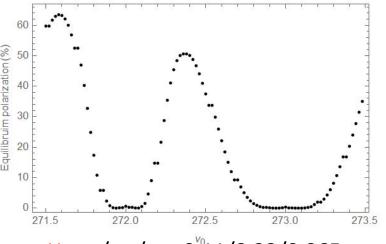
•••

104.0

104.5







H, vx/ vy/ vz=0.11/0.22/0.065

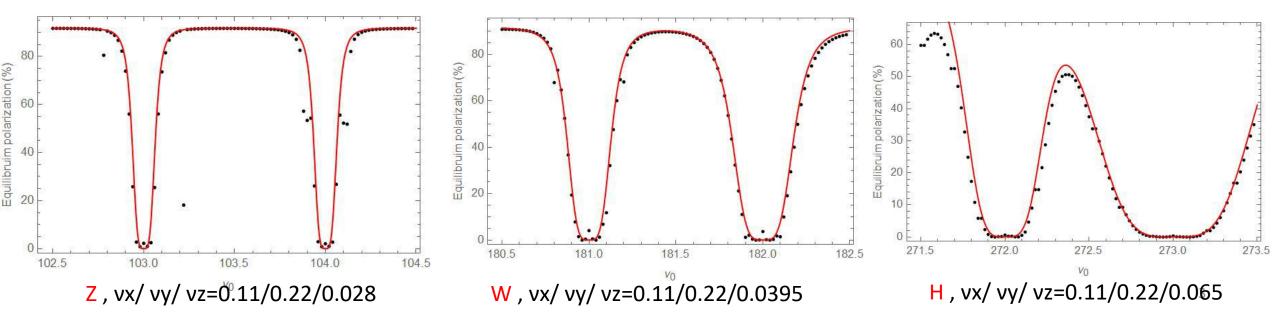
理论解释1

Derbenev 和Kondratenko的理论,拟合:

Ya.S. Derbenev, A.M. Kondratenko and A.N. Skrinsky, Par. Ace. 1979, Vol. 9, pp. 247-266.

$$\frac{\tau_{sk}}{\tau_{dep}} = \frac{11v^2}{18} \sum_{k} \frac{|\omega_k|^2}{(\nu - k)^4}$$

- 可以实现上升沿和下降沿的匹配;
- 但是,无法解释整数共振处的"凸起";



理论解释2

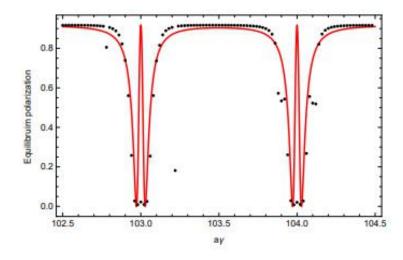
• Yokoya 的理论,拟合:

$$\frac{\tau_{sk}}{\tau_{dep}} \propto |b_{k,s}|^2 \frac{(\nu - k)^2}{((\nu - k)^2 - \nu_s^2)^2}$$

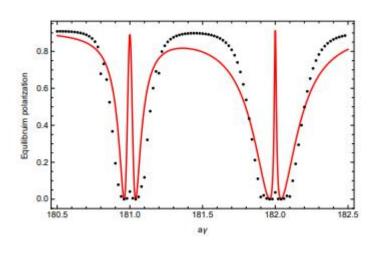
- 整数共振处有"凸起",但很大;
- 但是,无法实现上升沿,下降沿的匹配;

K. Yokoya, KEK Intenal 85-7, (1985).

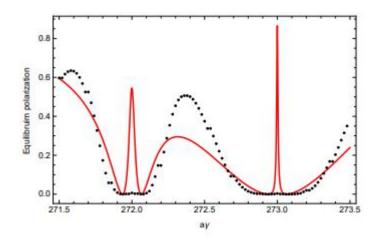
猜想:边带共振强度|b_{k,s}|^2,在整数自旋共振附近时,会随着离整数共振之间的距离发生变化,有没有可能消除掉分子项(nu-k)^2。



 $\frac{Z}{v}$, vx/ vy/ vz=0.11/0.22/0.028



W, vx/ vy/ vz=0.11/0.22/0.0395



H, vx/ vy/ vz=0.11/0.22/0.065

理论推导: 理想水平储存环

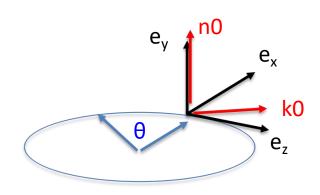
$$P \approx \frac{P_{\infty}}{1 + \frac{\tau_{sk}}{\tau_{dep}}} \qquad P_{\infty} = \frac{-8}{5\sqrt{3}} \frac{\oint ds \left\langle \frac{1}{|\rho|^3} \hat{b} \cdot (\hat{n} - \gamma \frac{\partial \hat{n}}{\partial \gamma}) \right\rangle}{\oint ds \left\langle \frac{1}{|\rho|^3} \right\rangle} \approx 92\% \qquad \frac{\tau_{sk}}{\tau_{dep}} \approx \frac{11}{18} \frac{\oint ds \left\langle \frac{1}{|\rho|^3} |\gamma \frac{\partial \hat{n}}{\partial \gamma}|^2 \right\rangle}{\oint ds \left\langle \frac{1}{|\rho|^3} \right\rangle}$$

- 关键在于自旋轨道耦合函数: $\gamma \frac{\partial \hat{\mathbf{n}}}{\partial \gamma}$
- 在理想水平储存环中: $\frac{d\vec{S}}{d\theta} = \left[\vec{\Omega}_0 + \vec{\omega}\right] \times \vec{S}$, Ω_0 为设计轨道上的自旋进动矢量,w来自于betatron+synchrotron
- w=0时,三个单位正交解n₀, m₀, l₀, 引入k₀=m₀+il₀;
- Yokoya给出:

$$\gamma \frac{\partial \vec{n}}{\partial \gamma}(\theta) = \frac{1}{2} Re[\vec{k}_0^{\star}(\theta) \cdot (D_x + D_{-x} + D_y + D_{-y} + D_s + D_{-s})]$$

• 我们只考虑纵向振荡的贡献:

$$\gamma \frac{\partial \vec{n}}{\partial \gamma}(\theta) = \frac{1}{2} Re[\vec{k}_0^{\star}(\theta) \cdot (D_s + D_{-s})]$$



理论推导: 理想水平储存环

• 我们只考虑纵向振荡的贡献:

$$\gamma \frac{\partial \vec{n}}{\partial \gamma}(\theta) = \frac{1}{2} Re[\vec{k}_0^{\star}(\theta) \cdot (D_s + D_{-s})]$$

$$D_{\pm s}(\theta) = \frac{i}{1 - e^{i2\pi(\nu \pm \nu_s)}} e^{\mp i\nu_s \theta}$$

$$\times \int_{\theta}^{\theta + 2\pi} \vec{k}_0(\theta') \left[\frac{\vec{e}_y}{\rho_x} - \frac{\vec{e}_x}{\rho_y} - (1 + a\gamma)(\eta_x G_x \vec{e}_y - \eta_y G_y \vec{e}_x) \right]_{\theta'} e^{\pm i\nu_s \theta'} R d\theta'$$

- $\pm \vec{r}_0 \parallel \vec{e}_y$, $\vec{k}_0 \perp \vec{e}_{y'}$, $\rho_y -> \infty$, $\eta_y = 0$: $D_{\pm s} = 0$
- 即在理想水平储存环中,纵向振荡对退极化没有影响。

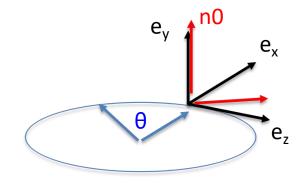
 ν_s :纵向工作点;

R:储存环的平均半径;

 ρ_x, ρ_y : 局部水平, 垂直偏转半径;

 η_x, η_y : 局部水平, 垂直色散函数;

 G_x, G_y : 局部水平, 垂直四极铁梯度;



理论推导: 水平储存环+磁铁误差

• 自旋运动方程:

$$rac{dec{S}}{d heta} = \left[ec{\Omega}_0 + \Deltaec{\Omega} + ec{\omega}
ight] imes ec{S}$$

$$\Delta \vec{\Omega} = -R(1 + a\gamma) \left[(x_{COD} \cdot G - t_x) \vec{e}_y + (y_{COD} \cdot G + t_y) \vec{e}_x \right]$$

- w=0时,运动方程的解: $\vec{n}=\vec{n}_0+\Delta\vec{n}$, $\vec{k}=\vec{k}_0+\Delta\vec{k}=\vec{m}+i\vec{l}_c$
- 扰动视为微扰:

$$\gamma \frac{\partial \vec{n}}{\partial \gamma}(\theta) = \frac{1}{2} Re[\vec{k}^{\star}(\theta) \cdot (\widetilde{D}_s + \widetilde{D}_{-s})|_{\vec{k}_0 \to \vec{k}}]$$

$$\widetilde{D}_{\pm s}(\theta) = \frac{i}{1 - e^{i2\pi(\nu \pm \nu_s)}} e^{\mp i\nu_s \theta}$$

$$\times \int_{\theta}^{\theta + 2\pi} \vec{k}(\theta') \left[\frac{\vec{e}_y}{\rho_x} - \frac{\vec{e}_x}{\rho_y} - (1 + a\gamma)(\eta_x G_x \vec{e}_y - \eta_y G_y \vec{e}_x) \right]_{\theta'} e^{\pm i\nu_s \theta'} R d\theta'$$

• 对于CEPC: $(1 + a\gamma) \gg 1$, $\eta_y << \eta_x$:

$$\widetilde{D}_{\pm s}(\theta) \approx \frac{i}{1 - e^{i2\pi(\nu \pm \nu_s)}} e^{\mp i\nu_s \theta}$$

$$\times \int_{\theta}^{\theta + 2\pi} \left[-(1 + a\gamma)(\eta_x G_{\mathbf{x}} \vec{e_y}) \right]_{\theta'} \cdot \vec{k}(\theta') e^{\pm i\nu_s \theta'} R d\theta'$$

理论推导: 水平储存环+磁铁误差

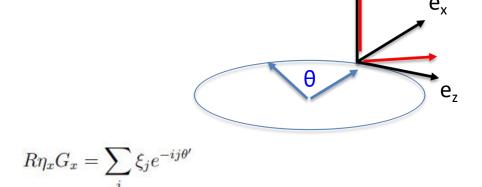
• k 在 e_v 方向的分量:

$$\begin{split} \vec{k} &= \vec{k}_0 + \alpha \vec{n}_0 + \beta \vec{k}_0^{\star} \\ \vec{e}_y \cdot \vec{k}(\theta') &= -\int_{-\infty}^{\theta'} \underline{\Delta \Omega_x} e^{i\nu(\Phi(\theta'') - \theta'')} e^{i\nu\theta''} d\theta'' \\ &= -\int_{-\infty}^{\theta'} \sum_k \omega_k e^{-ik\theta''} e^{i\nu\theta''} d\theta'' \\ &= i \sum \frac{\omega_k e^{i(\nu - k)\theta'}}{\nu - k} \end{split}$$

$$\Delta\Omega_x e^{i\nu\Phi(\theta'')-\theta''} = \sum_k \omega_k e^{-ik\theta''}$$

• 得到D±s:

$$\widetilde{D}_{\pm s}(\theta) \approx \frac{e^{\mp i\nu_s \theta} (1 + a\gamma)}{1 - e^{i2\pi(\nu \pm \nu_s)}} \times \sum_{k} \frac{\omega_k}{\nu - k} \int_{\theta}^{\theta + 2\pi} (R\eta_x G_x) e^{i(\nu - k \pm \nu_s)\theta'} d\theta'$$



理论推导: 水平储存环+磁铁误差

• v-k <<1 时: 其中经过计算,对于CEPC, $\xi_0 = 0.96 \approx 1$,且j≠0时, $\xi_j \ll 1$

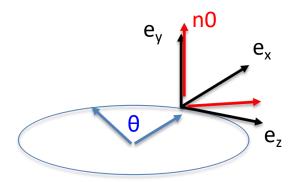
$$\frac{\widetilde{D}_{+s} + \widetilde{D}_{-s}}{2}|_{\theta} = i(1 + a\gamma) \sum_{k} \frac{\omega_{k}}{\nu - k} \sum_{j} \xi_{j} \frac{\nu - k - j}{(\nu - k - j)^{2} - \nu_{s}^{2}} e^{i(\nu - k - j)\theta}$$

$$\gamma \frac{\partial \vec{n}}{\partial \gamma}(\theta) = Re[\vec{k}^{\star} \cdot e^{i(\nu - k)\theta} \cdot i(1 + a\gamma) \sum_{k} \frac{\omega_{k}}{\nu - k} \xi_{0} \frac{\nu - k}{(\nu - k)^{2} - \nu_{s}^{2}}]$$

$$\frac{\partial \vec{n}}{\partial \gamma}|^{2} \approx (1 + a\gamma)^{2} \left[\sum_{k} \frac{\omega_{k}^{2} \xi_{0}^{2}}{((\nu - k)^{2} - \nu_{s}^{2})^{2}} \right]$$
整数共振项

• 以整数自旋共振为中心的一阶纵向边带自旋共振:

$$\begin{split} P &\approx \frac{P_{\infty}}{1 + \frac{\tau_{sk}}{\tau_{dep}}} \\ &\approx \frac{P_{\infty}}{1 + \frac{11}{18} \frac{\oint ds \left\langle \frac{1}{|\rho|^3} |\frac{\partial \hat{n}}{\gamma \partial \gamma}|^2 \right\rangle}{\oint ds \left\langle \frac{1}{|\rho|^3} \right\rangle}} \\ &\approx \frac{P_{\infty}}{1 + (1 + a\gamma)^2 \sum_k \frac{\omega_k^2 \xi_0^2}{((\nu - k)^2 - \nu_s^2)^2}} \end{split}$$



拟合结果

CEPC中, (v-k=0)与 (v-k±v_s=0)叠加,其中(v-k±v_s=0)对退极化有主要贡献;

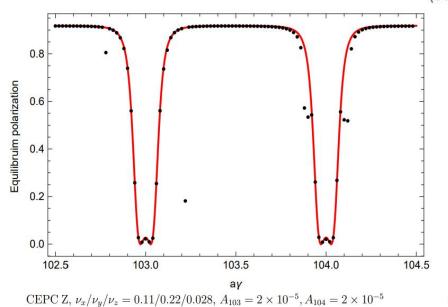
$$P \approx \frac{P_{\infty}}{1 + \frac{\tau_{sk}}{\tau_{dep}}}$$

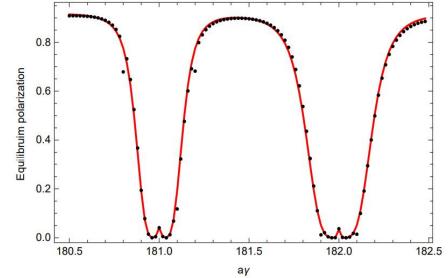
|v-k|>>v_s,近似为DK的理论公式,上升,下降沿吻合的原因:

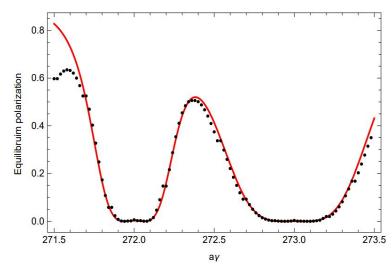
$$\frac{\tau_{sk}}{\tau_{dep}} = \sum_{k} \frac{A_k}{((\nu - k)^2 - \nu_s^2)^2} \longrightarrow \frac{\tau_{sk}}{\tau_{dep}} = \frac{11v^2}{18} \sum_{k} \frac{|\omega_k|^2}{(\nu - k)^4}$$

|v-k| < v_s时, P∞不再等于92%:

$$P_{\infty} = \frac{-8}{5\sqrt{3}} \frac{\oint ds \left\langle \frac{1}{|\rho|^3} \widehat{b} \cdot (\widehat{n} - \gamma \frac{\partial \widehat{n}}{\partial \gamma}) \right\rangle}{\oint ds \left\langle \frac{1}{|\rho|^3} \right\rangle}$$

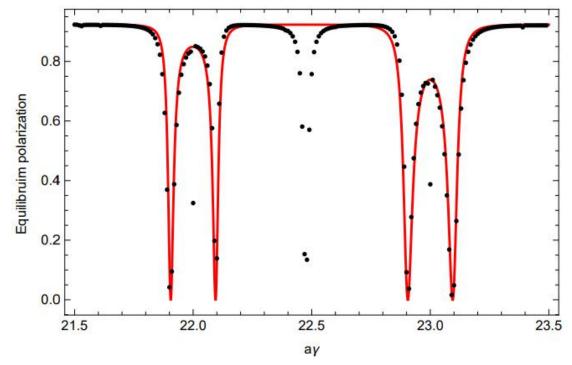






拟合结果

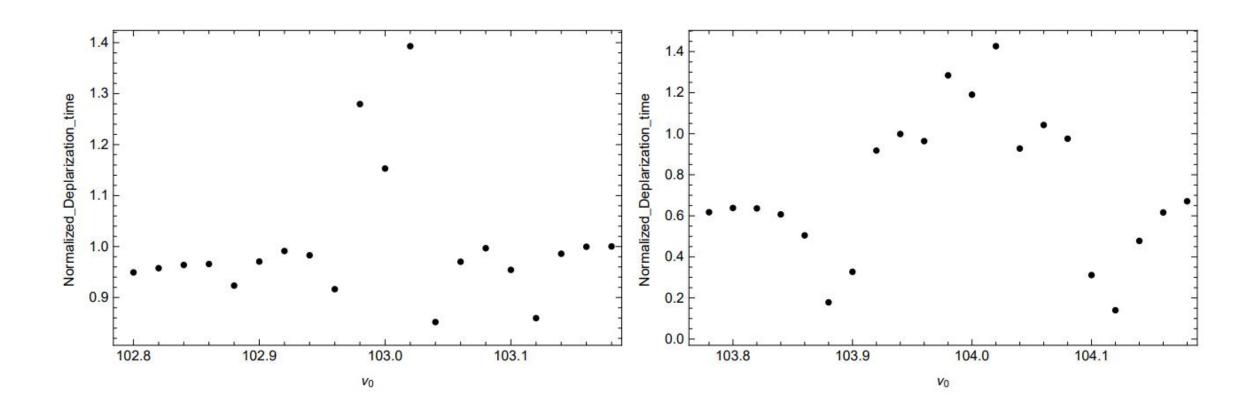
- (v-k=0) 与 (v-k±v_s=0) 不叠加的情况:
 - Other case: C=2112m, E=10GeV, Δ Y=10 μ m;
 - v-k=0 贡献可以忽略, v-k \pm v_s=0贡献有主要贡献;



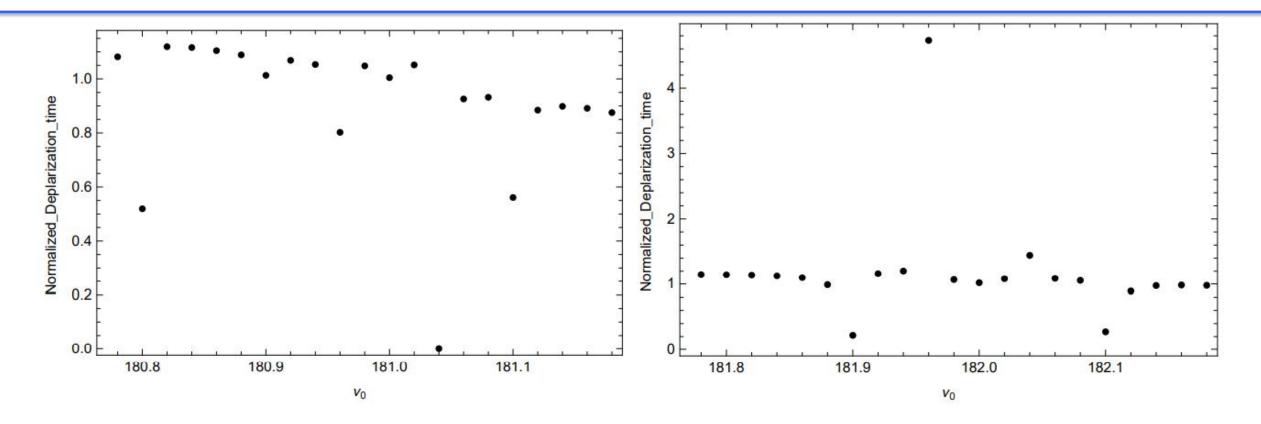
$$rac{ au_{sk}}{ au_{dep}} = \sum_{k} rac{A_k}{((
u - k)^2 -
u_s^2)^2}$$

 $\nu_x/\nu_y/\nu_z = 0.392/0.475/0.094, A_{22} = 0.7 \times 10^{-5}, A_{23} = 2 \times 10^{-5}$

Back up



Back up



Back up

