

Jet Tomography in high-energy nuclear collisions

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Intersections, Tsinghua U., 2022. 8. 26–28

热烈祝贺清华大学
物理系复系40周年！



Outline

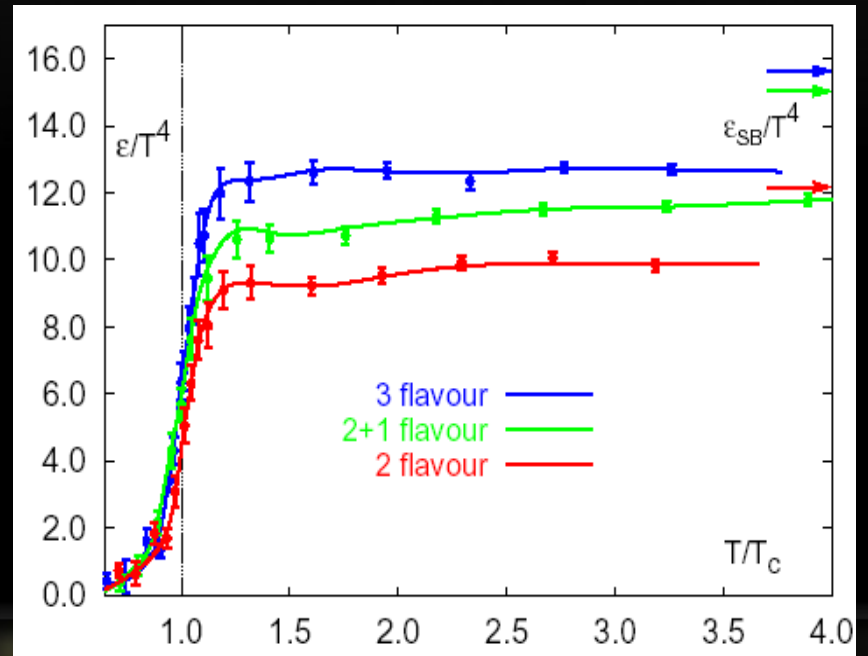
- Introduction
- Leading hadron production
- Full jet observables
 - 1) splitting scales
 - 2) W +jet
 - 3) heavy flavor jets
- Summary

Deconfinement and QGP

It would be interesting to explore new phenomena by distributing high energy or high nuclear density over a relatively large volume.

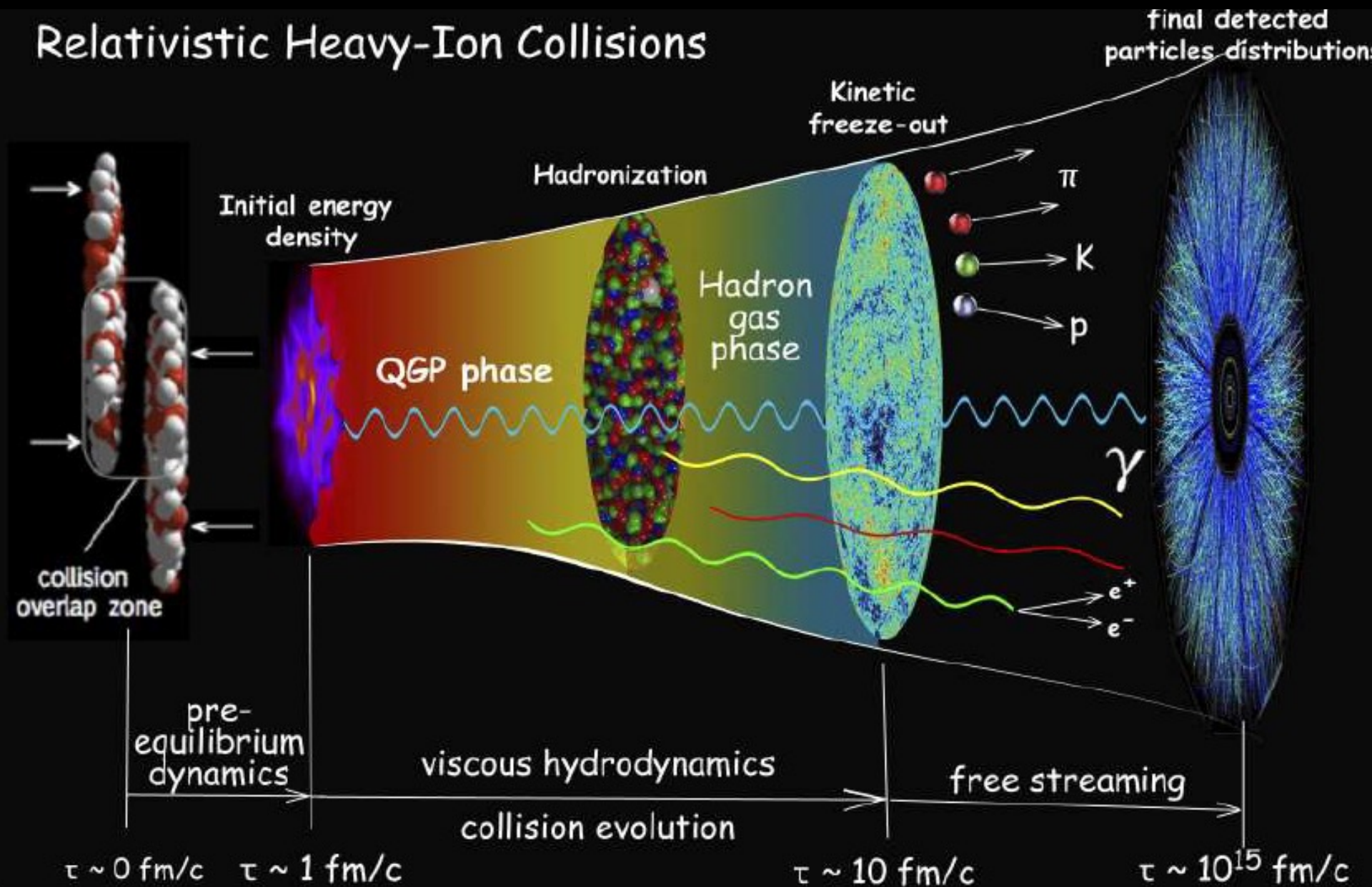
T. D. Lee (1978)

Lattice QCD predicts phase of thermal QCD matter with sharp rise in number of degrees of freedom near $T_c=170\text{MeV}$.



The Little Bang

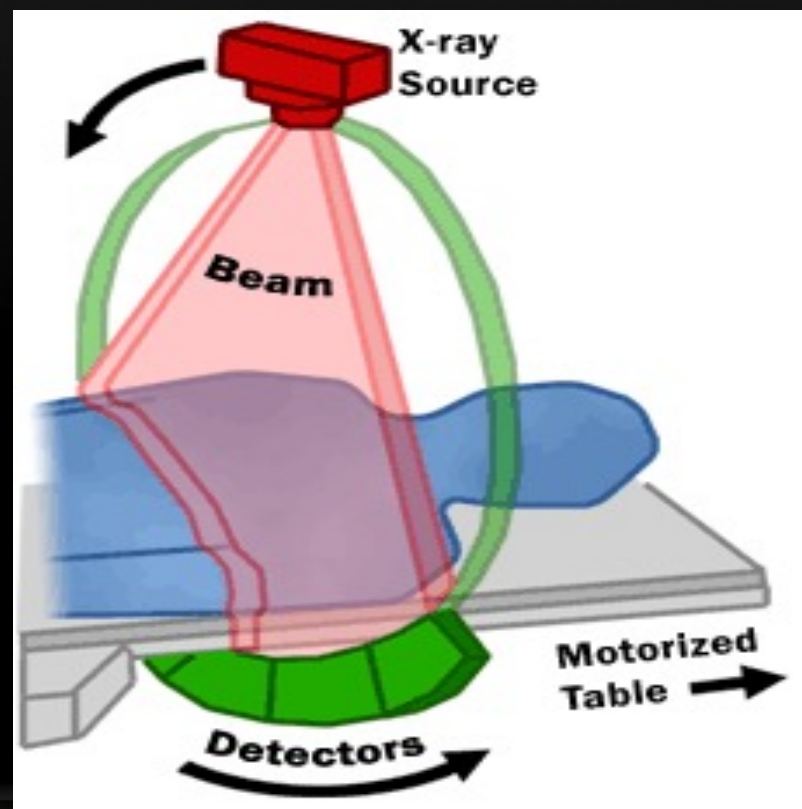
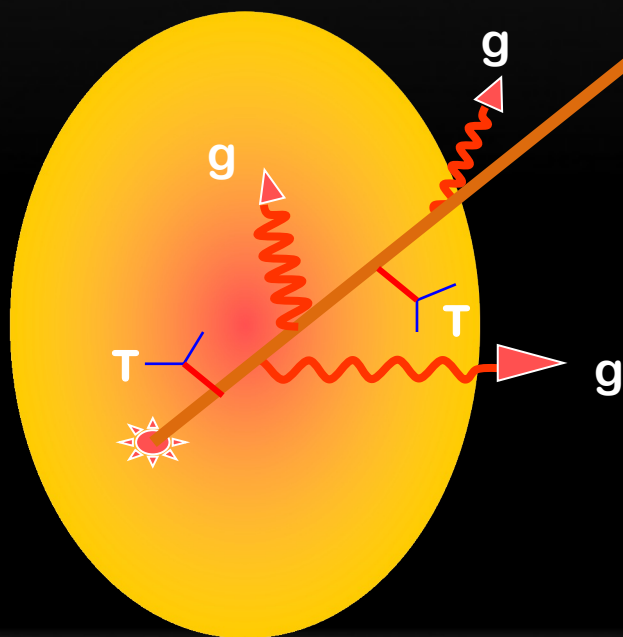
Relativistic Heavy-Ion Collisions



Jet quenching

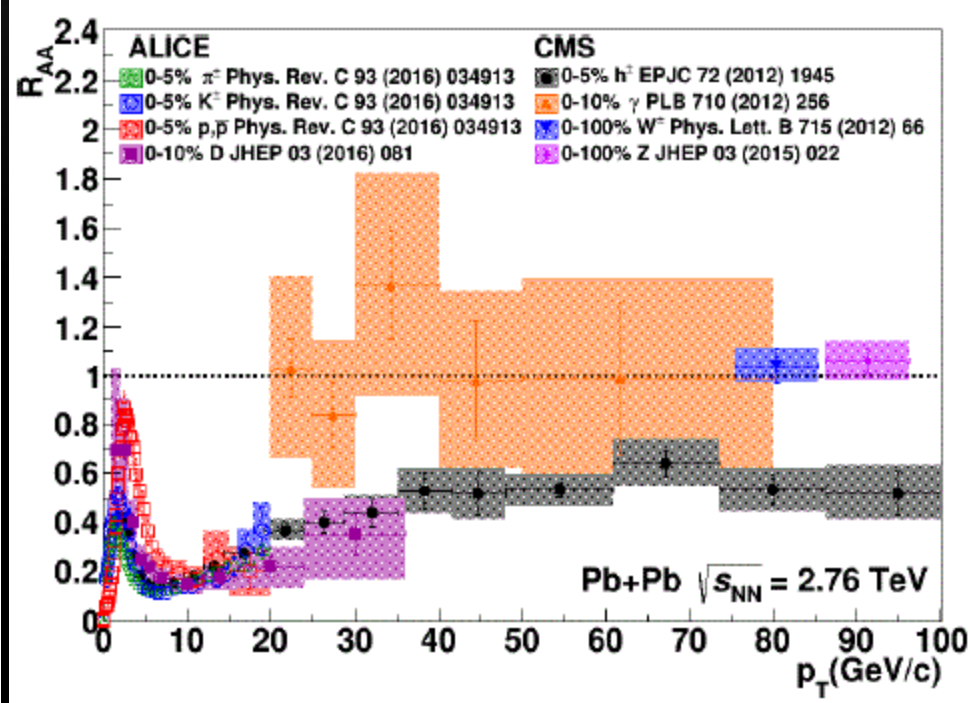
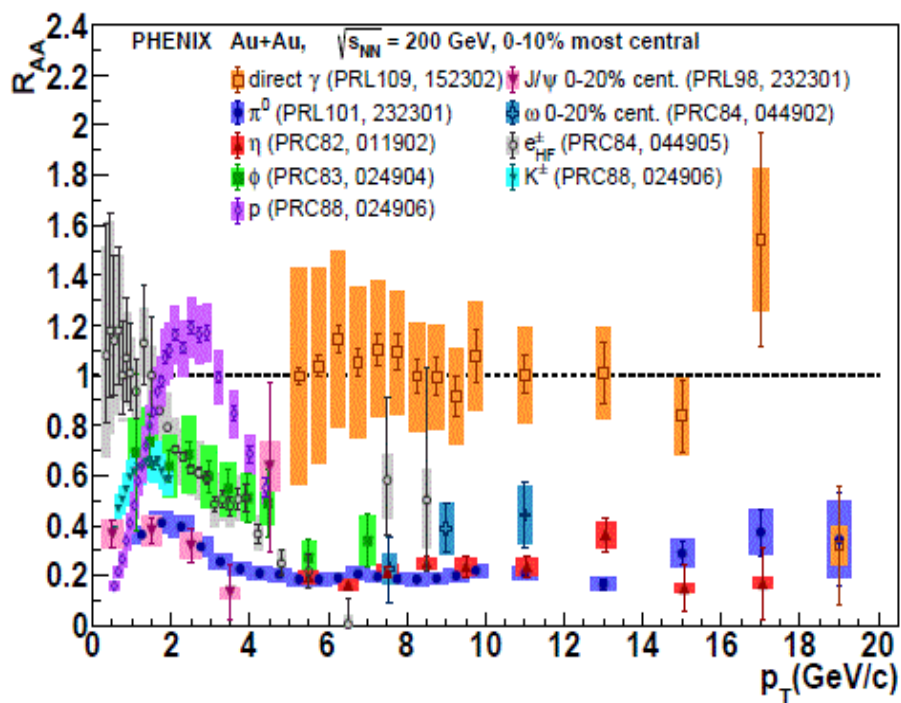
Parton energy has been proposed as an excellent probe of the hot/dense matter created at HIC.

Single Hadron Tomography



Jet quenching at RHIC and LHC

$$R_{AA} = \frac{\text{Yield}_{AA} / \langle N_{\text{binary}} \rangle_{AA}}{\text{Yield}_{pp}}$$



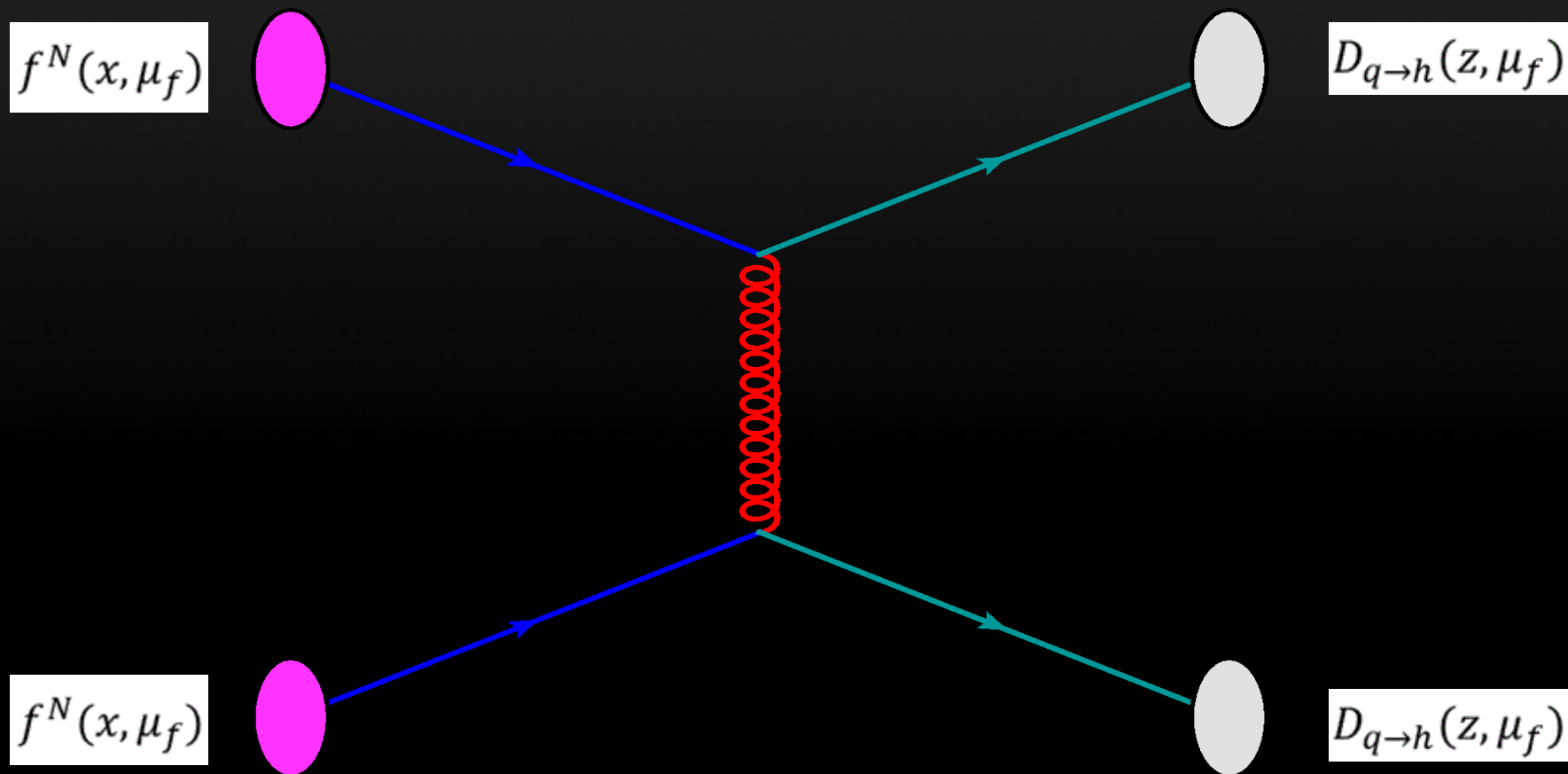
Fingerprints of jet quenching

领头强子

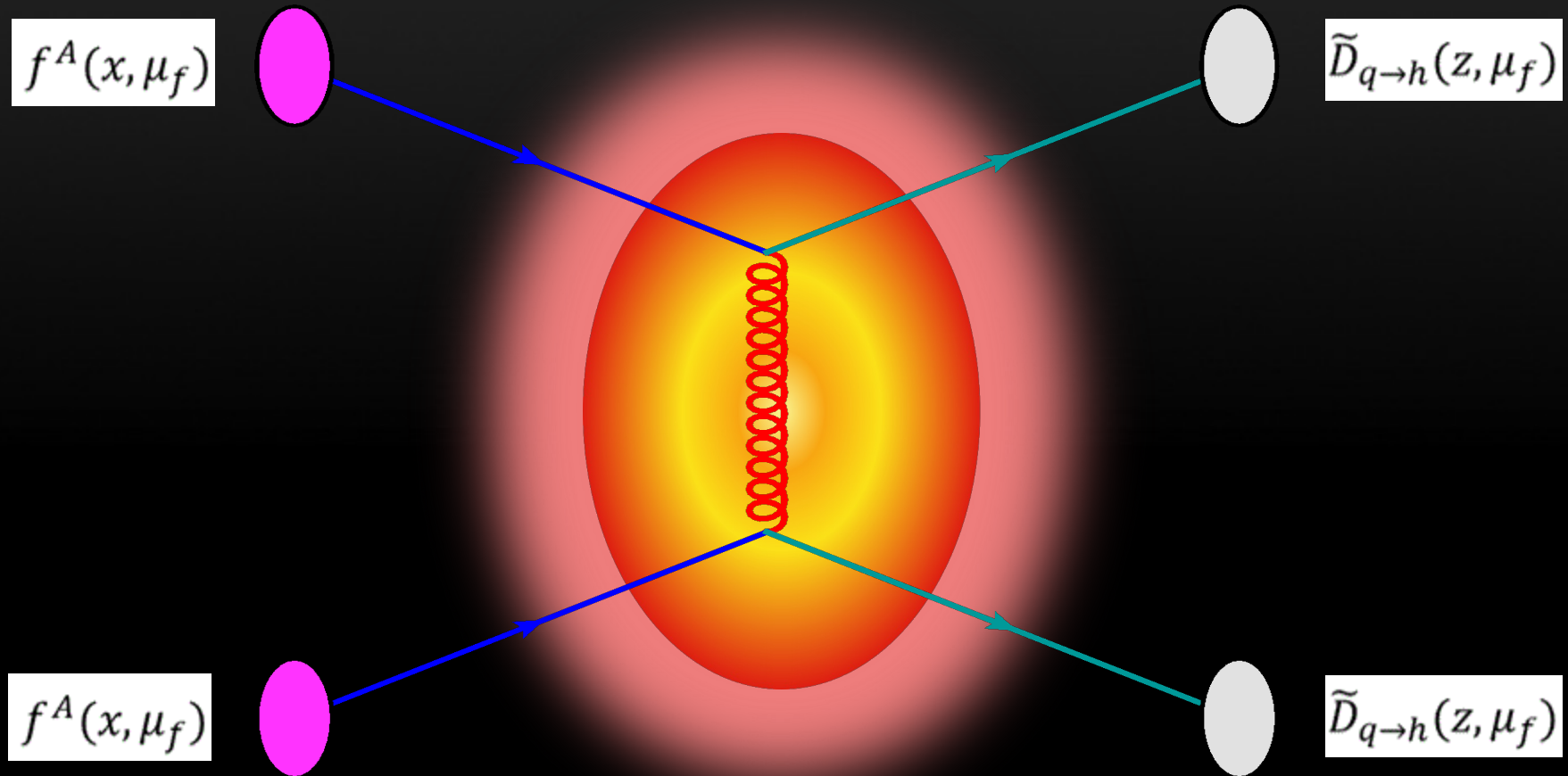


整体喷注

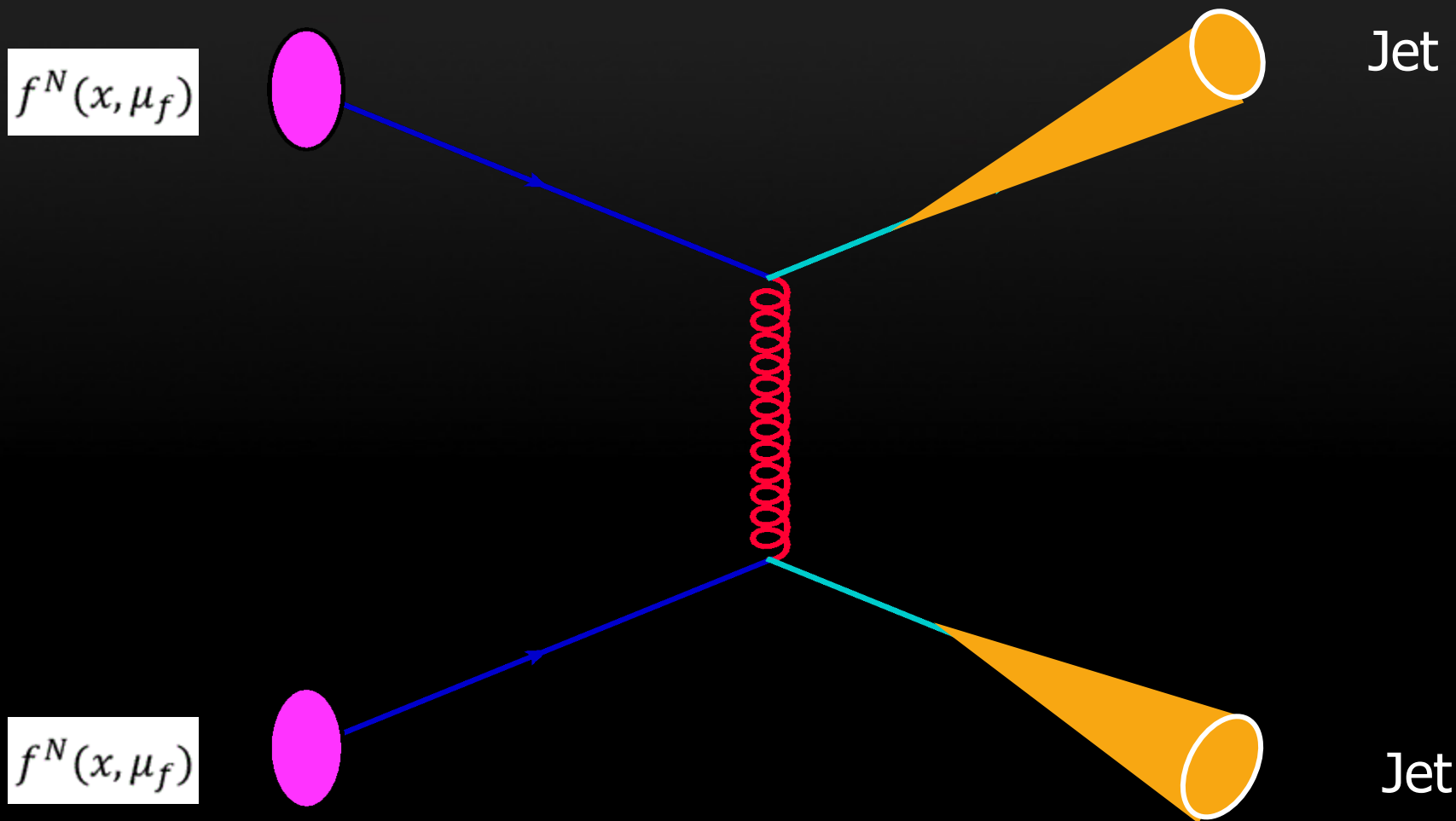
Leading hadron production



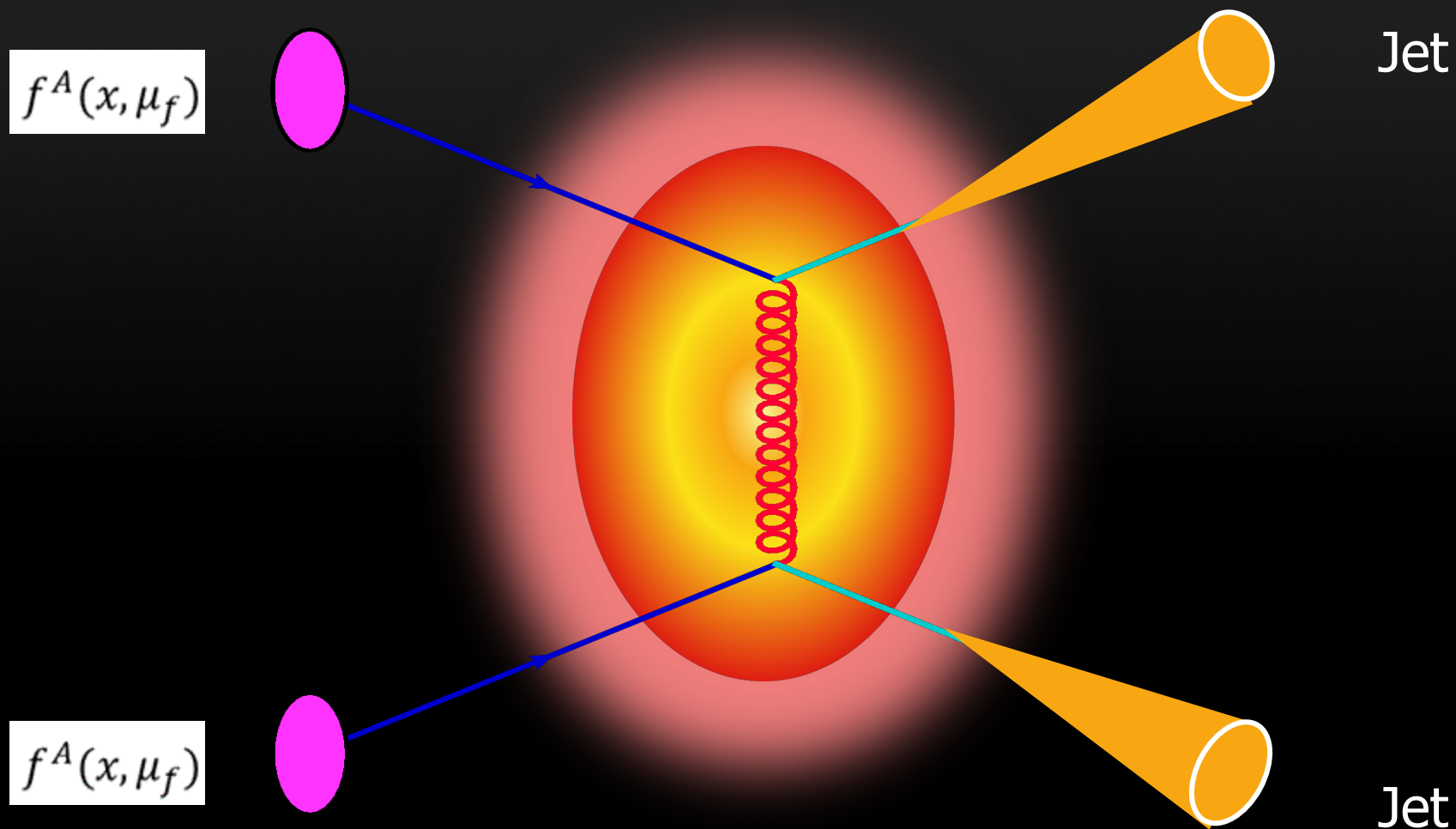
Leading hadron production



Full jets



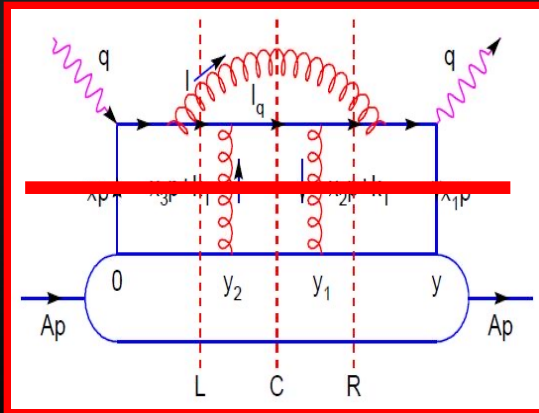
Full jets



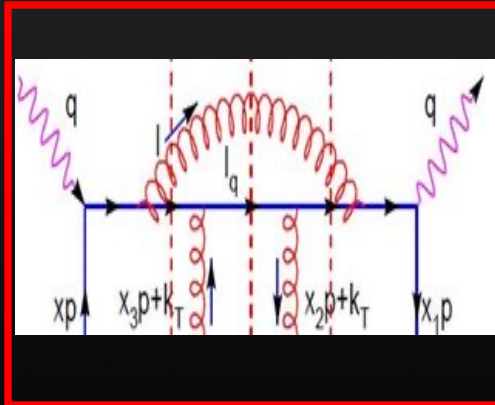
Large p_T hadrons in HIC

Jet quenching with higher twist approach

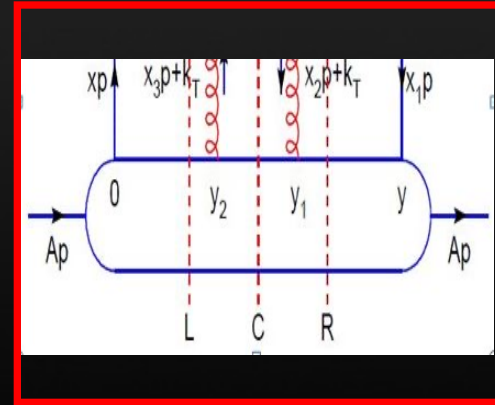
The parton energy loss due to medium-induced gluon radiation has been calculated with higher twist approach.



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⊗



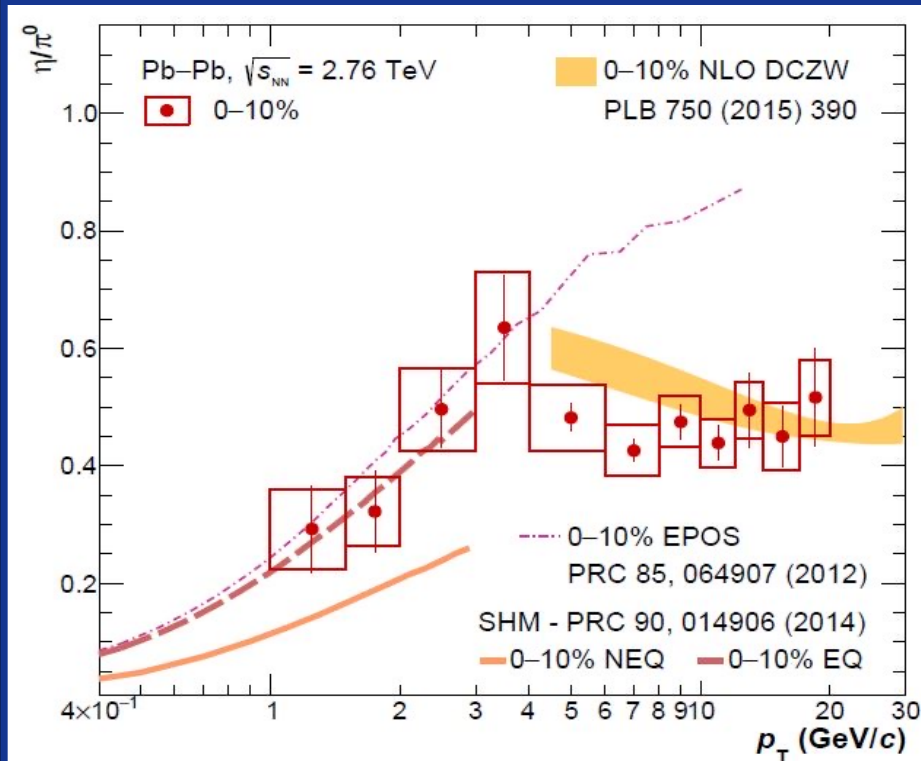
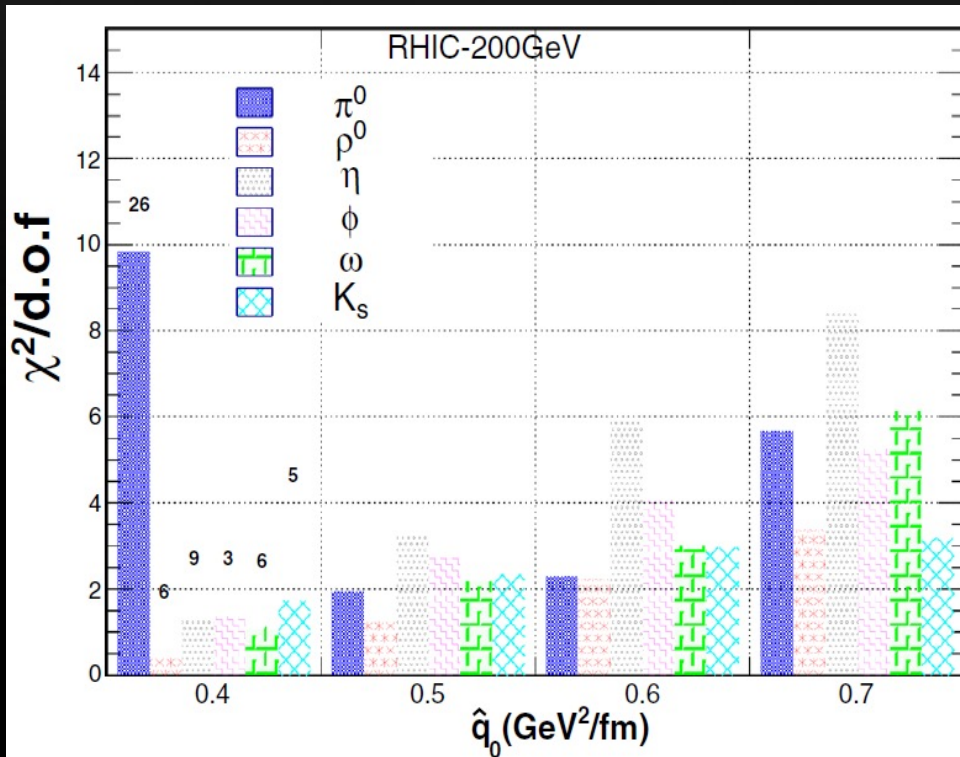
$$\begin{aligned} \tilde{D}_{q \rightarrow h}(z_h, \mu^2) &\equiv D_{q \rightarrow h}(z_h, \mu^2) + \int_0^{\mu^2} \frac{d\ell_T^2}{\ell_T^2} \frac{\alpha_s}{2\pi} \int_{z_h}^1 \frac{dz}{z} [\Delta\gamma_{q \rightarrow qg}(z, x, x_L, \ell_T^2) D_{q \rightarrow h}(z_h/z) \\ &+ \Delta\gamma_{q \rightarrow gq}(z, x, x_L, \ell_T^2) D_{g \rightarrow h}(z_h/z)] , \end{aligned}$$

$$\begin{aligned} \frac{1}{N_{\text{bin}}^{AB}(b)} \frac{d\sigma_{AB}^h}{dy d^2p_T} &= \sum_{abcd} \int dx_a dx_b f_{a/A}(x_a, \mu^2) f_{b/B}(x_b, \mu^2) \\ &\times \frac{d\sigma}{d\hat{t}}(ab \rightarrow cd) \frac{\langle \tilde{D}_c^h(z_h, Q^2, E, b) \rangle}{\pi z_c} + \mathcal{O}(\alpha_s^3). \end{aligned}$$

X Guo, X N Wang, PRL(2001); X Guo, X N Wang, NPA (2001);
BWZ, X N Wang, NPA(2003) ; BWZ, E Wang, X N Wang, PRL (2004)

- Yields of 6 species of hadrons and their ratios in Au+Au and Pb+Pb are calculated.

ALICE, PRC 98(2018)044901



W Dai, X Chen, BWZ, et al, PLB (2015); EPJC (2017); PRC(2018)

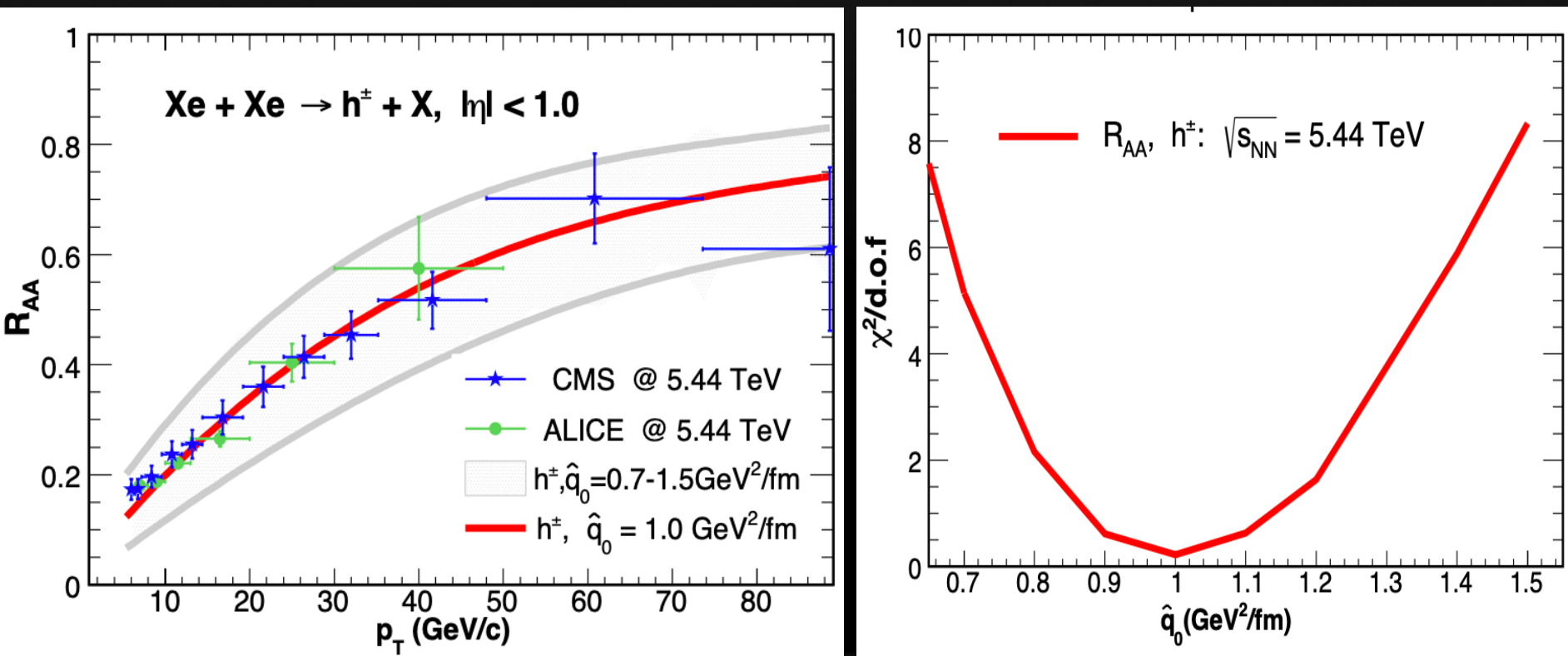
G Ma, W Dai, BWZ, E Wang, EPJC (2019)

Charged hadrons in Xe+Xe at NLO

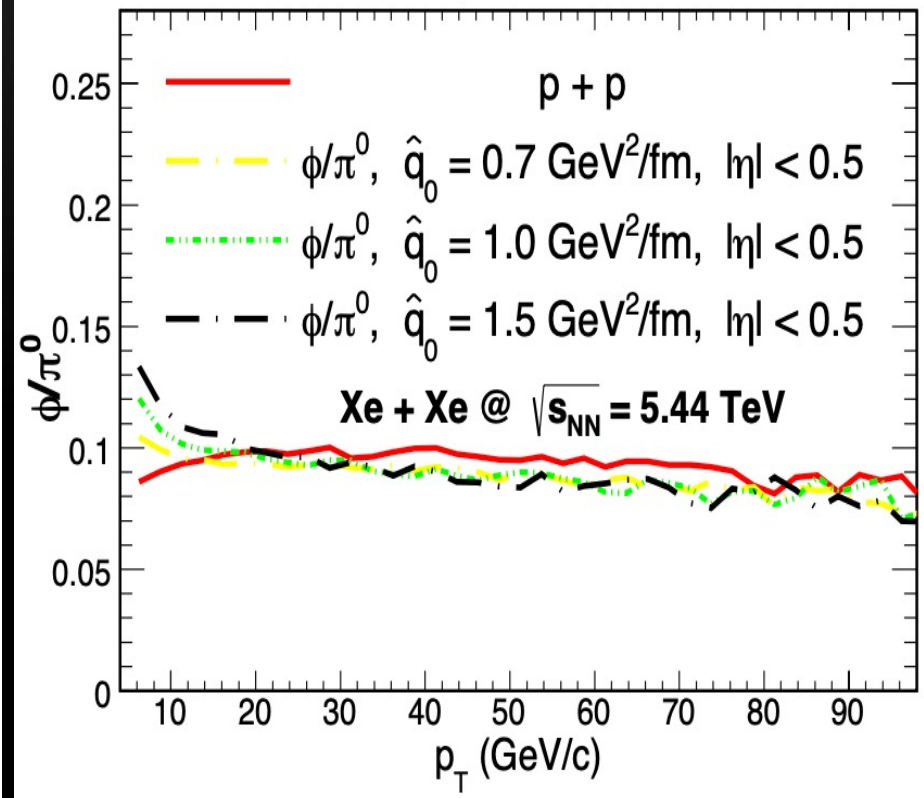
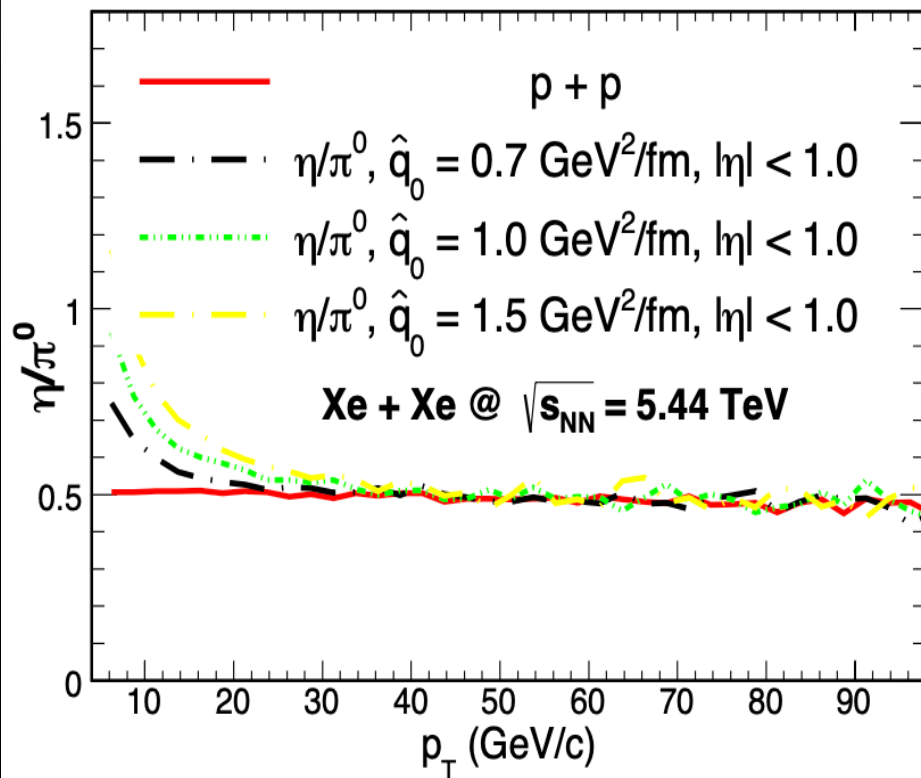
16

- Charged hadron productions in Xe+Xe has been studied.
- Jet transport coefficient is extracted by fitting the data.

$$(dE/dx)_{\text{rad}} = -C_2 \hat{q} L$$

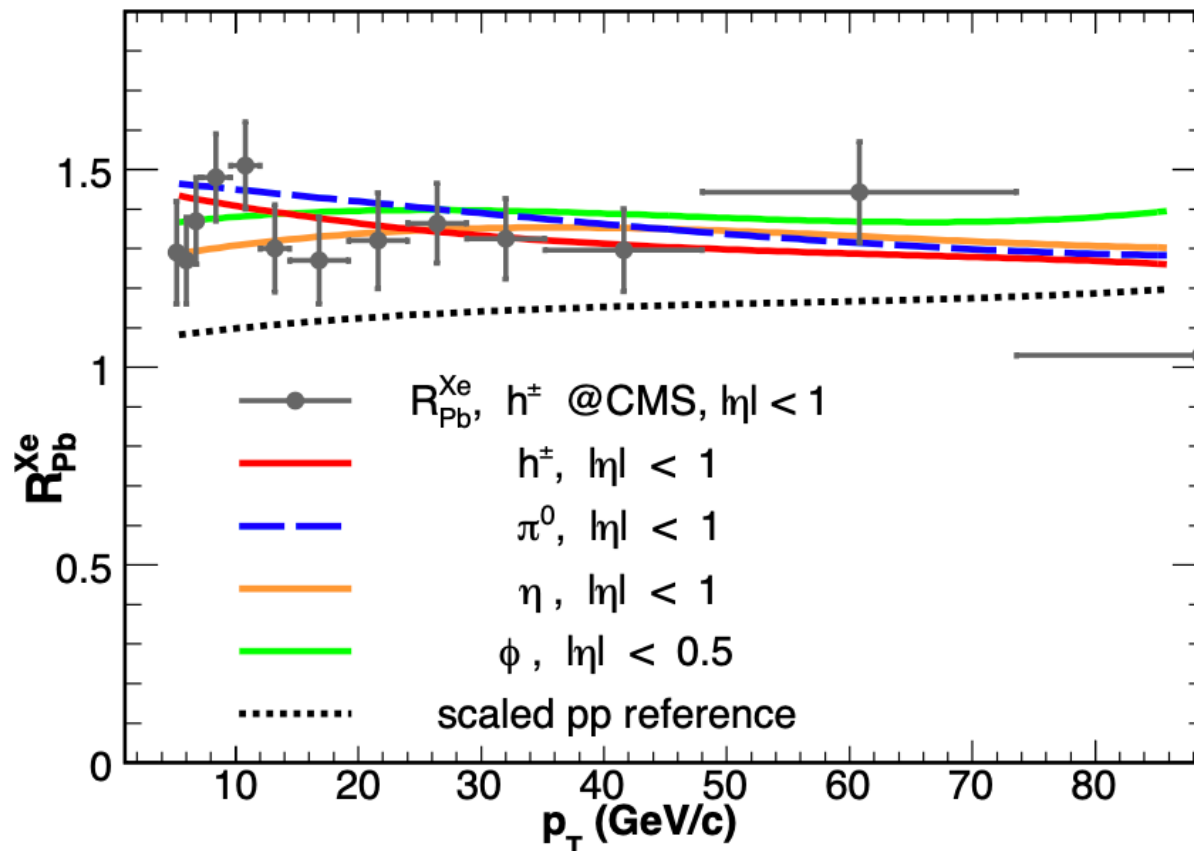


Particle ratios in Xe+Xe at NLO

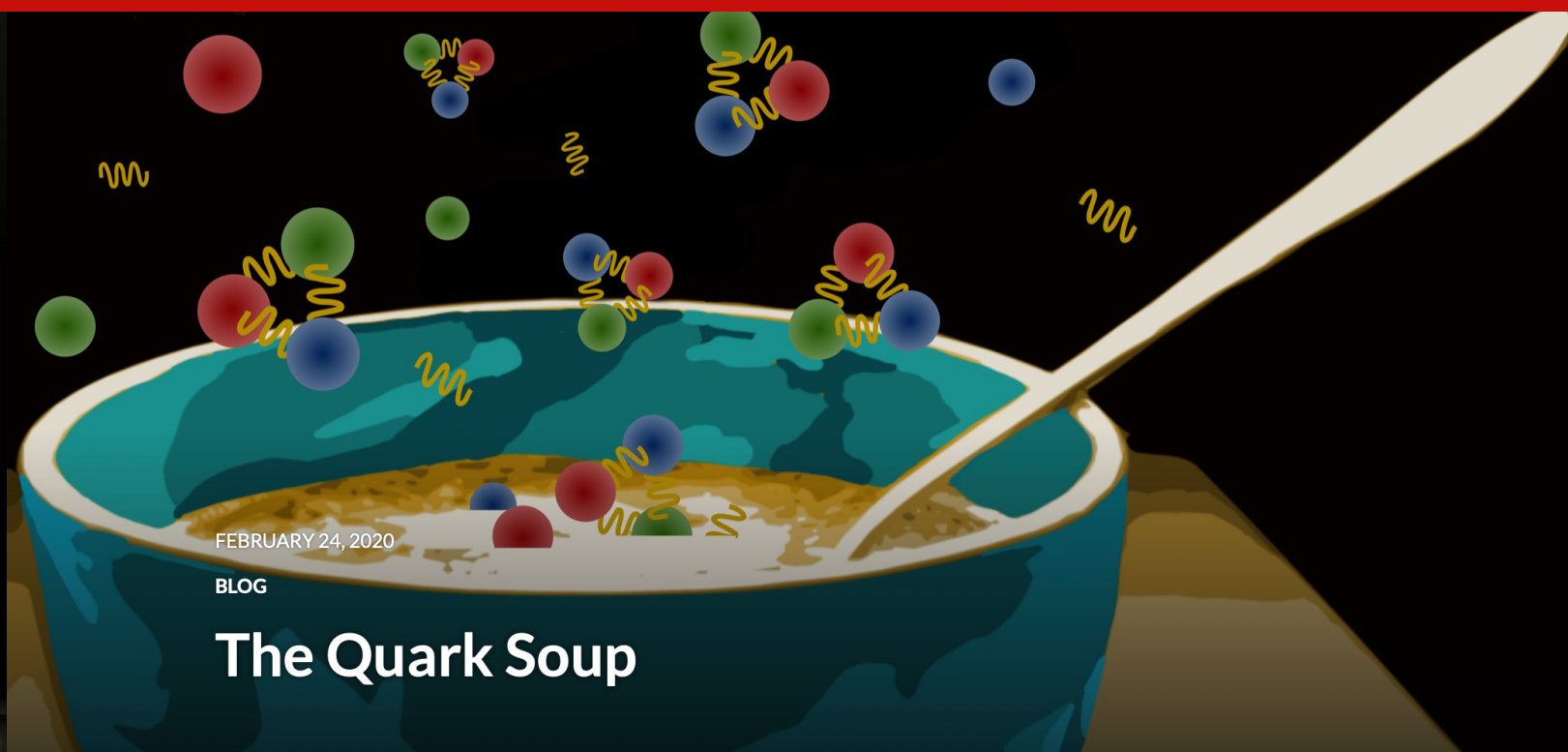


System size dependence of energy loss

$$R_{\text{Pb}}^{\text{Xe}}(p_{\text{T}}) = \frac{dN^{\text{XeXe}}/dp_{\text{T}}}{dN^{\text{PbPb}}/dp_{\text{T}}} \frac{T_{\text{PbPb}}}{T_{\text{XeXe}}}$$

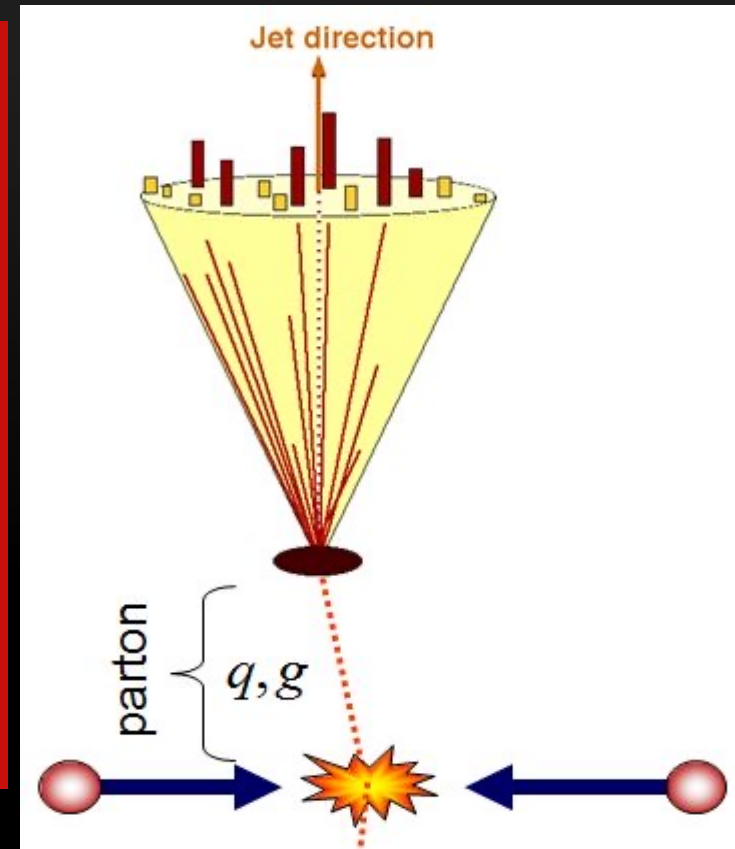


Jets in quark soup



What is a jet?

- A jet is a spray of final-state particles roughly moving in the same direction and defined by jet finding algorithms.
- At LO pQCD, jet \approx parton.
- In pQCD local-parton-hadron duality (LPHD) is used
- Jet: more precise and powerful



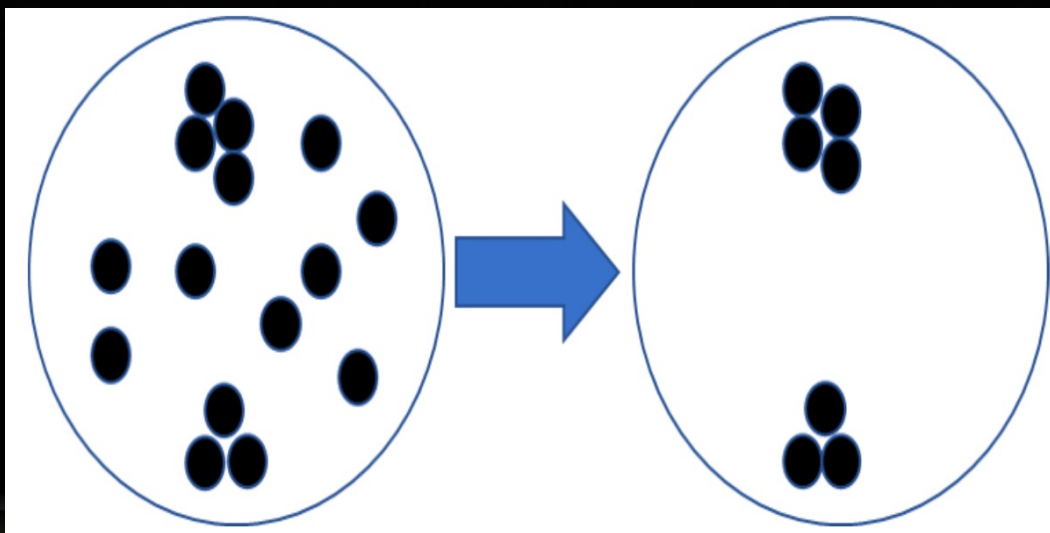
$$E_T = \sum_{i \in \text{jet}} E_{T,i}$$

$$y = \sum_{i \in \text{jet}} y_i E_{T,i} / E_T$$

$$\phi = \sum_{i \in \text{jet}} \phi_i E_{T,i} / E_T$$

$$R_{ij} = \sqrt{(y_i - y_j)^2 + (\phi_i - \phi_j)^2}$$

Reclustered large radius jets



remove soft radiation

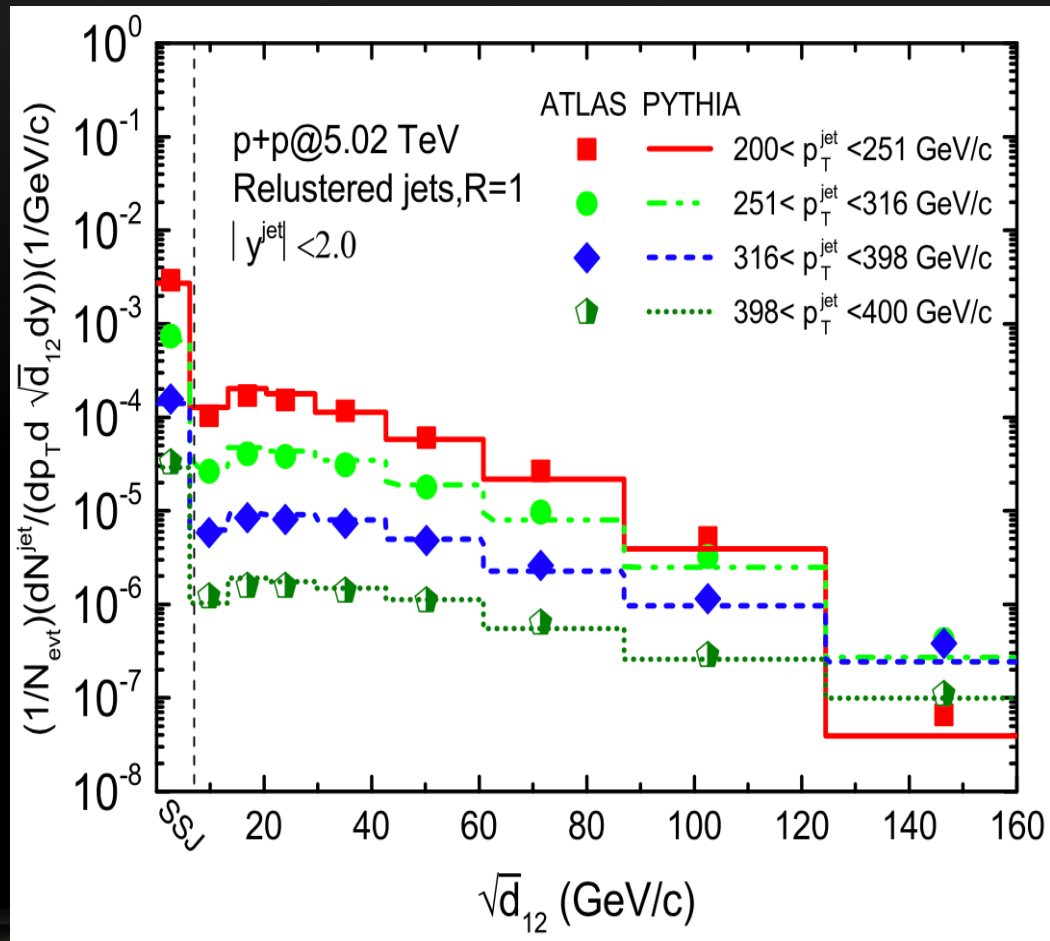
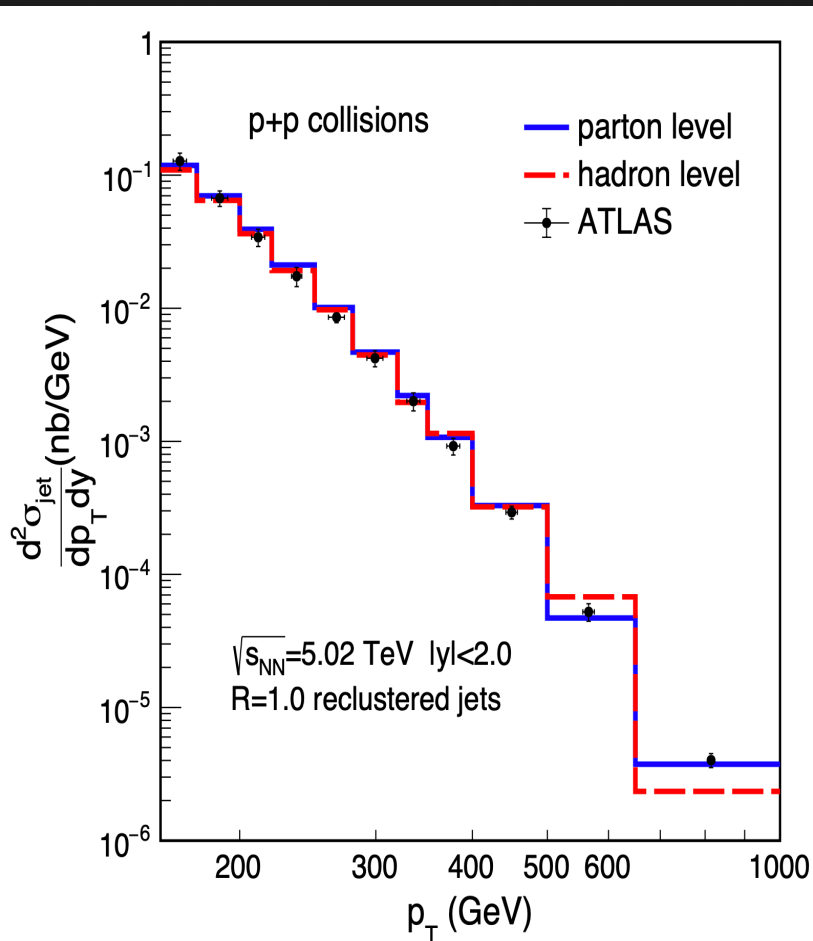
splitting scale

$$d_{12} = \min(p_{T1}^2, p_{T2}^2) \cdot \Delta R_{12}^2 / R^2,$$

$$\Delta R_{12} = \sqrt{\Delta \phi_{12}^2 + \Delta y_{12}^2}$$

Reclustered LR jets in p+p

$$\sqrt{d_{12}} = \sqrt{\min(p_{T,1}^2, p_{T,2}^2) * \Delta R_{12}^2}$$



Linear Boltzmann Transport Model

- Elastic scattering:

$$\begin{aligned}
 p_1 \cdot \partial f_1(p_1) &= - \int dp_2 dp_3 dp_4 (f_1 f_2 - f_3 f_4) |M_{12 \rightarrow 34}|^2 \\
 &\quad \times (2\pi)^4 \delta^4(P_1 + P_2 - P_3 - P_4) \\
 dp_i &\equiv \frac{d^3 p_i}{2E_i (2\pi)^3}, \quad |M_{12 \rightarrow 34}|^2 = C g^2 (s^2 + u^2) / (t + \mu^2)^2 \\
 f_i &= 1 / (e_i^{p \cdot u / T} \pm 1) \quad (i = 2, 4), \quad f_i = (2\pi)^3 \delta^3(\vec{p} - \vec{p}_i) \delta^3(\vec{x} - \vec{x}_i) \quad (i = 1, 3)
 \end{aligned}$$

X N Wang, Y Zhu, PRL(2013); He, Luo, Wang, Zhu, PRC (2015)

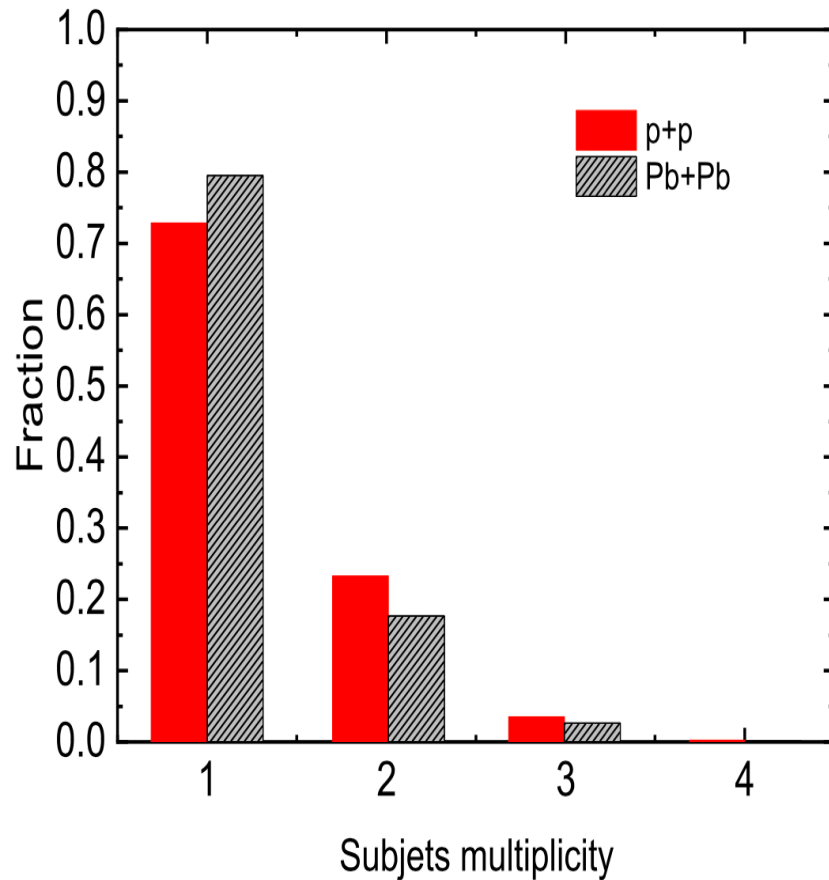
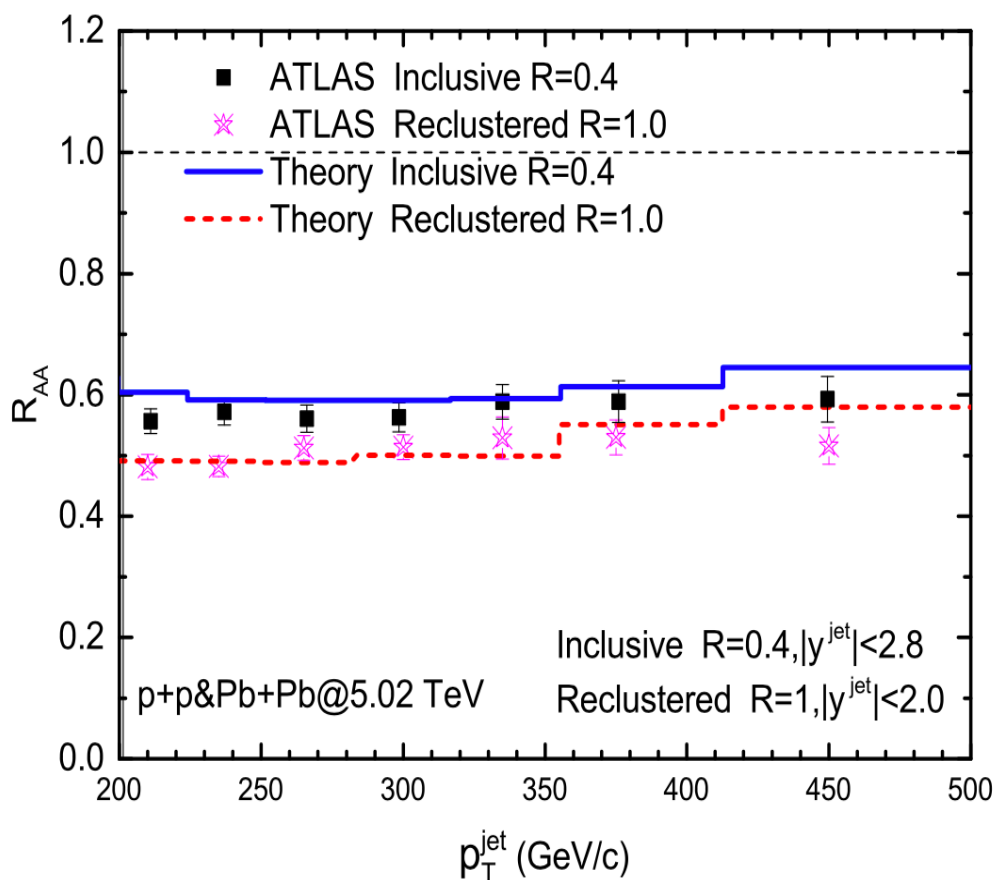
- Inelastic scattering by the higher twist approach:

$$\frac{dN_g}{dx dk_{\perp}^2 dt} = \frac{2\alpha_s P(x) \hat{q}}{\pi k_{\perp}^4} \sin^2 \left(\frac{t - t_i}{2\tau_f} \right) \left(\frac{k_{\perp}^2}{k_{\perp}^2 + x^2 M^2} \right)^4$$

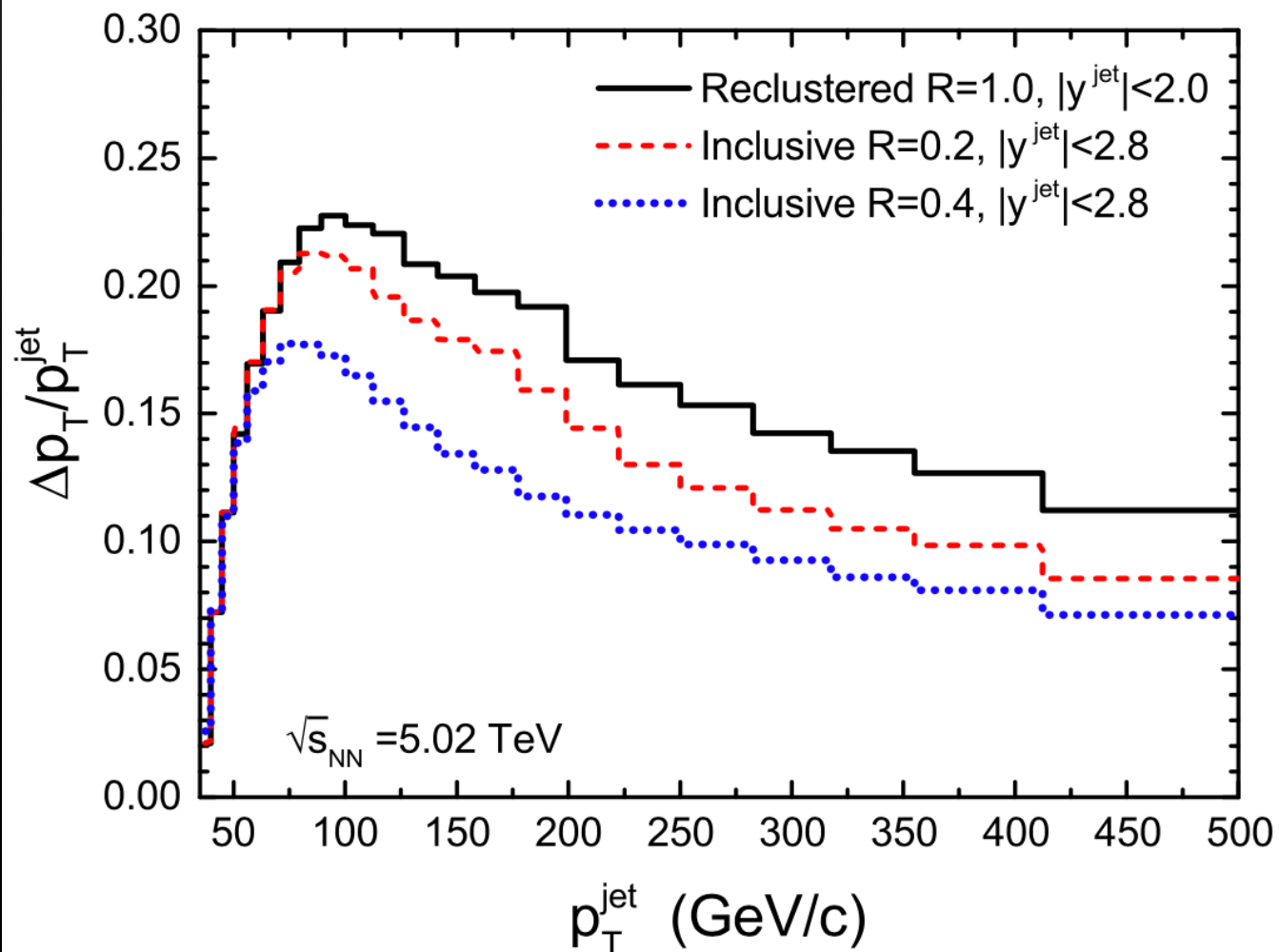
Guo, X N Wang, PRL(2002); BWZ, X Wang, NPA(2003);
 BWZ, E Wang, X N Wang, PRL (2004); Majumder, PRD(2012)

Nuclear modifications

- Nuclear suppression of reclustered LR jets at $R=1.0$ is larger than that of inclusive jets with $R=0.4$.

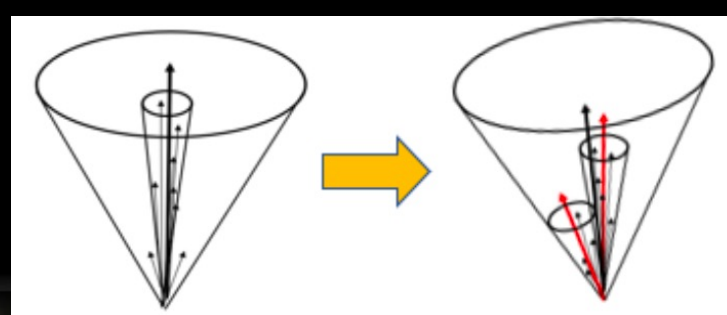
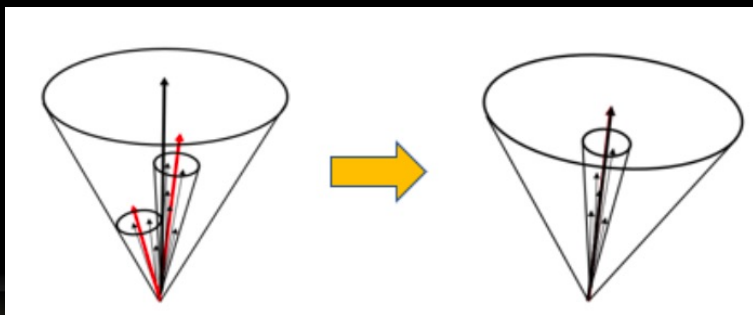
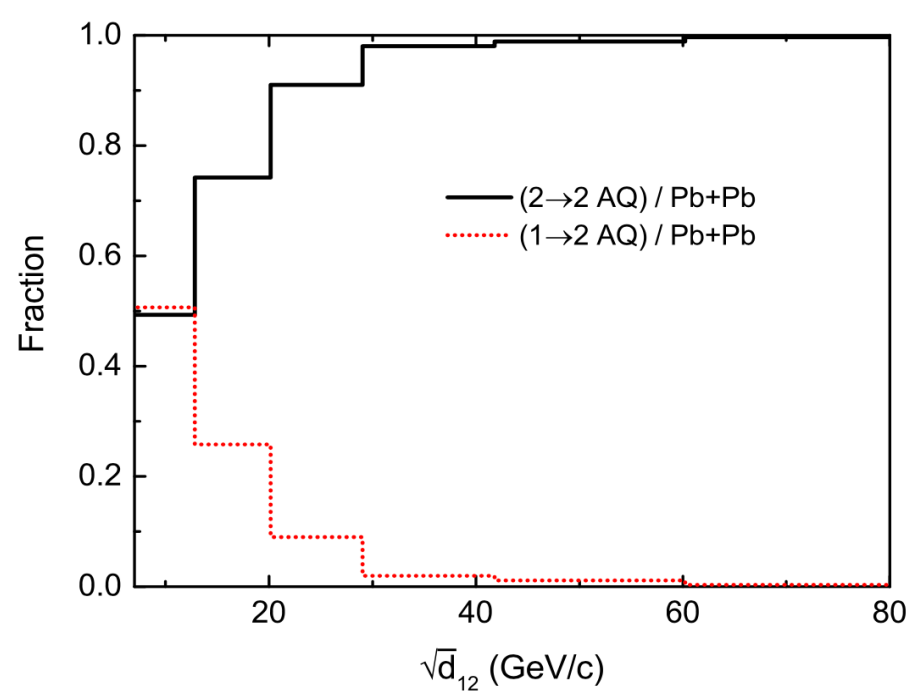
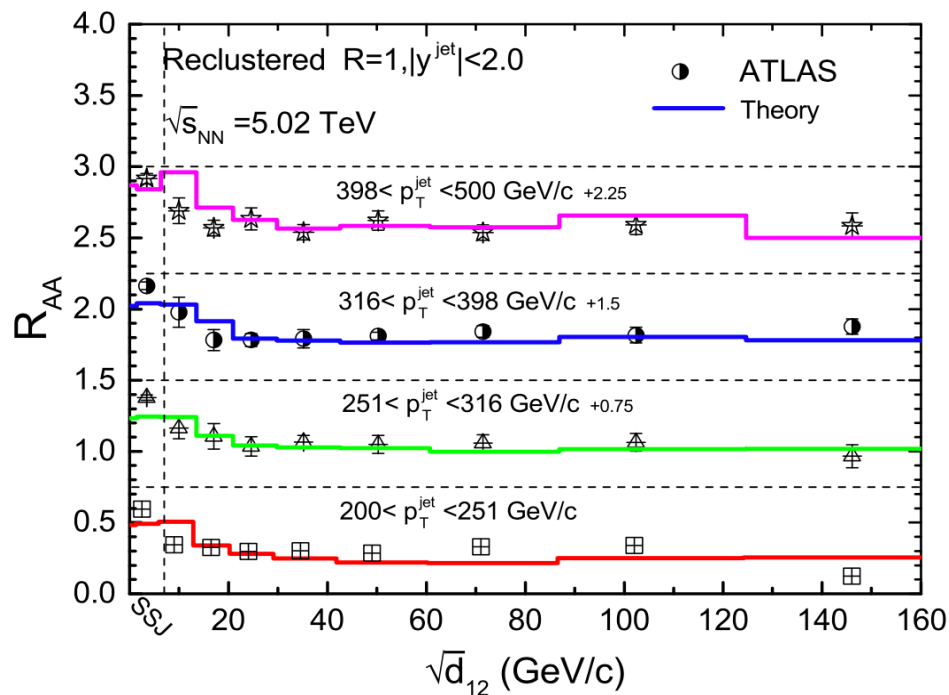


Energy loss fraction

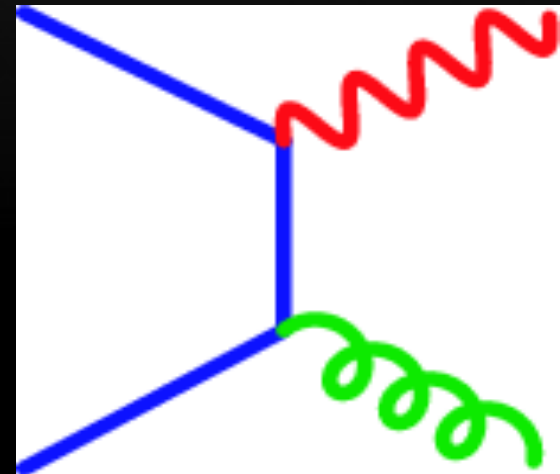
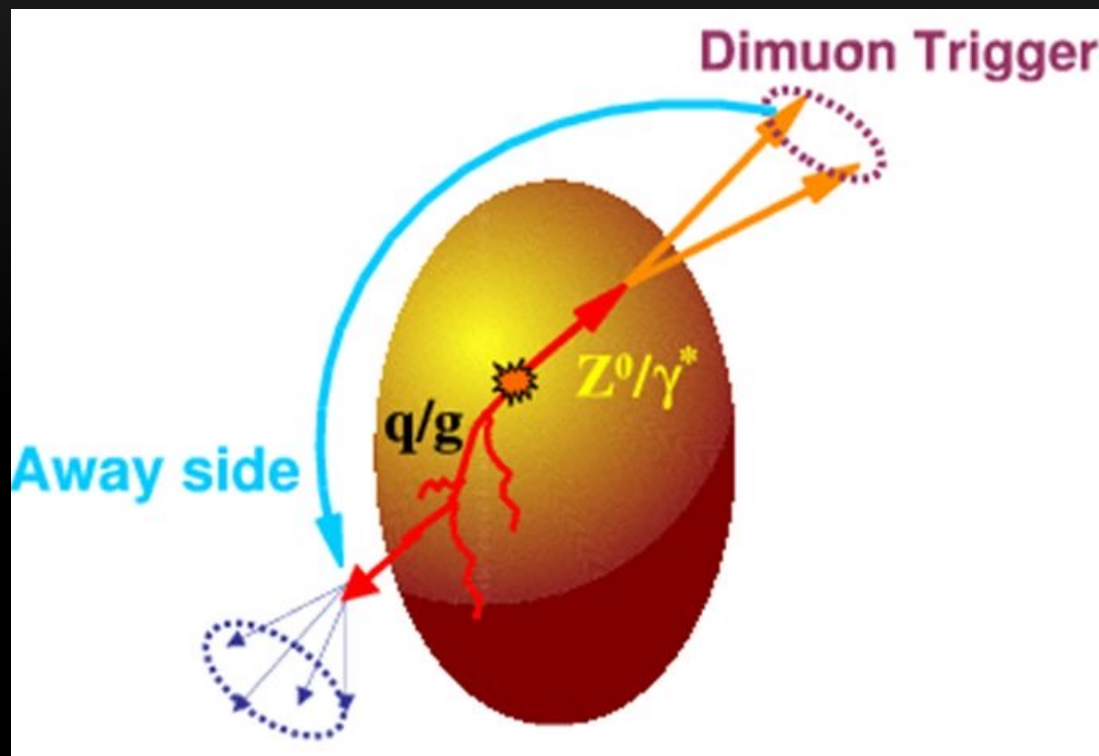


Energy loss of reclustered jets

$$\sqrt{d_{12}} = \sqrt{\min(p_{T,1}^2, p_{T,2}^2) * \Delta R_{12}^2}$$



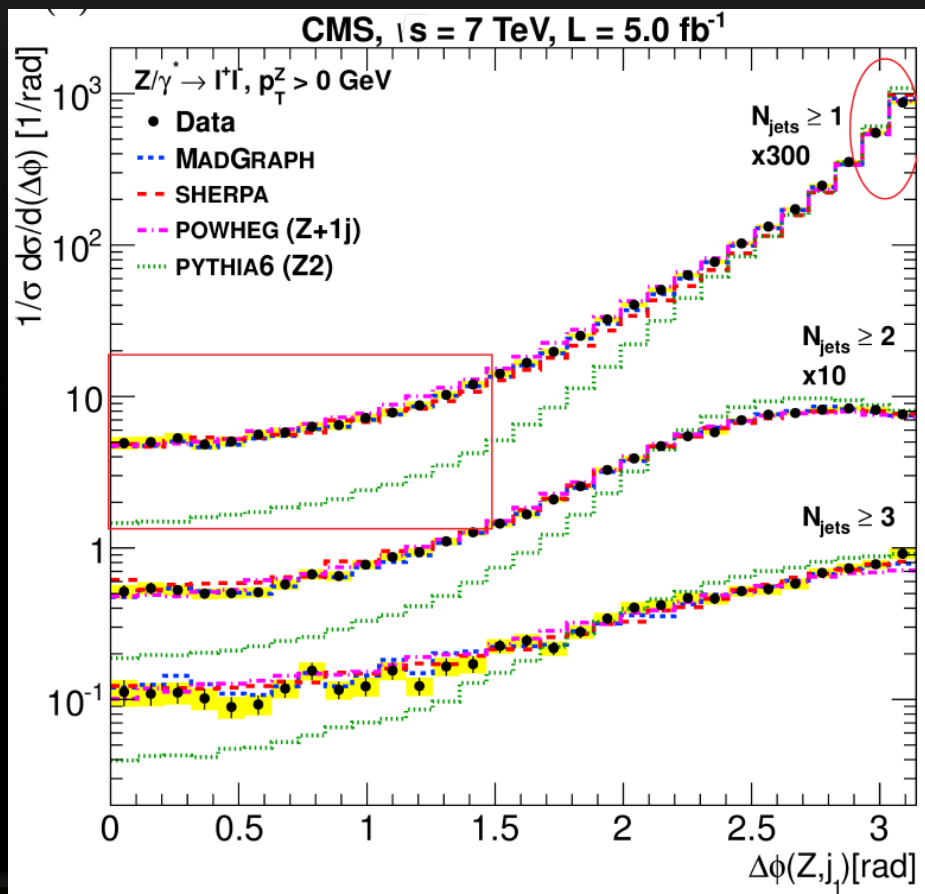
Tagged jet production in HIC



“golden channel”

Angular correlation in Z+jet

- NLO calculations fail at angular difference $\sim \pi$;
- LO+PS calculations fail at small angular difference.
- Z+jet in A+A: NLO+PS+Eloss



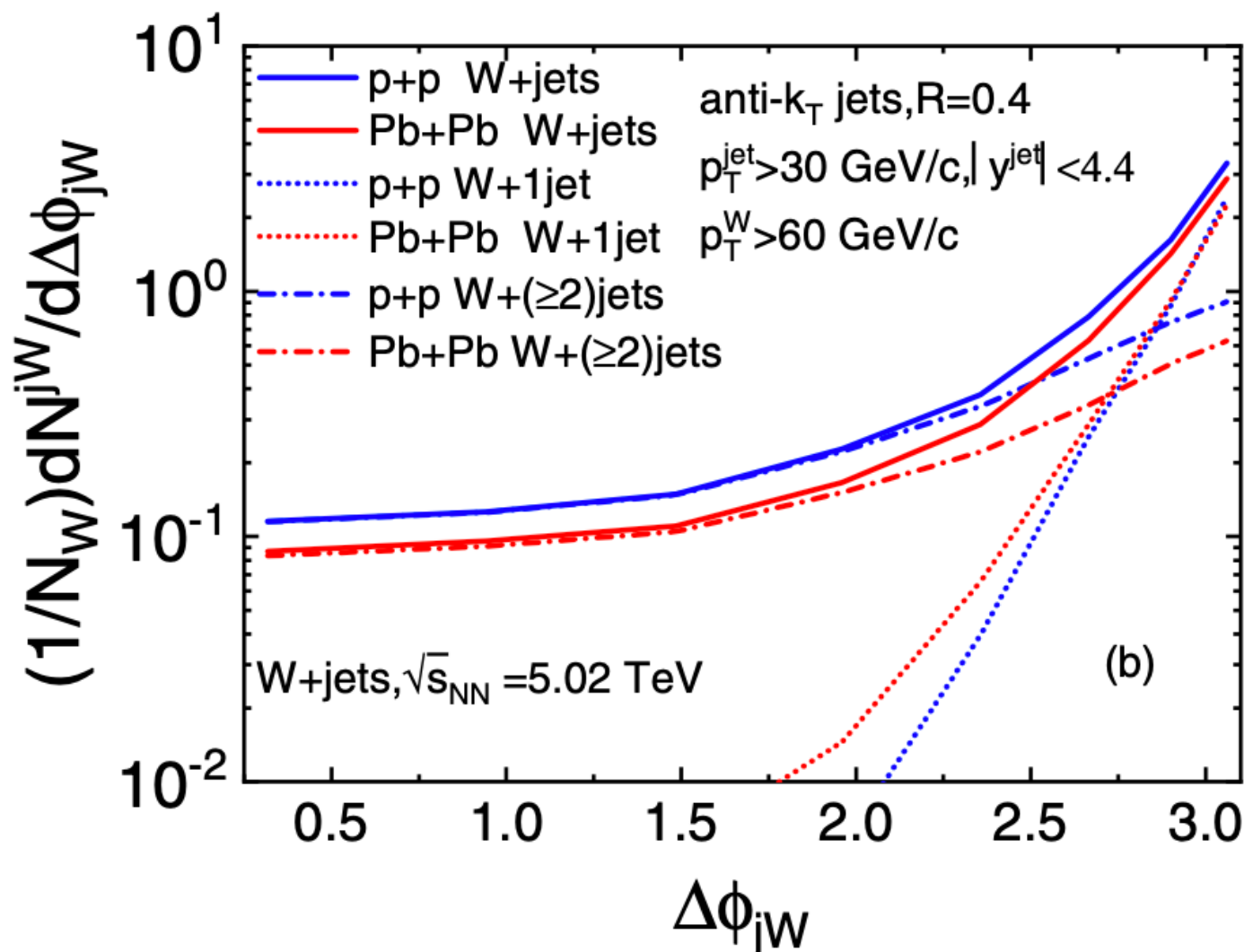
CMS, PLB 722
(2013) 238

Chen, Qin, Wei,
Xiao, Zhang,
PLB (2017)

**NLO +PS
+Eloss**

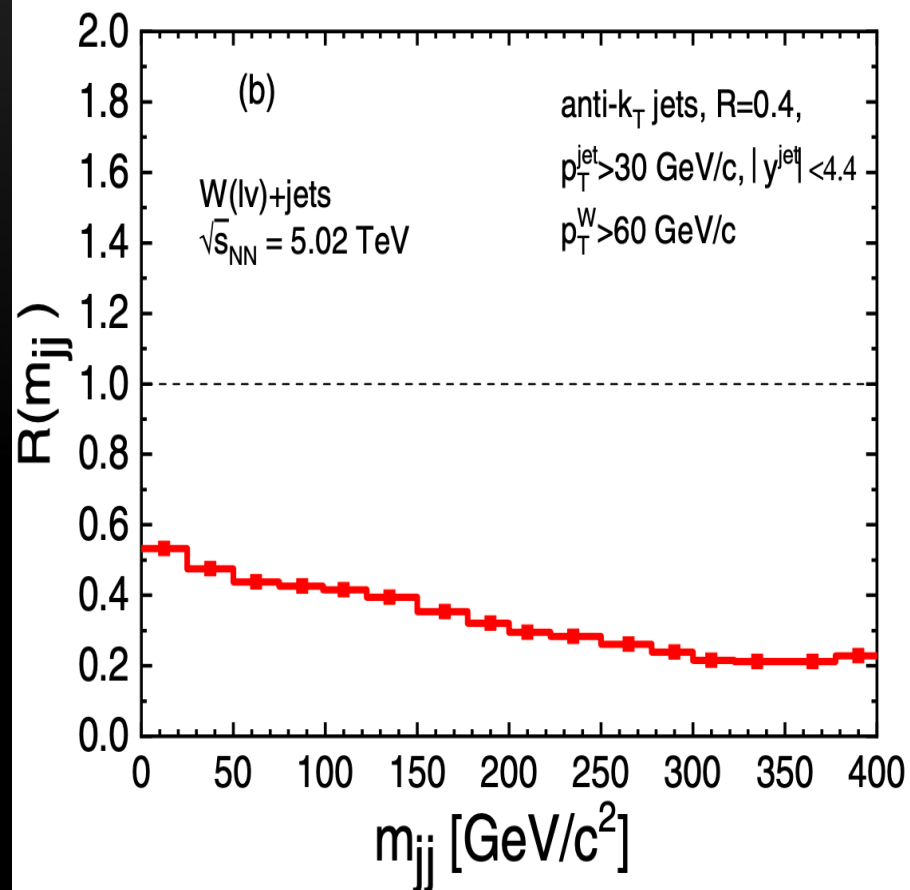
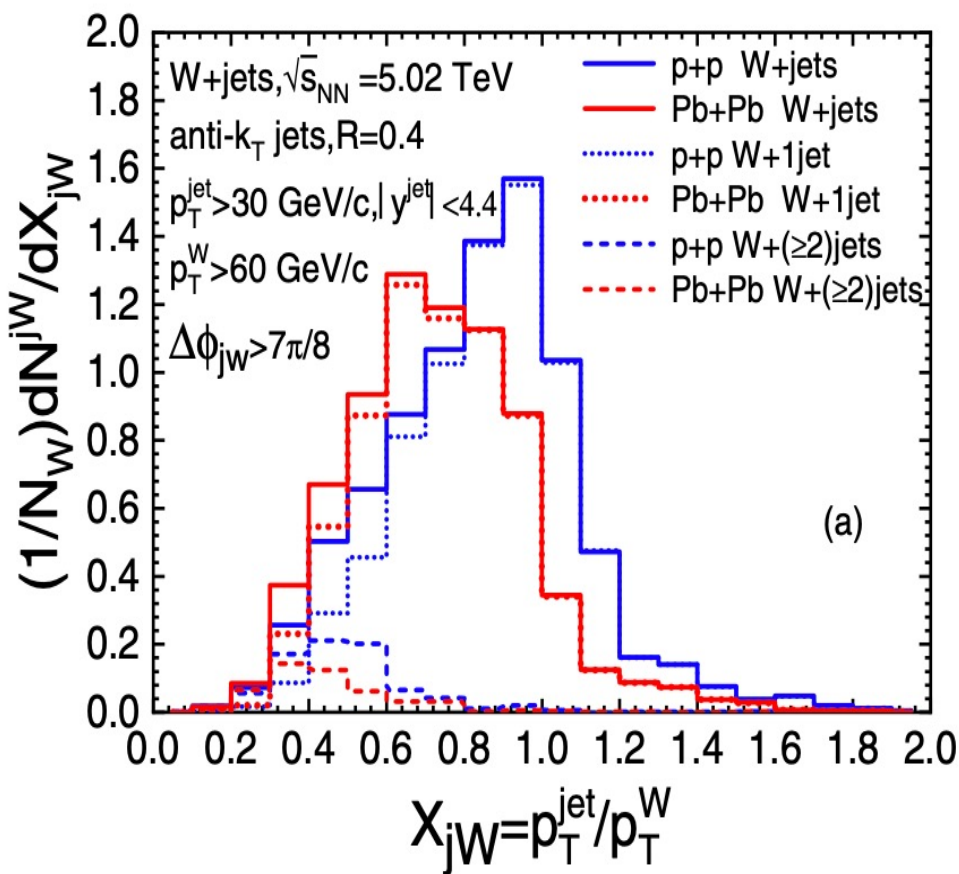
S L Zhang, T Luo, X N Wang, BWZ, PRC (2018)

W+jet in HIC at the LHC



W+jet in HIC at the LHC

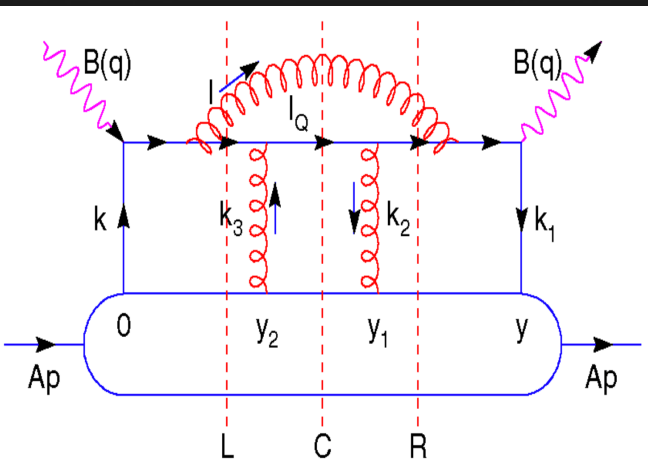
$$m_{jj} = [(E^L + E^{\text{Sub } L})^2 - (\vec{p}^L + \vec{p}^{\text{Sub } L})^2]^{1/2}$$



Heavy flavor jets

Heavy quark energy loss

- Heavy quark energy loss will be suppressed due to dead-cone effect relative to light quark.

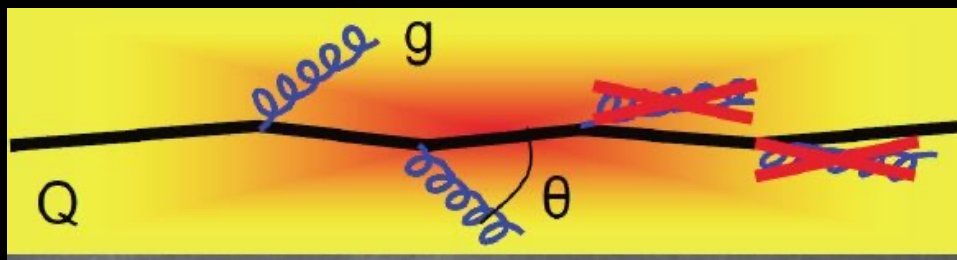


$$\frac{dN_g}{dx dk_{\perp}^2 dt} = \frac{2\alpha_s P(x) \hat{q}}{\pi k_{\perp}^4} \sin^2 \left(\frac{t - t_i}{2\tau_f} \right) \left(\frac{k_{\perp}^2}{k_{\perp}^2 + x^2 M^2} \right)^4$$

$$k_{\perp} = \omega \theta$$

$$\omega = xE$$

$$f_{Q/q} = \left(1 + \frac{\theta_0^2}{\theta^2} \right)^{-4}$$



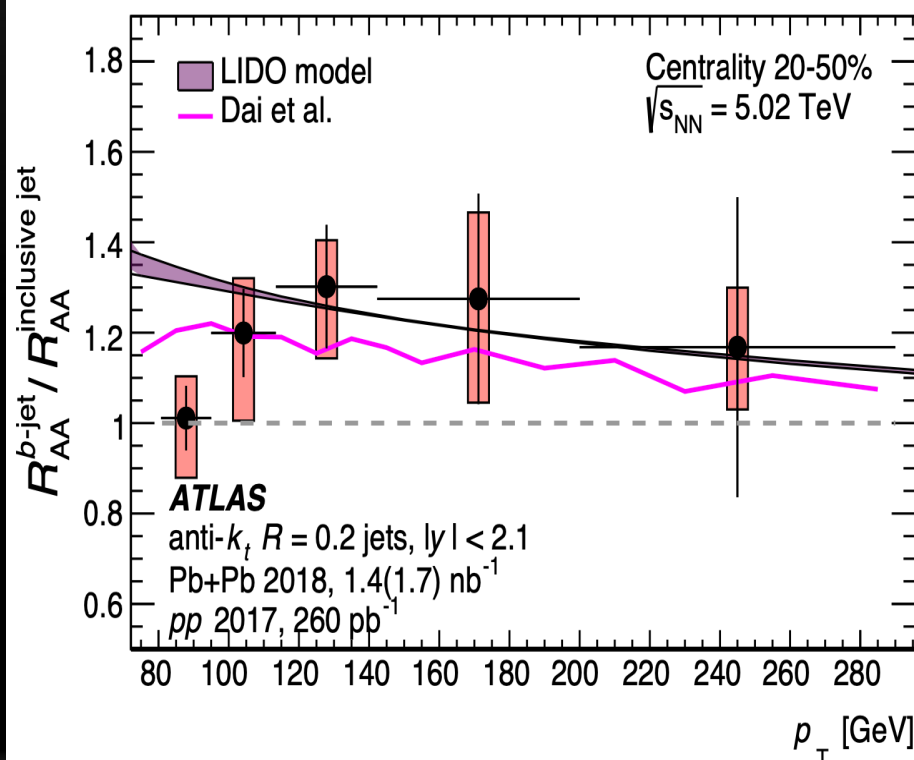
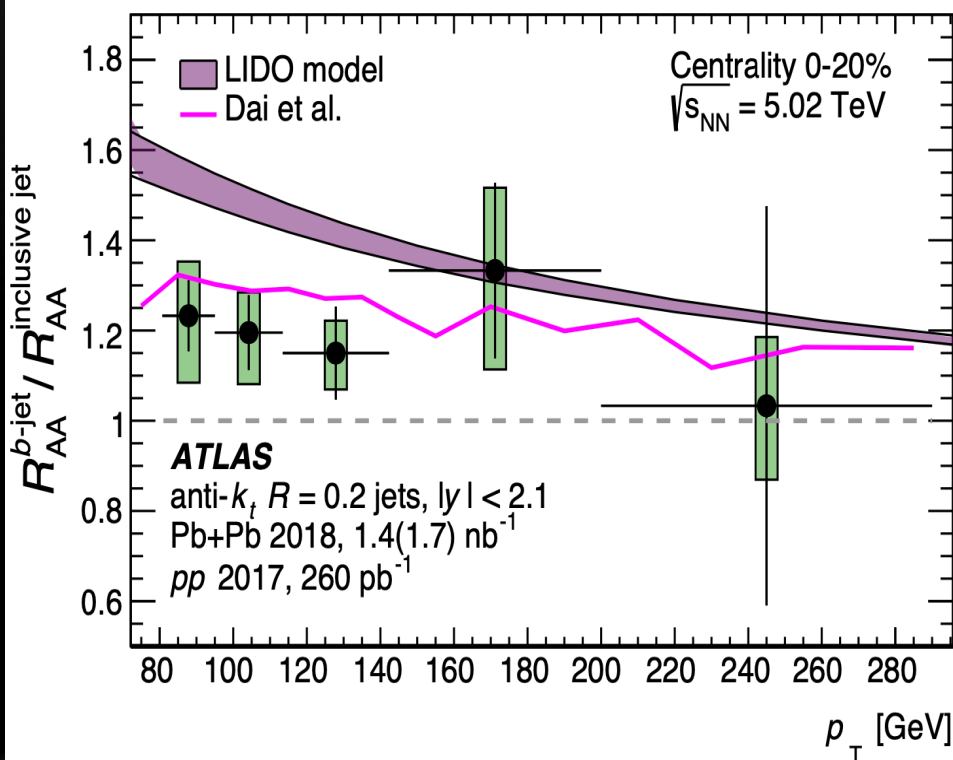
BWZ, E Wang, X N Wang, PRL (2004); NPA (2015)

Dokshitzer, Kharzeev, PLB (2011); Djordjevic, Gyulassy, PRC (2013)

Suppression of HF jets

- Heavy flavor jet should be less suppressed as compared to inclusive jets due to dead-cone effect.

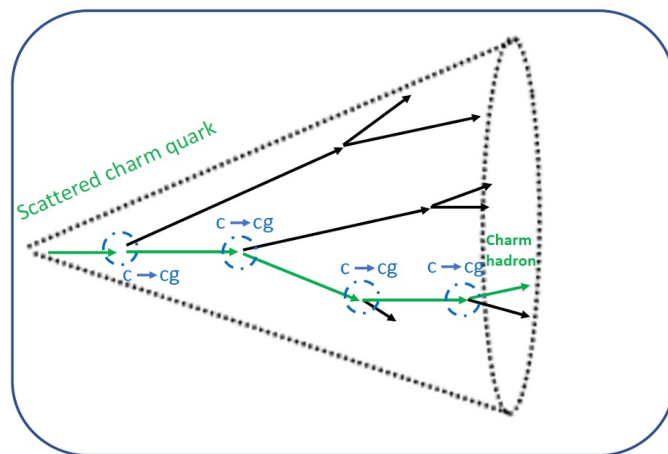
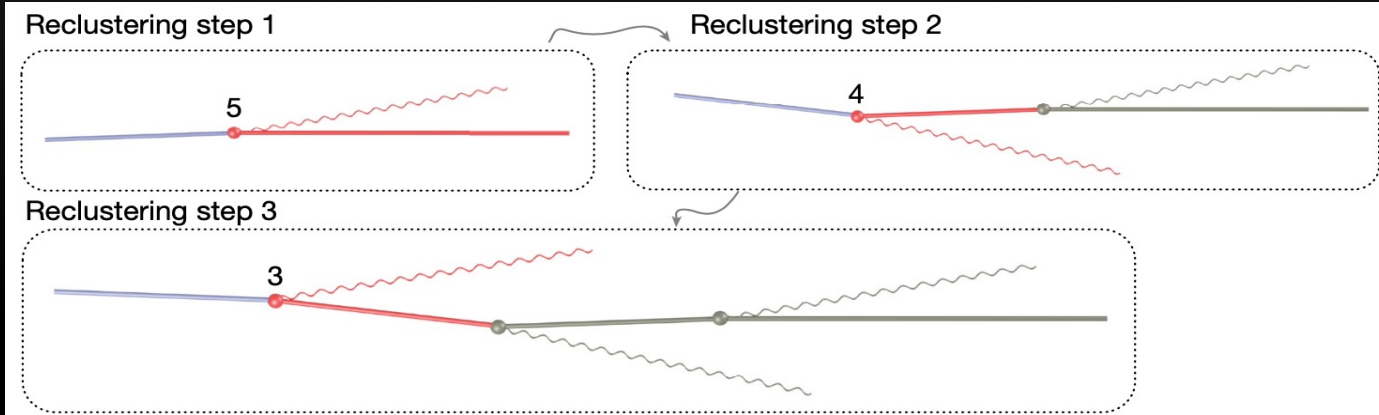
ATLAS, arXiv: 2204.13530



Dead-cone effect in vacuum

- A direct observation of dead-cone effect in p+p is made with an iterative declustering techniques by ALICE.

$$dP_{HQ} \simeq \frac{\alpha_s C_F}{\pi} \frac{d\omega}{\omega} \frac{k_{\perp}^2 dk_{\perp}^2}{(k_{\perp}^2 + \omega^2 \theta_0^2)^2} = dP_0 \left(1 + \frac{\theta_0^2}{\theta^2}\right)^2$$

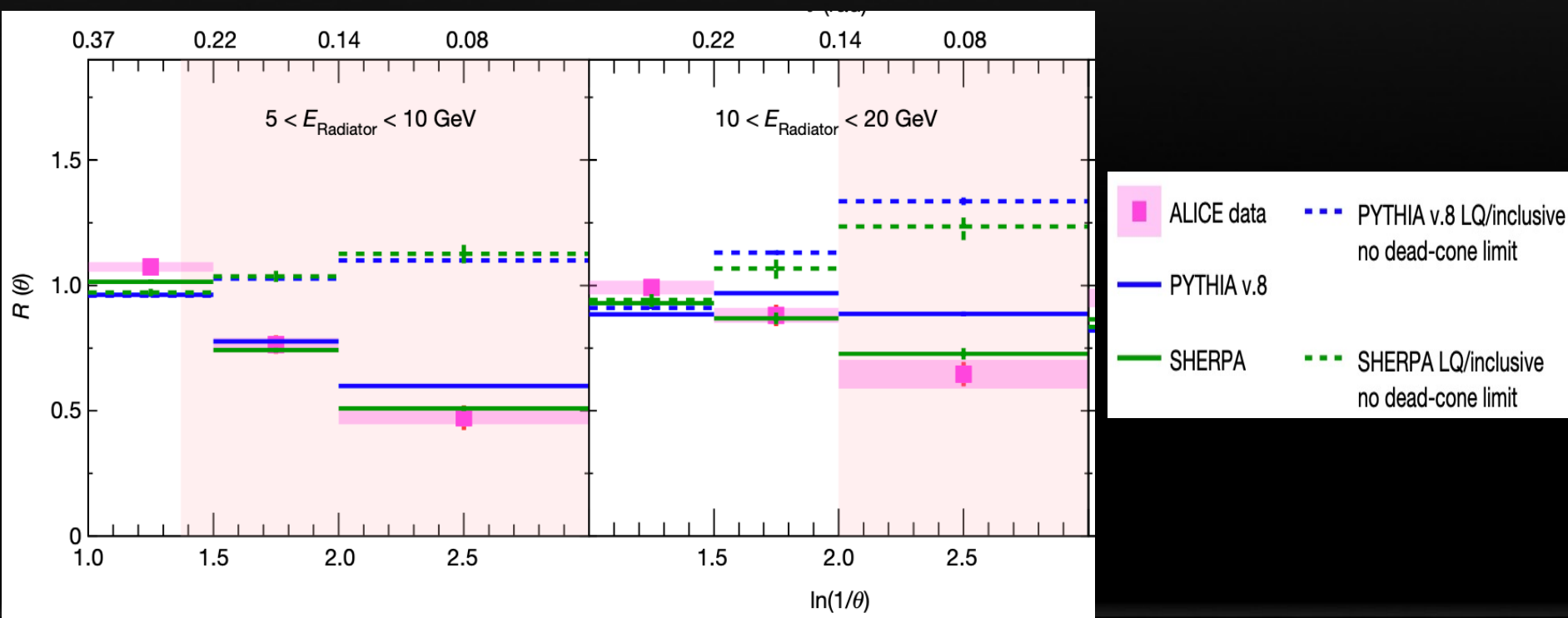


ALICE, Nature
605 (2022) 440

Dead-cone effect in vacuum

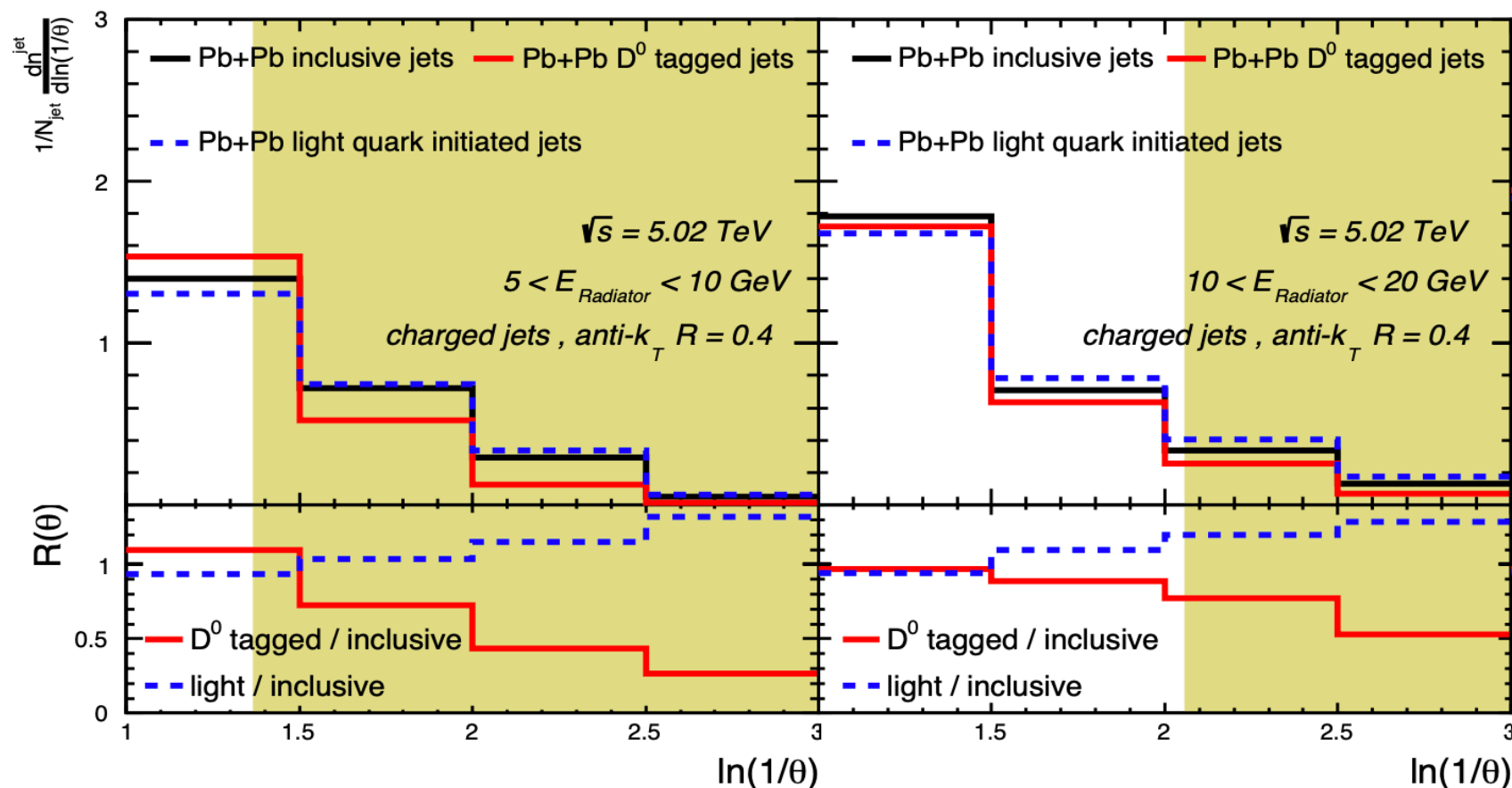
- A direct observation of dead-cone effect in p+p is made with an iterative declustering techniques by ALICE.

$$R(\theta) = \frac{1}{N^{D^0 \text{ jets}}} \frac{dn^{D^0 \text{ jets}}}{d\ln(1/\theta)} / \frac{1}{N^{\text{inclusive jet}}} \frac{dn^{\text{inclusive jet}}}{d\ln(1/\theta)}$$

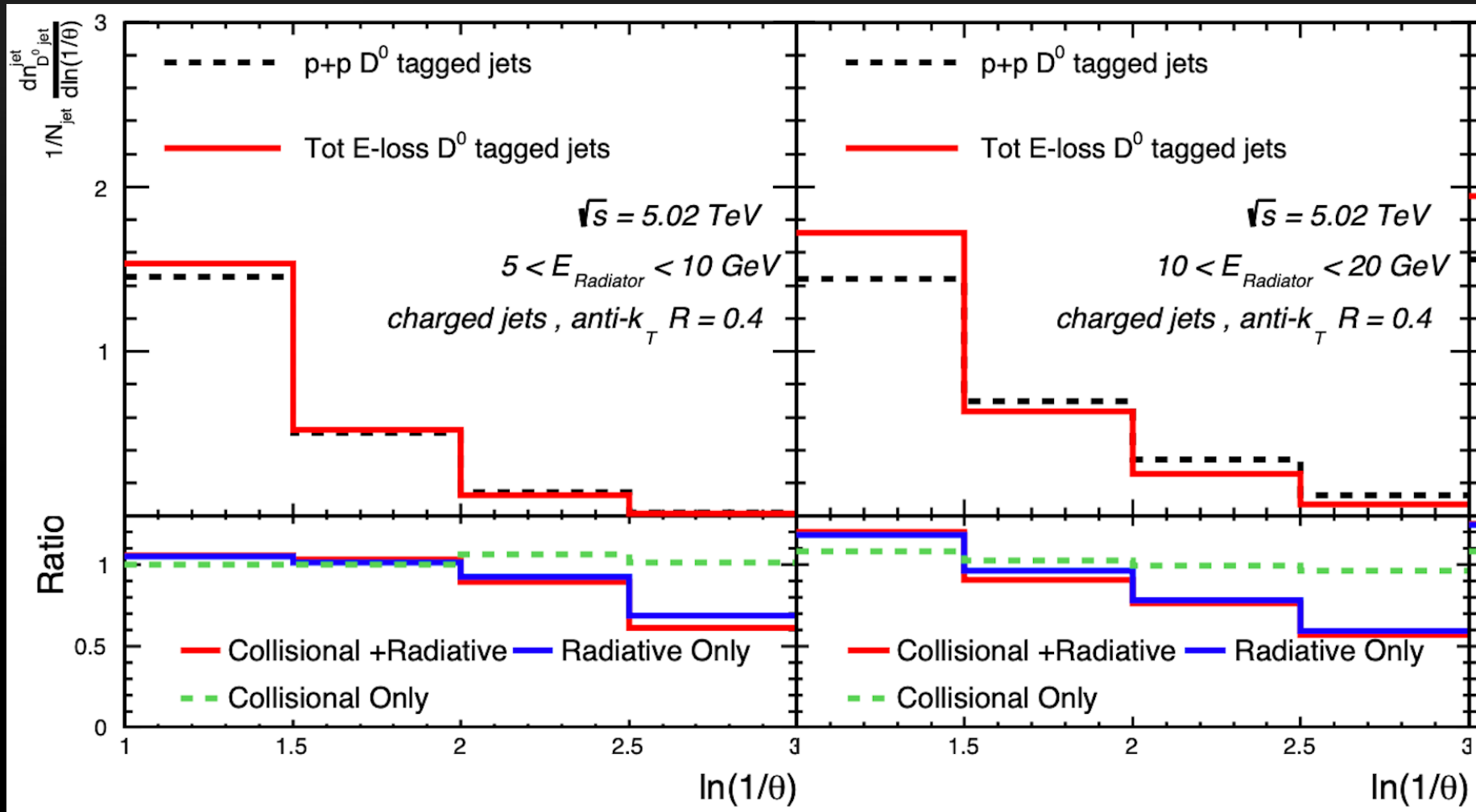


Dead-cone effect in A+A

$$R(\theta) = \frac{1}{N^{D^0 \text{ jets}}} \frac{dn^{D^0 \text{ jets}}}{d\ln(1/\theta)} / \frac{1}{N^{\text{inclusive jet}}} \frac{dn^{\text{inclusive jet}}}{d\ln(1/\theta)}$$



Dead-cone effect in A+A



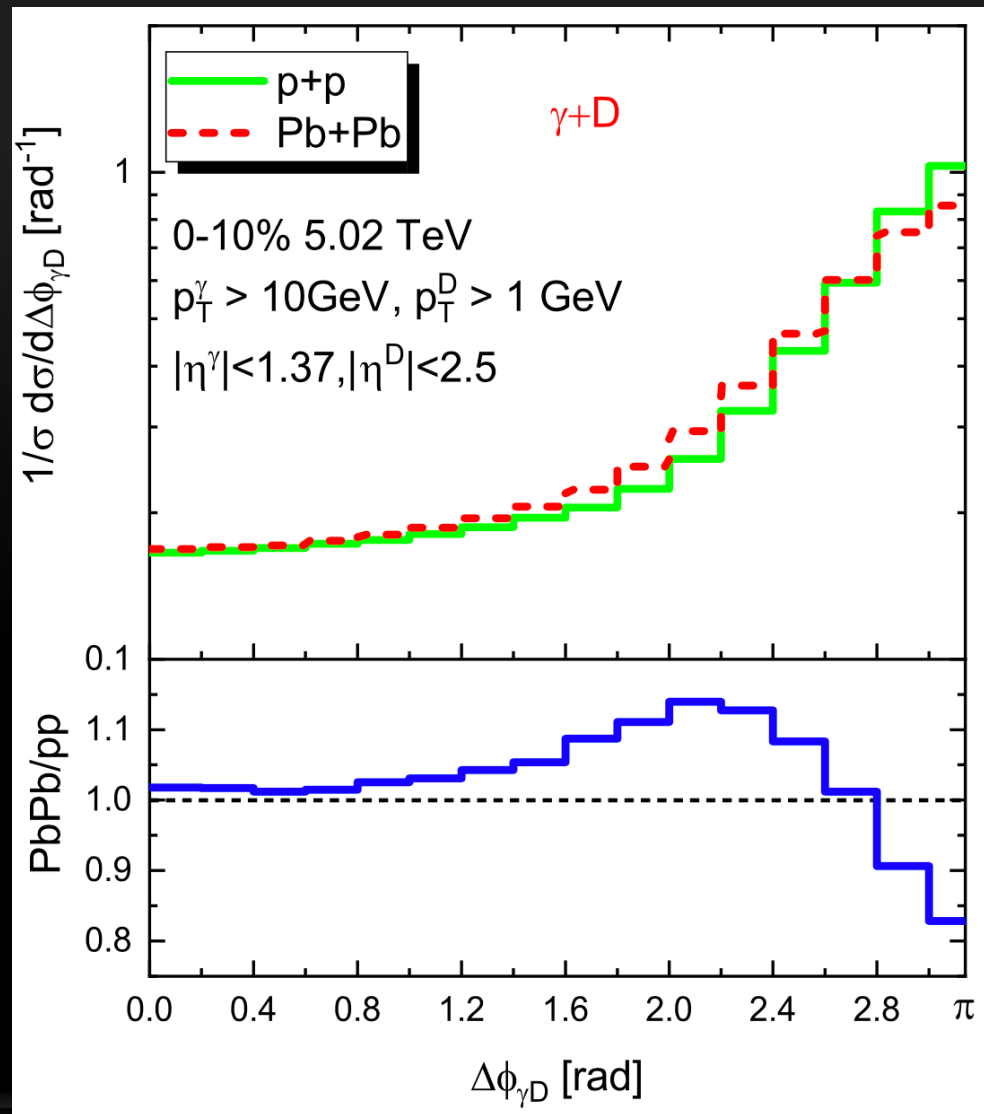
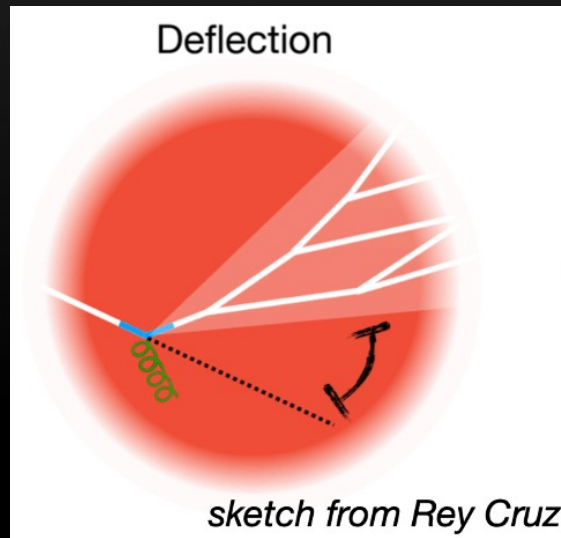
Mean value of emission angle

E_{Radiator}	Inclusive jets	D ⁰ jets	
	$\langle\theta\rangle_{\text{jets}}$	$\langle\theta\rangle_{\text{jets}}$	
5 – 10 GeV	0.31	0.34	pp
	0.36	0.36	AA
10 – 20 GeV	0.40	0.37	pp
	0.45	0.42	AA
20 – 35 GeV	0.47	0.42	pp
	0.49	0.47	AA

Mean value of emission angle

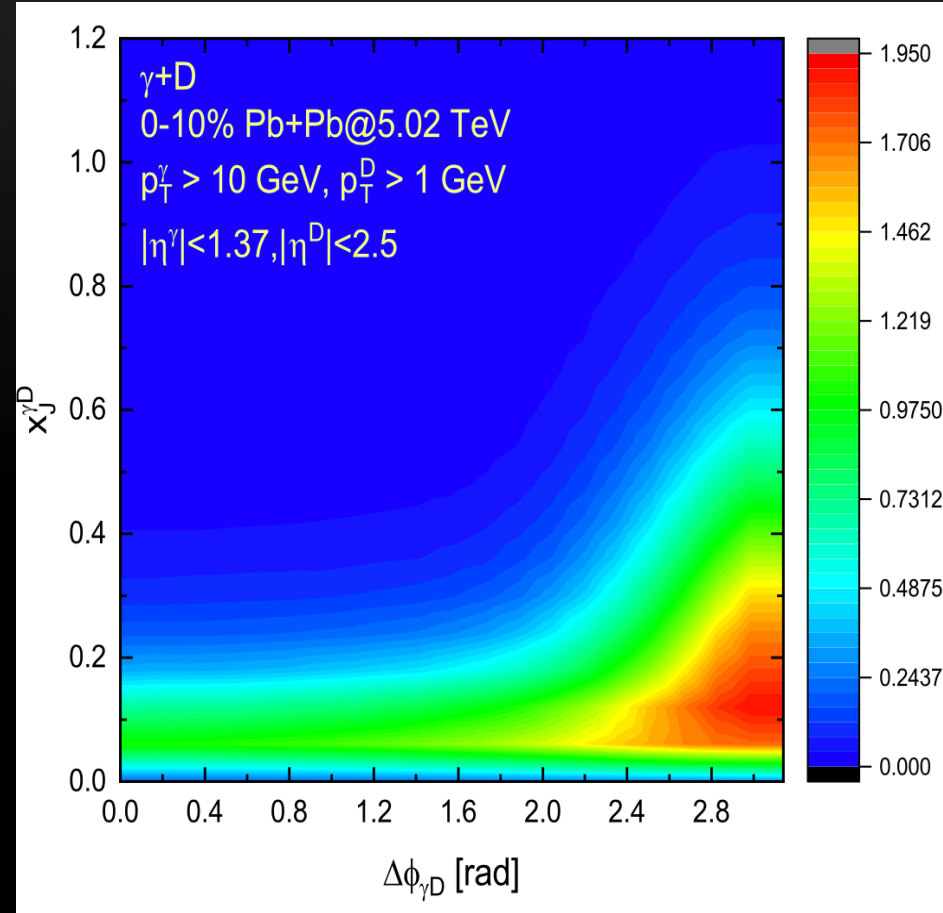
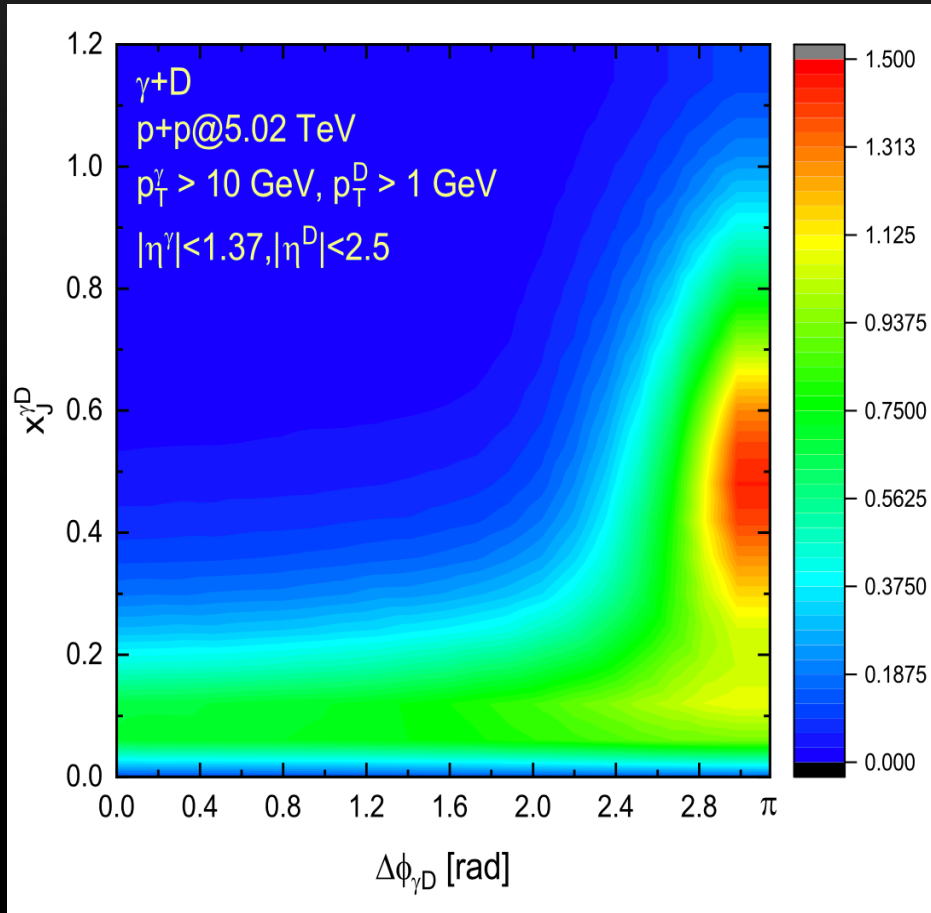
E_{Radiator}	Inclusive jets		D^0 jets		
	$\langle\theta\rangle_{\text{spl}}$	N_{spl}	$\langle\theta\rangle_{\text{spl}}$	N_{spl}	
5 – 10 GeV	0.227	1.358	0.277	1.233	pp
	0.256	1.405	0.280	1.280	AA
10 – 20 GeV	0.220	1.810	0.244	1.510	pp
	0.254	1.757	0.263	1.600	AA
20 – 35 GeV	0.232	2.040	0.232	1.822	pp
	0.249	1.977	0.251	1.860	AA

Momentum deflection of gamma+D



Momentum deflection of gamma+D

$$x_J^{\gamma D} = p_T^D / p_T^{\gamma}$$



Summary

- A systematic study of identified mesons at NLO in HIC has been made.
- A framework of combining NLO+PS for initial hard production with parton energy loss in the QGP has been developed.
- Our calculations provide nice descriptions of experimental data on full jet observables.
- A direct observation of dead-cone effect of jet quenching is possible by studying the emission angle of heavy flavor tagged jets in HIC.

Backup