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# Impact of globally spin-aligned vector mesons on the search for the chiral magnetic effect in heavy-ion collisions

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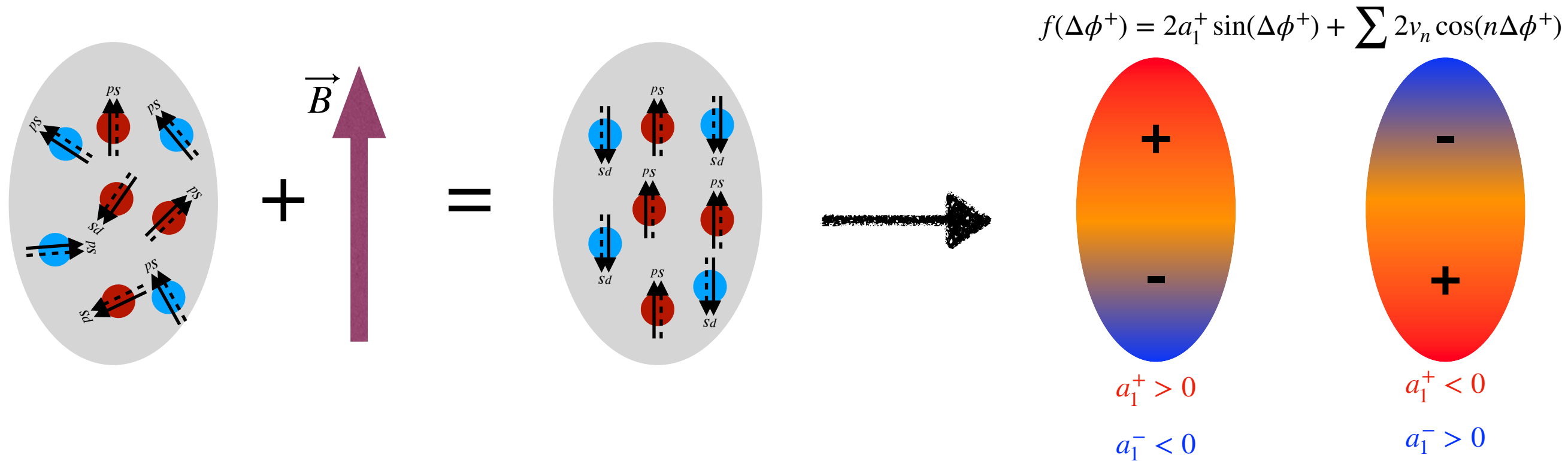
Physics Letters B 839 (2023) 137777

Diyu Shen, Jinhui Chen, Aihong Tang and Gang Wang



# The chiral magnetic effect

**Chirality imbalance + external magnetic field = electrical current**



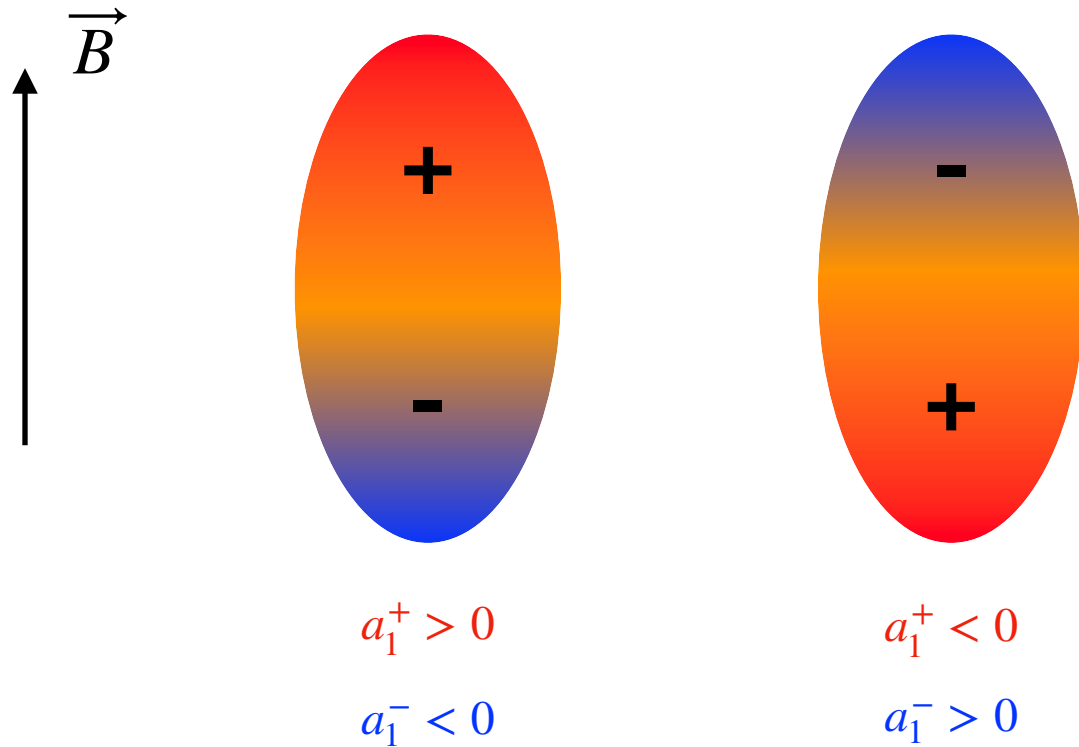
- Electrical current along the  $\vec{B}$  field generates charge separation.
- Charge dipole gives non-zero  $a_1$  ( $-a_1$ ) to positively (negatively) charged particles, while the  $a_1$  fluctuates event by event so on average  $\langle a_1 \rangle = 0$ .
- Observables are developed in experiment to measure the  $a_1$  and/or charge separation.

D.E. Kharzeev, J. Liao, Nat. Rev. Phys. 3 (2021) 55–63

# The chiral magnetic effect

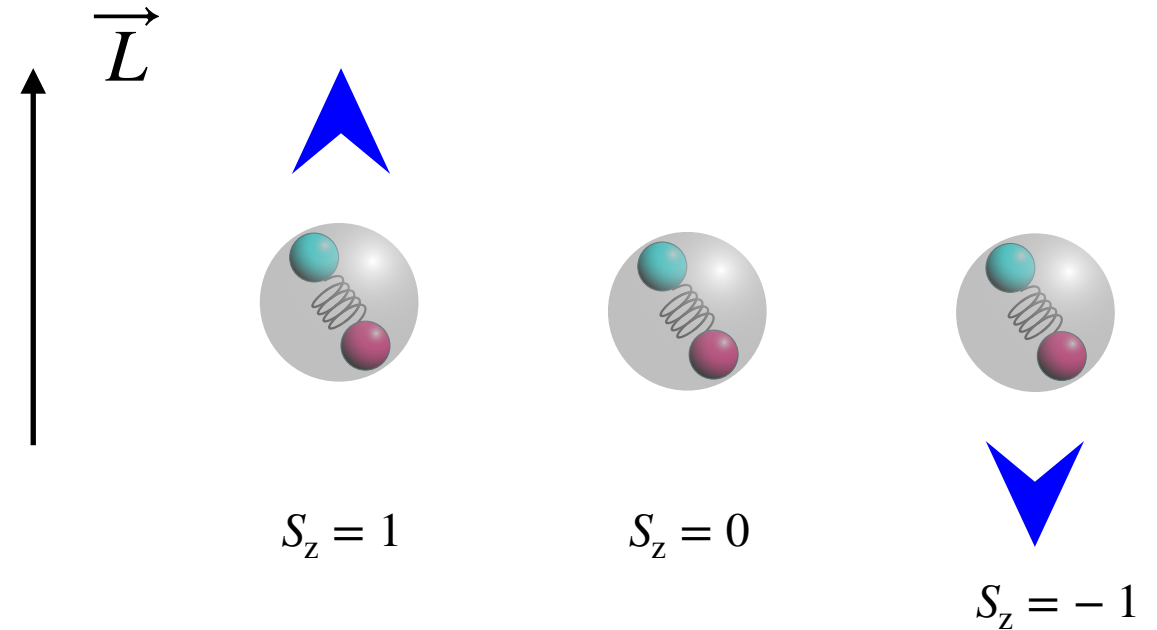
Z.T. Liang et al., Physics Letters B 629 (2005) 20–26

## CME



- Charge separation along  $\vec{B}$  field, characterized by non-zero  $a_1$ .

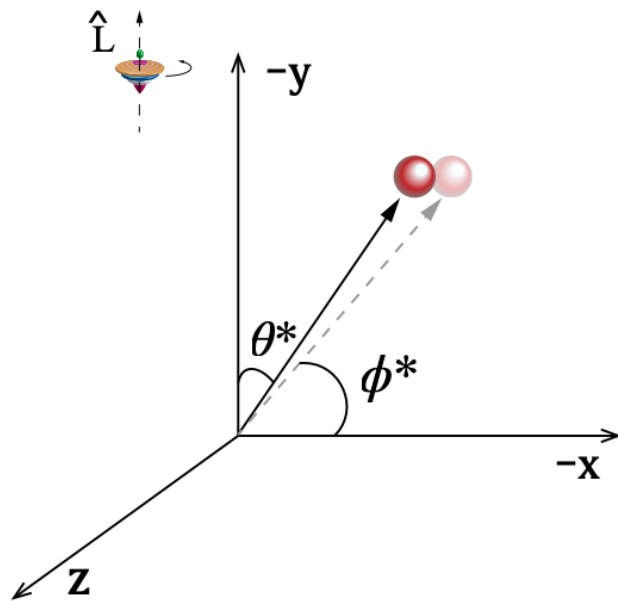
## Global spin alignment of vector meson



- Spin state along orbital angular momentum, characterized by  $\rho_{00}$  in spin density matrix.

Directions of  $\vec{B}$  and  $\vec{L}$  are correlated, both are perpendicular to reaction plane

# Global spin alignment

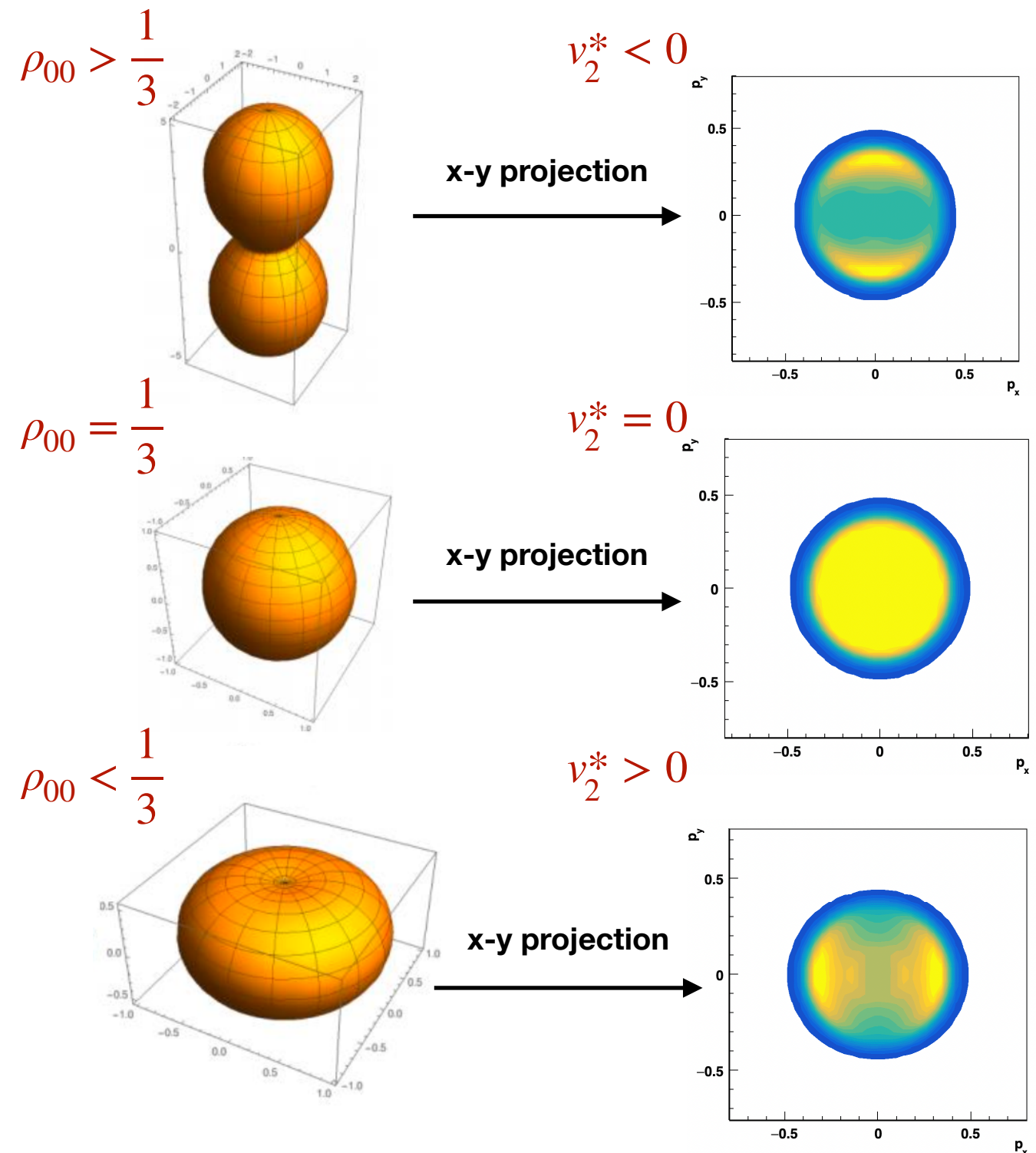


$$\frac{dN}{d \cos \theta^*} = \frac{3}{4} [(1 - \rho_{00}) + (3\rho_{00} - 1) \cos^2 \theta^*]$$

$$\frac{dN}{d\phi^*} = \frac{1}{2\pi} \left[ 1 - \frac{1}{2}(3\rho_{00} - 1) \cos 2\phi^* \right]$$

$$\therefore v_2^* = -\frac{3\rho_{00} - 1}{4}$$

$\rho_{00} \neq 1/3$  means non-zero “elliptic flow” of decay products in CMS.



# Global spin alignment to the $\Delta\gamma$ observable

$$\gamma_{112} = \langle \cos(\phi_\alpha + \phi_\beta - 2\Psi_{RP}) \rangle \quad \text{S.A. Voloshin, Phys. Rev. C 70 (2004) 057901}$$

$$\begin{aligned} \gamma_{112}^{\text{OS}} &= \langle \cos(\phi_+ + \phi_- - 2\Psi_{RP}) \rangle \\ &= \langle \cos \Delta\phi_+ \rangle \langle \cos \Delta\phi_- \rangle + \frac{N_\rho}{N_+ N_-} \text{Cov}(\cos \Delta\phi_+, \cos \Delta\phi_-) - \langle \sin \Delta\phi_+ \rangle \langle \sin \Delta\phi_- \rangle - \frac{N_\rho}{N_+ N_-} \text{Cov}(\sin \Delta\phi_+, \sin \Delta\phi_-) \end{aligned}$$


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$$\langle ab \rangle = \langle a \rangle \langle b \rangle + \text{Cov}(a, b)$$

In center-of-mass system, the covariance between  $\pi^+$  and  $\pi^-$  from  $\rho$  decays is

$$\text{Cov}(\cos \phi_+^*, \cos \phi_-^*) = - \langle \cos^2 \phi_+^* \rangle + \langle \cos \phi_+^* \rangle^2 = -\frac{1}{2} + \frac{1}{8}(3\rho_{00} - 1),$$

$$\text{Cov}(\sin \phi_+^*, \sin \phi_-^*) = - \langle \sin^2 \phi_+^* \rangle + \langle \sin \phi_+^* \rangle^2 = -\frac{1}{2} - \frac{1}{8}(3\rho_{00} - 1).$$

Where we used  $\phi_+^* = \phi_-^* + \pi$  and  $\frac{dN}{d\phi^*} = \frac{1}{2\pi} \left[ 1 - \frac{1}{2}(3\rho_{00} - 1) \cos 2\phi^* \right]$ ,

# Global spin alignment to the $\Delta\gamma$ observable

$$\text{Therefore, } \Delta\gamma^* = \gamma^{*\text{OS}} - \gamma^{*\text{SS}} = \frac{N_\rho}{N_+N_-} \frac{3\rho_{00} - 1}{4} = -\frac{N_\rho}{N_+N_-} v_2^*$$

The  $\Delta\gamma$  is proportional to the “elliptic flow” of decay products in  $\rho$  rest frame.

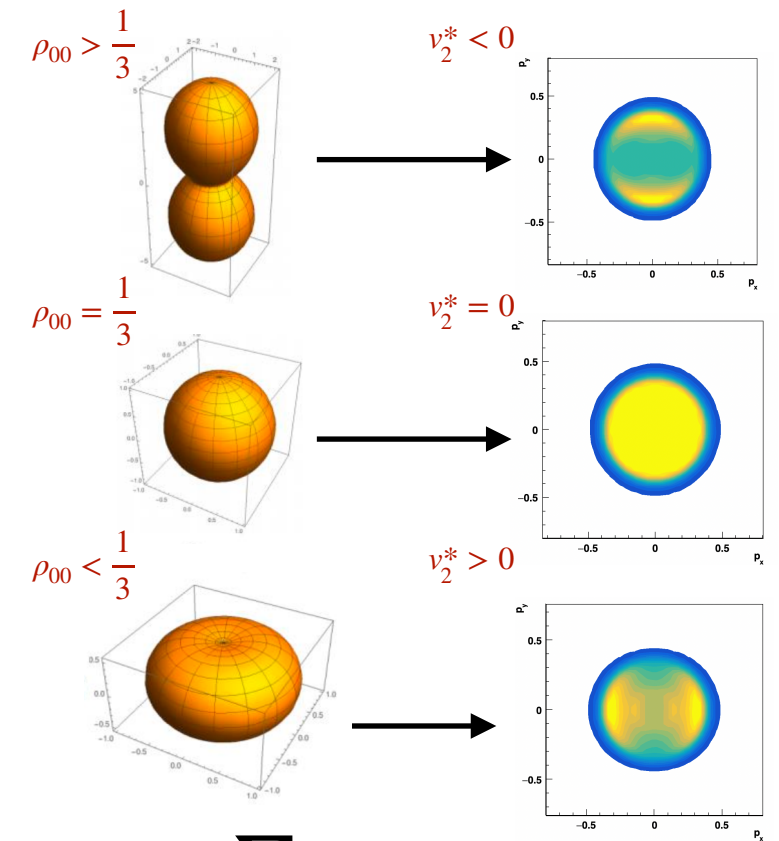
In lab frame,

$$\text{Cov}(\cos \phi_+, \cos \phi_-) = \underbrace{f_c}_{\text{Boost factor in plane}} \text{Cov}(\cos \phi_+^*, \cos \phi_-^*) = f_c \left[ -\frac{1}{2} + \frac{1}{8}(3\rho_{00} - 1) \right]$$

$$\text{Cov}(\sin \phi_+, \sin \phi_-) = \underbrace{f_s}_{\text{Boost factor out of plane}} \text{Cov}(\sin \phi_+^*, \sin \phi_-^*) = f_s \left[ -\frac{1}{2} - \frac{1}{8}(3\rho_{00} - 1) \right]$$

$$\Delta\gamma_{112} = \frac{N_\rho}{N_+N_-} \left[ \frac{1}{8}(f_c + f_s)(3\rho_{00} - 1) - \frac{1}{2}(f_c - f_s) \right] \sim \frac{N_\rho}{N_+N_-} \left[ A_1(3\rho_{00} - 1) + A_2 v_2^\rho \right]$$

**A linear dependence of  $\Delta\gamma_{112}$  on  $\rho_{00}$ .**



$$f_c = f_0 + \sum a_n (v_2^\rho)^n$$

$$f_s = f_0 + \sum b_n (v_2^\rho)^n$$

# Global spin alignment to the R correlator

Definition: N. Magdy, Phys. Rev. C 97 (2018) 061901

$$R_{\Psi_2}(\Delta S) \equiv \frac{N(\Delta S_{\text{real}})}{N(\Delta S_{\text{shuffled}})} / \frac{N(\Delta S_{\text{real}}^\perp)}{N(\Delta S_{\text{shuffled}}^\perp)},$$

$$\Delta S = \langle \sin \Delta \phi_+ \rangle - \langle \sin \Delta \phi_- \rangle,$$

$$\Delta S^\perp = \langle \cos \Delta \phi_+ \rangle - \langle \cos \Delta \phi_- \rangle,$$

$$\text{Cov}(\langle \sin \Delta \phi_+ \rangle, \langle \sin \Delta \phi_- \rangle)$$

$$\sigma^2(\Delta S_{\text{real}}) = f_s \left[ \sigma_s^2 + \frac{N_\rho}{N_+ N_-} \left( 1 + \frac{3\rho_{00} - 1}{4} \right) \right],$$

$$\sigma^2(\Delta S_{\text{shuffled}}) = f_s \sigma_s^2,$$

$$\sigma^2(\Delta S_{\text{real}}^\perp) = f_c \left[ \sigma_c^2 + \frac{N_\rho}{N_+ N_-} \left( 1 - \frac{3\rho_{00} - 1}{4} \right) \right],$$

$$\sigma^2(\Delta S_{\text{shuffled}}^\perp) = f_c \sigma_c^2,$$

$$\frac{S_{\text{concavity}}}{\sigma_R^2} = \frac{1}{\sigma^2(\Delta S_{\text{real}})} - \frac{1}{\sigma^2(\Delta S_{\text{shuffled}})} - \frac{1}{\sigma^2(\Delta S_{\text{real}}^\perp)} + \frac{1}{\sigma^2(\Delta S_{\text{shuffled}}^\perp)}.$$

$$\text{Sign}(S_{\text{concavity}}) = \text{Sign} \left[ -\frac{N_\rho}{2N_+ N_-} (3\rho_{00} - 1) \right]$$

$$\Delta \sigma_R^2 = \sigma^2(\Delta S_{\text{real}}) - \sigma^2(\Delta S_{\text{shuffled}}) - \sigma^2(\Delta S_{\text{real}}^\perp) + \sigma^2(\Delta S_{\text{shuffled}}^\perp).$$

$$\Delta \sigma_R^2 = \frac{N_\rho}{N_+ N_-} \left[ \frac{1}{4} (f_c + f_s) (3\rho_{00} - 1) + (f_c - f_s) \right].$$

**A linear dependence of  $\Delta \sigma_R^2$  on  $\rho_{00}$ .**

# Global spin alignment to the signed balance function

## Signed balance function

A. H. Tang, Chin. Phys. C 44 054101

$$\begin{aligned}\Delta B_y &\equiv \left[ \frac{N_{y(+ -)} - N_{y(++ )}}{N_+} - \frac{N_{y(- +)} - N_{y(-- )}}{N_-} \right] \\ &\quad - \left[ \frac{N_{y(- +)} - N_{y(++ )}}{N_+} - \frac{N_{y(+ -)} - N_{y(-- )}}{N_-} \right] \\ &= \frac{N_+ + N_-}{N_+ N_-} [N_{y(+ -)} - N_{y(- +)}],\end{aligned}$$

$$r \equiv \sigma(\Delta B_y) / \sigma(\Delta B_x).$$

Assuming all particles have same pT, we will have

$$\sigma^2(\Delta B_y) \approx \frac{64M^2}{\pi^4} \left( \frac{4}{9M} + 1 + \frac{4}{3}v_2 \right) \sigma^2(\Delta S_{\text{real}}),$$

$$\sigma^2(\Delta B_x) \approx \frac{64M^2}{\pi^4} \left( \frac{4}{9M} + 1 - \frac{4}{3}v_2 \right) \sigma^2(\Delta S_{\text{real}}^\perp).$$

$$\begin{aligned}\Delta \sigma^2(\Delta B) &\equiv \sigma^2(\Delta B_y) - \sigma^2(\Delta B_x) \\ &\approx c_1 + c_2(3\rho_{00} - 1),\end{aligned}$$

**A linear dependence of  $\Delta \sigma^2(\Delta B)$  on  $\rho_{00}$ .**

## Setups of toy model

- Spectrum of primordial pion

$$\frac{dN_{\pi^\pm}}{dm_T^2} \propto \frac{1}{e^{m_T/T_{BE}} - 1},$$

- Spectrum of  $\rho$

$$\frac{dN_\rho}{dm_T^2} \propto \frac{e^{-(m_T - m_\rho)/T}}{T(m_\rho + T)},$$

- 195 pairs of  $\pi^+\pi^-$  with 33 from  $\rho$  decays
- $v_2$  and  $v_3$  of primordial pions are set to zero.
- Spin alignment effect is introduced by sampling decay products according to

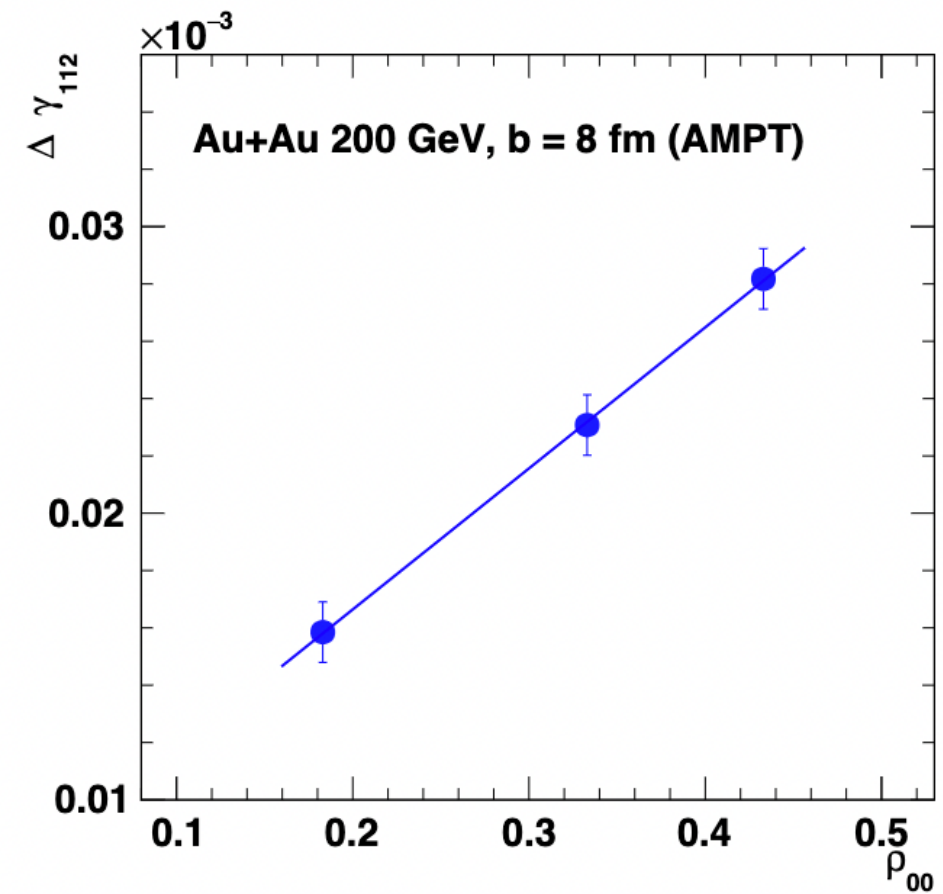
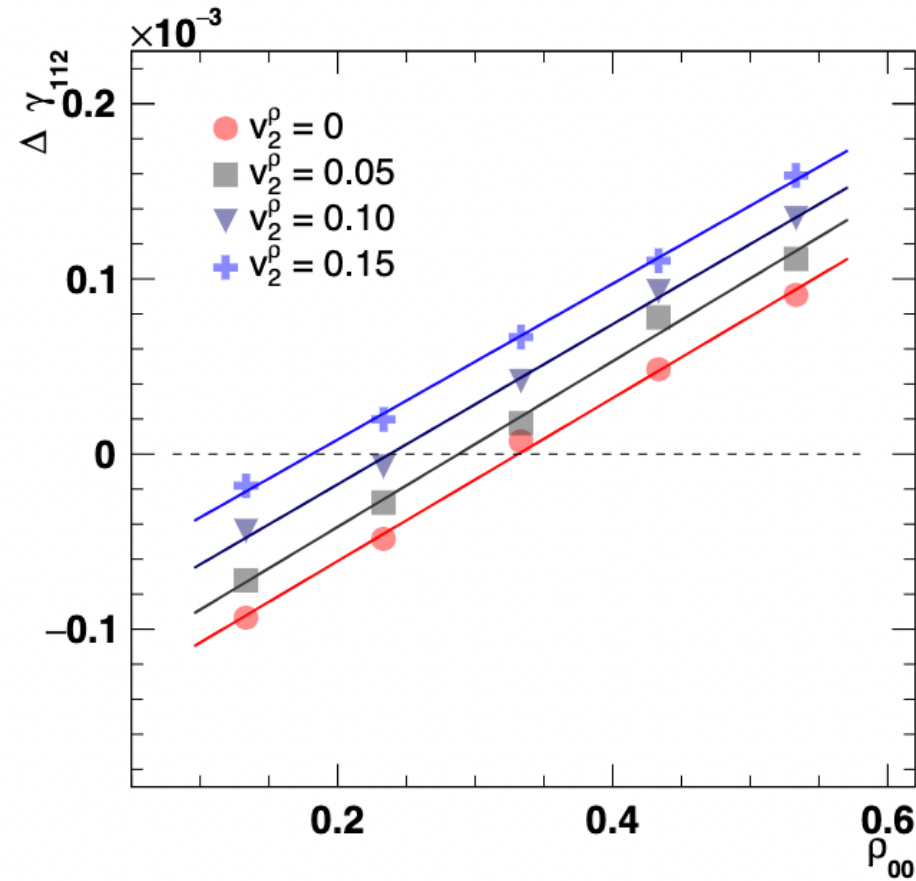
$$\frac{dN}{d\cos\theta^*} = \frac{3}{4} [(1 - \rho_{00}) + (3\rho_{00} - 1) \cos^2 \theta^*]$$

## Setups of AMPT

- String melting version
- AuAu 200 GeV with impact parameter  $b \sim 8$  fm
- Spin alignment effect is introduced by sampling decay products according to

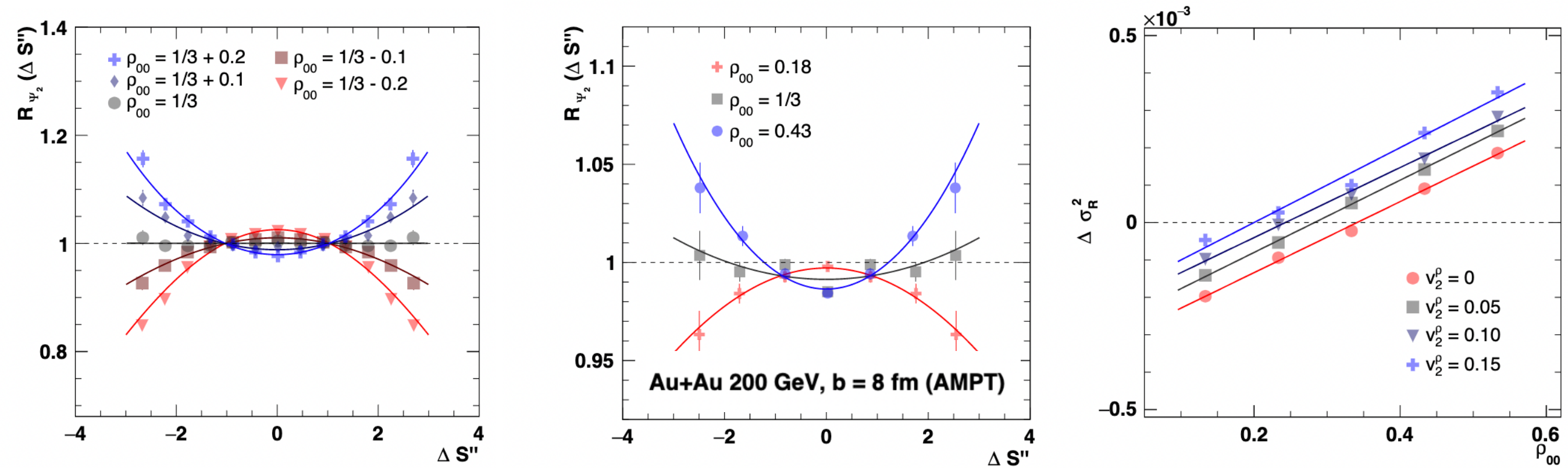
$$\frac{dN}{d\cos\theta^*} = \frac{3}{4} [(1 - \rho_{00}) + (3\rho_{00} - 1) \cos^2 \theta^*]$$

# Model simulations



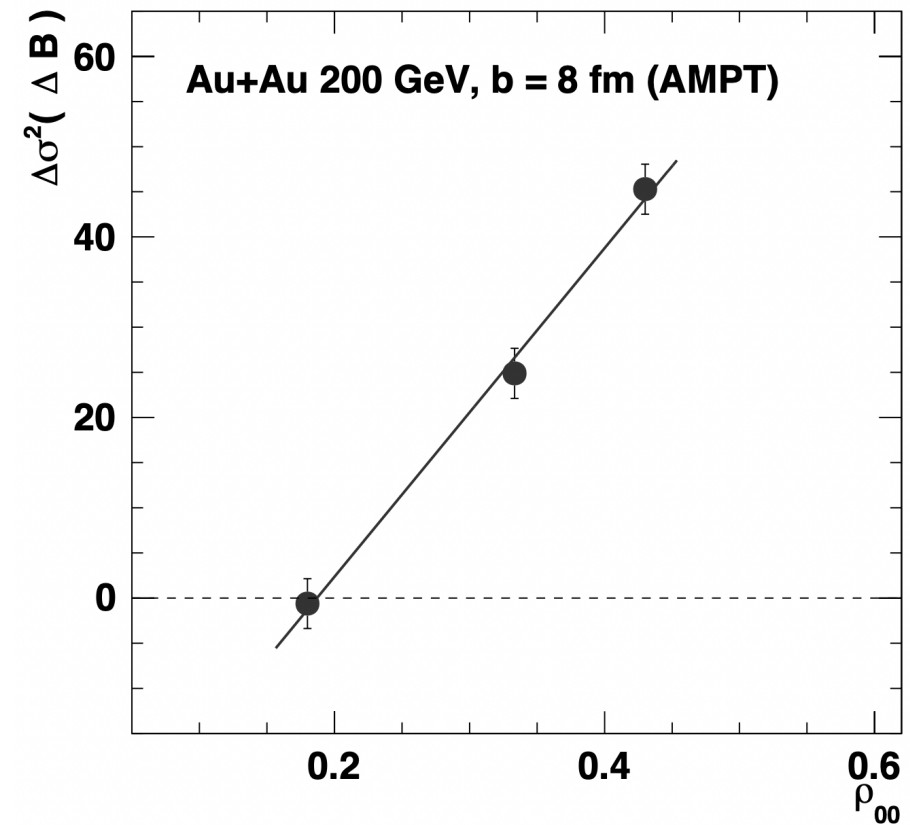
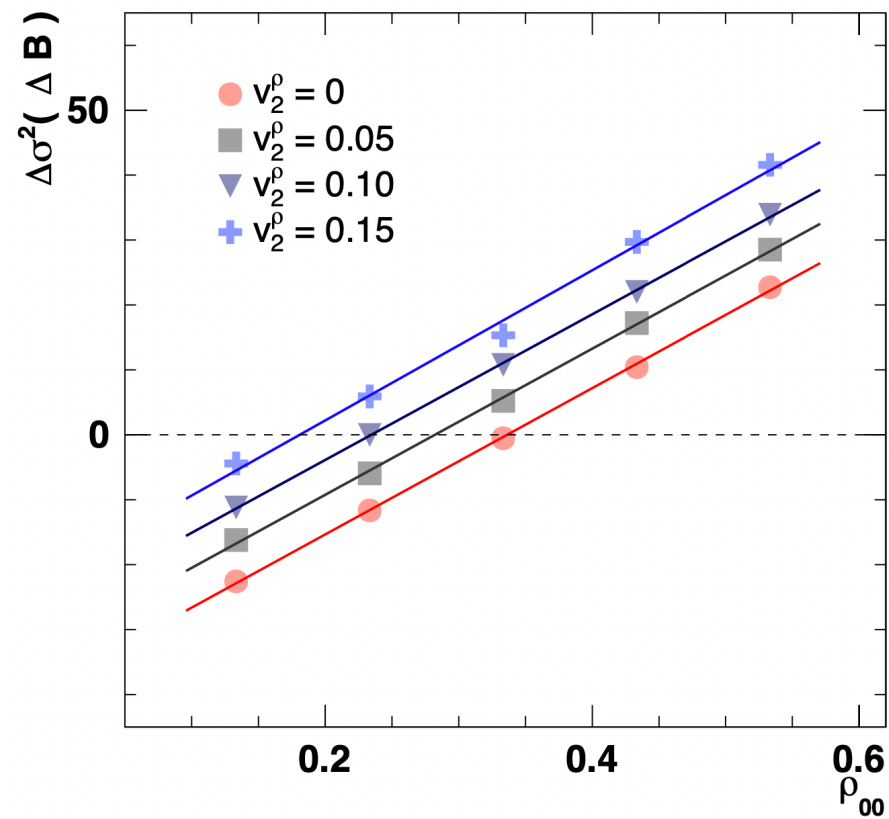
**A linear dependence of  $\Delta \gamma$  as a function of  $\rho_{00}$  has been observed, slope and intercept depend on spectra and flow of  $\rho$  mesons in models.**

# Model simulations



$R_{\Psi_2}(\Delta S)$  has similar  $\rho_{00}$  dependence as gamma,  $\Delta \sigma_R^2$  is also a linear function.

# Model simulations



**Signed balance function is also sensitive to  $\rho_{00}$ , the  $\Delta\sigma^2(\Delta B)$  is also a linear function.**

# Summary

- The  $\Delta\gamma_{112}$ ,  $R_{\Psi_2}(\Delta S)$  and signed balance function  $r_{\text{lab}}$  are both influenced by the spin alignment  $\rho_{00}$  of vector mesons.
- If the  $\rho_{00}$  is smaller (larger) than  $1/3$ , it gives negative (positive) signal to the  $\Delta\gamma_{112}$ ,  $R_{\Psi_2}(\Delta S)$  and  $r_{\text{lab}}$ .
- The spin alignment of  $\rho_0$  meson is critical for the CME measurements.

