

Spectra and flow of magnetised lepton pairs

Chowdhury Aminul Islam

7th International Conference on Chirality, Vorticity and
Magnetic Field

University of Chinese Academy of Sciences
Yanqi Lake Campus

19.07.2023

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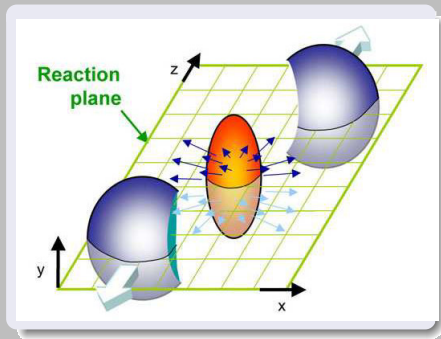
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Noncentral Heavy Ion collisions:

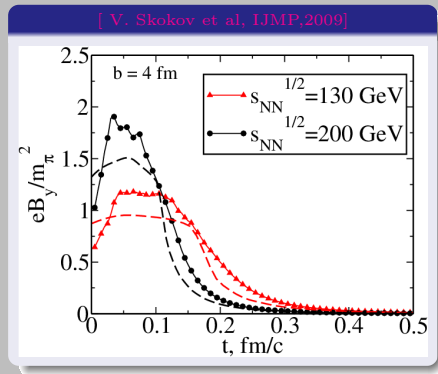


Pictorial representation of noncentral heavy ion collisions.

Production of a strong magnetic field in HICs

- A very strong magnetic field ($\approx m_\pi^2$ at RHIC and $\approx 10 m_\pi^2$ at LHC) is generated in the direction perpendicular to the reaction plane, due to the relative motion of the ions themselves.

$$(m_\pi^2 = 1.96 \times 10^{-2} \text{ GeV}^2 \approx 10^{18} \text{ Gauss})$$

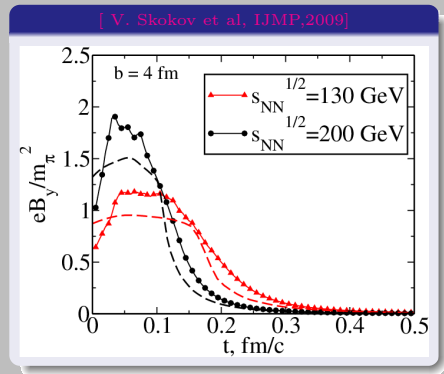


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- A comparison with other terrestrial strengths: Earth $\approx 10^{-18} m_\pi^2$, usual laboratory $\approx 10^{-13} m_\pi^2$, max.
- A magnetar: $\approx 10^{-5} - 10^{-3} m_\pi^2$.

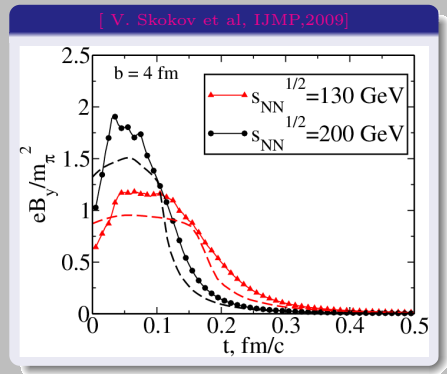


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- The presence of an external field in the medium subsequently requires modification of the present theoretical tools.

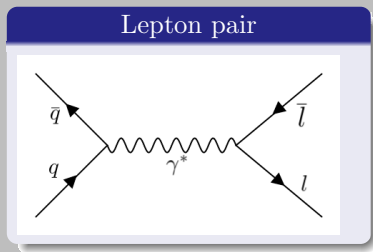


EM probes, particularly leptons:

- Photons and leptons (virtual photons) can probe the interior of a QCD medium.
- They are produced from multiple stages.
- They can be used to extract information from the hot and dense matter.
- We are particularly interested in the thermal leptons.

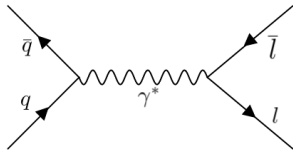
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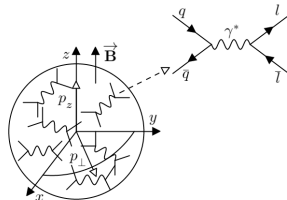


Lepton pair from the magnetised medium:

Lepton pair



A magnetised medium

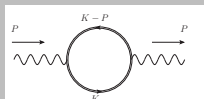


- A straightforward question is to ask whether the dilepton production will be affected by the magnetic field.

Methodology:

- This dilepton rate (DR) is given by [H. A. Weldon PRD 42, 2384].

$$\frac{dN}{d^4X d^4P} \equiv \frac{dR}{d^4P} = \frac{\alpha_{\text{EM}}}{12\pi^3} \frac{1}{P^2} \frac{1}{e^{p_0/T} - 1} \sum_{f=u,d} \frac{1}{\pi} \text{Im} \Pi_{\mu,f}^{\mu}(P). \quad (1)$$



- We can use the fermionic propagator, depending on the scenarios, to write down the one loop electromagnetic (EM) polarization tensor

$$\Pi_f^{\mu\nu}(P) = -iN_c q_f^2 \int \frac{d^4K}{(2\pi)^4} \text{Tr} [\gamma^\mu S_f^B(K) \gamma^\nu S_f^B(K-P)]. \quad (2)$$

Fermionic propagator in presence of eB :

- Schwinger propagator in momentum space as

$$S_f^{(B)}(K) = \exp\left(-\frac{k_\perp^2}{|q_f B|}\right) \sum_{\ell=0}^{\infty} (-1)^\ell \frac{D_\ell(K, q_f B)}{k_\parallel^2 - 2\ell|q_f B| - m_f^2 + i\epsilon}, \quad (3)$$

where

$$\begin{aligned} D_\ell(K, q_f B) = & (k_\parallel + m_f) \left\{ L_\ell\left(\frac{2k_\perp^2}{|q_f B|}\right) \left[\mathbb{1} - i\gamma^1 \gamma^2 \text{sgn}(q_f B)\right] - L_{\ell-1}\left(\frac{2k_\perp^2}{|q_f B|}\right) \left[\mathbb{1} + i\gamma^1 \gamma^2 \text{sgn}(q_f B)\right] \right\} \\ & + 4k_\perp L_{\ell-1}^1\left(\frac{2k_\perp^2}{|q_f B|}\right). \end{aligned} \quad (4)$$

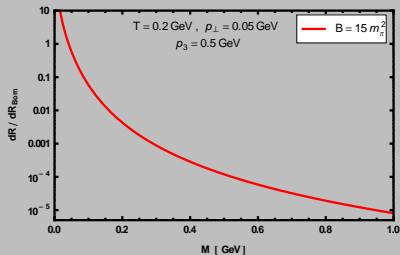
- m_f and q_f are the mass and charge of the fermion of flavor f , respectively; ℓ denotes the Landau level index.

Work done so far:

- So far dilepton rate from a magnetised medium has been calculated in different articles using different techniques. [A. Bandyopadhyay et al, Snigdha Ghosh et al, N. Sadooghi et al, X. Wang et al].
- Many of the calculations used different approximations, particularly either strong or weak magnetic field approximations.
- For arbitrary magnetic field, either parallel (p_z) or perpendicular (p_\perp) component taken to be zero.

One of the very fast:

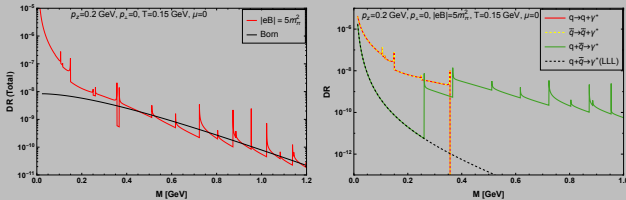
- The rate was estimated in the LLL approx (strong eB) [A. Bandyopadhyay, CAI, M. G. Mustafa, PRD, 2016].



Computational novelty in the present effort:

- We have relaxed all approximations related to the field and external momentum.
- It is easy to grasp because of the simplicity in ITF.
- There are similar works which talk about the ellipticity of the lepton pairs as well. [X. Wang and I. Shovkovy, 2022, 2023](#); [✉ Talk by X. Wang](#)

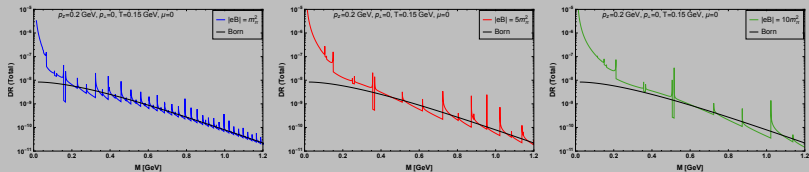
Effect of magnetic field on DR:



- In the left panel we have the plot of DR as a function of invariant mass for $eB = 5 m_\pi^2$ with p_T being zero. In the right panel the contribution coming from different processes are shown separately.

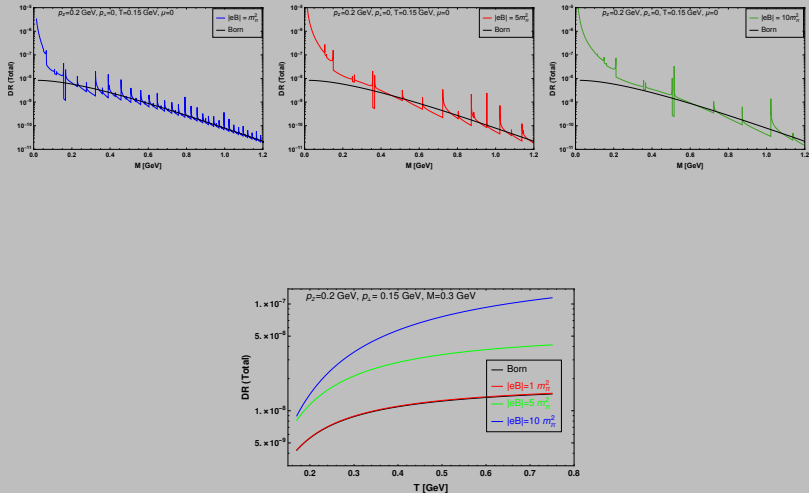
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Effect of magnetic field on DR:



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Effect of magnetic field on DR:



Rate to spectrum:

- The expression of dilepton rate,

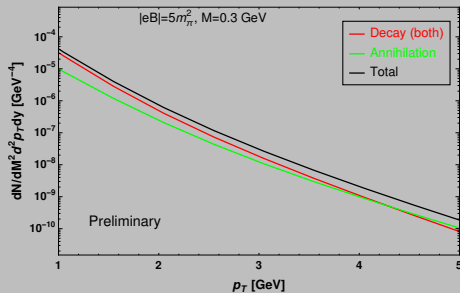
$$\frac{dN}{d^4X d^4P} \equiv \frac{dR}{d^4P} = \frac{\alpha_{\text{EM}}}{12\pi^3} \frac{1}{P^2} \frac{1}{e^{p_0/T} - 1} \sum_{f=u,d} \frac{1}{\pi} \text{Im} \Pi_{\mu,f}^{\mu}(P). \quad (5)$$

- For convenience we make a variable transformation from $(p_0, p_x, p_y, p_z) \rightarrow (M, p_T, \phi, y)$

- The Born rate becomes,

$$\begin{aligned} \frac{dN}{d^4x dM dM p_T dp_T dy d\phi} &= \frac{5N_c \alpha_{\text{EM}}^2}{108\pi^4} \frac{1}{\sqrt{M_T^2 \sinh^2 y + p_T^2}} \frac{T}{e^{\frac{M_T \cosh y}{T}} - 1} \left(1 + \frac{2m_f^2}{M^2} \right) \\ &\times \log \left(\frac{\left[\exp \left(-\frac{M_T \cosh y + \mu}{T} \right) + e^{-\omega - /T} \right] \left[e^{-\mu/T} + e^{-\omega + /T} \right]}{\left[\exp \left(-\frac{M_T \cosh y + \mu}{T} \right) + e^{-\omega + /T} \right] \left[e^{-\mu/T} + e^{-\omega - /T} \right]} \right), \end{aligned} \quad (6)$$

Dilepton spectrum



A. Das, A. Bandyopadhyay, CAI, R. Chatterjee (ongoing)

Summary:

- We have been able to calculate the dilepton rate for arbitrary strength of magnetic field for the most general case.
- We could break down the rate into the contributions coming from different processes and showed that it gets enhanced as compared to the Born rate in presence of eB , particularly at the lower end of the invariant mass.
- We are trying to obtain the spectrum in presence of eB .
- We expect that the enhancement in the rate will be as well reflected in the spectrum and in turn could affect the corresponding flow.

Thank You

A few words on the lifetime of the field

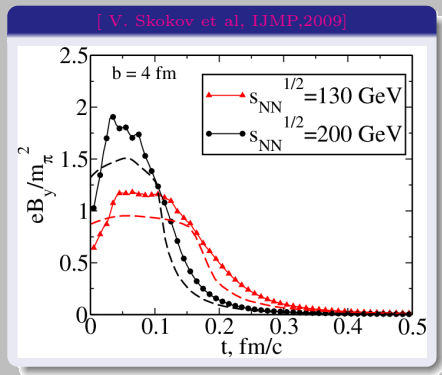
- The very high initial magnitude of this magnetic field then decreases very fast, being inversely proportional to the square of time(?).

[A. Bzdak et al, PRL, 2013;

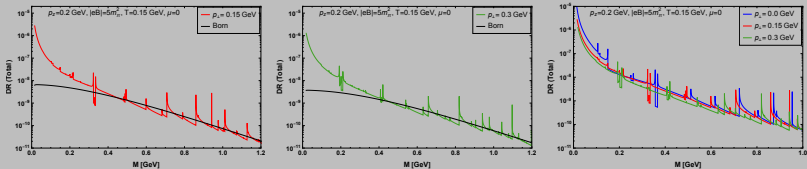
K. Tuchin, PRC, 2013]

[Z. Wang et al, PRC, 2022;

STAR, PRC, 2022]



Effect of the transverse momentum:



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