



Measurement of the Higgs boson in decays to bosons using the ATLAS detector

Dongshuo Du

University of Science and Technology of China

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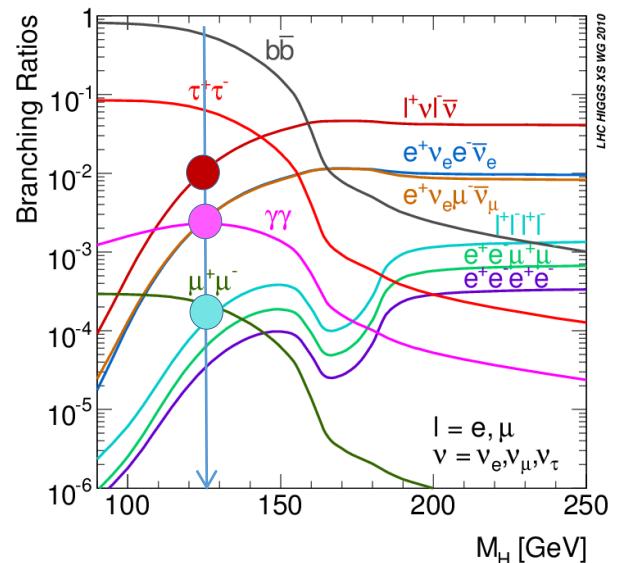
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辽宁师范大学, 2022/08/08 - 2022/08/11

Outline

- Overview of Higgs physics in diboson final states
- Higgs boson cross sections measurement
 - Measurement of Simplified Template Cross Section (STXS) in $H \rightarrow WW^*$, $H \rightarrow ZZ^*$ and $H \rightarrow \gamma\gamma$ decay channels
 - Fiducial inclusive and differential cross-section measurements in $H \rightarrow ZZ^*$ and $H \rightarrow \gamma\gamma$ decay channels
- Summary

Higgs Physics In Diboson Final States

- The most prolific decay: $H \rightarrow b\bar{b}$ (58%), very hard to observe
- The branching ratios (BR) of $H \rightarrow WW^*(\rightarrow l\bar{l}l\bar{l}) / \gamma\gamma / ZZ^*(\rightarrow 4l)$ are 1.0%, 0.23% and 0.012%, respectively
- The final states (e, μ, γ) are very sensitive & leave a clean signature in the ATLAS detector
 - $H \rightarrow ZZ^* \rightarrow 4l$ and $H \rightarrow \gamma\gamma$: the most sensitive channels for observation
 - $H \rightarrow ZZ^* \rightarrow 4l$, S/B > 2
 - $H \rightarrow \gamma\gamma$, S/B ~5%
 - $H \rightarrow WW^* \rightarrow l\bar{l}l\bar{l}$ (high yield with fair S/B)
 - Larger branching ratio • Clean signature
 - High background
- Powerful tool for measuring the Higgs properties



Simplified Template Cross Section (STXS)

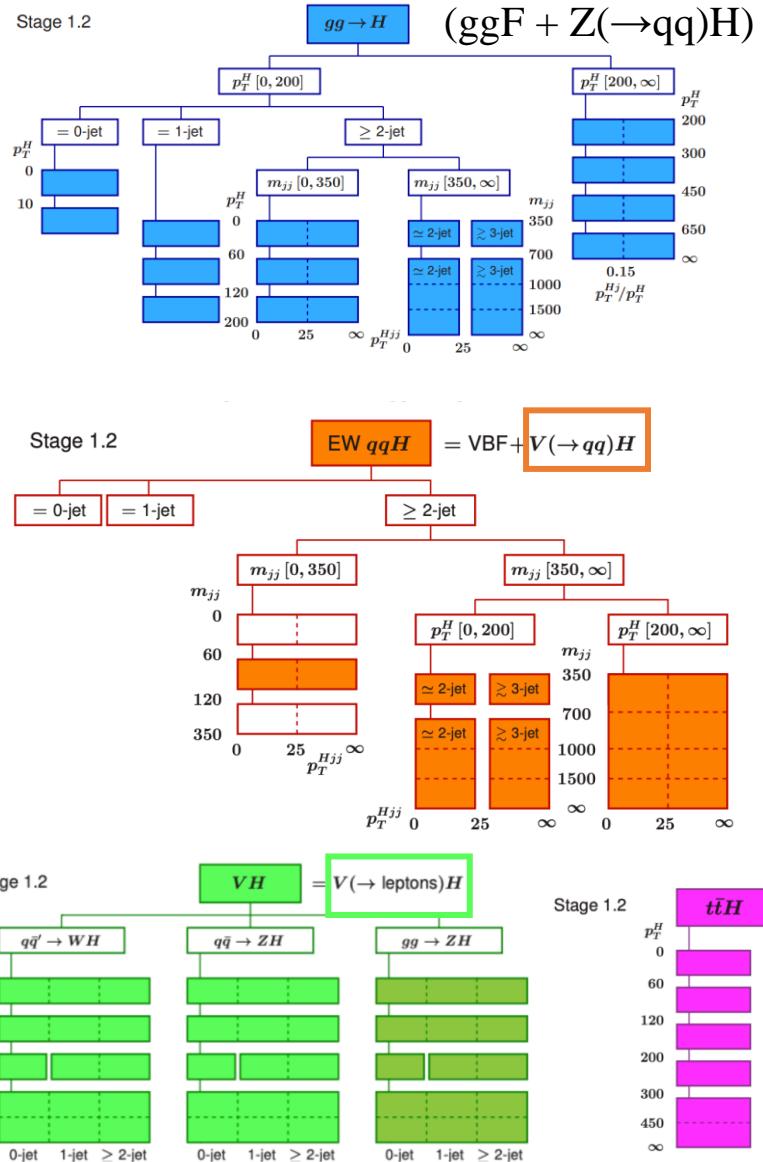
- The aim with STXS method:
 - Improve sensitivity of measurements
 - Reduce their dependence on the theory
 - Isolation of possible BSM effects

- STXS framework provides different stages (e.g. stage 0, stage 1, stage 1.2) with increasingly fine granularity, [more details](#)

- Categorizing events into bins of key (truth) variables (p_T^H , N_{jets} , m_{jj}) in different production modes (ggH , qqH , VH and tth)

- STXS well-suited to combine different decay channels

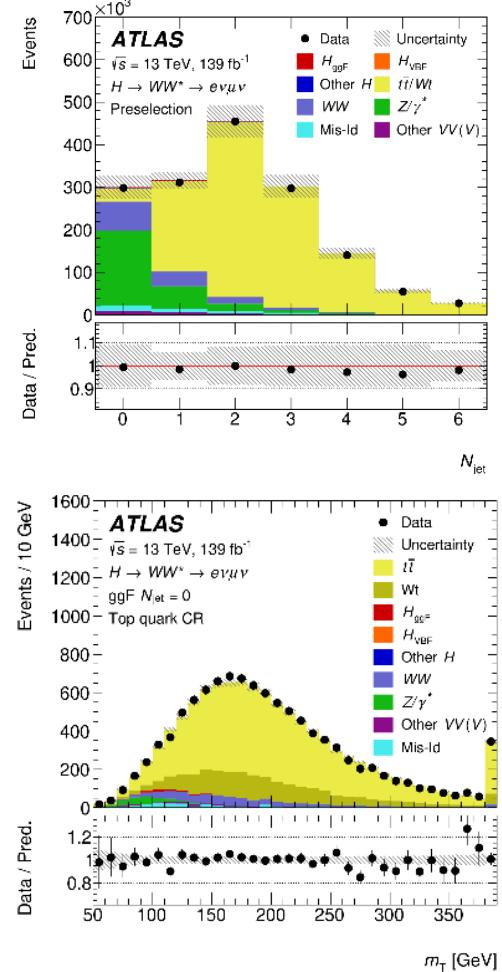
- Higgs boson properties measured with 139 fb^{-1} ($\sqrt{s} = 13 \text{ TeV}$) for Higgs boson rapidity $|y_H| < 2.5$



H \rightarrow WW*: Analysis Strategy

- Signal: different flavour ($e\mu + \mu e$) opposite charge leptons + MET
- Events split in 4 analysis categories based on N_{jets} ^(*)
 - ggF: $N_{\text{jets}} = 0, 1, \geq 2$, cut based
 - m_T used as discriminant variable
 - VBF: $N_{\text{jets}} \geq 2$, “deep” neural network (DNN) based
 - DNN used as discriminant variable
- Main background:
 - Non-resonant qqWW, top and $Z \rightarrow \tau\tau$
 - ggF: qqWW, top and $Z \rightarrow \tau\tau$ normalized by control regions (CR)
 - VBF: top and $Z \rightarrow \tau\tau$ normalized by CRs
 - Background with mis-identified leptons estimated by data-driven fake factor method

Submitted to PRD



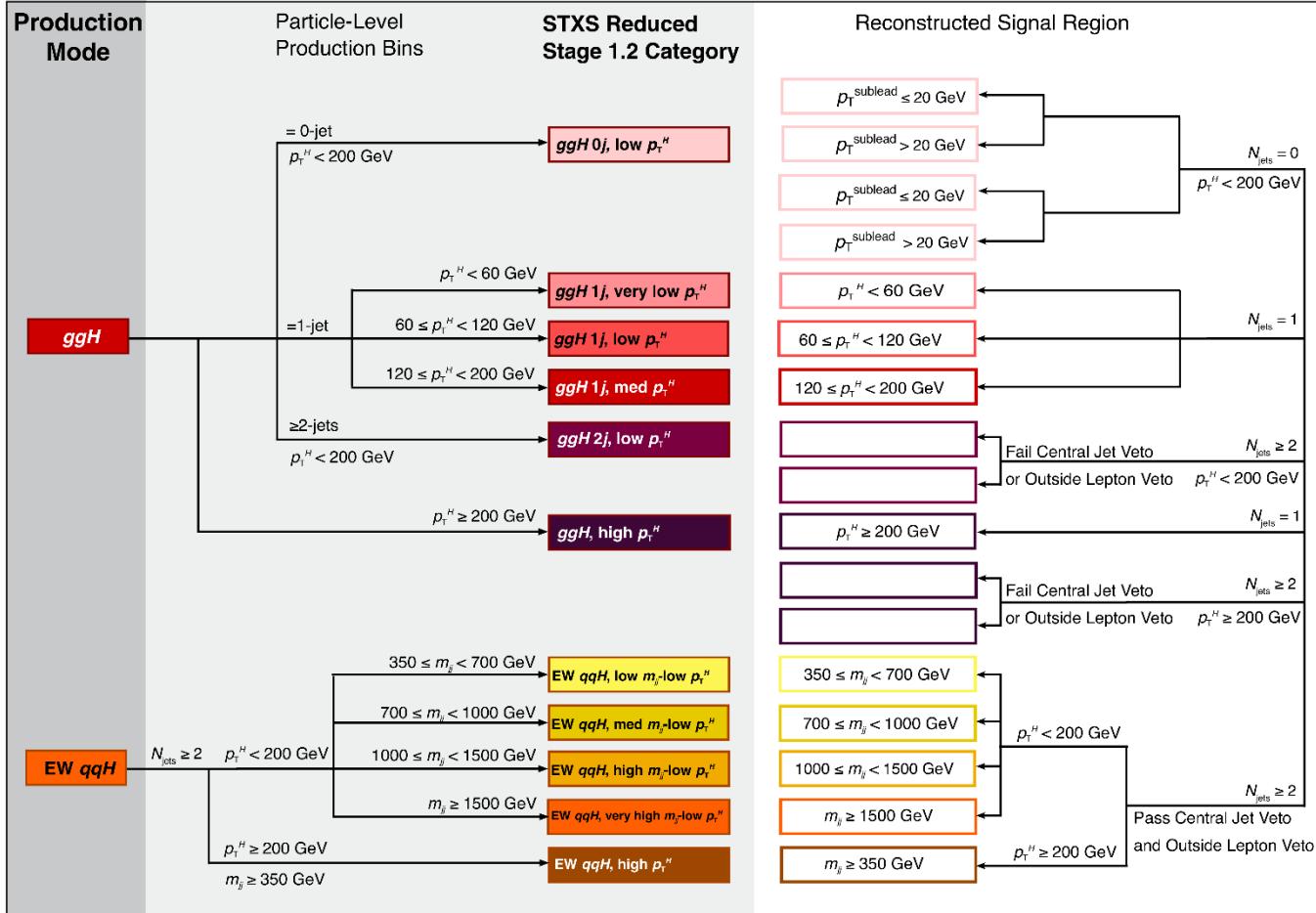
(*) Full event selection in the backup

$$m_T = \sqrt{(E_T^{\ell\ell} + E_T^{\text{miss}})^2 - |\vec{p}_T^{\ell\ell} + \vec{E}_T^{\text{miss}}|^2}, \quad E_T^{\ell\ell} = \sqrt{|\vec{p}_T^{\ell\ell}|^2 + m_{\ell\ell}^2}$$

H \rightarrow WW*: STXS

$H \rightarrow WW^* \rightarrow e\nu\mu\nu$ ATLAS $\sqrt{s} = 13$ TeV, 139 fb^{-1}

Submitted to PRD



- Split by p_T^H , N_{jets} , m_{jj} into 11 categories
- 17 signal regions (SR)
- 27 control regions (CR)

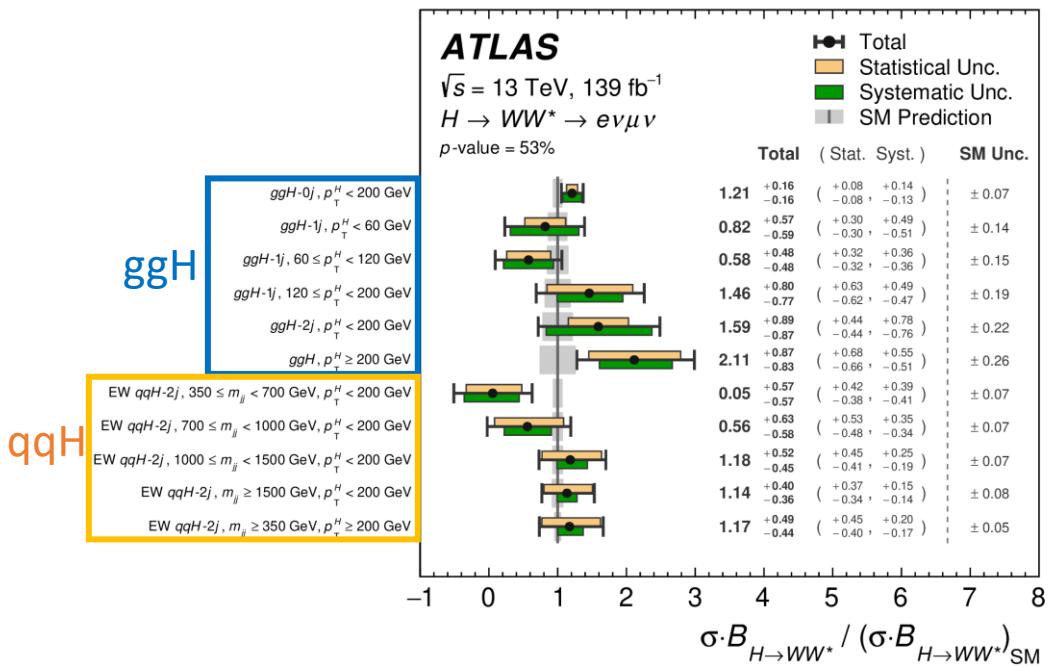
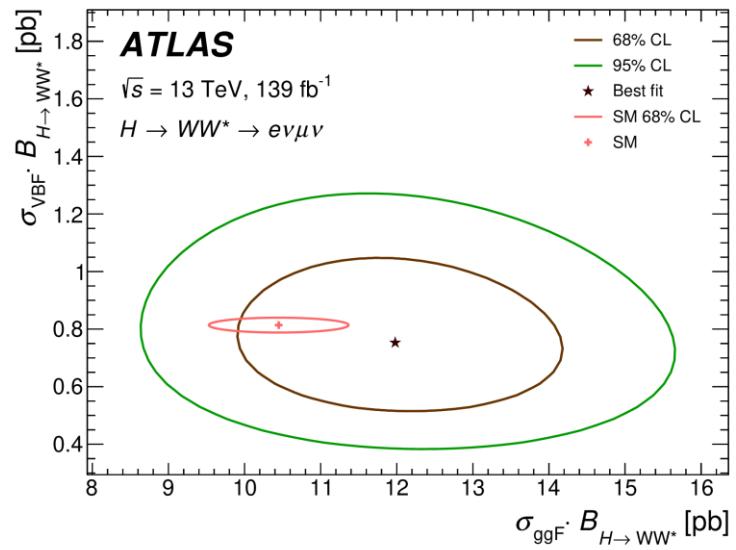
- This analysis based on the reduced stage 1.2 category to ensure sensitivity for all measurements.
- CRs split similar to SRs where statistics allow

$$p_T^H = |\mathbf{p}_T^{\ell\ell} + \mathbf{E}_T^{\text{miss}}|$$

H \rightarrow WW*: Results

Submitted to PRD

Coupling results



$$\begin{aligned} \sigma_{\text{ggF}} \cdot B_{H \rightarrow WW^*} &= 12.0 \pm 1.4 \text{ pb} \\ &= 12.0 \pm 0.6 \text{ (stat.)}^{+0.9}_{-0.8} \text{ (exp. syst.)}^{+0.6}_{-0.5} \text{ (sig. theo.)} \pm 0.8 \text{ (bkg. theo.) pb} \\ \sigma_{\text{VBF}} \cdot B_{H \rightarrow WW^*} &= 0.75^{+0.19}_{-0.16} \text{ pb} \\ &= 0.75 \pm 0.11 \text{ (stat.)}^{+0.07}_{-0.06} \text{ (exp. syst.)}^{+0.12}_{-0.08} \text{ (sig. theo.)}^{+0.07}_{-0.06} \text{ (bkg. theo.) pb,} \end{aligned}$$

➤ Compatible with SM predictions within 1 σ

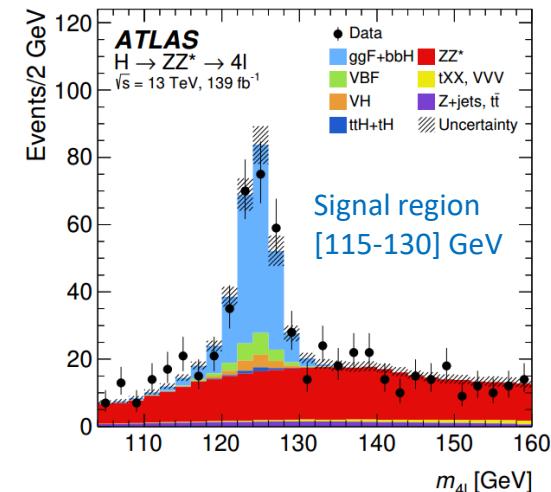
- Extracted by profile likelihood fit: 17 SRs (m_{jj}/DNN) + 27 CRs
- ggH uncertainties limited by both stat. + syst. uncertainty
- qqH uncertainties limited by statistical uncertainty at high m_{jj} / p_T^H

$H \rightarrow ZZ^* \rightarrow 4\ell$: Four-lepton Invariant Mass Distribution

Analysis features:

- Small branching fraction (0.0124% at $m_H = 125$ GeV)
- Best final S/B ratio, better than 2:1
- Clean signature with fully reconstructed final state
- good mass resolution = 1-2%

Signal: 4 leptons (4e, 4 μ , 2e2 μ and 2 μ 2e)

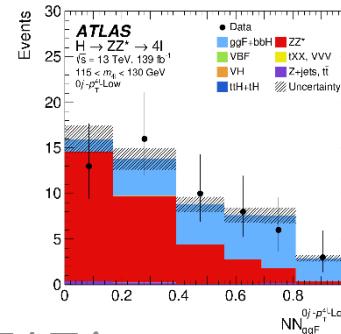


[Eur. Phys. J. C 80, 957 \(2020\).](#)

Backgrounds:

- Dominant background: Non-resonant $ZZ^*/Z\gamma^*$, estimated from data sideband (105–160 GeV)
- Reducible background with non-prompt leptons for $Z+\text{jets}$, $t\bar{t}\text{bar}$ constrained by dedicated CR.

Cross sections extracted by fit to the NN observable

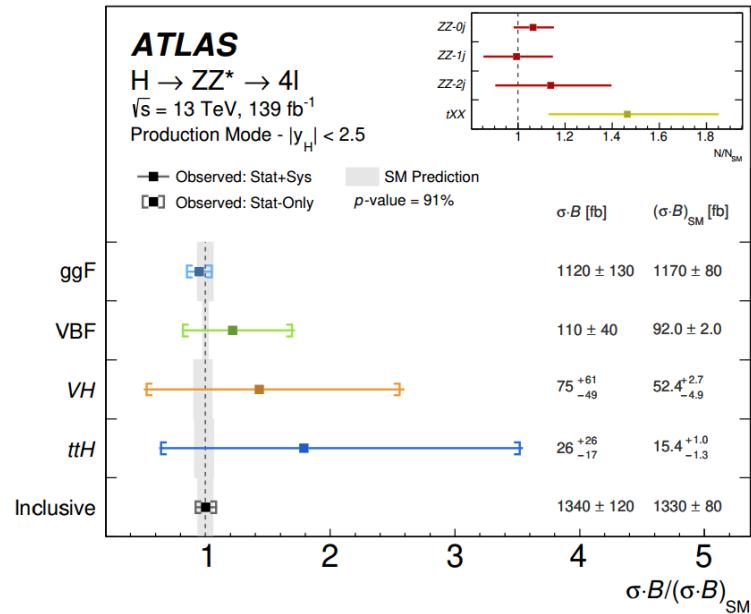


ggF 0jet category
with $p_T^H < 10$ GeV

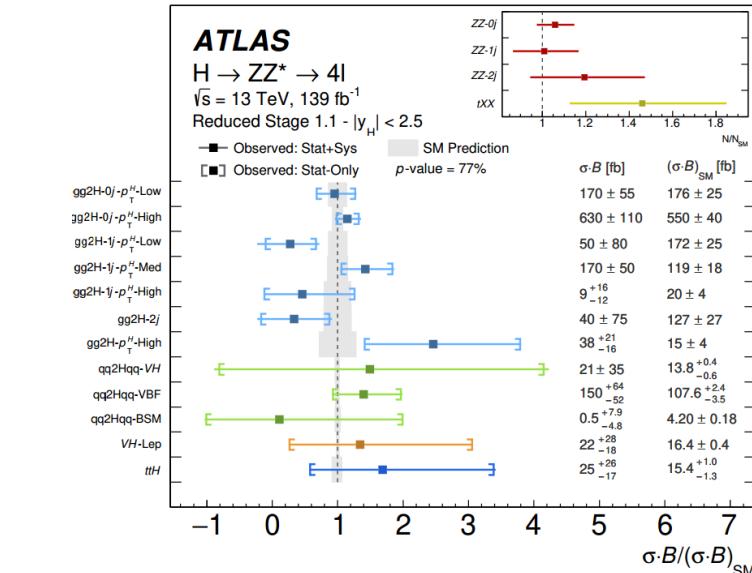
H \rightarrow ZZ* \rightarrow 4 ℓ : STXS

[Eur. Phys. J. C 80, 957 \(2020\).](#)

Production Mode Stage



Reduced Stage 1.1



- The inclusive H \rightarrow ZZ* production cross-section for $|y_H| < 2.5$:
 $1.34 \pm 0.12 \text{ pb}$
- Good agreement with the SM predictions:
 $1.33 \pm 0.08 \text{ pb}$
- 12 pois measured simultaneous
 - ggF/VBF: 7/3 categories separately
 - VH-lep/ttH: One bin from each mode
- Due to finer categorization, results are statistically limited.

$H \rightarrow \gamma\gamma$: Diphoton Invariant Mass Distribution

Analysis features:

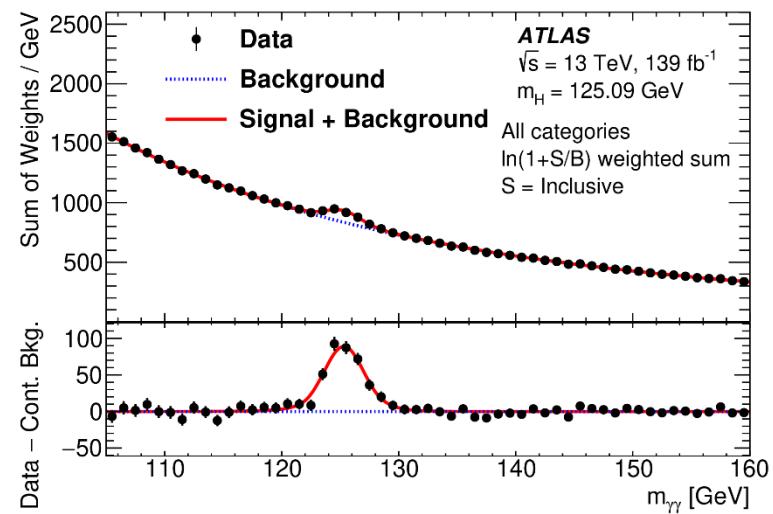
- Fairly high branching fraction wrt $H \rightarrow ZZ^* \rightarrow 4\ell$:
~20 times larger
- **Excellent performance of photon reconstruction and identification**
- **Final states are fully reconstructable**
- **Good mass resolution = 1-2%**

Signal: diphoton

Backgrounds:

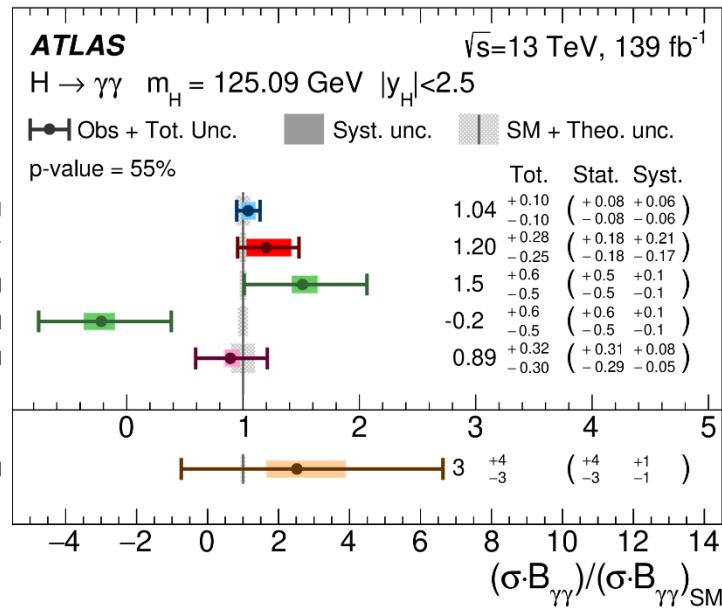
- SM diphoton production (>75%) or $\gamma + \text{jets}$ (~20%), $\text{jet} + \text{jets}$ (<5%) estimated from data sideband

[Submitted to JHEP](#)



Signal + background model fit for all categories with $m_{\gamma\gamma}$, and the residual plot after subtracting the backgrounds

H $\rightarrow\gamma\gamma$: Production Mode Cross Sections



[Submitted to JHEP](#)

- The contribution from the bbH process is included in the gg
- Good agreement with SM prediction.
- $\sigma_{tH}^{Obs (exp)} < 10 (6.8) \sigma_{tH}^{SM}$ at 95% CL
 - Unique sensitivity to the sign of the top-quark Yukawa coupling

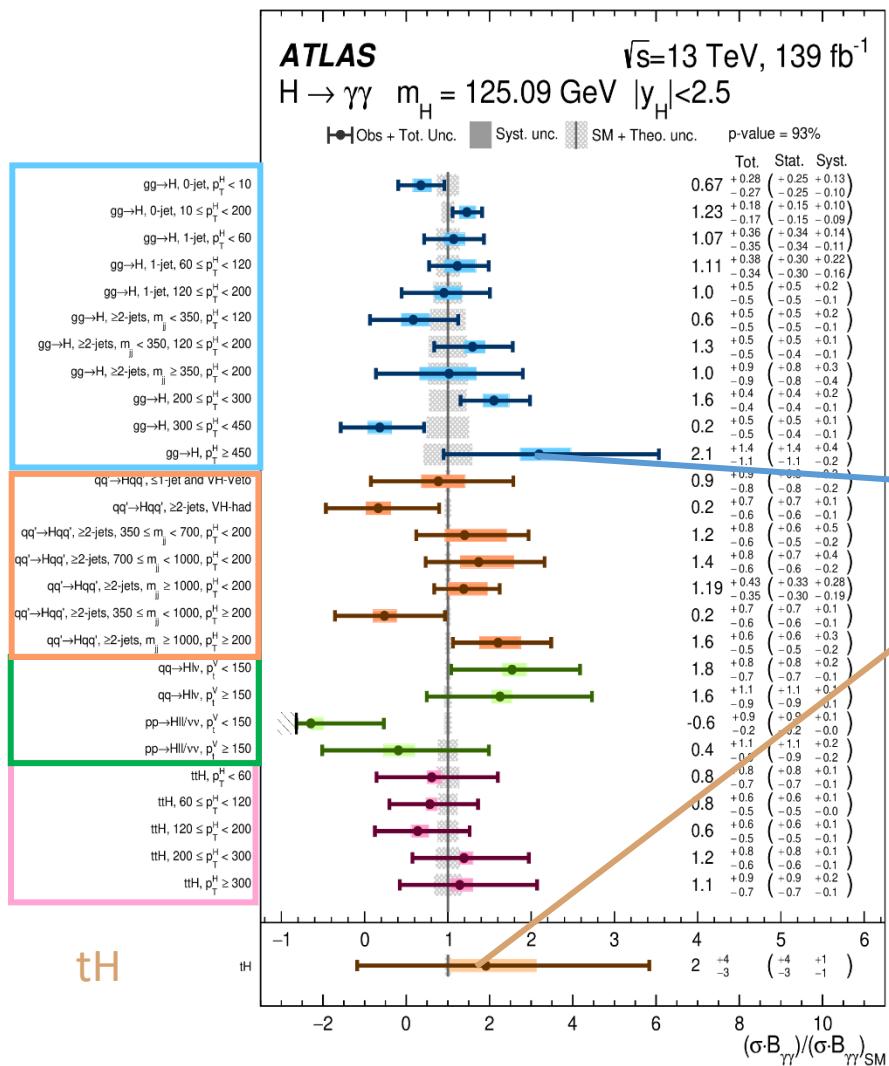
- The overall Higgs boson signal strength:

$$\mu = 1.04^{+0.10}_{-0.09} = 1.04 \pm 0.06 \text{ (stat.)}^{+0.06}_{-0.05} \text{ (theory syst.)}^{+0.05}_{-0.04} \text{ (exp. syst.)}.$$

Good agreement with SM prediction

H \rightarrow $\gamma\gamma$: STXS

Submitted to JHEP



- More statistics with Full Run2 dataset, finer binning especially for non-ggH production mode.
- Cross-sections of 28 STXS regions are measured
 - First to measure $t t H$ differentially
 - Access to the very high p_T^H region
 - Access to $t H$ production mode ($t H q b$ and $t H W$)
- Very good agreement with SM prediction with a p-value of 93%.
- Due to finer categorization, measurements are statistically limited.

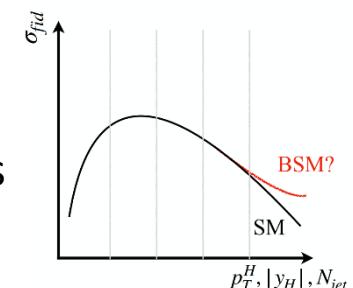
Differential and Fiducial Cross Sections

- Fiducial region: The fiducial selection closely match the detector level
 - Minimise model-dependent acceptance extrapolations
 - Allow for easy comparison of physics models today and in the future

$$\sigma^{\text{fid}} = \sigma^{\text{total}} \times A \times BR = \frac{N_s}{C \times L_{\text{int}}}$$

N_s : number of observed signal events, $C = \frac{N_{\text{Rec}}}{N_{\text{Fid}}}$: the correction factor for detector efficiency and resolution effect, L_{int} : Integrated luminosity

- **Inclusive fiducial cross-section:** No attempt to separate Higgs production/decay modes
→ compare with best available predictions in the detector phase space
- **Differential cross-section:** Different differential cross-section measurement sensitive to the different physics:
 - The Higgs p_T distribution can test couplings to b and c quarks, and high p_T can tests new physics.
 - The jet multiplicity is sensitive to different production mechanisms and the theoretical modelling of high-quark and gluon emission.

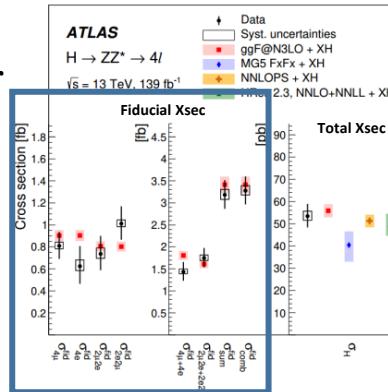


- Combination needs extrapolation to the full phase space due to the fiducial regions are different (acceptance + BR)

H \rightarrow ZZ* \rightarrow 4 ℓ : Inclusive and Differential Cross Section

[Eur. Phys. J. C 80, 942 \(2020\)](#)

- H \rightarrow ZZ* \rightarrow 4 ℓ decay mode provides excellent resolution for Higgs kinematic variables
- Fitting 4 ℓ distribution in each final state (**right**) or **differential bin** (**bottom**) to extract the measured cross section



Inclusive fiducial cross section:

$$\sigma_{\text{fid}} = 3.28 \pm 0.30 \text{ (stat.)} \pm 0.11 \text{ (syst.) fb}$$

$$\sigma_{\text{fid},SM} = 3.41 \pm 0.18 \text{ fb}$$

10% level precision and Good agreement with LHCXSWG prediction

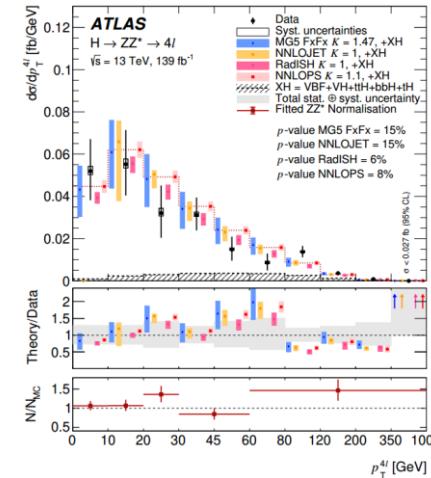
Differential cross section with 20 observables:

- 9 Higgs kinematic-related variables: like p_T^H
- 7 Jet-related variables: like N_{jets}
- 4 Higgs boson and jet-related variables

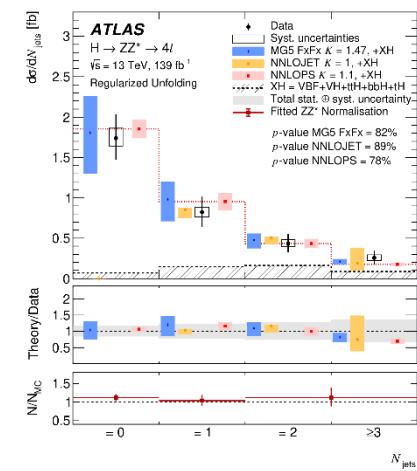
Measured cross sections in good agreement with SM predictions within uncertainties

Limited by statistical uncertainties

Low p_T^H test couplings to b and c quarks and High p_T^H test BSM



Sensitive to production mode



Predictions for comparison

H $\rightarrow\gamma\gamma$: Inclusive and Differential Cross Section

- Excellent performance of photon reconstruction and identification
→ **good resolution for Higgs variables**
- The H $\rightarrow \gamma\gamma$ signal yields are extracted from fits to the diphoton invariant mass spectrum

Inclusive fiducial cross section:

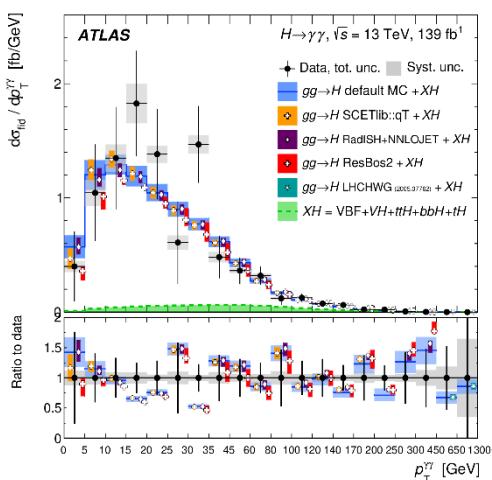
$$\sigma_{\text{fid}} = 67 \pm 6 \text{ fb}$$

$$\sigma_{\text{fid},SM} = 64.2 \pm 3.4 \text{ fb}$$

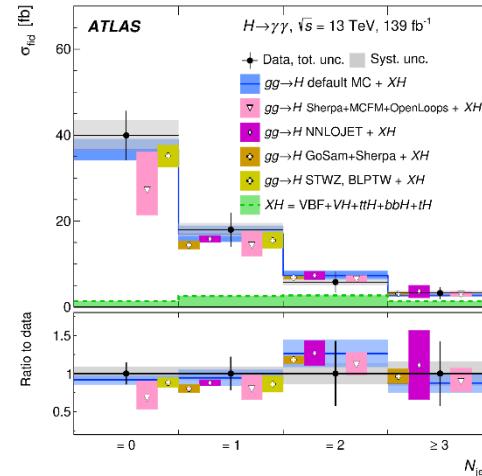
The measurement has precision at 9% level

Differential cross-section still limited by statistical uncertainties

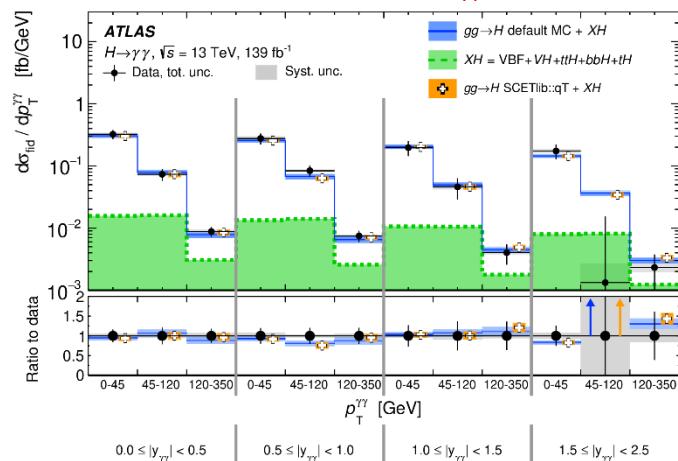
Low p_T^H test couplings to b and c quarks and High p_T^H test BSM



N_{jets} → test modelling of radiations at high pT and production modes



Double-differential cross section:
p_T^{γγ} in bins of |y_{γγ}|



Measured cross sections in good agreement with SM predictions within uncertainties

Accepted by JHEP

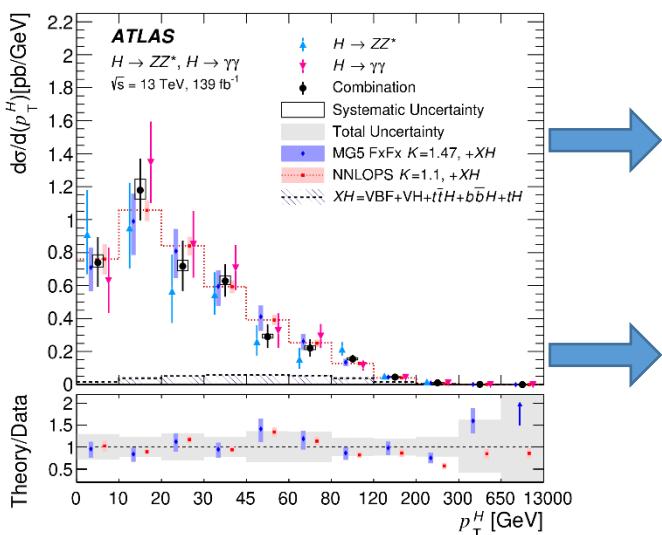
Dongshuo Du (USTC)

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More details in Fabio's talk

Combination: $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ$

- Combination between $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ$ can improve the measurement precision.
 - Based on the individual channel
 - The uncertainties from same source are correlated
 - Both channels are extrapolated to the full phase space to do the combination
- The Combined differential result still dominated by the statistical uncertainty



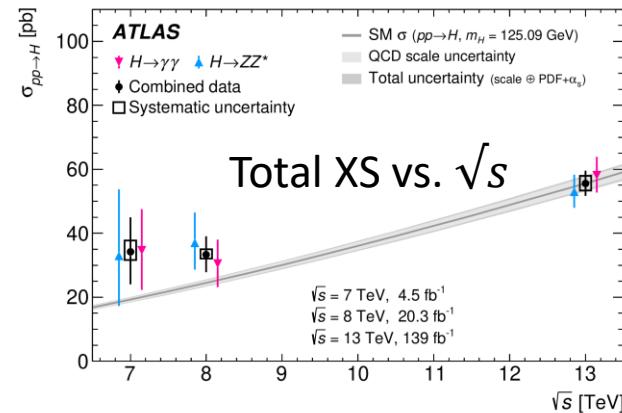
p_T^H precision:
 20-30% up to 300 GeV
 ~60% 300-650 GeV
 >100% 650-1200 GeV

The differential XS distribution of p_T^H can be used to indirectly constrain the Yukawa coupling of charm and bottom quark (k_c and k_b).

[More details in Tao's talk](#)

Submitted to JHEP

Measured total XS: $55.5^{+4.0}_{-3.8} \text{ pb}$
 SM prediction: $55.6 \pm 2.8 \text{ pb}$



With more data, the precision can be further improved

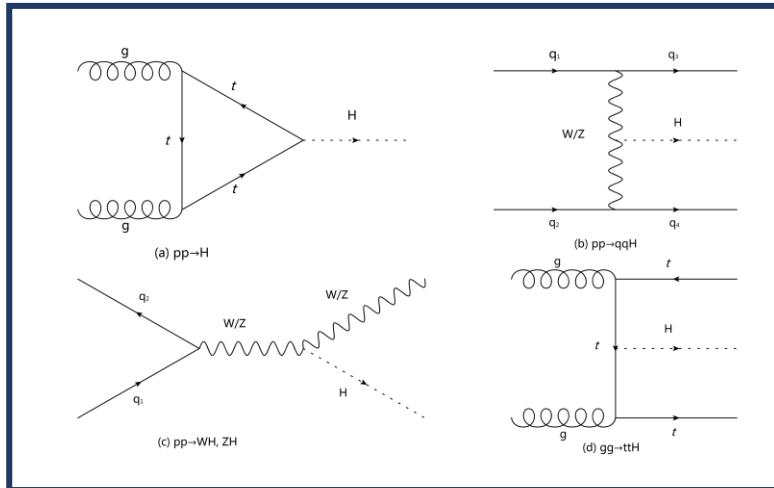
Summary

- $H \rightarrow ZZ^*$, $H \rightarrow \gamma\gamma$, $H \rightarrow WW^*$ channels investigated with 139 fb^{-1} of data collected with the ATLAS detector @13TeV
 - Inclusive, STXS and differential cross section measurements are presented
 - All the measurements are in agreement with the SM predictions
 - The precision of the fiducial cross section measurement in both $H \rightarrow ZZ^*$ and $H \rightarrow \gamma\gamma$ channel about 10%.
 - All the differential cross section measurement still dominant by the statistical uncertainty
- **Run3 has just started with much more to come. Stay tuned!**

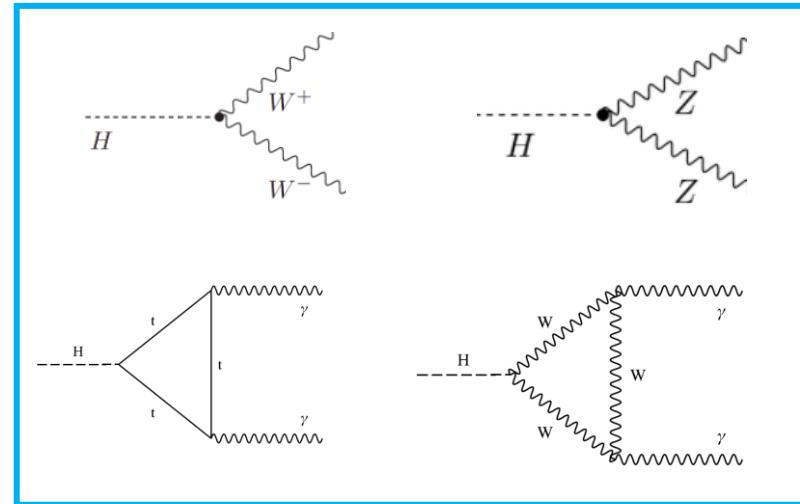
Thank you for your attention

Back up

Higgs Physics In Diboson Final States

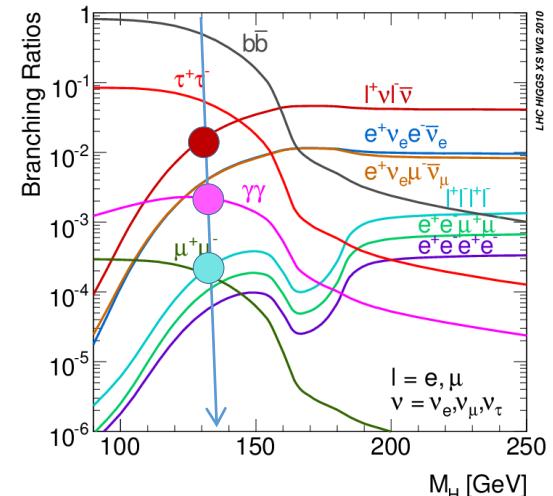


Production modes



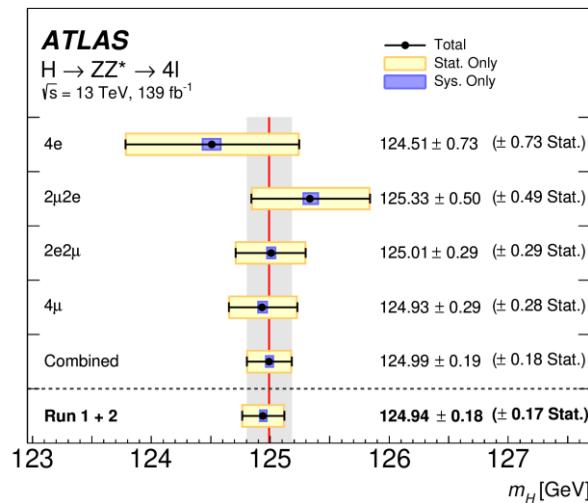
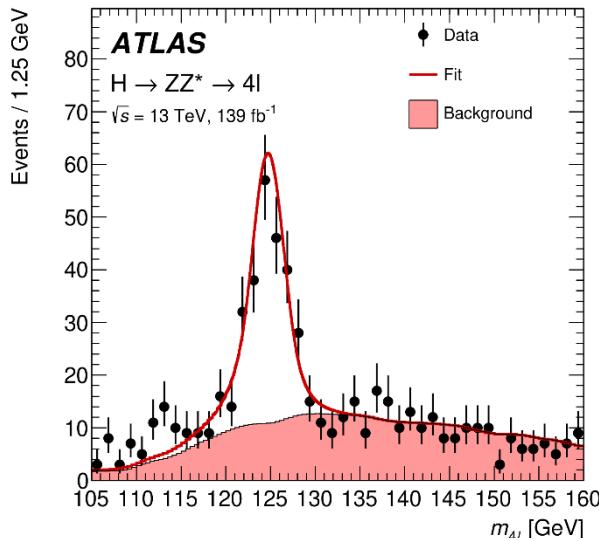
Higgs decay to diboson

- The most prolific decay: $H \rightarrow bb$ (58%), very hard to observe
- The branching ratios (BR) of $H \rightarrow WW^*$ ($\rightarrow ll\nu\bar{\nu}$) / ZZ^* ($\rightarrow 4l$) / $\gamma\gamma$ are 1.0%, 0.012% and 0.23%, respectively
- The final states (e, μ, γ) are very sensitive & leave a clean signature in the ATLAS detector



Higgs mass measurement

[Submitted to Physics Letters B](#)



- First mass measurement of Higgs boson with HZZ channel by using Full Run2 data
 $m_H = 124.99 \pm 0.18(\text{stat.}) \pm 0.04(\text{syst.}) \text{ GeV.}$
- m_H is extracted by performing a simultaneous fit to the four subchannels (4μ , $2e2\mu$, $2\mu2e$, and $4e$) in the $m_{4\ell}$ range between 105 and 160 GeV.
- The results are still limited by statistical uncertainty
- Run1+Run2 combined results with [new correlation scheme](#)(Systematic uncertainties largely reduced)

$$m_H = 124.94 \pm 0.17(\text{stat.}) \pm 0.03(\text{syst.}) \text{ GeV,}$$

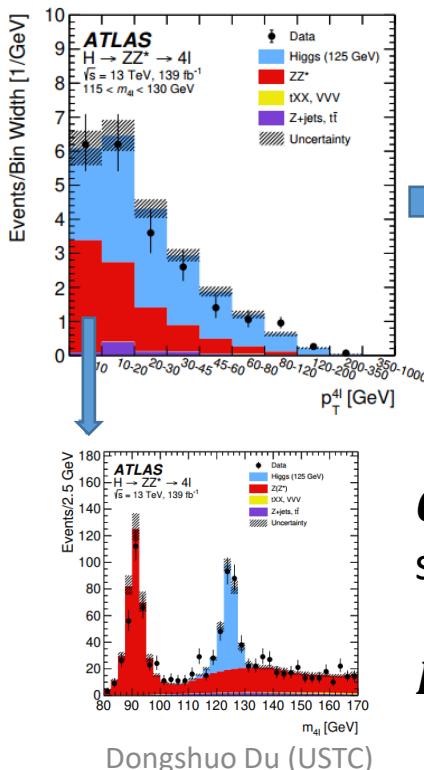
- The most precise measurement of m_H in the $H \rightarrow ZZ^* \rightarrow 4\ell$ channel by the ATLAS Collaboration

More details in [Han Li's talk](#)

More about the Fiducial and differential cross sections measurement

Fiducial region is defined to closely match to experimental acceptance

To obtain the data distributions in fiducial region, the **response matrix $R_{\text{truth_reco}}$** is used to accounting for the bin migrations

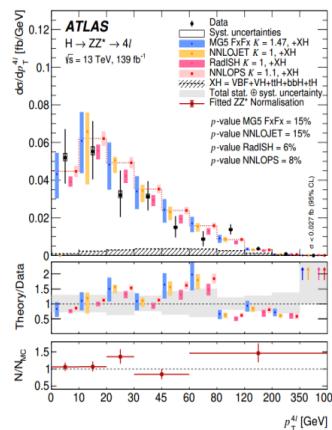


Unfolding

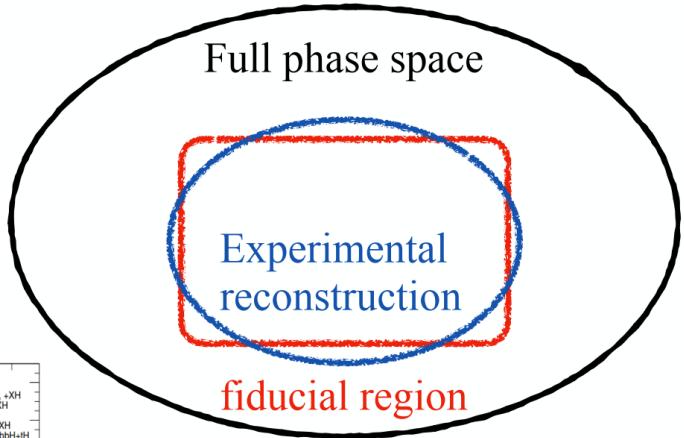
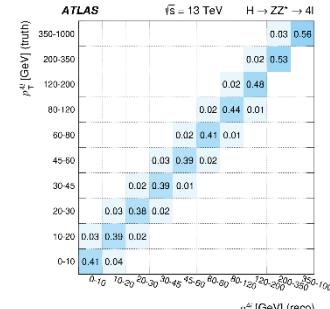
$$N_r^{(H)} = \frac{1}{C_r^{\text{fid}}} \left[\sum_t L \times (\sigma_t \times B_{\gamma\gamma}) \times R_{t,r} \right],$$

C_r^{fid} : Corrects for events that pass selection but outside of the fiducial region.

$R_{t,r}$ is the response matrix.



Response Matrix: few percent migration



Main improvement for HWW analysis

Big picture:

1. Total integrated luminosity
2. Additional consider ggF 2jet analysis
3. STXS also measured in this analysis

More specific:

- ① Ttbar nnlo reweighting to correct the mismodeling of the leading-lepton P_T due to missing higher-order corrections
- ② B-tagging change from MV2C10 to DL1r
- ③ DNN to replace the BDT for VBF analysis
- ④ The central jet veto is applied for jets with $pT > 30$ GeV, compared to $pT > 20$ GeV used previously, reducing the theory modelling uncertainty

...

Main improvement for H $\gamma\gamma$

- New categorization; $\sim 20\%$ improvement in sensitivity.
- Increased granularity; 28 STXS categories from 101 reconstruction categories

Categorization

The previous analyses performed such categorization sequentially, giving higher priority to production processes with lower cross sections, the new strategy considers all production processes simultaneously.

1. Signal multi-class BDT: splits signal into STXS regions, aiming to simultaneously reduce the uncertainties and correlations between the classes.
2. Background rejection binary BDT: in each of the signal multi-class BDT output classes a dedicated BDT is trained. This BDT aims to separate signal from the backgrounds.

Multi-class BDT: Training variables

[Submitted to JHEP](#)

$\eta_{\gamma_1}, \eta_{\gamma_2}, p_T^{\gamma\gamma}, y_{\gamma\gamma},$
 $p_{T,jj}^\dagger, m_{jj}$, and $\Delta y, \Delta\phi, \Delta\eta$ between j_1 and j_2 ,

$p_{T,\gamma\gamma j_1}, m_{\gamma\gamma j_1}, p_{T,\gamma\gamma jj}^\dagger, m_{\gamma\gamma jj}$
 $\Delta y, \Delta\phi$ between the $\gamma\gamma$ and jj systems,
minimum ΔR between jets and photons,

invariant mass of the system comprising all jets in the event,

dilepton p_T , di- e or di- μ invariant mass (leptons are required to be oppositely charged),

E_T^{miss} , p_T and transverse mass of the lepton + E_T^{miss} system,

p_T, η, ϕ of top-quark candidates, $m_{t_1 t_2}$

Number of jets \dagger , of central jets ($|\eta| < 2.5$) \dagger , of b -jets \dagger and of leptons,

p_T of the highest- p_T jet, scalar sum of the p_T of all jets,

scalar sum of the transverse energies of all particles ($\sum E_T$), E_T^{miss} significance,

$|E_T^{\text{miss}} - E_T^{\text{miss}}(\text{primary vertex with the highest } \sum p_{T,\text{track}}^2)| > 30 \text{ GeV}$

Top reconstruction BDT of the top-quark candidates,

$\Delta R(W, b)$ of t_2 ,

$\eta_{j_F}, m_{\gamma\gamma j_F}$

Average number of interactions per bunch crossing.

Binary BDT: Training variables

| STXS classes | Variables |
|---|--|
| Individual STXS classes from $gg \rightarrow H$ $qq' \rightarrow Hqq'$ $qq \rightarrow H\ell\nu$ $pp \rightarrow H\ell\ell$ $pp \rightarrow H\nu\bar{\nu}$ | All multiclass BDT variables, $p_T^{\gamma\gamma}$ projected to the thrust axis of the $\gamma\gamma$ system ($p_{Tt}^{\gamma\gamma}$), $\Delta\eta_{\gamma\gamma}, \eta^{\text{Zepp}} = \frac{\eta_{\gamma\gamma} - \eta_{jj}}{2},$ $\phi_{\gamma\gamma}^* = \tan\left(\frac{\pi - \Delta\phi_{\gamma\gamma} }{2}\right) \sqrt{1 - \tanh^2\left(\frac{\Delta\eta_{\gamma\gamma}}{2}\right)},$ $\cos\theta_{\gamma\gamma}^* = \left \frac{(E^{\gamma_1} + p_z^{\gamma_1}) \cdot (E^{\gamma_2} - p_z^{\gamma_2}) - (E^{\gamma_1} - p_z^{\gamma_1}) \cdot (E^{\gamma_2} + p_z^{\gamma_2})}{m_{\gamma\gamma} + \sqrt{m_{\gamma\gamma}^2 + (p_T^{\gamma\gamma})^2}} \right $ Number of electrons and muons. |
| all $t\bar{t}H$ and tHW STXS classes combined | p_T, η, ϕ of γ_1 and γ_2 , p_T, η, ϕ and b -tagging scores of the six highest- p_T jets, $E_T^{\text{miss}}, E_T^{\text{miss}}$ significance, E_T^{miss} azimuthal angle, Top reconstruction BDT scores of the top-quark candidates, p_T, η, ϕ of the two highest- p_T leptons. |
| $tHqb$ | $p_T^{\gamma\gamma}/m_{\gamma\gamma}, \eta_{\gamma\gamma},$ p_T , invariant mass, BDT score and $\Delta R(W, b)$ of t_1 , p_T, η of t_2 , p_T, η of j_F , Angular variables: $\Delta\eta_{\gamma\gamma t_1}, \Delta\theta_{\gamma\gamma t_2}, \Delta\theta_{t_1 j_F}, \Delta\theta_{t_2 j_F}, \Delta\theta_{\gamma\gamma j_F}$ Invariant mass variables: $m_{\gamma\gamma j_F}, m_{t_1 j_F}, m_{t_2 j_F}, m_{\gamma\gamma t_1}$ Number of jets with $p_T > 25$ GeV, Number of b -jets with $p_T > 25$ GeV*; Number of leptons*, E_T^{miss} significance* |

[Submitted to JHEP](#)

Main improvement for HZZ

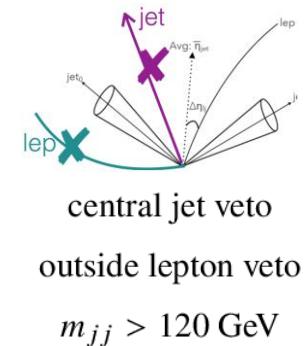
- Improved **lepton isolation** to mitigate the impact of pile-up
- Improved **jet reconstruction** using a particle flow algorithm
- Additional event categories for the classification of Higgs boson candidates
- New discriminants to enhance the sensitivity to distinguish the various production modes of the SM Higgs boson
- Use of data sidebands to constrain the dominant ZZ^* background process

H \rightarrow WW*: analysis

| Category | $N_{\text{jet},(p_T>30 \text{ GeV})} = 0 \text{ ggF}$ | $N_{\text{jet},(p_T>30 \text{ GeV})} = 1 \text{ ggF}$ | $N_{\text{jet},(p_T>30 \text{ GeV})} \geq 2 \text{ ggF}$ | $N_{\text{jet},(p_T>30 \text{ GeV})} \geq 2 \text{ VBF}$ |
|--|---|---|--|--|
| Preselection | Two isolated, different-flavor leptons ($\ell = e, \mu$) with opposite charge $p_T^{\text{lead}} > 22 \text{ GeV}, p_T^{\text{sublead}} > 15 \text{ GeV}$ $m_{\ell\ell} > 10 \text{ GeV}$ | | | |
| | $p_T^{\text{miss}} > 20 \text{ GeV}$ | | | |
| Background rejection | $N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$ $\Delta\phi_{\ell\ell, E_T^{\text{miss}}} > \pi/2$ $p_T^{\ell\ell} > 30 \text{ GeV}$ | | | |
| | $m_{\tau\tau} < m_Z - 25 \text{ GeV}$ $\max(m_T^\ell) > 50 \text{ GeV}$ | | | |
| $H \rightarrow WW^* \rightarrow e\nu\mu\nu$ topology | $m_{\ell\ell} < 55 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 1.8$ | | | |
| | fail central jet veto or fail outside lepton veto | | | |
| Discriminating fit variable | $ m_{jj} - 85 > 15 \text{ GeV}$ or $\Delta y_{jj} > 1.2$ | | | |
| | m_T | | | |
| DNN | | | | |

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[Reconstructed with the
collinear approximation
method](#)



H \rightarrow WW*: analysis

| CR | $N_{\text{jet},(p_T>30 \text{ GeV})} = 0 \text{ ggF}$ | $N_{\text{jet},(p_T>30 \text{ GeV})} = 1 \text{ ggF}$ | $N_{\text{jet},(p_T>30 \text{ GeV})} \geq 2 \text{ ggF}$ | $N_{\text{jet},(p_T>30 \text{ GeV})} \geq 2 \text{ VBF}$ |
|---|--|---|--|--|
| $N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$ | | | | |
| $qq \rightarrow WW$ | $\Delta\phi_{\ell\ell, E_T^{\text{miss}}} > \pi/2$ $p_T^{\ell\ell} > 30 \text{ GeV}$ $55 < m_{\ell\ell} < 110 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 2.6$ | $m_{\ell\ell} > 80 \text{ GeV}$ $ m_{\tau\tau} - m_Z > 25 \text{ GeV}$ $\max(m_T^\ell) > 50 \text{ GeV}$ | $m_{\tau\tau} < m_Z - 25 \text{ GeV}$ | $m_{T2} > 165 \text{ GeV}$ |
| | | | fail central jet veto or fail outside lepton veto | |
| | | | $ m_{jj} - 85 > 15 \text{ GeV}$ or $\Delta y_{jj} > 1.2$ | |
| | | | | |
| | | | | |
| | | | | |
| $t\bar{t}/Wt$ | $N_{b\text{-jet},(20 < p_T < 30 \text{ GeV})} > 0$ $\Delta\phi_{\ell\ell, E_T^{\text{miss}}} > \pi/2$ $p_T^{\ell\ell} > 30 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 2.8$ | $N_{b\text{-jet},(p_T>30 \text{ GeV})} = 1$ $N_{b\text{-jet},(20 < p_T < 30 \text{ GeV})} = 0$ | $N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$ | $N_{b\text{-jet},(p_T>20 \text{ GeV})} = 1$ |
| | | | $m_{\tau\tau} < m_Z - 25 \text{ GeV}$ | |
| | | | $\max(m_T^\ell) > 50 \text{ GeV}$ | $m_{\ell\ell} > 80 \text{ GeV}$ $\Delta\phi_{\ell\ell} < 1.8$ $m_{T2} < 165 \text{ GeV}$ |
| | | | fail central jet veto or fail outside lepton veto | central jet veto outside lepton veto |
| | | | $ m_{jj} - 85 > 15 \text{ GeV}$ or $\Delta y_{jj} > 1.2$ | |
| | | | | |
| Z/γ^* | $N_{b\text{-jet},(p_T>20 \text{ GeV})} = 0$ $m_{\ell\ell} < 80 \text{ GeV}$ no p_T^{miss} requirement | $m_{\ell\ell} < 55 \text{ GeV}$ | $m_{\ell\ell} < 70 \text{ GeV}$ | |
| | | | | |
| | | | | |
| | | | | |
| | | | | |
| | | | | |

Submitted to PRD

VBF ≥ 2 Jet: Deep Neural Network (DNN)

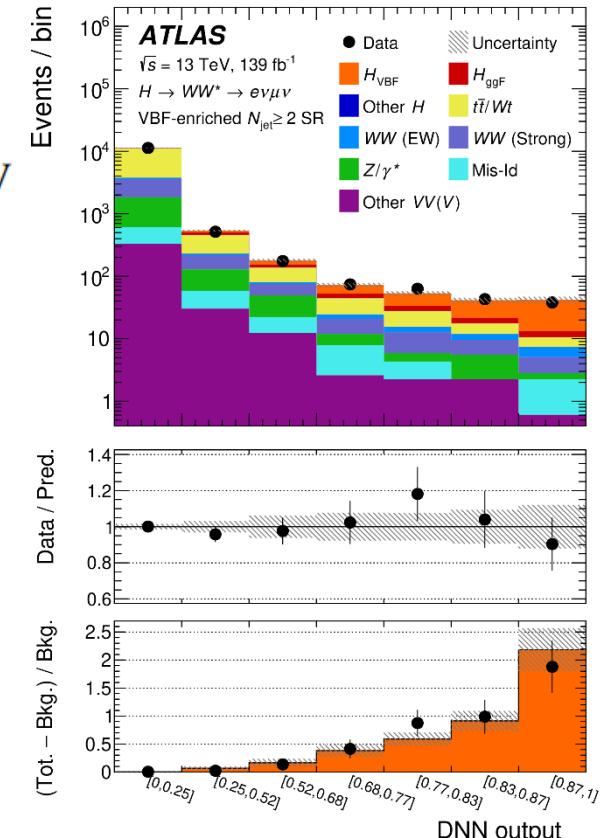
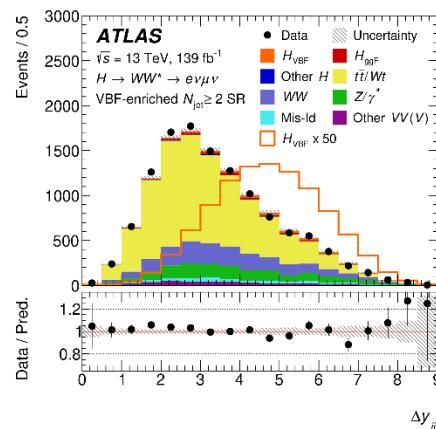
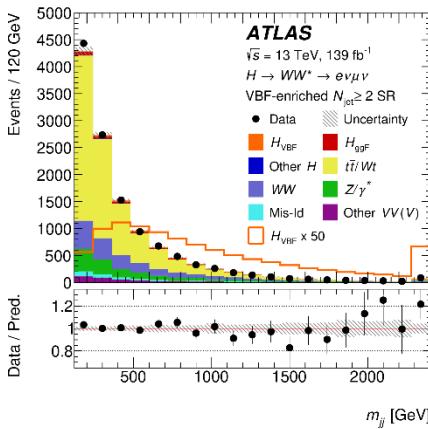
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DNN inputs:

$\Delta\phi_{\ell\ell}, m_{\ell\ell}, m_T, \Delta y_{jj}, m_{jj}, p_T^{\text{tot}}, \eta_\ell^{\text{centrality}}, m_{\ell 1j1}, m_{\ell 1j2}, m_{\ell 2j1}, m_{\ell 2j2}, p_T^{\text{jet}_1}, p_T^{\text{jet}_2}, p_T^{\text{jet}_3}$, and MET significance

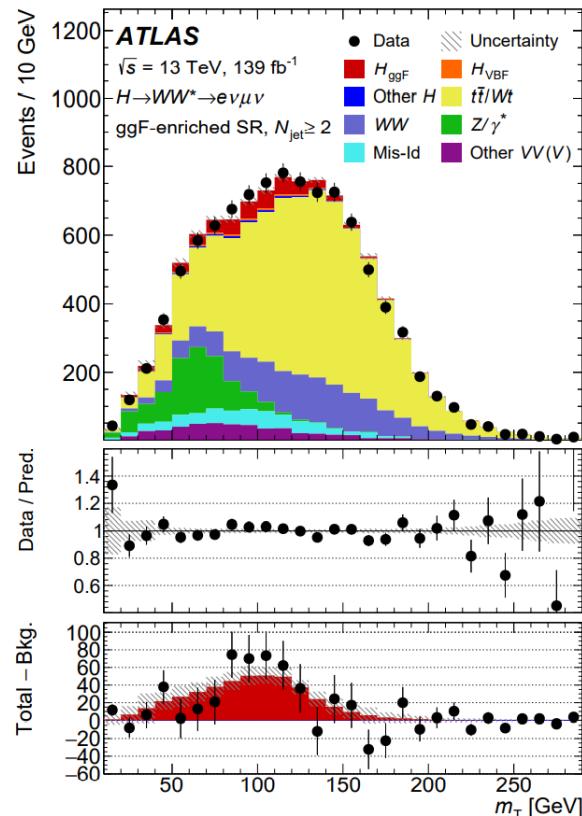
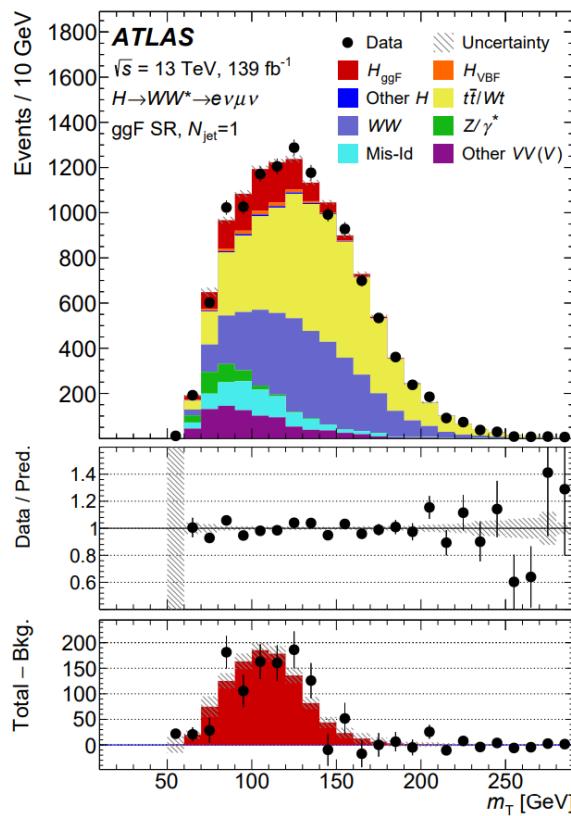
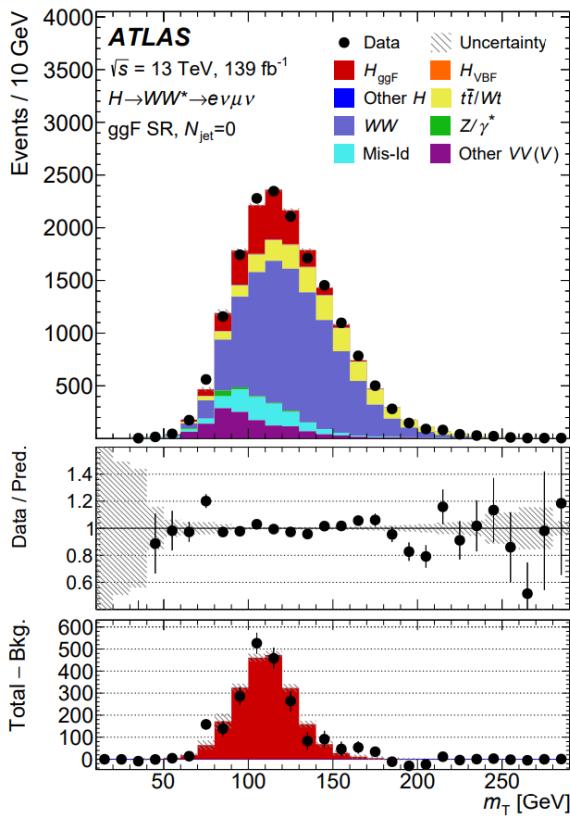
Training done after preselection, b veto && $m_{\tau\tau} < m_Z - 25 \text{ GeV}$

Training target: VBF VS all bkg(ggF, ttbar, ww*..)



Two most powerful variables

H \rightarrow WW*: analysis



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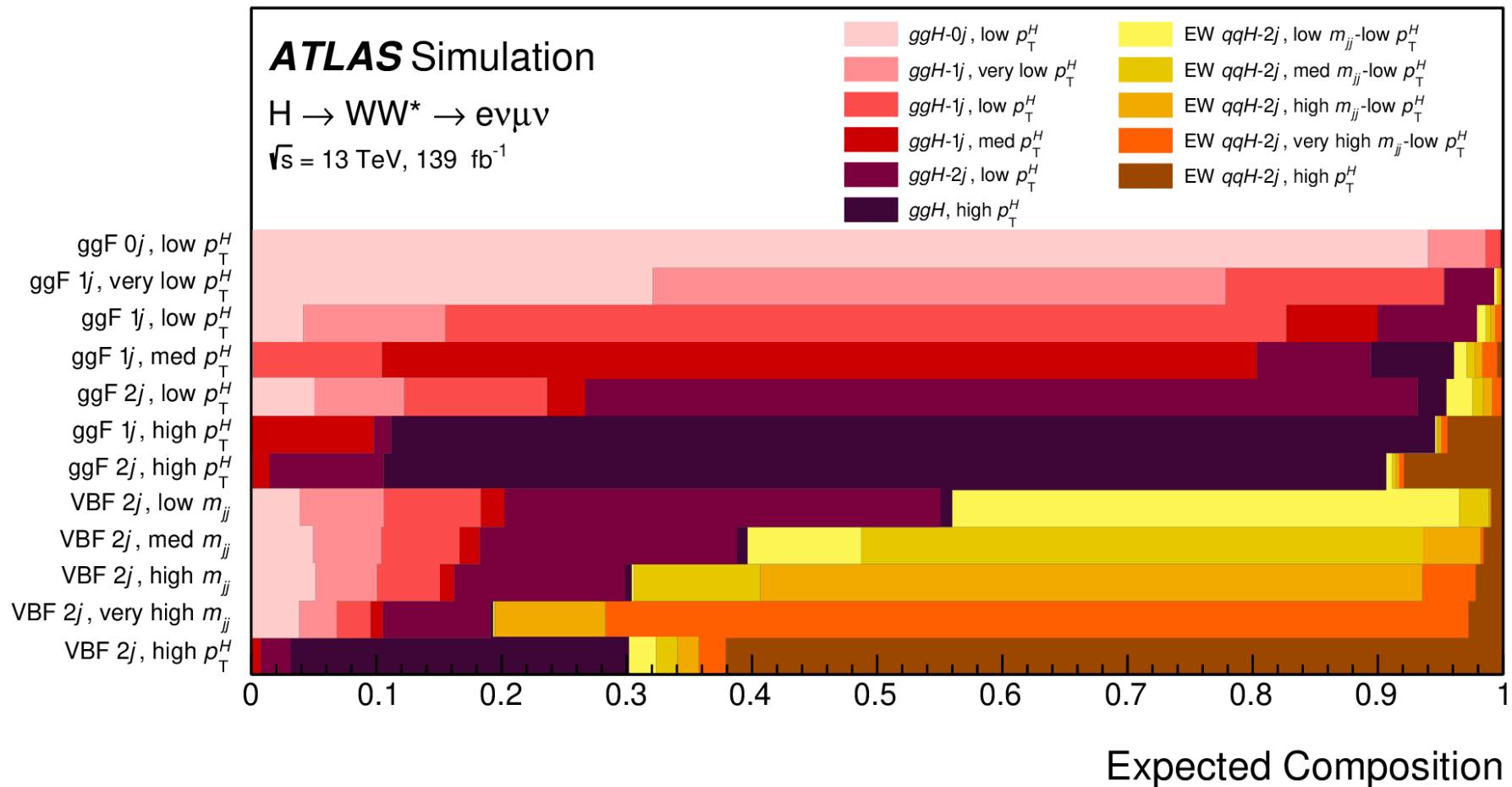
H \rightarrow WW*: Break down

| Source | $\frac{\Delta\sigma_{\text{ggF+VBF}} \cdot \mathcal{B}_{H \rightarrow WW^*}}{\sigma_{\text{ggF+VBF}} \cdot \mathcal{B}_{H \rightarrow WW^*}} [\%]$ | $\frac{\Delta\sigma_{\text{ggF}} \cdot \mathcal{B}_{H \rightarrow WW^*}}{\sigma_{\text{ggF}} \cdot \mathcal{B}_{H \rightarrow WW^*}} [\%]$ | $\frac{\Delta\sigma_{\text{VBF}} \cdot \mathcal{B}_{H \rightarrow WW^*}}{\sigma_{\text{VBF}} \cdot \mathcal{B}_{H \rightarrow WW^*}} [\%]$ |
|--------------------------------|--|--|--|
| Data statistical uncertainties | 4.6 | 5.1 | 15 |
| Total systematic uncertainties | 9.5 | 11 | 18 |
| MC statistical uncertainties | 3.0 | 3.8 | 4.9 |
| Experimental uncertainties | 5.2 | 6.3 | 6.7 |
| Flavor tagging | 2.3 | 2.7 | 1.0 |
| Jet energy scale | 0.9 | 1.1 | 3.7 |
| Jet energy resolution | 2.0 | 2.4 | 2.1 |
| E_T^{miss} | 0.7 | 2.2 | 4.9 |
| Muons | 1.8 | 2.1 | 0.8 |
| Electrons | 1.3 | 1.6 | 0.4 |
| Fake factors | 2.1 | 2.4 | 0.8 |
| Pileup | 2.4 | 2.5 | 1.3 |
| Luminosity | 2.1 | 2.0 | 2.2 |
| Theoretical uncertainties | 6.8 | 7.8 | 16 |
| ggF | 3.8 | 4.3 | 4.6 |
| VBF | 3.2 | 0.7 | 12 |
| WW | 3.5 | 4.2 | 5.5 |
| Top | 2.9 | 3.8 | 6.4 |
| Z $\tau\tau$ | 1.8 | 2.3 | 1.0 |
| Other VV | 2.3 | 2.9 | 1.5 |
| Other Higgs | 0.9 | 0.4 | 0.4 |
| Background normalizations | 3.6 | 4.5 | 4.9 |
| WW | 2.2 | 2.8 | 0.6 |
| Top | 1.9 | 2.3 | 3.4 |
| Z $\tau\tau$ | 2.7 | 3.1 | 3.4 |
| Total | 10 | 12 | 23 |

H \rightarrow WW*: STXS

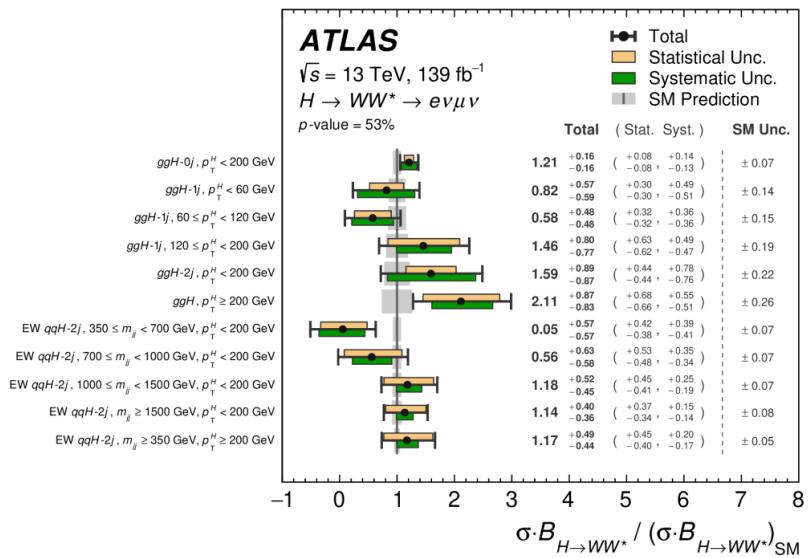
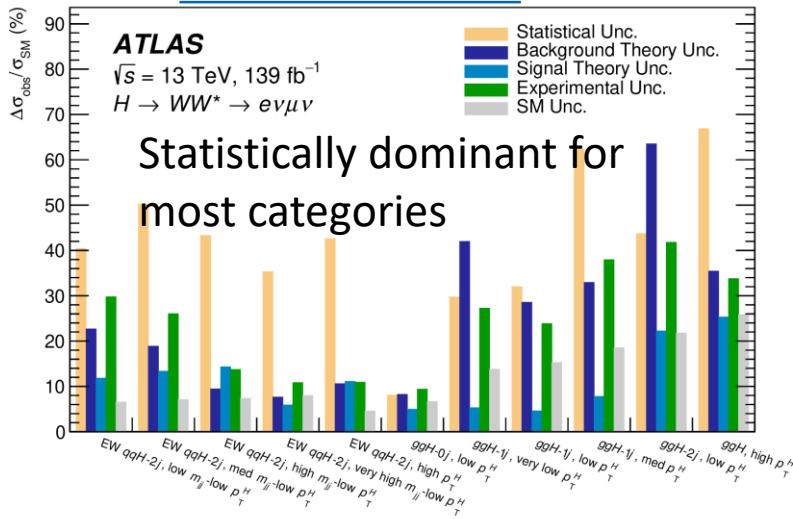
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Reconstructed Signal Region

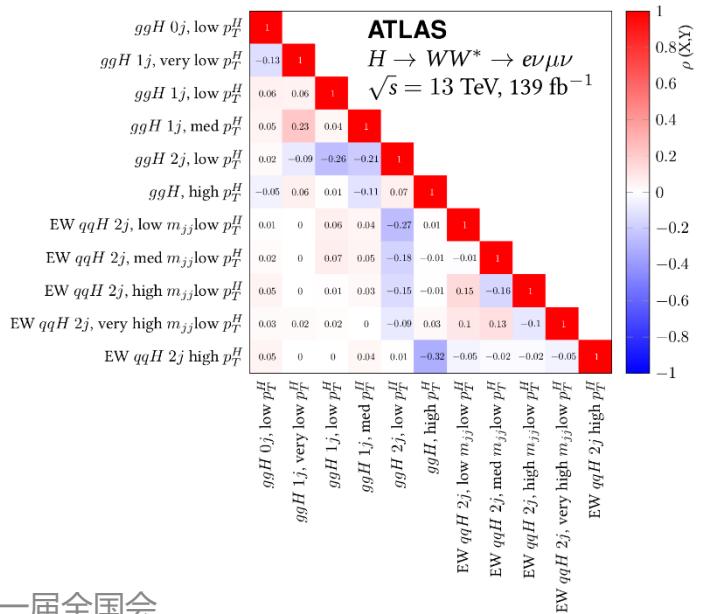


H \rightarrow WW*: STXS

Submitted to PRD



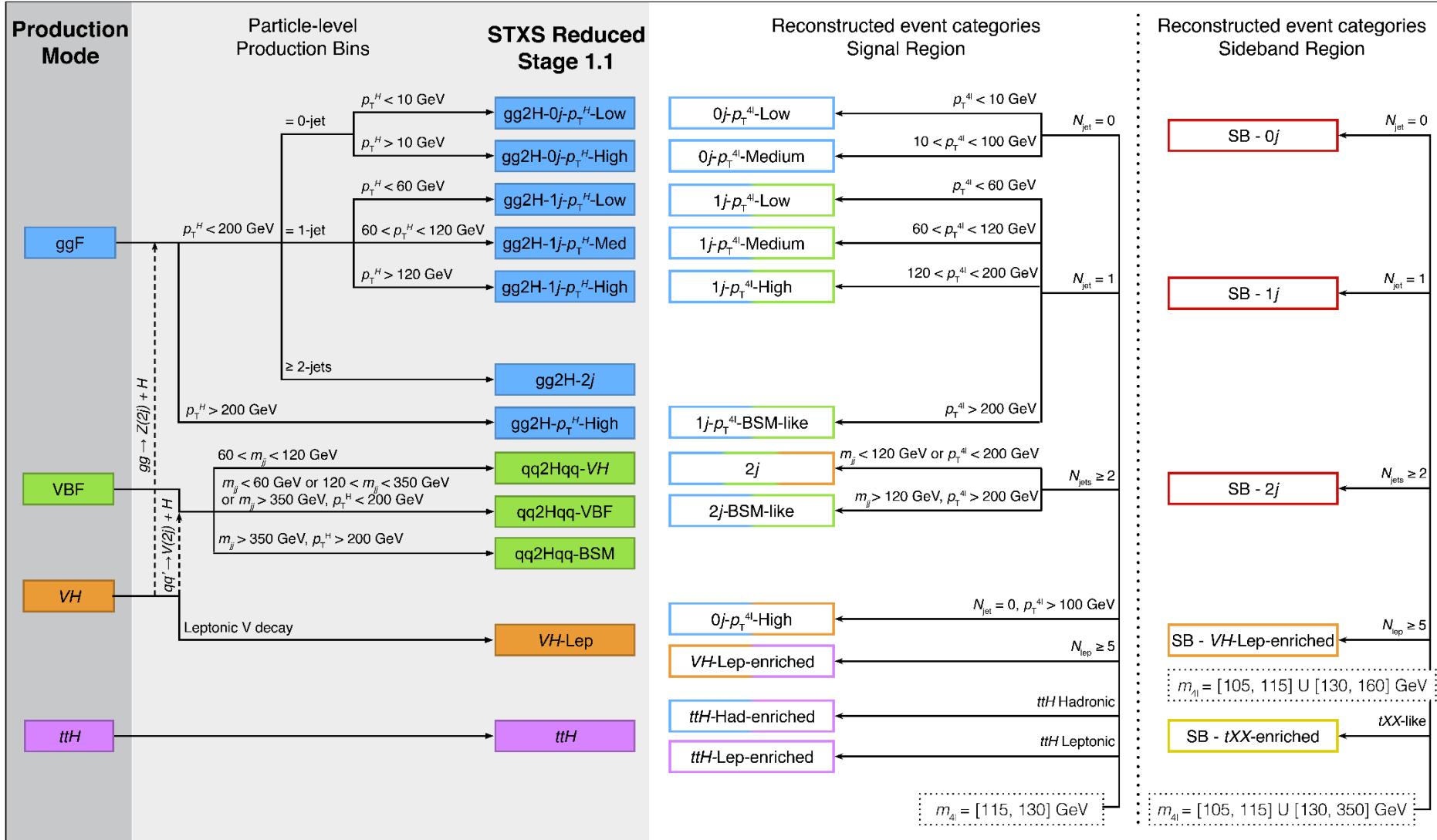
| STXS bin ($\sigma_i \cdot \mathcal{B}_{H \rightarrow WW^*}$) | Value [fb] | Uncertainty [fb] | | | | SM prediction [fb] | |
|---|------------|------------------|--------------|--------------|--------------|--------------------|----------------|
| | | Total | Stat. | Exp. Syst. | Sig. Theo. | | |
| ggH-0j, low p_T^H $p_T^H < 200 \text{ GeV}$ | 7100 | +950 -910 | +480 -470 | +570 -530 | +320 -260 | +490 -480 | 5870 ± 390 |
| ggH-1j, very low p_T^H $p_T^H < 60 \text{ GeV}$ | 1140 | +800 -820 | +420 -410 | +380 -380 | +80 -70 | +570 -600 | 1400 ± 190 |
| ggH-1j, low p_T^H $60 \leq p_T^H < 120 \text{ GeV}$ | 540 | +470 -470 | +310 -310 | +230 -230 | +42 -47 | +270 -280 | 970 ± 150 |
| ggH-1j, med p_T^H $120 \leq p_T^H < 200 \text{ GeV}$ | 230 | +130 -120 | +100 -100 | +60 -60 | +10 -10 | +50 -50 | 160 ± 30 |
| ggH-2j, low p_T^H $p_T^H < 200 \text{ GeV}$ | 1610 | +900 -890 | +440 -440 | +430 -420 | +300 -150 | +640 -650 | 1010 ± 220 |
| ggH, high p_T^H $p_T^H \geq 200 \text{ GeV}$ | 260 | +100 -100 | +80 -80 | +40 -40 | +40 -20 | +40 -40 | 122 ± 31 |
| EW qqH-2j, low m_{jj} -low p_T^H $350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}$ | 6 | +63 -62 | +46 -42 | +31 -34 | +11 -14 | +24 -26 | 109 ± 7 |
| EW qqH-2j, med m_{jj} -low p_T^H $700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}$ | 31 | +35 -33 | +30 -27 | +15 -14 | +8 -7 | +11 -10 | 56 ± 4 |
| EW qqH-2j, high m_{jj} -low p_T^H $1000 \leq m_{jj} < 1500 \text{ GeV}, p_T^H < 200 \text{ GeV}$ | 60 | +26 -23 | +23 -21 | +7 -7 | +9 -5 | +5 -5 | 51 ± 4 |
| EW qqH-2j, very high m_{jj} -low p_T^H $m_{jj} \geq 1500 \text{ GeV}, p_T^H < 200 \text{ GeV}$ | 57 | +20 -18 | +18 -17 | +5 -5 | +3 -3 | +4 -4 | 50 ± 4 |
| EW qqH-2j, high p_T^H $m_{jj} \geq 350 \text{ GeV}, p_T^H \geq 200 \text{ GeV}$ | 37 | +16 -14 | +14 -13 | +4 -3 | +4 -3 | +3 -3 | 32 ± 1 |



H \rightarrow ZZ*: STXS

Eur. Phys. J. C 80, 957 (2020).

ATLAS $\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$



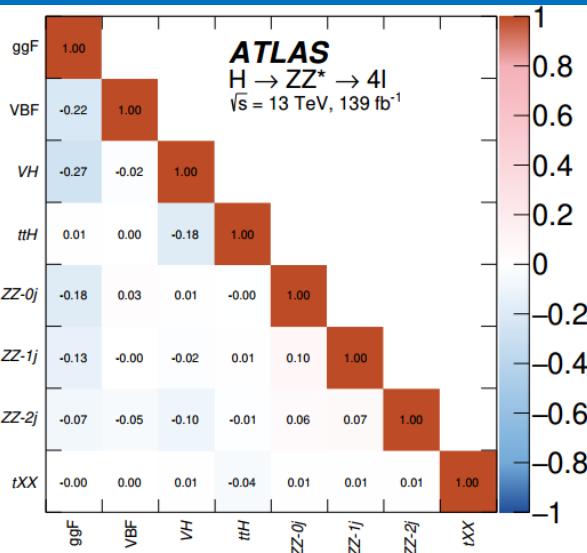
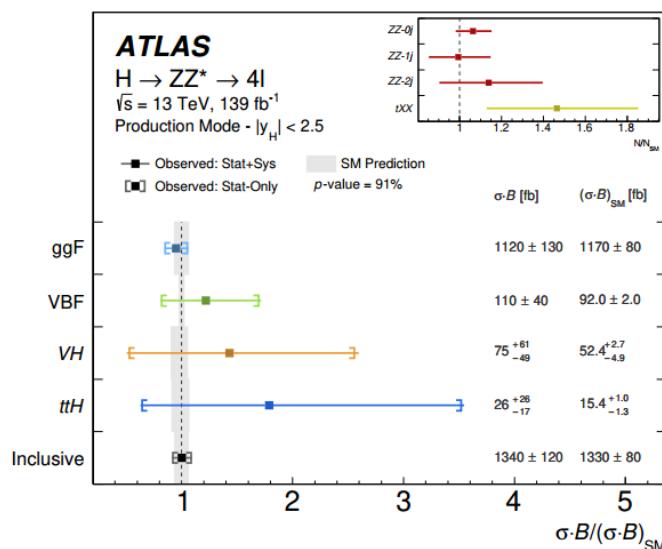
H \rightarrow ZZ*: STXS

[Eur. Phys. J. C 80, 957 \(2020\).](#)

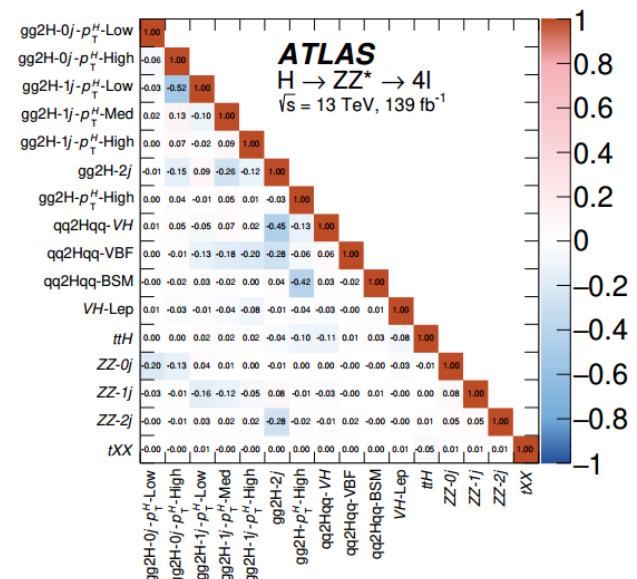
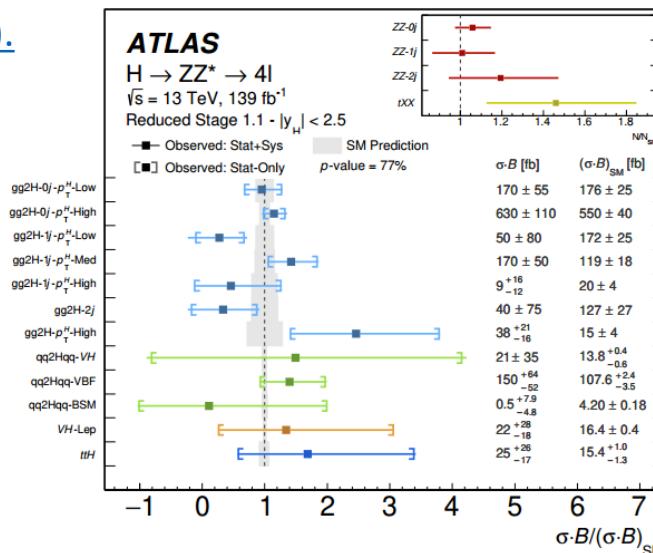
- Two types of NNs are used: feed-forward multilayer perceptron (MLP) and recurrent (RNN)
- The input variables used to train the MLP, and the two RNNs for the four leptons and the jets

| Category | Processes | MLP | Lepton RNN | Jet RNN | Discriminant |
|--|---------------|---|-----------------------|-----------------|--|
| 0j- $p_T^{4\ell}$ -Low 0j- $p_T^{4\ell}$ -Med | ggF, ZZ* | $p_T^{4\ell}, D_{ZZ^*}, m_{12}, m_{34}, \cos \theta^* , \cos \theta_1, \phi_{ZZ}$ | p_T^ℓ, η_ℓ | - | NN _{ggF} |
| 1j- $p_T^{4\ell}$ -Low | ggF, VBF, ZZ* | $p_T^{4\ell}, p_T^j, \eta_j, \Delta R_{4\ell j}, D_{ZZ^*}$ | p_T^ℓ, η_ℓ | - | NN _{VBF} for NN _{ZZ} < 0.25 NN _{ZZ} for NN _{ZZ} > 0.25 |
| 1j- $p_T^{4\ell}$ -Med | ggF, VBF, ZZ* | $p_T^{4\ell}, p_T^j, \eta_j, E_T^{\text{miss}}, \Delta R_{4\ell j}, D_{ZZ^*}, \eta_{4\ell}$ | p_T^ℓ, η_ℓ | - | NN _{VBF} for NN _{ZZ} < 0.25 NN _{ZZ} for NN _{ZZ} > 0.25 |
| 1j- $p_T^{4\ell}$ -High | ggF, VBF | $p_T^{4\ell}, p_T^j, \eta_j, E_T^{\text{miss}}, \Delta R_{4\ell j}, \eta_{4\ell}$ | p_T^ℓ, η_ℓ | - | NN _{VBF} |
| 2j | ggF, VBF, VH | $m_{jj}, p_T^{4\ell jj}$ | p_T^ℓ, η_ℓ | p_T^j, η_j | NN _{VBF} for NN _{VH} < 0.2 NN _{VH} for NN _{VH} > 0.2 |
| 2j-BSM-like | ggF, VBF | $\eta_{ZZ}^{\text{Zepp}}, p_T^{4\ell jj}$ | p_T^ℓ, η_ℓ | p_T^j, η_j | NN _{VBF} |
| VH-Lep-enriched | VH, ttH | $N_{\text{jets}}, N_{b\text{-jets}, 70\%}, E_T^{\text{miss}}, H_T$ | p_T^ℓ | - | NN _{ttH} |
| ttH-Had-enriched | ggF, ttH, tXX | $p_T^{4\ell}, m_{jj}, \Delta R_{4\ell j}, N_{b\text{-jets}, 70\%}$ | p_T^ℓ, η_ℓ | p_T^j, η_j | NN _{ttH} for NN _{tXX} < 0.4 NN _{tXX} for NN _{tXX} > 0.4 |

H \rightarrow ZZ*: STXS



[Eur. Phys. J. C](https://doi.org/10.1140/epjc/s10050-020-08570-0)
80, 957 (2020).



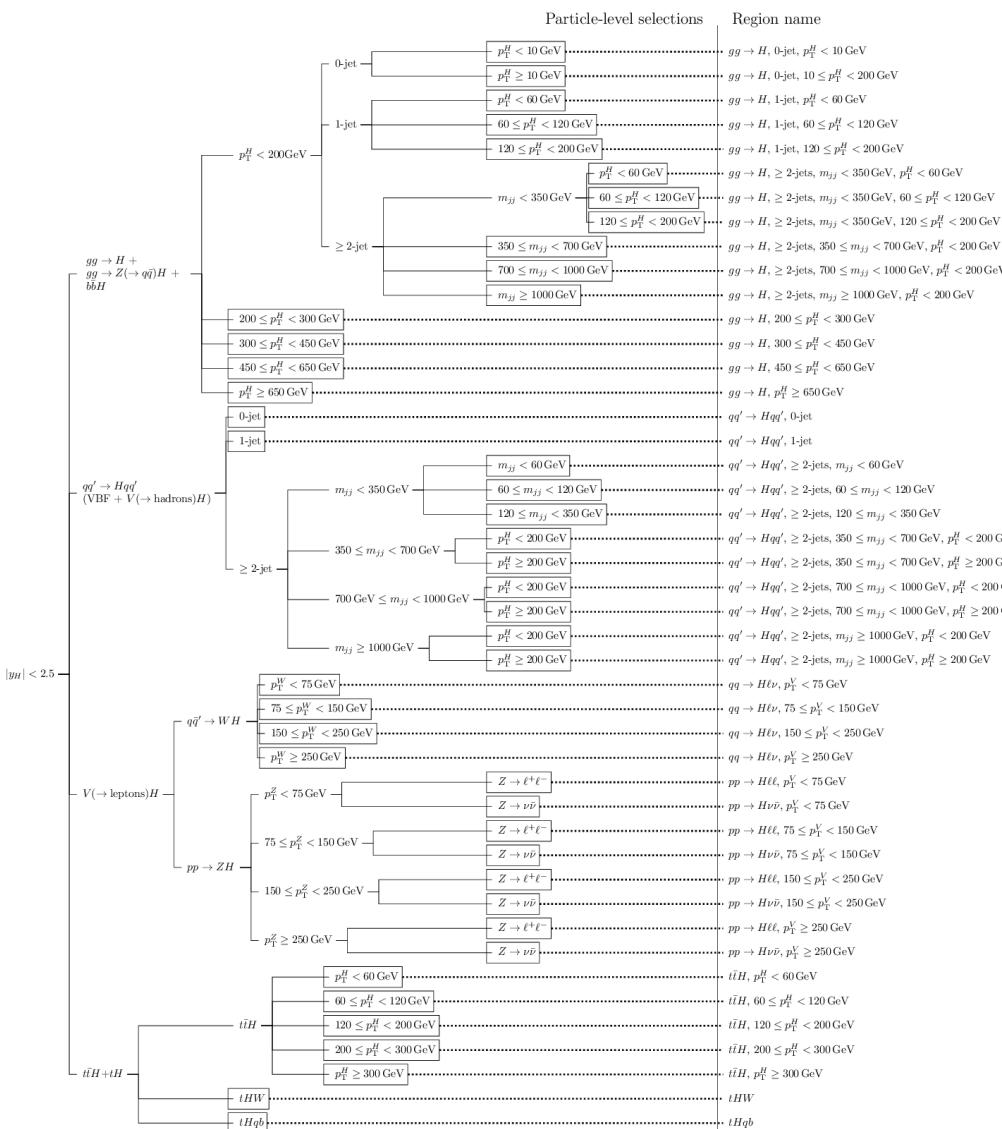
H \rightarrow ZZ*: STXS

| Measurement | Experimental uncertainties [%] | | | | Theory uncertainties [%] | | | | |
|---|--------------------------------|----------------------|--------------------|------------------|--------------------------|-------|--------|------|------|
| | Lumi. | $e, \mu,$ pile-up | Jets, flav. tag | Reducible bkg | Background | | Signal | | |
| | Inclusive cross-section | | | | | | | | |
| | 1.7 | 2.5 | 0.5 | < 0.5 | 1 | < 0.5 | < 0.5 | 1 | 2 |
| Production mode cross-sections | | | | | | | | | |
| ggF | 1.7 | 2.5 | 1 | < 0.5 | 1.5 | < 0.5 | 0.5 | 1 | 2 |
| VBF | 1.7 | 2 | 4 | < 0.5 | 1.5 | < 0.5 | 1 | 5 | 7 |
| VH | 1.9 | 2 | 4 | 1 | 6 | < 0.5 | 2 | 13.5 | 7.5 |
| ttH | 1.7 | 2 | 6 | < 0.5 | 1 | 0.5 | 0.5 | 12.5 | 4 |
| Reduced Stage-1.1 production bin cross-sections | | | | | | | | | |
| gg2H-0j- p_T^H -Low | 1.7 | 3 | 1.5 | 0.5 | 6.5 | < 0.5 | < 0.5 | 1 | 1.5 |
| gg2H-0j- p_T^H -High | 1.7 | 3 | 5 | < 0.5 | 3 | < 0.5 | < 0.5 | 0.5 | 5.5 |
| gg2H-1j- p_T^H -Low | 1.7 | 2.5 | 12 | 0.5 | 7 | < 0.5 | < 0.5 | 1 | 6 |
| gg2H-1j- p_T^H -Med | 1.7 | 3 | 7.5 | < 0.5 | 1 | < 0.5 | < 0.5 | 1.5 | 5.5 |
| gg2H-1j- p_T^H -High | 1.7 | 3 | 11 | 0.5 | 2 | < 0.5 | < 0.5 | 2 | 7.5 |
| gg2H-2j | 1.7 | 2.5 | 16.5 | 1 | 12.5 | 0.5 | < 0.5 | 2.5 | 10.5 |
| gg2H- p_T^H -High | 1.7 | 1.5 | 3 | 0.5 | 3.5 | < 0.5 | < 0.5 | 2 | 3.5 |
| qq2Hqq-VH | 1.8 | 4 | 17 | 1 | 4 | 1 | 0.5 | 5.5 | 8 |
| qq2Hqq-VBF | 1.7 | 2 | 3.5 | < 0.5 | 5 | < 0.5 | < 0.5 | 6 | 10.5 |
| qq2Hqq-BSM | 1.7 | 2 | 4 | < 0.5 | 2.5 | < 0.5 | < 0.5 | 3 | 8 |
| VH-Lep | 1.8 | 2.5 | 2 | 1 | 2 | 0.5 | < 0.5 | 1.5 | 3 |
| ttH | 1.7 | 2.5 | 5 | 0.5 | 1 | 0.5 | < 0.5 | 11 | 3 |

[Eur. Phys. J. C
80, 957 \(2020\).](#)

- The luminosity uncertainty, which is measured to be 1.7% and increases for the VH signal processes due to the simulation-based normalisation of the background.

$H \rightarrow \gamma\gamma$: 45 STXS analysis region



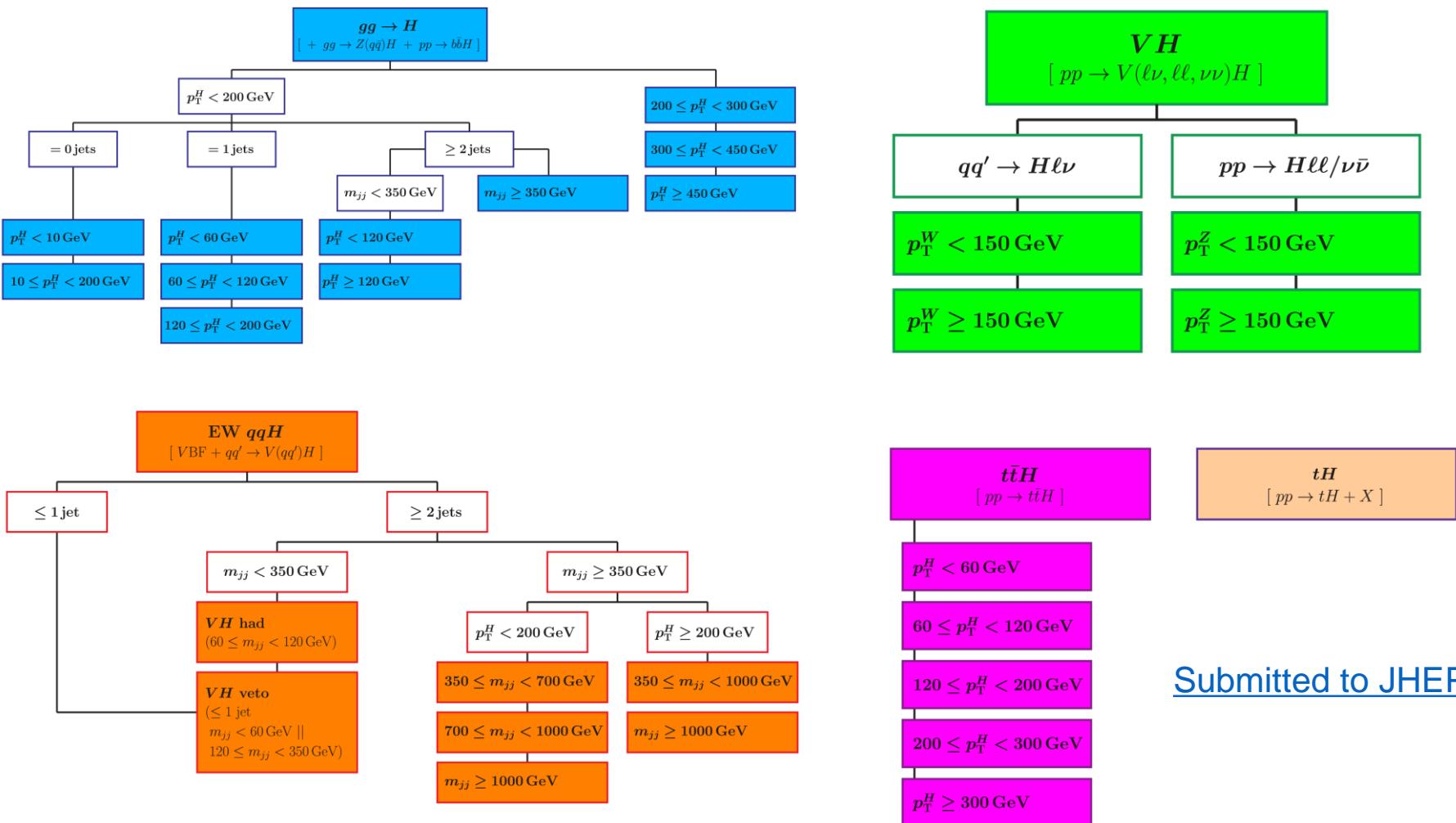
Submitted to JHEP

$H \rightarrow \gamma\gamma$:101 analysis categories

| Category | Function | N_{data} | N_{spur} | Wald | Category | Function | N_{data} | N_{spur} | Wald |
|---|----------|-------------------|-------------------|------|--|----------|-------------------|-------------------|------|
| $gg \rightarrow H$ | | | | | | | | | |
| 0-jet, $p_T^H < 10 \text{ GeV}$ | ExpPoly2 | 191623 | 64.8 | | $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{High-purity}$ | Exp | 18 | 0.189 | ✓ |
| 0-jet, $p_T^H \geq 10 \text{ GeV}$ | ExpPoly2 | 349266 | 50.4 | | $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{Med-purity}$ | Exp | 84 | 0.513 | ✓ |
| 1-jet, $p_T^H < 60 \text{ GeV}, \text{High-purity}$ | ExpPoly2 | 32644 | 20.7 | | $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{Low-purity}$ | Exp | 595 | 0.721 | |
| 1-jet, $p_T^H < 60 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 85229 | 24.9 | | $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{High-purity}$ | Exp | 19 | 0.110 | ✓ |
| 1-jet, $60 \leq p_T^H < 120 \text{ GeV}, \text{High-purity}$ | Exp | 26236 | 23.7 | | $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{Med-purity}$ | Exp | 411 | 0.193 | |
| 1-jet, $60 \leq p_T^H < 120 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 56669 | 21.3 | | $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{High-purity}$ | Exp | 23 | 1.30 | ✓ |
| 1-jet, $120 \leq p_T^H < 200 \text{ GeV}, \text{High-purity}$ | ExpPoly2 | 1570 | 1.48 | | $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H \geq 200 \text{ GeV}, \text{Med-purity}$ | Exp | 56 | 0.329 | ✓ |
| 1-jet, $120 \leq p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 6163 | 5.33 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, p_T^H < 60 \text{ GeV}, \text{High-purity}$ | ExpPoly2 | 8513 | 1.51 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, p_T^H < 60 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 31163 | 13.6 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, p_T^H < 60 \text{ GeV}, \text{Low-purity}$ | ExpPoly2 | 120357 | 15.7 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, 60 \leq p_T^H < 120 \text{ GeV}, \text{High-purity}$ | ExpPoly2 | 7582 | 2.26 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, 60 \leq p_T^H < 120 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 48362 | 6.21 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, 120 \leq p_T^H < 200 \text{ GeV}, \text{High-purity}$ | ExpPoly2 | 728 | 0.004 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 350 \text{ GeV}, 120 \leq p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | PowerLaw | 3007 | 0.983 | | | | | | |
| $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 432 | 0.487 | | | | | | |
| $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 3084 | 1.33 | | | | | | |
| $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Low-purity}$ | Exp | 7999 | 5.78 | | | | | | |
| $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 302 | 0.560 | | | | | | |
| $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | Exp | 1033 | 1.44 | | | | | | |
| $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Low-purity}$ | Exp | 3187 | 4.32 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 113 | 0.192 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | Exp | 332 | 0.804 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Low-purity}$ | PowerLaw | 1020 | 1.09 | | | | | | |
| $200 \leq p_T^H < 300 \text{ GeV}, \text{High-purity}$ | Exp | 420 | 1.68 | | | | | | |
| $200 \leq p_T^H < 300 \text{ GeV}, \text{Med-purity}$ | Exp | 2296 | 0.714 | | | | | | |
| $300 \leq p_T^H < 450 \text{ GeV}, \text{High-purity}$ | Exp | 25 | 0.407 | ✓ | | | | | |
| $300 \leq p_T^H < 450 \text{ GeV}, \text{Med-purity}$ | Exp | 186 | 0.259 | | | | | | |
| $300 \leq p_T^H < 450 \text{ GeV}, \text{Low-purity}$ | Exp | 422 | 0.121 | | | | | | |
| $450 \leq p_T^H < 650 \text{ GeV}, \text{High-purity}$ | Exp | 15 | 0.138 | ✓ | | | | | |
| $450 \leq p_T^H < 650 \text{ GeV}, \text{Med-purity}$ | Exp | 25 | 0.391 | ✓ | | | | | |
| $450 \leq p_T^H < 650 \text{ GeV}, \text{Low-purity}$ | Exp | 109 | 0.031 | | | | | | |
| $p_T^H \geq 650 \text{ GeV}$ | Exp | 14 | 0.448 | ✓ | | | | | |
| $qq' \rightarrow Hqq'$ | | | | | | | | | |
| 0-jet, High-purity | Exp | 176 | 0.180 | | | | | | |
| 0-jet, Med-purity | ExpPoly2 | 3238 | 4.73 | | | | | | |
| 0-jet, Low-purity | ExpPoly2 | 133314 | 49.7 | | | | | | |
| 1-jet, High-purity | Exp | 19 | 0.125 | ✓ | | | | | |
| 1-jet, Med-purity | Exp | 187 | 0.361 | | | | | | |
| 1-jet, Low-purity | PowerLaw | 1040 | 1.97 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 60 \text{ GeV}, \text{High-purity}$ | Exp | 17 | 0.499 | ✓ | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 60 \text{ GeV}, \text{Med-purity}$ | Exp | 157 | 0.489 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} < 60 \text{ GeV}, \text{Low-purity}$ | PowerLaw | 1978 | 1.29 | | | | | | |
| $\geq 2\text{-jets}, 60 \leq m_{jj} < 120 \text{ GeV}, \text{High-purity}$ | Exp | 53 | 0.165 | ✓ | | | | | |
| $\geq 2\text{-jets}, 60 \leq m_{jj} < 120 \text{ GeV}, \text{Med-purity}$ | Exp | 329 | 0.520 | | | | | | |
| $\geq 2\text{-jets}, 60 \leq m_{jj} < 120 \text{ GeV}, \text{Low-purity}$ | PowerLaw | 709 | 1.15 | | | | | | |
| $\geq 2\text{-jets}, 120 \leq m_{jj} < 350 \text{ GeV}, \text{High-purity}$ | Exp | 214 | 1.08 | | | | | | |
| $\geq 2\text{-jets}, 120 \leq m_{jj} < 350 \text{ GeV}, \text{Med-purity}$ | ExpPoly2 | 1671 | 1.07 | | | | | | |
| $\geq 2\text{-jets}, 120 \leq m_{jj} < 350 \text{ GeV}, \text{Low-purity}$ | PowerLaw | 11195 | 6.34 | | | | | | |
| $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 25 | 0.162 | ✓ | | | | | |
| $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | Exp | 260 | 0.443 | | | | | | |
| $\geq 2\text{-jets}, 350 \leq m_{jj} < 700 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Low-purity}$ | Exp | 753 | 1.17 | | | | | | |
| $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 25 | 0.670 | ✓ | | | | | |
| $\geq 2\text{-jets}, 700 \leq m_{jj} < 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | Exp | 166 | 0.713 | | | | | | |
| $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 48 | 1.47 | ✓ | | | | | |
| $\geq 2\text{-jets}, m_{jj} \geq 1000 \text{ GeV}, p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | Exp | 142 | 0.270 | | | | | | |
| $t\bar{t}H$ | | | | | | | | | |
| $p_T^H < 60 \text{ GeV}, \text{High-purity}$ | Exp | 35 | 0.040 | | | | | | |
| $p_T^H < 60 \text{ GeV}, \text{Med-purity}$ | Exp | 96 | 0.192 | | | | | | |
| $60 \leq p_T^H < 120 \text{ GeV}, \text{High-purity}$ | Exp | 34 | 0.038 | | | | | | |
| $60 \leq p_T^H < 120 \text{ GeV}, \text{Med-purity}$ | Exp | 74 | 0.274 | | | | | | |
| $120 \leq p_T^H < 200 \text{ GeV}, \text{High-purity}$ | Exp | 39 | 0.018 | | | | | | |
| $120 \leq p_T^H < 200 \text{ GeV}, \text{Med-purity}$ | Exp | 37 | 0.057 | | | | | | |
| $200 \leq p_T^H < 300 \text{ GeV}$ | Exp | 23 | 0.261 | | | | | | |
| $p_T^H \geq 300 \text{ GeV}$ | Exp | 19 | 0.180 | ✓ | | | | | |
| tH | | | | | | | | | |
| $tHqb, \text{High-purity}$ | Exp | 17 | 0.371 | ✓ | | | | | |
| $tHqb, \text{Med-purity}$ | Exp | 19 | 0.320 | ✓ | | | | | |
| $tHqb, \text{BSM} (\kappa_t = -1)$ | Exp | 14 | 0.496 | ✓ | | | | | |
| tHW | Exp | 38 | 0.070 | | | | | | |
| Low-purity top | Exp | 500 | 0.870 | | | | | | |

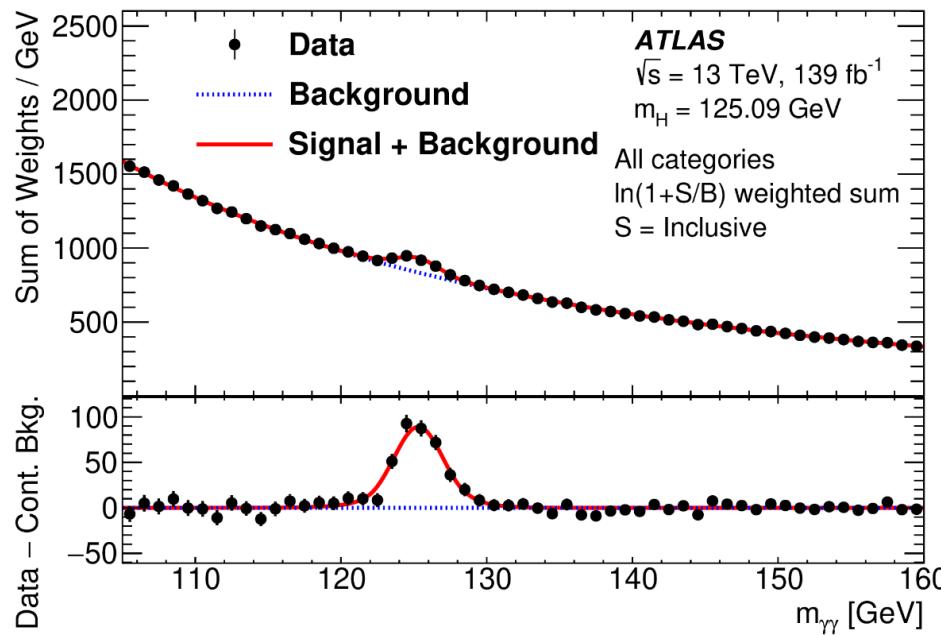
Submitted to JHEP

$H \rightarrow \gamma\gamma$: Measured STXS (28 pois)



Submitted to JHEP

$H \rightarrow \gamma\gamma$: Production Mode Cross Sections



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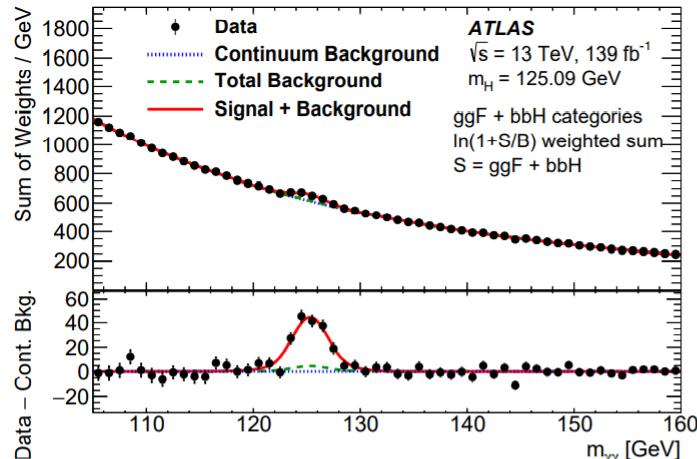
The events in each category are weighted by $\ln(1+S/B)$

The fitted signal-plus-background pdfs from all categories are also weighted and summed, shown as the solid line.

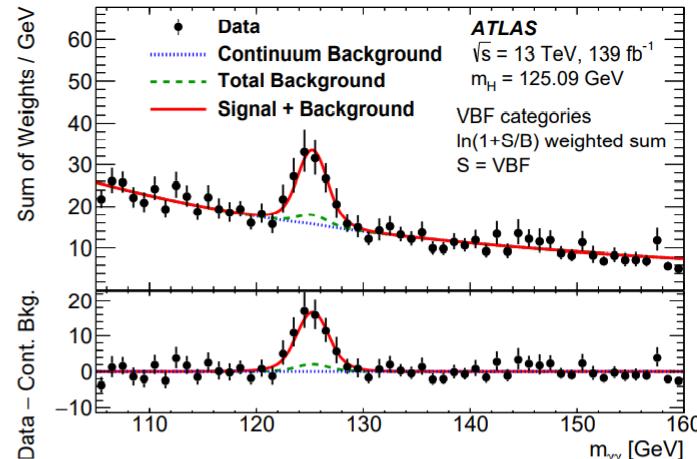
- This choice of event weight is designed to enhance the contribution of events from categories with higher signal-to-background ratio in a way that approximately matches the impact of these events in the categorized analysis of the data
- Shape description: a double-sided Crystal Ball (DSCB) function, a Gaussian distribution around the peak region, continued by power-law tails at lower and higher $m_{\gamma\gamma}$ values.

H $\rightarrow\gamma\gamma$: Production Mode Cross Sections

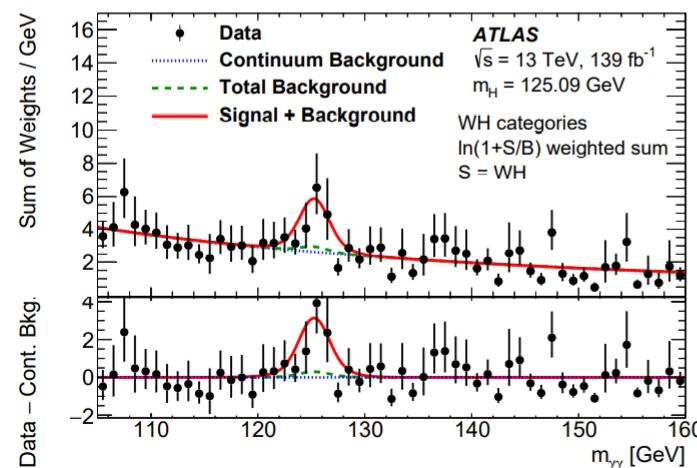
Submitted to JHEP



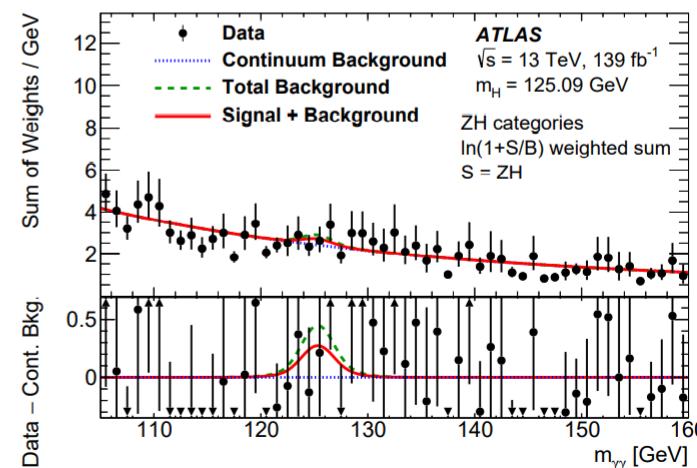
(a) ggF+ $b\bar{b}H$



(b) VBF



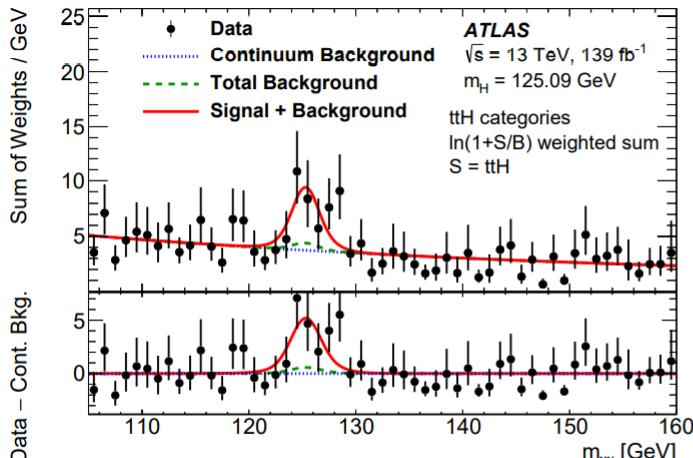
(c) WH



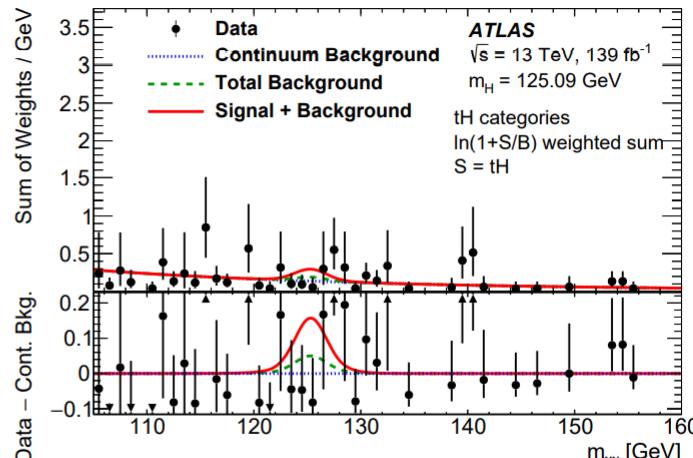
(d) ZH

$H \rightarrow \gamma\gamma$: Production Mode Cross Sections

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(e) $t\bar{t}H$



(f) tH

In this calculation, only Higgs boson events from the targeted production processes are considered as signal events.

Higgs boson events from other processes as well as the continuum background events are considered as background.

$H \rightarrow \gamma\gamma$: Production Mode Cross Sections (breakdown)

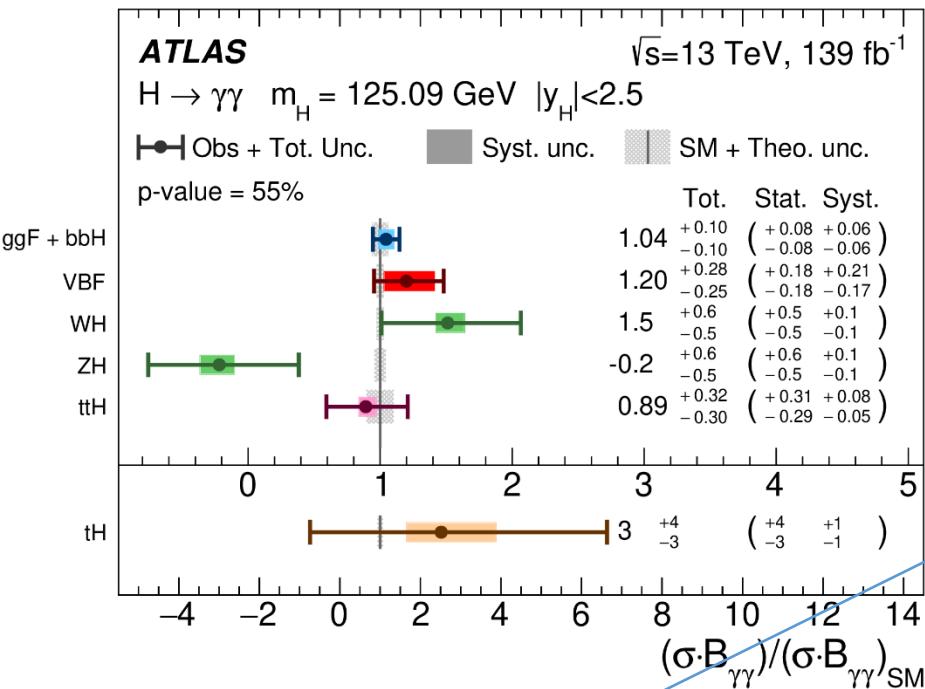
| Uncertainty source | $\Delta\mu [\%]$ |
|---|------------------|
| Theory uncertainties | |
| Higher-Order QCD Terms | ± 3.8 |
| Branching Ratio | ± 3.0 |
| Underlying Event and Parton Shower | ± 2.5 |
| PDF and α_s | ± 2.1 |
| Matrix Element | ± 1.0 |
| Modeling of Heavy Flavor Jets in non- $t\bar{t}H$ Processes | $< \pm 1$ |
| Experimental uncertainties | |
| Photon energy resolution | ± 2.8 |
| Photon efficiency | ± 2.6 |
| Luminosity | ± 1.8 |
| Pile-up | ± 1.5 |
| Background modelling | ± 1.3 |
| Photon energy scale | $< \pm 1$ |
| Jet/ E_T^{miss} | $< \pm 1$ |
| Flavour tagging | $< \pm 1$ |
| Leptons | $< \pm 1$ |
| Higgs boson mass | $< \pm 1$ |

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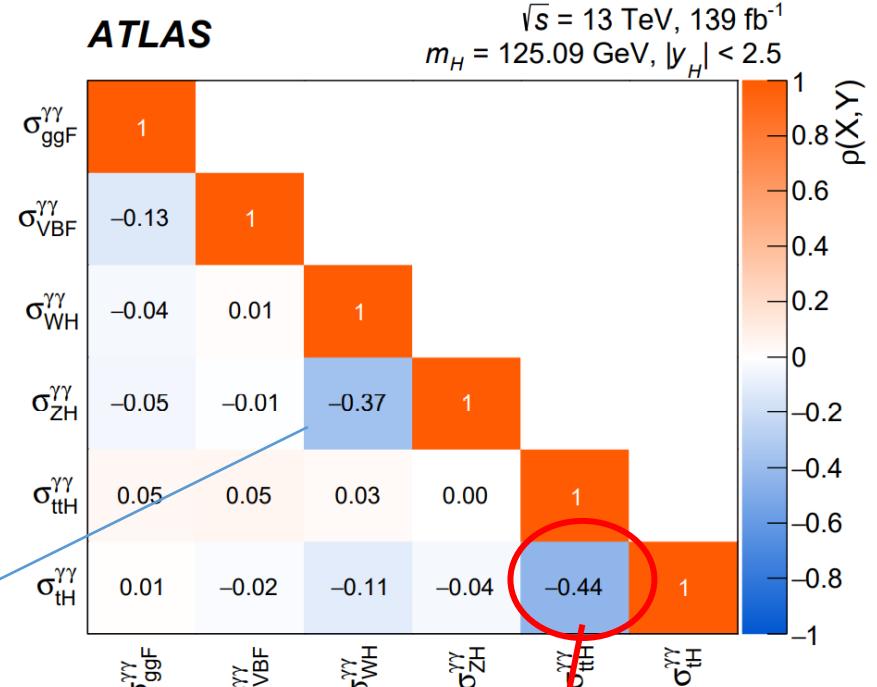
Inclusive case

H $\rightarrow\gamma\gamma$: Production Mode XSs

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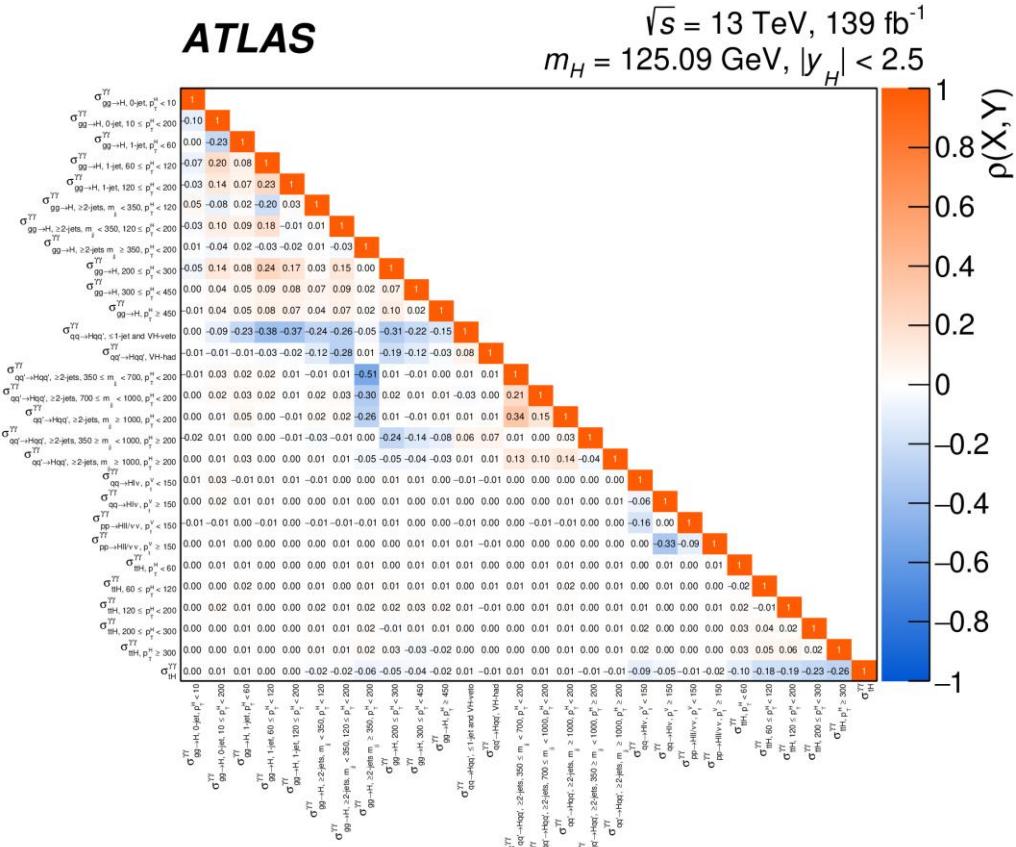
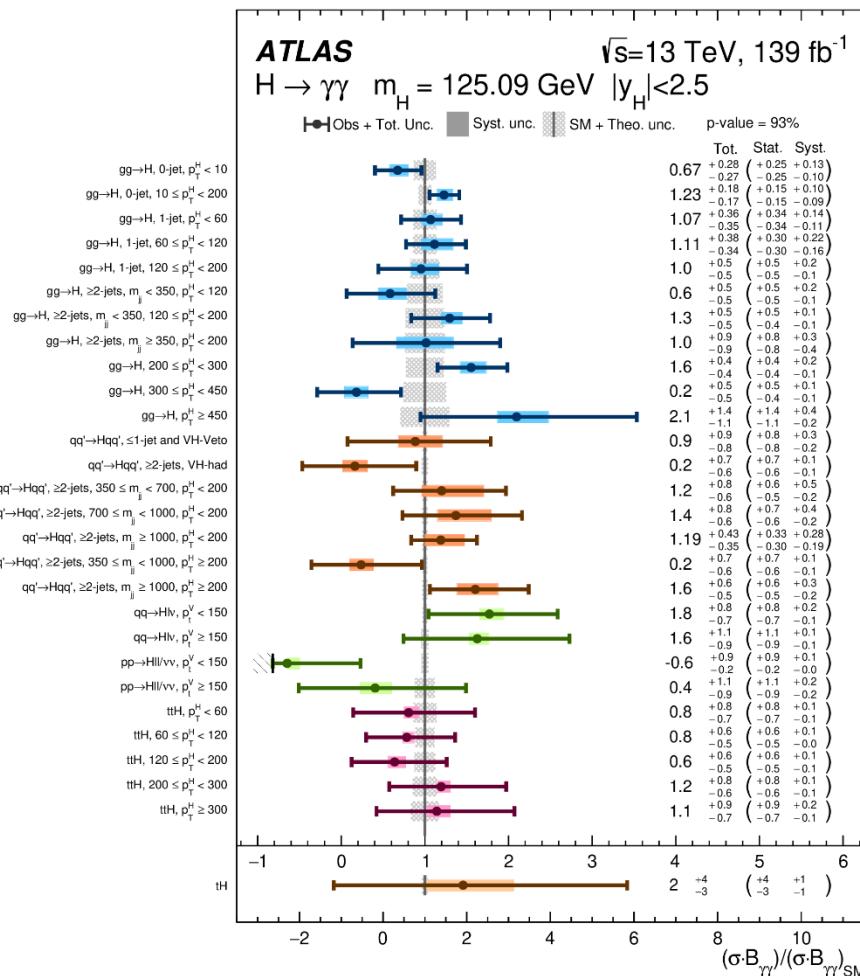
Due to contamination by $qq \rightarrow H\nu\nu$ events
 in the categories targeting the $pp \rightarrow H\nu\nu$
 process.



Due to the ttH contamination in the
 categories targeting tH

$H \rightarrow \gamma\gamma$: STXS (28 pois)

Submitted to JHEP



$H \rightarrow \gamma\gamma$: STXS (28 pois)

Table 1: Best-fit values and uncertainties for the production cross-section times $H \rightarrow \gamma\gamma$ branching ratio ($\sigma_i \times B_{\gamma\gamma}$) in each STXS region. The values for the $gg \rightarrow H$ process also include the contributions from $b\bar{b}H$ production. The total uncertainties are decomposed into statistical (Stat.) and systematic (Syst.) uncertainties. The uncertainties for the $pp \rightarrow H\ell\ell/\nu\bar{\nu}, p_T^V < 150$ GeV region are truncated at the value for which the model pdf becomes negative. SM predictions are also shown for each quantity with their total uncertainties.

| STXS region ($\sigma_i \times B_{\gamma\gamma}$) | Value [fb] | Uncertainty [fb] | | | SM prediction [fb] |
|--|---------------|------------------|----------------|----------------|---------------------------|
| | | Total | Stat. | Syst. | |
| $gg \rightarrow H, 0\text{-jet}, p_T^H < 10$ GeV | 10 | +4 -4 | +4 -4 | +2 -1 | 15^{+2}_{-2} |
| $gg \rightarrow H, 0\text{-jet}, p_T^H \geq 10$ GeV | 58 | +9 -8 | +7 -7 | +5 -4 | 47^{+4}_{-4} |
| $gg \rightarrow H, 1\text{-jet}, p_T^H < 60$ GeV | 16 | +5 -5 | +5 -5 | +2 -2 | 15^{+2}_{-2} |
| $gg \rightarrow H, 1\text{-jet}, 60 \leq p_T^H < 120$ GeV | 11 | +4 -3 | +3 -3 | +2 -2 | 10^{+1}_{-1} |
| $gg \rightarrow H, 1\text{-jet}, 120 \leq p_T^H < 200$ GeV | 1.6 | +0.9 -0.9 | +0.9 -0.8 | +0.4 -0.2 | $1.7^{+0.3}_{-0.3}$ |
| $gg \rightarrow H, \geq 2\text{-jets}, m_{jj} < 350$ GeV, $p_T^H < 120$ GeV | 4 | +4 -3 | +3 -3 | +1 -1 | 7^{+1}_{-1} |
| $gg \rightarrow H, \geq 2\text{-jets}, m_{jj} < 350$ GeV, $120 \leq p_T^H < 200$ GeV | 2.8 | +1.0 -1.0 | +1.0 -1.0 | +0.3 -0.2 | $2.1^{+0.5}_{-0.5}$ |
| $gg \rightarrow H, \geq 2\text{-jets}, m_{jj} \geq 350$ GeV, $p_T^H < 200$ GeV | 2 | +2 -2 | +2 -2 | +1 -1 | $2.0^{+0.5}_{-0.5}$ |
| $gg \rightarrow H, 200 \leq p_T^H < 300$ GeV | 1.6 | +0.4 -0.4 | +0.4 -0.4 | +0.2 -0.1 | $1.0^{+0.2}_{-0.2}$ |
| $gg \rightarrow H, 300 \leq p_T^H < 450$ GeV | 0.04 | +0.13 -0.11 | +0.12 -0.11 | +0.03 -0.03 | $0.24^{+0.06}_{-0.06}$ |
| $gg \rightarrow H, p_T^H \geq 450$ GeV | 0.09 | +0.06 -0.05 | +0.06 -0.05 | +0.02 -0.01 | $0.04^{+0.01}_{-0.01}$ |
| $qq' \rightarrow Hqq', \leq 1\text{-jet and } VH\text{-veto}$ | 6 | +6 -5 | +6 -5 | +2 -1 | $6.6^{+0.2}_{-0.2}$ |
| $qq' \rightarrow Hqq', VH\text{-had}$ | 0.19 | +0.85 -0.73 | +0.83 -0.71 | +0.17 -0.17 | $1.16^{+0.04}_{-0.04}$ |
| $qq' \rightarrow Hqq', \geq 2\text{-jets}, 350 \leq m_{jj} < 700$ GeV, $p_T^H < 200$ GeV | 1.5 | +0.9 -0.7 | +0.7 -0.6 | +0.6 -0.3 | $1.22^{+0.04}_{-0.04}$ |
| $qq' \rightarrow Hqq', \geq 2\text{-jets}, 700 \leq m_{jj} < 1000$ GeV, $p_T^H < 200$ GeV | 0.8 | +0.5 -0.4 | +0.4 -0.3 | +0.2 -0.1 | $0.58^{+0.02}_{-0.02}$ |
| $qq' \rightarrow Hqq', \geq 2\text{-jets}, m_{jj} \geq 1000$ GeV, $p_T^H < 200$ GeV | 1.2 | +0.4 -0.4 | +0.3 -0.3 | +0.3 -0.2 | $1.00^{+0.03}_{-0.03}$ |
| $qq' \rightarrow Hqq', \geq 2\text{-jets}, 350 \leq m_{jj} < 1000$ GeV, $p_T^H \geq 200$ GeV | 0.04 | +0.12 -0.10 | +0.12 -0.10 | +0.02 -0.02 | $0.167^{+0.005}_{-0.005}$ |
| $qq' \rightarrow Hqq', \geq 2\text{-jets}, m_{jj} \geq 1000$ GeV, $p_T^H \geq 200$ GeV | 0.27 | +0.11 -0.09 | +0.10 -0.08 | +0.05 -0.04 | $0.166^{+0.005}_{-0.005}$ |
| $qq \rightarrow H\ell\nu, p_T^V < 150$ GeV | 1.4 | +0.6 -0.6 | +0.6 -0.6 | +0.1 -0.1 | $0.79^{+0.02}_{-0.02}$ |
| $qq \rightarrow H\ell\nu, p_T^V \geq 150$ GeV | 0.20 | +0.13 -0.11 | +0.13 -0.11 | +0.02 -0.01 | $0.121^{+0.005}_{-0.005}$ |
| $pp \rightarrow H\ell\ell/\nu\bar{\nu}, p_T^V < 150$ GeV | -0.29 | +0.40 -0.08 | +0.39 -0.08 | +0.07 -0.00 | $0.45^{+0.02}_{-0.02}$ |
| $pp \rightarrow H\ell\ell/\nu\bar{\nu}, p_T^V \geq 150$ GeV | 0.04 | +0.10 -0.08 | +0.10 -0.08 | +0.02 -0.02 | $0.09^{+0.01}_{-0.01}$ |
| $t\bar{t}H, p_T^H < 60$ GeV | 0.22 | +0.21 -0.18 | +0.21 -0.18 | +0.03 -0.01 | $0.27^{+0.04}_{-0.04}$ |
| $t\bar{t}H, 60 \leq p_T^H < 120$ GeV | 0.32 | +0.23 -0.20 | +0.23 -0.20 | +0.04 -0.02 | $0.40^{+0.05}_{-0.05}$ |
| $t\bar{t}H, 120 \leq p_T^H < 200$ GeV | 0.18 | +0.18 -0.15 | +0.17 -0.15 | +0.04 -0.02 | $0.29^{+0.04}_{-0.04}$ |
| $t\bar{t}H, 200 \leq p_T^H < 300$ GeV | 0.14 | +0.09 -0.07 | +0.09 -0.07 | +0.01 -0.01 | $0.12^{+0.02}_{-0.02}$ |
| $t\bar{t}H, p_T^H \geq 300$ GeV | 0.06 | +0.05 -0.04 | +0.05 -0.04 | +0.01 -0.01 | $0.06^{+0.01}_{-0.01}$ |
| tH | 0.4 | +0.8 -0.6 | +0.7 -0.6 | +0.2 -0.2 | $0.19^{+0.01}_{-0.02}$ |

Submitted to JHEP

Results for STXS

H \rightarrow ZZ* \rightarrow 4 ℓ : Inclusive and Differential Cross Section

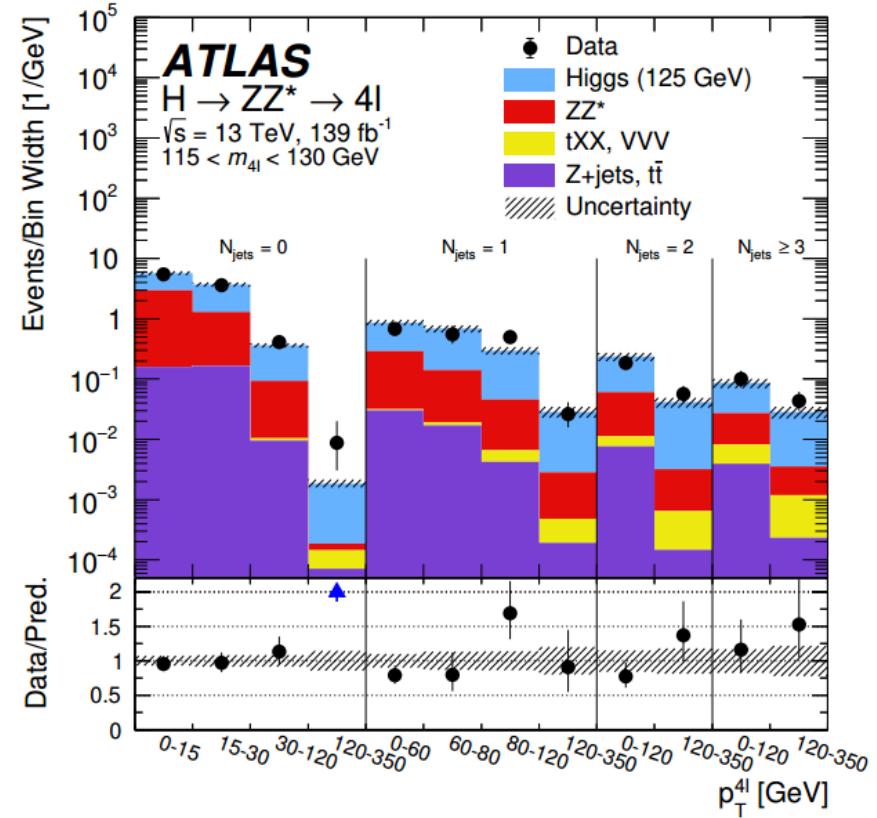
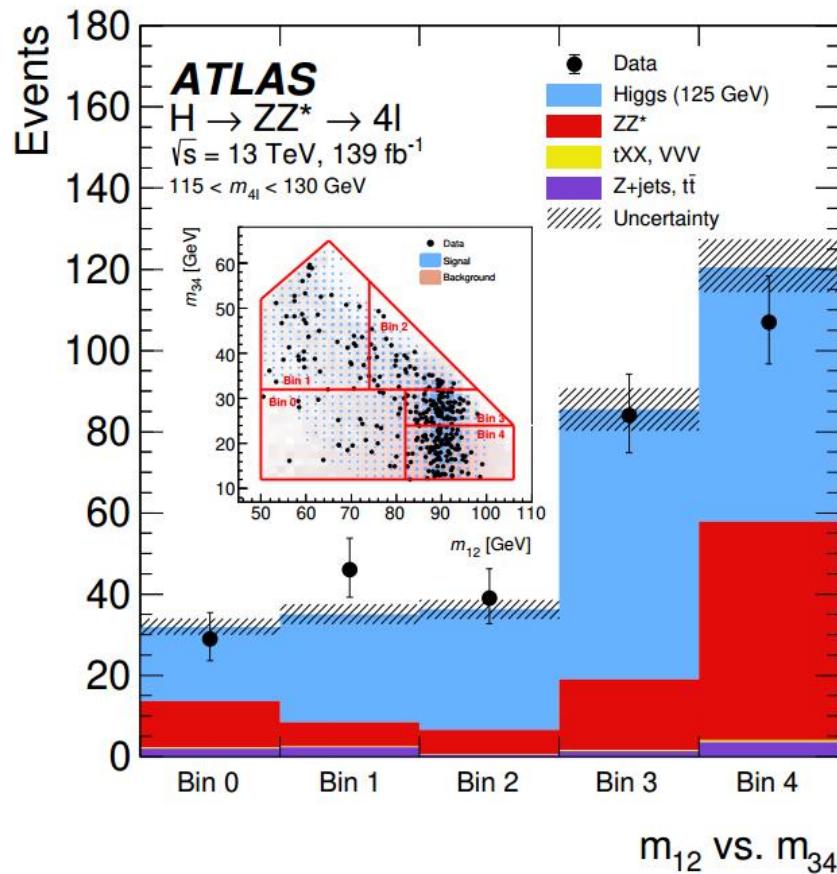
[Eur. Phys. J. C 80, 942 \(2020\)](#)

Matrix-based likelihood unfolding

| Higgs boson kinematic-related variables | |
|---|--|
| $p_T^{4\ell}$, $ y_{4\ell} $ | Transverse momentum and rapidity of the four-lepton system |
| m_{12} , m_{34} | Invariant mass of the leading and subleading lepton pair |
| $ \cos \theta^* $ | Magnitude of the cosine of the decay angle of the leading lepton pair in the four-lepton rest frame relative to the beam axis |
| $\cos \theta_1$, $\cos \theta_2$ | Production angles of the anti-leptons from the two Z bosons, where the angle is relative to the Z vector. |
| ϕ , ϕ_1 | Two azimuthal angles between the three planes constructed from the Z bosons and leptons in the Higgs boson rest frame. |
| Jet-related variables | |
| N_{jets} , $N_{b\text{-jets}}$ $p_T^{\text{lead. jet}}$, $p_T^{\text{sublead. jet}}$ | Jet and b -jet multiplicity Transverse momentum of the leading and subleading jet, for events with at least one and two jets, respectively. Here, the leading jet refers to the jet with the highest p_T in the event, while subleading refers to the jet with the second-highest p_T . |
| m_{jj} , $ \Delta\eta_{jj} $, $\Delta\phi_{jj}$ | Invariant mass, difference in pseudorapidity, and signed difference in ϕ of the leading and subleading jets for events with at least two jets |
| Higgs boson and jet-related variables | |
| $p_T^{4\ell j}$, $m_{4\ell j}$ | Transverse momentum and invariant mass of the four-lepton system and leading jet, for events with at least one jet |
| $p_T^{4\ell jj}$, $m_{4\ell jj}$ | Transverse momentum and invariant mass of the four-lepton system and leading and subleading jets, for events with at least two jets |

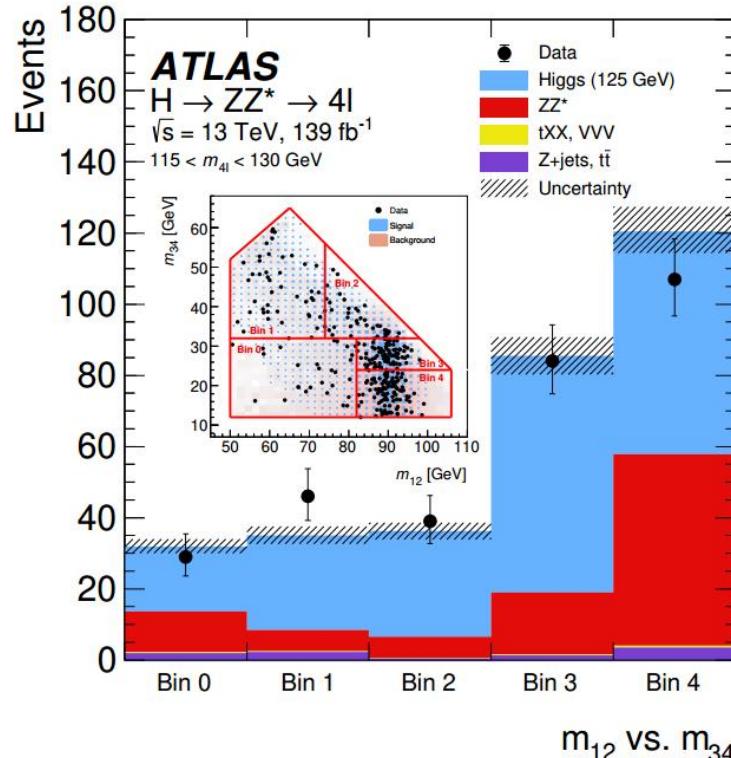
$H \rightarrow ZZ^* \rightarrow 4\ell$: Inclusive and Differential Cross Section

Eur. Phys. J. C 80, 942 (2020)



$H \rightarrow ZZ^* \rightarrow 4\ell$: Inclusive and Differential Cross Section

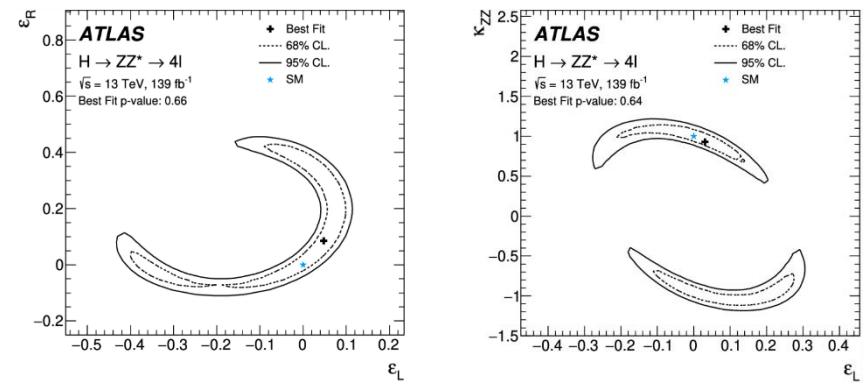
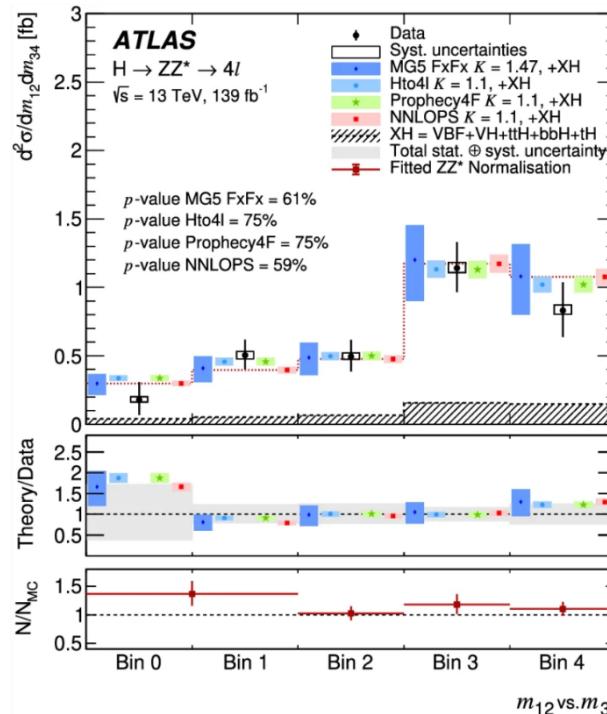
Eur. Phys. J. C 80, 942 (2020)



pseudo-observables ([arXiv:1504.04018](https://arxiv.org/abs/1504.04018)):

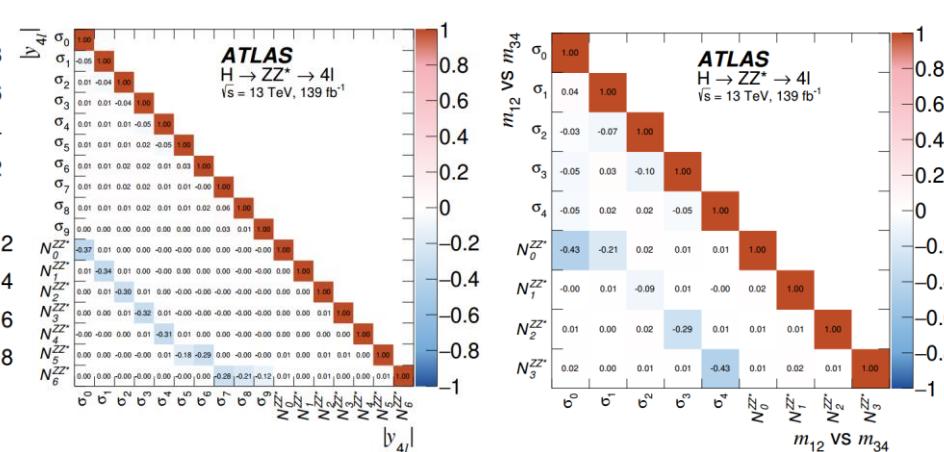
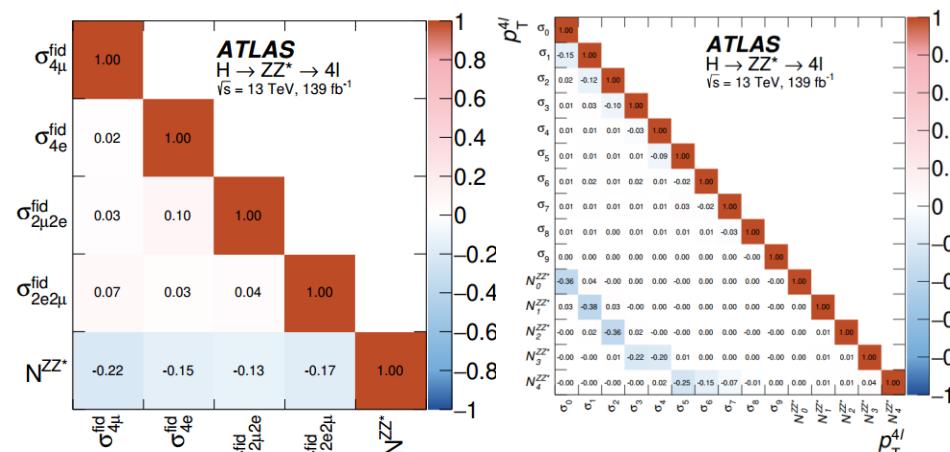
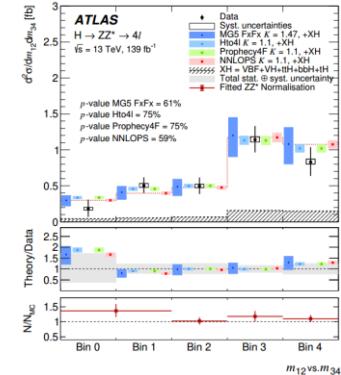
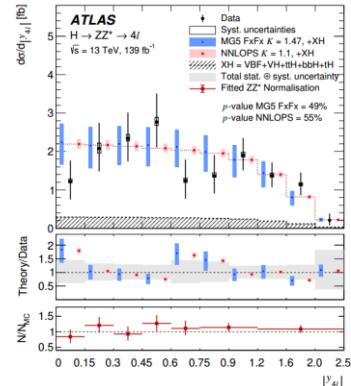
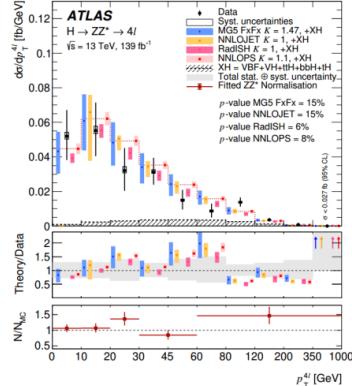
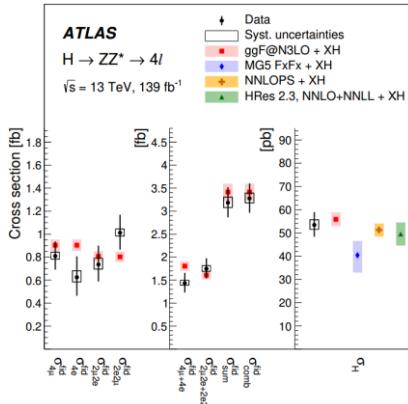
ϵ_L, ϵ_R : relating the couplings for left- and right-handed leptons to Higgs boson

κ_{ZZ} changes the coupling between H to Z boson



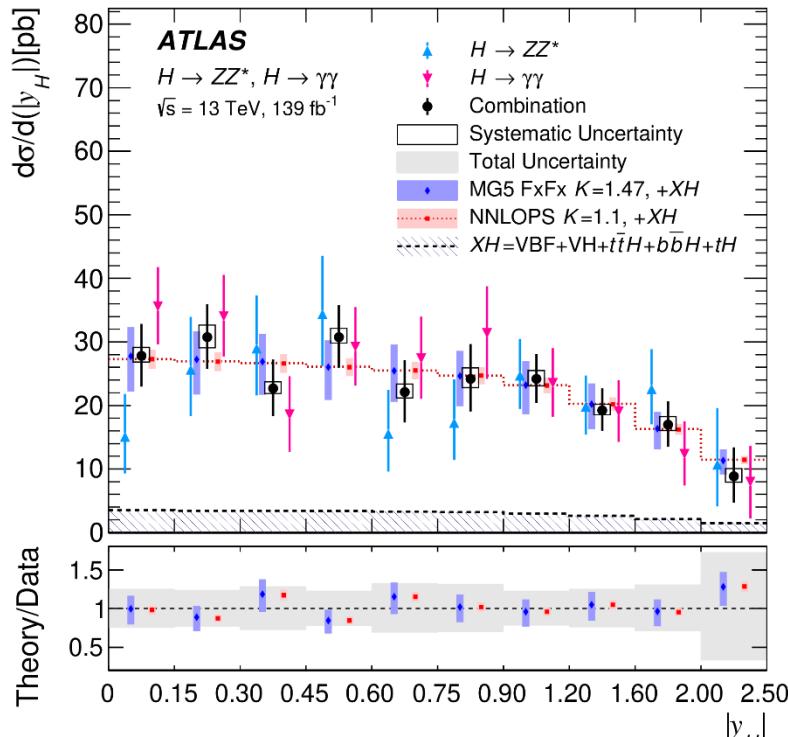
H \rightarrow ZZ* \rightarrow 4 ℓ : Inclusive and Differential Cross Section

Eur. Phys. J. C 80, 942 (2020)



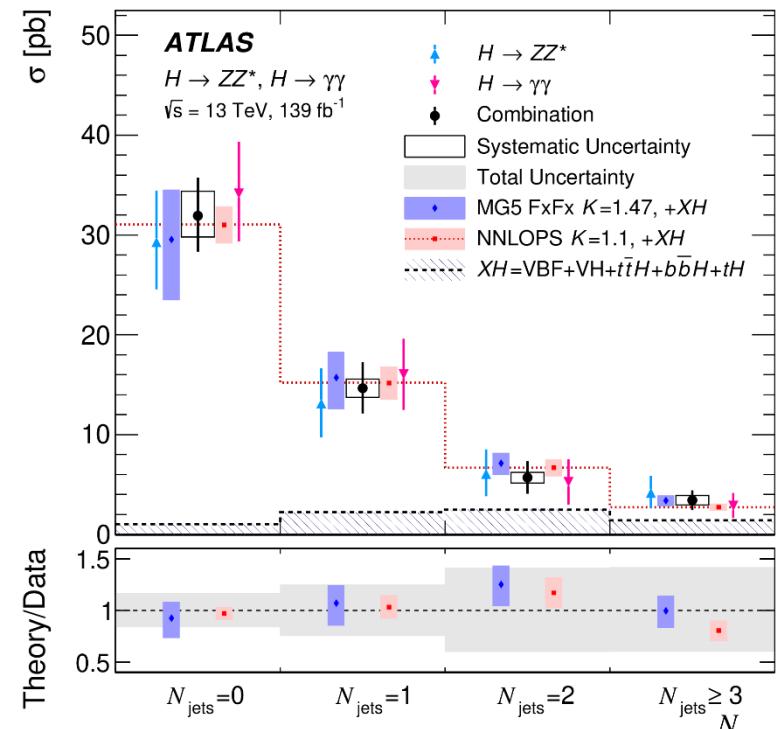
Combination: $H(\gamma\gamma)$ and $H(ZZ)$

Submitted to JHEP



y_H precision: 20-30% up to 2

Probes PDFs and pQCD modeling
of the ggF production mechanism

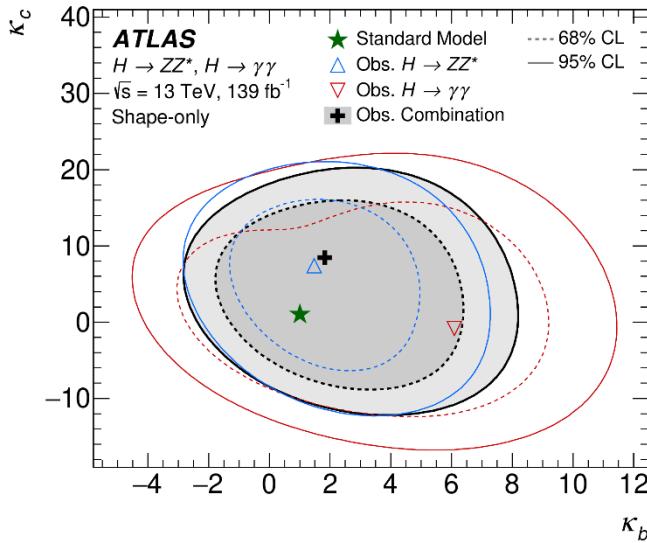


N_{jets} precision: 20-30%

Jet distributions can test pQCD modeling of
the ggF production mode, and contributions
of other production modes

Combination: H($\gamma\gamma$) and H(ZZ)

Submitted to JHEP



The combined observed limit on k_b less stringent than the individual ZZ* :

- Cross section quadratically depend on k_b
- Double minimum in the NLL scan
- k_b parameter being further from the SM expectation only with the H($\gamma\gamma$) channel

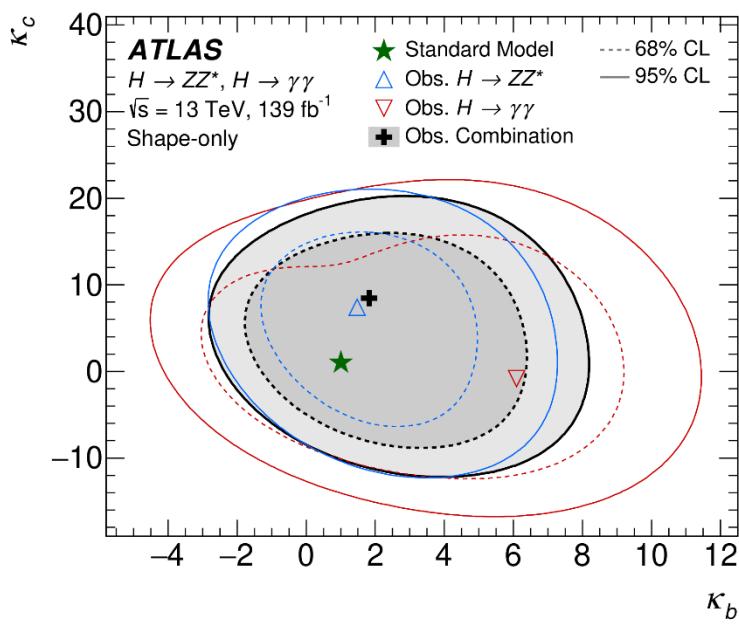
| Channel | Parameter | Observed | Expected |
|--|------------|-------------------------|-------------------------|
| | | 95% confidence interval | 95% confidence interval |
| $H \rightarrow ZZ^* \rightarrow 4\ell$ | κ_b | [−1.8, 6.4] | [−3.3, 9.3] |
| | κ_c | [−7.7, 18.3] | [−12.3, 19.2] |
| $H \rightarrow \gamma\gamma$ | κ_b | [−3.5, 10.2] | [−2.5, 8.0] |
| | κ_c | [−12.6, 18.3] | [−10.1, 17.3] |
| Combined | κ_b | [−2.0, 7.4] | [−2.0, 7.4] |
| | κ_c | [−8.6, 17.3] | [−8.5, 15.9] |

The results for one coupling modifier are obtained while fixing the other one to the SM expectation

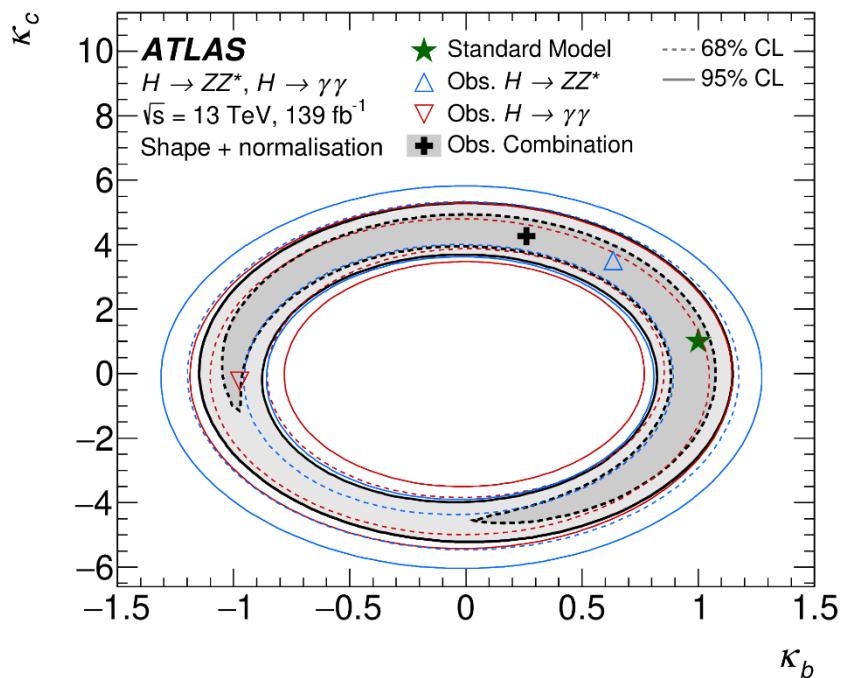
The combined observed limit on k_c similar to the individual ZZ*, but the 68% CL observed combined limits on k_c are worse than the results from the ZZ channel :

- The correlation between the k_b and k_c parameters
- Different best-fit k_b observations for ZZ and $\gamma\gamma$ channel
- Data fluctuations in some p_T^H bins

Interpretation with shape vs. shape + norm.



Shape only



Shape + normalization