Coherent elastic neutrino-nucleus scattering and quenching factor measurement

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arXiv: 2104.01811, Phys. Rev. D 104, 015005 (2021) arXiv: 2202.10622, accepted as a Letter in PRD

中国物理学会高能物理分会第十一届全国会员代表大会暨学术年会 大连,8/10/2022



- Introduction to CEvNS
- Measurement of Ge quenching factor
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Nuclear recoil



- Dark matter direct detection
- Coherent Elastic v-Nucleus Scattering (CEvNS)

Moment transfer $\longrightarrow q = \sqrt{2ME_R} \leq 1/R$ \leftarrow Nuclear radius Nuclear recoil energy $E_R \leq \frac{2E_v^2}{M+2E_v} \sim O(1) \text{ keV}$ for $E_v < 50 \text{ MeV}$ DM direct detection experiments \longrightarrow detection thresholds of O(1) keV

$CE\nu NS$ experiments

• π DAR source @ SNS

COHERENT first observed CE ν NS in 2017 at the 6.7 σ CL with a Csl detector

COHERENT, Science 357,1123 (2017)

Later confirmed in 2020 at more than 3σ CL with Ar detector

COHERENT, PRL 126, 012002 (2021)

Reactor antineutrino source

CONNIE uses a Si detector with 0.1 keV_{ee}threshold CONNIE, PRD 100, 092005 (2019)

CONUS uses a Ge detector with 0.3 keV_{ee} threshold

CONNIE, PRL 126, 041804 (2021)

NCC-1701 uses a Ge detector with 0.2 keV_{ee} threshold

Colaresi at al. , PRD 104, 072003 (2021) [2108.02880] Colaresi at al. , arXiv: 2202.09672



10¹⁰

0

2

4 Energy (MeV) 6

$CE\nu NS$ spectrum

• Differential cross section

$$\frac{d\sigma_{SM}}{dE_R} = \frac{G_F^2 M}{4\pi} q_W^2 \left(1 - \frac{ME_R}{2E_\nu^2}\right) F^2(\mathfrak{q})$$

• Event spectrum

$$\frac{dR}{dE_R} = N_T \int \frac{d\Phi}{dE_\nu} \frac{d\sigma}{dE_R} dE_\nu$$

Only a small portion of nuclear recoiling energy E_R will go into electronic ionization energy E_I , which is measured.

Quenching factor (QF):
$$Q \equiv E_I/E_R$$

Measured number of events:

$$N_i = t \int_{E_I^i}^{E_I^{i+1}} \eta \frac{dR}{dE_R} \left(\frac{1}{Q} - \frac{E_I}{Q} \frac{dQ}{dE_I}\right) dE_I$$

 $\frac{dE_R}{dE_I}$

Lindhard Model

Lindhard, Nielsen, Scharff and Thomsen, Mat. Fys. Medd. Dan. Vid. Selsk. 33 10 (1963)

$$Q(E_R) = \frac{k g(\epsilon)}{1 + k g(\epsilon)},$$

where $g(\epsilon) = 3 \epsilon^{0.15} + 0.7 \epsilon^{0.6} + \epsilon$

Dimensionless reduced energy

$$\epsilon = 11.5 Z^{-\frac{7}{3}} \left(\frac{E_R}{\text{keV}_{\text{nr}}}\right)$$

The slope of electronic stopping power $k=0.1333Z^{\frac{2}{3}}A^{-\frac{1}{2}}$

a larger k value leads to larger fraction of total energy going into electron.

Key approximations:

- The atomic binding energy of electrons is negligible.
- Energy transfers to electrons are small relative to energy transfers to atoms. Sorensen, PRD 91, 083509 (2015) [arXiv: 1412.3028]



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Measurement of Ge QF



The new dataset may be explained by the Lindhard model supplemented with the Migdal effect. Collar, Kavner, Lewis, PRD 103, 122003 (2021) [arXiv: 2102.10089]

Modified Lindhard Model

$$Q(E_R) = \frac{k g(\epsilon)}{1 + k g(\epsilon)} - \frac{q}{\epsilon}$$

Sorensen, PRD 91, 083509 (2015) [arXiv: 1412.3028]

- A small q value has no significant effect at the large recoil energy region.
- A positive q value allows a sharp cutoff in the energy given to electrons.
- A negative q value allows an enhancement in the energy given to electrons.



JL, Liu, Marfatia, PRD 104, 015005 (2021) [arXiv: 2104.01811]

New Physics in CEvNS

• A light vector Z'

$$\frac{d\sigma_{SM+Z'}}{dE_R} = \left(1 - \frac{q_{Z'}}{q_W}\right)^2 \frac{d\sigma_{SM}}{dE_R}$$
$$q_{Z'} = \frac{3\sqrt{2}\left(N+Z\right){g'}^2}{G_F\left(2ME_R + M_{Z'}^2\right)}$$

• A light scalar ϕ

$$\frac{d\sigma_{SM+\phi}}{dE_R} = \frac{d\sigma_{SM}}{dE_R} + \frac{d\sigma_{\phi}}{dE_R}$$
$$\frac{d\sigma_{\phi}}{dE_R} = \frac{G_F^2}{4\pi} q_{\phi}^2 \frac{2ME_R}{E_{\nu}^2} MF^2(\mathbf{q})$$
$$q_{\phi} = \frac{(14N+15.1Z) g_{\phi}^2}{\sqrt{2}G_F(2ME_R+M_{\phi}^2)}$$

10⁵ SM w/ Migdal : $k = 0.160, q = -13 \times 10^{-5}$ Z' w/ Lindhard : $g_{Z'} = 6.71 \times 10^{-5}$, $m_{Z'} = 32.1$ MeV 10⁴ - SM w/ Lindhard Counts/bin background 1000 $15 \text{ counts/keV}_{ee}/kg/day$ 100 10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 E_I (keV_{ee}) 10⁵ SM w/ Migdal : $k = 0.160, q = -10 \times 10^{-5}$ ----- ϕ w/Lindhard : $g_{\phi} = 6.04 \times 10^{-6}$, $m_{\phi} = 10.0$ MeV 10⁴ SM w/ Lindhard Counts/bin background 1000 100 10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 E_I (keV_{ee})

d = 20 m

P = 3.9 GW

t = 7 kg-year



- Both the light Z' and scalar cases with the standard Lindhard model can fit the SM spectrum for a given set of QF parameters k and q.
- This will lead to confusion in determining the nature of new physics.

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Suggestive evidence for coherent elastic neutrino-nucleus scattering from reactor antineutrinos arXiv: 2202.09672

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Solid (dotted) line shows the SM CEvNS prediction for the iron filter (Lindhard) QF.

Blue (black) line are based on the data from the iron filter (photo-neutron) measurement.

Quenching factor sensitivity



CEvNS and Quenching Factor

New Physics



- Assume the standard Lindhard QF model is valid.
- 1σ allowed region of light Z' is also allowed by COHERENT due to

$$q_{Z'} = 2q_W \implies \frac{g_{Z'}}{m_{Z'}} = \sqrt{\frac{\sqrt{2}G_F \left[N - (1 - 4\sin^2\theta_W)Z\right]}{3(N+Z)}}$$

New Physics



- Marginalize over the (k,q) of the modified Lindhard model.
- Constraints are qualitatively affected by the QF model.

Summary

- Recent direct QF measurement in Ge indicates a departure from the standard Lindhard model at low energies.
- This deviation can be parameterized by a negative q in the modified Lindhard model.
- NCC-1701 data provide an independent probe of Ge quenching factor, and the best-fit point is consistent with the direct QF measurements.
- A precise measurement of the QF is essential to detect new physics at CEvNS.
 Thanks!

Backup slides

Migdal Effect

"...it is usually assumed that the atomic electrons around the recoil nucleus immediately follow the motion of the nucleus. However, it takes some time for the electrons to catch up, which causes ionization and excitation of the recoil atom..."

Migdal, J.Phys.(USSR) 4, 449 (1941).

Ibe, Nakano, Shoji and Suzuki, JHEP 03, 194 (2018) [arXiv:1707.07258 [hep-ph]].



A value of P =50% are needed to obtain agreement of the data, being a factor of approximately seven above the integrated ionization probabilities calculated for Migdal shakeoff from atomic germanium in 1707.07258

Collar, Kavner, Lewis, PRD 103, 122003 (2021) [arXiv: 2102.10089]