

# Measurement of the W boson mass with the D0 detector

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# Motivation

- ❖ The Standard Model (SM) predicts a relationship between the W boson mass and other parameters of electroweak theory:

$$M_W = \sqrt{\frac{\pi\alpha}{\sqrt{2}G_F} \frac{1}{\sin\theta_W \sqrt{1-\Delta r}}}$$

- ❖ Contributions to  $M_W$  through radiative corrections  $\Delta r$ .

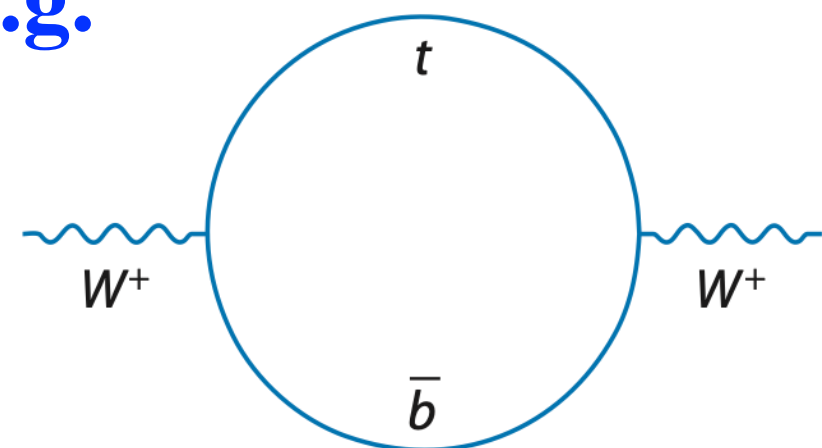
- ❖ Precisely test the electroweak theory at the loop level.

- ❖ In case of SM, the precise W mass and top mass measurements can predict the SM Higgs boson mass.
- ❖ By comparing the prediction and direct Higgs mass measurement, we can know how good is the SM prediction. If disagreement is big, we can infer contributions from theories beyond SM, such as SUSY.

**W mass related to Top quark mass:**

$$\Delta r \propto M_t^2$$

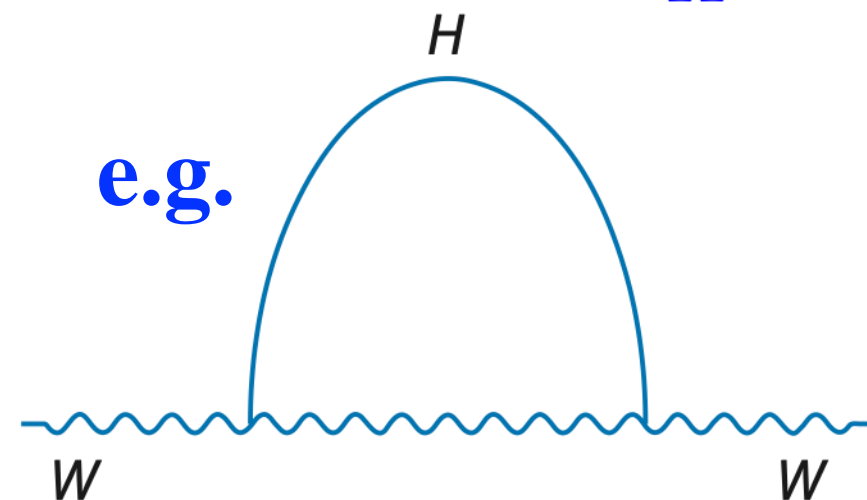
e.g.



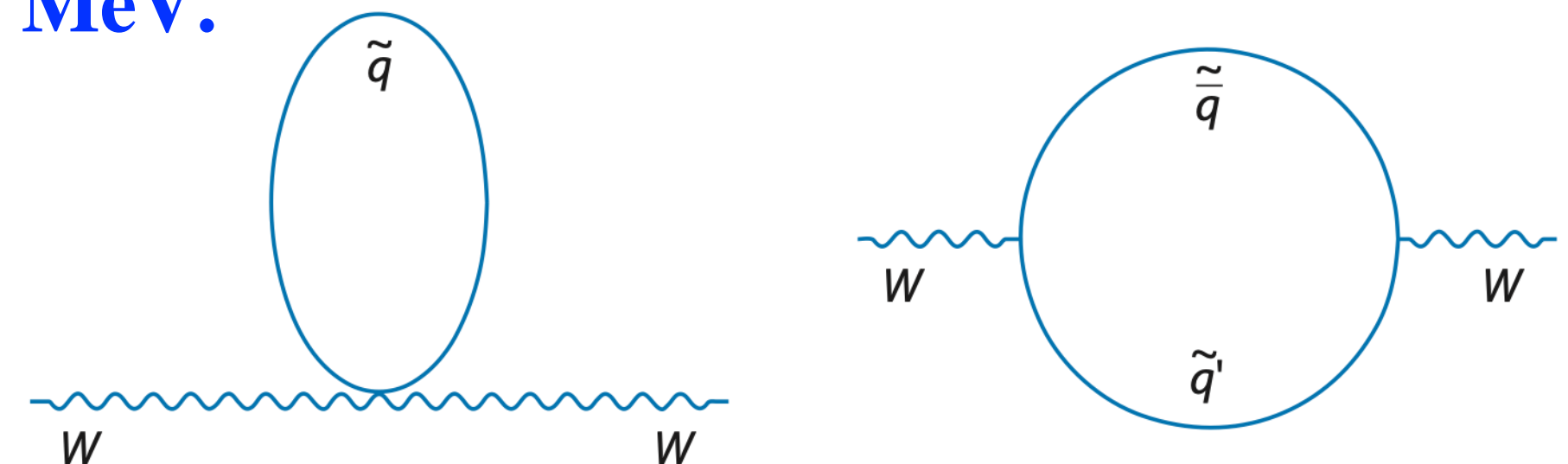
**W mass related to SM Higgs mass:**

$$\Delta r \propto \ln M_H$$

e.g.



**Beyond SM, contribution from SUSY particles can induce a total radiative correction to  $M_W$  of 100 to 200 MeV.**



# Compare predicted and measured

## The old state of the art in 2012

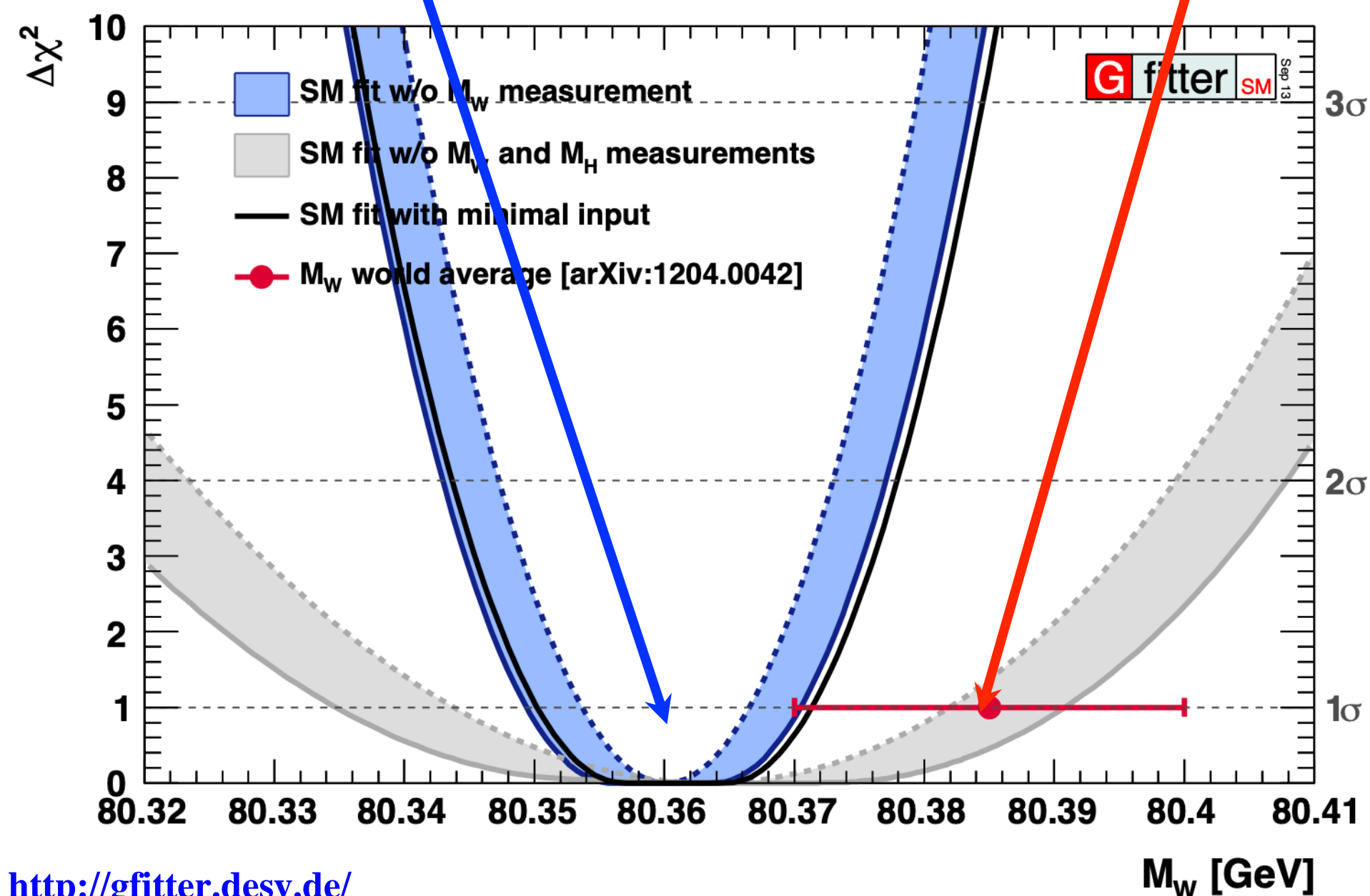
$M_W = 80359 \text{ MeV} \pm 11 \text{ MeV}$

*Comparing with the world average directly measured value in 2012:*

$M_W = 80385 \text{ MeV} \pm 15 \text{ MeV}$

Predicted

Measured

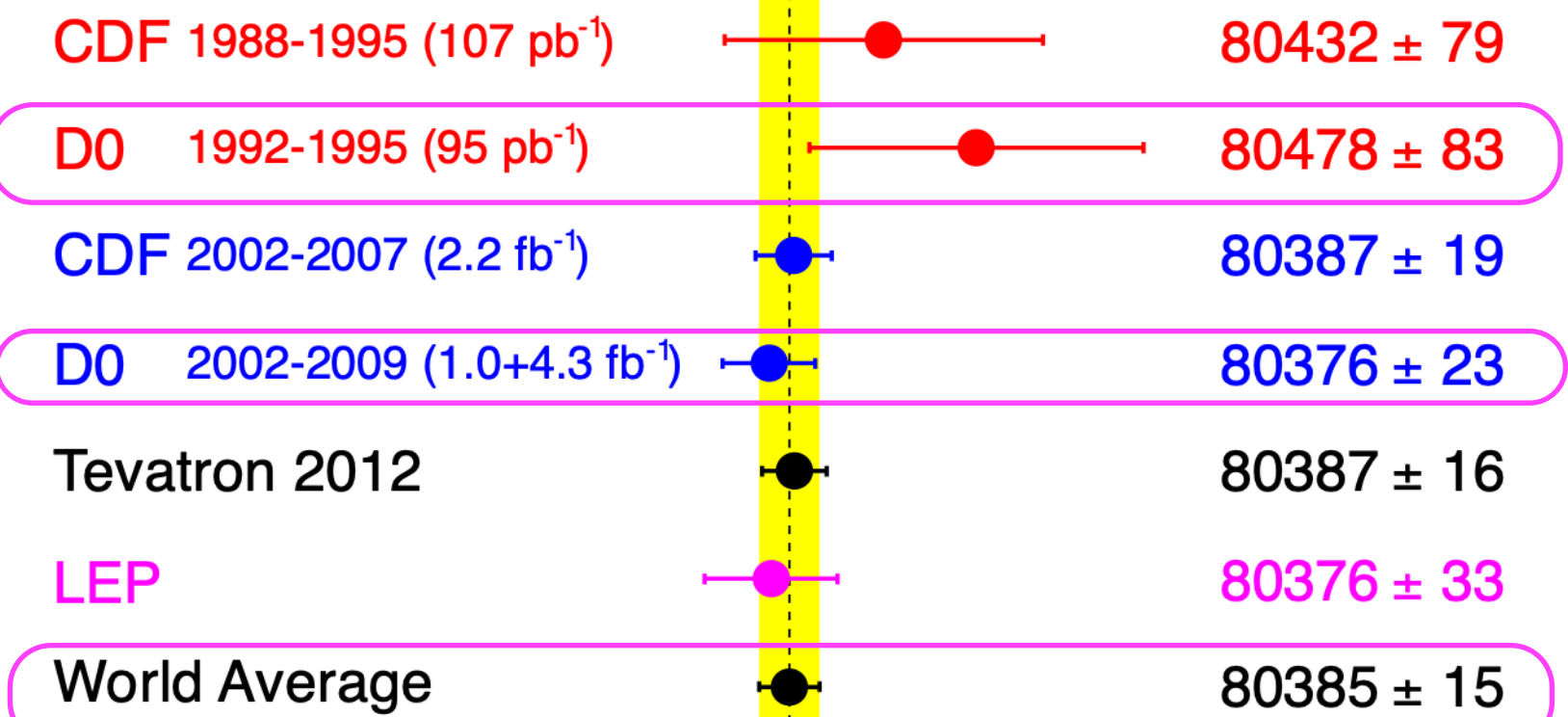


A **~1.3 sigma** difference between the two  $M_W$  central values.

Mass of the W Boson

Measurement

$M_W$  [MeV]

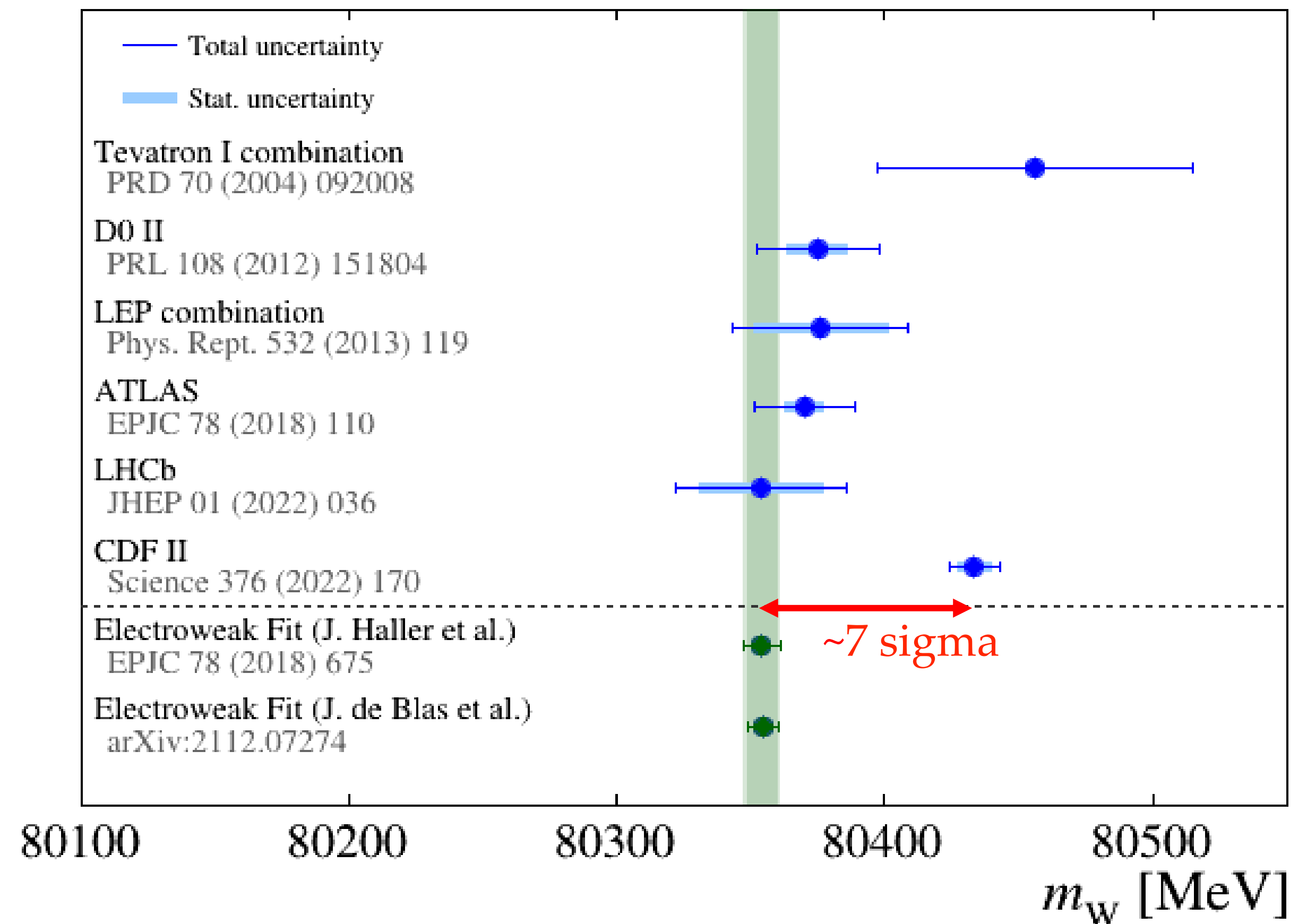
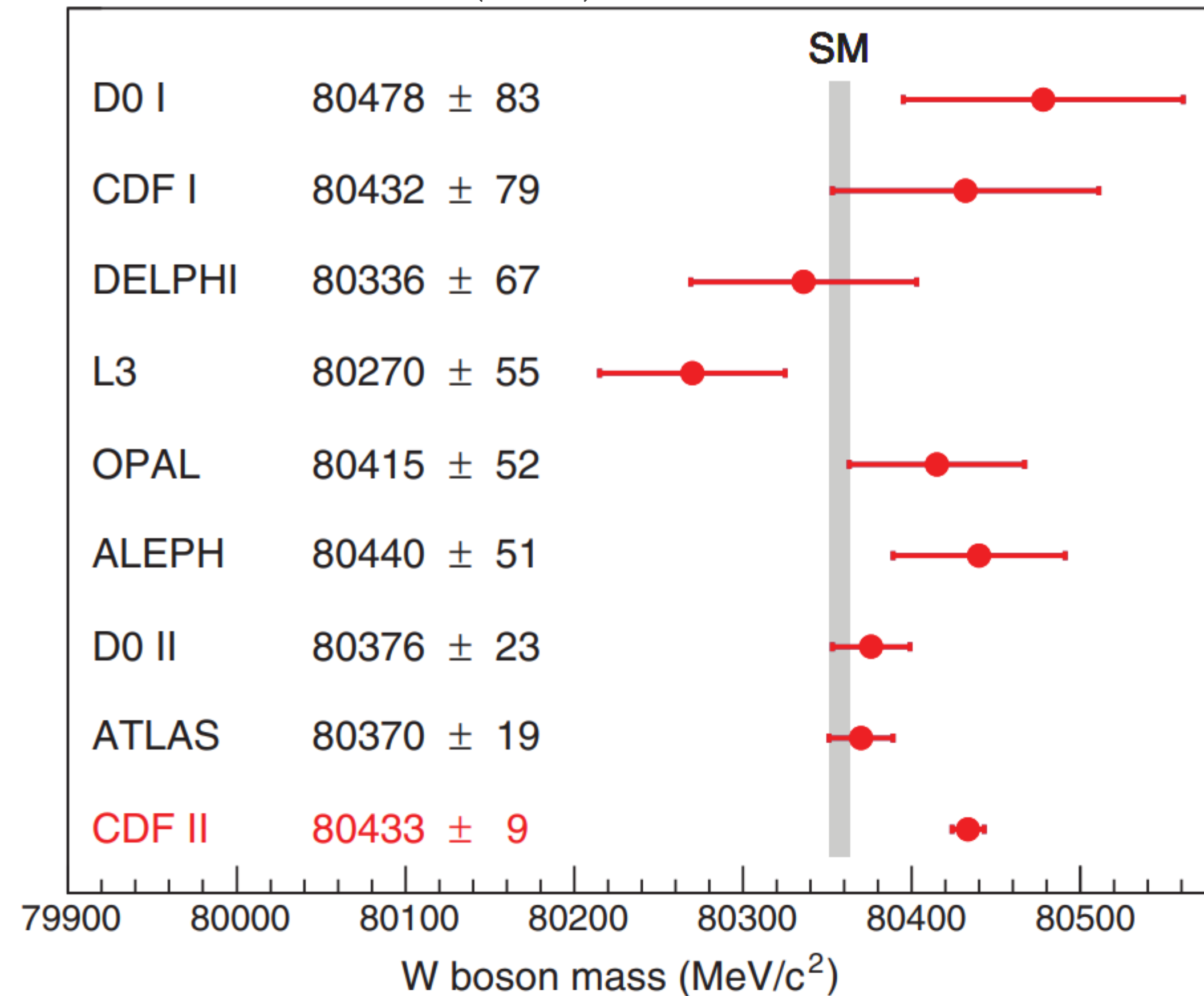


$M_W$  [MeV]

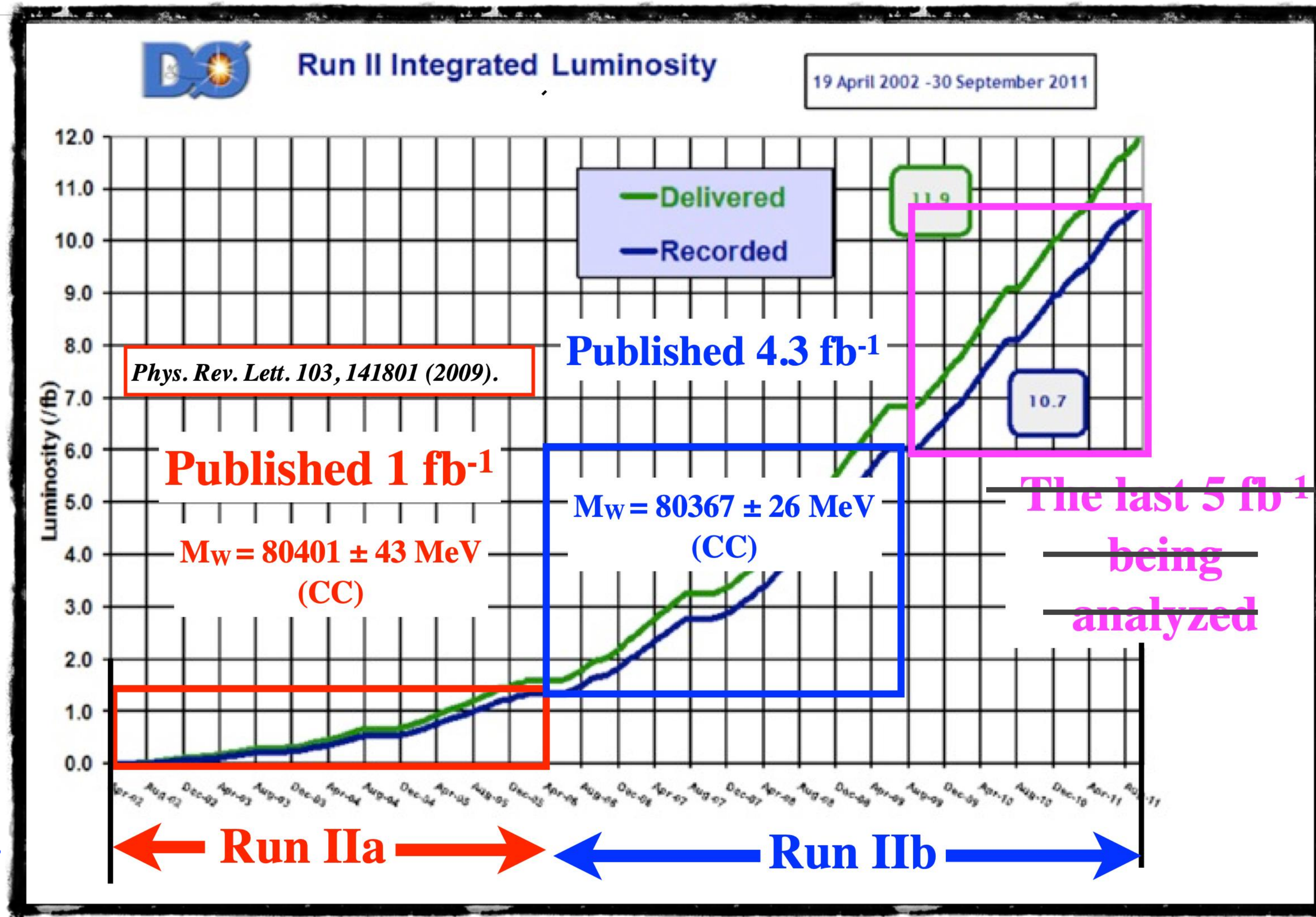
# Compare predicted and measured

## Current state of the art

Science 376, 170–176 (2022)



# The D0 data sets

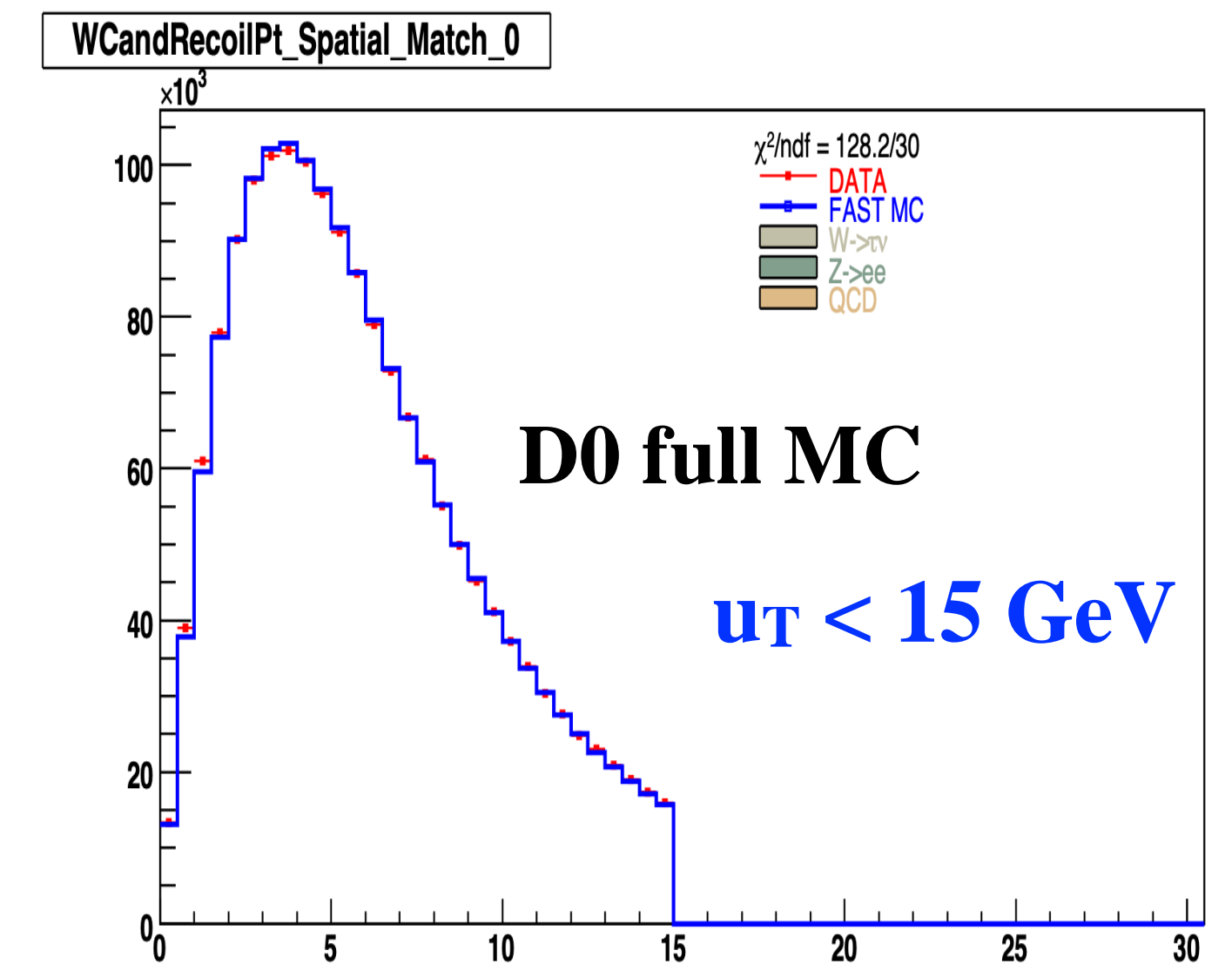
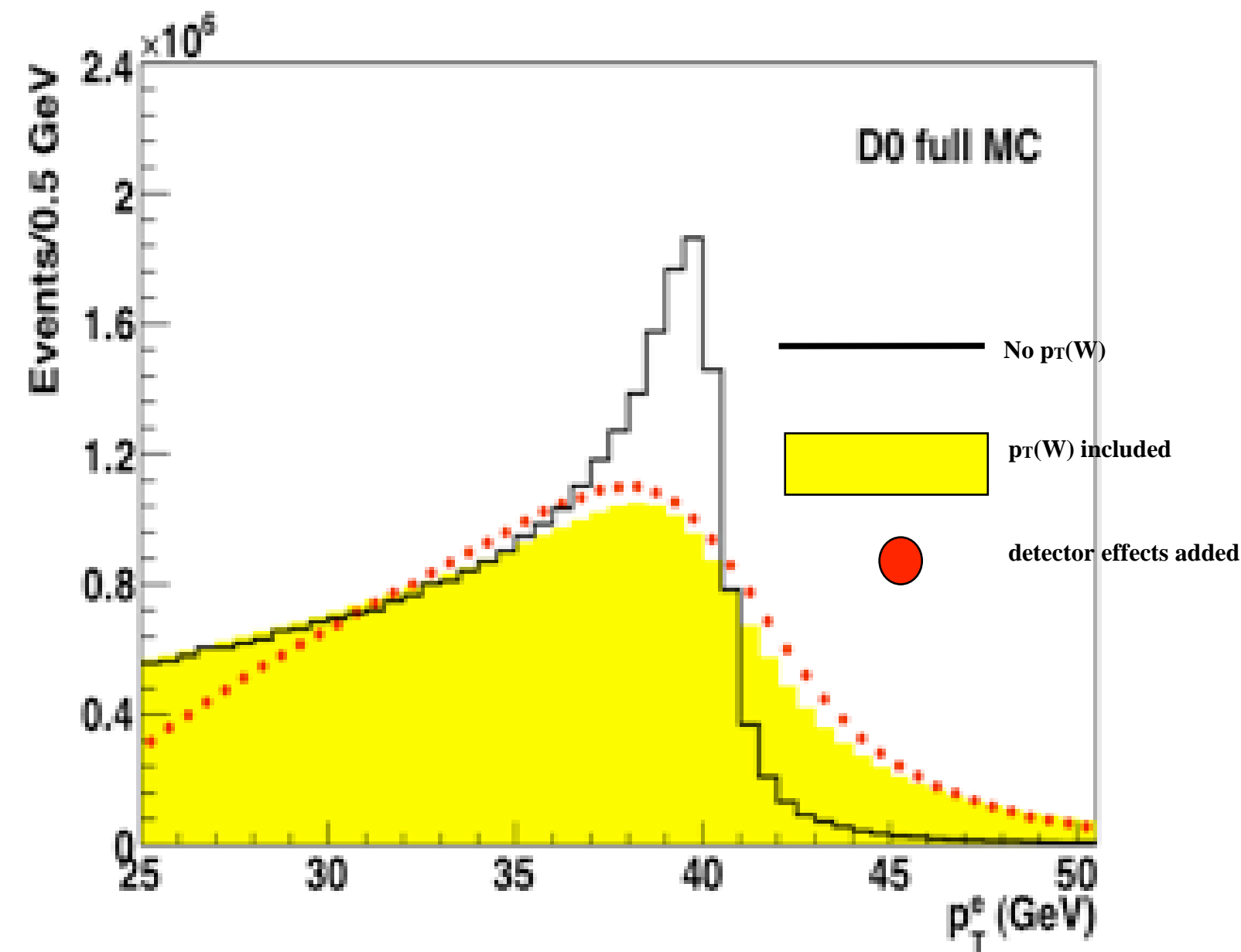


CC:  
Central  
Calorimeter

# Event reconstruction

## Event Selection:

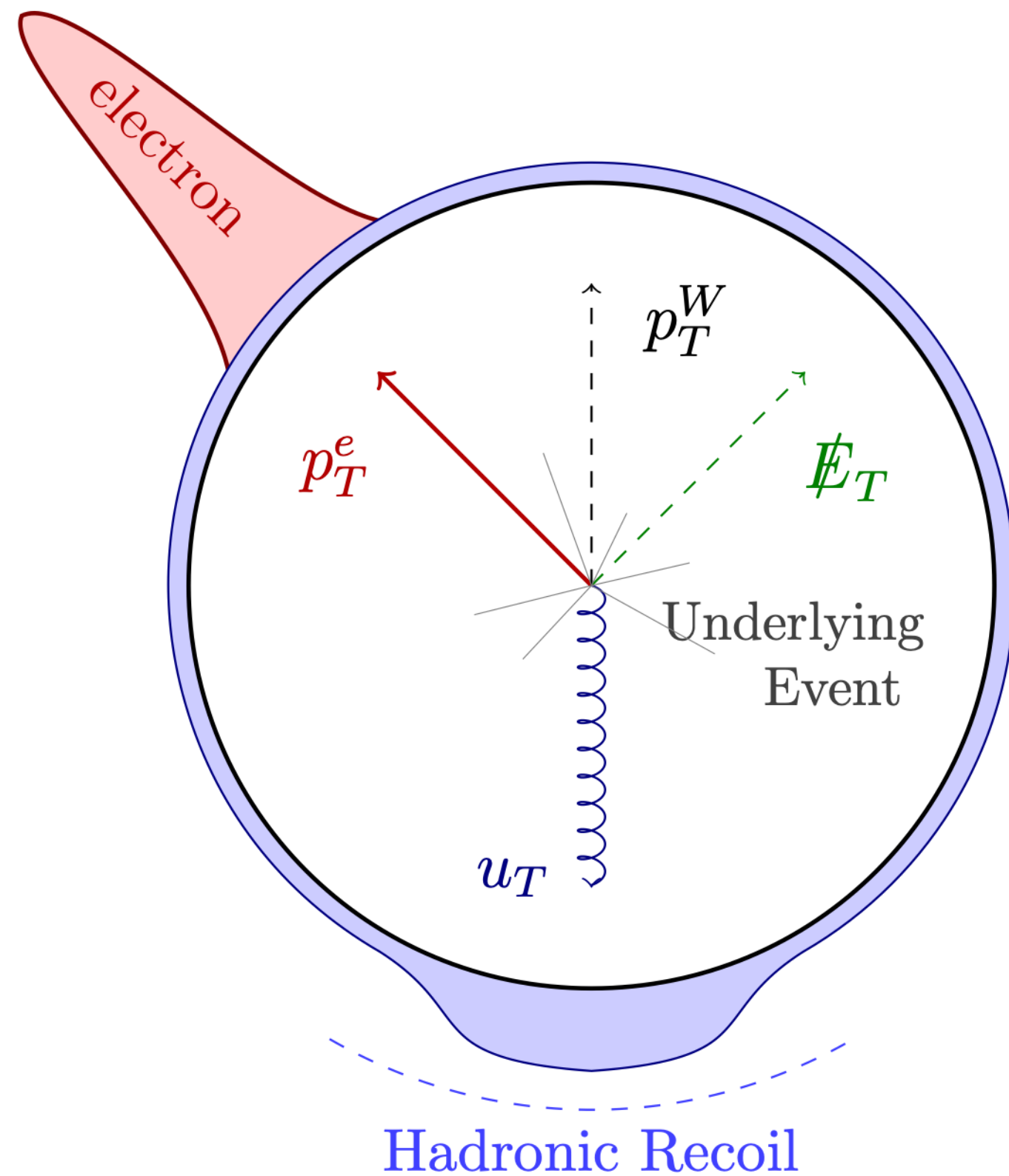
- $W \rightarrow e\nu$  events
- Central electrons:  $|\eta| < 1.05$
- $p_T(e) > 25$  GeV
- Missing  $E_T > 25$  GeV
- Hadronic recoil:  $u_T < 15$  GeV
- After selection:
- 1,677,394  $W \rightarrow e\nu$  candidates



Transverse boost of the W boson degrades the sharpness of the  $p_T(e)$  Jacobian edge. Requiring  $u_T < 15$  GeV is helpful, however, it also transfers certain recoil modeling uncertainty into  $p_T(e)$

# Analysis strategy

## A Typical $W \rightarrow ev$ Event in DØ Detector



## Reconstruct three observables:

$$M_T, P_T^l, \cancel{E}_T$$
$$M_T^W = \sqrt{2P_T^l \cancel{E}_T (1 - \cos \Delta\phi)}$$

using CC electrons with  $p_T > 25 \text{ GeV}$

## Using $Z \rightarrow ee$ events for detector calibration

**A Fast MC model to generate templates of the 3 observables with different W mass hypotheses. Fit the templates to the data to extract W mass.**

## The Fast MC model:

- Event Generator: Resbos(CTEQ6.6)+PHOTOS
- Parameterized Detector Model

# The Observables

Can directly reconstruct two variables:

Lepton  $p_T$  can be precisely measured, 0.01% precision.

$$\vec{P}_T^l$$

$$\vec{u}_T$$

Hadronic recoil: vectorial sum of the transverse energies of all the calorimeter cells outside the lepton reconstruction window.

- less precise, ~1% precision,
- low resolution,  $\Delta u_T > 3.5 \text{ GeV}$
- hadronic energy response is only ~ 65%

Calculate three observables to extract the W boson mass :

$$P_T^l$$

$$\cancel{E}_T$$

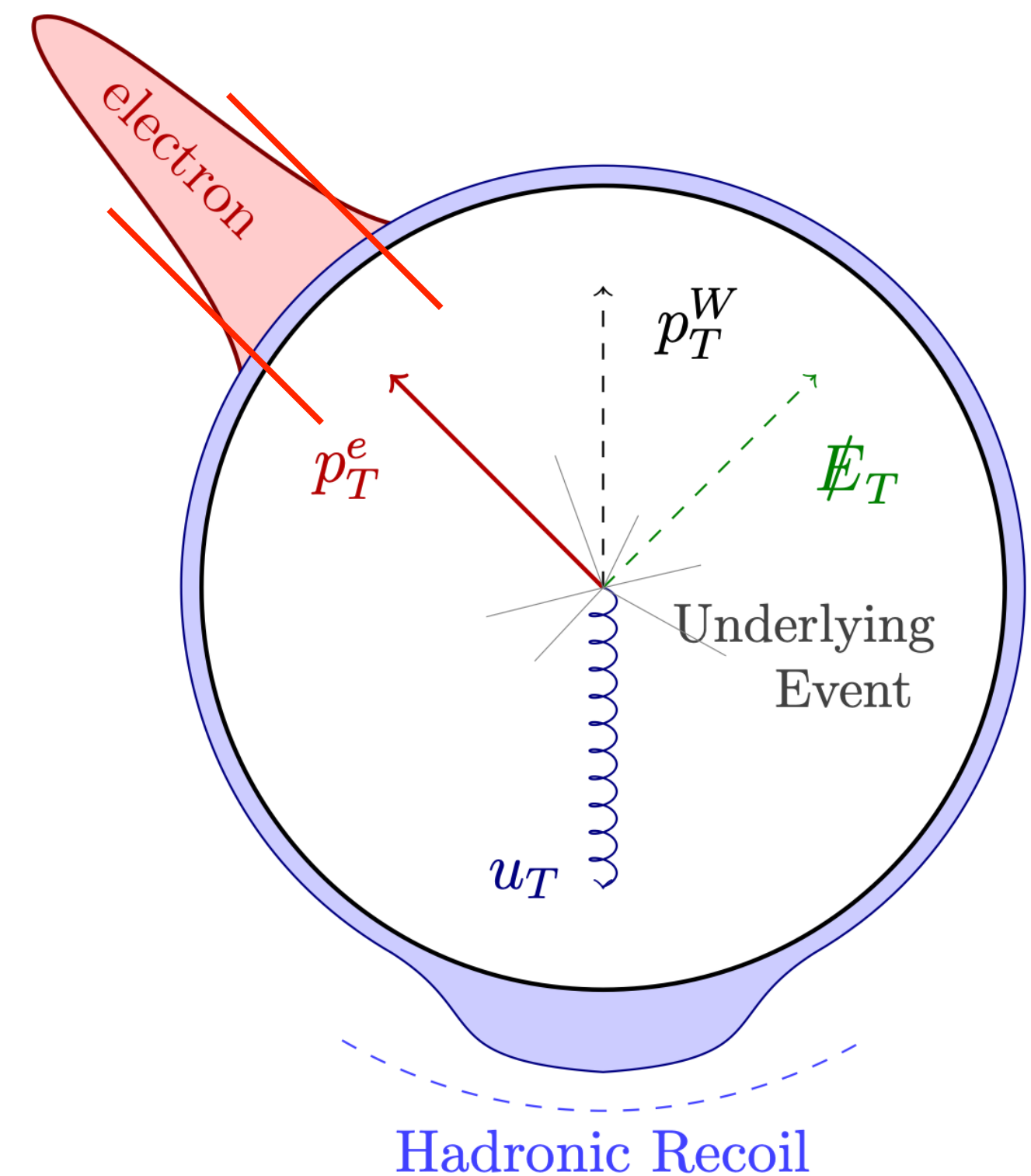
$$M_T$$

$$= |\vec{P}_T^l + \vec{u}_T|$$

$$\sim P_T^l + u_T \cos \Delta\phi_{(\text{lepton-recoil})}$$

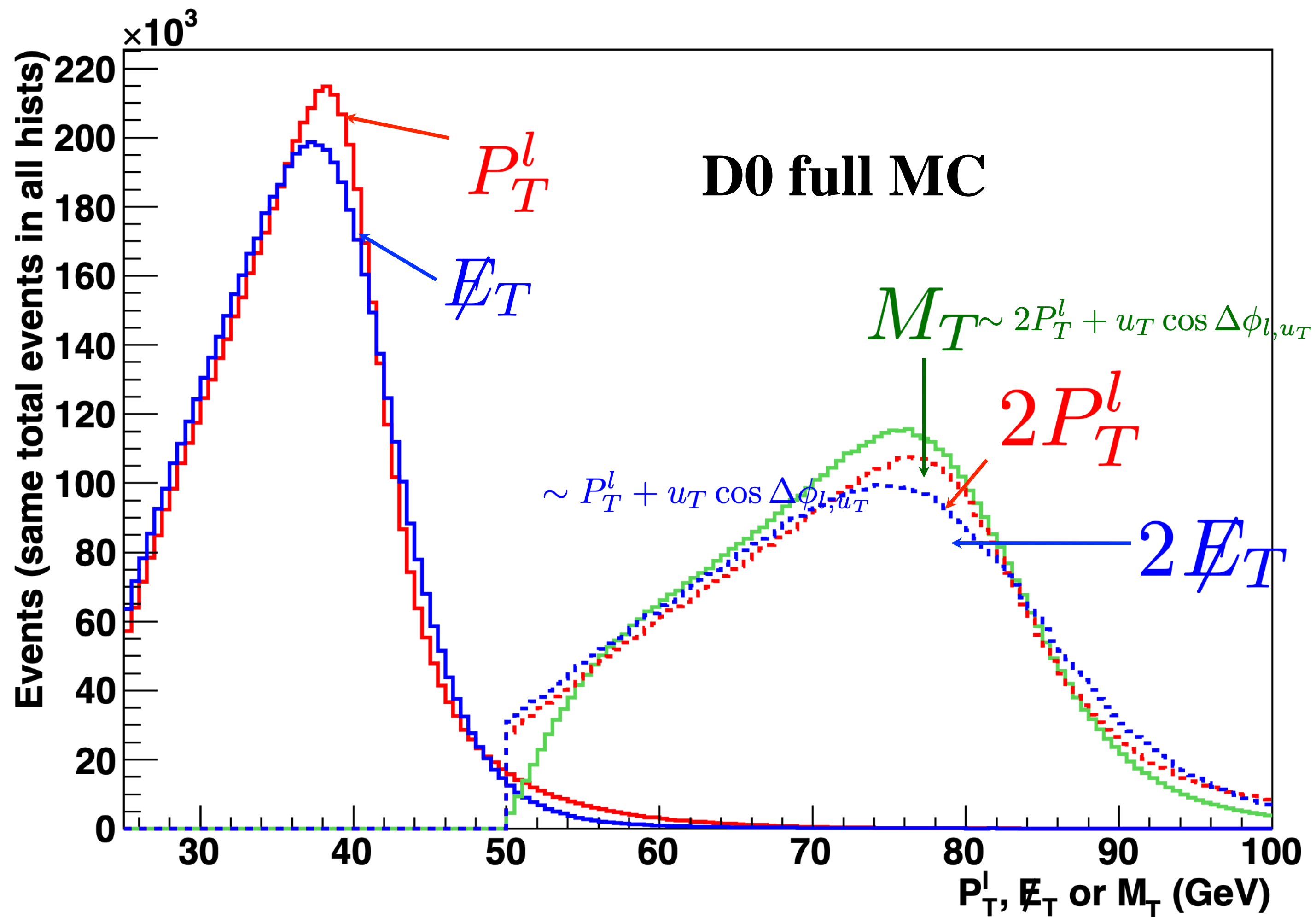
$$= \sqrt{(P_T^l + \cancel{E}_T)^2 - (\vec{P}_T^l + \vec{\cancel{E}}_T)^2}$$

$$\sim 2P_T^l + u_T \cos \Delta\phi_{\text{lepton-recoil}}$$



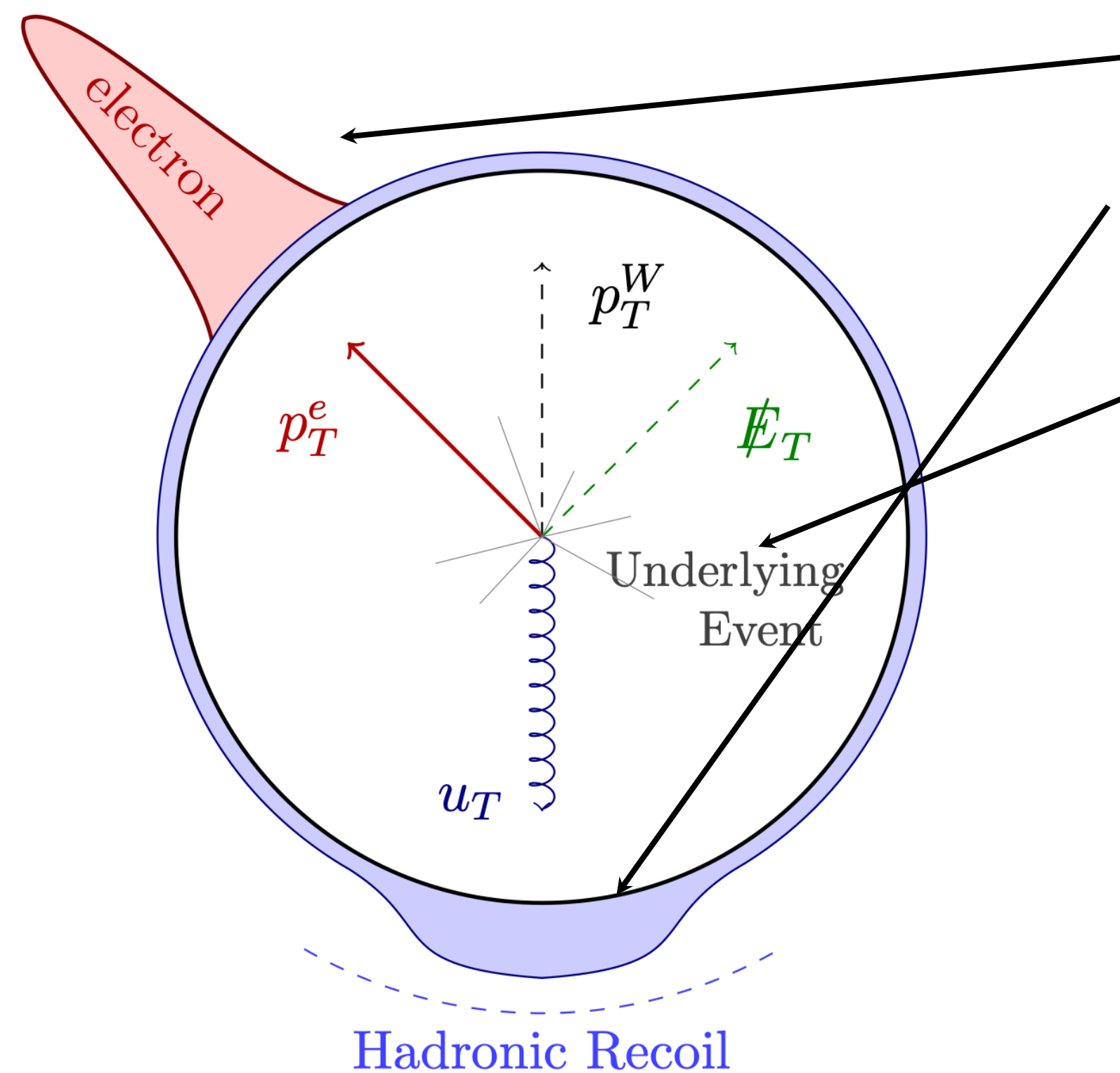


# Which observable is more powerful?



# Parametrized detector model

The parametrized detector model (PMCS) has to simulate:



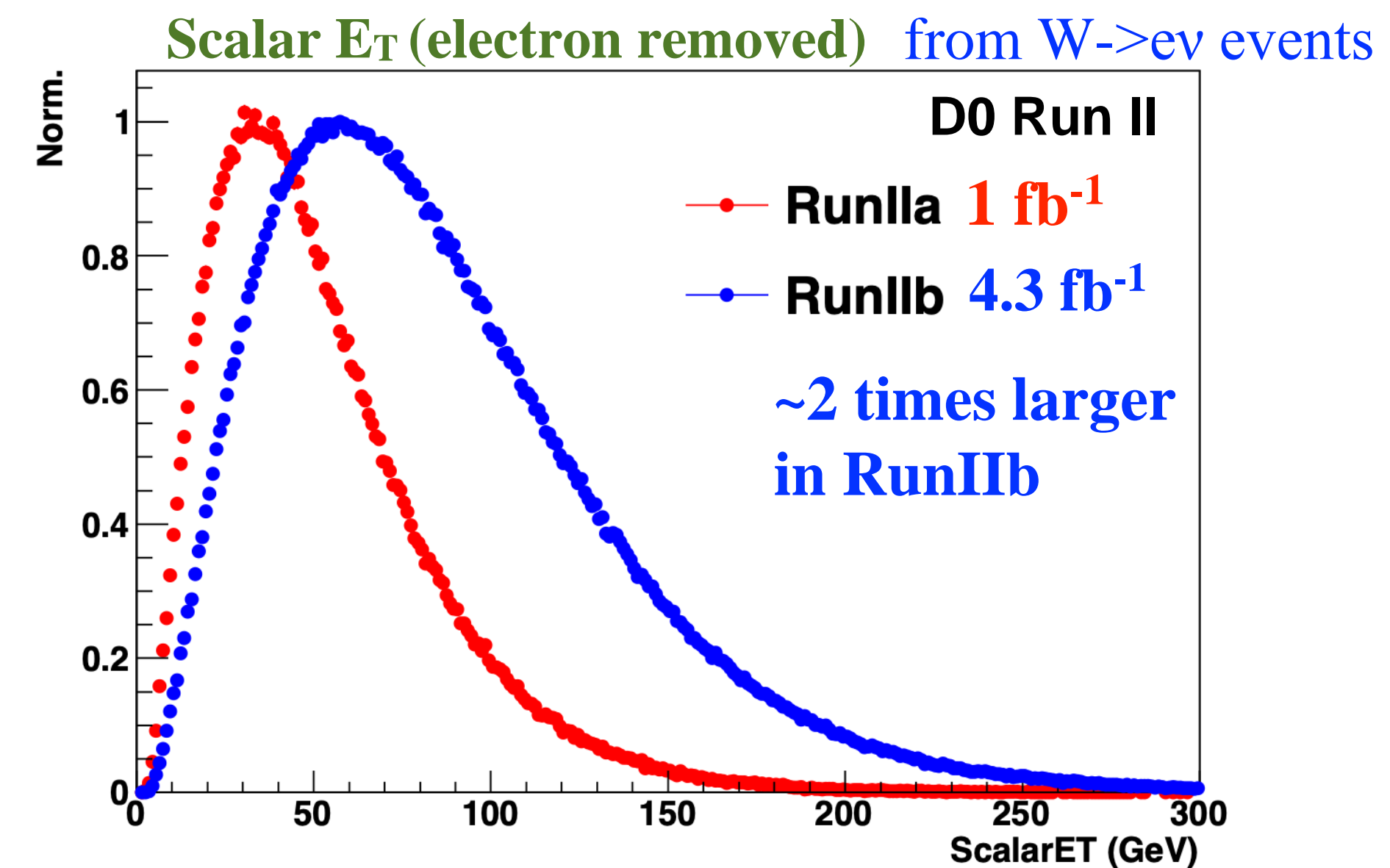
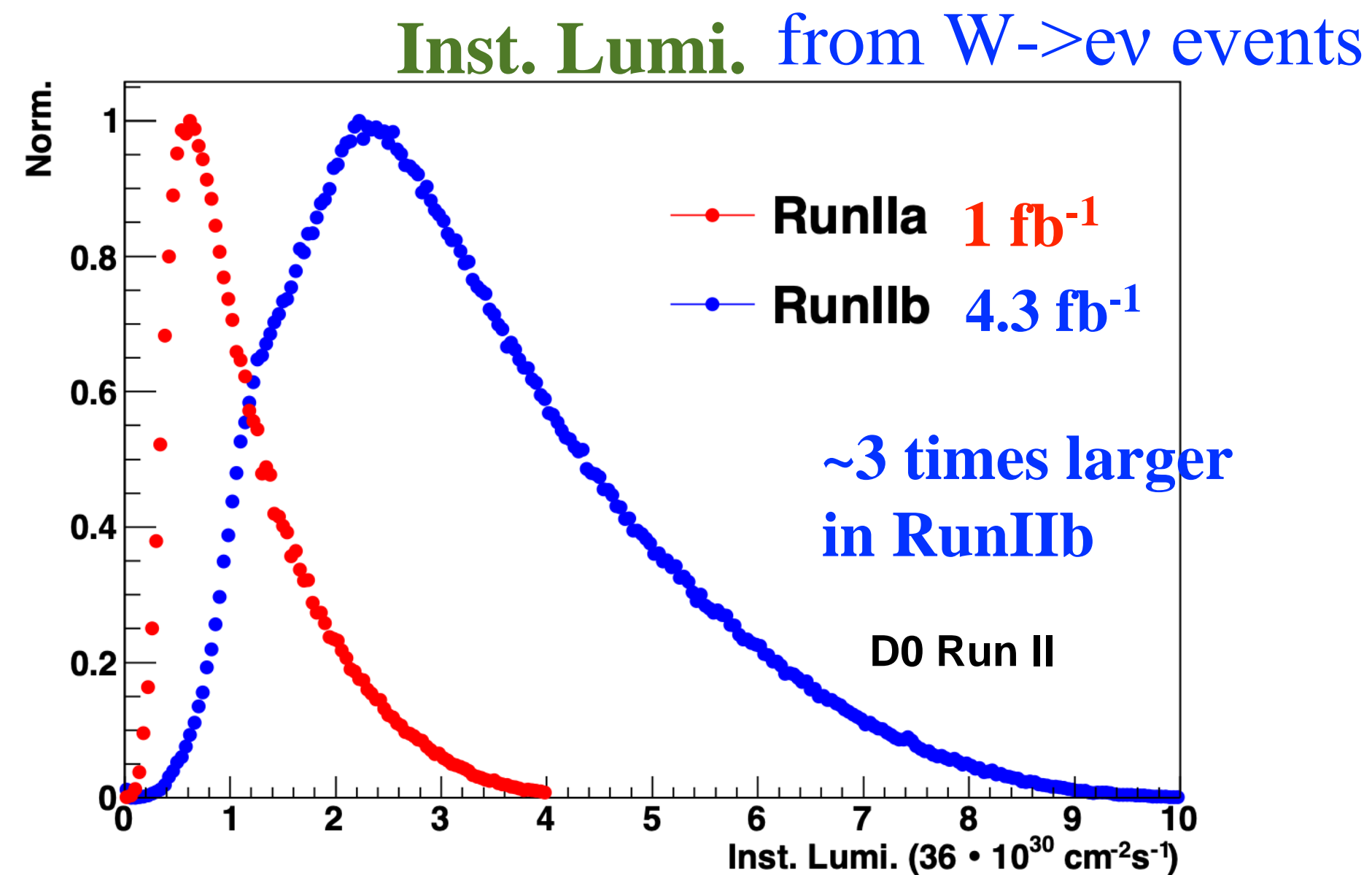
- Electron energy response and smearing
- Hadronic recoil energy response and smearing
- Underlying energy:
  - additional ppbar interactions (pileup):
    - average number of primary vertices:  $>4$
  - spectator parton interactions
- Event selection efficiency
- Background

# Electron energy model

## Correct/model non-linear energy responses:

- Correction of the energy loss due to dead material,
- Correction of the response decrease due to pileup
- Modeling underlying energy contamination from pileup and hadronic recoil

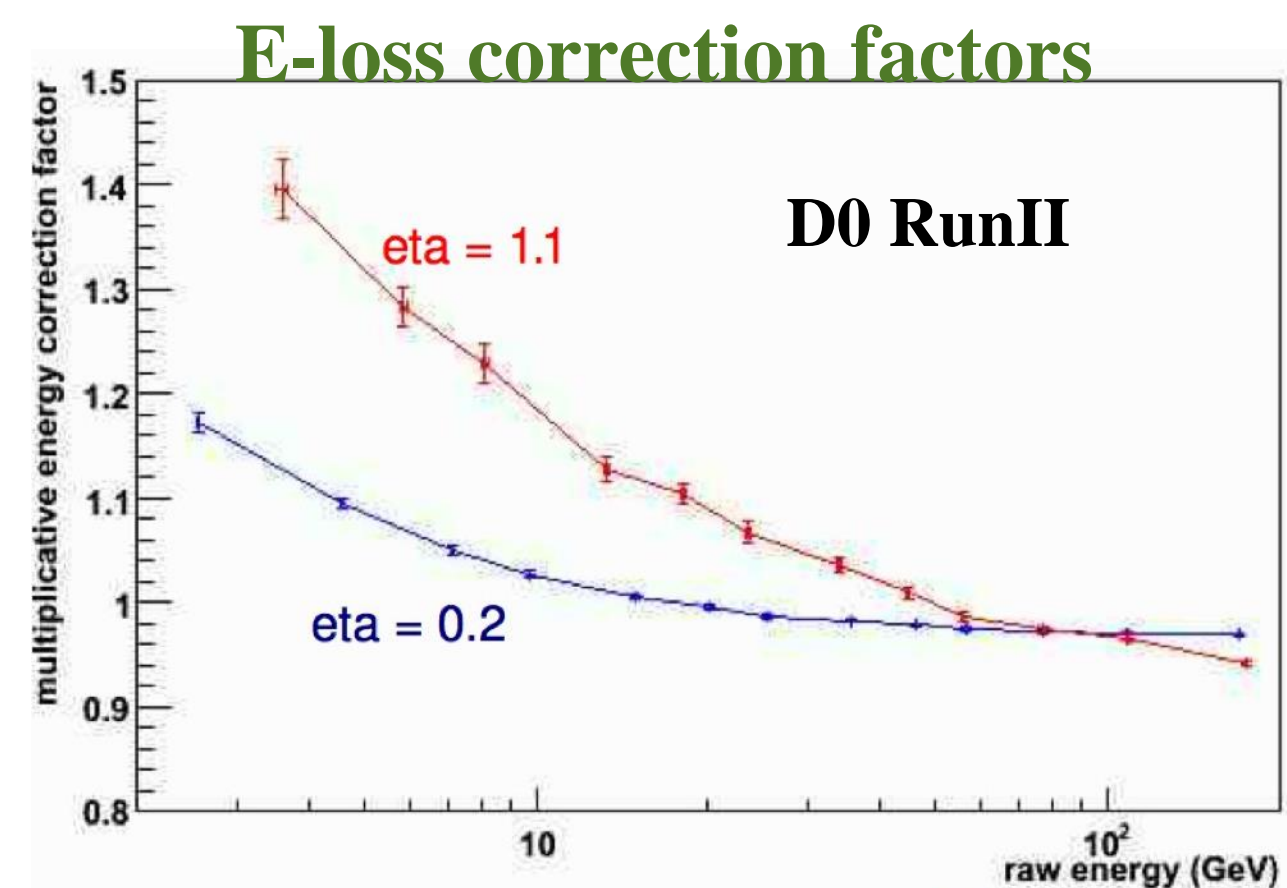
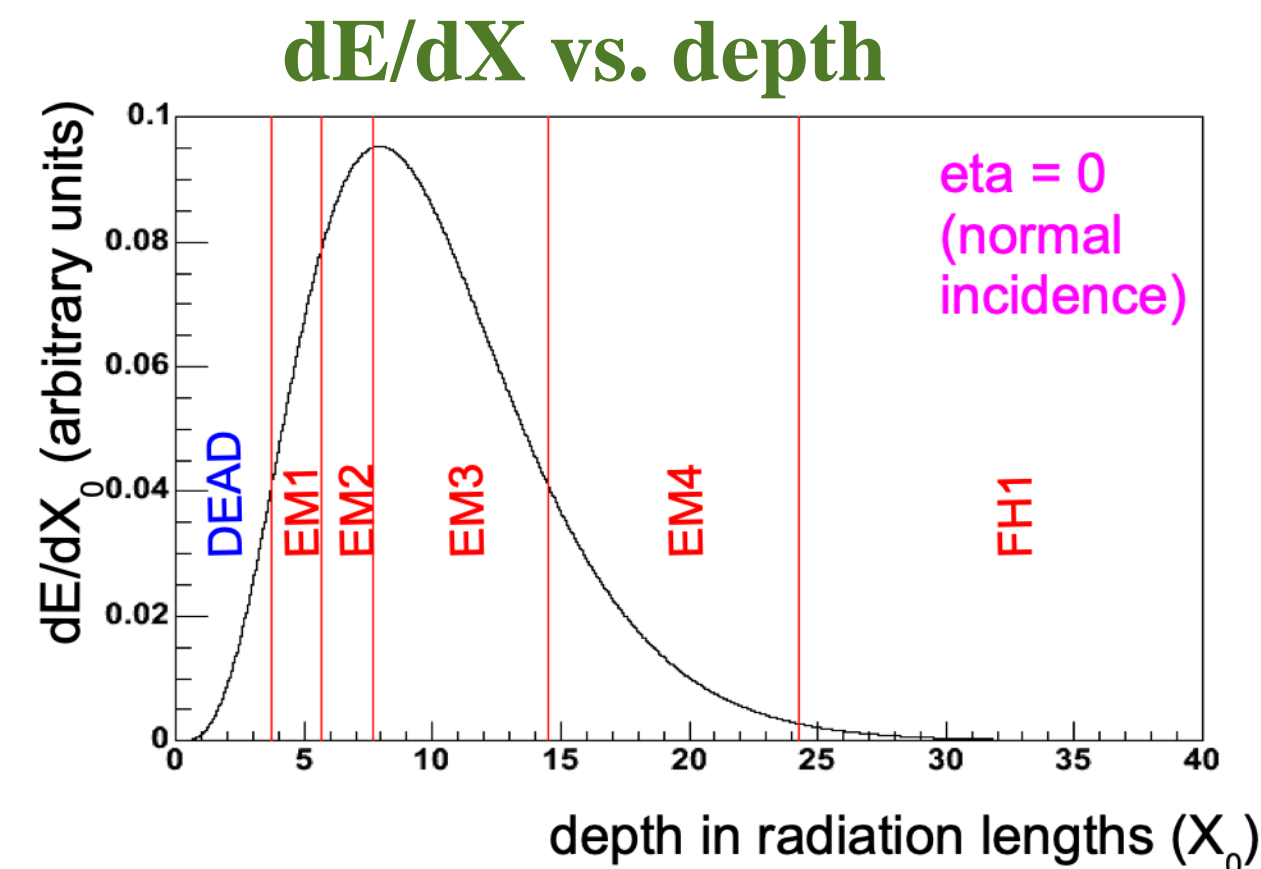
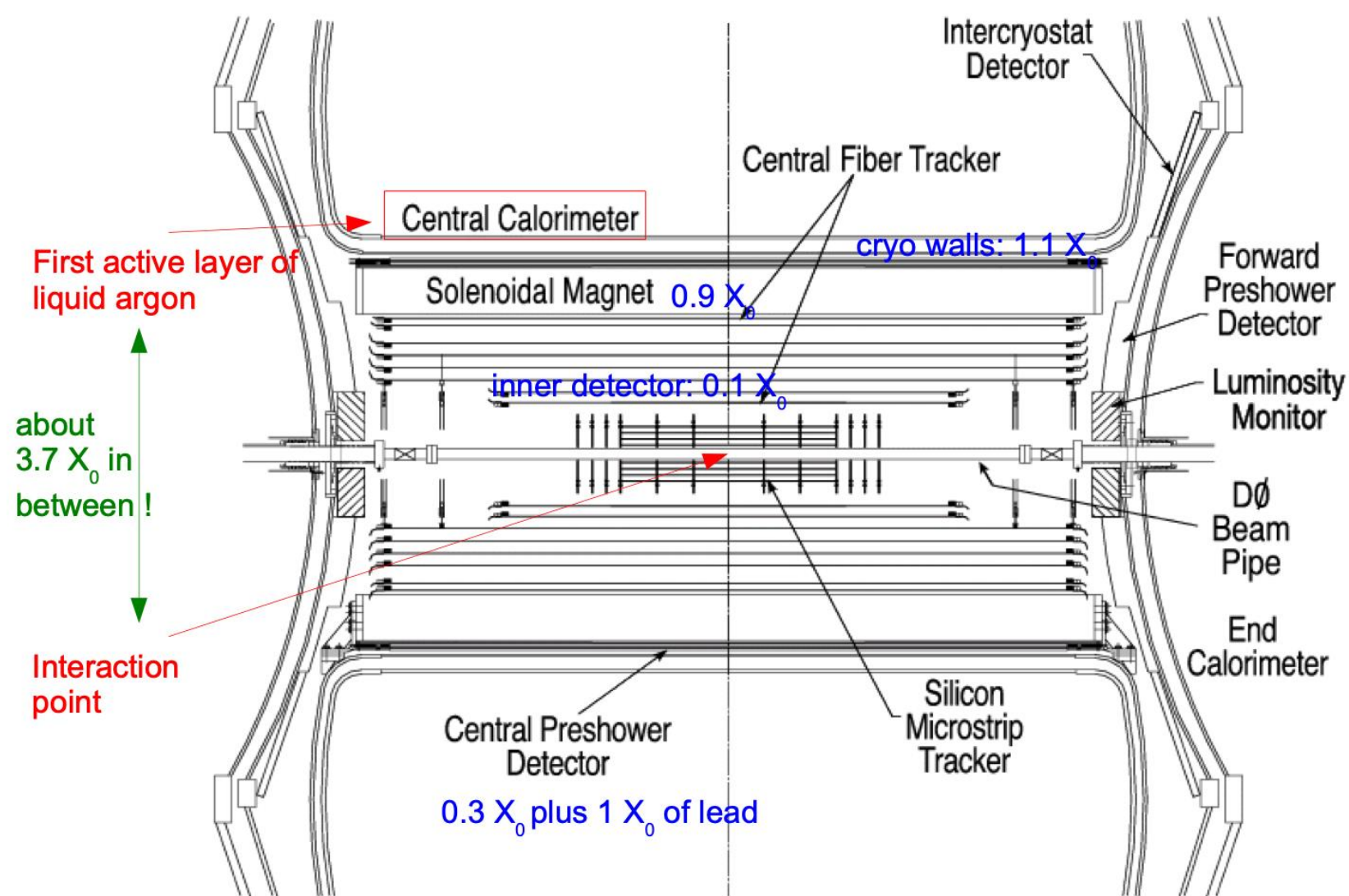
Final electron energy response is tuned using  $Z \rightarrow ee$  events assuming a linear response.



# Dead material, electron energy loss

## About $3.7 X_0$ dead material in front of EM calorimeter

- Electrons start to lose energy before flying to the EM calorimeter
- Depends on electron energy and incidence angle ( $\eta$ )

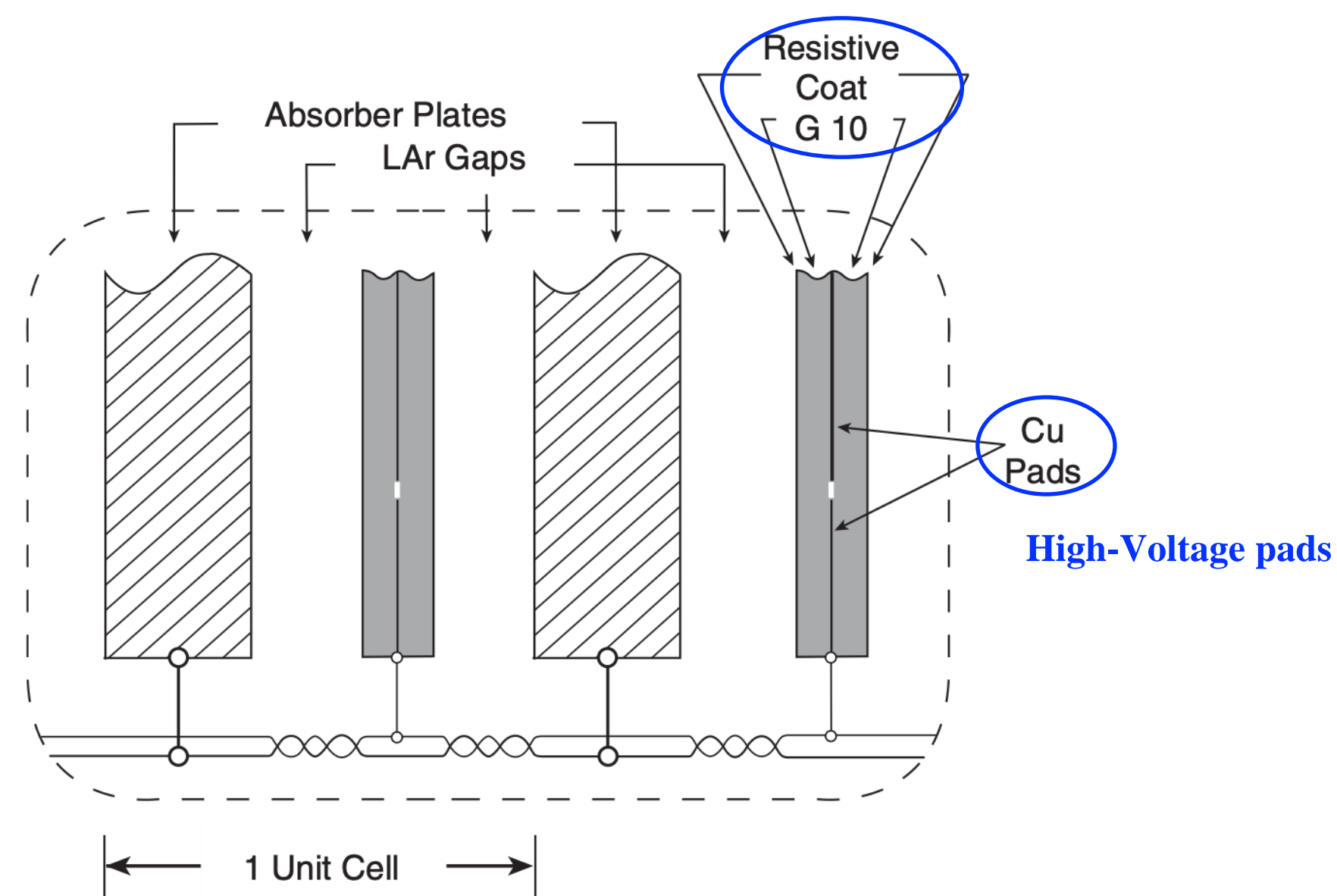


# Response reduction due to pileup

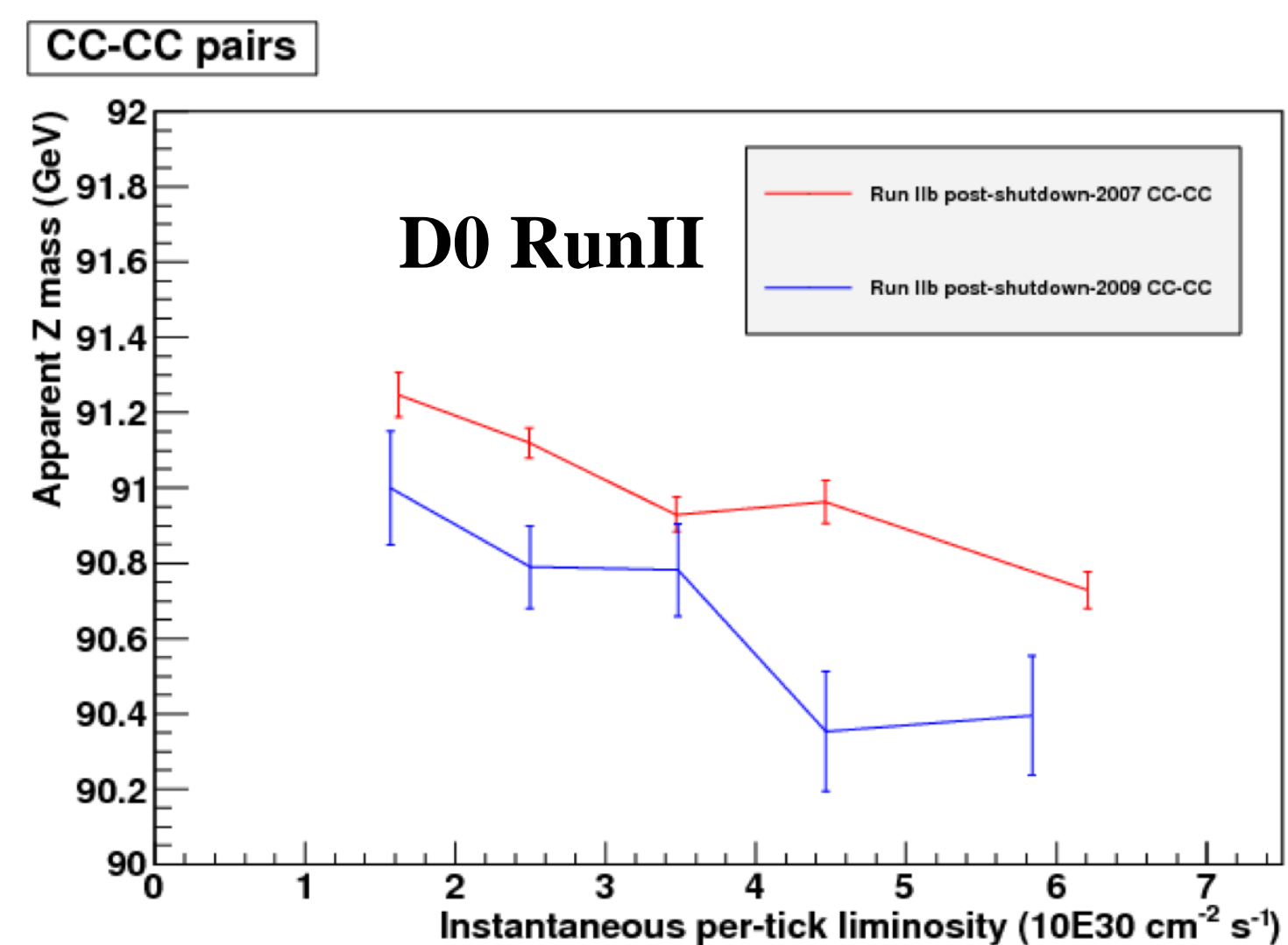
## Pileup causes reduction of energy response!

- Too much pileup creates high current in the readout
- The current that flows through resistive coat of the HV pads results in HV drops, thus, reduces the energy response

### Unit cell of the LAr calorimeter readout



### Fitted Z mass vs. Inst. Lumi.



# Electron energy model

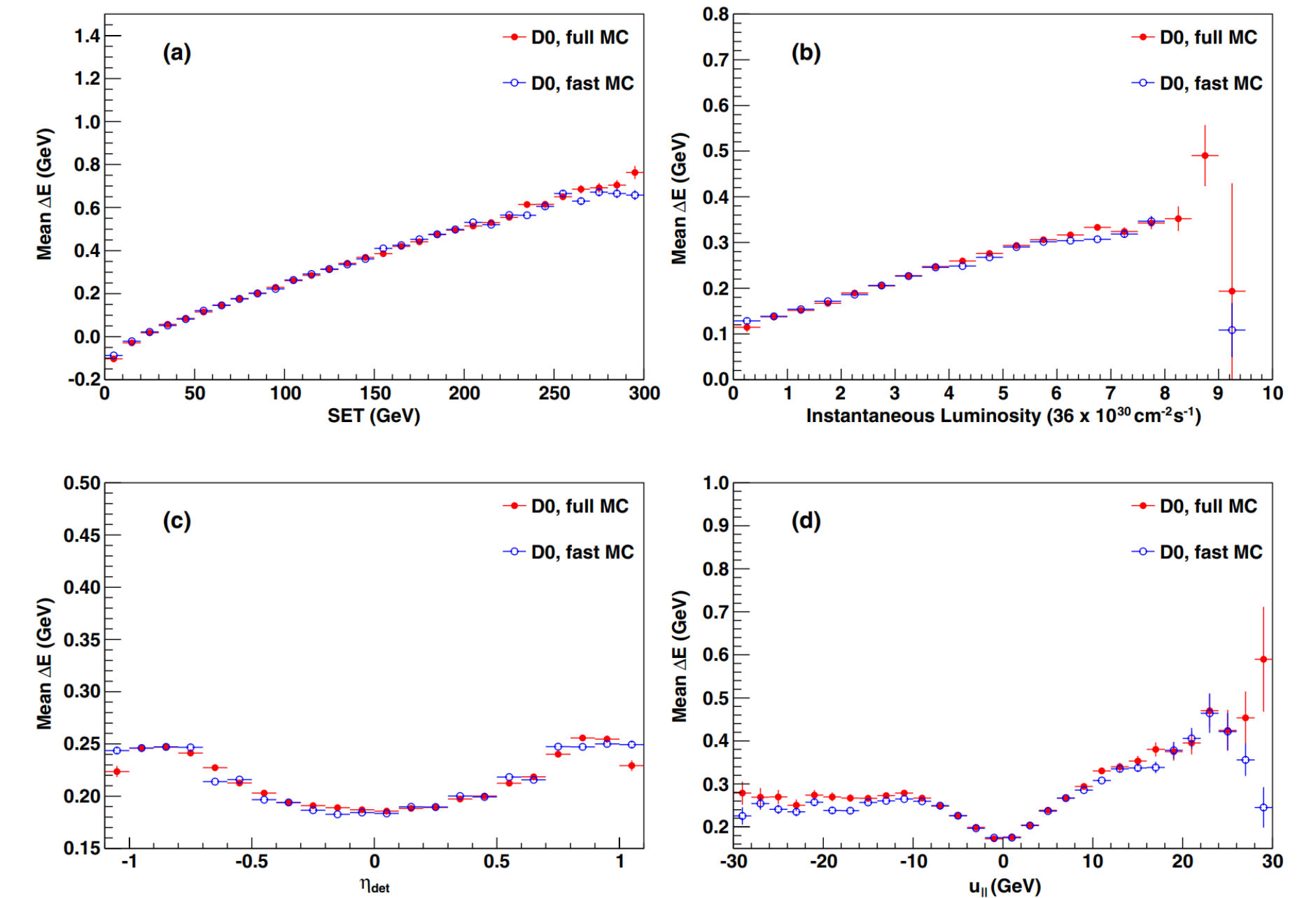
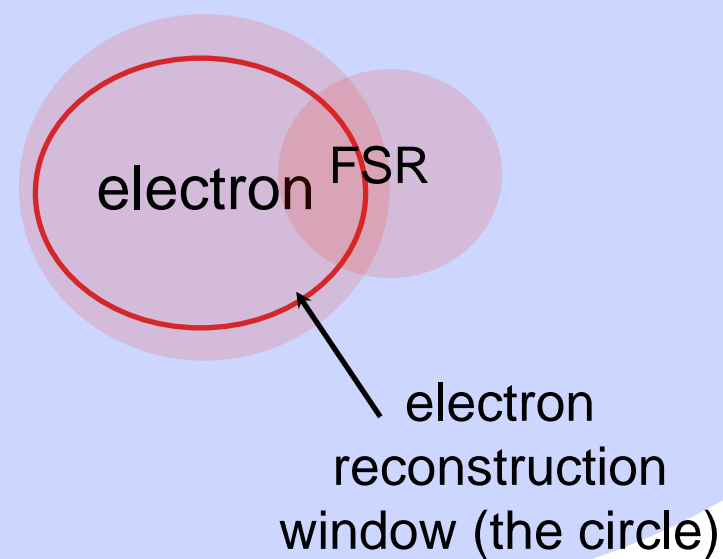
**Electron Model:**  $E_{reco} = R_{EM}(E_{true}) \otimes \sigma_{EM}(E_{true}) + \Delta E_{corr}$

Response
Resolution
Energy correction

## $\Delta E_{corr}$ Model:

1. Energy loss due to FSR
2. Recoil, spectator parton interactions and pileup contamination inside the electron reconstruction cone
3. Effects due to electronics noise subtraction and baseline subtraction (to subtract residue energy deposition from previous bunch crossings)

Hard Recoil, spectator parton interactions, and pile-up



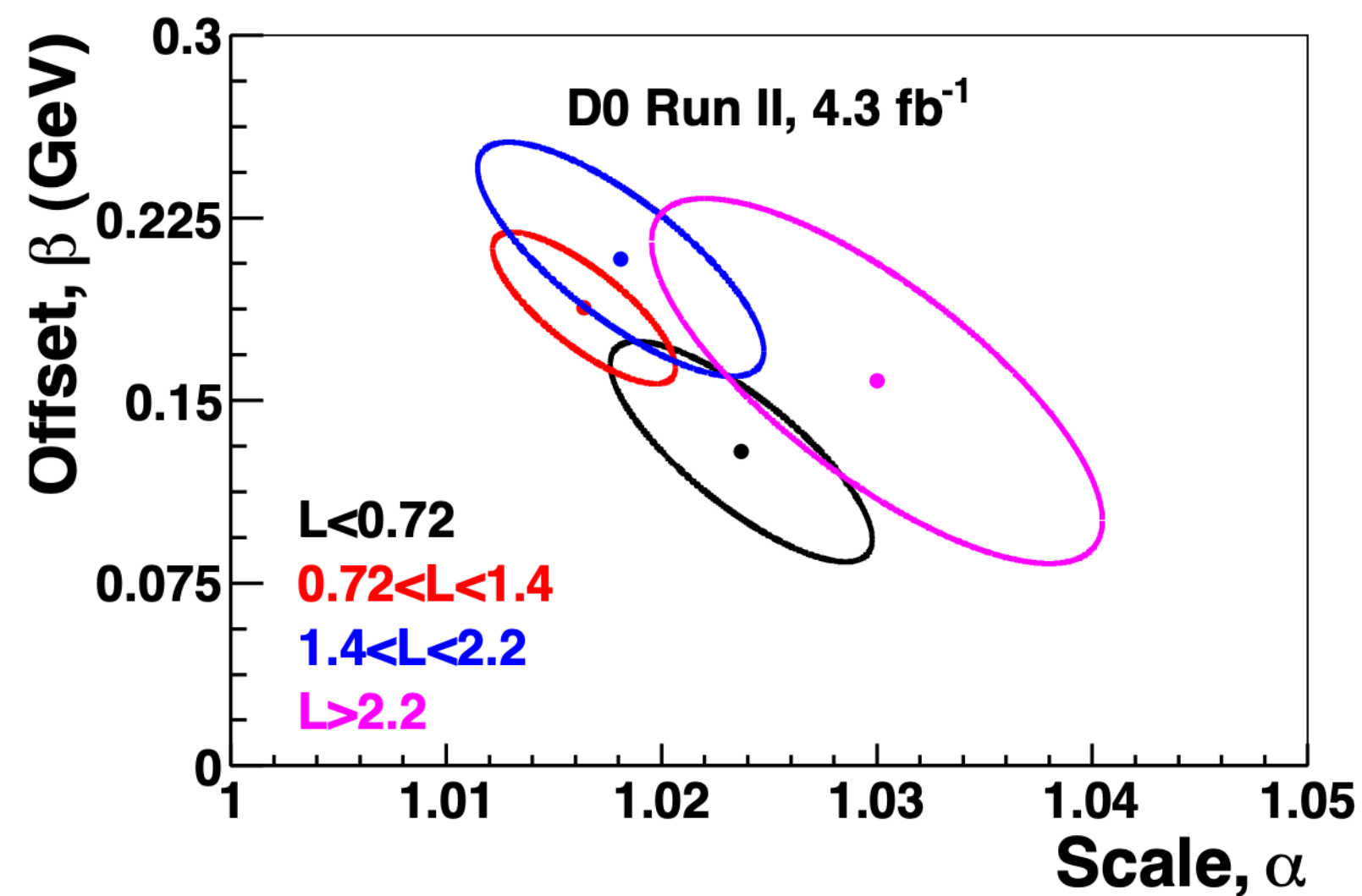
# Final electron energy scale

After the correction and modeling of the non-linear energy responses, the final electron energy response is calibrated using  $Z \rightarrow ee$  events assuming a linear response:

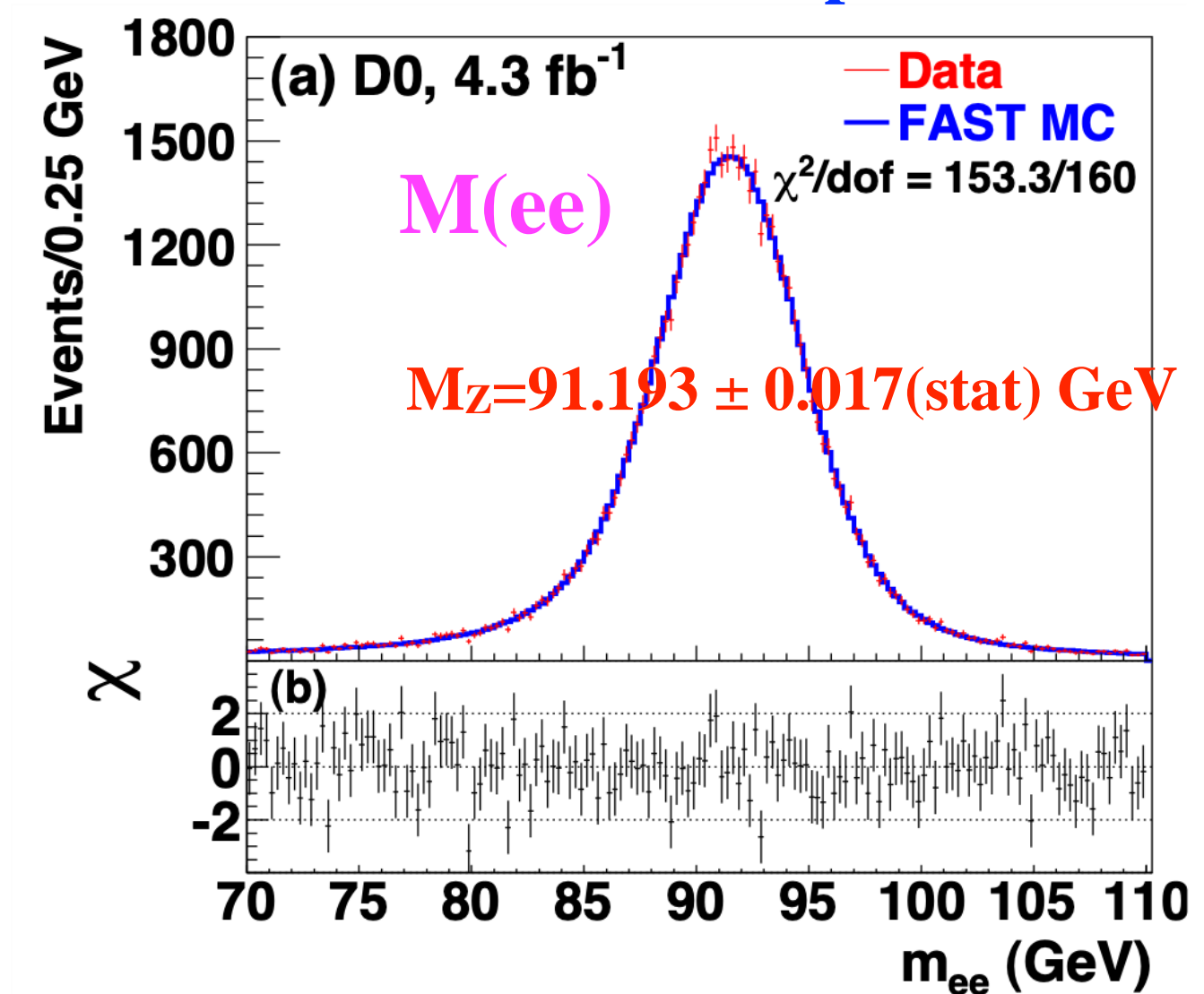
$$R_{EM}(E_{true}) = \alpha \cdot (E_{true} - \bar{E}_{true}) + \beta + \bar{E}_{true}$$

Essentially, measuring the ratio  $M_W/M_Z$ , limited by the  $Z \rightarrow ee$  statistics

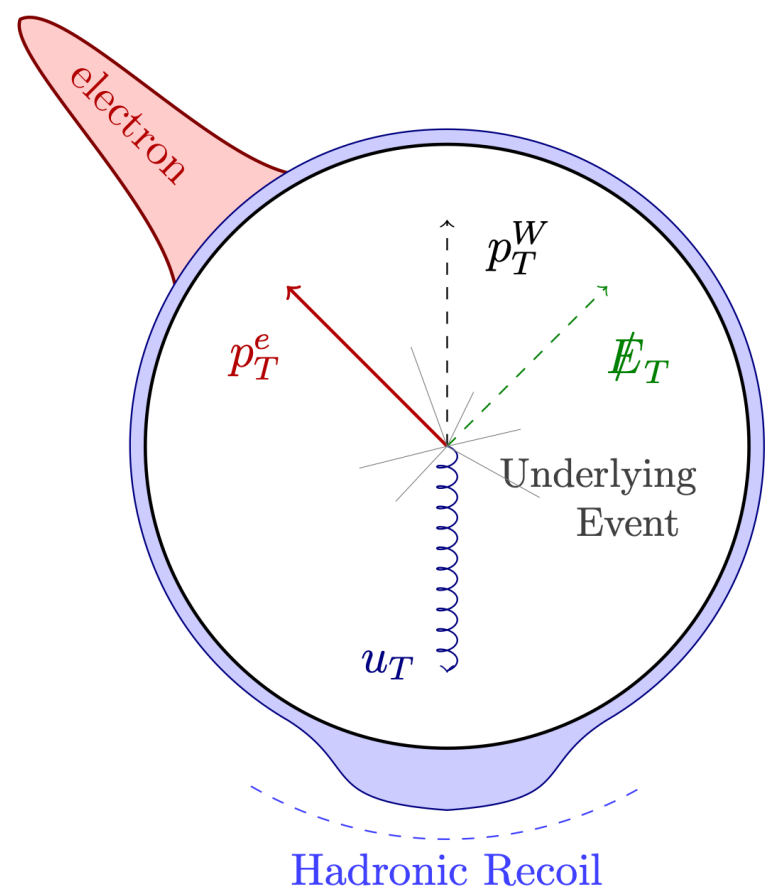
Scale and offset are determined in 4 inst. lumi. bins



Validation: it was calibrated using Z mass, now, we re-determine the Z mass using the calibration, we should find back the Z mass we put in.



# Hadronic recoil model



## Recoil Model:

$$\vec{u}_T = \vec{u}_T^{\text{Hard}} + \vec{u}_T^{\text{Soft}} + \vec{u}_T^{\text{Elec}} + \vec{u}_T^{\text{FSR}}$$

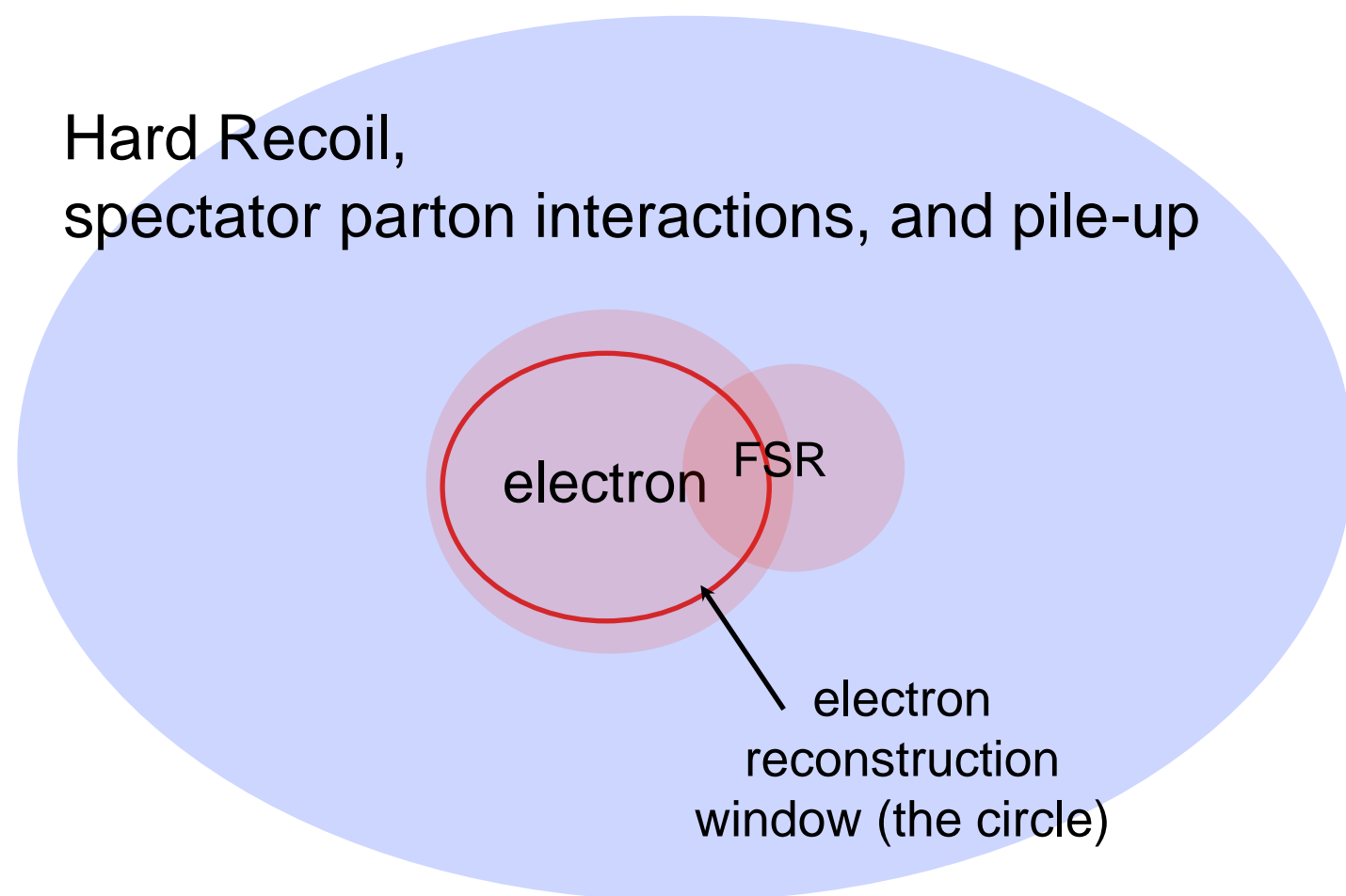
“pure” Hard Recoil  
balancing W or Z boson

Soft Recoil:  
pileup and  
spectator  
parton  
interactions

FSR photons that fly  
outside the electron  
reconstruction  
window.

Recoil energy that  
falls in the electron  
reconstruction  
window, as well as  
electron energy  
leakage to the recoil.

**In the same framework of  $\Delta E_{corr}$   
Modeling**  
**What has been added to (subtracted  
from) the electron has to be  
subtracted from (added to) the recoil.**



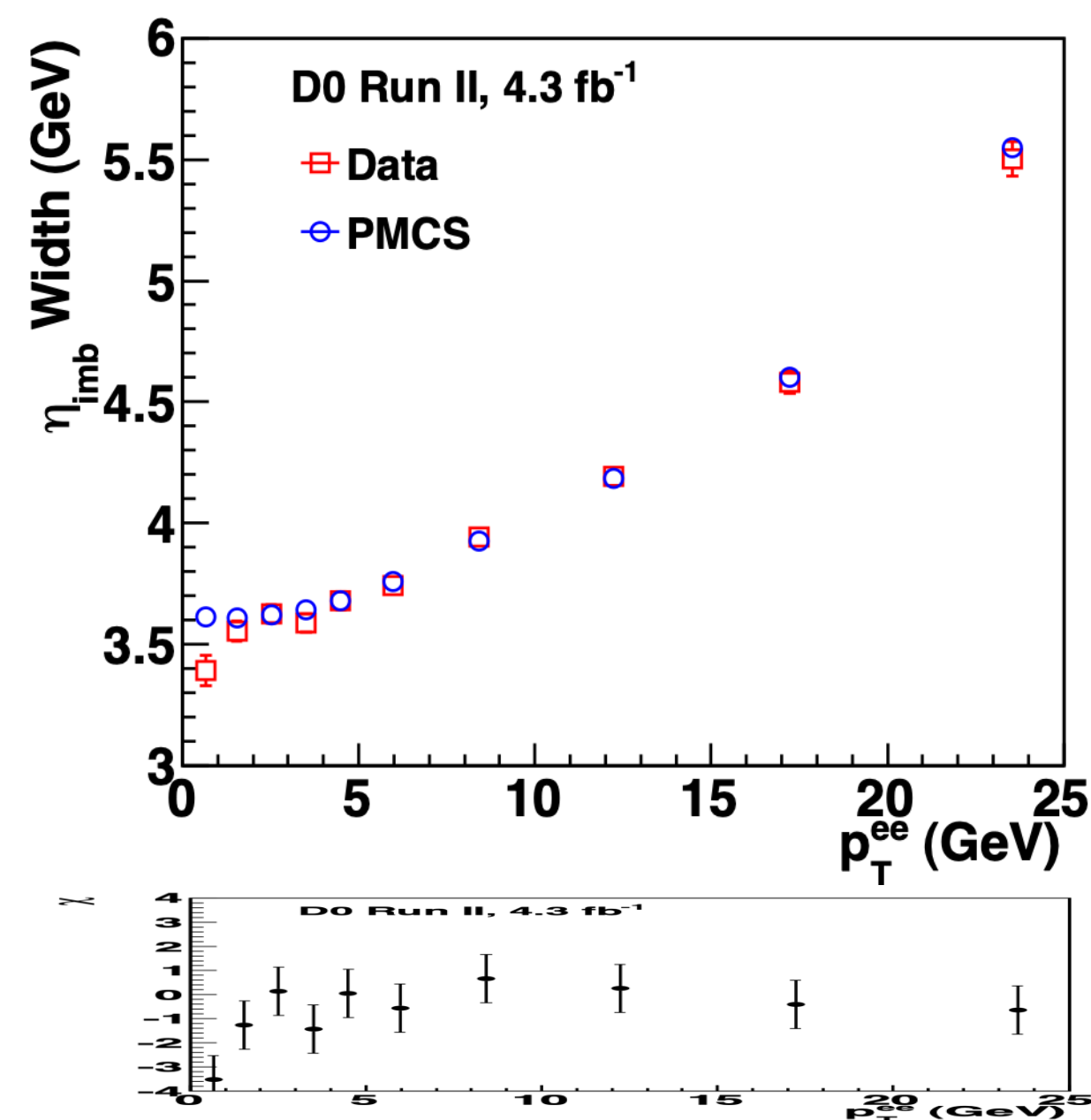
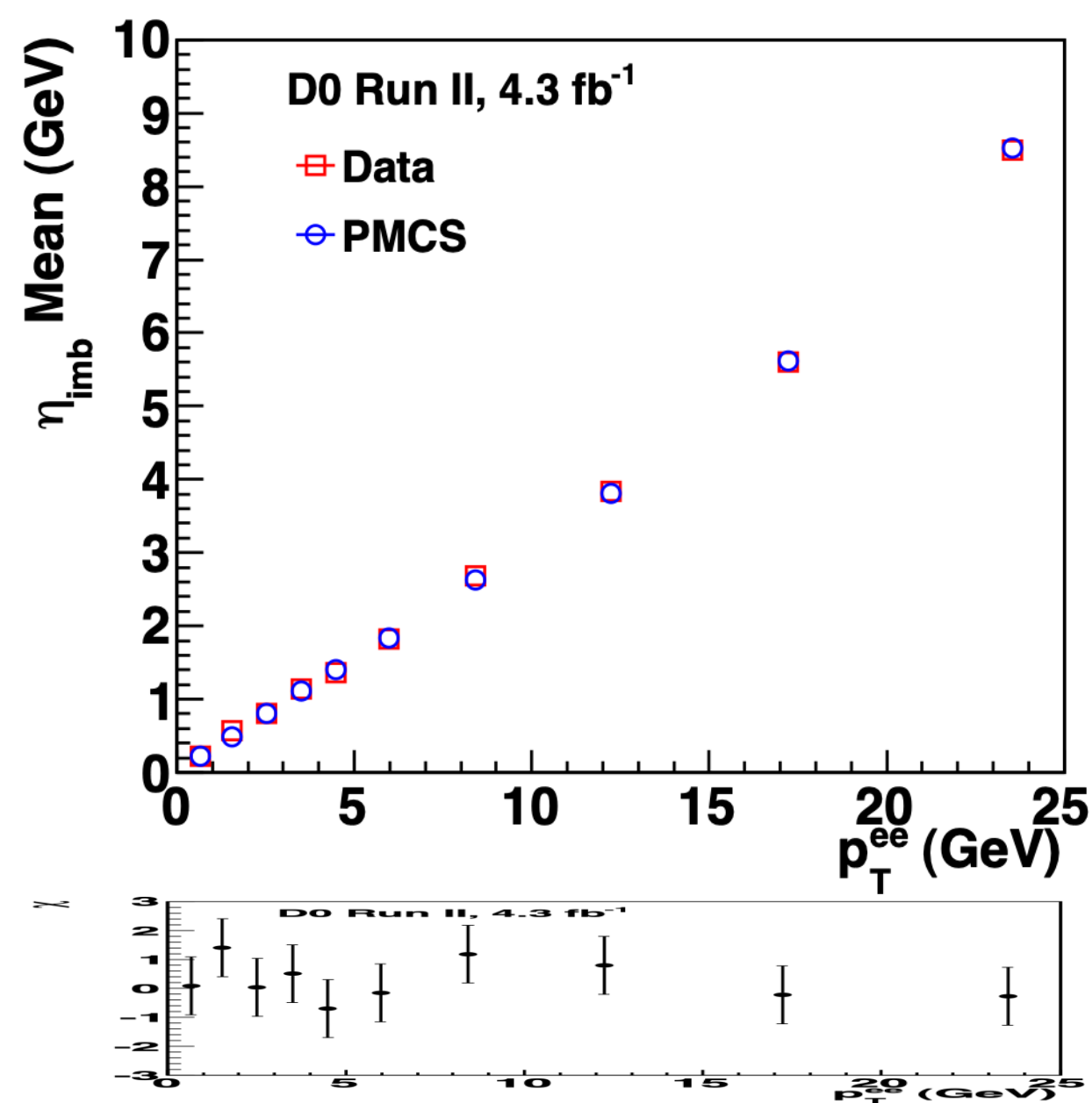


# Recoil fine tuning

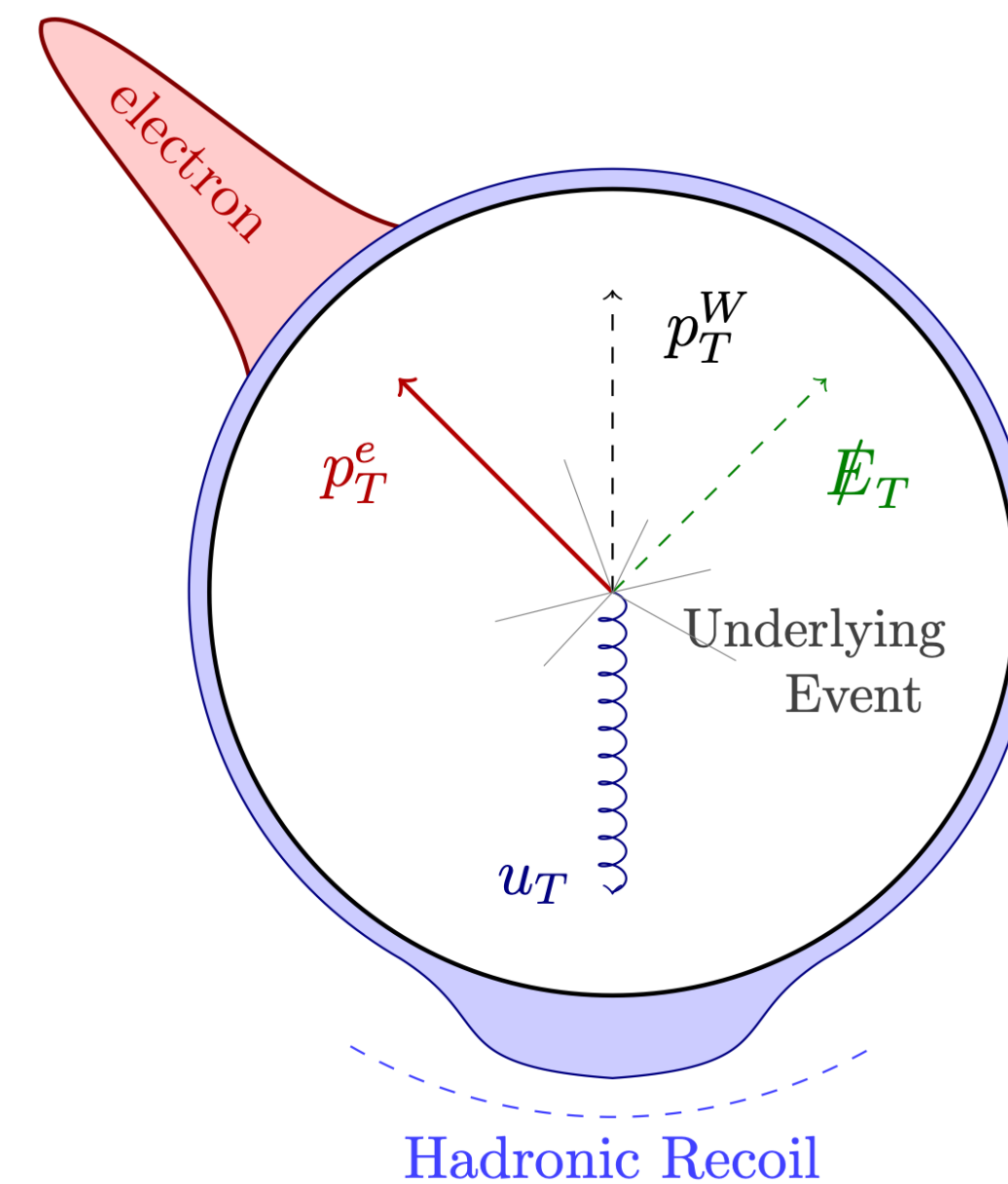
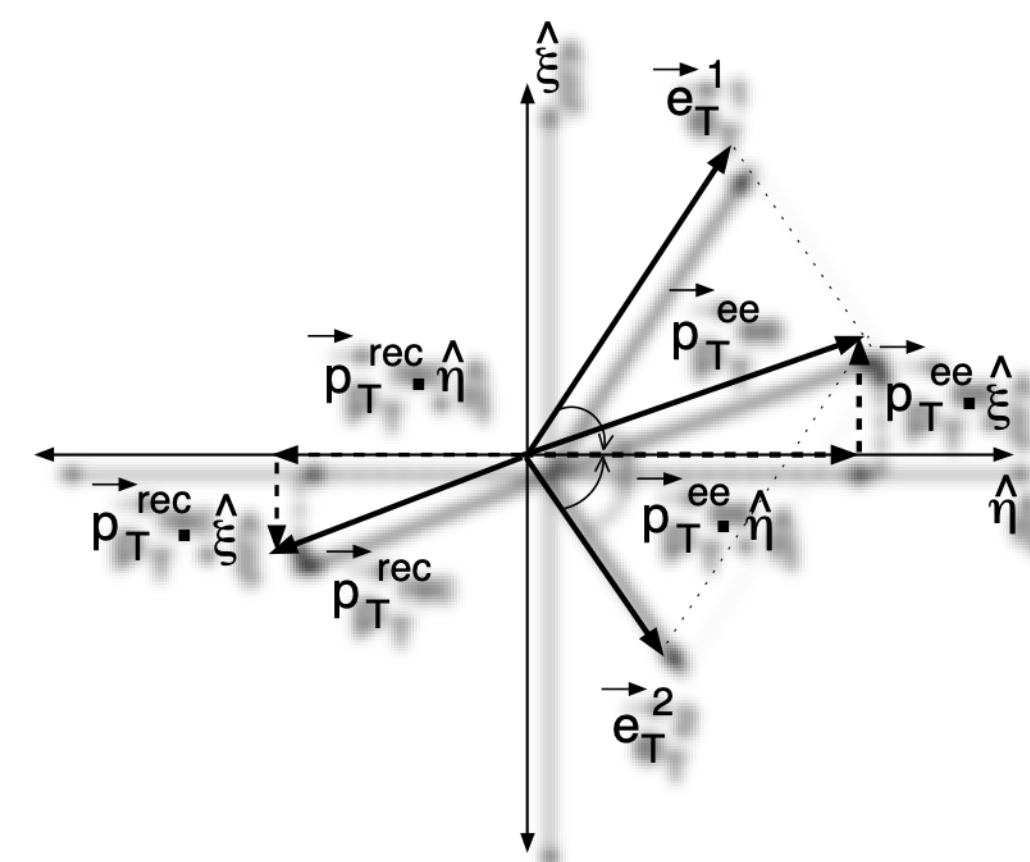
The recoil model is fine tuned using standard UA2 observables

$$\eta_{imb} = (\vec{p}_T^Z + \vec{u}_T) \cdot \hat{\eta}$$

- Use the mean of  $\eta$ -imbalance to tune the recoil response
- Use the width of  $\eta$ -imbalance to tune the recoil resolution



$\eta$ -axis: the bisector of two electron momenta of Z $\rightarrow$ ee events



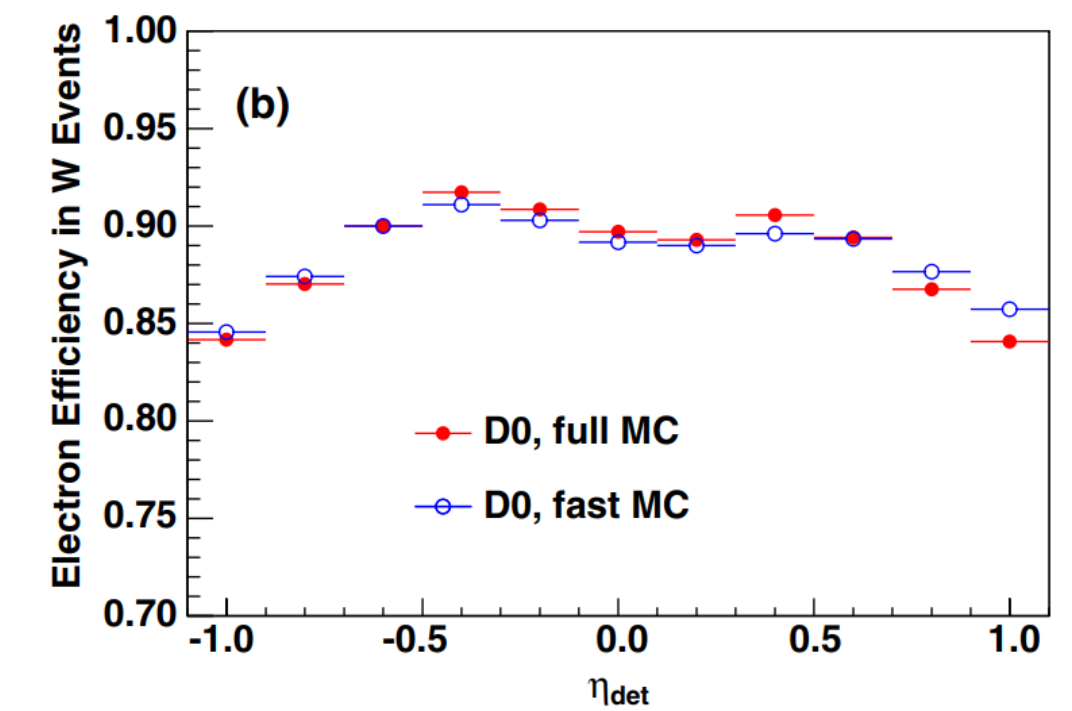
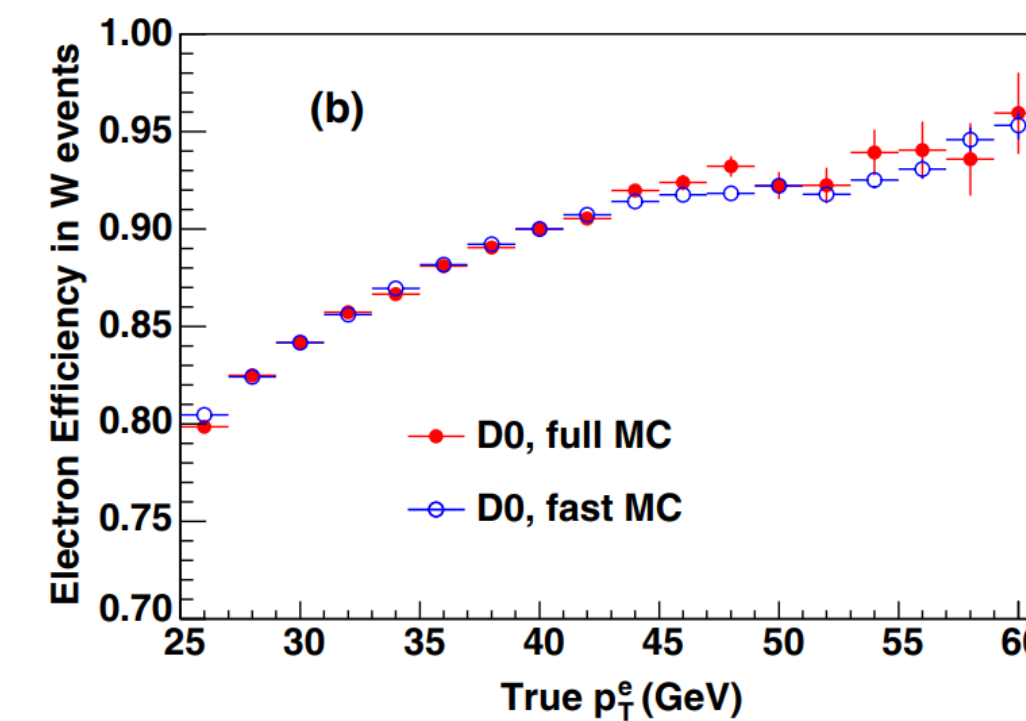
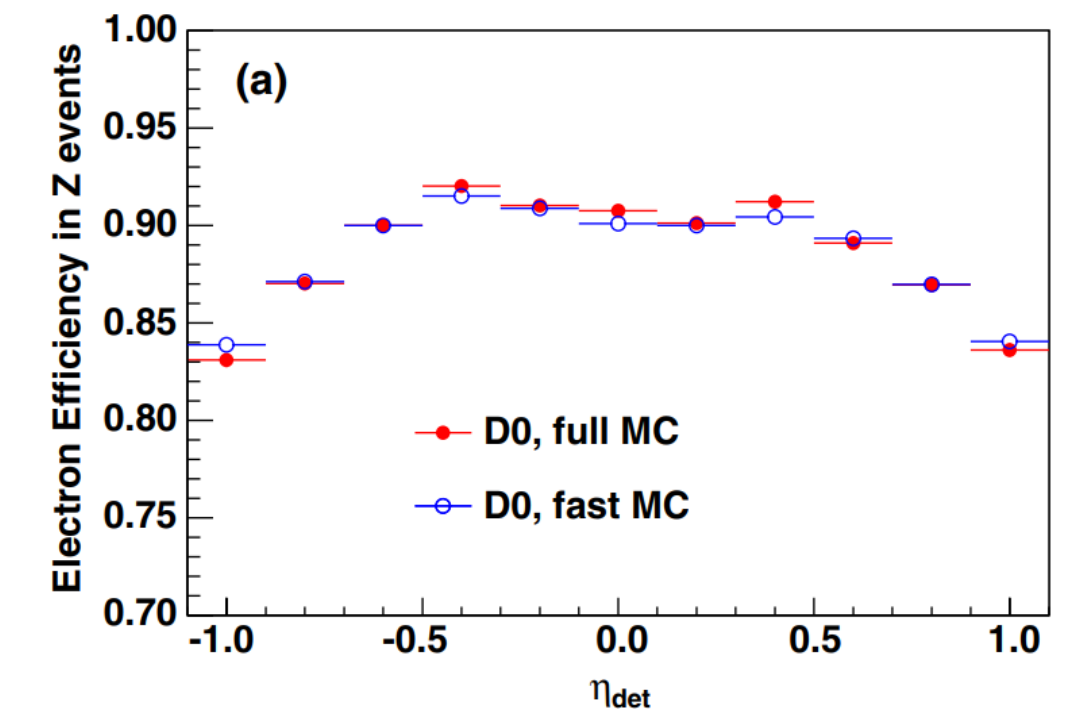
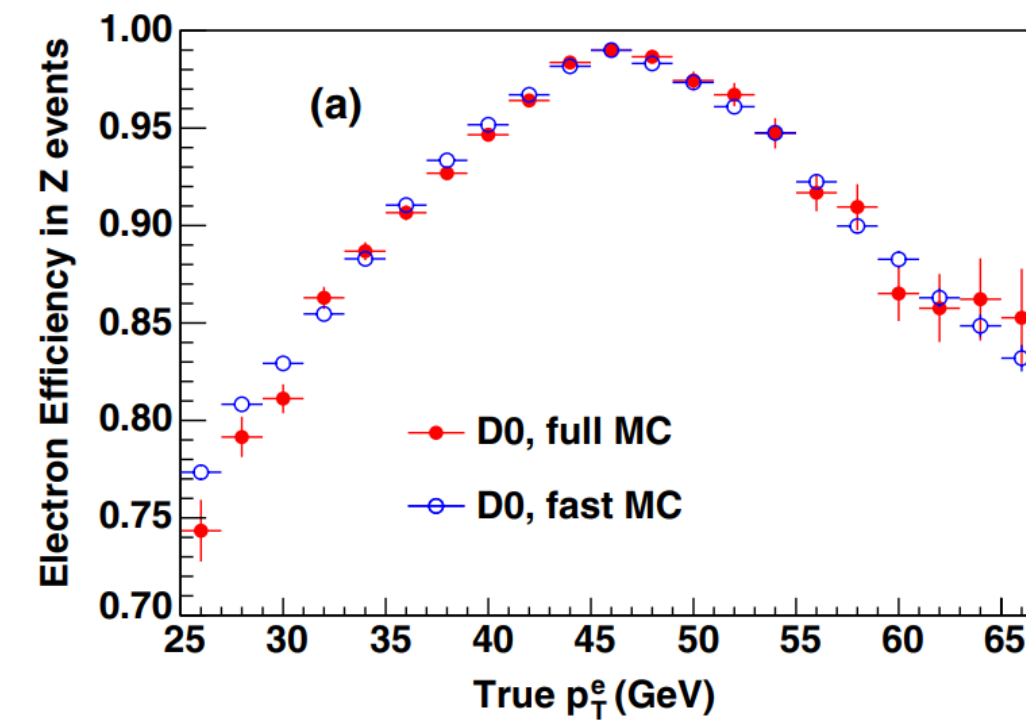
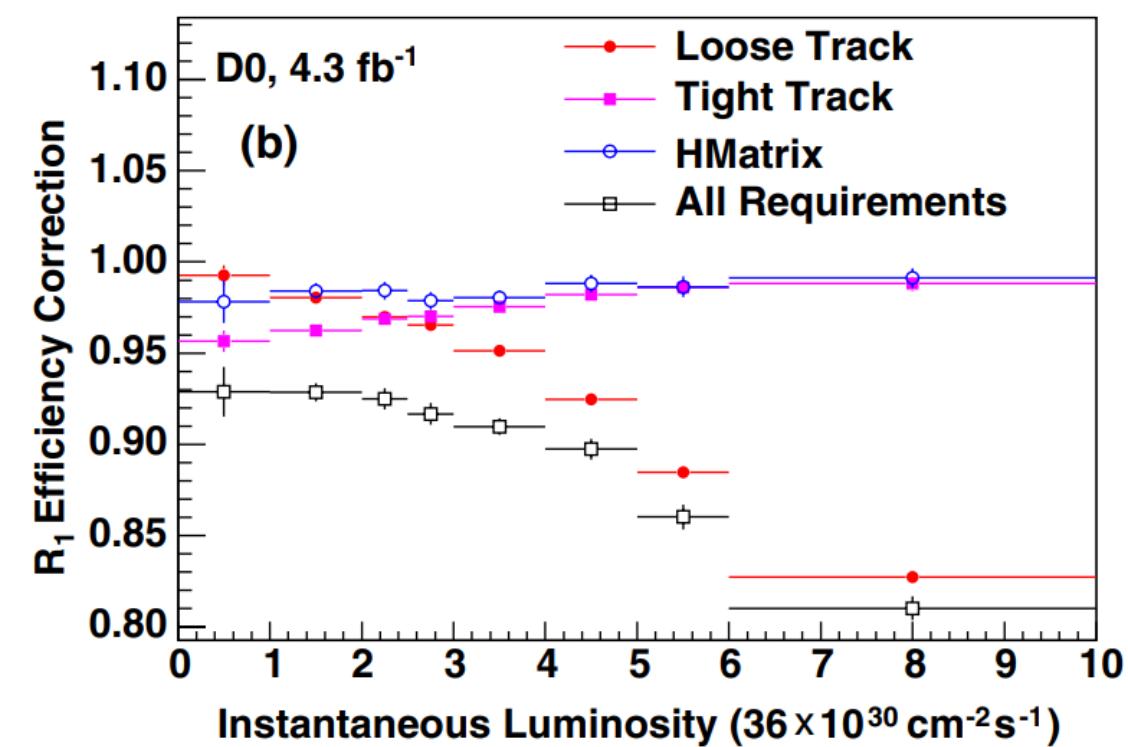
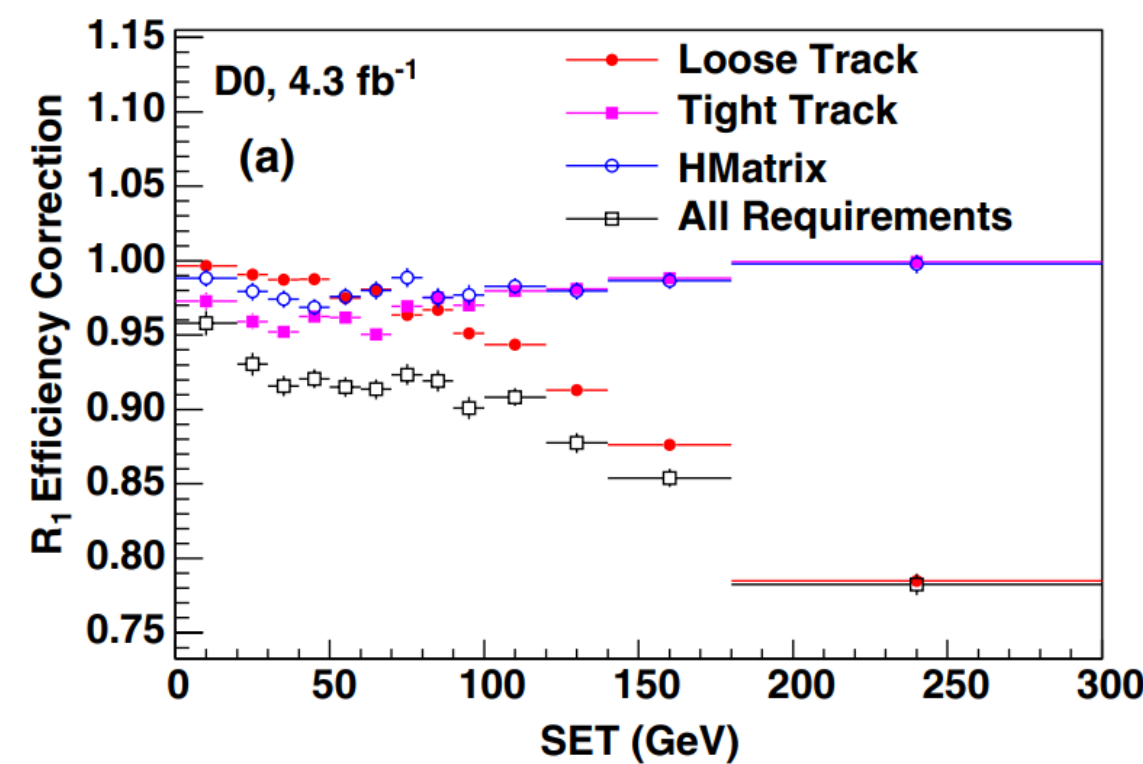
# Electron efficiency model

Efficiency modeling in the high inst. lumi. condition is challenging:

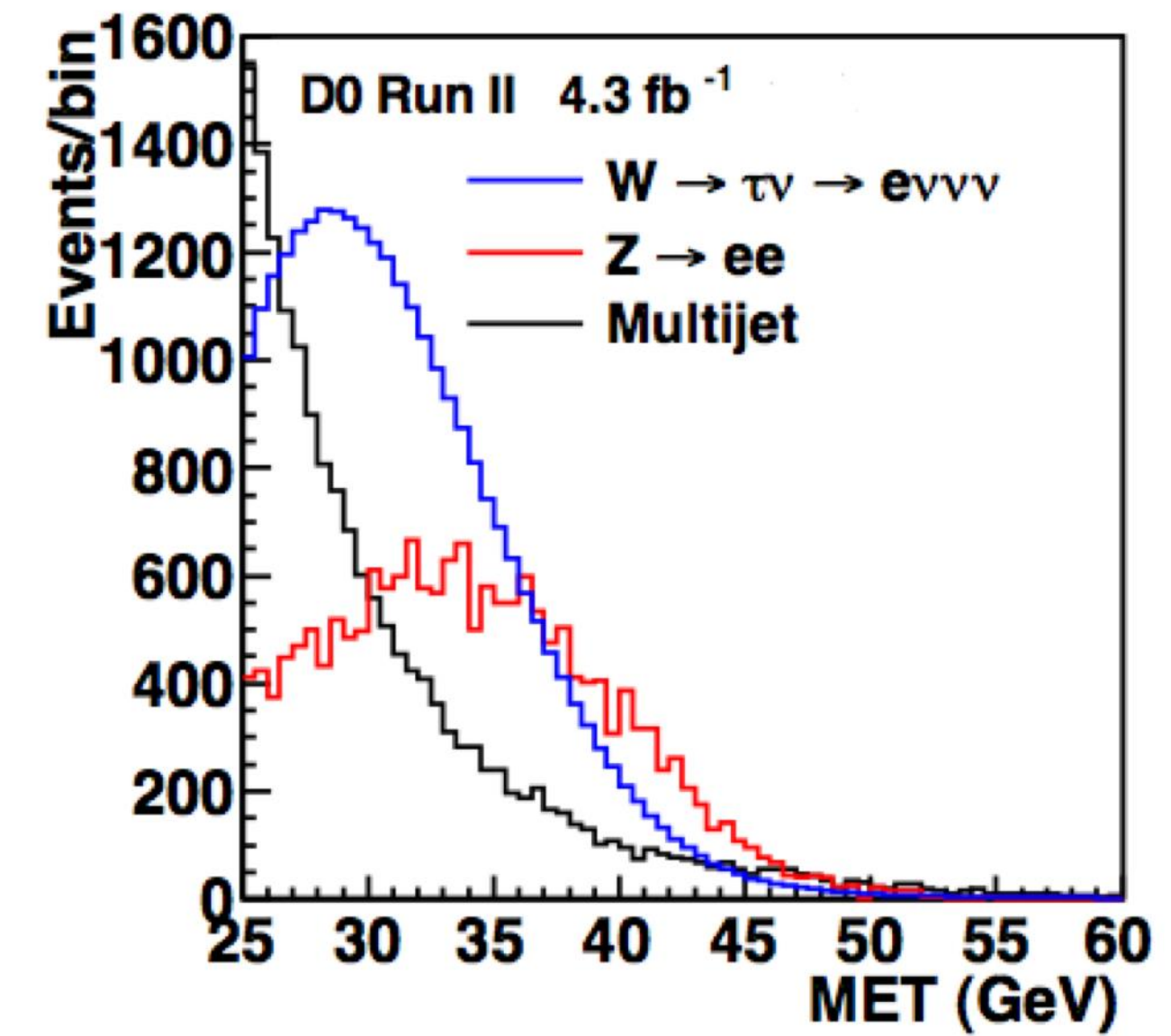
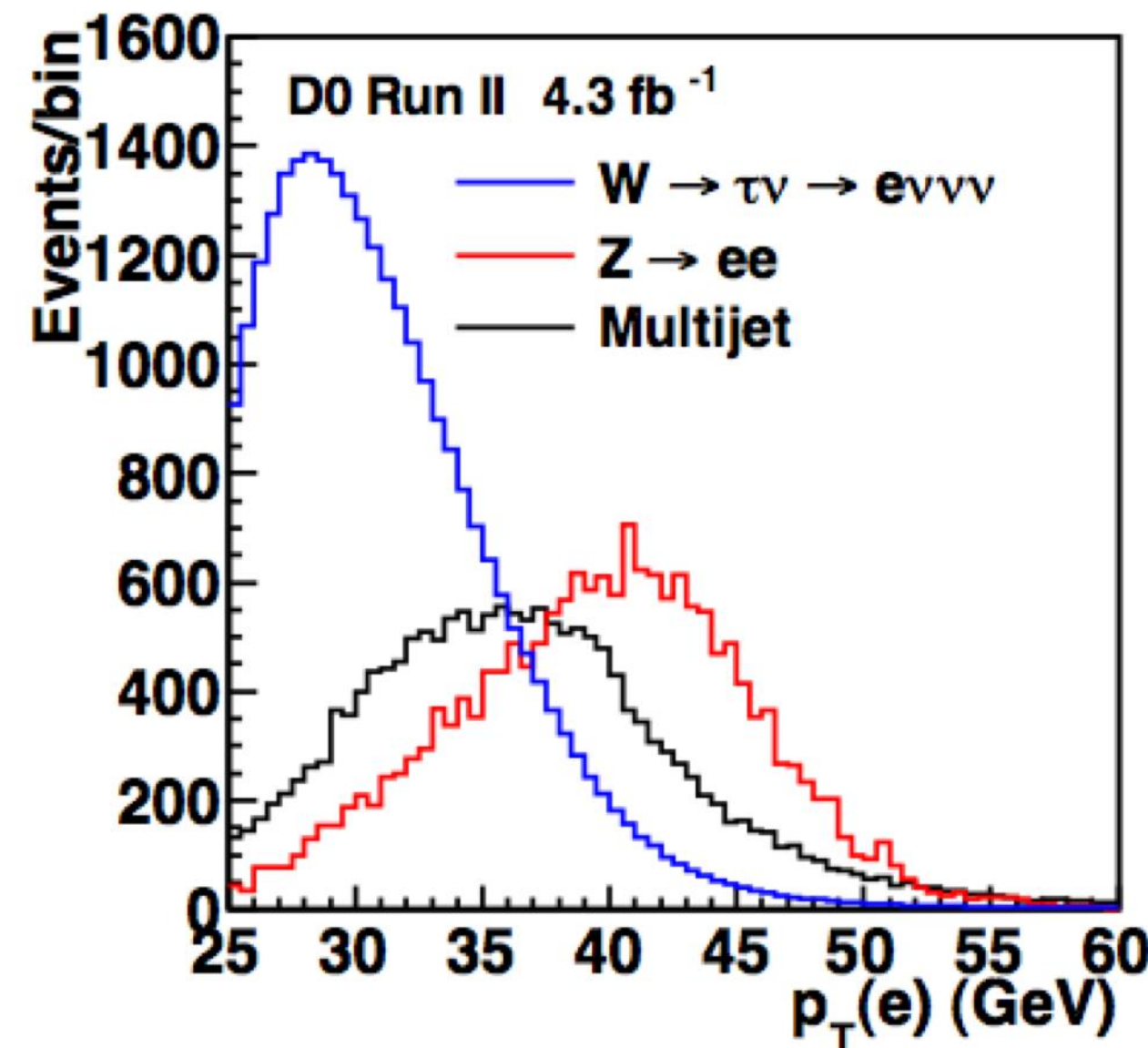
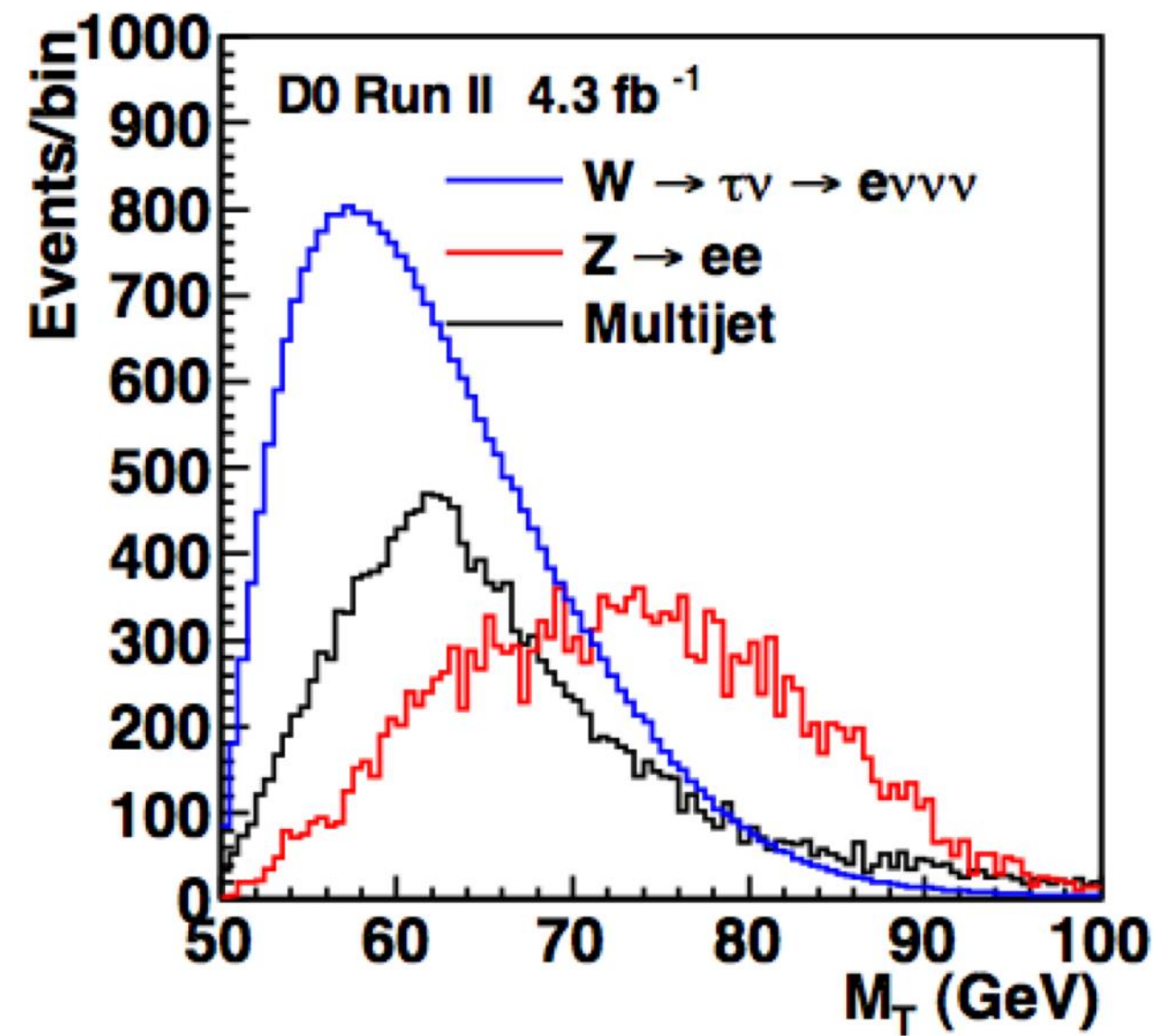
- pileup and hard recoil contaminate the electron reconstruction window,
- correlations with electron kinematics.

A two-step modeling:

- model the efficiency in a detailed simulation overlaid with pileup from collider data.
- check efficiency dependences using  $Z \rightarrow ee$  events comparing data and detailed simulation.



# Backgrounds



## Backgrounds:

**W->tau nu : 1.67%**

**Z->ee: 1.08%**

**Multijet (QCD): 1.02%**

Z->ee :

When the other electron falls in the cracks of the EM module  
 Estimated using electron + electron track

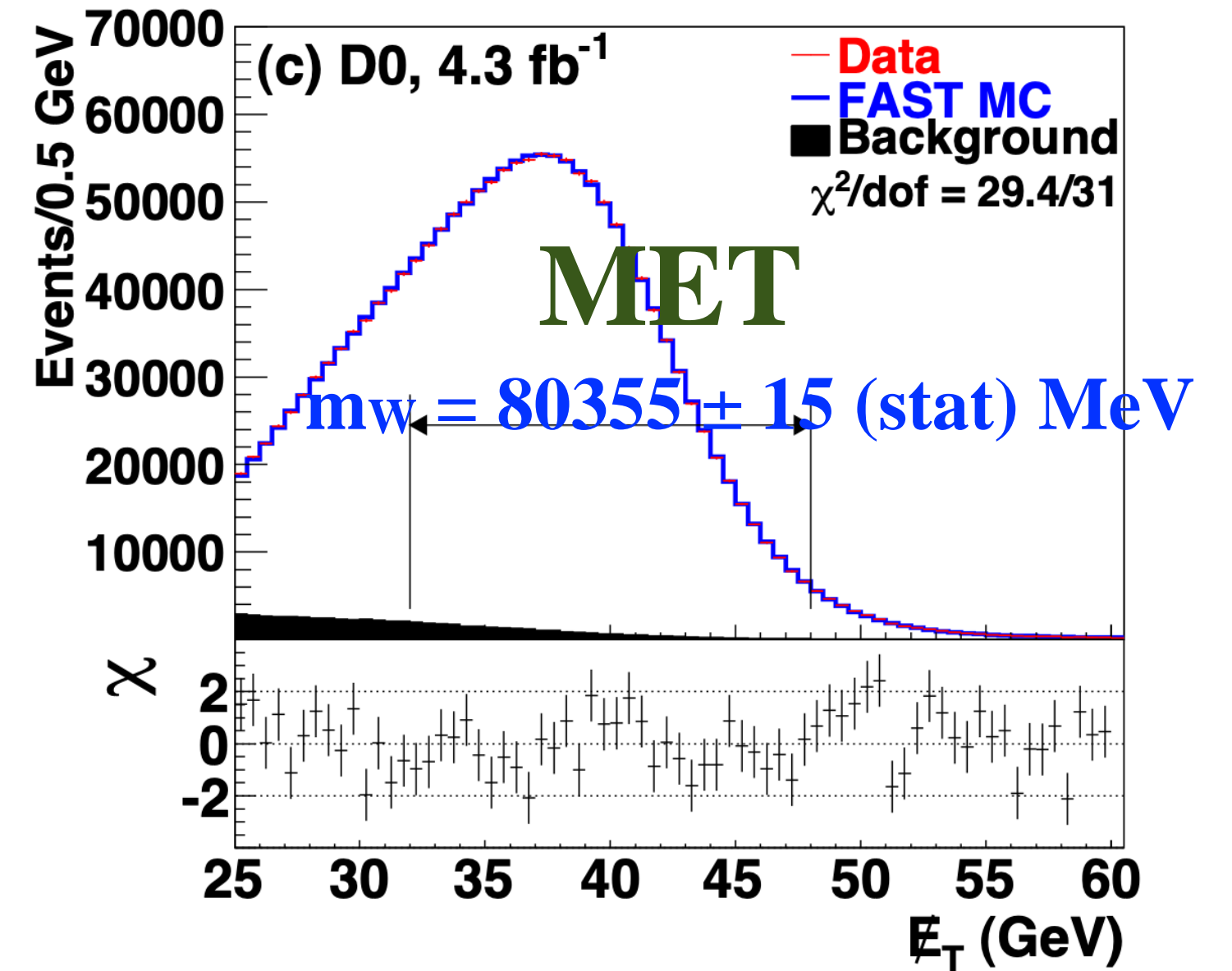
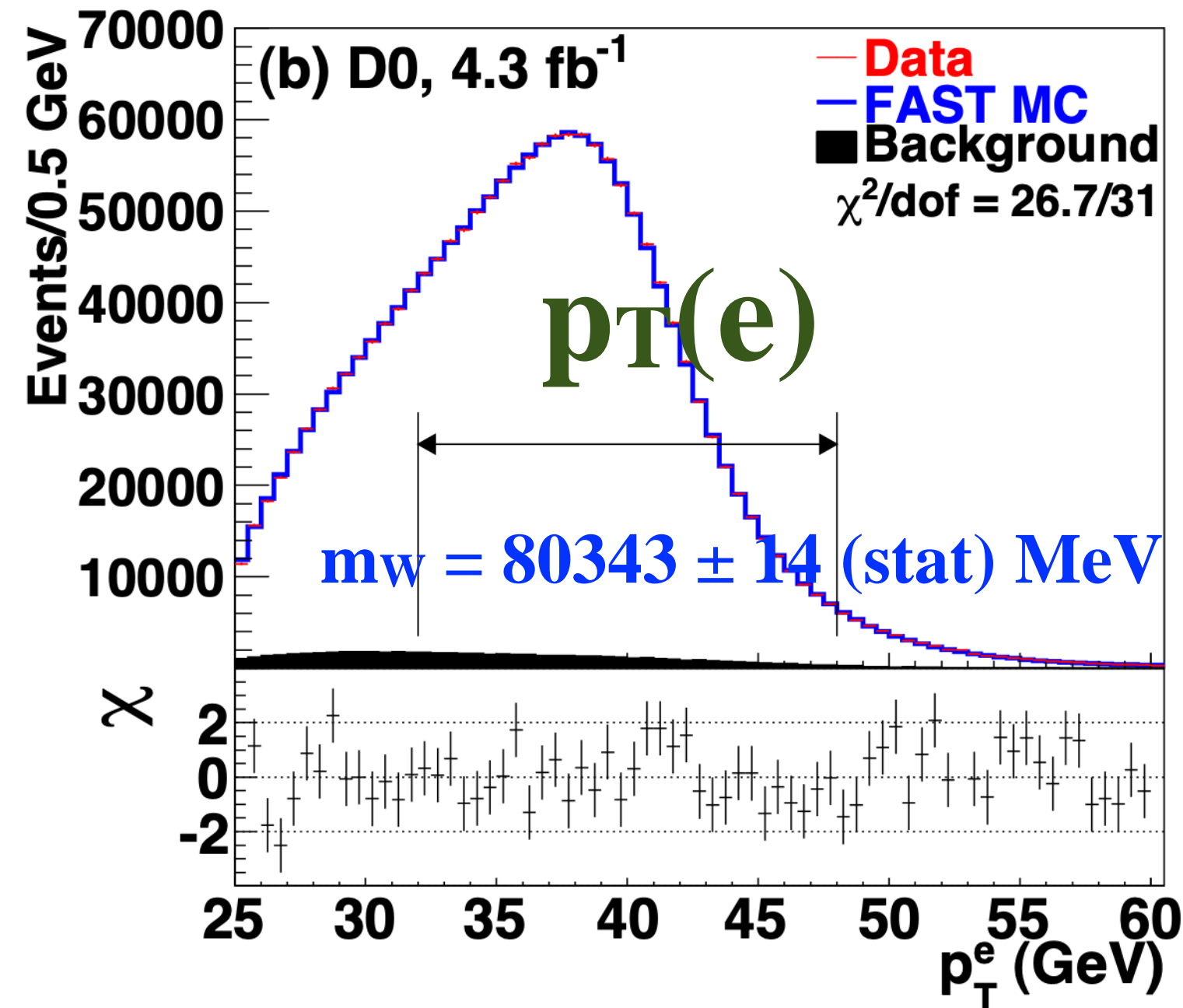
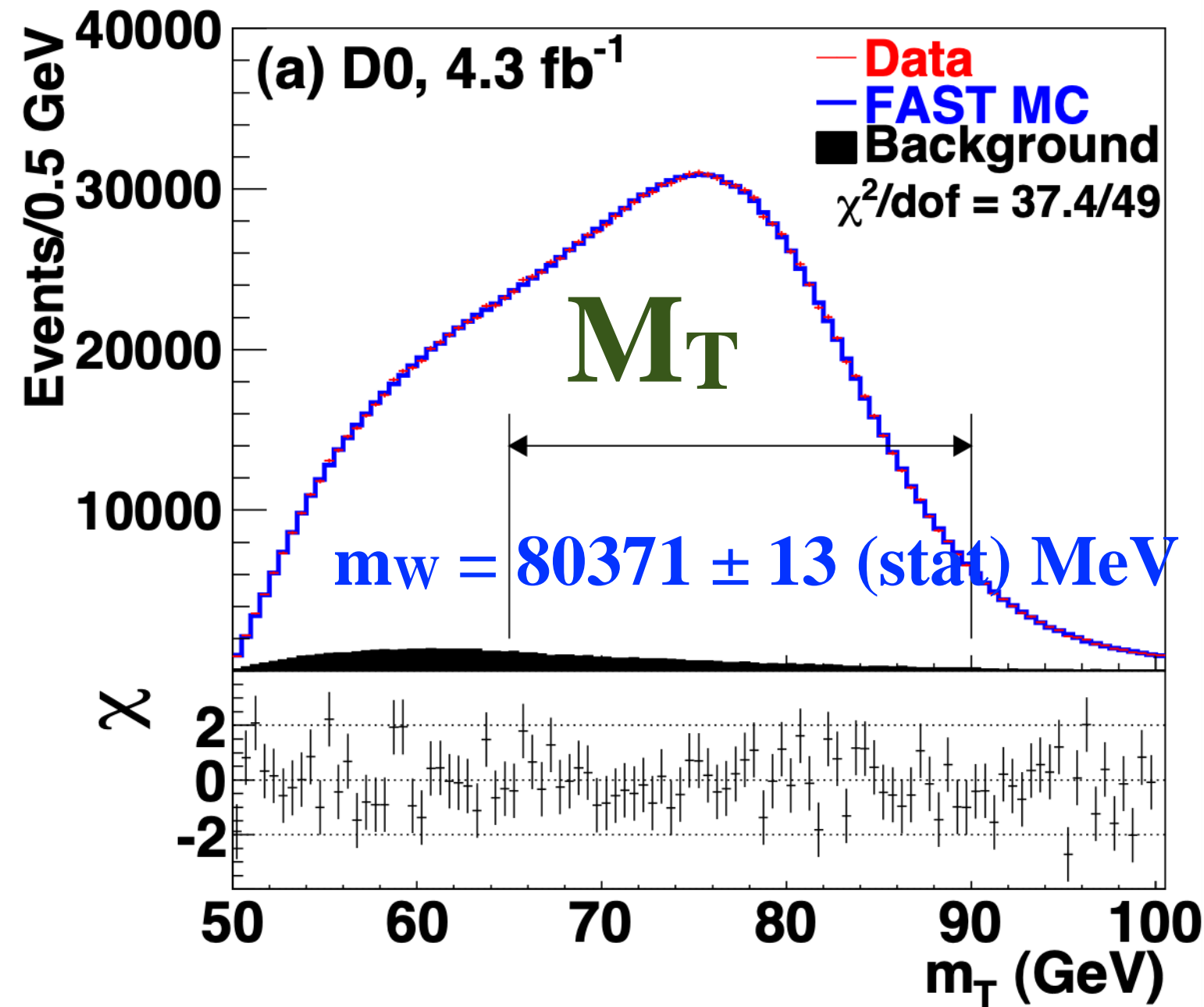
Multijet (QCD): matrix method, solving equation with loose cuts

$$N_L^{(i)} = N_W^{(i)} + N_{MJ}^{(i)},$$

$$N^{(i)} = \epsilon_e^{(i)} N_W^{(i)} + \epsilon_f^{(i)} N_{MJ}^{(i)},$$

# Results using 1/2 of the D0 full data set

1,677k W events



**The combined D0 5.3 fb<sup>-1</sup> result:**

5.3 fb<sup>-1</sup> = 1.0 fb<sup>-1</sup> (2009) + 4.3 fb<sup>-1</sup> (2012)

$$M_W = 80.375 \pm 0.011 \text{ (stat)} \pm 0.020 \text{ (syst)} \text{ GeV}$$

$$= 80.375 \pm 0.023 \text{ GeV.}$$

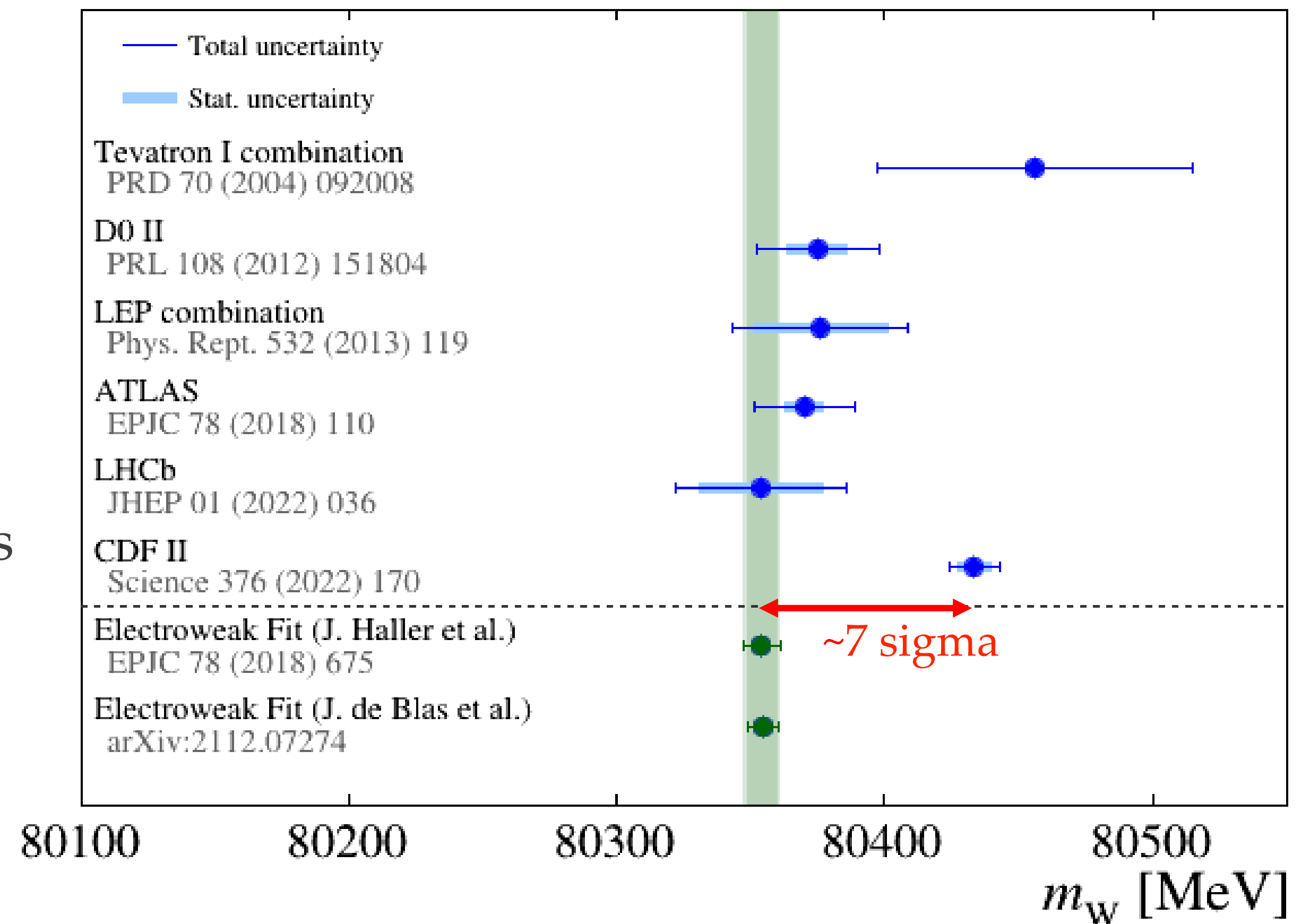
# Systematic uncertainties

D0 4.3 fb<sup>-1</sup>, e-channel

Source	$\sigma(m_W)$ MeV $m_T$	$\sigma(m_W)$ MeV $p_T^e$	$\sigma(m_W)$ MeV $E_T$
<b>Experimental</b>			
Electron Energy Scale	16	17	16
Electron Energy Resolution	2	2	3
Electron Energy Nonlinearity	4	6	7
$W$ and $Z$ Electron energy loss differences	4	4	4
Recoil Model	5	6	14
Electron Efficiencies	1	3	5
Backgrounds	2	2	2
<b>Experimental Total</b>	18	20	24
<b>W production and decay model</b>			
PDF	11	11	14
QED	7	7	9
Boson $p_T$	2	5	2
<b>W model Total</b>	13	14	17
<b>Total Systematic Uncertainty</b>	22	24	29

# Summary

- ❖ A highly motivated analysis:
  - ❖ Precisely examine the standard model, searching for evidences for new physics.
- ❖ CDF Run II gives the most precise measurement at the moment of an uncertainty 9 MeV with  $\sim 7$  sigma apart from the EW global fit
- ❖ For a better understanding of the new CDF results, the D0 measurement using 1/2 D0 data is reviewed in this talk, which gave a 23 MeV precision in 2012.
- ❖ Future results are expected from LHC experiments and future e<sup>+</sup>e<sup>-</sup> colliders.



## Backup slides

# Uncertainties and projection to the D0 full data set

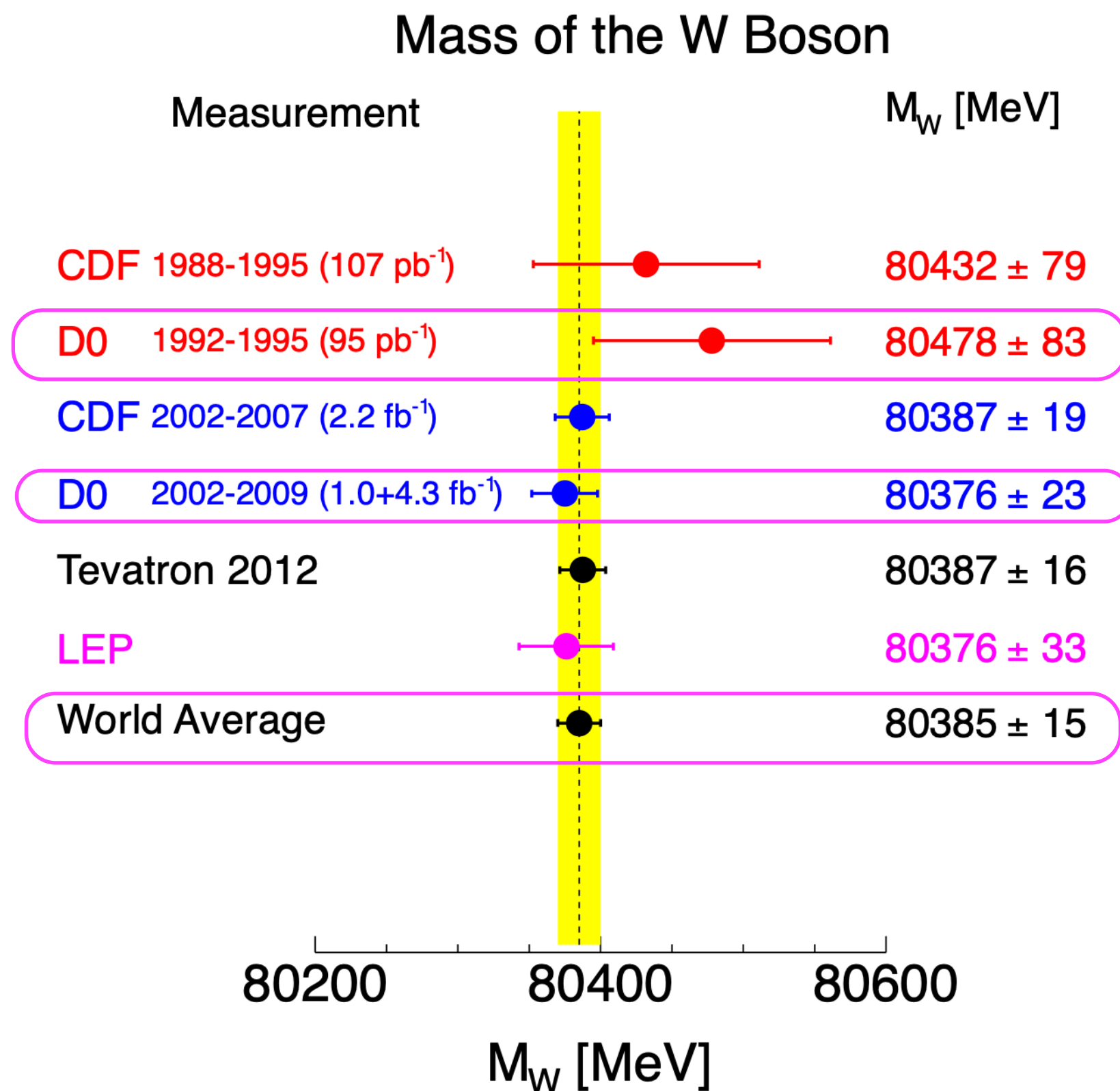
CC: Central Calorimeter  
EC: Endcap Calorimeter

Source (Unit in MeV)	Published (2009) 1 fb <sup>-1</sup> CC	Published (2012) 4.3 fb <sup>-1</sup> CC	Projection 10 fb <sup>-1</sup> CC	Projection 10 fb <sup>-1</sup> CC+EC
<b>Statistical</b>	<b>23</b>	<b>13</b>	<b>9</b>	<b>8</b>
Experimental syst.				
Electron energy scale	34	16	11	10
Electron energy resolution	2	2	2	2
Electron energy nonlinearity	4	4	2	2
W and Z electron energy loss differences	4	4	2	2
Recoil model	6	5	3	2
Electron efficiencies	5	1	1	1
Backgrounds	2	2	2	2
Exp. Syst. Subtotal	35	18	12	11
Theoretical syst.				
PDF	9	11	11	5
QED	7	7	3	3
Boson pT	2	2	2	2
Theo. Syst. Subtotal	12	13	12	6
<b>Systematic total</b>	<b>37</b>	<b>22</b>	<b>17</b>	<b>13</b>
<b>Total</b>	<b>44</b>	<b>26</b>	<b>19</b>	<b>15</b>



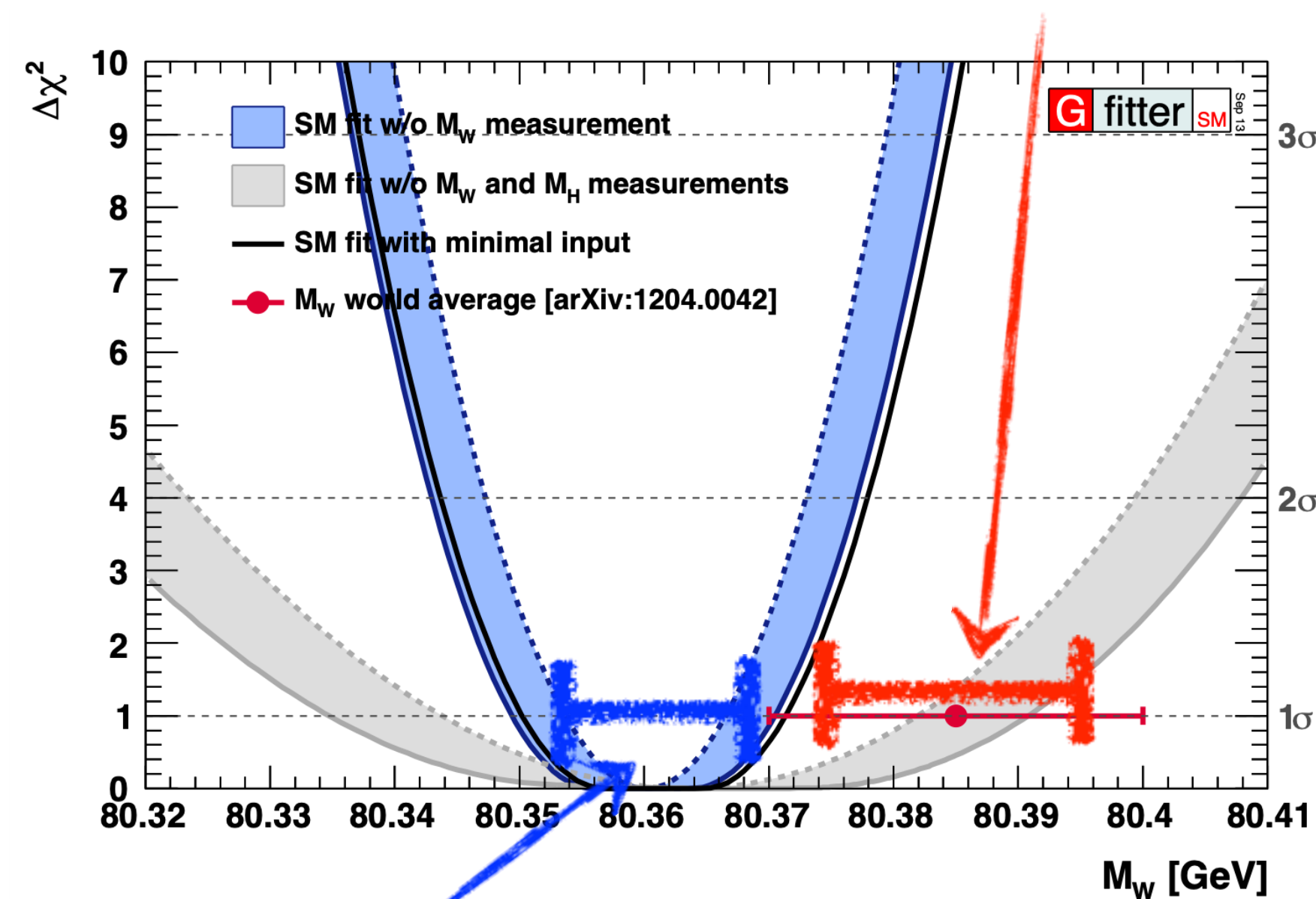
# World average and expected new results

Adding the expected new D0 results:



We will have

Measurement	$M_W$ [MeV]
CC (10 fb <sup>-1</sup> )	19 MeV
CC+EC (10 fb <sup>-1</sup> )	15 MeV
New world aver. :	14 MeV
	12 MeV



If the central value doesn't change in the future, we can expect a 2-sigma deviation apart from the SM. But if we expect better discrimination, we need CDF, ATLAS, CMS, and future e+e- colliders: CEPC, ILC, FCC-ee

The predicted W mass uncertainty can also be narrowed by improve the Top mass measurement.

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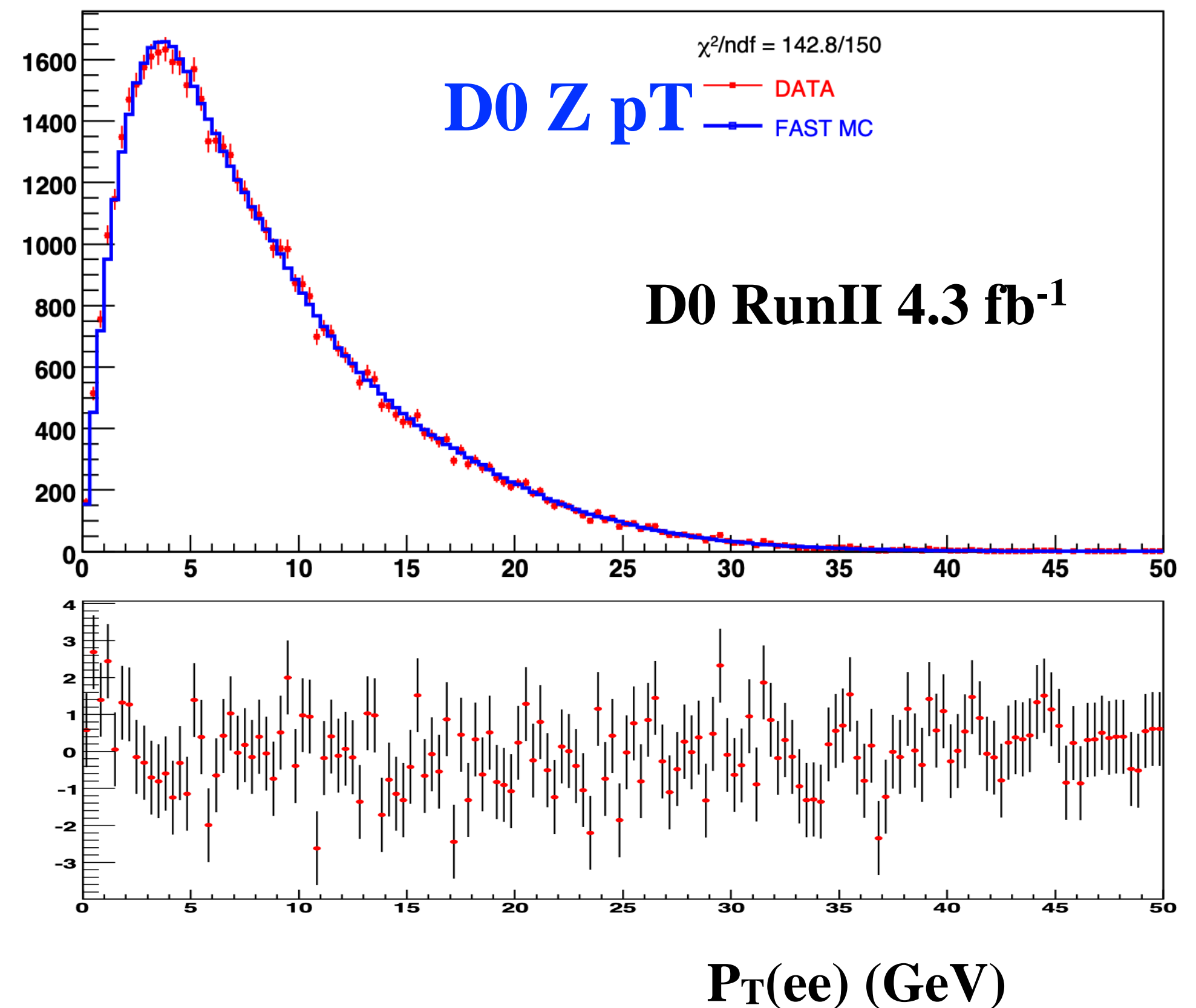
# Some words about the PDF Uncertainty

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- ❖ In principle, the transverse observables (e.g.  $M_T$ ,  $P_T(e)$ ) are insensitive to the uncertainties in the (longitudinal) parton distribution functions (PDF).
- ❖ However, our cuts on the leptons  $\eta$  ( $|\eta| < 1.0$ ) is not invariant under longitudinal boosts. Changes in PDFs can modify the shapes of the transverse observables under  $\eta$  cuts. Therefore, PDF uncertainties are introduced.
- ❖ Ways to reduce the PDF uncertainties:
  - ❖ Extending the  $\eta$  coverage as much as possible, including end-cap leptons:
    - ❖ Can reduce by a factor of two, need to understand the energy scale, pileup, and backgrounds for the end-cap leptons.
  - ❖ Reduce the PDF uncertainties by other measurements:
    - ❖ e.g.  $W$  charge asymmetry measurements.

# Theoretical modeling

- **Resbos:** Next-to-leading order event generator with next-to-next-to-leading logarithm resummation of soft gluons, gives the best boson pT description so far. [C. Balazs and C. P. Yuan, *Phys. Rev. D* 56, 5558 (1997).]
- **Photos:** generates up to two final state radiation photons. [P. Golonka and Z. Was, *Eur. Phys. J. C* 45, 97 (2006).]



# Model of W production and decay

Tool	Process	QCD	EW
RESBOS	$W, Z$	NLO	-
WGRAD	$W$	LO	complete $\mathcal{O}(\alpha)$ , Matrix Element, $\leq 1$ photon
ZGRAD	$Z$	LO	complete $\mathcal{O}(\alpha)$ , Matrix Element, $\leq 1$ photon
PHOTOS			QED FSR, $\leq 2$ photons

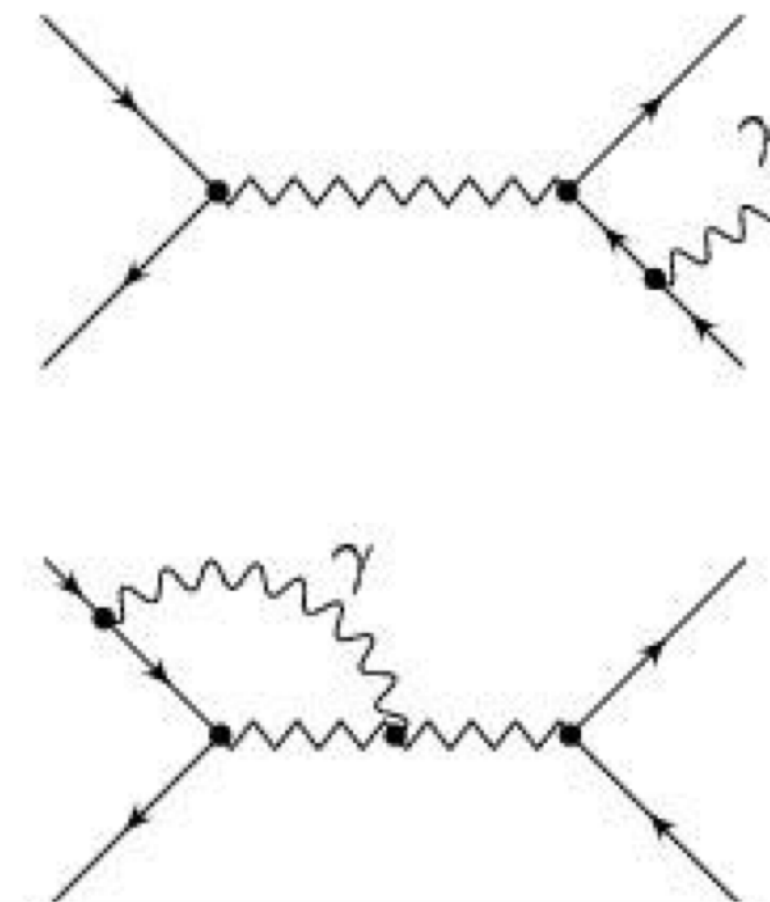
Our main generator is “**ResBos+Photos**”. The NLO QCD in **ResBos** allows us to get a reasonable description of the  $p_T$  of the vector bosons. The two leading EWK effects are the first FSR photon and the second FSR photon. **Photos** gives us a reasonable model for both.

We use **W/ZGRAD** to get a feeling for the effect of the full EWK corrections.

The final “QED” uncertainty we quote is **7/7/9 MeV** ( $m_T, p_T, MET$ ).

This is the sum of different effects; the two main ones are:

- Effect of full EWK corrections, from comparison of W/ZGRAD in “FSR only” and in “full EWK” modes (**5/5/5 MeV**).
- Very simple estimate of “quality of FSR model”, from comparison of W/ZGRAD in FSR-only mode vs **Photos** (**5/5/5 MeV**).



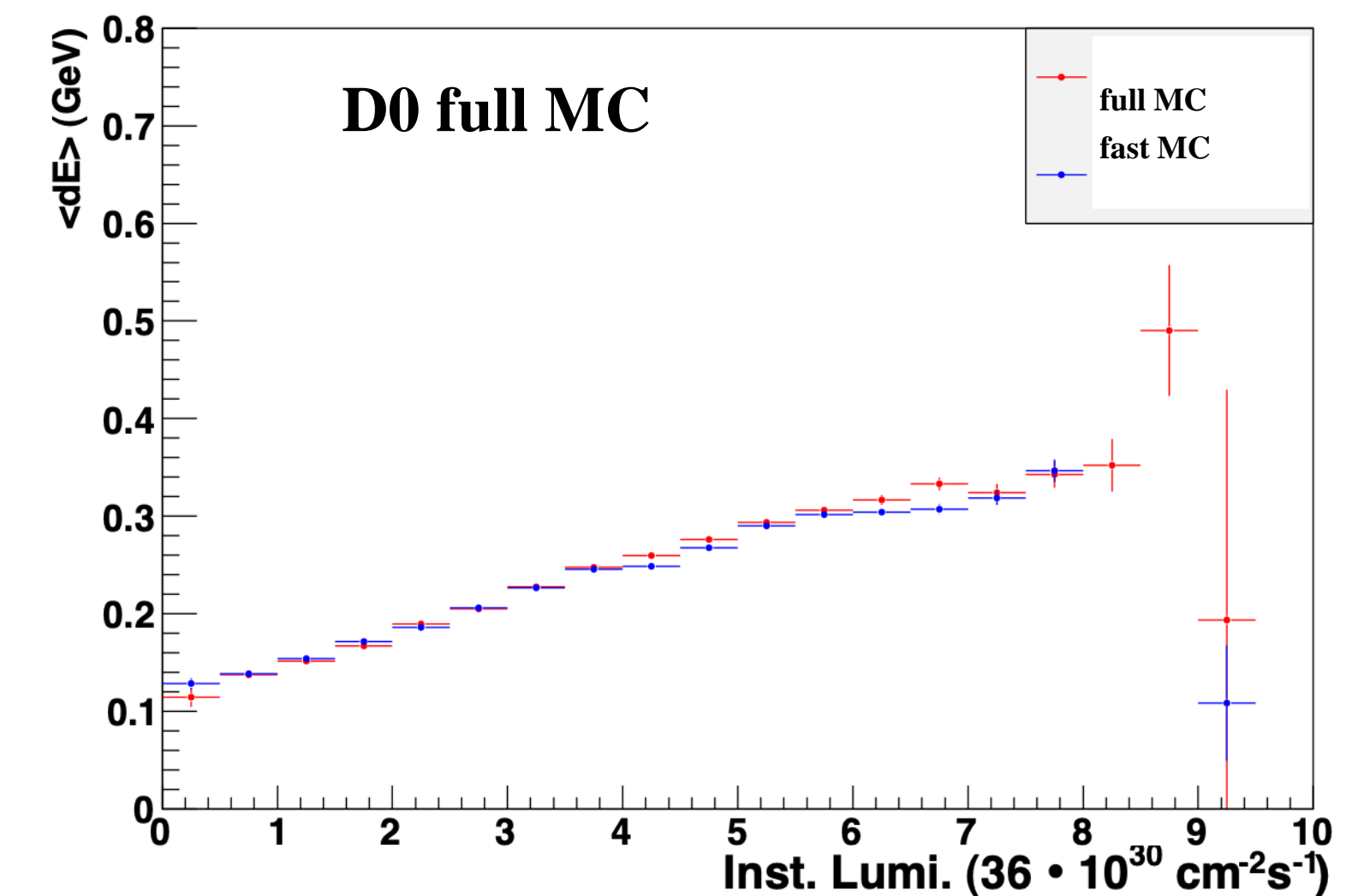
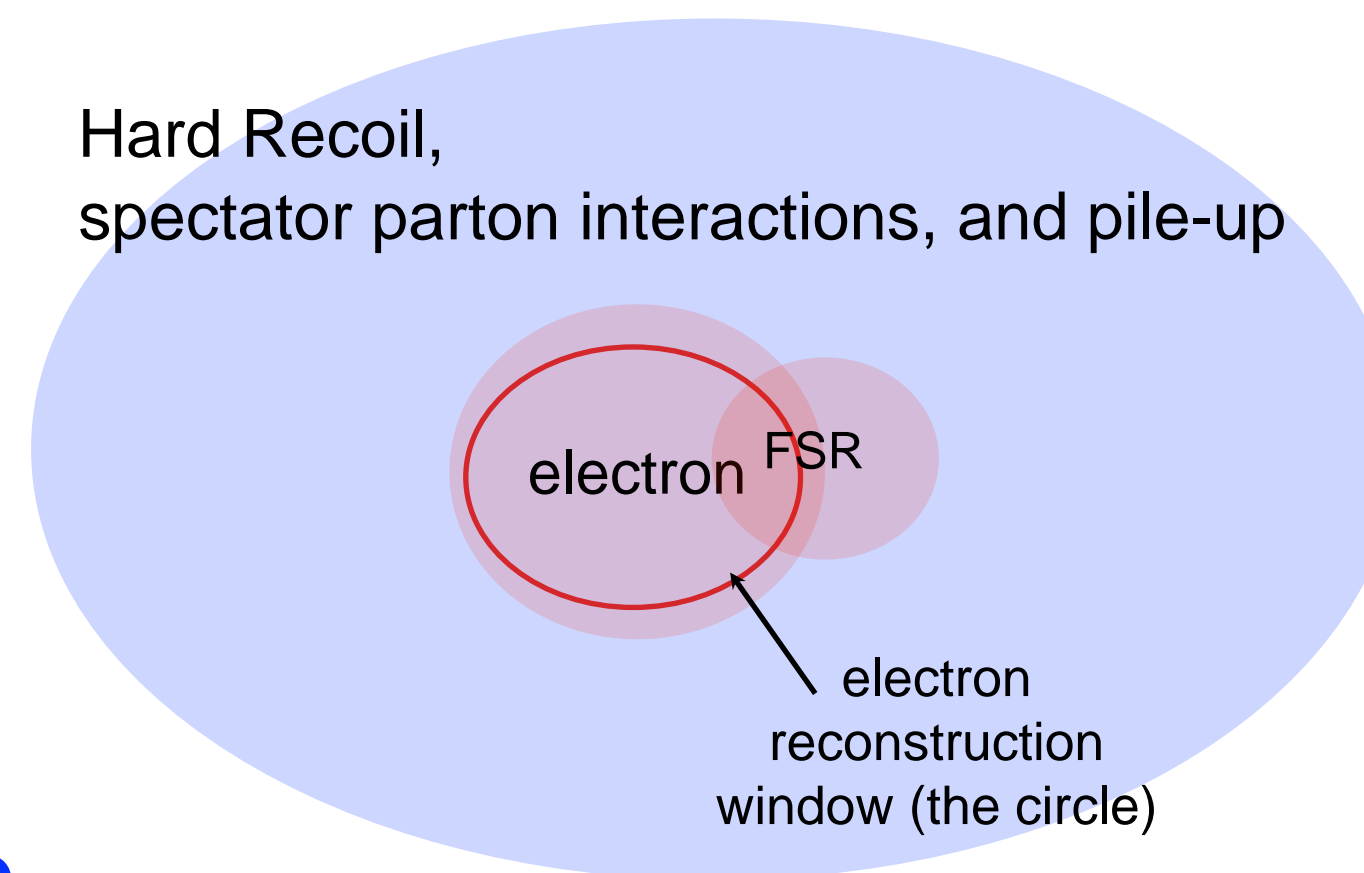
# Modeling the pileup contamination

**Electron Model:**  $E_{reco} = R_{EM}(E_{true}) \otimes \sigma_{EM}(E_{true}) + \Delta E_{corr}$

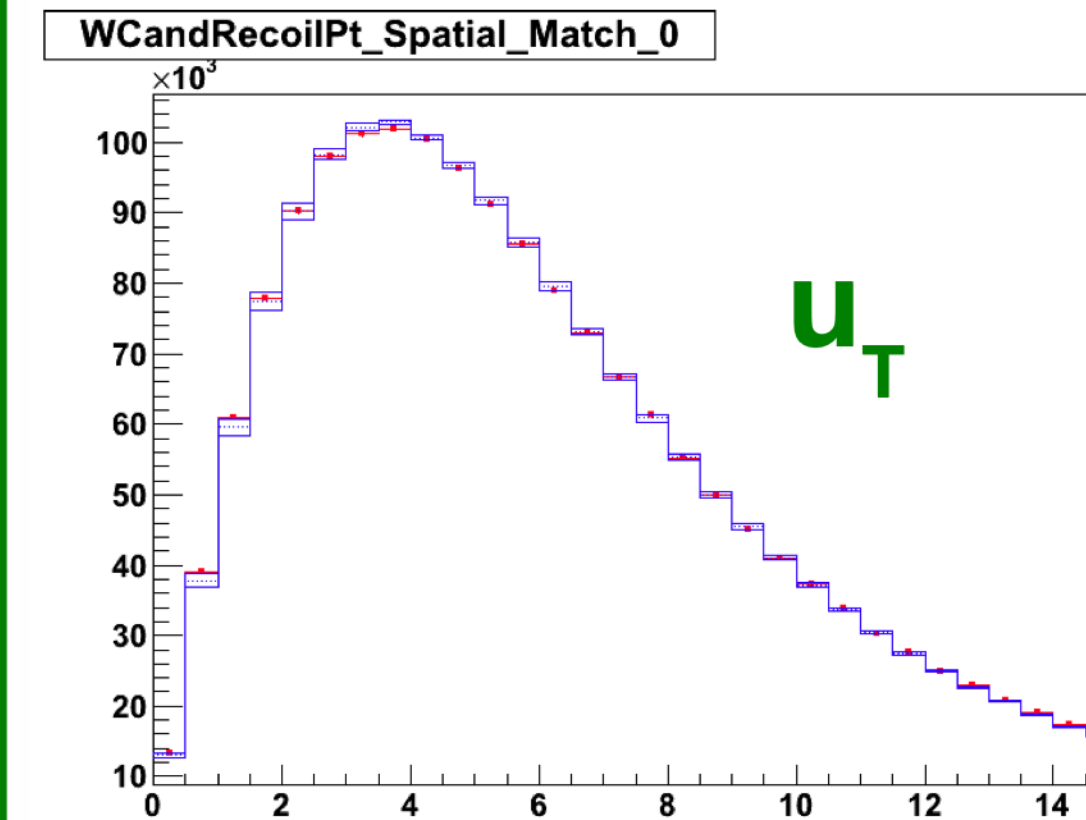
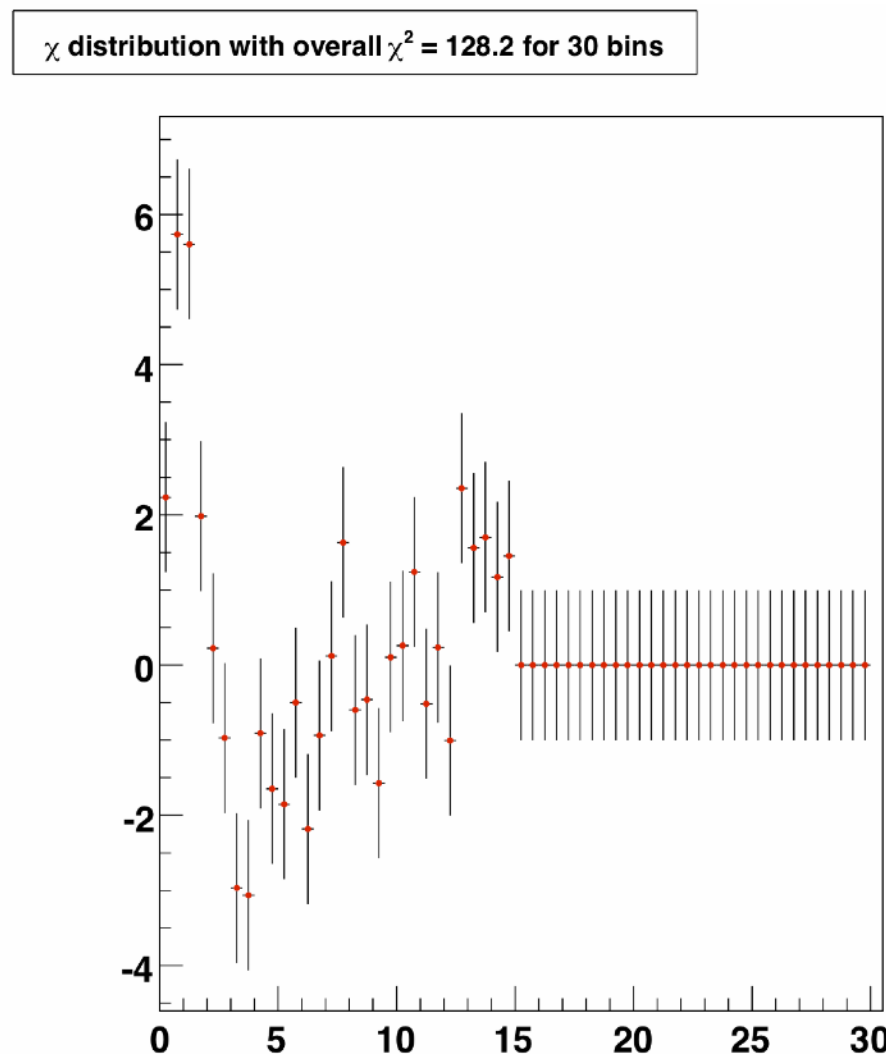
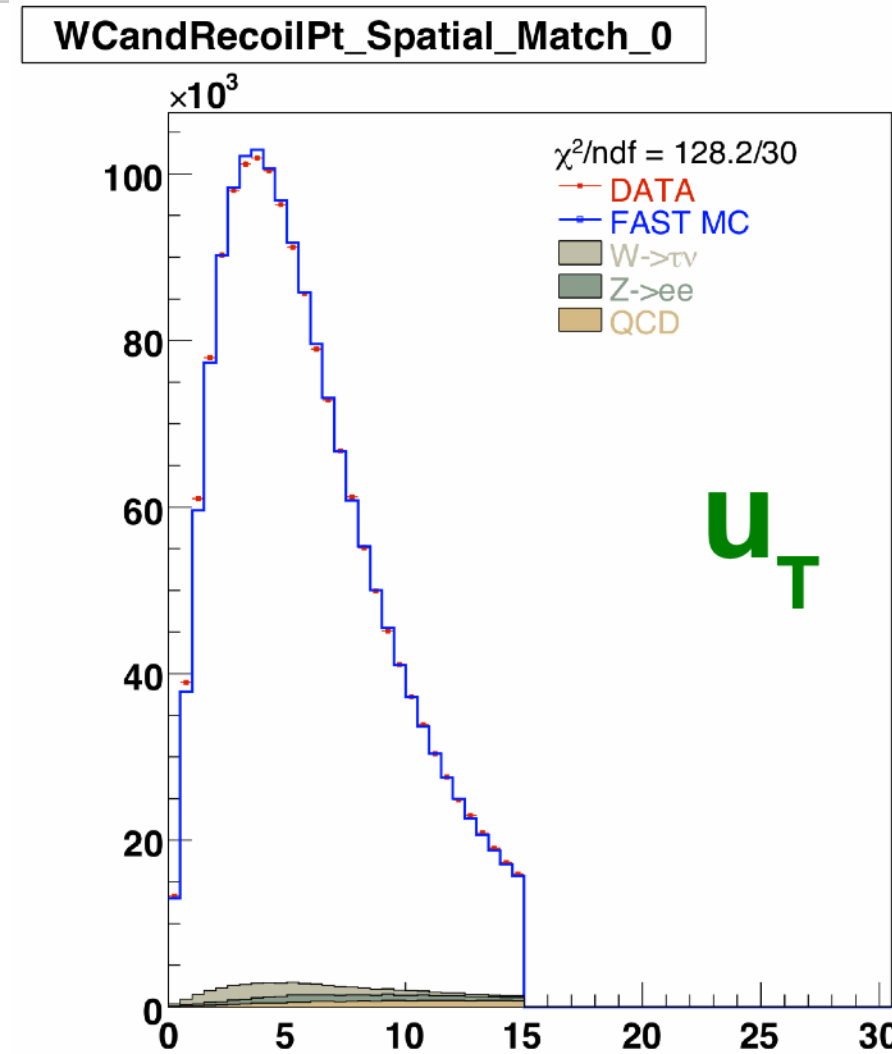
**Response**                      **Resolution**                      **Energy correction**

## $\Delta E_{corr}$ Model:

1. Energy loss due to FSR
2. Recoil, spectator partons interactions and pileup contamination inside the electron reconstruction cone
3. Effects due to electronics noise subtraction and baseline subtraction (to subtract residue energy deposition from previous bunch crossings)



# W data

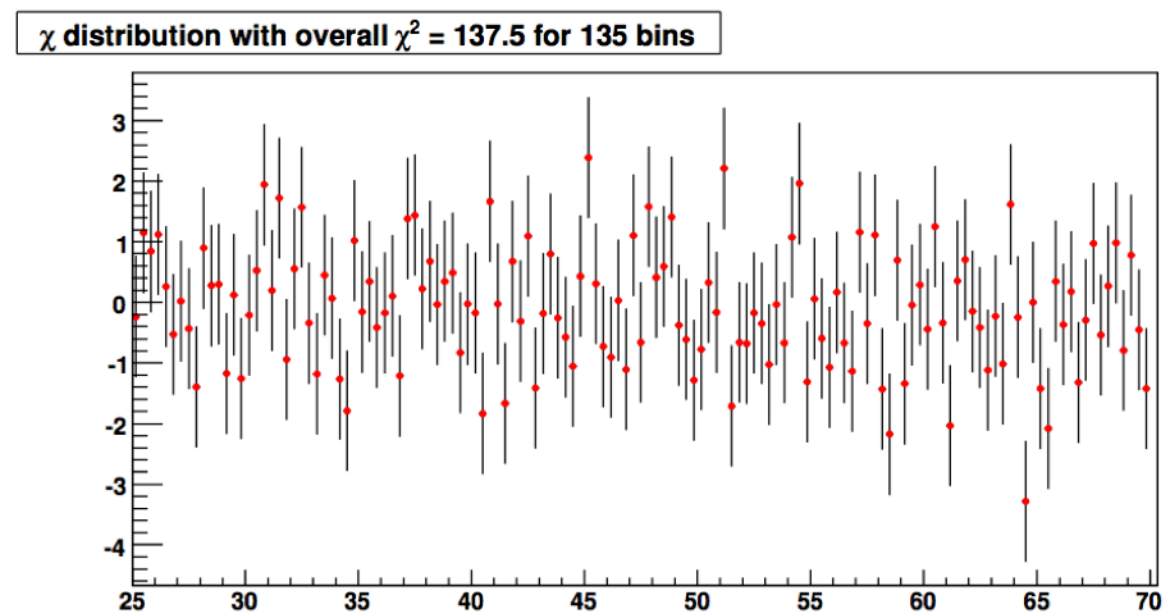
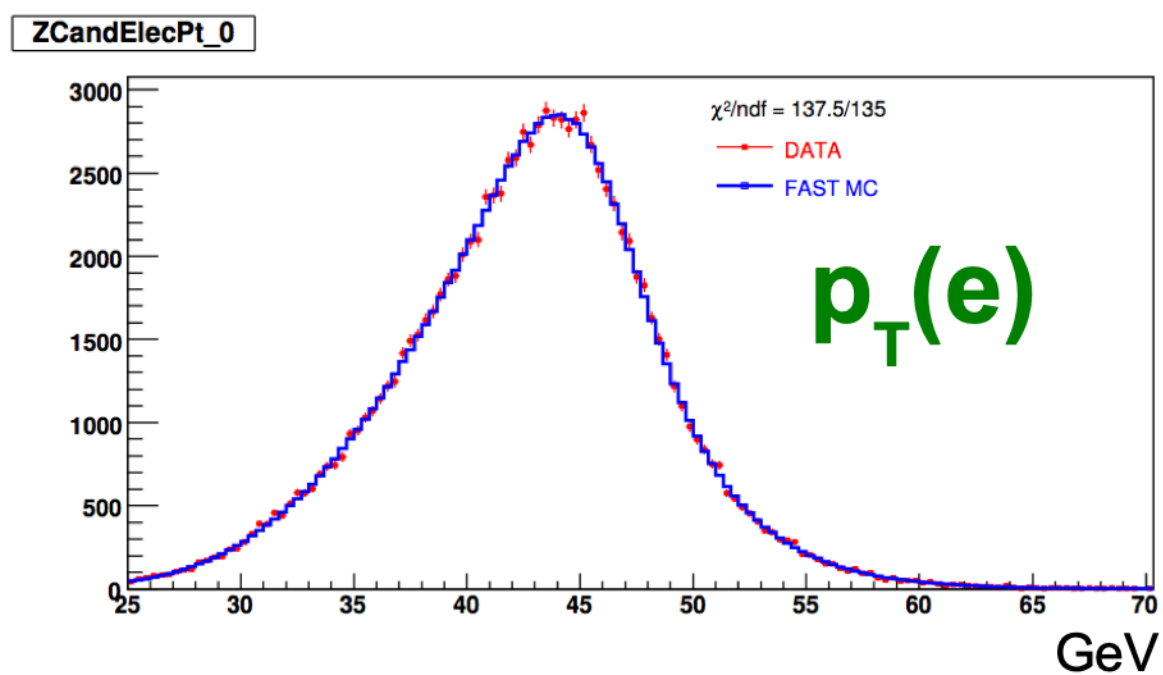
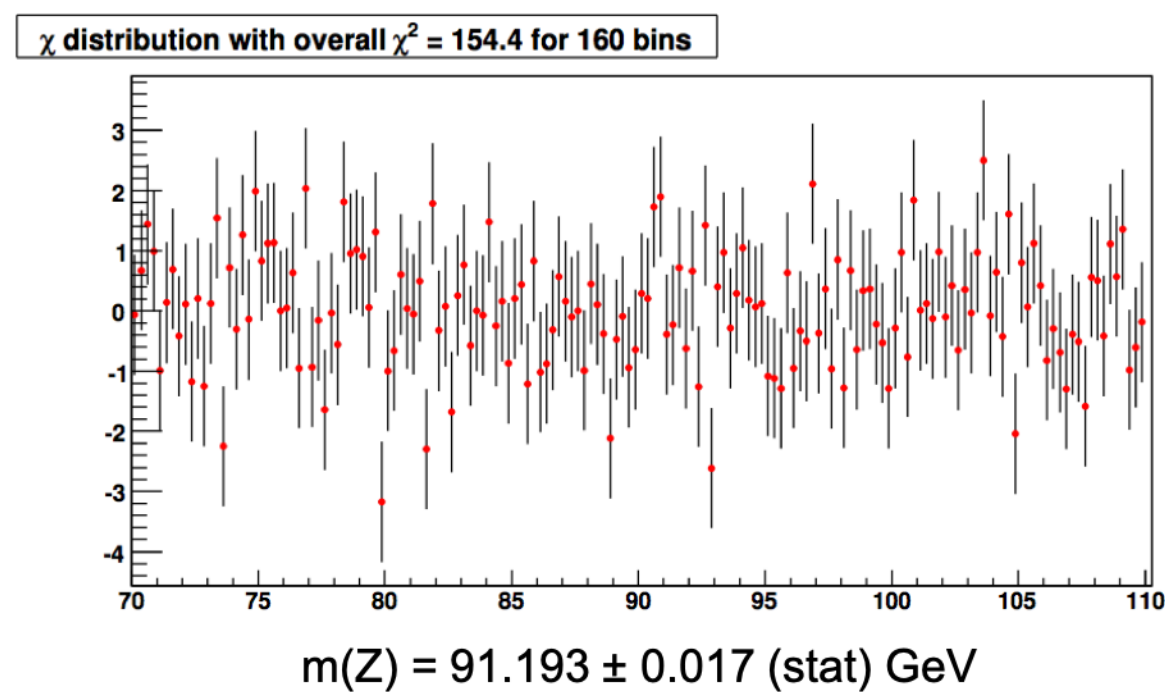
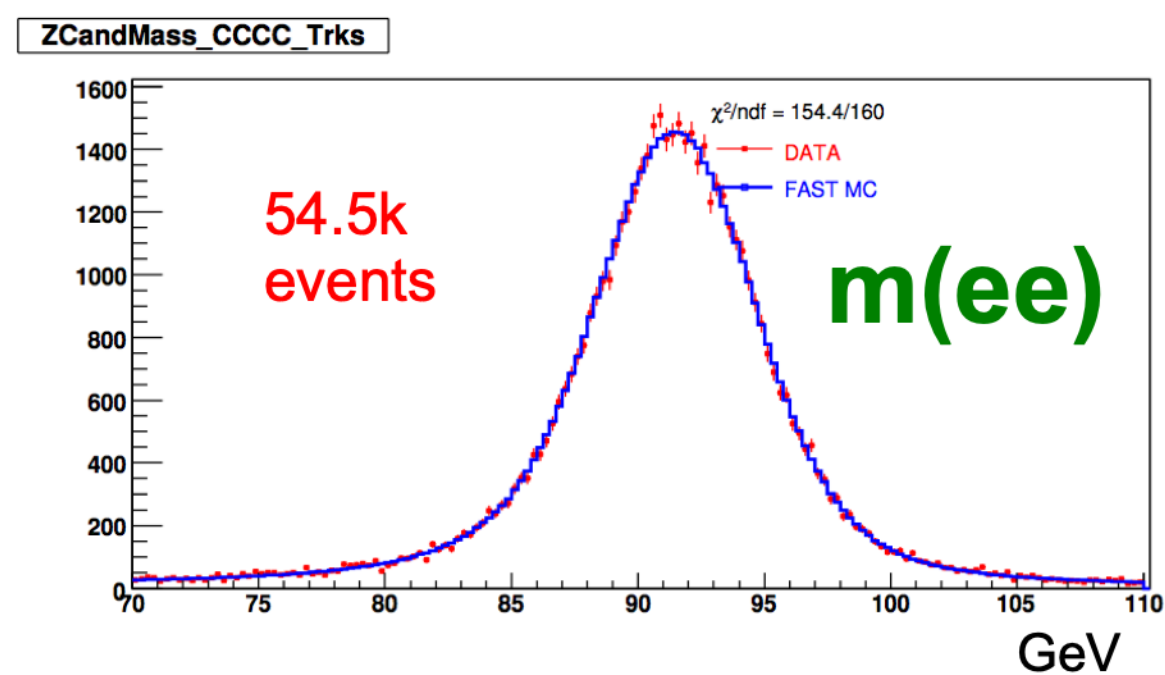


Here the error bars only reflect the finite statistics of the W candidate sample.

These are the same W candidates in the data. The blue band represents the uncertainties in the fast MC prediction due to the uncertainties in the recoil tune from the finite Z statistics.

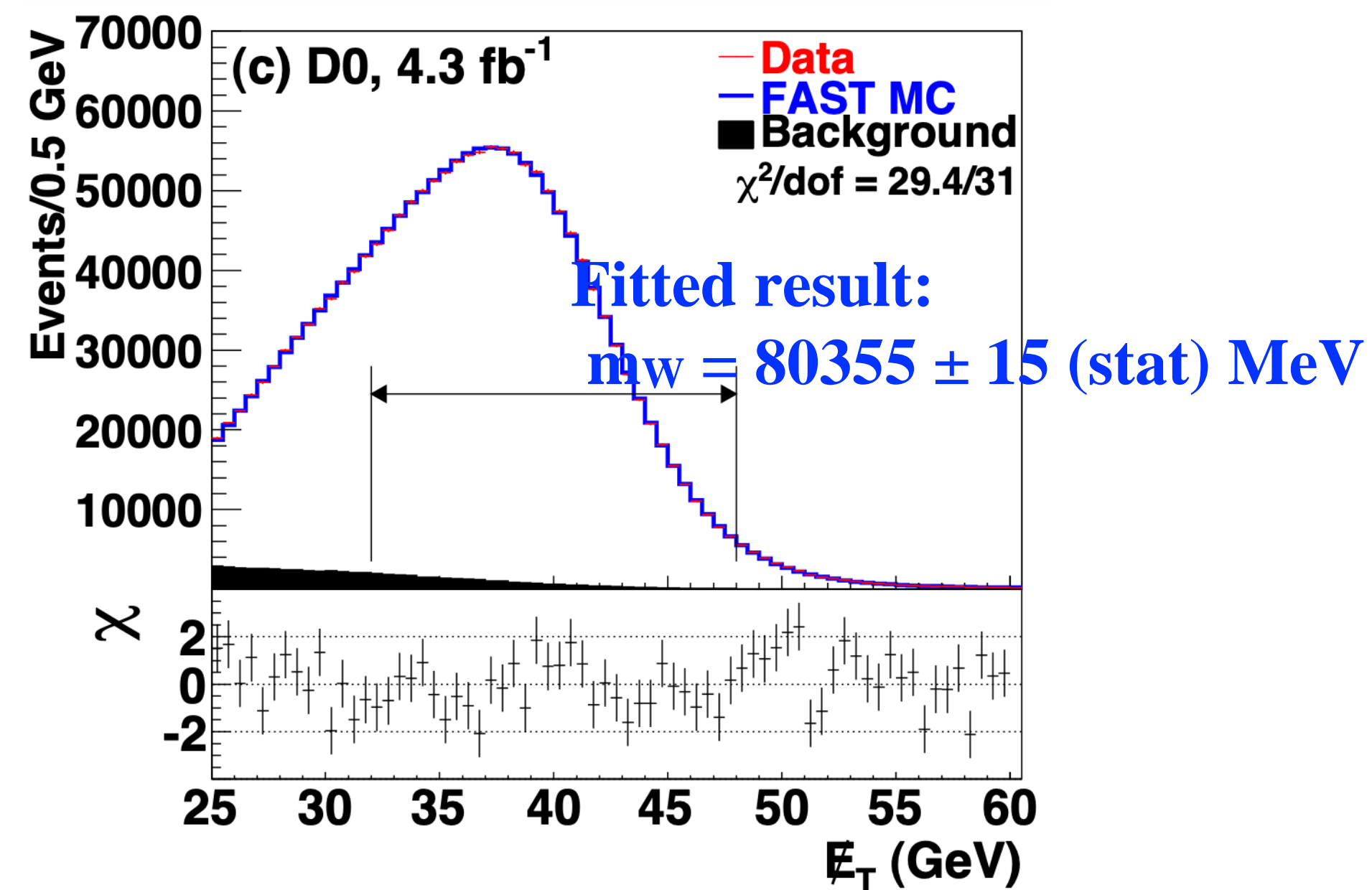
Good agreement between data and parameterised Monte Carlo.

# Z data

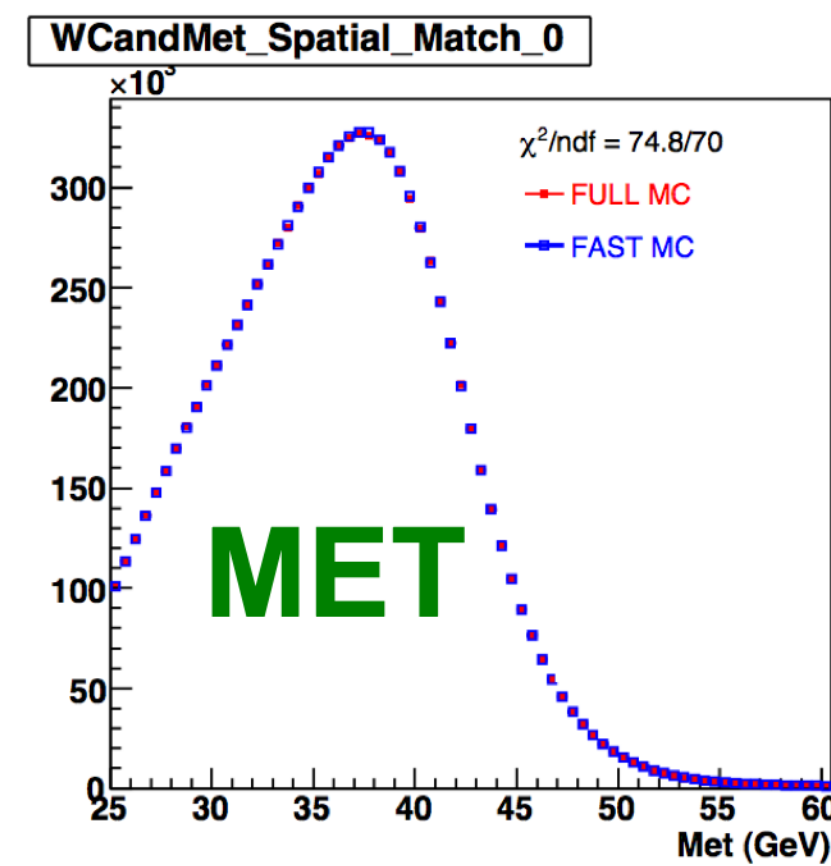
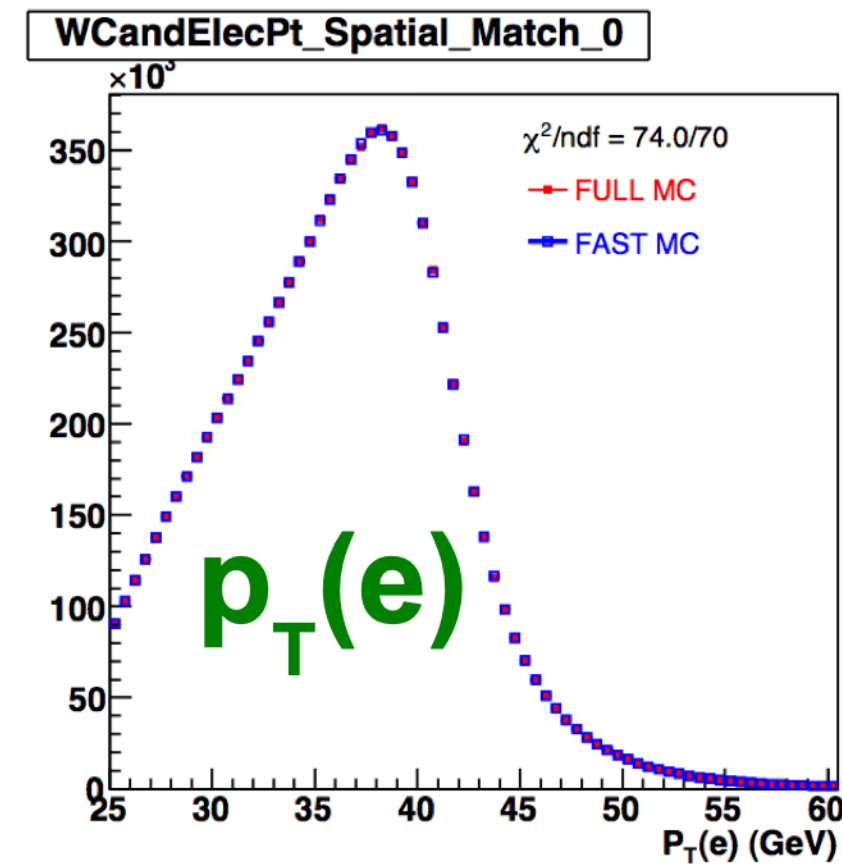
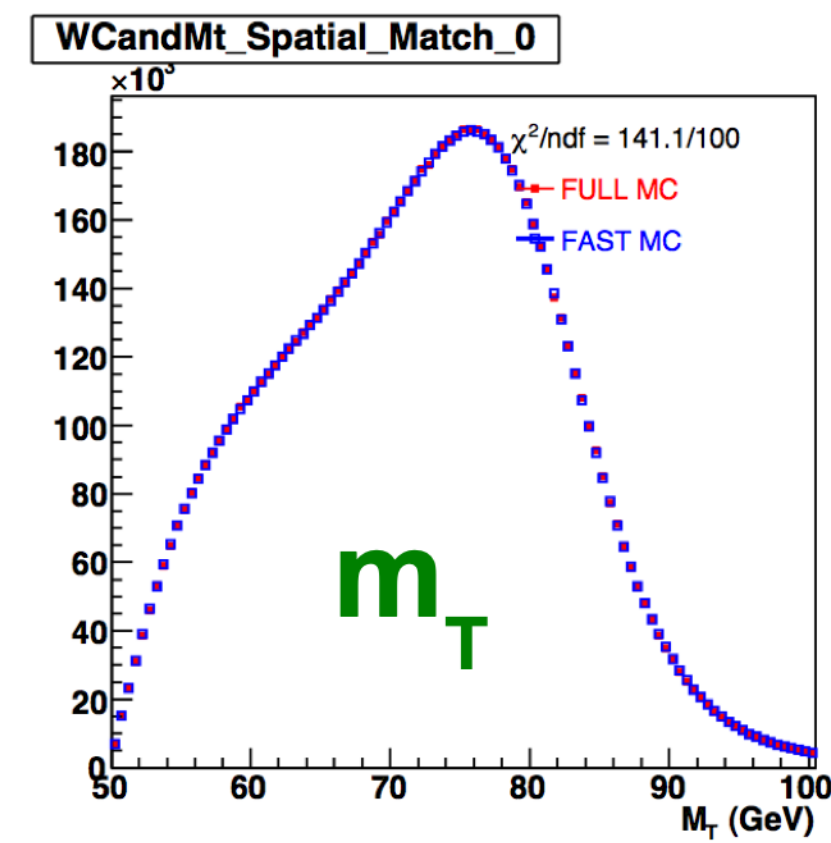


Good agreement between data and parameterised Monte Carlo.

# MET



# full MC closure test



9.8 M events  
after selection

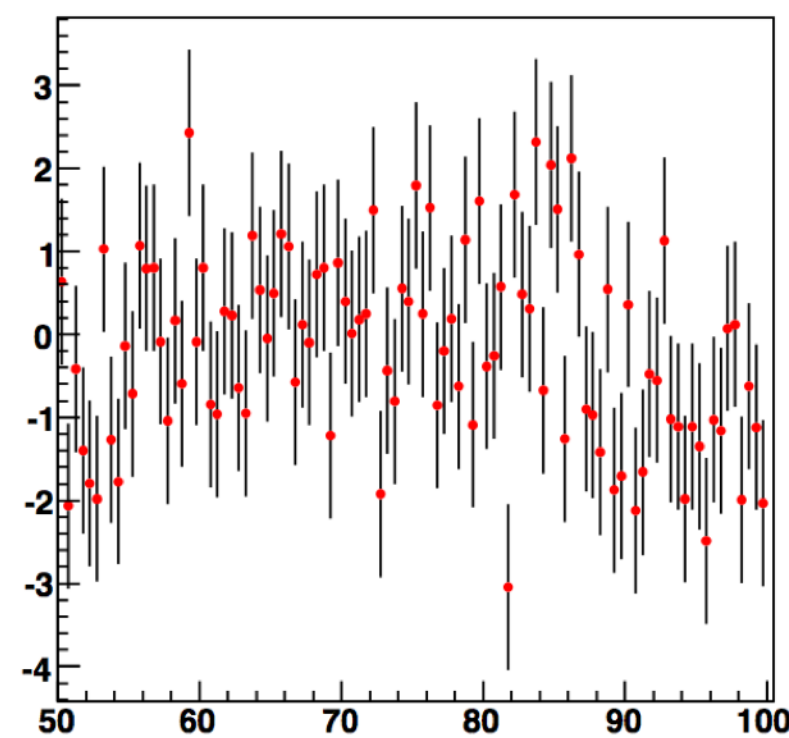
$m_T$ :  $80.448 \pm 0.005$  (stat) GeV

$p_T(e)$ :  $80.448 \pm 0.005$  (stat) GeV

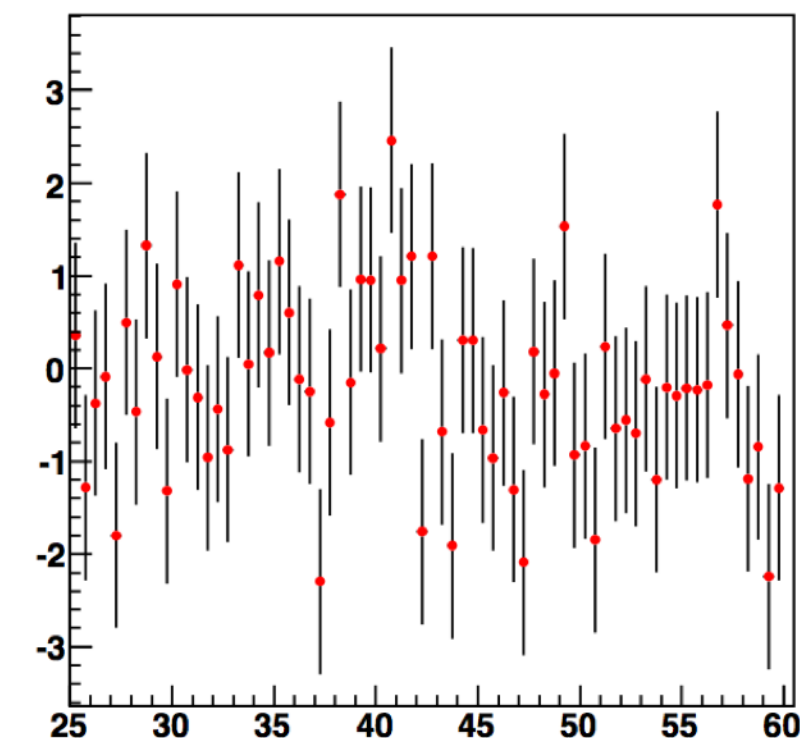
Met:  $80.455 \pm 0.006$  (stat) GeV

Input value: 80.450 GeV

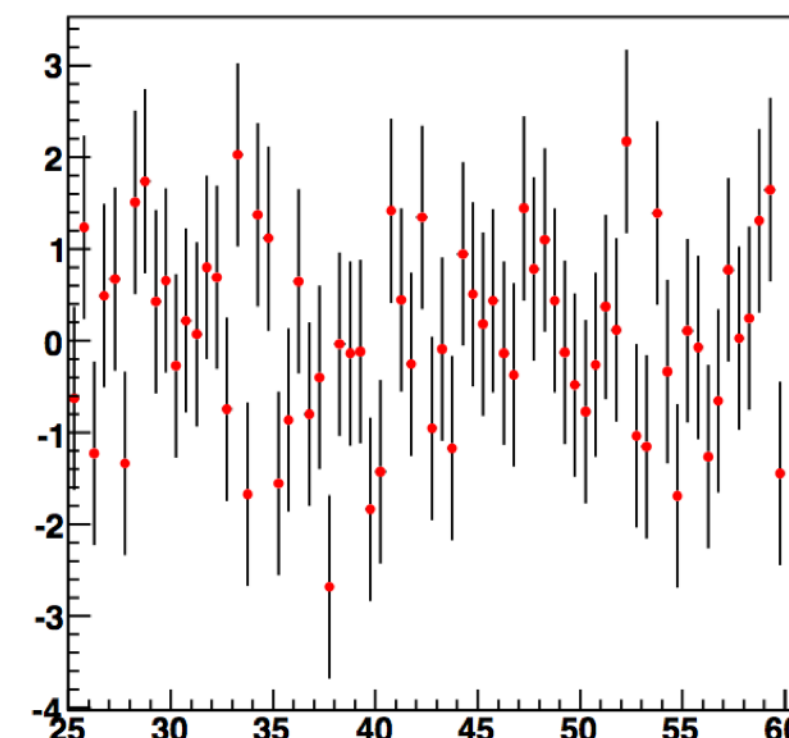
$\chi$  distribution with overall  $\chi^2 = 141.1$  for 100 bins



$\chi$  distribution with overall  $\chi^2 = 74.0$  for 70 bins



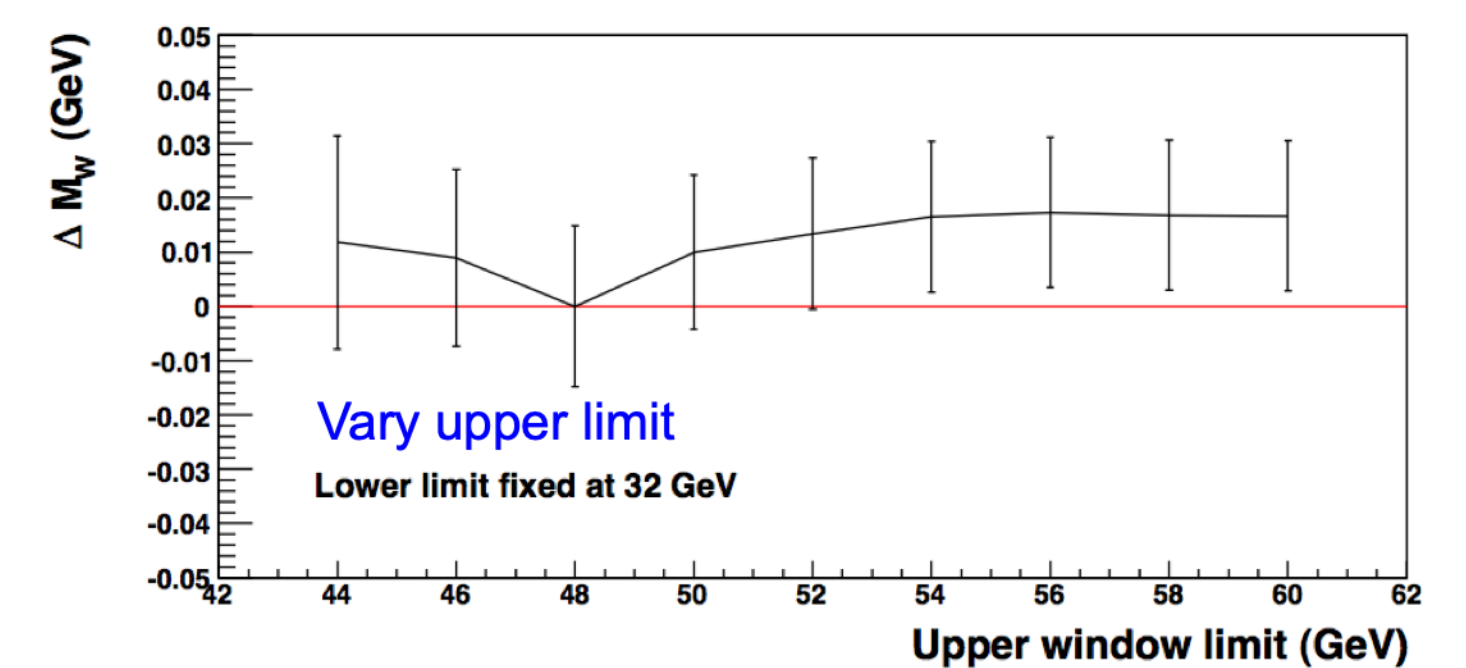
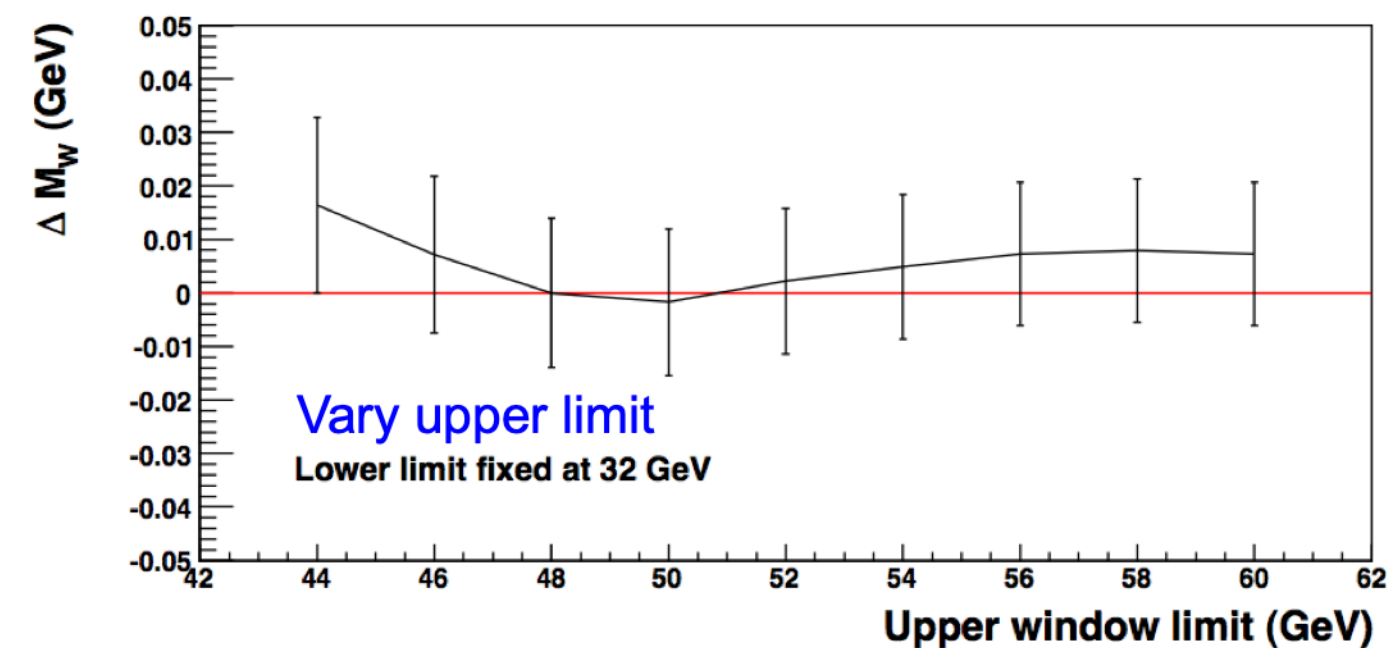
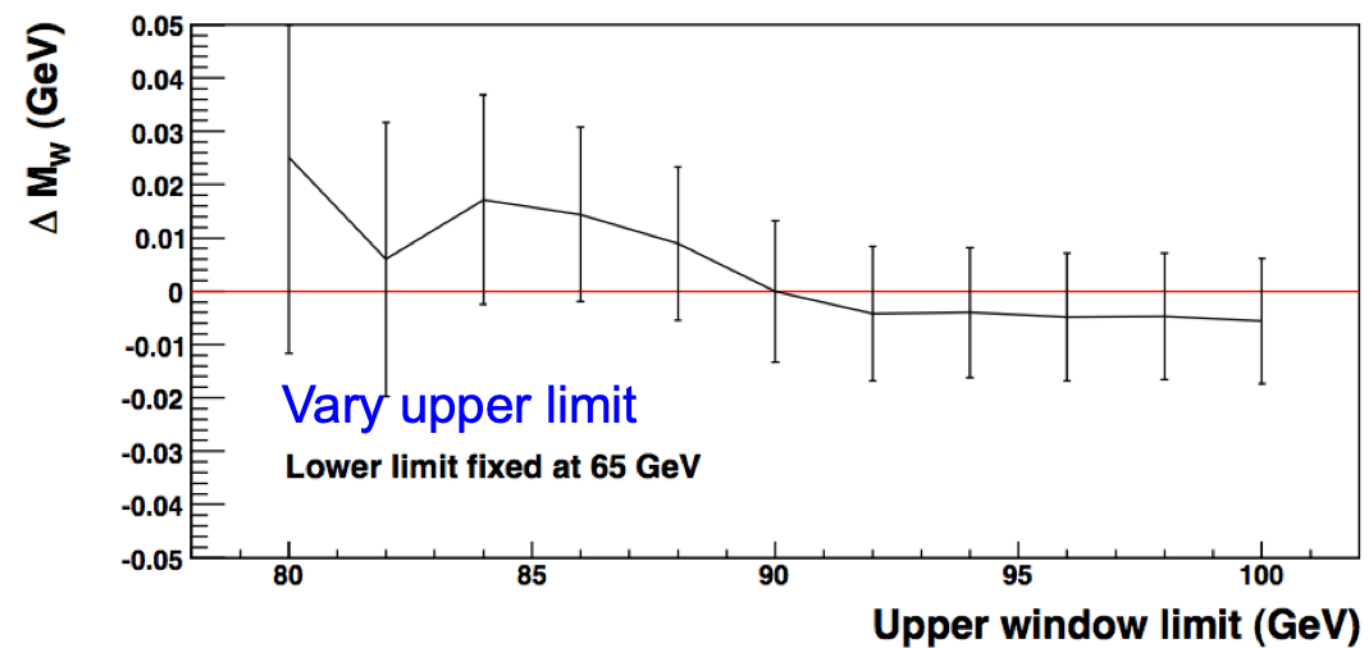
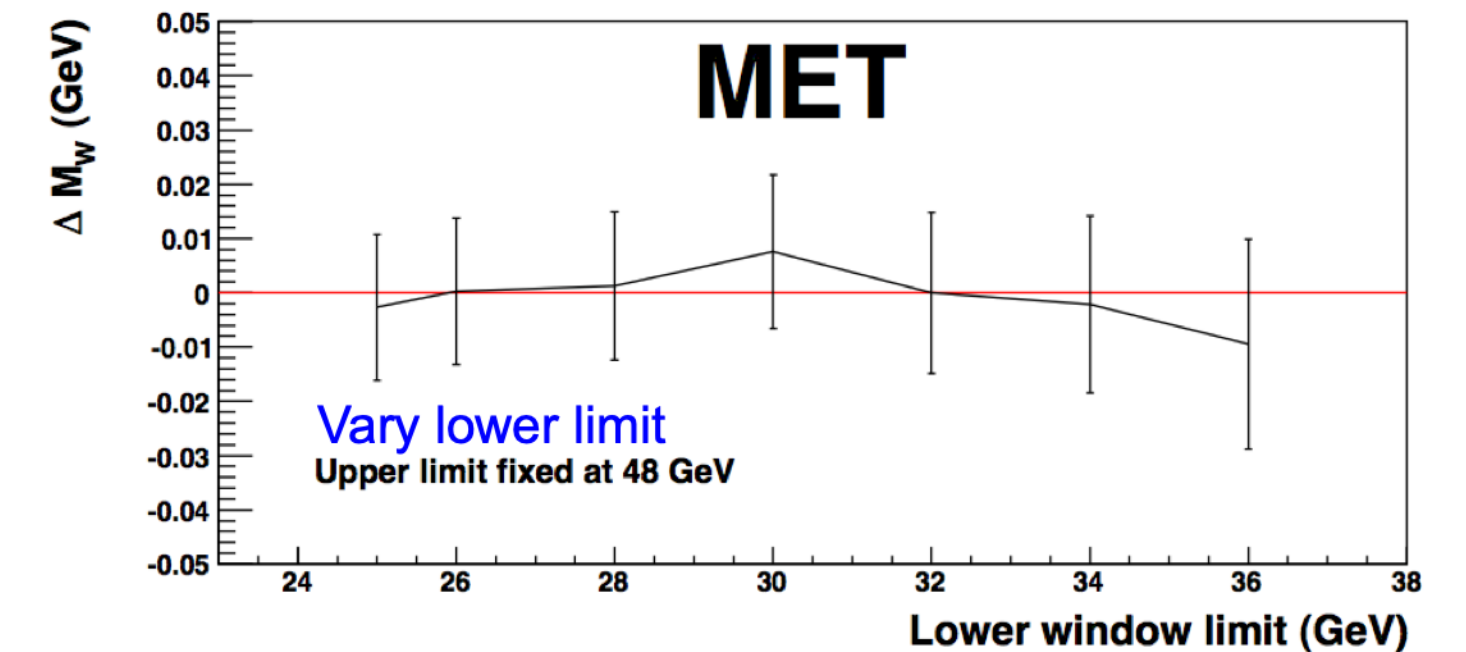
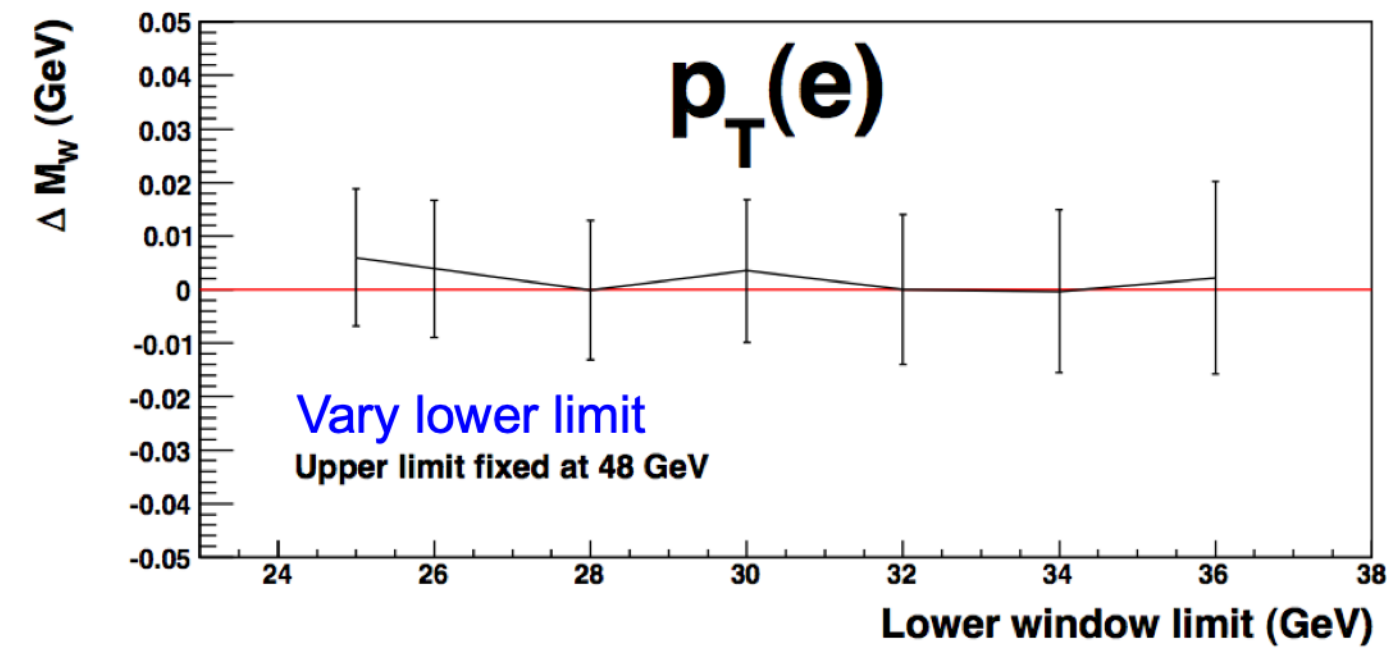
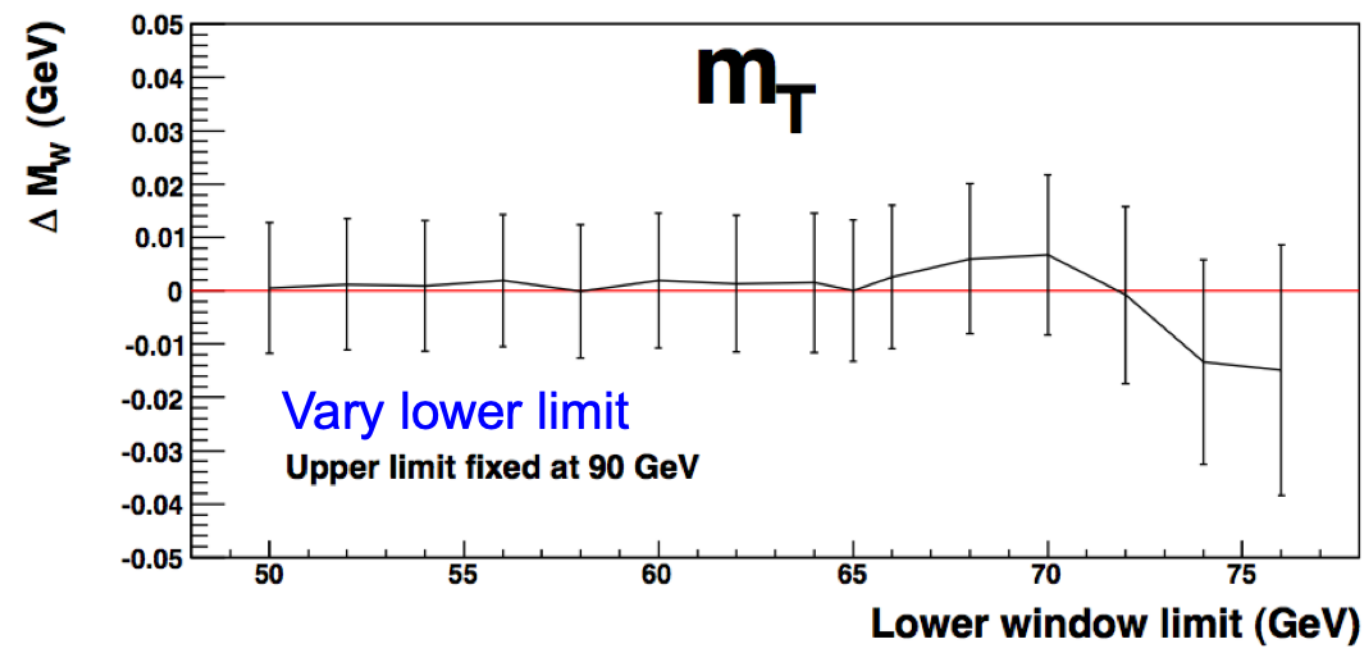
$\chi$  distribution with overall  $\chi^2 = 74.8$  for 70 bins





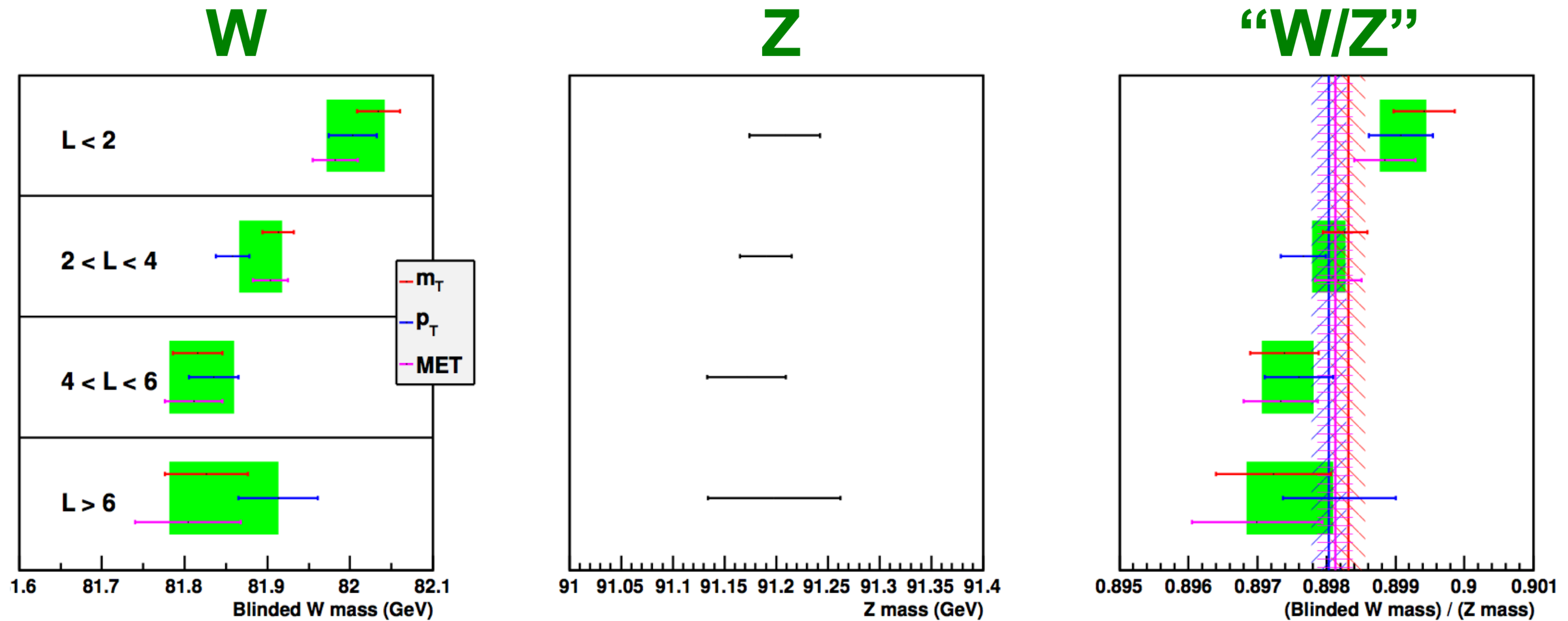
# Consistency Check:

## W mass fitting window check:



# Consistency Check:

Splitting Dataset according to Inst. Lumi.:

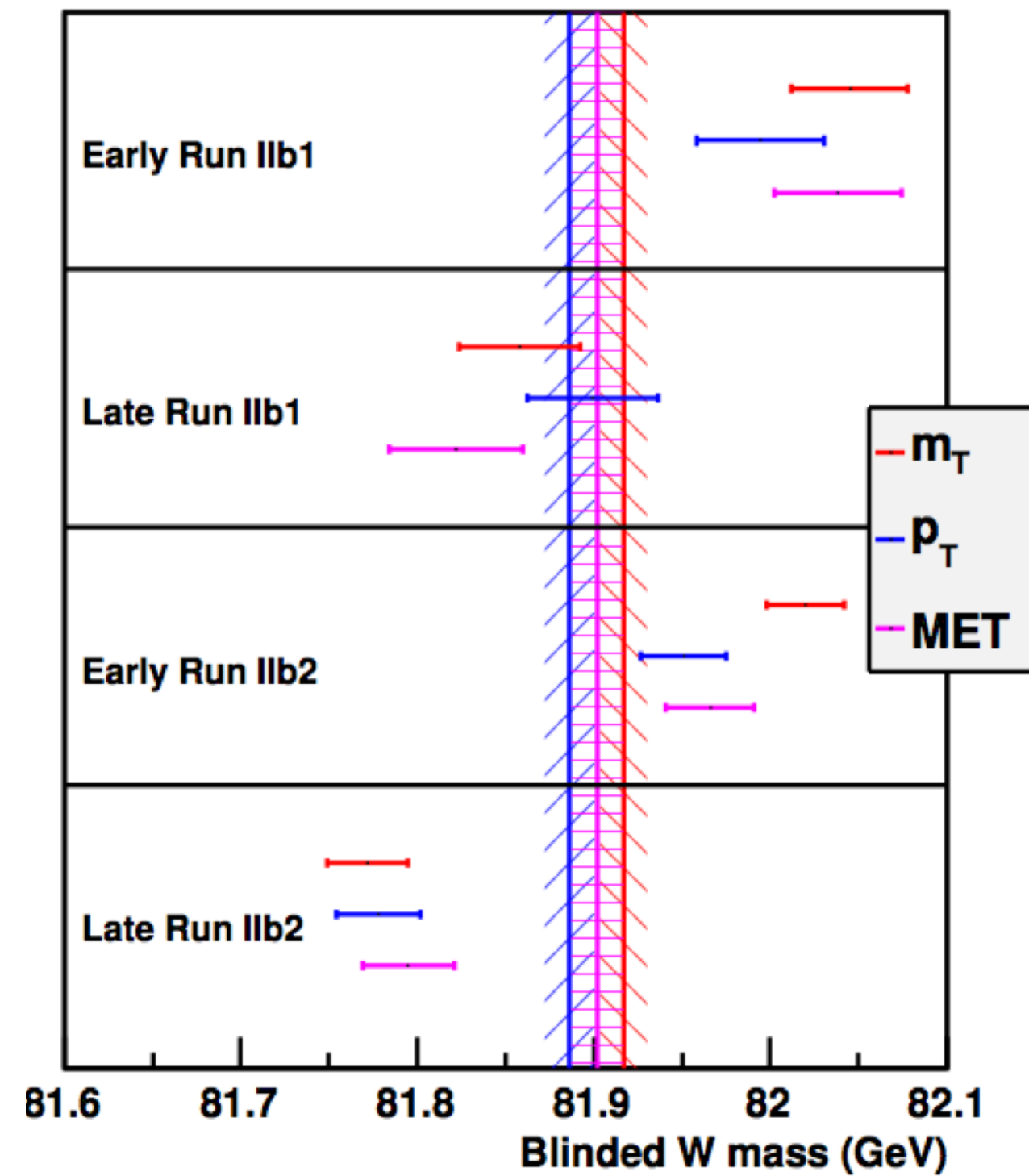


Green band is EM scale uncertainty.

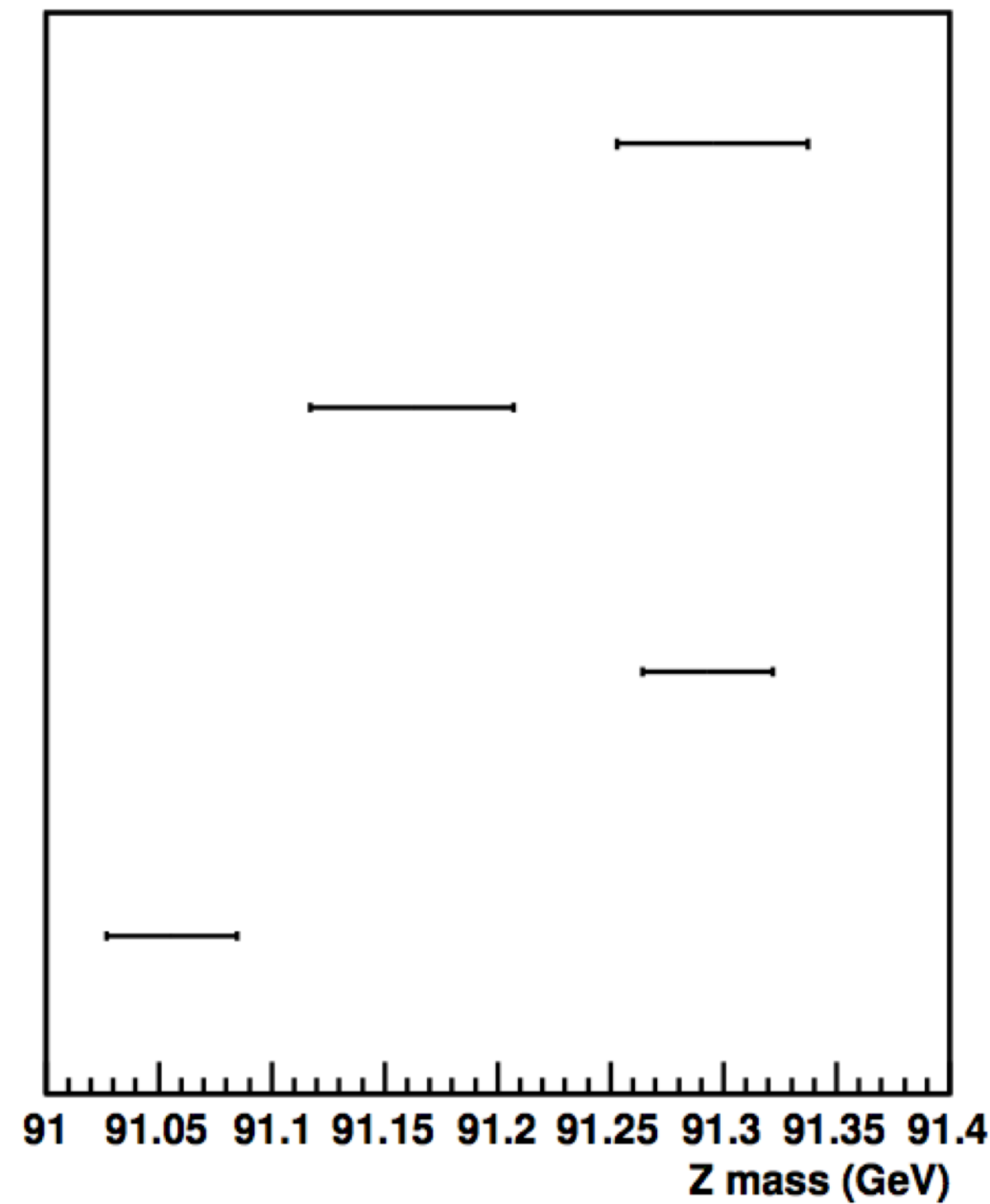
# Consistency Check

Splitting Dataset according to Run Periods:

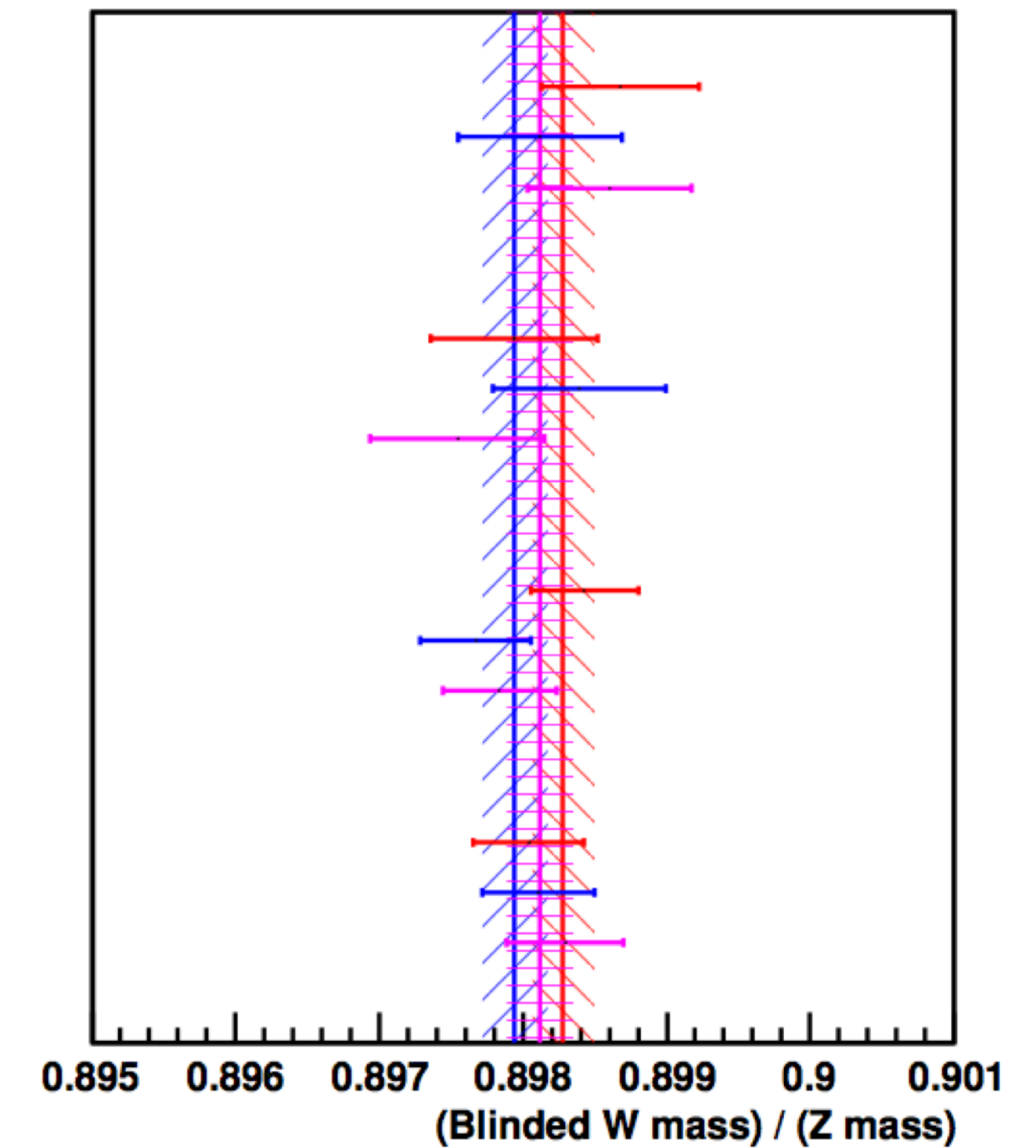
**W**



**Z**

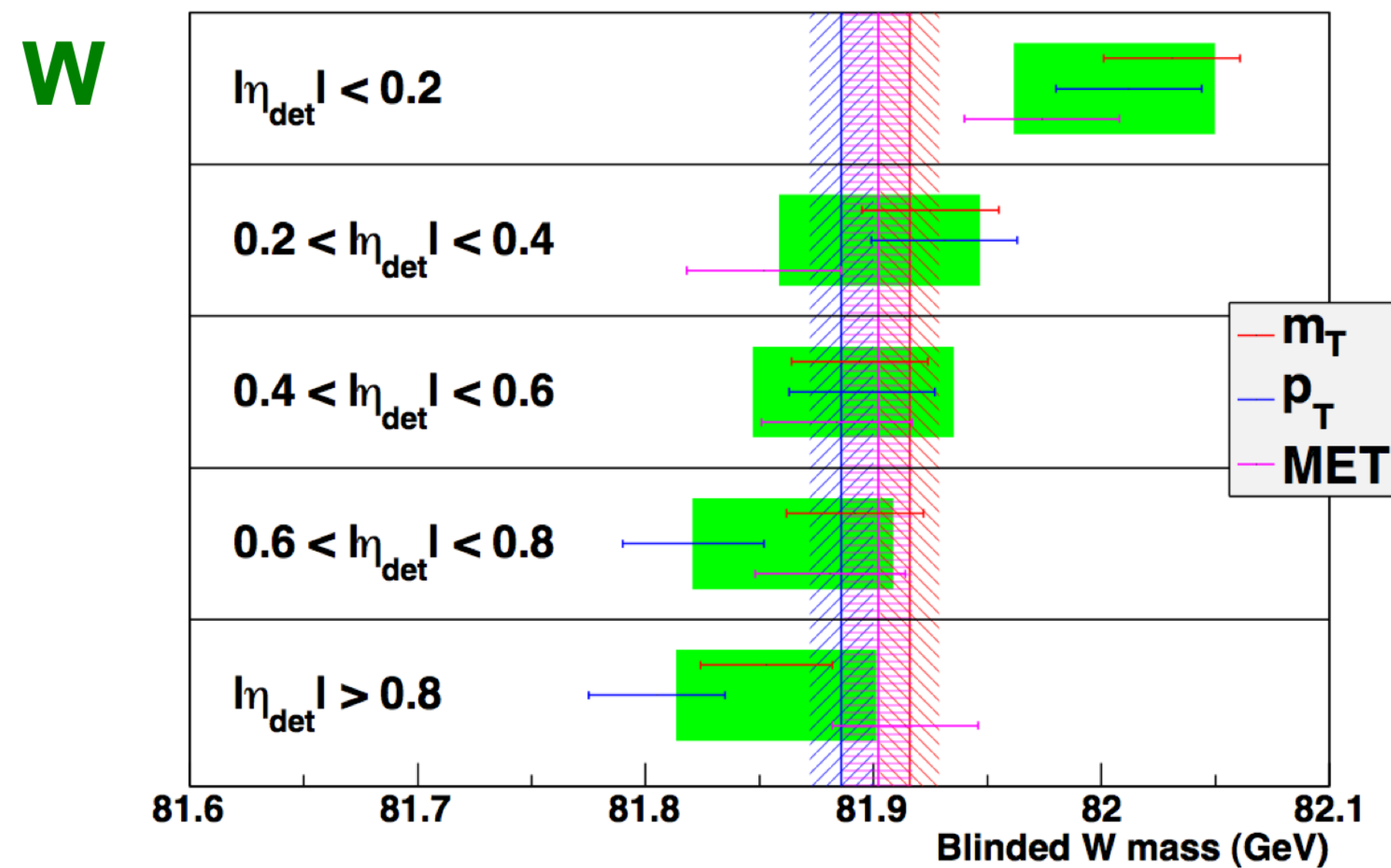


**“W/Z”**

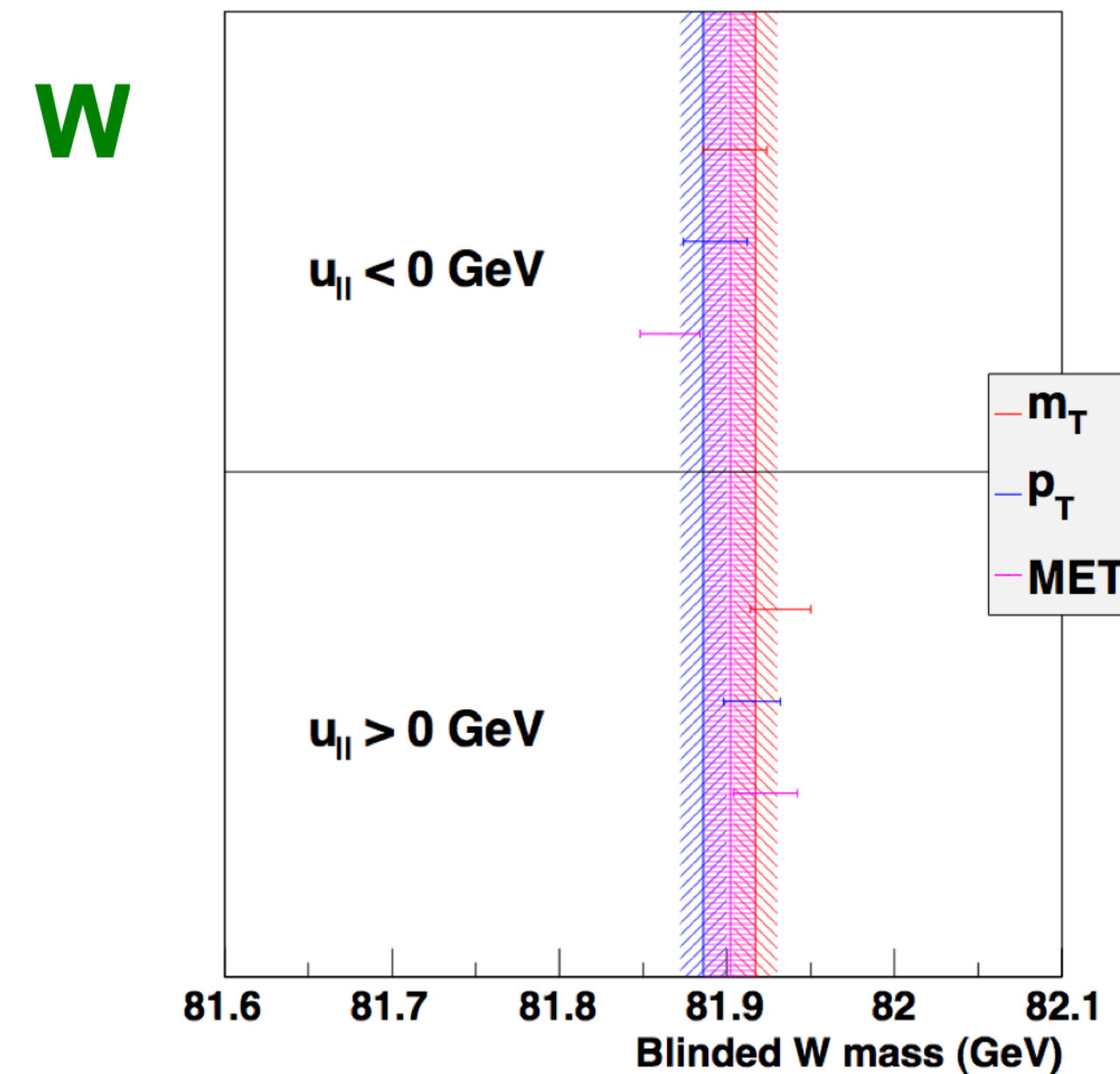


# Consistency Check:

Splitting Dataset according to detector Eta:



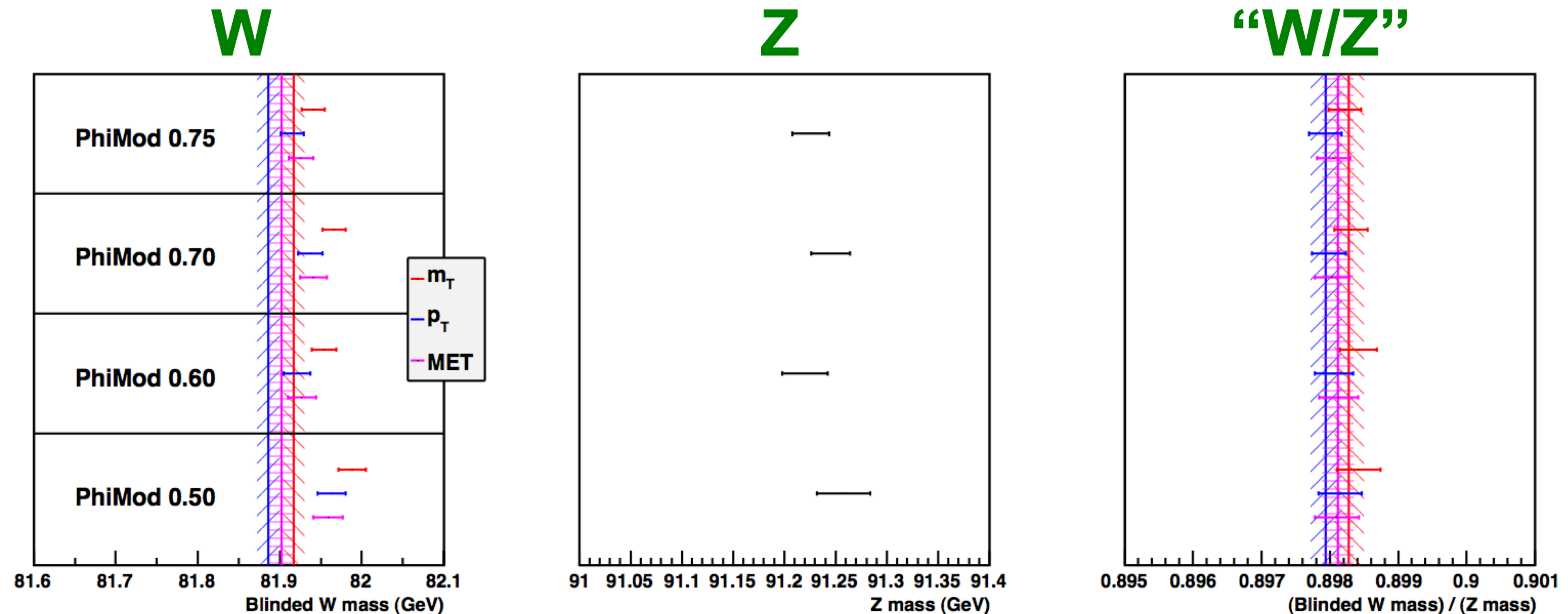
Splitting Dataset according to  $u_{||}$ :



$u_{||}$ : Hadronic Recoil vector  
projection to the electron's direction.

# Consistency Check:

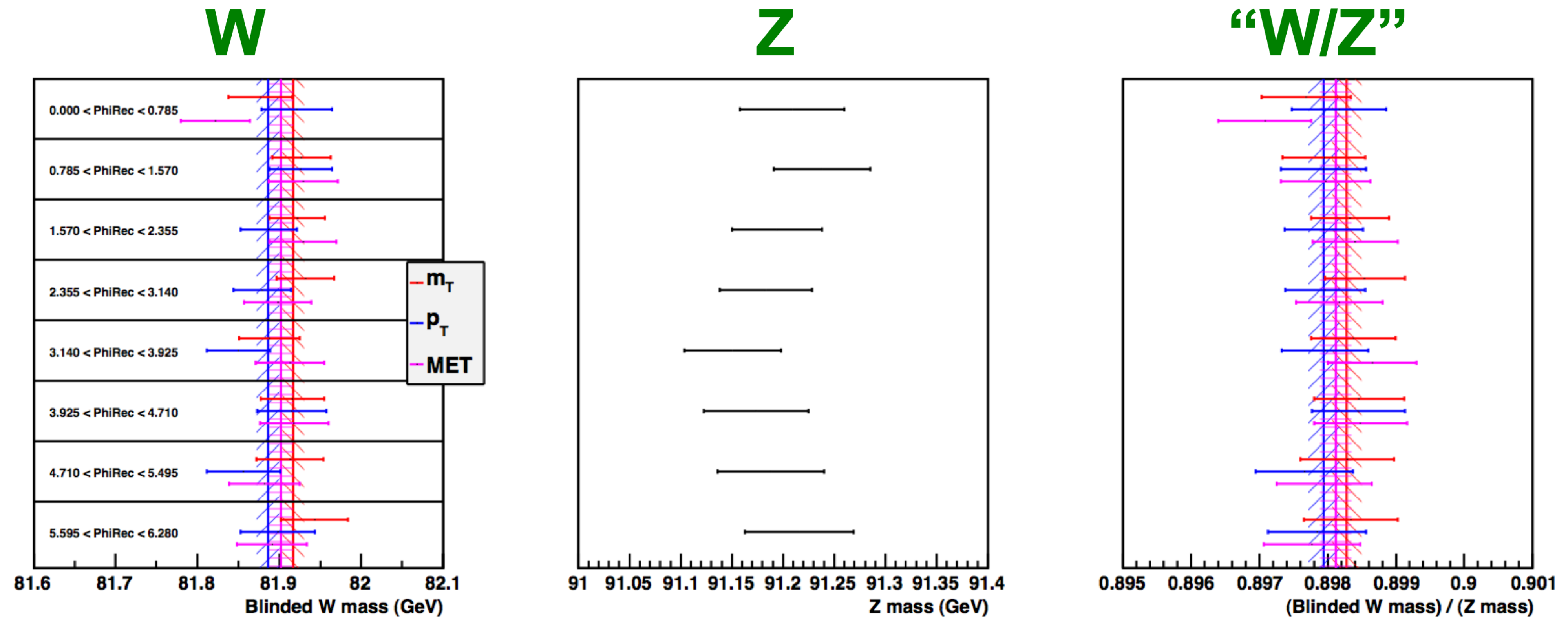
Splitting Dataset according to phi-mod:



An electron shooting to the edge of one calorimeter module in phi would be different in response than to the center of the module. Phi-Mod is the fractional position of one calorimeter module in phi.

# Consistency Check:

Splitting Dataset according to the Phi direction of the Hadronic Recoil vector:



# Combination

Results from the three observables are highly correlated:

$$\rho = \begin{pmatrix} \rho_{m_T m_T} & \rho_{m_T p_T^e} & \rho_{m_T \cancel{E}_T} \\ \rho_{m_T p_T^e} & \rho_{p_T^e p_T^e} & \rho_{p_T^e \cancel{E}_T} \\ \rho_{m_T \cancel{E}_T} & \rho_{p_T^e \cancel{E}_T} & \rho_{\cancel{E}_T \cancel{E}_T} \end{pmatrix} = \begin{pmatrix} 1.0 & 0.89 & 0.86 \\ 0.89 & 1.0 & 0.75 \\ 0.86 & 0.75 & 1.0 \end{pmatrix}$$

When we consider only the uncertainties which are allowed to decrease in the combination (e.g. not QED), we find that the MET measurement has negligible weight.

We therefore only retain  $p_T(e)$  and  $m_T$  for the combination:

**Run IIb 4.3 fb<sup>-1</sup> result:**

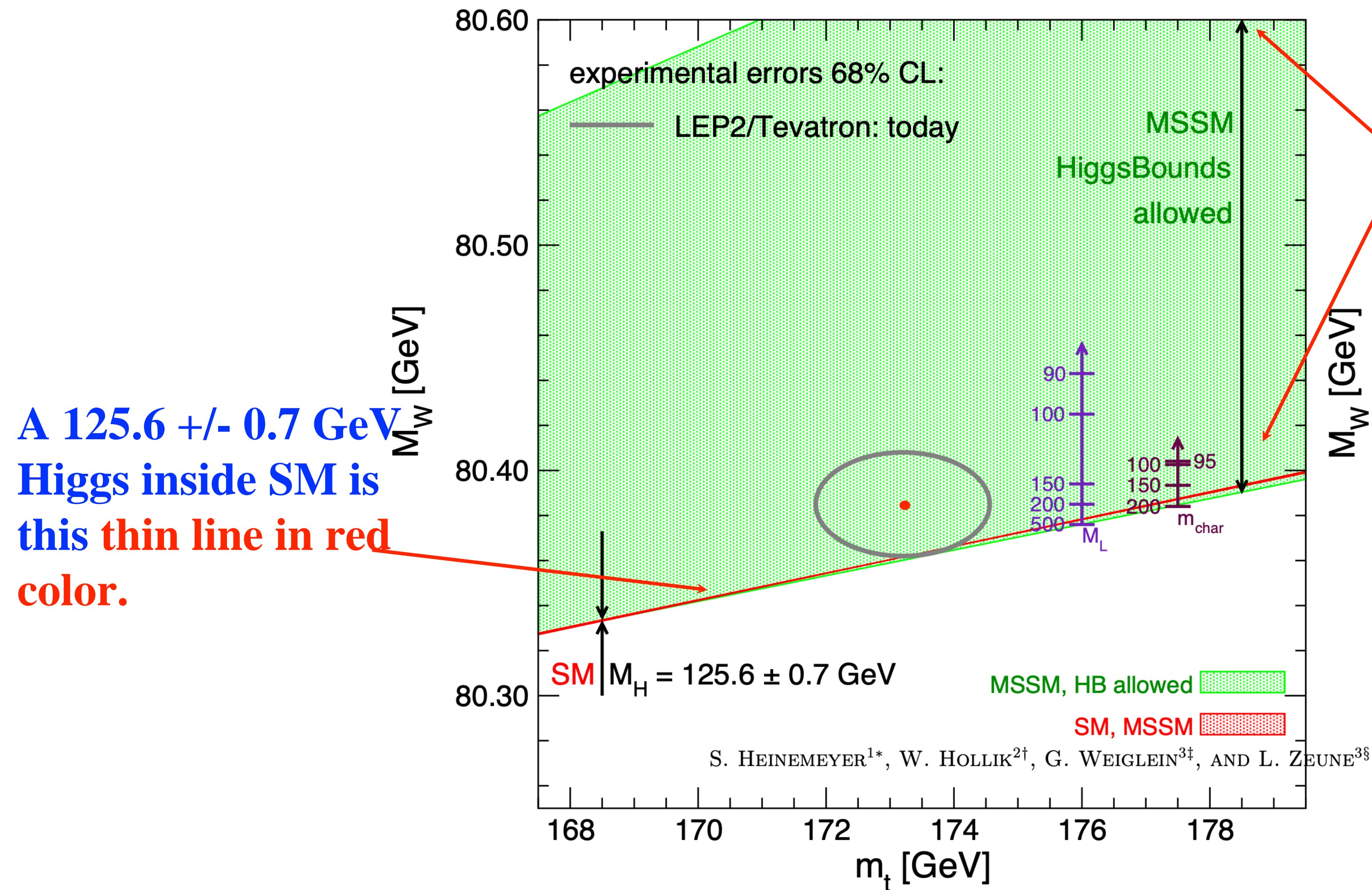
$$\begin{aligned} M_W &= 80.367 \pm 0.013 \text{ (stat)} \pm 0.022 \text{ (syst)} \text{ GeV} \\ &= 80.367 \pm 0.026 \text{ GeV.} \end{aligned}$$

Further combine with Run IIa 1 fb<sup>-1</sup> result, we obtain

**the new Run II 5.3 fb<sup>-1</sup> result:**

$$\begin{aligned} M_W &= 80.375 \pm 0.011 \text{ (stat)} \pm 0.020 \text{ (syst)} \text{ GeV} \\ &= 80.375 \pm 0.023 \text{ GeV.} \end{aligned}$$

# Help to understand the nature of the Higgs



A  $125.6 \pm 0.7$  GeV Higgs inside SM is this thin line in red color.

But a same  $125.6 \pm 0.7$  GeV Higgs inside MSSM is this wide band in green color!

- (1) The current W mass central value points to MSSM, not SM.
- (2) But precision is not good, not enough to distinguish SM or models beyond SM, such as MSSM.



