



Global analysis of QCD and SMEFT boosted with Machine Learnings

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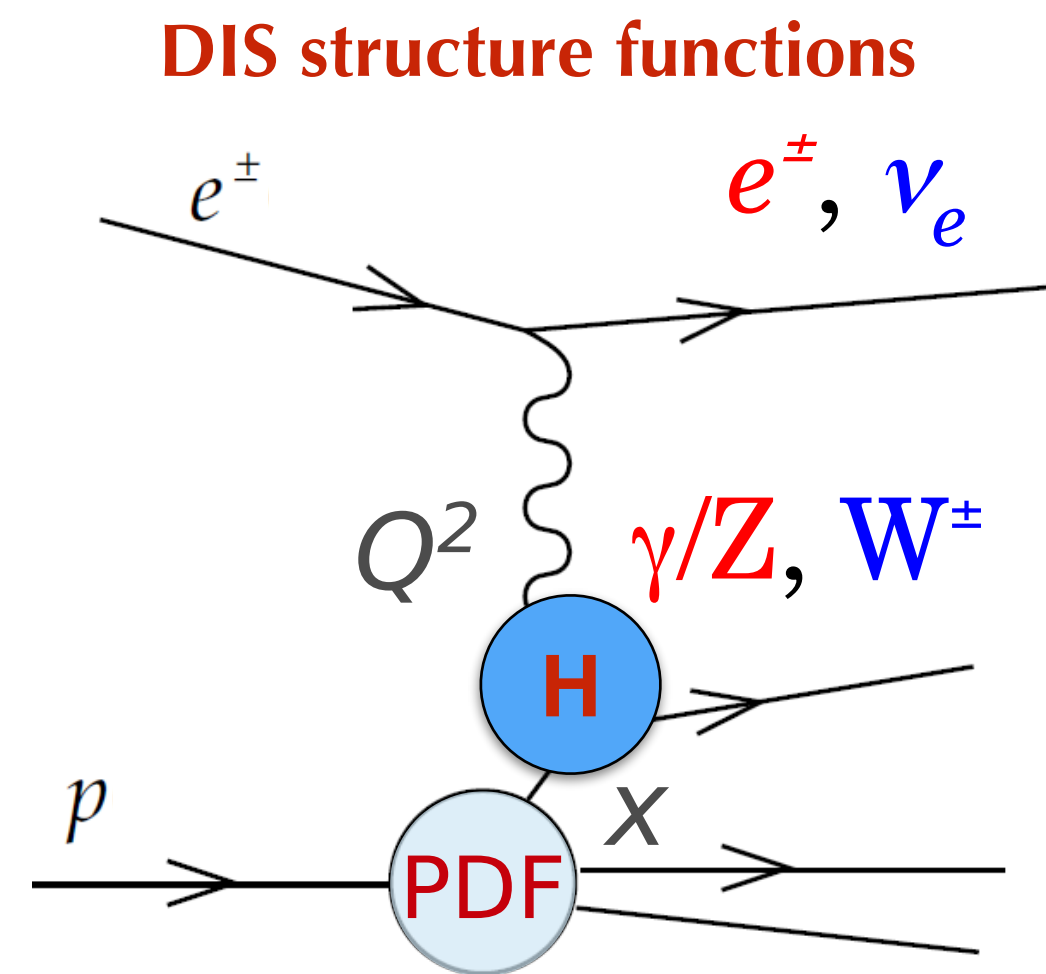


Outline

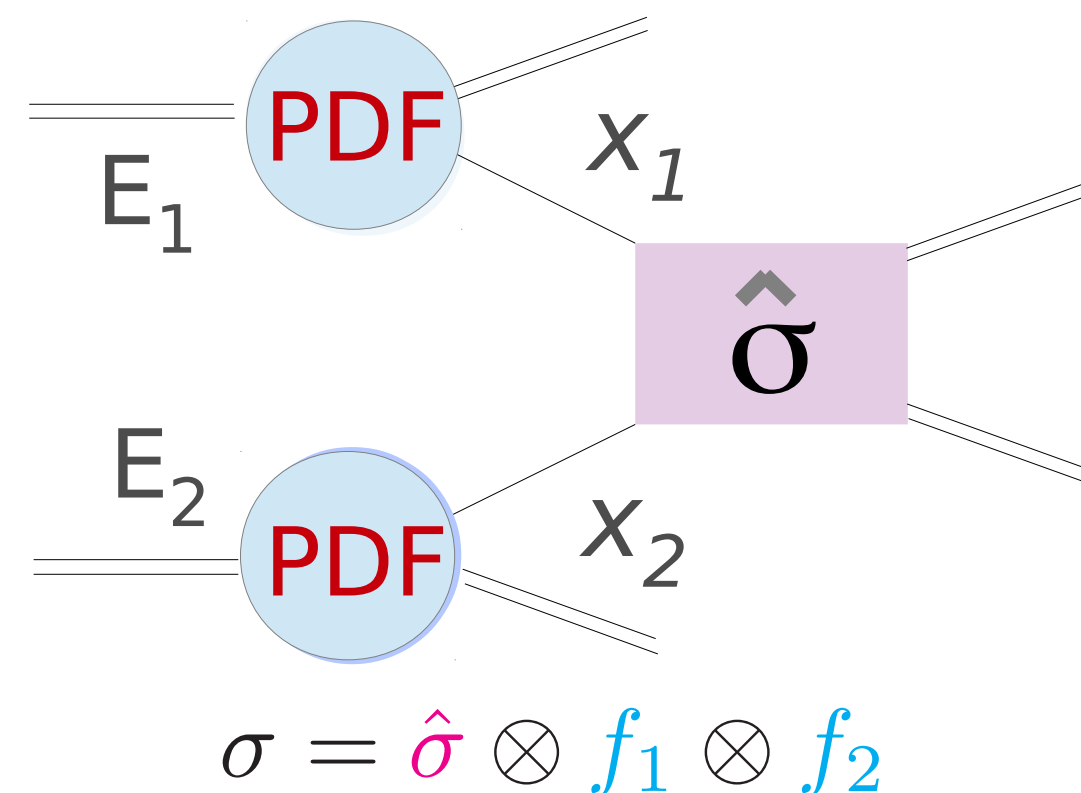
- ◆ 1. Introduction to PDFs for LHC
- ◆ 2. A framework of Global analysis boosted with machine learnings and applications
 - Understanding PDF uncertainties in W-mass direct measurements
 - Implications of PDF for searches of new physics at the LHC
- ◆ 3. Summary

QCD collinear factorization

- QCD collinear factorization ensures universal separation of long-distance and short-distance contributions in high energy scatterings involving initial state hadrons, and enables predictions on cross sections



hadron-hadron collision



$$F_2(x, Q^2) = \sum_{i=q, \bar{q}, g} \int_0^1 d\xi C_2^i(x/\xi, Q^2/\mu_r^2, \mu_f^2/\mu_r^2, \alpha_s(\mu_r^2)) \times f_{i/h}(\xi, \mu_f)$$

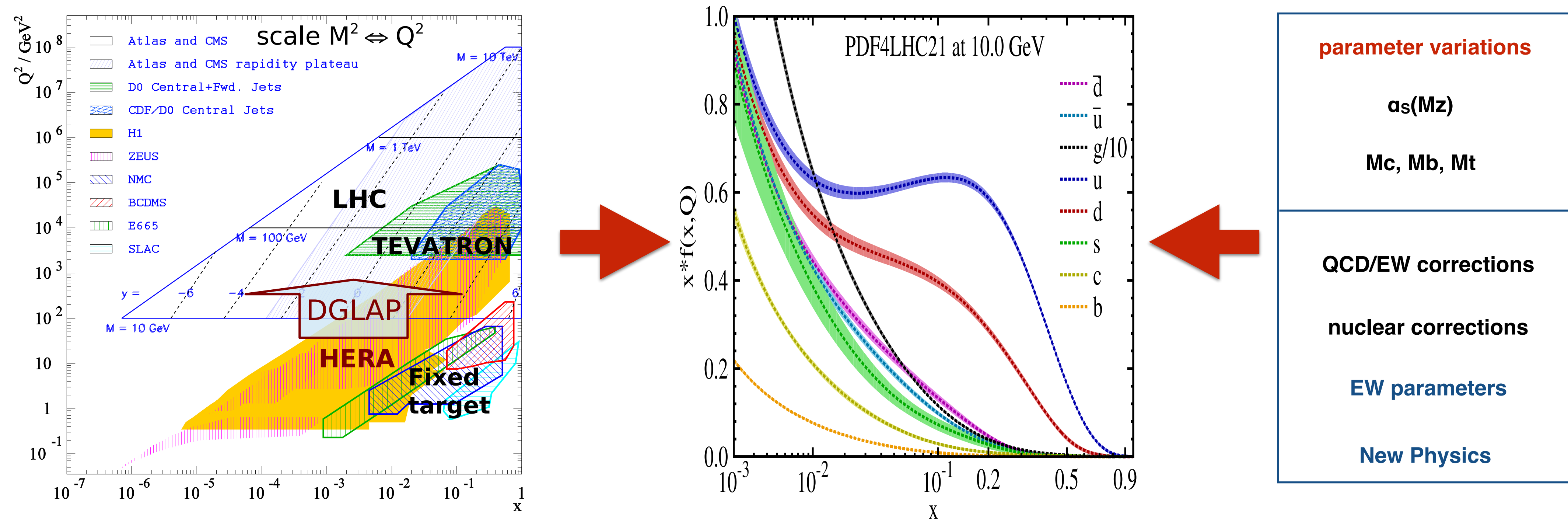
[Collins, Soper, Sterman, 1989]

- coefficient functions, hard scattering; infrared (IR) safe, calculable in pQCD, independent of the hadron
- PDFs, reveal inner structure of hadrons; non-perturbative (NP) origin, universality, e.g. DIS vs. pp collisions
- factorization scale μ_f
- runnings of $f_{i/h}$ with μ_f are governed by the DGLAP equation

choose $\mu_f = \mu_r = Q$, thus Q dependence (scaling violation) of F_2 are mostly from PDFs and thus are predicted by the DGLAP evolution

Global analysis of PDFs

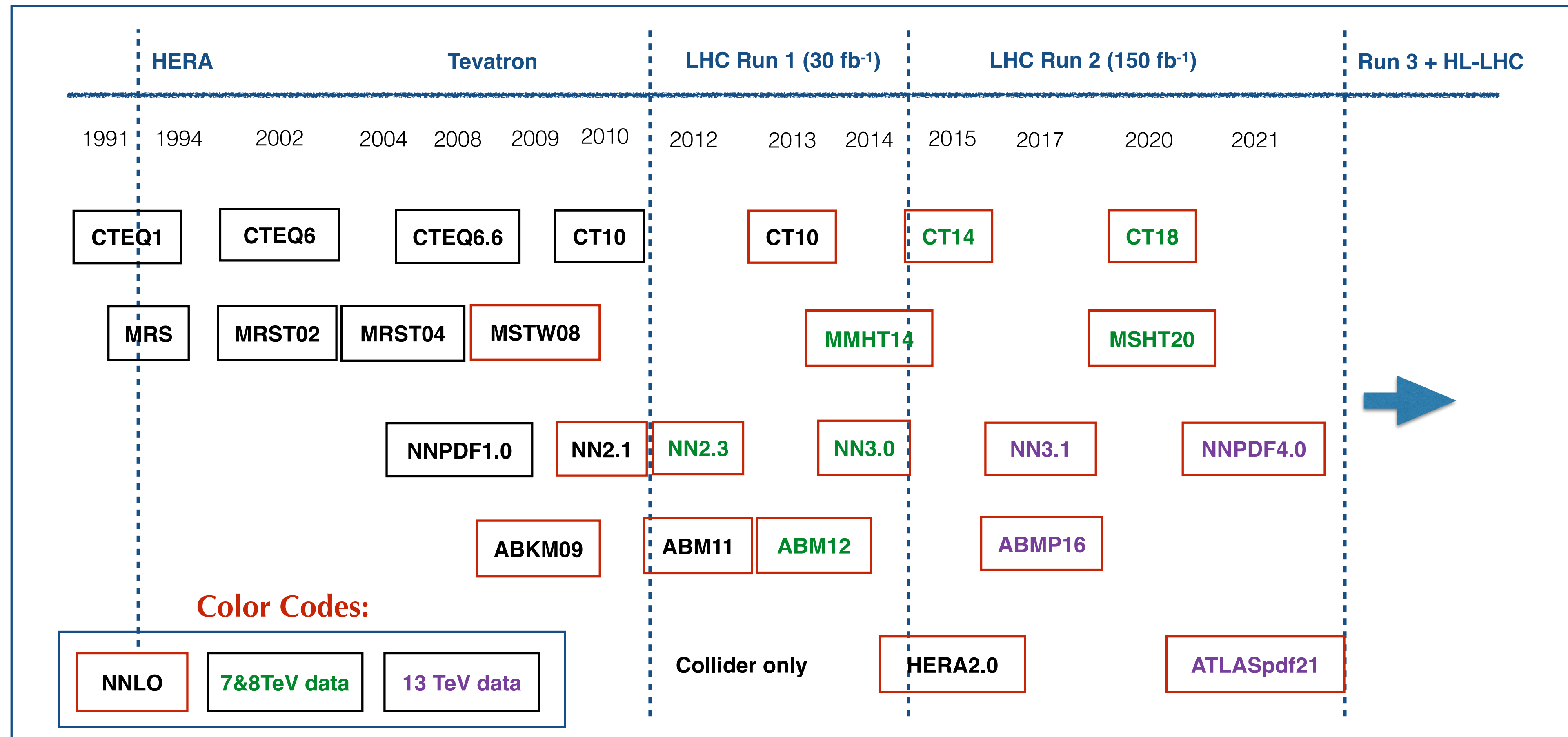
- PDFs are usually extracted from global analysis on variety of data, e.g., DIS, Drell-Yan, jets and top quark productions at fixed-target and collider experiments, with increasing weight from LHC, together with SM QCD parameters
[see JG, Harland-Lang, Rojo 1709.04922 for review article]



- diversity of the analysed data are important to ensure flavor separation and to avoid theoretical/experimental bias; extensions to include EW parameters and possible new physics for a self-consistent determination
- alternative approach from lattice QCD simulations, for various PDF moments or PDFs directly calculated in x -space with large momentum effective theory or pseudo-PDFs [2004.03543]

Major analysis groups

- PDFs provided by several major analysis groups (CT, MSHT, NNPDF, ABM, HERAPDF, ATLASpdf, CJ, JAM...) using slightly different heavy-quark schemes, selections of data, and methodologies

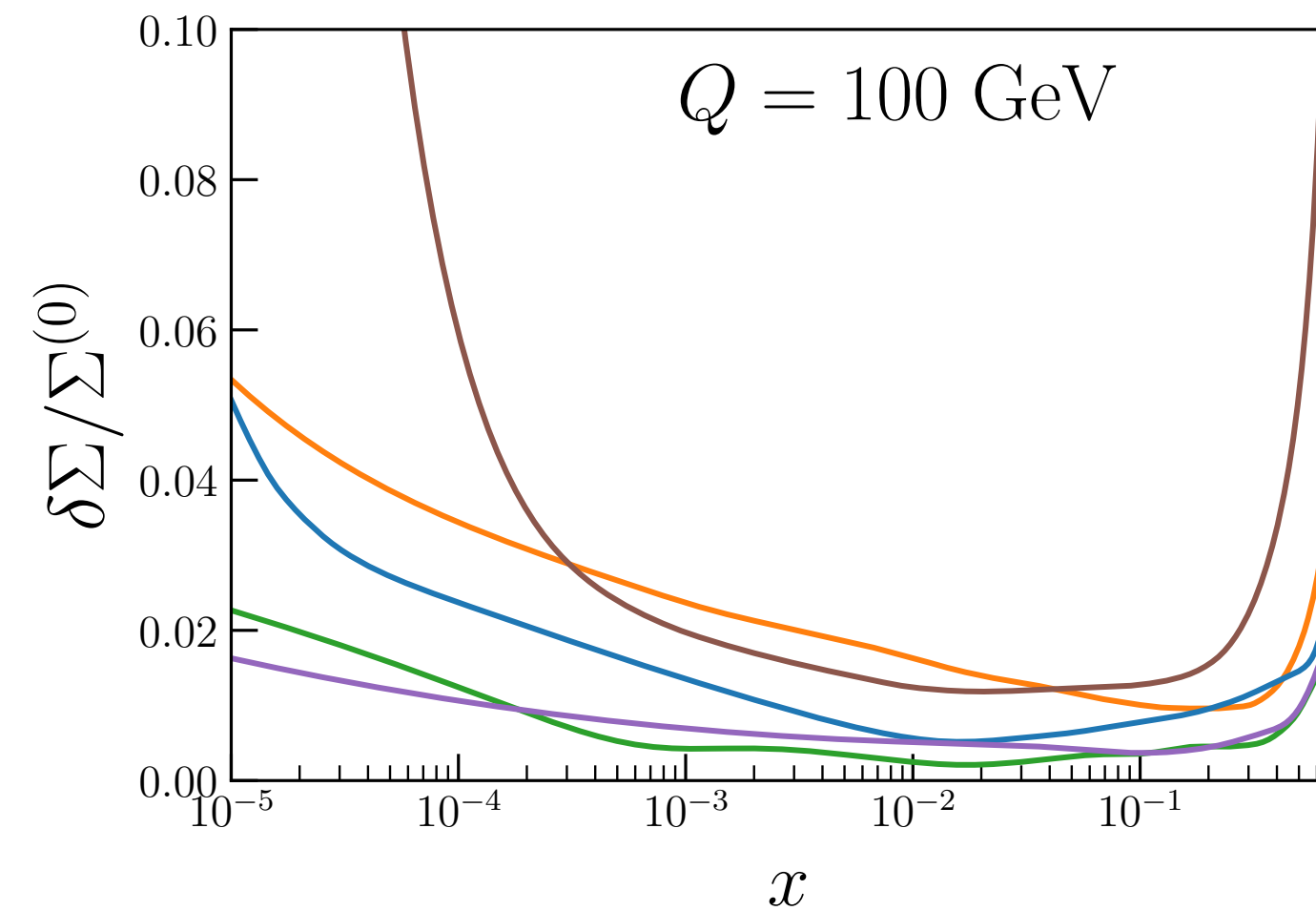
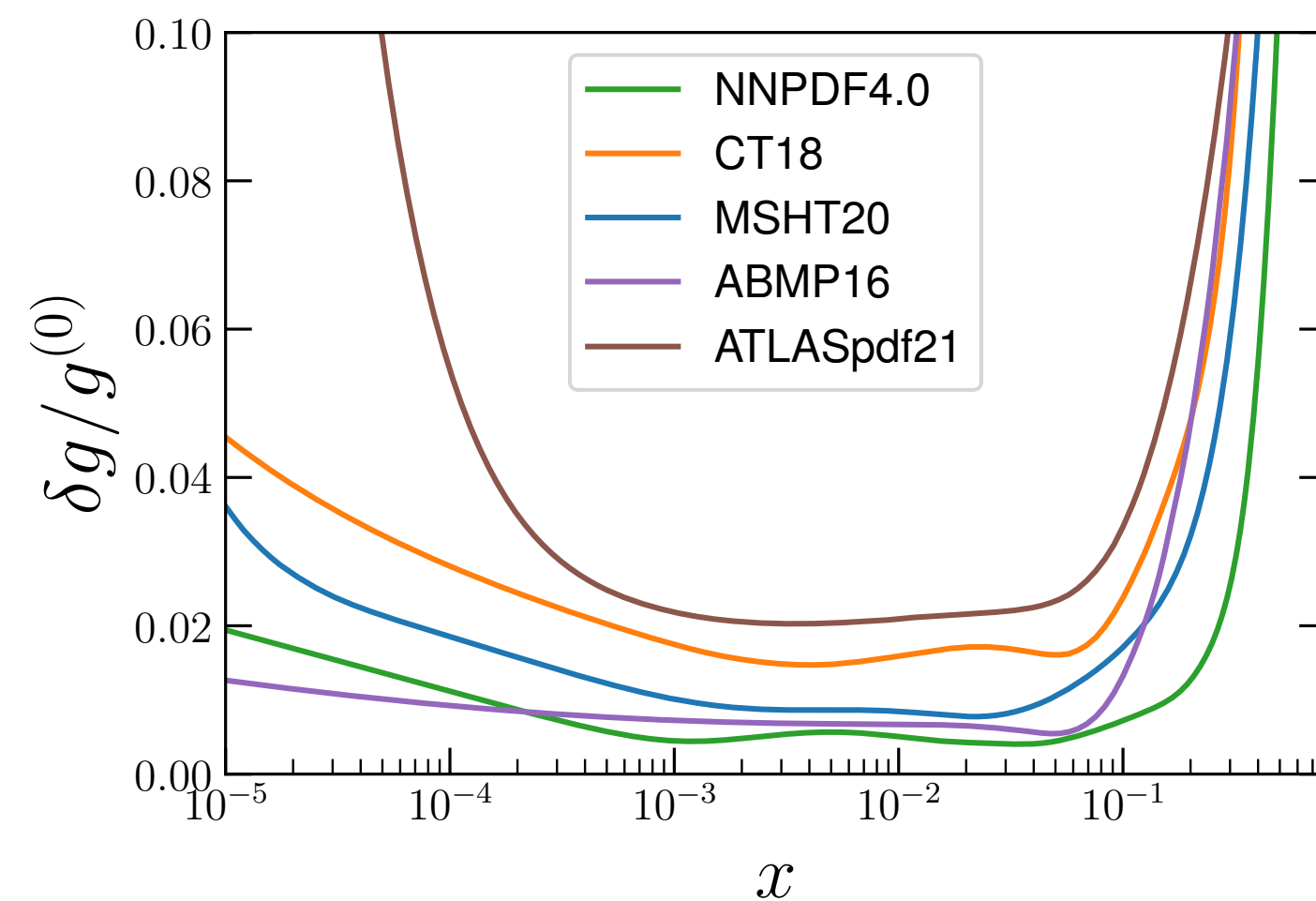
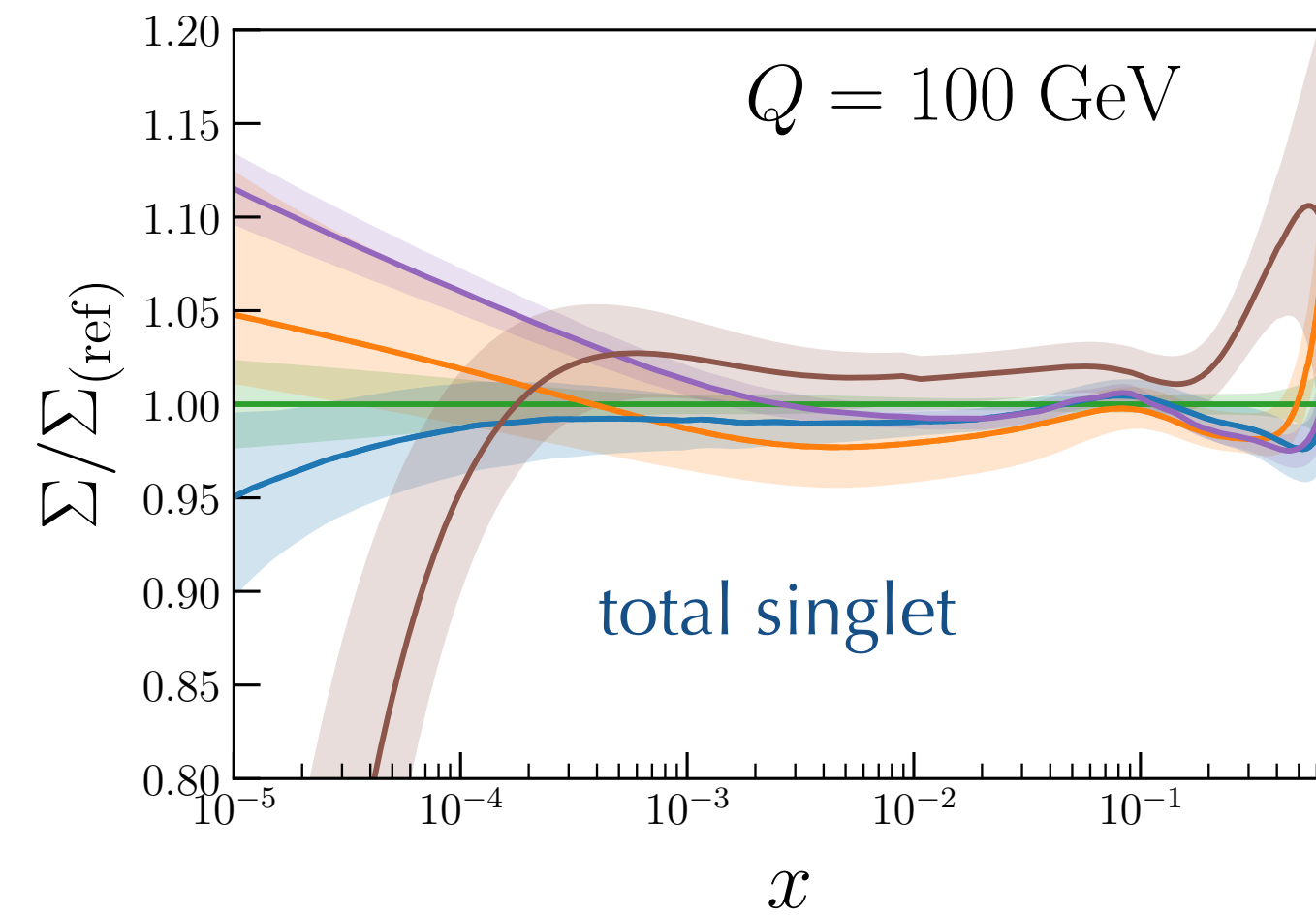
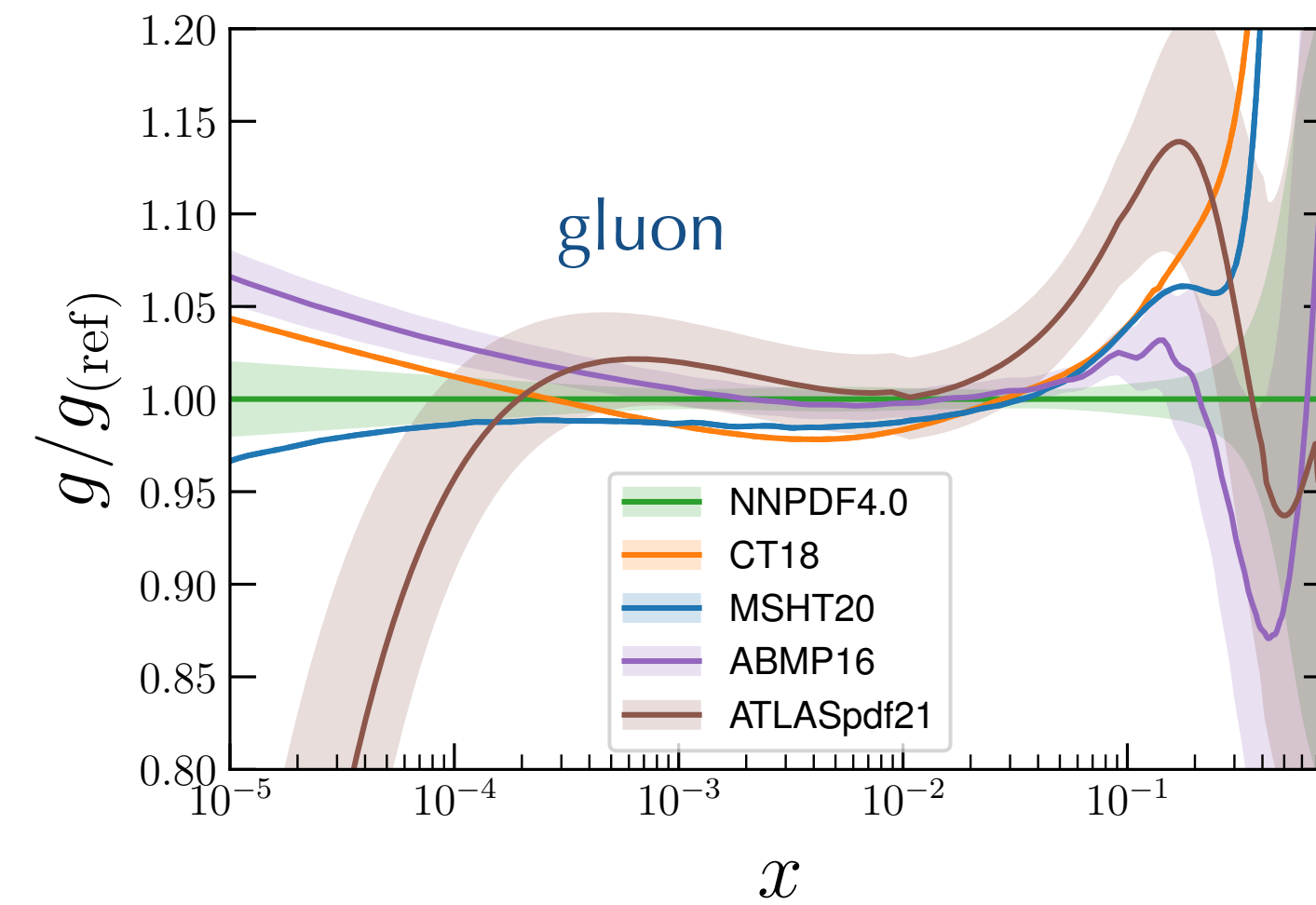


must have as many independent analyses as possible to have a faithful determination of PDFs and their uncertainties; state of the art PDFs are extracted at NNLO in QCD (+NLO EW) and with numerous LHC data

PDF comparisons

- Many ongoing efforts on comparisons and understanding of differences of up-to-date PDFs, in order to have a faithful determination of PDFs and its uncertainties

[Snowmass 2021, 2203.13923]



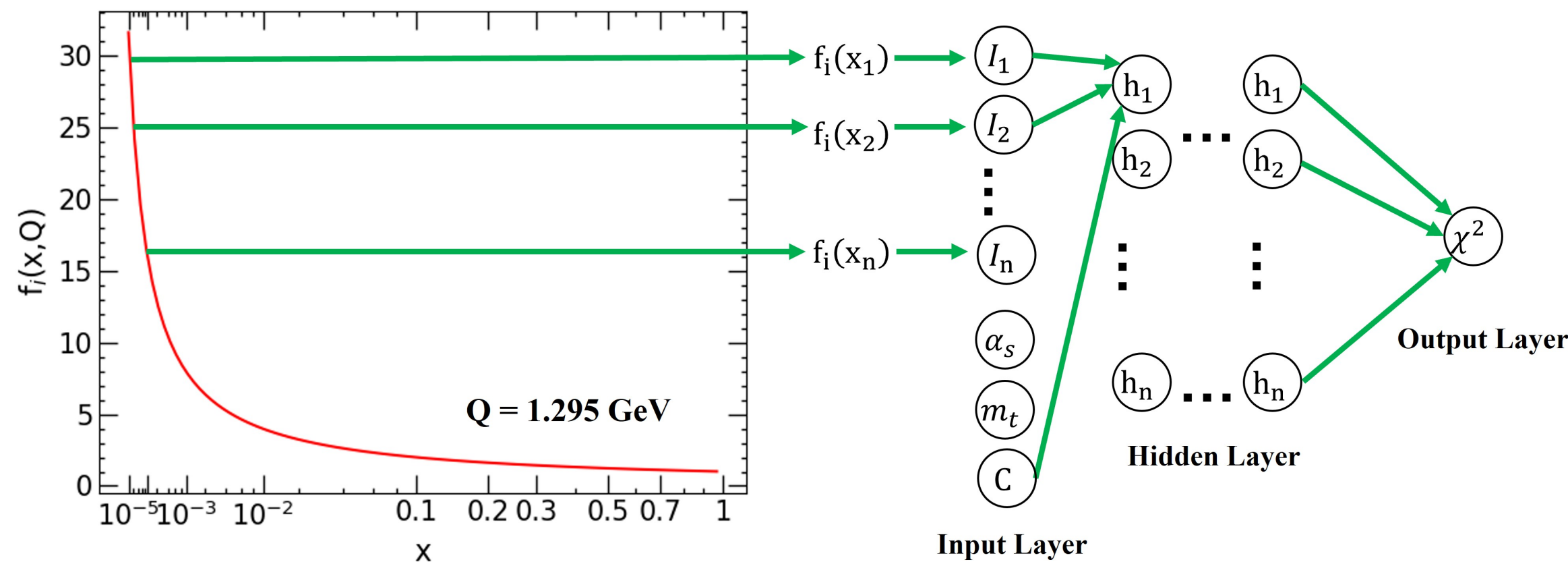
- general agreement between different groups (NN4.0, CT18, MSHT20, ABMP16, ATLAS21) over the range of x in 10^{-4} to 10^{-1} within uncertainties
- gluon: notable differences at $x \sim 0.2$, with 2σ for NN vs. CT&MSHT; singlet: ATLASpdf deviate at $x < 10^{-4}$ due to $Q^2 > 10 \text{ GeV}^2$ applied on HERA data, and at $x > 0.2$ due to lack of fixed-target data
- NN and ABMP show uncertainty of $\sim 1\text{-}2\%$ in constrained region mostly due to methodologies; CT18 being conservative among all fits; ATLAS unc. blow up in unconstrained region

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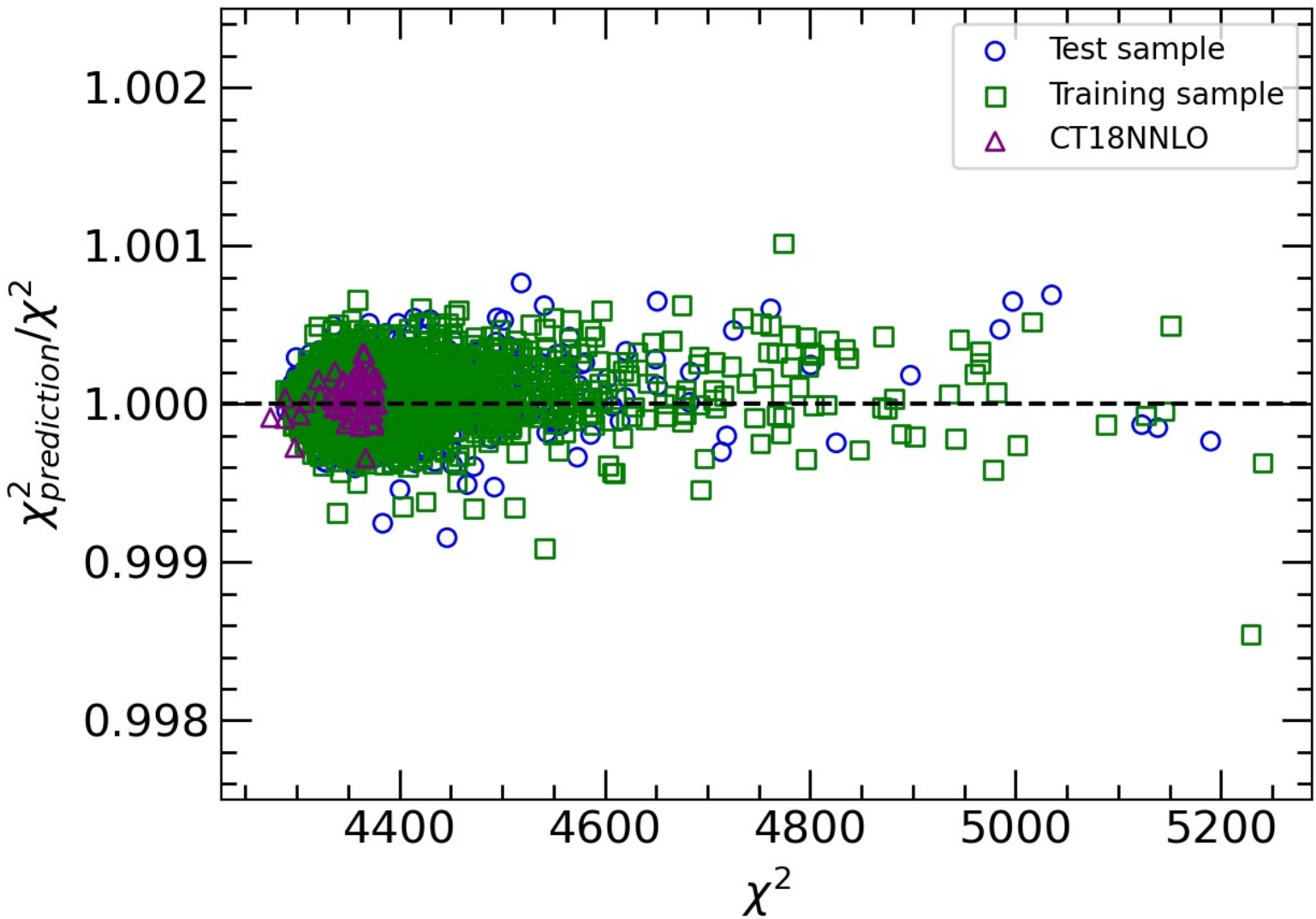
Global analysis with Machine Learnings

- ✦ We developed an efficient framework of using neural networks and machine learnings to modeling dependencies of the log-likelihood functions (χ^2) or cross sections on parton distributions; the computing efficiency improved by orders of magnitudes comparing to traditional method **[JG, DY Liu, CL Sun 2201.06586]**

sketch of NN architecture



modeling accuracy



- ✦ multi-layer NNs with discretized PDFs as input; trained with a large sample from CT18 MC PDFs
- ✦ reproduction accuracy of χ^2 better than one per mille
- ✦ almost costless comparing to traditional methods that requires extensive calculations of cross sections

computing efficiency per parameter point

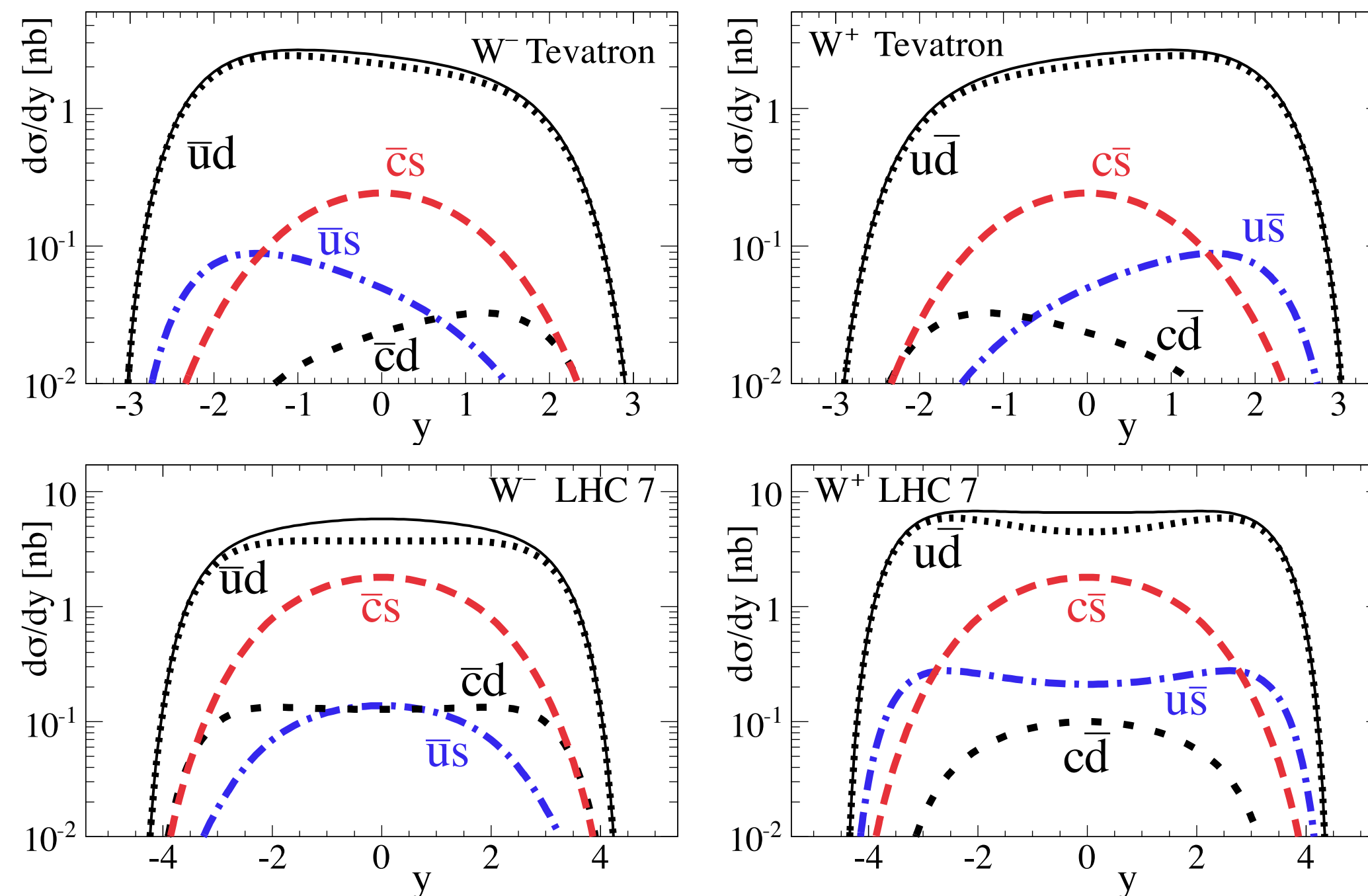
cost \ target			
	χ^2	σ	$f(x, Q)$
method			
NNs	0.70 ms	0.41 ms	0.37 ms
traditional	$10^7(300)$ ms	$10^6(30)$ ms	20(2) ms

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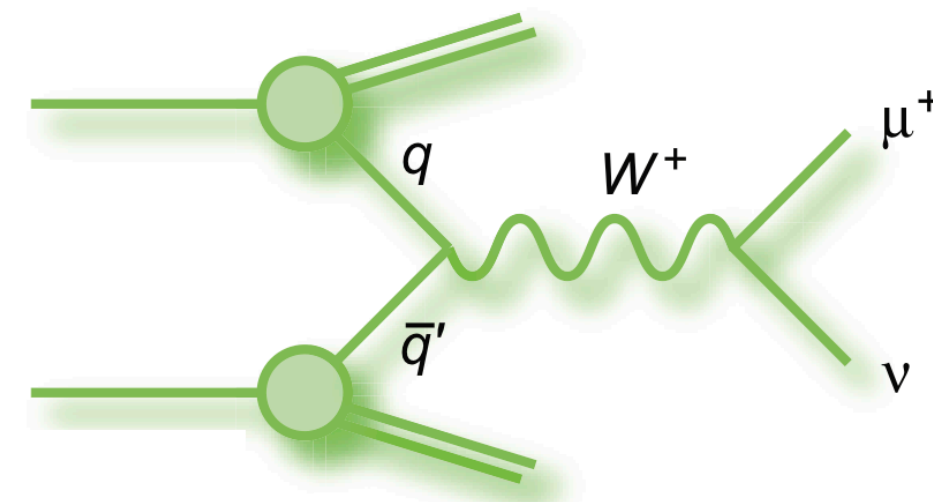
Question 1: PDF uncertainties in W-mass direct measurements

- PDFs are key inputs for precision programs at hadron colliders, e.g., direct measurements on the W boson mass and the weak mixing angle

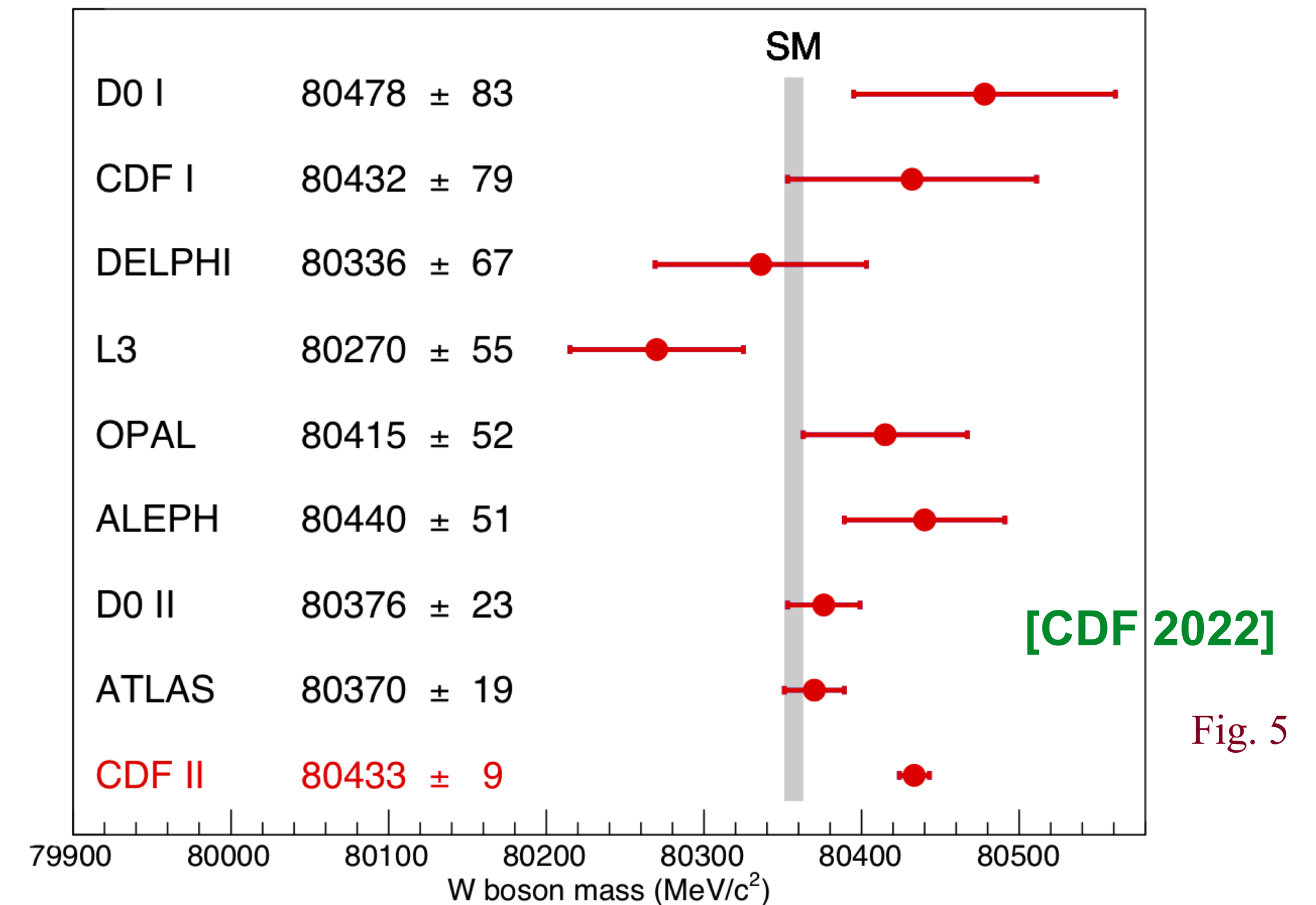
W boson rapidity distribution [1203.1290]



Drell-Yan production



W boson mass from different experiments



SM expectation: $M_W = 80,357 \pm 4_{\text{inputs}} \pm 4_{\text{theory}}$ (PDG 2020)

LHCb measurement: $M_W = 80,354 \pm 23_{\text{stat}} \pm 10_{\text{exp}} \pm 17_{\text{theory}} \pm 9_{\text{PDF}}$ [JHEP 2022, 36 (2022)]

PDF unc. of CDF / ATLAS / LHCb: 3.9 / 8 / 9 MeV

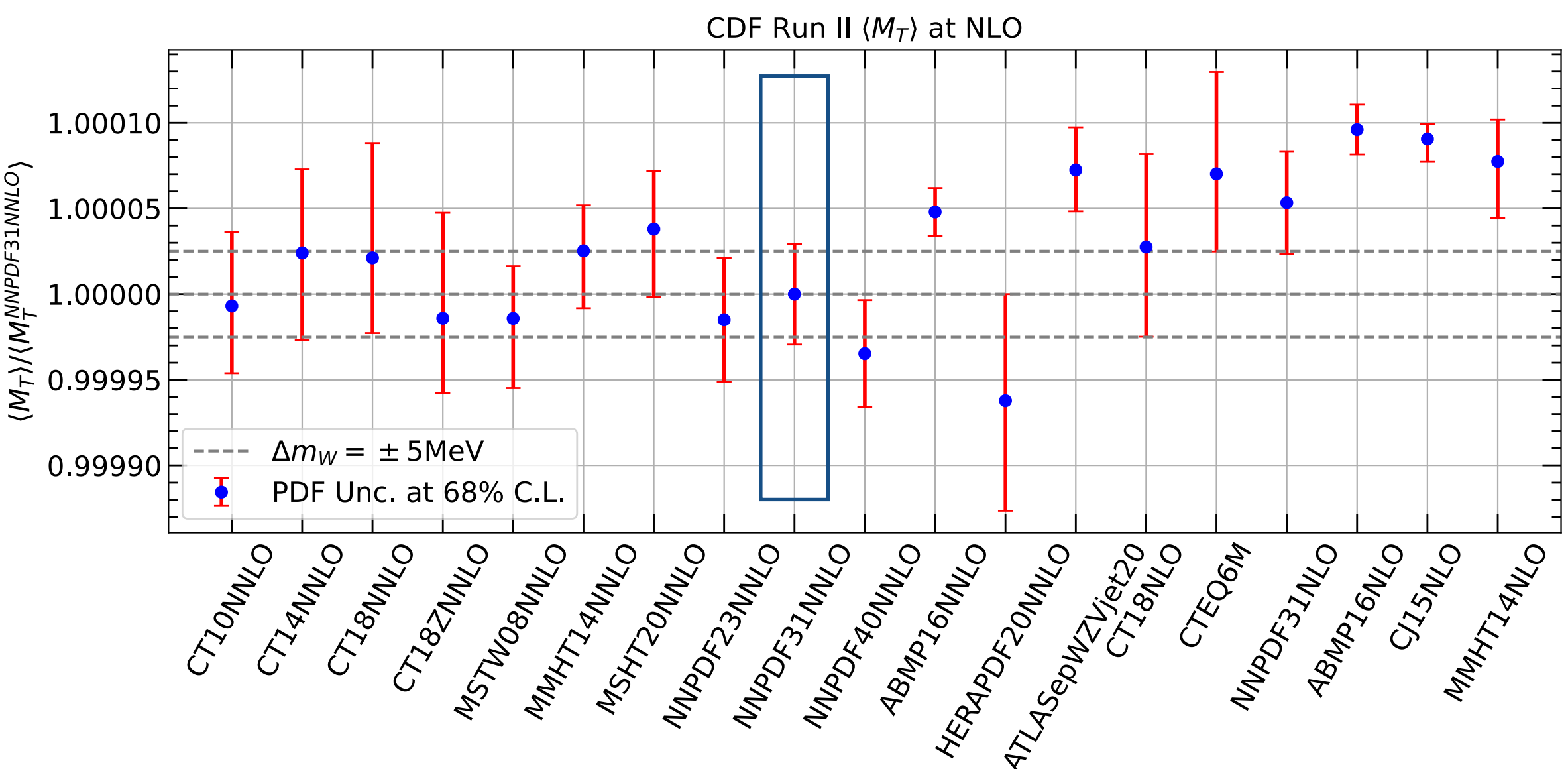
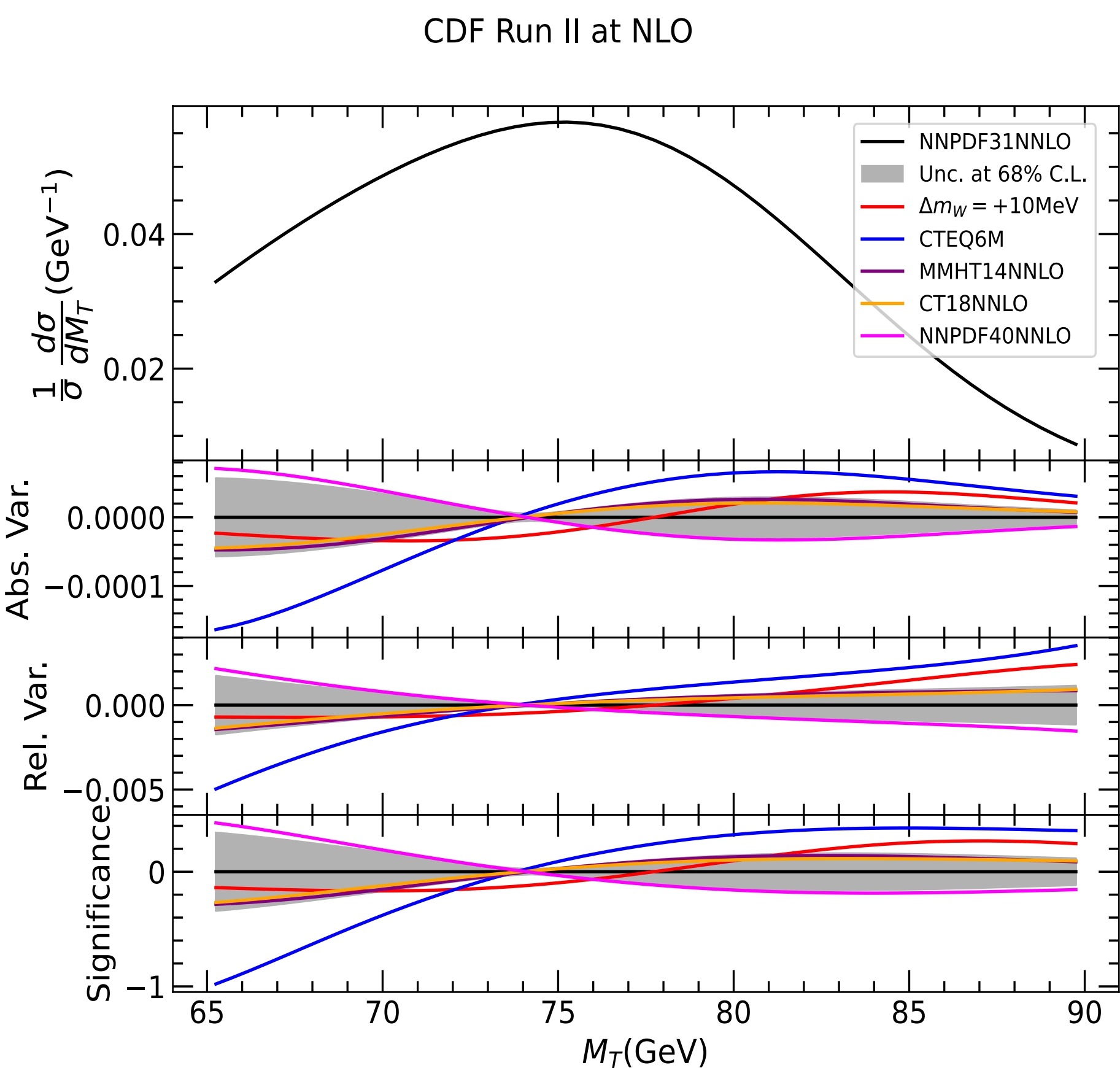
PDF variations can not explain CDF discrepancy

- ✦ We estimate shift of extracted W boson mass induced by variation of PDFs, and the associated PDF uncertainty for a variety of PDFs, focusing on the kinematic variable of transverse mass at CDF

[JG, DY Liu, KP Xie, 2205.03942]

mean value of m_T

normalized m_T distribution
PDF var. vs. M_W var.



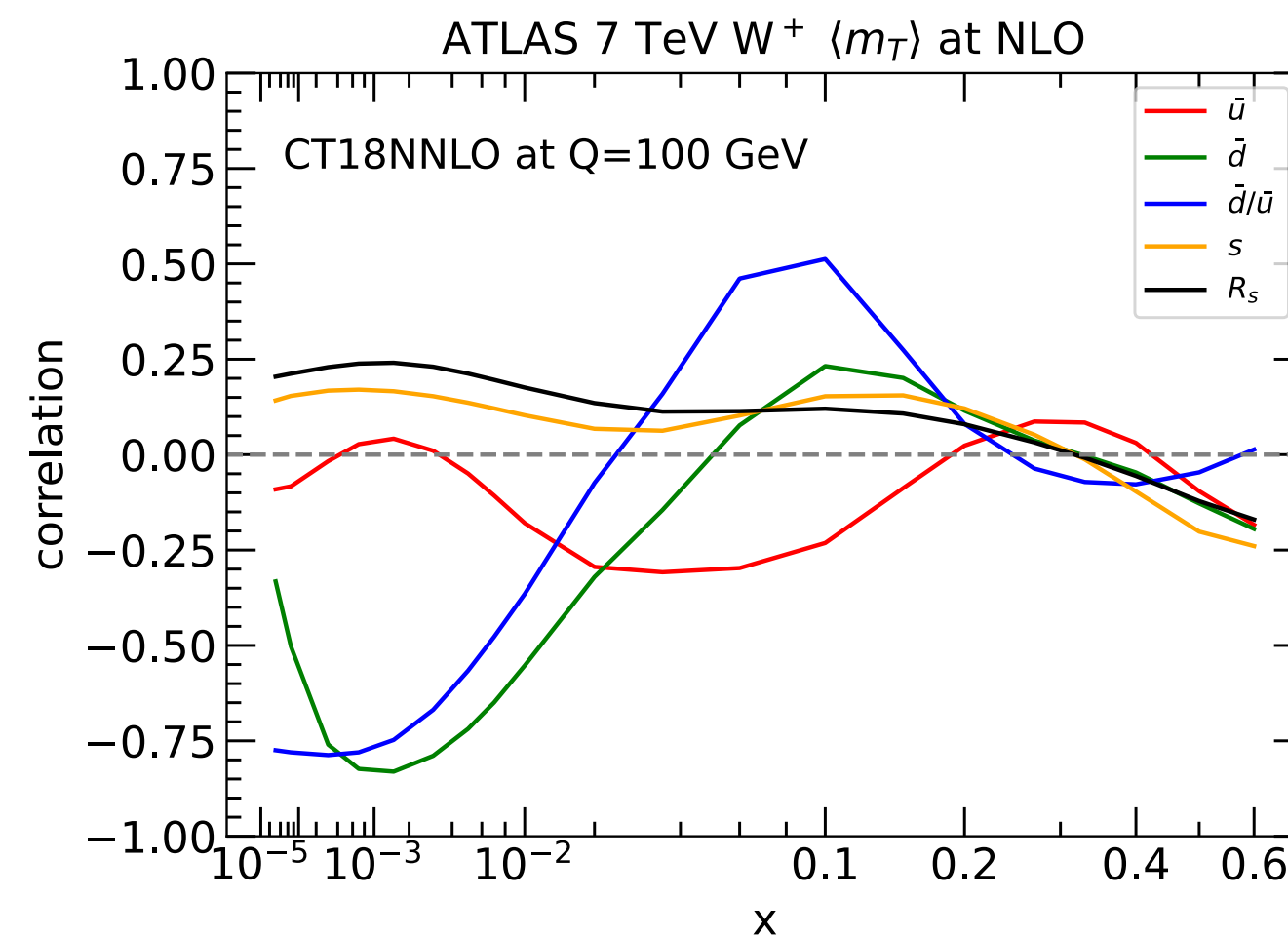
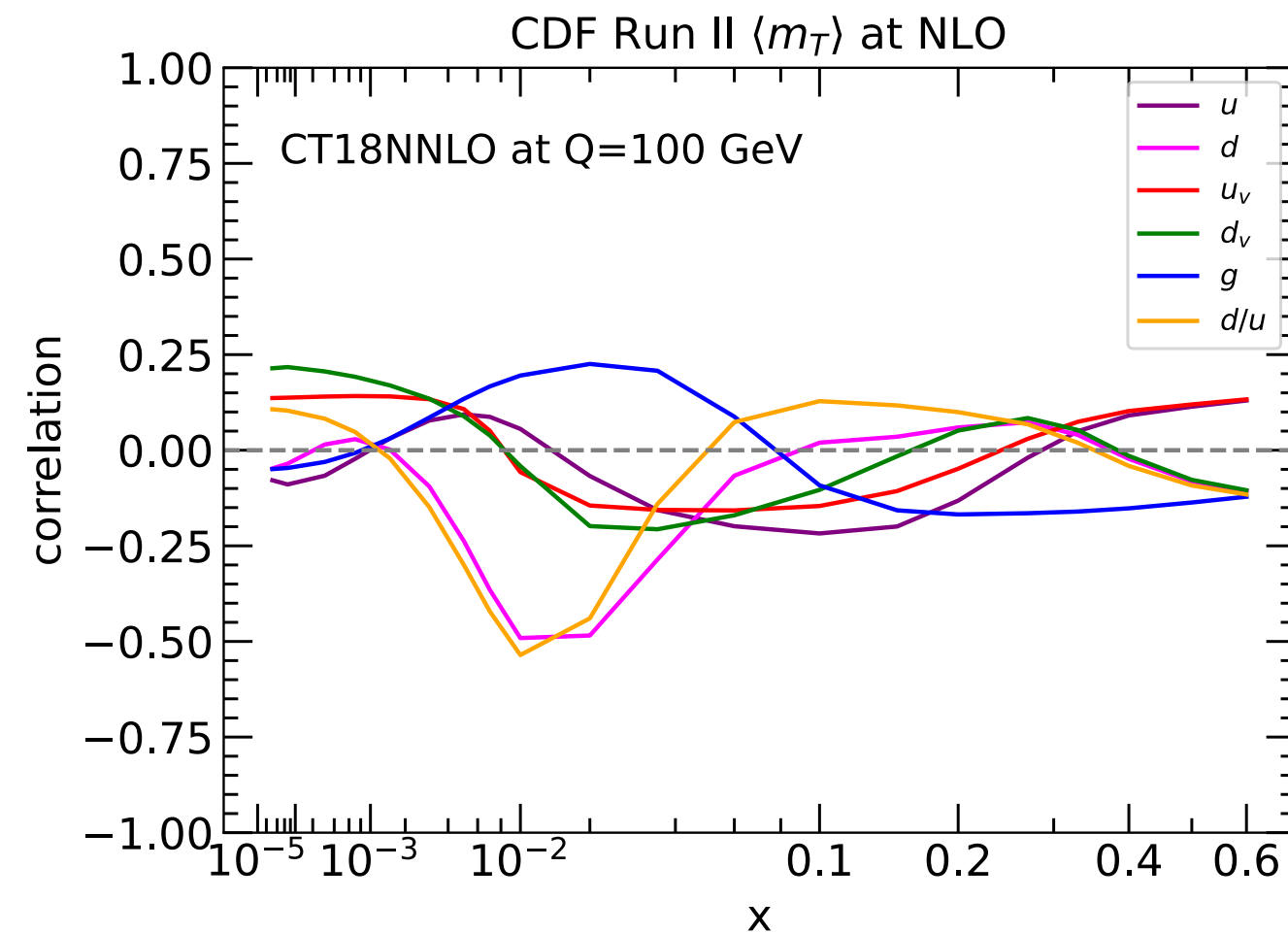
estimated shift and PDF unc. of W mass

δM_W in MeV	sta.	NNPDF3.1	CT18	MMHT2014	NNPDF4.0	MSHT2020
$\langle M_T \rangle$ (LO)	—	$0^{+8.3}_{-8.3}$	$-1.0^{+8.3}_{-11.4}$	$-3.3^{+7.4}_{-4.2}$	$+7.8^{+5.1}_{-5.1}$	$-3.1^{+6.7}_{-5.7}$
χ^2 fit (LO)	8.0	$0^{+7.6}_{-7.6}$	$-1.0^{+5.4}_{-8.6}$	$-3.3^{+6.1}_{-3.0}$	$+8.0^{+3.7}_{-3.7}$	$-3.0^{+5.0}_{-4.0}$
$\langle M_T \rangle$ (NLO)	—	$0^{+5.9}_{-5.9}$	$-4.2^{+8.8}_{-13.3}$	$-5.0^{+6.7}_{-5.3}$	$+6.9^{+6.2}_{-6.2}$	$-7.6^{+7.9}_{-6.7}$
χ^2 fit (NLO)	8.0	$0^{+4.2}_{-4.2}$	$-4.3^{+5.4}_{-10.1}$	$-5.1^{+4.8}_{-3.4}$	$+7.1^{+4.5}_{-4.5}$	$-7.8^{+5.7}_{-4.5}$
CDF	9.2	$0^{+3.9}_{-3.9}$	—	—	—	—

PDF variations can not explain CDF discrepancy

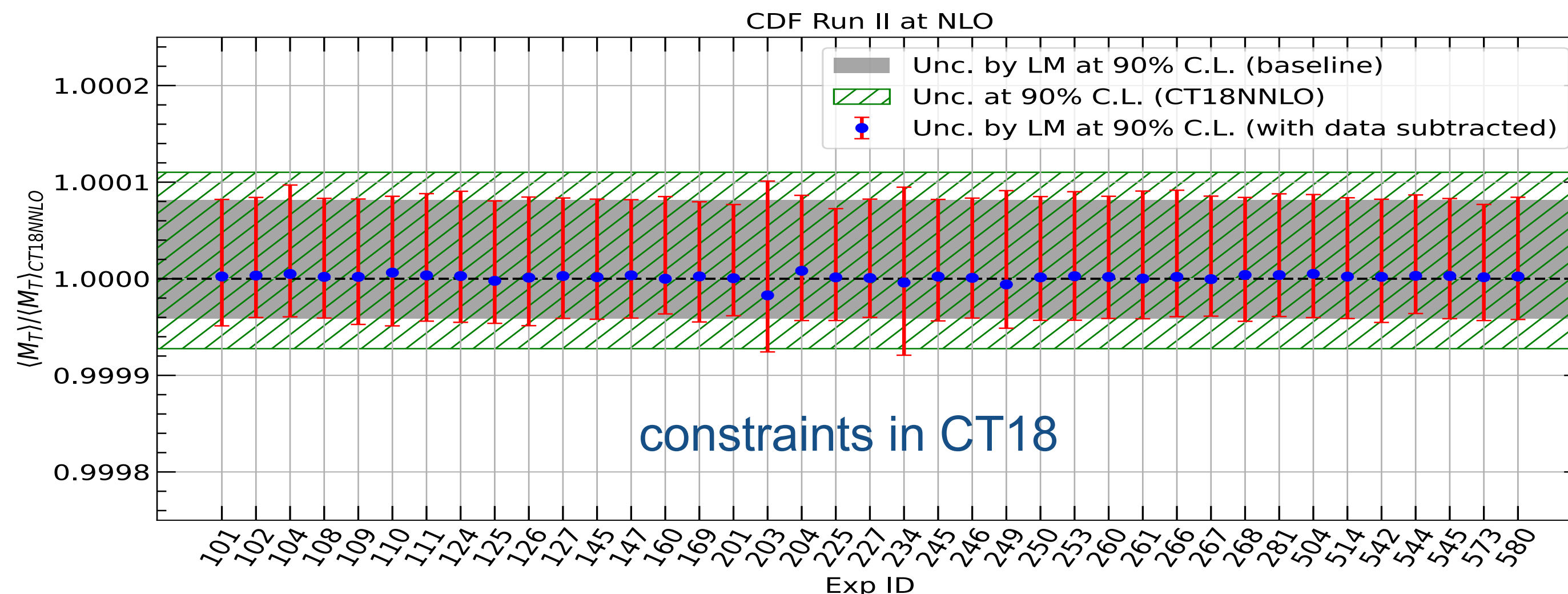
- ✦ We carry out a series of Lagrange multiplier scans to identify the constraints on the transverse mass distribution (using mean M_T) imposed by individual data sets in the CT18 global analysis

PDF induced correlations



[JG, DY Liu, KP Xie, 2205.03942]

- ✦ m_T at CDF (ATLAS) is mostly sensitive to the d-quark (d-bar-quark) at $x \sim 0.01$ (0.001); CDF and ATLAS are largely uncorrelated



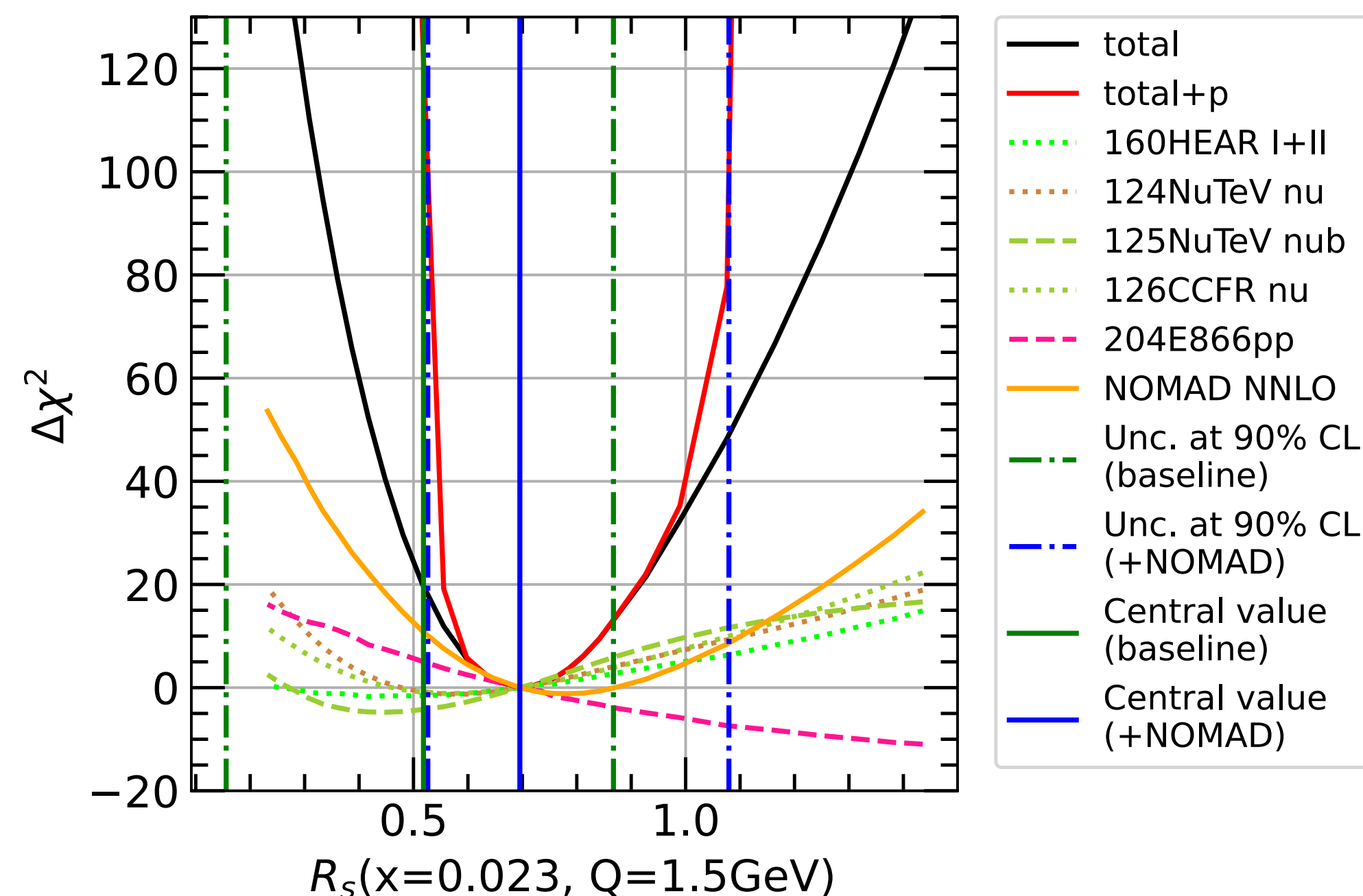
- ✦ m_T at CDF is largely constrained by the DIS and Drell-Yan data on deuteron target, the Tevatron lepton charge asymmetry data; at ATLAS also the CMS charge asymmetry data

Strangeness is moderately suppressed

- ✦ We include NOMAD data into a global analysis of PDFs (CT18 as the baseline), and analysis its impact to PDFs, especially focusing on strange-quark PDF and strange to light sea-quark ratio $R_s=(s+sb)/(ub+db)$

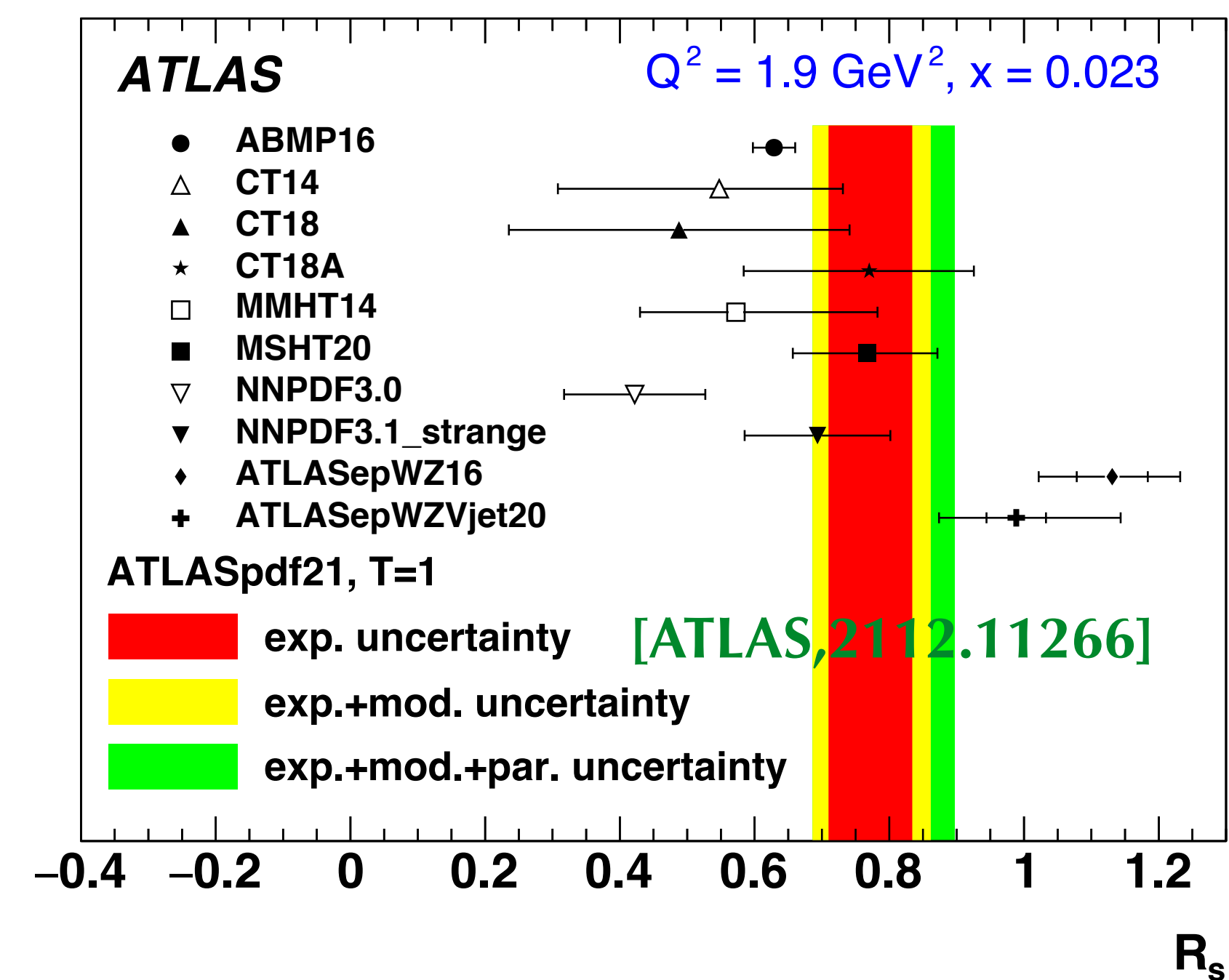
[JG, DY Liu, CL Sun 2201.06586]

LM scans on R_s



- ✦ NOMAD prefers larger s-PDF comparing to NuTeV and CCFR dimuon; leads to increase of R_s , from 0.5 to 0.7
- ✦ reduction of PDF uncertainty by more than 30%

recent ATLAS measurement



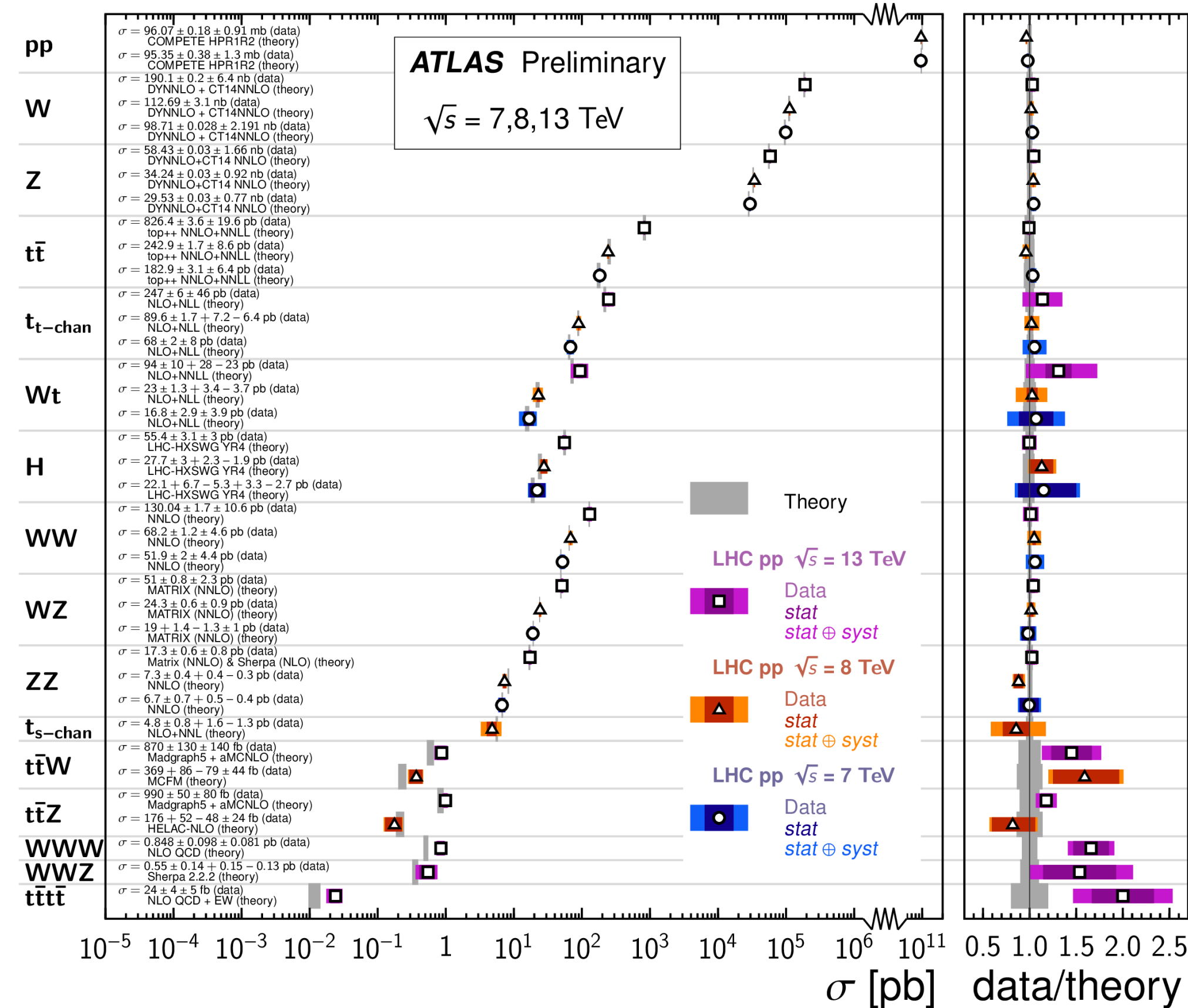
- ✦ tensions between dimuon data ($R_s \sim 0.5$) and LHC data ($R_s \sim 1$) exist for years; now relieved
- ✦ most recent ATLAS data shows $R_s \sim 0.8$

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Question 2: PDF bias in searches of new physics at the LHC

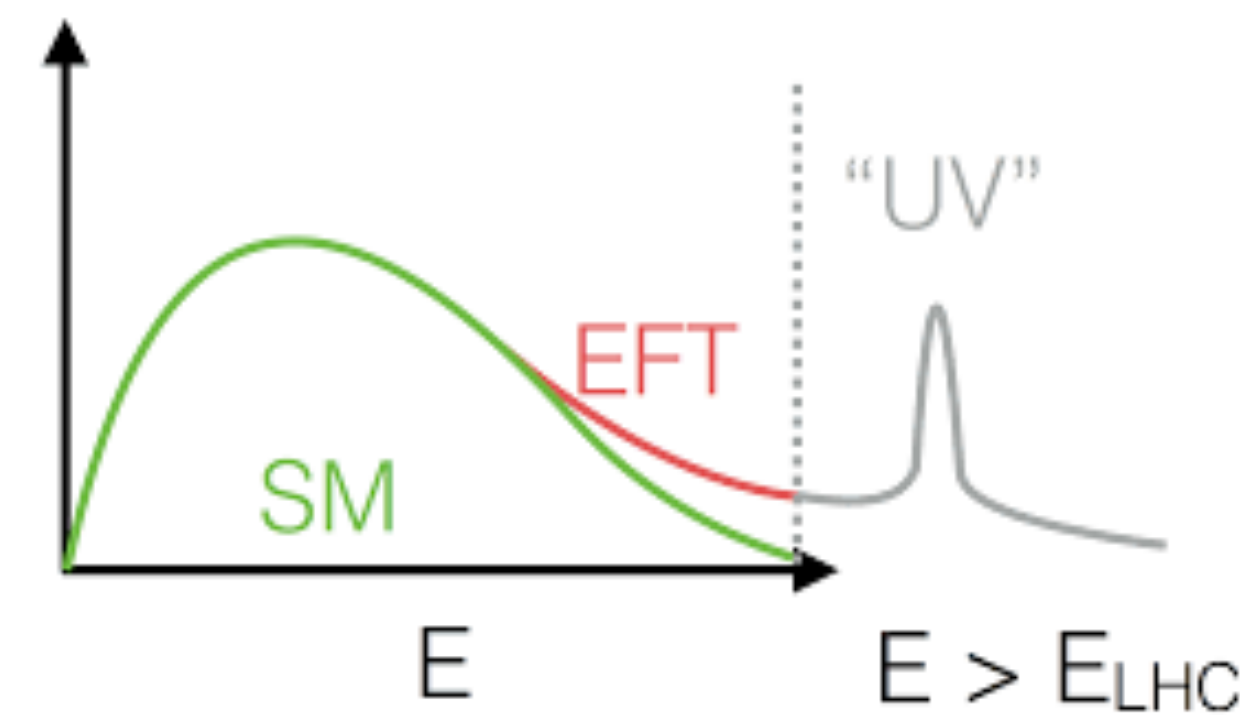
- PDFs are key inputs for searches of new physics beyond the SM at hadron colliders, especially non-resonance signatures hiding in high mass tails taking SM effective theory (SMEFT) as an example

Standard Model Total Production Cross Section Measurements



SM cross sections from ATLAS

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i^{(6)} O_i^{(6)}}{\Lambda^2}$$



conventionally, constraints on NP are determined using PDFs extracted from similar data sets but with pure SM assumptions

$$\frac{d\sigma}{dp_T} = f_1(x_1) \otimes f_2(x_2) \otimes \frac{d\hat{\sigma}_{SM}}{dp_T} [1 + O(\alpha_s) + O(\alpha_{EW}) + O(\frac{1}{\Lambda^2})]$$

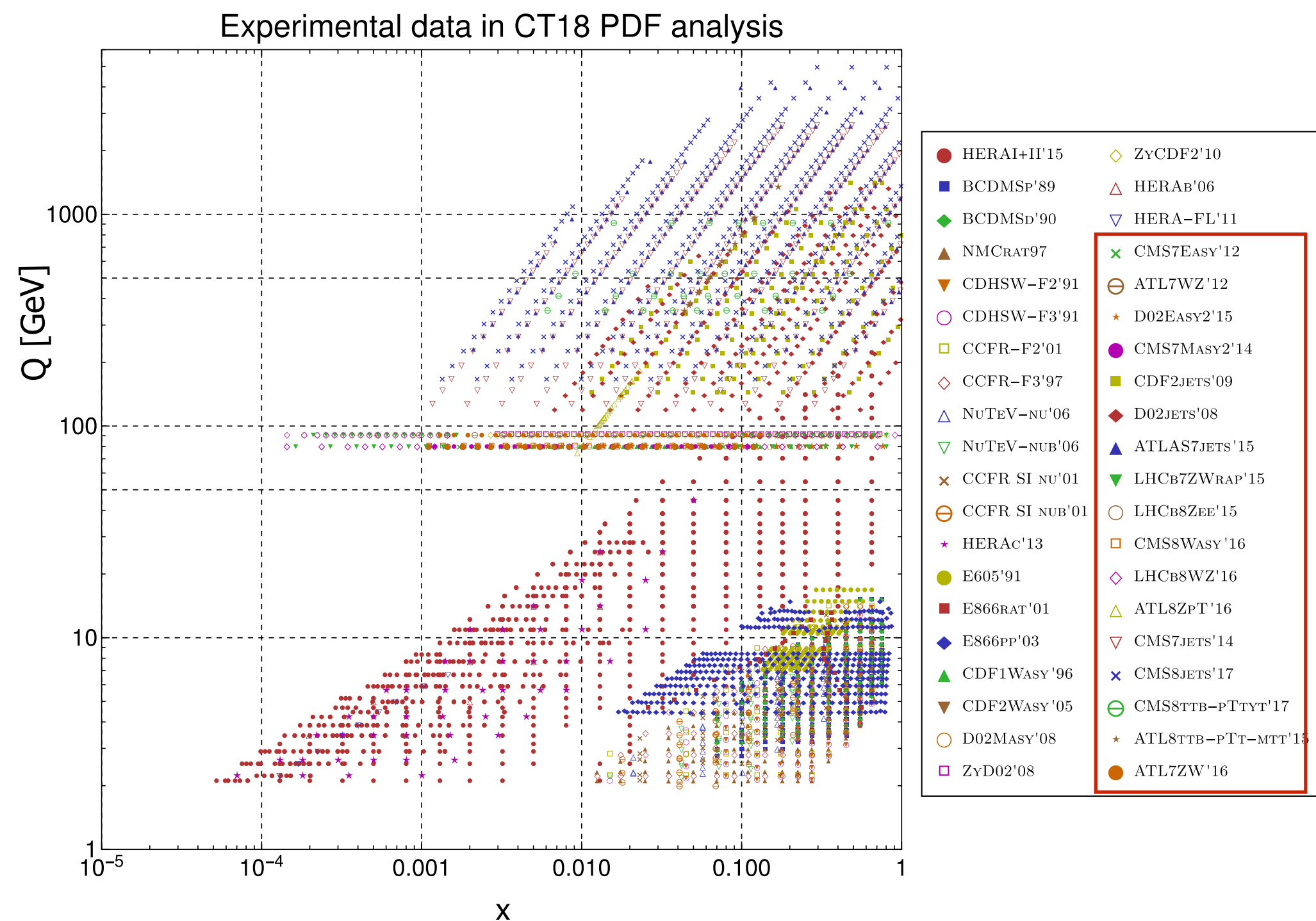
Joint fit of PDFs and BSM

- Based on our framework of neural networks we performed a joint fit of PDFs and new physics beyond the SM (PDF+BSM); the later are described by EFT couplings of operators in SMEFT

full data set (CT18 as baseline)

[JG, DY Liu+, 2022]

subset of data (top pair, jet) directly constrain EFT



focusing on BSM relevant for top pair and jet production that are both sensitive to gluon PDFs

Experiments	$\sqrt{s}(\text{TeV})$	$\mathcal{L}(\text{fb}^{-1})$	observable	N_{pt}
*† LHC(Tevatron)	7/8/13(1.96)	—	$t\bar{t}$ total cross section [16–21]	8
*† ATLAS $t\bar{t}$	8	20.3	1D dis. in $p_{T,t}$ or $m_{t\bar{t}}$ [22]	15
*† CMS $t\bar{t}$	8	19.7	2D dis. in $p_{T,t}$ and y_t [23]	16
CMS $t\bar{t}$	8	19.7	1D dis. in $m_{t\bar{t}}$ [24]	7
*† ATLAS $t\bar{t}$	13	36	1D dis. in $m_{t\bar{t}}$ [25]	7
*† CMS $t\bar{t}$	13	35.9	1D dis. in $m_{t\bar{t}}$ [26]	7
*† CDF II inc. jet	1.96	1.13	2D dis. in p_T and y [27]	72
*† D0 II inc. jet	1.96	0.7	2D dis. in p_T and y [28]	110
*† ATLAS inc. jet	7	4.5	2D dis. in p_T and y [29]	140
*† CMS inc. jet	7	5	2D dis. in p_T and y [30]	158
* CMS inc. jet	8	19.7	2D dis. in p_T and y [31]	185
† CMS dijet	8	19.7	3D dis. in $p_T^{\text{ave.}}$, y_b and y^* [32]	122
† CMS inc. jet	13	36.3	2D dis. in p_T and y [8]	78

theoretical calculations

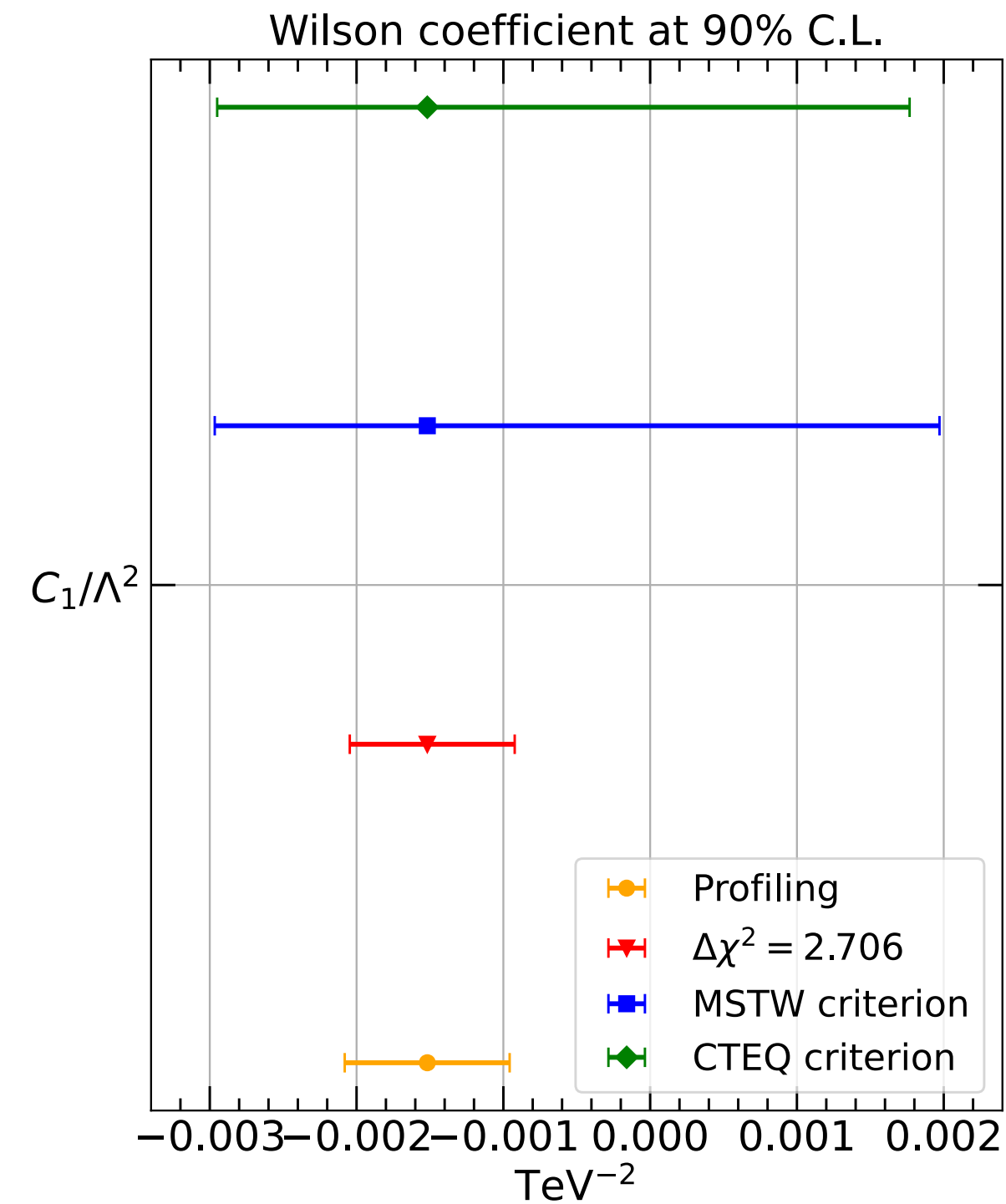
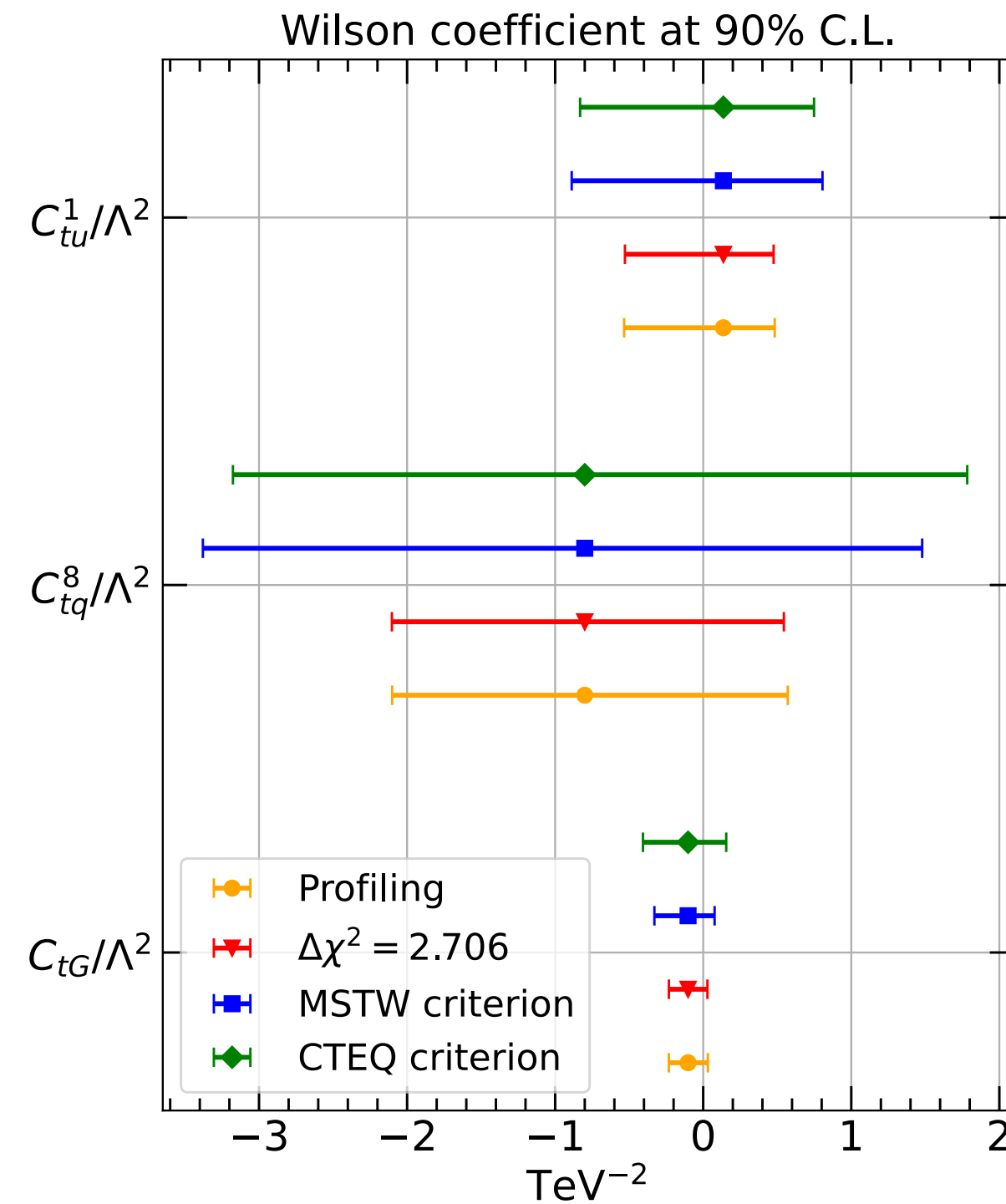
$$\frac{d\sigma}{d\mathcal{O}} = \frac{d\sigma_{\text{SM}}}{d\mathcal{O}} + \sum_i \frac{d\tilde{\sigma}_i}{d\mathcal{O}} \frac{C_i}{\Lambda^2} + \sum_{i,j} \frac{d\tilde{\sigma}_{ij}}{d\mathcal{O}} \frac{C_i C_j}{\Lambda^4}$$

observable	μ_0	SM QCD	SM EW	SMEFT QCD	th. unc.
$t\bar{t}$ total	m_t	NNLO+NNLL	no	NLO	$\mu_{F,R}$ var.
$t\bar{t}$ p_T dist.	$m_T/2$	NNLO	NLO	NLO	$\mu_{F,R}$ var.
$t\bar{t}$ $m_{t\bar{t}}$ dist.	$H_T/4$	NNLO(+NLP)	NLO	NLO	$\mu_{F,R}$ var.
$t\bar{t}$ 2D dist.	$H_T/4$	NNLO	no	NLO	no
inc. jet	$p_{T,j}$	NNLO	NLO	NLO	0.5% uncor.
dijet	m_{jj}	NNLO	NLO	NLO	0.5% uncor

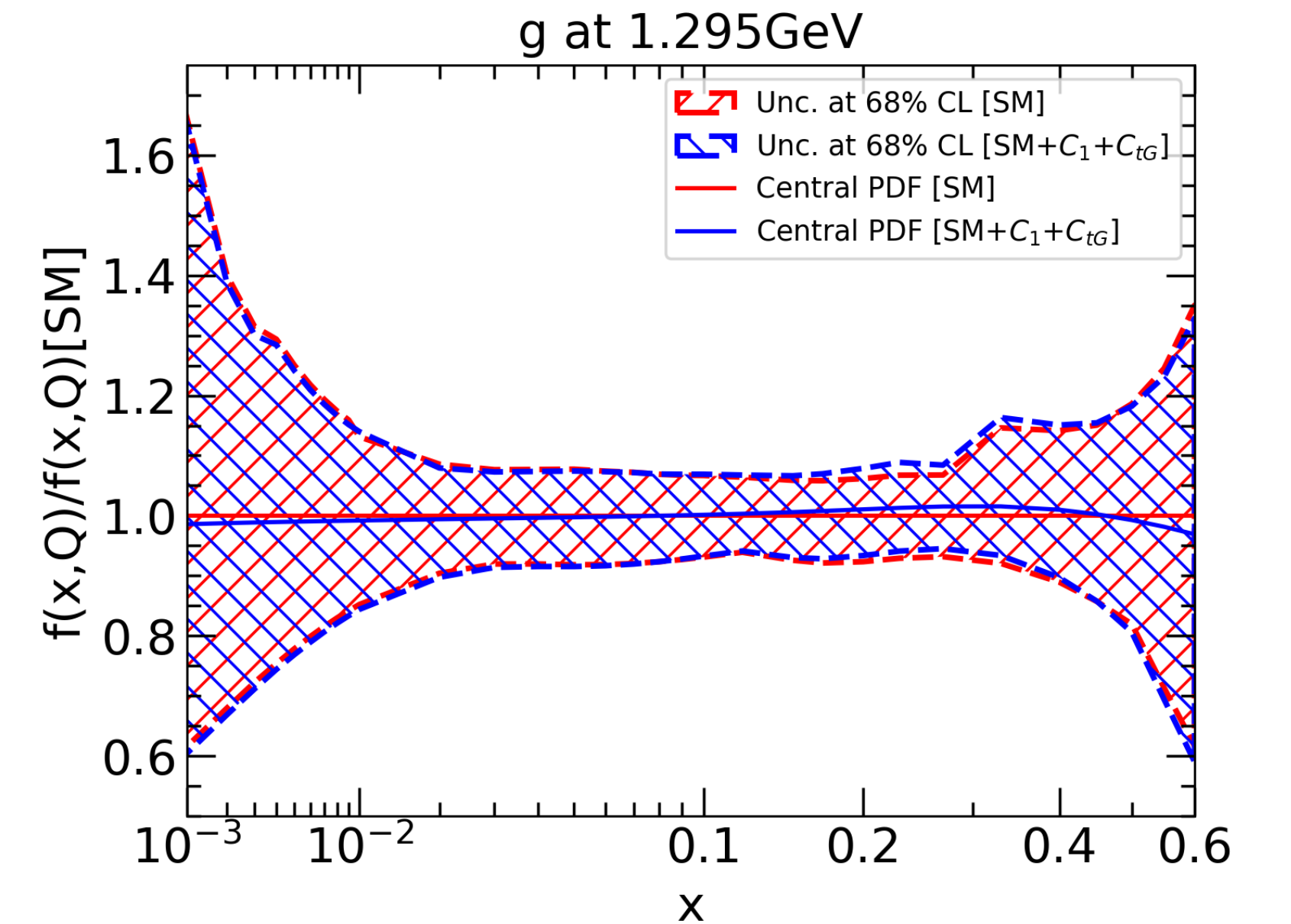
Correlations of PDFs and BSM are mild

- Unbiased results on four-quark and gluonic operators are obtained using global data sets with different tolerance criteria; current correlations between EFT and PDFs in global analyses are found to be mild

SMEFT coefficients extracted



gluon PDFs from SM/SM+EFT fits



$$O^1_{tu} = \sum_{i=1}^2 (\bar{t}\gamma_\mu t) (\bar{u}_i\gamma^\mu u_i) ,$$

$$O^1_{td} = \sum_{i=1}^3 (\bar{t}\gamma^\mu t) (\bar{d}_i\gamma_\mu d_i) ,$$

$$O_1 = 2\pi \left(\sum_{i=1}^3 \bar{q}_{Li}\gamma_\mu q_{Li} \right) \left(\sum_{j=1}^3 \bar{q}_{Lj}\gamma^\mu q_{Lj} \right)$$

$$O_{tG} = ig_s (\bar{Q}_{L,3}\tau^{\mu\nu}T^A t) \tilde{\varphi} G^A_{\mu\nu} + \text{h.c.} ,$$

$$O^8_{tq} = \sum_{i=1}^2 (\bar{Q}_i\gamma_\mu T^A Q_i) (\bar{t}\gamma^\mu T^A t) ,$$

SMEFT Operators

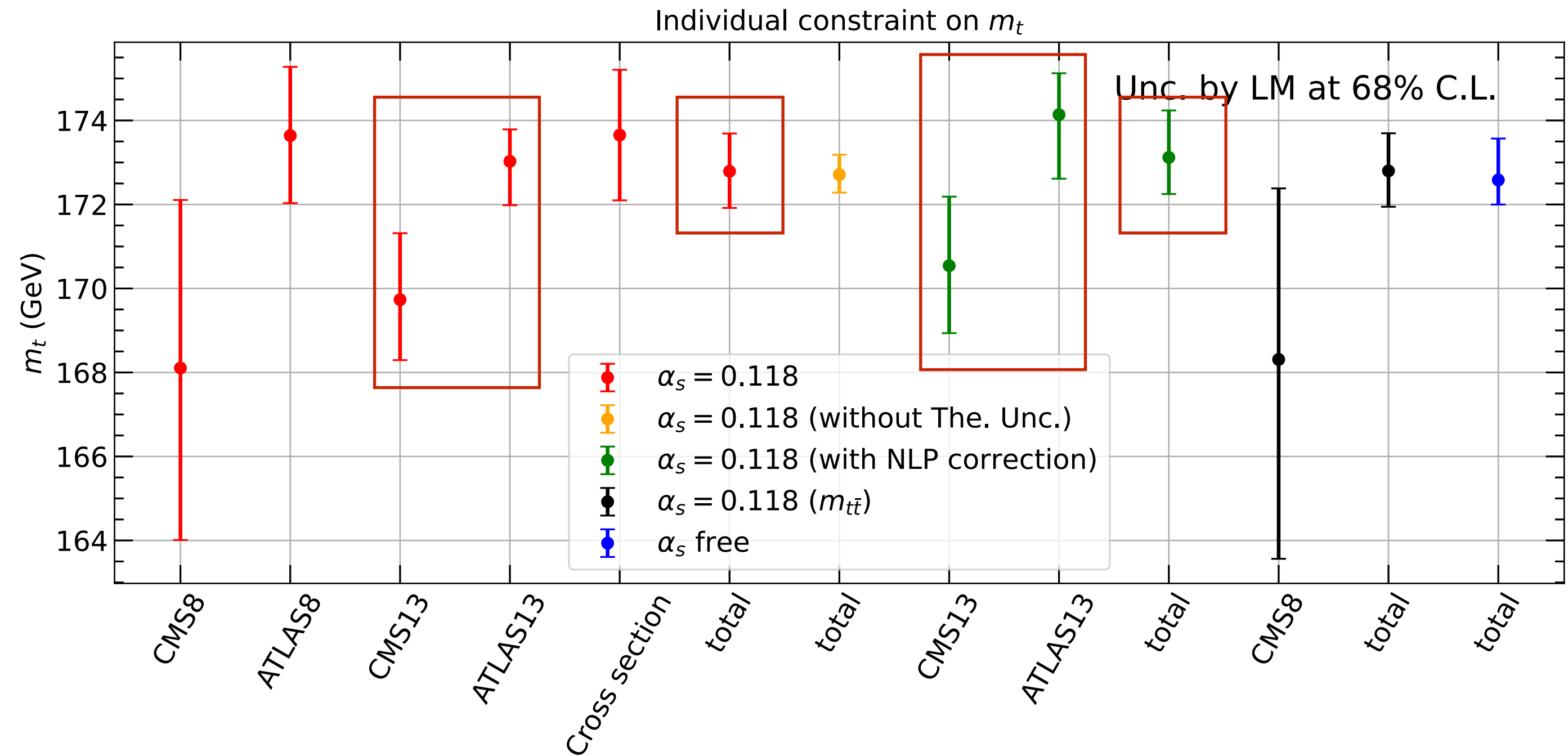
PDFs are essentially unchanged indicating weak correlations between PDF and BSM, and robustness of global analysis of QCD due to variety of experimental data

Top-quark mass from global analysis

- ◆ Meanwhile we extract the top-quark mass from the global analyses including the same data sets on total and differential cross sections of top-quark pair production **[JG, DY Liu+, 2022]**

setup	nominal	nominal	nominal	PDF fixed	no the. unc.	NLP
$\alpha_s(M_Z)$	free	0.1162	0.118	0.118	0.118	0.118
$m_t(\text{GeV})$	$172.58^{+0.99}_{-0.58}$	$172.58^{+0.98}_{-0.58}$	$172.79^{+0.90}_{-0.87}$	$172.79^{+0.90}_{-0.86}$	$172.71^{+0.43}_{-0.43}$	$173.12^{+1.12}_{-0.87}$

PDG world average
 $172.7 \pm 0.3 \text{ GeV}$



top-quark pole mass extracted at various conditions

- ◆ top-mass variations induce changes on total cross sections and kinematic bins close to threshold
- ◆ individual data sets from LHC are more or less consistent with preference on the central values lower for CMS than ATLAS
- ◆ theoretical unc. (beyond NNLO) are dominant; Coulomb corrections lead to larger top mass but can not account for differences between ATLAS and CMS

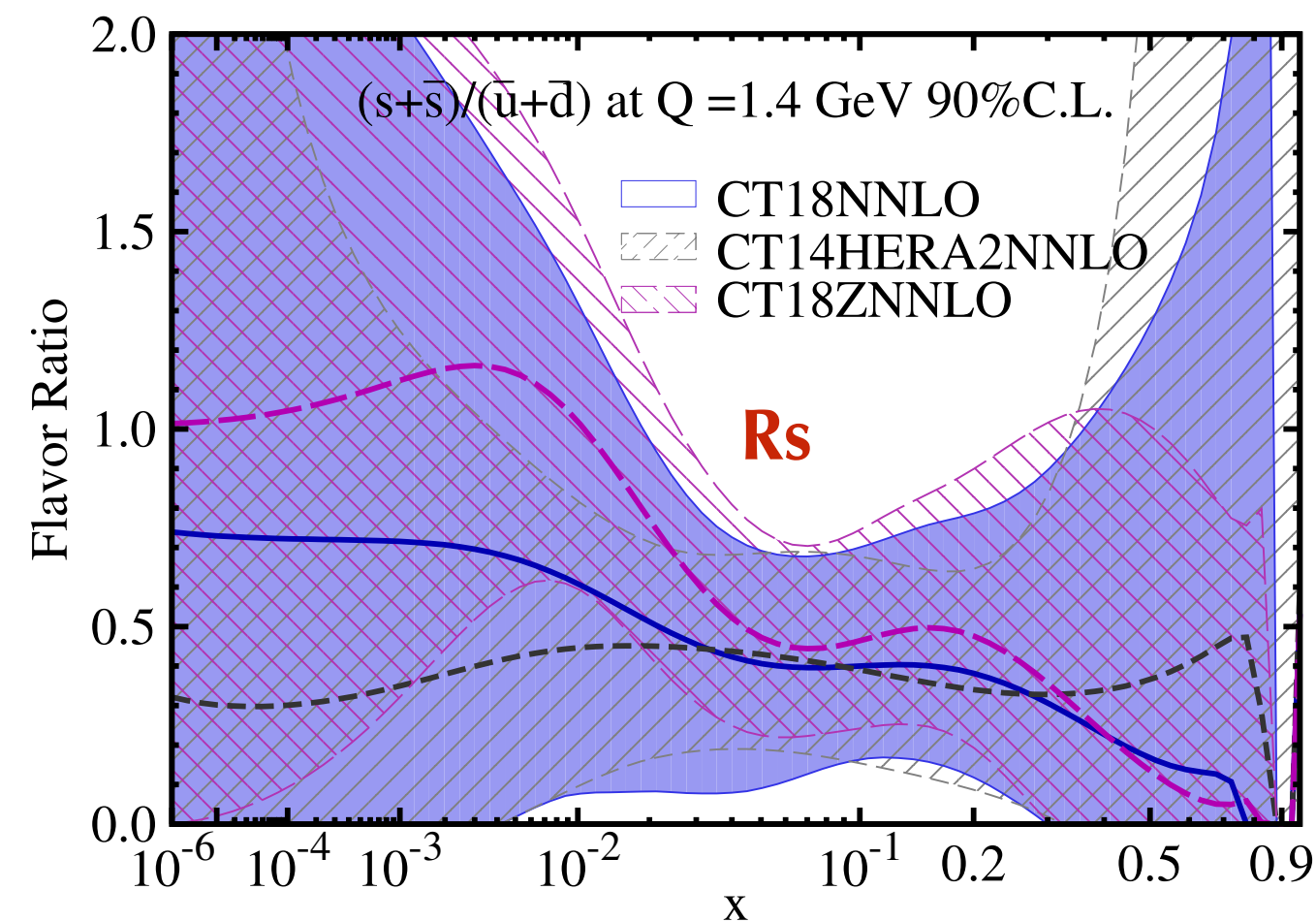
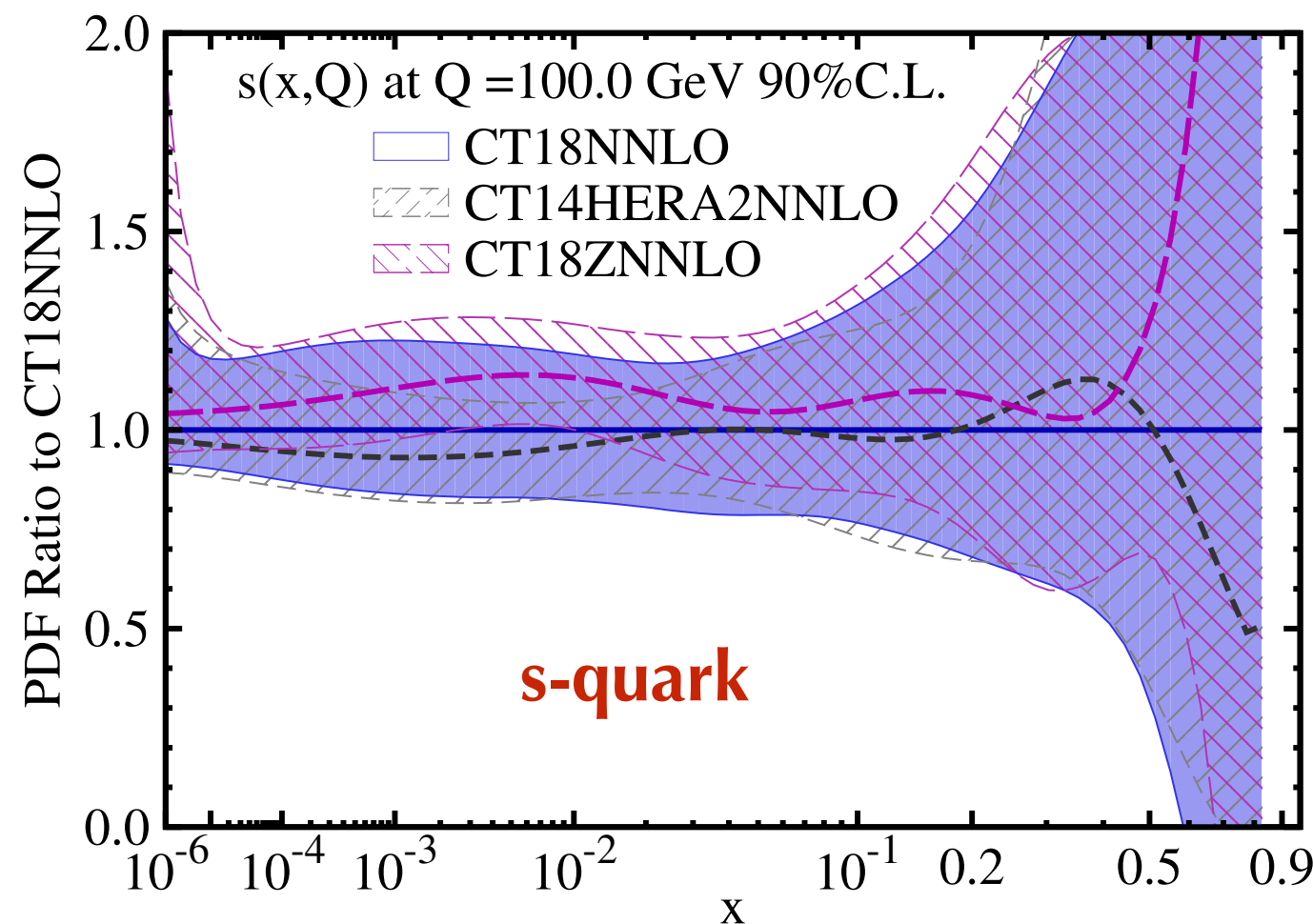
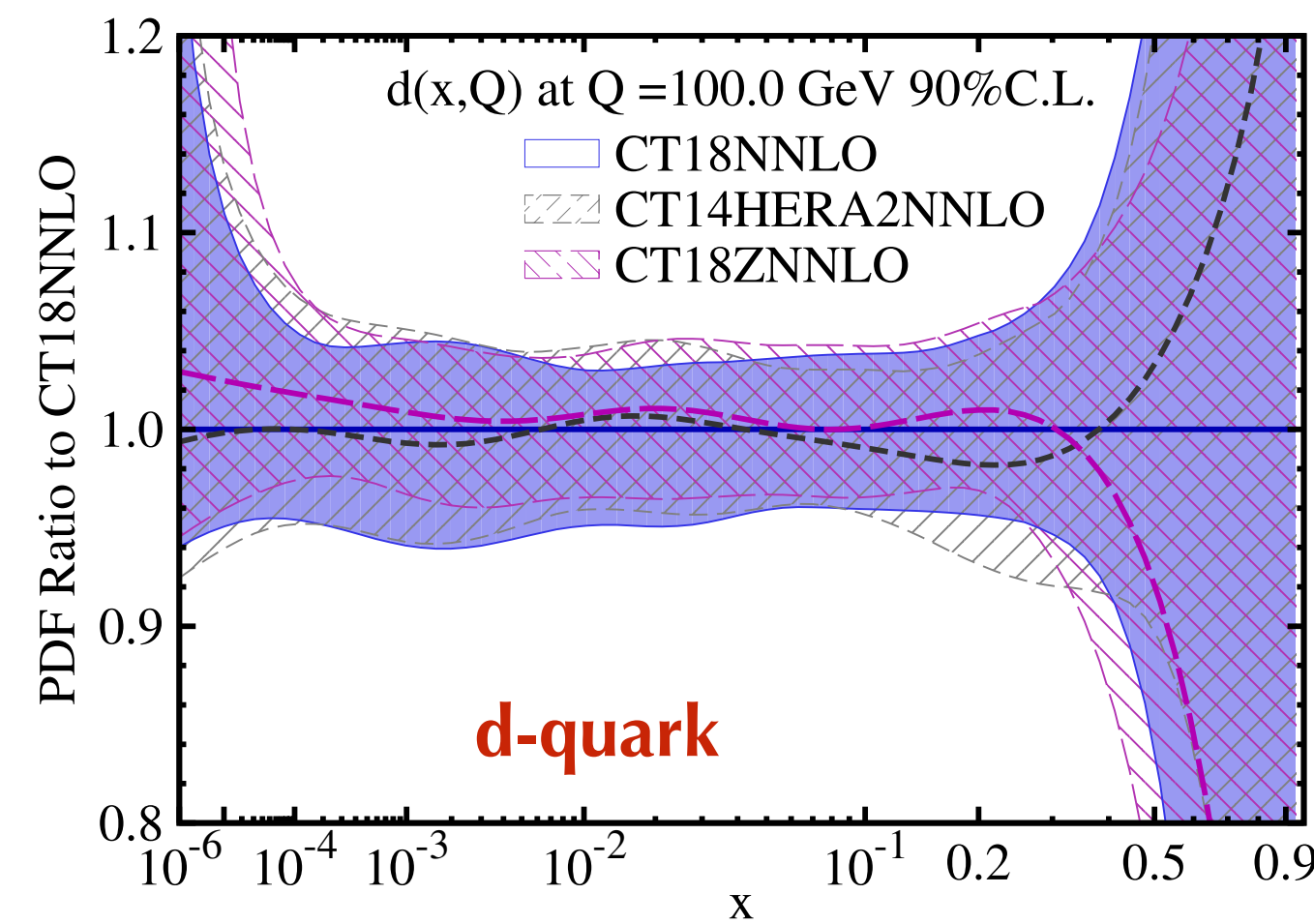
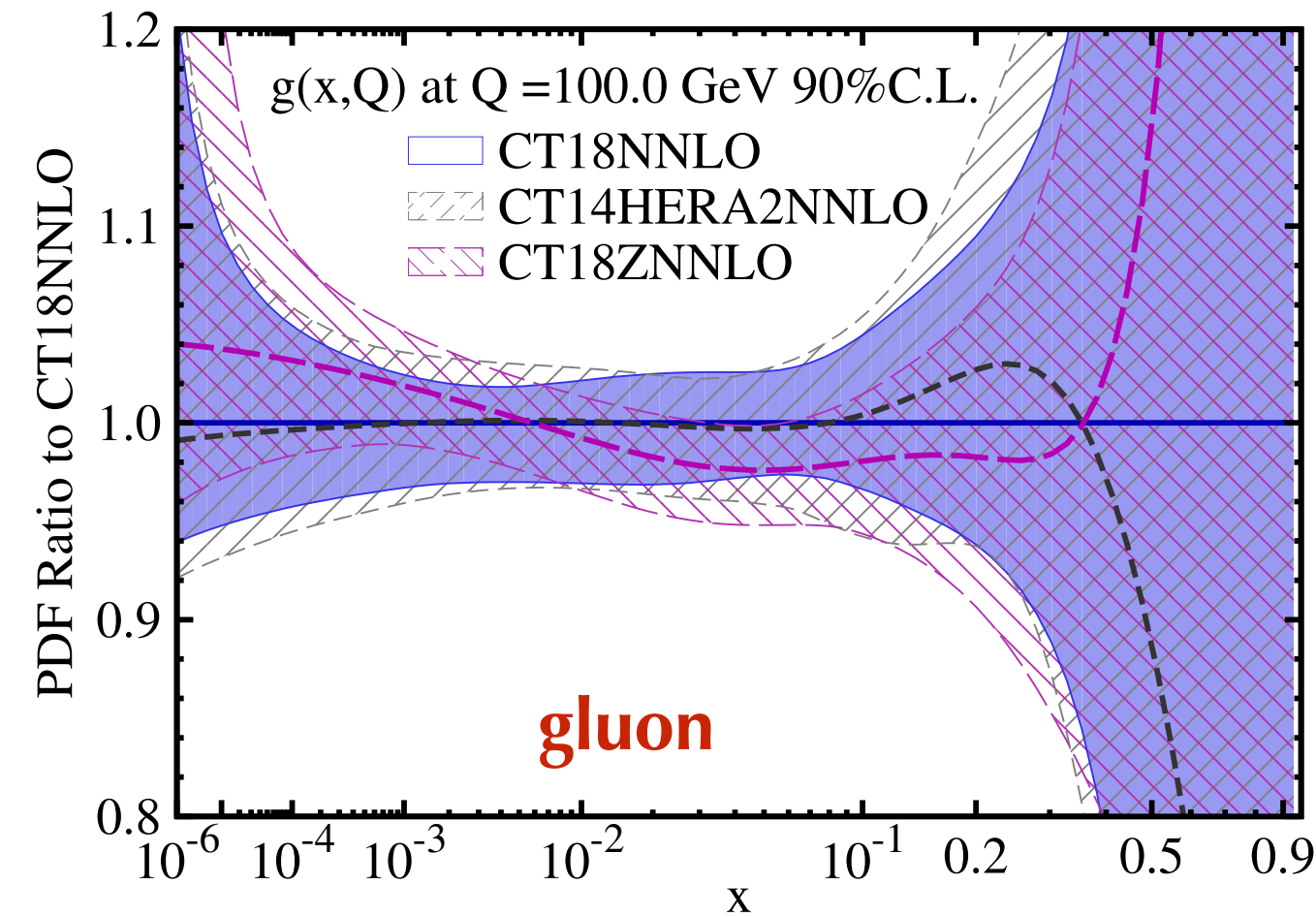
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Summary

- ✦ We developed a framework of global analysis for efficient evaluation on uncertainties of QCD inputs and BSM parameters using machine learnings and neural networks
- ✦ PDF uncertainties are one of the dominant theoretical uncertainties in direct measurements of the W-boson mass; variations due to PDFs are much smaller than discrepancies of the CDF measurement
- ✦ Strange-quark PDFs are slightly suppressed at $x \sim 0.02$ as now supported by both DIS and LHC data
- ✦ Correlations between gluon PDFs and BSM effects in global analyses with top-quark pair and jet production data are mild, indicating robustness of the global analyses of QCD
- ✦ Top-quark mass determinations from measurements of total and differential cross sections are consistent with those determinations based on kinematic reconstructions

CTEQ-TEA PDFs

- CT18 PDFs show moderate reductions of PDF uncertainties due to new LHC data sets, and agree with previous CT14 within uncertainties; alternative fits CT18Z/A/X for evaluation of certain systematic effects



- CT18 vs CT14: gluon unc. reduced everywhere (jets, Z pT, top); d-quark unc. reduced at $x \sim 0.2$ (LHCb W/Z); s-quark almost unchanged

- ATLAS 7 TeV W/Z data are not included in CT18 fit but in CT18A; CT18X uses a x -dependent scale in DIS to mimic small- x resummations

- CT18Z includes both variations, differences wrt. CT18 are most significant in s-quark and gluon/sea-quarks

[CT18, 1912.10053]