天文探测暗物质粒子

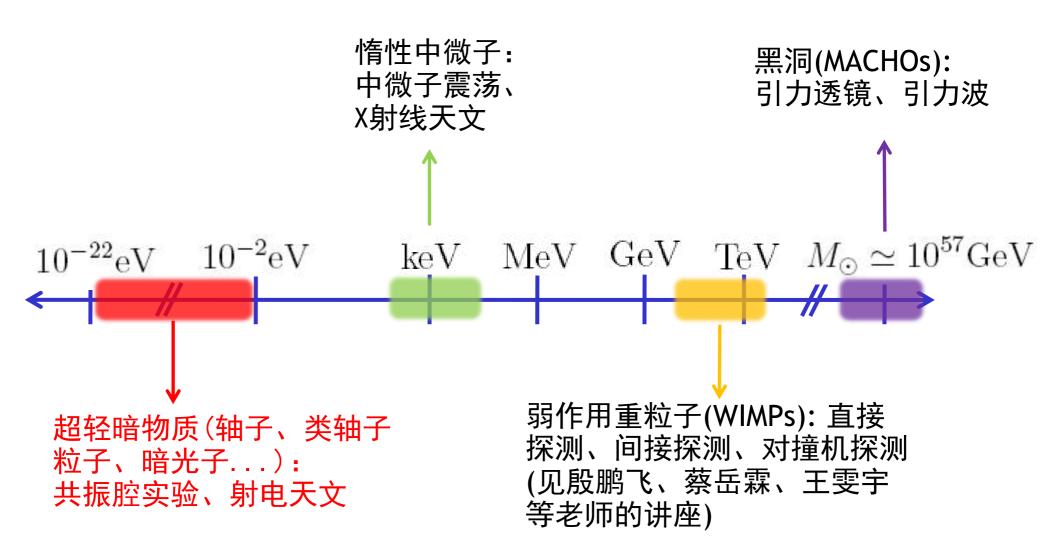
袁强

中国科学院紫金山天文台

2022.7.20-21

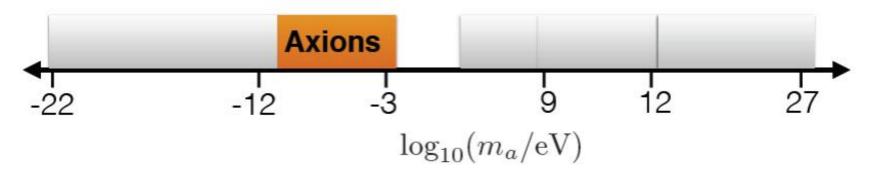
2022年理论物理前沿讲习班: 暗物质与新物理暑期学校

理论学家眼中的暗物质



暗物质的实验探测极具挑战性,天文观测在多种形式暗物质 候选体探测中起着非常重要的作用

轴子和类轴子粒子



Invented to solve strong CP problem in Standard Model

$$\mathcal{L}_{\rm axion} = -\left(\bar{\theta} + \frac{a}{f_a}\right) \frac{g^2}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu} \qquad \text{(neutron EDM} \propto \bar{\theta}\text{)}$$

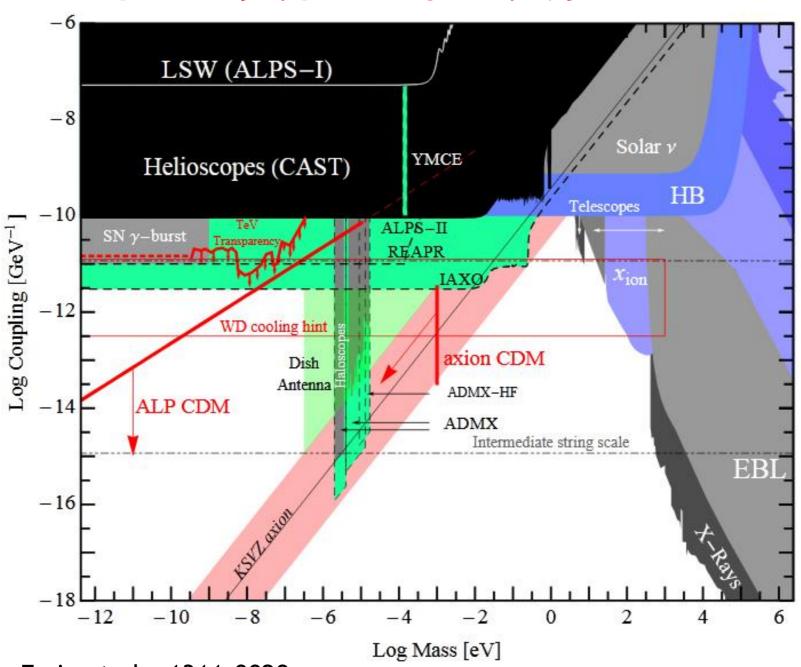
QCD effect gives axion mass

$$m_a pprox rac{f_\pi}{f_a} m_\pi pprox 10^{-9} \ \mathrm{eV} \left(rac{10^{16} \ \mathrm{GeV}}{f_a}
ight)$$

Axion couples with electromagnetic field

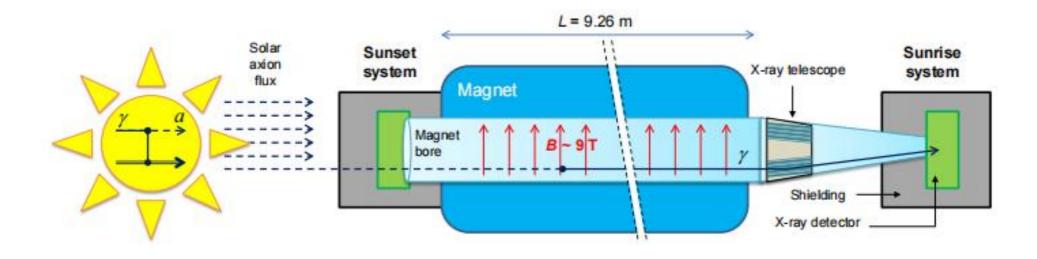
$$\mathcal{L} = -rac{1}{4} rac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} ilde{F}^{\mu
u} \qquad rac{g_{a\gamma\gamma}}{f_a} \propto rac{lpha_{\mathsf{EM}}}{f_a}$$

轴子/类轴子探测实验汇总



Essig et al., 1311.0029

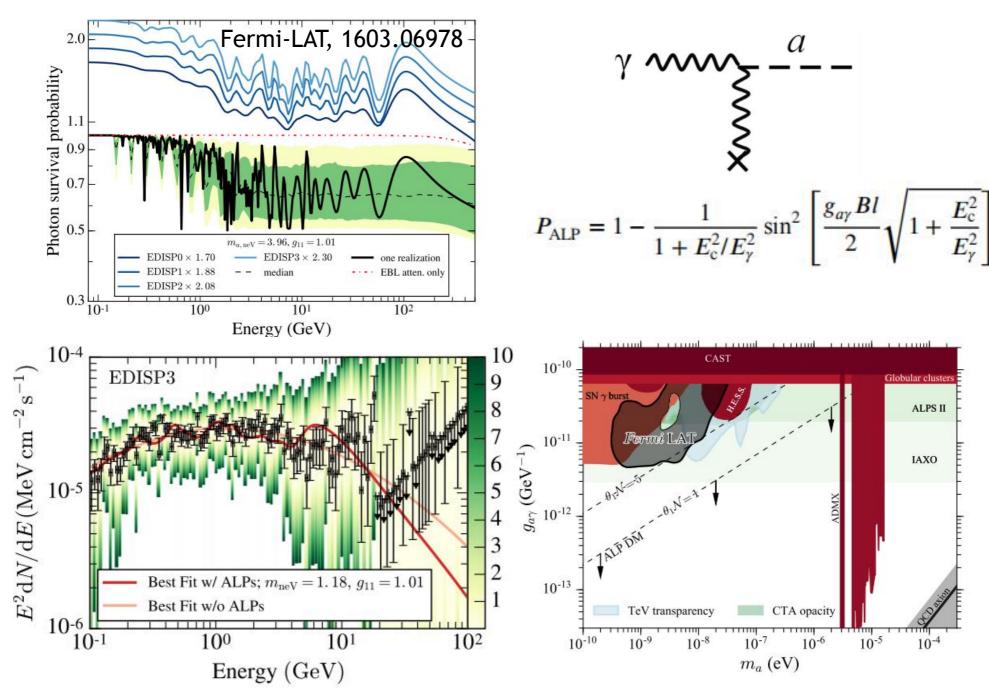
轴子太阳望远镜CAST



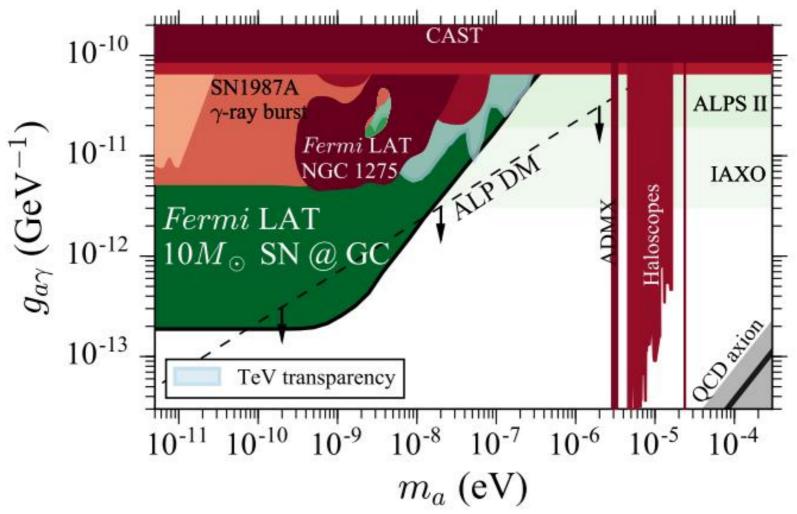
- ➤ 太阳内部存在大量的等离子体和光子,温度大约1-10 keV
- > 光子在带电粒子库伦场中通过Primarkoff过程转化成为轴子

$$\frac{d\Phi_{a}}{dE} = 6.02 \times 10^{10} \text{ cm}^{-2} \text{ s}^{-1} \text{ keV}^{-1} g_{10}^{2} E^{2.481} e^{-E/1.205}$$

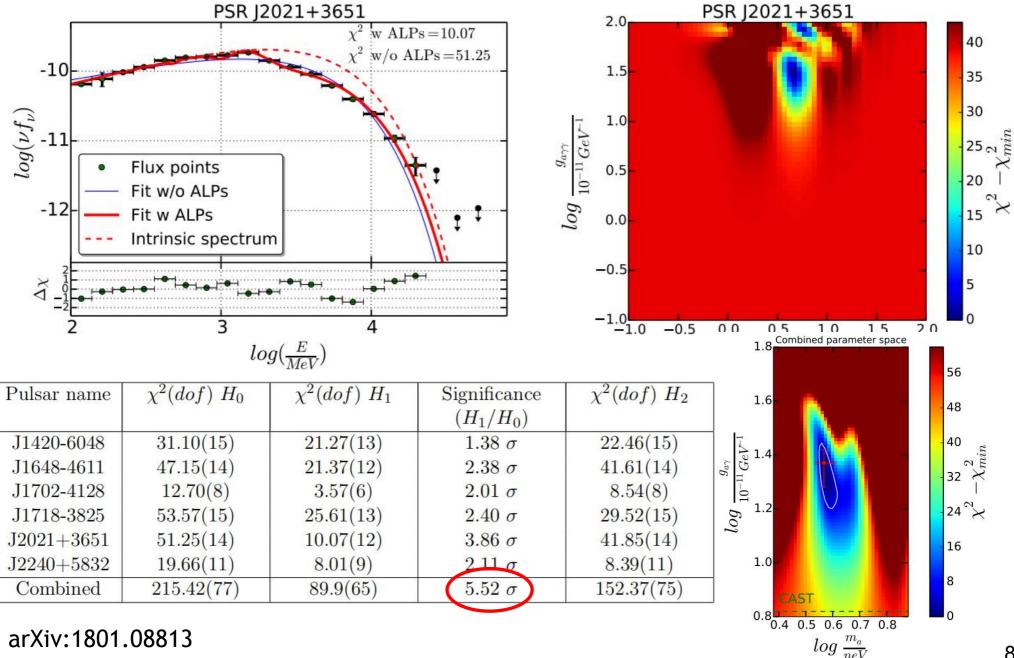
- ➤ 在地球上加设磁场将太阳轴子再次转化成为X射线光子
- ➤ 通过X射线探测器加以探测



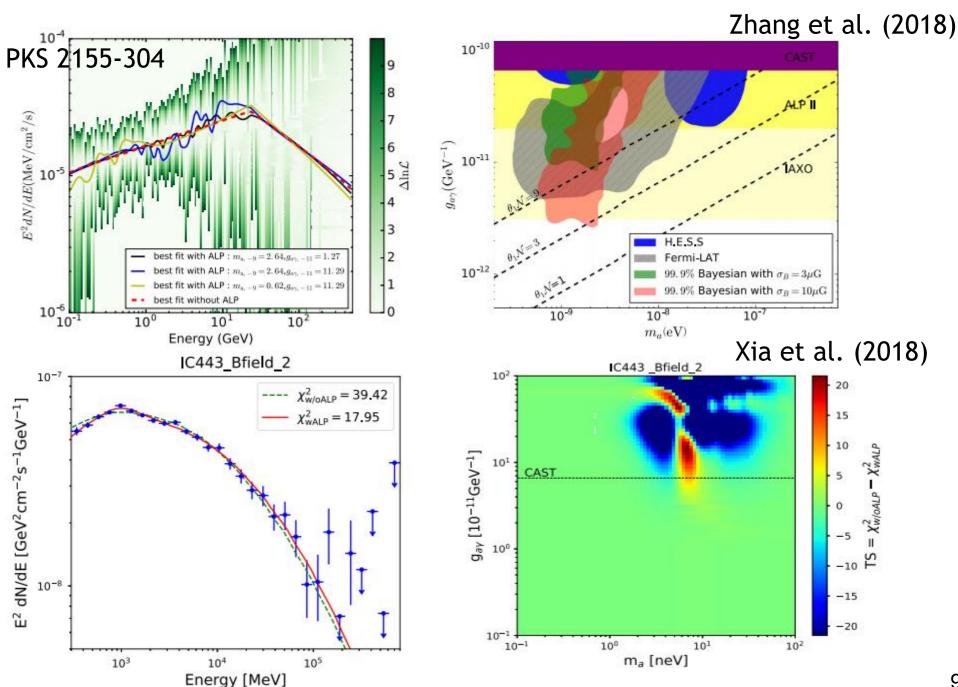
arXiv:1609.02350



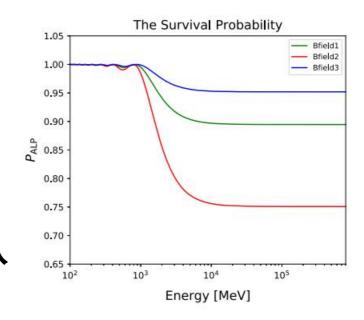
假设银河系中发生超新星爆发,Fermi-LAT将可以非常灵敏 地探测来自轴子-光子振荡产生的伽马射线



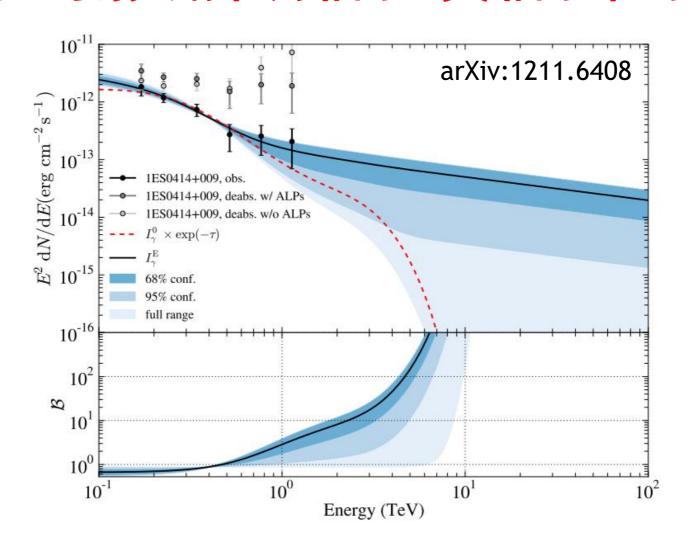
arXiv:1801.08813



- 轴子-光子振荡在伽马射线能谱上的 影响主要是"压低"和"不规则波动"
- "压低"效应通常会和源的内禀性质以及传播吸收效应等简并,难以证认

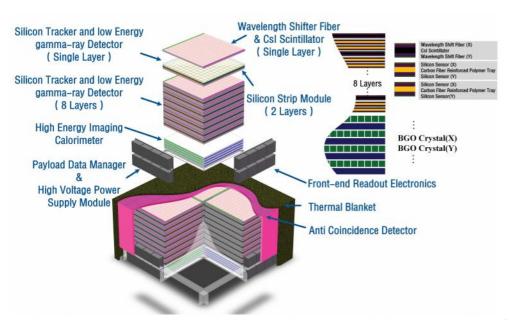


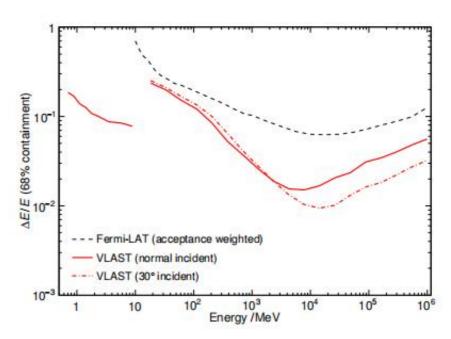
- "不规则波动"对能量分辨率和数据统计量要求很高, Fermi卫星统计量大但能量分辨率较差(~10%), DAMPE卫星能量分辨率很高(~1%)但统计量较低
- > 天体磁场模型也是一项主要的系统误差来源



探测河外辐射源在超高能伽马射线波段的"反常超出"可能是一个行之 有效的方法:能量越高由红外光学背景导致的吸收不透明度越高,光 子-轴子转换可以有效增加伽马射线传播距离 doi: 10.15940/j.cnki.0001-5245.2022.03.002

甚大面积伽马射线空间望远镜计划*

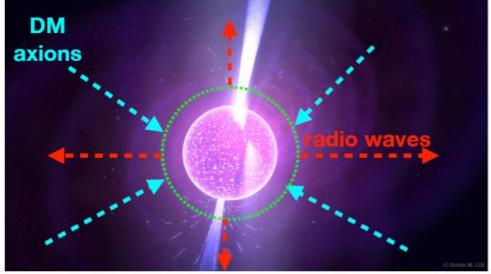




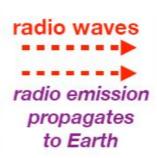
有效接受度达10 m²sr的旗舰型空间探测设施

中子星射电谱线观测轴子

NS with strong B-field and surrounding plasma

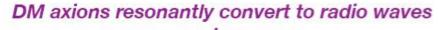


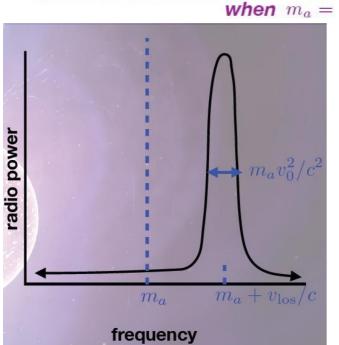
射电望远镜的宽频覆盖可以有效探索地面共振腔实验的盲区

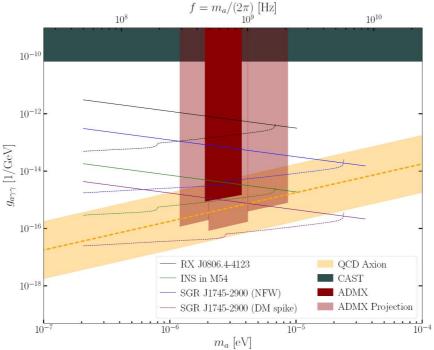




Narrow radio line detectable at Earth with $f=m_a/(2\pi)$.







Pshirkov+ JETP 2009; Huang+ PRD 2018; Hook+ PRL 2018

中子星射电谱线观测轴子

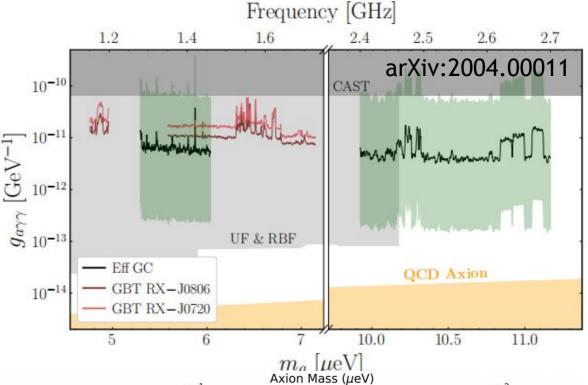
- \succ 信号特征:近似单一频率的线谱辐射, $f=m_a$
- 信号强度依赖于中子星磁场、周期、轴子-光子耦合、轴子暗物质密度等

$$\begin{split} \overline{S}_{\nu_{i}} &= \frac{F}{\Delta \nu} = 3.8 \times 10^{-6} \text{ Jy} \left(\frac{100 \text{ pc}}{d}\right)^{2} \left(\frac{16 \text{ kHz}}{\Delta \nu}\right) \\ &\times \left(\frac{d\mathcal{P}/d\Omega}{5.7 \times 10^{9} \text{W}}\right) \int_{\nu_{i,\text{min}}}^{\nu_{i,\text{max}}} \frac{d\nu}{\sqrt{2\pi}\sigma_{0}} e^{-\frac{(\nu - m_{a})^{2}}{2\sigma_{0}^{2}}}, \\ \frac{d\mathcal{P}}{d\Omega} &\simeq 5.7 \times 10^{9} \text{ W} \left(\frac{g_{a\gamma\gamma}}{10^{-12} \text{ GeV}^{-1}}\right)^{2} \left(\frac{r_{\text{NS}}}{10 \text{ km}}\right)^{5/2} \left(\frac{m_{a}}{\text{GHz}}\right)^{4/3} \\ &\times \left(\frac{B_{0}}{10^{14} \text{ G}}\right)^{5/6} \left(\frac{P}{\text{sec}}\right)^{7/6} \left(\frac{\rho_{\text{DM}}^{\infty}}{0.45 \text{ GeV cm}^{-3}}\right) \left(\frac{M_{\text{NS}}}{M_{\odot}}\right)^{1/2} \\ &\times \left(\frac{200 \text{ km s}^{-1}}{\nu_{0}}\right) \frac{3 \left(\hat{\mathbf{m}} \cdot \hat{\mathbf{r}}\right)^{2} + 1}{\left|3 \cos \theta \, \hat{\mathbf{m}} \cdot \hat{\mathbf{r}} - \cos \theta_{\text{m}}\right|^{7/6}}, \end{split}$$

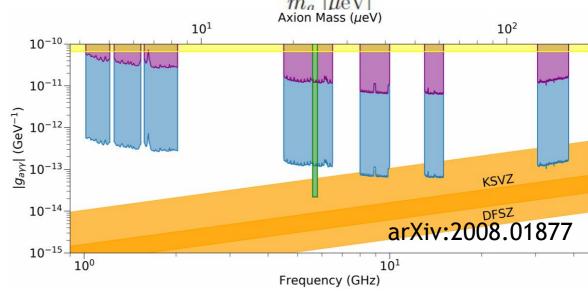
▶ 孤立、射电暗弱、强磁场中子星是优选目标

中子星射电谱线观测轴子

GBT和Effelsberg观测银心磁星和近邻孤立中子星

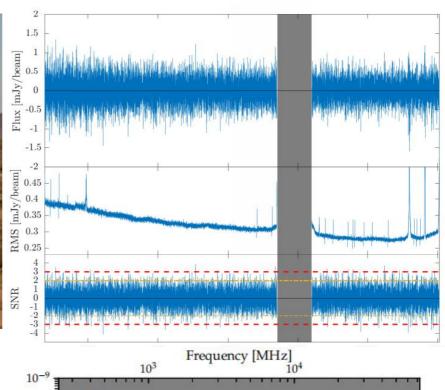


➤ VLA观测银心磁星PSR J1745-2900

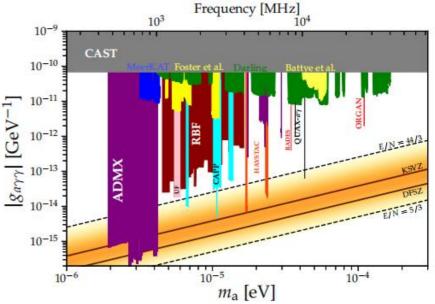


MeerKAT观测孤立中子星



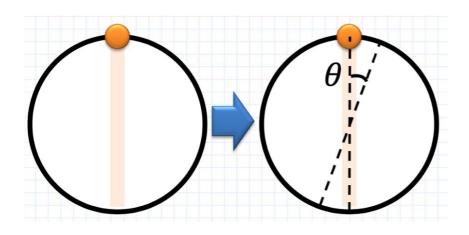


- ➤ MeerKAT UHF band: 580-1051 MHz covers the gap between ADMX and RBF
- Reaches better sensitivity than other observations at overlapping frequency



射电偏振辐射探测轴子

arXiv:1901.10981



$$\omega_{\pm} \approx k \pm \frac{1}{2} g \left(\frac{\partial \varphi}{\partial t} + \nabla \varphi \cdot \frac{\mathbf{k}}{k} \right)$$

线偏振光在轴子背景场 中传播,会由于双折射 效应产生偏振角的转动

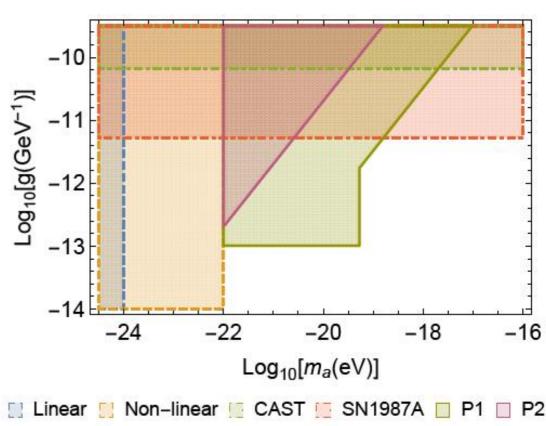


FIG. 1: Projected sensitivities to detect the CAB, using linearly polarized pulsar light as a probe, in the two benchmark scenarios: P_1 and P_2 .

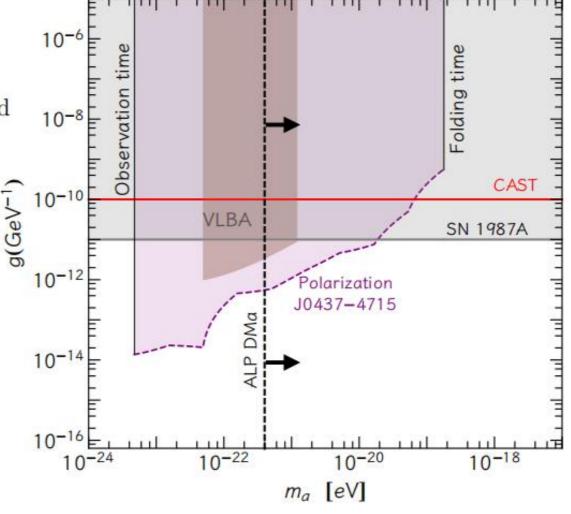
射电偏振辐射探测轴子

arXiv:1902.02695

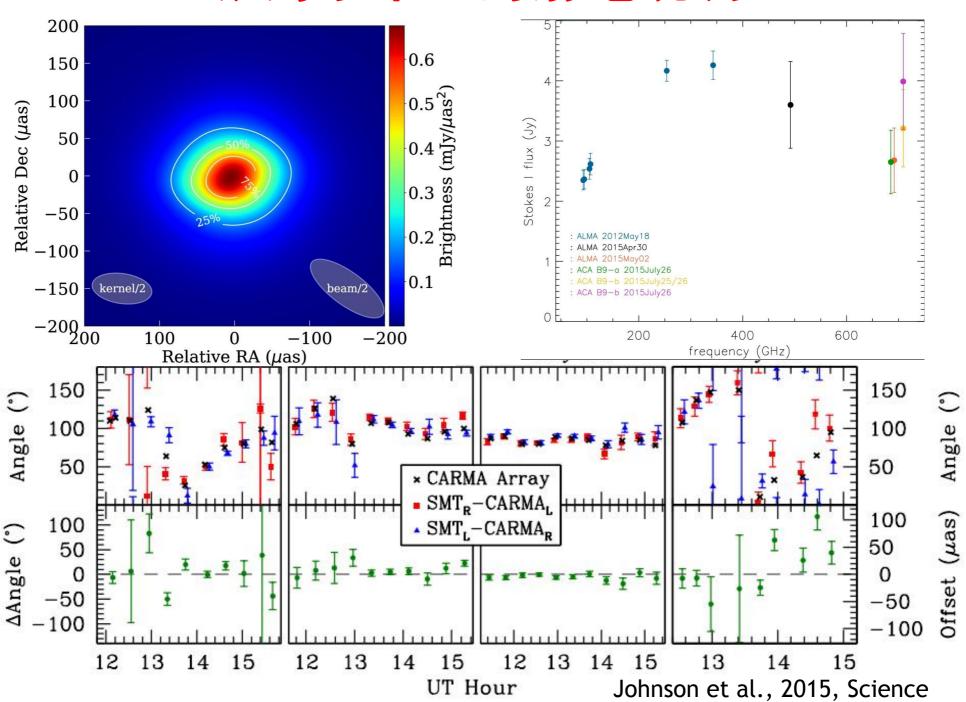
$$\theta(t,T) \sim 1.4 \times 10^{-2} \sin(m_a t + \delta)$$

$$\left(\frac{g}{10^{-12} \,\text{GeV}^{-1}}\right) \frac{10^{-22} \,\text{eV}}{m_a} \text{ rad}$$

Parkes对PSR J0437-4715 的偏振观测对超轻ALP给 出严格限制

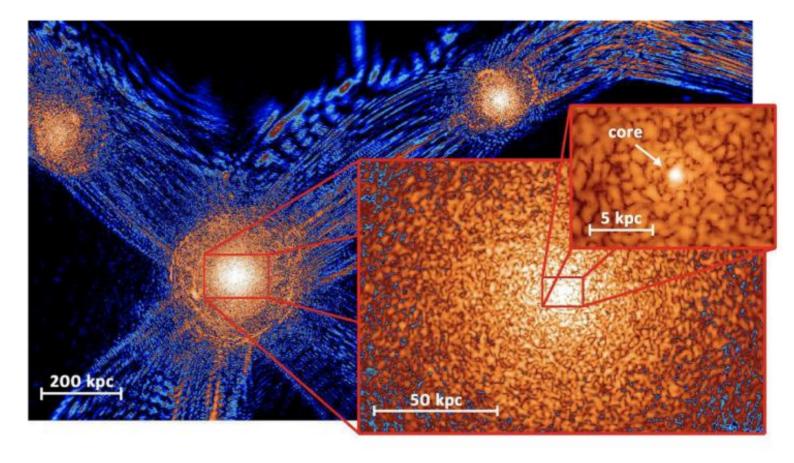


银河系中心的射电观测



银心暗物质分布

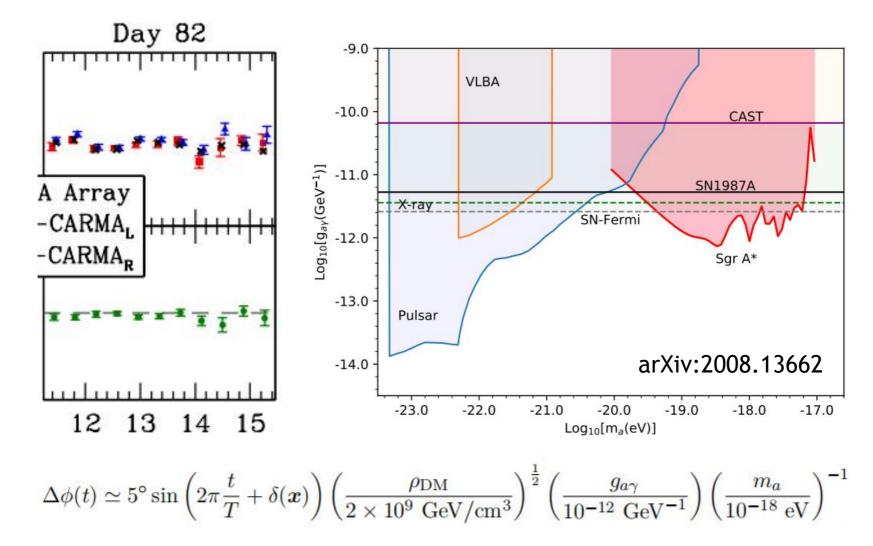
质量非常轻的 暗物质波动性 明显,影响其 在小尺度上的 结构分布



$$\rho_{\rm DM} = \begin{cases} 190 \times \left(\frac{m_a}{10^{-18} {\rm eV}}\right)^{-2} \left(\frac{r_c}{1 {\rm pc}}\right)^{-4} M_{\odot} {\rm pc}^{-3}, & \text{for } r < r_c \\ \frac{\rho_0}{r/R_g (1+r/R_g)^2}, & \text{for } r > r_c \end{cases}$$

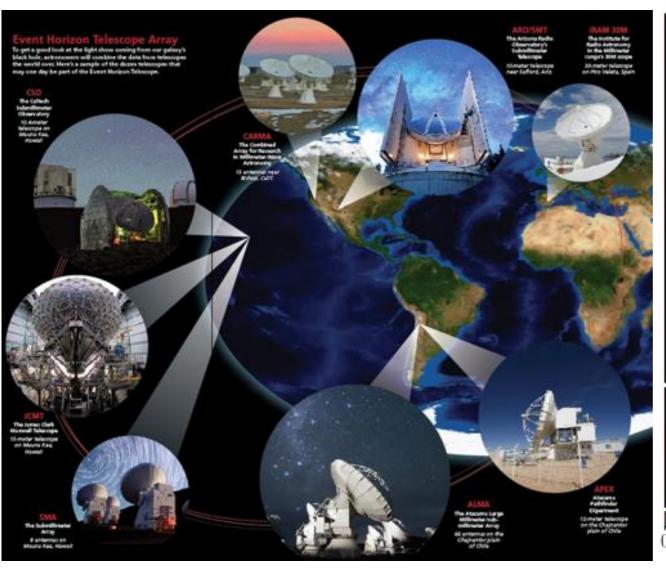
Soliton core + NFW

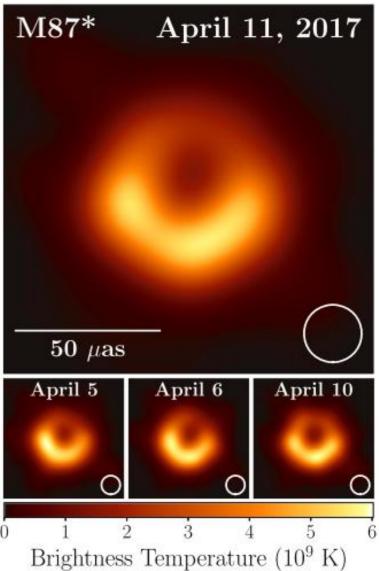
银心射电观测限制轴子暗物质模型



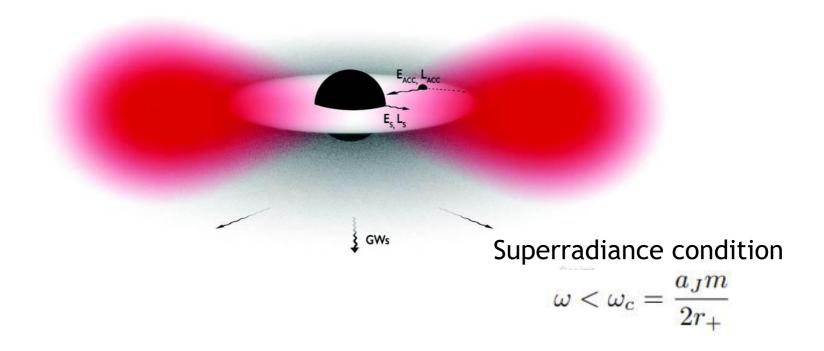
- VLBI对银心的偏振观测发现偏振度和偏振角呈复杂变化,反映了 吸积盘的不稳定性
- ▶ 利用其相对稳定的一天观测结果(Day 82)可以对ALP模型给出限制

事件视界望远镜



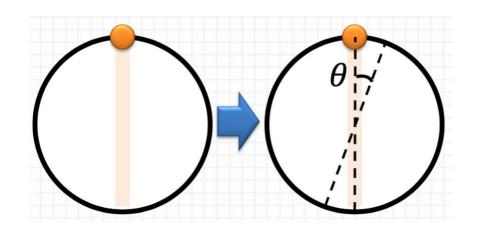


黑洞超辐射



- 满足超辐射条件的玻色子在克尔黑洞外围将增长形成玻色云; 类似于氢原子外的电子云,也称作引力原子
- ▶ 玻色云的演化(增长、坍缩、外流等)消耗黑洞自转能
- > 玻色云的总能量可以很大,和黑洞质量相比拟

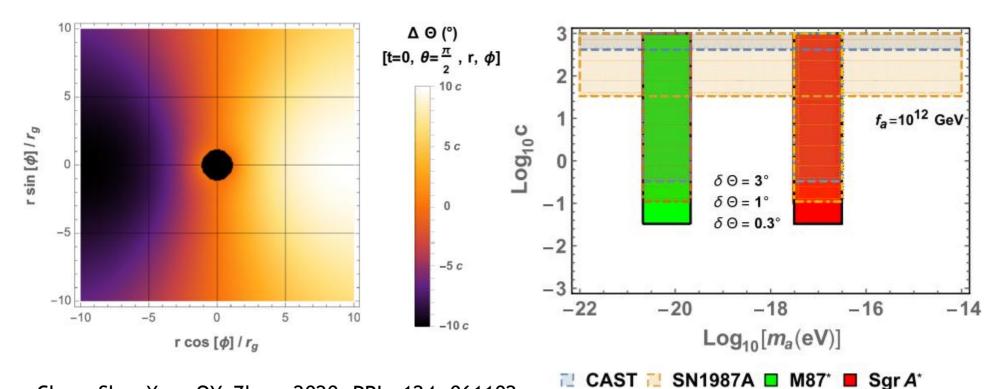
EHT偏振探测轴子/类轴子



$$\Delta\Theta = g_{a\gamma} \Delta a(t_{\text{obs}}, \mathbf{x}_{\text{obs}}; t_{\text{emit}}, \mathbf{x}_{\text{emit}})$$

$$= g_{a\gamma} \int_{\text{emit}}^{\text{obs}} ds \ n^{\mu} \ \partial_{\mu} a$$

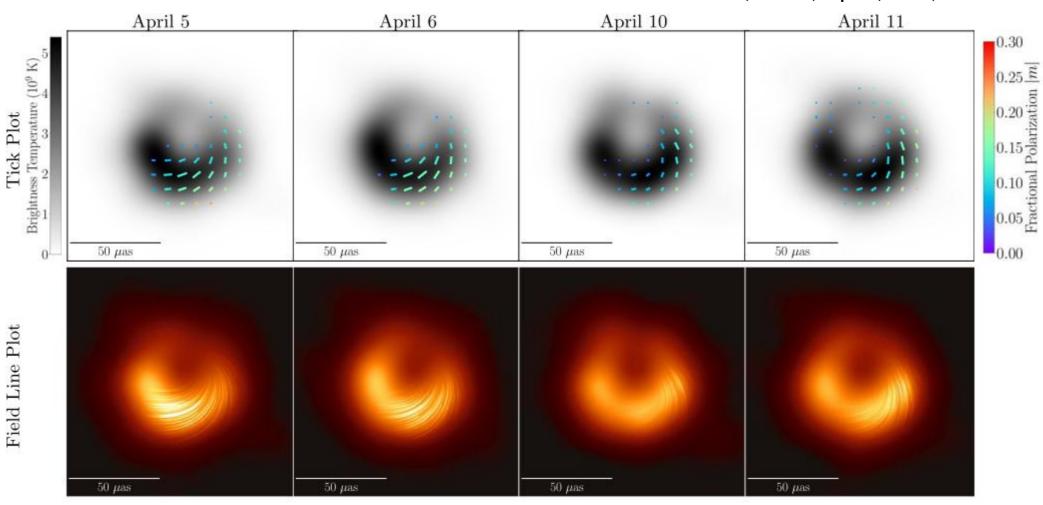
$$= g_{a\gamma} [a(t_{\text{obs}}, \mathbf{x}_{\text{obs}}) - a(t_{\text{emit}}, \mathbf{x}_{\text{emit}})],$$



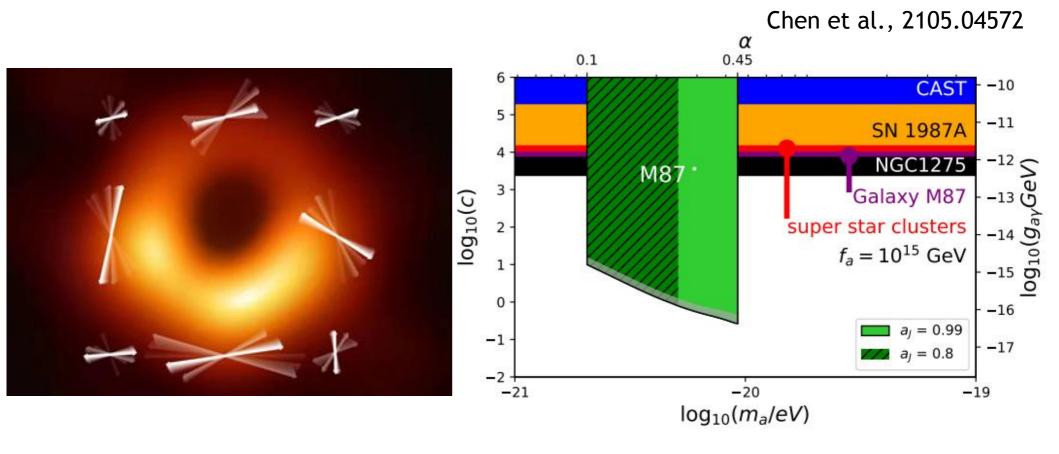
Chen, Shu, Xue, QY, Zhao, 2020, PRL, 124, 061102

EHT M87*偏振成像观测

EHT Coll., 2021, ApJL, 910, L12



EHT偏振探测轴子/类轴子



- ▶ 时间上周期性变化;方位角周期性变化
- ➤ EHT超高分辨的偏振成像正好可以用于探测此类效应

超轻"模糊"暗物质(fuzzy dark matter)

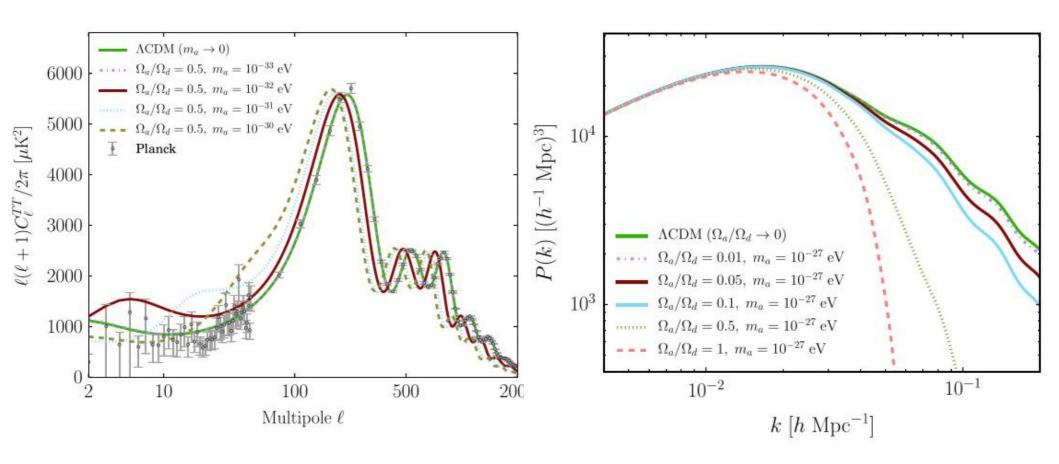
- 超轻的玻色暗物质以玻色-爱因斯坦凝聚状态存在,在大尺度 上可以类似冷暗物质
- ➤ 质量非常轻的粒子(~10⁻²² eV)更像一个波,波长相当于矮星系的尺度(kpc),抹平小尺度结构(Hu et al., 2000),解决冷暗物质模型的"小尺度危机"

$$i(\partial_t + \frac{3}{2}\frac{\dot{a}}{a})\psi = (-\frac{1}{2m}\nabla^2 + m\Psi)\psi$$

$$i\partial_t \psi = (-\frac{1}{2m}\nabla^2 + m\Psi)\psi, \quad \nabla^2 \Psi = 4\pi G\delta\rho$$

$$r_{Jh} \sim 3.4 (c_{10}/f_{10})^{1/3} m_{22}^{-2/3} M_{10}^{-1/9} (\Omega_m h^2)^{-2/9} \text{kpc}$$

FDM的宇宙学效应

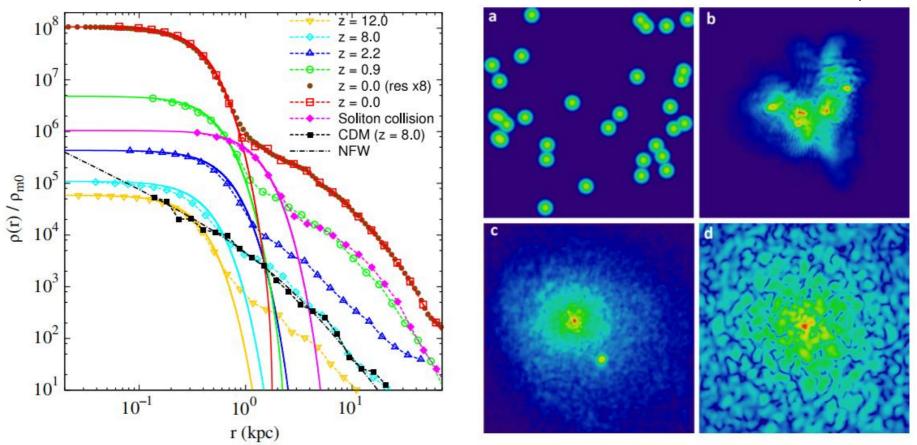


超轻暗物质($w\sim -1$)影响宇宙的膨胀率,从而影响CMB和物质功率谱;物质功率谱上主要体现为对小尺度结构的压低

Marsh, 2016, Phys. Rept.

FDM小尺度结构

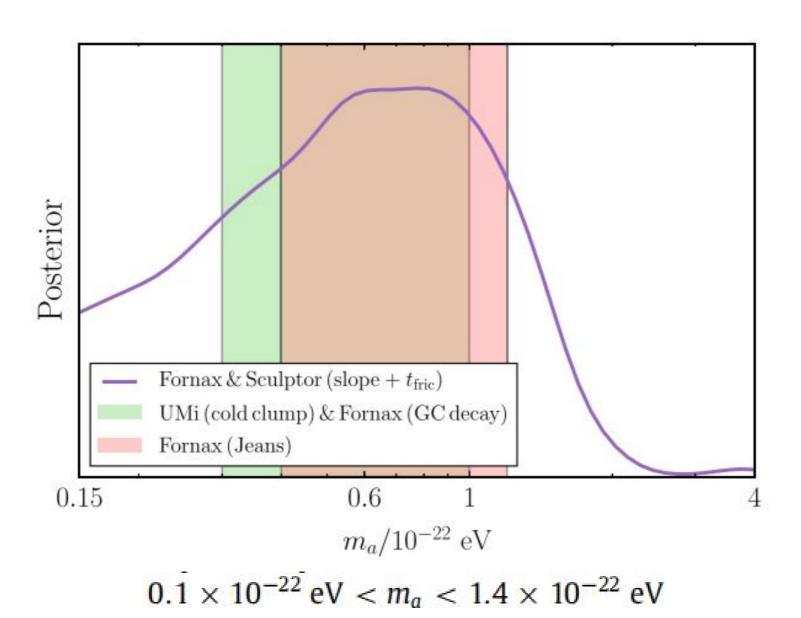
Schive et al., 2014, PRL



FDM形成的暗晕在小尺度表现为soliton core结构,大尺度和冷暗物质类似,可以解决冷暗物质模型的"小尺度危机":

- The missing satellites problem [267,268]: CDM predicts more small Milky Way satellites than are observed.
- The too-big-to-fail problem [269]: CDM predicts more massive satellites that should contain stars than are observed.
- The cusp-core problem [270]: many observed low-mass systems contain flat central density profiles, not NFW cusps.

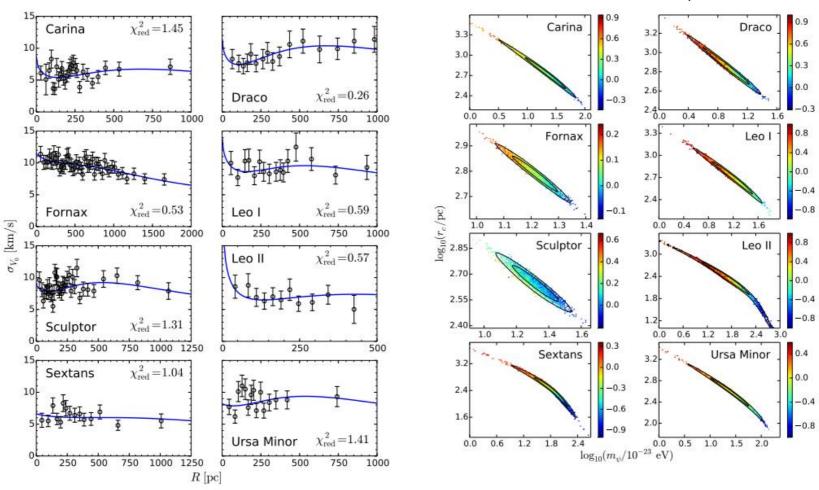
FDM解释矮星系结构



Marsh, 2016, Phys. Rept.

FDM解释矮星系结构

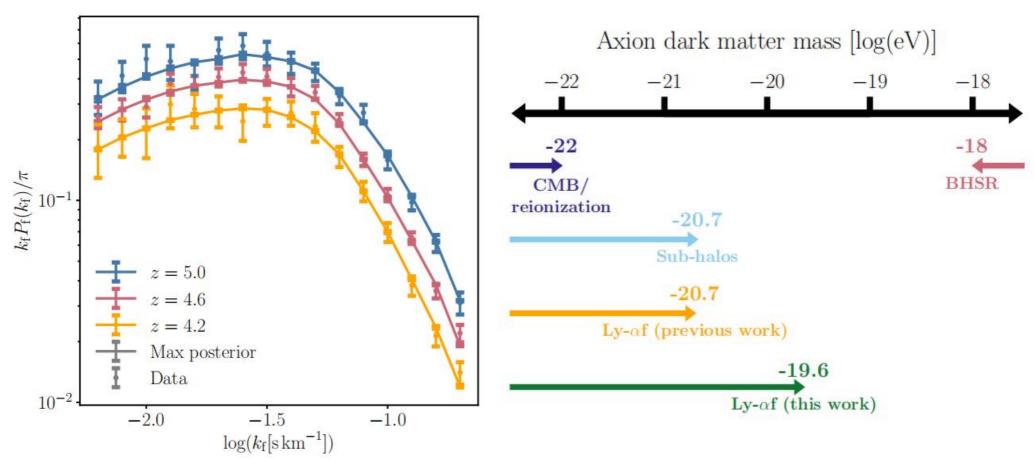
Chen et al., arXiv:1606.09030



pressure against gravity. Here we apply Jeans analysis assuming a soliton core profile to the kinematic data of eight classical dSphs so as to constrain m_{ψ} , and obtain $m_{\psi} = 1.18^{+0.28}_{-0.24} \times 10^{-22} \,\text{eV}$ and $m_{\psi} = 1.79^{+0.35}_{-0.33} \times 10^{-22} \,\text{eV}$ (2 σ) using the observational data sets of Walker et al. (2007) and Walker et al. (2009b), respectively. We

Lya森林的限制

arXiv:2007.12705



We present a new bound on the ultra-light axion (ULA) dark matter mass m_a , using the Lymanalpha forest to look for suppressed cosmic structure growth: a 95% lower limit $m_a > 2 \times 10^{-20} \,\text{eV}$.

Lyα森林和矮椭球星系的结果稍微存在矛盾

Ultralight scalars as cosmological dark matter

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Department of Physics, Columbia University, New York, NY 10027

Jeremiah P. Ostriker

Department of Astronomy, Columbia University, New York, NY 10027 and
Department of Astrophysical Sciences, Princeton University, Princeton, NJ 08544

Scott Tremaine and Edward Witten Institute for Advanced Study, Princeton, NJ 08540

In summary, the hypothesis that the principal component of the ubiquitous dark matter is an ultra-light axion is an attractive and testable alternative to CDM, having no serious inconsistencies with current data if the particle mass $m \gtrsim 10^{-22}\,\mathrm{eV}$. There are significant and attractive observational consequences if the mass is in the range $1-10\times 10^{-22}\,\mathrm{eV}$. There is tension with observations of the Lyman- α forest, which favor masses $10-20\times 10^{-22}\,\mathrm{eV}$ or higher. More sophisticated calculations of reionization and of structure formation in FDM are required to determine whether this variety of constraints is consistent with observations.

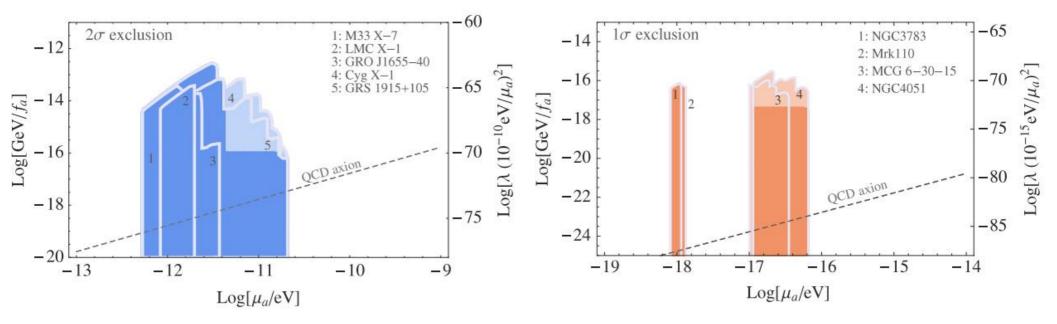
黑洞自旋和超轻暗物质

- 满足一定条件的玻色粒子会通过提取黑洞转动能形成玻色云, 称为超辐射(Penrose过程)
- \triangleright 考虑玻色子的自相互作用,玻色云密度增长至某临界值后会发生坍缩(bosenova),最大容许玻色云粒子数变为 $N_{bosenova}$
- 通过测量黑洞自旋可以对该过程给出限制

$$N_{
m max} \simeq rac{GM_{
m bh}^2}{\mu} \Delta a_*,$$

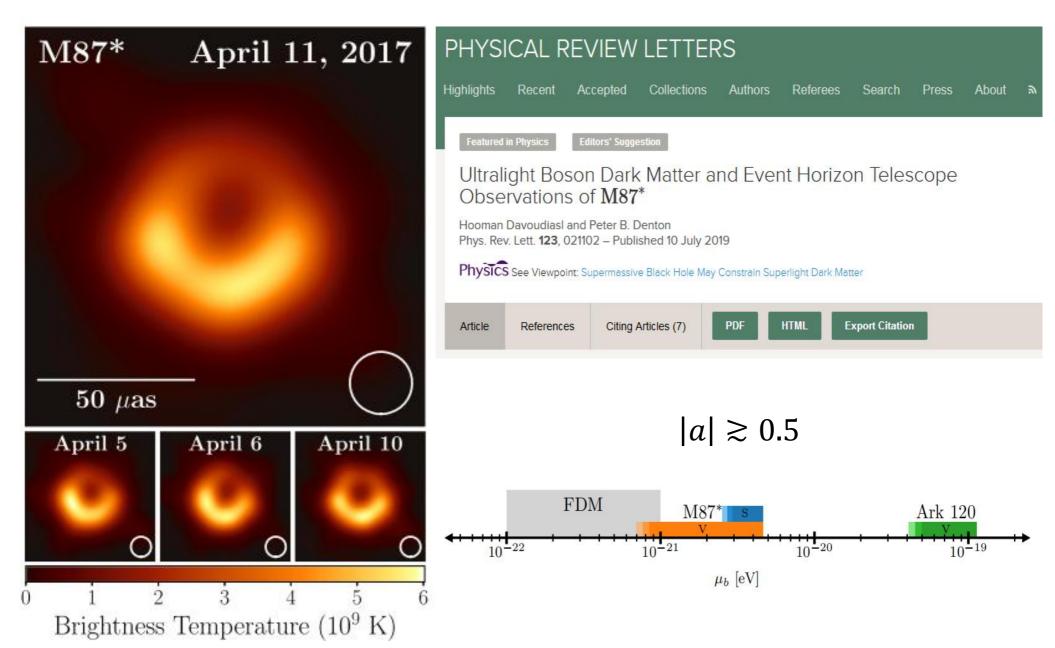
$$\Gamma \tau_{\rm bh} > \ln N_{\rm max}$$

$$\Gamma \tau_{\rm bh}(N_{\rm bosenova}/N_{\rm max}) > \ln N_{\rm bosenova}$$



Marsh, 2016, Phys. Rept.

黑洞自旋和超轻暗物质



黑洞自旋和超轻暗物质

TABLE I: SMBHs with masses and spins measured with various method.

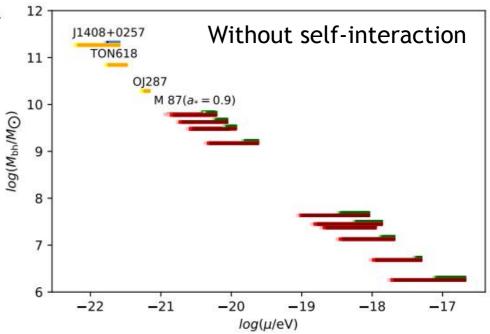
	Object	$M_{\rm bh}~(10^8{ m M}_\odot)$	Spin	Refs.
	Mrk 110	$0.251^{+0.061}_{-0.061}$	$0.96^{+0.03}_{-0.07}$	[38, 39]
	Mrk 335	$0.142^{+0.037}_{-0.037}$	> 0.91	[38, 40]
	NGC 3783	$0.298^{+0.054}_{-0.054}$	> 0.98	[38, 41]
	NGC 4051	$0.019^{+0.008}_{-0.008}$	> 0.99	[38] 42
	NGC 4151	$0.457^{+0.057}_{-0.047}$	> 0.90	[43] 44
	NGC 5506	$0.051^{+0.022}_{-0.012}$	$0.93^{+0.04}_{-0.04}$	[45] 46
UV/optical	J1152 + 0702	32.3+4.8	$0.998^{+0.000}_{-0.032}$	[22]
	J1158 - 0322	$31.6^{+4.7}_{-3.4}$	$0.898^{+0.036}_{-0.035}$	[22]
	J0941 + 0443	$44.7^{+5.4}_{-4.8}$	$0.998^{+0.000}_{-0.032}$	[22]
	J0303 + 0027	$63.1_{-6.9}^{+7.7}$	$0.998^{+0.000}_{-0.032}$	[22]
	J0927 + 0004	$15.8^{+1.9}_{-2.0}$	$0.998^{+0.000}_{-0.036}$	[22]
Most massive	OJ 287	THE CONTRACT WAS A SECOND STREET	0.381 ± 0.004	[33]
	Ton 618	660	> 0.60	[34]
	SDSS J140821.67	1060	0.07	1051
	+025733.2	1960	> 0.97	[35]

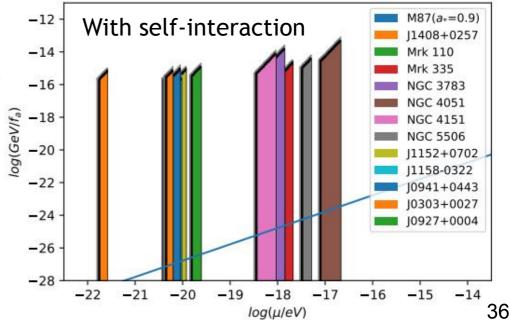
$$\Gamma_s = \frac{1}{24} a_* r_g^8 \mu^9,$$

$$\frac{dN}{dt} = \Gamma N.$$

$$\Gamma_v = 4 a_* r_g^6 \mu^7,$$

Zu et al., 2020, EPJP, 135, 709

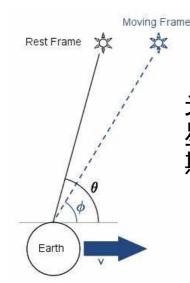




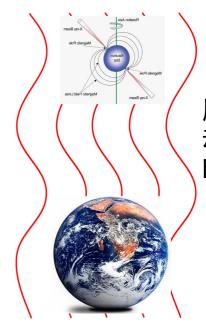
恒星位置和脉冲星计时探测超轻暗物质



质量非常轻(<10⁻²⁰ eV)的暗物质表现出宏观的波动性,其和物质的相互作用与引力波导致的时空扰动相似,可以在天体位置或者脉冲到达时间上留下印记

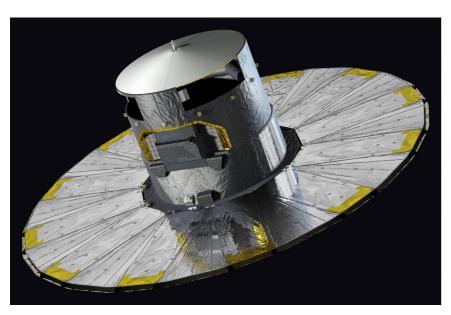


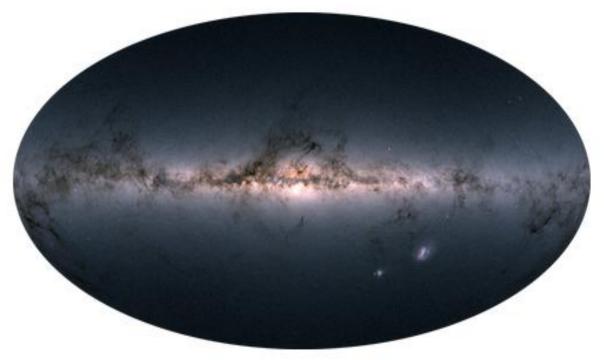
光行差导致恒 星位置产生周 期性额外摆动



脉冲星或地球运 动导致脉冲到达 时间超前或滞后

Gaia高精度天体位置测量





- Gaia将以非常高的精度(~100 μas)观测大量(~109)恒星的位置 和运动
- 将带来银河系结构和动力学、恒星物理、地外行星、基础物理等领域的革命性突破

超轻暗光子导致天体位置变化

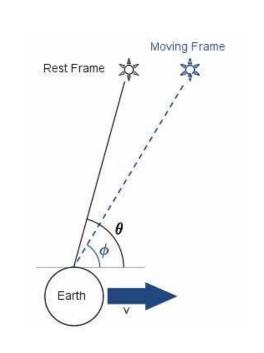
- ▶ 考虑一类暗光子模型:它们以微弱的耦合强度和重子数B或者 重子-轻子之差B-L (dark charge)不为零的物体相互作用
- ➤ Gaia卫星在暗光子场中会感受到一个周期性振荡的加速度

$$a(t, x) \simeq \epsilon e^{\frac{q}{m}} m_A A_0 \cos(m_A t - k \cdot x)$$

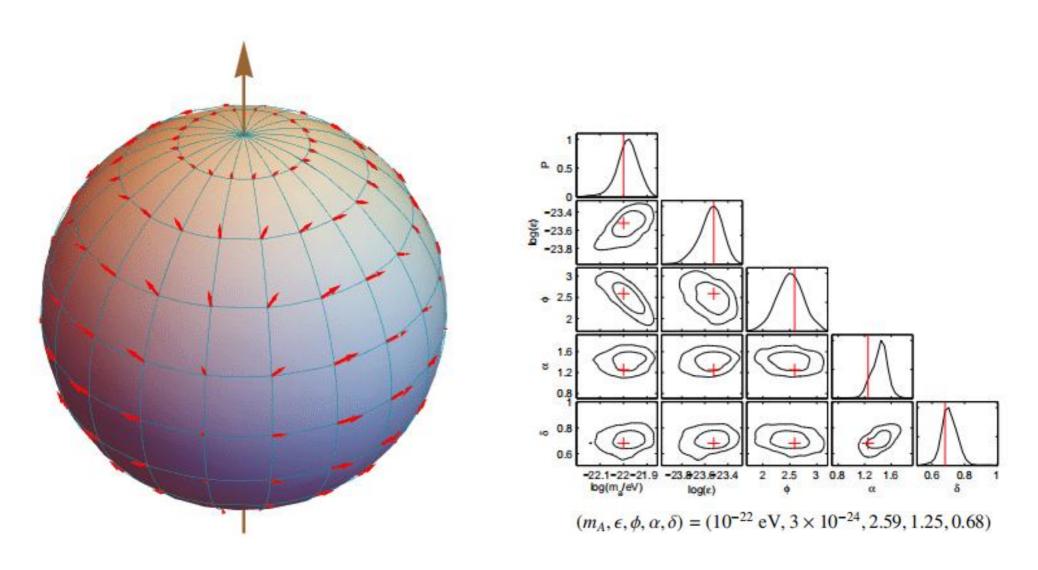
该加速度导致卫星产生一项额外的周期性运动,由于光行差现象使得远处的恒星视位置发生改变

$$\Delta v(t, \mathbf{x}) \simeq \epsilon e \frac{q}{m} \mathbf{A_0} \sin(m_A t - \mathbf{k} \cdot \mathbf{x}).$$
 $\Delta \theta \simeq -\Delta v \sin \theta$

大量恒星的运动会形成一个规律的图案, 反映暗光子场的性质

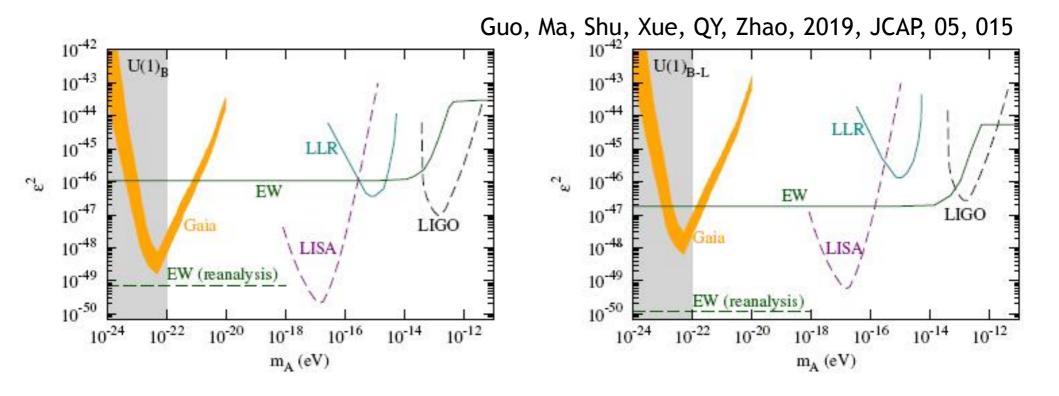


模拟和参数重建



Guo, Ma, Shu, Xue, QY, Zhao, 2019, JCAP, 05, 015

Gaia对超轻暗光子的灵敏度



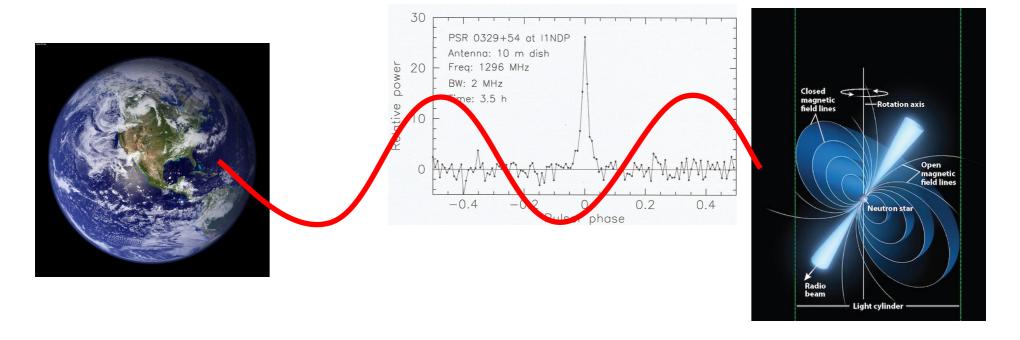
- 对m_A<10⁻²¹ eV区间将达到比实验室探测更好的灵敏度
- 很低质量区间(m_A<10⁻²² eV)灵敏度变差是由于长周期振荡和恒星自行的简并;高质量区间灵敏度变差是因为暗光子场振幅变小

PTA探测超轻暗光子

脉冲星和地球在暗光子场中均会产生振荡,导致脉冲到达时间出现残差:

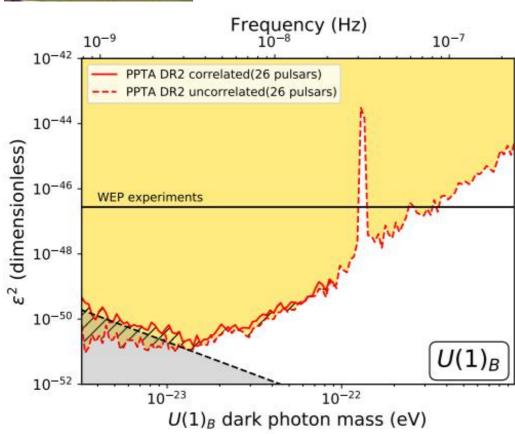
$$\delta x_{e,p}(t) \simeq -\frac{\epsilon eq}{m_A m} A_0^{e,p} \cos \left[m_A(t-t_0) + \alpha_{e,p} \right] \qquad \qquad \Delta t_r^d(t) = \frac{\left| d + \delta x_p \left(t - \frac{|d|}{v(t)} \right) - \delta x_e(t) \right| - |d|}{v(t)}$$

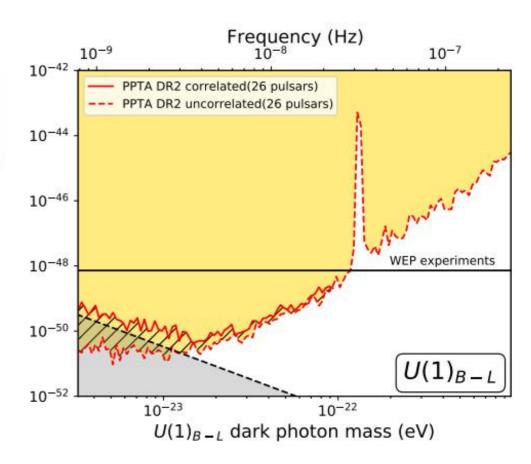
$$\simeq \frac{n_p \cdot \Delta x(t)}{v(t)},$$



暗光子模型参数限制







Xue et al., with PPTA (2022, Phys. Rev. Res., 4, L012022)

天文观测的发展趋势

测量准:空间分辨率高、时间分辨率高、能量分辨率高

样本多:大视场、大深度、高采样

> 手段全: 全波段、多信使

天文精测在暗物质粒子探测中发挥的重要作用 越来越得以凸显!

总结 & 展望

- 暗物质候选体形形色色,跨越非常大的能段,实验探测具有很大挑战
- 天文观测是通常粒子物理实验方法的重要补充,而且更为重要的是天文观测可以覆盖比实验室探测更宽的质量范围,探索更多的可能性
- 未来的天文学观测将更倾向于大数据、高精度、高分辨、 全波段,将在暗物质探测领域大有可为
- → 中国的机会: FAST、LHAASO、SKA、CSST、ngEHT、eXTP、VLAST