

量子计算与高能核物理交叉前沿讲习班 @ 华南师范大学, 广州  
2022年11月13日至28日

# Quantum Sensing and searches for new particles

## —— Sapphire: Spin Amplifier for Particle PHysics REsearches



Xinhua Peng

Nov. 22, 2022

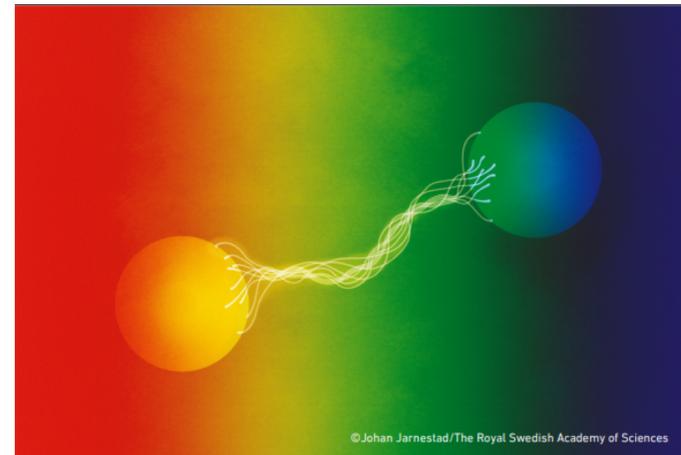
Univ. of Sci. & Tech of China (USTC)

CAS Key Laboratory of Microscale Magnetic Resonance

# Nobel Prize in Physics 2022

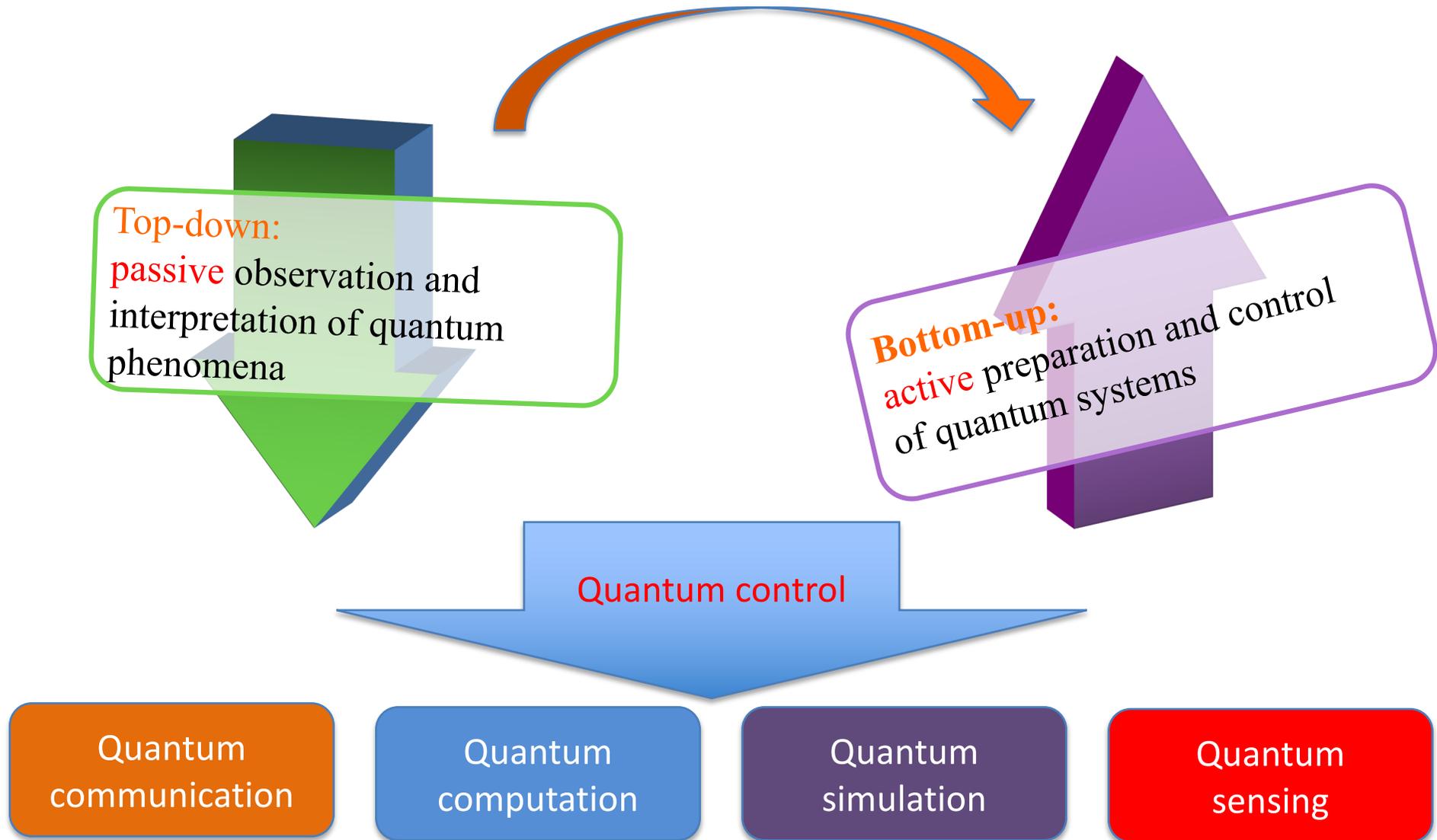


*"for experiments with entangled photons, establishing the violation of Bell inequalities and pioneering quantum information science"*



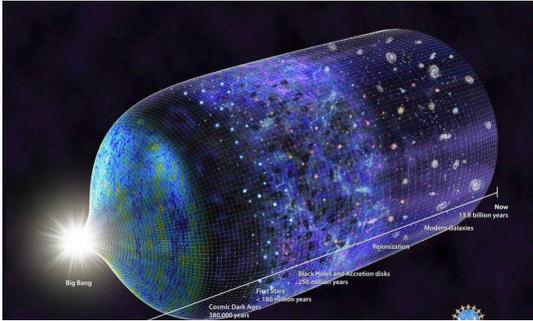
Alain Aspect, John Clauser and Anton Zeilinger have each conducted groundbreaking experiments using entangled quantum states, where two particles behave like a single unit even when they are separated. Their results have cleared the way for new technology based upon quantum information.

# The second quantum revolution

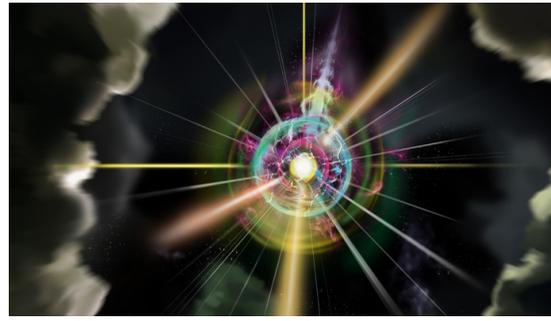


**Break the limits of classical technologies**

# Quantum precision measurement



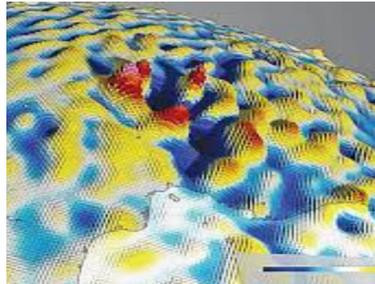
Dark energy



Monopoles



Dark matter



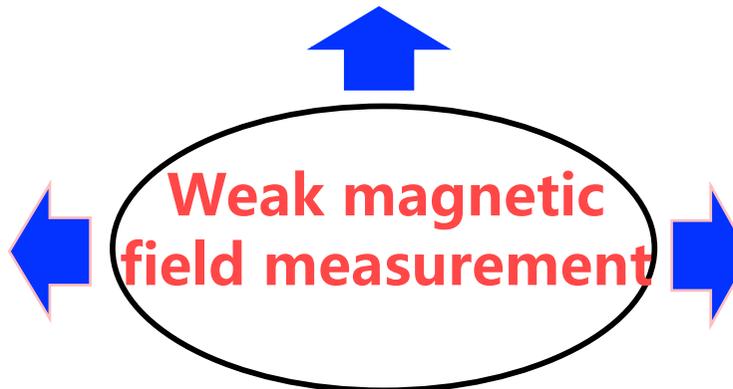
Navigation



Defence security

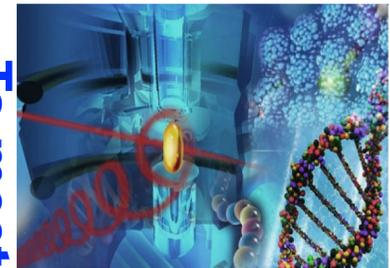
To meet the major national strategic needs

Facing the world's frontier of science and technology

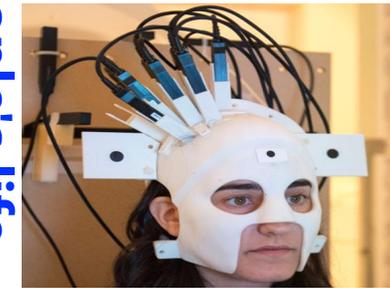


Quantum precision measurement

To meet people's life and health needs

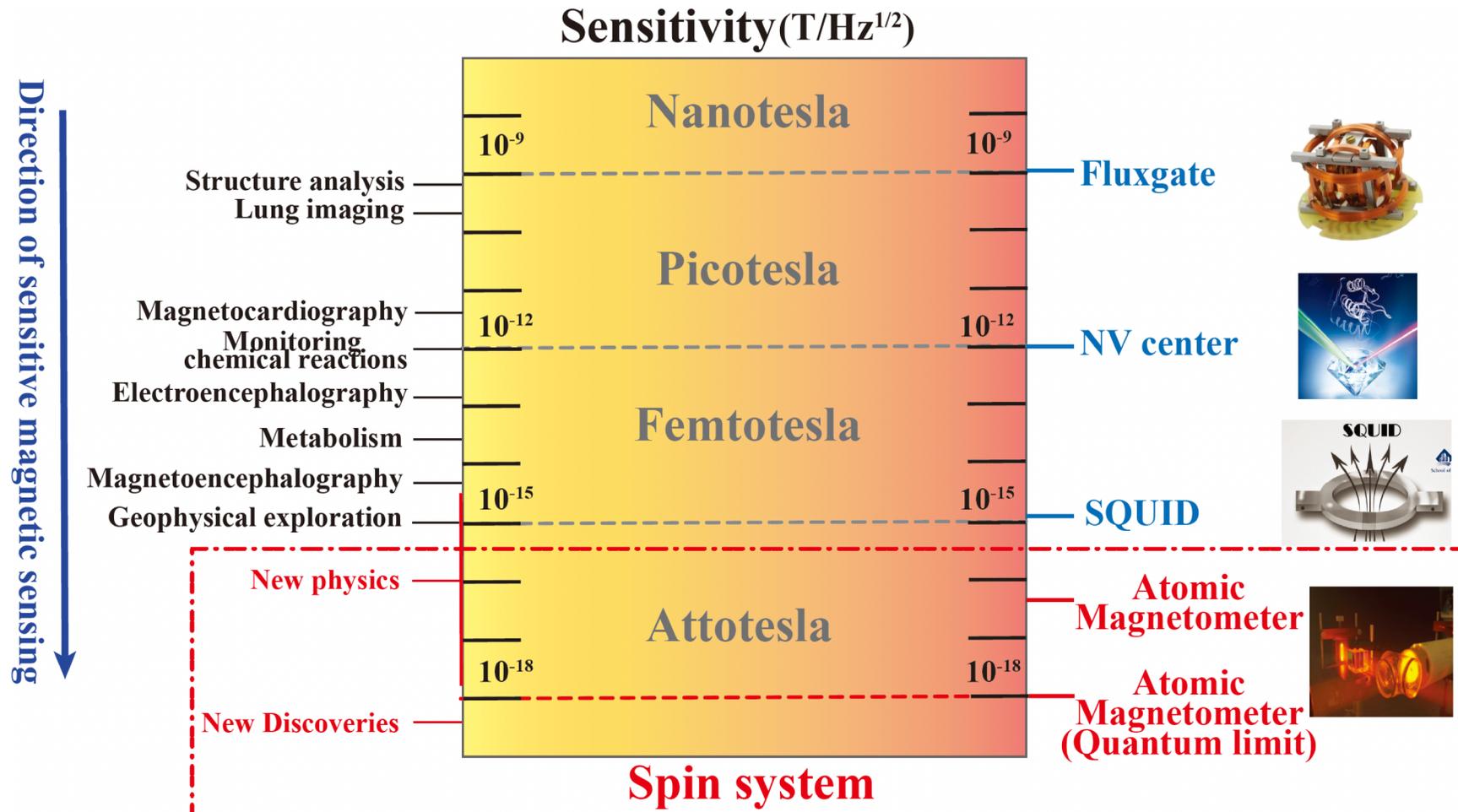


MRI



MEG/MCG

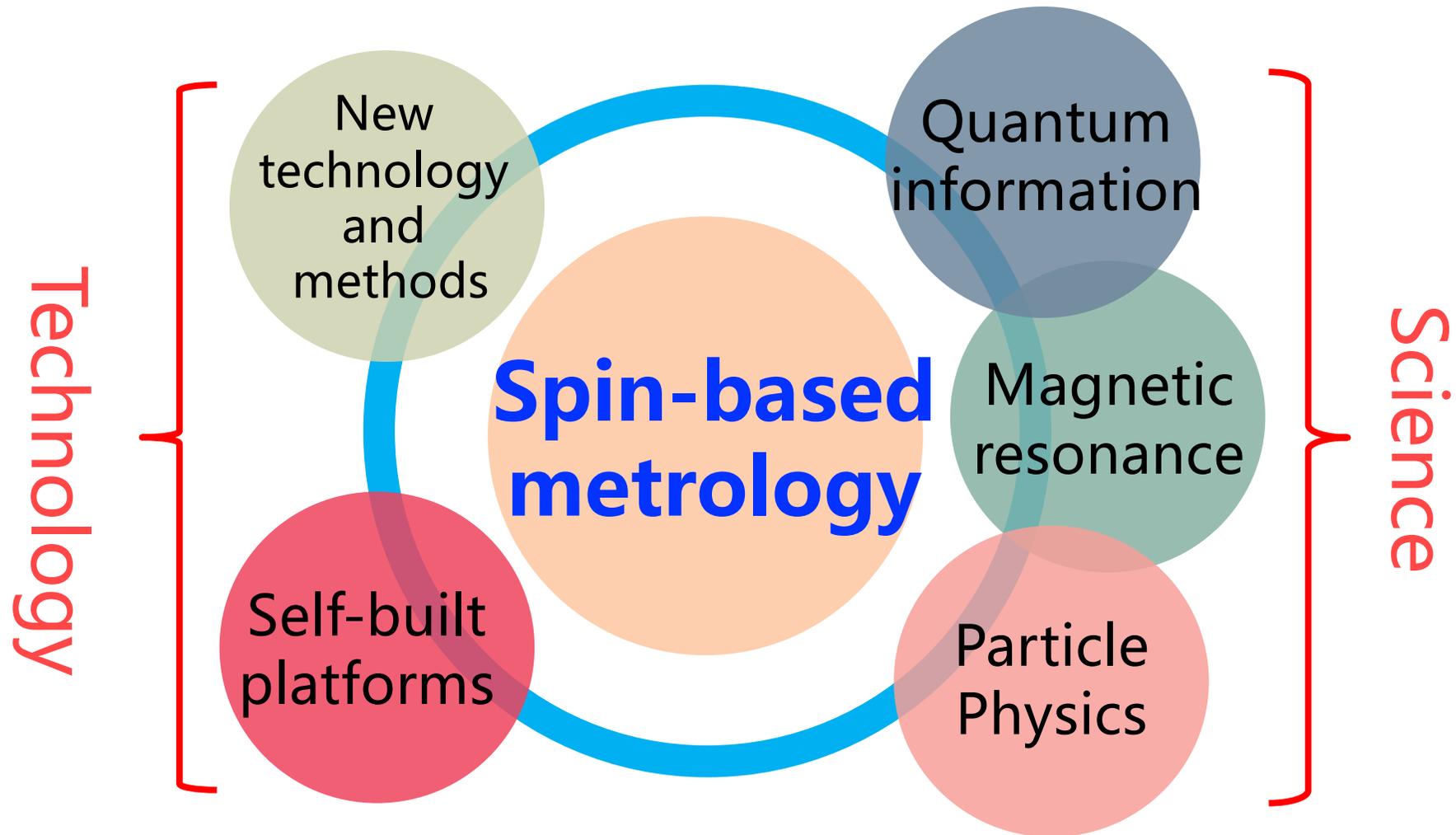
# Magnetic phenomena and their strength



**Quantum sensors make new discovery possible !**

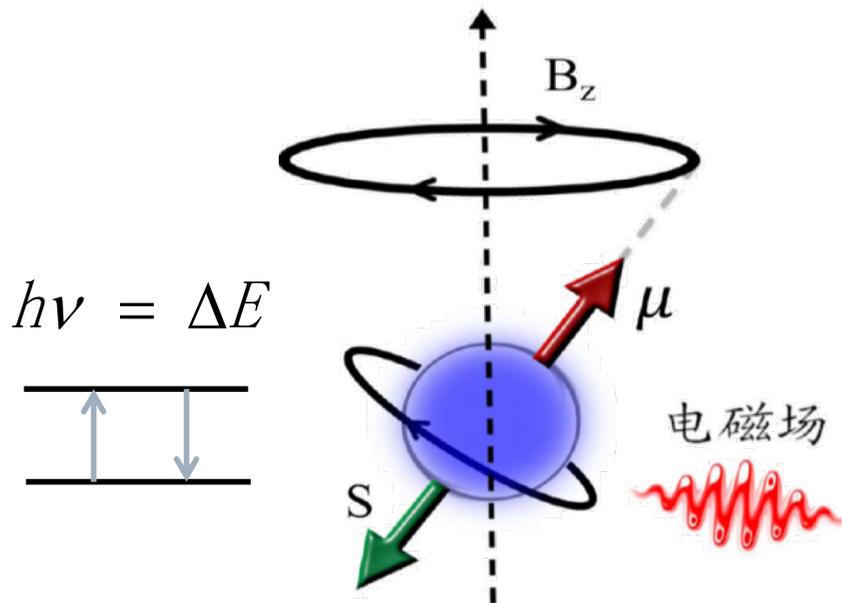
# Our research

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# Spin-based quantum precision measurement

**Principle :** The nucleus or electron spins placed in an external magnetic field, can absorb and release the electromagnetic radiation of the corresponding frequency, then the magnetic resonance phenomenon occurs.



Lamor precession

$$\omega = \gamma B$$

# Outline

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- Dark Matter and Axion: What is it and Why now?
  - Quantum sensors: How to see it?
    - Noble-gas spin systems: intrinsic systems vs. Floquet systems
    - Noble-gas spin amplification effect
    - Noble-gas spin Masing effect
  - Dark matter searches with spin-based-amplifier magnetometers
    - Axion dark matter
    - Exotic interactions mediated by new particles (the fifth force)
-

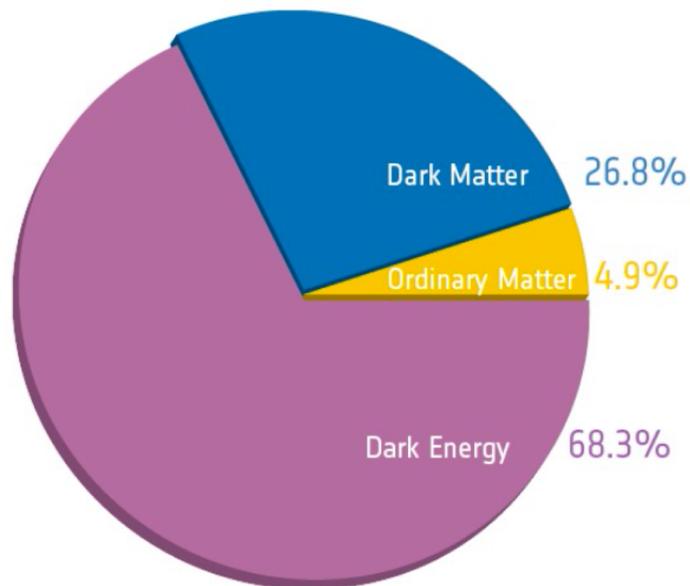
# Outline

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-

# Dark matter

## What is the nature of dark matter?

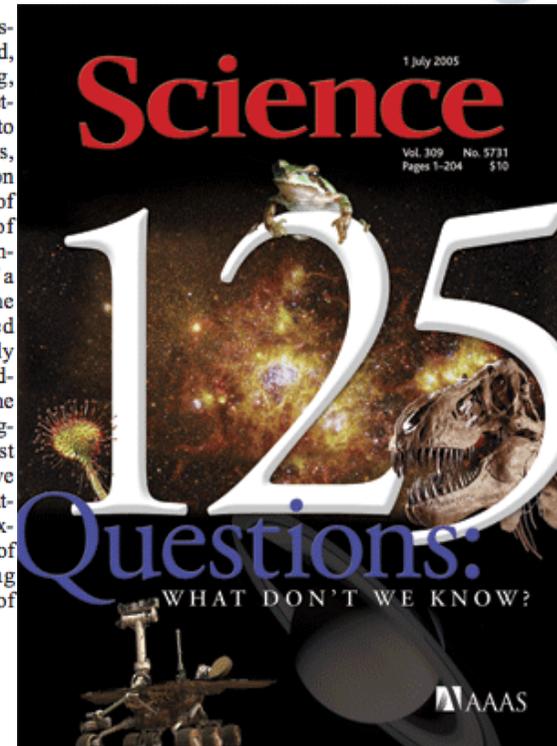


- What kind of microscopic particles are made of?
- How do particles interact?

WHAT DON'T WE KNOW?

## What Is the Universe Made Of

Every once in a while, cosmologists are dragged, kicking and screaming, into a universe much more unsettling than they had any reason to expect. In the 1500s and 1600s, Copernicus, Kepler, and Newton showed that Earth is just one of many planets orbiting one of many stars, destroying the comfortable Medieval notion of a closed and tiny cosmos. In the 1920s, Edwin Hubble showed that our universe is constantly expanding and evolving, a finding that eventually shattered the idea that the universe is unchanging and eternal. And in the past few decades, cosmologists have discovered that the ordinary matter that makes up stars and galaxies and people is less than 5% of everything there is. Grappling with this new understanding of



# Axions and axion-like particles



Steven Weinberg

VOLUME 40, NUMBER 4

PHYSICAL REVIEW LETTERS

23 JANUARY 1978

## A New Light Boson?

Steven Weinberg

*Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138*

(Received 6 December 1977)

It is pointed out that a global U(1) symmetry, that has been introduced in order to preserve the parity and time-reversal invariance of strong interactions despite the effects of instantons, would lead to a neutral pseudoscalar boson, the "axion," with mass roughly of order 100 keV to 1 MeV. Experimental implications are discussed.

PHYSICAL REVIEW D VOLUME 30, NUMBER 1

1 JULY 1984

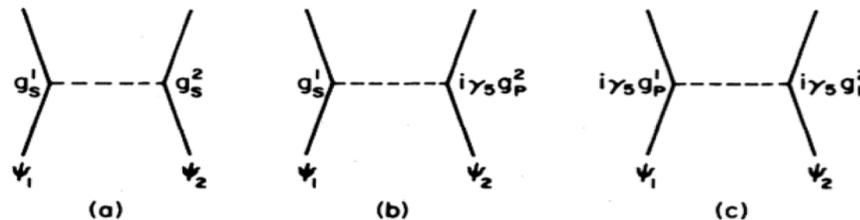
## New macroscopic forces?

J. E. Moody\* and Frank Wilczek

*Institute for Theoretical Physics, University of California, Santa Barbara, California 93106*

(Received 17 January 1984)

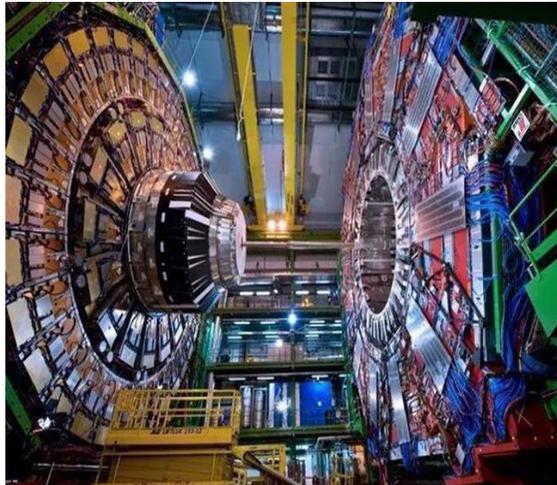
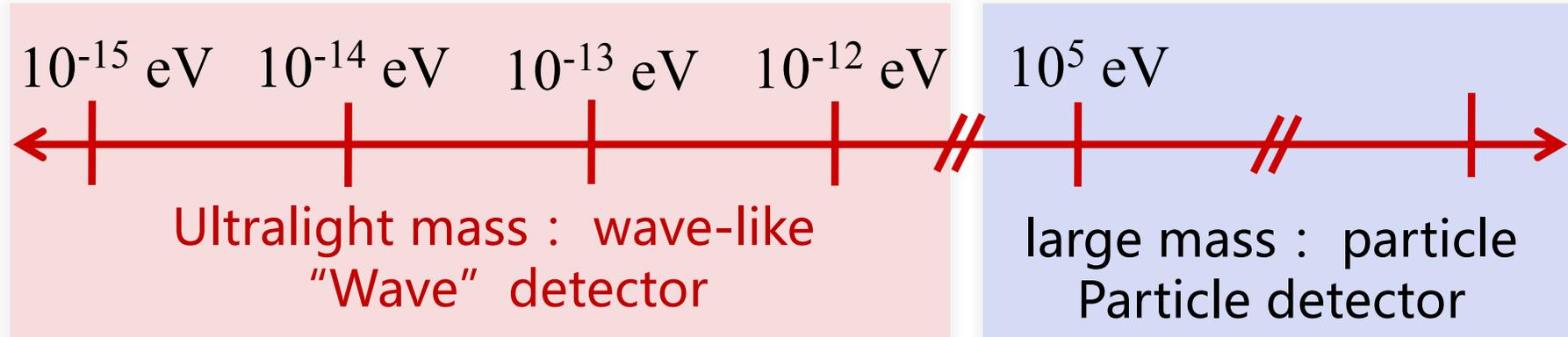
The forces mediated by spin-0 bosons are described, along with the existing experimental limits. The mass and couplings of the invisible axion are derived, followed by suggestions for experiments to detect axions via the macroscopic forces they mediate. In particular, novel tests of the T-violating axion monopole-dipole forces are proposed.



Frank Wilczek



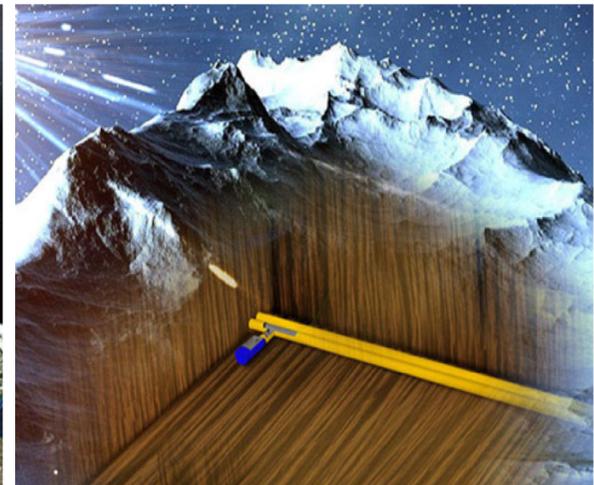
# How to see it?



Collider exp.  
(LHC, BESIII)

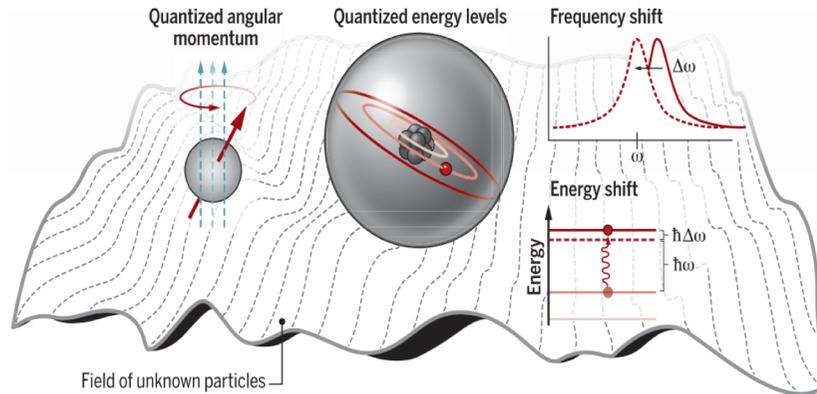


Space exp.  
(DAMPE, AMS)



Underground exp.  
(PandaX, CDEX)

# Quantum sensors: sensitive low-energy tools



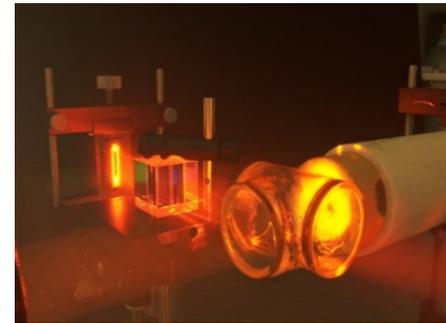
**Dark matter** is very weakly coupled to **atomic molecular systems**, causing small energy shifts

Science 357, 990 (2017)

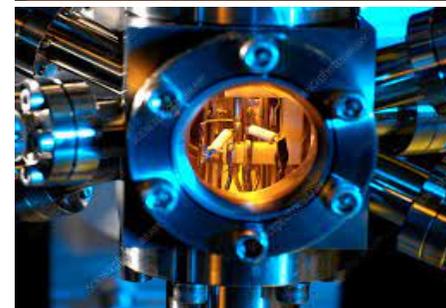
Rev. Mod. Phys. 90,025008 (2018)

- **Sensitive** : quantum-noise limits
  - **Tabletop** : small size and cheap
  - **Networks** : low false-alarm rate
- Beyond the astrophysical limits !**

## Quantum sensors



**Magnetic field:  
Magnetometers  
(or spin sensor)**



**Time standard:  
Optical clocks**



**Gravity:  
Interferometry**

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-

# How to see it?

## The nonrelativistic Hamiltonian

$$\mathcal{H} = \sqrt{\frac{\epsilon_0}{\mu_0}} g_{a\gamma\gamma} \int a \mathbf{E} \cdot \mathbf{B} dV + g_{\text{aff}} \hbar c \nabla a \cdot \hat{\mathbf{S}} + \sqrt{\epsilon_0 (\hbar c)^3} g_{\text{EDM}} a \hat{\mathbf{S}} \cdot \mathbf{E}$$

**Couplings** axion-photon

axion-fermion

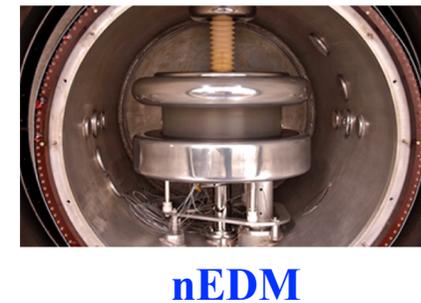
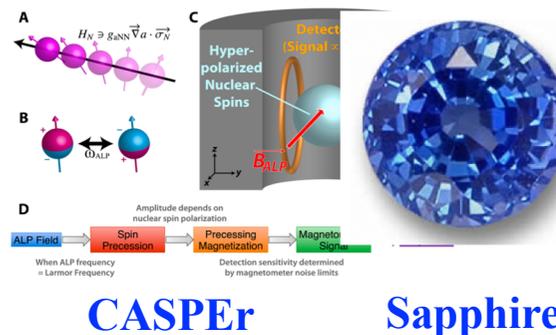
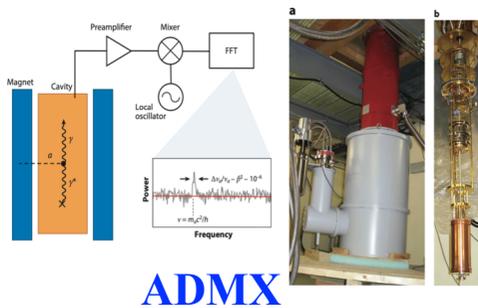
axion-gluon

**Effects**

axions is converted to real photons

axions generate pseudomagnetic field on spins

axions generate oscillating EDM



# Quantum spin metrology

---

$$\omega = \gamma B$$

- Alkali-metal atoms      K , Rb , Cs
- Hyperpolarized noble gas       $^3\text{He}$  ,  $^{129}\text{Xe}$

Precision limit:  $\delta\omega = \frac{1}{P\gamma\sqrt{NT_2t}}$

- Polarization:  $P$
- Number of spins:  $N$
- Spin relaxation:  $T_2$

Why we choose alkali metal atoms/noble gas ?

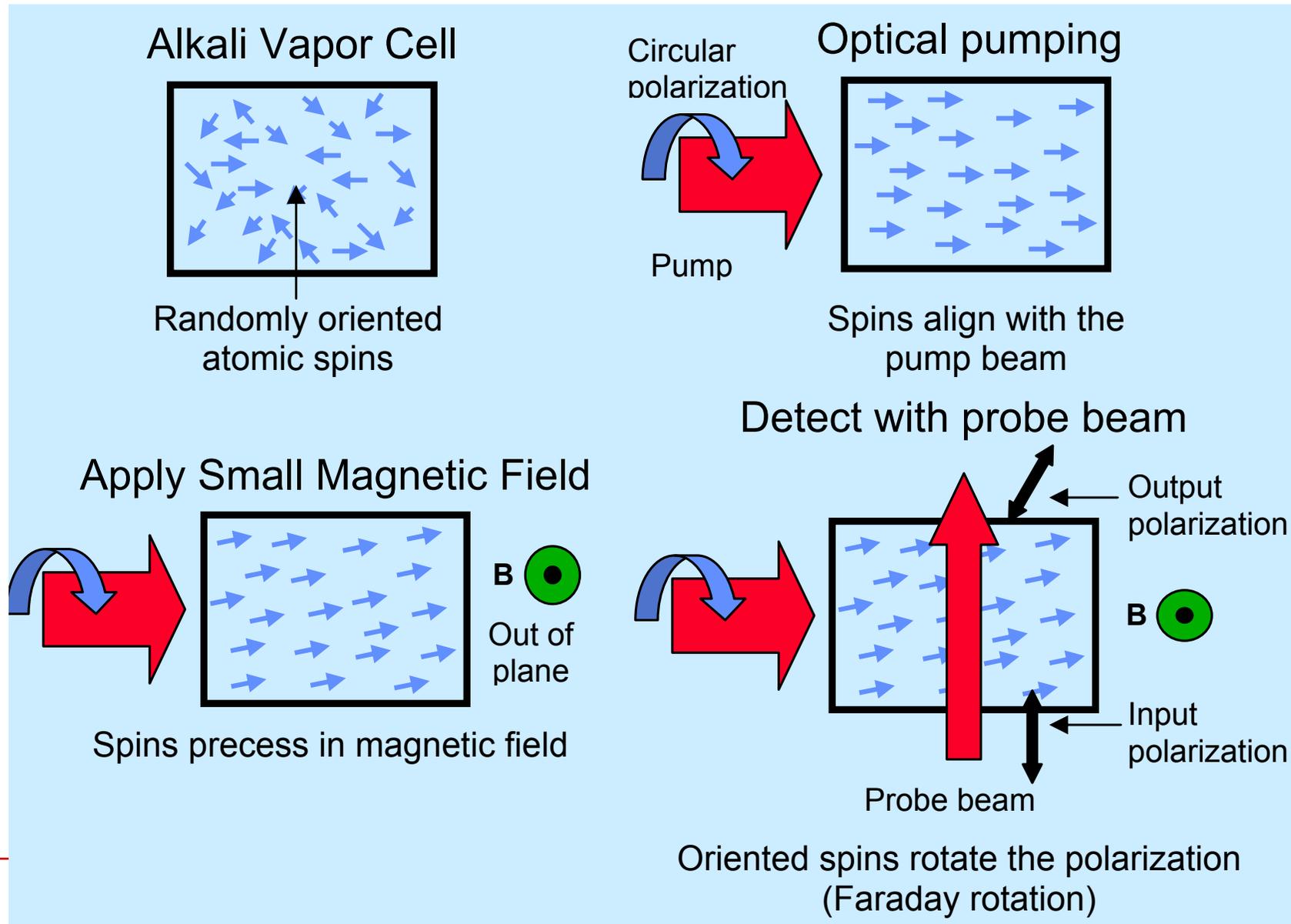
- Alkali metal atoms: Well-developed pump and probe techniques
  - Noble gas: Long coherence time (  $>10$  s )
  - Spin exchange interactions via collisions (Allow information transfer between them)
-

## Coherence time ( $T_2^*$ )

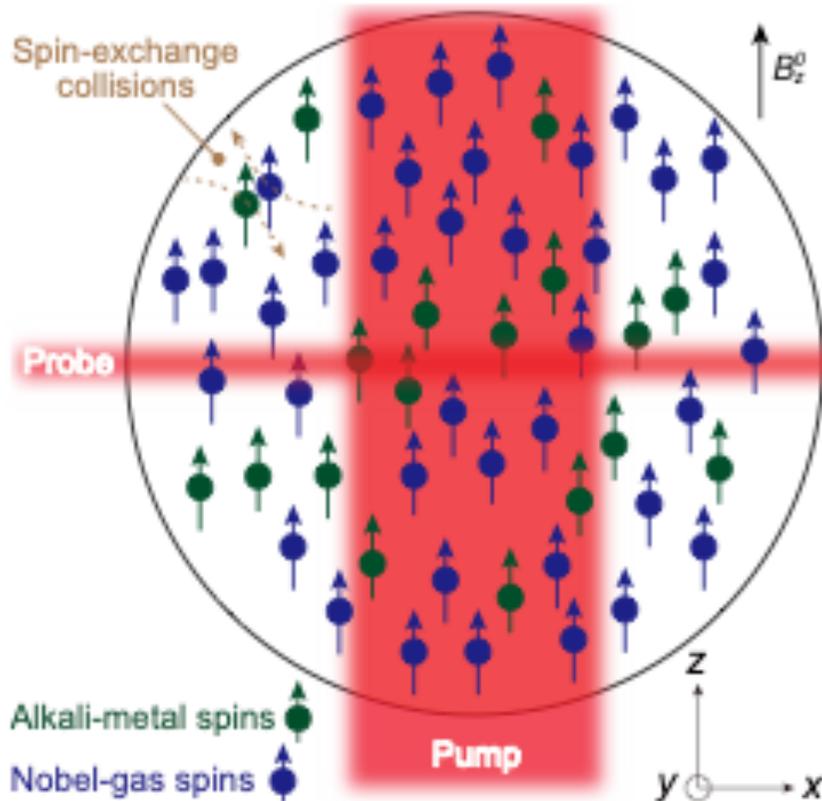
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- $^3\text{He}$ :  $I=1/2$ ,  $T_2 > 1$  h, abundance: 0.01%
- $^{129}\text{Xe}$ :  $I=1/2$ ,  $T_2 \sim 20$  s, abundance: 26.4%
- $^{21}\text{Ne}$ :  $I=3/2$ , abundance: 21%
- $^{83}\text{Kr}$ :  $I=9/2$ , abundance: 83%
- $^{131}\text{Xe}$ :  $I=3/2$ ,  $T_2 \sim 5$  s, abundance: 21.2%

# Atomic magnetometer



# Spin dynamics



$$\beta M_0^e \mathbf{P}^e \ll \beta M_0^n \mathbf{P}^n$$

Coupled Bloch equations

$$\frac{\partial \mathbf{P}^e}{\partial t} = \frac{\gamma_e}{Q} (B_z^0 \hat{z} + \mathbf{B}_a + \beta M_0^n \mathbf{P}^n) \times \mathbf{P}^e + \frac{P_0^e \hat{z} - \mathbf{P}^e}{T_e Q},$$

$$\frac{\partial \mathbf{P}^n}{\partial t} = \gamma_n (B_z^0 \hat{z} + \mathbf{B}_a) \times \mathbf{P}^n + \frac{P_0^n \hat{z} - \mathbf{P}^n}{\{T_{2n}, T_{2n}, T_{1n}\}}.$$

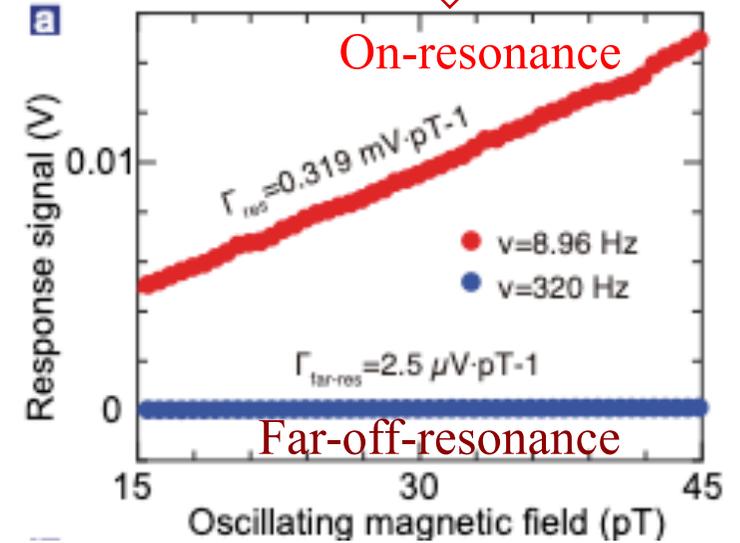
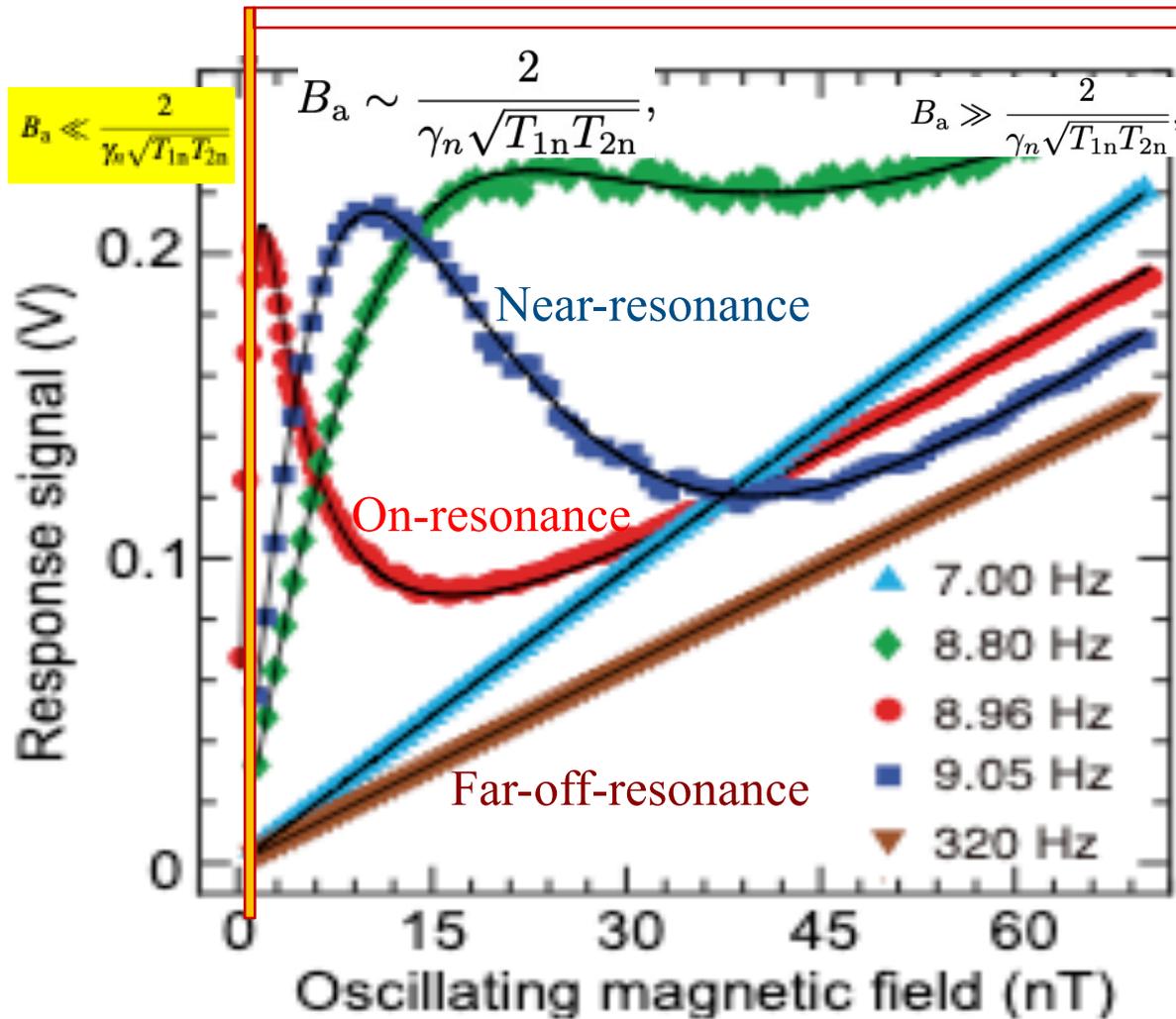
Polarized noble-gas nuclear spins generate an effective field experienced by alkali-metal atoms due to the Fermi-contact interactions

$$\mathbf{B}_{\text{eff}}^n = \beta M_0^n \mathbf{P}^n$$

Fermi-contact factor  $\kappa_0 \sim 540$  (Rb+Xe)

$$\beta = 8\pi\kappa_0/3$$

# Spin dynamics



Slope: Response to the magnetic field

At least two orders of magnitude improvement

Linear Nonlinear

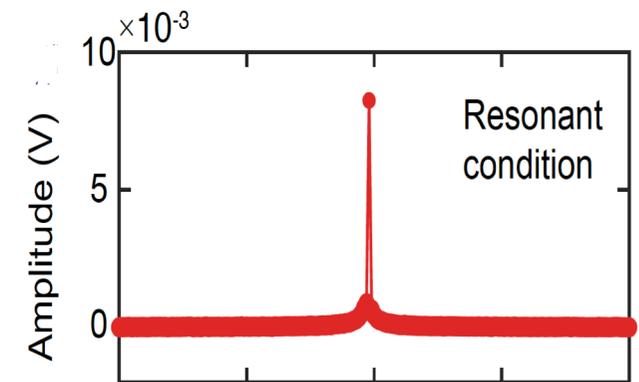
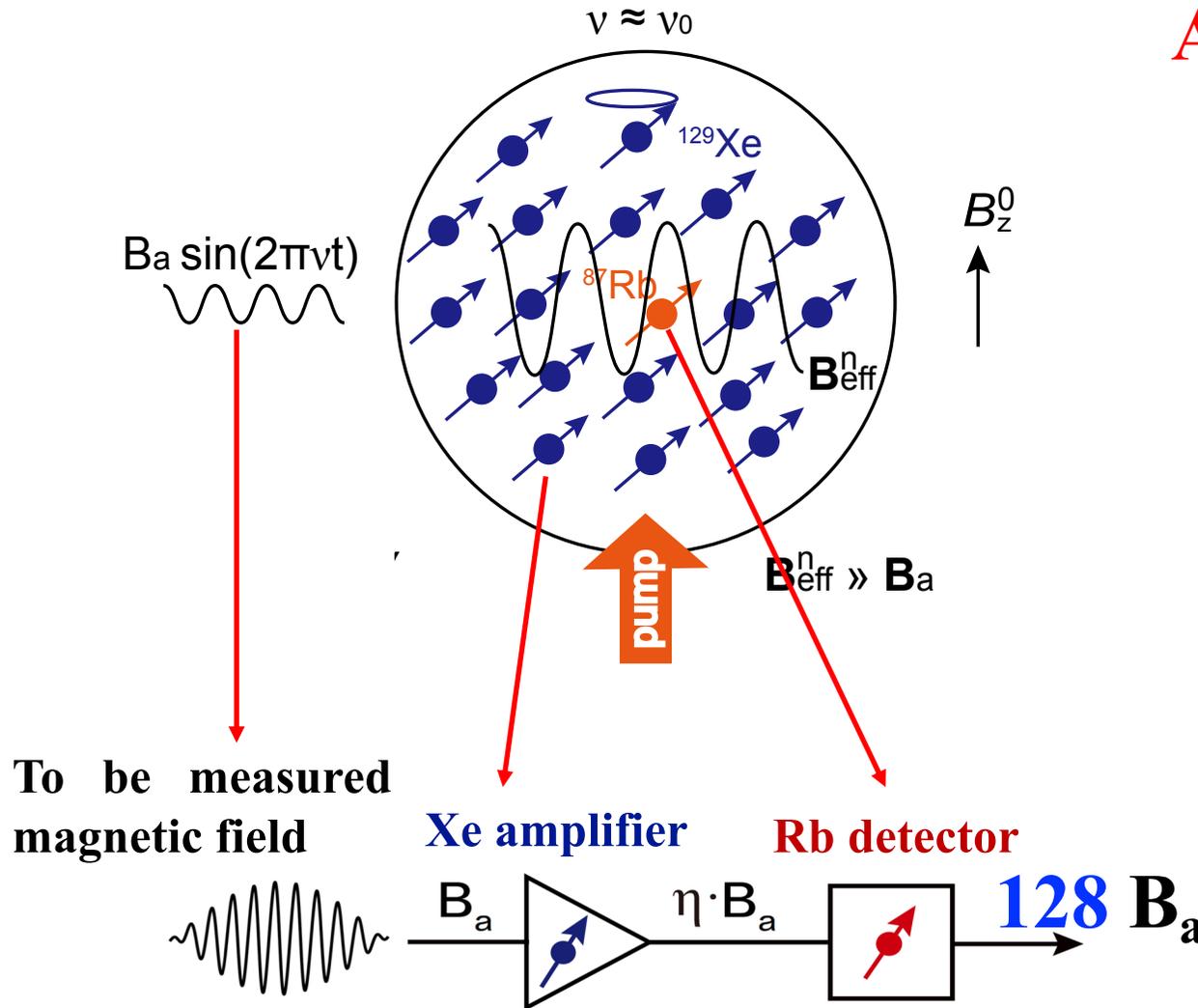
Saturated

# Amplification-assisted magnetometer

**Amplification factor:**

$$\eta = \frac{1}{2} \beta M_0^n P_0^n \gamma_n T_{2n} \gg 1$$

$$\delta B = \frac{1}{\eta P \gamma \sqrt{N T_2 t}}$$

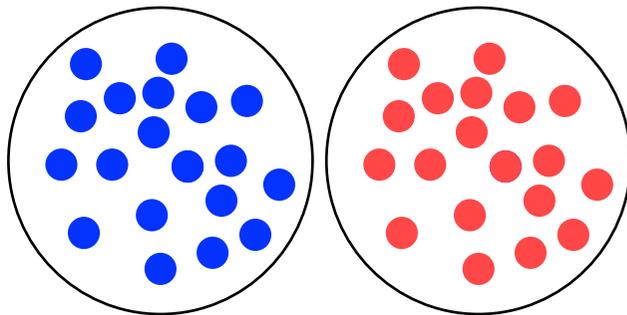


# Spin amplifier: Amplification of magnetic field

## Magnetic interaction

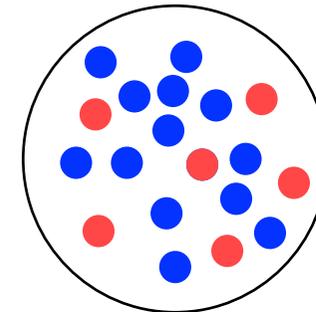
$$\mathbf{B}(\mathbf{r}) = \underbrace{\frac{\mu_0}{4\pi} \frac{3\hat{r}(\hat{r} \cdot \mathbf{m}) - \mathbf{m}}{r^3}}_{\text{"remote"}} + \underbrace{\frac{2\mu_0\mathbf{m}}{3} \delta(\mathbf{r})}_{\text{"contact"}}$$

External readout (dipole field)



Nuclear spins    Alkali-metal spins

**Fermi contact readout**

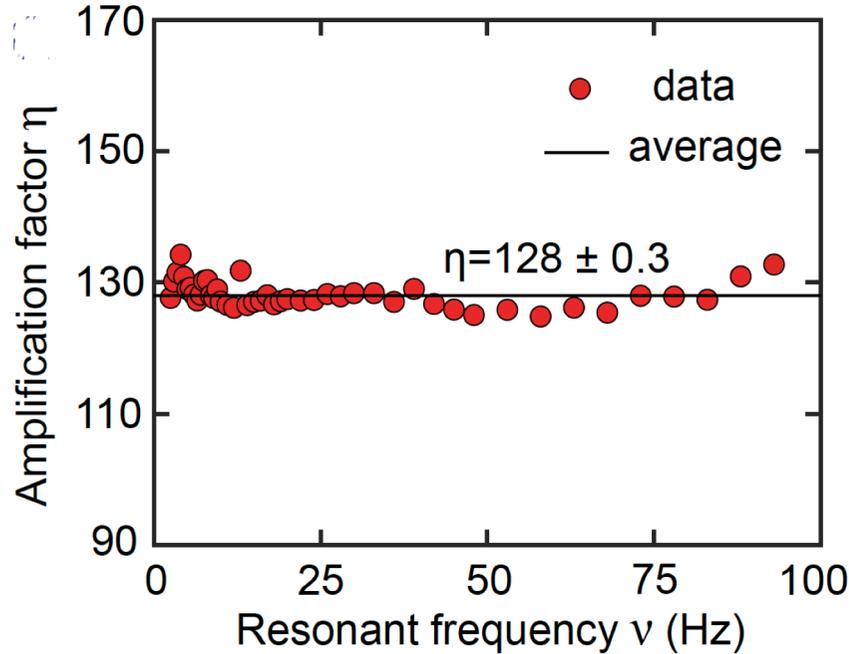


Noble-gas + alkali-metal  
overlapping spins

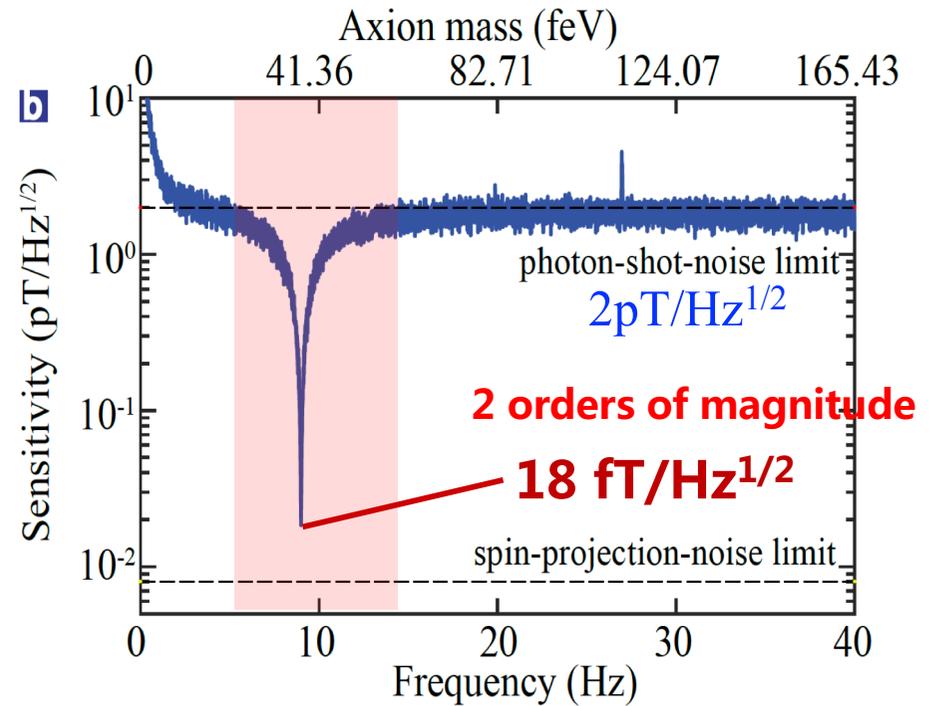
- Large Fermi-contact enhancement ( $\sim 600$  for Rb-Xe)
- Without spin polarization loss
- Continuous measurement

Jiang et al., [Nature Physics](#) **17**, pages 1402–1407 (2021)

# Ultrasensitive magnetic field sensing



Magnetic field is amplified by a factor of more than 100!

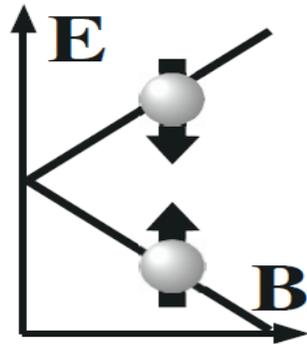


Femtotesla-level sensitivity!

$$1\text{fT} = 10^{-15}\text{T}$$

Jiang et al., *Nature Physics* **17**, pages 1402–1407 (2021)

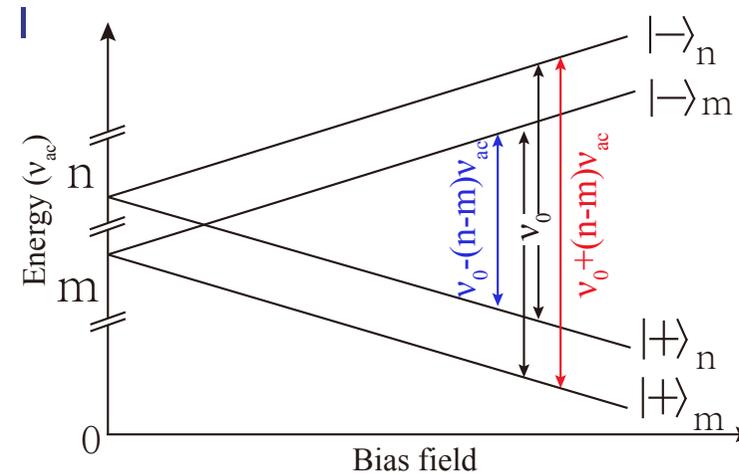
# Periodically driven (Floquet) spin systems



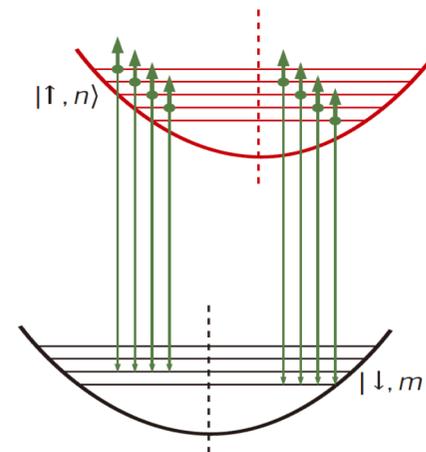
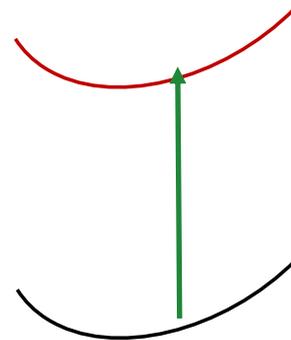
Two-level spin system

Extended to a series of Floquet states

## Floquet levels/Floquet states

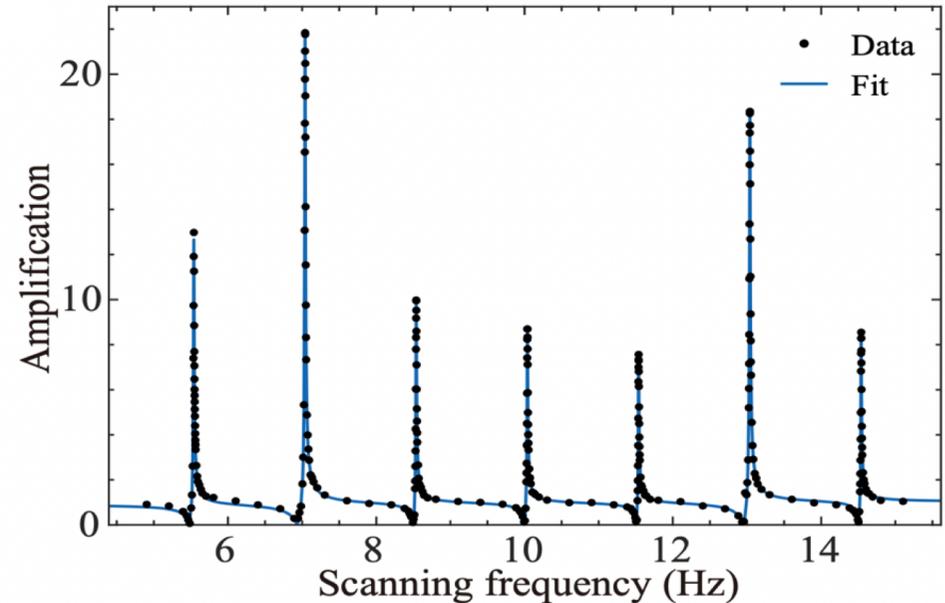
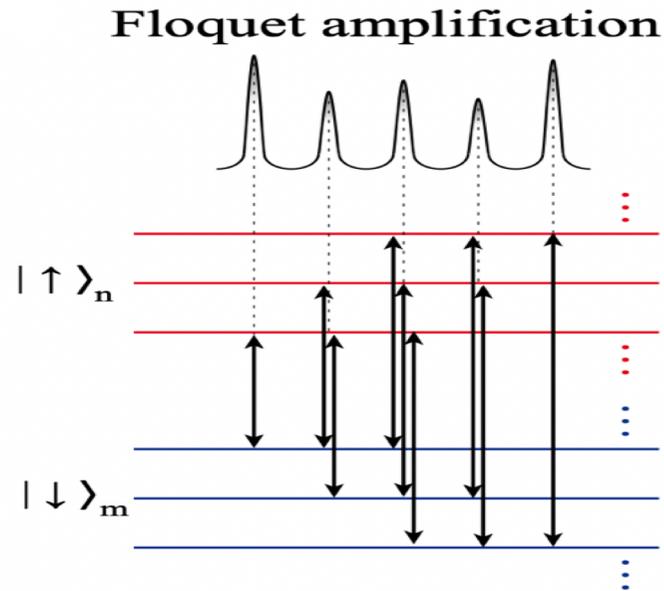


One transition



Many Floquet transitions

# Floquet spin amplification

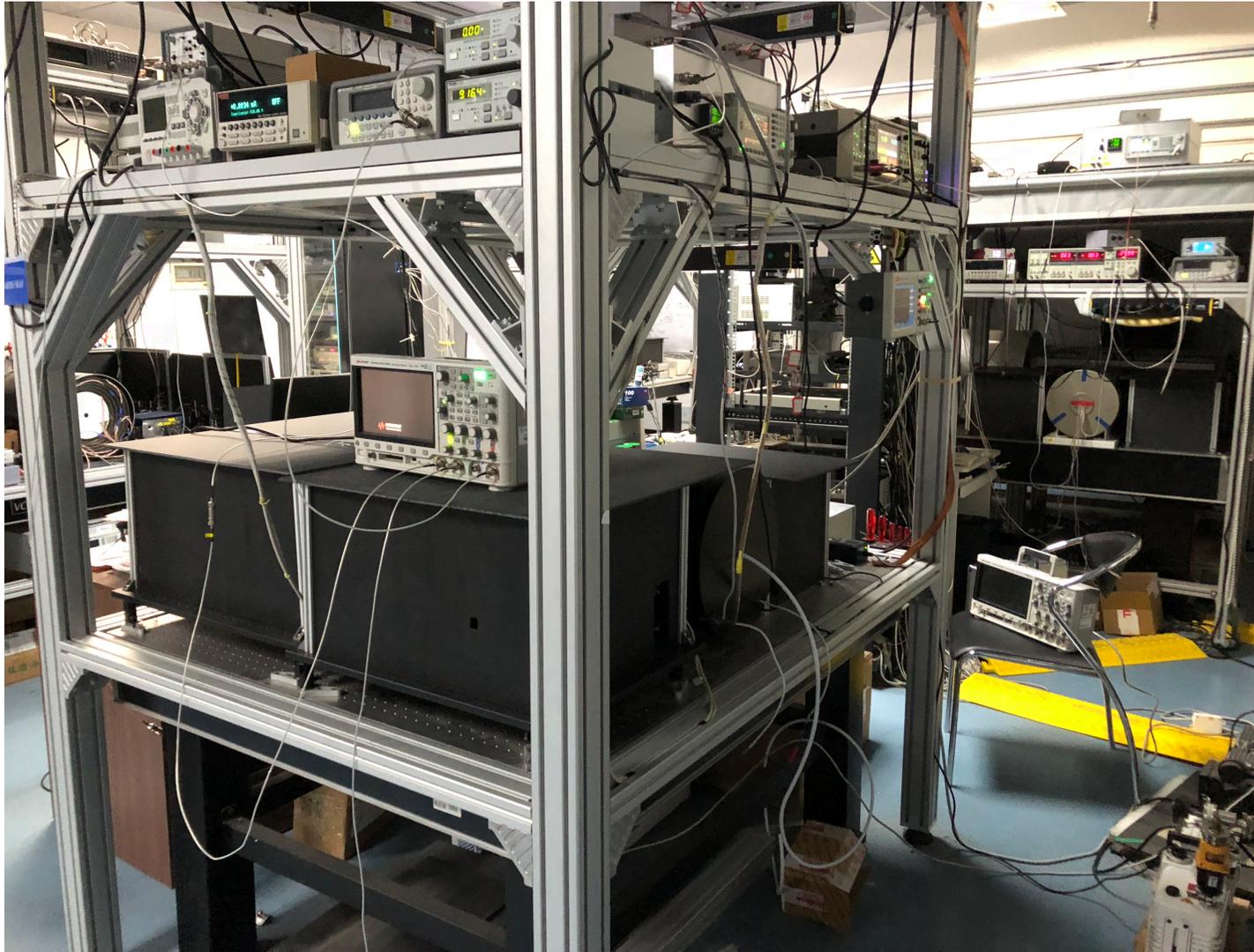


**Amplification factor:**

$$\eta_{k,0}(u) = \frac{4\pi}{3} \kappa_0 M^n P_0^n \gamma_n T_{2n} J_k^2(u)$$

10 times improvement in  
the detection range,  
 $\text{fT}/\text{Hz}^{1/2}$  detect sensitivity

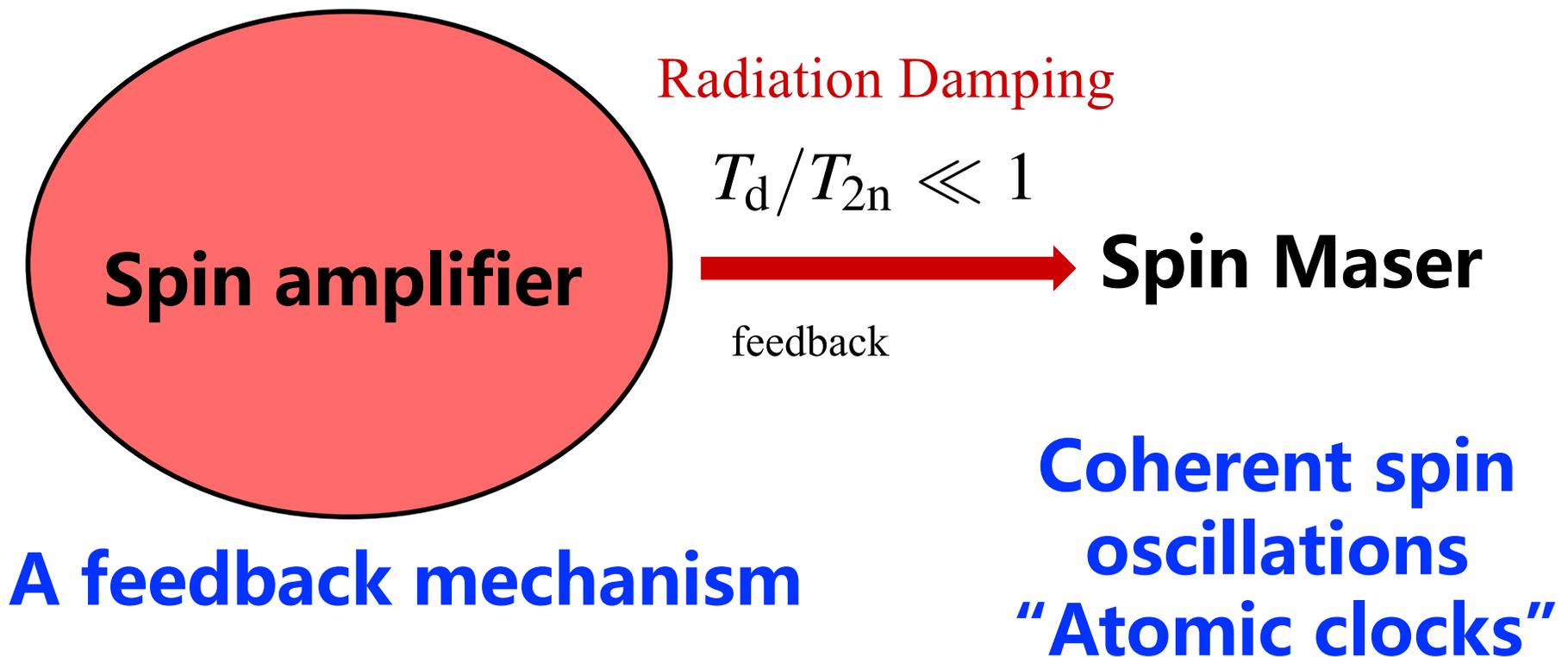
# 飞特斯拉 ( fT ) 级别的自旋传感装置



# Noble-gas masing effect

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Spin amplification combines with a “cavity-like” feedback



# Spin amplification combines with a feedback

## Nobel Prize: Maser and atomic clock



C. Townes



G. Basov



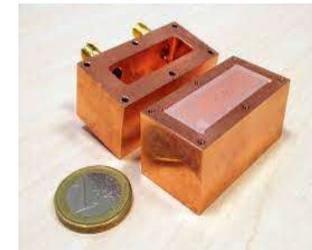
M. Prokhorov



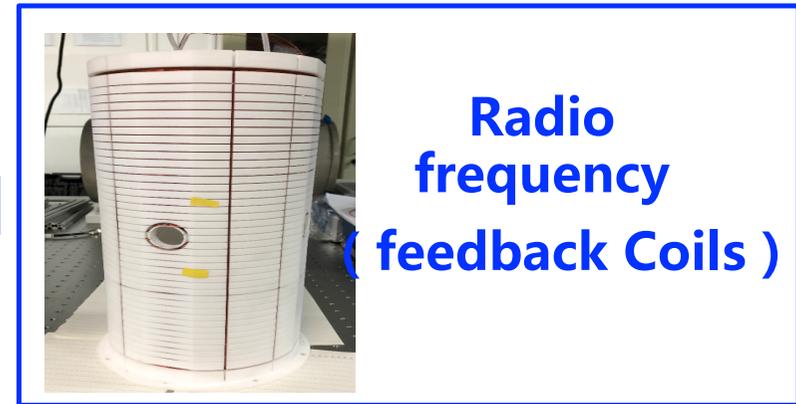
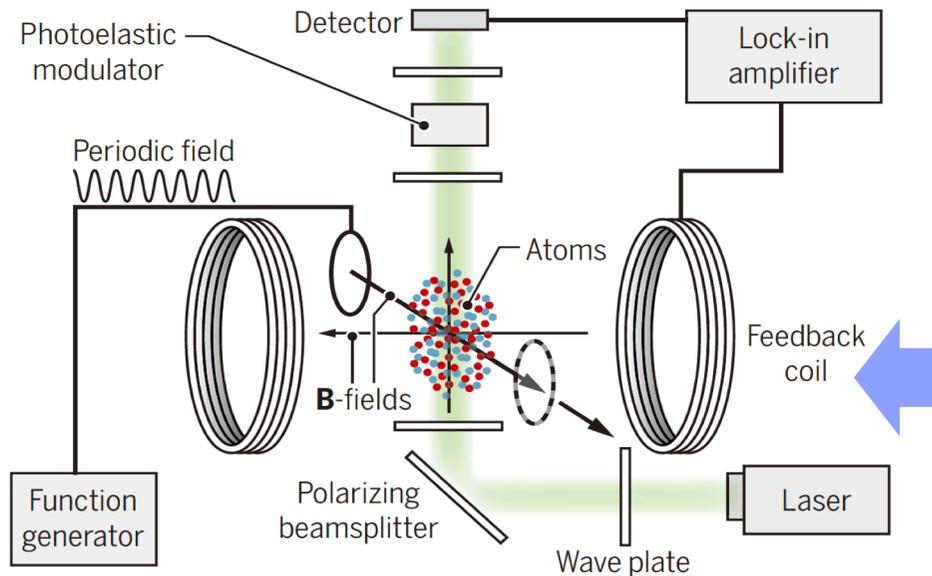
N. F. Ramsey



**Optical frequency**  
( **Optical cavity** )



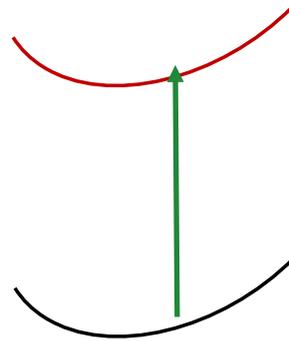
**Microwave frequency**  
( **Cavity** )



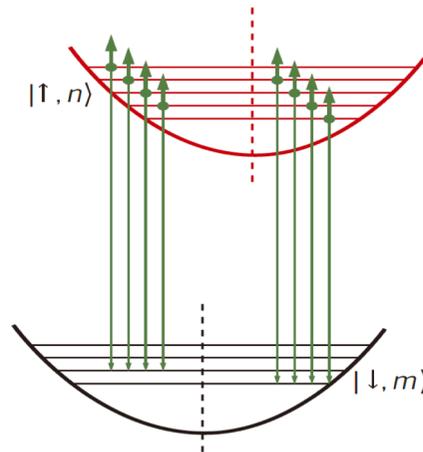
**Radio frequency**  
( **feedback Coils** )

# Periodically driven (Floquet) spin maser

Two-level spin system

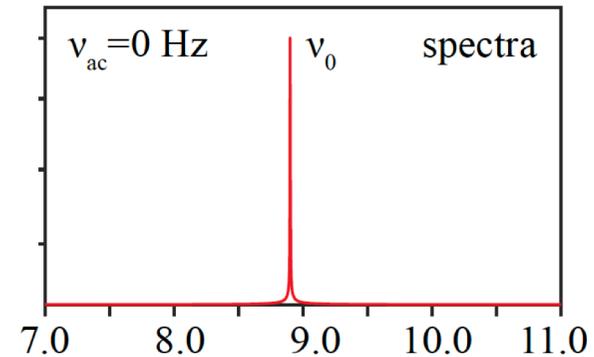


Pseudo-magnetic field from particles

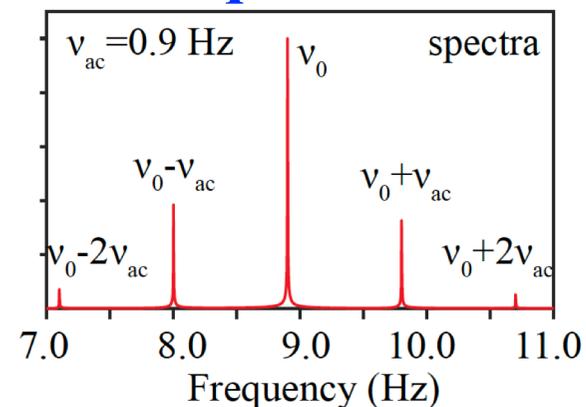


Extended to a series of Floquet states

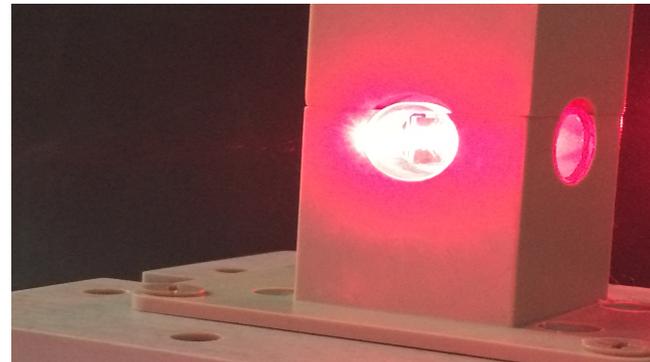
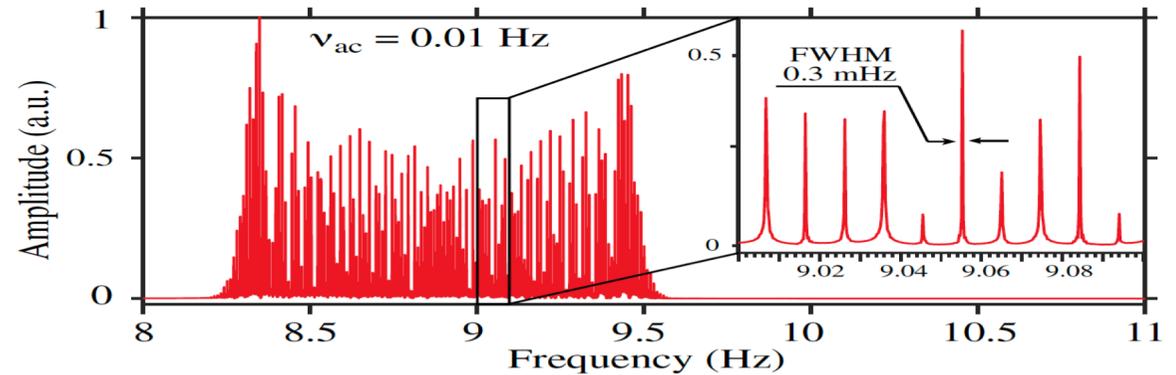
**traditional maser**



**Floquet maser**



# Periodically driven (Floquet) spin maser



频率梳~130  
条跃迁线

ScienceAdvances

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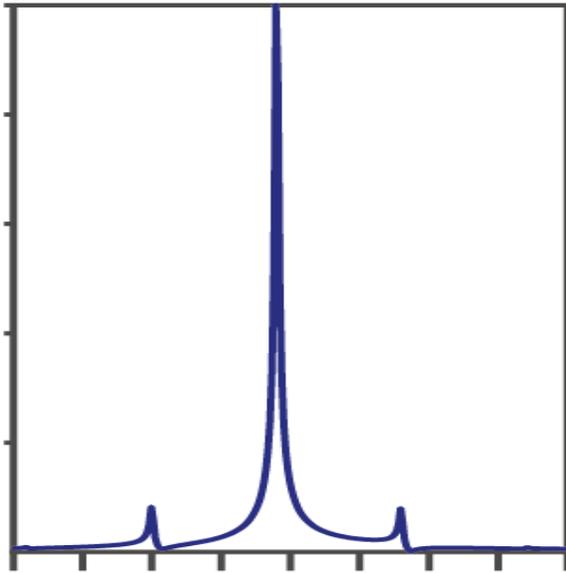
Floquet maser



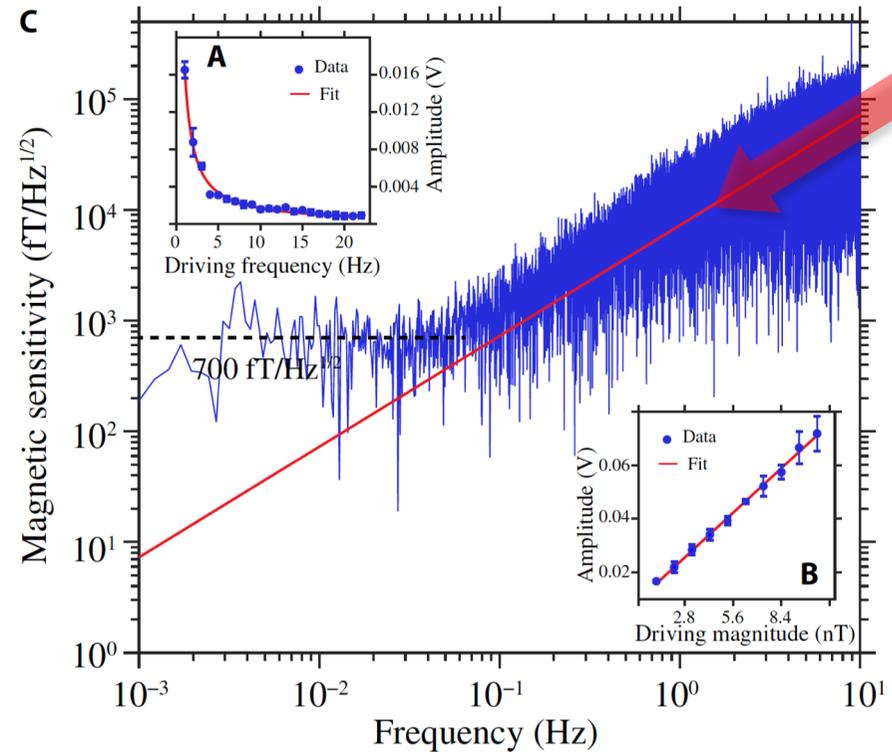
[Min Jiang](#)<sup>1,2,3</sup>, [Haowen Su](#)<sup>1,2,3</sup>, [Ze Wu](#)<sup>1,2,3</sup>, [Xinhua Peng](#)<sup>1,2,3,\*</sup> and [Dmitry Budker](#)<sup>4,5,6</sup>

# Floquet-maser-based sensing

1-1000mHz: axion-mass  $10^{-17}\text{eV}-10^{-14}\text{eV}$



$$S_1 \propto J_1\left(\frac{\gamma B_{ac}}{U_{ac}}\right) \approx \frac{\gamma}{2U_{ac}} B_{ac}$$



Best sensor for 1-100mHz !



## Floquet maser



Min Jiang<sup>1,2,3</sup>, Haowen Su<sup>1,2,3</sup>, Ze Wu<sup>1,2,3</sup>, Xinhua Peng<sup>1,2,3,\*</sup> and Dmitry Budker<sup>4,5,6</sup>

# Science

## A masing ladder

A maser that amplifies emission of periodically modulated quantum states has uses in metrology

**PHYS**  **ORG**

### Extending maser techniques to Floquet systems

techniques to Floquet systems. In their paper published in the journal *Science Advances*, the group describes their approach to creating a new type of maser by amplifying radio frequencies in Floquet systems. Ren-Bao Liu, with the Chinese

**physicsworld** RESEARCH UPDATE

New Floquet maser is very good at detecting low frequency magnetic fields

《Science》 Perspectives reports:

“... demonstrate a new type of maser...

Conceivable applications of this work include precision clocks and detection of ultralight dark matter particles such as axions”

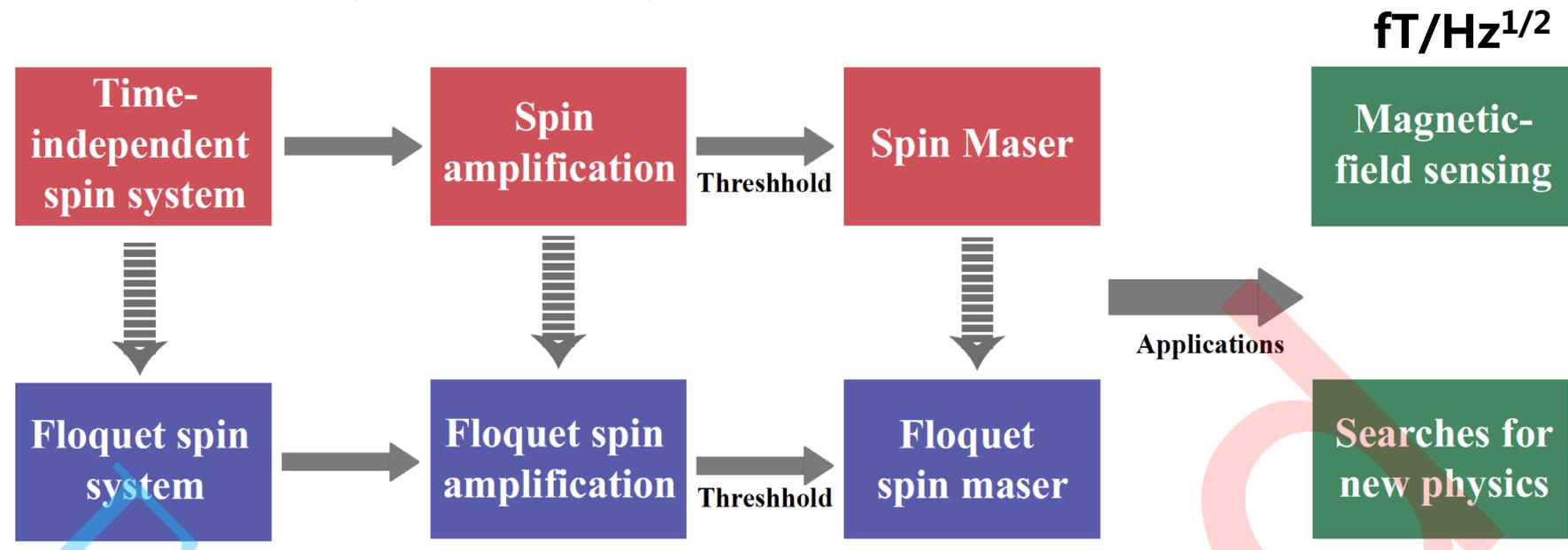
# Spin amplification

SCIENCE CHINA

Information Sciences

## Review of noble-gas spin amplification via the spin-exchange collisions

Haowen Su<sup>1,2</sup>, Min Jiang<sup>1,2</sup> and Xinhua Peng<sup>1,2</sup>



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-

# "Sapphire" project

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## Spin Amplifier for Particle PHysics Research



*Can I Afford a Blue Sapphire?*



# Sapphire project ( “蓝宝石” 计划 )

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## Spin Amplifier for Particle PHysics Research



**直接探测：**  
暗物质候选粒子  
轴子、暗光子

**“蓝宝石”计划：**  
自旋放大、fT灵敏度、  
阵列探测、桌面式

**间接探测：**  
新奇相互作用  
新粒子作为传播子

# Direct detection for ALP

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Goal: to detect an oscillating nuclear-spin-dependent energy shift caused by the weak coupling between Spin and Axion

“pseudo-magnetic” field:

$$\vec{B}_{\text{ALP}} \propto g_{\text{aNN}} \cos(m_{\text{ALP}} t) \vec{v}$$

Resonance frequency probes the mass  $m_{\text{ALP}}$   
Amplitude  $B_{\text{ALP}}$  probes the coupling value of  $g_{\text{aNN}}$

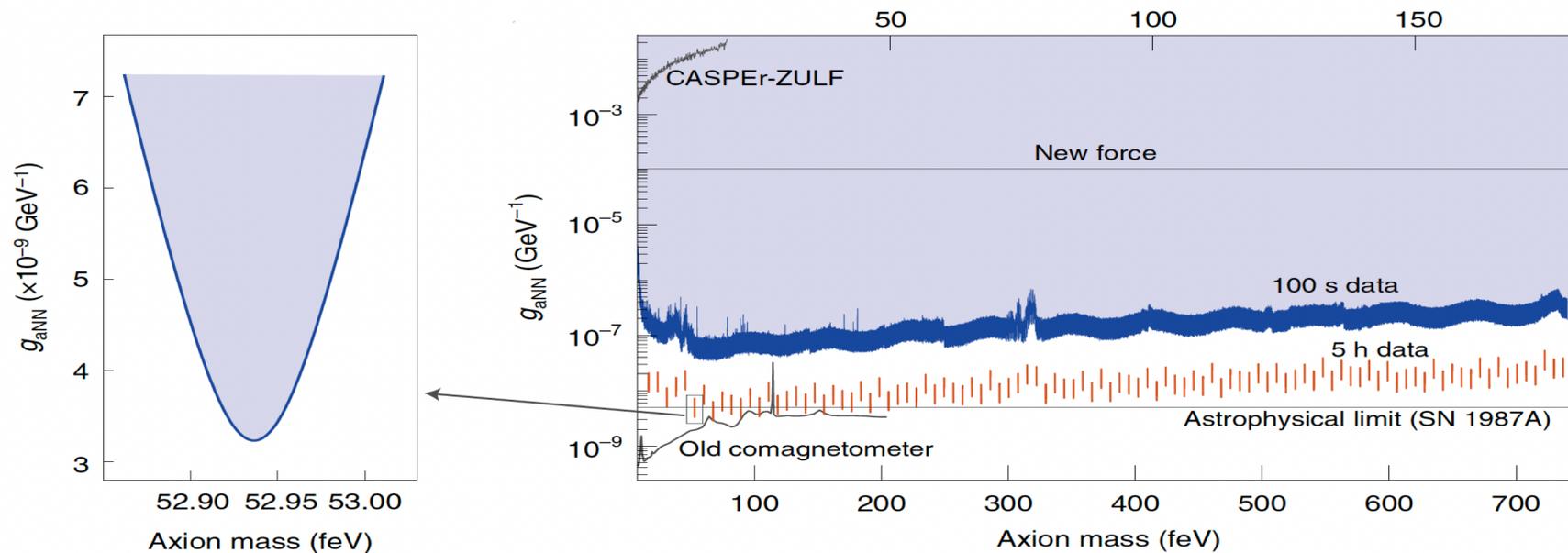
More sensitive detector  $\rightarrow$  Smaller coupling  $g_{\text{aNN}}$   
Lower-frequency detector  $\rightarrow$  Lighter mass  $m_{\text{ALP}}$

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# Beyond the astrophysical limits

## Search for axion-like dark matter with spin-based amplifiers

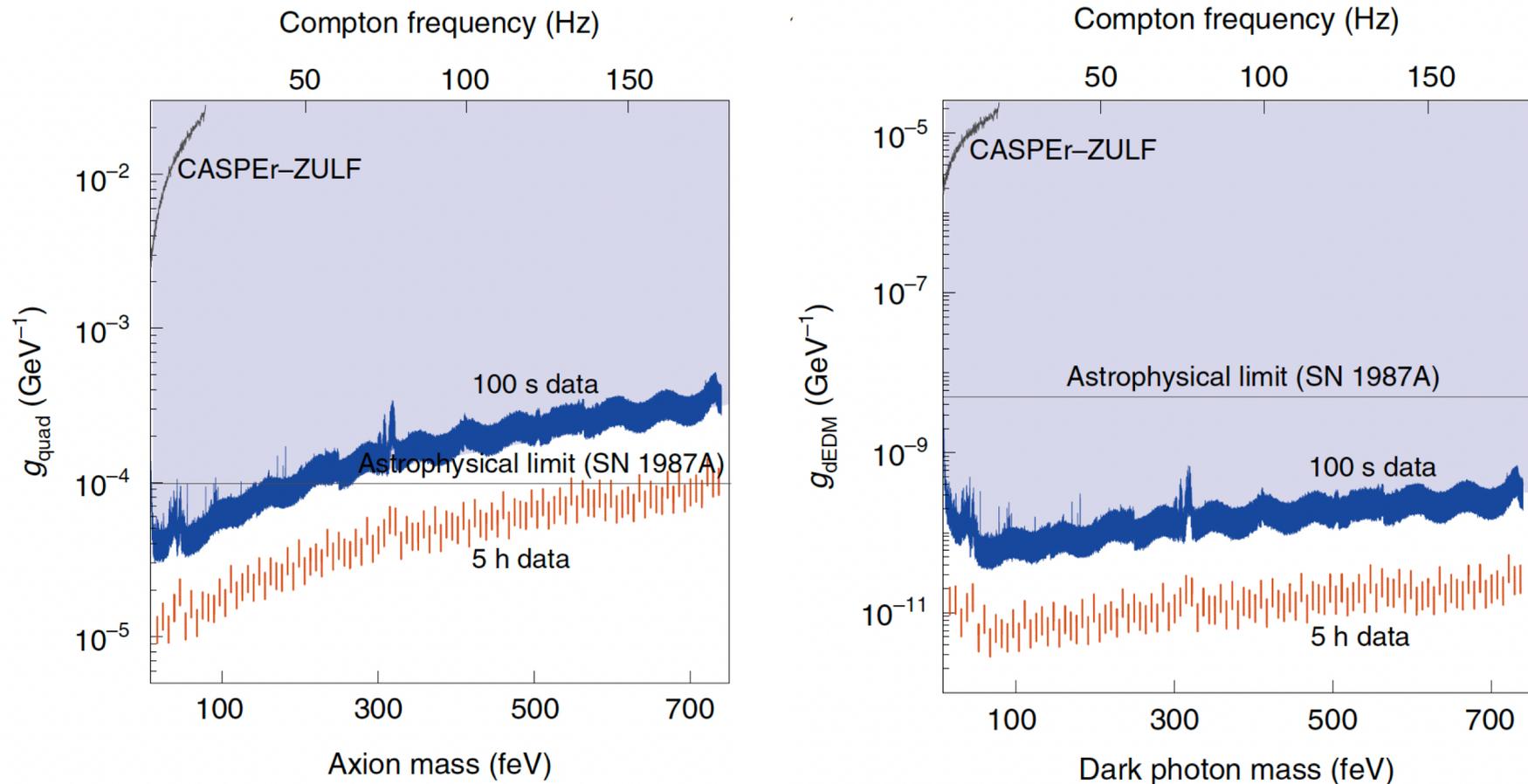
Min Jiang<sup>1,2,3,7</sup>, Haowen Su<sup>1,2,3,7</sup>, Antoine Garcon<sup>4,5</sup>, Xinhua Peng<sup>1,2,3</sup> and Dmitry Budker<sup>4,5,6</sup>



## New constraints on axion-neutron coupling

Jiang et al., *Nature Physics* 17, pages 1402–1407 (2021)

# Beyond the astrophysical limits



At least five orders of magnitude improvement,  
beyond the astrophysical limits

Jiang et al., [Nature Physics](#) **17**, pages 1402–1407 (2021)

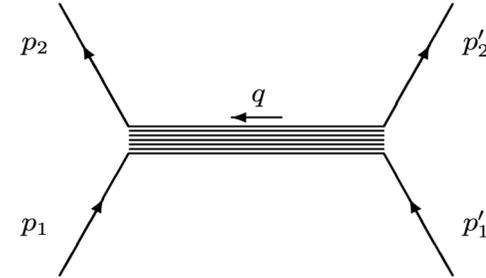
# Exotic interactions beyond the standard model



B. A. Dobrescu



I. Mocioiu



Extend axion to new mediator bosons and lead to 16 interactions

$$\begin{aligned} \mathcal{V}_1 &= \frac{1}{r} y(r) , \\ \mathcal{V}_2 &= \frac{1}{r} \vec{\sigma} \cdot \vec{\sigma}' y(r) , \\ \mathcal{V}_3 &= \frac{1}{m^2 r^3} \left[ \vec{\sigma} \cdot \vec{\sigma}' \left( 1 - r \frac{d}{dr} \right) - 3 \left( \vec{\sigma} \cdot \hat{r} \right) \left( \vec{\sigma}' \cdot \hat{r} \right) \left( 1 - r \frac{d}{dr} + \frac{1}{3} r^2 \frac{d^2}{dr^2} \right) \right] y(r) \\ \mathcal{V}_{4,5} &= -\frac{1}{2m r^2} \left( \vec{\sigma} \pm \vec{\sigma}' \right) \cdot \left( \vec{v} \times \hat{r} \right) \left( 1 - r \frac{d}{dr} \right) y(r) , \\ \mathcal{V}_{6,7} &= -\frac{1}{2m r^2} \left[ \left( \vec{\sigma} \cdot \vec{v} \right) \left( \vec{\sigma}' \cdot \hat{r} \right) \pm \left( \vec{\sigma} \cdot \hat{r} \right) \left( \vec{\sigma}' \cdot \vec{v} \right) \right] \left( 1 - r \frac{d}{dr} \right) y(r) , \\ \mathcal{V}_8 &= \frac{1}{r} \left( \vec{\sigma} \cdot \vec{v} \right) \left( \vec{\sigma}' \cdot \vec{v} \right) y(r) , \end{aligned}$$

$$\begin{aligned} \mathcal{V}_{9,10} &= -\frac{1}{2m r^2} \left( \vec{\sigma} \pm \vec{\sigma}' \right) \cdot \hat{r} \left( 1 - r \frac{d}{dr} \right) y(r) , \\ \mathcal{V}_{11} &= -\frac{1}{m r^2} \left( \vec{\sigma} \times \vec{\sigma}' \right) \cdot \hat{r} \left( 1 - r \frac{d}{dr} \right) y(r) , \\ \mathcal{V}_{12,13} &= \frac{1}{2r} \left( \vec{\sigma} \pm \vec{\sigma}' \right) \cdot \vec{v} y(r) , \\ \mathcal{V}_{14} &= \frac{1}{r} \left( \vec{\sigma} \times \vec{\sigma}' \right) \cdot \vec{v} y(r) , \\ \mathcal{V}_{15} &= -\frac{3}{2m^2 r^3} \left\{ \left[ \vec{\sigma} \cdot \left( \vec{v} \times \hat{r} \right) \right] \left( \vec{\sigma}' \cdot \hat{r} \right) + \left( \vec{\sigma} \cdot \hat{r} \right) \left[ \vec{\sigma}' \cdot \left( \vec{v} \times \hat{r} \right) \right] \right\} \\ &\quad \times \left( 1 - r \frac{d}{dr} + \frac{1}{3} r^2 \frac{d^2}{dr^2} \right) y(r) , \\ \mathcal{V}_{16} &= -\frac{1}{2m r^2} \left\{ \left[ \vec{\sigma} \cdot \left( \vec{v} \times \hat{r} \right) \right] \left( \vec{\sigma}' \cdot \vec{v} \right) + \left( \vec{\sigma} \cdot \vec{v} \right) \left[ \vec{\sigma}' \cdot \left( \vec{v} \times \hat{r} \right) \right] \right\} \left( 1 - r \frac{d}{dr} \right) y(r) \end{aligned}$$

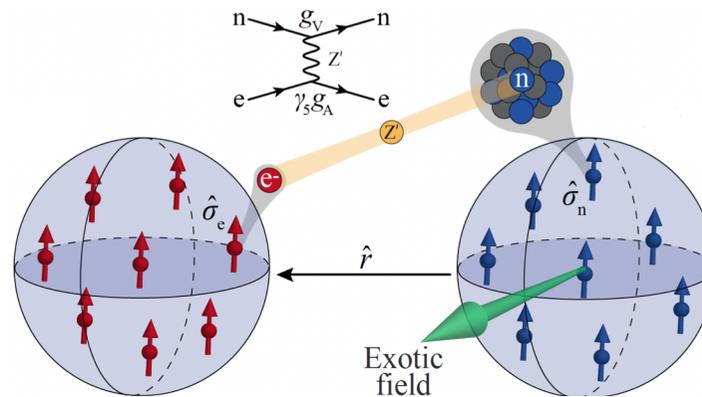
B. Dobrescu, I. Mocioiu, JHEP 11, 005 (2006)

# How to see it ?

Electron,  
Proton,  
Neutron,  
nucleon

“Spin source”

“Spin sensor”



Noble-gas  
spin amplifier

xeon-129:  
~70% neutron

Energy shift



Exotic (pseudo-magnetic) field

$$-\boldsymbol{\mu}_{\text{Xe}} \cdot \mathbf{B}_j^{\text{exo}} = V_j$$

$$\mathbf{B}_{12+13}^{\text{exo}} = -f_{12+13} \frac{\hbar}{8\pi |\boldsymbol{\mu}_{\text{Xe}}|} \iiint \rho(\hat{r})(\mathbf{v}) \left(\frac{1}{r}\right) e^{-r/\lambda} d\mathbf{r}$$

Coupling to be measured

More sensitive detector



Smaller coupling  $f$

# Searched exotic interactions in our study

$$\mathcal{V}_1 = \frac{1}{r} y(r) ,$$

$$\mathcal{V}_2 = \frac{1}{r} \vec{\sigma} \cdot \vec{\sigma}' y(r) ,$$

$$\mathcal{V}_3 = \frac{1}{m^2 r^3} \left[ \vec{\sigma} \cdot \vec{\sigma}' \left( 1 - r \frac{d}{dr} \right) - 3 (\vec{\sigma} \cdot \hat{r}) (\vec{\sigma}' \cdot \hat{r}) \left( 1 - r \frac{d}{dr} + \frac{1}{3} r^2 \frac{d^2}{dr^2} \right) \right] y(r)$$

$$\mathcal{V}_{4,5} = -\frac{1}{2m r^2} (\vec{\sigma} \pm \vec{\sigma}') \cdot (\vec{v} \times \hat{r}) \left( 1 - r \frac{d}{dr} \right) y(r) ,$$

$$\mathcal{V}_{6,7} = -\frac{1}{2m r^2} \left[ (\vec{\sigma} \cdot \vec{v}) (\vec{\sigma}' \cdot \hat{r}) \pm (\vec{\sigma} \cdot \hat{r}) (\vec{\sigma}' \cdot \vec{v}) \right] \left( 1 - r \frac{d}{dr} \right) y(r) ,$$

$$\mathcal{V}_8 = \frac{1}{r} (\vec{\sigma} \cdot \vec{v}) (\vec{\sigma}' \cdot \vec{v}) y(r) ,$$

$$\mathcal{V}_{9,10} = -\frac{1}{2m r^2} (\vec{\sigma} \pm \vec{\sigma}') \cdot \hat{r} \left( 1 - r \frac{d}{dr} \right) y(r) ,$$

$$\mathcal{V}_{11} = -\frac{1}{m r^2} (\vec{\sigma} \times \vec{\sigma}') \cdot \hat{r} \left( 1 - r \frac{d}{dr} \right) y(r) ,$$

$$\mathcal{V}_{12,13} = \frac{1}{2r} (\vec{\sigma} \pm \vec{\sigma}') \cdot \vec{v} y(r) ,$$

$$\mathcal{V}_{14} = \frac{1}{r} (\vec{\sigma} \times \vec{\sigma}') \cdot \vec{v} y(r) ,$$

$$\mathcal{V}_{15} = -\frac{3}{2m^2 r^3} \left\{ [\vec{\sigma} \cdot (\vec{v} \times \hat{r})] (\vec{\sigma}' \cdot \hat{r}) + (\vec{\sigma} \cdot \hat{r}) [\vec{\sigma}' \cdot (\vec{v} \times \hat{r})] \right\} \times \left( 1 - r \frac{d}{dr} + \frac{1}{3} r^2 \frac{d^2}{dr^2} \right) y(r) ,$$

$$\mathcal{V}_{16} = -\frac{1}{2m r^2} \left\{ [\vec{\sigma} \cdot (\vec{v} \times \hat{r})] (\vec{\sigma}' \cdot \vec{v}) + (\vec{\sigma} \cdot \vec{v}) [\vec{\sigma}' \cdot (\vec{v} \times \hat{r})] \right\} \left( 1 - r \frac{d}{dr} \right) y(r)$$

V3 exp.

polarized electron/proton  
(rubidium vapor)

force mediator: axion

V11 exp.

polarized electron/proton  
(rubidium vapor)

force mediator: Z' bosons

V4,5 and V12,13 exp.

Unpolarized nucleon  
(BGO crystal)

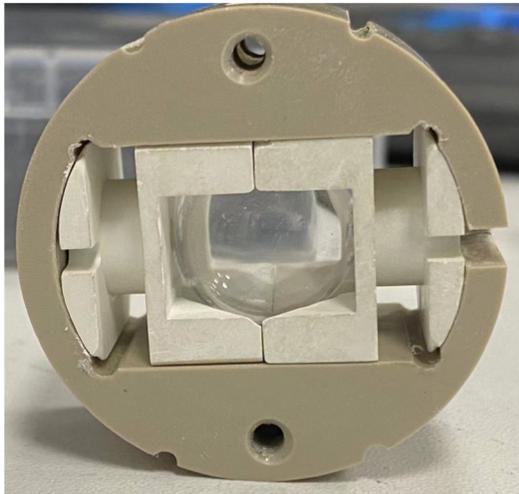
# Various spin sources

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## **I: electron/proton (Rb) vapor**

high polarization

Short force range (small size)



Electron/proton-neutron coupling

## **II: BGO crystal**

high nucleon density

Non-magnetic effect



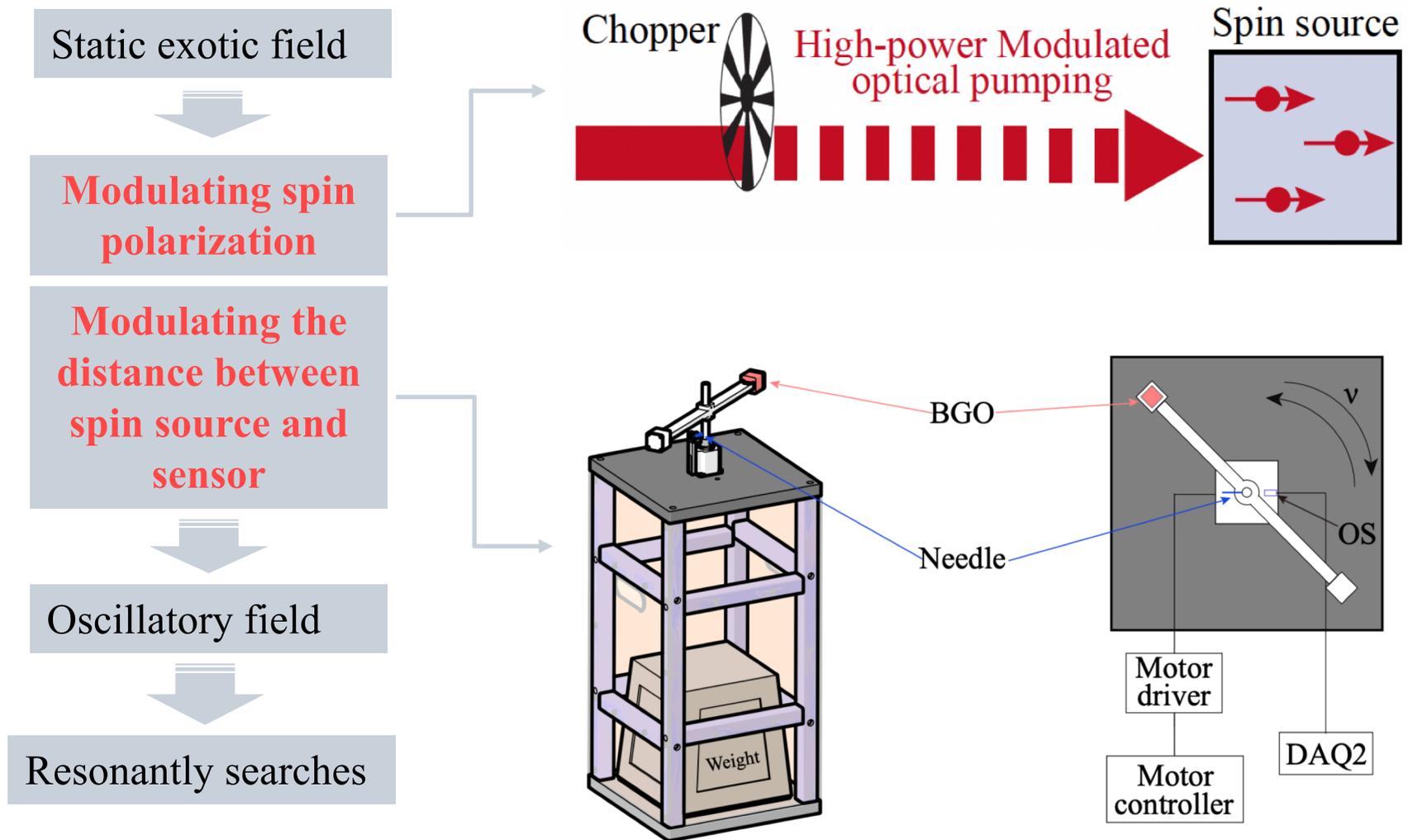
Nucleon-neutron coupling

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**Develop spin sources for different exotic interaction searches**

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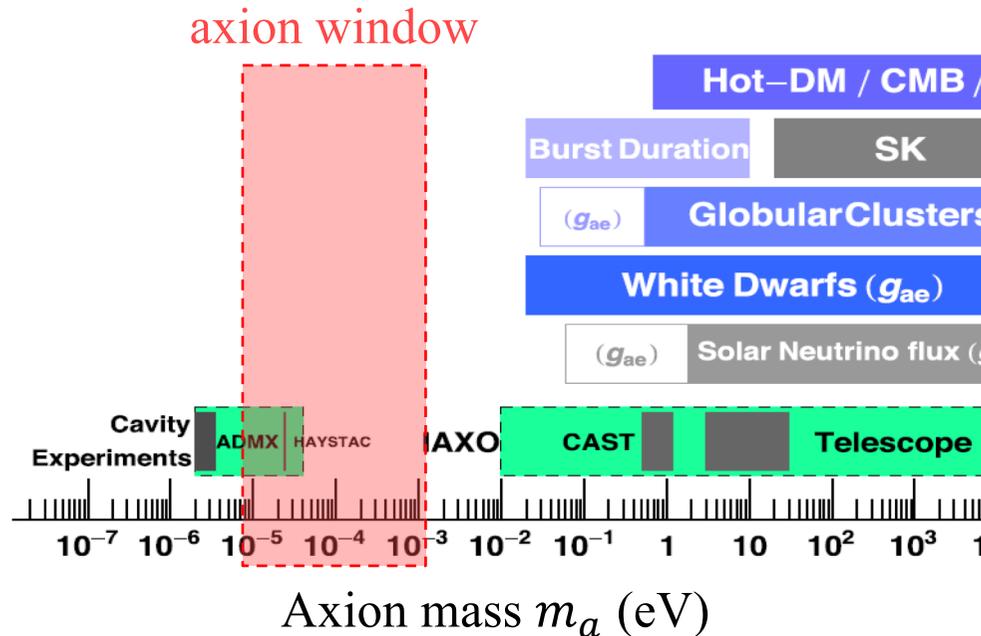
# Resonant detection with spin amplifiers



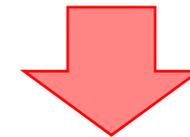
# Searches for an axion-mediated interaction

Axion-mediated dipole-dipole interaction

$$V_{pp} = -\frac{g_p^1 g_p^2}{4} \left[ (\hat{\sigma}_1 \cdot \hat{\sigma}_2) \left( \frac{m_a}{r^2} + \frac{1}{r^3} \right) - (\hat{\sigma}_1 \cdot \hat{r})(\hat{\sigma}_2 \cdot \hat{r}) \left( \frac{m_a^2}{r} + \frac{3m_a}{r^2} + \frac{3}{r^3} \right) \right] \frac{e^{-m_a r}}{4\pi m_1 m_2}$$



High-temperature QCD  
SMASH model  
Axion string networks

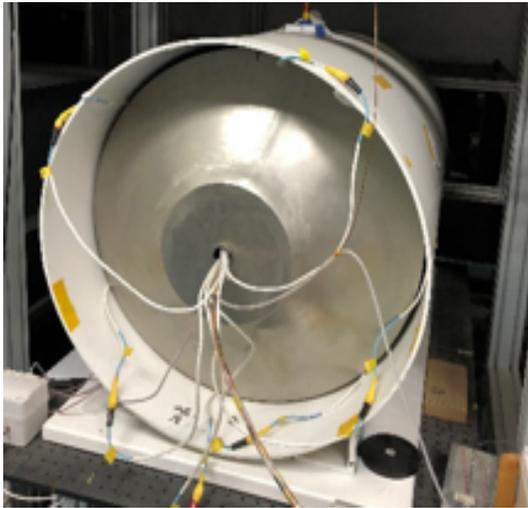


Axion window  
(  $10 \mu\text{eV} \sim 1 \text{meV}$ , )  
(  $\lesssim 2\text{cm}$  force range )

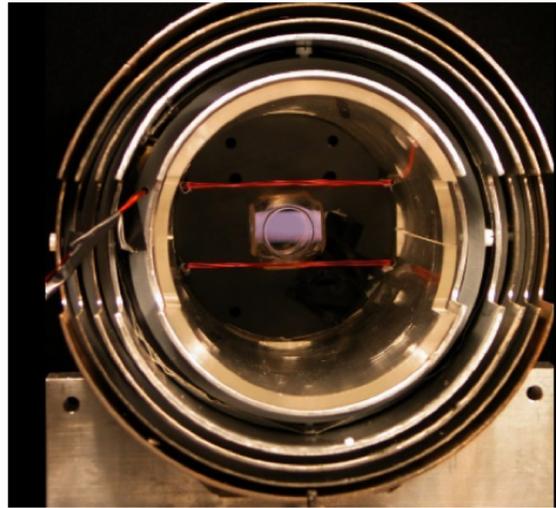
Spurious dipole field: Search remains challenging in the axion window

# How to shield the spurious ordinary field?

High-permeability materials:  $\mu$ -metal



Meter scale



Decimeter scale



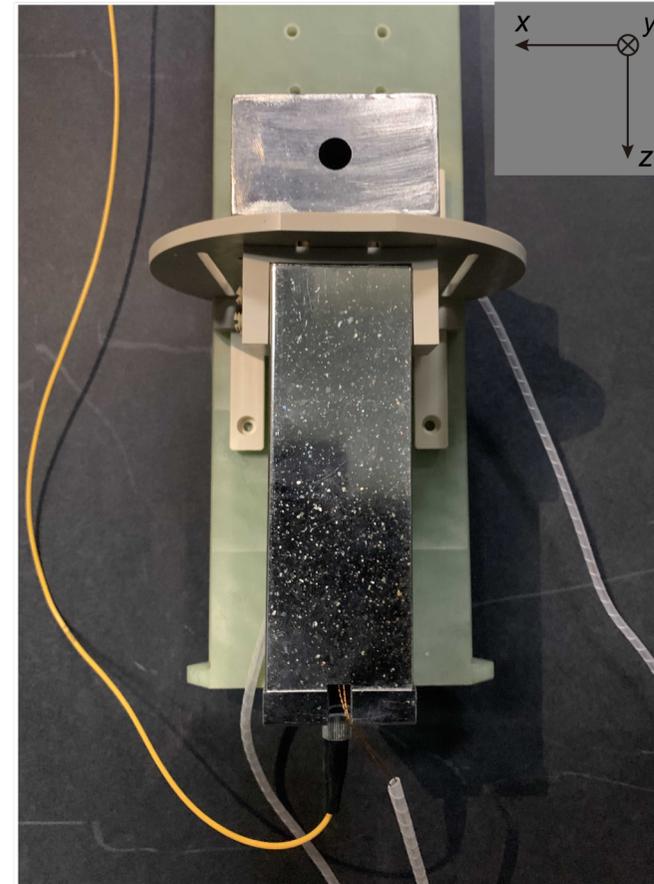
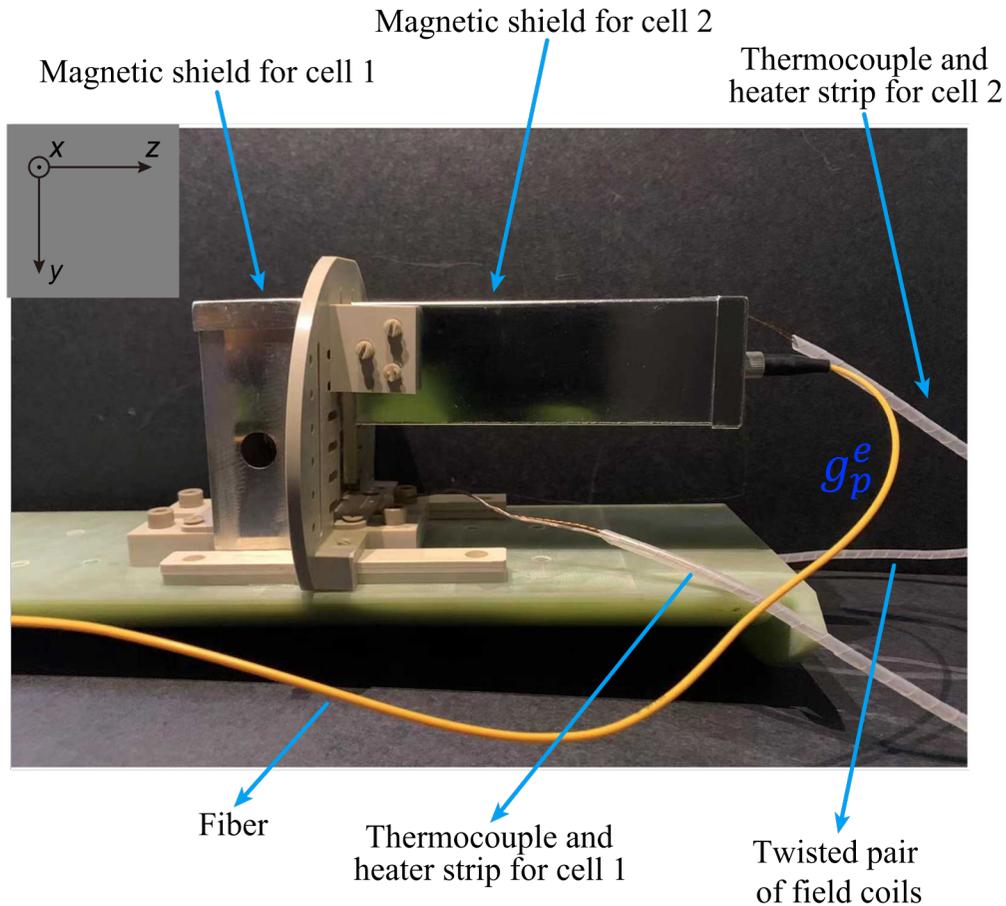
Centimeter scale

Axion window

Mass range  $\longrightarrow$  Force range  
( $10 \mu\text{eV} \sim 1 \text{meV}$ )  $\longrightarrow$  ( $0.2 \text{mm} \sim 20\text{mm}$ )

Y. Wang, H. Su, **M. Jiang**<sup>†</sup>, X. Peng<sup>†</sup> *et al.*, PRL 129, 051801 (2022)

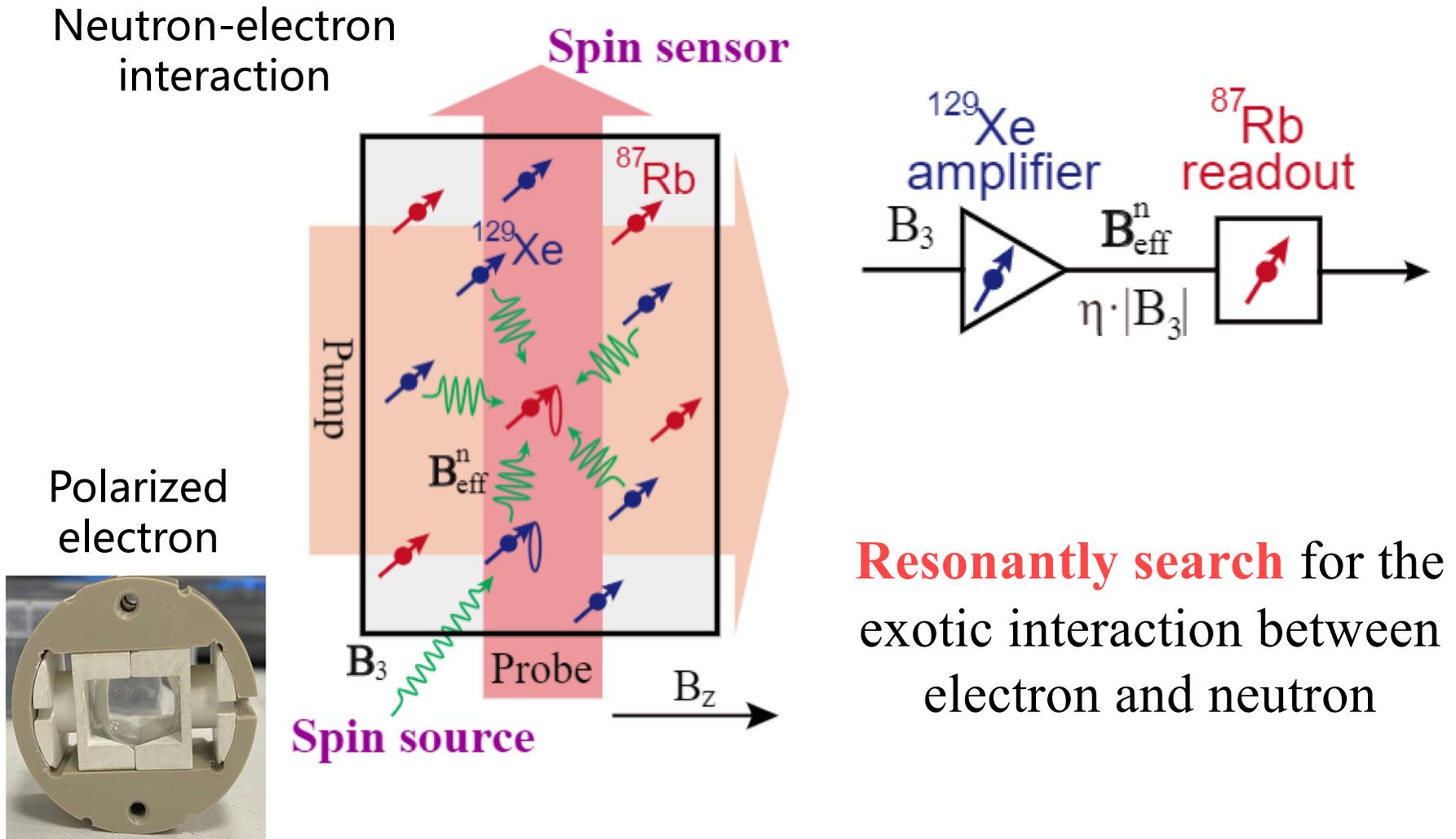
# How to shield the spurious ordinary field?



Spurious dipole field is suppressed at least  $\sim 10^4$

Y. Wang, H. Su, **M. Jiang**<sup>†</sup>, X. Peng<sup>†</sup> *et al.*, PRL 129, 051801 (2022)

# Search for an axion-mediated interaction



Y. Wang, H. Su, **M. Jiang**<sup>†</sup>, X. Peng<sup>†</sup> *et al.*, PRL 129, 051801 (2022)

# Search for an axion-mediated interaction

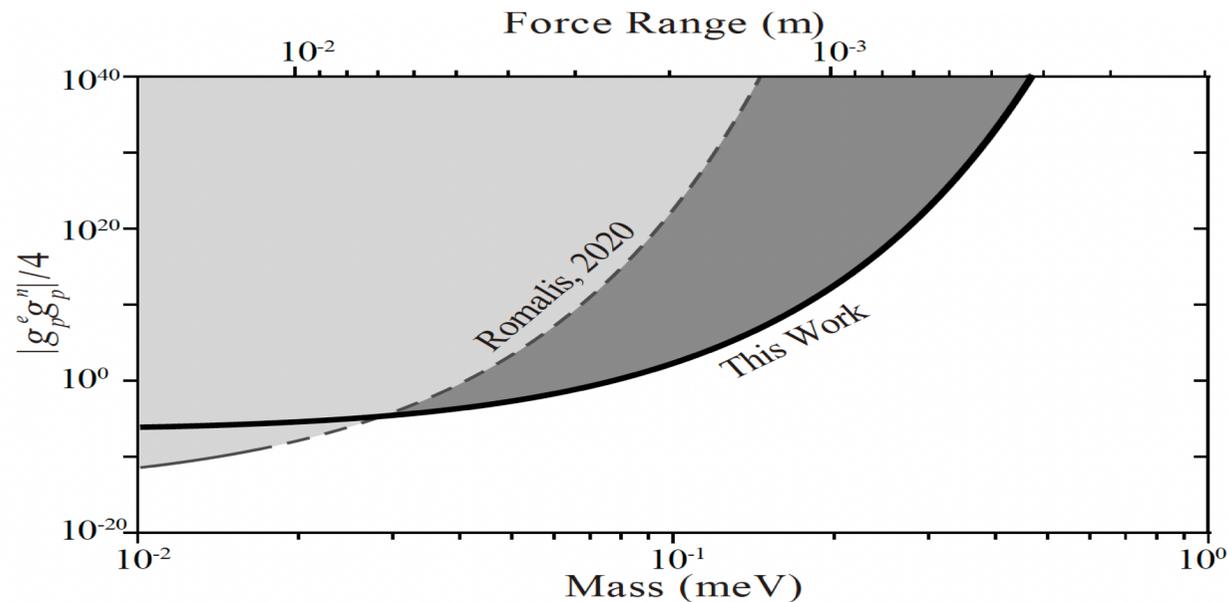
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## Limits on Axions and Axionlike Particles within the Axion Window Using a Spin-Based Amplifier

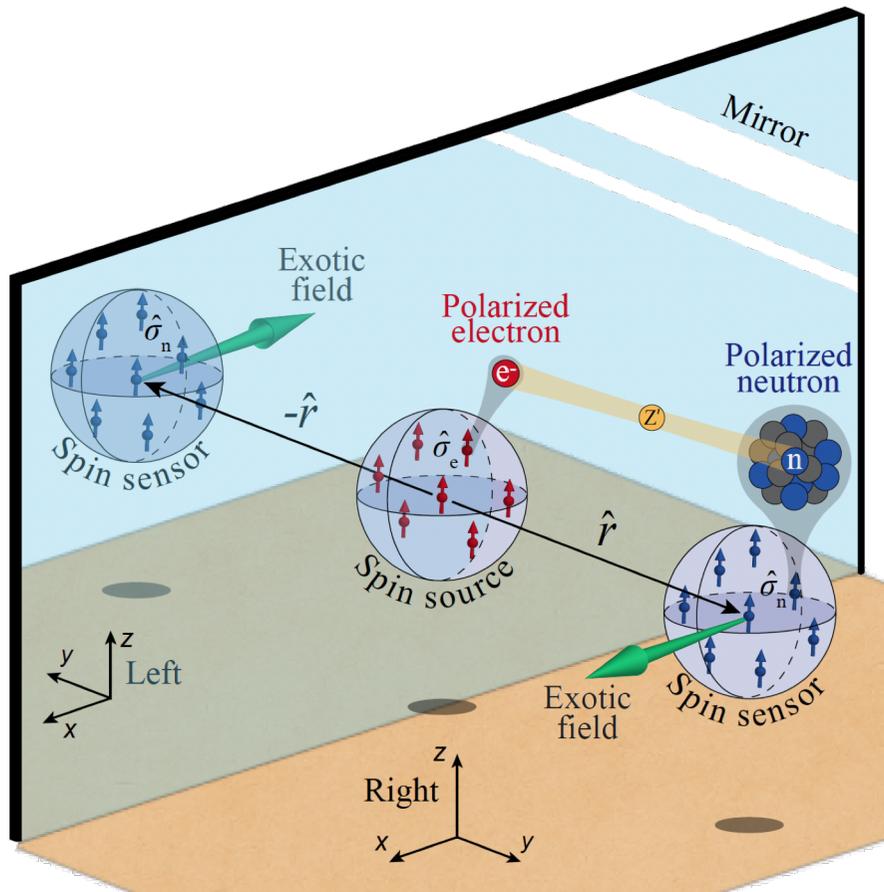
Yuanhong Wang, Haowen Su, Min Jiang, Ying Huang, Yushu Qin, Chang Guo, Zehao Wang, Dongdong Hu, Wei Ji, Pavel Fadeev, Xinhua Peng, and Dmitry Budker  
Phys. Rev. Lett. **129**, 051801 – Published 25 July 2022



The most stringent constraints on  $g_p^e g_p^n$  within the axion window

Y. Wang, H. Su, **M. Jiang**<sup>†</sup>, X. Peng<sup>†</sup> *et al.*, PRL 129, 051801 (2022)

# Search for exotic parity-violation interactions



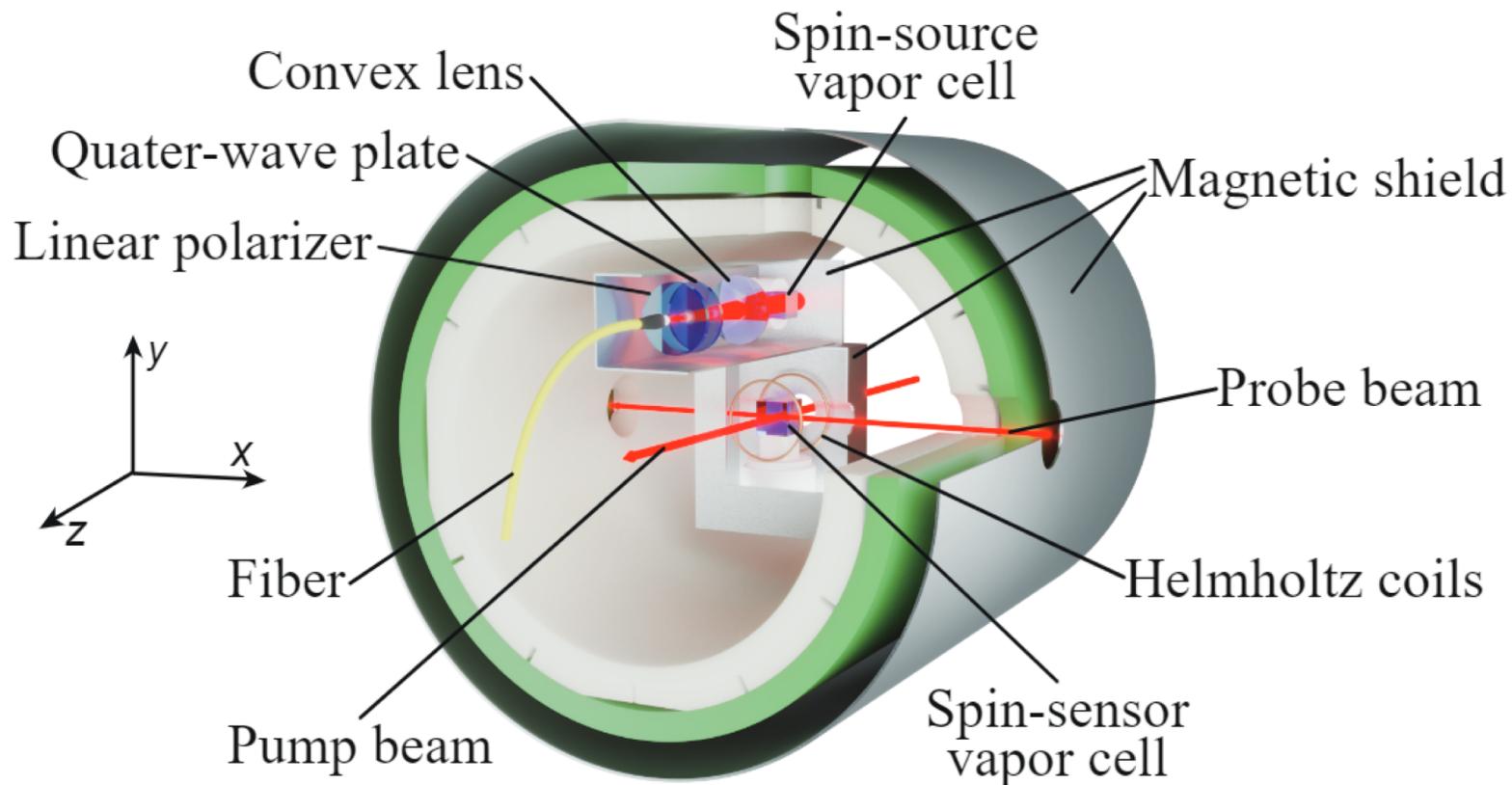
$$\begin{array}{c}
 n \rightarrow \begin{array}{c} g_V \\ \text{---} \\ Z' \\ \text{---} \\ \gamma_5 g_A \end{array} \rightarrow n \\
 e \rightarrow \begin{array}{c} \gamma_5 g_A \\ \text{---} \\ Z' \\ \text{---} \\ g_V \end{array} \rightarrow e \\
 + \\
 n \rightarrow \begin{array}{c} \gamma_5 g_A \\ \text{---} \\ Z' \\ \text{---} \\ g_V \end{array} \rightarrow n \\
 e \rightarrow \begin{array}{c} g_V \\ \text{---} \\ Z' \\ \text{---} \\ \gamma_5 g_A \end{array} \rightarrow e
 \end{array}$$

Parity-odd spin-spin interaction mediated by  $Z'$  boson

arXiv:2205.07222 (2022); under review in Science Advances

# Search for exotic parity-violation interactions

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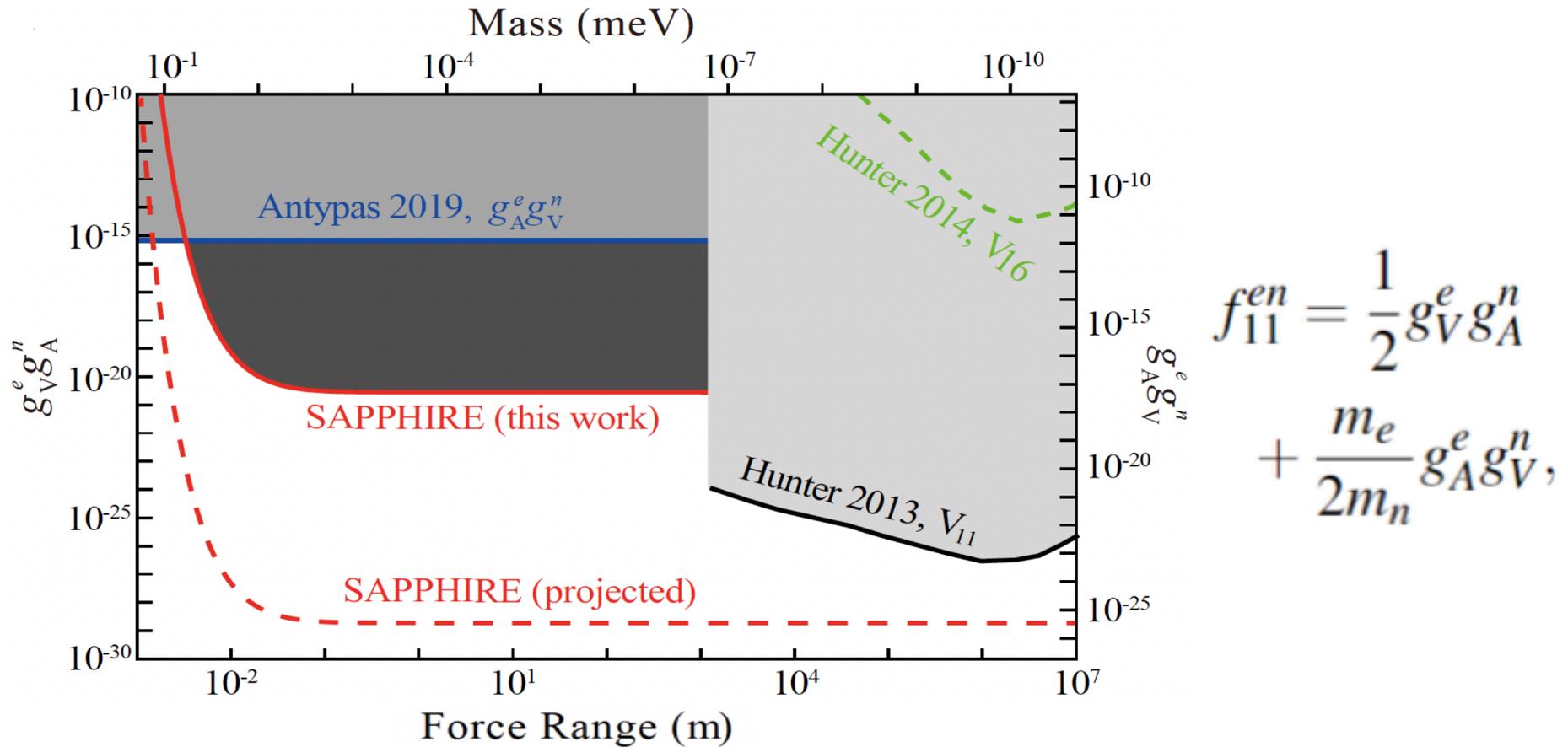


## Experimental setup

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arXiv:2205.07222 (2022); under review in Science Advances

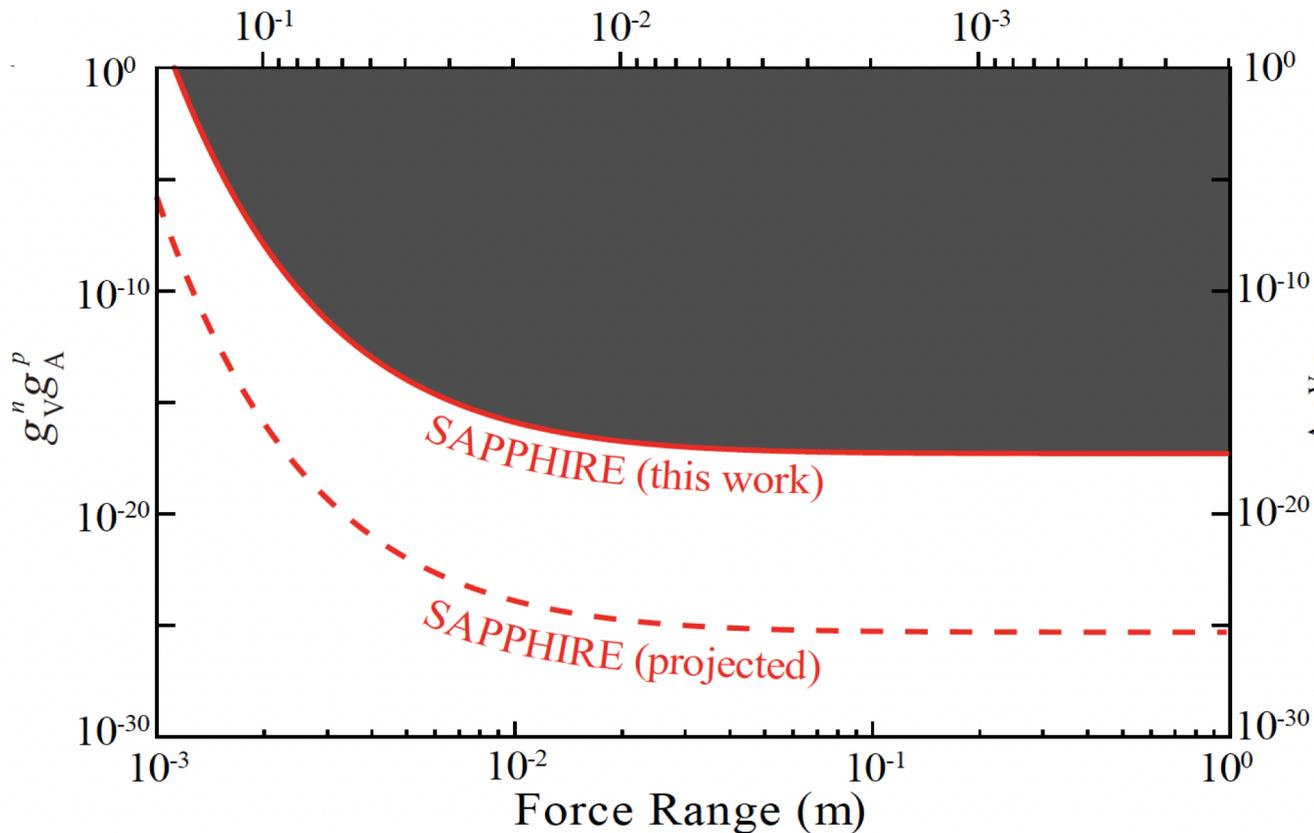
# Constraints on electron-neutron coupling



A five-order-of-magnitude improvement

arXiv:2205.07222 (2022); under review in Science Advances

# Constraints on proton-neutron coupling



$$f_{11}^{np} = \frac{m_e}{2m_p} g_A^n g_V^p + \frac{m_e}{2m_n} g_V^n g_A^p$$

$p$  – valence protons within the Rb nuclei in the source

$n$  – valence neutrons within the Xe nuclei in the sensor

arXiv:2205.07222 (2022); under review in Science Advances

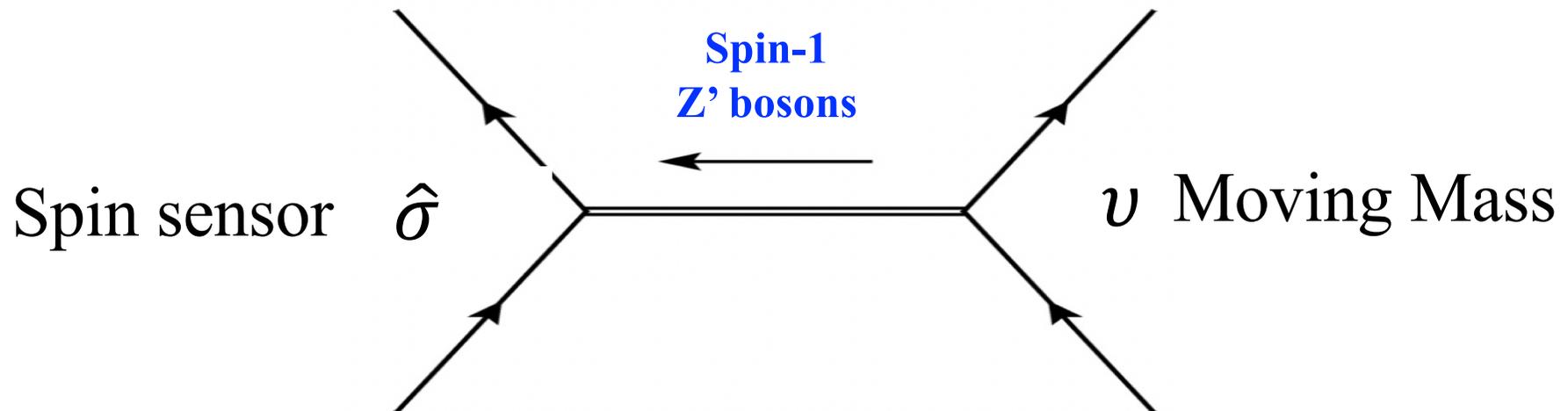
# Search for exotic spin-dependent interactions

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$$V_{4+5} = -f_{4+5} \frac{\hbar^2}{8\pi mc} [\hat{\sigma} \cdot (\mathbf{v} \times \hat{r})] \left( \frac{1}{\lambda r} + \frac{1}{r^2} \right) e^{-r/\lambda},$$

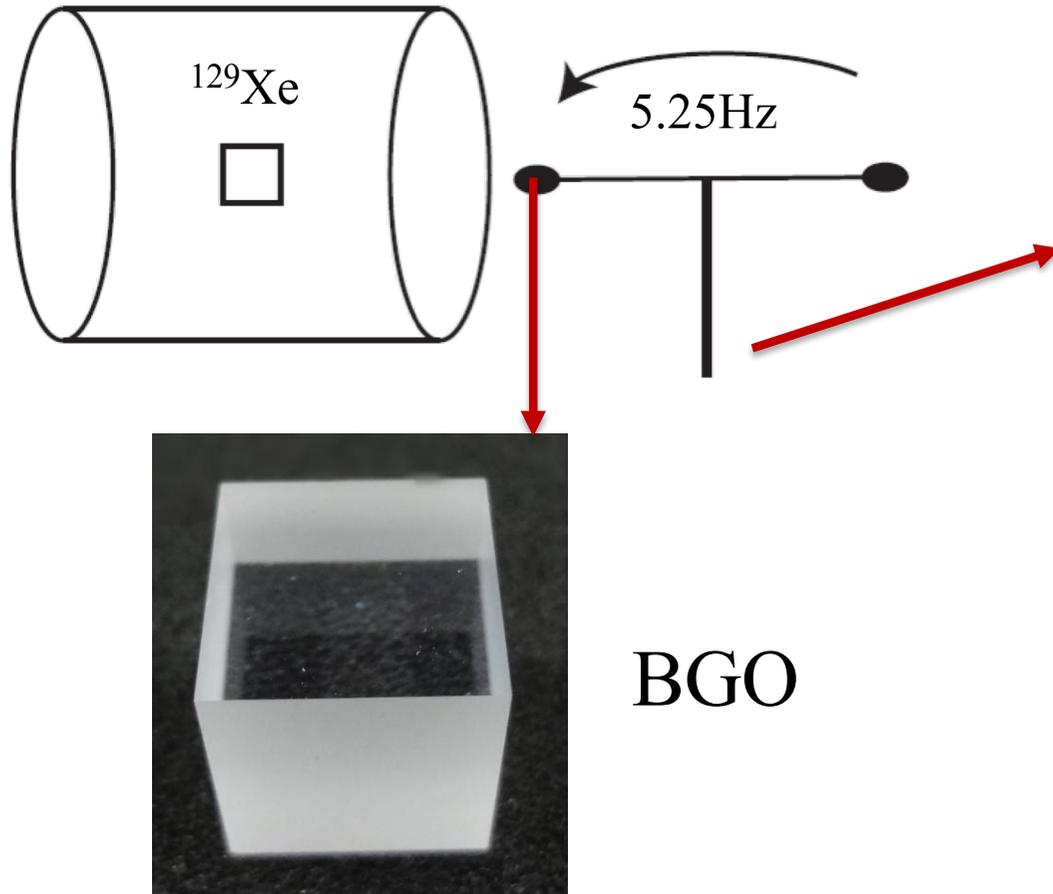
$$V_{12+13} = f_{12+13} \frac{\hbar}{8\pi} (\hat{\sigma} \cdot \mathbf{v}) \left( \frac{1}{r} \right) e^{-r/\lambda},$$

Velocity dependence

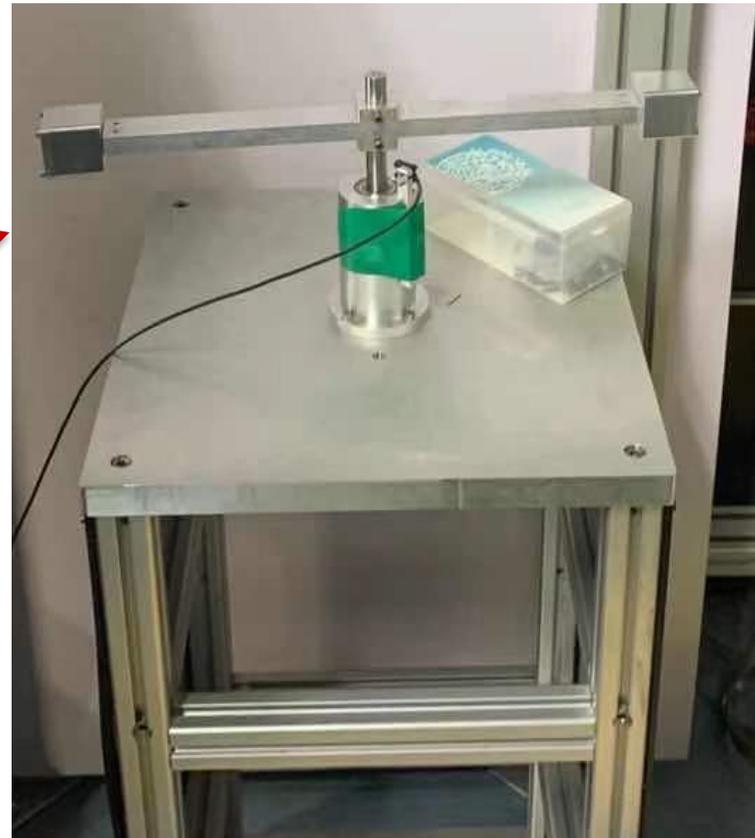


# Search setup

Spin-mass coupling

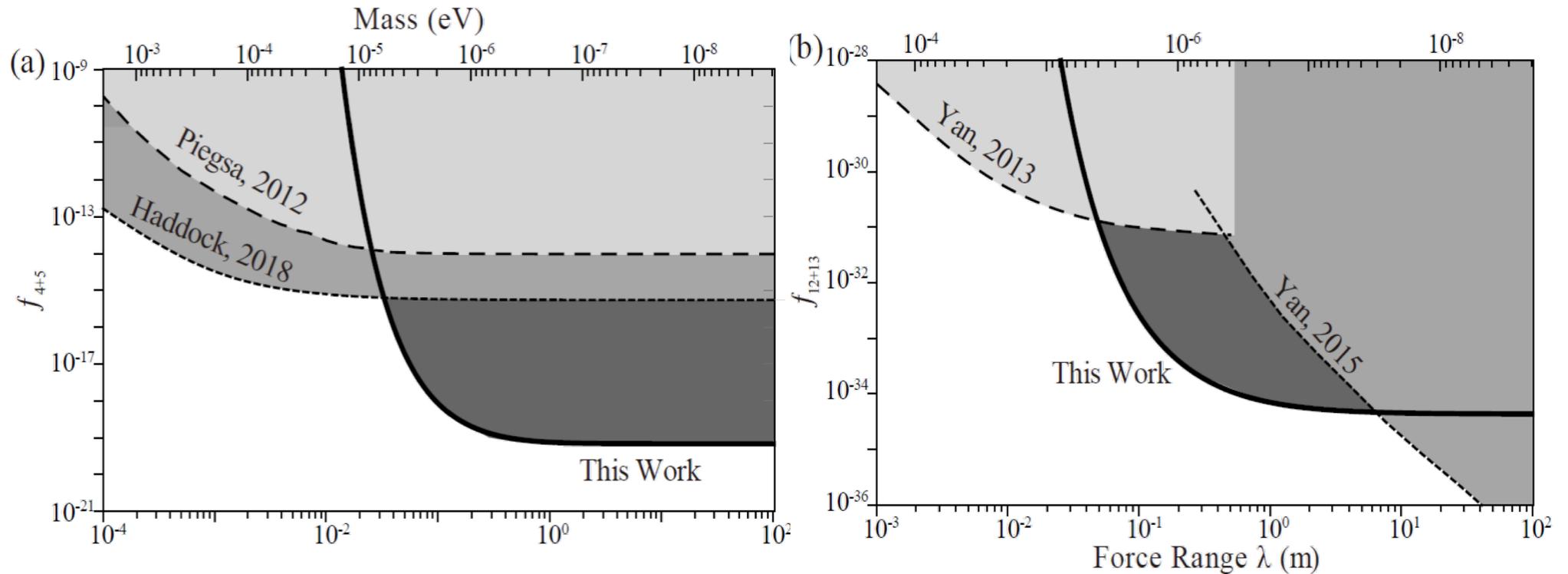


Rotator



H. Su\*, Y. Wang\*, M. Jiang#, X. Peng# *et al.* [Science Advances](#) 7, eabi9535 (2021).

# Constraints on spin-dependent interactions



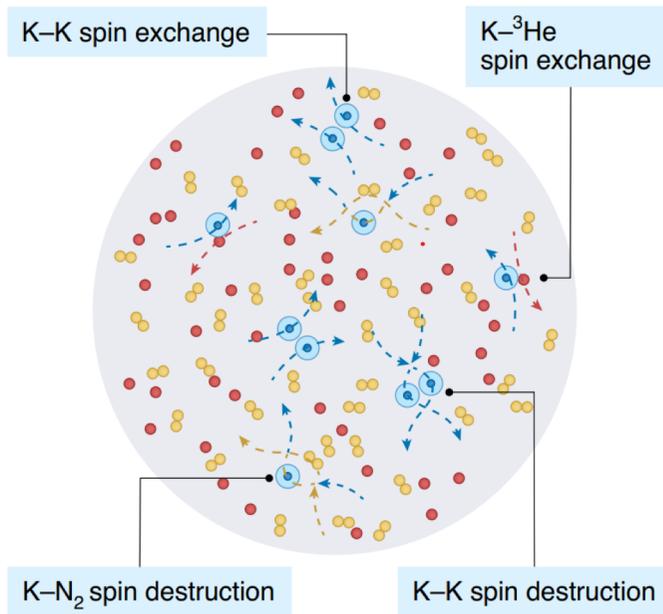
**At least 2 orders of magnitude improvement on constraints on  $Z'$  boson**

H. Su, Y. Wang, M. Jiang<sup>†</sup>, X. Peng<sup>†</sup> *et al.* [Science Advances](#) 7, eabi9535 (2021).



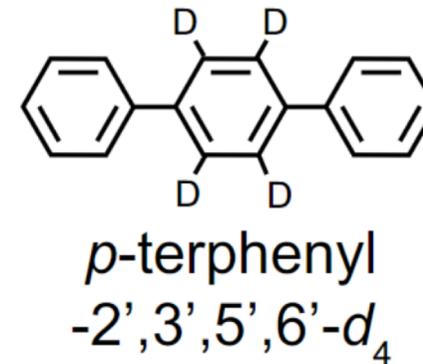
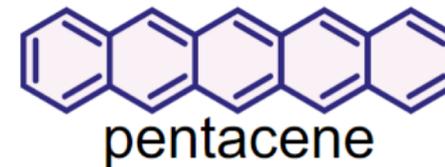
# SAPPHIRE projected sensitivity

## $^3\text{He}$ -K spin amplifier



4 orders of magnitude improvement

## Solid-state spin source



$10^{14} \text{ cm}^{-3}$   $\rightarrow$   $10^{18} \text{ cm}^{-3}$

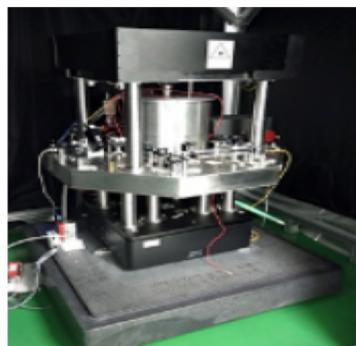
**8 orders of magnitude improvements are ongoing!**

# 正在开展: aT级别弱磁测量新方法和技术

国际公开报道最高指标

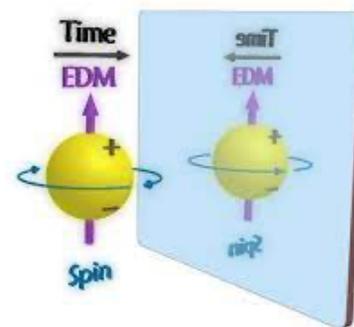


普林斯顿160aT/Hz<sup>1/2</sup>



北航680aT/Hz<sup>1/2</sup>

亟需突破



aT测量技术：开启科学新大门

**核心挑战：弱磁探测响应机理，厘清经典/量子噪声机制及抑制技术**



**新方法**

**基于原子自旋新效应：**  
自旋量子放大  
自旋协同效应  
多体量子相变等

**新技术**



**发展超低噪声器件及噪声抑制技术：**  
高性能原子气室设计和加工  
(高压、长寿命、低漏率)  
超导屏蔽技术、超稳激光等

# 正在开展：“蓝宝石”计划及暗物质测量网络

- ✓ 研究趋势：从单个传感器到阵列式；从单一类型到多种类型
- 提高灵敏度  $\propto 1/\sqrt{N}$       显著降低报警率  $\propto p^N$       结构信息

## “蓝宝石”计划发展规划



USTC : 1台  
2019-2020年



USTC : 4台  
2021年

4个城市 : 5台  
2022年

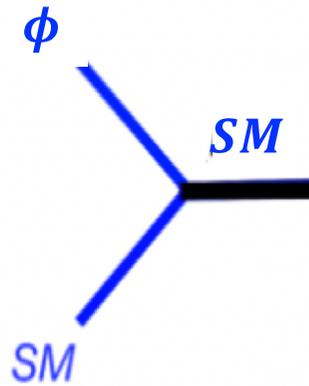
空间站、月球与深空  
2022-

**建成5台探测器阵列（合肥，杭州，苏州，哈尔滨）  
已完成2个月数据采集和分析，即将公布结果！**

基金委原创探索项目（2021-2024）

# Summary

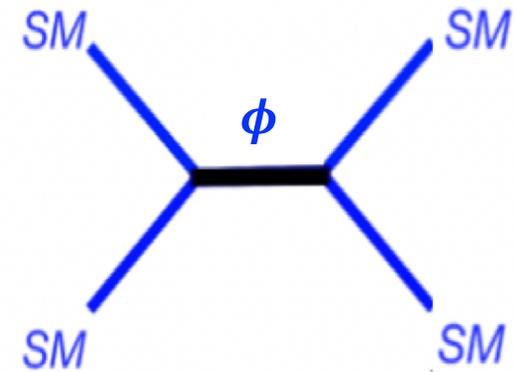
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Search for axions,  
axion-like  
particles, dark  
photon



Develop sensitive  
spin amplifier for  
particle physics  
(**SAPHIRE**)



Exotic fifth force  
between SM  
particles

# Thanks to .....

---

Prof. Jiangfeng Du  
Prof. Zhengguo Zhao

## Collaborators:

Prof. Dmitry Budker  
Prof. Jing Shu  
Dr. Yifan Chen  
Dr. Yue Zhao  
Dr. Xiao Xue

.....

## Funding:

USTC, CAS, NNSFC.....



Dmitry Budker

## NMR group (especially low-field NMR group):

Dr. Min Jiang  
Haowen Su  
Yuanhong Wang  
Qing Li  
Shiming Song  
Ze Wu  
Minxiang Xu  
Yushu Qing

.....



Min Jiang

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Thank you for your attention!