



Axion-like particle solution to muon $g-2$ and its test at Z-factory

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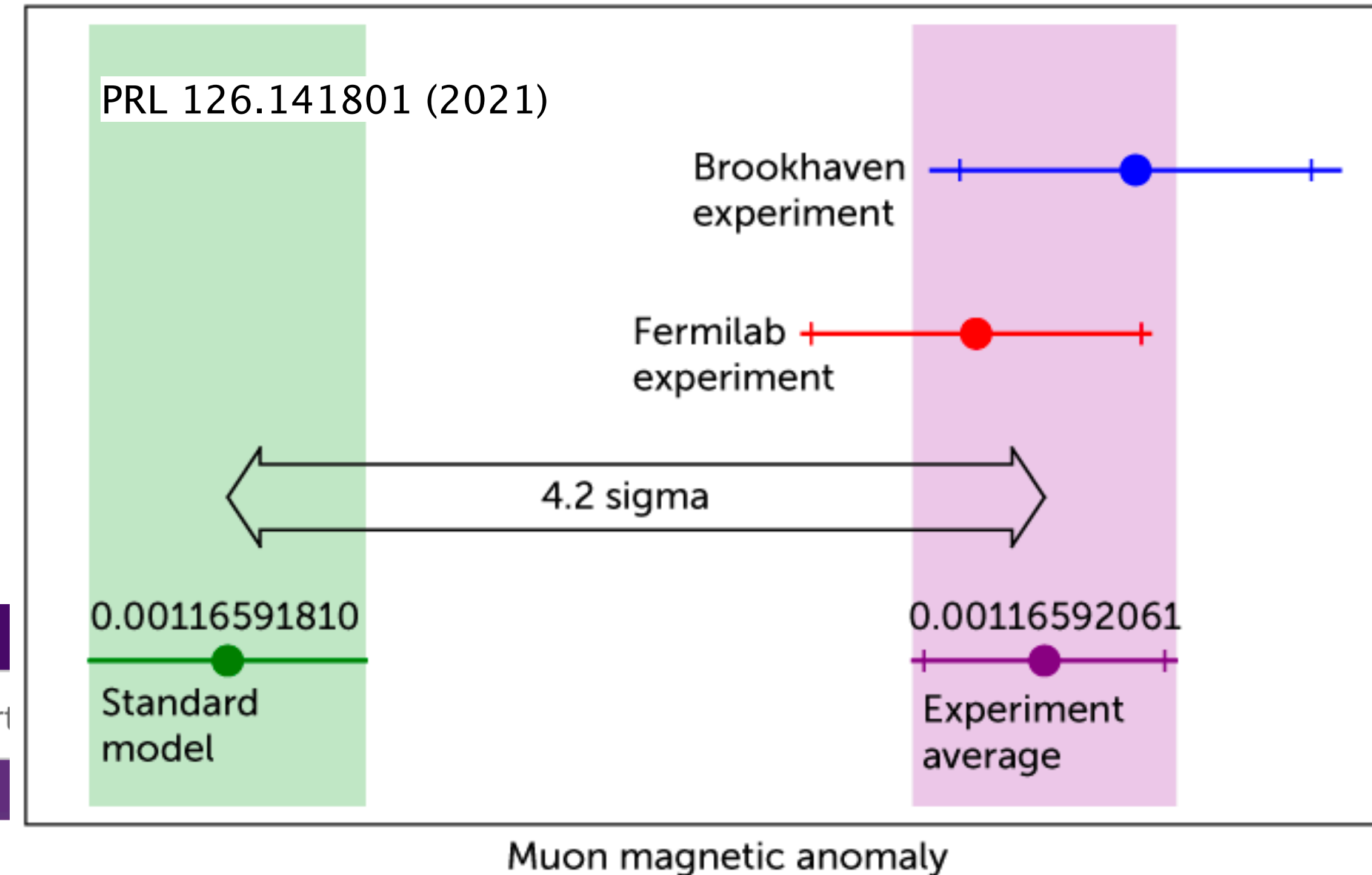
Collaborated with: Xiaolin Ma, Lian-tao Wang, Xiao-ping Wang
Preliminary results

Joint Workshop of the CEPC Physics, Software and New Detector Concept
2022-05-24

The Muon g-2 anomaly

- Positive value and a 4.2 σ (Fermilab + Brookhaven)

$$\Delta a_{\mu} = a_{\mu}^{\text{exp}} - a_{\mu}^{\text{th}} = (25.1 \pm 5.9) \times 10^{-10}$$



Physics

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The Era of Anomalies

May 14, 2020 • *Physics* 13, 79

Particle physicists are faced with a growing list of “anomalies”—experimental results that conflict with the standard model but fail to overturn it for lack of sufficient evidence.

The (pseudo)scalar answer to g-2

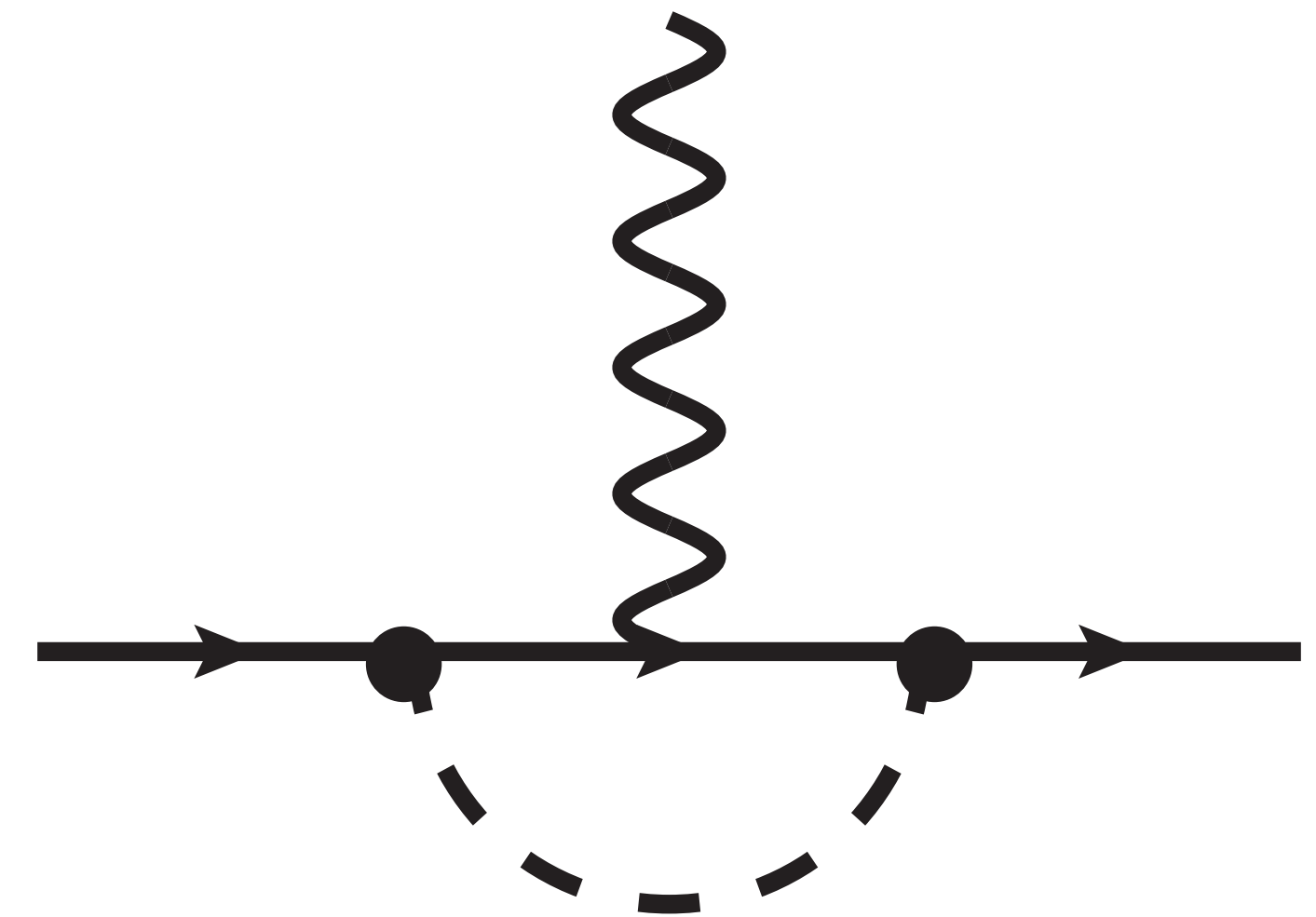
- The (pseudo)scalar Yukawa coupling to lepton

$$\mathcal{L}_{\text{yuk}} = \phi \bar{\ell} (g_R + i g_I \gamma_5) \ell$$

- The 1-loop contribution to g-2

$$\Delta a_\ell = \frac{1}{8\pi^2} \int_0^1 dx \frac{(1-x)^2 ((1+x)g_R^2 - (1-x)g_I^2)}{(1-x)^2 + x (m_\phi/m_\ell)^2}$$

- For scalar, $\Delta a_\ell > 0$
- For (pseudo)scalar, $\Delta a_\ell < 0$



The pseudo-scalar solution

- Further requiring photon coupling

$$\mathcal{L} = \frac{1}{4} g_{a\gamma\gamma} \tilde{F}F + iy_{a\psi} a \bar{\psi} \gamma_5 \psi$$

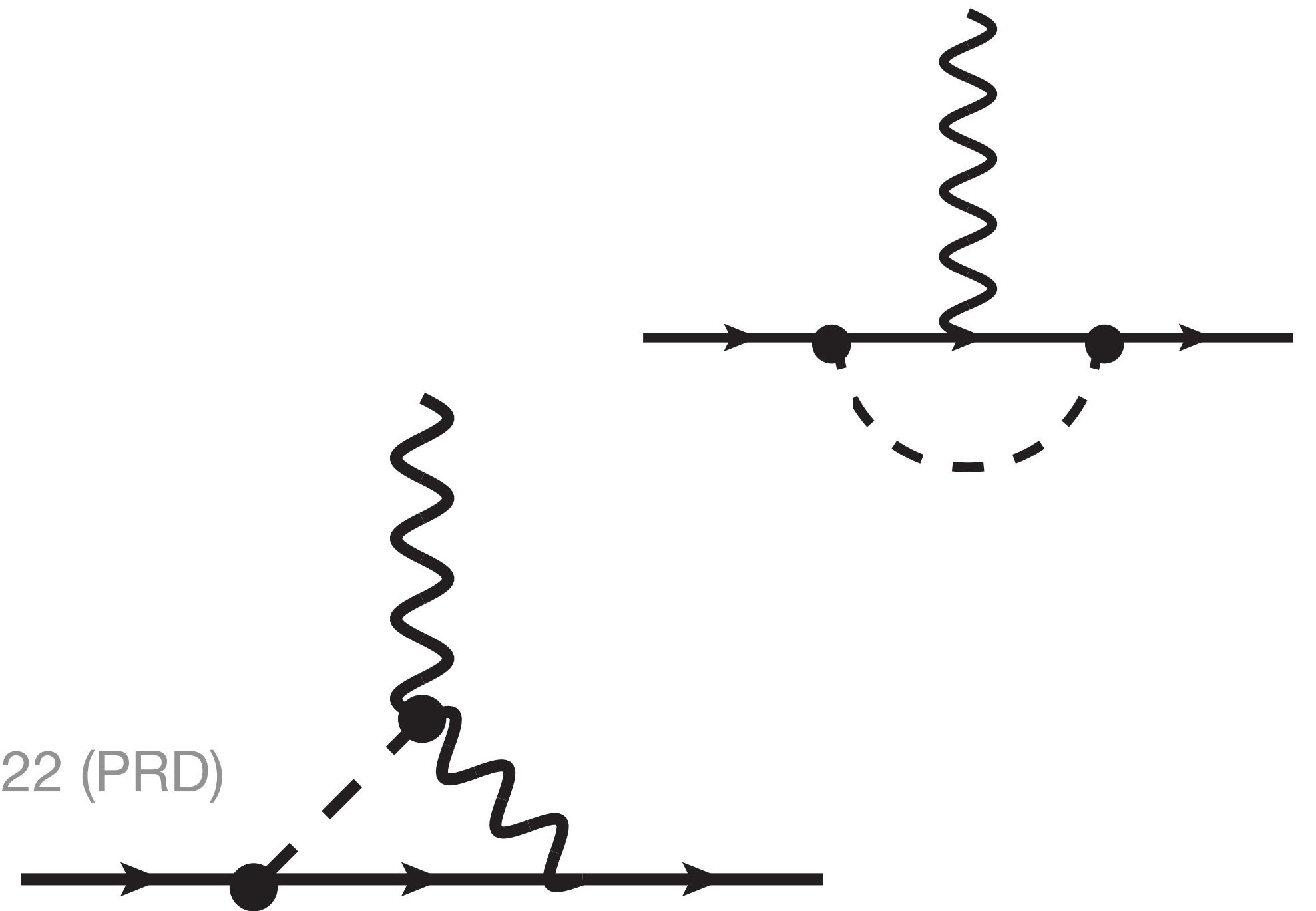
- The 1-loop BZ contribution

$$\Delta a_\ell^{\text{BZ}} \approx \frac{m_\ell}{4\pi^2} g_{a\gamma\gamma} y_{a\ell} \log \frac{\Lambda}{m_a}$$

Marciano et al, 1607.01022 (PRD)

- Assumes $g_{a\gamma\gamma}$ remains essentially constant throughout the integration over virtual photon-loop momentum

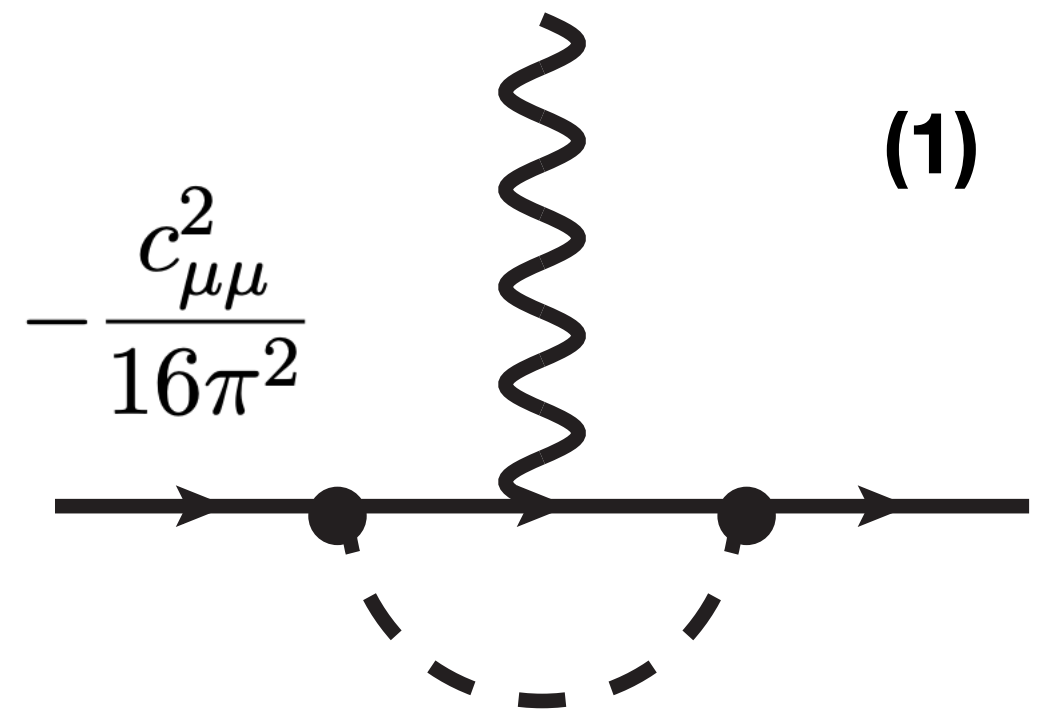
- $g_{a\gamma\gamma}$ and $y_{a\ell}$ can adjust its sign to give positive result



A complete calculation for axion-like particle

- The axion-like particle Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_f \frac{C_{ff}}{2} \frac{\partial^\mu a}{f_a} \bar{f} \gamma_\mu \gamma_5 f + \frac{g^2}{16\pi^2} C_{WW} \frac{a}{f_a} W_{\mu\nu}^i \tilde{W}^{\mu\nu,i} + \frac{g^2}{16\pi^2} C_{BB} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu}$$

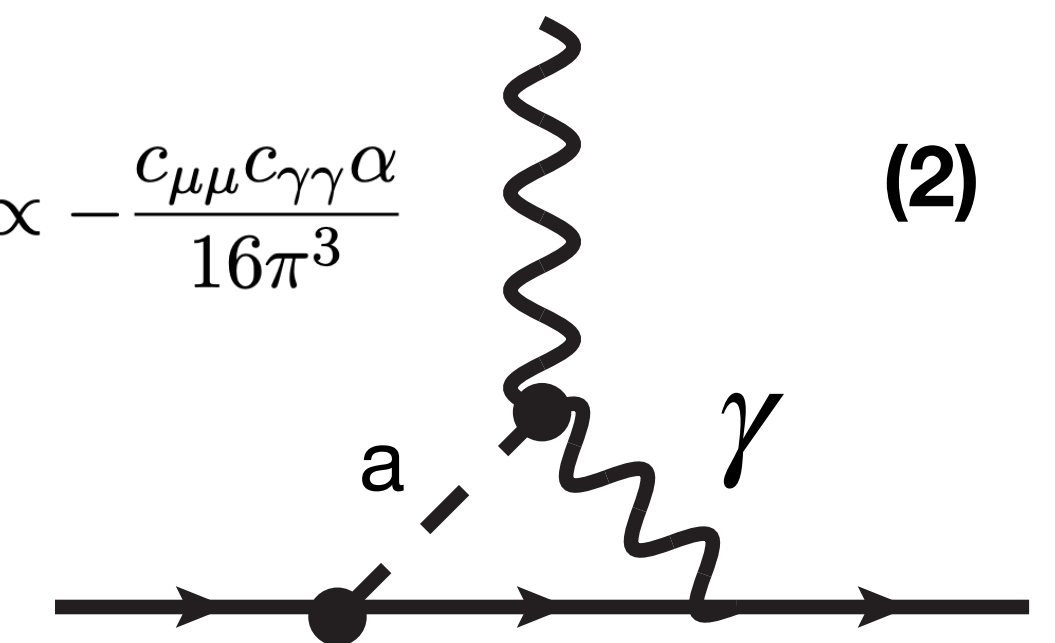
$$\Delta a_\mu^{(1)} \propto -\frac{c_{\mu\mu}^2}{16\pi^2}$$


(1)

- The 3rd diagram subtlety:

Bauer et al, 1708.00443 (JHEP)

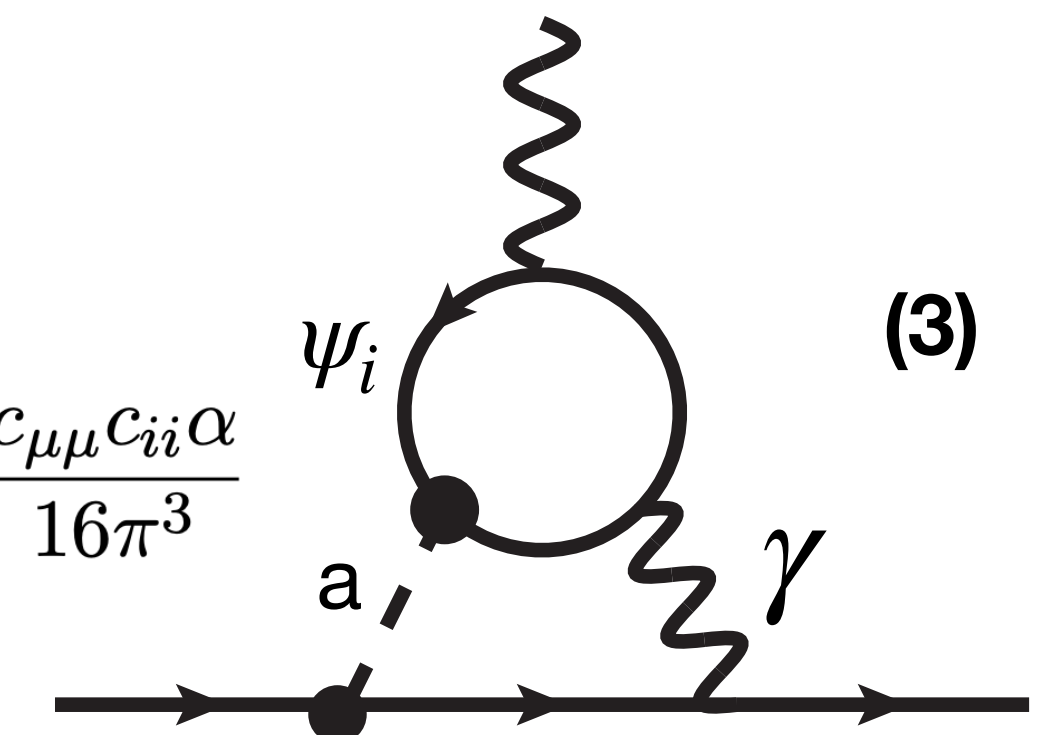
- Previously, calculate Barr-Zee diagram for on-shell ALP and photon and use the result of 2nd diagram

$$\Delta a_\mu^{(2)} \propto -\frac{c_{\mu\mu} c_{\gamma\gamma\alpha}}{16\pi^3}$$


(2)

- Recently, [calculate off-shell ALP and photon](#)

Buen-Abad et al, 2104.03267 (JHEP)

$$\Delta a_\mu^{(3)} \propto -\frac{c_{\mu\mu} c_{ii\alpha}}{16\pi^3}$$


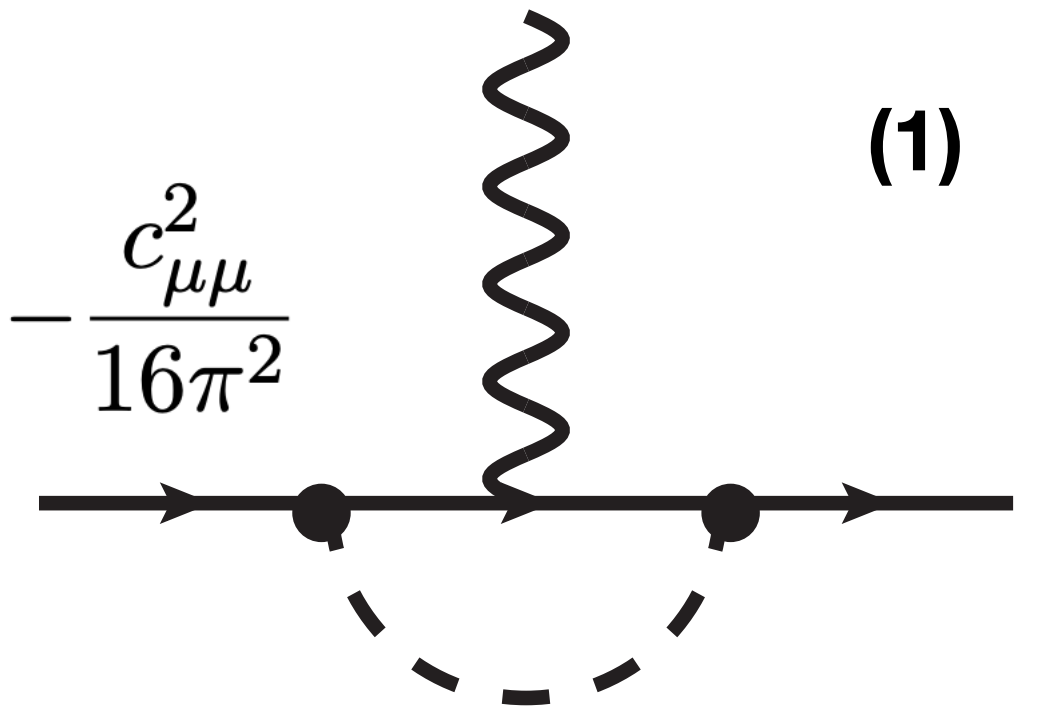
(3)

A complete calculation for axion-like particle

- The axion-like particle Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_f \frac{C_{ff}}{2} \frac{\partial^\mu a}{f_a} \bar{f} \gamma_\mu \gamma_5 f + \frac{g^2}{16\pi^2} C_{WW} \frac{a}{f_a} W_{\mu\nu}^i \tilde{W}^{\mu\nu,i} + \frac{g^2}{16\pi^2} C_{BB} \frac{a}{f_a} B_{\mu\nu} \tilde{B}^{\mu\nu}$$

$$\Delta a_\mu^{(1)} \propto -\frac{c_{\mu\mu}^2}{16\pi^2}$$

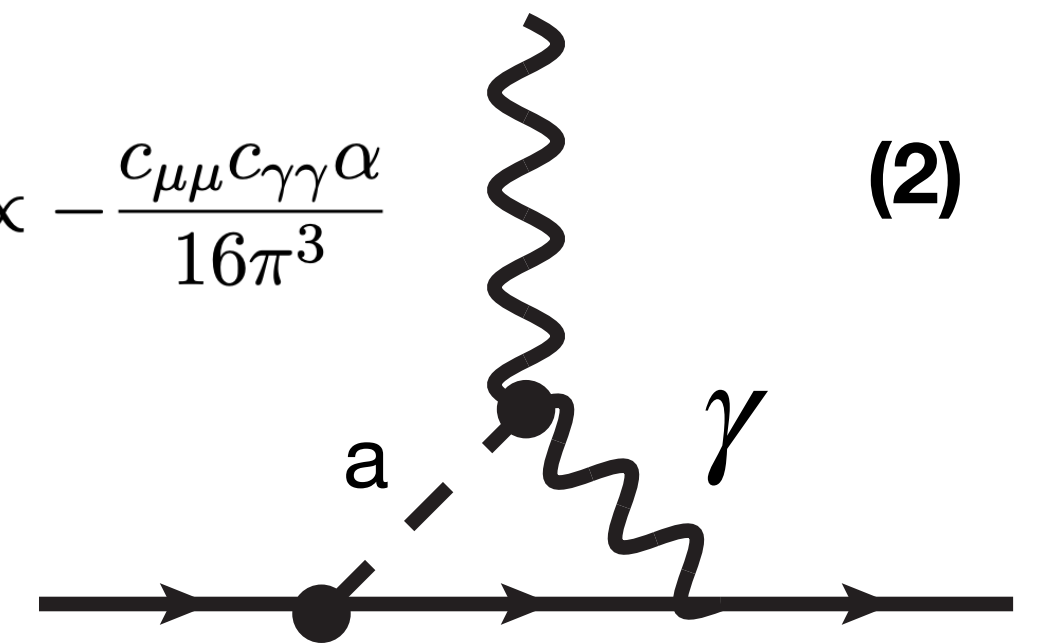


- The 3rd diagram subtlety:

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- Previously, calculate Barr-Zee diagram for on-shell ALP and photon and use the result of 2nd diagram

$$\Delta a_\mu^{(2)} \propto -\frac{c_{\mu\mu} c_{\gamma\gamma} \alpha}{16\pi^3}$$



- Recently, calculate off-shell ALP and photon

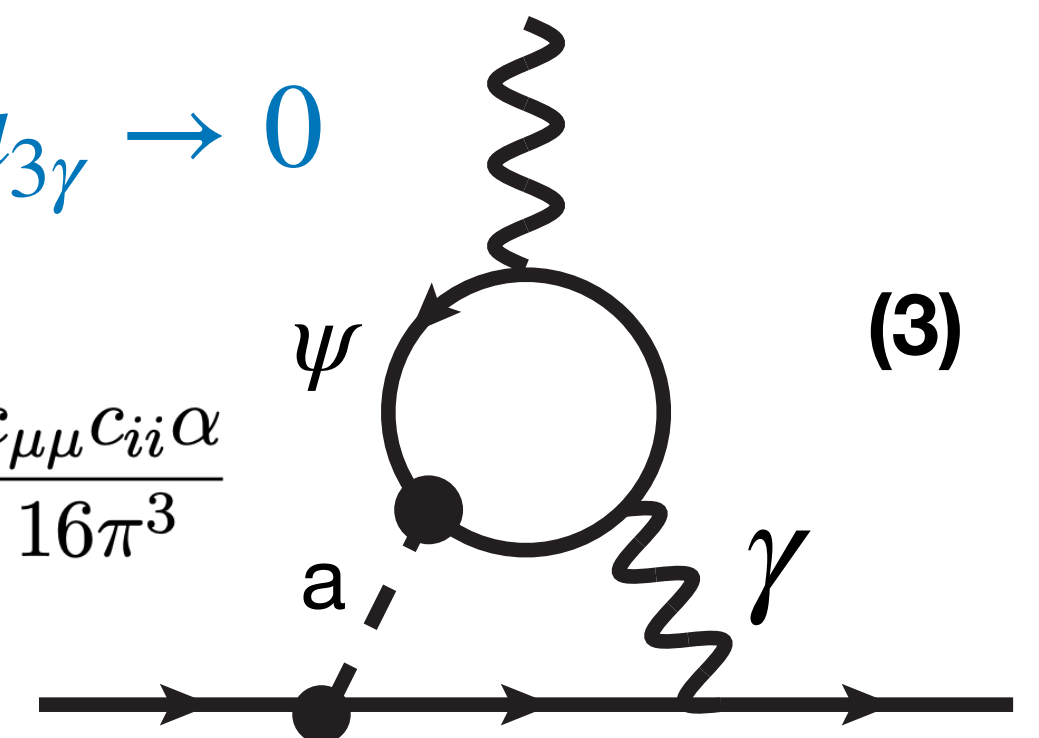
Buen-Abad et al, 2104.03267 (JHEP)

$$\Delta a_{2\gamma} = -\frac{m_\mu^2 c_{\mu\mu} c_{\gamma\gamma} \alpha}{8\pi^3 f_a^2} \cdot h_\gamma(x, \mu) \quad x \equiv \frac{m_a^2}{m_\mu^2}$$

$$m_\psi \rightarrow \infty, \Delta a_{3\gamma} \rightarrow 0$$

$$\Delta a_{3\gamma} = -\frac{c_{\mu\mu}^2 \alpha m_\mu^2}{8\pi^3 f_a^2} \left[H_\gamma(\text{two loop}) + h_\gamma(\text{counter term}) \right]$$

$$\Delta a_\mu^{(3)} \propto -\frac{c_{\mu\mu} c_{ii} \alpha}{16\pi^3}$$



- Different sign for $c_{\mu\mu}$ and $c_{\gamma\gamma}$ is needed

A complete calculation for axion-like particle

- The axion-like particle Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_f \frac{C_{ff}}{2} \frac{\partial^\mu a}{f_a} \bar{f} \gamma_\mu \gamma_5 f + \frac{\alpha C_{\gamma\gamma}}{4\pi} \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\alpha C_{\gamma Z}}{2\pi s_w c_w} \frac{a}{f_a} F_{\mu\nu} \tilde{Z}^{\mu\nu} + \frac{\alpha C_{ZZ}}{4\pi s_w^2 c_w^2} \frac{a}{f_a} Z_{\mu\nu} \tilde{Z}^{\mu\nu}$$

- Our calculation

$$\Delta a_{2Z} = \frac{\alpha c_{\gamma Z} c_{\mu\mu} m_\mu^2 (4s_w^2 - 1)}{32\pi^3 c_w^2 s_w^2 f_a^2} \cdot h_Z(x, y, \mu)$$

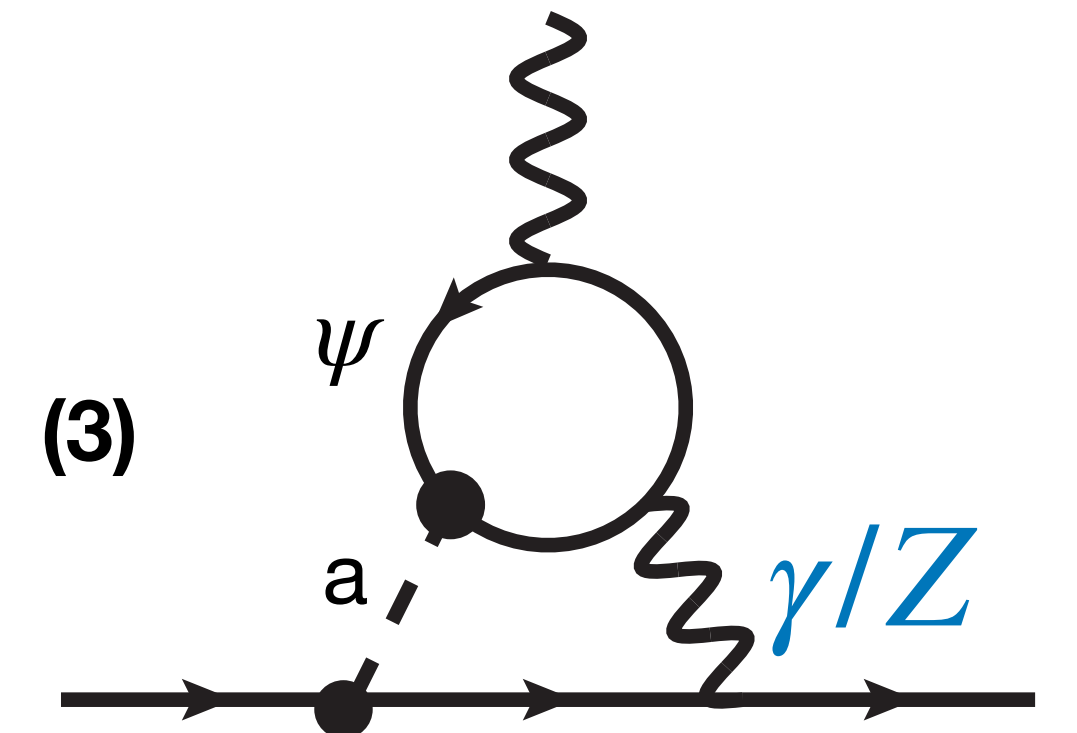
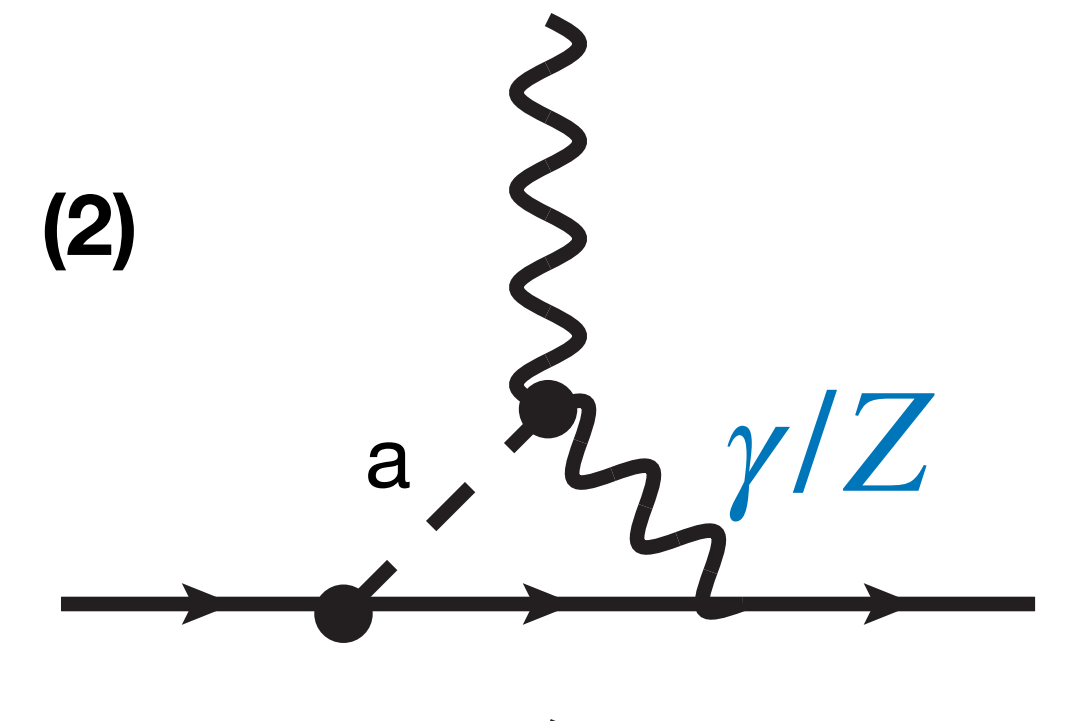
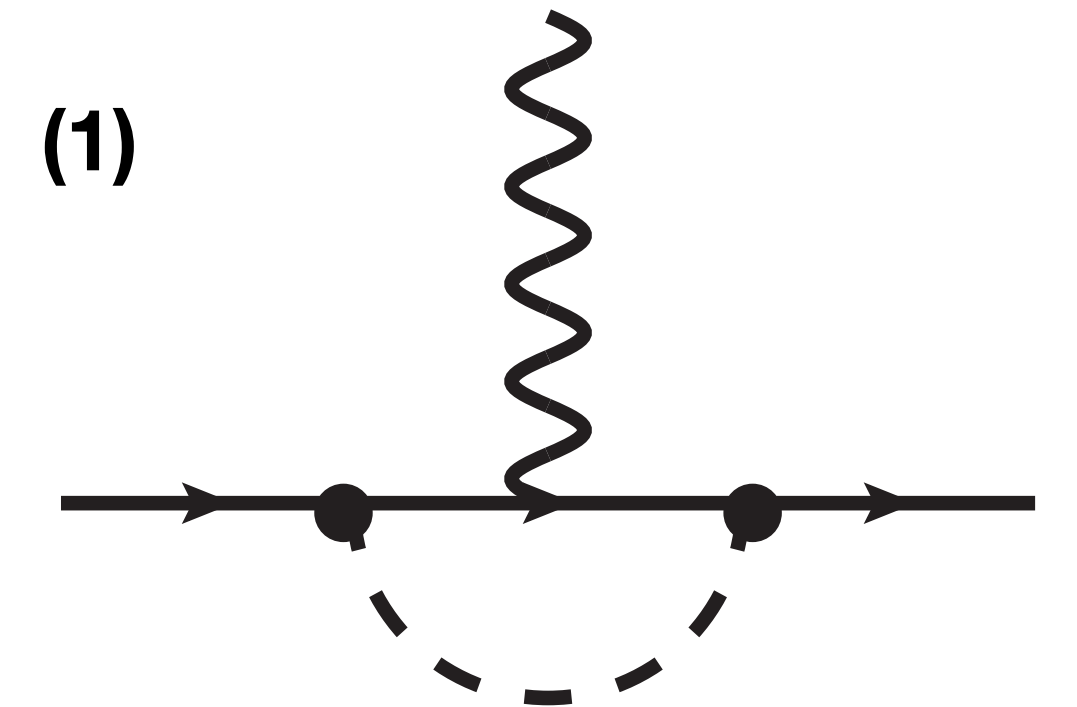
$$m_\psi \rightarrow \infty, \Delta a_3 \rightarrow 0$$

$$\Delta a_{3Z} = -\frac{\alpha c_{\mu\mu}^2 m_\mu^2 (4s_w^2 - 1)^2}{128\pi^3 c_w^2 s_w^2 f_a^2} [H_Z(\text{two loop}) + h_Z(\text{counter term})]$$

$$x \equiv \frac{m_a^2}{m_\mu^2}, \quad y \equiv \frac{m_Z^2}{m_\mu^2}$$

- Calculation in another basis

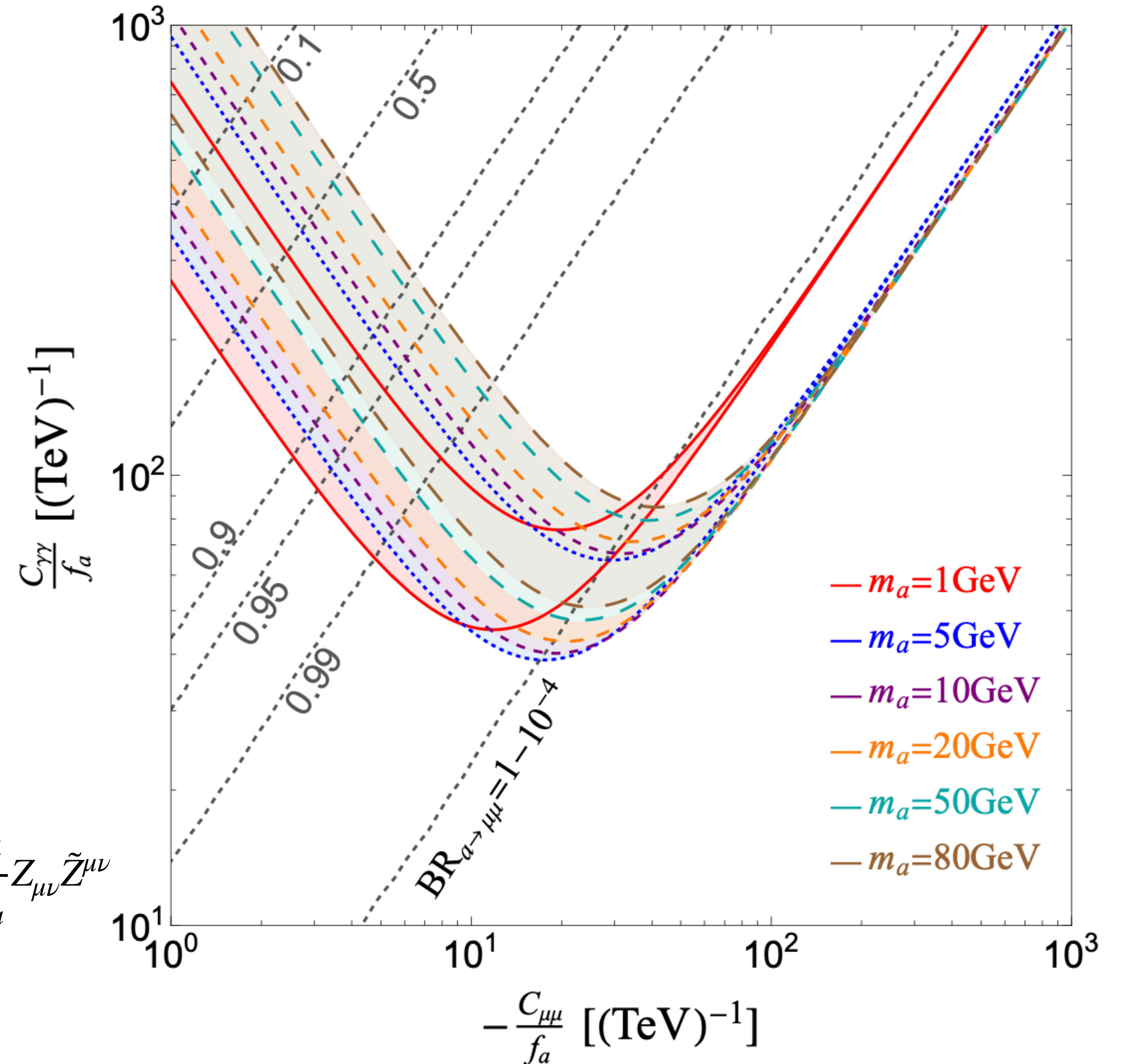
$$\frac{c_{\mu\mu}}{2} \frac{\partial^\mu a}{f_a} \bar{\mu} \gamma_\mu \gamma_5 \mu \rightarrow -\frac{ic_{\mu\mu} m_\mu}{f_a} a \bar{\mu} \gamma_5 \mu + \frac{\alpha c_{\mu\mu}}{4\pi f_a} a F \tilde{F} + \frac{\alpha c_{\mu\mu} (s_w^2 Q - \frac{1}{2} T_3)}{2\pi c_w s_w f_a} a F^{\mu\nu} \tilde{Z}_{\mu\nu}$$



Muon $g-2$ solution and ALP decay BR

- Assume only coupling to Muon and Photon
- In $g-2$ solution region, mostly decay to $a \rightarrow \mu^+ \mu^-$
- The inclusion of Z diagram makes some difference for large m_a
- Exotic Z decay should happen

$$\mathcal{L}_{\text{eff}} = \sum_f \frac{C_{ff}}{2} \frac{\partial^\mu a}{f_a} \bar{f} \gamma_\mu \gamma_5 f + \frac{\alpha C_{\gamma\gamma}}{4\pi} \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\alpha C_{\gamma Z}}{2\pi s_w c_w} \frac{a}{f_a} F_{\mu\nu} \tilde{Z}^{\mu\nu} + \frac{\alpha C_{ZZ}}{4\pi s_w^2 c_w^2} \frac{a}{f_a} Z_{\mu\nu} \tilde{Z}^{\mu\nu}$$



Muon g-2 solution and ALP decay BR

- The BR($a \rightarrow \gamma\gamma, Z \rightarrow \gamma a$) needs to consider 1-loop contribution from fermion loop

$$\Gamma(Z \rightarrow \gamma a) = \frac{\alpha^2 (m_Z) m_Z^3}{96\pi^3 s_w^2 c_w^2 f_a^2} |C_{\gamma Z}^{\text{eff}}|^2 \left(1 - \frac{m_a^2}{m_Z^2}\right)^3,$$

$$\Gamma(a \rightarrow \mu^+ \mu^-) = \frac{m_a m_\mu^2}{8\pi f_a^2} |C_{\mu\mu}|^2 \sqrt{1 - \frac{4m_\mu^2}{m_a^2}},$$

$$\Gamma(a \rightarrow \gamma\gamma) = \frac{\alpha^2 m_a^3}{64\pi^3 f_a^2} |C_{\gamma\gamma}^{\text{eff}}|^2,$$

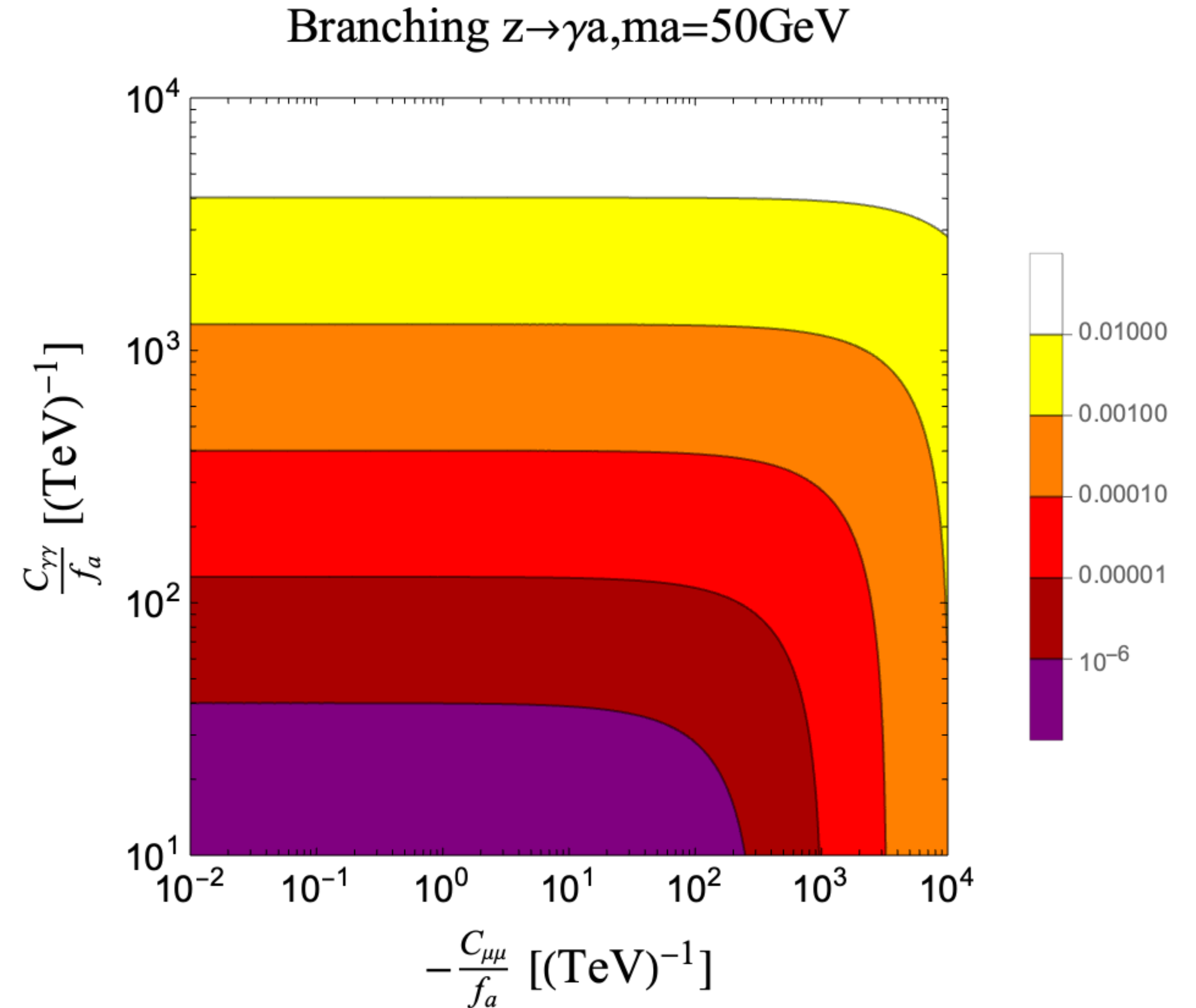
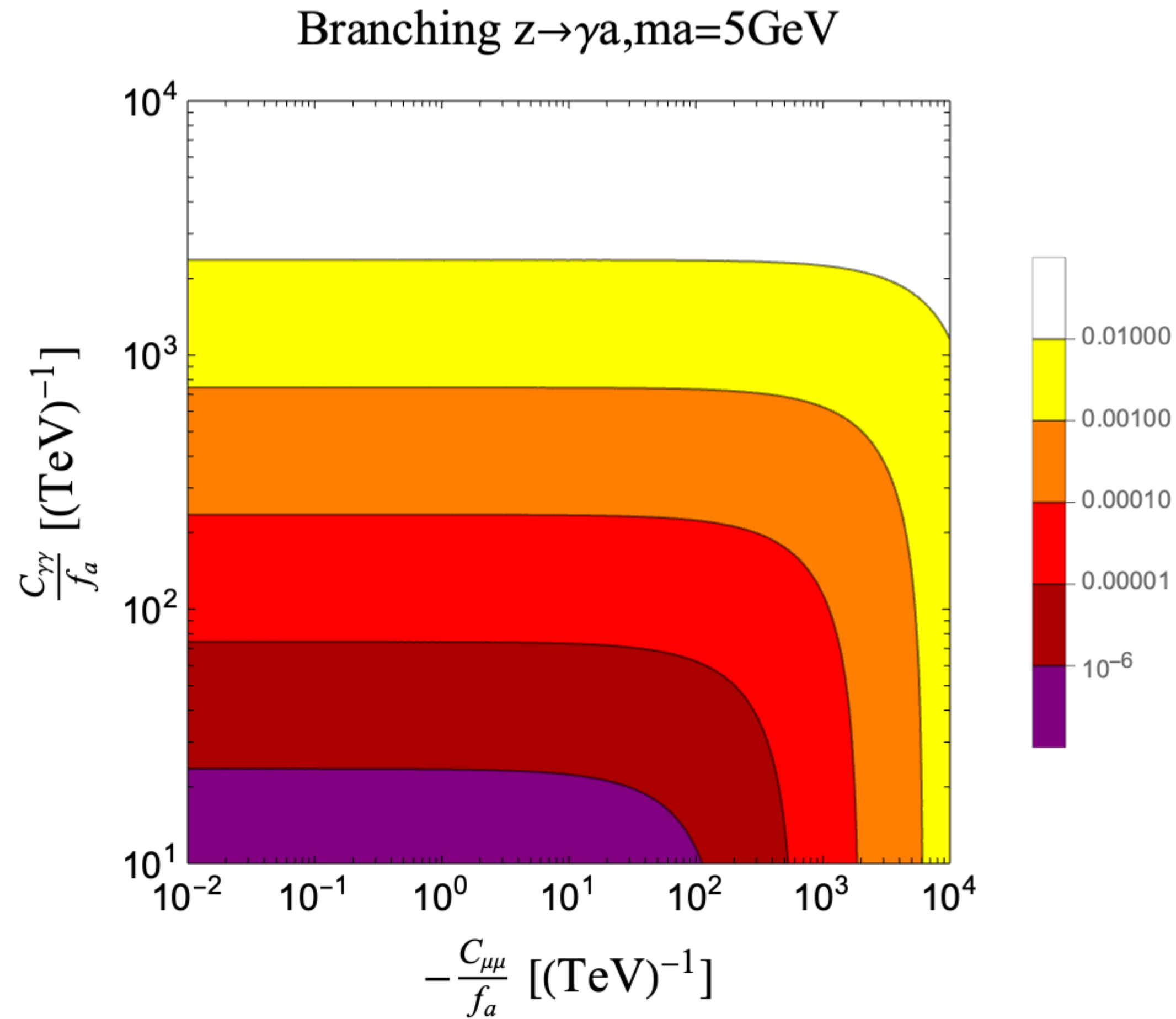
$$C_{\gamma\gamma}^{\text{eff}} = C_{\gamma\gamma} + C_{\mu\mu} \left[1 + \frac{m_\mu^2}{m_a^2} \cdot \mathcal{F}\left(\frac{m_a^2}{m_\mu^2}\right) \right] + \mathcal{O}(\alpha),$$

$$C_{\gamma Z}^{\text{eff}} = C_{\gamma Z} + C_{\mu\mu} \left(\frac{1}{4} - s_w^2 \right) \left[1 + \frac{m_\mu^2}{m_a^2 - m_Z^2} \cdot \left(\mathcal{F}\left(\frac{m_a^2}{m_\mu^2}\right) - \mathcal{F}\left(\frac{m_Z^2}{m_\mu^2}\right) \right) \right] + \mathcal{O}(\alpha)$$

$$\mathcal{F}(x) \equiv \ln^2 \left(\frac{\sqrt{x(x-4)} - x + 2}{2} \right)$$

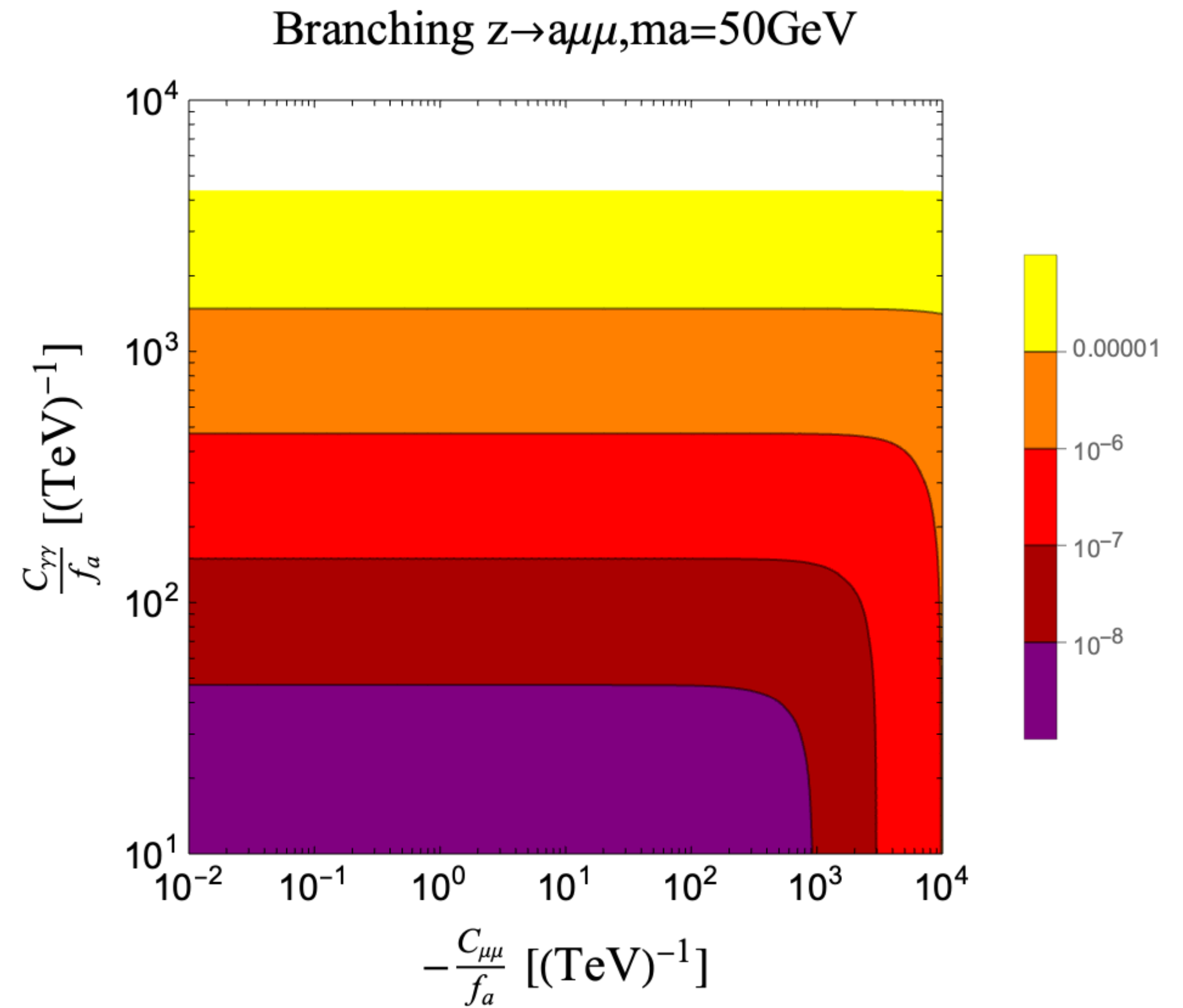
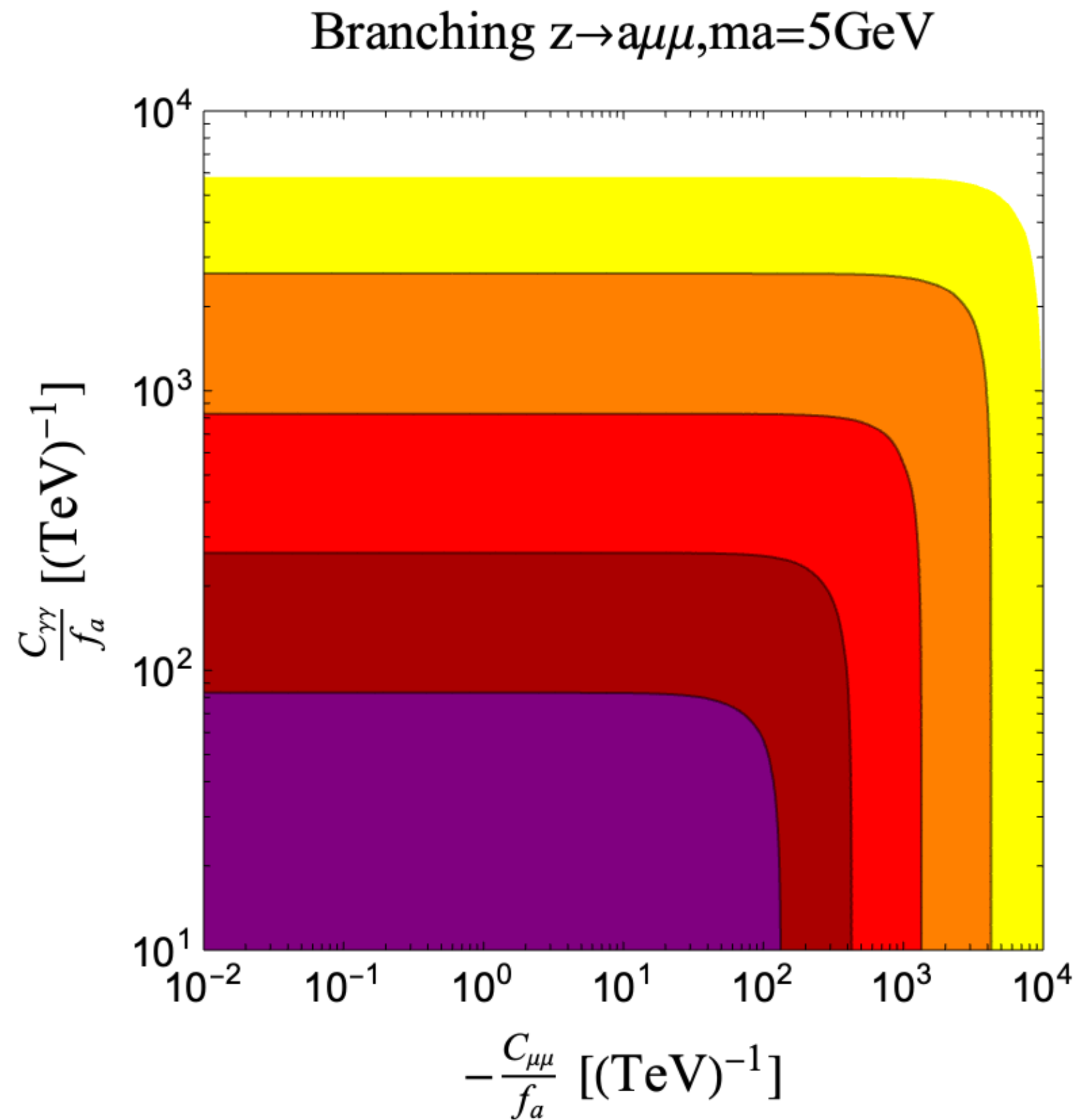
The relevant search at Z-factory (CEPC/FCC-ee)

- The exotic Z decay: $Z \rightarrow a + \gamma$ and $Z \rightarrow a + \mu^+ + \mu^-$



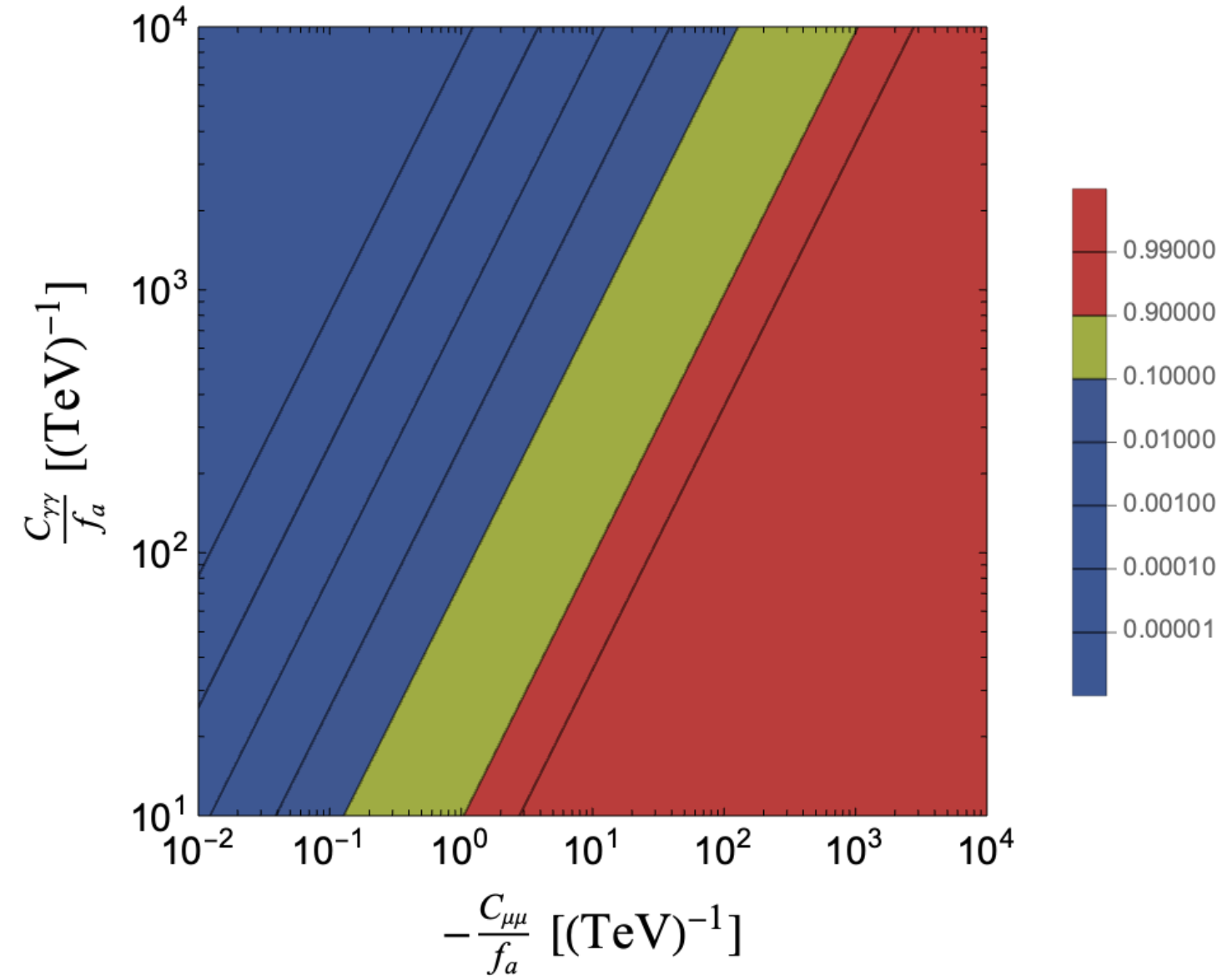
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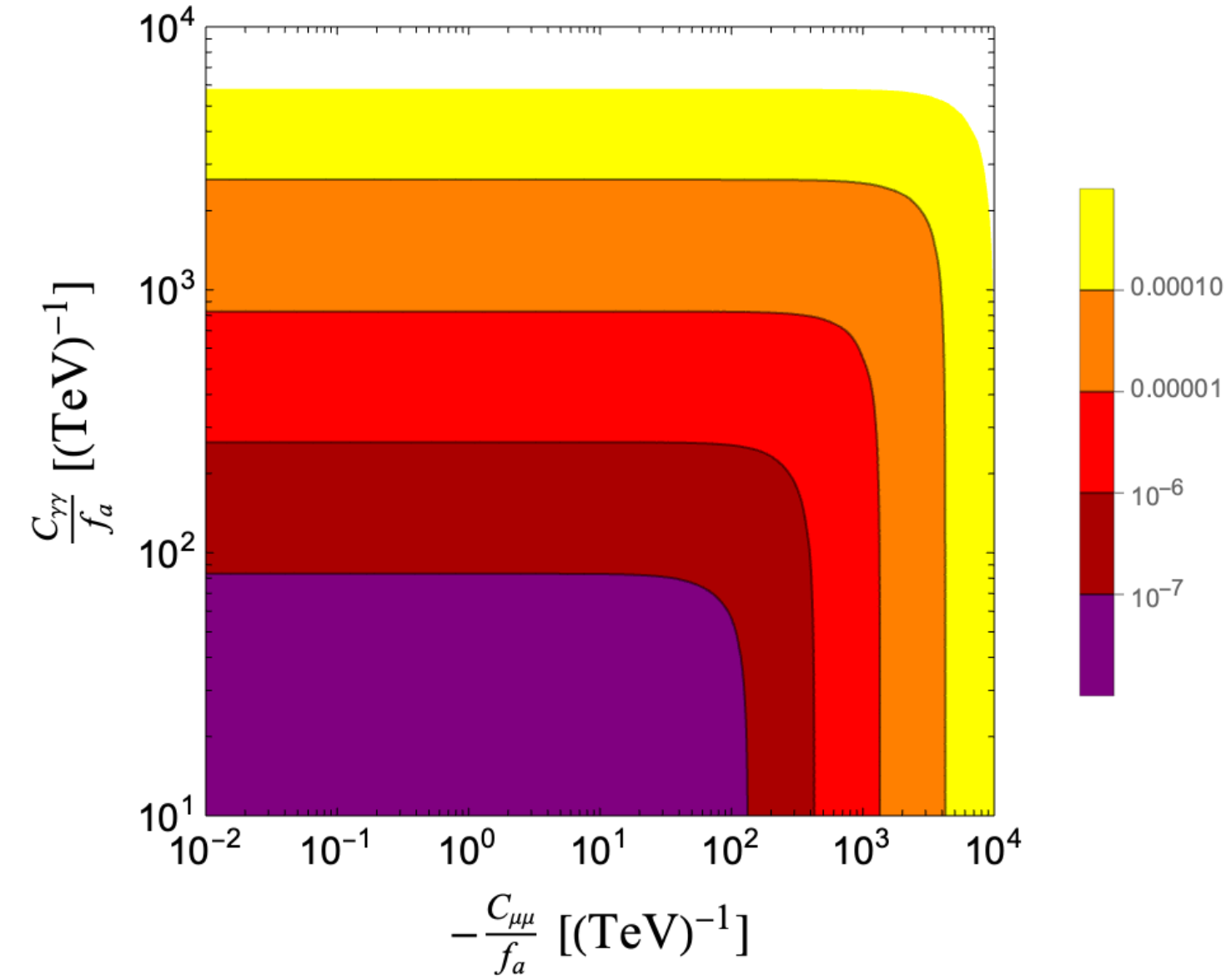


ALP $a \rightarrow \mu\mu$ branching ratio and effects

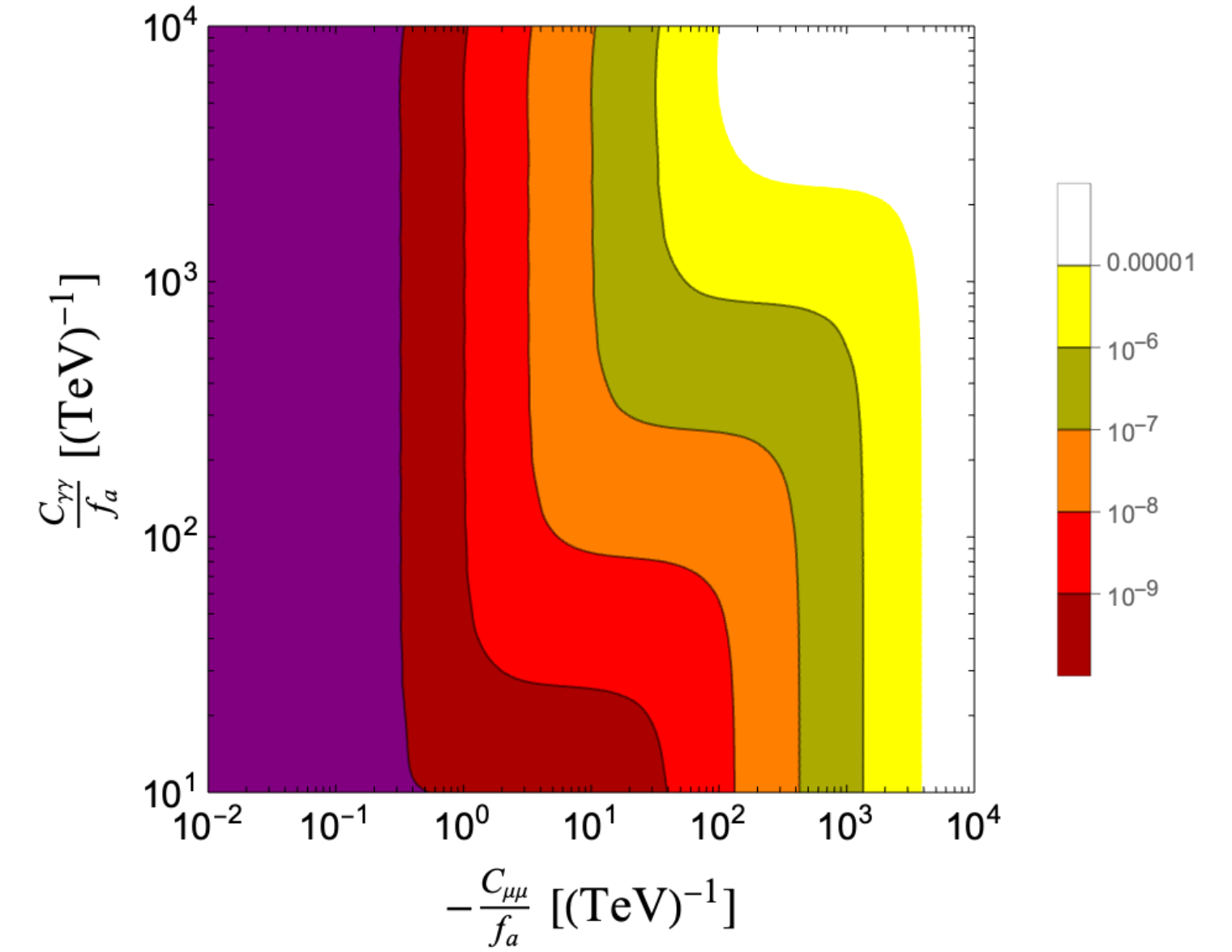
Branching $a \rightarrow \mu\mu, m_a=5\text{GeV}$



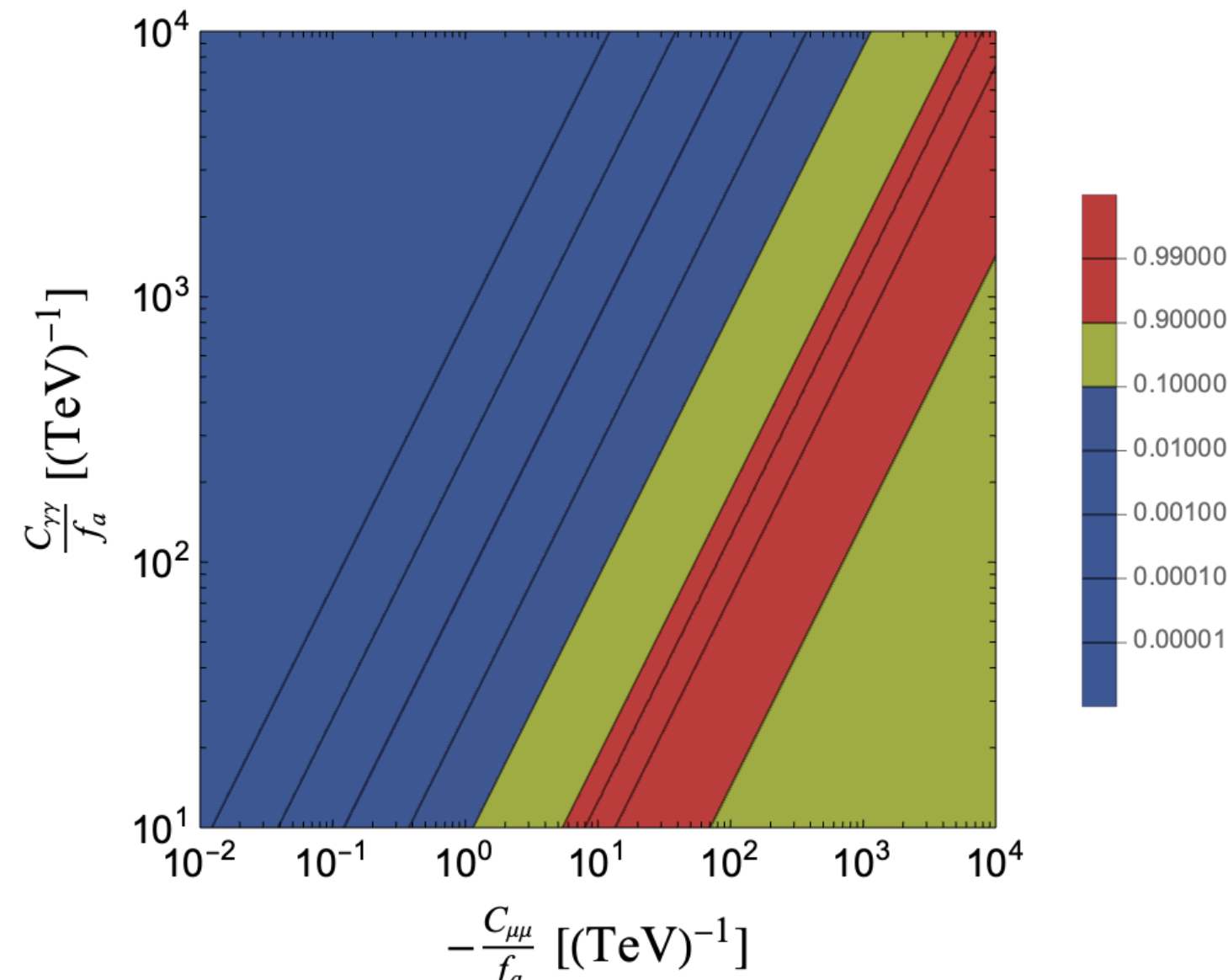
Branching $z \rightarrow a\mu\mu, m_a=5\text{GeV}$



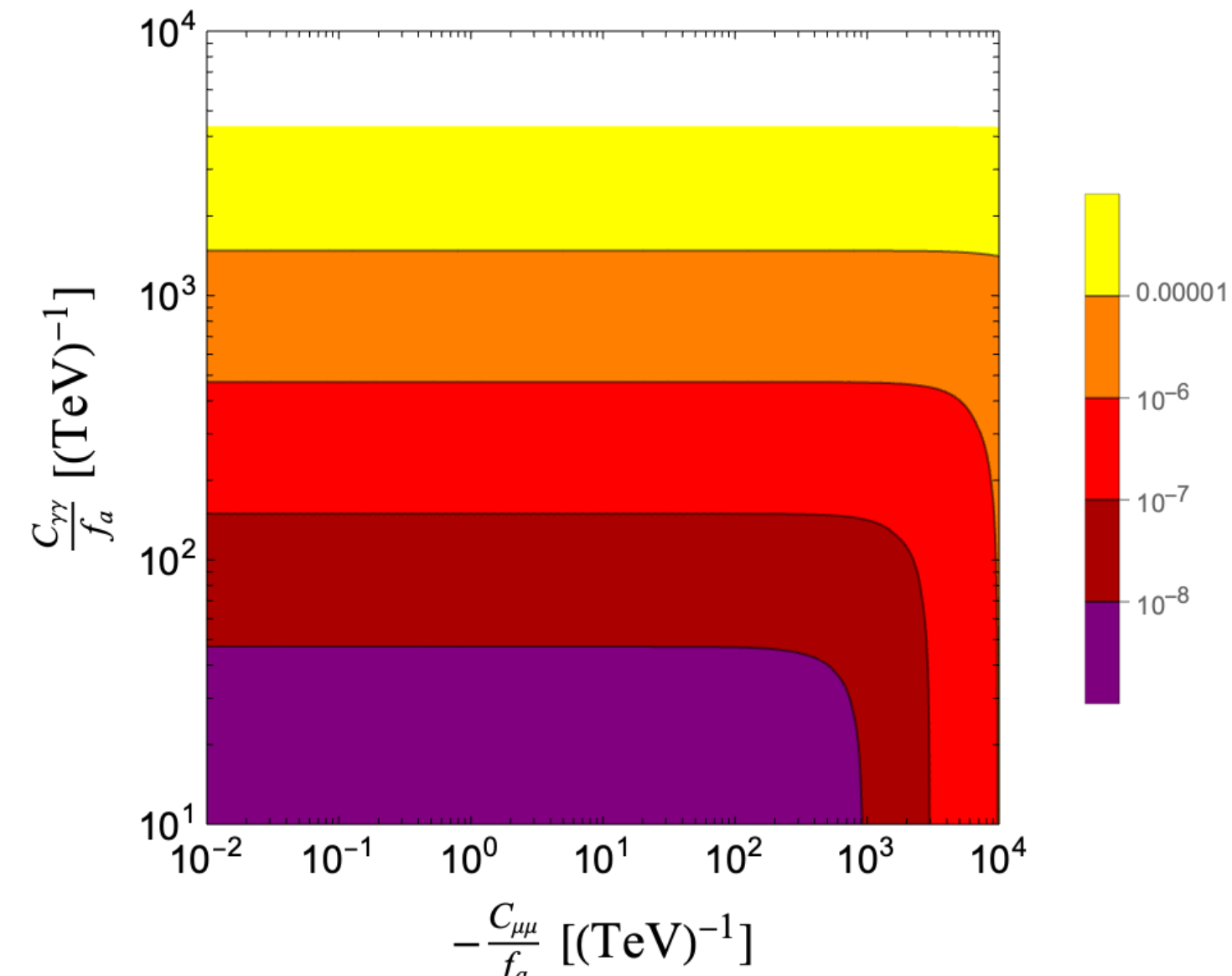
Branching $z \rightarrow 4\mu, m_a=5\text{GeV}$



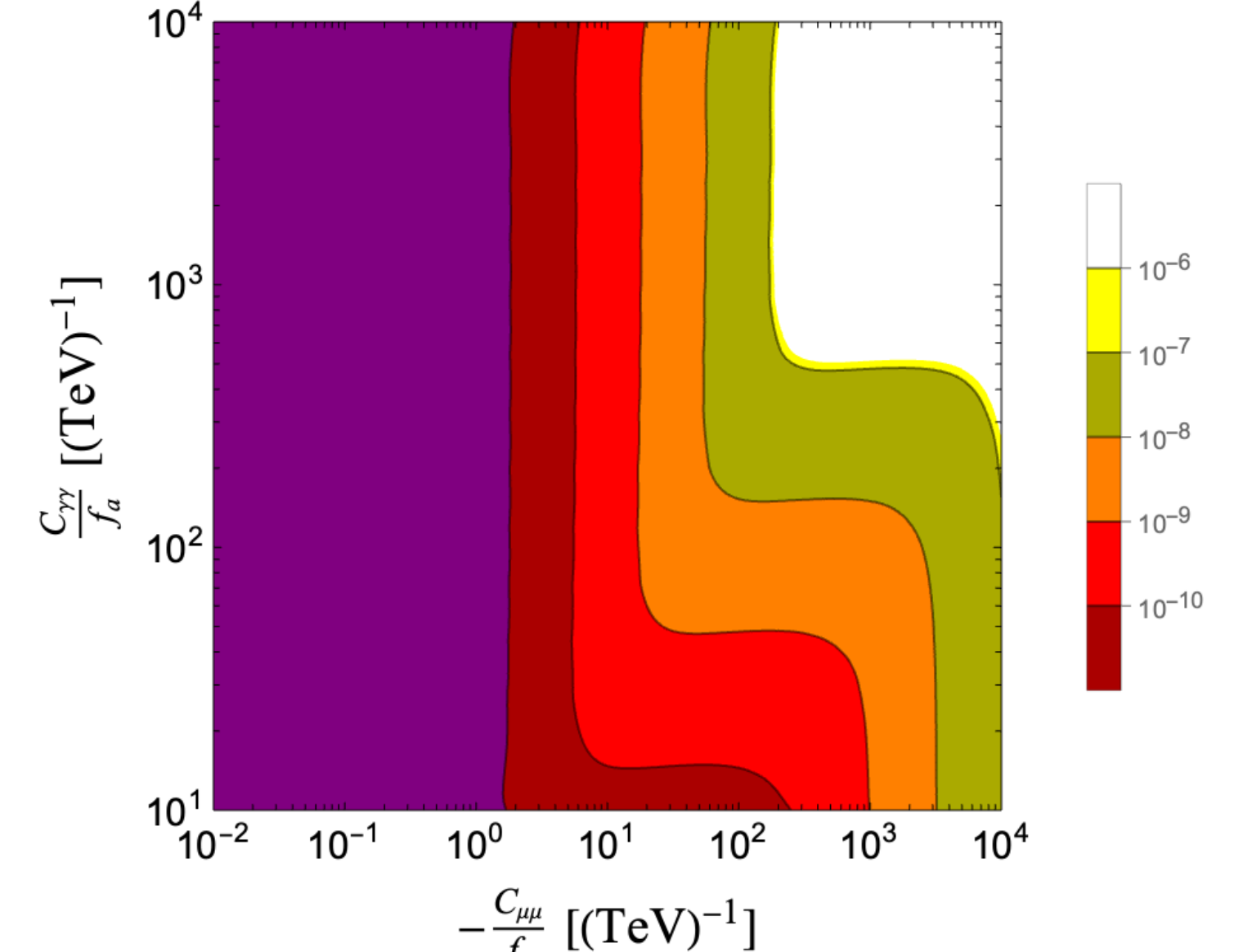
Branching $a \rightarrow \mu\mu, m_a=50\text{GeV}$



Branching $z \rightarrow a\mu\mu, m_a=50\text{GeV}$



Branching $z \rightarrow 4\mu, m_a=50\text{GeV}$



The existing constraints for ALP

- Constraining a - γ coupling only:

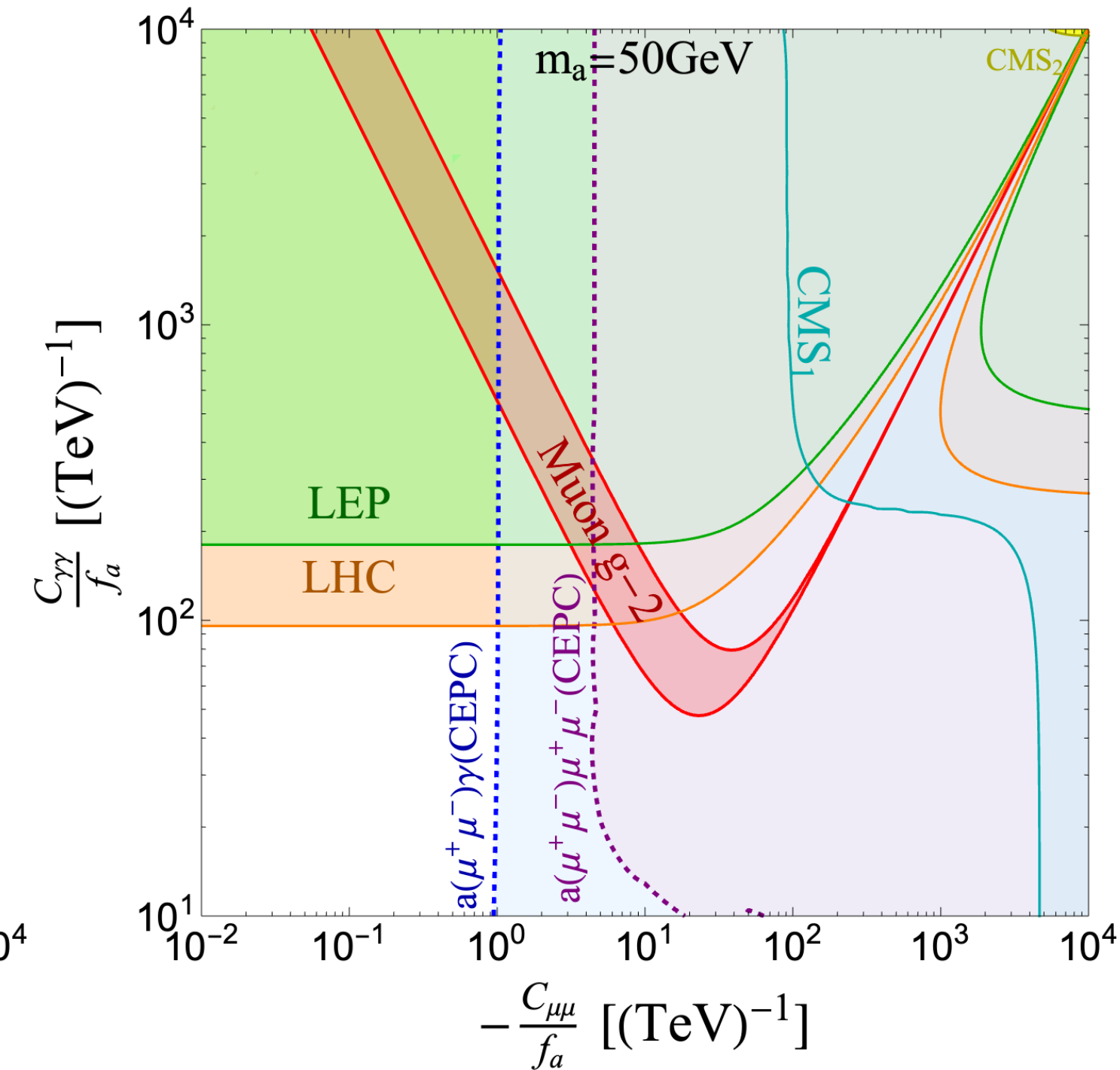
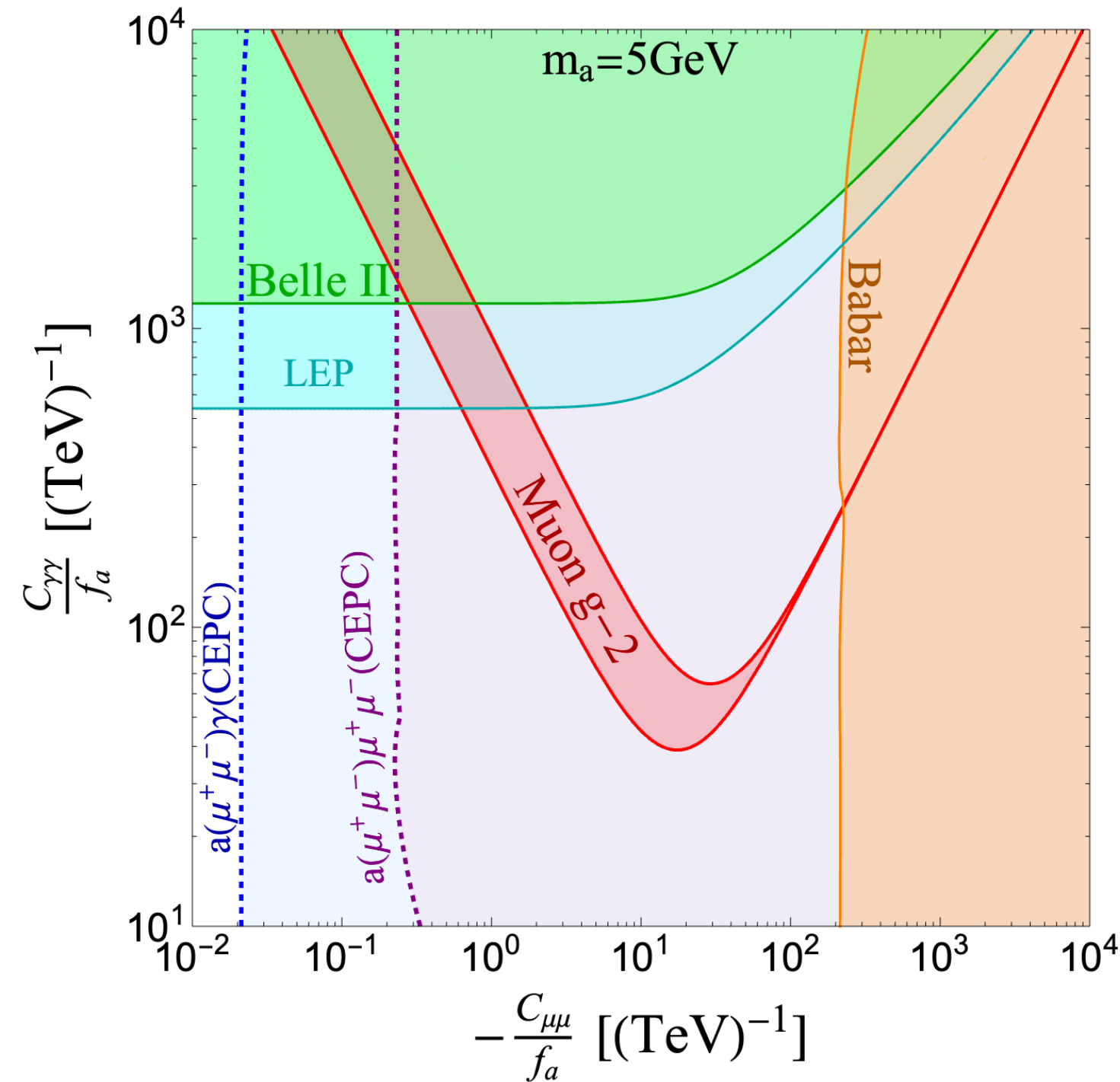
- Belle-II, LEP: $e^+e^- \rightarrow a\gamma \rightarrow (\gamma\gamma)\gamma$
- LHC: $pp \rightarrow a\gamma \rightarrow (\gamma\gamma)\gamma$

- Constraining a - μ coupling only:

- BaBar: recast $e^+e^- \rightarrow \mu^+\mu^-Z'$
- CMS(4μ): $pp \rightarrow \mu^+\mu^-\phi$

- Constraining both coupling

- CMS($\bar{t}t + 2\mu$): $pp \rightarrow \bar{t}t\phi \rightarrow \bar{t}t(\mu^+\mu^-)$



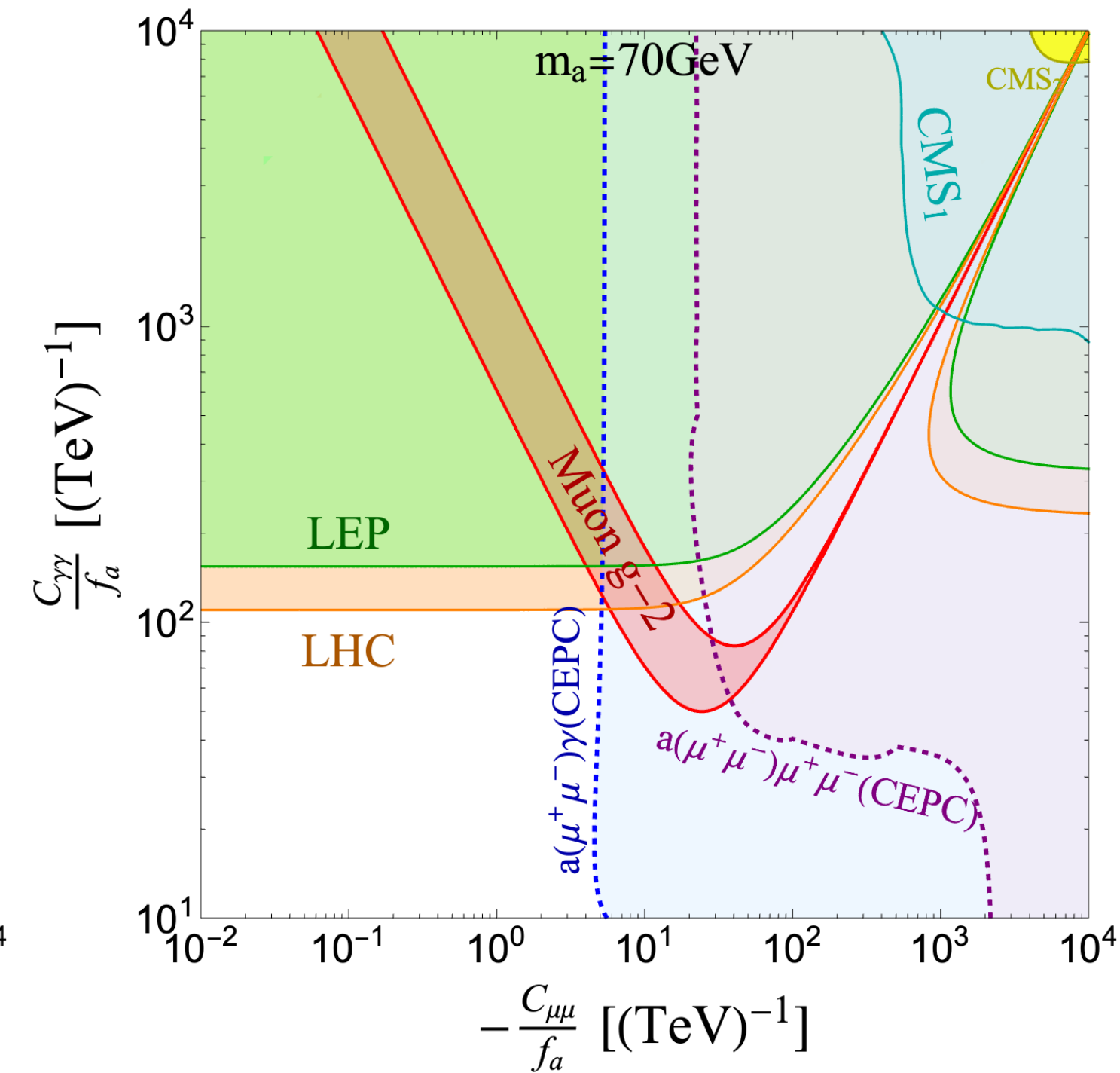
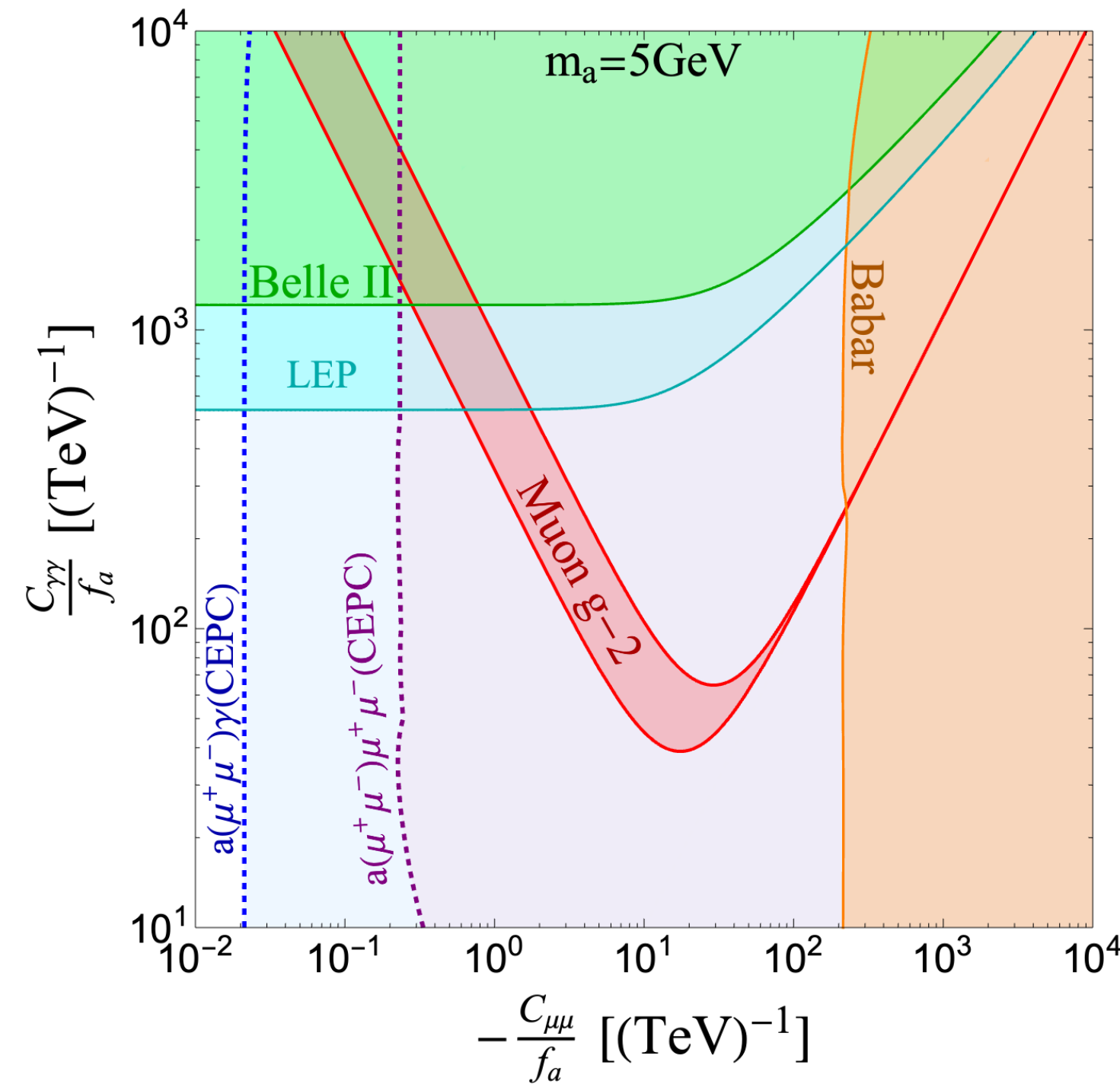
Preliminary results

The relevant search at Z-factory ($C_{WW} = 0$)

- The exotic Z decay: $Z \rightarrow a + \gamma$
and $Z \rightarrow a + \mu^+ + \mu^-$
- With ALP decay: $a \rightarrow \mu^+ \mu^-$
- Relevant SM background: $\mu^+ \mu^- \gamma$
and 4μ
- Cuts: $E_\gamma > 2\text{GeV}$, $p_T^\mu > 5\text{GeV}$,
 $|\eta| < 3$, $\Delta R_{\mu\mu} > 0.1$,

dimuon resolution:

$$\left| \frac{m_{\mu\mu}}{m_a} - 1 \right| < 0.19\%$$

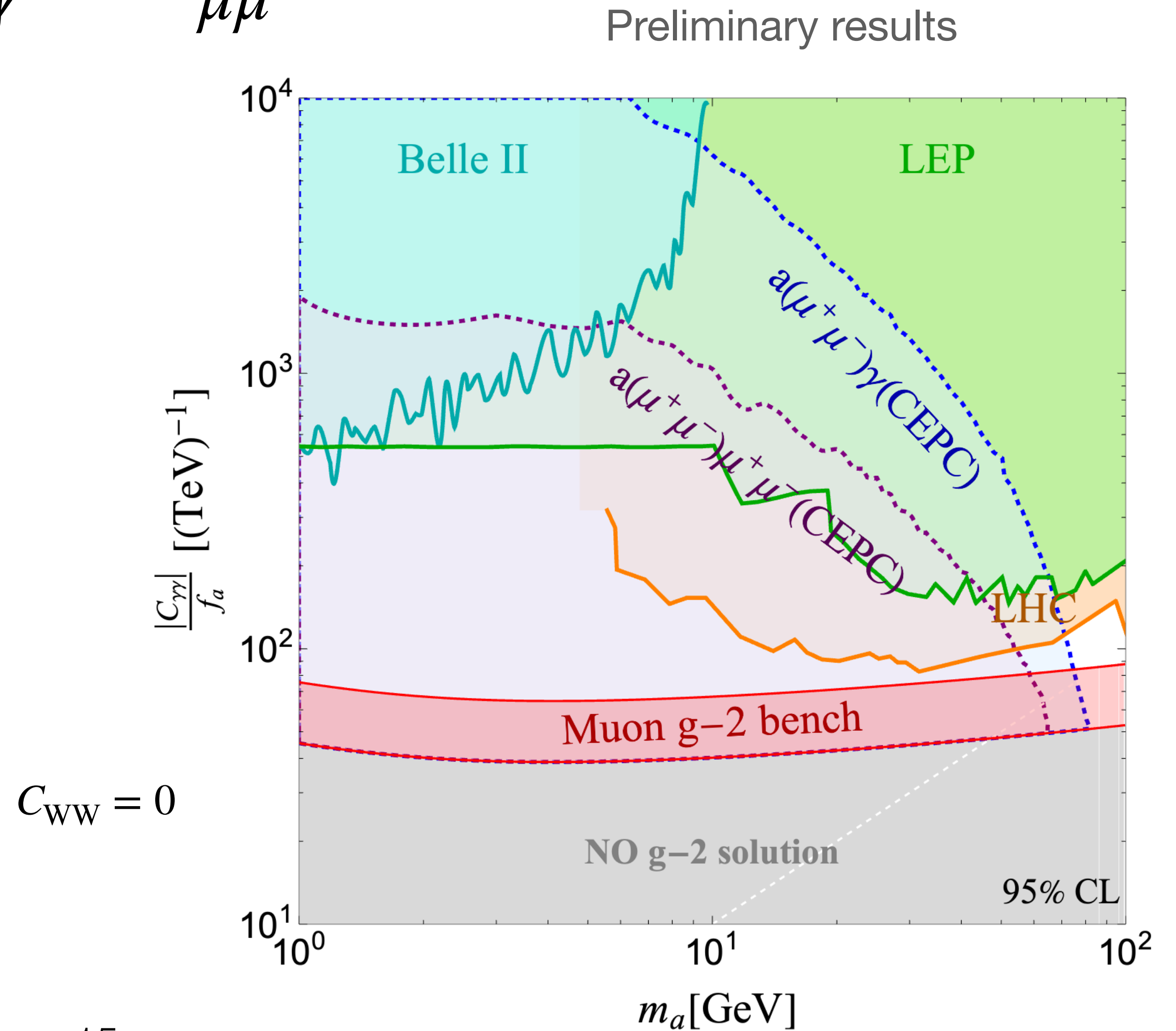
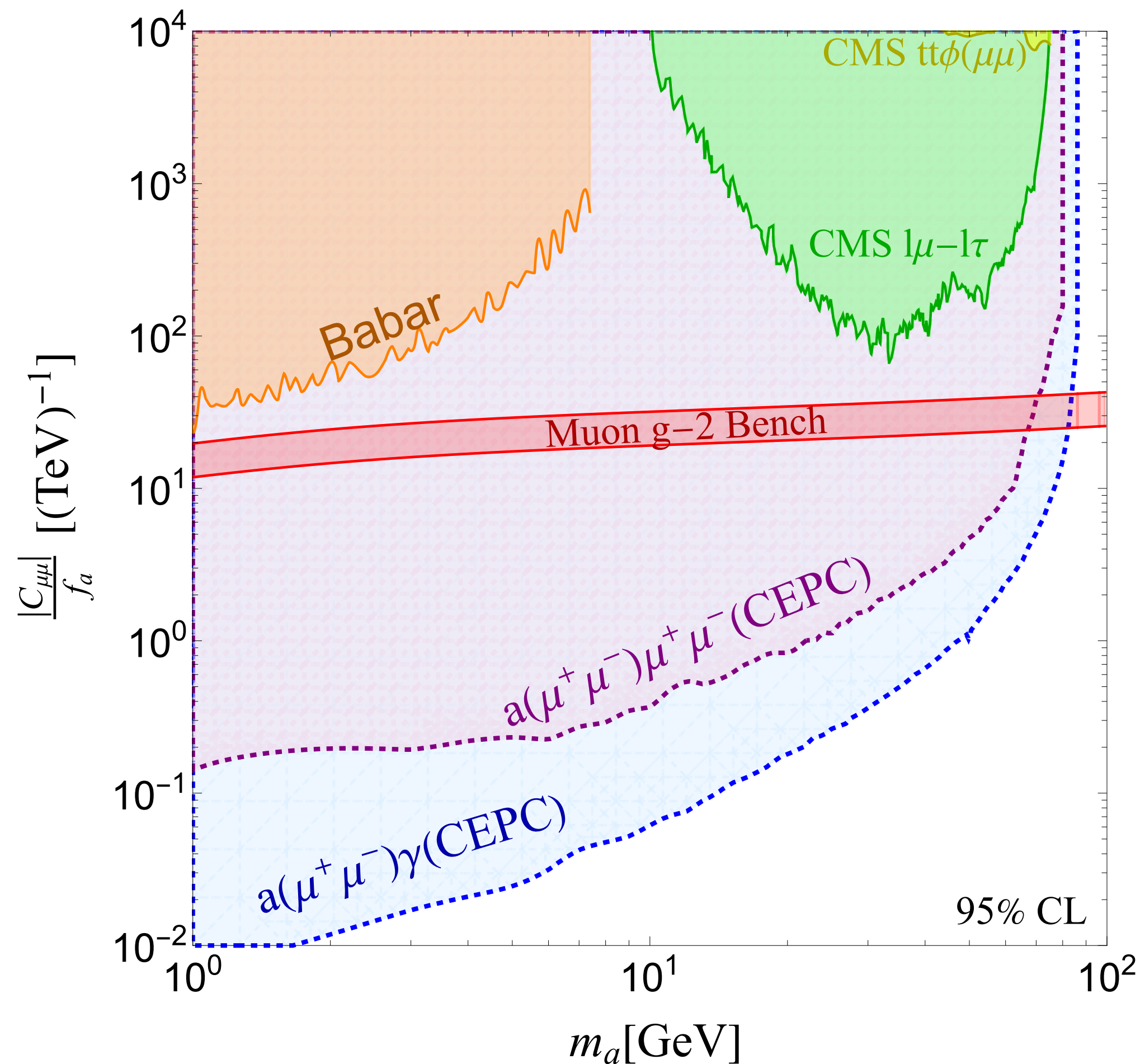


$C_{WW} = 0$

- Large m_a prefers decay $a \rightarrow \gamma\gamma$
- SM bkg: $\mu^+ \mu^- \gamma$ dominated by soft γ

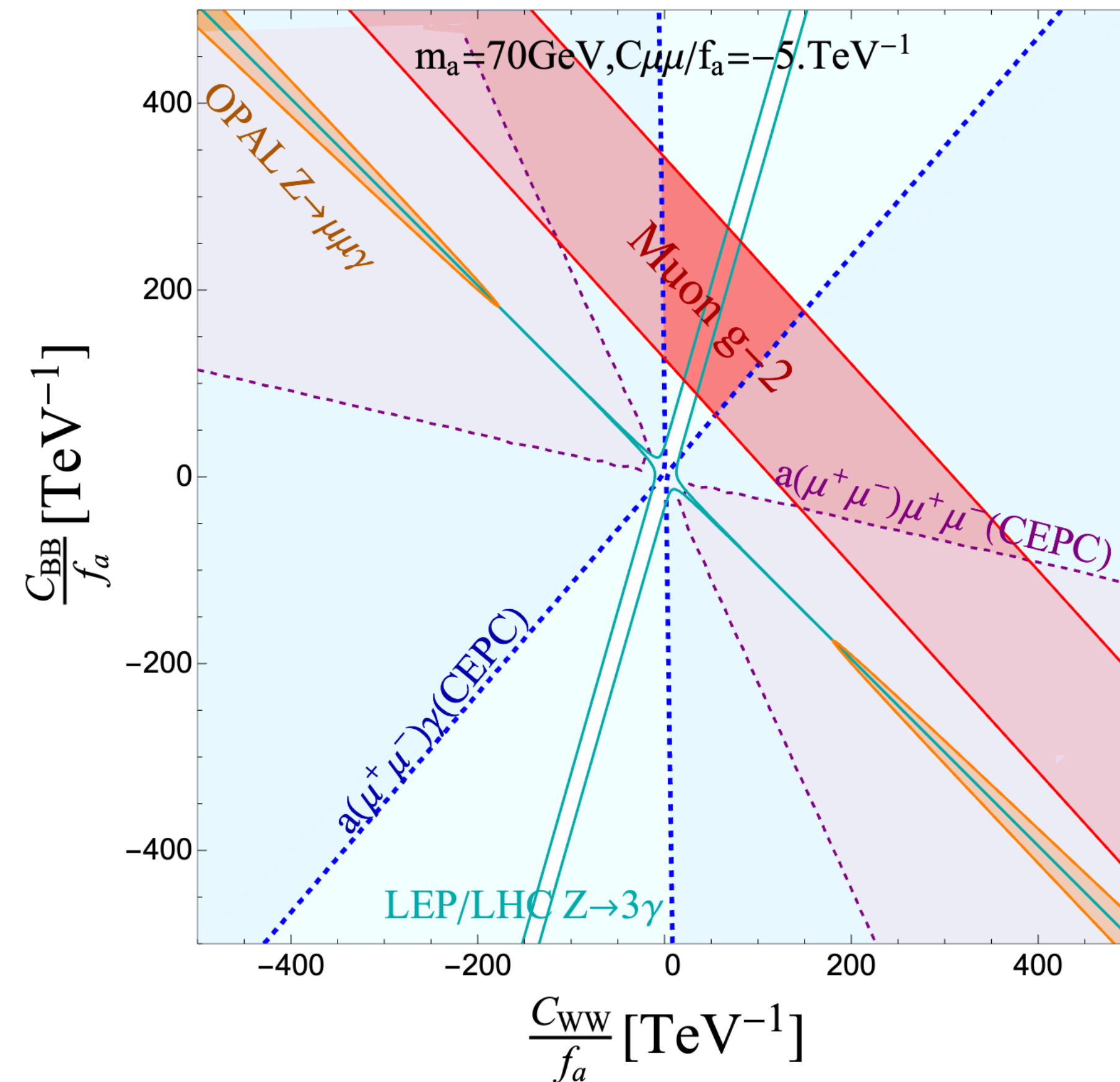
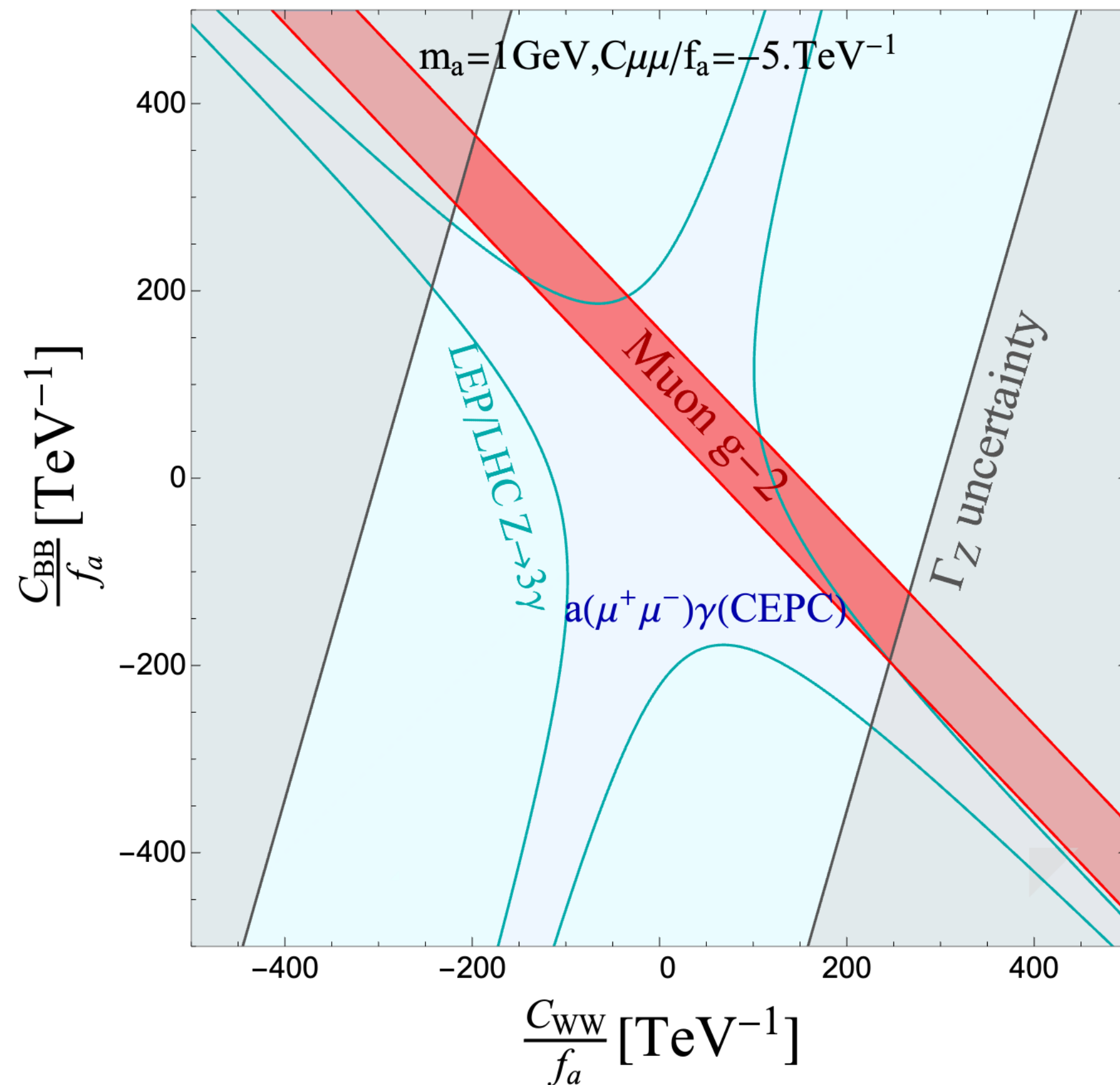
The relevant search at Z-factory ($C_{WW} = 0$)

- Marginalize one coupling $C_{\gamma\gamma}$ or $C_{\mu\mu}$



The relevant search at Z-factory ($C_{WW} \neq 0$)

- 4 parameters: $m_a, C_{\mu\mu}, C_{BB}, C_{WW}$
- Z-factory $a \rightarrow \mu^+ \mu^-$ search is sensitive to light ALP
- Z-factory $a \rightarrow \gamma\gamma$ search will be sensitive to heavy ALP (in progress) Preliminary results



Summary

- Muon g-2 experiments show 4.2σ discrepancy with SM
- ALP can provide a solution with couplings $C_{\mu\mu}$ and $C_{\gamma\gamma}$
- In UV model, $C_{\gamma\gamma}$ comes from C_{WW} and C_{BB} , leads to $C_{\gamma Z}$
- At Z-factory, it leads to exotic Z decay:
 - $Z \rightarrow a\gamma, a\mu^+\mu^-$
- Future Z factory can provide sensitivity covers most of the g-2 region for $m_a \lesssim m_Z$

Thank you!

Backup slides

$$x \equiv \frac{m_a^2}{m_\mu^2}, \quad y \equiv \frac{m_Z^2}{m_\mu^2}, \quad \Delta^2 \equiv \frac{m_\mu^2}{z(1-z)} \quad (15)$$

$$B(m_\mu^2, m_{a/Z}, m_\mu) = \sqrt{m_{a/Z}^2(m_{a/Z}^2 - 4m_\mu^2)} \ln \left(\frac{\sqrt{m_{a/Z}^2 - 4m_\mu^2} + m_{a/Z}}{2m_\mu} \right) / m_\mu^2 \quad (16)$$

- Special functions

$$H_\gamma = \int_0^1 dz \frac{\Delta^2}{2(\Delta^2 - m_a^2)} \left[-\frac{2m_a^2 + 2m_\mu^2}{3m_\mu^2} B(m_\mu^2, m_a, m_\mu) + \frac{2\Delta^2 + 2m_\mu^2}{3m_\mu^2} B(m_\mu^2, m_a, \Delta) \right. \\ \left. + \frac{\Delta^4}{3m_\mu^4} \ln \left(\frac{m_\mu^2}{\Delta^2} \right) + \frac{m_a^4}{3m_\mu^4} \ln \left(\frac{m_a^2}{m_\mu^2} \right) + \frac{2\Delta^2 - m_a^2}{3m_\mu^2} \right] \quad (17)$$

$$H_Z = \int_0^1 dz \frac{\Delta^2}{\Delta^2 - m_a^2} \left[-\frac{ma^2(m_a^2 + 2m_\mu^2)B(m_\mu^2, m_a, m_\mu)}{3(m_a^2 - m_Z^2)m_\mu^2} + \frac{m_a^6}{6m_\mu^4(m_a^2 - m_Z^2)} \ln \left(\frac{m_a^2}{m_\mu^2} \right) \right] \\ + \frac{\Delta^2}{\Delta^2 - m_Z^2} \left[\frac{m_Z^2(2m_\mu^2 + m_Z^2)B(m_\mu^2, m_Z, m_\mu)}{3m_\mu^2(m_a^2 - m_Z^2)} + \frac{m_Z^6}{6m_\mu^4(m_Z^2 - m_a^2)} \ln \left(\frac{m_Z^2}{m_\mu^2} \right) \right] \\ + \frac{\Delta^2}{3m_\mu^2} + \frac{\Delta^4(\Delta^2 + 2m_\mu^2)B(m_\mu^2, m_\mu, \Delta)}{3m_\mu^2(\Delta^2 - m_a^2)(\Delta^2 - m_Z^2)} + \frac{\Delta^8}{6m_\mu^4(\Delta^2 - m_a^2)(\Delta^2 - m_Z^2)} \ln \left(\frac{m_\mu^2}{\Delta^2} \right) \quad (18)$$

$$h_\gamma = \ln \left(\frac{\mu^2}{m_\mu^2} \right) + 2 - \frac{x^2}{6} \ln(x) + \frac{x}{3} + \frac{x+2}{3} \sqrt{x(x-4)} \ln \left(\frac{\sqrt{x-4} + \sqrt{x}}{2} \right) \quad (19)$$

$$h_Z = \ln \left(\frac{\mu^2}{m_Z^2} \right) + \frac{x(x+2)B(m_\mu^2, m_a, m_\mu)}{3(x-y)} + \frac{y(2+y)B(m_\mu^2, m_Z, m_\mu)}{3(y-x)} \\ + \frac{(x+3+y)}{3} - \frac{6x+y^3-6y}{6(x-y)} \ln \left(\frac{x}{y} \right) - \frac{x^2+xy-6+y^2}{6} \ln(x) \quad (20)$$