



遼寧師範大學

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# Searching for axionlike particles at future electron-positron colliders

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based on

H.-Y. Zhang, C.-X. Yue, Y.-C. Guo and **SY**, PRD 104 (2021) 9, 096008

C.-X. Yue, **SY**, H. Wang and N. Zhang, arXiv:2204.04702 accepted by PRD

Joint Workshop of the CEPC Physics, Software and New Detector  
Concept in 2022, May 23-25, 2022

# Outline

1. Introduction to axion-like particles (ALPs).
2. ALPs searches at colliders
3. Searching for ALPs at future  $e^+e^-$  colliders  
via light-by-light scattering
4. Searching for ALPs via decay  $Z \rightarrow a f \bar{f}$  at Z factories
5. Summary

# Introduction to axion-like particles

- The discovery of Higgs bring us into new territory of spin-0 particles.
- Axion have been postulated to address the strong CP problem, which is the pNGB associated to Peccei-Quinn symmetry, a global U(1).  $m_a \sim m_\pi f_\pi / f_a$
- Many extensions of the SM feature one or several spontaneously broken global U(1) symmetries, thus predicting axion-like particles (ALPs).
- ALPs: No direct relation between coupling and mass.

# General effective Lagrangian of ALPs

*bottom-up view*

SMEFT  $\mathcal{L} = \mathcal{L}_0 + \sum_i \frac{c_i}{\Lambda^{d-4}} \mathbf{O}_i$

Building Blocks:

SM fields:  $B_{\mu\nu}, W_{\mu\nu}, G_{\mu\nu}$

EW scalar doublet:  $\Phi(x) = \frac{v + h(x)}{\sqrt{2}} e^{i\vec{\pi}\vec{\sigma}/v}$

New pseudoscalar (GB):  $\frac{\partial_\mu a}{f_a}$

# General effective Lagrangian of ALPs

## Linear Effective Lagrangian

### NLO bosonic operators

$$\mathcal{O}_{\tilde{B}} = -B_{\mu\nu} \tilde{B}^{\mu\nu} \frac{a}{f_a}$$

$$\mathcal{O}_{\tilde{G}} = -G_{\mu\nu}^a \tilde{G}^{a\mu\nu} \frac{a}{f_a}$$

$$\mathcal{O}_{\tilde{W}} = -W_{\mu\nu} \tilde{W}^{\mu\nu} \frac{a}{f_a}$$

$$\mathcal{O}_{a\Phi} = i(\Phi^\dagger \overleftrightarrow{D}_\mu \Phi) \frac{\partial^\mu a}{f_a}$$

$$\mathcal{L}_{\text{eff}}^{D \leq 5} = \frac{1}{2} (\partial_\mu a) (\partial^\mu a) - \frac{M_a^2}{2} a^2 + \frac{\partial^\mu a}{\Lambda} \sum_F \bar{\psi}_F C_F \gamma_\mu \psi_F$$

$$+ g_s^2 C_{GG} \frac{a}{\Lambda} G_{\mu\nu}^A \tilde{G}^{\mu\nu,A} + g^2 C_{WW} \frac{a}{\Lambda} W_{\mu\nu}^A \tilde{W}^{\mu\nu,A} + g'^2 C_{BB} \frac{a}{\Lambda} B_{\mu\nu} \tilde{B}^{\mu\nu}$$

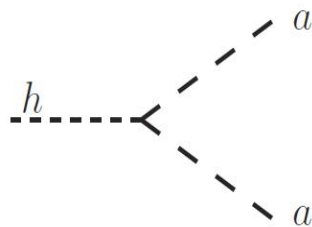
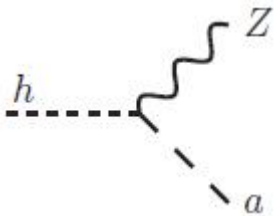
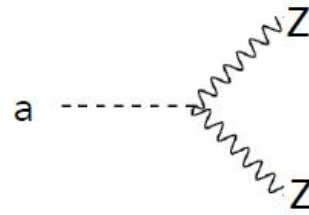
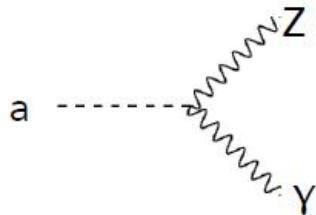
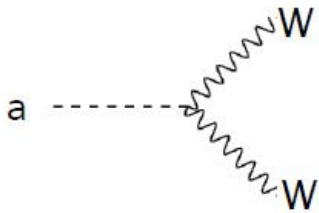
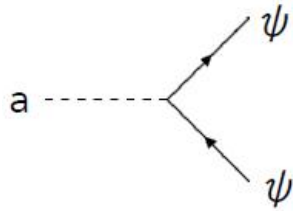
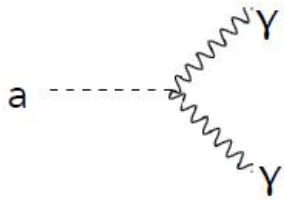
$$\mathcal{L}_{\text{eff}}^{D \geq 6} = \frac{c_{ah}}{f^2} (\partial_\mu a) (\partial^\mu a) \phi^\dagger \phi + \frac{c_{Zh}}{f^3} (\partial^\mu a) \left( \phi^\dagger i D_\mu \phi + \text{h.c.} \right) \phi^\dagger \phi + \dots$$

*H.Georgi, D.B.Kaplan & L.Randall, PLB169(1986)73-78*

*M.Bauer et al., JHEP12(2017),044*

*I.Brivio et al., EPJC77(2017),8,572 (including Nolinear Effective Lagrangian)*

# Vertices



# ALPs searches at colliders

## Production modes

- Resonant production

$$gg \rightarrow a \quad \gamma\gamma \rightarrow a$$

$$e^+e^- \rightarrow a \quad \text{strongly suppressed}$$

- Associated production

$$pp \rightarrow a W^\pm \quad pp \rightarrow a Z(\gamma)$$

$$pp \rightarrow a h \quad pp \rightarrow t\bar{t}a \quad pp \rightarrow a W^\pm \gamma$$

$$e^+e^- \rightarrow a Z(\gamma) \quad e^+e^- \rightarrow a h$$

$$e^+e^- \rightarrow e^+e^-a \quad e^+e^- \rightarrow \nu\bar{\nu}a$$

- Exotic SM decays

$$h \rightarrow Za \quad h \rightarrow aa \quad Z \rightarrow a\gamma$$

- Other modes

## Decay channels

- Stable ALPs  $\sim E$

- Long-Lived ALPs with a Displaced-Vertex

- Short lived ALP

$$a \rightarrow \gamma\gamma \quad a \rightarrow \ell^+\ell^- \quad a \rightarrow jj$$

$$a \rightarrow b\bar{b} \quad a \rightarrow VV \quad a \rightarrow t\bar{t}$$

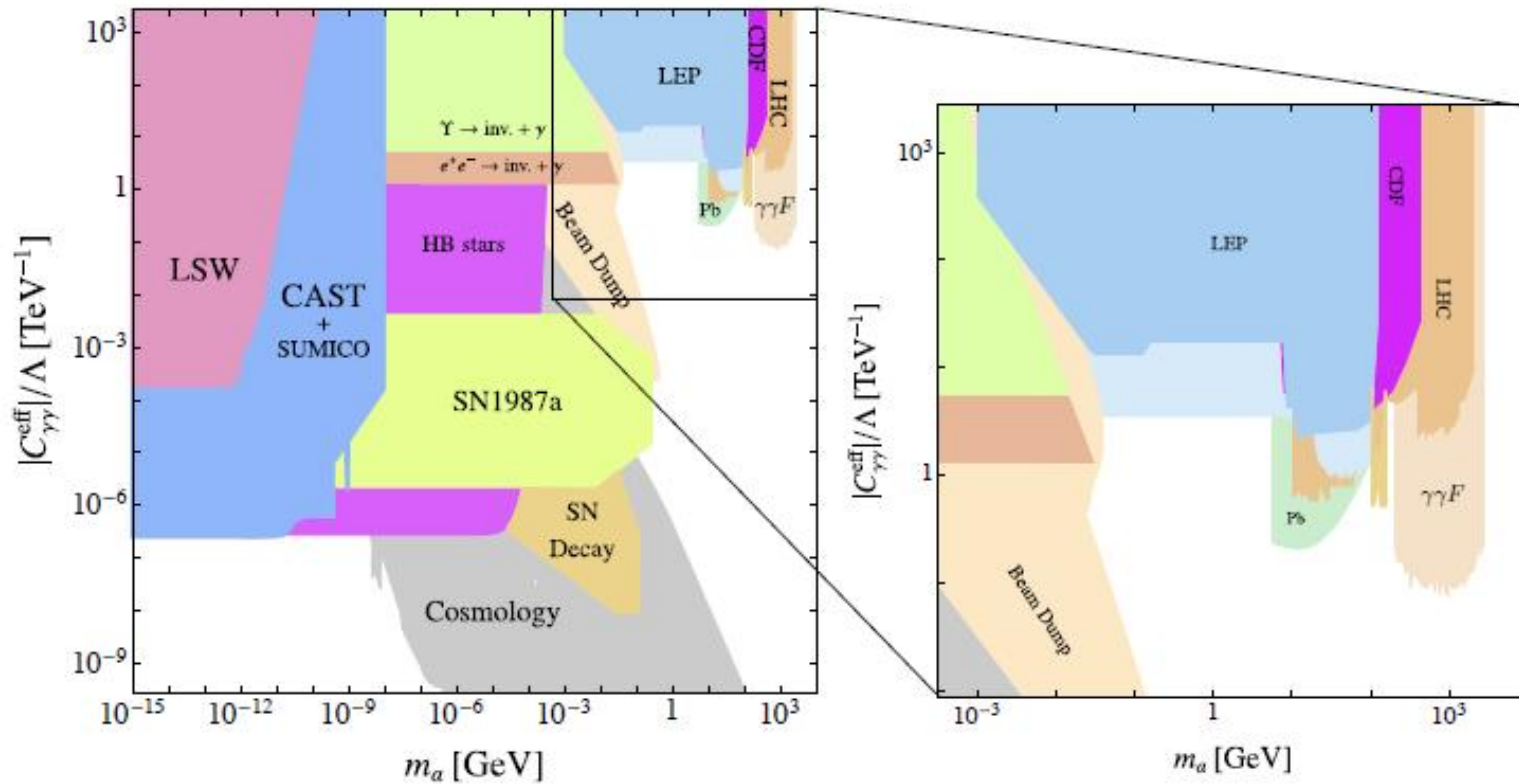
*I.Brivio et al., EPJC77(2017),8,572*

*M.Bauer et al., EPJC79(2019),1,74*

*CERN Yellow Rep. Monogr. Vol. 3 (2018)*

# Constraints: di-photon coupling

*M. Bauer et al., JHEP12(2017),044*



1. For light ALPs ( $\ll \text{MeV}$ ), cosmological and astrophysical measurements place very tight bounds on the coupling to photons.
2. For heavier ALPs, the limits are less stringent.



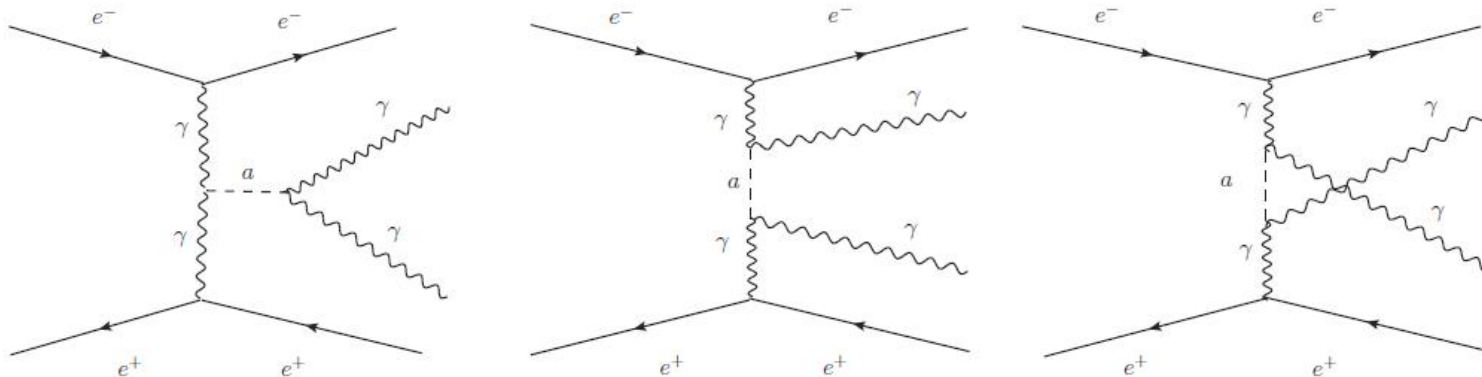
# Searching for ALPs at future e<sup>+</sup>e<sup>-</sup> colliders via light-by-light scattering

*H-Y Zhang, C-X Yue, Y-C Guo and SY*  
PRD 104 (2021) 9, 096008

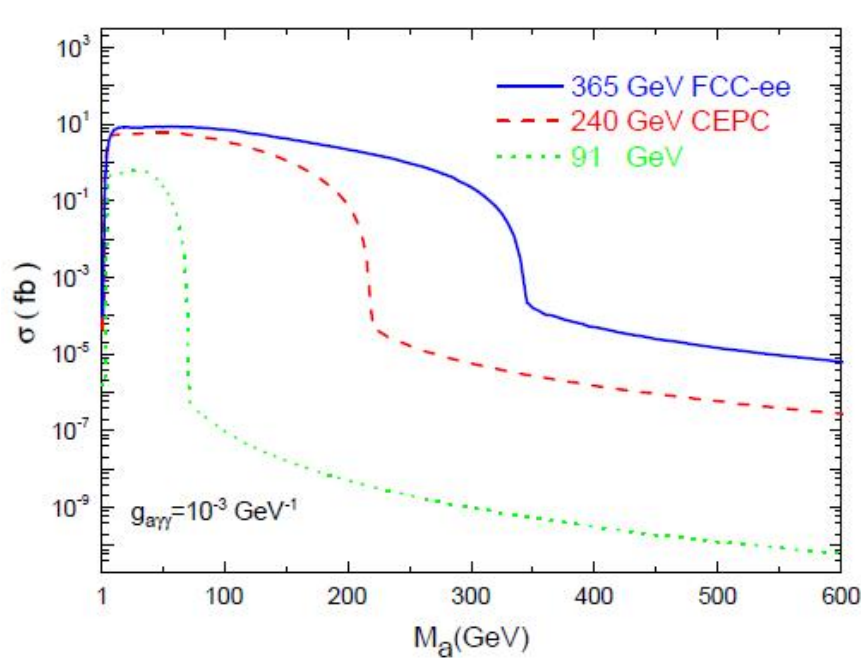
- The LHC generally is more sensitive to the heavy ALP searches by LBL scattering. The CLIC studies obtain a stronger bounds for TeV ALPs.

*C. Baldenegro et al., JHEP06,(2018)131 (LHC LBL)*  
*S.C. Inan and A.V. Kisselev, JHEP06(2020)183;*  
*Chin.Phys.C 45 (2021) 4, 043109 (CLIC LBL)*

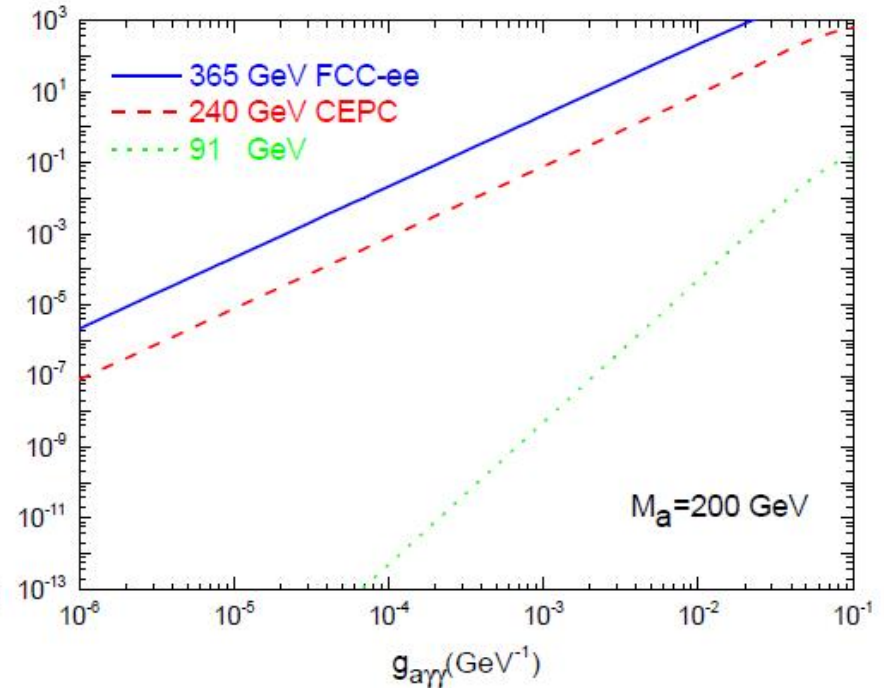
- It is interesting to study LBL at the CEPC and FCC-ee.



# Searching for ALPs at future e+e- colliders via light-by-light scattering



(a)



(b)

The cross section for LBL scattering

# Background for the LBL signal

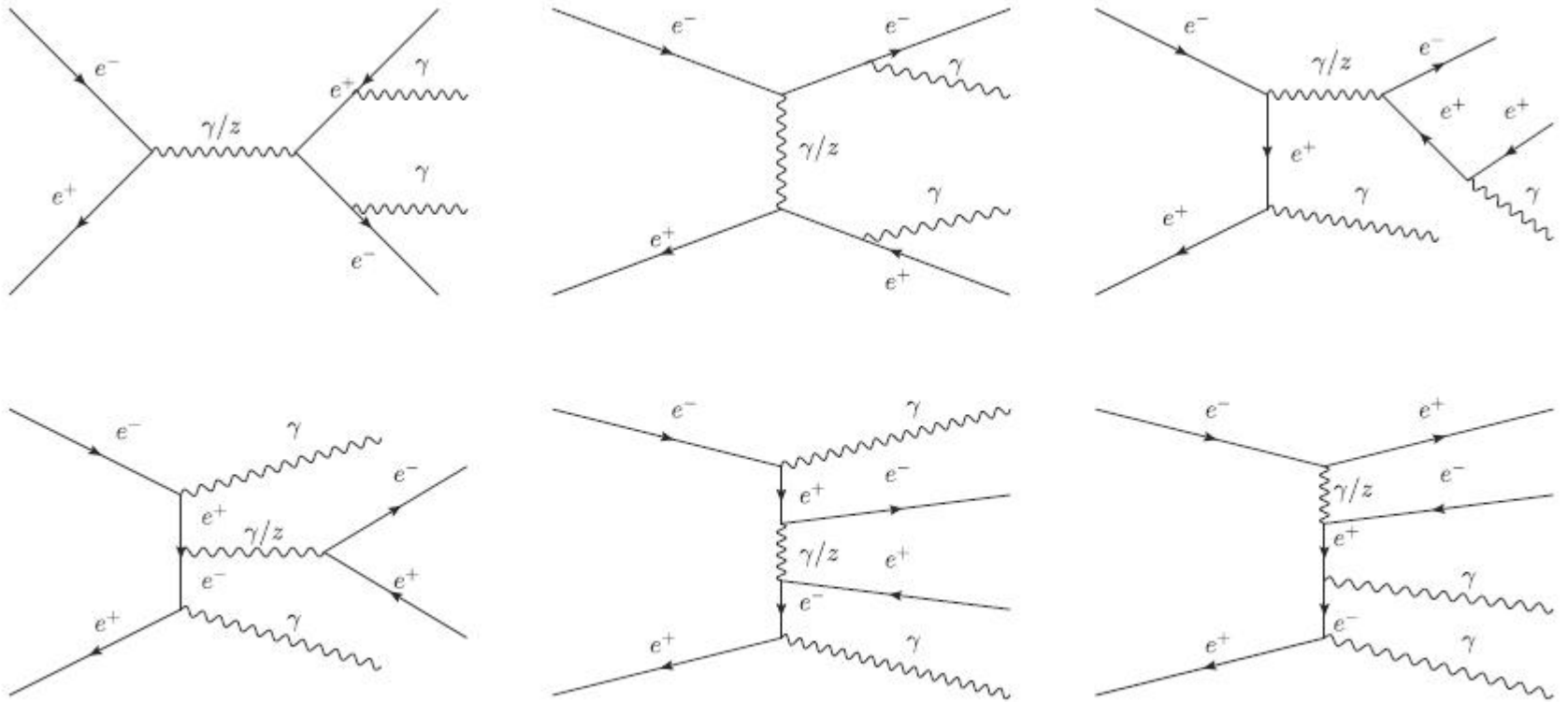


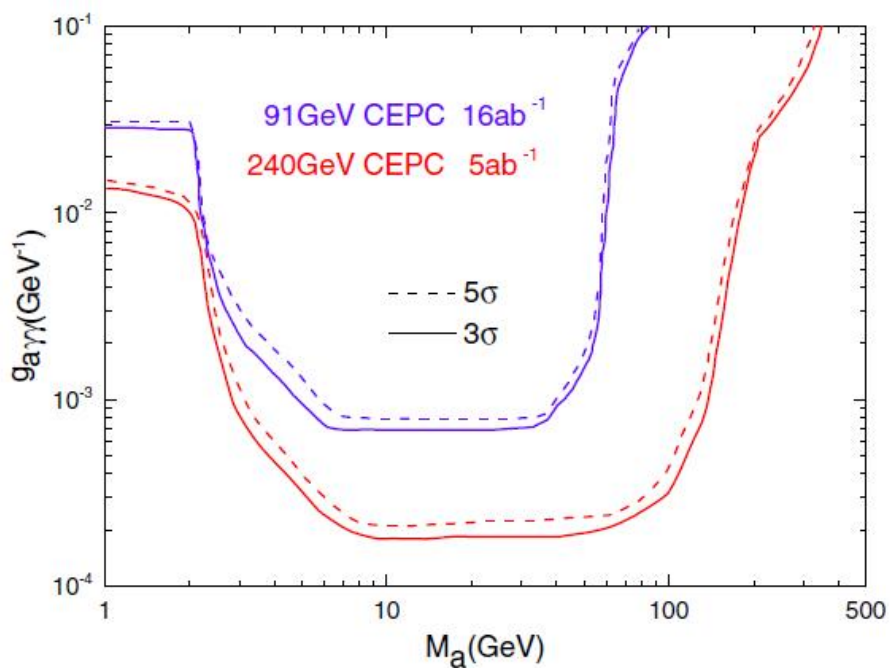
FIG. 3. The typical diagrams for the background of the process  $e^+e^- \rightarrow \gamma\gamma e^+e^-$ .

# Detecting ALPs at the CEPC

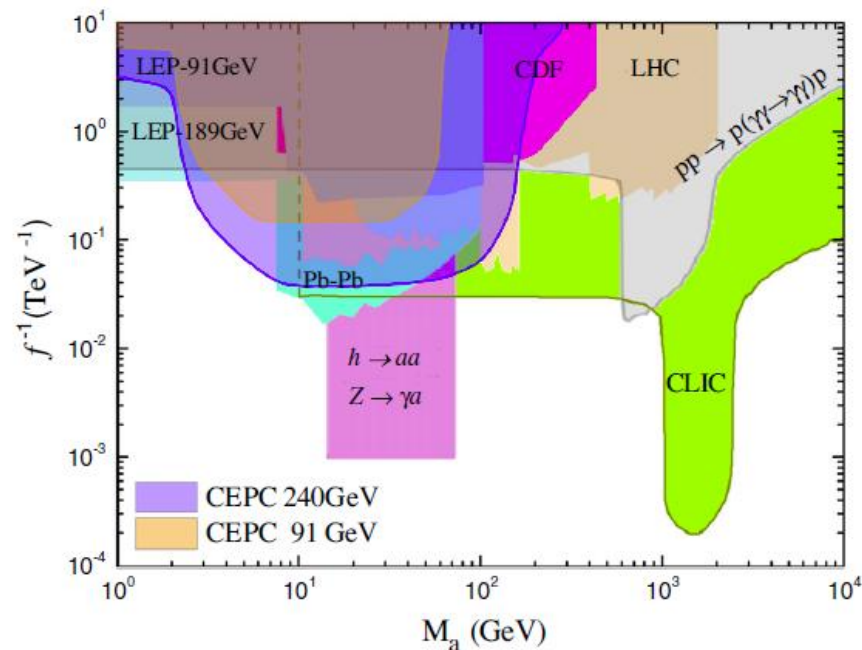
Cuts	$\sqrt{s} = 240 \text{ GeV}$	$\sqrt{s} = 91 \text{ GeV}$
Cut-1: Electron and positron pseudo-rapidity	$0.4 < \eta(e^+) < 2.4$ $-2.4 < \eta(e^-) < -0.4$	$-0.3 < \eta(e^+) < 0.9$ $-0.9 < \eta(e^-) < 0.3$
Cut-2: Angle between the ALP and the beam axis	$0.7 < \theta(\gamma\gamma) < 2.4$	$0.7 < \theta(\gamma\gamma) < 2.4$
Cut-3: Angular separation between electron-positron	$\Delta\theta(e^+e^-) < 2.9$	$\Delta\theta(e^+e^-) < 2.4$
Cut-4: Transverse momentum of reconstructed diphotons	$p_T(\gamma\gamma) > 45 \text{ GeV}$	$p_T(\gamma\gamma) > 20 \text{ GeV}$

CEPC @ $\sqrt{s} = 240 \text{ (91) GeV}$							
Cuts	Signal (fb)						Background (fb)
	$M_a = 6 \text{ GeV}$	$M_a = 8 \text{ GeV}$	$M_a = 10 \text{ GeV}$	$M_a = 50 \text{ GeV}$	$M_a = 100 \text{ GeV}$	$M_a = 160 \text{ GeV}$	$\gamma\gamma e^+e^-$
Basic cuts	3.4378(0.249)	4.8088(0.4796)	5.2928(0.5003)	5.9064(0.2432)	3.585	0.8021	67.0614(98.8986)
Cut 1	2.9865(0.0316)	3.932(0.1267)	4.138(0.1417)	4.5336(0.0977)	2.4778	0.4436	33.7026(40.928)
Cut 2	2.1714(0.0309)	3.0176(0.1264)	3.2819(0.1411)	3.1262(0.0904)	1.6993	0.3145	12.628(34.93)
Cut 3	2.1368(0.0226)	3.0383(0.1156)	3.2422(0.1297)	3.0238(0.0717)	1.6497	0.3052	9.042(8.396)
Cut 4	1.4(0.0226)	2.2984(0.1156)	2.5065(0.1297)	2.0519(0.0501)	0.8747	0.0392	3.3614(6.1921)

# Prospects for detecting ALP at CEPC



The  $3\sigma$  and  $5\sigma$  discovery curves in the  $M_a - g_{a\gamma\gamma}$  plane at the CEPC.



The 95% C.L. exclusion regions in the  $M_a - g_{a\gamma\gamma}$  plane at the CEPC and other colliders



# ALPs at the CEPC & FCC-ee

CEPC @ $\sqrt{s} = 240$ (91) GeV							
Cuts	Signal (fb)						Background (fb)
	$M_a = 6$ GeV	$M_a = 8$ GeV	$M_a = 10$ GeV	$M_a = 50$ GeV	$M_a = 100$ GeV	$M_a = 160$ GeV	$\gamma\gamma e^+e^-$
Basic cuts	3.4378(0.249)	4.8088(0.4796)	5.2928(0.5003)	5.9064(0.2432)	3.585	0.8021	67.0614(98.8986)
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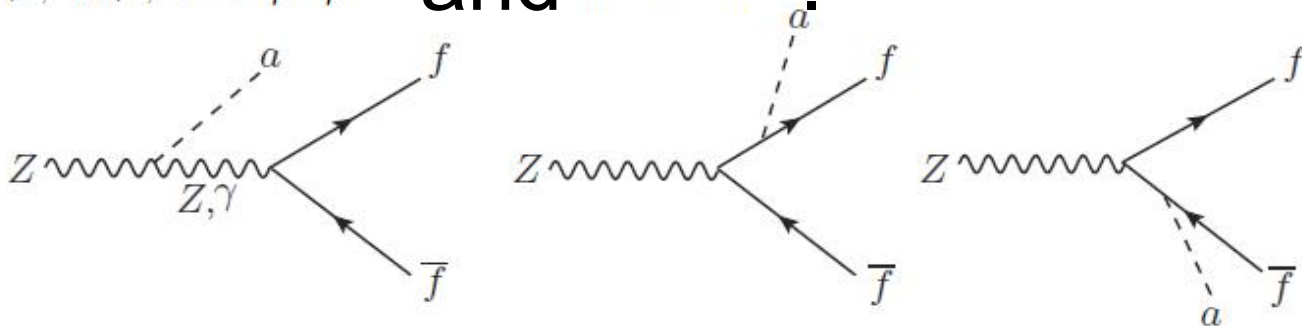
FCC-ee @ $\sqrt{s} = 365$ (91) GeV							
Cuts	Signal (fb)						Background (fb)
	$M_a = 6$ GeV	$M_a = 8$ GeV	$M_a = 10$ GeV	$M_a = 50$ GeV	$M_a = 100$ GeV	$M_a = 200$ GeV	$\gamma\gamma e^+e^-$
Basic cuts	2.9092(0.2483)	5.0074(0.4786)	6.5272(0.5001)	8.4206(0.2432)	7.1235	2.1737	54.203(98.8188)
Cut 1	2.1634(0.0311)	4.2978(0.1265)	5.3419(0.142)	4.5123( 0.0977)	4.9093	1.2593	29.233(41.0505)
Cut 2	1.3962(0.0307)	2.6956(0.1261)	3.6755(0.1416)	2.9963(0.0904)	3.1011	0.7942	8.3373(35.0206)
Cut 3	1.2374(0.0223)	2.5417(0.1152)	3.5173(0.1304)	2.8482 (0.0717)	2.9926	0.768	4.8137(8.4019)
Cut 4	0.9014(0.0222)	2.2243(0.115)	3.1819(0.1303)	2.5198(0.05)	2.5458	0.453	2.6445(6.1842)

# Searching for ALPs via exotic decay $Z \rightarrow a f \bar{f}$ at CEPC

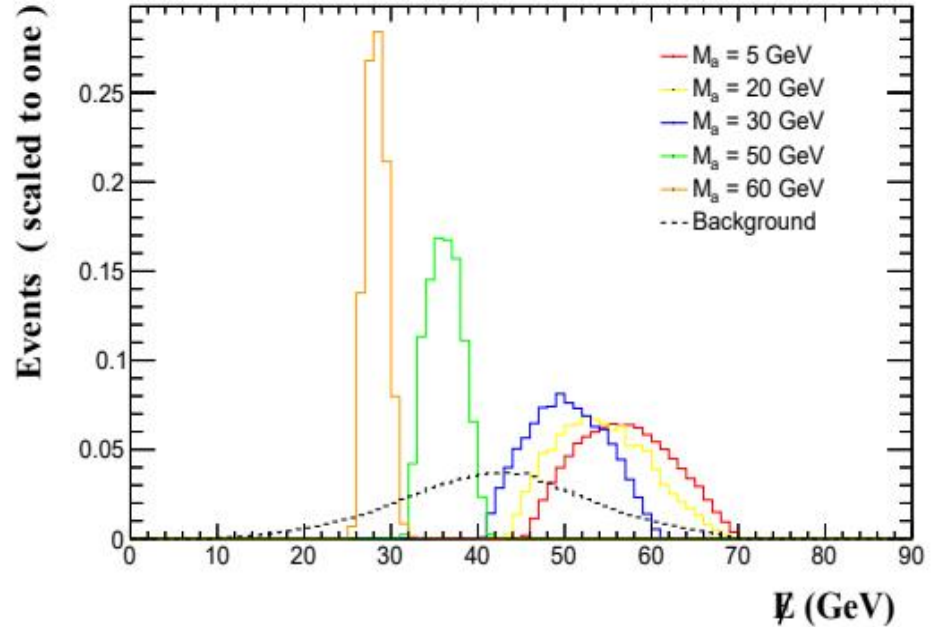
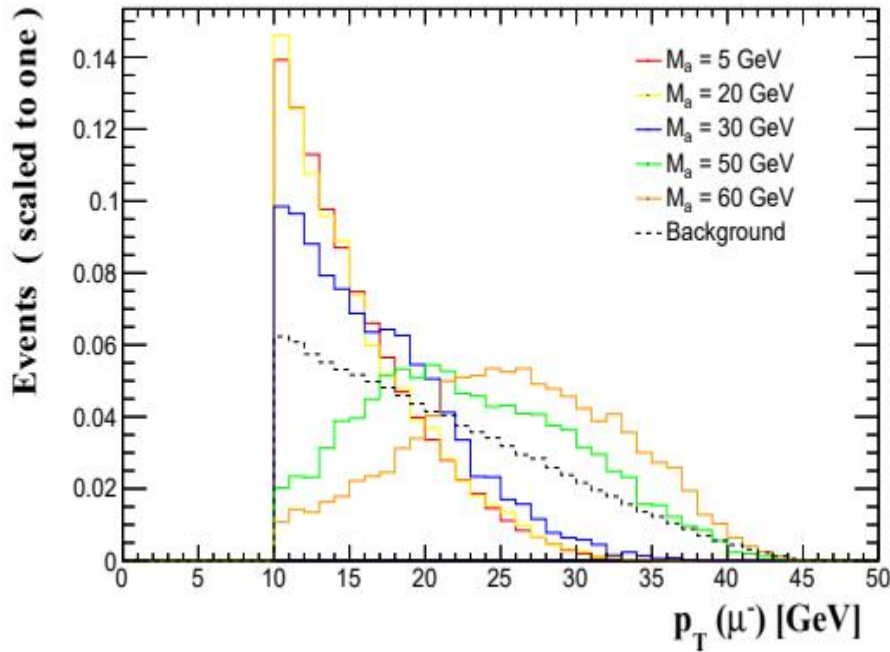
C.-X. Yue, **SY**, H. Wang and N. Zhang, arXiv:2204.04702 accepted by PRD

- CEPC and FCC-ee can produce up to  $10^{12}$  Z bosons.
- It is possible for the observations of rare decay of Z.
- Future Z factories is powerful for detecting dark sector models via exotic Z decay. J.Liu, *et al.*, PRD97,095044,2018.
- We focus on four types of exotic Z-decay signals

$$Z \rightarrow \mu^+ \mu^- \cancel{E}, bb \cancel{E}, e^+ e^- \mu^+ \mu^- \text{ and } e^+ e^- bb$$



$$Z \rightarrow \mu^+ \mu^- \cancel{E} \quad \text{and} \quad Z \rightarrow b\bar{b} \cancel{E}$$



The normalized distribution for kinematic variables for signal  $\mu^+ \mu^- \cancel{E}$  and BG.



$$Z \rightarrow \mu^+ \mu^- \cancel{E} \quad \text{and} \quad Z \rightarrow b\bar{b} \cancel{E}$$

Cuts	Cross sections for signal(background) (fb)				
	$m_a = 5\text{GeV}$	$m_a = 10\text{GeV}$	$m_a = 30\text{GeV}$	$m_a = 50\text{GeV}$	$m_a = 60\text{GeV}$
Basic cuts	0.3406(0.2602)	0.3177(0.2602)	0.2368(0.2602)	0.0358(0.2602)	0.0062(0.2602)
Cut 1-A	0.3404(0.0059)	0.3175(0.0072)	0.2335(0.0300)	0.0343(0.0342)	0.0058(0.0195)
$S/\sqrt{S+B}$	18.29	17.62	14.38	5.59	1.16

Cut1-A

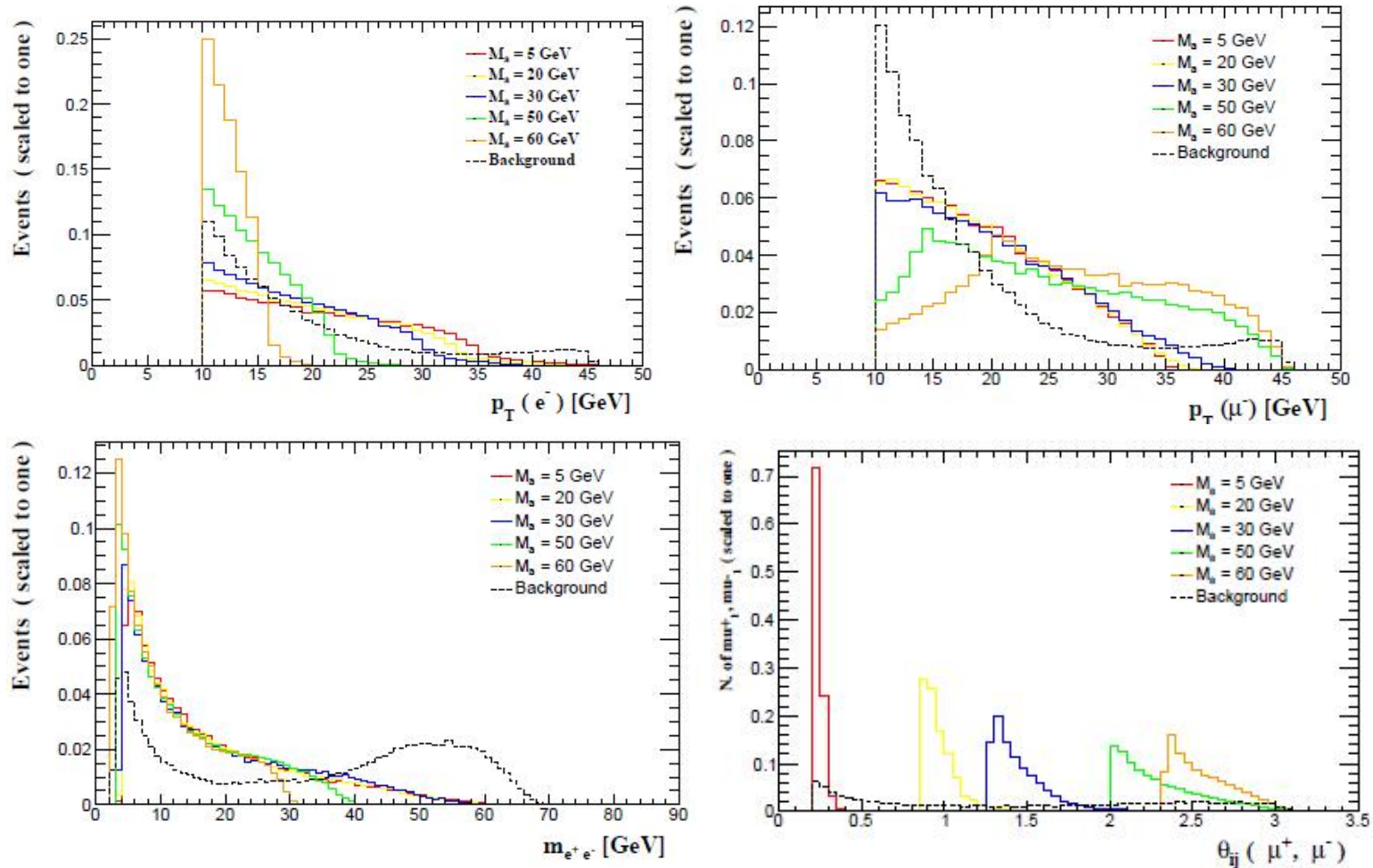
$$|m_{\mu^+\mu^-} - m_a| < 3 \text{ GeV}$$

Cuts	Cross sections for signal(background) (fb)				
	$m_a = 15\text{GeV}$	$m_a = 30\text{GeV}$	$m_a = 40\text{GeV}$	$m_a = 50\text{GeV}$	$m_a = 60\text{GeV}$
Basic cuts	0.0460(0.7172)	0.03345(0.7172)	0.0284(0.7172)	0.0142(0.7172)	0.0134(0.7172)
Cut 1-B	0.0449(0.0126)	0.0279(0.08387)	0.0199(0.1832)	0.0078(0.2115)	0.0015(0.1364)
$S/\sqrt{S+B}$	5.92	2.64	1.39	0.53	0.12

Cut1-B

$$|m_{b\bar{b}} - m_a| < 5 \text{ GeV}$$

$$Z \rightarrow e^+ e^- \mu^+ \mu^- \text{ and } Z \rightarrow e^+ e^- b\bar{b}$$



The normalized distribution for kinematic variables for signal  $e^+ e^- \mu^+ \mu^-$  and BG.

$$Z \rightarrow e^+e^-\mu^+\mu^- \text{ and } Z \rightarrow e^+e^-bb$$

Cuts	Cross sections for signal(background) (fb)				
	$m_a = 5\text{GeV}$	$m_a = 10\text{GeV}$	$m_a = 30\text{GeV}$	$m_a = 50\text{GeV}$	$m_a = 60\text{GeV}$
Basic cuts	1.5314(8.0284)	1.4735(8.0284)	1.1559(8.0284)	0.4615(8.0284)	0.1067(8.0284)
Cut 1-C	1.2659(3.4300)	1.2231(3.4300)	0.9318(3.4300)	0.4258(3.4300)	0.1065(3.4300)
Cut 2-C	1.2659(0.1743)	1.2215(0.1623)	0.9143(0.1689)	0.4066(0.6499)	0.1005(0.5635)
$S/\sqrt{S+B}$	33.36	32.83	27.78	12.51	3.90

Cut-C

$$m_{e^+e^-} < 30 \text{ GeV}$$

$$|m_{\mu^+\mu^-} - m_a| < 3 \text{ GeV}$$

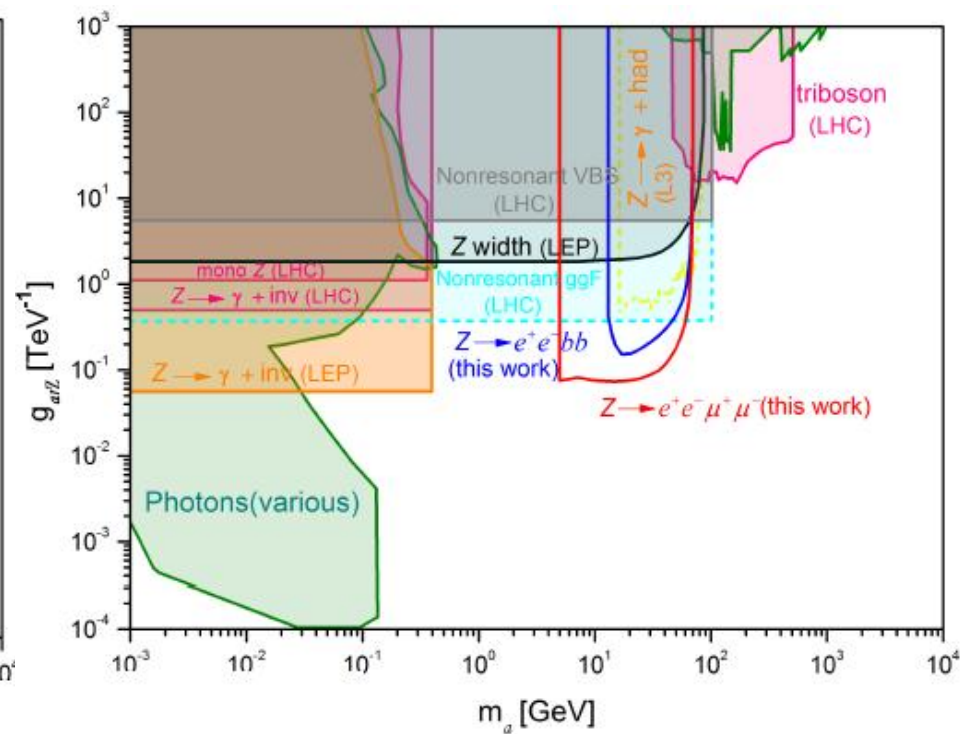
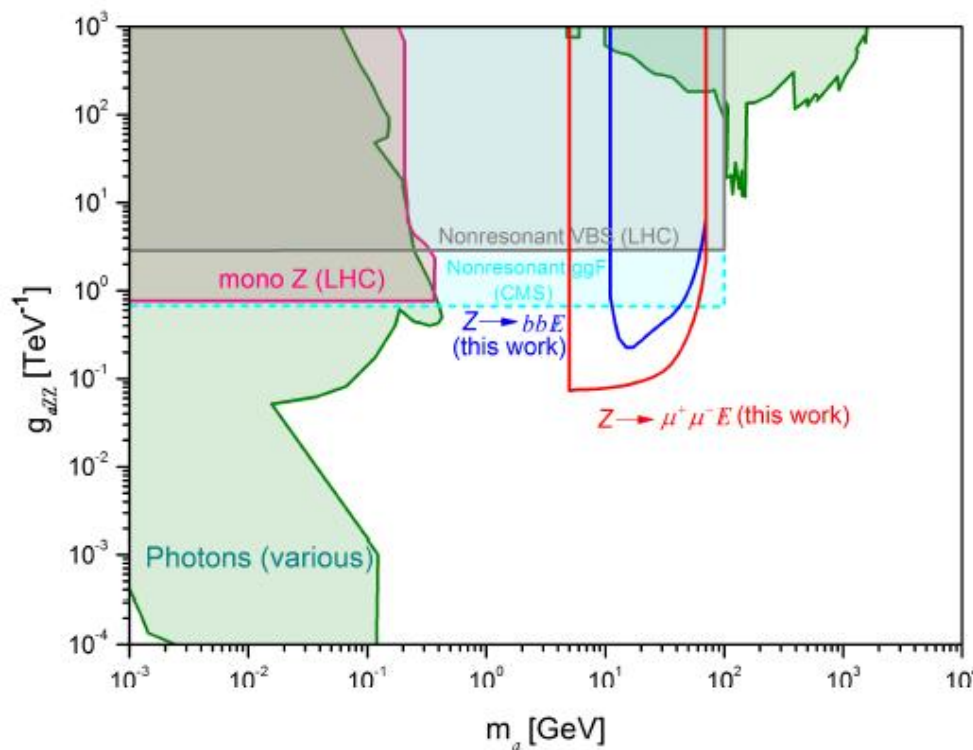
Cuts	Cross sections for signal(background) (fb)				
	$m_a = 15\text{GeV}$	$m_a = 30\text{GeV}$	$m_a = 40\text{GeV}$	$m_a = 50\text{GeV}$	$m_a = 60\text{GeV}$
Basic cuts	0.2058(2.7076)	0.2497(2.7076)	0.1874(2.7076)	0.1191(2.7076)	0.0354(2.7076)
Cut 1-D	0.1711(1.8332)	0.2251(1.8332)	0.1658(1.8332)	0.1103(1.8332)	0.0353(1.8332)
Cut 2-D	0.1694(0.1417)	0.1763(0.2081)	0.1106(0.3301)	0.0617(0.5102)	0.0160(0.4212)
$S/\sqrt{S+B}$	9.61	8.99	5.27	2.58	0.76

Cut-D

$$m_{e^+e^-} < 30 \text{ GeV}$$

$$|m_{bb} - m_a| < 5 \text{ GeV}$$

# Sensitivity bounds



Sensitivity bounds on  $g_{a\ell\ell}$  and  $m_a$  at 95% C.L. from exotic Z decays and other current exclusion regions.

# Summary

- ALPs have a much wider parameter space and hence generate rich phenomenology at colliders.
- CEPC and FCC-ee can provide a good environment to exploring ALPs.
- The studies on LBL scattering and exotic decay  $Z \rightarrow a\bar{f}f$  found that CEPC and FCC-ee might be more sensitive to the ALPs in light mass range than LHC. It is expected that the future  $e^+e^-$  colliders could discovery or exclude ALPs.



