

# Probing dark matter particles at CEPC

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[arXiv:1903.12114, ZL, Yong-Heng Xu, Yu Zhang]

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# Outline

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## CEPC sensitivities on different dark matter (DM) models

1

Millicharged DM

2

Vector portal DM

3

DM with EFT interactions

# Interaction Lagrangian for different Dirac DM models

**millicharged DM**  $\mathcal{L} = e\varepsilon A_\mu \bar{\chi} \gamma^\mu \chi$

[Feldman, ZL, Nath, hep-ph/0702123, 331 cites]

**vector-portal DM**  $\mathcal{L} = Z'_\mu \bar{\chi} \gamma^\mu (g_V^\chi - g_A^\chi \gamma_5) \chi + Z'_\mu \bar{f} \gamma^\mu (g_V^f - g_A^f \gamma_5) f$

## DM models with EFT interaction

**scalar (s)**  $\frac{1}{\Lambda_s^2} \bar{\chi} \chi \bar{\ell} \ell$

**vector**  $\frac{1}{\Lambda_V^2} \bar{\chi} \gamma_\mu \chi \bar{\ell} \gamma^\mu \ell$

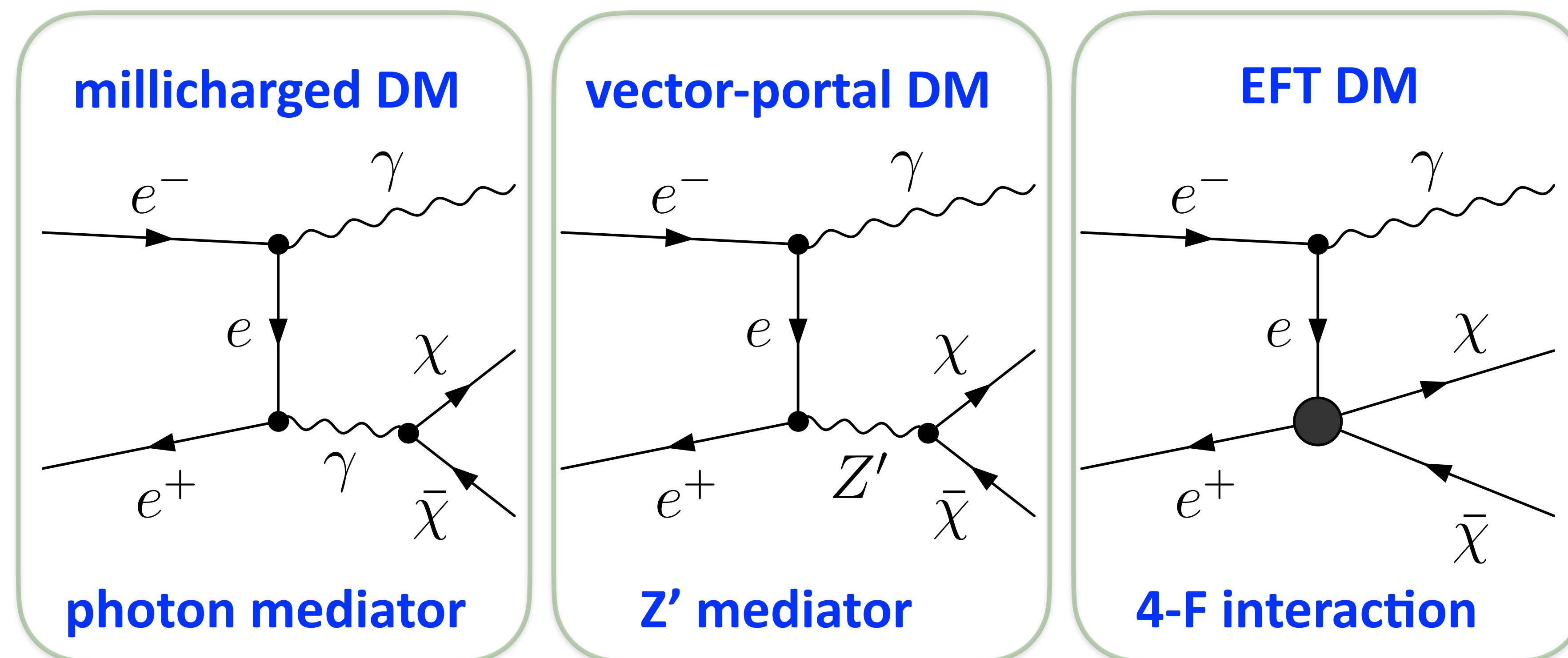
**scalar (t)**  $\frac{1}{\Lambda_t^2} \bar{\chi} \ell \bar{\ell} \chi$

**axial-vector**  $\frac{1}{\Lambda_A^2} \bar{\chi} \gamma_\mu \gamma_5 \chi \bar{\ell} \gamma^\mu \gamma_5 \ell$

[Fox, Harnik, Kopp, Tsai, 1103.0240; Chae, Perelstein, 1211.4008]

# Signal process: mono-photon final state at CEPC

new physics process:  $e^+e^- \rightarrow \bar{\chi}\chi\gamma$



maximum photon energy

$$E_\gamma < \frac{s - 4m_\chi^2}{2\sqrt{s}} \equiv E_\chi^m$$

# 3 CEPC running modes

H-mode: 240 GeV & 5.6/ab

Z-mode: 91.2 GeV & 16 /ab

WW-mode: 160 GeV & 2.6/ab

[CEPC Conceptual Design Report: Volume 2 - Physics & Detector, 1811.10545]

Operation mode	$\sqrt{s}$ (GeV)	$L$ per IP ( $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ )	Years	Total $\int L$ ( $\text{ab}^{-1}$ , 2 IPs)	Event yields
$H$	240	3	7	5.6	$1 \times 10^6$
$Z$	91.2	32 (*)	2	16	$7 \times 10^{11}$
$W^+W^-$	158–172	10	1	2.6	$2 \times 10^7$ (†)

# NP mono-photon xsec @ CEPC

millicharged DM

$$\frac{d\sigma}{dE_\gamma dz_\gamma} = \frac{8\alpha^3 \varepsilon^2 (1+2y)\beta_\chi}{3sE_\gamma} \left[ \frac{1+x(1+z_\gamma^2)}{1-z_\gamma^2} \right]$$

[Liu, Zhang, [1808.00983](#)]

vector-portal DM

$$\frac{d\sigma}{dE_\gamma dz_\gamma} = \frac{\alpha \left[ (g_V^f)^2 + (g_A^f)^2 \right] \left[ (g_V^\chi)^2 (1+2y) + (g_A^\chi)^2 (1-4y) \right] s_\gamma^2 \beta_\chi}{6\pi^2 s E_\gamma [(s_\gamma - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2]} \left[ \frac{1+x(1+z_\gamma^2)}{1-z_\gamma^2} \right]$$

$E_\gamma$  is the energy of the final state photon

$$z_\gamma \equiv \cos \theta_\gamma \quad s_\gamma = s - 2\sqrt{s}E_\gamma$$

$\theta_\gamma$  is the polar angle of the final state photon

$$x \equiv E_\gamma^2 / s_\gamma$$

$s$  is the center-of-mass energy square

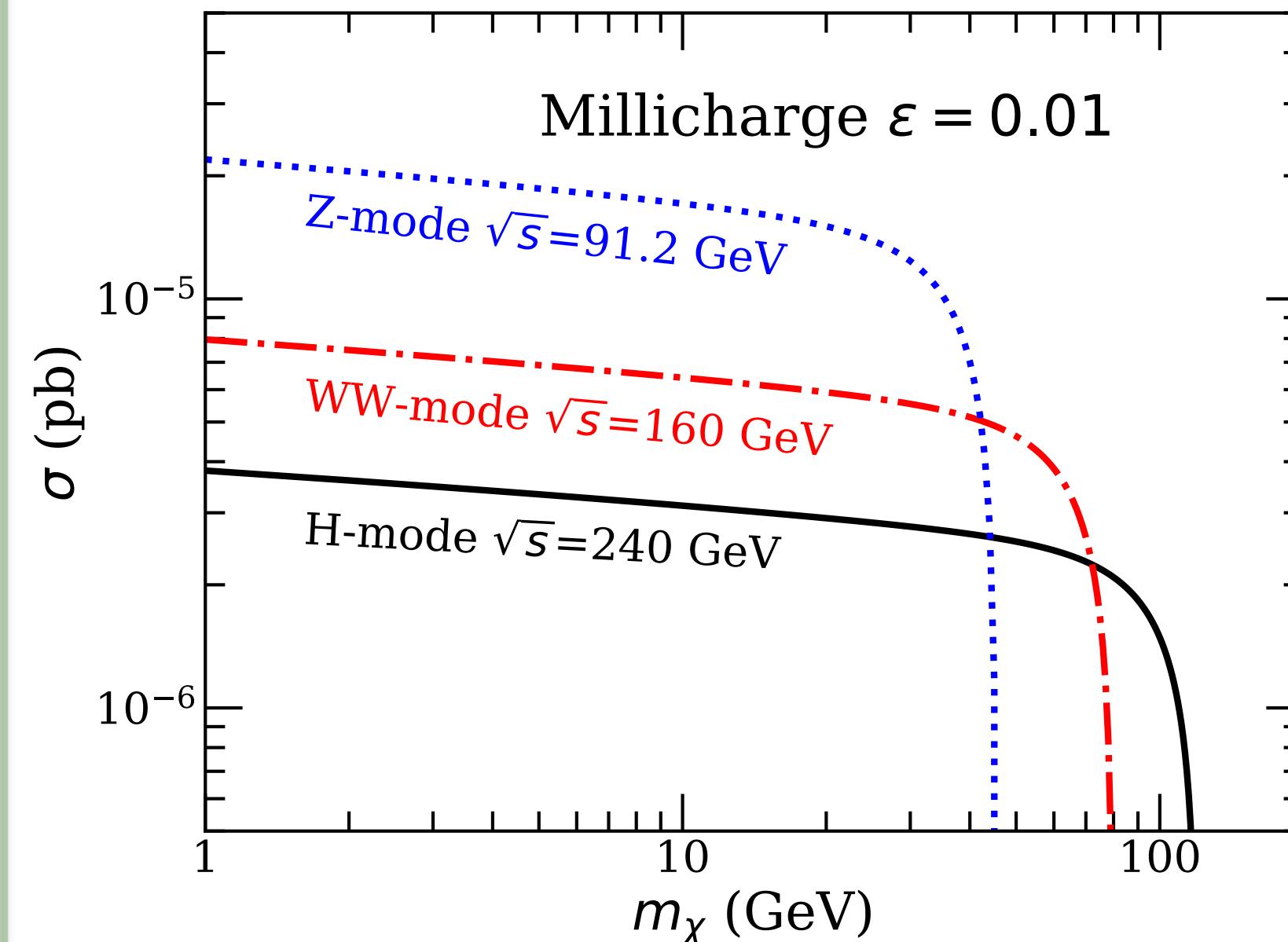
$$y \equiv m_\chi^2 / s_\gamma \quad \beta_\chi = (1-4y)^{1/2}$$

$m_\chi$  is the DM mass

# mono-photon signal (mediators)

detector cuts  $E_\gamma > 0.1$  GeV &  $|z_\gamma| < 0.99$

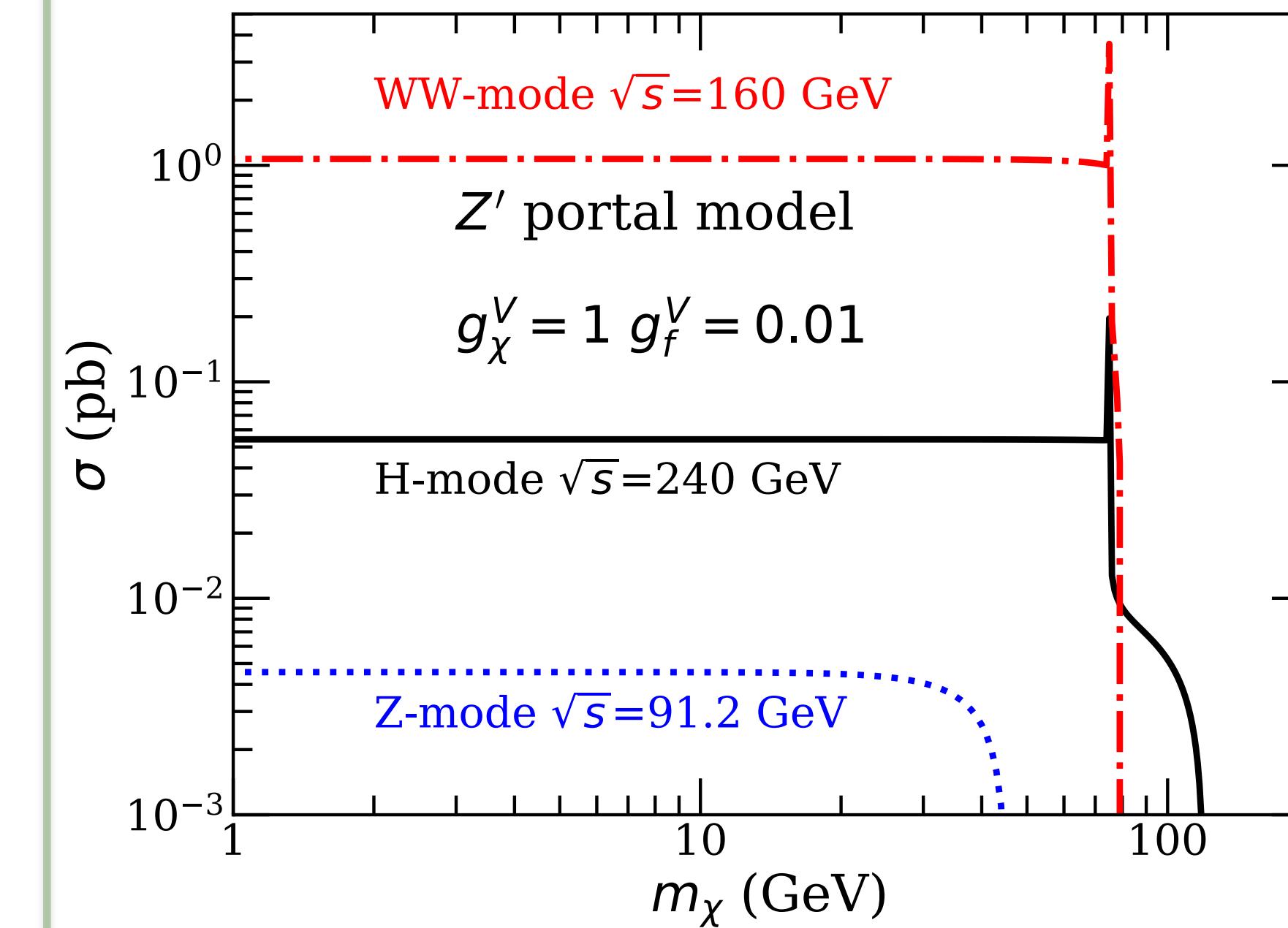
Millicharged DM  $\epsilon = 0.01$



Z>WW>H

vector-coupling

Z'-portal DM  $m_{Z'} = 150$  GeV

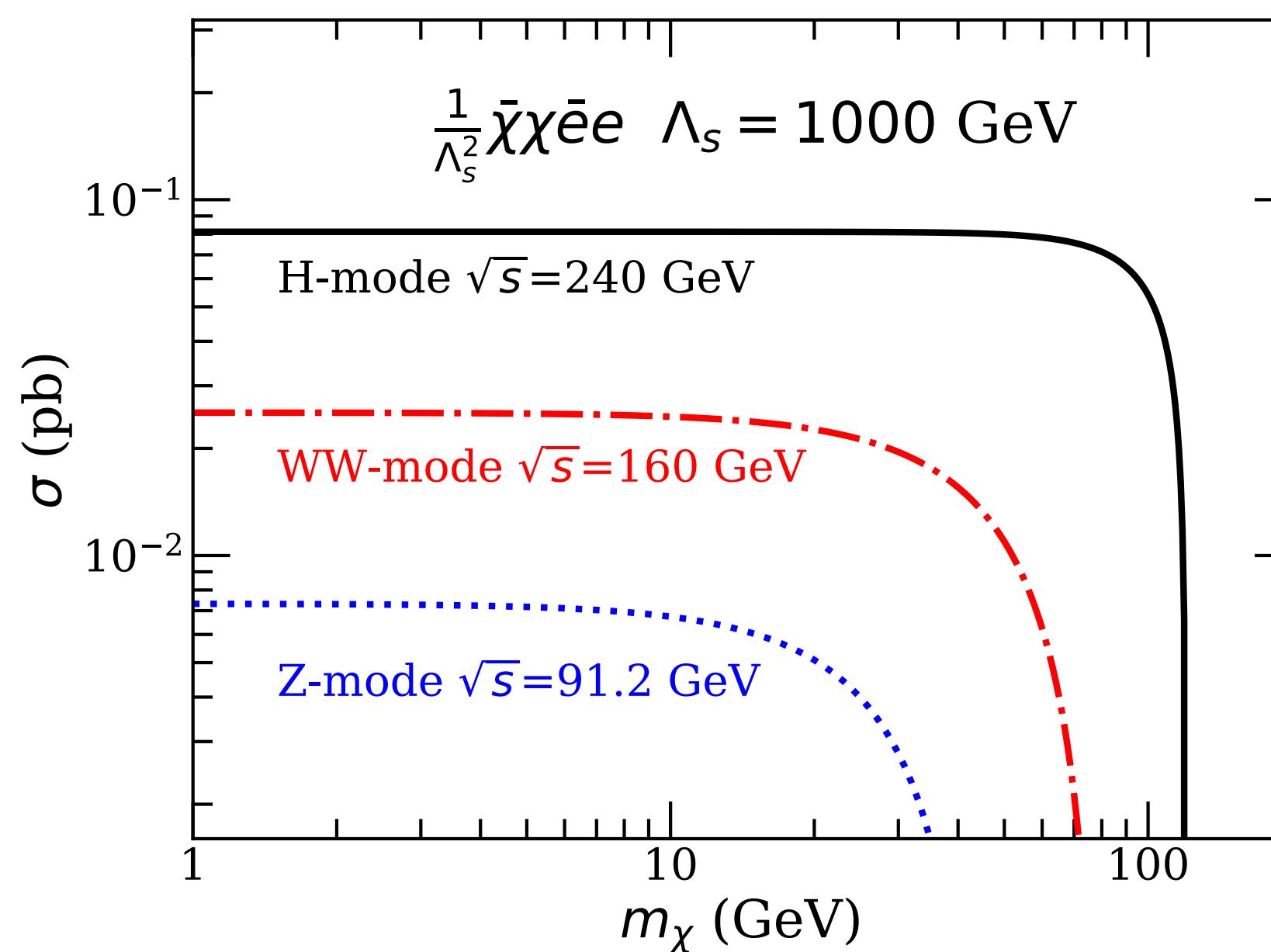


WW>H>Z

# mono-photon signal (EFT)

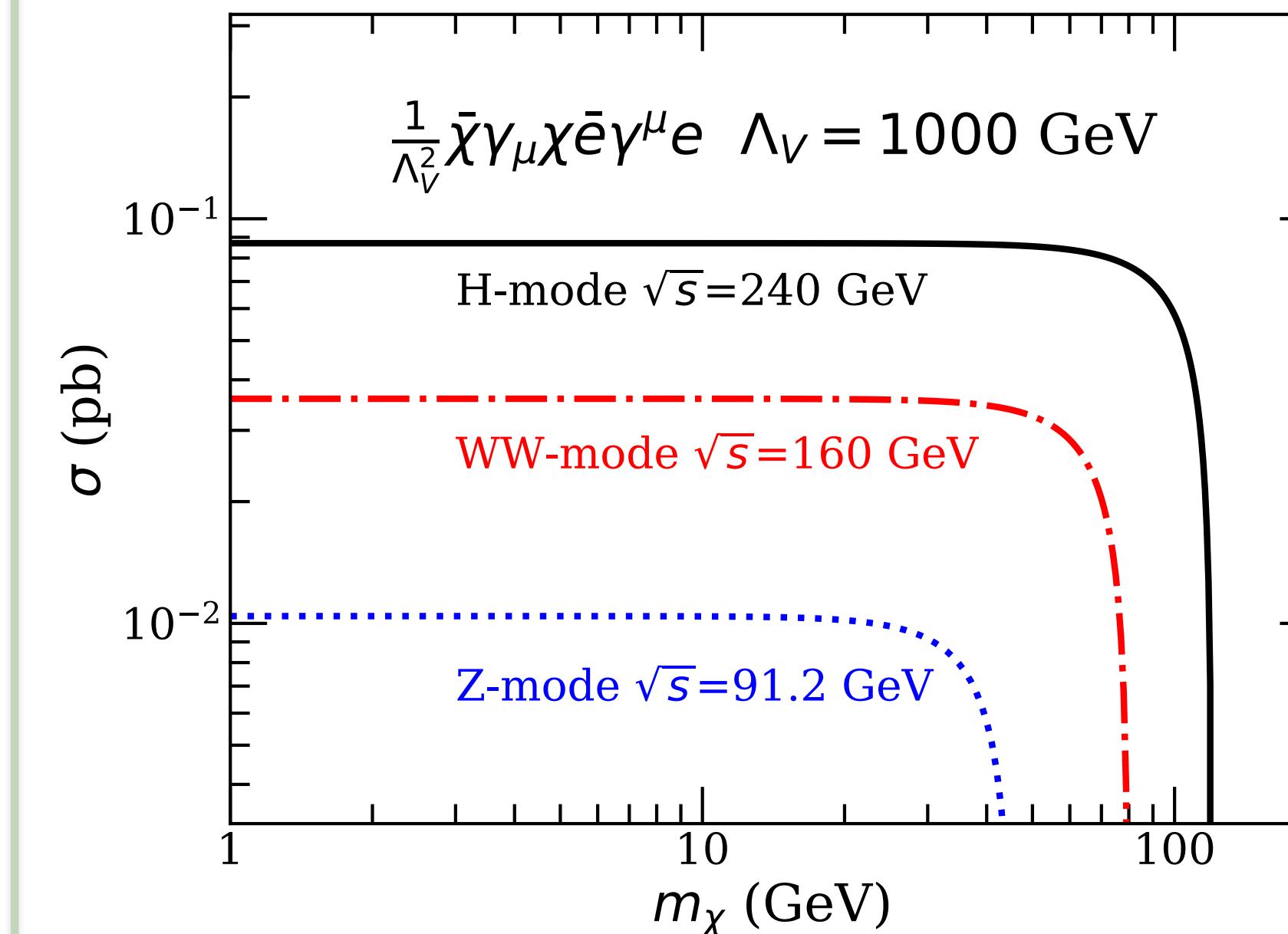
detector cuts  $E_\gamma > 0.1$  GeV &  $|z_\gamma| < 0.99$

scalar-EFT DM  $\Lambda = \text{TeV}$



H>WW>Z

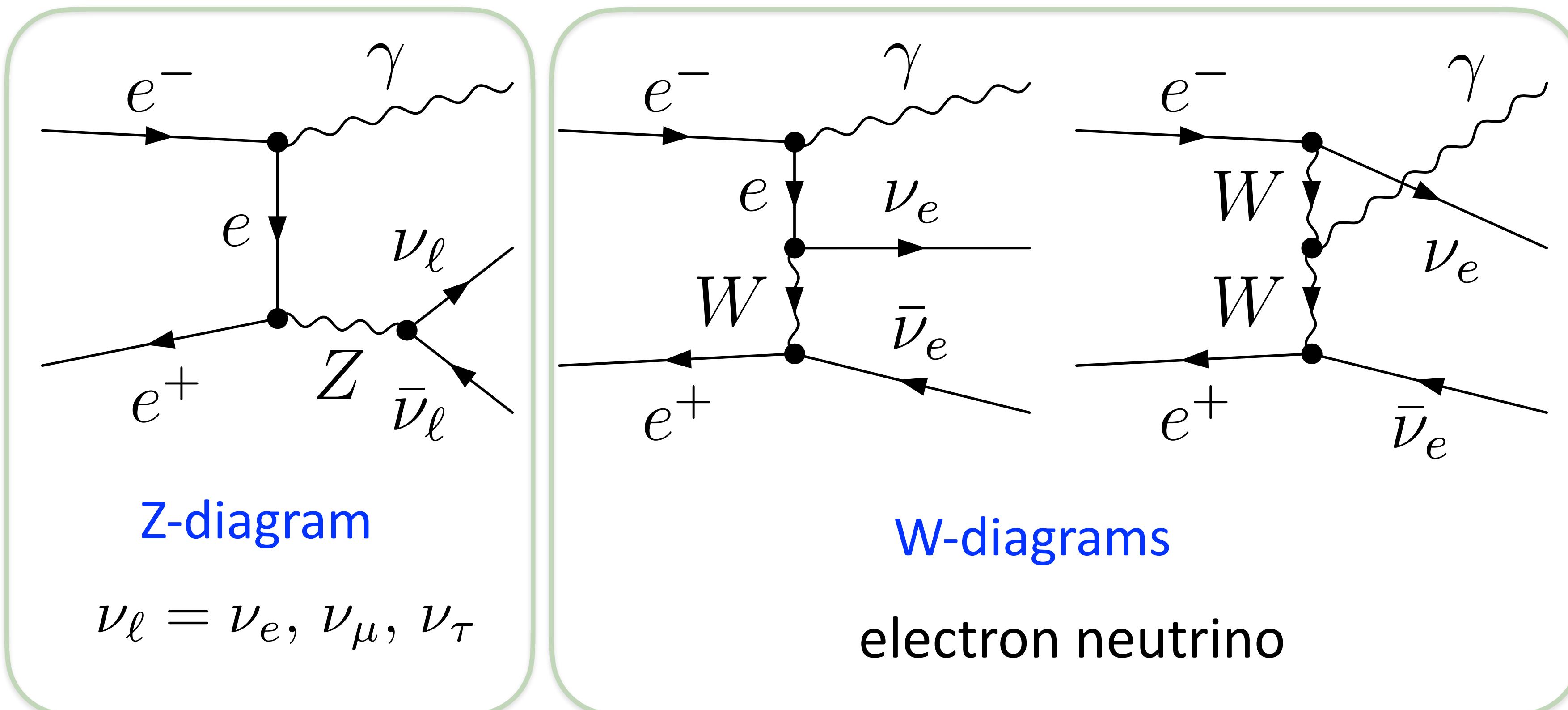
vector-EFT DM  $\Lambda = \text{TeV}$



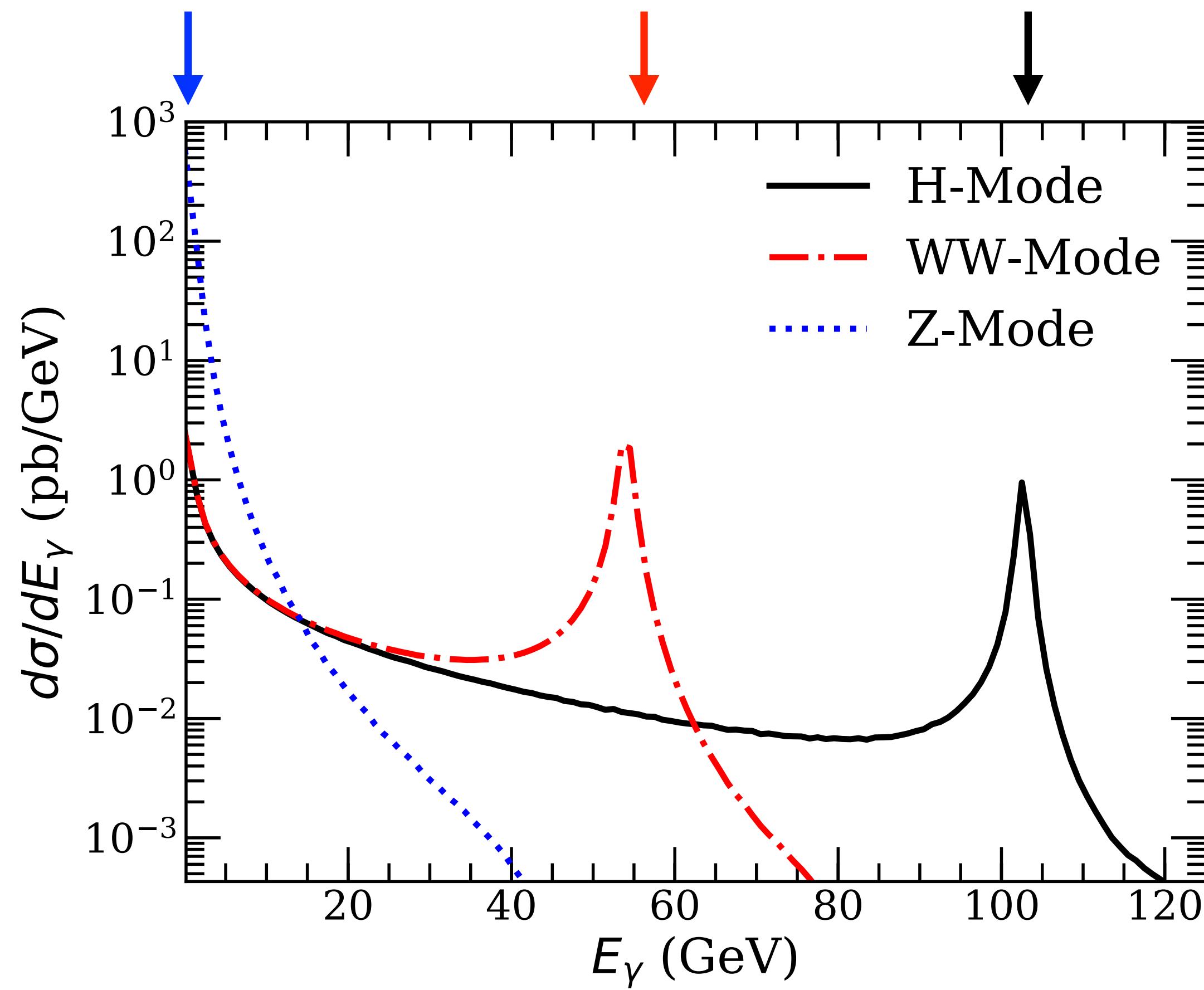
H>WW>Z

# Irreducible SM background

irreducible BG:  $e^+e^- \rightarrow \bar{\nu}\nu\gamma$



# Photon energy in irreducible BG



"Z resonance"

FWHM

veto

$$E_\gamma^Z \in (E_\gamma^Z \pm 5\Gamma_\gamma^Z)$$

$$E_\gamma^Z = \frac{s - m_Z^2}{2\sqrt{s}}$$

$$\Gamma_\gamma^Z = \frac{m_Z \Gamma_Z}{\sqrt{s}}$$

# Reducible SM background

Reducible BG:  $e^+e^- \rightarrow \gamma + X$  w/  $X = f\bar{f}, \gamma, \gamma\gamma$

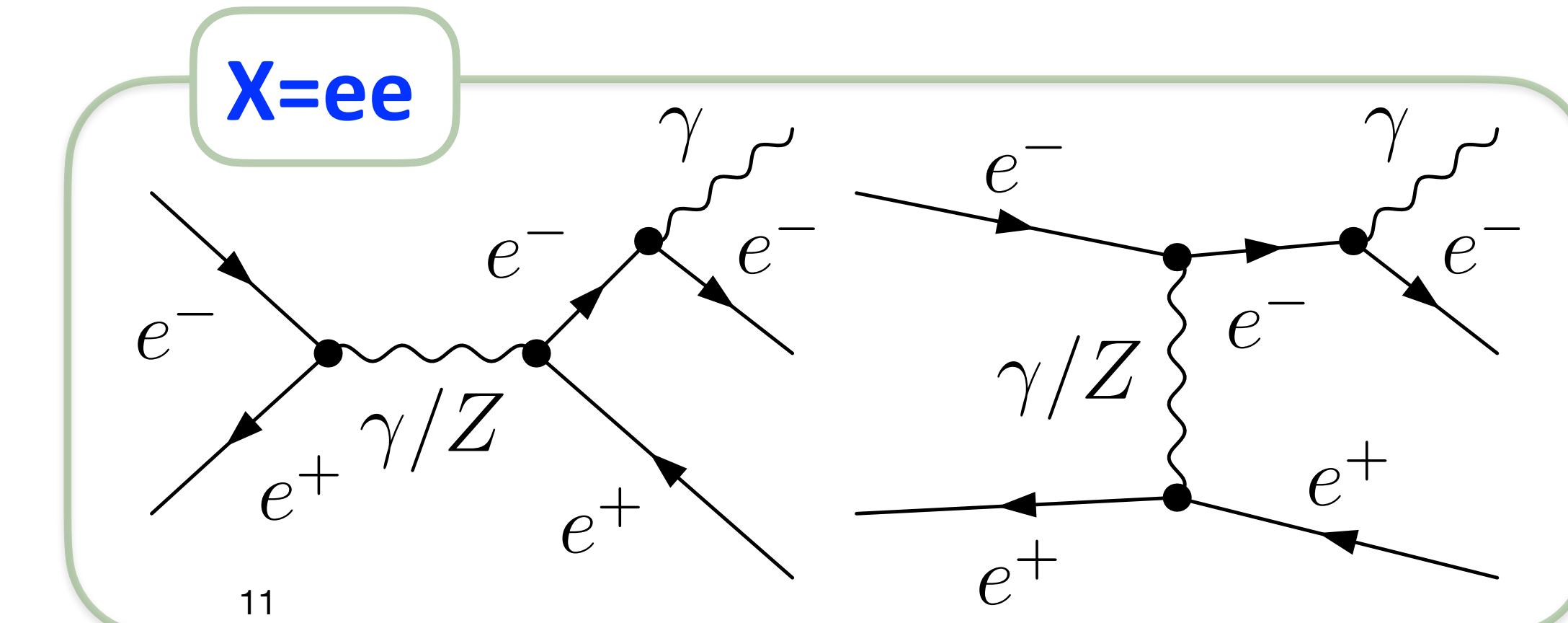
$X$  becomes invisible due to limited detection capability

[1811.10545] EMC coverage:  $|\cos \theta| < 0.99$  &  $E > 0.1$  GeV

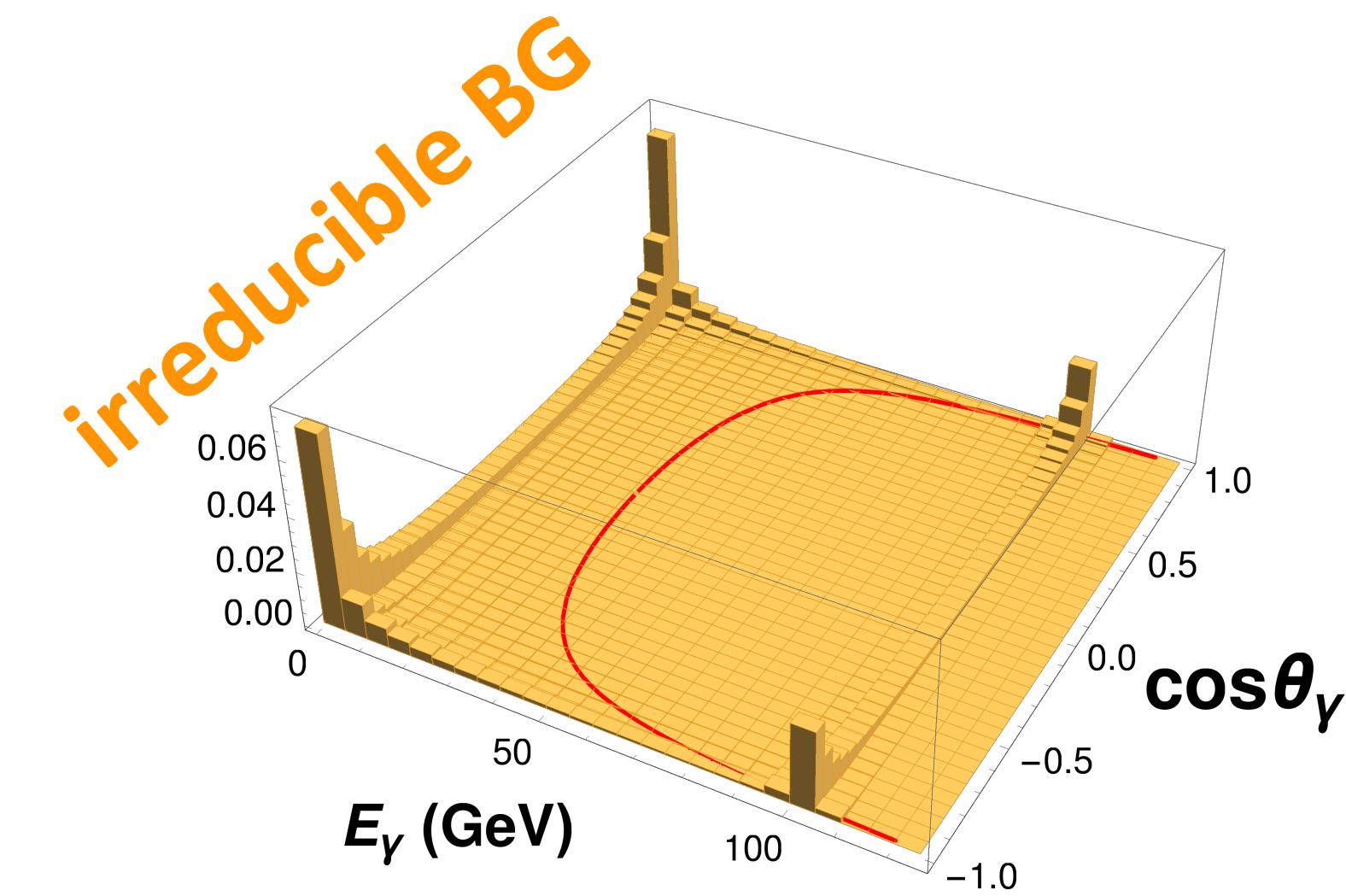
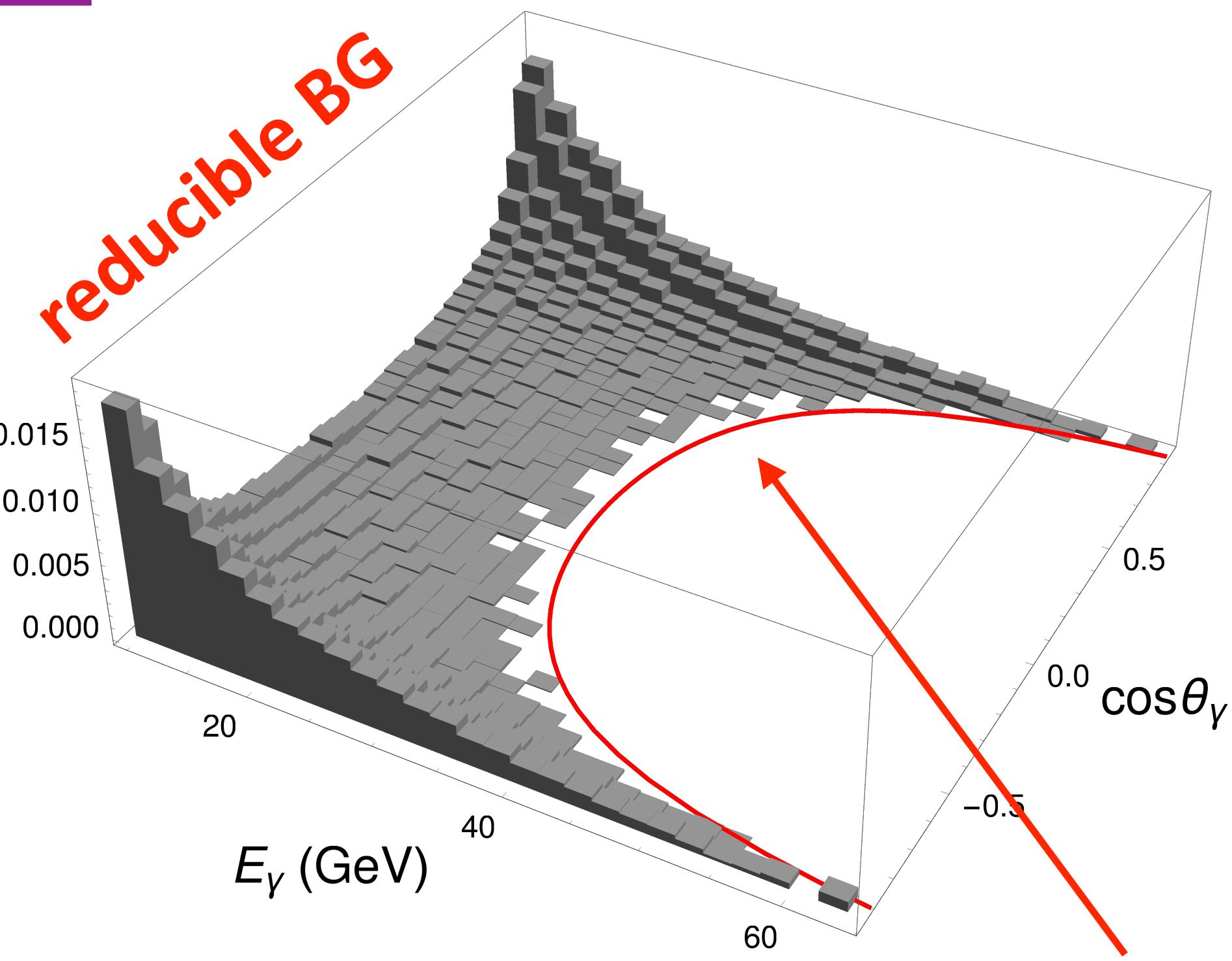
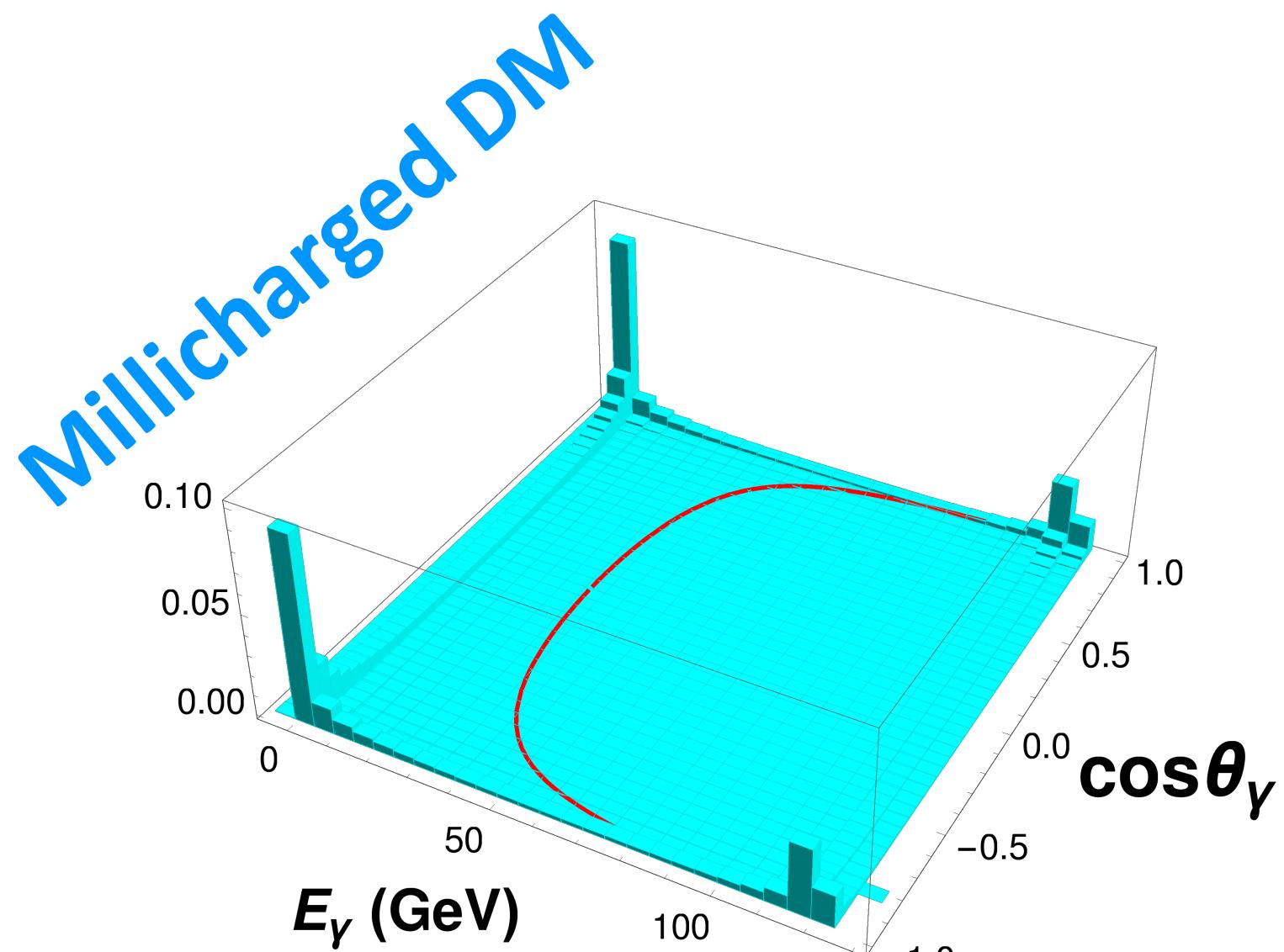
EMC energy resolution:  $\frac{\sigma(E)}{E} = \frac{10.1\%}{\sqrt{E/\text{GeV}}} \oplus 0.4\%$

EMC position resolution:  $\sim 0.1$  mm [S. Wang, thesis, 2017]

$X = \gamma$  is small  
 $X = \gamma\gamma, f\bar{f}$  are large



# Remove reducible BG



maximum photon energy

$$E_\gamma^m = \sqrt{s} \left[ 1 + \frac{\sin \theta_\gamma}{\sin \theta_b} \right]^{-1} \equiv E_B^m(\theta_\gamma)$$

$$\theta_{e^\pm} = \theta_b$$

$$|\cos \theta_b| = 0.99$$

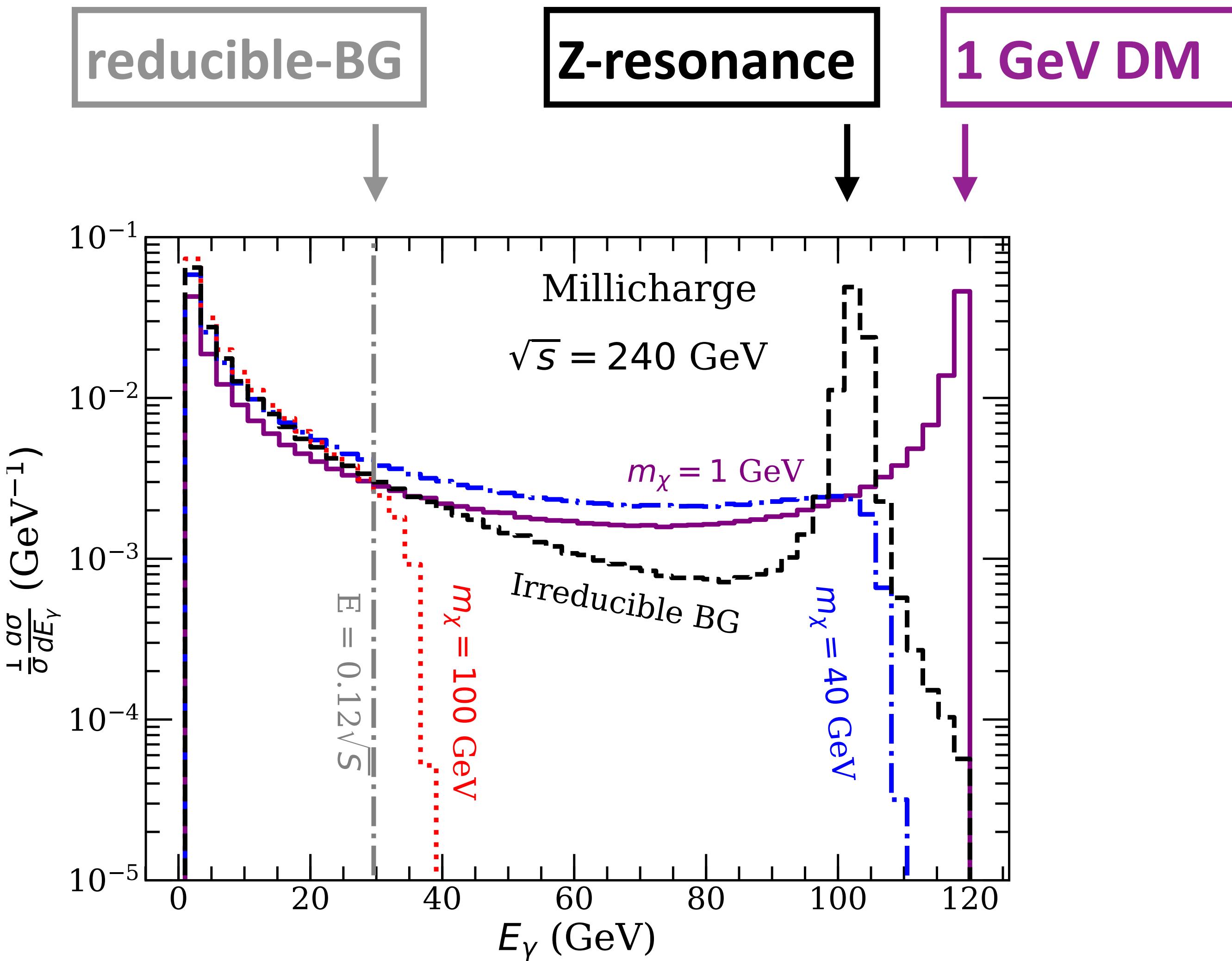
# Detector cuts

- $E_\gamma > 0.1 \text{ GeV}$  ← **EMC**
  - $|\cos \theta_\gamma| < |\cos \theta_b| = 0.99$  ← **EMC**
  - $E_\gamma < E_\chi^m = (s - 4m_\chi^2)/(2\sqrt{s})$  ← **DM**
  - veto  $E_\gamma \in (E_\gamma^Z(s) \pm 5\Gamma_\gamma^Z(s))$  ← **Z resonance**
  - $E_\gamma(\theta_\gamma) > E_B^m(\theta_\gamma) = \sqrt{s} \Big/ \left[ 1 + \frac{\sin \theta_\gamma}{\sin \theta_b} \right]$  ← **reducible BG**
- ↑  $(E_B^m)_{\min} \simeq 0.12\sqrt{s}$  when  $\theta_\gamma = \pi/2$

1

millicharged DM

# Energy distribution (H-mode)

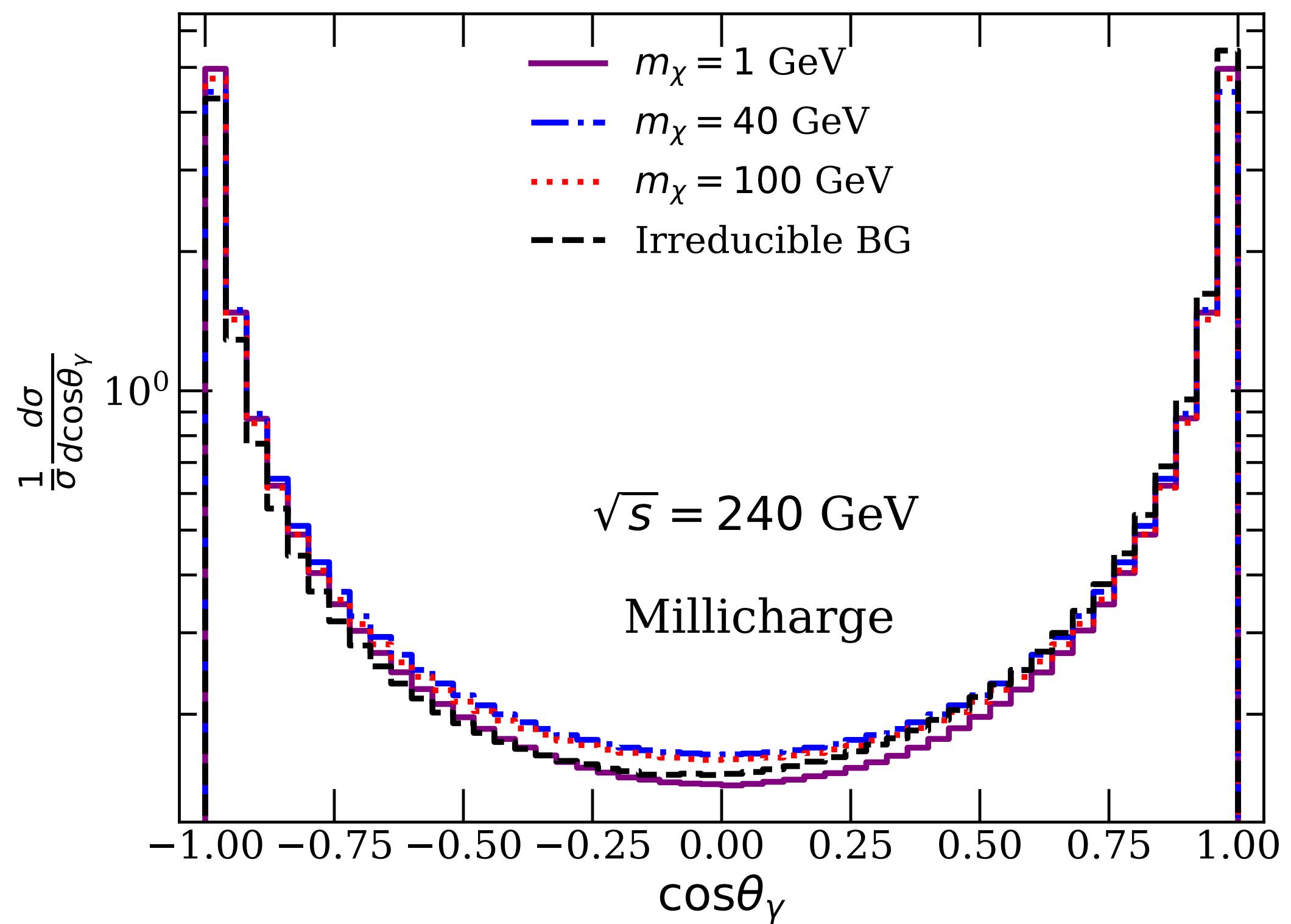


Millicharged DM with  
mass = 1, 40, 100 GeV  
&  $\epsilon = 0.01$

detector cuts  
 $E_\gamma > 0.1$  GeV &  $|z_\gamma| < 0.99$

to enhance sensitivity  
 $E_\gamma > E_\gamma^Z(s) + 5\Gamma_\gamma^Z(s)$   
for  $m_\chi < 25$  GeV

# Angular distribution (H-mode)



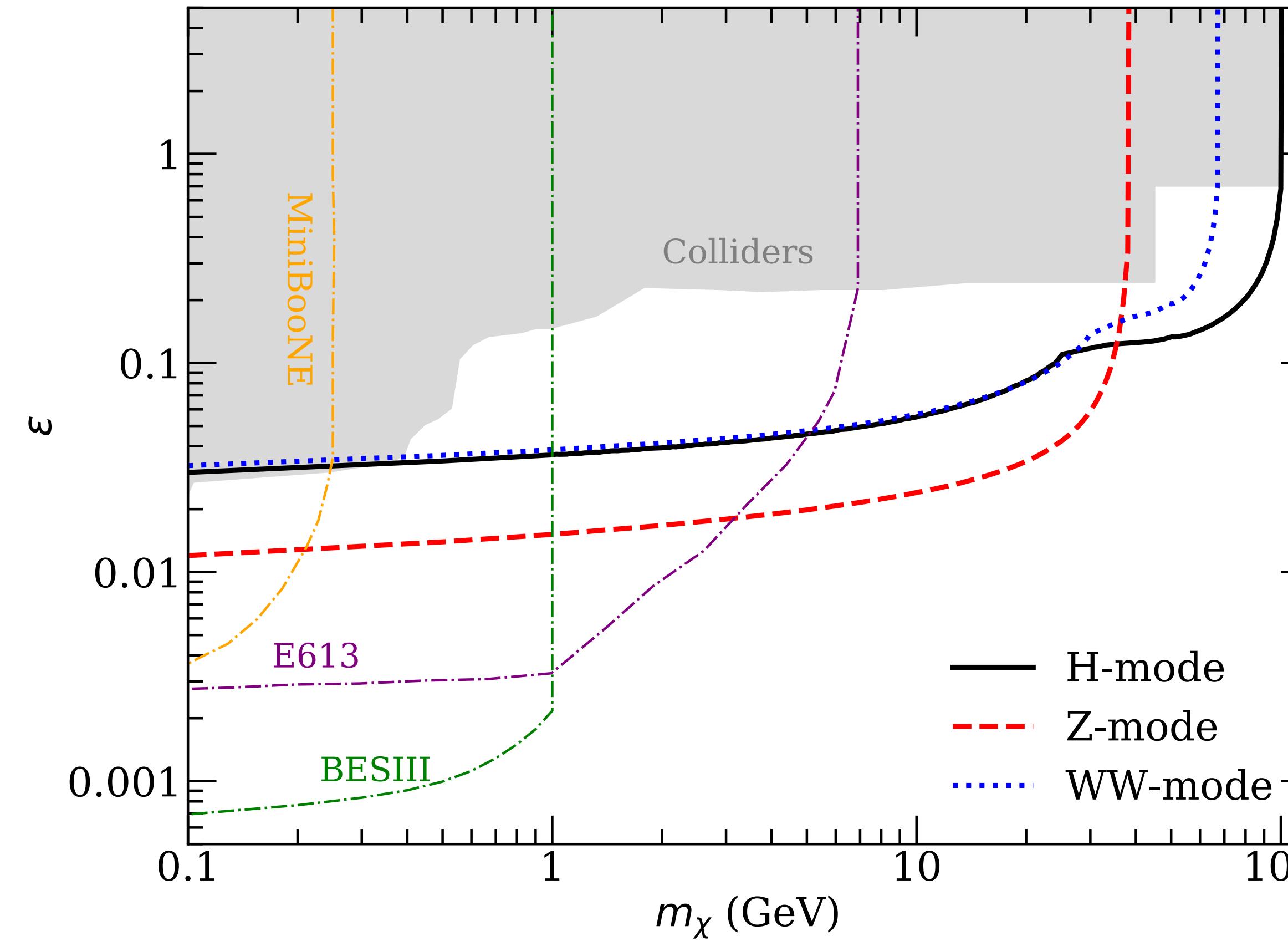
Millicharged DM with  
mass = 1, 40, 100 GeV  
&  $\varepsilon = 0.01$

detector cuts  
 $E_\gamma > 0.1 \text{ GeV} \text{ & } |z_\gamma| < 0.99$

similar distributions  $\implies$   
no further angular cut

# CEPC sensitivity on millicharged DM

[ZL, Y.-H. Xu, and Y. Zhang, 1903.12114]



95% CL bound  
 $S/\sqrt{B} = 2$

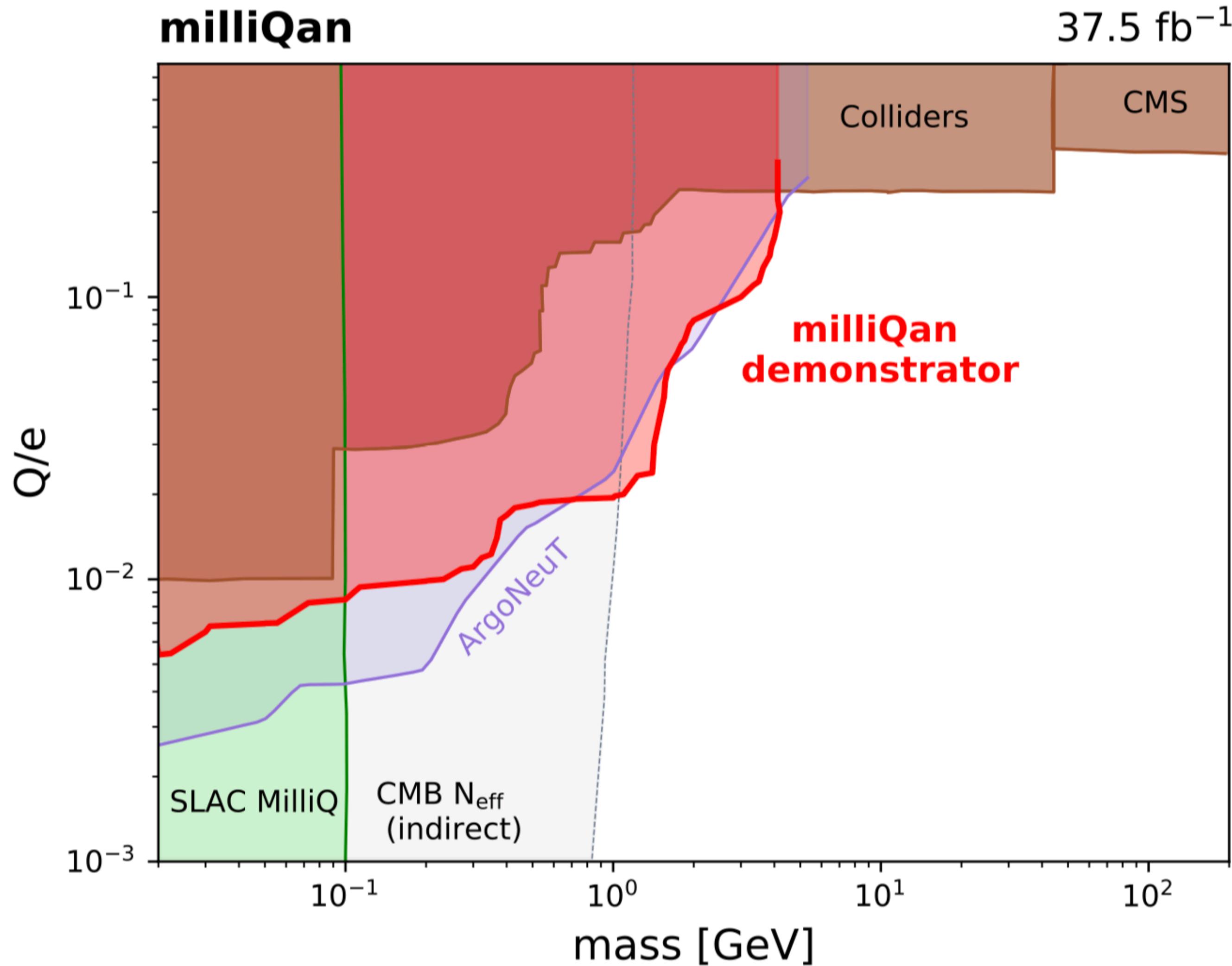
millicharged particle  $\chi$

$e \epsilon A_\mu^\gamma \bar{\chi} \gamma^\mu \chi$

Z-mode can probe  
 $\epsilon \sim 0.02$  for  $m = 5$  GeV

CEPC can probe unexplored millicharge values

# Recent limits from ArgoNeuT & milliQan demonstrator



[ArgoNeuT, 1911.07996]

ArgoNeuT  $10^{20}$  POT

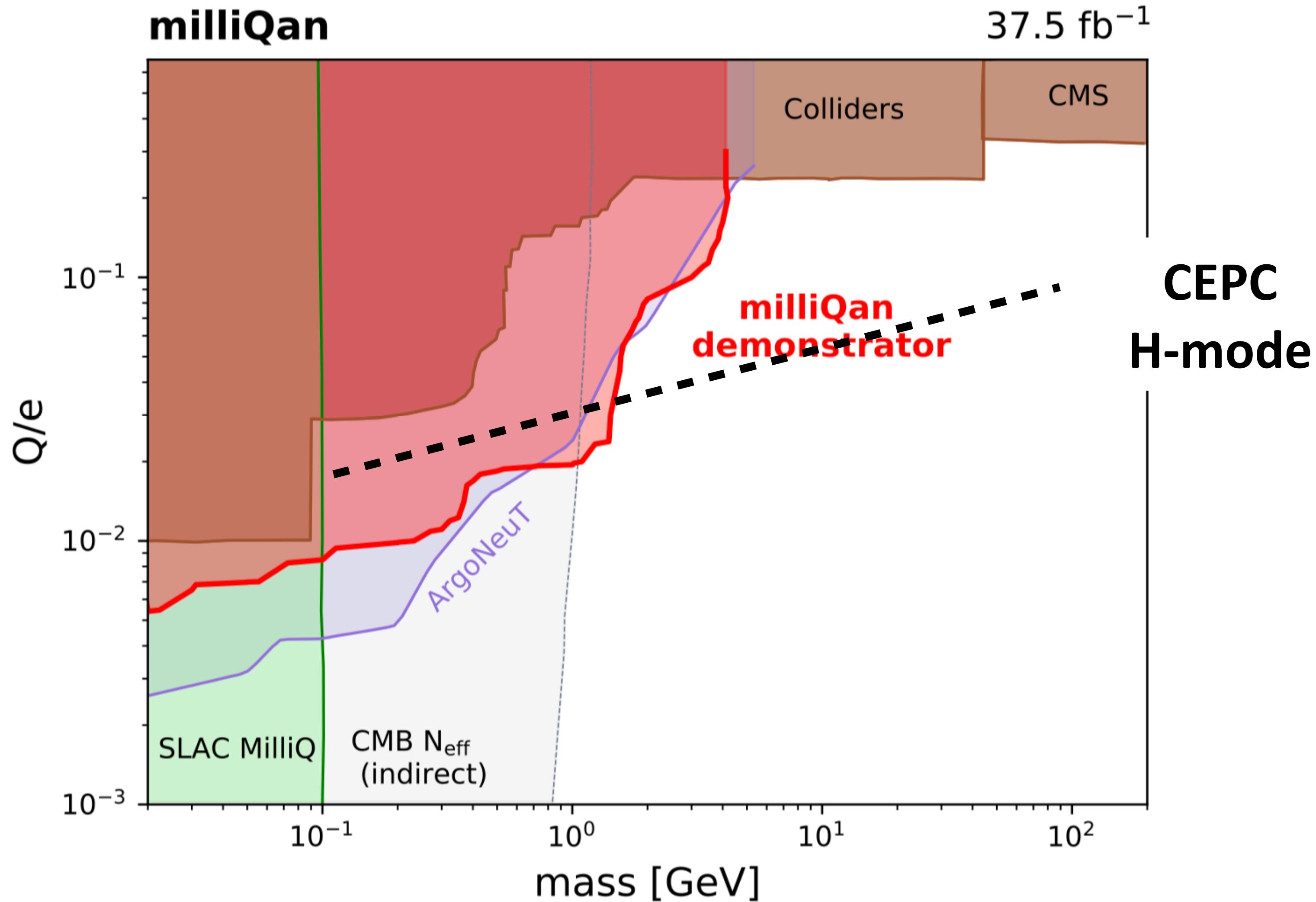
[milliQan demonstrator, 2005.06518]

37/fb in 2018

also limits from  
neutrino experiments

[Magill et al., 1806.03310]

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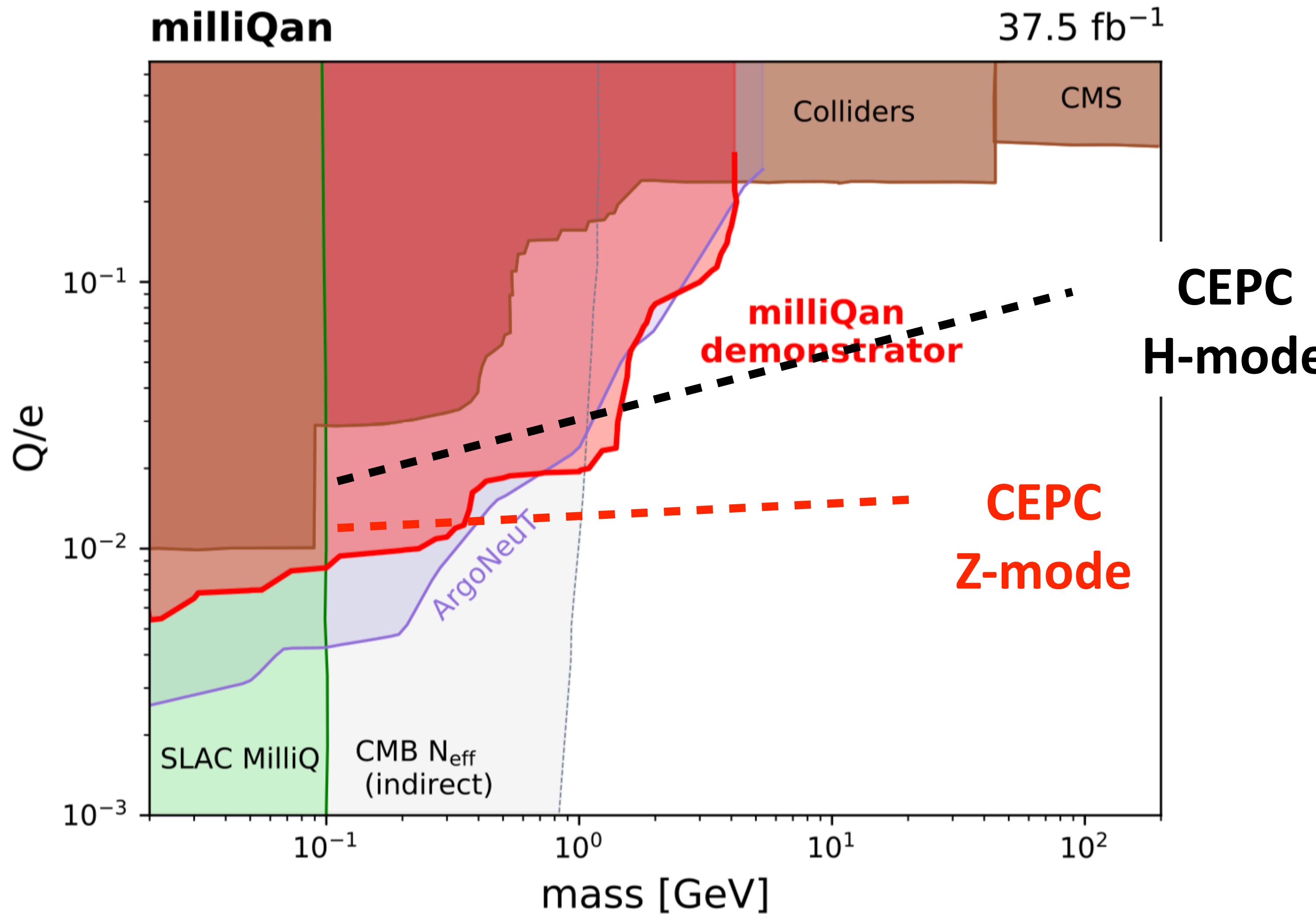
[milliQan demonstrator, 2005.06518]

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# Recent limits from ArgoNeuT & MilliQan demonstrator



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2

## Vector portal DM

# Vector-portal DM

$$\mathcal{L} = Z'_\mu \bar{\chi} \gamma^\mu (g_V^\chi - g_A^\chi \gamma_5) \chi + Z'_\mu \bar{f} \gamma^\mu (g_V^f - g_A^f \gamma_5) f$$

Z' mass = **150** GeV

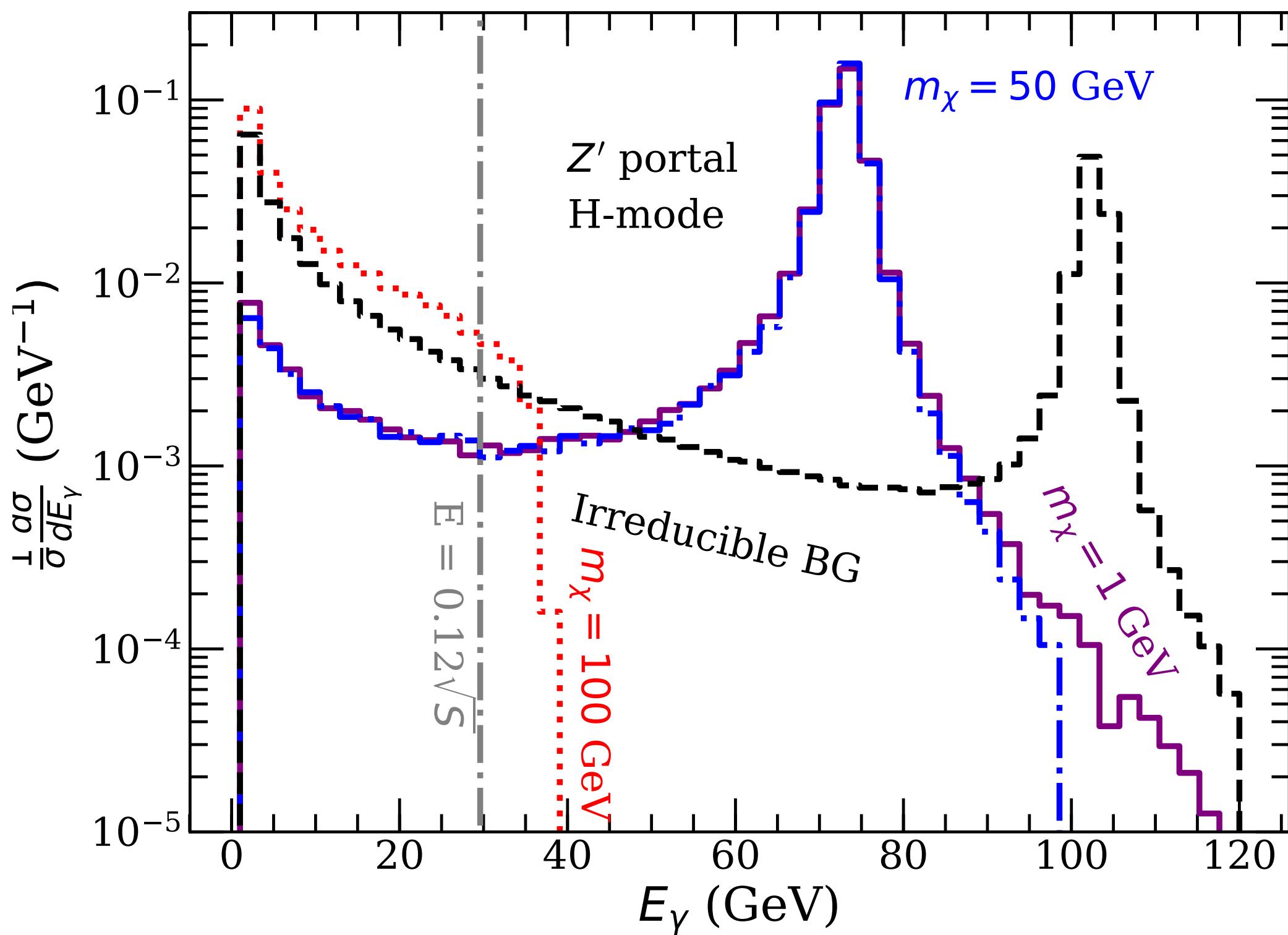
2 scenarios

- (1) **vector** coupling only
- (2) **axial-vector** coupling only

2 channels

- (1) monophoton (**invisible**)
- (2) di-muon (**visible**)

# Energy distribution (Z' portal DM)



$$E_\gamma > 0.1 \text{ GeV} \text{ & } |z_\gamma| < 0.99$$

$$m_{Z'} = 150 \text{ GeV}$$

“Z’ resonance”

$$E_\gamma = \frac{s - m_{Z'}^2}{2\sqrt{s}}$$

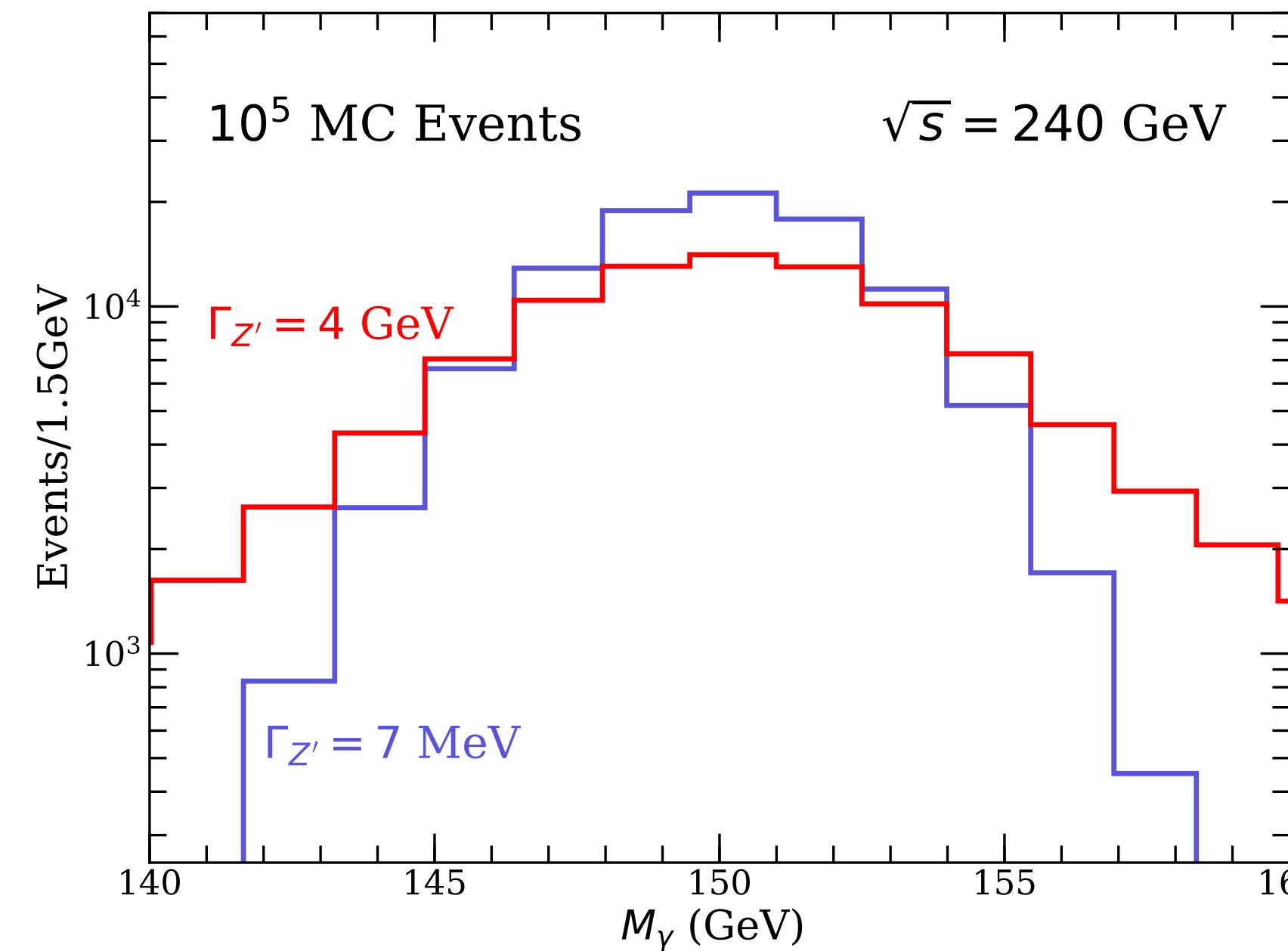
reconstruct Z’

$$M_\gamma = \sqrt{s - 2\sqrt{s}E_\gamma}$$

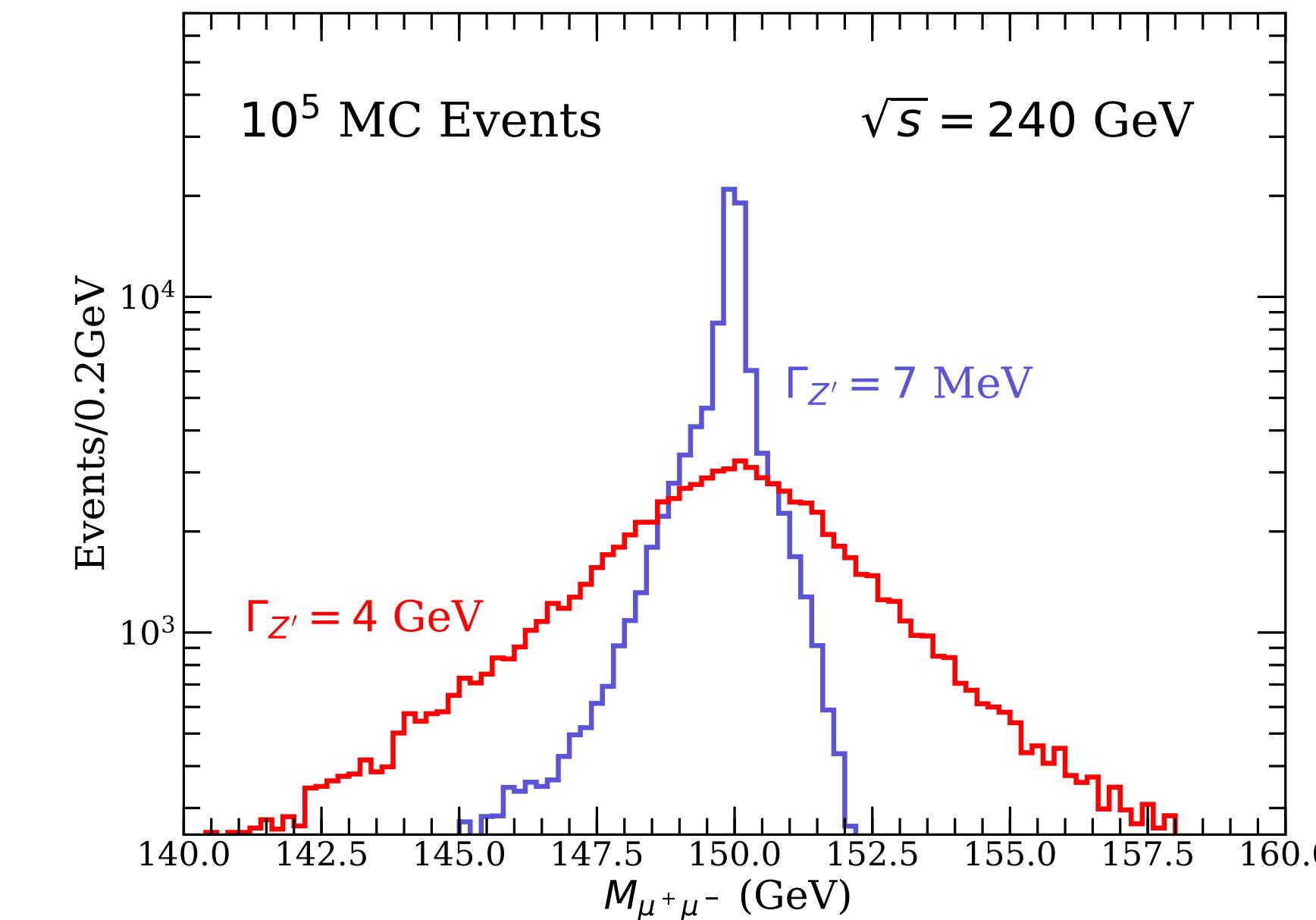
$$147 < M_\gamma/\text{GeV} < 153$$

# Invisible versus visible

## invisible: monophoton



## visible: dimuon



**dimuon**

$$\frac{\delta p_T}{p_T} = \frac{p_T}{10^5 \text{ GeV}} \bigoplus 0.1\% \text{ for } |\eta| < 1.0$$

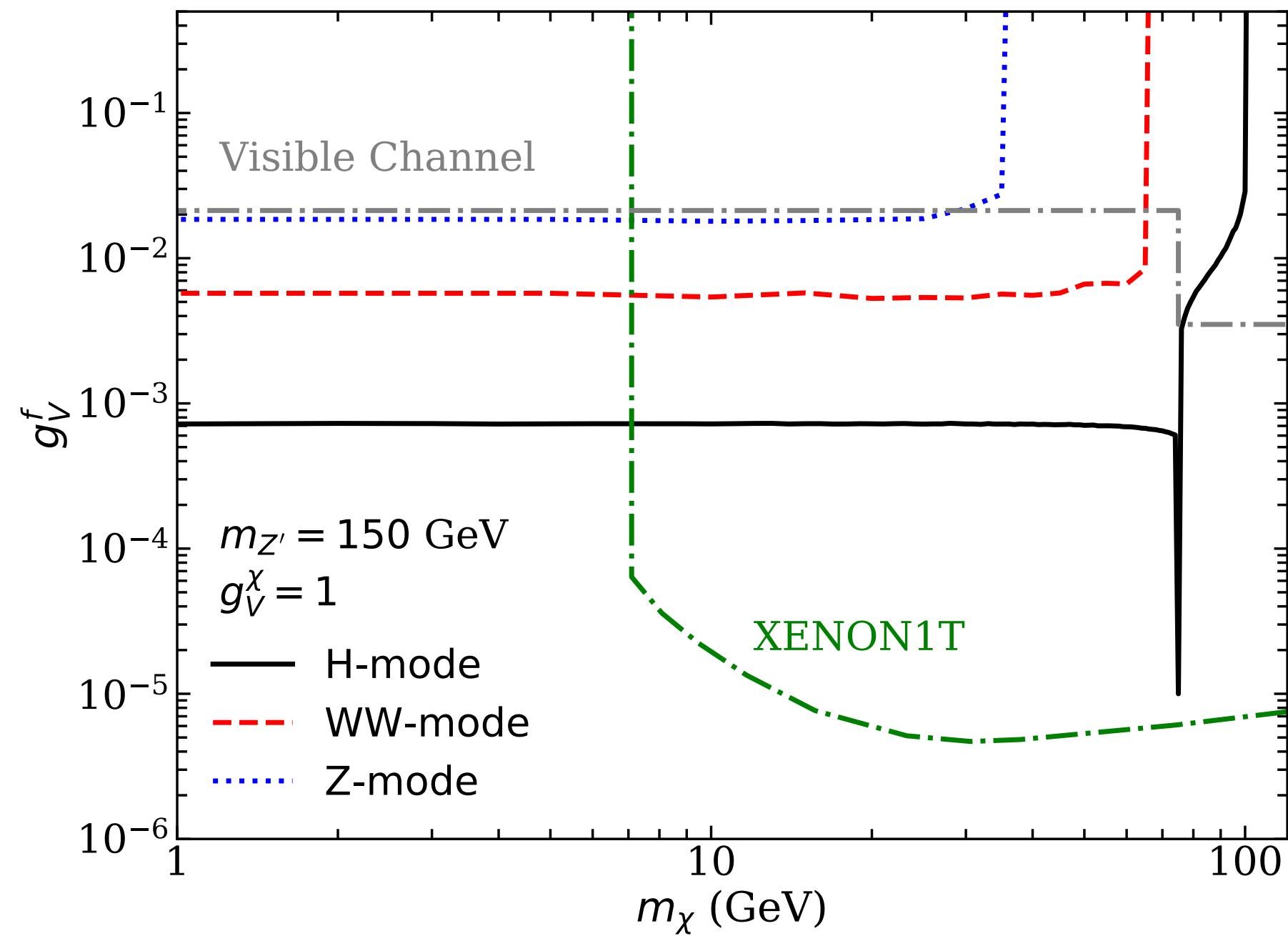
& 10 times larger for  $1.0 < |\eta| < 3.0$

[see also, e.g. Karliner, Low, Rosner, Wang, 1503.07209]

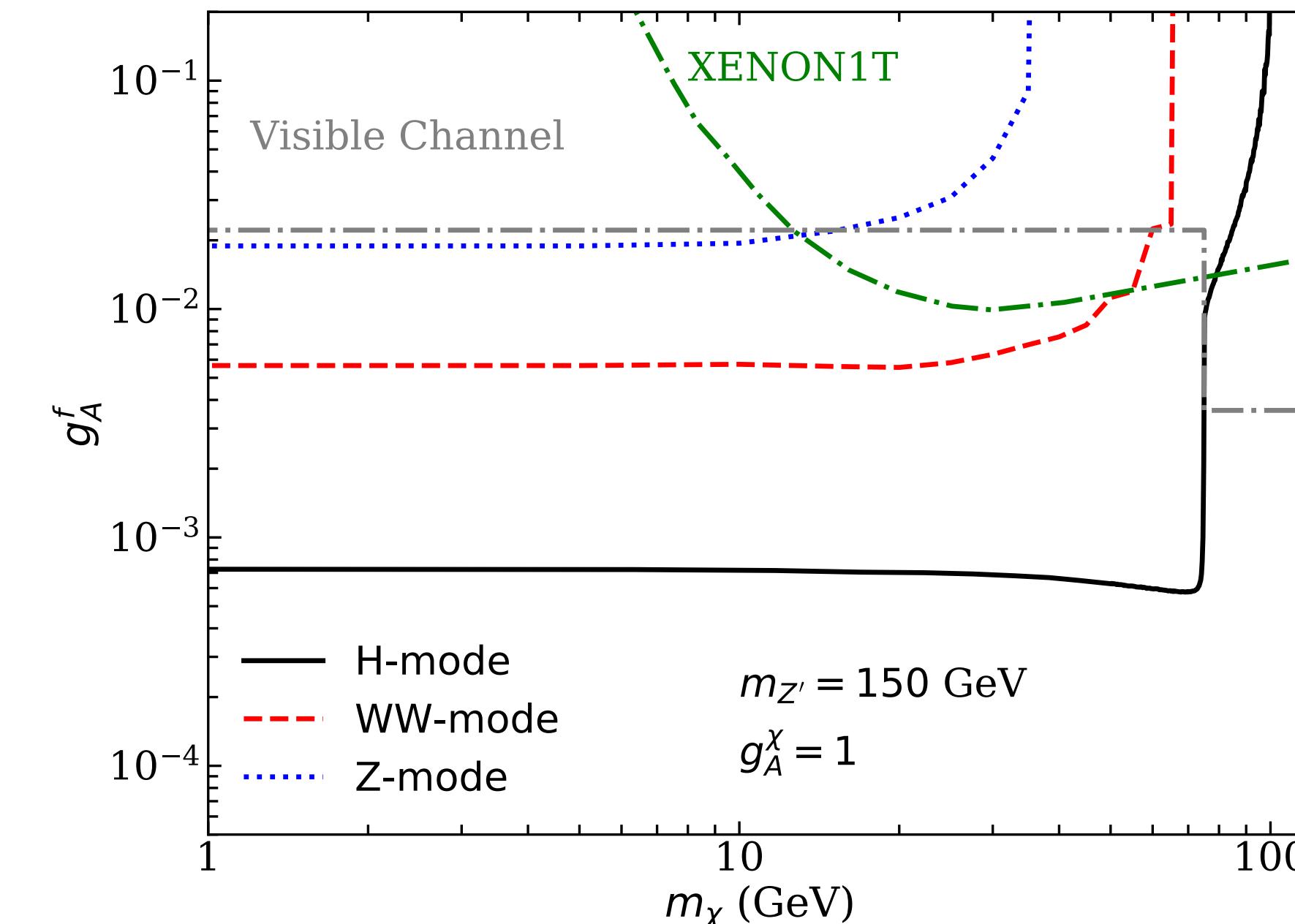
# CEPC sensitivity to Z'-portal DM

[ZL, Y.-H. Xu, and Y. Zhang, 1903.12114]

vector only



axial-vector only



strong coupling to  $\chi$ :  $g \sim 1$

DMDD: assume universal coupling to SM fermions

3

## DM with EFT interactions

# DM models with EFT interactions

**vector**

$$\frac{1}{\Lambda_V^2} \bar{\chi} \gamma_\mu \chi \bar{\ell} \gamma^\mu \ell$$

[Fox, Harnik, Kopp, Tsai, 1103.0240]

[Chae, Perelstein, 1211.4008]

**axial-vector**

$$\frac{1}{\Lambda_A^2} \bar{\chi} \gamma_\mu \gamma_5 \chi \bar{\ell} \gamma^\mu \gamma_5 \ell$$

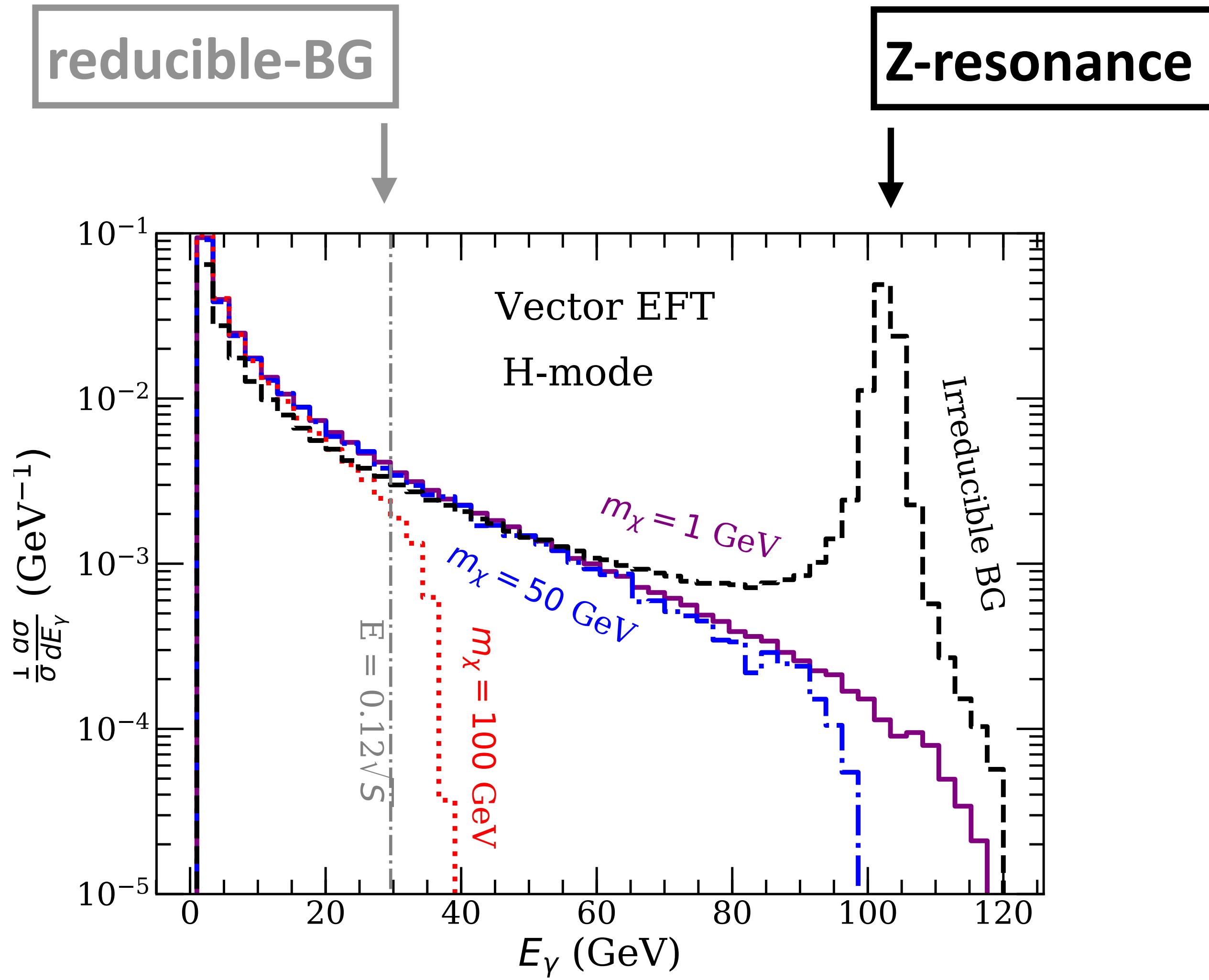
**scalar (s-channel)**

$$\frac{1}{\Lambda_s^2} \bar{\chi} \chi \bar{\ell} \ell$$

**scalar (t-channel)**

$$\frac{1}{\Lambda_t^2} \bar{\chi} \ell \bar{\ell} \chi$$

# Energy distribution (EFT DM)

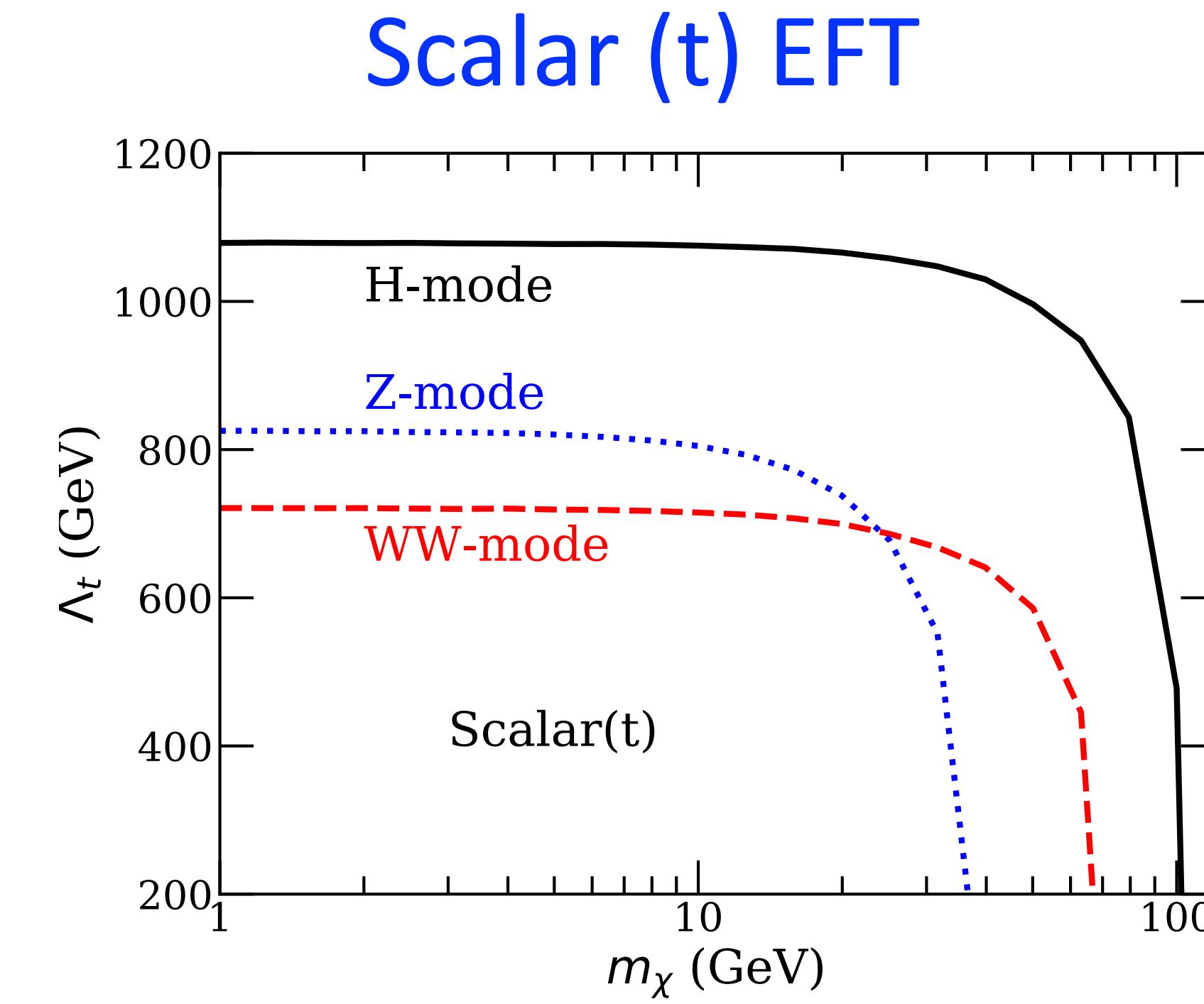
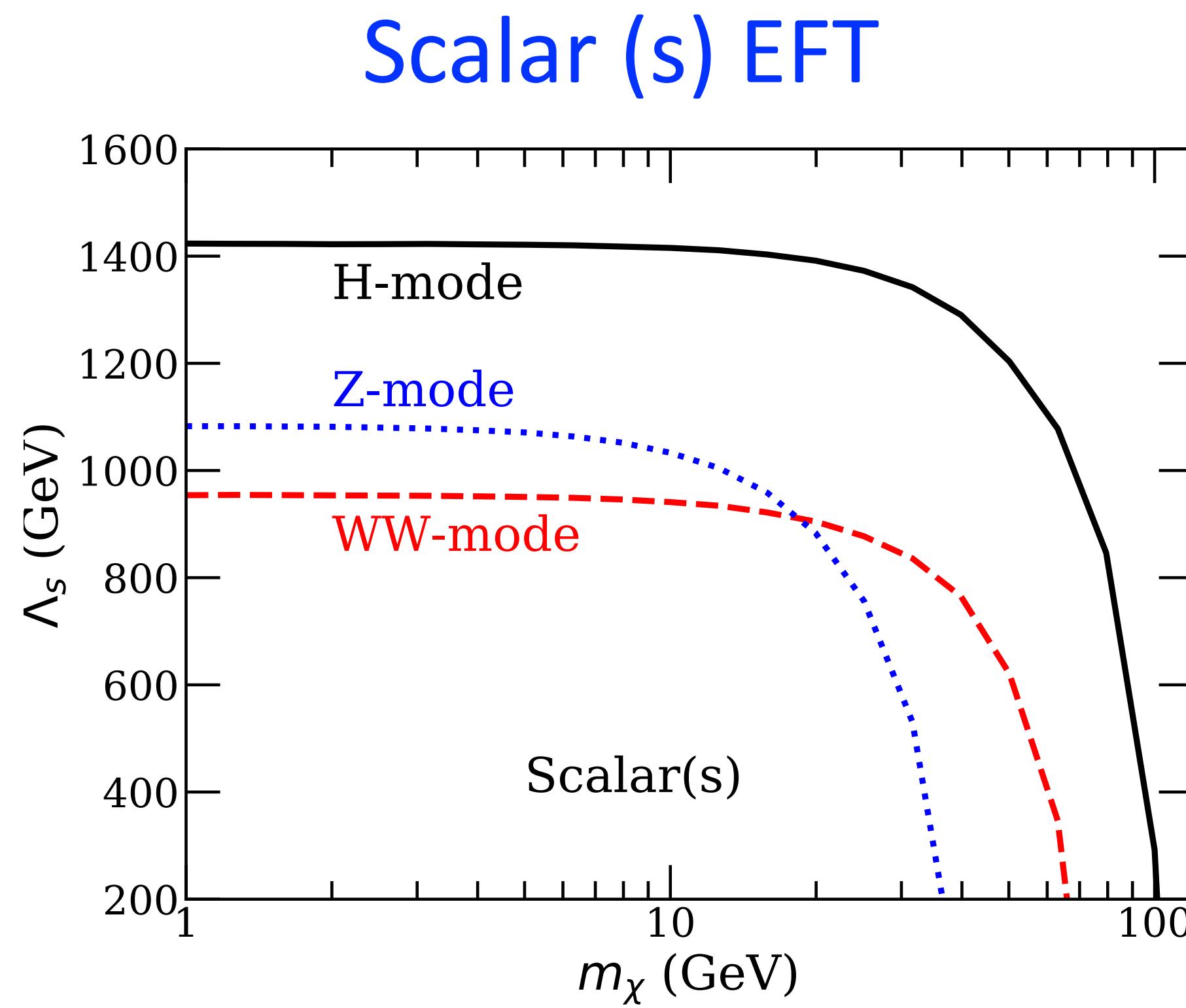


Vector EFT

detector cuts

$$E_\gamma > 0.1 \text{ GeV} \text{ & } |z_\gamma| < 0.99$$

# CEPC sensitivity to EFT DM (scalar cases)

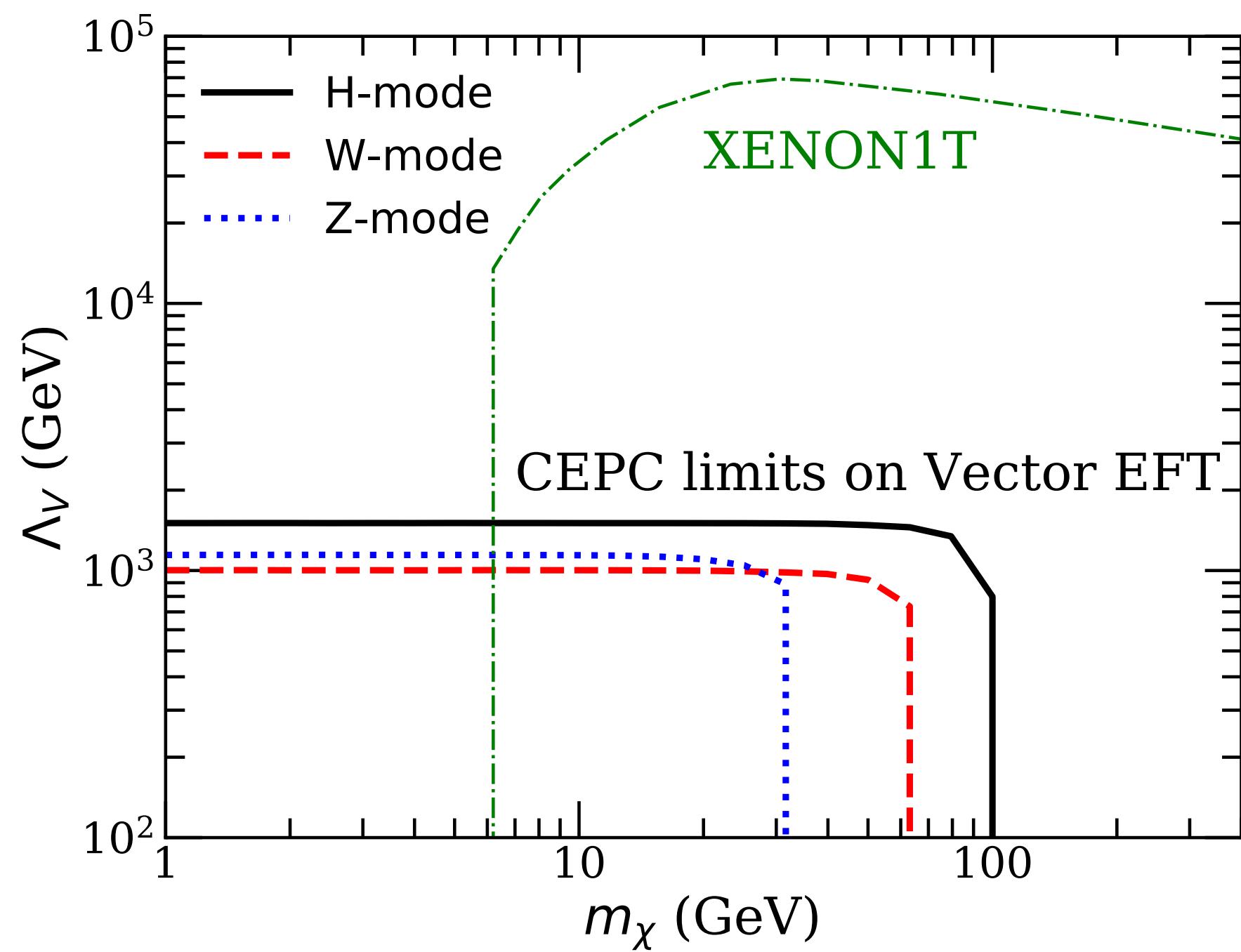


[ZL, Y.-H. Xu, and Y. Zhang, 1903.12114]

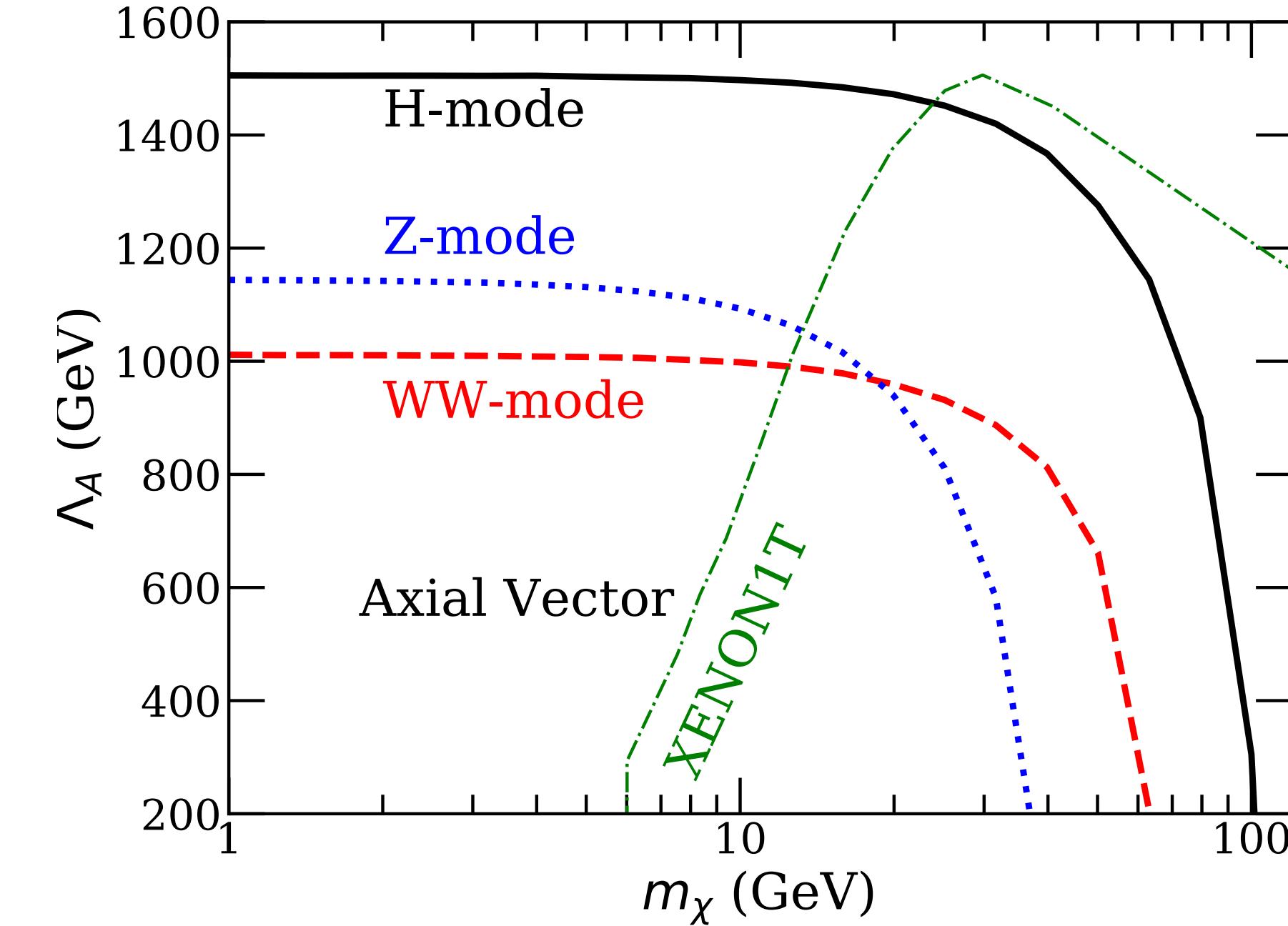
# CEPC sensitivity to EFT DM (V & A cases)

[ZL, Y.-H. Xu, and Y. Zhang, 1903.12114]

## Vector EFT



## Axial-vector EFT



DMDD: assume universal coupling to SM fermions

# Summary

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To optimize the CEPC sensitivity to different dark matter models, we propose a list of detector cuts to suppress the background

CEPC can probe millicharge DM that is currently unexplored

CEPC can probe the parameter space of vector-portal DM models and EFT DM models that are unconstrained by DMDD

[[arXiv:1903.12114](https://arxiv.org/abs/1903.12114), ZL, Yong-Heng Xu, Yu Zhang]

# slides for Millicharged DM

# Electric charge is quantized in SM

	mass →	charge →	spin →	
QUARKS				
u	$\approx 2.3 \text{ MeV}/c^2$	2/3	1/2	up
c	$\approx 1.275 \text{ GeV}/c^2$	2/3	1/2	charm
t	$\approx 173.07 \text{ GeV}/c^2$	2/3	1/2	top
g	0	0	1	gluon
H	$\approx 126 \text{ GeV}/c^2$	0	0	Higgs boson
d	$\approx 4.8 \text{ MeV}/c^2$	-1/3	1/2	down
s	$\approx 95 \text{ MeV}/c^2$	-1/3	1/2	strange
b	$\approx 4.18 \text{ GeV}/c^2$	-1/3	1/2	bottom
$\gamma$	0	0	1	photon
e	$0.511 \text{ MeV}/c^2$	-1	1/2	electron
$\mu$	$105.7 \text{ MeV}/c^2$	-1	1/2	muon
$\tau$	$1.777 \text{ GeV}/c^2$	-1	1/2	tau
Z	$91.2 \text{ GeV}/c^2$	0	1	Z boson
$\nu_e$	$< 2.2 \text{ eV}/c^2$	0	1/2	electron neutrino
$\nu_\mu$	$< 0.17 \text{ MeV}/c^2$	0	1/2	muon neutrino
$\nu_\tau$	$< 15.5 \text{ MeV}/c^2$	0	1/2	tau neutrino
GAUGE BOSONS				
W	$80.4 \text{ GeV}/c^2$	$\pm 1$	1	W boson

integer multiple of that of d-quark

Stringent experimental limits, e.g.

$$Q_p - Q_e < (0.8 \pm 0.8) \times 10^{-21} e$$

[Marinelli et al. 1984]

# But BSM particles can naturally possess a millicharge

Kinetic mixing

[Holdom, Phys.Lett. 166B, 196 (1986)]

[Foot & He, Phys. Lett. 267B, 509 (1991)]

$$\mathcal{L} = -\frac{\delta}{2} B_{\mu\nu} X^{\mu\nu} + X_\mu J_X^\mu$$

Stueckelberg mass mixing

[Kors & Nath hep-ph/0402047]

[Feldman, ZL, Nath, hep-ph/0603039]

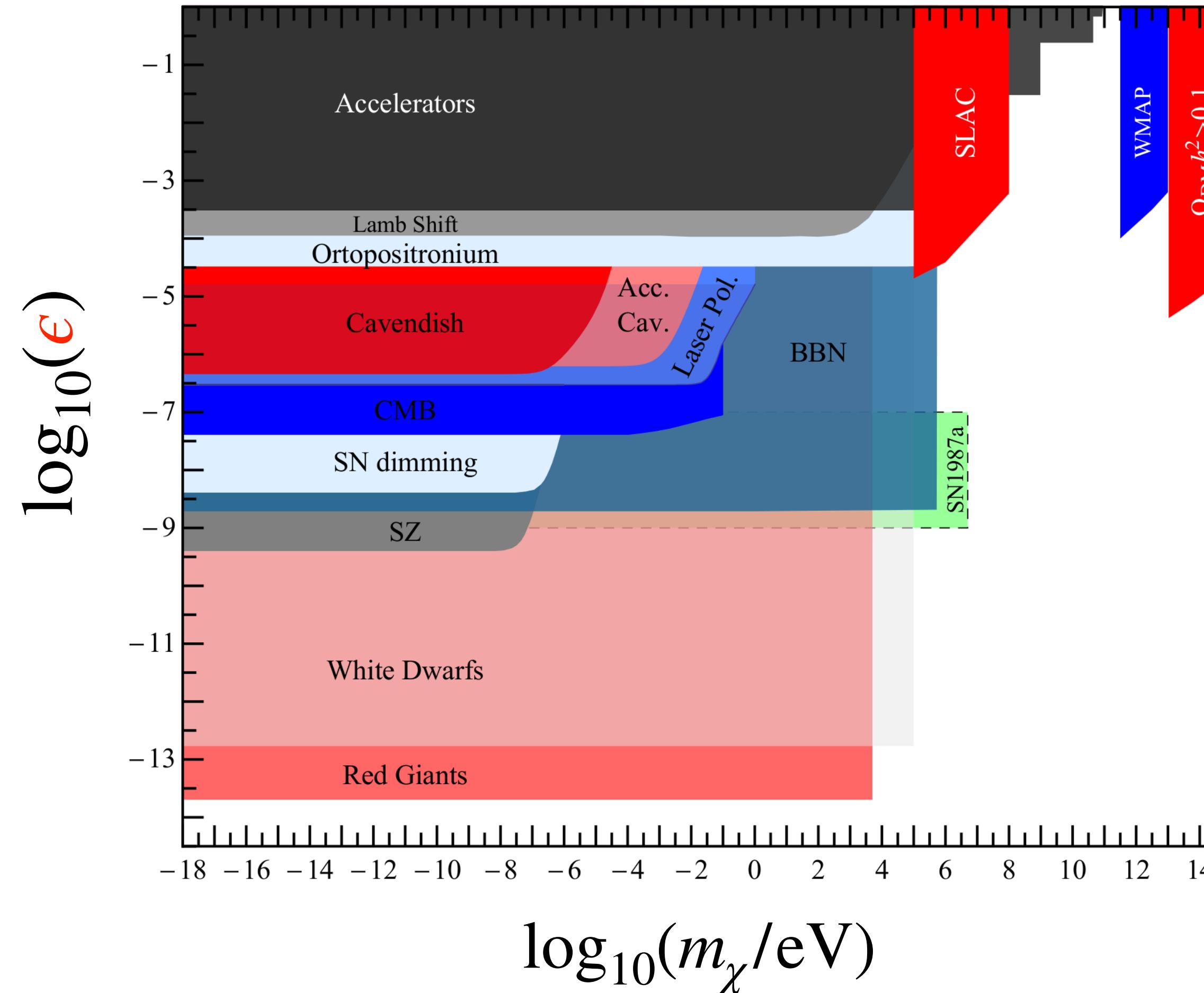
$$\mathcal{L} = -\frac{1}{2}(\partial_\mu \sigma + M_1 X_\mu + M_2 B_\mu)^2 + X_\mu J_X^\mu$$

Both kinetic mixing & Stueckelberg mass mixing

[Feldman, ZL, Nath, hep-ph/0702123, 331 cites]

# Constraints on millicharged particles in BSM

[Jaeckel & Ringwald, 1002.0329]

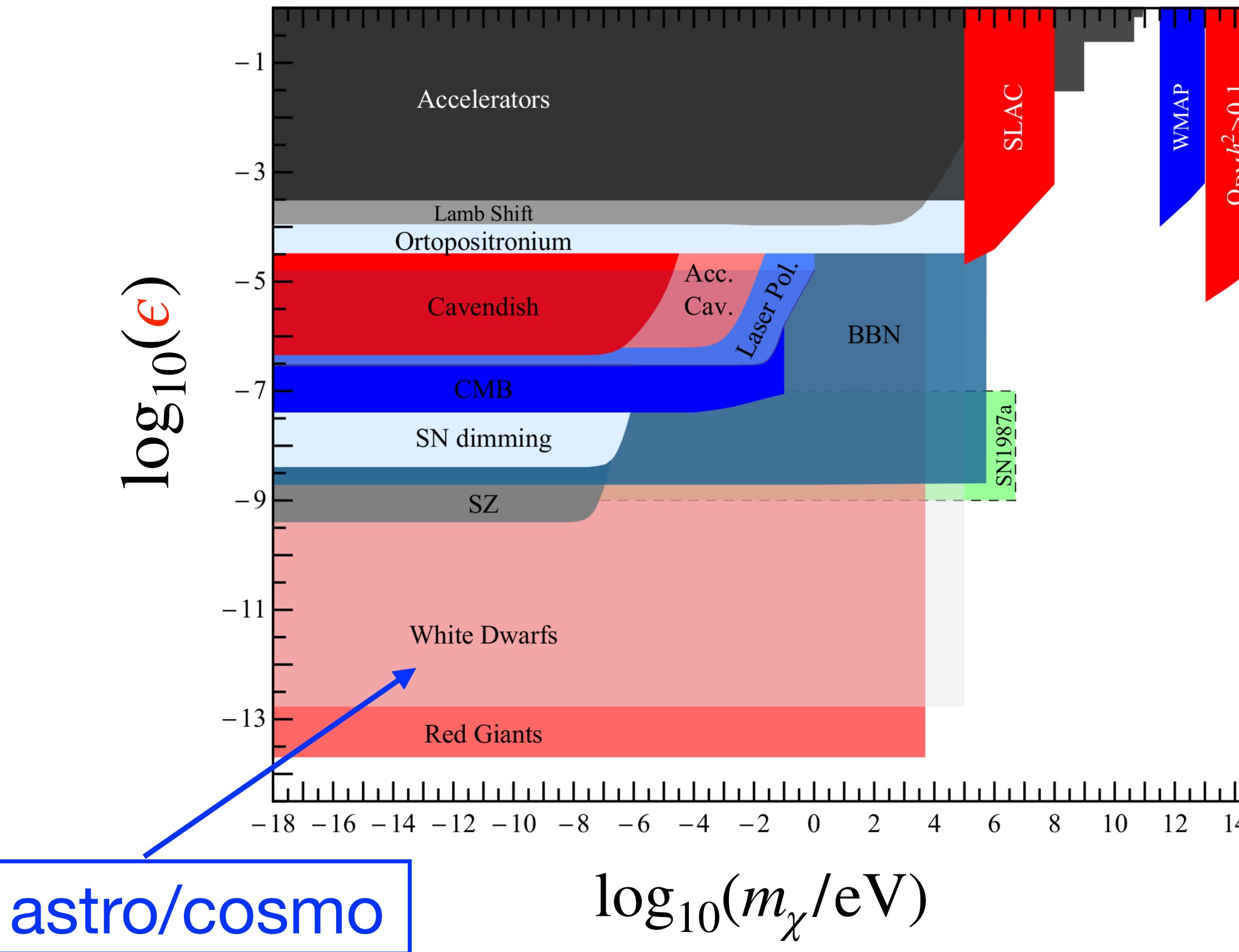


millicharged particle  $\chi$

$$e \, \epsilon \, A_\mu^\gamma \bar{\chi} \gamma^\mu \chi$$

# Constraints on millicharged particles in BSM

[Jaeckel & Ringwald, 1002.0329]

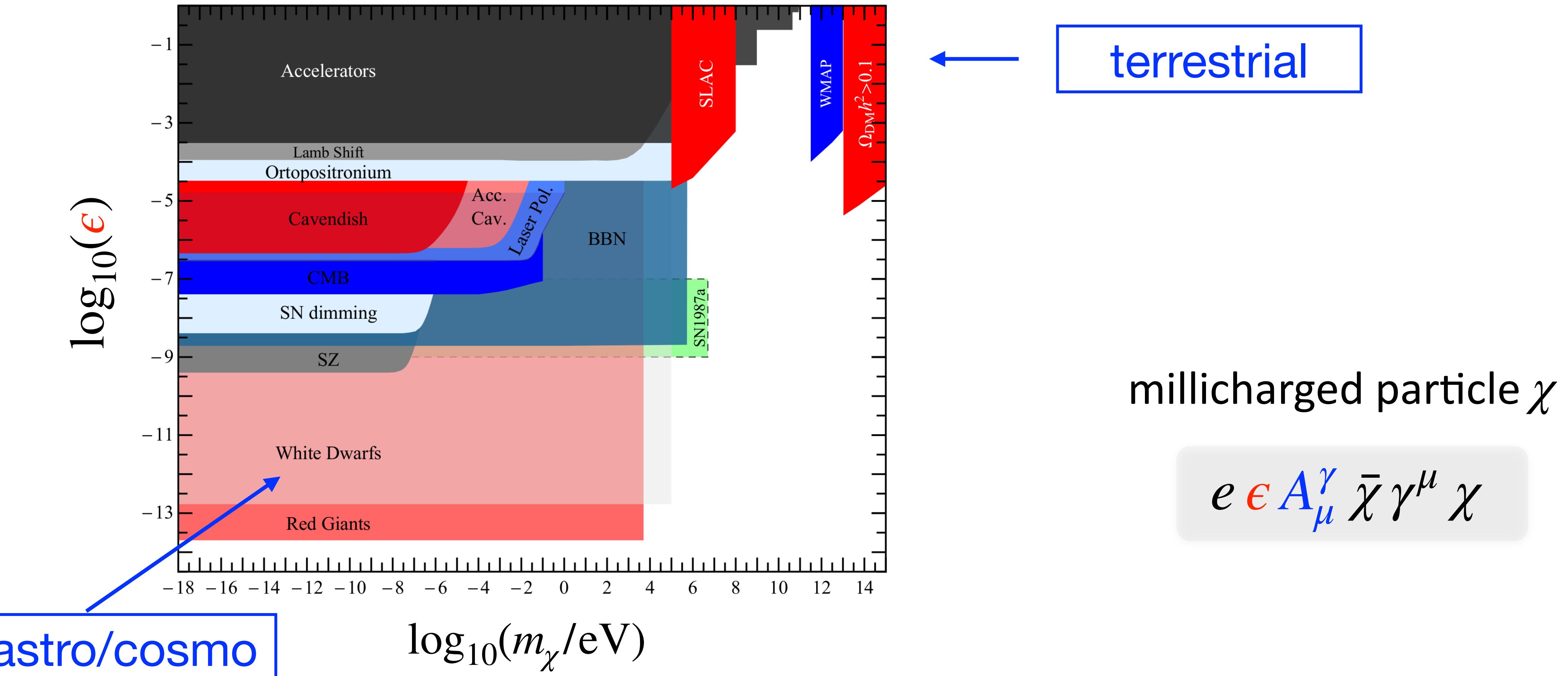


millicharged particle  $\chi$

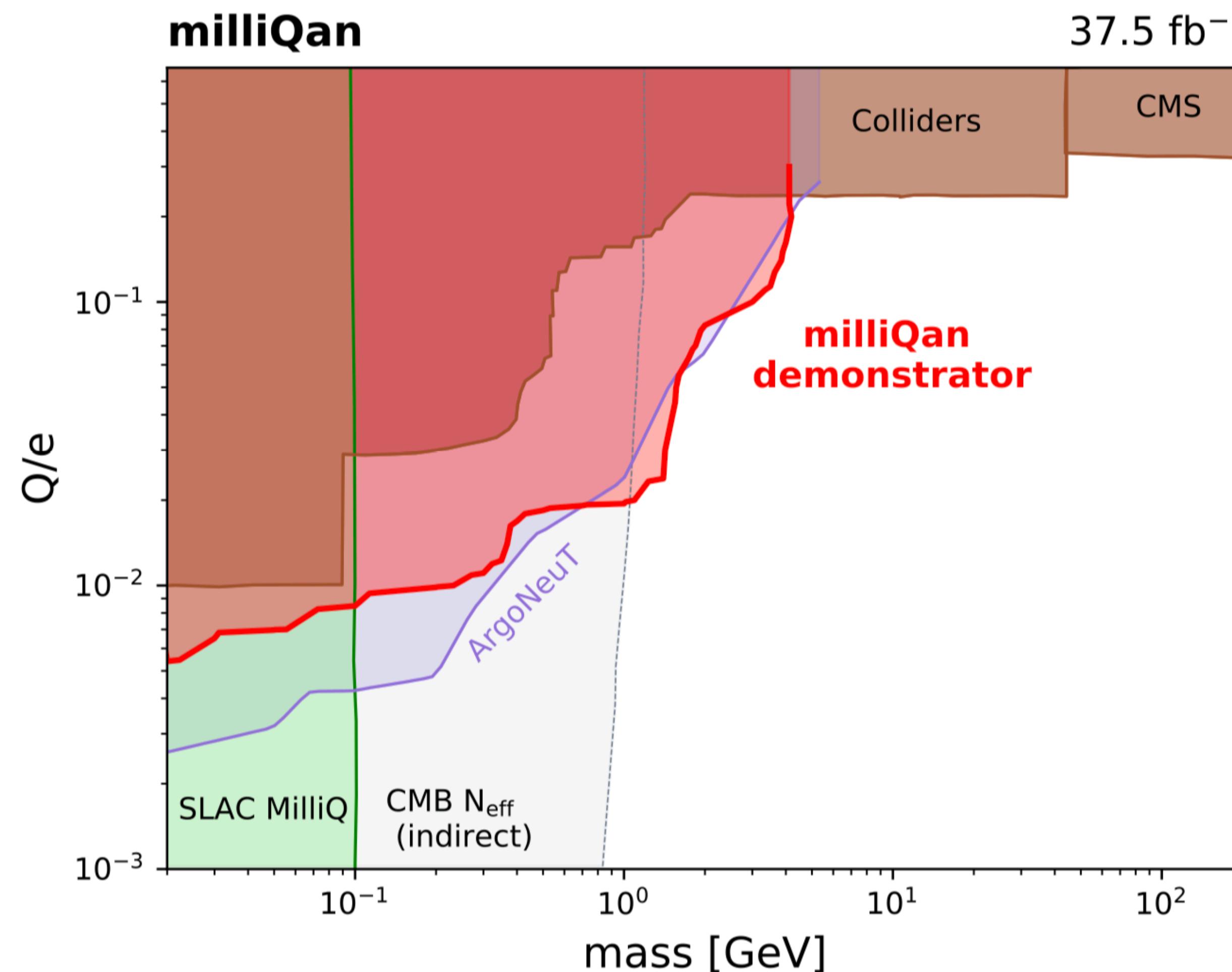
$$e \epsilon A_\mu^\gamma \bar{\chi} \gamma^\mu \chi$$

# Constraints on millicharged particles in BSM

[Jaeckel & Ringwald, 1002.0329]



# Recent limits from ArgoNeuT & milliQan demonstrator



[ArgoNeuT, 1911.07996]

ArgoNeuT  $10^{20}$  POT

[milliQan demonstrator, 2005.06518]

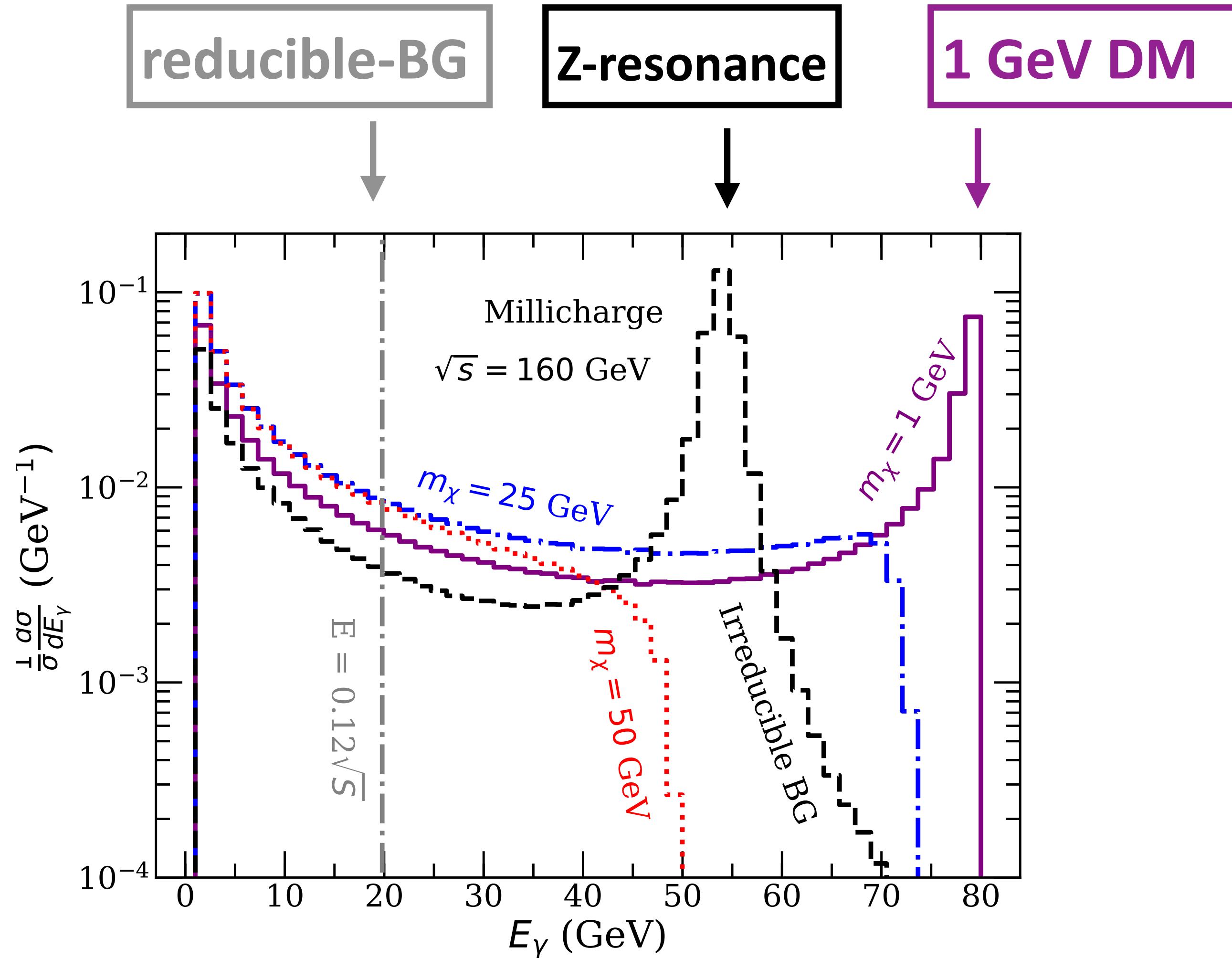
37/fb in 2018

also limits from  
neutrino experiments

[Magill et al., 1806.03310]

# Energy distributions for Millicharged DM

# Energy distribution (WW-mode)



Millicharged DM  
w/  $m = 1, 25, 50$   
GeV &  $\epsilon = 0.01$

detector cuts

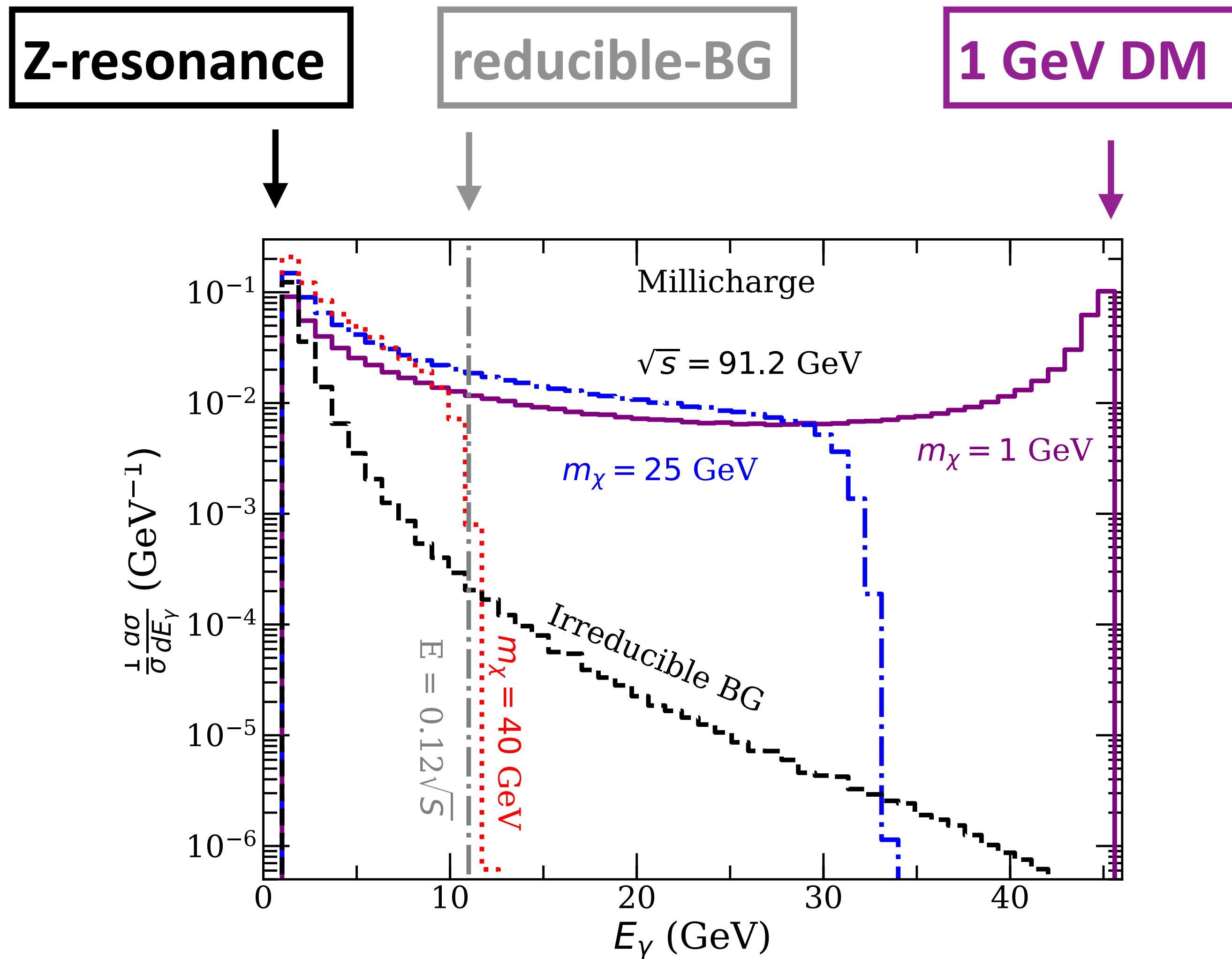
$$E_\gamma > 0.1 \text{ GeV} \text{ & } |z_\gamma| < 0.99$$

to enhance sensitivity

$$E_\gamma > E_\gamma^Z(s) + 5\Gamma_\gamma^Z(s)$$

$$\text{for } m_\chi < 30 \text{ GeV}$$

# Energy distribution (Z-mode)



Millicharged DM w/  
 $m = 1, 25, 40 \text{ GeV}$  &  
 $\epsilon = 0.01$

detector cuts

$E_\gamma > 0.1 \text{ GeV}$  &  $|z_\gamma| < 0.99$