

The Multi-lepton Anomalies at the LHC and implications

Bruce Mellado

The Institute for Collider Particle Physics and iThemba LABS



INSTITUTE FOR
COLLIDER
PARTICLE
PHYSICS



UNIVERSITY OF THE WITWATERSRAND

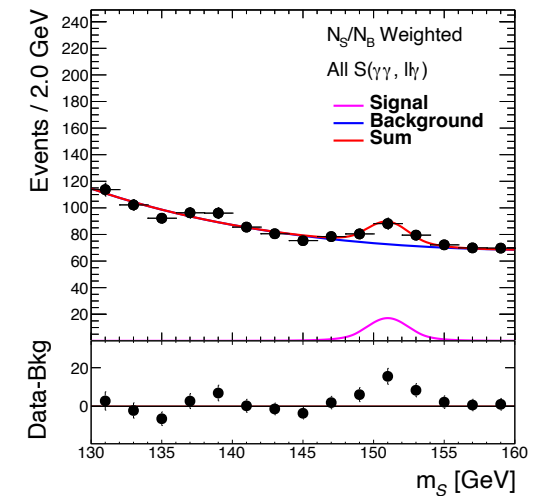
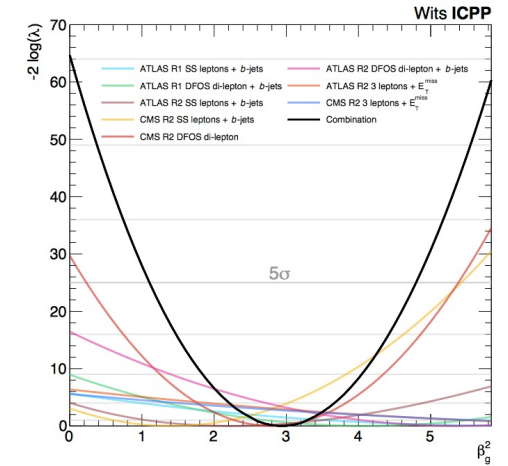


核物理与核技术国家重点实验室(北京大学)
State Key Laboratory of Nuclear Physics and Technology, Peking University

Higgs Potential 2022, 26/07/22

Outline

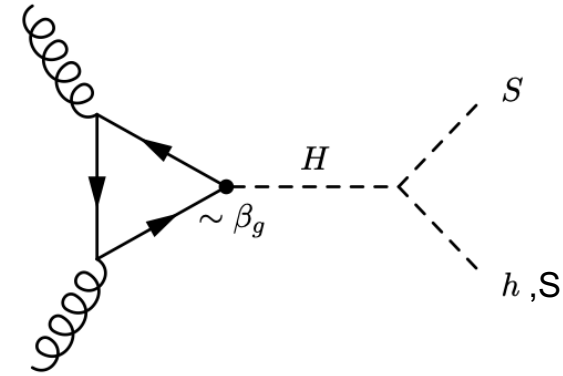
- The simplified model
- The multilepton anomalies
 - Methodology
 - The anatomy
- Impact on Higgs physics
- A possible candidate of S



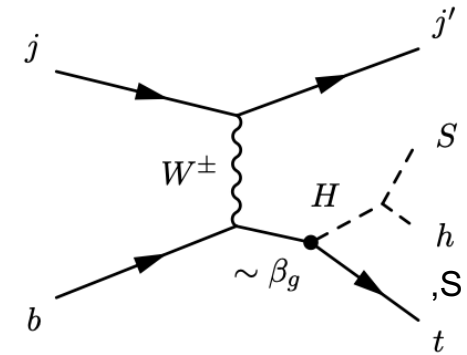
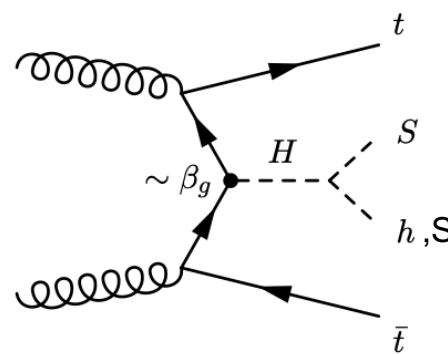
The Simplified Model and 2HDM+S

The simplified Model (from Run I)

- 1. The starting point of the hypothesis is the existence of a boson, H, that contains Higgs-like interactions, with a mass in the range 250-280 GeV**
- 2. In order to avoid large quartic couplings, incorporate a mediator scalar, S, that interacts with the SM and Dark Matter.**
- 3. Dominance of $H \rightarrow Sh, SS$ decay over other decays**



(a) Gluon fusion (ggF).



$$\mathcal{L}_{\text{int}} \supset -\beta_g \frac{m_t}{v} t\bar{t}H + \beta_V \frac{m_V^2}{v} g_{\mu\nu} V^\mu V^\nu H$$

$$\mathcal{L}_{HhS} = -\frac{1}{2} v \left[\lambda_{hhS} hhS + \lambda_{hSS} hSS + \lambda_{HHS} HHS + \lambda_{HSS} HSS + \lambda_{HhS} HhS \right],$$

The Decays of H

- In the general case, H can have couplings as those displayed by a Higgs boson in addition to decays involving the intermediate scalar and Dark Matter

$$H \rightarrow WW, ZZ, q\bar{q}, gg, Z\gamma, \gamma\gamma, \chi\chi$$
$$+ H \rightarrow SS, Sh, hh$$

Dominant decays

Diboson decay

$$H \rightarrow h(+X), S(+X)$$

The 2HDM+S

Eur. Phys. J. C (2016) 76:580

Introduce singlet real scalar, S.

2HDM potential, $\mathcal{V}(\Phi_1, \Phi_2)$

2HDM+S potential

$$\begin{aligned} &= m_1^2 \Phi_1^\dagger \Phi_1 + m_2^2 \Phi_2^\dagger \Phi_2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + \text{h.c.}) \\ &+ \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\ &+ \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 \\ &+ \frac{1}{2} \lambda_5 \left[(\Phi_1^\dagger \Phi_2)^2 + \text{h.c.} \right] \\ &+ \left\{ \left[\lambda_6 (\Phi_1^\dagger \Phi_1) + \lambda_7 (\Phi_2^\dagger \Phi_2) \right] \Phi_1^\dagger \Phi_2 + \text{h.c.} \right\} \end{aligned}$$

$$\begin{aligned} &\mathcal{V}(\Phi_1, \Phi_2) + \frac{1}{2} m_{S_0}^2 S^2 + \frac{\lambda_{S_1}}{2} \Phi_1^\dagger \Phi_1 S^2 \\ &+ \frac{\lambda_{S_2}}{2} \Phi_2^\dagger \Phi_2 S^2 + \frac{\lambda_{S_3}}{4} (\Phi_1^\dagger \Phi_2 + \text{h.c.}) S^2 \\ &+ \frac{\lambda_{S_4}}{4!} S^4 + \mu_1 \Phi_1^\dagger \Phi_1 S + \mu_2 \Phi_2^\dagger \Phi_2 S \\ &+ \mu_3 \left[\Phi_1^\dagger \Phi_2 + \text{h.c.} \right] S + \mu_S S^3. \end{aligned}$$

Out of considerations of simplicity, assume S to be Higgs-like, which is not too far fetched.

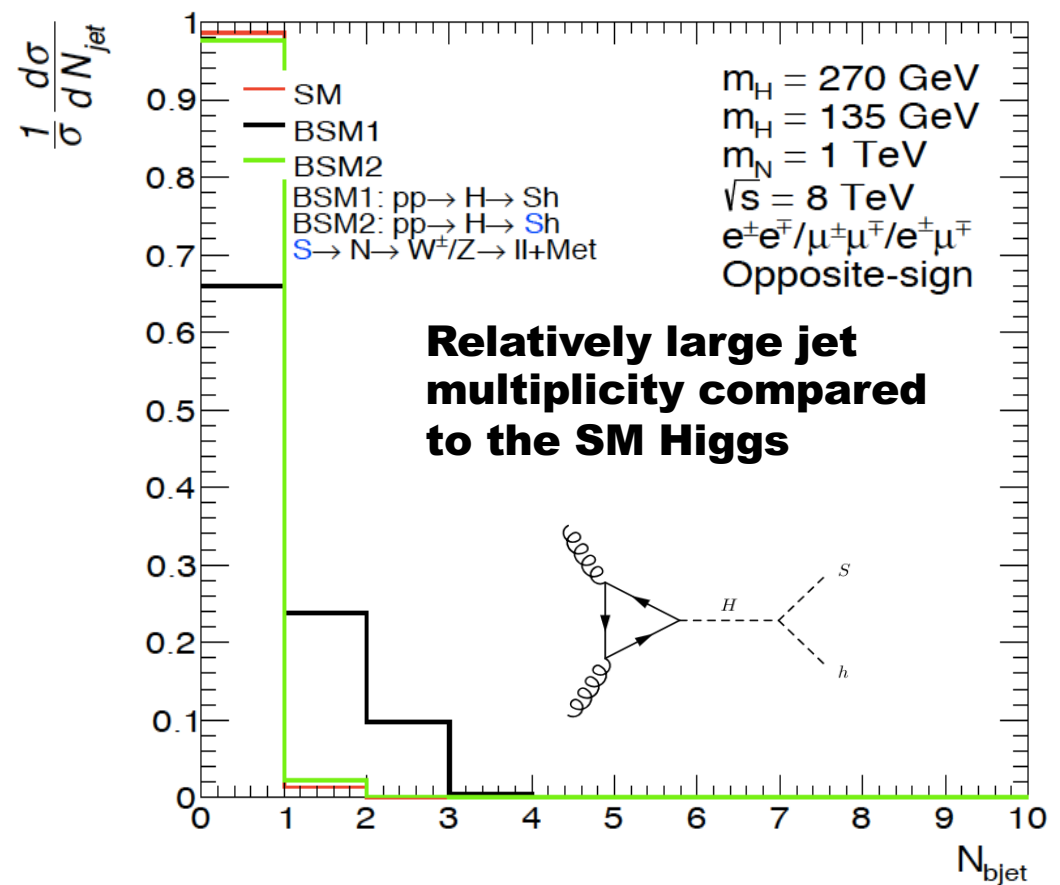
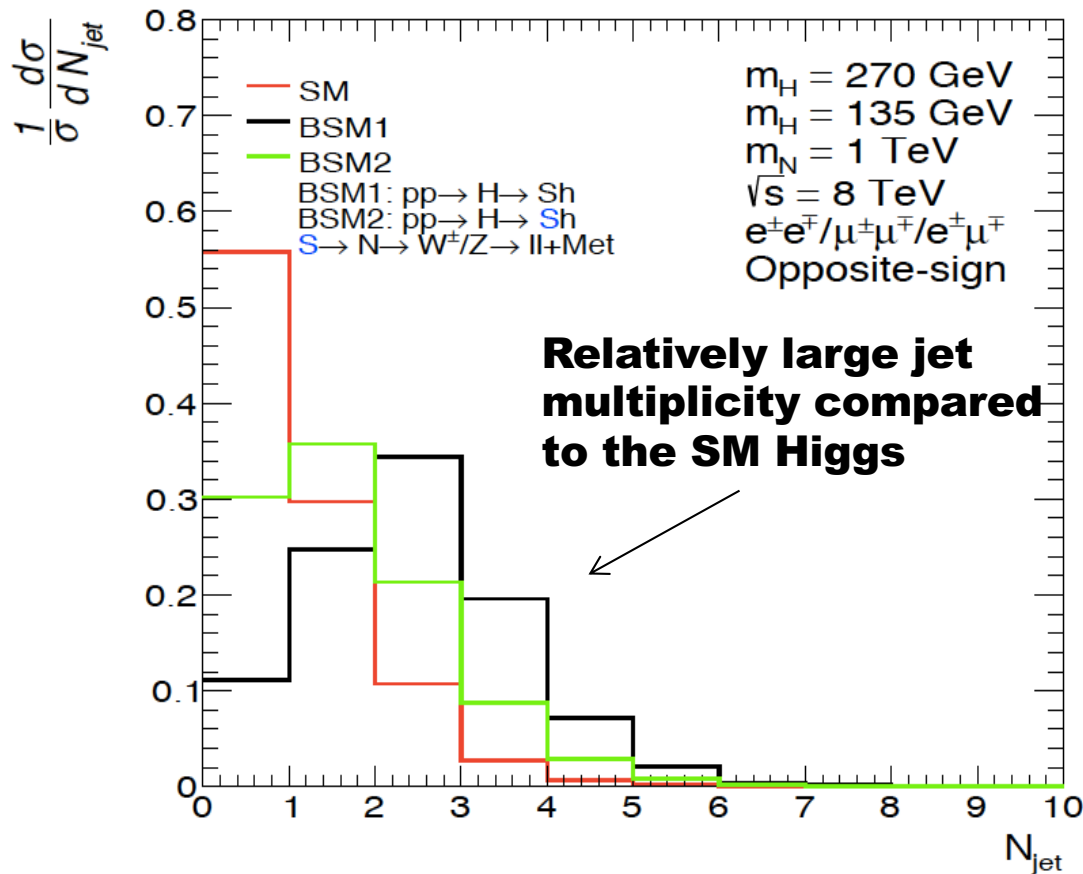
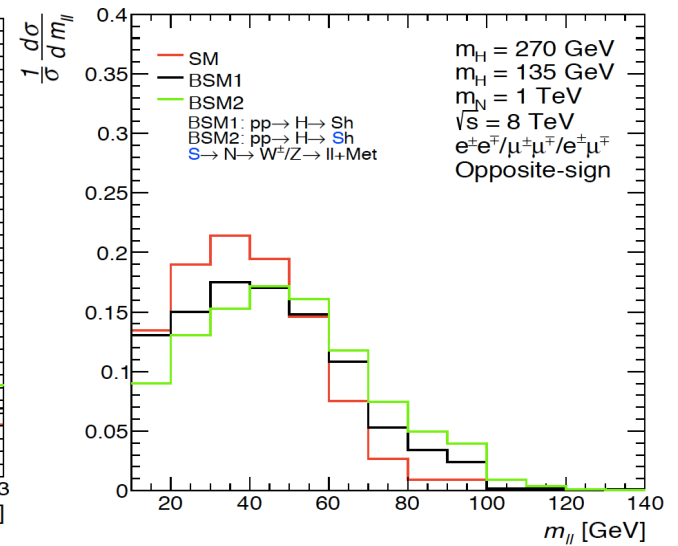
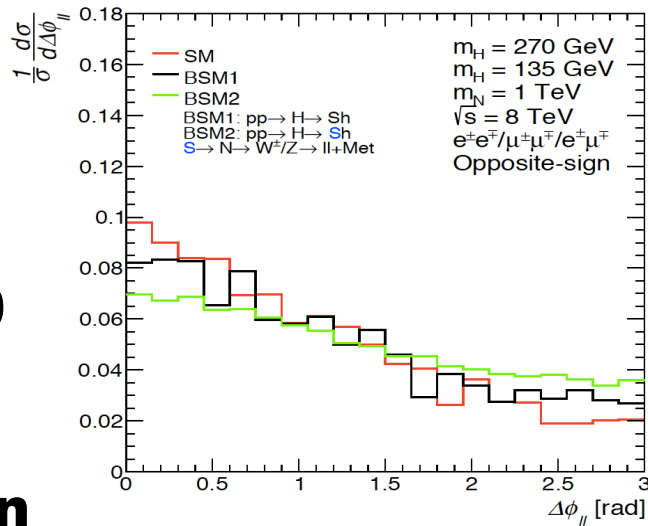
Multi-lepton final states

It is paramount to remark that the excesses are seen in final states that were predicted 2015/2016 on the basis of a simplified model and not the result of scan of the available phase-space. Additionally, the parameters of the model were fixed then leaving only one degree of freedom: normalization. Thus, no look-elsewhere effects in parameter or phase-space

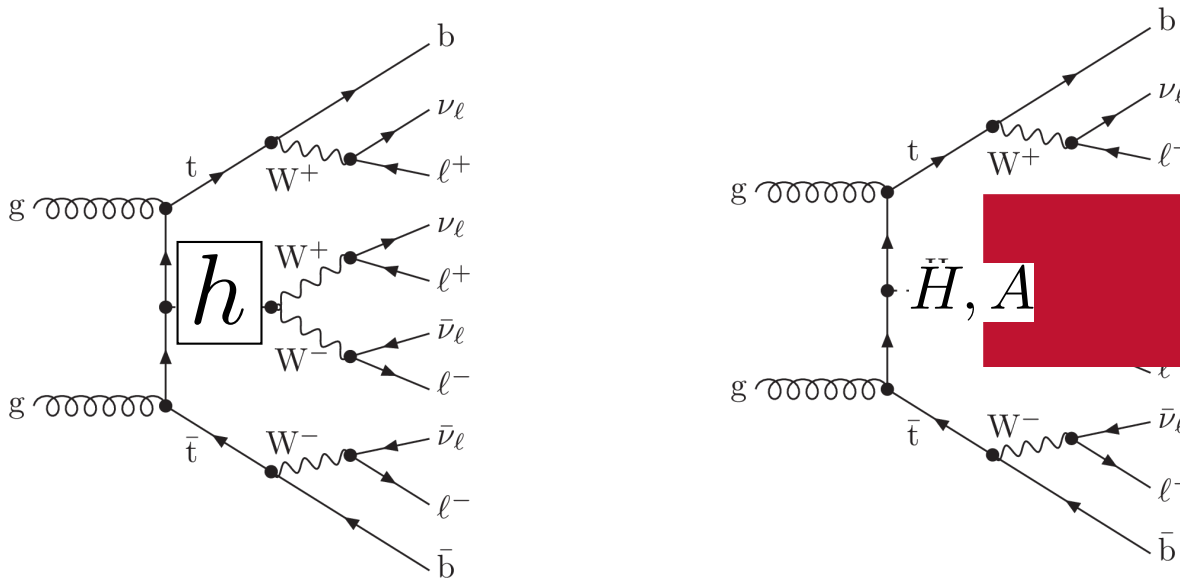
$$pp \rightarrow H \rightarrow Sh$$

$$\rightarrow \ell^+ \ell^- + X$$

Expect di-leptons ($m_{ll} < 100$ GeV) with jets and b-jets with rates comparable to that of the SM Higgs boson

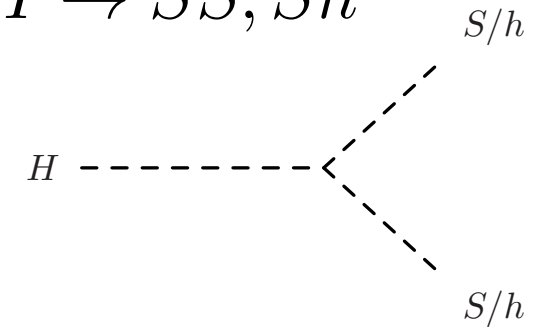


Top associated Higgs production (Multi-lepton final states)



$$A \rightarrow t\bar{t}, ZH$$

$$H \rightarrow SS, Sh$$



Reduced cross-section of $t\bar{t}H+tH$ is compensated by di-boson, (SS, Sh) decay and large $\text{Br}(S \rightarrow WW)$. Production of same sign leptons, three leptons is enhanced. Enhanced tH cross-section

Produces SS 2l, 3l with b-jets, including 3 b-jets

Explains anomalously large $t\bar{t}W+t\bar{t}h+4t$ cross-sections seen by ATLAS and CMS

Methodology

(to avoid biases and look-else-where effects)

Based Higgs p_T , hh, tth, VV in Run 1
Eur. Phys. J. C (2016) 76:580

Model defined and predictions made for
multilepton excesses

Multi-lepton excesses in Run 1 and few
Run 2 results available in 2017

J.Phys.G 45 (2018) 11, 115003

Model parameters fixed in 2017 with
 $m_H=270$ GeV, $m_S=150$ GeV,
S treated as SM Higgs-like,
dominance of $H \rightarrow Sh, SS$

Fixed final states and phase-space
defined by fixed model parameters.
NO tuning, NO scanning

Update same final states with
more data in Run 2

Study new final states where
excesses predicted and data
available in Run 1 and Run 2
(e.g., SS0b, 3l0b, ZW0b)

J.Phys. G46 (2019) no.11, 115001

JHEP 1910 (2019) 157

Chin.Phys.C 44 (2020) 6, 063103

Physics Letters B 811 (2020) 135964

Eur.Phys.J.C 81 (2021) 365

Combination of fit results (2019)

- **Simultaneous fit for all measurements:**
- **To the right: (-2 log) profile likelihood ratio for each individual result and the combination of them all**
- **The significance for each fit is calculated as**

$$\sqrt{-2 \log \lambda(0)}$$

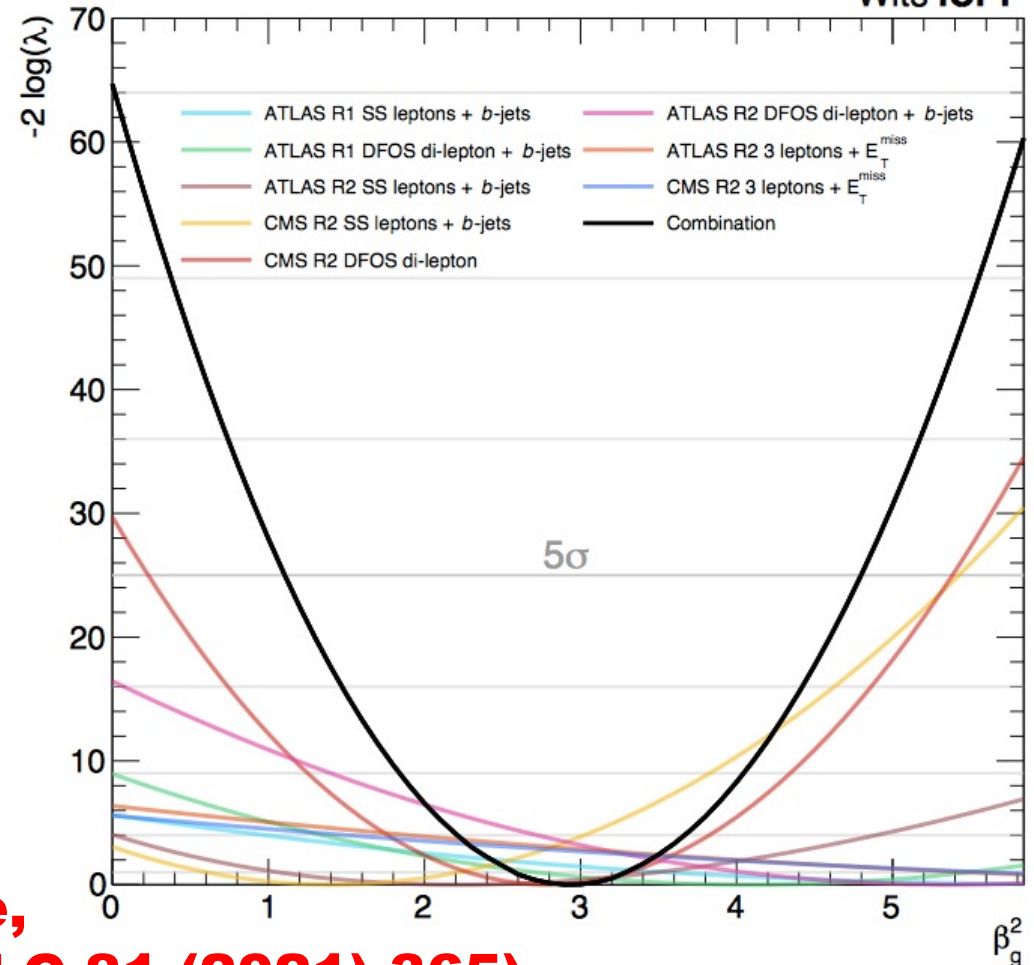
- **Best-fit: $\beta_g^2 = 2.92 \pm 0.35$**
- **Corresponds to 8.04σ**

Excesses have been growing since, and new have emerged (Eur.Phys.J.C 81 (2021) 365)

Interpretation: Measure of the inability of current MC tools to describe multiple-lepton data and how a simplified model with $H \rightarrow S_h$ is able to capture the effect with one parameter

JHEP 1910 (2019) 157

Wits ICPP



Anatomy of the multi-lepton anomalies

Final state	Characteristic	Dominant SM process	Significance
l^+l^- + jets, b-jets	$m_{ll} < 100$ GeV, dominated by 0b-jet and 1b-jet	tt+Wt	$>5\sigma$
l^+l^- + full-jet veto	$m_{ll} < 100$ GeV	WW	$\sim 3\sigma$
$l^\pm l^\pm$ & $l^\pm l^\pm l$ + b-jets	Moderate H_T	ttW, 4t	$>3\sigma$
$l^\pm l^\pm$ & $l^\pm l^\pm l$ et al., no b-jets	In association with h	Wh, (WWW)	4.2σ
Z($\rightarrow l^+l^-$)+l	$p_{TZ} < 100$ GeV	ZW	$>3\sigma$

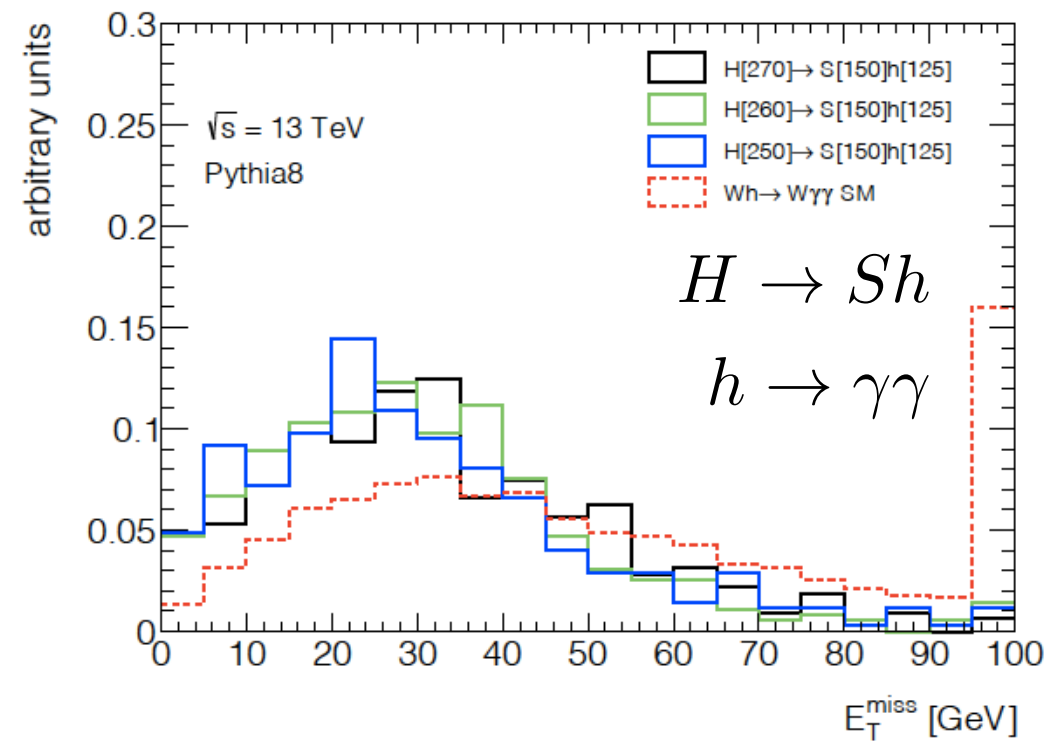
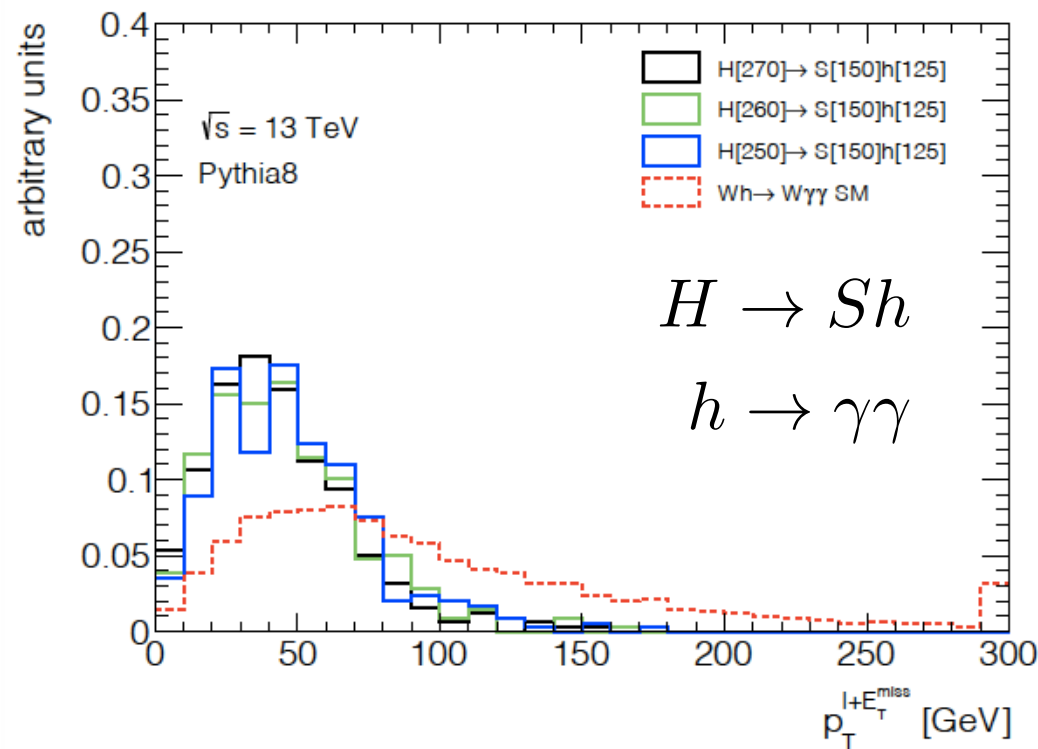
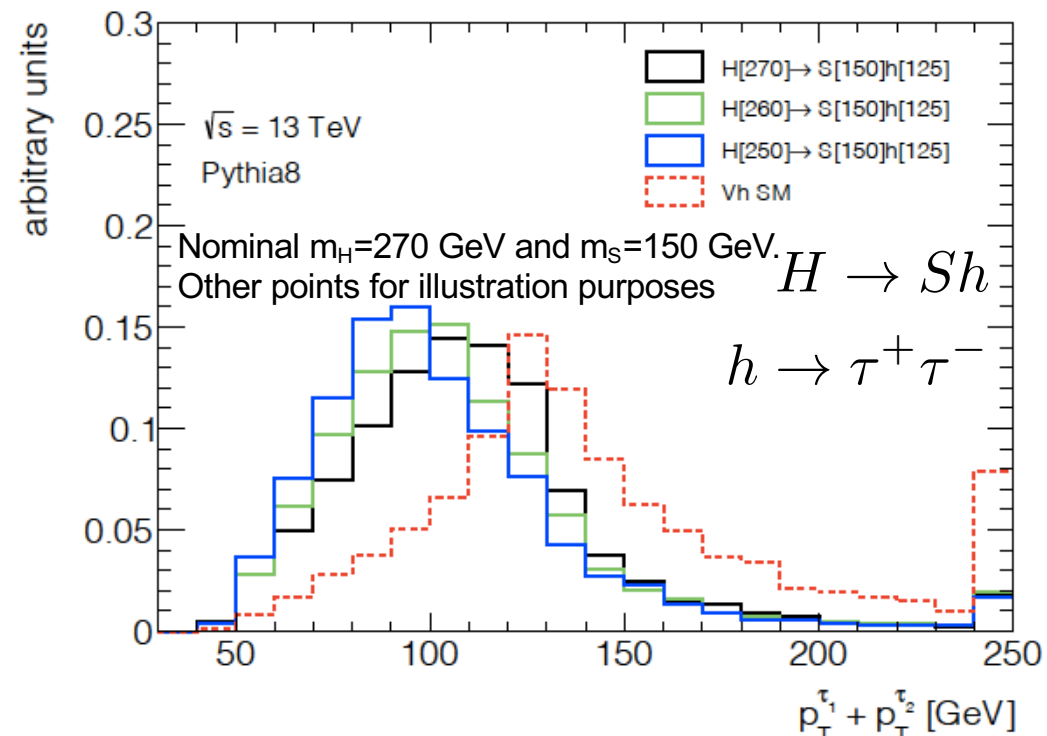
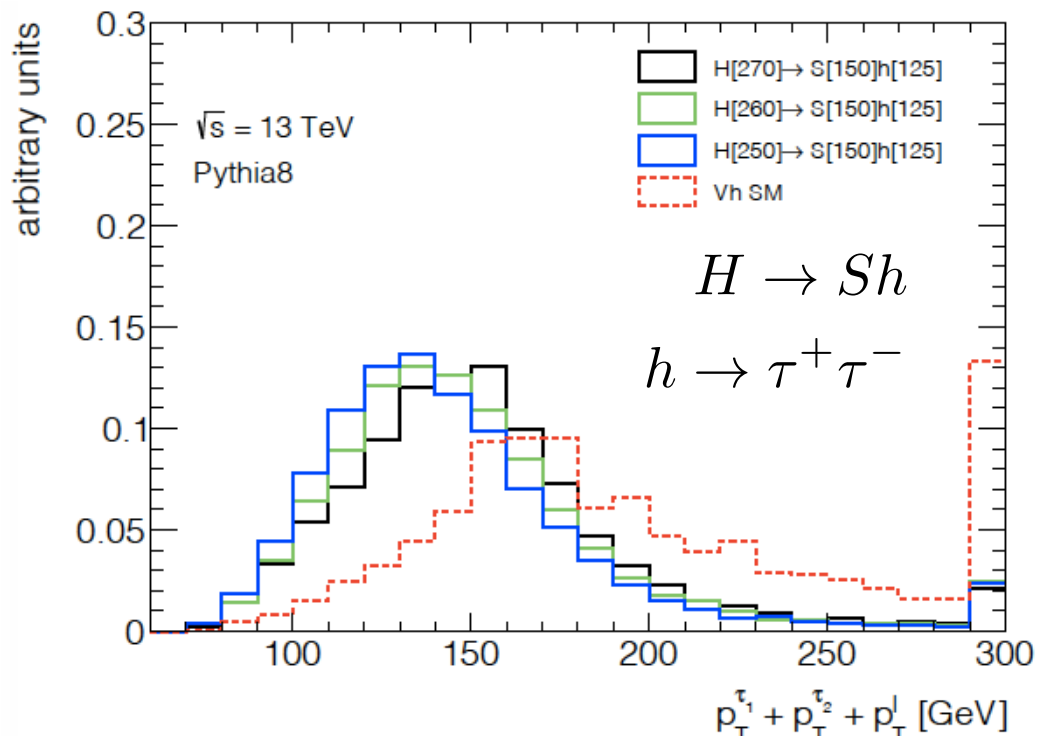
Anomalies cannot be explained by mismodelling of a particular process, e.g. ttbar production alone.

Impact on Higgs Physics

The presence of a BSM signal of the type $H \rightarrow Sh$ would lead to:

- The presence of extra leptons in association with h . Affects the Wh measurement (Eur.Phys.J.C 81 (2021) 365)**
- Distortion of Higgs p_T and rapidity (under study)**

No tuning of model parameters performed. Look at fixed corners of the phase-space fixed with parameters of 2017.



Survey of LHC results on Vh ($V=W,Z$) production (Eur.Phys.J.C 81 (2021) 365)

The BSM ($H \rightarrow Sh$) signal appears at low p_{Th} and the SM signal is prevalent at larger p_{Th} (no tuning of parameters)

Include those results from ATLAS and CMS where no requirements on p_{Th} (or correlated observables) is not done or used in an MVA.

Those results where the final state is treated more “inclusively” display elevated signal strengths for Wh production:

$$\mu(Wh) = 2.41 \pm 0.37$$

This represents a 3.8σ deviation from the SM value of 1. BSM signal normalization less than expected from multilepton excesses assuming $Br(H \rightarrow Sh)=100\%$. Indicates that $Br(H \rightarrow SS) > Br(H \rightarrow Sh)$

Higgs decay	Ref.	Experiment	\sqrt{s}, \mathcal{L} TeV, fb ⁻¹	Final state	Category	μ	Used in combination	Comments				
WW	66	ATLAS	7, 4.5 8, 20.3	2 ℓ	DFO5 2j	$2.2^{+2.0}_{-1.9}$	✓	$m_{\ell_0 \ell_2}$ used as input BDT discriminating variable				
					SS 1j	$8.4^{+4.3}_{-3.8}$	✓		2 ℓ combination: $\mu = 3.7^{+1.9}_{-1.5}$			
					SS 2j	$7.6^{+5.0}_{-3.4}$	✓					
				1SFOS	$-2.9^{+2.7}_{-2.1}$	x						
				0SFOS	$1.7^{+1.9}_{-1.4}$	✓						
	67	ATLAS	13, 36.1	3 ℓ	1SFOS	$2.3^{+1.2}_{-1.0}$	✓	1SFOS channel uses $m_{\ell_0 \ell_2}$ in the BDT but excess driven by 0SFOS				
					0SFOS							
					68	CMS	7, 4.9 8, 19.4	2 ℓ	DFO5 2j	$0.39^{+1.97}_{-1.87}$	✓	Discrepancy at low $m_{\ell\ell}$
								3 ℓ	0+1SFOS	$0.56^{+1.27}_{-0.95}$	✓	
								69	CMS	13, 35.9	2 ℓ	DFO5 2j
3 ℓ	0+1SFOS	$2.23^{+1.76}_{-1.53}$	✓									
$\tau\tau$	70	ATLAS	8, 20.3	1 ℓ	$\ell + \tau_h \tau_h$	1.8 ± 3.1	✓	BDT based on $p_T^{\tau_1} + p_T^{\tau_2}$ Split $p_T^{\ell_1} + p_T^{\ell_2} + p_T^{\tau}$ at 130 GeV				
				2 ℓ	$e^{\pm} \mu^{\pm} + \tau_h$	1.3 ± 2.8	✓					
	71	CMS	7, 4.9 8, 19.7	1 ℓ	$\ell + \tau_h \tau_h$	-0.33 ± 1.02	x					
				2 ℓ	$e^{\pm} \mu^{\pm} + \tau_h$		x					
	72	CMS	13, 35.9	1 ℓ	$\ell + \tau_h \tau_h$	$3.39^{+1.68}_{-1.54}$	✓					
				2 ℓ	$e^{\pm} \mu^{\pm} + \tau_h$							
$\gamma\gamma$	73	ATLAS	7, 5.4 8, 20.3	$\ell\nu$	One-lepton	1.0 ± 1.6	x	$E_T^{miss} > 70 - 100$ GeV				
				$\ell\nu, \nu\nu$	E_T^{miss}			$p_{T\gamma}^{\gamma\gamma} > 70$ GeV				
				jj	Hadronic							
	74	CMS	7, 5.1 8, 19.7	$\ell\nu$	One-lepton		x	Split E_T^{miss} at 45 GeV				
				$\ell\nu, \nu\nu$	E_T^{miss}	$-0.16^{+1.16}_{-0.79}$	x	$E_T^{miss} > 70$ GeV				
				jj	Hadronic			$p_{T\gamma}^{\gamma\gamma} > 13m_{\gamma\gamma}/12$				
	75	ATLAS	13, 139	$\ell\nu$	One-lepton	$2.41^{+0.71}_{-0.70}$	✓	$p_T^{\ell+E_T^{miss}} < 150$ GeV				
				$\ell\nu, \nu\nu$	E_T^{miss}	$2.64^{+1.16}_{-0.99}$	x	$p_T^{\ell+E_T^{miss}} > 150$ GeV				
				jj	Hadronic		x	$E_T^{miss} > 75$ GeV				
	76	CMS	13, 35.6	$\ell\nu$	One-lepton	$0.76^{+0.95}_{-0.83}$	x	$60 < m_{jj} < 120$ GeV				
				$\ell\nu, \nu\nu$	E_T^{miss}		x	$m_{jj} \in [0, 60] \cup [120, 350]$ GeV				
				jj	Hadronic	$3.16^{+1.84}_{-1.72}$	✓	Superseded by full Run 2 result				
77	CMS	13, 137	$\ell\nu$	One-lepton	$3.0^{+1.5}_{-1.3}$	x	$E_T^{miss} > 85$ GeV					
			$\ell\nu, \nu\nu$	E_T^{miss}		x	$p_T^{\gamma\gamma}/m_{\gamma\gamma}$ not used					
			jj	Hadronic	$5.1^{+2.5}_{-2.3}$	✓	$p_T^V < 75$ GeV					
78	ATLAS	13, 139	$\ell\nu$	One-lepton	$1.31^{+1.42}_{-1.12}$	✓	$p_T^{\gamma\gamma}/m_{\gamma\gamma}$ used in BDT					
			$\ell\nu, \nu\nu$	E_T^{miss}	$0.89^{+0.89}_{-0.91}$	x						
			jj	Hadronic		x						
ZZ	79	CMS	13, 137.1	$\ell\ell\ell + \ell\nu$	Lep-enriched	$1.44^{+1.17}_{-0.93}$	x	Number of jets used in MVA				
				$\ell\ell\ell + q\bar{q}$	2j			m_{jj} used in MVA				
				$\ell\ell\ell + \ell\nu$	Lep-low p_T^{Δ}	$3.21^{+2.49}_{-1.85}$	✓	$p_T^{\Delta} < 150$ GeV				
79	CMS	13, 137.1	$\ell\ell\ell + \ell\nu$	Lep-high p_T^{Δ}	$0.00^{+1.57}_{-0.60}$	x	$p_T^{\Delta} > 150$ GeV					
			$\ell\ell\ell + q\bar{q}$	2j	$0.57^{+1.20}_{-0.57}$	x	$60 < m_{jj} < 120$ GeV					

CMS *Preliminary*

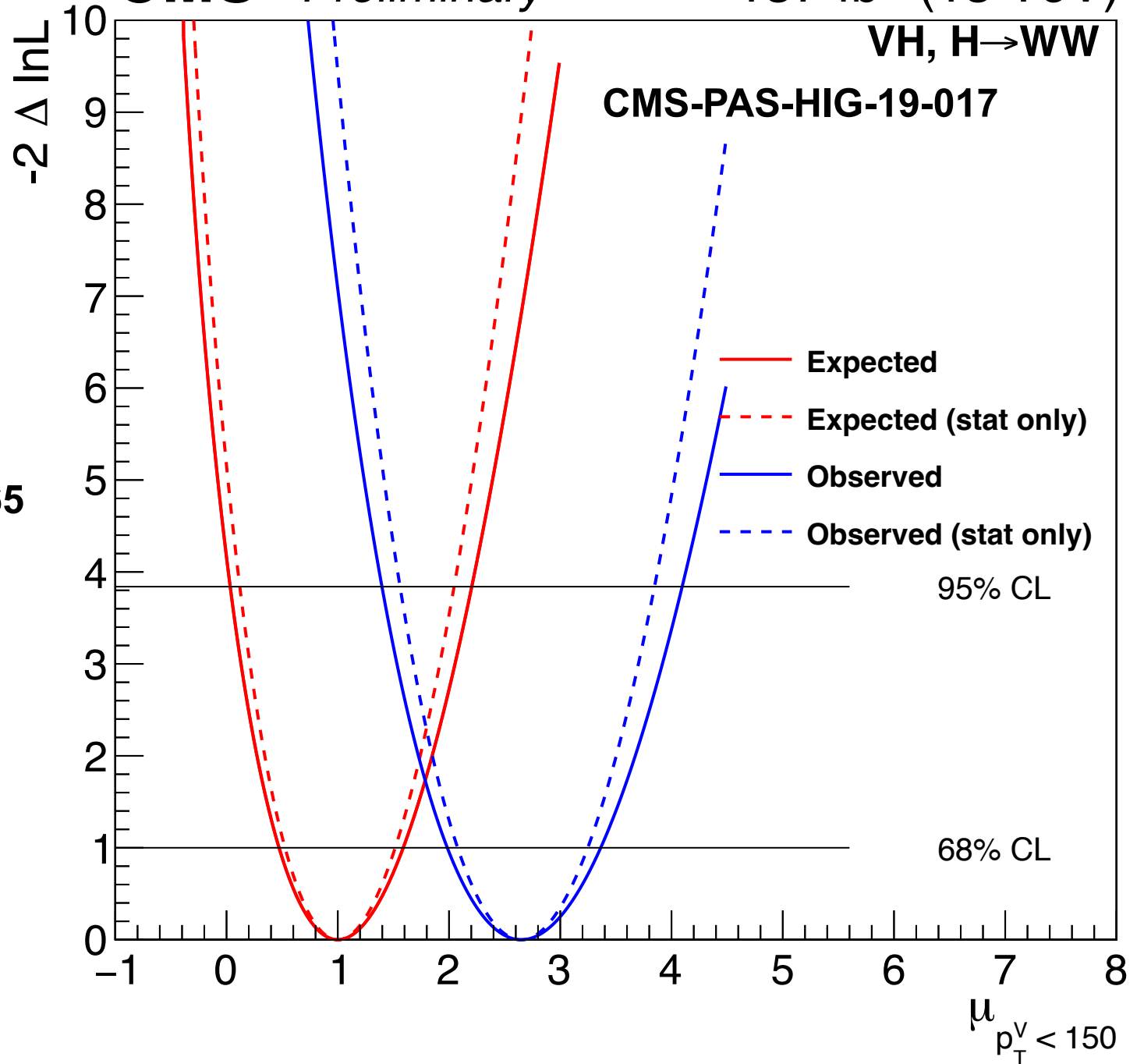
137 fb⁻¹ (13 TeV)

VH, H→WW

CMS-PAS-HIG-19-017

New results from CMS in the measurement of Vh , $h \rightarrow WW$ add to the anomalies reported in *Eur.Phys.J.C* 81 (2021) 365

Deviation from the SM becomes stronger with $p_{TV} < 150$ GeV

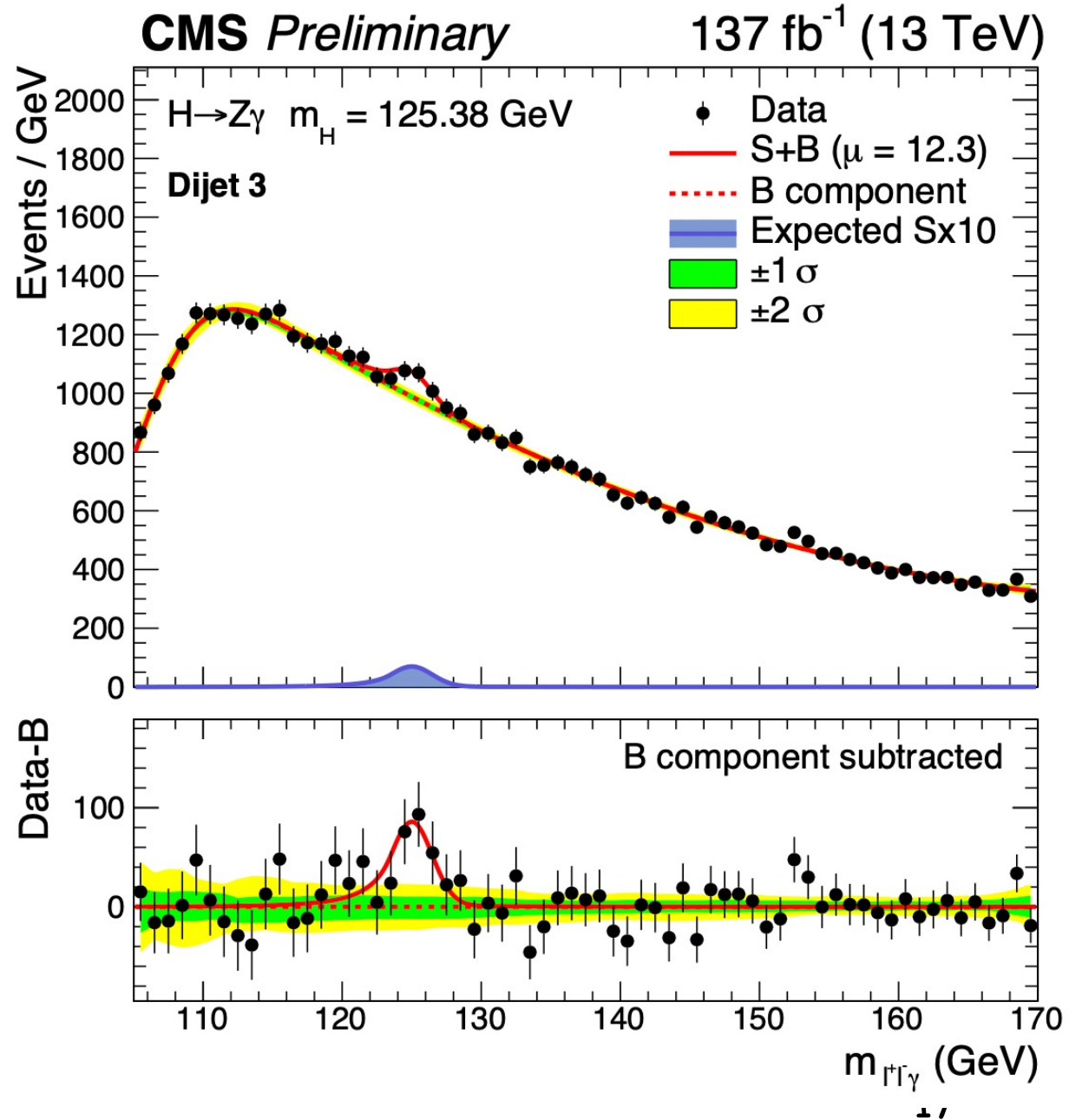


CMS Physics Analysis Summary

Search for the Higgs boson decay to $Z\gamma$ in proton-proton collisions at $\sqrt{s} = 13$ TeV

The CMS Collaboration

CMS observes what appears to be an upward fluctuation of the $h \rightarrow Z\gamma$ in the di-jet bin optimized for the measurement of Wh production. The Signal strength deviates from unity by 3.2σ .

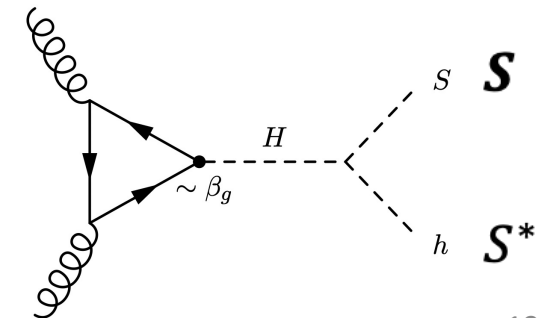


**A possible
candidate of S**

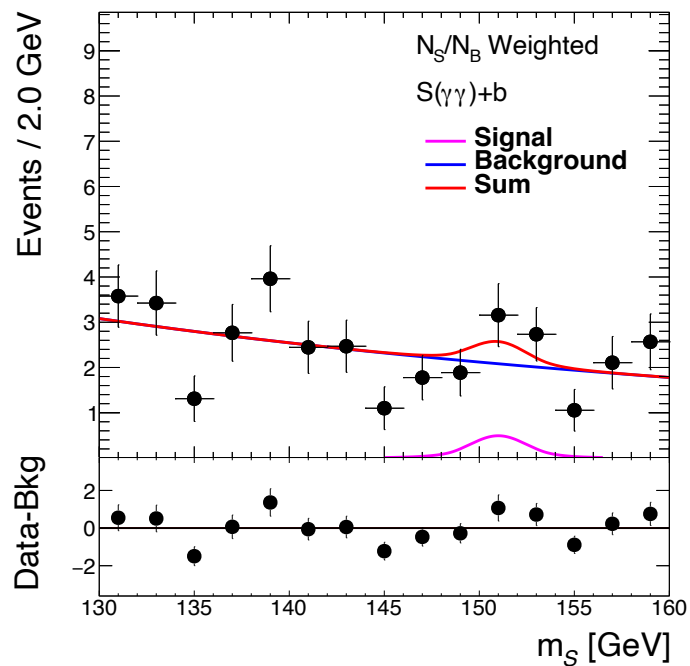
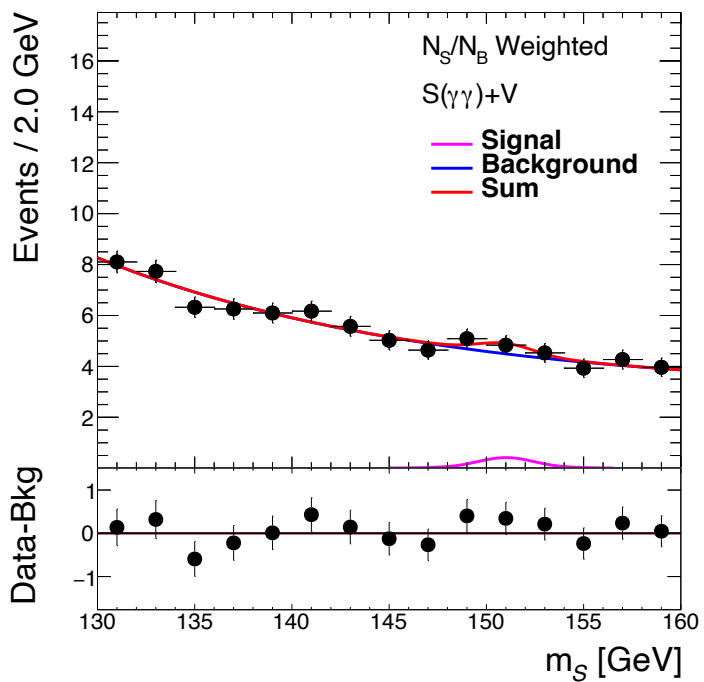
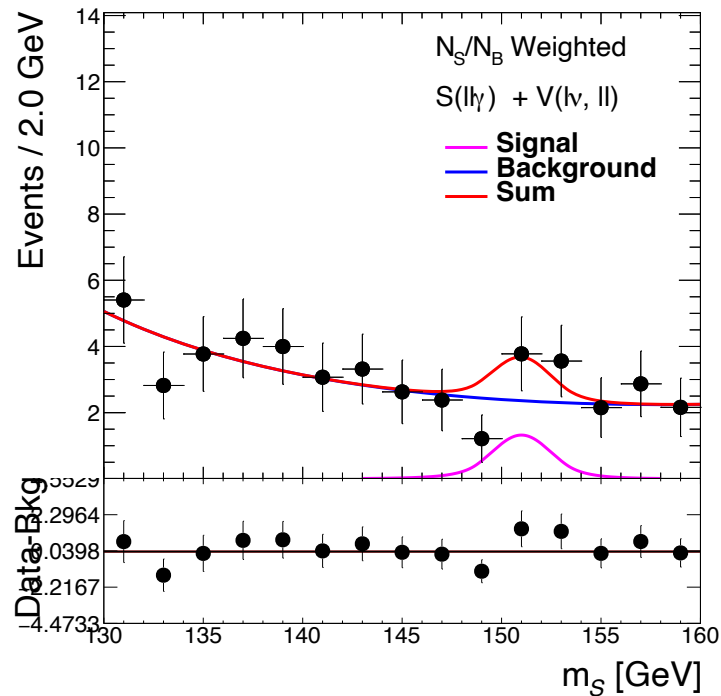
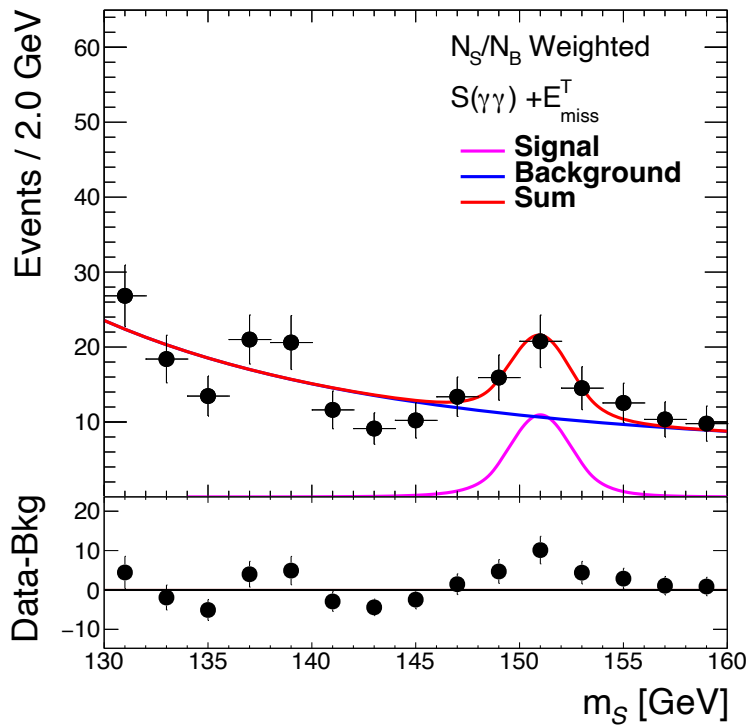
Procedure

(avoiding “cherry picking”)

- ❑ **Setting a well-defined procedure is essential to the integrity of a search. Scanning nullifies significance**
- ❑ **From the di-lepton anomalies: $m_h < m_s < 170$ GeV**
 - ❑ **It is critical that search be localized and motivated**
- ❑ **Focus on $\gamma\gamma$ and $Z\gamma$ decays**
- ❑ **As per the model that described the multi-lepton anomalies, we select final state according to di-boson signatures. S is produced via the decay of something heavier and not directly**
 - ❑ **Re-use Higgs boson data**
 - ❑ **Remove VBF and boosted topologies**
 - **Related to direct production**

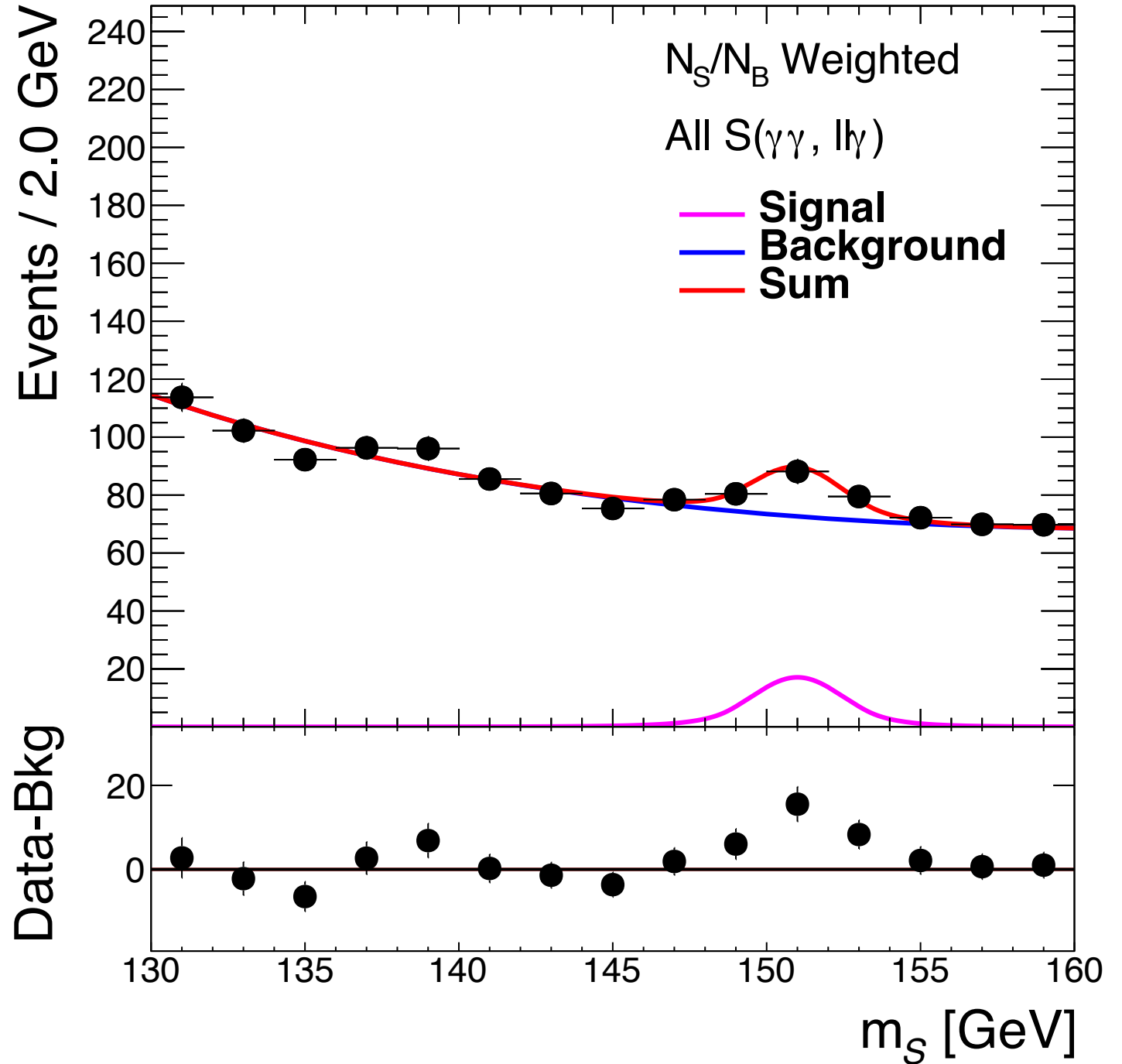


From Run 1 multi-lepton excesses model-dependent prediction of $m_s = 150 \pm 5$ GeV

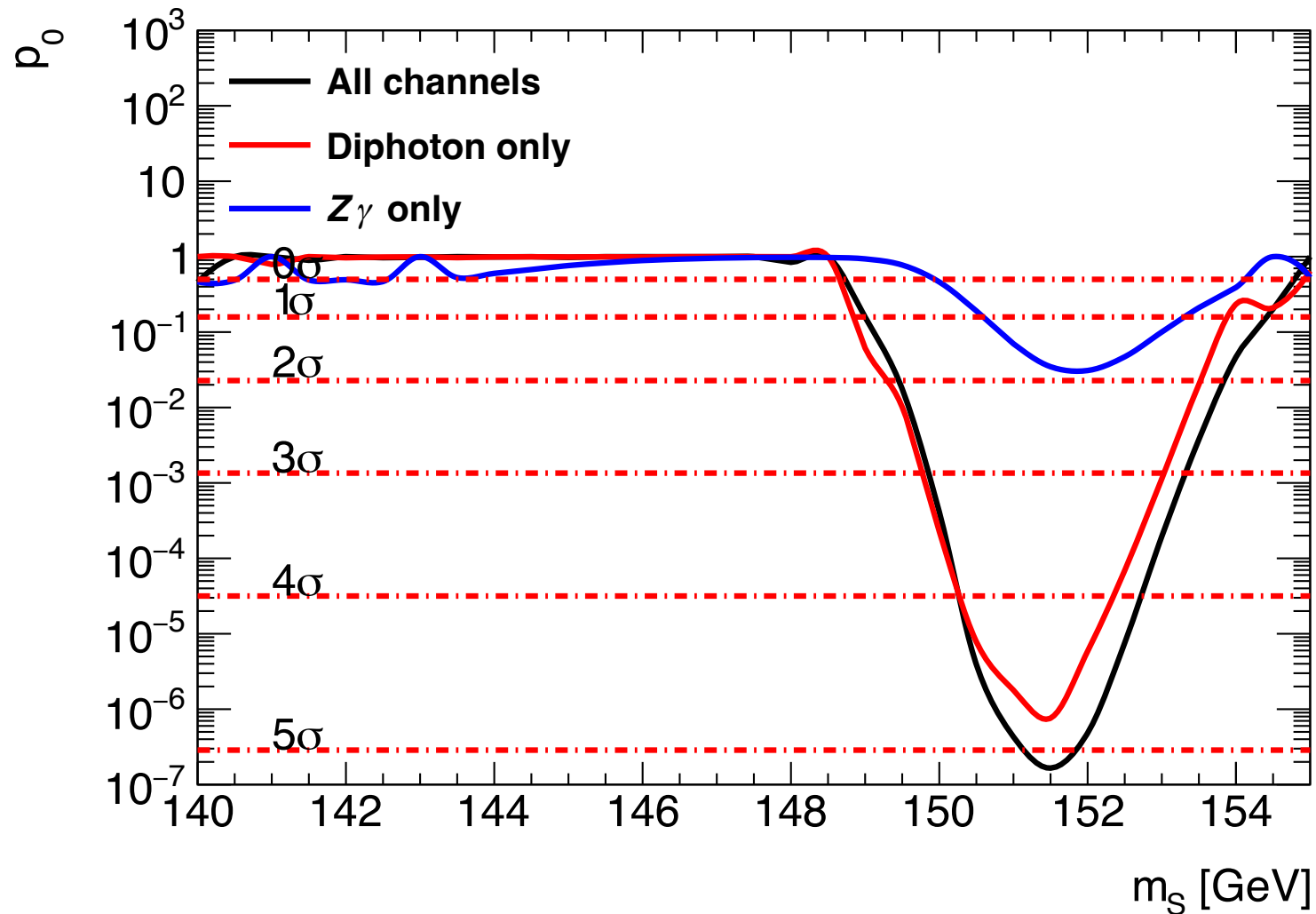


Analysis of publicly available di-photon and $Z\gamma$ spectra in associated production gives global 4.8σ excess around 151 GeV. Fiducial yields consistent with $H \rightarrow SS^*$ hypothesis with $m_H = 270$ GeV (see above)

Excess not seen in $S \rightarrow ZZ \rightarrow 4\ell, \ell = e, \mu$



Result is obtained with public results from the LHC experiments. Using a simplified model and two degrees of freedom global significance drops to $\sim 4\sigma$ (under review).



CMS Physics Analysis Summary

<https://cds.cern.ch/record/2803738/files/HIG-20-013-pas.pdf>

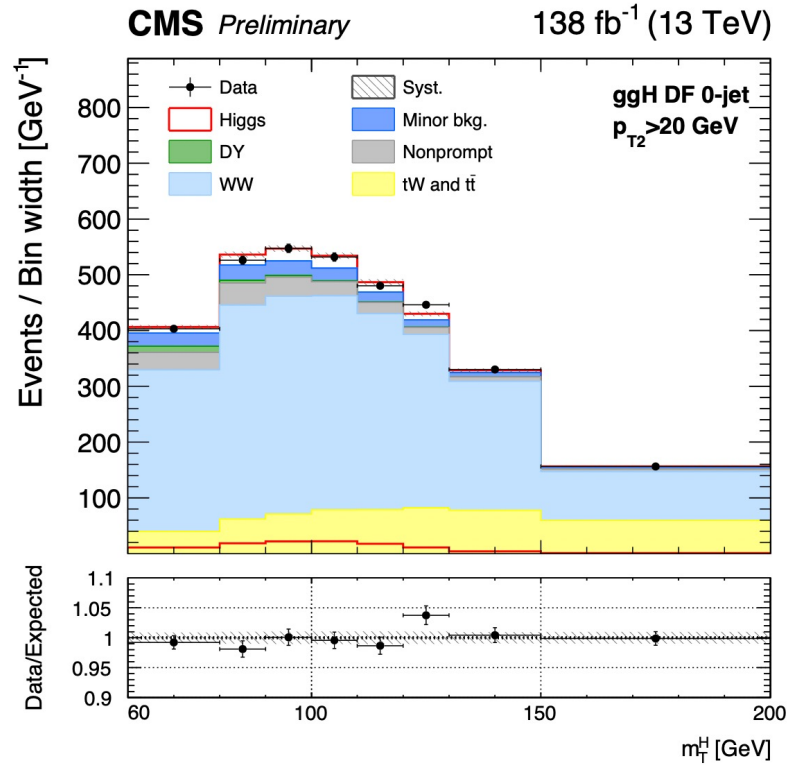
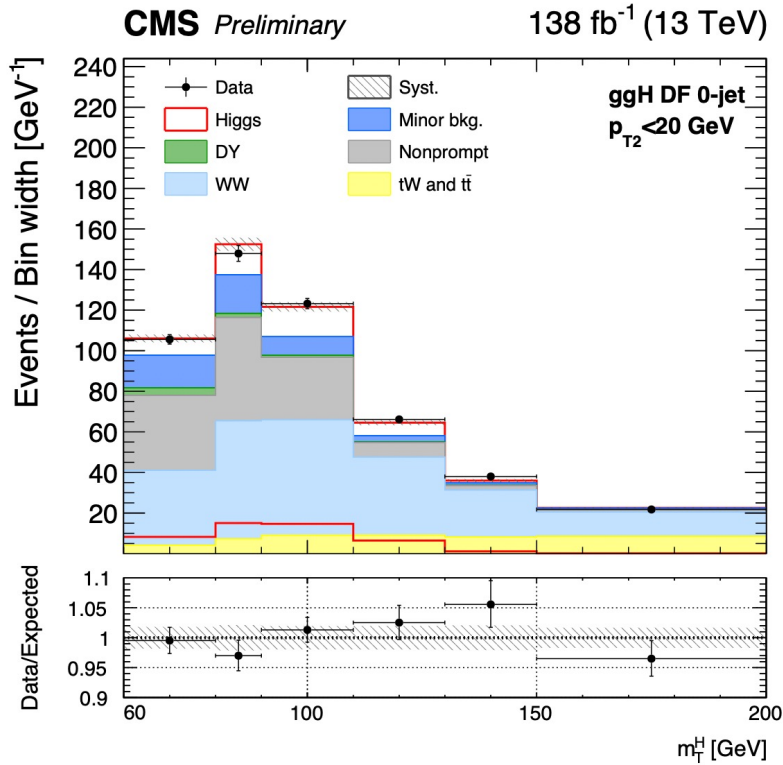
Contact: cms-pag-conveners-higgs@cern.ch

2022/03/11

Recent results with $ll+MET$ not included above

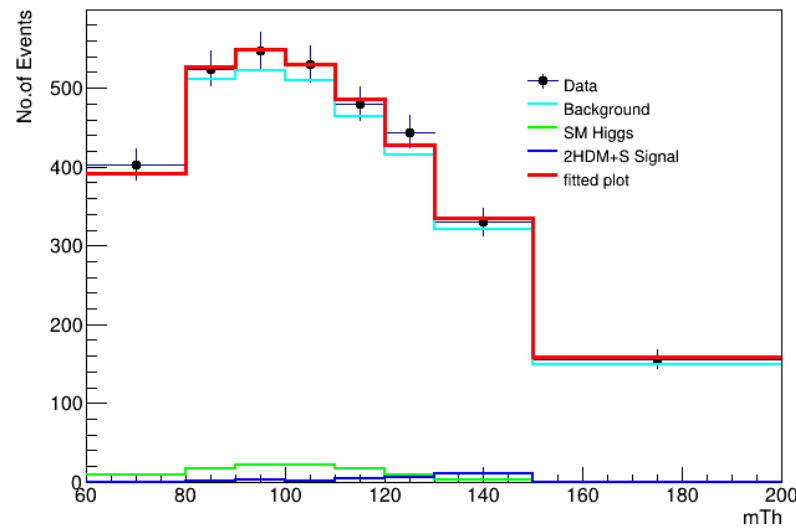
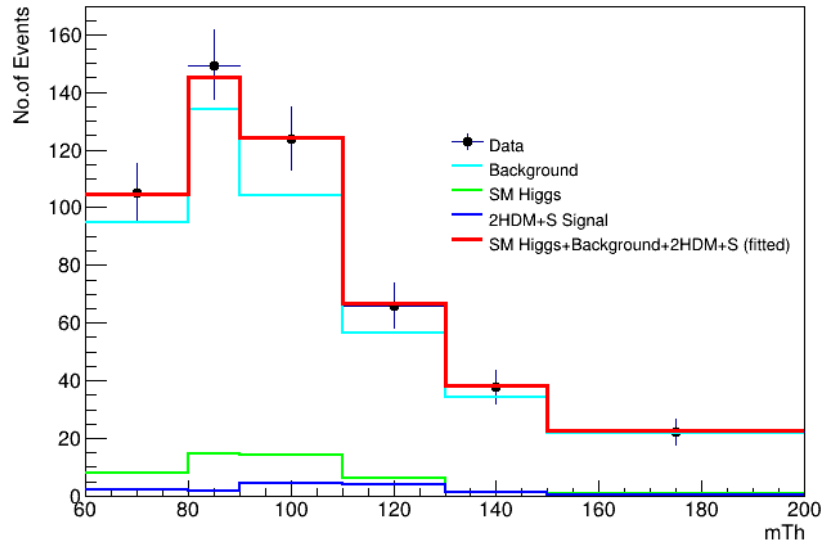
Measurements of properties of the Higgs boson in the W boson pair decay channel in proton-proton collisions at
 $\sqrt{s} = 13$ TeV

The CMS Collaboration



Contribution from a 2HDM+S signal is given by the dark blue histograms in the bottom two plots.

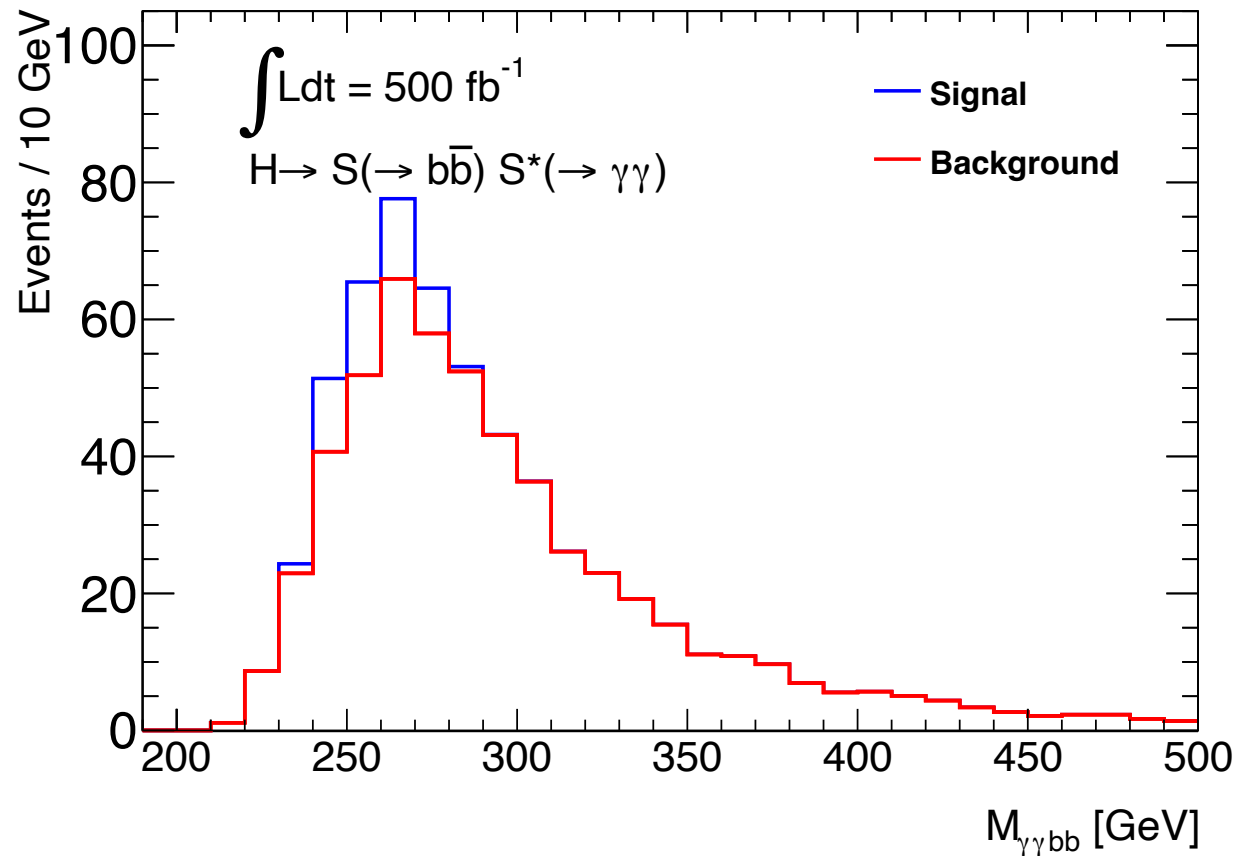
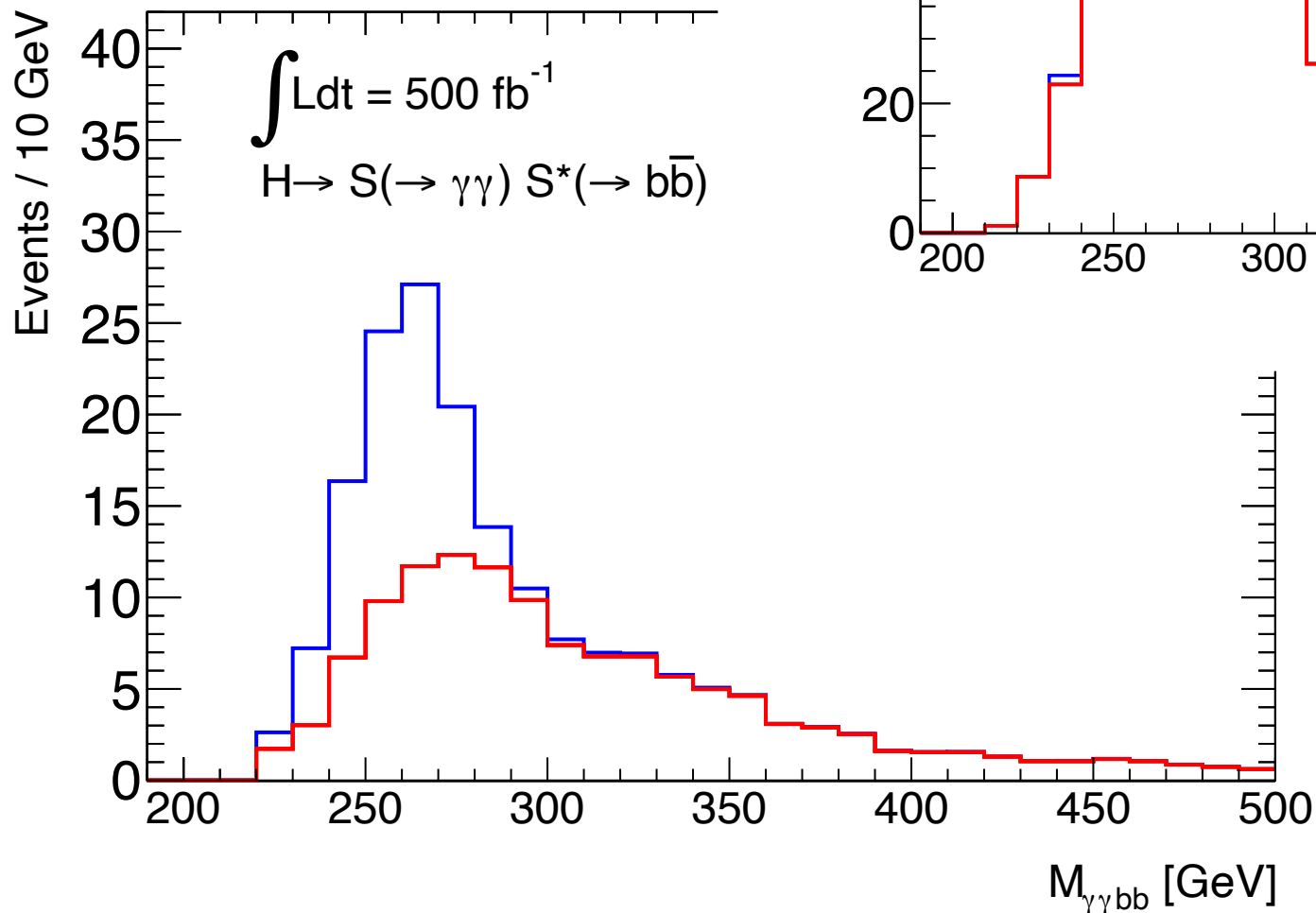
The BSM signal normalization of compatible for the two bins in p_{T12} .



Signal extraction in the 1j bin and combination with other excesses in progress.

Data favors an admixture of a signal with $m_S = 150$ GeV

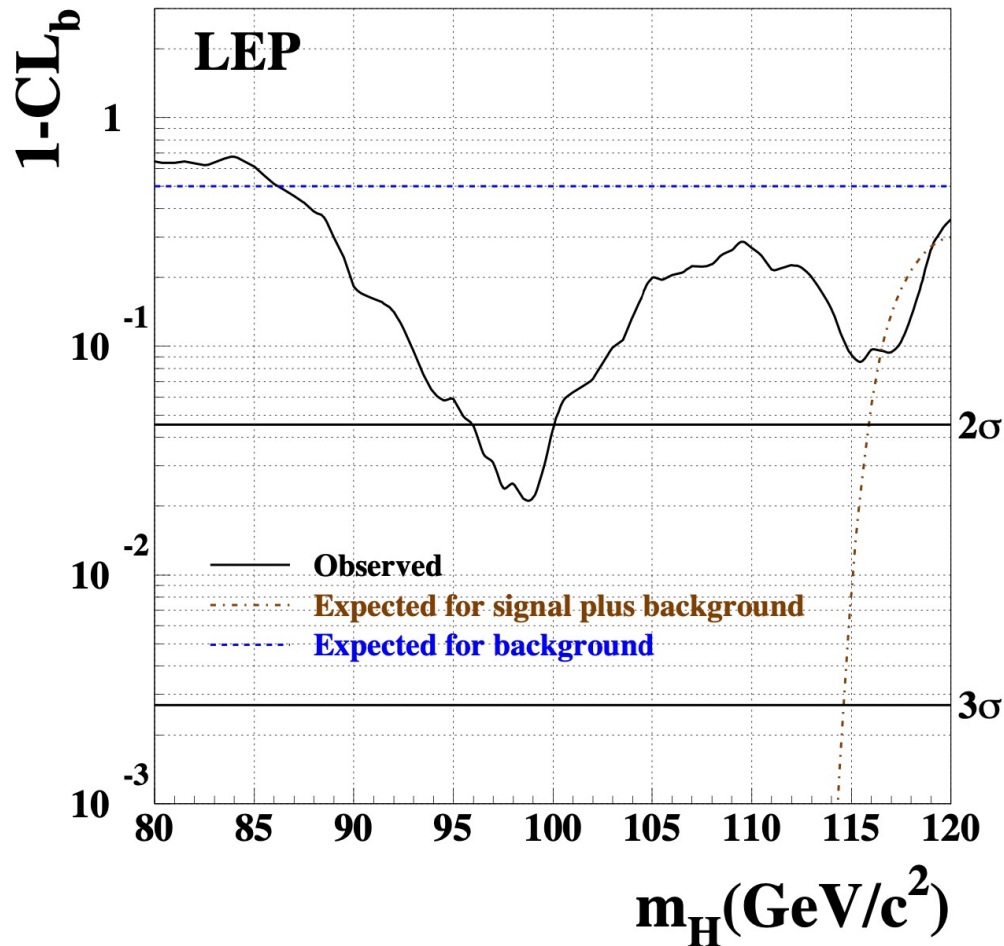
Abovementioned excess further motivates searches for bosons in asymmetric $\gamma\gamma b\bar{b}$ configurations not performed before at the LHC



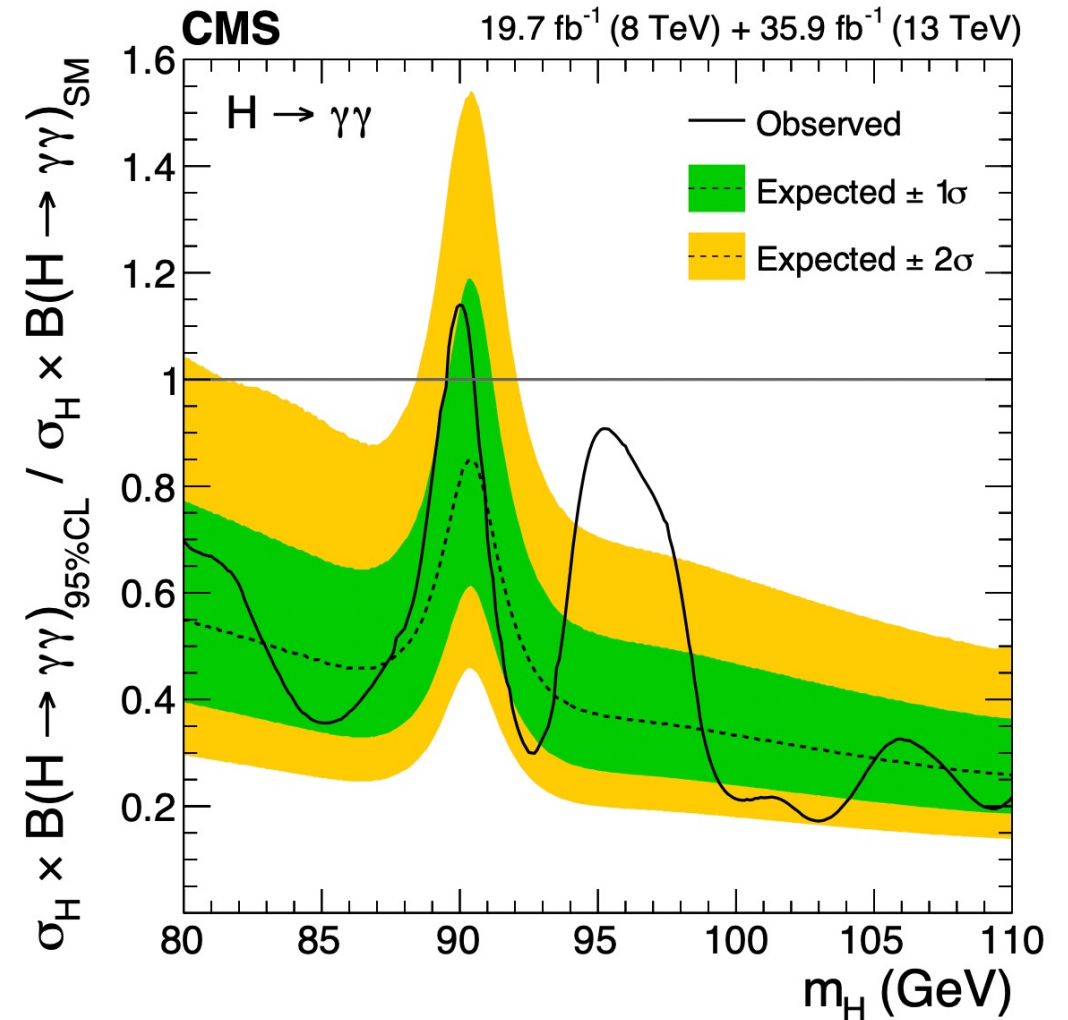
Expect more than 7σ significance for one experiment with the Run 2 + Run 3 data sets.

Some tantalizing results around 96 GeV from LEP and CMS, not contradicted by ATLAS. Interesting to see what the full Run 2 data set has to say. Further motivates asymmetric searches $H \rightarrow SS'$...

LEP, Phys. Lett. B 565 (2003) 61–75



CMS, Phys. Lett. B 793 (2019) 320–347



Oliver Fischer^{†,1}, Bruce Mellado^{†,2,3},
 Stefan Antusch⁴, Emanuele Bagnaschi⁵, Shankha Banerjee⁶, Geoff Beck²,
 Benedetta Belfatto^{7,8}, Matthew Bellis⁹, Zurab Berezhiani^{10,11}, Monika
 Blanke^{12,13}, Bernat Capdevila^{14,15}, Kingman Cheung¹⁶, Andreas
 Crivellin^{5,6,17}, Nishita Desai¹⁸, Bhupal Dev¹⁹, Rohini Godbole²⁰, Tao Han²¹,
 Philip Harris^{22, 23}, Martin Hoferichter²⁴, Matthew Kirk^{25,26}, Suchita
 Kulkarni²⁷, Clemens Lange²⁸, Kati Lassila-Perini²⁹, Zhen Liu³⁰, Farvah
 Mahmoudi^{6,31}, Claudio Andrea Manzari^{5,17}, David Marzocca³², Biswarup
 Mukhopadhyaya³³, Antonio Pich³⁴, Xifeng Ruan², Luc Schnell^{35, 36}, Jesse
 Thaler^{22, 23}, and Susanne Westhoff³⁷

¹*Department of Mathematical Sciences, University of Liverpool, Liverpool, L69 7ZL, UK*

²*School of Physics and Institute for Collider Particle Physics, University of the Witwatersrand, Johannesburg, Wits 2050, South Africa.*

³*iThemba LABS, National Research Foundation, PO Box 722, Somerset West 7129, South Africa.*

⁴*Department of Physics, University of Basel, Klingelbergstr. 82, CH-4056 Basel, Switzerland*

⁵*Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland*

⁶*CERN Theory Division, CH-1211 Geneva 23, Switzerland*

⁷*Dipartimento di Fisica "E. Fermi", Università di Pisa, Largo Bruno Pontecorvo 3, I-56127 Pisa, Italy*

⁸*INFN Sezione di Pisa, Largo Bruno Pontecorvo 3, I-56127 Pisa, Italy*

⁹*Siena College, 515 Loudon Road, Loudonville, NY 12211-1462, United States*

¹⁰*Dipartimento di Fisica e Chimica, Università di L'Aquila, 67100 Coppito, L'Aquila, Italy*

¹¹*INFN, Laboratori Nazionali del Gran Sasso, 67100 Assergi, L'Aquila, Italy*

¹²*Institute for Astroparticle Physics (IAP), Karlsruhe Institute of Technology, Hermann-von-Helmholtz-Platz 1, D-76344 Eggenstein-Leopoldshafen, Germany*

¹³*Institute for Theoretical Particle Physics (TTP), Karlsruhe Institute of Technology, Engesserstrasse 7, D-76128 Karlsruhe, Germany*

Outlook and Conclusions

- **Discrepancies in multi-lepton final states at LHC w.r.t. current MCs are not statistical fluctuations**
 - **They appear in corners of the phase-space dominated by different processes** ($Wt/tt/4t$, VV , ttV , Vh , WWW)
 - **Hard to explain with MC mismodelling**
 - **Discrepancies interpreted with simplified model where $H \rightarrow SS$, Sh is treated as SM Higgs-like and one parameter is floated**
- **Features of the Higgs data from LHC agree with predictions the simplified model used here**
- **Analysis of $\gamma\gamma/Z\gamma$ spectra in associated production gives a large global excess around 151 GeV**
 - **Motivates $H \rightarrow SS^*$, $Sh \rightarrow \gamma\gamma bb$ searches, where asymmetric configurations play an important role**