



第八届中国LHC物理研讨会 The 8th China LHC Physics Workshop

2022年11月23日-27日，南京汤山颐尚酒店
Nov. 23-27, 2022, Nanjing Tangshan Yishang Hotel

Exotic hadrons at



Xiao-Rui Lyu (吕晓睿)

University of Chinese Academy Sciences

(on behalf of the LHCb collaboration)



Outline

- Introduction
- $\chi_{c1}(3872) \rightarrow \pi^+ \pi^- J/\psi$
- Open-charm tetraquark $T_{c\bar{s}0}^a(2900)$ in $B \rightarrow DD_s \pi$
- $X(3960)$ in $B^+ \rightarrow D_s^+ D_s^- K^+$
- Neutral $Z_{cs}(4000)/T_{\phi s1}^\theta(4000)$ in $B^0 \rightarrow J/\psi \phi K_S$
- Searches for exotics in
 - ✓ $B^- \rightarrow \Lambda_c^+ \Lambda_c^- K^-$
 - ✓ $B^+ \rightarrow K^+ \eta J/\psi$
- Pentaquark $P_{\psi s}^\Lambda(4338)$ in $B^- \rightarrow J/\psi \Lambda \bar{p}$
- Summary

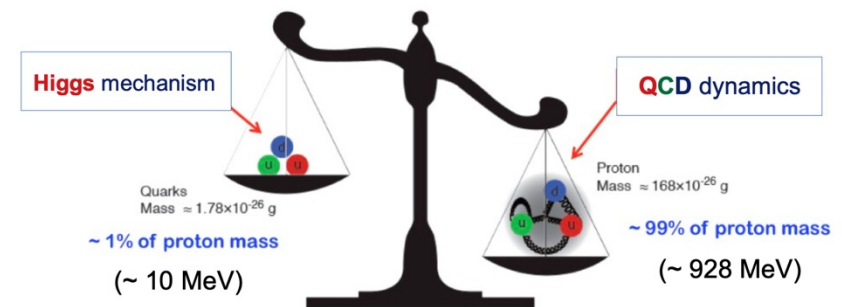
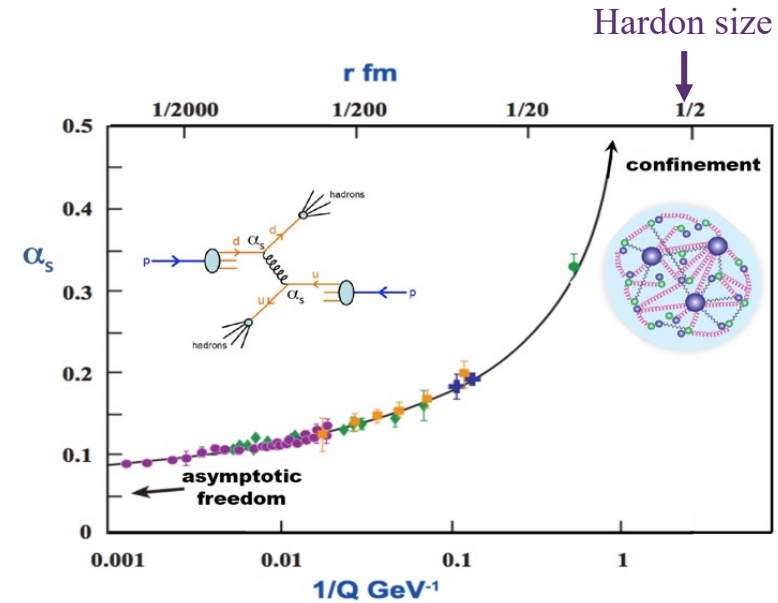
more details, see talks by:

Shuangli Yang and Zhihong Shen today afternoon;

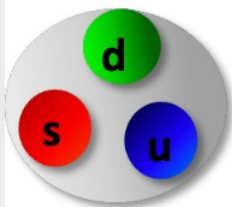
Quan Zou yesterday afternoon

Introduction

- Quarks and gluons not isolated in nature.
 - Formation of colorless bound states: “**Hadrons**”
 - **1-fm scale** size of hadrons?
- Hadron spectroscopy provides opportunities to study QCD in the non-perturbative region
 - Extensive and precise spectroscopy combined with a thorough theoretical analysis, will add substantially to our knowledge of QCD
- Complex exotic hadrons can reveal new or hidden aspects of the dynamics of strong interactions
 - Predicted in quark model
 - Recent results show strong evidence for their existence



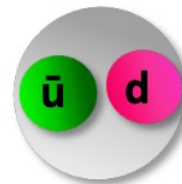
Different types of hadrons to be explored



Baryons are red-blue-green triplets

$\Lambda = usd$

Mesons are color-anticolor pairs

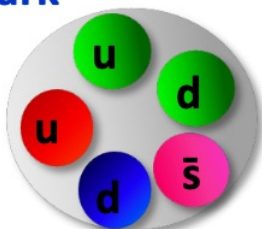


$\pi = \bar{u}d$

Other possible combinations of quarks and gluons :

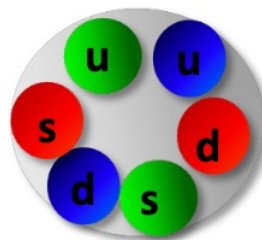
Pentaquark

$S = +1$
Baryon



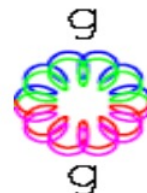
H di-Baryon

Tightly bound
6 quark state



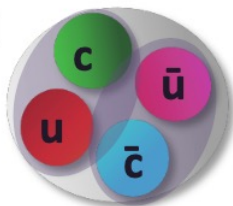
Glueball

Color-singlet multi-gluon bound state



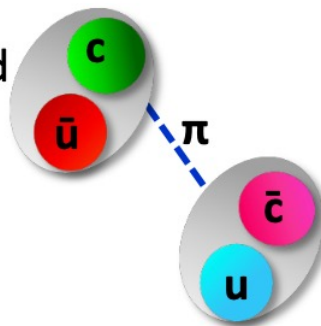
Tetraquark

Tightly bound
diquark &
anti-diquark

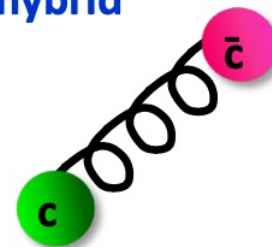


Molecule

loosely bound
meson-
antimeson
"molecule"



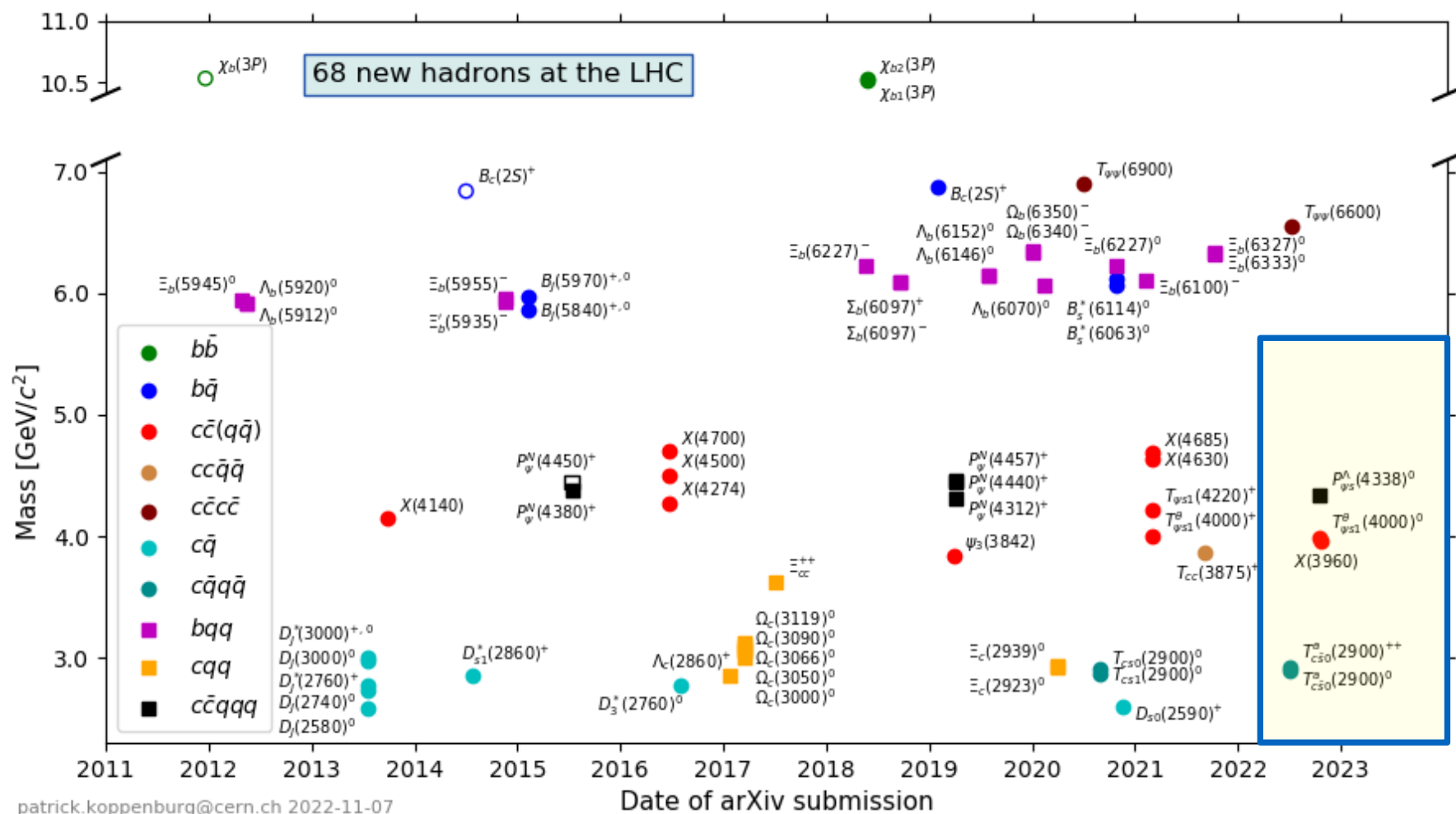
$q\bar{q}$ -gluon hybrid
mesons



Heavy hadron Spectroscopy



an example: LHC as a Large Hadron Discovery Factory



The LHCb detector and physics

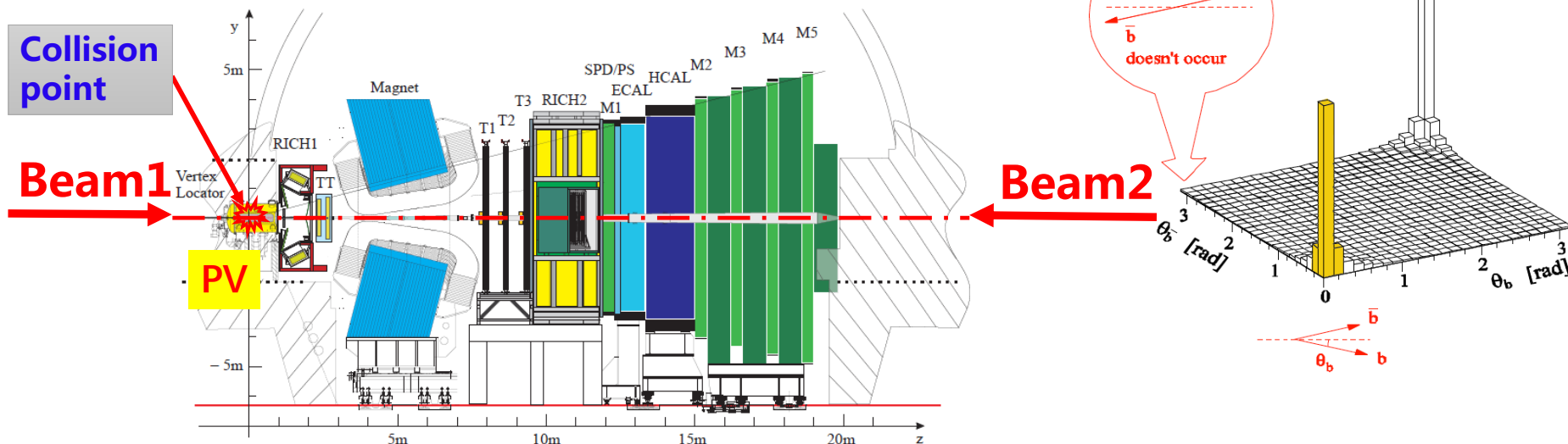


JINST 3 (2008) S08005

Int. J. Mod. Phys. A 30 (2015) 1530022

Pseudorapidity coverage

$$2 < \eta < 5$$



Excellent vertex and IP, decay time resolution:

- $\sigma(\text{IP}) \approx 20 \mu\text{m}$ for high- p_T tracks
- $\sigma(\tau) \approx 45 \text{ fs}$ for $B_s^0 \rightarrow J/\psi\phi$ and $B_s^0 \rightarrow D_s^- \pi^+$ decays

Very good momentum resolution:

- $\delta p/p \approx 0.5\% - 1\%$ for $p \in (0, 200) \text{ GeV}$
- $\sigma(m_B) \approx 24 \text{ MeV}$ for two-body decays

Hadron and Muon identification

- $\epsilon_{K \rightarrow K} \approx 95\%$ for $\epsilon_{\pi \rightarrow K} \approx 5\%$ up to 100 GeV
 - $\epsilon_{\mu \rightarrow \mu} \approx 97\%$ for $\epsilon_{\pi \rightarrow \mu} \approx 1 - 3\%$
- Data good for analyses
- $> 99\%$

Designed for CP violation and heavy flavor studies

X(3872) before LHC

$\chi_{c1}(3872)$: B factory era

X(3872)

$I^G(J^P) = ?^?(?^?)$

PDG2004

OMITTED FROM SUMMARY TABLE

Seen by CHOI 03 in $B \rightarrow K \pi^+ \pi^- J/\psi(1S)$ decays as a narrow peak in the invariant mass distribution of the $\pi^+ \pi^- J/\psi(1S)$ final state, but not seen in the $\gamma \chi_{c1}$ final state of these decays. Possibly absent in the invariant mass spectrum of the final state $\pi^+ \pi^- J/\psi(1S)$ in $e^+ e^-$ collisions. Interpretation as a 1^{--} charmonium state not favored.

Quantum numbers are not established.

X(3872) MASS

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
3872.0 ± 0.6 ± 0.5	36	CHOI	03 BELL	$B \rightarrow K \pi^+ \pi^- J/\psi$

X(3872) WIDTH

VALUE (MeV)	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<2.3	90	36	CHOI	03 BELL	$B \rightarrow K \pi^+ \pi^- J/\psi$

X(3872) DECAY MODES

Mode	Fraction (Γ_i/Γ)
Γ_1 $e^+ e^-$	
Γ_2 $\pi^+ \pi^- J/\psi(1S)$	seen
Γ_3 $\gamma \chi_{c1}$	

X(3872)

$I^G(J^P) = 0^?(?^+)$

PDG2012

Seen by CHOI 03 in $B \rightarrow K \pi^+ \pi^- J/\psi(1S)$ decays as a narrow peak in the invariant mass distribution of the $\pi^+ \pi^- J/\psi(1S)$ final state, but not seen in the $\gamma \chi_{c1}$ final state of these decays. Possibly absent in the invariant mass spectrum of the final state $\pi^+ \pi^- J/\psi(1S)$ in $e^+ e^-$ collisions. Interpretation as a 1^{--} charmonium state not favored. Isovector hypothesis excluded by AUBERT 05B and CHOI 11. A helicity amplitude analysis of the $X(3872) \rightarrow J/\psi \pi^+ \pi^-$ decay gives two possible J^P assignments: $J^P = 1^{++}$ and 2^{-+} (ABULENCIA 07E and CHOI 11). A study of the 3π invariant mass distribution in $J/\psi \omega$ decays slightly favors $J^P = 2^{-}$ (DEL-AMO-SANCHEZ 10B).

See our note on "Developments in Heavy Quarkonium Spectroscopy".

X(3872) MASS FROM $J/\psi X$ MODE

LHCb measured its mass using 2010 data (34.7 pb⁻¹)

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
3871.68 ± 0.17 OUR AVERAGE				
3871.95 ± 0.48 ± 0.12	0.6k	AALJ	12H LHCb	$p \bar{p} \rightarrow J/\psi \pi^+ \pi^- X$
3871.85 ± 0.27 ± 0.19	~170	¹ CHOI	11 BELL	$B \rightarrow K \pi^+ \pi^- J/\psi$
3873 ± 1.8 ± 1.6	±1.3	² DEL-AMO-SA.10B	BABR	$B \rightarrow \omega J/\psi K$
3871.61 ± 0.16 ± 0.19	6k	^{2,3} AALTONEN	09AU CDF2	$p \bar{p} \rightarrow J/\psi \pi^+ \pi^- X$
3871.4 ± 0.6 ± 0.1	93.4	AUBERT	08Y BABR	$B^+ \rightarrow K^+ J/\psi \pi^+ \pi^-$
3868.7 ± 1.5 ± 0.4	9.4	AUBERT	08Y BABR	$B^0 \rightarrow K_S^0 J/\psi \pi^+ \pi^-$
3871.8 ± 3.1 ± 3.0	522	^{2,4} ABAZOV	04F D0	$p \bar{p} \rightarrow J/\psi \pi^+ \pi^- X$
... We do not use the following data for averages, fits, limits, etc. ...				
3868.6				Mass: close to the $\bar{D}^{*0} D^0$ threshold
3871.3				Width: narrow
3873.4				C parity: +
3871.3				Isospin: 0
3872.0				
3836				

X(3872) during LHC Era

Produced in b -hadron decays: $B, B_s, \Lambda_b, \dots$ and prompt production in $e^+e^-/pp/p\bar{p}$ and heavy ion collision;

PDG2022

$\chi_{c1}(3872)$

$$J^G(J^{PC}) = 0^+(1^{++})$$

also known as $X(3872)$

This state shows properties different from a conventional $q\bar{q}$ state.
A candidate for an exotic structure. See the review on non- $q\bar{q}$ states.

$\chi_{c1}(3872)$ MASS FROM $J/\psi X$ MODE

Mass resolution: $0.06 \text{ MeV}/c^2$

VALUE (MeV)	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
3871.65 ± 0.06	0.06	OUR AVERAGE			
$3871.64 \pm 0.06 \pm 0.01$	19.8k	1	AAIJ	20S	LHCB $B^+ \rightarrow J/\psi \pi^+ \pi^- K^+$
$3871.9 \pm 0.7 \pm 0.2$	20		ABLIKIM	14	BES3 $e^+e^- \rightarrow J/\psi \pi^+ \pi^- \gamma$
$3871.95 \pm 0.48 \pm 0.12$	0.6k		AAIJ	12H	LHCB $pp \rightarrow J/\psi \pi^+ \pi^- X$
$3871.85 \pm 0.27 \pm 0.19$	170	2	CHOI	11	BELL $B \rightarrow K \pi^+ \pi^- J/\psi$
$3873 \pm 1.8 \pm 1.3$	27	3	DEL-AMO-SA...	10B	BABR $B \rightarrow \omega J/\psi K$
$3871.61 \pm 0.16 \pm 0.19$	6k	3,4	AALTONEN	09AU	CDF2 $p\bar{p} \rightarrow J/\psi \pi^+ \pi^- X$
$3871.4 \pm 0.6 \pm 0.1$	93.4		AUBERT	08Y	BABR $B^+ \rightarrow K^+ J/\psi \pi^+ \pi^-$

$\chi_{c1}(3872)$ WIDTH

Width: $1.19 \pm 0.21 \text{ MeV}/c^2$

VALUE (MeV)	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
1.19 ± 0.21	OUR AVERAGE		Error includes scale factor of 1.1.		
$1.39 \pm 0.24 \pm 0.10$	15.6k	¹	AAIJ	20AD	LHCB $pp \rightarrow J/\psi \pi^+ \pi^- X$
$0.96 \pm 0.19 \pm 0.21$	4.2k	²	AAIJ	20S	LHCB $B^+ \rightarrow J/\psi \pi^+ \pi^- K^+$
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●					
<2.4	90		ABLIKIM	14	BES3 $e^+e^- \rightarrow J/\psi \pi^+ \pi^- \gamma$
<1.2	90		CHOI	11	BELL $B \rightarrow K \pi^+ \pi^- J/\psi$
<3.3	90		AUBERT	08Y	BABR $B^+ \rightarrow K^+ J/\psi \pi^+ \pi^-$
<4.1	90	69	AUBERT	06	BABR $B \rightarrow K \pi^+ \pi^- J/\psi$
<2.3	90	36	³ CHOI	03	BELL $B \rightarrow K \pi^+ \pi^- J/\psi$

... We do not use the following data for averages, fits, limits, etc. ...

¹ Using $\chi_{c1}(3872)$ produced in inclusive b -hadron decays. Breit-Wigner parametrization.

² Using Breit-Wigner parametrization. Partially overlapping dataset with that of AAJ 20AD.

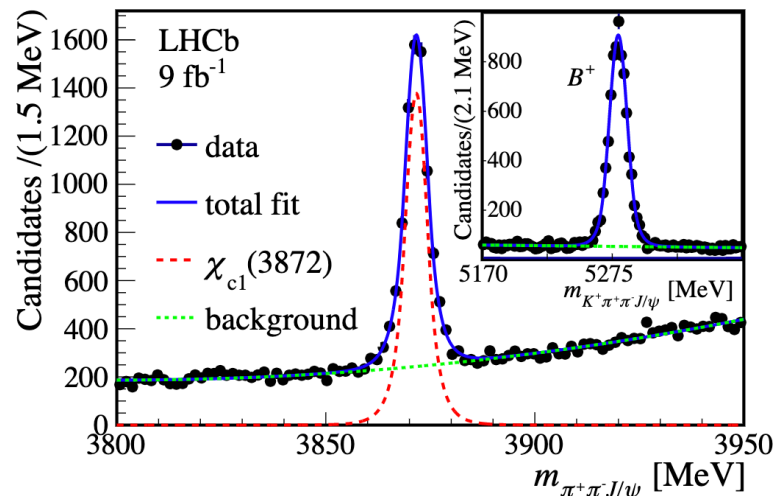
³ Superseded by CHOI 11.

	Mode	Fraction (Γ_i / Γ)
Γ_1	e^+e^-	$< 2.8 \times 10^{-6}$
Γ_2	$\pi^+\pi^- J/\psi(1S)$	$(3.8 \pm 1.2)\%$
Γ_3	$\pi^+\pi^-\pi^0 J/\psi(1S)$	not seen
Γ_4	$\omega\eta_c(1S)$	$< 33\%$
Γ_5	$\omega J/\psi(1S)$	$(4.3 \pm 2.1)\%$
Γ_6	$\phi\phi$	not seen
Γ_7	$D^0\bar{D}^0\pi^0$	$(49^{+18}_{-20})\%$
Γ_8	$\bar{D}^{*0}D^0$	$(37 \pm 9)\%$
Γ_9	$\gamma\gamma$	$< 11\%$
Γ_{10}	$D^0\bar{D}^0$	$< 29\%$
Γ_{11}	D^+D^-	$< 19\%$
Γ_{12}	$\pi^0\chi_{c2}$	$< 4\%$
Γ_{13}	$\pi^0\chi_{c1}$	$(3.4 \pm 1.6)\%$
Γ_{14}	$\pi^0\chi_{c0}$	$< 70\%$
Γ_{15}	$\pi^+\pi^-\eta_c(1S)$	$< 14\%$
Γ_{16}	$\pi^+\pi^-\chi_{c1}$	$< 7 \times 10^{-3}$
Γ_{17}	$p\bar{p}$	$< 2.4 \times 10^{-5}$
▼ Radiative decays		
Γ_{18}	γD^+D^-	$< 4\%$
Γ_{19}	$\gamma\bar{D}^0D^0$	$< 6\%$
Γ_{20}	$\gamma J/\psi$	$(8 \pm 4) \times 10^{-3}$
Γ_{21}	$\gamma\chi_{c1}$	$< 9 \times 10^{-3}$
Γ_{22}	$\gamma\chi_{c2}$	$< 3.2\%$
Γ_{23}	$\gamma\psi(2S)$	$(4.5 \pm 2.0)\%$

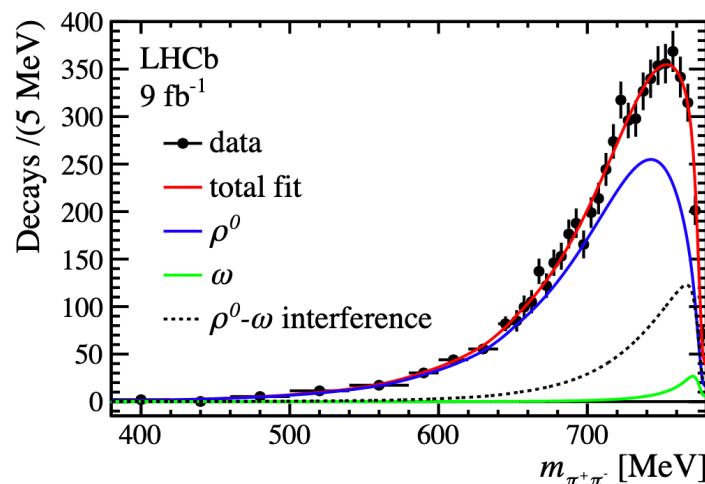
- Loosely $D^0\bar{D}^{0*}$ bound state?
- Mixture of $\chi_{c1}(2P)$ and $D^0\bar{D}^{0*}$?

Observation of sizeable ω contribution in $X(3872) \rightarrow \pi^+ \pi^- J/\psi$

arXiv:2204.12597



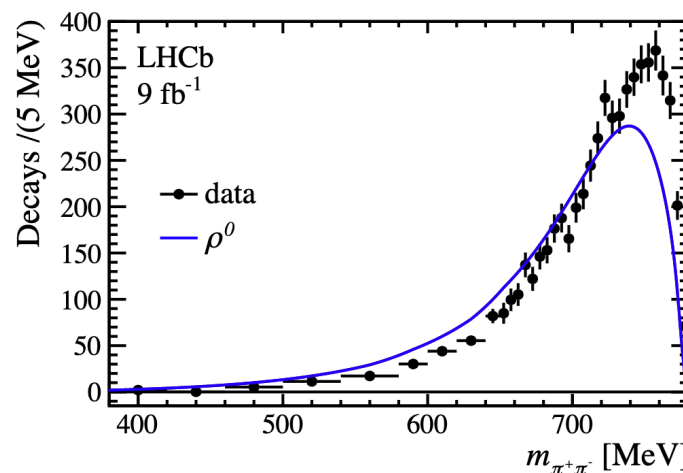
$B^+ \rightarrow K^+ X(3872)$ studied with RUN 1+2 data



$$\frac{g_{\chi_{c1}(3872) \rightarrow \rho^0 J/\psi}}{g_{\chi_{c1}(3872) \rightarrow \omega J/\psi}} = \sqrt{\frac{\mathcal{B}(\omega \rightarrow \pi^+ \pi^-)}{\mathcal{R}_{\omega/\rho}^0}} = 0.29 \pm 0.04.$$

$$\frac{g_{\psi(2S) \rightarrow \pi^0 J/\psi}}{g_{\psi(2S) \rightarrow \eta J/\psi}} = \sqrt{\frac{\mathcal{B}(\psi(2S) \rightarrow \pi^0 J/\psi)}{\mathcal{B}(\psi(2S) \rightarrow \eta J/\psi)}} \frac{p_{\eta}^3}{p_{\pi^0}^3} = 0.045 \pm 0.001$$

The isospin violating ρ^0 contribution, quantified for the first time with proper subtraction of the ω contribution, is an order of magnitude too large for $X(3872)$ to be a pure charmonium state.



Open flavor tetraquark

- The $D_{s0}^*(2317)^+ (D_s^+ \pi^0)$ state was observed in 2003.
- It is argued to contain some **tetraquark component** in several theoretical descriptions whose $I = 1$ partners can exist in the $D_s^+ \pi^\pm$ final states.
- [*Cheng & Hou*](#): It would be astonishing if a doubly charged resonance is found. [PLB 566, 193 (2003)]

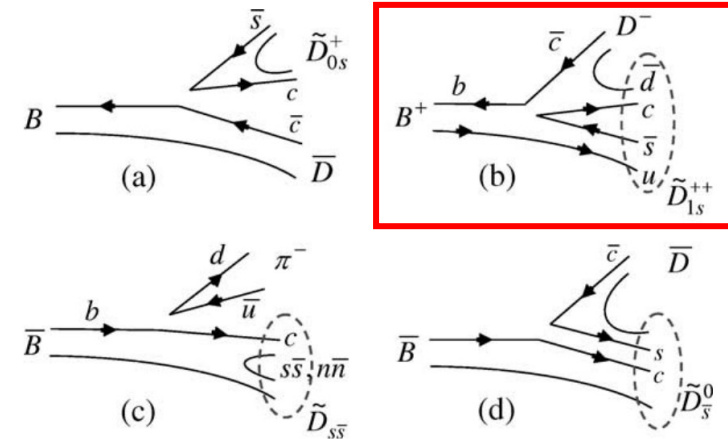
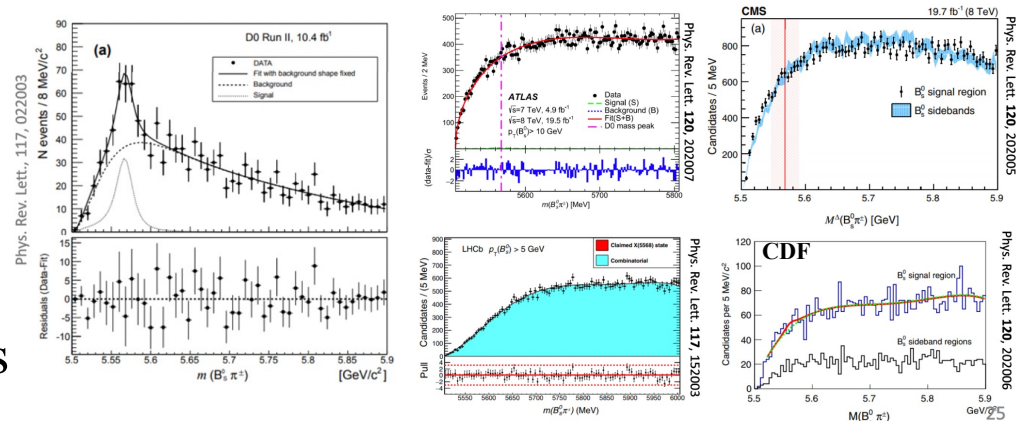


Fig. 2. Diagrams for (a) $B \rightarrow \bar{D} \bar{D}_{0s}^+$, (b) $B^+ \rightarrow D^- \bar{D}_{1s}^{++}$ ($B \rightarrow \bar{D} \bar{D}_{1s}$), (c) $\bar{B} \rightarrow \pi^- \bar{D}_{S\bar{S}}, \pi^- \bar{D}$, (d) $B \rightarrow D \bar{D}_S^0$.

- D0 claimed evidence for the X(5568) in decaying to $B_s \pi^+$, interpreted as tetraquark state $[b\bar{s}u\bar{d}]$
- But not seen in other experiments



Observation of a doubly charged tetraquark $T_{c\bar{s}0}(2900)^{++}$ [$c\bar{s}u\bar{d}$] and its neutral partner $T_{c\bar{s}0}(2900)^0$ [$c\bar{s}u\bar{d}$]

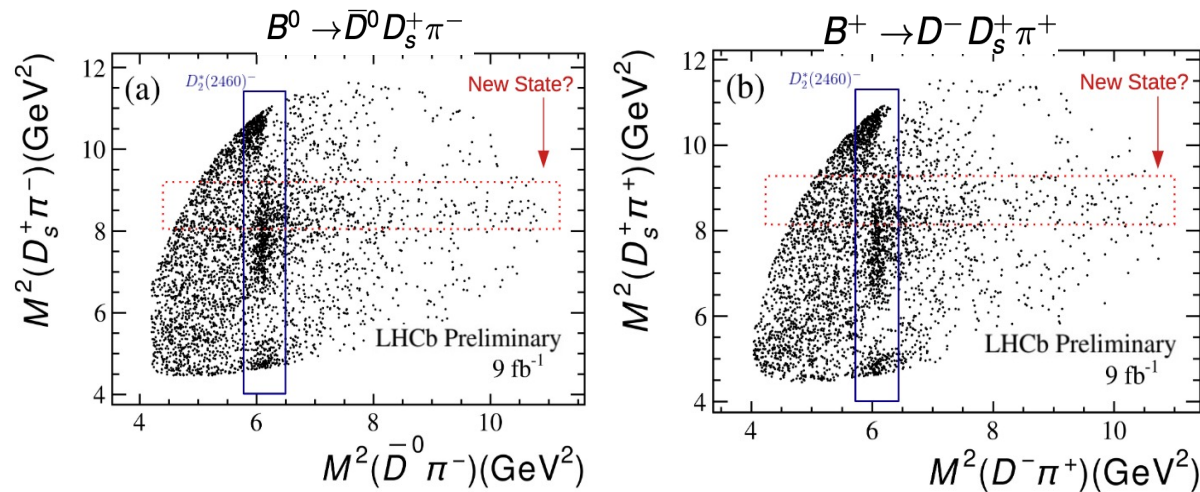
LHCb-PAPER-2022-026
LHCb-PAPER-2022-027

- First simultaneous amplitude analysis of $B^+ \rightarrow D^- D_s^+ \pi^+$ & $B^0 \rightarrow \bar{D}^0 D_s^+ \pi^-$ with RUN 1+2 9 fb⁻¹ data

- $D_s \pi$ mass spectra well described by adding $J^P = 0^+$ ($> 7.5 \sigma$) $T_{c\bar{s}0}^a(2900) > 9 \sigma$

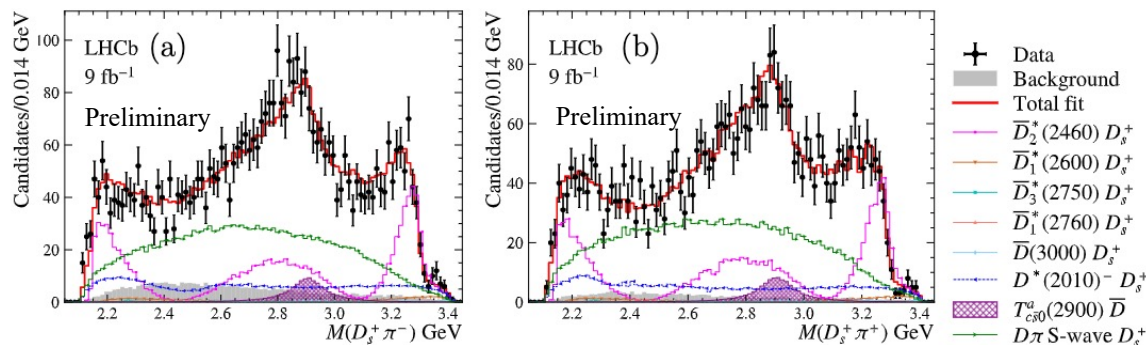
$$M = 2.908 \pm 0.011 \pm 0.020 \text{ GeV}$$

$$\Gamma = 0.136 \pm 0.023 \pm 0.011 \text{ GeV}$$



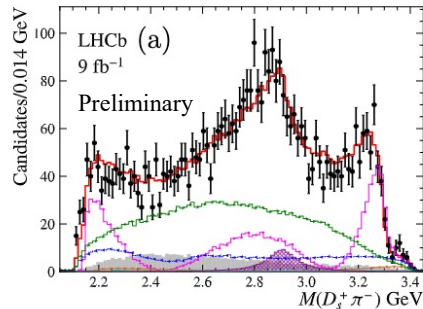
- Separate resonance fits:

- $T_{c\bar{s}}^{a0}$: $m = 2892 \pm 14 \pm 15 \text{ MeV}$,
 $\Gamma = 119 \pm 26 \pm 12 \text{ MeV}$;
- $T_{c\bar{s}}^{a++}$: $m = 2921 \pm 17 \pm 19 \text{ MeV}$,
 $\Gamma = 137 \pm 32 \pm 14 \text{ MeV}$

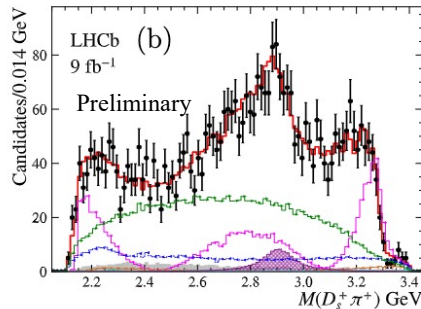


$T_{c\bar{s}0}^a(2900)$ and $X_{0,1}(2900)$

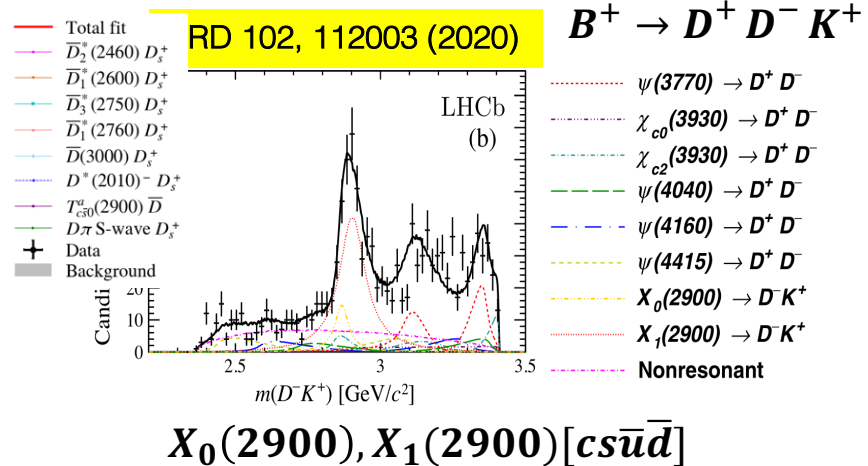
LHCb-PAPER-2022-026 ; LHCb-PAPER-2022-027



$T_{c\bar{s}0}^a(2900)^0 [c\bar{s}u\bar{d}]$



$T_{c\bar{s}0}^a(2900)^{++} [c\bar{s}u\bar{d}]$



$X_0(2900), X_1(2900) [c\bar{s}u\bar{d}]$

	Mass (GeV)	Width (GeV)	J^P
$T_{c\bar{s}0}^a(2900)^0$ & $T_{c\bar{s}0}^a(2900)^{++}$	$2.908 \pm 0.011 \pm 0.020$	$0.136 \pm 0.023 \pm 0.020$	0^+
$X_0(2900)/T_{cs0}(2900)$	$2.866 \pm 0.007 \pm 0.002$	$0.057 \pm 0.012 \pm 0.004$	0^+
$X_1(2900)/T_{cs1}(2900)$	$2.904 \pm 0.005 \pm 0.001$	$0.110 \pm 0.011 \pm 0.004$	1^-

• $T_{c\bar{s}0}^a(2900)$ v.s. $X_0(2900)$

- ✓ Similar mass, but width and flavor contents are different.
- ✓ $T_{c\bar{s}1}^a(2900)?$
- ✓ $T_{c\bar{s}0}^a(2900)^{++} \rightarrow D^+ K^+?$
- ✓ $T_{c\bar{s}0}^a(2900)^+ \rightarrow D_s^+ \pi^0, D_s^+ \pi^+ \pi^-?$

- **no isospin relation:** $[c\bar{s}u\bar{d}]$ v.s. $[c\bar{s}u\bar{d}]$
- **U-spin relation:** $[c\bar{s}u\bar{d}]$ v.s. $[c\bar{d}u\bar{s}]$
- $T_{c\bar{s}0}^a(2900)$ mass and width larger than $T_{cs0}(2900)$

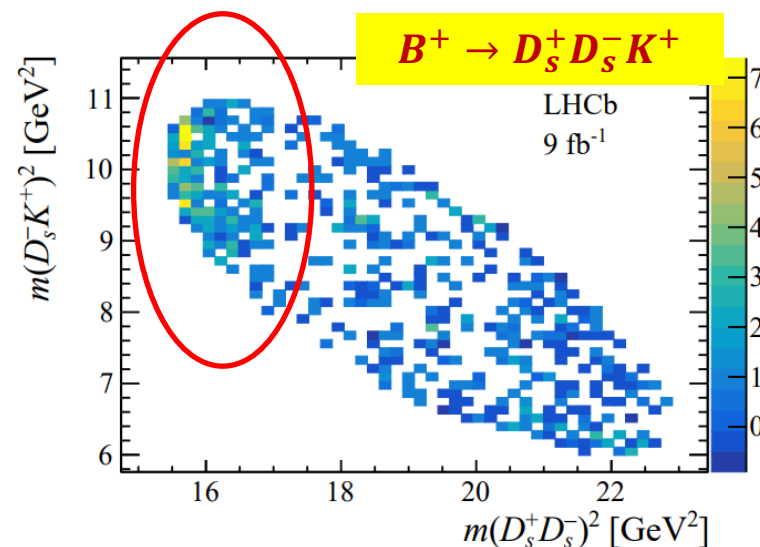
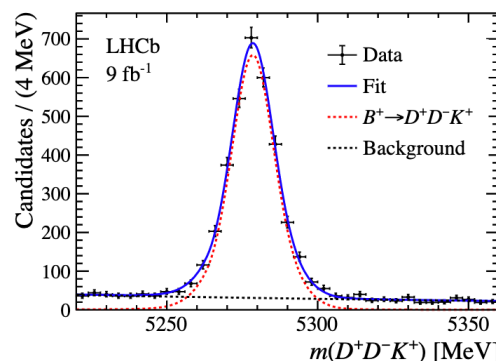
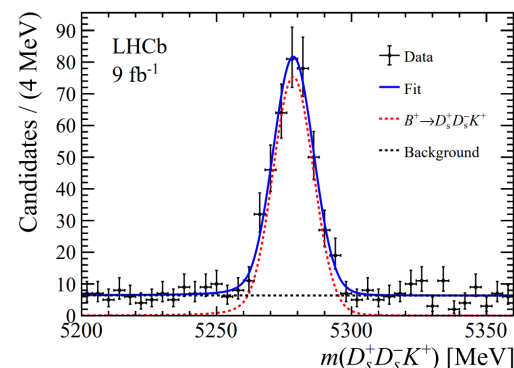
Study on $B^+ \rightarrow D_s^+ D_s^- K^+$

arXiv:2211.05034

arXiv:2210.15153

- Relative measurement of branching fractions :

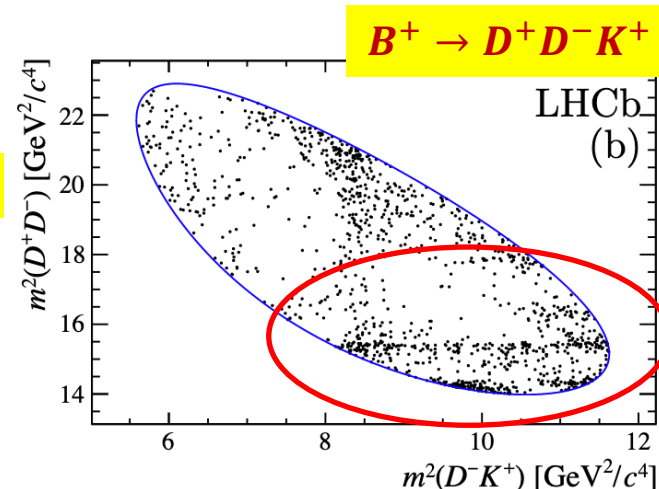
$$\frac{\mathcal{B}(B^+ \rightarrow D_s^+ D_s^- K^+)}{\mathcal{B}(B^+ \rightarrow D^+ D^- K^+)} = 0.525 \pm 0.033 \pm 0.027 \pm 0.034$$



- $D_s^+ D_s^-$ near-threshold enhancement is seen
- Similar to the $\chi_{c0,2}(3930)$ observed in $B^+ \rightarrow D^+ D^- K^+$?

PRL125, 242001 (2020)

Resonance	Mass (GeV/ c^2)	Width (MeV)
$\chi_{c0}(3930)$	$3.9238 \pm 0.0015 \pm 0.0004$	$17.4 \pm 5.1 \pm 0.8$
$\chi_{c2}(3930)$	$3.9268 \pm 0.0024 \pm 0.0008$	$34.2 \pm 6.6 \pm 1.1$

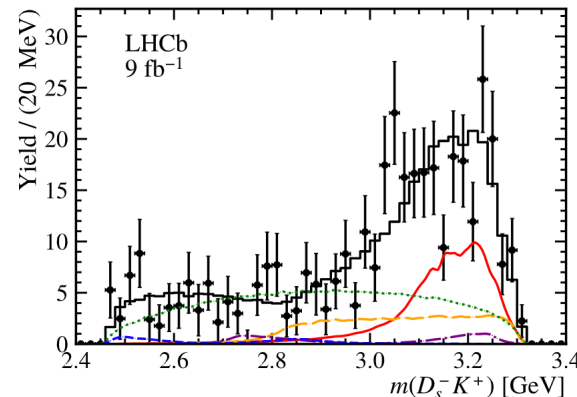
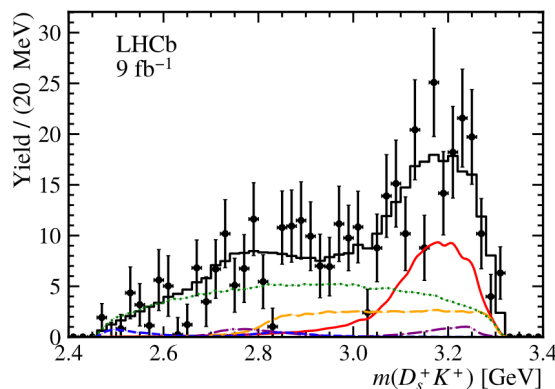
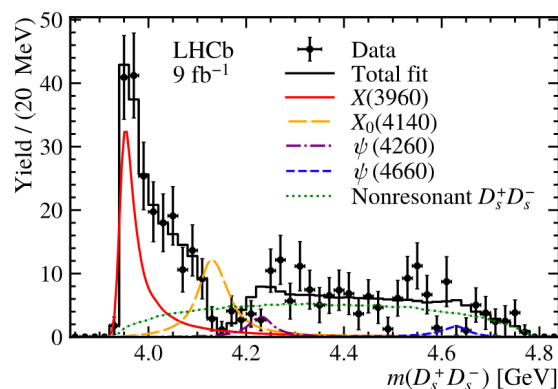


New $[c\bar{c}s\bar{s}]$ state in $D_s^+ D_s^-$

arXiv:2211.05034

arXiv:2210.15153

- Near threshold structure $X(3960)$: 12σ , $J^P = 0^{++}$
- $X_0(4140)$ accounts for the dip around 4.14 GeV



Components	J^{PC}	Mass (MeV)	Width (MeV)	Fit fraction (%)	Significance (σ)
$X(3960)$	0^{++}	$3956 \pm 5 \pm 10$	$43 \pm 13 \pm 8$	$25.4 \pm 7.7 \pm 5.0$	12.6 (14.6)
$X_0(4140)$	0^{++}	$4133 \pm 6 \pm 6$	$67 \pm 17 \pm 7$	$16.7 \pm 4.7 \pm 3.9$	3.8 (4.1)
$\psi(4260)$	1^{--}	4230 (fixed)	55 (fixed)	$3.6 \pm 0.4 \pm 3.2$	3.2 (3.6)
$\psi(4660)$	1^{--}	4633 (fixed)	64 (fixed)	$2.2 \pm 0.2 \pm 0.8$	3.0 (3.2)
NR	S-wave	—	—	$46.1 \pm 13.2 \pm 11.3$	3.1 (3.4)

dip at 4.14 GeV
via interference

$$\frac{\mathcal{B}(X \rightarrow D^+ D^-)}{\mathcal{B}(X \rightarrow D_s^+ D_s^-)} = 0.29 \pm 0.09 \pm 0.10 \pm 0.08$$

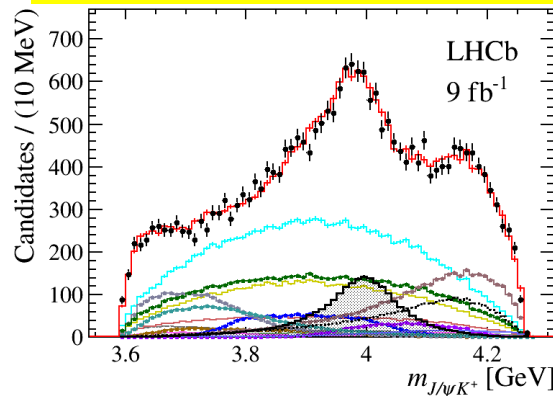
Resonance	Mass (GeV/ c^2)	Width (MeV)
$\chi_{c0}(3930)$	$3.9238 \pm 0.0015 \pm 0.0004$	$17.4 \pm 5.1 \pm 0.8$

If $X(3960)$ and $\chi_{c0}(3930)$ the same particle, the ratio (<1) indicates intrinsic $[s\bar{s}]$ components. Hence, they could be exotic $[c\bar{c}s\bar{s}]$ tetraquark state.

Z_{cs} [$c\bar{c}u\bar{s}$] states

- Charged Z_{cs} states are observed at BESIII and LHCb: $Z_{cs}(3985)$, $Z_{cs}(4000)$, $Z_{cs}(4220)$
- It is natural to search for the neutral isospin partners

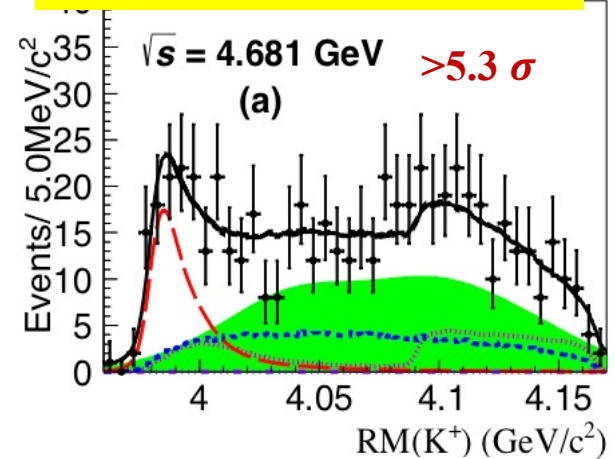
LHCb, PRL127, 082001 (2021)



All $Z_{cs}(1^+)$				$25 \pm 5^{+11}_{-12}$
$Z_{cs}(4000)$	15 (16)	$4003 \pm 6^{+4}_{-14}$	$131 \pm 15 \pm 26$	$9.4 \pm 2.1 \pm 3.4$
$Z_{cs}(4220)$	5.9 (8.4)	$4216 \pm 24^{+43}_{-30}$	$233 \pm 52^{+97}_{-73}$	$10 \pm 4^{+10}_{-7}$

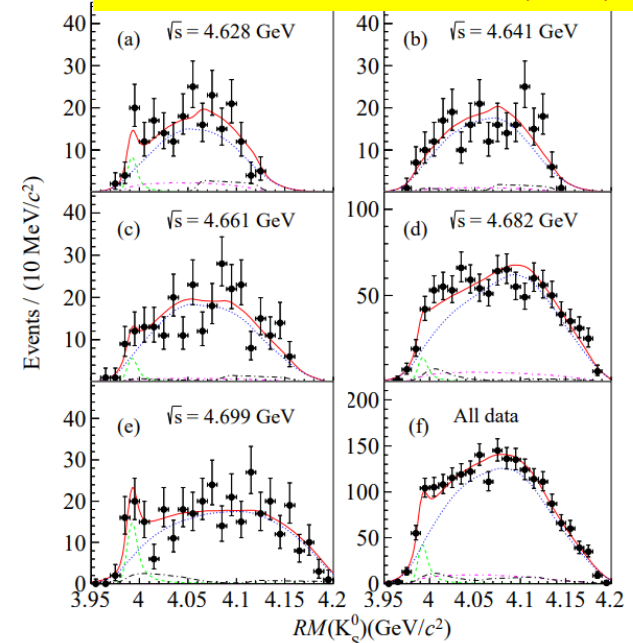
	Mass (MeV/c ²)	Width (MeV)
$Z_{cs}(3985)^0$	$3992.2 \pm 1.7 \pm 1.6$	$7.7^{+4.1}_{-3.8} \pm 4.3$
$Z_{cs}(3985)^+$	$3985.2^{+2.1}_{-2.0} \pm 1.7$	$13.8^{+8.1}_{-5.2} \pm 4.9$

BESIII, PRL 126, 102001 (2021)



$$e^+e^- \rightarrow K_S^0 (D_s^- D^{*+} + D_s^{*-} D^+)$$

BESIII, PRL 129, 112003 (2022)

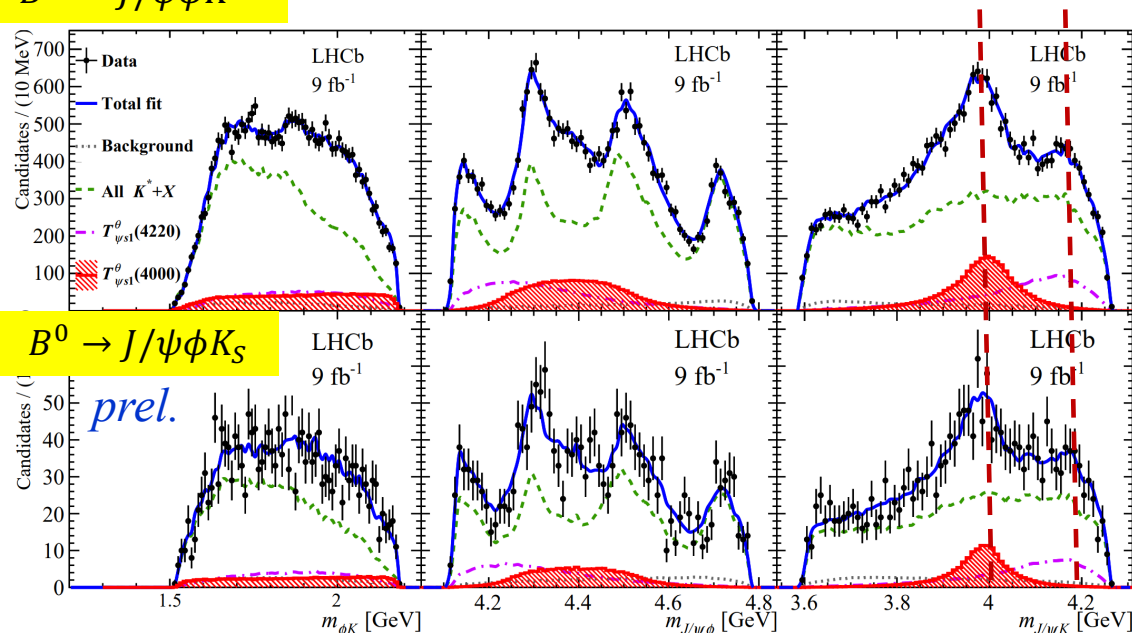


$T_{\psi s 1}^{\theta}(4000)^0$ in $B^0 \rightarrow J/\psi \phi K_S^0$

$B^+ \rightarrow J/\psi \phi K^+$

LHCb-PAPER-2022-040

preliminary results



simultaneous fits to $B^0 \rightarrow J/\psi \phi K_S$ and $B^+ \rightarrow J/\psi \phi K^+$, assuming isospin symmetry for all the intermediate states, except for the charged and neutral $T_{\psi s 1}^{\theta}(4000)$ states.

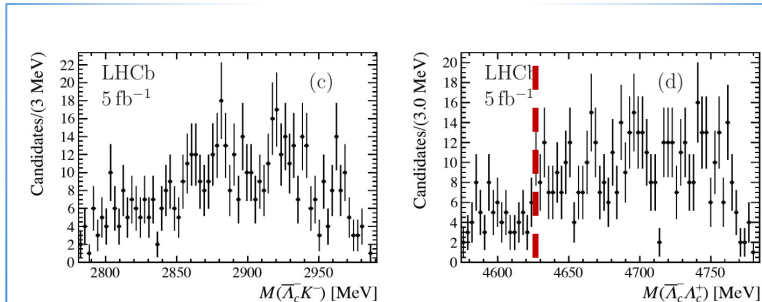
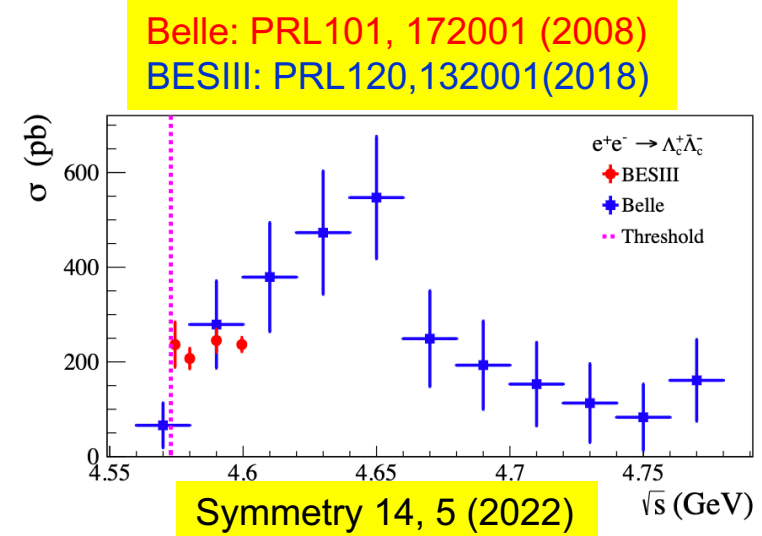
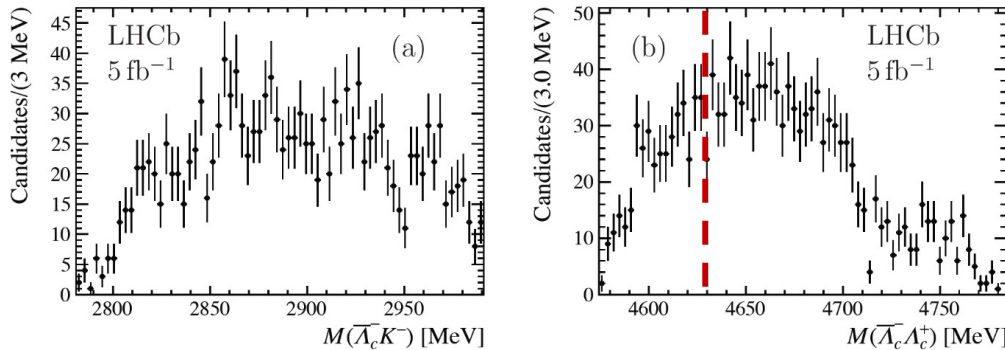
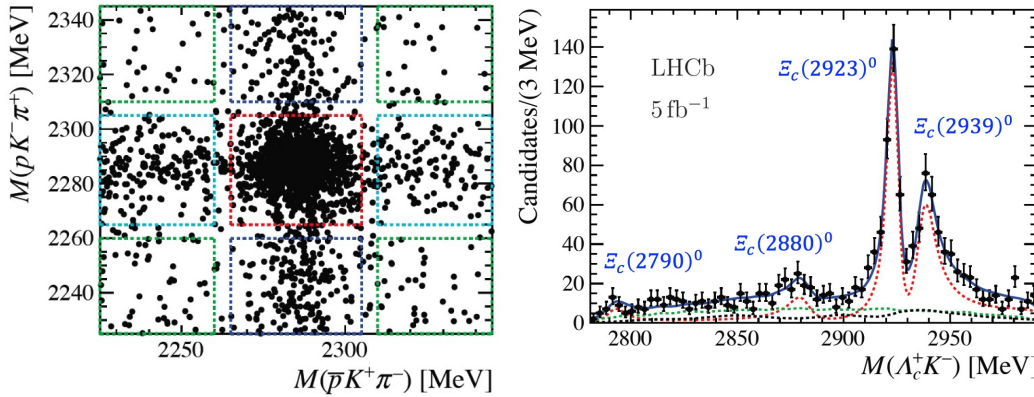
- Consistent with being isospin partners: $\Delta m = -12.1^{+11.1+6.0}_{-10.2-4.2} \text{ MeV}/c^2$
- Significance is 4.0σ without isospin symmetry for $T_{\psi s 1}^{\theta}(4000)$, while 5.4σ with isospin symmetry constraints

	J^P	Mass (MeV/c^2)	Width (MeV)	Fit fraction
<i>prel.</i> $T_{\psi s 1}^{\theta}(4000)^0 \rightarrow J/\psi K_S^0$	1^+	$3991.3^{+11.7+8.5}_{-10.4-16.7}$	$104.8^{+29.3+17.1}_{-25.3-23.3}$	$7.9 \pm 2.5^{+3.0}_{-2.8}$
$Z_{cs}^+ / T_{\psi s 1}^{\theta}(4000)^+ \rightarrow J/\psi K^+$	1^+	$4003 \pm 6^{+4}_{-14}$	$131 \pm 15 \pm 26$	$9.4 \pm 2.1 \pm 3.4$

Search for exotics in $B^- \rightarrow \Lambda_c^+ \bar{\Lambda}_c^- K^-$

arXiv:2211.00812

- $\Lambda_c^+ \bar{\Lambda}_c^-$ and $\Lambda_c^- K^-$ systems are good places to search for exotics
 - Near threshold enhancement of $Y(4630) \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-$ is observed at Belle.
- However, the Breit-Wigner line shape is not supported by BESIII



vetoing $2.90 \text{ GeV} < M(\Lambda_c^+ K^-) < 2.97 \text{ GeV}$

No obvious structures are seen in $\Lambda_c^+ \bar{\Lambda}_c^-$ and $\Lambda_c^- K^-$ systems

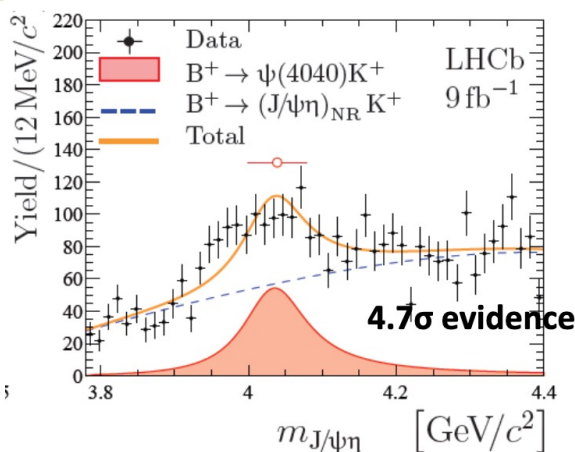
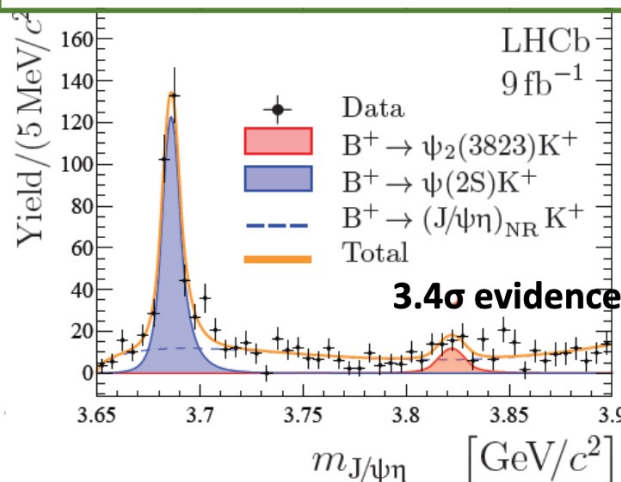
Charmonium(-like) states in $B^+ \rightarrow K^+ \eta J/\psi$

JHEP 22, 046 (2022)

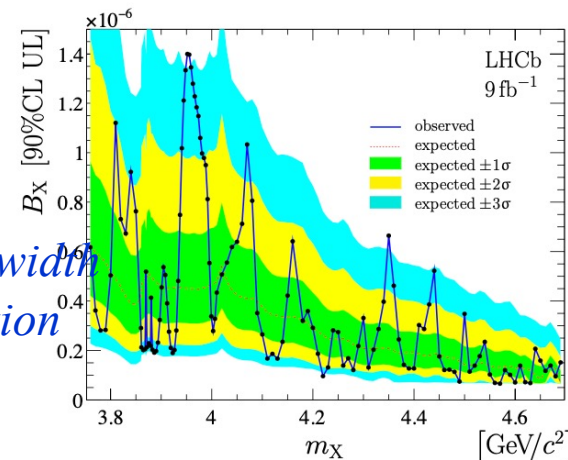
- With RUN 1+2 LHCb data, evidence is found for $\psi_2(3823) \rightarrow \eta J/\psi$ and $\psi(4040) \rightarrow \eta J/\psi$
- Searches for the known XYZ states and provides UL on BF in $B^+ \rightarrow K^+ X$

$$F_X \equiv \frac{\mathcal{B}(B^+ \rightarrow XK^+) \times \mathcal{B}(X \rightarrow J/\psi \eta)}{\mathcal{B}(B^+ \rightarrow \psi(2S)K^+) \times \mathcal{B}(\psi(2S) \rightarrow J/\psi \eta)},$$

$$B_X \equiv \mathcal{B}(B^+ \rightarrow XK^+) \times \mathcal{B}(X \rightarrow J/\psi \eta)$$



*narrow width
assumption*



Upper limit at 90% CL
 F_X [10^{-2}] B_X [10^{-7}]

$\psi(3770)$	2.2	4.6
$\psi_3(3842)$	2.9	6.1
$\psi(4160)$	4.2	8.7
$\psi(4415)$	4.6	9.6

R(3760)	2.0	4.1
R(3790)	3.2	6.7
$Z_c(3900)^0$	2.1	4.3
$\psi(4230)$	1.9	3.9
$\psi(4360)$	6.0	12.4
$\psi(4390)$	11.6	24.1

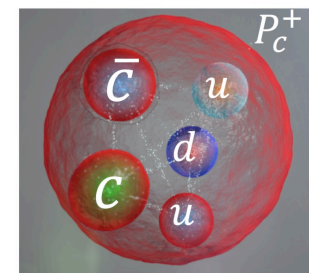
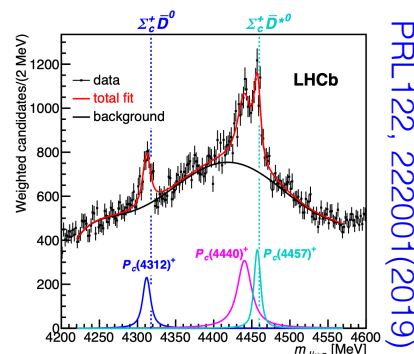
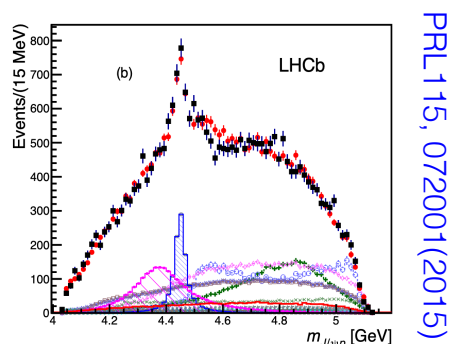
$Z_c(4430)^0$	6.1	12.7
X'_C	1.9	3.9

$$F_{\psi_2(3823)} = (5.95^{+3.38}_{-2.55}) \times 10^{-2}, \quad B_{\psi_2(3823)} = (1.25^{+0.71}_{-0.53} \pm 0.04) \times 10^{-6},$$

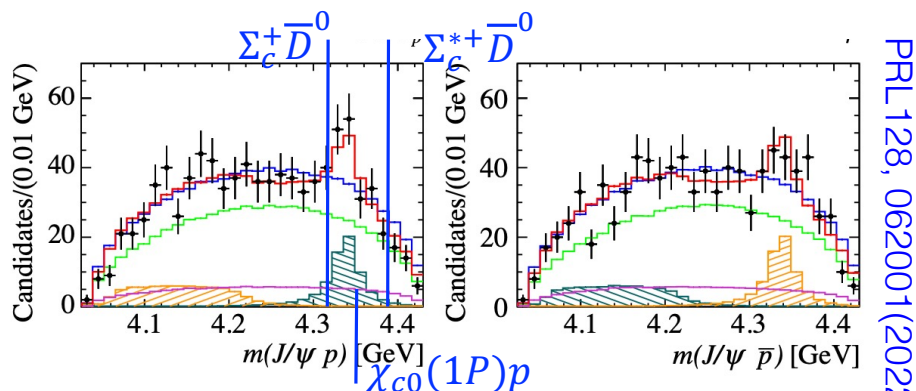
$$F_{\psi(4040)} = (40.6 \pm 11.2) \times 10^{-2}, \quad B_{\psi(4040)} = (8.53 \pm 2.35 \pm 0.30) \times 10^{-6}$$

Pentaquark states at LHCb

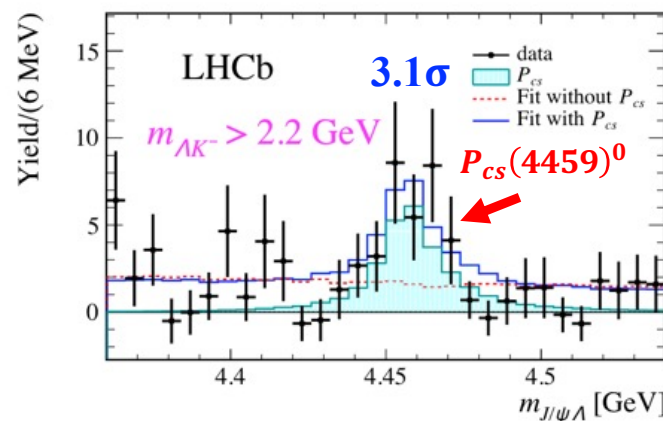
- Observation of $[\bar{c}uud]$ pentaquarks: $P_{\psi}^N(4312)^+$, $P_{\psi}^N(4440)^+$, $P_{\psi}^N(4457)^+$ in $\Lambda_b^0 \rightarrow J/\psi p K^-$ decays; near thresholds of $\Sigma_c^+ \bar{D}^0$, $\Sigma_c^+ \bar{D}^{*0}$, J^P not determined



- Evidence of $[\bar{c}uud]$ pentaquark: $P_{\psi}^N(4337)^+$ in $B_s^0 \rightarrow J/\psi p \bar{p}$ decays



- Evidence for $[\bar{c}cuds]$ pentaquark candidate with strangeness: $P_{\psi s}^{\Lambda}(4459)^0$ in $\Xi_b^- \rightarrow J/\psi \Lambda K^-$ decays, near threshold of $\Xi_c^0 \bar{D}^{*0}$

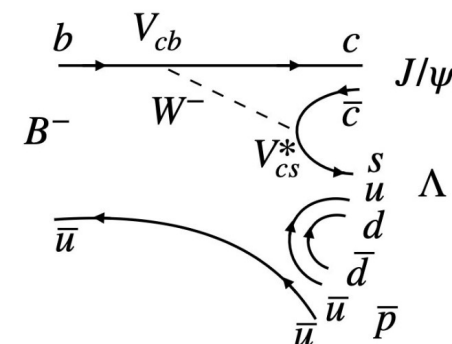


Sci. Bull. 66, 1278(2021)

Observation of the hidden-charm strange pentaquark [$c\bar{c}uds$]

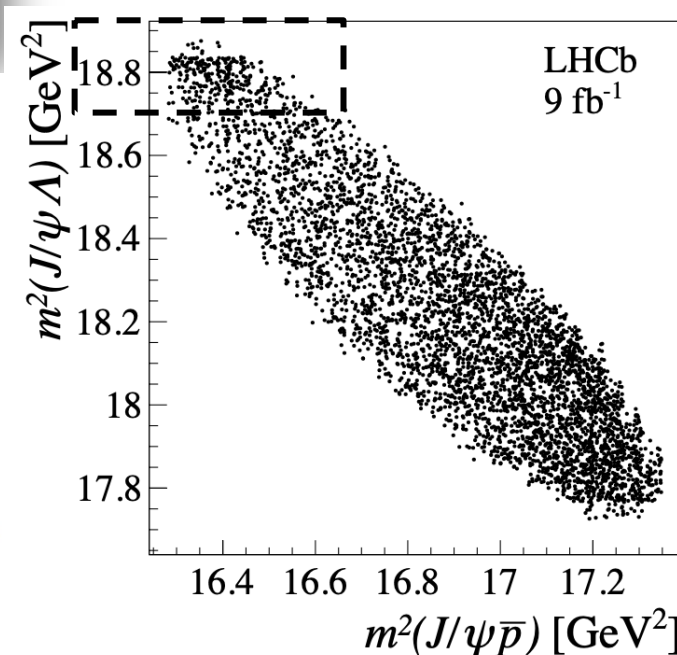
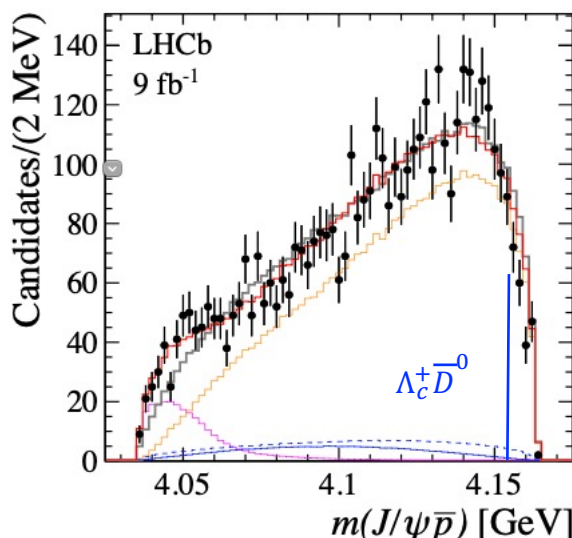
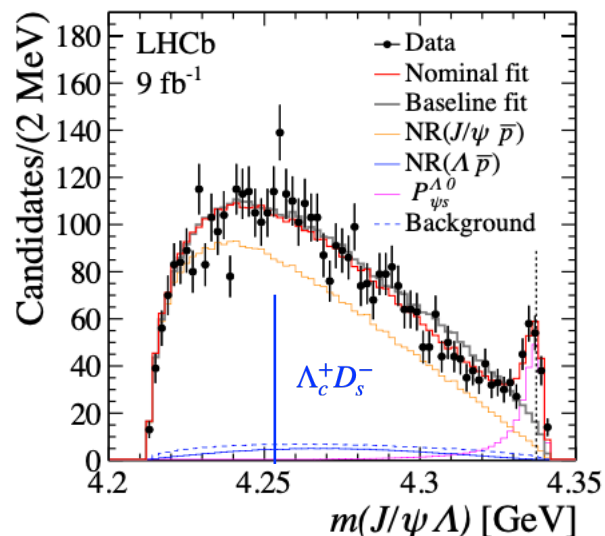
arXiv:2210.10346

- narrow structure in $J/\psi\Lambda$ in $B^- \rightarrow J/\psi\Lambda\bar{p}$, with 9 fb^{-1} LHCb data
- amplitude analysis is performed
- $P_{\psi s}^{\Lambda}(4338)^0 \rightarrow J/\psi\Lambda$ observed with significance larger than 10σ
- $J^P = \frac{1}{2}^-$ preferred and close to $\Xi_c^+ D^-$ threshold
 - 0.8 MeV above $\Xi_c^+ D^-$;
 - 2.9 MeV above $\Xi_c^0 \bar{D}^0$



$$M_{P_{cs}} = 4338.2 \pm 0.7 \pm 0.4 \text{ MeV}$$

$$\Gamma_{P_{cs}} = 7.0 \pm 1.2 \pm 1.3 \text{ MeV}$$



no evidence for $P_{\psi s}^{\Lambda}(4255)^0$ and $P_{\psi}^N(4152)^+$

Summary

- 3 tetraquark states and 1 pentaquark state observed at LHCb in 2022
 - $T_{c\bar{s}0}^a(2900)^{0,++}$ [$c\bar{s}u\bar{d}$] in $B \rightarrow DD_s \pi$
 - first doubly charged tetraquark candidate
 - $X(3960)$ [$c\bar{c}s\bar{s}$] in $B^+ \rightarrow D_s^+ D_s^- K^+$
 - $T_{\psi s1}^\theta(4000)^0$ [$c\bar{c}d\bar{s}$] in $B^0 \rightarrow J/\psi \phi K_S^0$
 - isospin partner of $Z_{cs}(4000)$
 - Observation of new pentaquark state $P_{\psi s}^\Lambda(4338)$ [$c\bar{c}uds$]
 - first open-strange pentaquark candidate
- Observation of $\chi_{c1}(3872) \rightarrow \omega J/\psi$ in $\chi_{c1}(3872) \rightarrow \pi^+ \pi^- J/\psi$
- More data in Run 3 are desired for further improvement on understanding their natures





Thank you!

谢谢!

Exotic hadron naming convention



LHCb-PUB-2022-013, arXiv:2206.15233

Minimal quark content	Current name	$I^{(G)}, J^{P(C)}$	Proposed name
$c\bar{c}$	$\chi_{c1}(3872)$	$I^G = 0^+, J^{PC} = 1^{++}$	$\chi_{c1}(3872)$
$c\bar{c}u\bar{d}$	$Z_c(3900)^+$	$I^G = 1^+, J^P = 1^+$	$T_{\psi 1}^b(3900)^+$
$c\bar{c}u\bar{d}$	$X(4100)^+$	$I^G = 1^-$	$T_{\psi}(4100)^+$
$c\bar{c}u\bar{d}$	$Z_c(4430)^+$	$I^G = 1^+, J^P = 1^+$	$T_{\psi 1}^b(4430)^+$
$c\bar{c}(s\bar{s})$	$\chi_{c1}(4140)$	$I^G = 0^+, J^{PC} = 1^{++}$	$\chi_{c1}(4140)$
$c\bar{c}u\bar{s}$	$Z_{cs}(4000)^+$	$I = \frac{1}{2}, J^P = 1^+$	$T_{\psi s 1}^{\theta}(4000)^+$
$c\bar{c}u\bar{s}$	$Z_{cs}(4220)^+$	$I = \frac{1}{2}, J^P = 1^?$	$T_{\psi s 1}(4220)^+$
$c\bar{c}c\bar{c}$	$X(6900)$	$I^G = 0^+, J^{PC} = ?^{?+}$	$T_{\psi\psi}(6900)$
$cs\bar{u}\bar{d}$	$X_0(2900)$	$J^P = 0^+$	$T_{cs 0}(2900)^0$
$cs\bar{u}\bar{d}$	$X_1(2900)$	$J^P = 1^-$	$T_{cs 1}(2900)^0$
$cc\bar{u}\bar{d}$	$T_{cc}(3875)^+$		$T_{cc}(3875)^+$
$bb\bar{u}\bar{d}$	$Z_b(10610)^+$	$I^G = 1^+, J^P = 1^+$	$T_{\tau 1}^b(10610)^+$
$c\bar{c}uud$	$P_c(4312)^+$	$I = \frac{1}{2}$	$P_{\psi}^N(4312)^+$
$c\bar{c}uds$	$P_{cs}(4459)^0$	$I = 0$	$P_{\psi s}^{\Lambda}(4459)^0$

□ P for pentaquarks, T for tetraquarks

□ P states: i.e. $P_{\psi}^N(4312)^+$, $P_{\psi s}^{\Lambda}(4459)^0$

Superscript: denotes isospin: $\Lambda, N, \Sigma, \Delta$ for $I = 0, 1/2, 1, 3/2$

Subscript: Υ, ψ, ϕ for hidden beauty, charm, strangeness; b, c, s for open flavor quantum numbers