

# Observation of doubly-charged tetraquark candidate at LHCb (CLHCP 2022)

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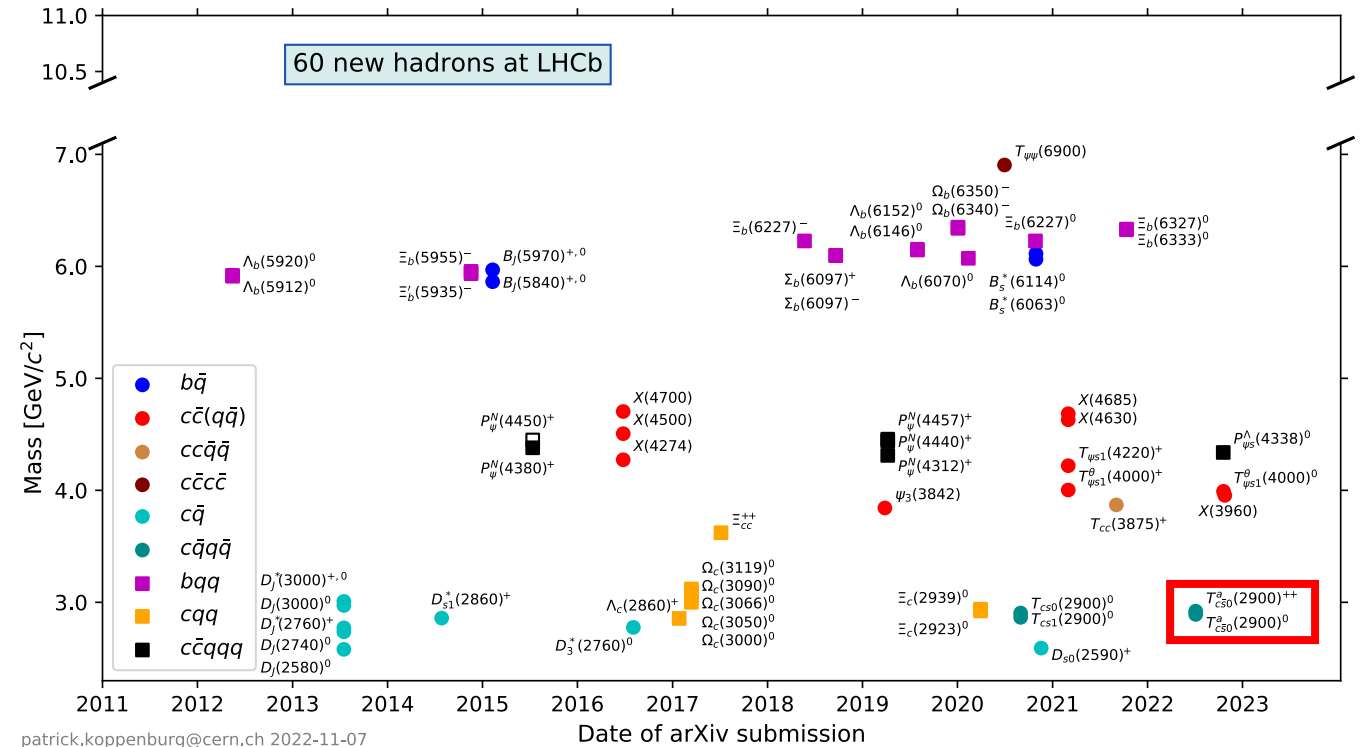
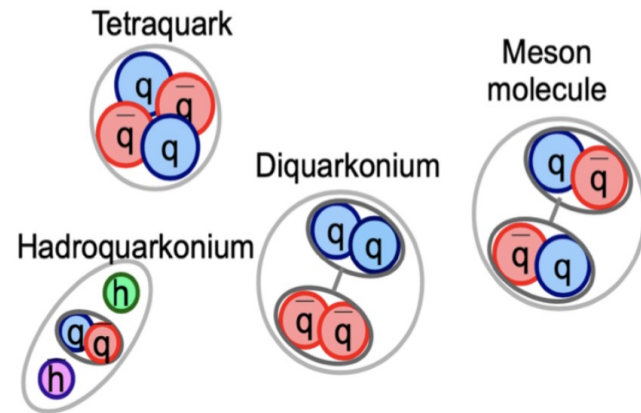
**On behalf of the LHCb Collaboration**

*Nov. 25<sup>th</sup>, 2022*



**中国科学院大学**  
University of Chinese Academy of Sciences

- The existence of mesonic exotic state has been discussed since 1964<sup>1</sup>.
- In 2003, Belle Collaboration reported the first observation of  $\chi_{c1}(3872)$ .
- Many mesonic exotic states are observed in the past two decades.
- A series of **theoretical models** are established to describe these states.



**Masses and discovery date for states observed at LHCb.**  
Hollow markers indicate superseded states.<sup>2,3</sup>

## ● 60 new hadrons observed at LHCb!

- 15 tetraquark candidates.
- 5 pentaquark candidates.<sup>[4]</sup>

<sup>1</sup> M. Gell-Mann, A schematic model of baryons and mesons, Phys. Lett. 8 (1964) 214.

<sup>2</sup> <https://www.nikhef.nl/~pkoppenb/particles.html>

<sup>3</sup> Exotic hadron naming convention: <https://arxiv.org/abs/2206.15233>

<sup>4</sup>  $P_c(4450)^+$  resolved into  $P_c(4440)^+$  and  $P_c(4457)^+$ .

## ● Why $D_s^+ \pi^\pm$ states?

- The  $D_{s0}^*(2317)^+$  ( $D_s^+ \pi^0$ ) state was observed in 2003.
- It is thought to have some **tetraquark component** in several theoretical descriptions, whose  $I = 1$  partners can exist in the  $D_s^+ \pi^\pm$  final states.
- **Prof. Hai-Yang:** It would be astonishing if a doubly charged resonance is found.

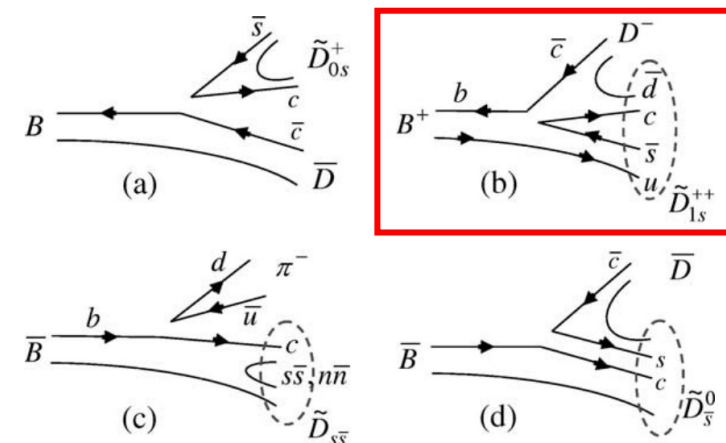
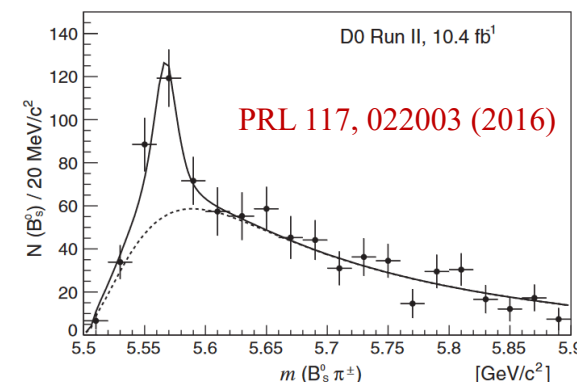


Fig. 2. Diagrams for (a)  $B \rightarrow \bar{D} \tilde{D}_{0s}^+$ , (b)  $B^+ \rightarrow D^- \tilde{D}_{1s}^{++}$  ( $B \rightarrow \bar{D} \tilde{D}_{1s}$ ), (c)  $\bar{B} \rightarrow \pi^- \tilde{D}_{s\bar{s}}, \pi^- \tilde{D}$ , (d)  $B \rightarrow D D_s^0$ .

PLB 566 (2003) 193–200

- Evidence of  $X(5568)^+$  ( $B_s^0 \pi^\pm$ ) claimed by D0 Collaboration in 2016.
- **Not been confirmed** by the other experiments.
- Natural to search for  $D_s^+ \pi^\pm$  resonances as predicted in the diquark-antidiquark model.



$X(5568)^+$

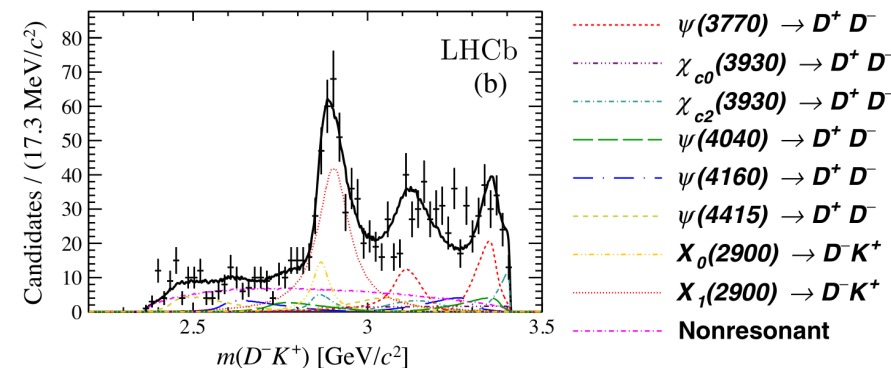
## ● Why $D_s^+ \pi^\pm$ states?

- In 2020, the  $X_0(2900)^0$  and  $X_1(2900)^0$  ( $c\bar{s}u\bar{d}$ ), have been observed at LHCb in  $D^+ K^-$  final states.
- They are candidates to be the first observed tetraquarks consisting of four different flavors.

PRL 125, 242001 (2020)

PRD 102, 112003 (2020)

➤ Would there be exotic states decaying into  $D_s^+ \pi^-$  ( $c\bar{s}u\bar{d}$ ) or  $D_s^+ \pi^+$  ( $c\bar{s}u\bar{d}$ ) final states?



## ● Amplitude analyses of $B^0 \rightarrow \bar{D}^0 D_s^+ \pi^-$ and $B^+ \rightarrow D^- D_s^+ \pi^+$ decays.

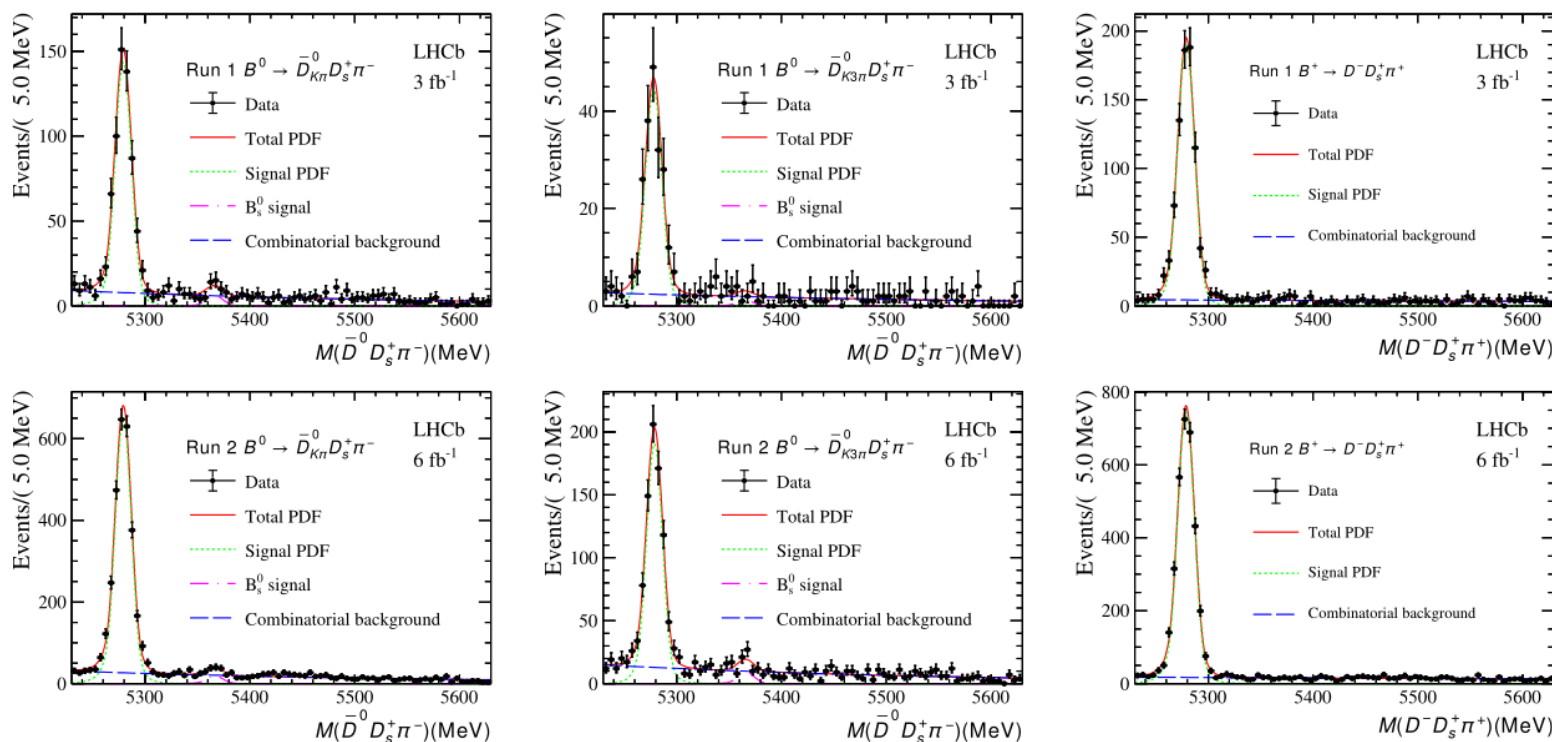
- Only **excited**  $\bar{D}^*$  states with natural spin-parity expected to contribute.
- Ideal channels to search for **possible exotic states** decaying to  $D_s^+ \pi^\pm$  final states.

- **Can be explored by LHCb datasets with large statistic!**

◆  $B^0 \rightarrow \bar{D}^0 D_s^+ \pi^-$  &  $B^+ \rightarrow D^- D_s^+ \pi^+$

LHCb-PAPER-2022-026  
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◆ Using all the Run 1 and Run 2 datasets.



◆  $\bar{D}^0 \rightarrow K^+ \pi^-, K^+ \pi^- \pi^- \pi^+$

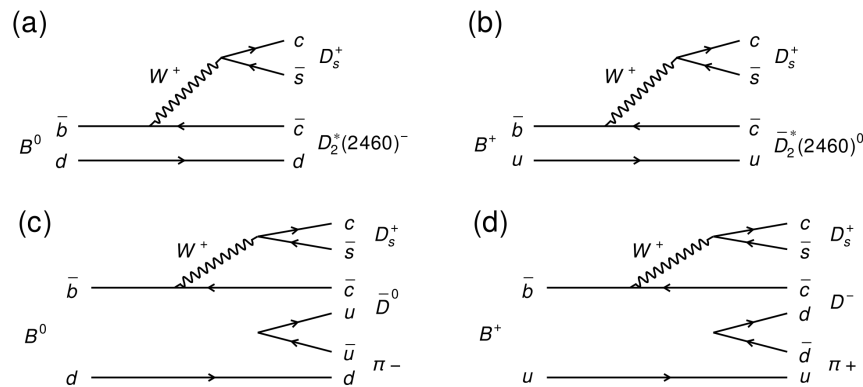
◆  $D^- \rightarrow K^+ \pi^- \pi^-$

◆  $D_s^+ \rightarrow K^+ K^- \pi^+$

✓ In  $B$  signal region:

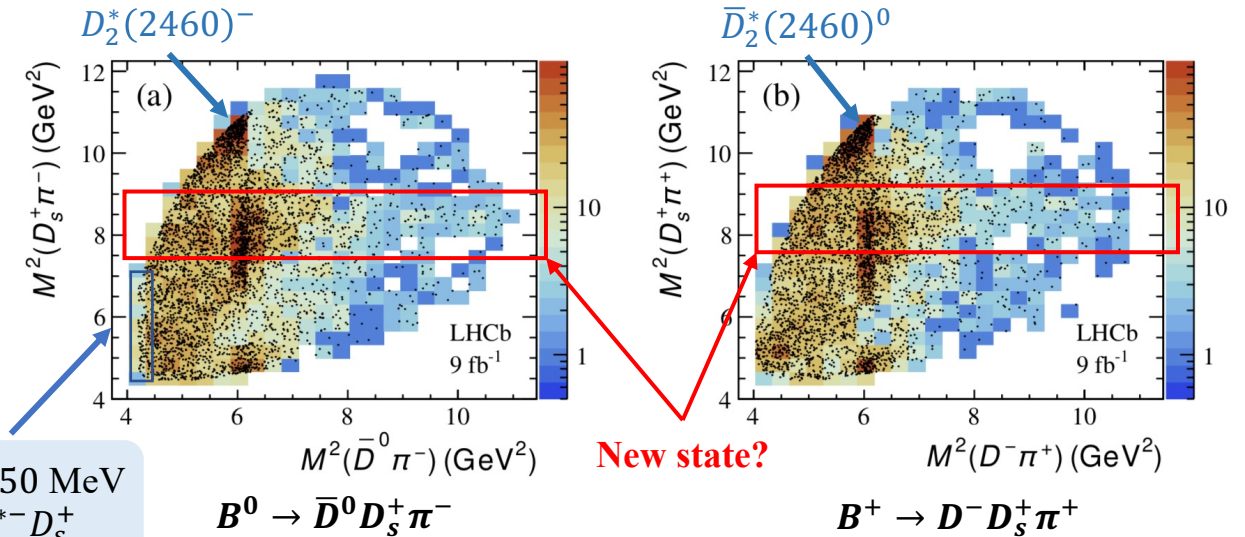
✓ **4009**  $B^0 \rightarrow \bar{D}^0 D_s^+ \pi^-$  signals, purity **90.7%**.

✓ **3750**  $B^+ \rightarrow D^- D_s^+ \pi^+$  signals, purity **95.2%**.



Feynman diagrams

$m(\bar{D}^0 \pi^-) > 2050 \text{ MeV}$   
Veto  $B^0 \rightarrow D^{*-} D_s^+$



- ✓ Clear spin-2 vertical band at the  $\bar{D}\pi$  mass squared around  $6 \text{ GeV}^2$ .
- ✓ A faint horizontal band at the  $D_s\pi$  mass squared around  $8.5 \text{ GeV}^2$ .
- Very similar features in the Dalitz plots of two channels, which are isospin-related.
- It suggests the feasibility to perform **simultaneous fit** of the two channels.
  - The amplitude parameters of all components are set to be the same of the two channels.

- Unbinned maximum likelihood fits performed with TF-PWA

$$P(x; \Theta) = f_{\text{sig}} \cdot P_{\text{sig}}^{\text{norm}}(x; \Theta) + f_{\text{bkg}} \cdot P_{\text{bkg}}^{\text{norm}}(x),$$

Fractions determined from mass fits

- Background modelled from upper sideband with extrapolating into signal regions

- Signal PDF: 
$$P_{\text{sig}}^{\text{norm}}(x; \Theta) = \frac{\epsilon(x) |\mathcal{A}(x; \Theta)|^2}{I_{\text{sig}}(\Theta)}.$$

Normalization factor

- Efficiencies obtained from full simulation with corrections for data-simulation difference

- Amplitude model: 
$$\mathcal{A}(x; \Theta) = \sum c_i \cdot \mathcal{A}_i(x; \Theta_i),$$

Angular distribution + line shape (RBW etc.)

● The  $\bar{D}\pi$  candidates with natural spin-parity on PDG:

Resonance	$J^P$	Mass ( GeV)	Width ( GeV)	Comments
$\bar{D}^*(2007)^0$	$1^-$	$2.00685 \pm 0.00005$	$< 2.1 \times 10^{-3}$	Width set to be 0.1 MeV
$D^*(2010)^-$	$1^-$	$2.01026 \pm 0.00005$	$(8.34 \pm 0.18) \times 10^{-5}$	
$\bar{D}_0^*(2300)$	$0^+$	$2.343 \pm 0.010$	$0.229 \pm 0.016$	#
$\bar{D}_2^*(2460)$	$2^+$	$2.4611 \pm 0.0007$	$0.0473 \pm 0.0008$	#
$\bar{D}_1^*(2600)^0$	$1^-$	$2.627 \pm 0.010$	$0.141 \pm 0.023$	#
$\bar{D}_3^*(2750)$	$3^-$	$2.7631 \pm 0.0032$	$0.066 \pm 0.005$	#
$\bar{D}_1^*(2760)^0$	$1^-$	$2.781 \pm 0.022$	$0.177 \pm 0.040$	#
$\bar{D}_J^*(3000)^0$	$?^?$	$3.214 \pm 0.060$	$0.186 \pm 0.080$	# $J^P = 4^+$ is assumed

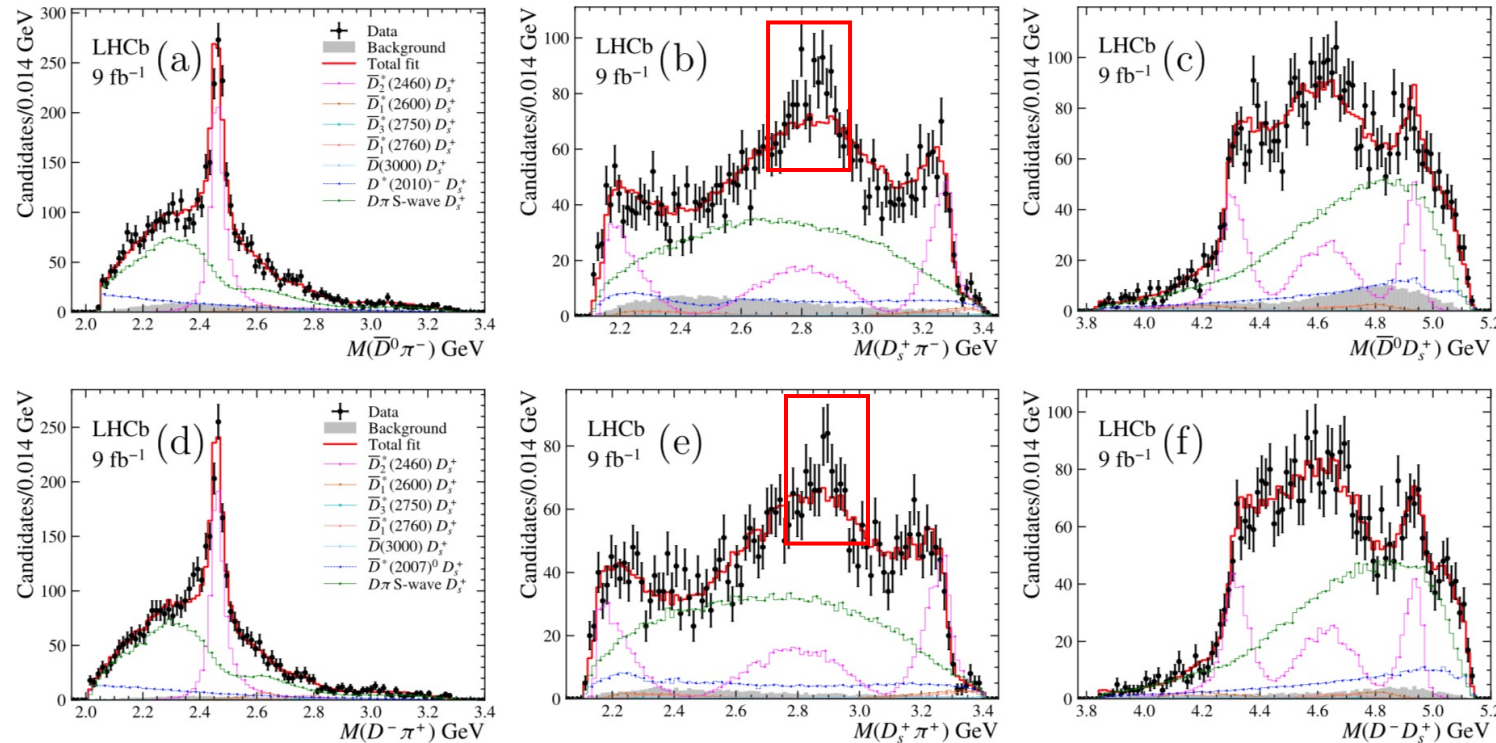
◆ Fit strategy – **simultaneous fit**

- Include all the  $D^*$  and  $D^{**}$  states with natural spin-parity.
- $\bar{D}\pi$  S-wave component: quasi-model-independent description (QMI) spline points.\*
- All parameters, except  $\bar{D}^*(2007)^0$  and  $D^*(2010)^-$  are shared.

\*11 spline points at [1.9, 2.0, 2.1, 2.2, 2.3, 2.4, 2.5, 2.6, 2.7, 2.9, 3.4] GeV



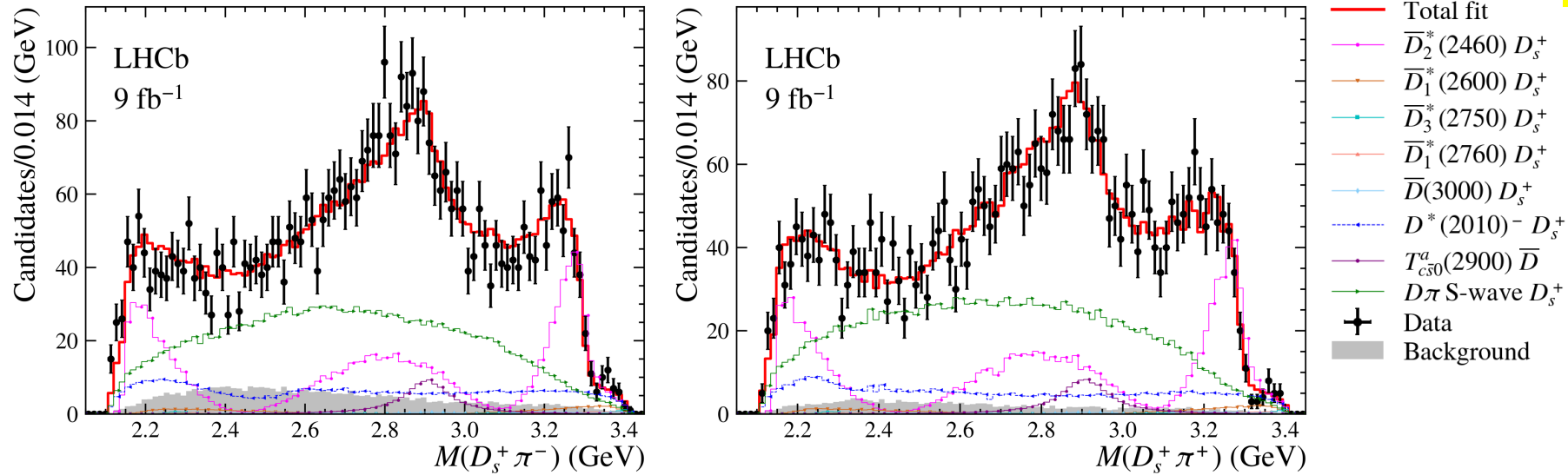
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LHCb-PAPER-2022-027



## ✓ Fit result

- ✓  $M(\bar{D}\pi)$  and  $M(\bar{D}D_s^+)$  well described.
- ✓  $D_1^*(2600)$ ,  $D_1^*(2760)$ ,  $D(3000)$  **not significant**, however, **still included** conservatively.
- ✓ Spin-parity of  **$D(3000)$**  favors  $4^+$ .
- ✓ Further new  $\bar{D}^{**}$  state with spin-parity up to  $4^+$  tested to be **disfavored**.

**Peaking structures not well described near  $M(D_s^+\pi) = 2.9 \text{ GeV}$  !**



◆ **Two  $D_s^+ \pi$  exotic states with shared parameters are added.**

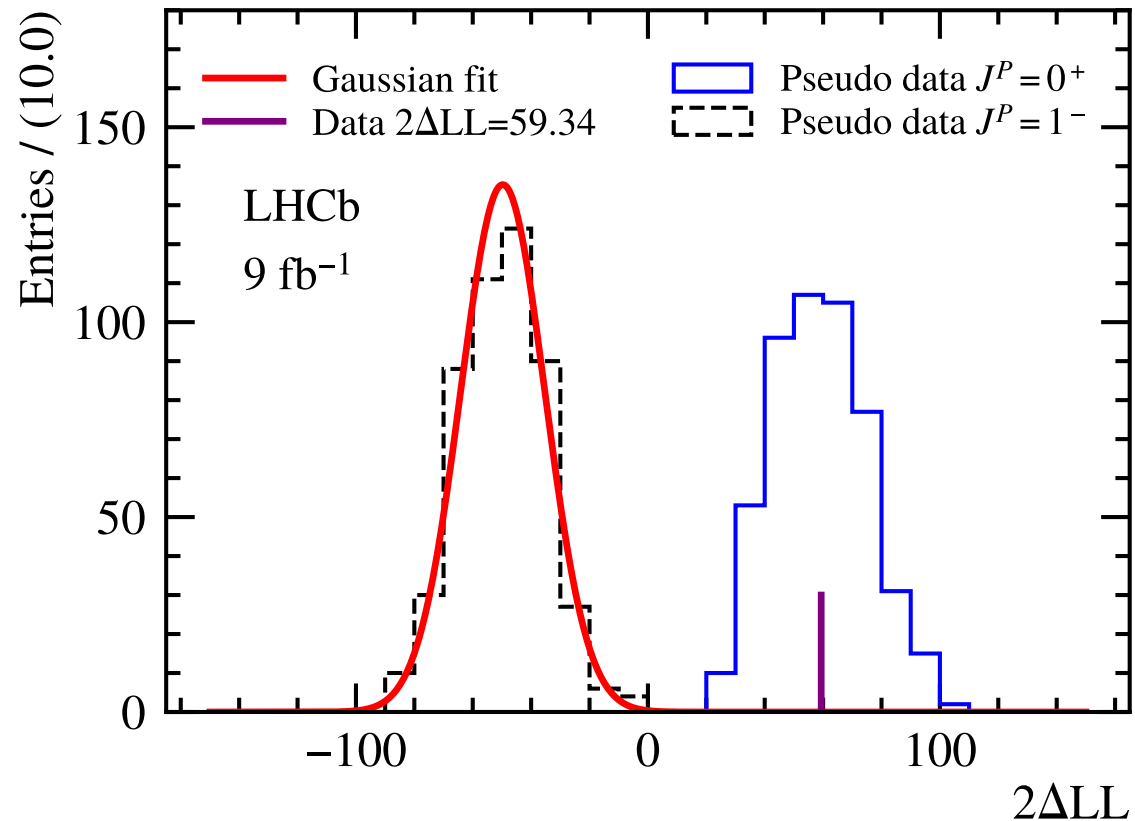
- ✓  $J^P$  up to  $3^+$  are tested,  **$0^+$**  produces the best likelihood.
- ✓ Significance greater than  **$9\sigma$** .
- ✓ Mass and width are measured:

$$M = 2.908 \pm 0.011 \pm 0.020 \text{ GeV}$$

$$\Gamma = 0.136 \pm 0.023 \pm 0.013 \text{ GeV}$$

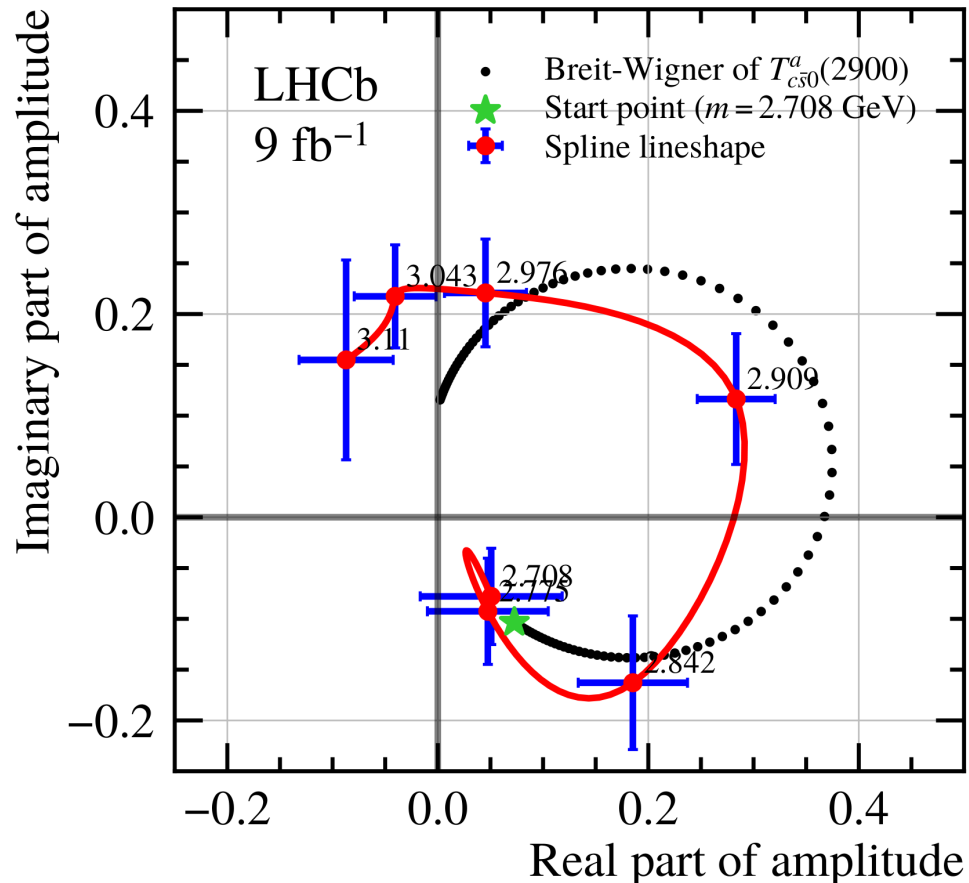
✓ **Named\*** as  $T_{c\bar{s}0}^a(2900)^0 (D_s^+ \pi^-)$  and  $T_{c\bar{s}0}^a(2900)^{++} (D_s^+ \pi^+)$

\* Exotic hadron naming convention: <https://arxiv.org/abs/2206.15233>



## ◆ Spin test

- ✓ Generate pseudoexperiments.
- ✓ Spin-parity favored  **$0^+$**  over  **$1^-$**  with a significance about  **$7.5\sigma$** .
- ✓  $0^+$  is also significantly preferred when exotics not constrained by isospin.

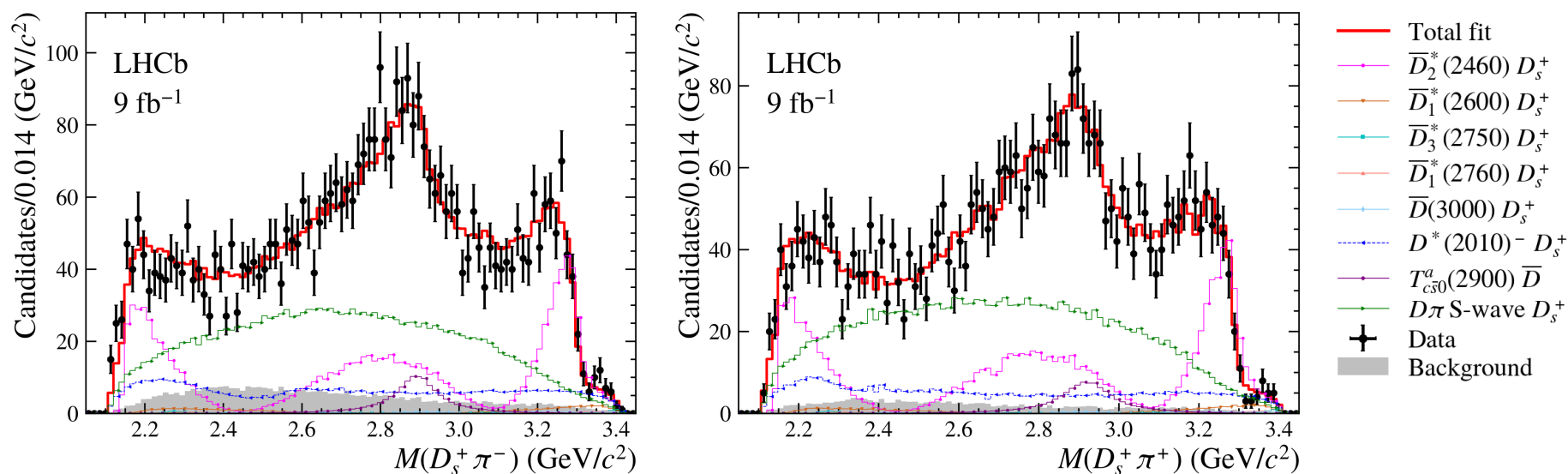


## ◆ Argand diagram

- ✓ Replace the RBW of  $T_{c\bar{s}0}^a$  with spline points.
- ✓ The 7 points from  $m - \frac{3}{2}\Gamma$  to  $m + \frac{3}{2}\Gamma$ , where  $m$  and  $\Gamma$  is the fitted values of  $T_{c\bar{s}0}^a$  mass and width.
- ✓ Lineshape **consistent** with RBW.

## ◆ Separate $T_{c\bar{s}0}^a$ parameters

- ✓  $-\ln\mathcal{L}$  improved by **2.8**, with **4** free parameters added.
- ✓ Masses and widths are **consistent** with each other.
- ✓ **Isospin triplet!**



	Mass (GeV)	Width (GeV)	Significance
$T_{c\bar{s}0}^a(2900)^0$	$2.892 \pm 0.014 \pm 0.015$	$0.119 \pm 0.026 \pm 0.013$	<b>8.0<math>\sigma</math></b>
$T_{c\bar{s}0}^a(2900)^{++}$	$2.921 \pm 0.017 \pm 0.020$	$0.137 \pm 0.032 \pm 0.017$	<b>6.5<math>\sigma</math></b>

- First observation of a **doubly charged mesonic** exotic state, together with its **neutral partner**.
  - ✓ Belong to the same isospin triplet.
  - ✓ Spin-parity:  $0^+$ .
  - ✓ Minimum quark content:  $T_{c\bar{s}0}^a(2900)^{++}: [c\bar{s}u\bar{d}]; T_{c\bar{s}0}^a(2900)^0: [c\bar{s}u\bar{d}]$
  - ✓ Similar mass with  $X_0(2900)$  ( $cs\bar{u}\bar{d}$ ), but width and flavor contents are different.

## ● Next step

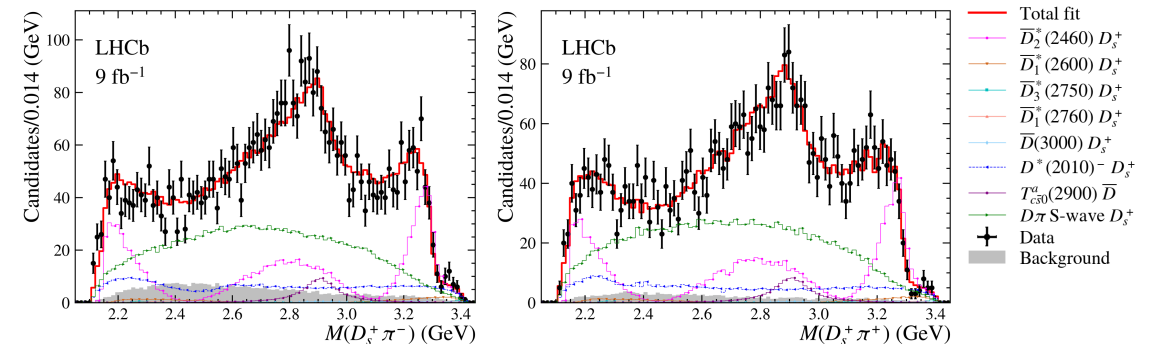
- Several ongoing  $B \rightarrow DDh$  analyses with LHCb Run 1 and Run 2 datasets.

•  $B: B^{0,+}, B_s^0, \Lambda_b^0 \dots; D: D^{0,+}, D^{*+}, D_s^+, \Lambda_c^+ \dots; h: K^+, \pi^+, p, \Lambda \dots$

- LHC Run 3 data taking started recently.

✓ More statistic, detailed analysis on  $B \rightarrow \bar{D}D_s\pi$

✓  $B \rightarrow \bar{D}D_s\pi^0, B \rightarrow \bar{D}D_s\pi\pi, B_s \rightarrow \bar{D}D_sK\pi \dots$



# Thank you!

# Back Up

## Forward single arm detector

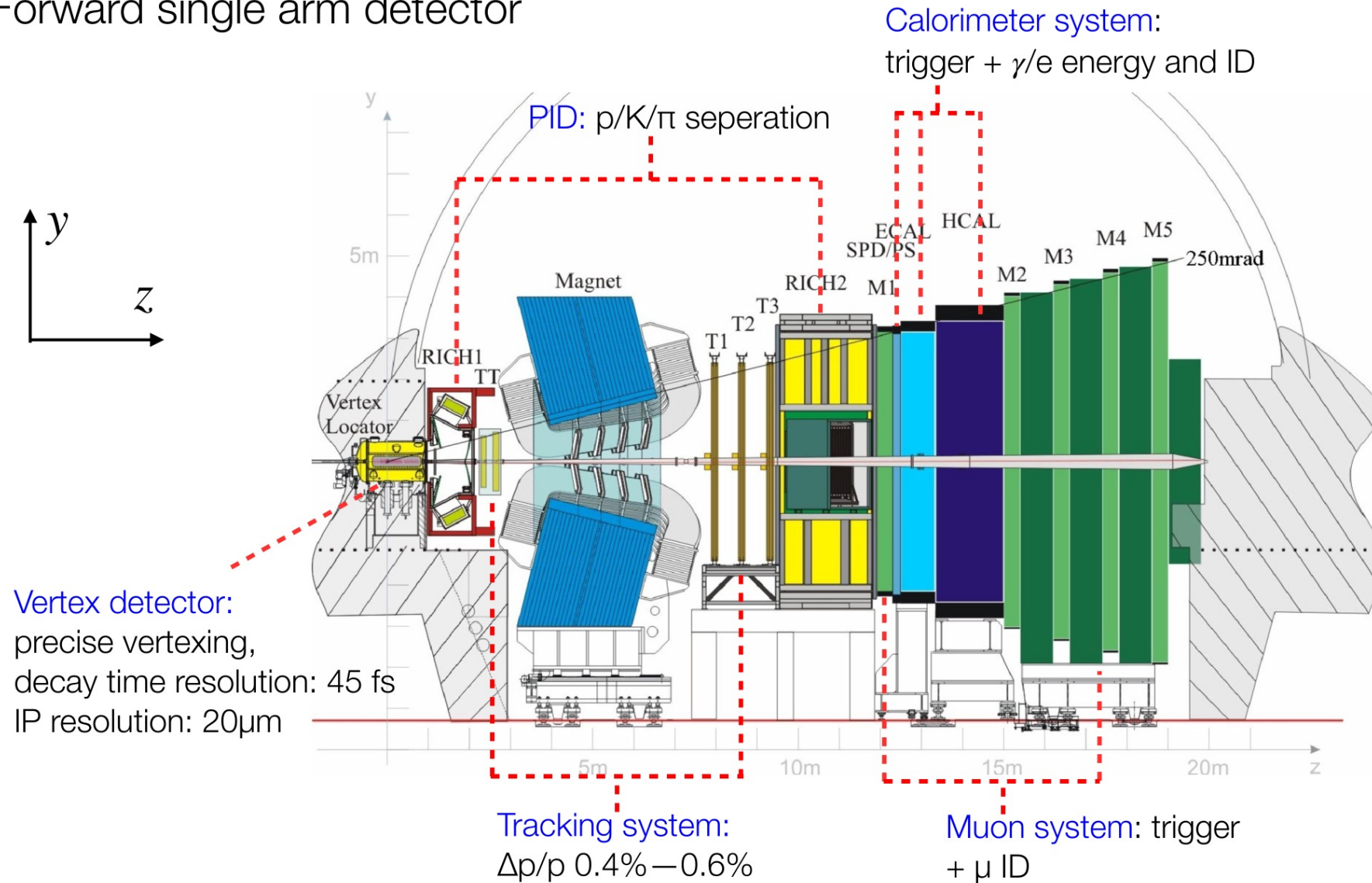




Table 2: Signal and background yields inside the  $B$  mass signal window, together with the signal purity, split by run period and decay mode. The uncertainties shown are statistical.

Decay	Parameter	Run 1	Run 2
$B^0 \rightarrow \bar{D}_{K\pi}^0 D_s^+ \pi^-$	Signal yield	$564 \pm 26$	$2534 \pm 55$
	Total candidates	633	2753
	Purity	89.1%	92.1%
$B^0 \rightarrow \bar{D}_{K3\pi}^0 D_s^+ \pi^-$	Signal yields	$177 \pm 14$	$734 \pm 31$
	Total candidates	199	835
	Purity	88.9%	87.9%
$B^+ \rightarrow D^- D_s^+ \pi^+$	Signal yield	$766 \pm 29$	$2984 \pm 57$
	Total candidates	797	3143
	Purity	96.1%	94.9%

## ● Amplitude formula

- For decay  $B \rightarrow D(1)D_s^+(2)\pi(3)$ , the amplitude is

$$\mathcal{A} = \sum_a C_a \cdot f_a(m_{12}^2) \cdot T_a(\theta_{12}) + \sum_b C_b \cdot f_b(m_{23}^2) \cdot T_b(\theta_{23}) + \sum_c C_c \cdot f_c(m_{13}^2) \cdot T_c(\theta_{13}),$$

- $C$  is complex coefficient
- $f(m^2)$  is lineshape function
- $T(\theta)$  is the angular term

## ● Maximum likelihood fit (cFit\*)

- $\text{PDF}(x; \Theta) = |A(x; \Theta)|^2$ , where  $x$  is the data point,  $\Theta$  is the parameter set.
- $-2\ln L = -2 \sum_i^N \ln \text{PDF}(x; \Theta)$

- For resonance, we use the relativistic Breit-Wigner lineshape

$$f_{BW} = q(m)^{L_1} p(m)^{L_2} \frac{F_1(m, L_1) F_2(m, L_2)}{m_0^2 - m^2 - im_0 \Gamma(m)}$$

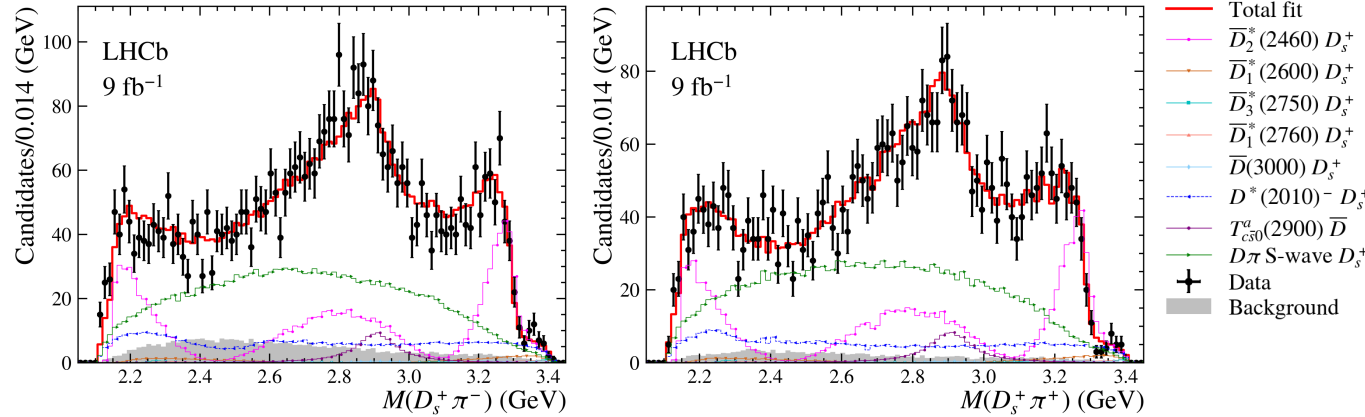
with running width

$$\Gamma(m) = \Gamma_0 \left( \frac{q(m)}{q_0} \right)^{2L_2+1} \frac{m_0}{m} F_2^2(m, L_2)$$

- The Blatt-Weisskopf form factors

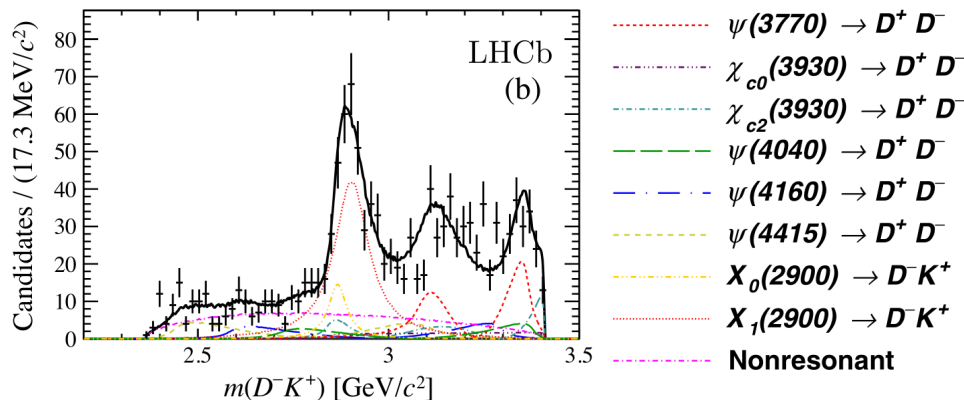
$$F(m, L) = \begin{cases} 1 & L = 0 \\ \sqrt{\frac{1+z^2(m)}{1+z_0^2}} & L = 1 \\ \sqrt{\frac{9+3z^2(m)+z^4(m)}{9+3z_0^2+z_0^4}} & L = 2 \\ \sqrt{\frac{225+45z^2(m)+6z^4(m)+z^6}{225+45z_0^2+6z_0^4+z_0^6}} & L = 3 \end{cases}$$

\*Detailed implementations in back-up slide



$T_{c\bar{s}0}^a(2900)^0 [c\bar{s}u\bar{d}]$

$T_{c\bar{s}0}^a(2900)^{++} [c\bar{s}u\bar{d}]$



$X_0(2900), X_1(2900) [c\bar{s}u\bar{d}]$

PRD 102, 112003 (2020)

## ● $T_{c\bar{s}0}^a(2900)$ & $X_0(2900)$

- ✓ Similar mass, but width and flavor contents are different.
- ✓  $T_{c\bar{s}1}^a(2900)?$
- ✓  $T_{c\bar{s}0}^a(2900)^{++} \rightarrow D^+ K^+?$
- ✓  $T_{c\bar{s}0}^a(2900)^+ \rightarrow D_s^+ \pi^0, D_s^+ \pi^+ \pi^-?$

**More statistic needed!**

	Mass (GeV)	Width (GeV)
$T_{c\bar{s}0}^a(2900)$	$2.908 \pm 0.011 \pm 0.020$	$0.136 \pm 0.023 \pm 0.020$
$X_0(2900)$	$2.866 \pm 0.007 \pm 0.002$	$0.057 \pm 0.012 \pm 0.004$
$X_1(2900)$	$2.904 \pm 0.005 \pm 0.001$	$0.110 \pm 0.011 \pm 0.004$