# Charmless two-body B decays in perturbative QCD factorization approach 

Shan Cheng<br>Hunan University

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Jian Chai, SC, Yao-hui Ju, Da-cheng Yan, Cai-dian Lü and Zhen-jun Xiao


## Overview

(1) Motivation

- Significances
- Experiment promotions
- Theory progresses
(2) Introduction of PQCD approach
- Three scale factorization frame
- Progresses towards to NLO
(3) PQCD updates of $B \rightarrow P P, P V, V V$ decays
(4) Conclusion


## Motivation: Signficiance

- Asymmetry between matter and antimatter indicates the interaction with CPV, Heavy flavour physics (HFP) provides many processes with CPV, but is not enough and new mechanism of CPV is imminent.
- In the post Higgs Era, the precise testing of SM and searching of NP are the core tasks of particle physics.
- HFP plays am important role in both two targets.
- B meson decays is the vanguard of HFP study.
- Timeline of B physics
$\dagger$ 1973, Kaboyashi \& Maskawa proposed a $3 \times 3$ unitary matrix (4 parameters) of quark mixing to accommodate CPV,
$\dagger$ 1977, CFS-E288 at FermiLab discovered $\Upsilon$ meson ( $b \bar{b}$ ), Lederman,
$\dagger$ 1981, Bigi \& Santa pointed out the expectation of large CPV in $B^{0}$ decay according to CKM theory,
$\dagger$ 1987, Oddone proposed the construction of $B$ factories to study CPV,
$\dagger$ 1999, BABAR and Belle started running; 2001(04), $A_{C P}(t, f)\left(A_{C P}\right)$ in $B^{0}$ decays,
$\dagger$ 2009, LHCb played in to the game; 2013(20), $A_{C P}\left(A_{C P}(t, f)\right)$ in $B_{s}$ decays, 2012, $A_{C P}$ in $B^{+}$decays; 2019, $\delta A_{C P}$ in $D$ decays,
$\dagger$ Anomalies: $R_{K^{(*)}}, R_{D}, P_{5}^{\prime}, B_{s} \rightarrow \mu \mu,\left|V_{u b}\right|,\left|V_{c b}\right|$


## Motivation: Experiment promotions

- Success running of $B$ factories(1999-2008), CPV in $B^{0}$ decays, anomalies
- LHC(2009-2023), 3-5 order more events, CPV in $B^{+}, B_{s}, D$ decays, anoma.
- SuperKEKB(2018-2026), $8 \times 10^{35} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}, 80$ times larger than that of KEKB, $\dagger$ reduce the uncertainty of $\left|V_{u b}\right|$ from $2.5 \%$ to $1.2 \%$ in $B \rightarrow \pi / \nu, \rho / \nu$ $\dagger$ improve the accuracy of $A_{C P}, \delta A_{C P}$ in $B \rightarrow K^{*} \pi, K \rho, K^{*} \rho$ to that in $K \pi$ channel $\dagger B \rightarrow V V: f_{\perp}$, relative phase $\phi_{\|}, \phi_{\perp}, \delta_{0}$ and helicity CP parameters $\mathcal{A}_{\mathrm{CP}}^{0}, \mathcal{A}_{\mathrm{C} P}^{\perp}$ $\dagger$ CKM triangle angle $\alpha$ from $B \rightarrow \pi^{0} \pi^{0}$, and et al. [CERN Yellow Rep. Monogr, 2019] $\ddagger$ The first measurements of $B^{+} \rightarrow \rho^{+} \rho^{0}, B^{0} \rightarrow K^{0} \pi^{0}$ have been released [Belle-II, 2021].
- HL-LHC(2027-2033), $\mathcal{L}=23(300) \mathrm{fb}^{-1}$ in phase $1(2), 2$ order larger than it at LHC, $\dagger$ reduce the uncertainty of $\left|V_{u b}\right|$ from $2.4 \%$ to $0.7 \%(0.4 \%)$ in $B \rightarrow \pi \pi, \pi \rho, \rho \rho$ $\dagger$ improve the measurement of $C_{\pi^{+} \pi^{-}}, S_{\pi^{+} \pi^{-}}$by three times (one order)
$\dagger$ CKM triangle angle $\alpha$ from $B \rightarrow \rho \rho, \rho \pi$, and et al. [E. Kou et al. [Belle-II], 2019]
- More precise study of $B$ decays from the theoretical side is imperative.


## Motivation: Theoretical progresses

- High precision calculation of two-body $B$ decays
$\dagger B \rightarrow \pi \pi$ decay is studied from LCSRs [Khodjamirian 2001,03,05] the high order \& power corrections of $B \rightarrow P, V$ form factors LQCD [HPQCD 2013] [Bharucha 2016, Wang 15,16,20, Lü 19, Beneke 17, Gubernari 19, Cheng 17,19]
$\dagger$ NF: $\sim F_{B \rightarrow M_{2}} \otimes f_{M_{1}}$ [Khodjamirian 2001,03,05]
GF: pQCD corrections from $O_{i=1,2}$ and $O_{i=3,10} \quad$ [Ali 1998,99]
QCDF: VC to $\mathcal{M}_{t, p}+$ correction to spectator scattering, full NNLO
[Benele 2010, Bell 15, 20, Huber 16, Beneke 06,07, Jain 07] [Lu 2022]
$\dagger$ SCET: introduces different fields in different energy regions, simple kinematics but complicated dynamics [Bauer 2001, Chay 04, Becher 15], QCDF/SCET [Beneke 2015]
- To eliminate the end-point singularity emerged in collinear factorization, the PQCD approach is established by picking up the $k_{T}$ of valence quarks.
$\dagger B \rightarrow M$ FFs and the annihilation amp. are both calculable [Keum 2001, Lü 01]
$\dagger$ LO $\left(\mathcal{O}\left(\alpha_{s}\right)\right) B \rightarrow P P, P V, V V$ decays [Xiao 2007; Lü̈ 02; Li 05, Li 06, Zou 15], [Hua 2021]
$\dagger$ partially NLO $\left(\mathcal{O}\left(\alpha_{s}^{2}\right)\right)$ : factorizable amplitudes [Cheng \& Xiao 2021]
to effective operators [Mishima 2003, Li 05]
to hard scattering [Li 2012, Cheng 14], [Li 13, Cheng 15,15, Hua 18], [Li 14, Liu 15,16]
- A systematical study of $B$ decays with high accuracy is in time.


## PQCD: Three scale factorization frame

New physics: $\mathcal{L}_{N P}$

Electroweak scale $\left(m_{W}\right): \mathcal{L}_{E W}+\mathcal{L}_{D>4}$
Heavy quark scale $\left(m_{b}\right): \mathcal{L}_{\text {eff }}=-\frac{\downarrow}{\sqrt{\sqrt{2}}} \begin{aligned} & \downarrow \\ & \downarrow\end{aligned} V_{\mathrm{CKM}} \sum_{i} C_{i}(\mu) O_{i}(\mu)+\mathcal{L}_{\text {eff }, D>6}$
Hadron scale $\left(\Lambda_{\mathrm{QCD}}\right)$ : LCDAs, PDF, PDA

- Derive the effective Hamiltonian by integrating over $m_{W}$ [Buchalla 1996]
$\dagger$ Product of two charged currents is expanded by a series of local operators $O_{i}$ with the weighted coefficients $C_{i}$
- Dynamics at the scale $\mathcal{O}\left(m_{W}\right)$ is absorbed into Wilson coefficients $C_{i}(\mu)$
$\dagger C_{i}$ is obtained by matching the $\mathcal{L}_{\text {eff }}$ with the full theory of weak decays
[Ma 80, Inami\&Lim 81, Clements 83]
- The rest go into the four fermion effective operators $O_{i}(\mu)$
- The key is to calculate the hadron matrix element $\left\langle M_{1} M_{2}\right| O_{i}|B\rangle$


## PQCD: Three scale factorization frame



(a)


Effective tree, bullent denotes $O_{i}$

- Integrating over the $m_{W}$ and $m_{b}$
- Weak phase difference between charged and FCNC of $b$ decays



## PQCD: Three scale factorization frame

- Diagrams at scales $\mathcal{O}\left(\Lambda_{\mathrm{QCD}}\right)-\mathcal{O}\left(m_{b}\right)$ : Hadron matrix element $\left\langle M_{1} M_{2}\right| O_{i}|B\rangle$
- Factorization: detach the hard kernel $\mathcal{H} O_{i}$ at scale $\mathcal{O}\left(m_{b}\right)$ from the hadron wave function $\Phi B, M_{1}, M_{2}$ mesons at scale $\mathcal{O}\left(\Lambda_{\mathrm{QCD}}\right)$
- Prediction power: $\mathcal{H}$ is calculated perturbatively order by order, $\Phi$ s are universal
different strong phases

- End-point singularities appear in diagrams (a,b,e,f)


## PQCD: Three scale factorization frame

- Introduce $\mathbf{k}_{T}$ to regularize the end-point singularity [Huang 1991]
- Enriches the study of hadron DAs, TMD definition with Wilson link, observables
- Scales of transversal momentum and the large logarithms [borrowed from H.N Li]

$\dagger$ Multiple scales and hence large single logarithms in $\mathcal{H}$ and $\Phi$ from QCD corrections
$\dagger$ Double logs in the soft-collinear regions $\alpha_{s}(\mu) \ln ^{2}\left(\mathrm{k}_{T}^{2} / m_{B}^{2}\right)$


## PQCD: Three scale factorization frame

- In order to repair the perturbative expansion, do resummation by using RGE
- $k_{T}$ resummation for $\mathcal{H}$ and obtain $S\left(x_{i}, b_{i}, Q\right)$ [Botts 1989, Li 92]
$\dagger$ decreases the inverse power of the momentum transfer in the divergence amplitude
$\dagger$ exhibits high suppression for large transversal distances (small $k_{T}$ ) interactions
- Integrating over $k_{T}$, large $\log \ln ^{2}\left(x_{i}\right)$ when intermediate gluon is on shell
- threshold resummation for $\Phi$ and obtain $S_{t}\left(x_{i}, Q\right)$ [Li 1999]
$\dagger$ suppresses the small $x_{i}$ regions
$\dagger$ repairs the self-consistency between $\alpha_{s}(t)$ and hard $\log \ln \left(x_{1} x_{3} Q^{2} / t^{2}\right)$
$\ddagger$ dynamics with $k_{T}<\sqrt{Q \Lambda}$ is organized into $S(x, b, Q)$
$\ddagger$ dynamics in small $x$ is suppressed by $S_{t}(x, Q)$


$$
\mathcal{M}\left(B \rightarrow M_{1} M_{2}\right)=\sum_{i} C_{i}\left(m_{W}, t\right) \otimes \mathcal{H}_{i}(t, b) \otimes \phi(x, b) \operatorname{Exp}\left[-s\left(p^{+}, b\right)-\int_{1 / b}^{t} \frac{d \bar{\mu}}{\bar{\mu}} \gamma_{\phi}\left(\alpha_{s}(\bar{\mu})\right)\right]
$$

## PQCD: Progresses towards to NLO

$\mathcal{M}\left(B \rightarrow M_{1} M_{2}\right)=\sum_{i} C_{i}\left(m_{W}, t\right) \otimes \mathcal{H}_{i}(t, b) \otimes \phi(x, b) \operatorname{Exp}\left[-s\left(p^{+}, b\right)-\int_{1 / b}^{t} \frac{d \bar{\mu}}{\bar{\mu}} \gamma_{\phi}\left(\alpha_{s}(\bar{\mu})\right)\right]$

- The NLO QCD/QED corrections to $C_{i}$ has been finished [Buchalla, 1996, Rev. Mod. Phys]

$$
\begin{aligned}
C_{1}\left(m_{W}\right) & =\frac{11}{2} \frac{\alpha_{s}\left(m_{W}\right)}{4 \pi} \\
C_{2}\left(m_{W}\right) & =1-\frac{11}{6} \frac{\alpha_{s}\left(m_{W}\right)}{4 \pi}-\frac{35}{18} \frac{\alpha_{\mathrm{em}}}{4 \pi} \\
C_{3}\left(m_{W}\right) & =-\frac{\alpha_{s}\left(m_{W}\right)}{24 \pi}\left[E_{0}\left(\frac{m_{t}^{2}}{m_{W}^{2}}\right)-\frac{2}{3}\right] \\
& +\frac{\alpha_{\mathrm{em}}}{6 \pi} \frac{1}{\sin ^{2} \theta_{W}}\left[2 B_{0}\left(\frac{m_{t}^{2}}{m_{W}^{2}}\right)+C_{0}\left(\frac{m_{t}^{2}}{m_{W}^{2}}\right)\right] \\
C_{4}\left(m_{W}\right) & =-\frac{\alpha_{s}\left(m_{W}\right)}{8 \pi}\left[E_{0}\left(\frac{m_{t}^{2}}{m_{W}^{2}}\right)-\frac{2}{3}\right],
\end{aligned}
$$

| $\mu(\mathrm{GeV})$ | 1.0 | 2.0 | 3.0 | 4.0 | 4.98 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $a_{s}(t)$ | 0.63, 0.47 | 0.39, 0.30 | 0.32, 0.25 | 0.29, 0.23 | 0.26, 0.21 |
| $c_{1}(\mu)$ | -0.27, -0.51 | -0.61, -0.31 | -0.85, -0.24 | $-1.05,-0.20$ | $-0.83,-0.17$ |
| $C_{2}(\mu)$ | 1.12, 1.28 | 1.33, 1.15 | 1.50, 1.11 | 1.66, 1.09 | 1.48, 1.07 |
| $C_{3}(\mu)$ | 0.01, 0.04 | 0.03, 0.02 | 0.05, 0.02 | 0.06, 0.01 | 0.05, 0.01 |
| $C_{4}(\mu)$ | -0.03, -0.09 | $-0.06,-0.05$ | -0.08, -0.04 | -0.10, -0.04 | $-0.07,-0.03$ |
| $C_{5}(\mu)$ | 0.01, 0.0.2 | 0.02, 0.01 | 0.02, 0.01 | $0.02,0.01$ | 0.02, 0.01 |
| $C_{6}(\mu)$ | -0.0.3, -0.13 | -0.09, -0.07 | -0.15, -0.05 | -0.20, -0.04 | -0.14, -0.04 |
| $c_{7}(\mu)$ | 0.00, -0.00 | 0.00, -0.00 | 0.00, -0.00 | 0.00, -0.00 | 0.00, -0.00 |
| $C_{8}(\mu)$ | 0.00, 0.00 | 0.00, 0.00 | 0.00, 0.00 | 0.00, 0.00 | 0.00, 0.00 |
| $c_{g}(\mu)$ | $-0.01,-0.01$ | -0.00, -0.01 | -0.01, -0.01 | -0.01, -0.01 | -0.01, -0.01 |
| $\mathrm{Cr}_{10}(\mu)$ | 0.00, 0.01 | 0.01, 0.00 | 0.01, 0.00 | 0.001,0.00 | 0.01, 0.00 |
| LO NLO |  |  |  |  |  |

$\dagger$ Inami-Lim functions $B, C, D, E$ from box, $Z, \gamma, g$ penguin diagrams, respectively
$\dagger$ Scale running from $m_{W}$ to $\mathcal{O}\left(m_{b}\right)$ by evolution matrix: $C_{i}(\mu)=U\left(\mu, m_{W}\right) C_{i}\left(m_{W}\right)$

## PQCD: Progresses towards to NLO

- The NLO corrections to ME $\left\langle M_{1} M_{2}\right| O_{i}|B\rangle$

$\dagger$ Vertex of effective operator in $\mathcal{M}_{a, b}$ Completed in collinear factorization [Beneke 2001, Mishima 03, Li 05]
$\dagger B \rightarrow P$ transition form factors in $\mathcal{M}_{a, b}$ Done up to twist three [Li 2012, Cheng 14]
$\dagger$ EM form factors in $\mathcal{M}_{\boldsymbol{e}, f}$ Done up to twist three for PP, PV [Li 2010,12, Cheng 14,15, Hua 18]
$\dagger$ Scalar form factor in $\mathcal{M}_{e, f}$ with helicity flip Done up to twist three [Cheng 15]
$\dagger$ Glauber gluon correction to $\mathcal{M}_{c, d}$ Done for $M=\pi$ [Li 2014, Liu 2015,16]
$\dagger$ TMD wave functions [Ji 2004, Bacchetta 2008, Collins 2011,14] [Li 2004, 2014, 2015]
$\dagger$ NLO corrections to spectator annihilation amplitudes are still missing


## $B \rightarrow P P, P V, V V$ decays: Amplitudes

- General decomposition of Wilson coefficients for each certain effective weak vertex

| Weak vertex | Typical amplitudes | Wilson coefficients |
| :---: | :---: | :---: |
| $[s, s, s],[d, d, d]$ | $\mathcal{E}^{\mathrm{LL}} / \mathcal{A}^{\mathrm{LL}}, \mathcal{E}_{N F}^{\mathrm{LL}} / \mathcal{H}_{N F}^{\mathrm{LL}}$ | $a_{3}+a_{4}-\frac{a_{9}+a_{10}}{2}, \quad C_{3}+C_{4}-\frac{C_{9}+C_{10}}{2}$ |
| spectator meson $M_{3}$ | $\mathcal{E}^{\mathrm{LR}} / \mathcal{H}^{\mathrm{LR}}, \quad \mathcal{E}_{N F}^{\mathrm{LR}} / \mathcal{A}_{N F}^{\mathrm{LR}}$ | $a_{5}-\frac{a_{7}}{2}, \quad C_{5}-\frac{c_{7}}{2}$ |
|  | $\mathcal{E}^{\mathrm{SP}} / \mathcal{A}^{\mathrm{SP}}, \quad \varepsilon_{N F}^{\mathrm{SP}} / \mathcal{A}_{N F}^{\mathrm{SP}}$ | $a_{6}-\frac{a_{8}}{2}, \quad C_{6}-\frac{c_{8}}{2}$ |
| $[d, s, s], \quad[s, d, d]$ | $\mathcal{E}^{\mathrm{LL}} / \mathcal{H}^{\mathrm{LL}}, \quad \mathcal{E}_{N F}^{\mathrm{LL}} / \mathcal{H}_{N F}^{\mathrm{LL}}$ | $a_{4}-\frac{a_{10}}{2}, \quad C_{3}-\frac{c_{9}}{2}$ |
| , $q_{2}, q_{3}$ | $\mathcal{E}^{\mathbf{L R}} / \mathcal{H}^{\mathbf{L R}}, \quad \mathcal{E}_{N F}^{\mathrm{LR}} / \mathcal{H}_{N F}^{\mathrm{LR}}$ | $a_{6}-\frac{a_{8}}{2}, \quad C_{5}-\frac{c_{7}}{2}$ |
| $\downarrow \quad[s, s, d], \quad[d, d, s]$ | $\mathcal{E}^{\mathrm{LL}} / \mathcal{F}^{\mathrm{LL}}, \quad \mathcal{E}_{N F}^{\mathrm{LL}} / \mathcal{H}_{N F}^{\mathrm{LL}}$ | $a_{3}-\frac{a_{9}}{2}, \quad C_{4}-\frac{c_{10}}{2}$ |
| emission meson $M_{2}$ | $\mathcal{E}^{\mathrm{LR}} / \mathcal{H}^{\mathrm{LR}}, \quad \mathcal{E}_{N F}^{\mathrm{LR}} / \mathcal{A}_{N F}^{\mathrm{LR}}$ | $a_{5}-\frac{a_{7}}{2}, \quad C_{6}-\frac{C_{8}}{2}$ |
| $[u, u, s], \quad[u, u, d]$ | $\mathcal{E}^{\mathrm{LL}} / \mathcal{H}^{\mathrm{LL}}, \quad \mathcal{E}_{N F}^{\mathrm{LL}} / \mathcal{F}_{N F}^{\mathrm{LL}}$ | $a_{2}, \quad C_{2}$ |
|  | $\mathcal{E}^{\mathrm{LR}} / \mathcal{F}^{\mathrm{LR}}, \quad \mathcal{E}_{N F}^{\mathrm{LR}} / \mathcal{H}_{N F}^{\mathrm{LR}}$ | $a_{3}+a_{9}, \quad C_{4}+C_{10}$ |
|  | $\mathcal{E}^{\mathrm{SP}} / \mathcal{H}^{\mathrm{SP}}, \quad \mathcal{E}_{N F}^{\mathrm{SP}} / \mathcal{H}_{N F}^{\mathrm{SP}}$ | $a_{5}+a_{7}, \quad C_{6}+C_{8}$ |
| $[s, u, u], \quad[d, u, u]$ | $\mathcal{E}^{\mathrm{LL}} / \mathcal{F}^{\mathrm{LL}}, \quad \mathcal{E}_{N F}^{\mathrm{LL}} / \mathcal{H}_{N F}^{\mathrm{LL}}$ | $a_{1}, \quad C_{1}$ |
|  | $\mathcal{E}^{\mathrm{LR}} / \mathcal{F}^{\mathbf{L R}}, \quad \mathcal{E}_{N F}^{\mathrm{LR}} / \mathcal{A}_{N F}^{\mathrm{LR}}$ | $a_{4}+a_{10}, \quad C_{3}+C_{9}$ |
|  | $\mathcal{E}^{\mathrm{SP}} / \mathcal{A}^{\mathrm{SP}}, \quad \varepsilon_{N F}^{\mathrm{SP}} / \mathcal{H}_{N F}^{\mathrm{SP}}$ | $a_{6}+a_{8}, \quad C_{5}+C_{7}$ |

- ie. Decay amplitude of $B^{+} \rightarrow \pi^{+} K^{0}$ at NLO

$$
\begin{aligned}
\mathcal{M}\left(B^{+} \rightarrow \pi^{+} K^{0}\right)= & \frac{G_{F}}{\sqrt{2}} V_{u b}^{*} V_{u s}\left[a_{1} \mathcal{A}_{\pi}^{\mathrm{LL}}+C_{1} \mathcal{A}_{N F, \pi}^{\mathrm{LL}}+\mathcal{M}_{B \rightarrow K^{\circ} \pi^{+}}^{(\mathrm{q}, \mathrm{u})}\right]+\frac{G_{F}}{\sqrt{2}} V_{c b}^{*} V_{c s} \mathcal{M}_{B \rightarrow K^{0} \pi^{+}}^{(\mathrm{ql,c}}-\frac{G_{F}}{\sqrt{2}} V_{t b}^{*} V_{t s}\left[\left(a_{4}-\frac{a_{10}}{2}\right) \mathcal{E}_{\pi}^{\mathrm{LL}}\right. \\
& +\left(a_{6}-\frac{a_{8}}{2}\right) \mathcal{E}_{\pi}^{\mathrm{SP}}+\left(C_{3}-\frac{C_{9}}{2}\right) \mathcal{E}_{N F, \pi}^{\mathrm{LL}}+\left(C_{5}-\frac{C_{7}}{2}\right) \mathcal{E}_{N F, \pi}^{\mathrm{LR}}+\left(a_{4}+a_{10}\right) \mathcal{A}_{\pi}^{\mathrm{LL}}+\left(a_{6}+a_{8}\right) \mathcal{A}_{\pi}^{\mathrm{SP}} \\
& \left.+\left(C_{3}+C_{9}\right) \mathcal{A}_{N F, \pi}^{\mathrm{LL}}+\left(C_{5}+C_{7}\right) \mathcal{A}_{N F, \pi}^{\mathrm{LR}}+\mathcal{M}_{B \rightarrow K^{0} \pi^{+}}^{(\mathrm{qLI}, \mathrm{t}}+\mathcal{M}_{B \rightarrow K^{\circ} \pi^{+}}^{(\mathrm{mp})}\right],
\end{aligned}
$$

$\triangle$ the glauber gluon corrections and TMD wave functions are not taken into account in this work

## $B \rightarrow P P, P V, V V$ decays: Amplitudes

- Operator decomposition of $B \rightarrow P P$ decays

Tree/color-favoured tree emission

QCD Penguin

Color-suppressed tree emission
$\mathbf{P e m}_{\mathrm{em}}$ : Electroweak penguin

E: tree annihilation amplitude

| Topology | Channel |
| :---: | :---: |
| $\left\{\mathbf{P}, \mathrm{T}, \mathrm{C}, \mathrm{E}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{0} K^{+}, \eta_{q} K^{+}$ |
| $\left\{\mathbf{T}, \mathrm{P}, \mathrm{C}, \mathrm{E}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{+} \eta_{q}$ |
| $\left\{\mathbf{T}, \mathrm{C}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{+} \pi^{0}$ |
| $\left\{\mathbf{P}, \mathrm{E}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{+} K^{0}, \eta_{s} K^{+}, K^{+} K^{0}$ |
| $\left\{\mathbf{P}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{+} \eta_{s}$ |
| $\left\{\mathbf{T}, \mathrm{P}, \mathrm{E}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{+} \pi^{-}$ |
| $\left\{\mathbf{P}, \mathrm{T}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{-} K^{+}$ |
| $\left\{\mathbf{C}, \mathbf{E}, \mathbf{P}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{0} \pi^{0}, \pi^{0} \eta_{q}, \eta_{q} \eta_{q}$ |
| $\left\{\mathbf{P}, \mathrm{C}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\pi^{0} K^{0}, \eta_{q} K^{0}$ |
| $\left\{\mathbf{P}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $\eta_{s} K^{0}, K^{0} \bar{K}^{0}, \pi^{0} \eta_{s}, \eta_{s} \eta_{s}, \eta_{q} \eta_{s}$ |
| $\left\{\mathbf{E}, \mathrm{P}, \mathrm{P}_{\mathrm{ew}}\right\}$ | $K^{+} K^{-}$ |

## $B \rightarrow P P, P V, V V$ decays: Numerics

- Main uncertainties of PQCD calculation: high order QCD corrections \& LCDAs $\downarrow$
$\dagger$ characterized by the variation in the factorization scale
$\dagger$ minimized by setting $\mu_{t}$ as the largest virtuality in hard scattering
$\dagger$ two-loop expression for the strong coupling
- Input parameters of meson LCDAs

| Meson | $\pi^{ \pm} / \pi^{0}$ | $K^{ \pm} / K^{0}$ | $\eta_{q}$ | $\eta_{z}$ |
| :---: | :---: | :---: | :---: | :---: |
| $m / \mathrm{GeV}[108]$ | $0.140 / 0.135$ | $0.494 / 0.498$ | 0.104 | 0.705 |
| $f / \mathrm{GeV}$ | $0.130[108]$ | $0.156[108]$ | $0.125[114]$ | $0.177[114]$ |
| $m_{0} / \mathrm{GeV}$ | 1.400 | $1.892[112]$ | 1.087 | 1.990 |
| $a_{1}$ | 0 | $0.076 \pm 0.004[113]$ | 0 | 0 |
| $a_{2}$ | $0.270 \pm 0.047[14]$ | $0.221 \pm 0.082[113]$ | $0.250 \pm 0.150[115]$ | $0.250 \pm 0.150[115]$ |
| Meson | $\rho^{ \pm} / \rho^{0}$ | $K^{+ \pm} / K^{* 0}$ | 0 | $\phi$ |
| $m / \mathrm{GeV}[108]$ | 0.775 | 0.892 | 0.783 | 0.019 |
| $f^{\\| l} / \mathrm{GeV}[9]$ | $0.210 / 0.213$ | 0.204 | 0.197 | 0.233 |
| $f^{\perp} / \mathrm{GeV}$ | $0.144 / 0.146[116]$ | $0.159[9]$ | $0.162[9]$ | 0 |
| $a_{1}^{\\| l}$ | 0 | $0.060 \pm 0.040[117]$ | 0 | 0 |
| $a_{1}^{\perp}$ | $0.040 \pm 0.030[117]$ | $0.150 \pm 0.120[117]$ | $0.230 \pm 0.080[117]$ |  |
| $a_{2}^{\\| l}$ | $0.180 \pm 0.037[116]$ | $0.160 \pm 0.090[117]$ | $0.140 \pm 0.120[117]$ | $0.140 \pm 0.070[117]$ |
| $a_{2}^{\perp}$ | $0.137 \pm 0.030[116]$ | $0.100 \pm 0.080[117]$ |  | 0 |

default scale 1 GeV

## $B \rightarrow P P, P V, V V$ decays: Numerics

- Anatomy of NLO corrections to $\mathcal{B}$ and $\mathcal{A}_{\mathrm{CP}}$ of $\pi \pi, \pi K$ modes

| Mode | LO | +VC | +QL | +MP | $+\mathcal{F}^{\text {NLO }}$ | PDG [108] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathcal{B}\left(B^{+} \rightarrow \pi^{+} \pi^{0}\right)$ | 3.58 | 3.89 | ... | ... | $4.18_{-0.97}^{+1.32}$ | $5.5 \pm 0.4$ |
| $\mathcal{A l}_{C P}$ | -0.05 | 0.09 | ... | ... | $0.08_{-0.09}^{+0.09}$ | $3 \pm 4$ |
| $\mathcal{B}\left(B^{0} \rightarrow \pi^{+} \pi^{-}\right)$ | 6.97 | 6.82 | 6.92 | 6.76 | $7.31_{-1.72}^{+2.38}$ | $5.12 \pm 0.19$ |
| $C_{\pi^{+} \pi^{-}}$ | -23.4 | -27.6 | -13.8 | -13.3 | $-12.8_{-3.3}^{+3.5}$ | $-32 \pm 4$ |
| $S_{\pi^{+} \pi^{-}}$ | -31.1 | -35.5 | -46.4 | -37.0 | $-36.4_{-1.5}^{+1.5}$ | $-65 \pm 4$ |
| $\mathcal{B}\left(B^{0} \rightarrow \pi^{0} \pi^{0}\right)$ | 0.14 | 0.29 | 0.30 | 0.22 | $0.23_{-0.05}^{+0.07}$ | $1.59 \pm 0.26$ |
| $C_{\pi^{0} \pi^{0}}$ | -3.1 | 60.1 | 73.6 | 77.6 | $80.2_{-6.7}^{+5.2}$ | $33 \pm 22$ |
| $\mathcal{B}\left(B^{+} \rightarrow \pi^{+} K^{0}\right)$ | 17.0 | 20.8 | 28.0 | 19.4 | $20.3_{-4.4}^{+6.3}$ | $23.7 \pm 0.8$ |
| $\mathcal{A}_{C P}$ | -1.19 | -0.95 | -0.06 | -0.08 | $-0.08_{-0.09}^{+0.08}$ | $-1.7 \pm 1.6$ |
| $\mathcal{B}\left(B^{+} \rightarrow \pi^{0} K^{+}\right)$ | 10.0 | 12.75 | 16.76 | 11.92 | $12.3_{-2.7}^{+3.8}$ | $12.9 \pm 0.5$ |
| $\mathcal{A}_{C P}$ | -10.9 | -5.20 | 2.26 | 2.48 | $2.28_{-1.74}^{+1.61}$ | $3.7 \pm 2.1$ |
| $\mathcal{B}\left(B^{0} \rightarrow \pi^{-} K^{+}\right)$ | 14.3 | 18.0 | 23.9 | 16.4 | $17.1_{-3.7}^{+5.2}$ | $19.6 \pm 0.5$ |
| $\mathcal{H c}_{C P}$ | -15.2 | -14.2 | -4.16 | -5.42 | $-5.43_{-2.34}^{+2.24}$ | $-8.3 \pm 0.4$ |
| $\mathcal{B}\left(B^{0} \rightarrow \pi^{0} K^{0}\right)$ | 5.90 | 8.12 | 10.4 | 6.99 | $7.38_{-1.50}^{+2.11}$ | $9.9 \pm 0.5$ |
| $C_{x^{0} K^{0}}$ | -2.62 | -7.31 | -6.57 | -7.97 | $-7.70_{-0.13}^{+0.21}$ | $0 \pm 13$ |
| $S_{\pi^{0} K^{0}}$ | 70.1 | 73.5 | 71.6 | 71.9 | $71.9_{-0.6}^{+0.6}$ | $58 \pm 17$ |

$\dagger \mathcal{B}: Q L$ and MP corrections cancel, VC and NLO ffs do not have a significant effect
$\dagger$ NLO corrections change asymmetry parameters more significantly than $\mathcal{B}$
$\dagger \mathrm{VC}(\mathrm{QL})$ flips the sign of the direct CPV of $\pi^{+} \pi^{0}$ and $\pi^{0} \pi^{0}\left(\pi^{0} K^{+}\right)$modes $\mathcal{A}_{C P}\left(B^{+} \rightarrow K^{+} \pi^{0}\right)-\mathcal{A}_{C P}\left(B^{+} \rightarrow K^{+} \pi^{-}\right)=7.71_{-2.92}^{+2.74}(\mathrm{PQCD})$ vs $12.0 \pm 2.4$ (Data)
$\dagger$ Color-suppressed modes $\left(\pi^{0} \pi^{0}, \pi^{0} K^{0}\right)$ are more sensitive to NLO corrections.
$\dagger$ PQCD shows a large direct CPV in $\pi^{-} K^{+}, \pi^{+} \pi^{-}$modes in 2000 (LO), which are confirmed by BABRA and Belle afterward.

## $B \rightarrow P P, P V, V V$ decays: Numerics

- Updated PQCD results for the branching ratios of $B \rightarrow P P$ decays (in units of $10^{-6}$ )

| Mode | PQCD | SCET1 [125] | SCET2 [125] | QCDF [127] | PDG [108] |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $B^{+} \rightarrow \pi^{+} K^{0}$ | $20.3{ }_{-4.4-0.1}^{+6.0 .1}$ | ... | $\cdots$ | $21.7_{-9.1}^{+13.4}$ | $23.7 \pm 0.8$ |
| $B^{+} \rightarrow \pi^{0} K^{+}$ | $12.3{ }_{-2.7-0.1}^{+3.8+0.1}$ | ... | ... | $12.5{ }_{-4.8}^{+6.8}$ | $12.9 \pm 0.5$ |
| $B^{+} \rightarrow \eta^{\prime} K^{+}$ | $52.0{ }_{-10.8}^{+15.0+2.7}$ | $69.5 \pm 28.4$ | $69.3 \pm 27.7$ | $74.5{ }_{-31.6}^{+63.6}$ | $70.4 \pm 2.5$ |
| $B^{+} \rightarrow \eta K^{+}$ | $6.688_{-1.60-0.96}^{+2.26+1.85}$ | $2.7 \pm 4.8$ | $2.3 \pm 4.5$ | $2.2{ }_{-1.3}^{+20}$ | $2.4 \pm 0.4$ |
| $B^{+} \rightarrow K^{+} K^{0}$ | $1.56{ }_{-0.34-0.02}^{+0.48+0.02}$ | ... | ... | $1.8{ }_{-0.7}^{+1.1}$ | $1.31 \pm 0.17$ |
| $B^{+} \rightarrow \pi^{0} \pi^{+}$ | $4.18_{-0.94-0.22}^{+1.30+0.22} 4.45$ | ... | $\cdots$ | $5.9{ }_{-1.6}^{+2.6}$ | $5.5 \pm 0.4$ |
| $\eta_{q^{-}} \eta_{s} \text { mixing } B_{B^{+} \rightarrow \pi^{+} \eta^{\prime}}$ | $2.00_{-0.42}^{+0.57}+0.31$ | $2.4 \pm 1.3$ | $2.8 \pm 1.3$ | $3.8{ }_{-0.6}^{1.6}$ | $2.7 \pm 0.9$ |
| $\checkmark \eta_{q}-\eta_{s^{-}} \eta_{g} \text { mixingés } \rightarrow \pi^{+} \eta$ | $2.62_{-0.57-0.40}^{+0.78+0.45}$ | $4.9 \pm 2.0$ | $5.0 \pm 2.1$ | $5.0{ }_{-0.9}^{+1.5}$ | $4.02 \pm 0.27$ |
| [Fan 2012] $B^{0} \rightarrow \pi^{-} K^{+}$ | 17.1-3.7.7-0.1 | ... | . $\cdot$ | $19.3+11.4$ | $19.6 \pm 0.5$ |
| $B^{0} \rightarrow \pi^{0} K^{0}$ | $7.388_{-1.50-0.04}^{+2.11+0.03}$ | ... | ... | $8.6{ }_{-3.6}^{+5.4}$ | $9.9 \pm 0.5$ |
| $B^{0} \rightarrow \eta^{\prime} K^{0}$ | $52.3{ }_{-10.8-0.3}^{+14.9+2.1}$ | $63.2 \pm 26.3$ | $62.2 \pm 25.4$ | $70.9{ }_{-29.8}^{+59.1}$ | $66 \pm 4$ |
| $B^{\circ} \rightarrow \eta K^{0}$ | $4.63_{-1.09-0.79}^{+1.57+1.51}$ | $2.4 \pm 4.4$ | $2.3 \pm 4.4$ | $1.5{ }_{-1.1}^{+1.7}$ | $1.23_{-0.24}^{+0.27}$ |
| $B^{0} \rightarrow K^{0} K^{0}$ | $1.48_{-0.33-0.000}^{+0.47+0.01}$ | ... | . $\cdot$ | $2.1{ }_{-0.8}^{+1.3}$ | $1.21 \pm 0.16$ |
| $B^{0} \rightarrow K^{+} K^{-}$ | $0.046_{-0.039-0.008}^{+0.058+0.09}$ | ... | ... | $0.1 \pm 0.04$ | $0.078 \pm 0.015$ |
| $B^{0} \rightarrow \pi^{+} \pi^{-}$ | $7.31_{-1.68-0.36}^{+235+0.38} 5.35$ | ... | $\ldots$ | $7.0{ }_{-1.0}^{+0.8}$ | $5.12 \pm 0.19$ |
| $\checkmark$ Glauber gluoh effect $\pi^{0} \pi^{0}$ |  | ... | ... | $1.1_{-0.5}^{+1.2}$ | $1.59 \pm 0.26$ |
| [Liu 2014] $\quad B^{0} \rightarrow \pi^{0} \eta^{\prime}$ | $0.20_{-0.03-0.01}^{+0.05+0.02}$ | $2.3 \pm 2.8$ | $1.3 \pm 0.6$ | $0.42_{-0.15}^{+0.28}$ | $1.2 \pm 0.6$ |
| $B^{0} \rightarrow \pi^{0} \eta$ | $0.20_{-0.04-0.01}^{+0.05+0.02}$ | $0.88 \pm 0.68$ | $0.68 \pm 0.62$ | $0.36_{-0.11}^{+0.13}$ | $0.41 \pm 0.17$ |
| $B^{0} \rightarrow \eta \eta$ | $0.37_{-0.07-0.07}^{+0.09+0.08}$ | $0.69 \pm 0.71$ | $1.0 \pm 1.5$ | $0.32_{-0.08}^{+0.15}$ | $<1$ |
| $B^{0} \rightarrow \eta^{\prime}$ | $0.29_{-0.05-0.06}^{+0.07}$ | $1.0 \pm 1.6$ | $2.2 \pm 5.5$ | $0.36{ }_{-0.13}^{+0.27}$ | $<1.2$ |
| $B^{0} \rightarrow \eta^{\prime} \eta^{\prime}$ | $0.42_{-0.07-0.11}^{+0.09+0.13}$ | $0.57 \pm 0.73$ | $1.2 \pm 3.7$ | $0.22_{-0.08}^{+0.16}$ | <1.7 |

$\dagger$ NLO corrections play an important role in penguin dominated models $\pi K, \eta^{\prime} K$ and pure annihilation mode $K^{0} K^{0}$
$\dagger$ PQCD predicted $\mathcal{B}\left(B_{s} \rightarrow \pi^{+} \pi^{-}\right) \sim 6 \times 10^{-6}$ in 2007 (LO), confirmed by CDF in 2011

## $B \rightarrow P P, P V, V V$ decays: Numerics

- Updated PQCD results for the CPV of $B \rightarrow P P$ decays (in units of $10^{-2}$ )



## $B \rightarrow P P, P V, V V$ decays: Numerics

- Updated PQCD results for the branching ratios of $B^{+} \rightarrow P V$ decays (in units of $10^{-6}$ )

$\dagger$ NLO corrections play an important role in $\phi, \omega$ involved modes, $\omega-\phi$ mixing ?


## $B \rightarrow P P, P V, V V$ decays: Numerics

- Updated PQCD results for the CPV of $B^{+} \rightarrow P V$ decays (in units of $10^{-2}$ )

|  | Mode | PQCD | SCET1 [128] | SCET2 [128] | QCDF [127] | PDG [108] |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $B^{+} \rightarrow \eta^{\prime} K^{++}$ | $1.54{ }_{-8.16-9.74}^{+9.05}$ | $2.7{ }_{-19.5}^{\text {27.4 }}$ | ${ }^{2.6}{ }_{-32.9}^{+26.7}$ | $65.5+635.9$ | $-26 \pm 27$ |
|  | $B^{+} \rightarrow \eta K^{++}$ | $-34.5+2.5+0.08$ | -2.6 ${ }_{5}^{+5.4}$ | -1.9 ${ }^{+3.4}$ | -9.7 ${ }_{8.0}^{+7.3}$ | $2 \pm 6$ |
| large CPV predictions | $B^{+} \rightarrow K^{+} \omega$ | $31.5{ }_{-1.1}^{+0.6+0.7}$ | $11.6_{-20.4}^{+18.2}$ | $12.3{ }_{-173}^{16.6}$ | $22.1{ }_{-18,2}$ | $-2 \pm 4$ |
|  | $B^{+} \rightarrow \pi^{+} K^{* 0}$ | $-0.944_{-0.29}^{+0.26+0.04}$ | 0 | 0 | $0.4{ }_{-4.2}^{\text {+4. }}$ | $-4 \pm 9$ |
|  | $B^{+} \rightarrow \pi^{0} K^{+}$ | $-0.011_{-4.87-1.26}^{+4.40+12}$ | $-17.8{ }_{-24.7}^{+30.4}$ | $-12.9{ }_{-12.2}^{+12.0}$ | $1.6{ }_{-4.2}^{+11.5}$ | $-39 \pm 21$ |
|  | $B^{+} \rightarrow K^{+} \rho^{0}$ | $58.7{ }_{-4.0-2.8}^{+4.3+3}$ | $9.2 .{ }_{-16.1}^{+15.2}$ | $16.0{ }_{-225}^{+20.5}$ | 45.4-36.2 | $37 \pm 10$ |
|  | $B^{+} \rightarrow K^{0} \rho^{+}$ | $0.99_{-0.01-0.18}^{+0.01+0.13}$ | 0 | 0 | $0.3{ }_{-0.3}^{+0.5}$ | $-3 \pm 15$ |
| large CPV in rare deca | $B^{+} \rightarrow K^{+} \bar{K}^{* 0}$ | $21.3{ }_{-5.7-1.4}^{+6.1 .2}$ | ${ }_{-3.6}{ }_{-5.3}^{+6.1}$ | $-4.4{ }_{-4.1}^{+4.1}$ | $-8.9{ }_{-2,6}^{+3.0}$ | $12 \pm 10$ |
|  | $B^{+} \rightarrow K^{+} \phi$ | $-1.93_{-0.60-0.42}^{+0.66+0.66}$ | 0 | 0 | $0.6{ }_{-0.1}^{+0.1}$ | $2.4 \pm 2.8$ |
|  | $B^{+} \rightarrow \pi^{+} \phi$ | 0.0 | $\cdots$ | ... | 0.0 | $1 \pm 5$ |
|  | $B^{+} \rightarrow \pi^{+} \omega$ | -29.8.0.4.9+0.8 | $0.5{ }_{-19.6}^{+19.1}$ | $2.3{ }_{-13.2}^{+13.4}$ | ${ }_{-13.2}{ }_{-10.9}$ | $-4 \pm 5$ |
|  | $B^{+} \rightarrow \pi^{+} \rho^{0}$ | 14.9 -0.4+0.6 | $-10.8{ }_{-12.7}^{+13.7}$ | ${ }_{-19.2}+15.6$ | $-9.88_{-10.5}^{+11.9}$ | $0.9 \pm 1.9$ |
|  | $B^{+} \rightarrow \pi^{0} \rho^{+}$ | $-7.31_{-0.02}^{+0.06+0.03}$ | $15.5{ }_{-19.0}^{+17.0}$ | $12.33_{-10}^{+9.4}$ | $9.7_{-10.8}^{+8.3}$ | $2 \pm 11$ |
| $\eta_{q}-\eta_{s}$ mixing | $B^{+} \rightarrow \eta^{\prime} \rho^{+}$ | $29.0{ }_{-0.4-0.1}^{+0.4+0.0}$ | $-19.8{ }_{-37.6}^{+66.6}$ | $-21.7{ }_{-24.3}^{+1359}$ | $1.4{ }_{-11.9}^{+14.0}$ | $26 \pm 17$ |
| $? \eta_{q^{-}} \eta_{s^{-}} \eta_{g}$ mixing | $B^{+} \rightarrow \eta \rho^{+}$ | -13.0.0.0.1+15 | -6.6 ${ }_{2}^{+21.5}$ | -9.1+16.78 | -8.5 ${ }_{5}^{+6.5}$ | $11 \pm 11$ |

$\dagger$ The measured direct CPV in $B \rightarrow P V$ is significantly larger than that in $B \rightarrow P P$
$\dagger$ It is hard to measure $B \rightarrow P V$ decays precisely $\Leftarrow$ vector meson is not stable
$\dagger$ Three-body $B$ decays along with intermediate $B \rightarrow P V$ decays, but difficult to resolve

## $B \rightarrow P P, P V, V V$ decays: Numerics

- Updated PQCD results for the branching ratios of $B^{+} \rightarrow V V$ decays (in units of $10^{-6}$ )
isospin symmetry
smallness of $\mathcal{B}\left(\rho^{0} \rho^{0}\right)$
$\quad \downarrow$
$\mathcal{B}\left(\rho^{+} \rho^{-}\right) \sim 2 \mathcal{B}\left(\rho^{+} \rho^{0}\right)$

| Mode | PQCD | SCET [130] | QCDF [127,131] | PDG [108] |
| :---: | :---: | :---: | :---: | :---: |
| $B^{+} \rightarrow \rho^{+} K^{* 0}$ | $9.40_{-1.34-0.95}^{+1.43+1.5}$ | $8.93 \pm 3.18$ | $9.2_{-5.5}^{+3.8}$ | $9.2 \pm 1.5$ |
| $B^{+} \rightarrow \rho^{0} K^{++}$ | $6.25_{-0.84-0.53}^{+1.12+0.59}$ | $4.64 \pm 1.37$ | $5.5_{-2.5}^{+1.4}$ | $4.6 \pm 1.1$ |
| $B^{+} \rightarrow \omega K^{*+}$ | $5.48_{-1.36-0.66}^{+1.52+0.8}$ | $5.56 \pm 1.60$ | $3.0_{-1.5}^{+2.5}$ | $<7.4$ |



```
PQCD: ~ 1.6
Data: ~ 1
```

$B^{0} \rightarrow \omega K^{* 0} \quad 5.93_{-0.73-1.55}^{+0.89+1.74}$
$B^{0} \rightarrow \phi K^{* 0} \quad 11.8_{-1.3-1}^{+1.6+1}$
$\begin{array}{llll}B^{0} \rightarrow K^{* 0} \bar{K}^{* 0} & 0.38_{-0.06-0.01}^{+0.09+0.02} & 0.48 \pm 0.16 & 0.6_{-0.3}^{+0.2}\end{array}$
$B^{0} \rightarrow K^{*+} K^{*-} \quad 0.17_{-0.02-0.03}^{+0.02+0.05}$

| $P Q C D: \mathcal{B}\left(\rho^{0} \rho^{0}\right) \sim 2 \mathcal{B}\left(\pi^{0} \pi^{0}\right)^{B^{0} \rightarrow \rho^{+} \rho^{-}} \begin{aligned} & B^{0} \rightarrow \rho^{0} \rho^{0}\end{aligned}$ | $0.17_{-0.02-0.03}^{+0.02+0.05}$ | $\cdots$ | $0.16{ }_{-0.1}^{+0.1}$ | $<2.0$ |
| :---: | :---: | :---: | :---: | :---: |
|  | $22.7{ }_{-4.8-0.6}^{+6,3+0.6}$ | $27.7 \pm 4.1$ | $25.5{ }_{-3.0}^{+2.8}$ | $27.7 \pm 1.9$ |
|  | $0.000 .54_{-0.11-0.04}^{+0.16+0.04}$ | $1.00 \pm 0.29$ | - $\begin{array}{r}0.99_{-0.5}^{+1.9} \\ \hline\end{array}$ | $0.96 \pm 0.15$ |
| Data : $\mathcal{B}\left(\rho^{0} \rho^{0}\right) \sim \mathcal{B}\left(\pi^{0} \pi^{0}\right) /{ }^{2}{ }_{B^{0} \rightarrow \rho^{0} \omega}$ |  | $0.59 \pm 0.19$ | incoo.s.jisbed | $0 \times 1.6$ |
| $B^{0} \rightarrow \rho^{0} \phi$ | $0.019_{-0.004-0.001}^{+0.005+0.002}$ | $\sim 0.002$ | ... | $<3.3$ |
| $B^{0} \rightarrow \omega \omega$ | $1.21_{-0.19-0.24}^{+0.24+0.31}$ | $0.39 \pm 0.13$ | $0.7_{-0.4}^{+1.1}$ | $1.2 \pm 0.4$ |
| $B^{0} \rightarrow \omega \phi$ | $0.018_{-0.004-0.005}^{+0.005+0.005}$ | $\sim 0.002$ | $\cdots$ | $<0.7$ |
| $B^{0} \rightarrow \phi \phi$ | $0.029_{-0.002-0.006}^{+0.002+0.006}$ | ... | $\cdots$ | $<0.027$ |

$\dagger$ NLO corrections play an important role in rare modes $\rho^{+} \phi, \rho^{0} \rho^{0}(\omega, \rho), \omega \omega(\phi)$

## $B \rightarrow P P, P V, V V$ decays: Numerics

$$
\begin{array}{r}
\sqrt{2} \mathcal{M}\left(B^{+} \rightarrow \pi^{+} \pi^{0}\right)=\mathcal{M}\left(B^{0} \rightarrow \pi^{+} \pi^{-}\right)-\mathcal{M}\left(B^{0} \rightarrow \pi^{0} \pi^{0}\right) \\
\sqrt{2} \mathcal{M}\left(B^{+} \rightarrow \pi^{+} \rho^{0}+\pi^{0} \rho^{+}\right)=\mathcal{M}\left(B^{0} \rightarrow \pi^{+} \rho^{-}+\pi^{-} \rho^{+}\right)-2 \mathcal{M}\left(B^{0} \rightarrow \pi^{0} \rho^{0}\right) \\
\sqrt{2} \mathcal{M}\left(B^{+} \rightarrow \rho^{+} \rho^{0}\right)=\mathcal{M}\left(B^{0} \rightarrow \rho^{+} \rho^{-}\right)-\mathcal{M}\left(B^{0} \rightarrow \rho^{0} \rho^{0}\right)
\end{array}
$$

- Updated PQCD results for the CPV of $B^{+} \rightarrow V V$ decays (in units of $10^{-2}$ )

| Mode | PQCD | SCET [130] | QCDF [127,131] | PDG [108] |
| :---: | :---: | :---: | :---: | :---: |
| $B^{+} \rightarrow \rho^{+} K^{* 0}$ | $0.58_{-0.12-0.18}^{+0.13+0.16}$ | $-0.56 \pm 0.61$ | $-0.3_{-1}^{+2}$ | $-1 \pm 16$ |
| $B^{+} \rightarrow \rho^{0} K^{*+}$ | $30.6_{-0.7-0.2}^{+0.5+0.1}$ | $29.3 \pm 31.0$ | $43_{-28}^{+13}$ | $31 \pm 13$ |
| $B^{+} \rightarrow \omega K^{*+}$ | $43.0_{-2.0-3.2}^{+1.7+3.8}$ | $24.3 \pm 27.1$ | $29 \pm 35$ | $\cdots$ |
| $B^{+} \rightarrow \phi K^{++}$ | $2.40_{-0.14-0.10}^{+0.14+0.13}$ | $-0.39 \pm 0.44$ | 0.05 | $-1 \pm 8$ |
| $B^{+} \rightarrow K^{*+} \bar{K}^{* 0}$ | $-26.8_{-2.4-2.0}^{+2.3+1.0}$ | $9.5 \pm 10.6$ | $\cdots$ | $\cdots$ |
| $B^{+} \rightarrow \rho^{0} \rho^{+}$ | $0.03_{-0.01-0.00}^{+0.00+0.00}$ | 0.0 | 0.06 | $-5 \pm 5$ |
| $B^{+} \rightarrow \rho^{+} \omega$ | $-25.9_{-1.9-1.2}^{+1.8+1.3}$ | $-13.6 \pm 16.1$ | $-8_{-4}^{+3}$ | $-20 \pm 9$ |
| $B^{+} \rightarrow \rho^{+} \phi$ | 0.0 | 0.0 | $\cdots$ | ... |
| $B^{0} \rightarrow \rho^{-} K^{*+}$ | $32.4_{-0.1-0.2}^{+0.1+0.1}$ | $20.6 \pm 23.3$ | $32_{-14}^{+2}$ | $21 \pm 15$ |
| $B^{0} \rightarrow \rho^{0} K^{* 0}$ | $-14.4_{-1.4-1.0}^{+1.2+0.9}$ | $-3.30 \pm 3.91$ | $-15 \pm 16$ | $-6 \pm 9$ |
| $B^{0} \rightarrow \omega K^{* 0}$ | $9.89_{-0.80-1.12}^{+0.96+1.59}$ | $3.66 \pm 4.05$ | $23_{-18}^{+10}$ | $45 \pm 25$ |
| $B^{0} \rightarrow \phi K^{* 0}$ | $0.86_{-0.06-0.06}^{+0.06+0.07}$ | $-0.39 \pm 0.44$ | $0.8{ }_{-0.5}^{+0.4}$ | $0 \pm 4$ |
| $B^{0} \rightarrow \rho^{+} \rho^{-}$ | $-1.85_{-0.11-0.00}^{+0.20+0.01}$ | $-7.68 \pm 9.19$ | $11_{-4}^{+11}$ | $C_{\rho^{+} \rho^{-}}=0 \pm 9$ |
|  | $-12.7_{-0.1-0.3}^{+0.1+0.4}$ | $\cdots$ | $-19_{-10}^{+9}$ | $S_{\rho^{+} \rho^{-}}=-14 \pm 13$ |
| $B^{0} \rightarrow \rho^{0} \rho^{0}$ | $74.6_{-1.9-2.3}^{+1.3+1.9}$ | $19.5 \pm 23.5$ | $-53_{-54}^{+26}$ | $C_{\rho^{0} \rho^{0}}=20 \pm 90$ |
|  | $1.38_{-0.03-1.93}^{+0.74+2.15}$ | ... | $16_{-49}^{+50}$ | $S_{\rho^{0} \rho^{0}}=30 \pm 70$ |

## $B \rightarrow P P, P V, V V$ decays: Numerics

- Updated PQCD results for the $f_{L}$ of $B^{+} \rightarrow V V$ decays (in units of $10^{-2}$ )

| Mode | $\mathrm{PQCD}_{\text {LO }}$ [51] | PQCD | SCET [130] | QCDF [127,131] | HFLAV [134] |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $B^{+} \rightarrow \rho^{+} K^{* 0}$ | $70.0 \pm 5.0$ | $76.6_{-1.4}^{+1.5}$ | $45.0 \pm 18.0$ | $48.0_{-40.0}^{+52.0}$ | $48 \pm 8$ |
| $B^{+} \rightarrow \rho^{0} K^{++}$ | $75.0_{-5.0}^{+4.0}$ | $80.0_{-1.5}^{+1.5}$ | $42.0 \pm 14.0$ | $67.0_{-48.0}^{+11.0}$ | $78 \pm 12$ |
| $B^{+} \rightarrow \omega K^{++}$ | $64.0 \pm 7.0$ | $77.4_{-0.9}^{+0.5}$ | $53.0 \pm 14.0$ | $67.0_{-39.0}^{+32.0}$ | $41 \pm 19$ |
| $B^{+} \rightarrow \phi K^{*+}$ | $57.0_{-5.9}^{+6.3}$ | $68.7_{-1.5}^{+1.3}$ | $51.0 \pm 16.4$ | $49.0_{-43.0}^{+51.0}$ | $50 \pm 5$ |
| $\boldsymbol{B}^{+} \rightarrow \boldsymbol{K}^{++} \bar{K}^{+0}$ | $74.0 \pm 7.0$ | $82.4{ }_{-1.1}^{+1.1}$ | $50.0 \pm 16.0$ | $45.0+38.0$ | $82_{-21}^{+15}$ |
| $B^{+} \rightarrow \rho^{0} \rho^{+}$ | $98.0 \pm 1.0$ | $96.9_{-0.1}^{+0.1}$ | $\sim 100$ | $96.0 \pm 2.0$ | $95 \pm 1.6$ |
| $B^{+} \rightarrow \rho^{+} \omega$ | $97.0 \pm 1.0$ | $96.3{ }_{-0.4}^{+0.3}$ | $97.0 \pm 1.0$ | $96.0{ }_{-3.0}^{+2.0}$ | $90 \pm 6$ |
| $B^{+} \rightarrow \rho^{+} \phi$ | $95.0 \pm 1.0$ | $81.3_{-1.8}^{+1.9}$ | $\sim 100$ | ... | ... |
| $B^{0} \rightarrow \rho^{-} K^{++}$ | $68.0_{-4.0}^{+5.0}$ | $75.7_{-1.4}^{+1.5}$ | $55 \pm 14$ | $53.0_{-32.0}^{+45.0}$ | $38 \pm 13$ |
| $B^{0} \rightarrow \rho^{0} K^{* 0}$ | $65.0_{-5.0}^{+4.0}$ | $71.0_{-1.3}^{+1.5}$ | $61.0 \pm 13.0$ | $39.0_{-31.0}^{+60.0}$ | $17.3 \pm 2.6$ |
| $B^{0} \rightarrow \omega K^{* 0}$ | $65.0 \pm 5.0$ | $77.7_{-0.9}^{+0.4}$ | $40.0 \pm 20.0$ | $58.0_{-17.0}^{+14.0}$ | $69 \pm 11$ |
| $B^{0} \rightarrow \phi K^{* 0}$ | $56.5_{-5.9}^{+5.8}$ | $69.5_{-1.5}^{+1.2}$ | $51.0 \pm 16.4$ | $50.0_{-44.0}^{+51.0}$ | $49.7 \pm 1.7$ |
| $B^{0} \rightarrow K^{* 0} \bar{K}^{* 0}$ | $58.0 \pm 8.0$ | $68.8_{-5.3}^{+5.3}$ | $50.0 \pm 16.0$ | $52.0_{-49.0}^{+48.0}$ | $74 \pm 5$ |
| $B^{0} \rightarrow K^{*+} K^{*-}$ | $-100.0$ | $-100.0$ | $\ldots$ | $-100.0$ | $\ldots$ |
| $B^{0} \rightarrow \rho^{+} \rho^{-}$ | $95.0 \pm 1.0$ | $93.8_{-0.1}^{+0.1}$ | $99.1 \pm 0.3$ | $92.0_{-3.0}^{+1.0}$ | $99.0_{-1.9}^{+2.1}$ |
| $B^{0} \rightarrow \rho^{0} \rho^{0}$ | $12.0_{-2.0}^{+16.0}$ | $80.9_{-1.9}^{+1.9}$ | $87.0 \pm 5.0$ | $92.0_{-37.0}^{+7.0}$ | $71_{-9}^{+8}$ |
| $B^{0} \rightarrow \rho^{0} \omega$ | $67.0_{-9.0}^{+8.0}$ | $74.2_{-0.1}^{+0.1}$ | $58.0 \pm 14.0$ | $52.0_{-44.0}^{+12.0}$ | ... |
| $B^{0} \rightarrow \rho^{0} \phi$ | $95.0 \pm 1.0$ | $81.3_{-1.8}^{+1.9}$ | $\sim 100$ | . | $\cdots$ |
| $B^{0} \rightarrow \omega \omega$ | $66.0_{-11.0}^{+10.0}$ | $88.4_{-0.8}^{+0.9}$ | $64.0 \pm 15.0$ | $94.0_{-20.0}^{+4.0}$ | ... |
| $B^{0} \rightarrow \omega \phi$ | $94.0_{-3.0}^{+2.0}$ | $80.8_{-1.4}^{+0.8}$ | $\sim 100$ | ... | $\ldots$ |
| $B^{0} \rightarrow \phi \phi$ | $97.0 \pm 1.0$ | $99.9_{-0.0}^{+0.0}$ | ... | ... | ... |

$\dagger$ PQCD showed $f_{L}$ in penguin dominated $B \rightarrow V V$ channels down by annihilation mechanism in 2002 (LO), before the "polarization puzzle" appeared.

## Conclusion

- The up-to-date PQCD predictions with including the current well-known NLO and sub-leading power corrections can explain most of the data.
$\dagger K \pi, K \rho, K \omega, K \phi$ and $K^{*} \rho, K^{*} \omega, K^{*} \phi$ channels $\checkmark \checkmark K^{*} \pi, K^{*} K$ channels $\checkmark$
$\dagger f_{L}$ of $K^{*} \rho, K^{*} \omega, K^{*} \phi$ channels is still larger than the HFLAV result LD effect in $B \rightarrow K^{*}$ transition ? NLO corrections to $B \rightarrow V$ form factors? width effect of the intermediate vector resonant (four-body decays) ?
$\dagger \eta^{(\prime)}$ involved channels do not consist well with data the large mixing mechanism $\eta_{q}-\eta_{s}-\eta_{g}$ provides a possible solution
$\dagger$ The CPV of charged (neutral) $B$ decays is (not) sensitive to the new added two power correction (heavy quark expansion), especially for the channels with at least one $\eta^{(\prime)}$ in the final state.


## Conclusion with opportunities and challenges of PQCD

- Corrections from 3 particle $B$ meson DAs and high twist light meson DAs
$\dagger$ interaction between largely off-shell gluon with three-particle configurations $\mathcal{O}\left(\Lambda / m_{B}\right)$ ?
- Complete NLO calculation for two-body $B$ meson decays
$\dagger$ vertex corrections, $B \rightarrow \rho$ type ff , tensor meson ffs , annihilation spectator amplitude $\ldots$
- Complete NLO calculation for the radiative and $P_{\text {EW }} B$ meson decays
$\dagger B$ meson distribution amplitude
- TMD wave functions of $B$ and $B_{c}$ mesons, $\Lambda_{b}$ baryon
- Systematic power counting with including $k_{T}$
- Sudakov factor of baryon and three particle configuration of meson
- Multibody $B$ decay, more observables, CPV sources, factorization formula
- Input of meson and dimeson DAs, optimal choice of factorization scale


## The End, Thanks.

