The W boson Mass and Muon g-2: Hadronic Uncertainties or New Physics?

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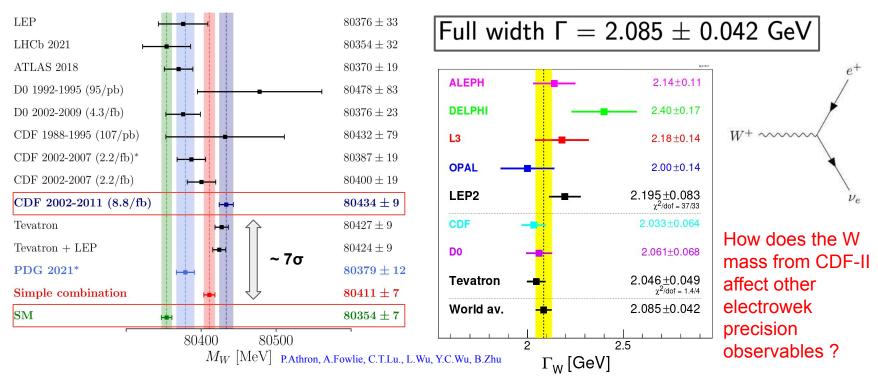
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- 1. Introduction
- 2. Electroweak fits of the W mass and muon g-2
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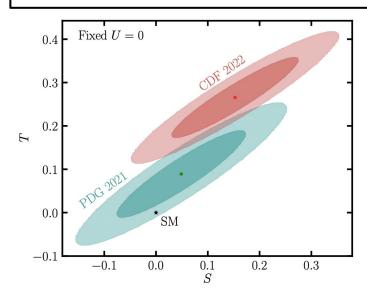
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The W boson mass and width measurements

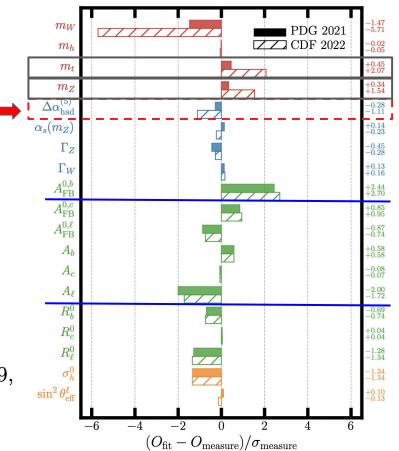


 $^{^*}$ Does not include $13.5\,\mathrm{MeV}$ shift in CDF 2002-2007 (2.2/fb)

The updatedc EW fits after the CDF-II W mass measurement



 $S = 0.06 \pm 0.10, T = 0.11 \pm 0.12, U = 0.13 \pm 0.09,$ $S = 0.14 \pm 0.08, T = 0.26 \pm 0.06 \text{ with } U = 0$

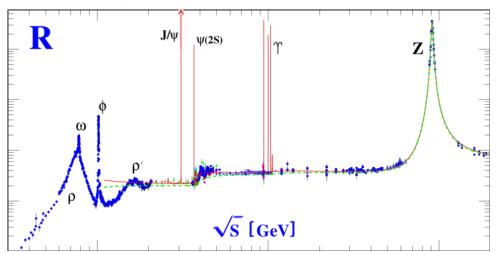


$$\alpha^{-1}(M_Z^2) = \alpha^{-1} \Big[1 - \Delta \alpha_{\rm lep}(M_Z^2) \\ - \Big[\Delta \alpha_{\rm had}^{(5)}(M_Z^2) - \Delta \alpha_{\rm top}(M_Z^2) \Big] \Big] \Big[\Delta \alpha_{\rm had}^{(5)} \big|_{e^+e^-} = 276.1(1.1) \times 10^{-4}. \Big]$$

Fine-structure constant

$$\Delta \alpha_{\text{had}}^{(5)}|_{e^+e^-} = 276.1(1.1) \times 10^{-4}.$$

$$e^+e^- \rightarrow \text{hadrons}$$

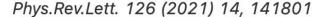


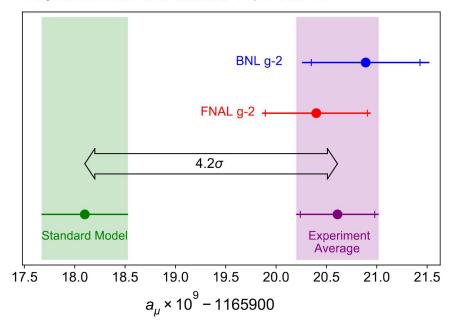
$$\Delta lpha_{
m had}^{(5)}(\emph{M}_{\it Z}^2) = rac{lpha \emph{M}_{\it Z}^2}{3\pi} P \int\limits_{s_{
m thr}}^{\infty} {
m d}s rac{\emph{R}_{
m had}(s)}{s(\emph{M}_{\it Z}^2-s)}$$

$$R_{\mathsf{had}}(s) = \frac{3s}{4\pi\alpha^2} \sigma(e^+e^- \to \mathsf{hadrons})$$

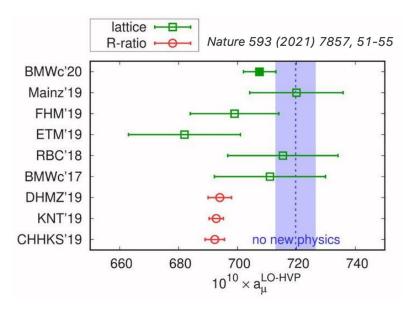
where
$$s_{\rm thr} = m_{\pi^0}^2$$

Muon g-2 excess from BNL and FNAL





Are lattice calculations consistent with the SM value?



The SM contributions to muon g-2

Contribution	Value ×10 ¹¹	_
Experiment (E821)	116 592 089(63)	Time-like :
${\text{HVP LO }(e^+e^-)}$	6931(40)	- Illie-like .
HVP NLO (e^+e^-)	-98.3(7)	m^2 t^{∞} d.
HVP NNLO (e^+e^-)	12.4(1)	$a_{\mu}^{\mathrm{HVP}} = \frac{m_{\mu}^2}{12\pi^3} \int_{m_{\pi^0}^2}^{\infty} \frac{\mathrm{d}s}{s} K(s) \sigma_{\mathrm{had}}(\sqrt{s}),$
HVP LO (lattice, udsc)	7116(184)	$a_{\mu} = \frac{12\pi^3}{12\pi^3} \int_{\mathbb{R}^2} s^{\mathbf{R}(3) \log_{\mathrm{had}}(\sqrt{3})},$
HLbL (phenomenology)	92(19)	-1 $Jm_{\pi^0}^{-1}$
HLbL NLO (phenomenology)	2(1)	
HLbL (lattice, <i>uds</i>)	79(35)	where m_{μ} and m_{π^0} are the muon and neutral pion
HLbL (phenomenology + lattice)	90(17)	masses, respectively, and $K(s)$ is the kernel function
QED	116 584 718.931(104)	
Electroweak	153.6(1.0)	
$HVP(e^+e^-, LO + NLO + NNLO)$	6845(40)	
HLbL (phenomenology + lattice + NLO)	92(18)	
Total SM Value	116 591 810(43)	
Difference: $\Delta a_{\mu} := a_{\mu}^{\rm exp} - a_{\mu}^{\rm SM}$	279(76)	Phys.Rept. 887 (2020) 1-166

What is common for W mass and muon g-2?

W mass:

$$M_W^2 = M_Z^2 \left\{ \frac{1}{2} + \sqrt{\frac{1}{4} - \frac{\pi \alpha}{\sqrt{2} G_\mu M_Z^2}} (1 + \Delta r)} \right\}$$

$$\Delta lpha_{
m had} = rac{M_Z^2}{4\pi^2 lpha} \int_{m_{\pi^0}^2}^{\infty} rac{{
m d}s}{M_Z^2 - s} \sigma_{
m had}(\sqrt{s}),$$



$$a_{\mu}^{\mathrm{SM}}=a_{\mu}^{\mathrm{QED}}+a_{\mu}^{\mathrm{EW}}+a_{\mu}^{\mathrm{HVP}}+a_{\mu}^{\mathrm{HLbL}}$$
 $a_{\mu}^{\mathrm{HVP}}=rac{m_{\mu}^{2}}{12\pi^{3}}\int_{m^{2}}^{\infty}rac{\mathrm{d}s}{s}K(s)\sigma_{\mathrm{had}}(\sqrt{s}),$

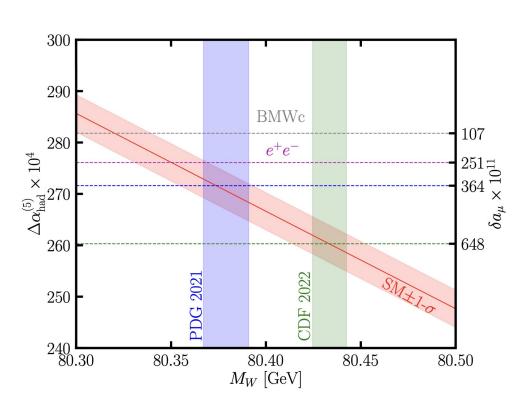
How many ways to determine $\Delta \alpha_{ m had}$ and $a_{\mu}^{ m HVP}$?

	$e^+e^- \to \text{hadrons}$	Lattice QCD	Electroweak fits	
$\Delta lpha_{ m had}$	Yes	Partial (only the results from low energy regions are reported from BMWc!)	Yes	
a_{μ}^{HVP}	Yes	Yes	No (the assumption of transformation is needed!)	

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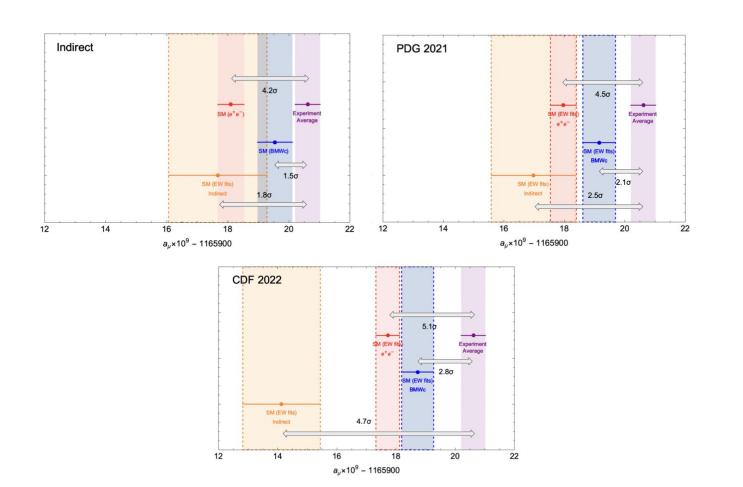
The relation between $\,\Deltalpha_{ m had}\,\,(\delta a_{\mu}\,\,)$ and M_W



EW fits table of the W mass and muon g-2

M_W Indirect					PDG 2021		CDF 2022			
$\Delta lpha_{ m had}$		BMWc	e^+e^-	${\bf Indirect}$	BMWc	e^+e^-	${\bf Indirect}$	BMWc	e^+e^-	Indirect
Innut	$M_W [{ m GeV}]$		-	-	80.379(12)	80.379(12)	80.379(12)	80.4335(94)	80.4335(94)	80.4335(94)
Input	$\Delta \alpha_{ m had} imes 10^4$	281.8(1.5)	276.1(1.1)	_	281.8(1.5)	276.1(1.1)	_	281.8(1.5)	276.1(1.1)	: <u>=</u>
Fitted	$\chi^2/{ m dof}$	18.32/15	16.01/15	15.89/14	23.41/16	18.74/16	17.59/15	74.51/16	62.58/16	47.19/15
	M_W [GeV]	80.348(6)	80.357(6)	80.359(9)	80.355(6)	80.361(6)	80.367(7)	80.375(5)	80.380(5)	80.396(7)
	$\Delta \alpha_{ m had} imes 10^4$	280.9(1.4)	275.9(1.1)	274.4(4.4)	280.3(1.4)	275.6(1.1)	271.7(3.8)	278.6(1.4)	274.7(1.0)	260.9(3.6)
	$\delta a_{\mu} imes 10^{11}$	-	=	294(166)	146(68)	264(59)	364(145)	188(68)	289(57)	648(137)
	Tension	-	-	1.8σ	2.1σ	4.5σ	2.5σ	2.8σ	5.1σ	4.7σ
	$\delta M_W \; [{ m MeV}]$	86(11)	77(11)	75(13)	79(11)	73(11)	67(12)	59(11)	54(11)	38(12)
	Tension	7.8σ	7.0σ	5.8σ	7.2σ	6.6σ	5.6σ	5.4σ	4.9σ	3.2σ

$$\delta M_W \equiv M_W^{\rm CDF} - M_W$$



The key observation

We demonstrate that the two anomalies pull the hadronic contributions in **opposite**directions by performing electroweak fits in which the hadronic contribution was allowed to float.

The fits show that including the g-2 measurement worsens the tension with the CDF measurement and conversely that adjustments that alleviate the CDF tension worsen the g-2 tension beyond 5σ .

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The singlet-triplet scalar leptoquark model

- We consider the singlet-triplet scalar leptoquark(LQ) mdoel to explain the muon g-2 and W boson mass.
- 2. The quantum numbers of the singlet and triplet scalar LQ are defined as

$$S_1(\overline{\bf 3},{\bf 1},1/3) \qquad S_3(\overline{\bf 3},{\bf 3},1/3)$$

3. The relevant Lagrangian terms include

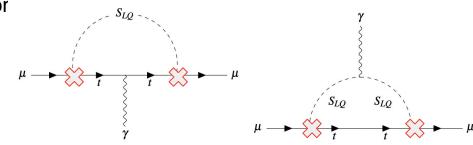
$$\mathcal{L}_{S_1 \& S_3} = \mathcal{L}_{\text{mix}} + \mathcal{L}_{\text{LQ}},$$

$$\mathcal{L}_{\text{mix}} = \lambda H^{\dagger} \left(\vec{\tau} \cdot \overrightarrow{S_3} \right) H S_1^* + \text{ h.c.}$$

$$\mathcal{L}_{\text{LQ}} = y_R^{ij} \bar{u}_{Ri}^C e_{Rj} S_1 + y_L^{ij} \bar{Q}_i^C i \tau_2 \left(\vec{\tau} \cdot \vec{S_3} \right) L_j + \text{h.c.}$$

^{*}Although a coupling between S_1 and the LH lepton and quark fileds is allowed, we don't cosider it here.

1. We The mixing between interaction eigenstates allows the physical mass eigenstates to have both LH and RH couplings to muons and induces chirality flipping enhancements in the one-loop muon g-2 correctior

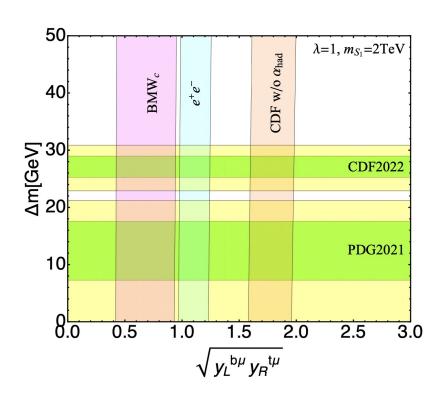


- arXiv: 2204.09031
- 2. Two scalar LQs with $Q = \frac{1}{3}$ can mix through the mixing interaction after EWSB

$$m_{S_{\pm}}^2 = \frac{m_{S_1}^2 + m_{S_3}^2}{2} \pm \frac{1}{2} \sqrt{\left(m_{S_1}^2 - m_{S_3}^2\right)^2 + 4\delta^2}$$
 where $\delta \equiv \lambda v^2/2$.

such mass splitting can generate an extra contribution to the T parameter and provide the shift in W mass from the SM predictuion.

The singlet-triplet scalar leptoquark model



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Conclusions

- We show that the W mass and muon g-2 anomalies pull the hadronic contributions in opposite directions, the new CDF W mass measurement indirectly increases the deviation in muon g-2.
- The singlet-triplet scalar leptoquark model can simultaneously explain both W mass and muon g-2 anomalies.
- 3. The results point to new physics that has large chirality flipping enhancements in the one-loop diagrams for muon g-2 and significant BSM contributions to the oblique T parameter that can be given through custodial symmetry violation.

Thank you for your attention

Back up

The time-like and space-like master formulae for $\Delta lpha_{ m had}$ and $a_{\mu}^{ m HVP}$

1. <u>Time-like</u>: Using for $e^+e^- \to \text{hadrons}$ data calculations

$$egin{align*} egin{align*} egin{align*}$$

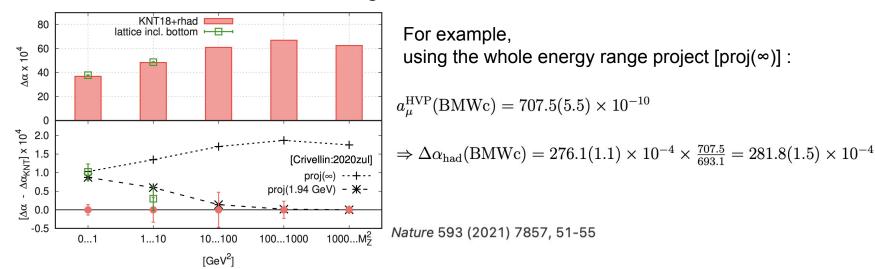
2. Space-like: Using for lattice QCD calculations

$$a_{\mu}^{\mathsf{HVP,LO}} = \left(\frac{\alpha}{\pi}\right)^2 \int_0^{\infty} \mathsf{d}s \, f(s) \hat{\Pi}(-s) \qquad \qquad \hat{\Pi}(s) = 4\pi^2 \Big[\Pi(s) - \Pi(0)\Big]$$

$$\Delta \alpha_{\mathsf{had}}^{(5)}(M_Z^2) = \frac{\alpha}{\pi} \hat{\Pi}(-M_Z^2) + \frac{\alpha}{\pi} \big(\hat{\Pi}(M_Z^2) - \hat{\Pi}(-M_Z^2)\big)$$

The problem to compare $\Delta \alpha_{ m had}$ form data-driven and lattice QCD

- 1. The $\Delta \alpha_{\rm had}$ is calculated at the scale MZ for five quark flavors from data-driven method with $\Delta \alpha_{\rm had}^{(5)}|_{e^+e^-} = 276.1(1.1) \times 10^{-4}$. KNT, DHMZ
- 2. However, we don't have enough information for $\Delta \alpha_{
 m had}$ from lattice QCD side.



The 3rd way to extract $\Delta \alpha_{ m had}$: Global Electrowek Fits

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GFitter

M_W			Indirect		PDG 2021			CDF 2022			
$\Delta\alpha_{\rm had}$		BMWc	e^+e^-	Indirect	BMWc	e^+e^-	Indirect	BMWc	e^+e^-	Indirect	
T4	M_W [GeV]	-	-	-	80.379(12)	80.379(12)	80.379(12)	80.4335(94)	80.4335(94)	80.4335(94)	
Input	$\Delta \alpha_{ m had} imes 10^4$	281.8(1.5)	276.1(1.1)	-	281.8(1.5)	276.1(1.1)		281.8(1.5)	276.1(1.1)	=	
	$\chi^2/{ m dof}$	18.32/15	16.01/15	15.89/14	23.41/16	18.74/16	17.59/15	74.51/16	62.58/16	47.19/15	
Fitted	$M_W [{ m GeV}]$	80.348(6)	80.357(6)	80.359(9)	80.355(6)	80.361(6)	80.367(7)	80.375(5)	80.380(5)	80.396(7)	
	$\Delta \alpha_{ m had} imes 10^4$	280.9(1.4)	275.9(1.1)	274.4(4.4)	280.3(1.4)	275.6(1.1)	271.7(3.8)	278.6(1.4)	274.7(1.0)	260.9(3.6)	
	$\delta a_{\mu} \times 10^{11}$	-	-	294(166)	146(68)	264(59)	364(145)	188(68)	289(57)	648(137)	
	Tension	-		1.8σ	2.1σ	4.5σ	2.5σ	2.8σ	5.1σ	4.7σ	
	$\delta M_W \; [{ m MeV}]$	86(11)	77(11)	75(13)	79(11)	73(11)	67(12)	59(11)	54(11)	38(12)	
	Tension	7.8σ	7.0σ	5.8σ	7.2σ	6.6σ	5.6σ	5.4σ	4.9σ	3.2σ	

Then, how to transform the information between $\Delta \alpha_{\rm had}$ and $a_{\mu}^{\rm HVP}$? Here we consider the whole energy range projection.

Three various projections between $\Delta lpha_{ m had}$ and $a_{\mu}^{ m HVP}$

1. According to *Crivellin:2020zul*, there are three different hypotheses for the projection between

$$\Delta lpha_{
m had}$$
 and $a_{\mu}^{
m HVP}$:

- (1) Low energy for the sum of exclusive channels: $m_{\pi_0} \leq \sqrt{s} \leq 1.937 \, \text{GeV}$,
- (2) Energy below the perturbative contributions : $m_{\pi_0} \leq \sqrt{s} \leq 11.199 \, \text{GeV or}$
- (3) The whole energy range : $m_{\pi_0} \leq \sqrt{s} \leq \infty$,

(<u>Hypothesis</u>: The part above the upper energy threshold is the same as data driven one and the uniform scaling is applied.)

2. Open questions: (1) Which projection should be preferred?

(The low energy projection agrees better with BMWc restuls.)

(2) Can we go beyond the uniform scaling (energy independent) hypothesis?

Using Global EW Fits to extract $\Delta \alpha_{\rm had}$: Low energy projection

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-	$M_W \hspace{1cm} ext{Indirect}$				PDG 2021		CDF 2022			
	$\Delta lpha_{ m had}$	BMWc	e^+e^-	Indirect	BMWc	e^+e^-	Indirect	BMWc	e^+e^-	Indirect
Input	M_W [GeV]	-		=	80.379(12)	80.379(12)	80.379(12)	80.4335(94)	80.4335(94)	80.4335(94)
	$\Delta \alpha_{ m had}^{(5)}(M_Z^2) imes 10^4$	277.4(1.2)	276.1(1.1)	=	277.4(1.2)	276.1(1.1)	-	277.4(1.2)	276.1(1.1)	_
Fitted	$\chi^2/{ m dof}$	16.28/15	16.01/15	15.89/14	19.51/16	18.74/16	17.59/15	65.07/16	62.58/16	47.19/15
	M_W [GeV]	80.355(6)	80.357(6)	80.359(9)	80.360(6)	80.361(6)	80.367(7)	80.379(5)	80.380(5)	80.396(7)
	$\Delta \alpha_{ m had} imes 10^4$	277.1(1.2)	275.9(1.1)	274.4(4.4)	276.8(1.1)	275.6(1.1)	271.7(3.8)	275.6(1.1)	274.7(1.0)	260.9(3.6)
	$\delta a_{\mu} imes 10^{11}$	-	_	438(396)	173(54)	306(54)	748(339)	306(54)	416(54)	1997(320)
	Tension	_	=	1.1σ	3.2σ	5.7σ	2.2σ	5.7σ	7.7σ	6.2σ
	$\delta M_W \; [{ m MeV}]$	79(11)	77(11)	75(13)	74(11)	73(11)	67(12)	55(11)	54(11)	38(12)
	Tension	7.2σ	7.0σ	5.8σ	6.7σ	6.6σ	5.6σ	5.0σ	4.9σ	3.2σ

For the case of low energy projection, $\Delta \alpha_{\rm had}$ is shrunk, but $a_{\mu}^{\rm HVP}$ is enlarged after the transformation compared with the whole energy range projection.