

Pulsar Polarization Arrays

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Based on work with Tao Liu, Xuzixiang Lou, arXiv: 2111.10615

Content

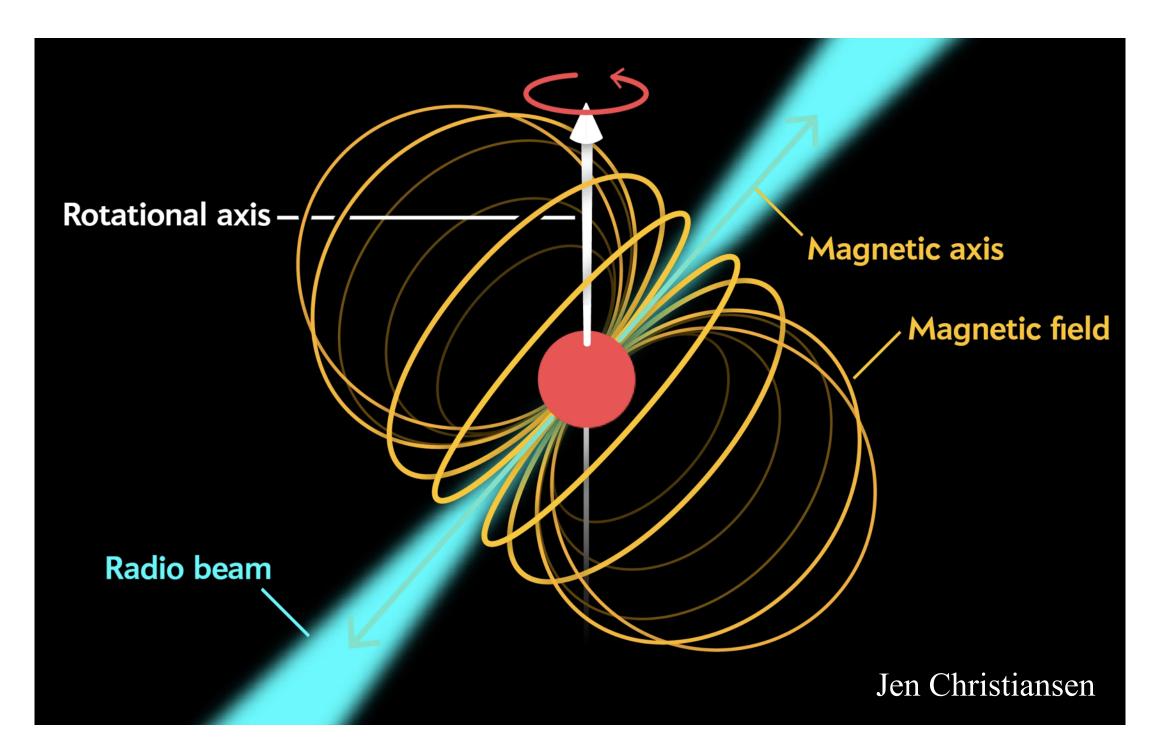
◆ Pulsar timing arrays (PTAs) and pulsar polarization arrays (PPAs)

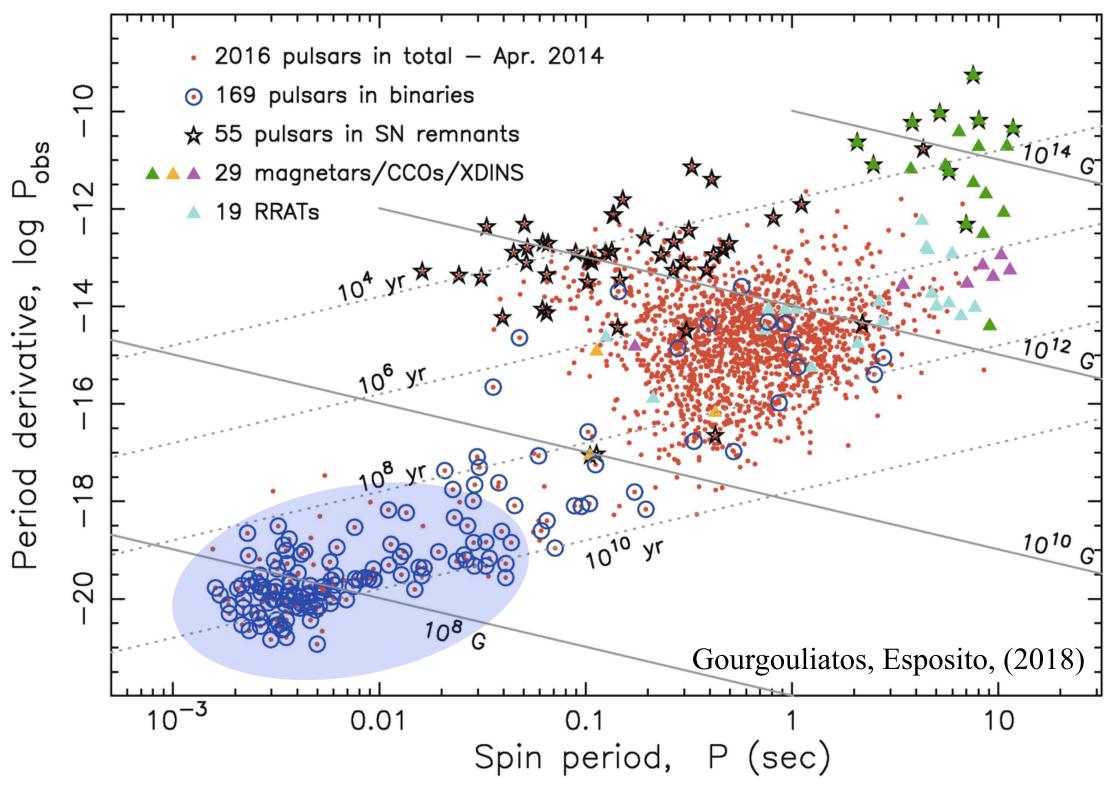
Ultralight axion-like dark matter (ALDM)

Ultralight ALDM detection by PPAs through cosmic birefringence

Pulsars

- ◆ Pulsars emit electromagnetic pulses with extraordinary regularity, with the period ranging from milliseconds to seconds. Up to now about 3000 pulsars observed in our galaxy
- Millisecond pulsars (MSPs) are especially stable due to mass and angular momentum transfer from a companion. Although emission mechanism not fully clarified, they play significant roles as astronomical clocks



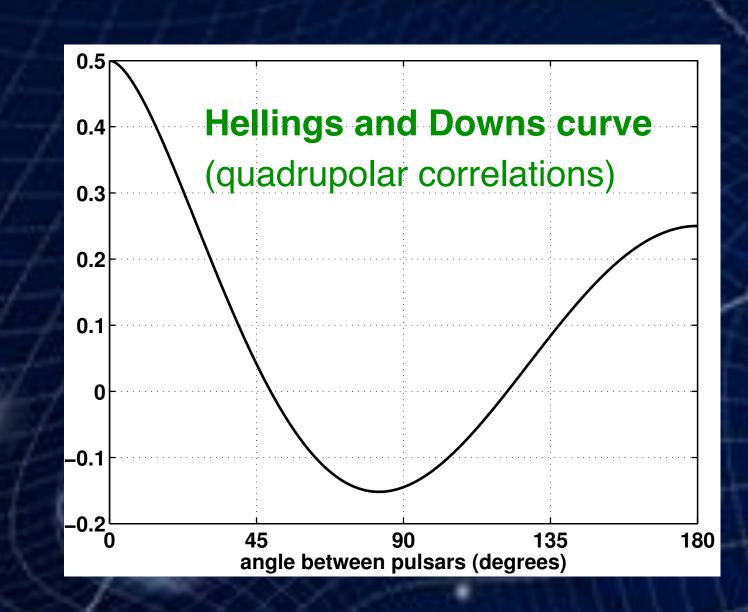


Pulsar timing arrays (PTAs)

- → Global PTA network monitors over 80 MSPs in the timespan of years. FAST/SKA may increase the number to the order of 1000
- * Given precise timing model of the expected time of arrival of the pulse, the measured time difference can be directly related to gravitational waves (GWs)

$$\Delta T(t) = \int_{-\infty}^{\infty} df \frac{1}{2} u^a u^b h_{ab}(f,\hat{n}) \, e^{i2\pi f(t+\hat{n}\cdot\vec{r}_2/c)} \, \frac{1}{i2\pi f} \, \frac{1}{1+\hat{n}\cdot\hat{u}} \left[1 - e^{-\frac{i2\pi fL}{c}(1+\hat{n}\cdot\hat{u})}\right]$$
 metric perturbation "Earth" "pulsar"

* PTAs with MSPs in different directions serve as galactic interferometers to measure nHZ GWs. Stochastic GW backgrounds (SGWBs) can be identified by the quadrupolar spatial correlations among pulsars

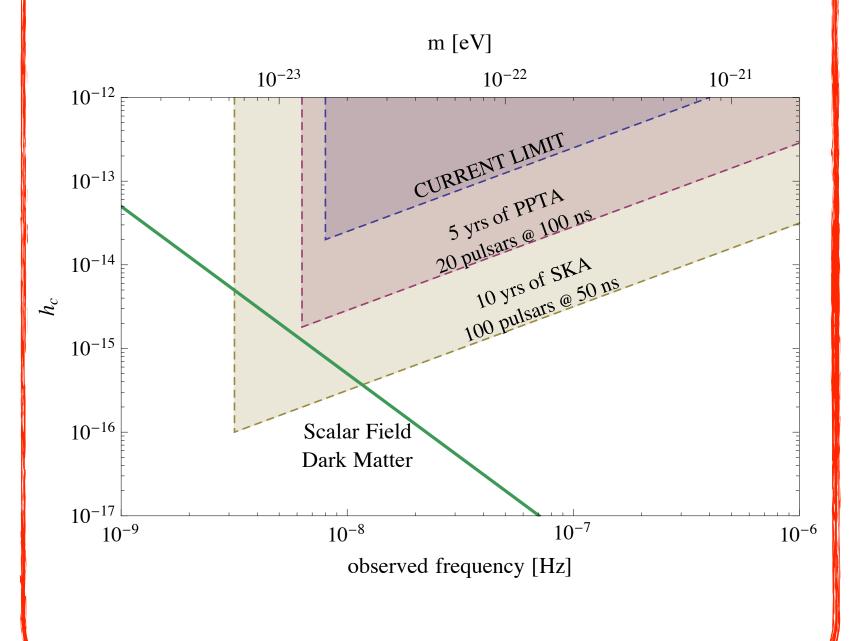


Pulsar timing arrays (PTAs)

Recently, the explorations on the PTA targets extended to dark matter (DM) physics

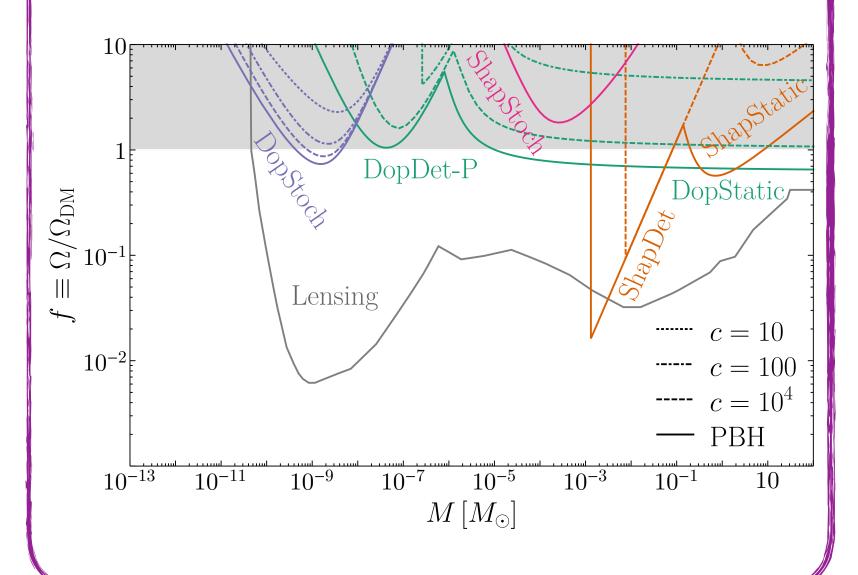
Oscillating gravitational potential induced by ultralight scalar DM

Khmelnitsky and Rubakov, 1309.5888; De Martino et al. 1705.04367; Porayko et al., 1810.03227



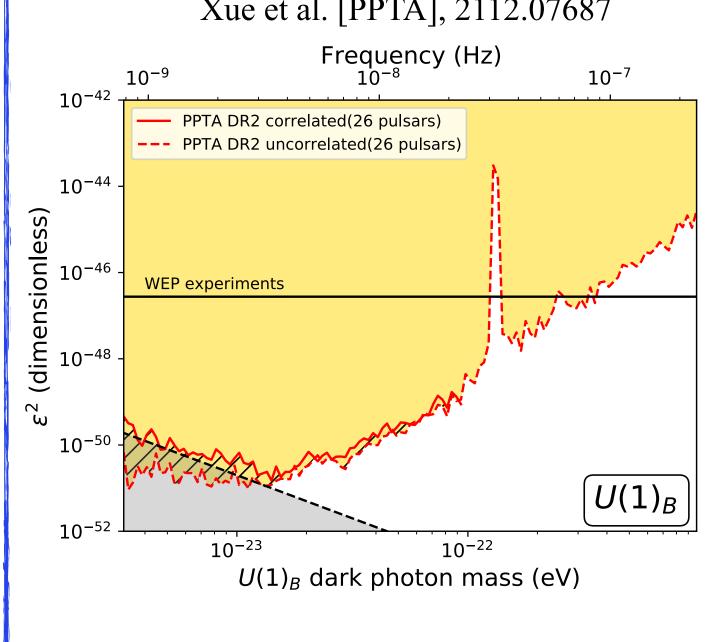
Doppler/Shapiro effects induced by transiting objects associated with DM substructure

Zurek et al., 1901.04490, 2005.03030, 2104.05717



Oscillating fifth force induced by ultralight dark photon DM

Xue et al. [PPTA], 2112.07687



Pulsar polarization arrays (PPAs)

◆ Pulsar radio emissions typically have strong linear polarization

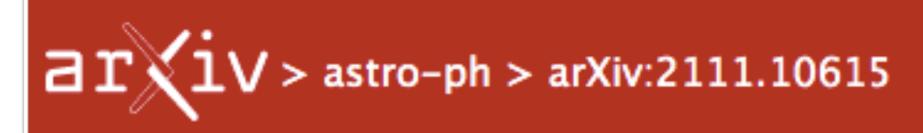
# NAME	P0 (s)		DM (cm^-3 pc)		RM (rad m^-2)	DIST (kpc)		
2 <u>J0006+1834</u> <u>cr</u> 3 <u>J0007+7303</u> <u>aa</u> 4 <u>J0011+08</u> <u>ds</u>	0.1153635682680 0.69374767047 0.3158731909 0.4+16 0.55287 0.78 0.240699038946	14 <u>cwp+17</u> 14 <u>cn95</u> 3 <u>awd+12</u> 0 <u>dsm+16</u> 11 <u>hlk+04</u>	218.6 11.41 * 24.9 30.405	6 wcp+18 55 bkk+16 0 * 0 dsm+16 13 bkk+16	* -20 * * -15.56	0 * 3 <u>npn+20</u> 0 * 0 * 10 <u>sbg+19</u>	6.357 0.860 1.400 5.399 1.776	ATNF Pulsar Catalo (Version: 1.66)	ogue
37	0.31835777079 0+00 0.0048654532114961 0.315 0.912 0.0018771818845850 0.19632137543003 0.19632137543006 0.19632137543006 0.19632137543006 0.19632137543006 0.19632137543006 0.19632137540006 0.19632137540006 0.19632137540006 0.19632137540006 0.19632137540006 0.19632137540006 0.19632137540006 0.19632137540006 0.1963211140006 0.196321140006 0.1963211140006 0.1963211140006 0.1963211140006 0.19632114	3 sjm+19 19 aab+21a 0 hrk+08 0 hrk+08 2 aaa+10b 16 ywml10 7 hlk+04 10 clm+05 13 hlk+04 5 mgf+19 15 bjd+06 8 hlk+04 5 psj+19 4 bjd+06 3 lbs+20	245.06 4.332772 75.65 37.7 13.76517 40.111 25.66 22.881 11.92577 19.7978 131.80 40.938 195.50 179.4 113.4	6 sjm+19 0 aab+21a 9 mss+20 3 mss+20 4 aaa+10b 3 hlk+04 3 hlk+04 12 bkk+16 4 srb+15 1 mgf+19 7 bjd+06 3 hlk+04 3 mss+20 1 bjd+06 2 antt94	-294 1.17 * * -38 -9.5 2 1.4 +6.6 21.6 25.2 -1.2 +122 126.9 57.7	2 npn+20 8 sbg+19 0 * 0 * 17 hmvd18 5 npn+20 4 h187 2 sbg+19 50 fdr15 1 mgf+19 50 hmvd18 4 h187 6 jsd+21 71 hmvd18 5 qmlg95	6.619 0.324 2.251 1.654 1.348 2.128 25.000 2.477 0.920 1.622 0.484 1.900 4.084 3.576 0.571	$RM = \frac{PA}{\lambda^2}$ rotation measure	pola — a wavele

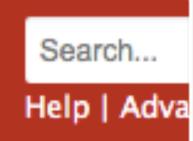
 Polarization profiles are usually obtained as a by-product of PTAs due to the crucial role played by polarization calibration for precise timing

polarization

angle

wavelength





Astrophysics > High Energy Astrophysical Phenomena

[Submitted on 20 Nov 2021]

Pulsar Polarization Arrays

Tao Liu, Xuzixiang Lou, Jing Ren

Pulsar timing arrays (PTAs) consisting of widely distributed and well-timed millisecond pulsars can serve as a galactic interferometer to measure gravitational waves. With the same data acquired for PTAs, we propose to develop pulsar polarization arrays (PPAs), to explore astrophysics and fundamental physics. As in the case of PTAs, PPAs are best suited to reveal temporal and spatial correlations at large scales that are hard to mimic by local noise. To demonstrate the physical potential of PPAs, we consider detection of ultralight axion-like dark matter (ALDM), through cosmic birefringence induced by its Chern-Simon coupling. Because of its tiny mass, the ultralight ALDM can be generated as a Bose-Einstein condensate, characterized by a strong wavy nature. Incorporating both temporal and spatial correlations of the signal, we show that PPAs have a potential to probe the Chern-Simon coupling up to $\sim 10^{-14}-10^{-17} {\rm GeV}^{-1}$, with a mass range $\sim 10^{-27}-10^{-21} {\rm eV}$.

Measured polarization angle:

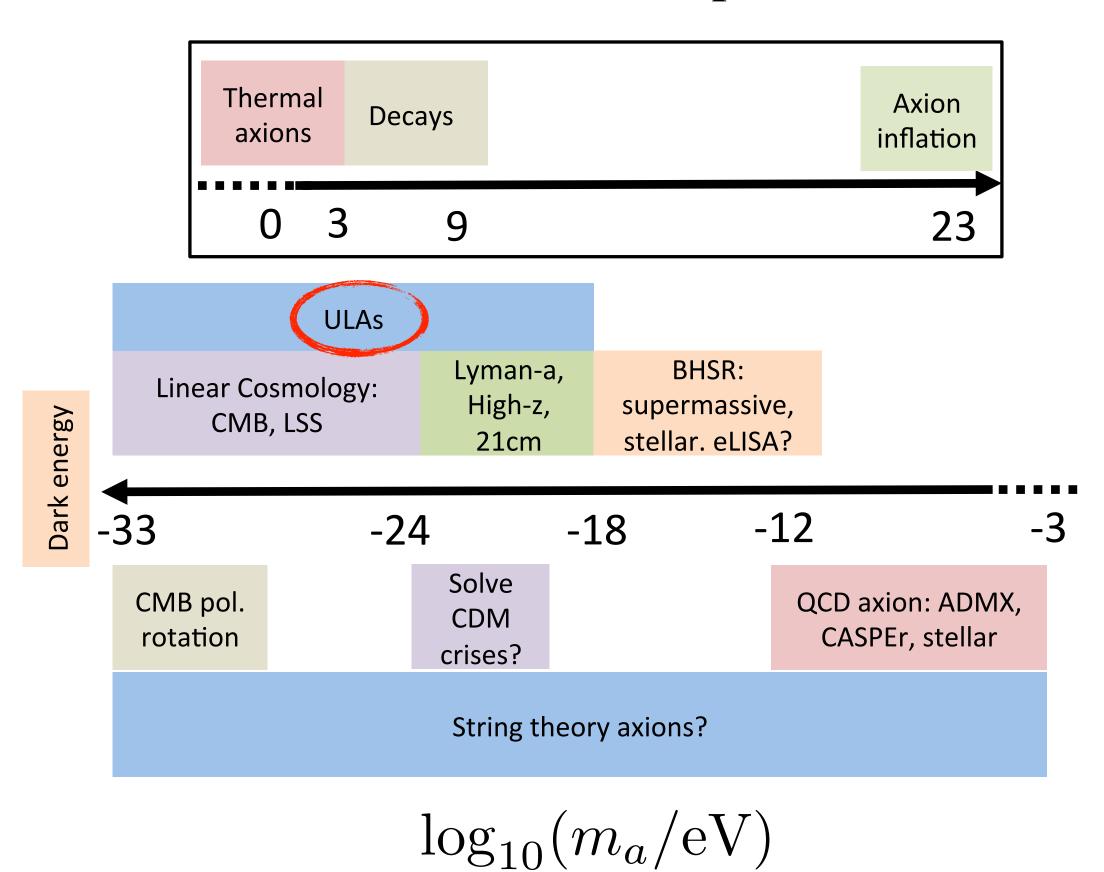
$$PA = PA_{source} + PA_{instru} + PA_{propa}$$

- Source: change of pulsar orientation/ magnetosphere; stochastic variation of single pulse profile (jitter noise)
- Instrument: related to PA calibration; radiometer noise
- **Propagation**: Faraday rotation caused by interstellar magnetic field/Earth's ionosphere; new physics effects??

Axion and axion-like particles (ALPs)

- QCD axion well motivated to solve the strong CP problem and also serve as an DM candidate; ALPs introduced in many BSM scenarios, e.g. string axions.
- ◆ Ultralight ALPs (m_a<10⁻¹⁸eV) can be generated as a Bose-Einstein condensate from misalignment; behave effectively as a classical scalar field
- ◆ Ultralight ALPs may serve as DM or DE during cosmic evolution. Here we focus on ALDM (e.g. fuzzy DM at m_a~10⁻²²eV, subdominant at m_a<10⁻²⁴eV)

ALPs landscapes



D. Marsh, Phys. Rept. 643, 1-79 (2016)

Physical properties of ALDM

◆ Local ALDM field made up by a large number of ALP classical fields with uncorrelated random phases Foster, Rodd, Safdi, PRD 97, no.12, 123006 (2018)

$$a(\mathbf{x},t) = \frac{\sqrt{\rho(\mathbf{x})}}{m_a} \sum_{d}^{\mathbf{density}} \frac{\mathbf{velocity \ element}}{\cos[\omega_a(t - \mathbf{v}_d \cdot \mathbf{x}) + \phi_d]}$$

- Non-relativistic limit: $\omega_a \approx m_a$; velocity distribution $f_{\mathbf{x}}(\mathbf{v})$ peaked around $|\mathbf{v}| \sim v_0 \sim 10^{-3} \text{c}$ (CDM velocity in our galaxy)
- Random nature: random amplitude α_d and phase ϕ_d (follow the Rayleigh and uniform distributions)
- ◆ ALPs effective action: gravitational interaction, coupling to SM particles...

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2} \nabla^{\mu} a \nabla_{\mu} a - \frac{1}{2} m_a^2 a^2 + \frac{g_{a\gamma}}{2} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{g_{af}}{2m_f} \nabla_{\mu} a (\bar{f} \gamma^{\mu} \gamma_5 f) + \dots \right]$$

Chern-Simons coupling

ALDM induced cosmic birefringence

When photons propagate in ALDM field *a*, the Chern-Simons coupling *g* corrects the dispersion relations of their two circular polarization modes in opposite way, yielding polarization angle (PA) rotation of linearly polarized photons

$$\omega_{\pm} \simeq k \pm g \left(\frac{\partial a}{\partial t} + \nabla a \cdot \frac{\mathbf{k}}{k} \right) \xrightarrow{\text{non-}} k \pm g \frac{\partial a}{\partial t}$$

$$\Delta \theta_{p,n} \approx \frac{g}{m_a} \sum_{d} \alpha_d \left\{ \sqrt{\rho_p f_p} \cos[m_a (t_n - L_p - \mathbf{v}_d \cdot \mathbf{x}_p) + \phi_d] - \sqrt{\rho_e f_e} \cos(m_a t_n + \phi_d) \right\} (\Delta v)^{\frac{3}{2}}.$$

- relies only on field profiles at two endpoints of photon traveling
- quasi-monochromatic oscillation around the frequency $\omega_a \approx m_a$
- no radio frequency dependence (FR increases with wavelength)

A variety of astrophysical light sources proposed before to constrain g

[Carroll et al., PRD, 1990; Antonucci, ARAA, 1993; Ivanov et al., JCAP, 2019; Fujita et al., PRL, 2019; Liu et al., PRD, 2020; Caputo et al., PRD, 2019; Chigusa et al., PLB, 2020; Chen et al., PRL, 2020 ...]

Yet, spatial correlations of the signal among individual sources not properly considered

ALDM induced signal for PPAs

PA rotation signal on PPAs with many pulsars

Liu, Lou, JR, 2021, [arXiv: 2111.10615]

• Construct a signal vector s from the time series of PA rotation for each pulsar in PPA

$$\mathbf{s} \equiv (\Delta \theta_{1,1}, ..., \Delta \theta_{1,N_1}, ..., \Delta \theta_{N,1}, ..., \Delta \theta_{N,N_N})^T$$

· Signal s follows a multivariable Gaussian distribution with zero mean and the covariance matrix

$$\Sigma_{p,n;q,m}^{(s)} \approx \frac{g^2}{m_a^2} \left\{ \rho_e \cos(m_a \Delta t) + \sqrt{\rho_p \rho_q} \cos[m_a (\Delta t - \Delta L)] \frac{\sin y_{pq}}{y_{pq}} + \dots \right\} \frac{y_{pq} = \Delta x/l_c}{l_c = 1/(m_a v_0)}$$
(coherence length)

Comparison	SGWB	ALDM (PTAs and PPAs)
Earth-term correlation	quadrupolar correlation	monopolar correlation (universal)
pulsar-term correlation	pulsar-terms suppressed in $L\gg 1/\omega$ limit; spatial correlations rapidly degrade	spatial correlations degrade slower $(L\gg l_c\gg 1/m_a)$; encode DM density dep.; enhanced at galactic center

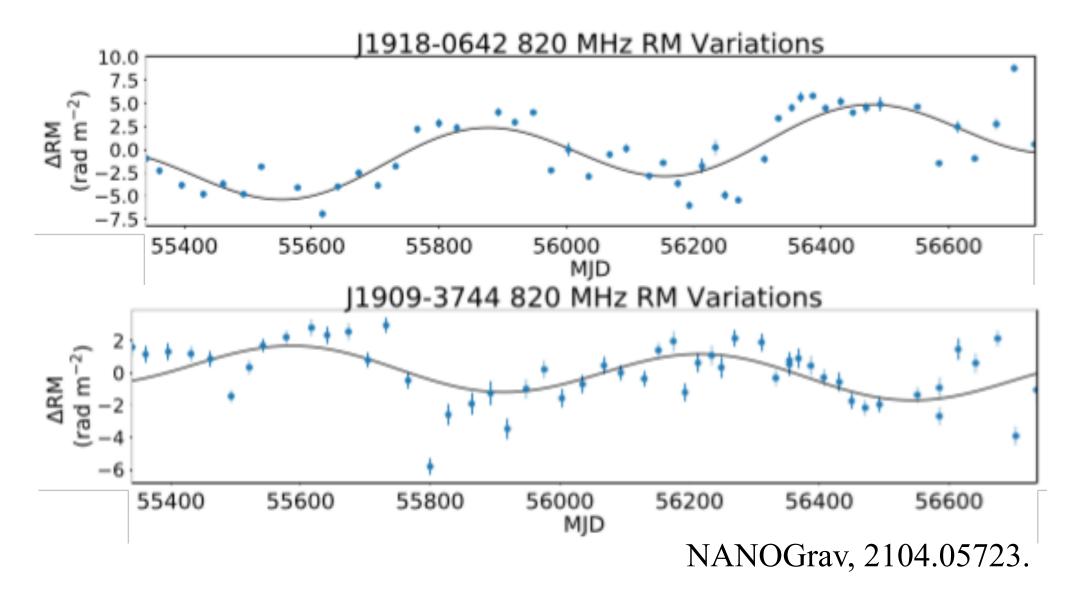
Pulsar correlations

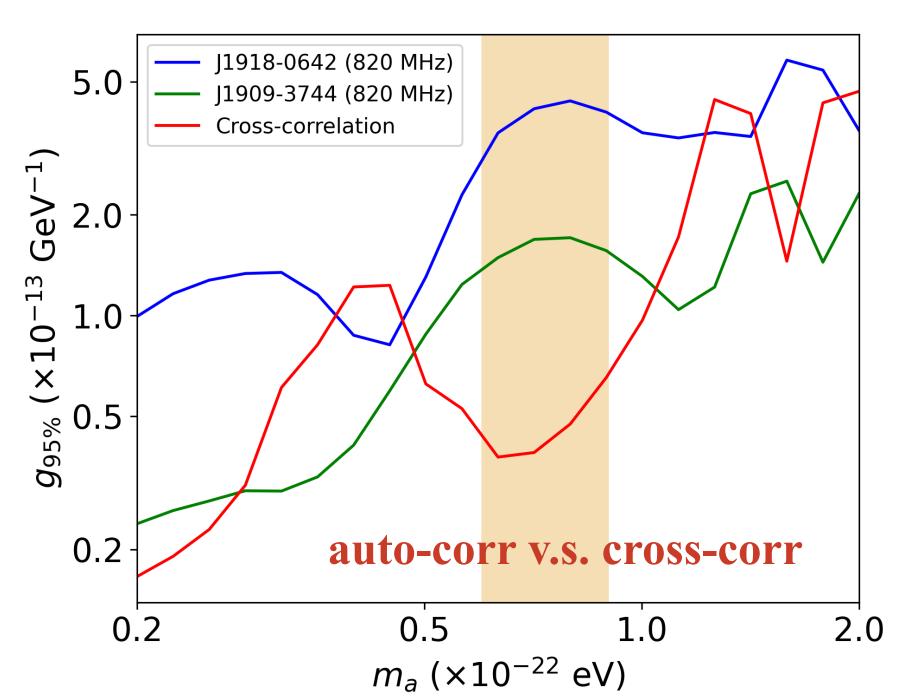
Considering background n dominated by white noise, data d = s + n follows a multivariate Gaussian distribution with zero mean and covariance matrix $\Sigma = \Sigma^{(s)} + B$

- Likelihood function: $\mathcal{L}(\theta|\mathbf{d}) = \det[2\pi\mathbf{\Sigma}]^{-1/2} \exp\left[-\frac{1}{2}\mathbf{d}^T \cdot \mathbf{\Sigma}^{-1} \cdot \mathbf{d}\right]$
- Exclusion TS on the coupling g: $q(g, m_a) \equiv 2[\ln \mathcal{L}(\hat{g}, m_a|\mathbf{d}) \ln \mathcal{L}(g, m_a|\mathbf{d})]$ (\hat{g} maximizes likelihood)

Both auto-correlation of individual pulsars and cross-correlation among different pulsars contribute. Cross-correlation highly valuable to distinguish spatial correlated and uncorrelated contributions

Detected oscillating PA rotation with period ~1yr





Pulsar correlations

Asimov form of exclusion TS for projection: averaging different data realizations under the background only assumption

$$\langle q \rangle = \text{Tr} \left[\mathbf{B} (\mathbf{\Sigma}^{-1} - \mathbf{B}^{-1}) \right] + \ln \det \mathbf{\Sigma} / \ln \det \mathbf{B} + \dots$$

$$\langle q \rangle \approx \frac{1}{2} \sum_{p,q} \frac{1}{\lambda_p \lambda_q} \text{Tr} \left(\mathbf{\Sigma}_{pq}^{(s)} \mathbf{\Sigma}_{qp}^{(s)} \right)$$
 (small signal limit; λ noise variance)

Auto-correlation

$$\operatorname{Tr}\left(\mathbf{\Sigma}_{pp}^{(s)}\mathbf{\Sigma}_{pp}^{(s)}\right) \sim \frac{g^4}{m_a^4} N_p^2 \left[\rho_e + \rho_p - 2\sqrt{\rho_e \rho_p} \cos(m_a L_p)\right]$$
of data points for *p*-th pulsar $\left[\frac{\sin y_{ep}}{y_{ep}}\right]^2$. $y_{ep} = L_p/l_c$

- N_p enhancement reflects temporal correlation, insensitive to distribution of sampled points
- Earth-pulsar correlation (3rd) suppressed if $L_p \gg l_c$

Cross-correlation

$$\operatorname{Tr}\left(\mathbf{\Sigma}_{pq}^{(s)}\mathbf{\Sigma}_{pq}^{(s)}\right) \sim \frac{g^4}{m_a^4} N_p N_q \left[\rho_e^2 + \rho_p \rho_q \frac{\sin^2 y_{pq}}{y_{pq}^2}\right] + 2\rho_e \sqrt{\rho_p \rho_q} \cos(m_a \Delta L) \frac{\sin y_{pq}}{y_{pq}} + f(y_{ep}, y_{eq}).$$

- Earth-terms cross-correlations (1st) universal
- Pulsar-terms cross-correlations (2nd) dominate for pulsars around galactic center

PPA projected sensitivity to ALDM

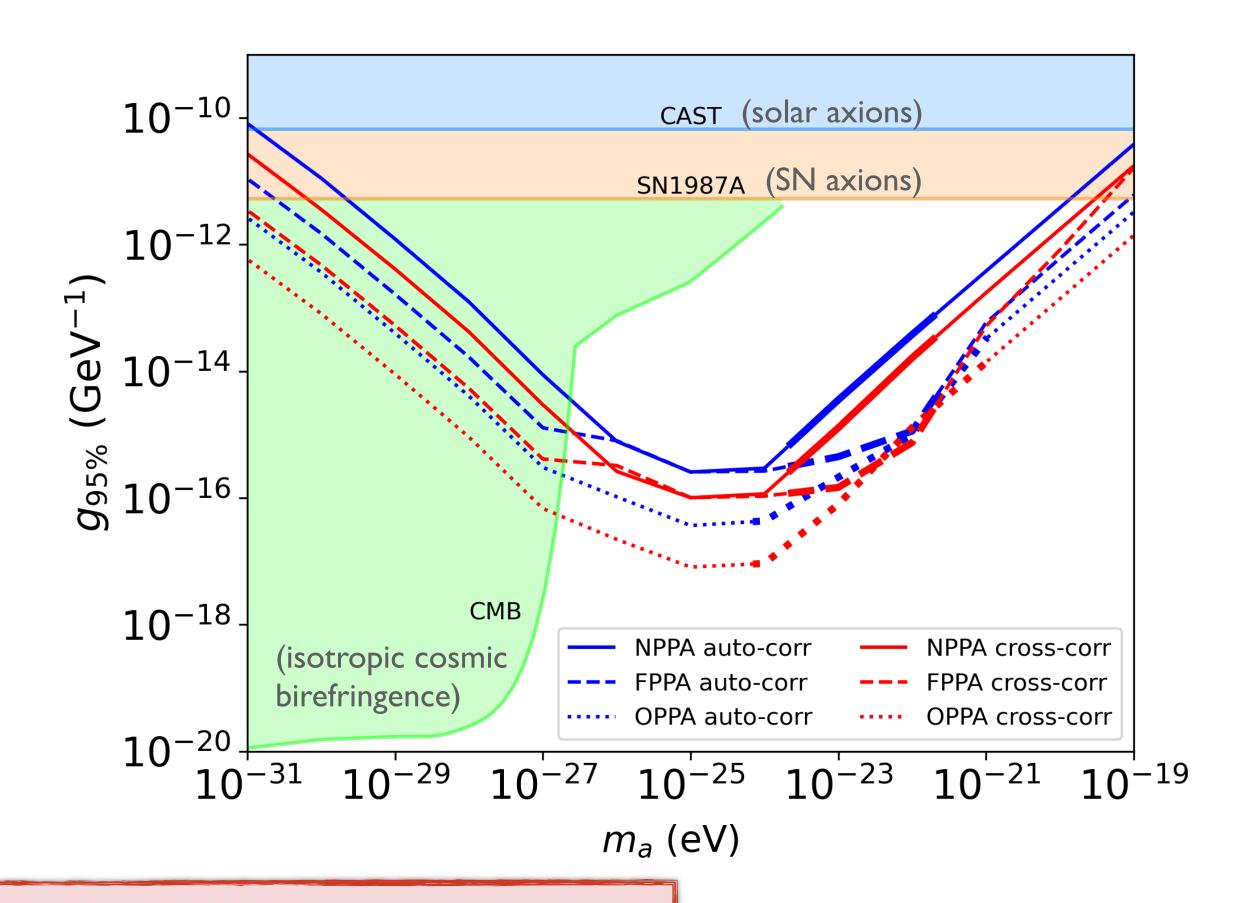
◆ Three PPA scenarios (different stage of PTAs)

NPPA: 100 MSPs around 1kpc (current PTAs)

FPPA: 100 MSPs in galactic bulge (*FAST/SKA era*)

OPPA: 1000 MSPs following ATNF (*FAST/SKA era*)

- ◆ Galactic ALDM density profile: soliton formed by quantum pressure $(r < l_c, \rho_c \propto m_a^2)$ + NFW $(r > l_c)$; CMB constraints on relic abundance imposed
- → Pulsar distance uncertainty (~20%) marginalized in constraints estimate



- · Cross-correlation limits are typically stronger than auto-correlation ones
- · As ma decreases, the limit decreases (fuzzy DM), stay flat and increases
- Projected PPA limits form a great complementarity with the existing bounds

Summary and Outlook

- We propose development of PPAs with the same data acquired for PTAs. PPAs can be used to detect ultralight ALDM induced cosmic birefringence with special temporal and spatial correlations. Projected limits on its Chern-Simons coupling form a great complementarity with existing bounds
- ◆ Real data analysis with PPTA polarization data Liu, Luu, JR, Shu, Xue, Zhao 2022, ongoing...
- ◆ Cross-correlation of oscillating signals induced by ALDM on both PPAs and PTAs Liu, JR, Xu, 2022, ongoing...
- More physical targets for PPAs as a new tool for exploring fundamental physics?
- Synergize PPAs with other observations to further distinguish different scenarios?

Thank You!